

Saturation model in UPCs

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Open Questions in Photon-Induced Interactions
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Exclusive vector meson production in UPCs

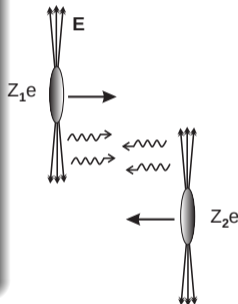
Advantages

- No net color charge transfer: at least two gluon exchange

$$\left. \frac{d\sigma^{\gamma^* A \rightarrow VA}}{dt} \right|_{t=0} = \frac{16\pi^3 \alpha_s^2 \Gamma_{ee}}{3\alpha_{em} M_V^5} \left[\mathbf{x}g(\mathbf{x}, Q^2) \right]^2$$

In practice connection to PDFs complicated (see Vadim's talk) [Ryskin, 1993](#)

- Measure J/ψ p_T = total momentum transfer Δ
= Fourier conjugate to \mathbf{b}
 \Rightarrow Access to spatial structure



UltraPeripheral heavy ion Collisions (UPC)

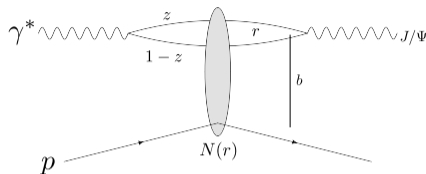
- At $|b_T| > R_1 + R_2$ one nucleus acts as a photon source

High energy (nuclear) DIS before the Electron Ion Collider!

Vector meson production at high energy/small x

High energy factorization:

- 1 $\gamma \rightarrow q\bar{q}$ splitting,
wave function $\Psi^\gamma(r, Q^2, z)$
- 2 $q\bar{q}$ dipole scatters elastically: $N(r, x, b)$
- 3 $q\bar{q} \rightarrow J/\Psi$,
wave function $\Psi^V(r, z)$



Diffractive scattering amplitude

$$\mathcal{A}^{\gamma^* p \rightarrow V p} \sim \int d^2 b d z d^2 r \Psi^\gamma \Psi^V(r, Q^2, z) e^{-i\mathbf{b} \cdot \mathbf{\Delta}} N(r, x, b)$$

- $N(r, x, b)$: universal dipole-proton scattering amplitude, resumming multiple scattering
Convenient degrees of freedom at small x

Phenomenology at the moment: LO (resumming $\alpha_s \ln 1/x$), NLO is coming...

H.M, J. Penttala, [arXiv:2104.02349](https://arxiv.org/abs/2104.02349) [hep-ph]

Target shape from coherent diffraction

Coherent cross section: target remains on the same quantum state

Average interaction of states that diagonalize the scattering matrix

- These states are $q\bar{q}$ dipoles with fixed size \mathbf{r} , probing fixed target configuration Ω
- Good, Walker, PRD 120, 1960

Coherent cross section

$$\frac{d\sigma^{\gamma^* A \rightarrow VA}}{dt} \sim |\langle \mathcal{A}^{\gamma^* A \rightarrow VA} \rangle_{\Omega}|^2$$

- Cross section probes average \mathbf{b} dependence of the scattering amplitude = target geometry

$$\langle \mathcal{A}^{\gamma^* P \rightarrow VP} \rangle_{\Omega} \sim \int d^2\mathbf{b} d^2\mathbf{z} d^2\mathbf{r} \psi^{\gamma^*} \psi^V(|\mathbf{r}|, z, Q^2) e^{-i\mathbf{b} \cdot \mathbf{\Delta}} \langle N(|\mathbf{r}|, \mathbf{x}, \mathbf{b}) \rangle_{\Omega}$$

Incoherent diffraction = target dissociation

Incoherent cross section

- Target final state $|f\rangle \neq$ initial state $|i\rangle$
- No net color charge transfer: rapidity gap between J/Ψ and target remnants

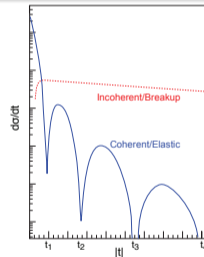
$$\begin{aligned}\sigma_{\text{incoherent}} &\sim \sum_{f \neq i} |\langle f | \mathcal{A} | i \rangle|^2 \\ &= \sum_f \langle i | \mathcal{A} | f \rangle^\dagger \langle f | \mathcal{A} | i \rangle - \langle i | \mathcal{A} | i \rangle^\dagger \langle i | \mathcal{A} | i \rangle\end{aligned}$$

