



# Jet quenching at RHIC and the LHC: a status report

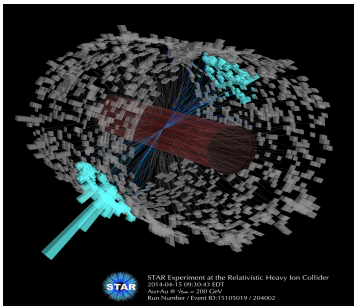
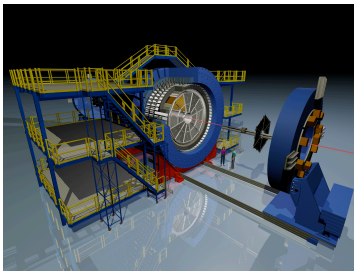


*biased*

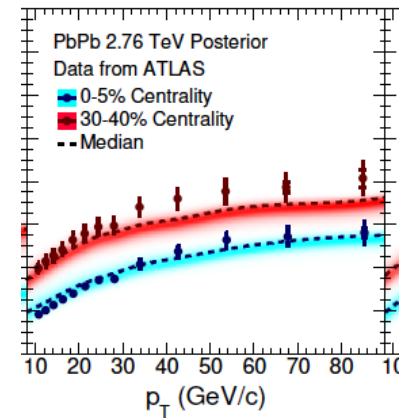
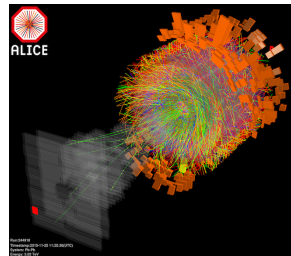
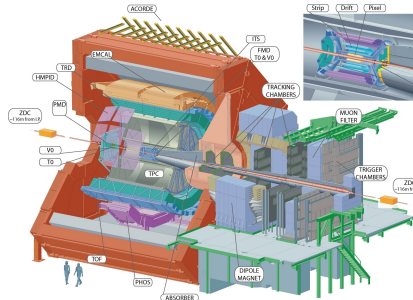
*Peter Jacobs*

*Lawrence Berkeley National Laboratory*

STAR@RHIC

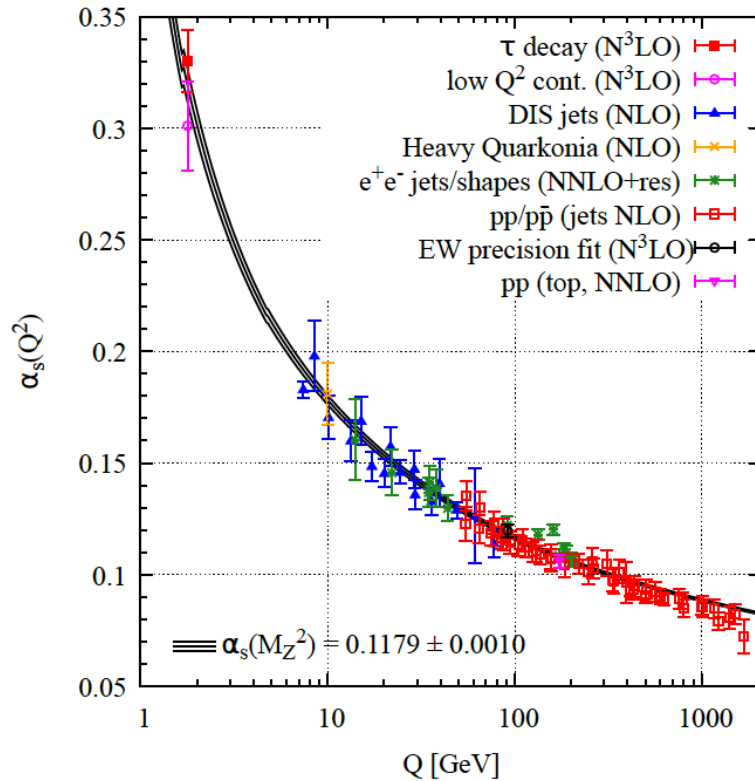
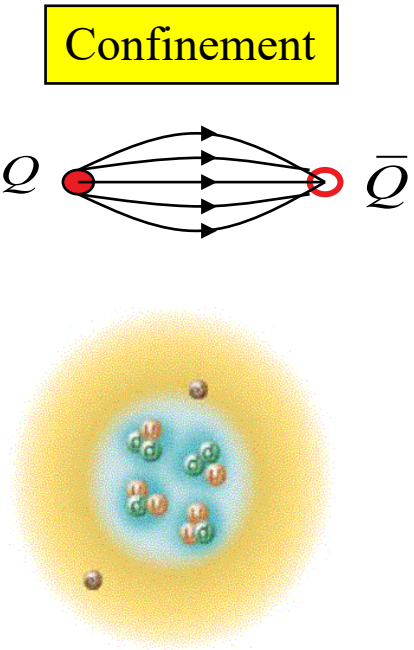