Average over initial states:

$$\sigma_{\text{incoherent}} \sim \langle |\mathcal{A}|^2 \rangle_\Omega - |\langle \mathcal{A} \rangle_\Omega|^2$$

Incoherent cross section = covariance of $\mathcal{A} \gamma^* A \rightarrow VA$

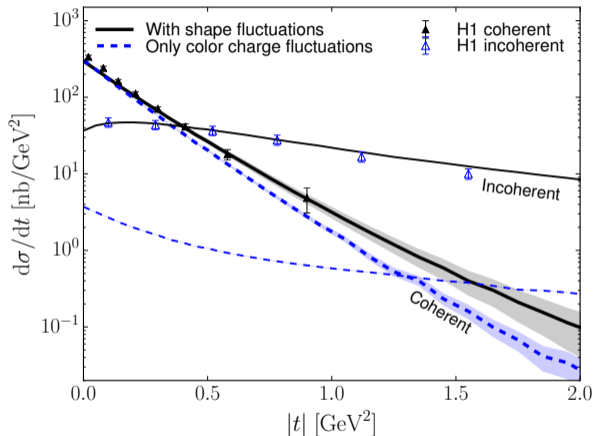
- Amount of event-by-event fluctuations in target configurations Ω



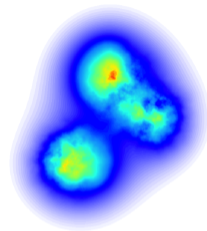
Miettinen, Pumplin, PRD 18, 1978,

Caldwell, Kowalski, PRC81 (2010) 025203

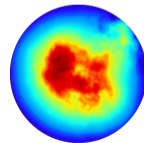
HERA data at small x : $\gamma + p \rightarrow J/\psi + p$ at $W = 75$ GeV ($x_{\mathbb{P}} \approx 10^{-3}$)



Fluctuations

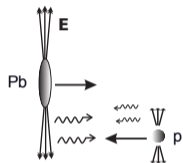


Round



HERA data well understood

(incoherent with fluctuating substructure) [H.M., B. Schenke, 1607.01711](#)



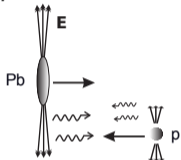
Ultraperipheral collisions in pA

What happens to the incoherent cross section at high W ?

Ultrapерipheral $p + A$:

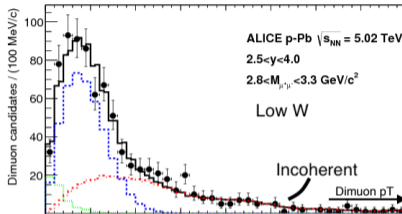
Photon flux $\sim Z^2$

$\gamma + p$ dominates



Low energy $\gamma - A$: coherent and incoherent visible

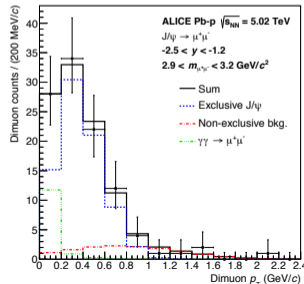
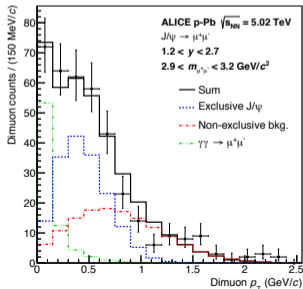
ALICE: 1406.7819



Larger COM energies:
incoherent $\rightarrow 0$ (?)
 \Rightarrow smoother proton?

ALICE:1809.03235

$x \sim 10^{-2} \rightarrow 10^{-5}$



Medium \rightarrow high energy

Can we understand the energy dependence?

Approach 1: parametrize the number of hot spots

Small- x gluon emissions increase the number of hot spots [Cepila, Contreras, Tapia Takaki, 1608.07559](#)

$$N_{hs}(x) \sim x^{p_1} (1 + p_2 \sqrt{x})$$



Approach 2: Solve small- x evolution equations

Evolve proton structure by solving evolution perturbatively

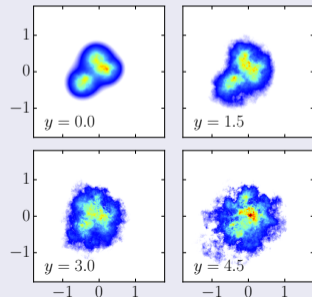
- BK eq. with impact parameter

[Berger, Stasto, 1106.5740](#), [Cepila, Contreras, Matas, 1812.02548](#)

- JIMWLK eq. [Schlichting, Schenke, 1407.8458](#), [H.M., Schenke, 1806.06783](#)

Fit HERA F_2 and exclusive data. [H.M, Schenke, 1806.06783](#)

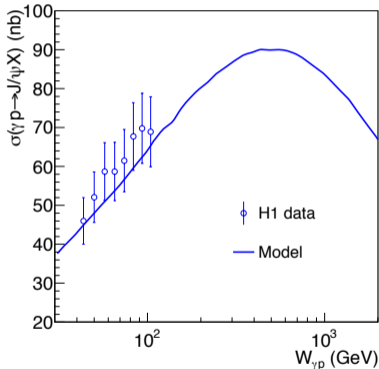
Difficulty: regulating confinement effects



Towards small x : $\gamma + p \rightarrow J/\psi + p^*$

Increasing # of hot spots with energy:
Smoother proton, less fluctuations

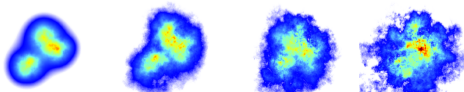
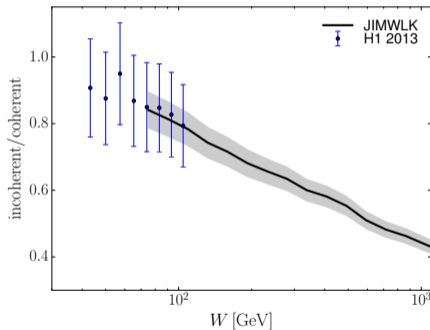
Cepila, Contreras, Tapia Takaki, 1608.07559



Heikki Mäntysaari (JYU)

JIMWLK evolution event-by-event
Includes also growing RMS size

H.M., Schenke, 1806.06783

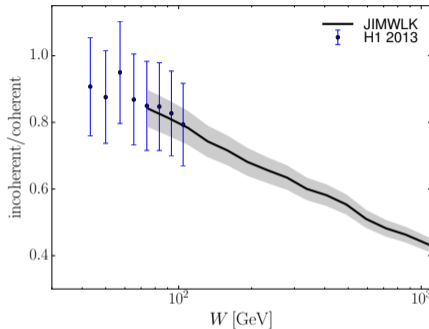
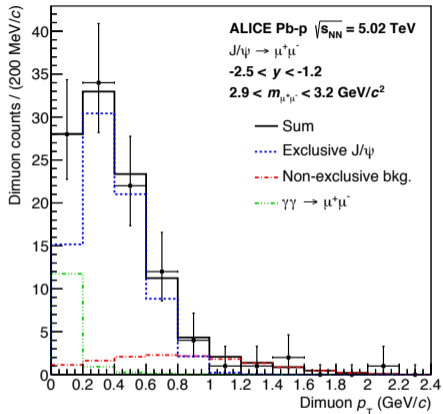


UPC

27.4.2021

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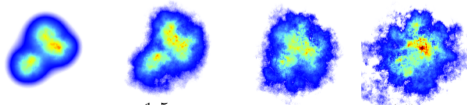
Small x in $\gamma + p \rightarrow J/\psi + p^*$



H.M, Schenke, 1806.06783

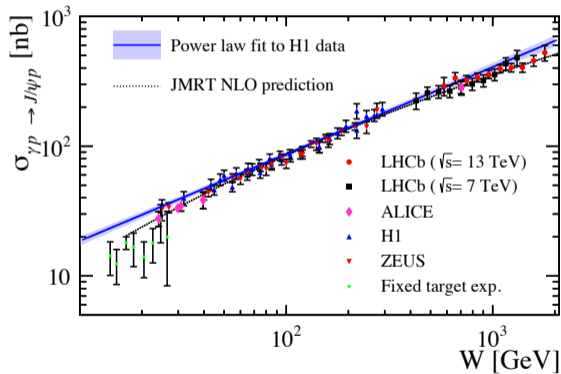
ALICE data qualitatively compatible with e-b-e fluctuating geometry evolution

- Final data at cross section level?
- Are we approaching the black disc limit?