ALICE@LHC



# Hallmark of QCD: running of $\alpha_s$

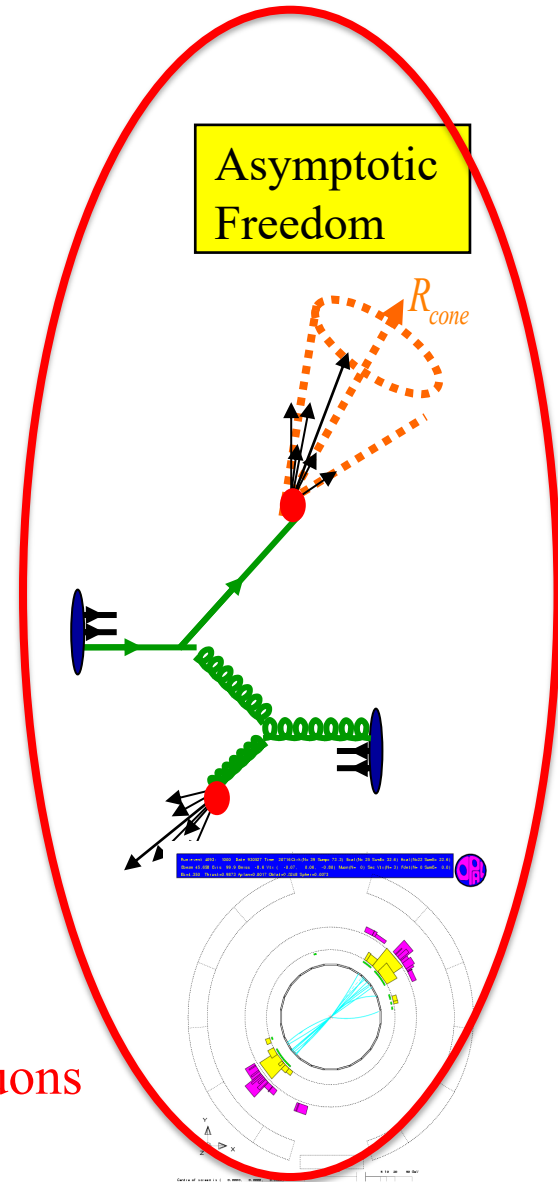
PDG2020



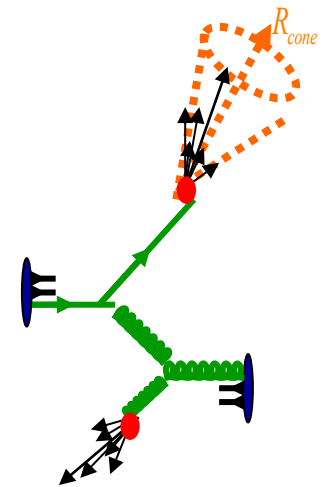
Low momentum

High momentum

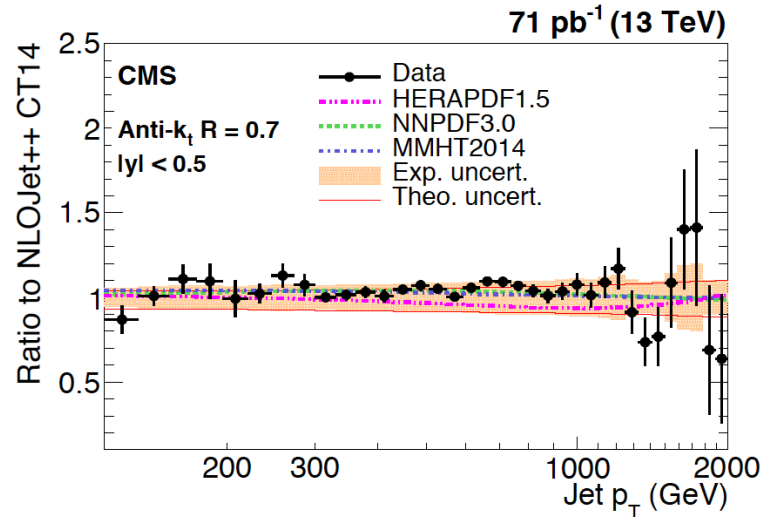
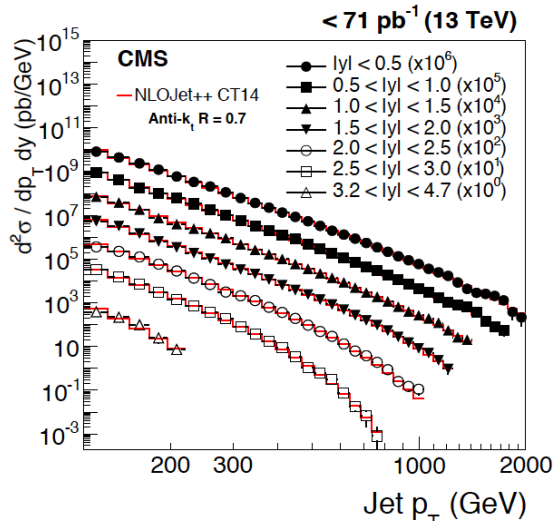
Scattering of energetic quarks and gluons  
→ jets