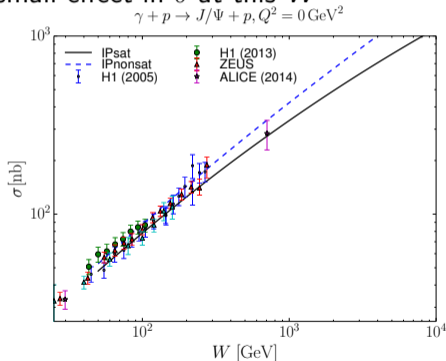
Saturation in $\gamma + p \rightarrow J/\psi + p$?

LHCb measurement in ultra peripheral $p + p$
 Generically saturation effects \Rightarrow deviations
 from the $\sigma \sim W^\delta$ behaviour

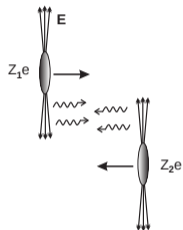


LHCb, 1806.04079

- Small hints at highest W ?
- Saturation effects expected to have only a small effect in δ at this W

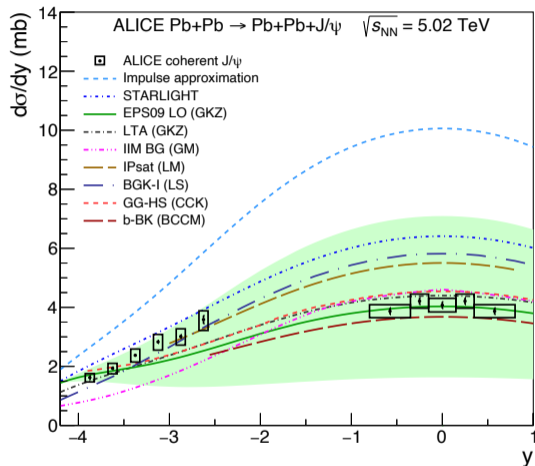


H.M, P. Zurita, 1804.05311



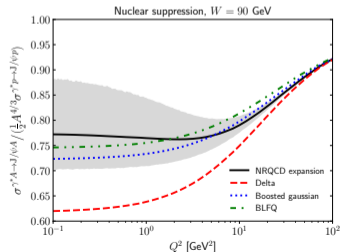
Ultraperipheral collisions in AA

Coherent diffraction today



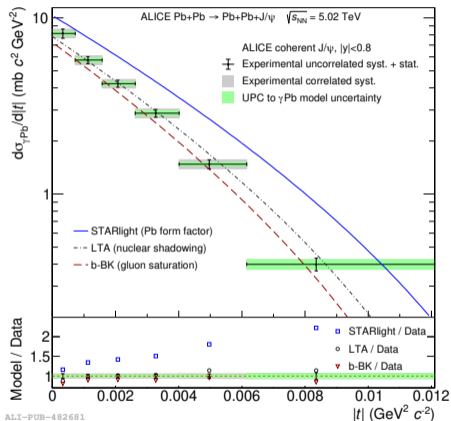
- Different calculations with saturation effects relatively successful
- Flat(?) rapidity dependence in $-2 < y < 0$???

Also: normalization uncertainty from J/ ψ wf, does not completely cancel in nuclear suppression ratio (=ratio to imp. approx.)



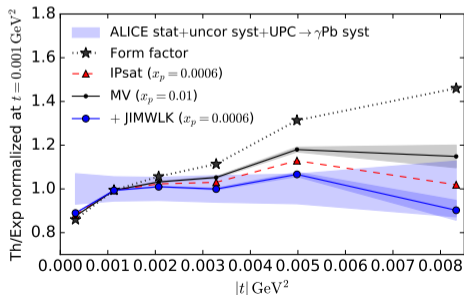
Measurements not statistic limited anymore

Coherent t spectra



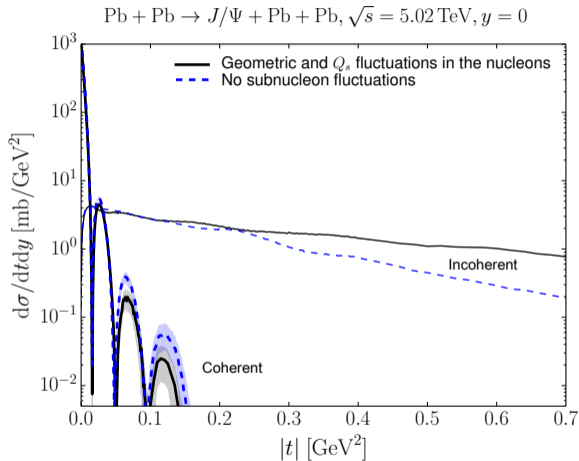
ALICE, 2101.04623

- First t spectra measured!
- (Would be interesting to see coh and incoh)
- Much steeper t spectra than FT of WS
- Steeper slope expected from saturation
- Especially the smallest t bin?



With T. Lappi, F. Salazar and B. Schenke

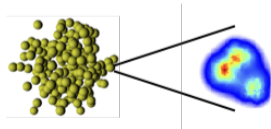
Accessing fluctuations at different scales



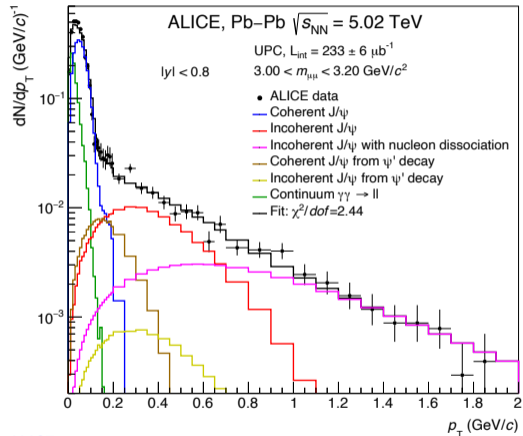
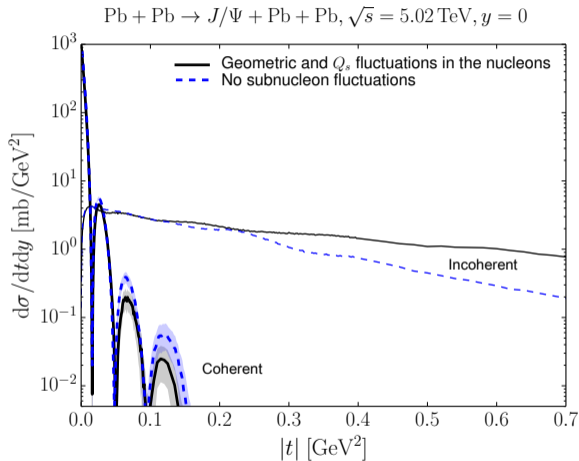
- $\sqrt{|t|}$ is conjugate to b_T
- Small $|t|$: fluctuations of nucleon positions
- Large $|t|$: fluctuations at subnucleon scale
- Incoherent slope changes at

$$|t| \approx 0.25 \text{ GeV}^2 \sim 0.4 \text{ fm}$$

which is the hot spots size from HERA



Accessing fluctuations at different scales



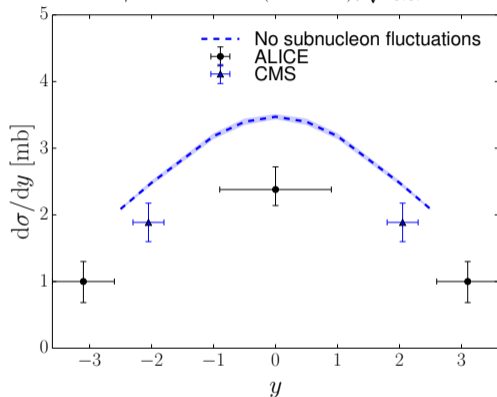
ALICE, 2101.04577

H. M., B. Schenke, arXiv:1703.09256

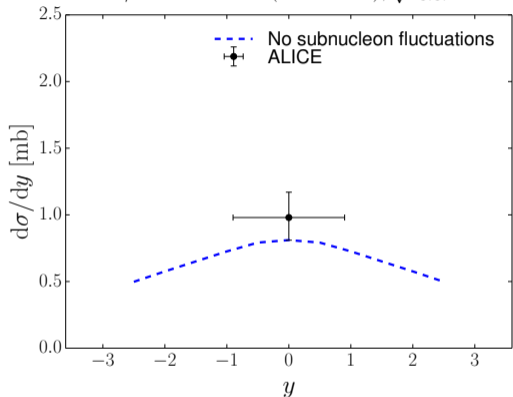
Qualitatively consistent with the ALICE data, but want cross section level spectra!