# Testing perturbative QCD: inclusive jet production in p+p collisions



*CMS, Eur. Phys. J C 76 (2016) 451*



## Magnificent achievement of QCD

- needed 30 years of development in theory, experiment, and algorithms to connect the two

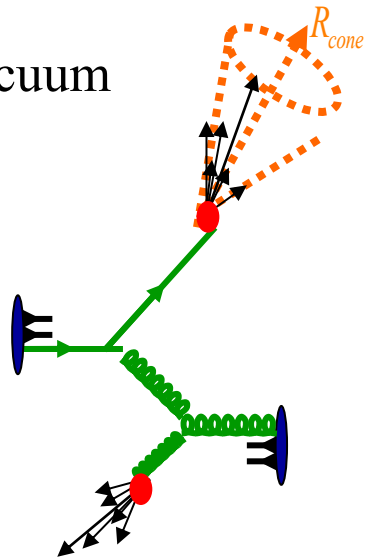
## Infrared and collinear-safe (IRC-safe) jet reconstruction algorithms:

- Integrate out all hadron degrees of freedom
- Same procedures applied to pQCD theory and experiment
- Enables direct, precise and improvable comparison of theory/experiment

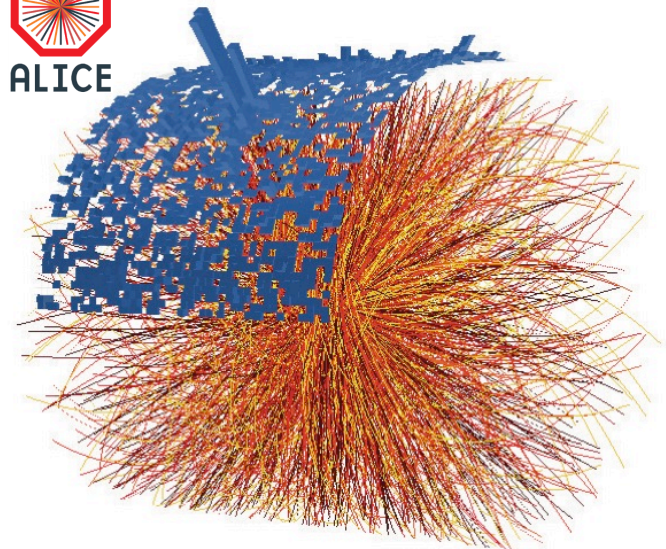
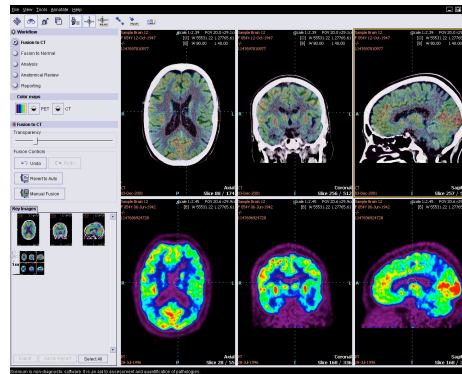
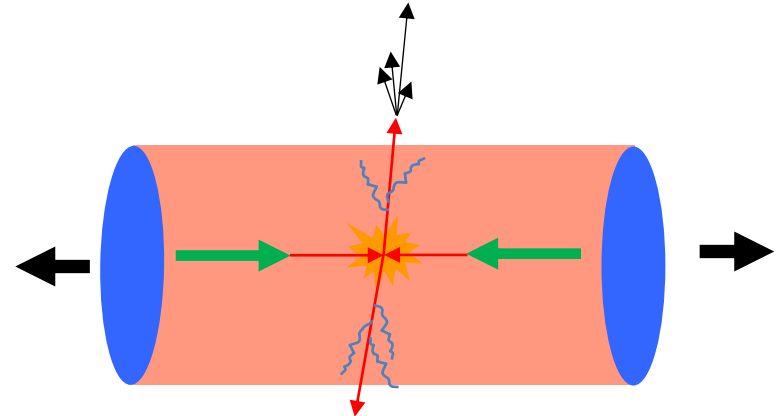
→ jets measure partons