Comparison to LHC data, no subnucleonic fluctuations

Pb + Pb $\rightarrow J/\Psi + \text{Pb} + \text{Pb}$ (coherent), $\sqrt{s_{NN}} = 2760$ GeV



Pb + Pb $\rightarrow J/\Psi + \text{Pb} + \text{Pb}^*$ (incoherent), $\sqrt{s_{NN}} = 2760$ GeV

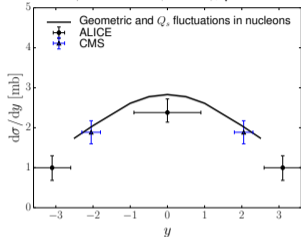


- Only fluctuations of nucleon positions from Woods-Saxon:
Coherent cross section overestimated and incoherent underestimated
- Overall normalization uncertainty from the J/Ψ wave function

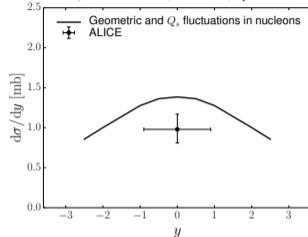
H.M., B. Schenke, arXiv:1703.09256

Comparison to LHC data, with subnucleon fluctuations

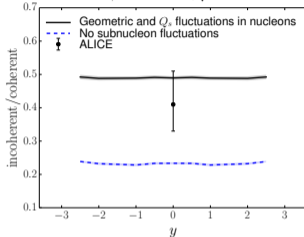
Pb + Pb $\rightarrow J/\psi + \text{Pb} + \text{Pb}$ (coherent), $\sqrt{s_{NN}} = 2760$ GeV



Pb + Pb $\rightarrow J/\psi + \text{Pb} + \text{Pb}^*$ (incoherent), $\sqrt{s_{NN}} = 2760$ GeV



Pb + Pb $\rightarrow J/\psi + \text{Pb} + \text{Pb}$, $\sqrt{s_{NN}} = 2760$ GeV



- Calculations with nucleon substructure consistent with the LHC data.
- Dependence on the J/ψ wave function mostly cancels in the ratio
- Note: no W dependent geometry included
- Missing: y and \sqrt{s} dependence of the incoherent cross section
How does the incoherent cross section and e-b-e fluctuations depend on $x_{\mathbb{P}}$?

- Vector meson production from HERA relatively well understood
- Lots of precise data from LHC (and RHIC!)
- Interesting observations include
 - Large nuclear suppression compared to impulse approximation
 - Disappearing incoherent cross section in high energy $\gamma + p$??
 - Total $\gamma + p \rightarrow J/\psi + p$ cross section: deviating from W^δ law from HERA???
 - Low- t part of the $\gamma + \text{Pb} \rightarrow J/\psi + \text{Pb}$ spectra????
- Wish list:
 - t spectra in $\gamma + p$ and $\gamma + A$ (coherent+incoherent ok)
 - y and \sqrt{s} dependence of the incoherent cross section
- Clear signatures of non-linear effects, but also need for precision theory
 - NLO (coming, see [H.M, J. Penttala, 2104.02349](#))
 - Vector meson wave function beyond phenomenological models
(see [T. Lappi, H.M, J. Penttala, 2006.02830](#); [Y. Li, P. Maris, J. Vary, 1704.06968](#))

Significant nuclear effects seen in $\gamma + A \rightarrow J/\psi + A$

Enhance non-linearities:

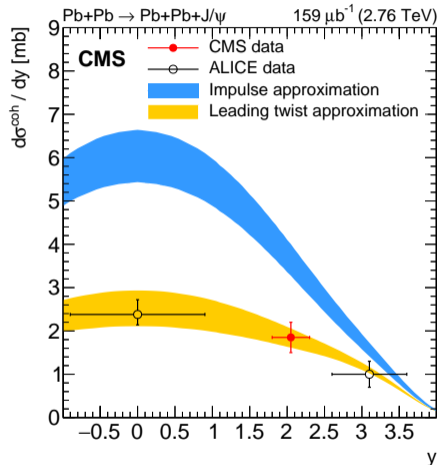
$$Q_{s,A}^2 \sim A^{1/3} x^{-\lambda}$$

(larger A is cheaper than smaller x)

The first UPC measurements:

clear nuclear suppression

Impulse approximation = scaled $\gamma + p$

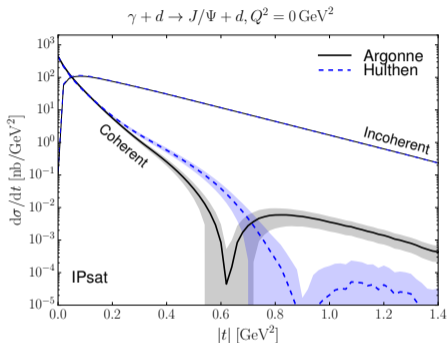


$$x_{\mathbb{P}} = M_V e^{\pm y} / \sqrt{s}$$

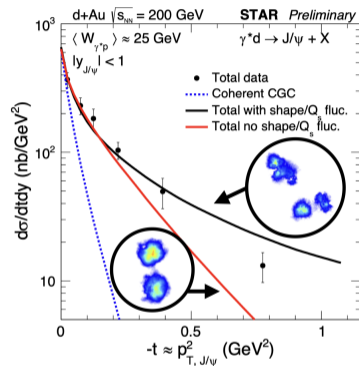
CMS, 1605.06966

Light ions

RHIC:
UPC in $d + Au$
(and $^3\text{He} + Au$)
 $x_{\mathbb{P}} \sim 10^{-2}$



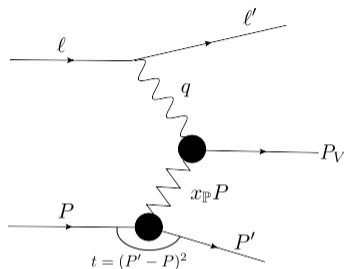
Distribution of small- x gluons in d :
Does it follow nucleon positions?
Details of the deuteron wf at small x



Nucleon substructure fluctuations in d
Preferred by STAR data (coh+incoh)

H.M., Schenke, 1910.03297; STAR 2009.04860

Diffractive scattering

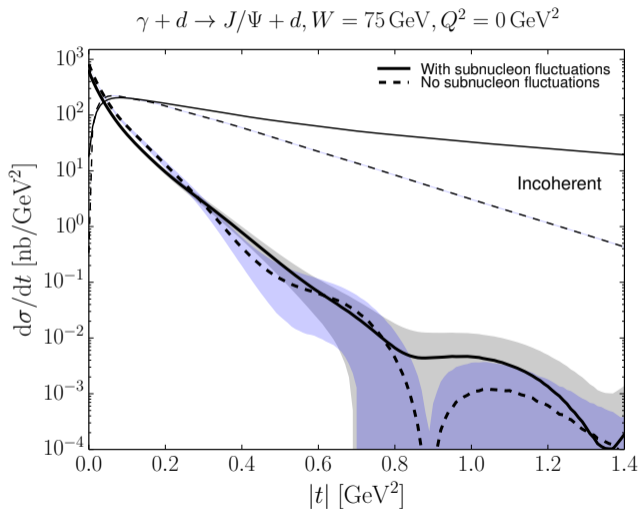


Diffractive events: no exchange of color charge (\rightarrow rapidity gap)

- Target remains intact: *coherent diffraction*, small $|t|$.
Probes average distribution of gluons.
- Target breaks up: *incoherent diffraction*, larger $|t|$.
Sensitive to fluctuations.

Target: proton or nucleus

Subnucleon fluctuations in deuteron



Dipole-proton scattering: IPsat model

An impact parameter dependent dipole amplitude

$$N(r, x, b) = 1 - \exp \left[-\frac{\pi^2}{2N_c} \alpha_s x g(x, \mu^2) T_p(b) r^2 \right]$$

- Fit to HERA data (F_2): initial condition for the DGLAP evolution of $xg(x, \mu^2)$ (Kowalski, Teaney 2003; Rezaeian et al, 2013, H.M, Zurita, 1804.05311)
- Proton profile T_p : Gaussian, width B_p

$$T_p(b) = \frac{1}{2\pi B_p} e^{-b^2/2B_p}$$

Adding color charge fluctuations: IP-Glasma

- Obtain saturation scale $Q_s(b_T)$ from IPsat (with constituent quarks)
- Sample color charges $\rho(b_T) \sim Q_s(b_T)$
- Solve Yang-Mills equations to obtain Wilson lines

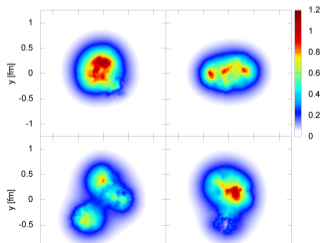
$$V(x_T) = P \exp \left(-ig \int dx^- \frac{\rho(x^-, x_T)}{\nabla^2 + m^2} \right)$$

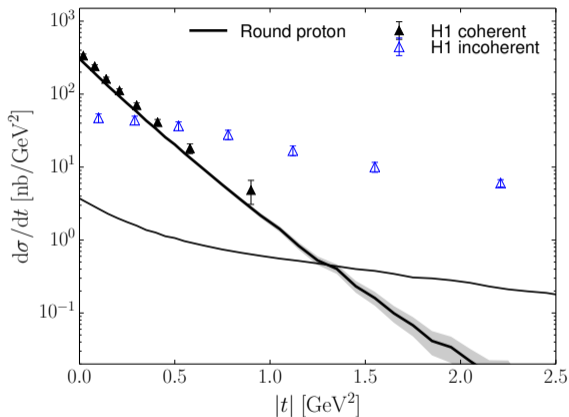
- Dipole amplitude: $N(x_T, y_T) = 1 - \text{Tr} V(x_T) V^\dagger(y_T) / N_c$
- Fix parameters B_{qc} , B_q and m with HERA data

Example configurations:

$$1 - \text{Re}(\text{Tr} V(x_T)) / N_c$$

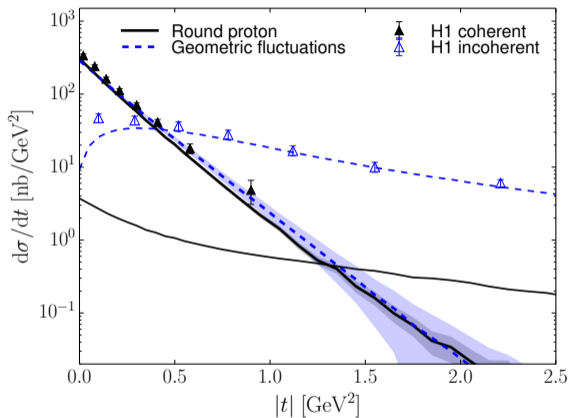
H.M., B. Schenke, arXiv:1603.04349





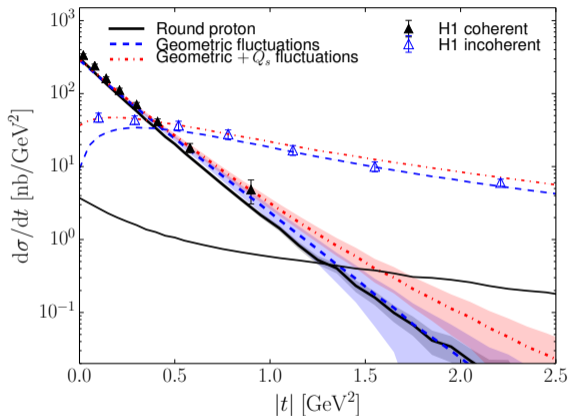
H.M., B. Schenke, in preparation and arXiv:1603.04349

- Color charge fluctuations alone are not enough



H.M., B. Schenke, arXiv:1603.04349

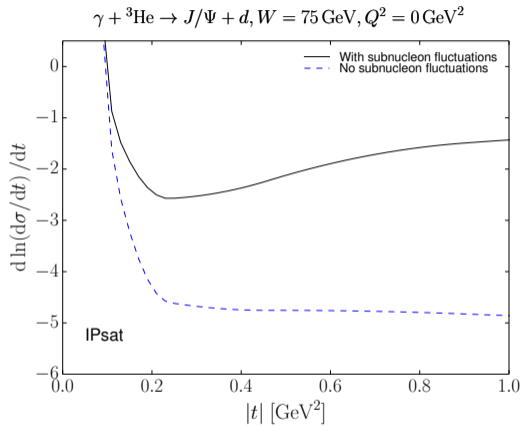
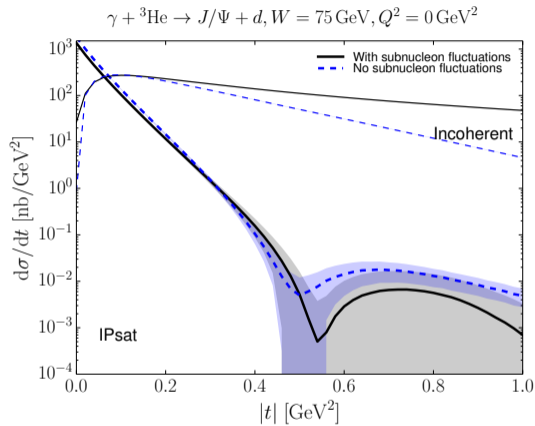
- Large geometric fluctuations are needed



H.M., B. Schenke, arXiv:1603.04349

- Q_s fluctuations improve description at small $|t|$

Similar results for ^3He



- Subnucleon fluctuations alter the incoherent $|t|$ slope, and significantly increase the incoherent cross section
- EIC will open a new window to the gluonic structure of light ions, also!