# Jets in QCD matter

Jets in vacuum

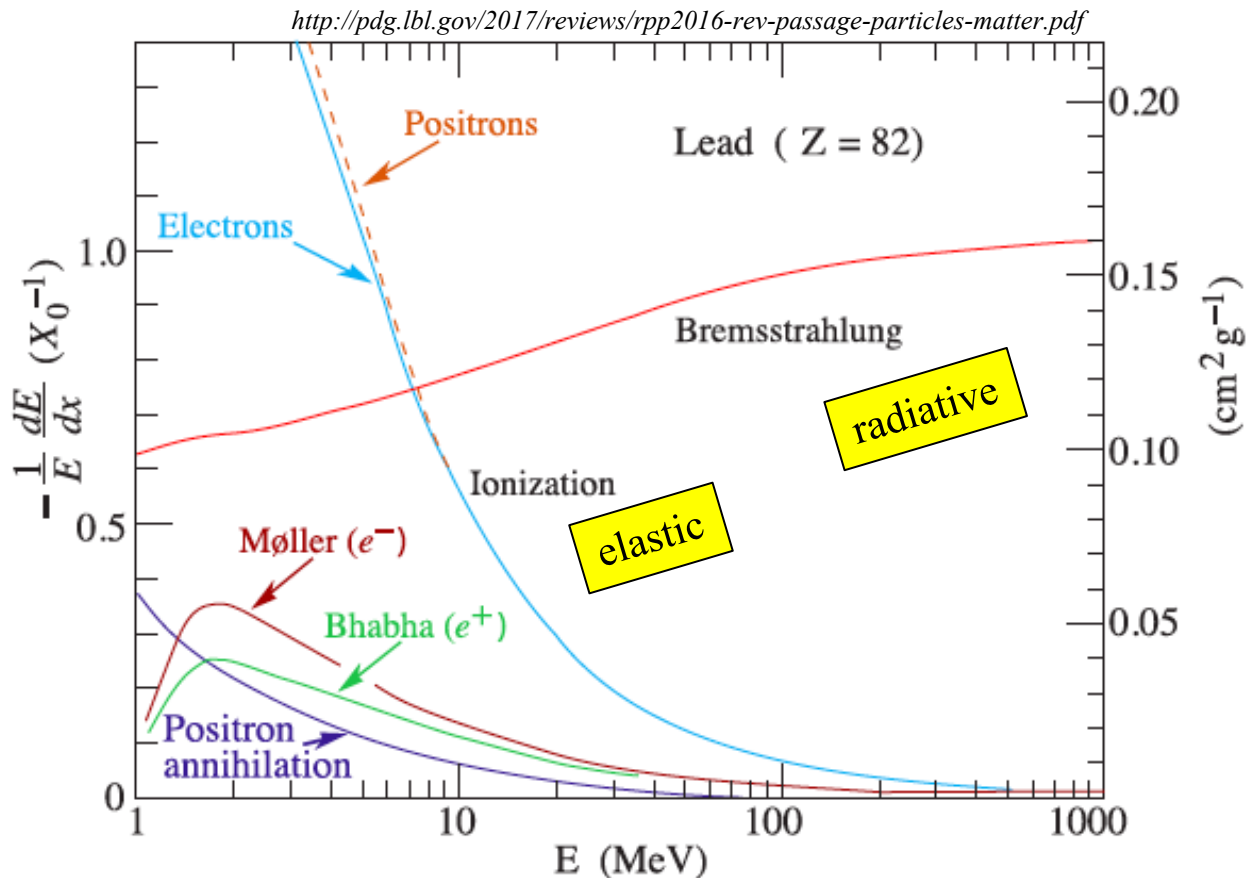


Jets in nuclear collisions



# Energy loss in QED

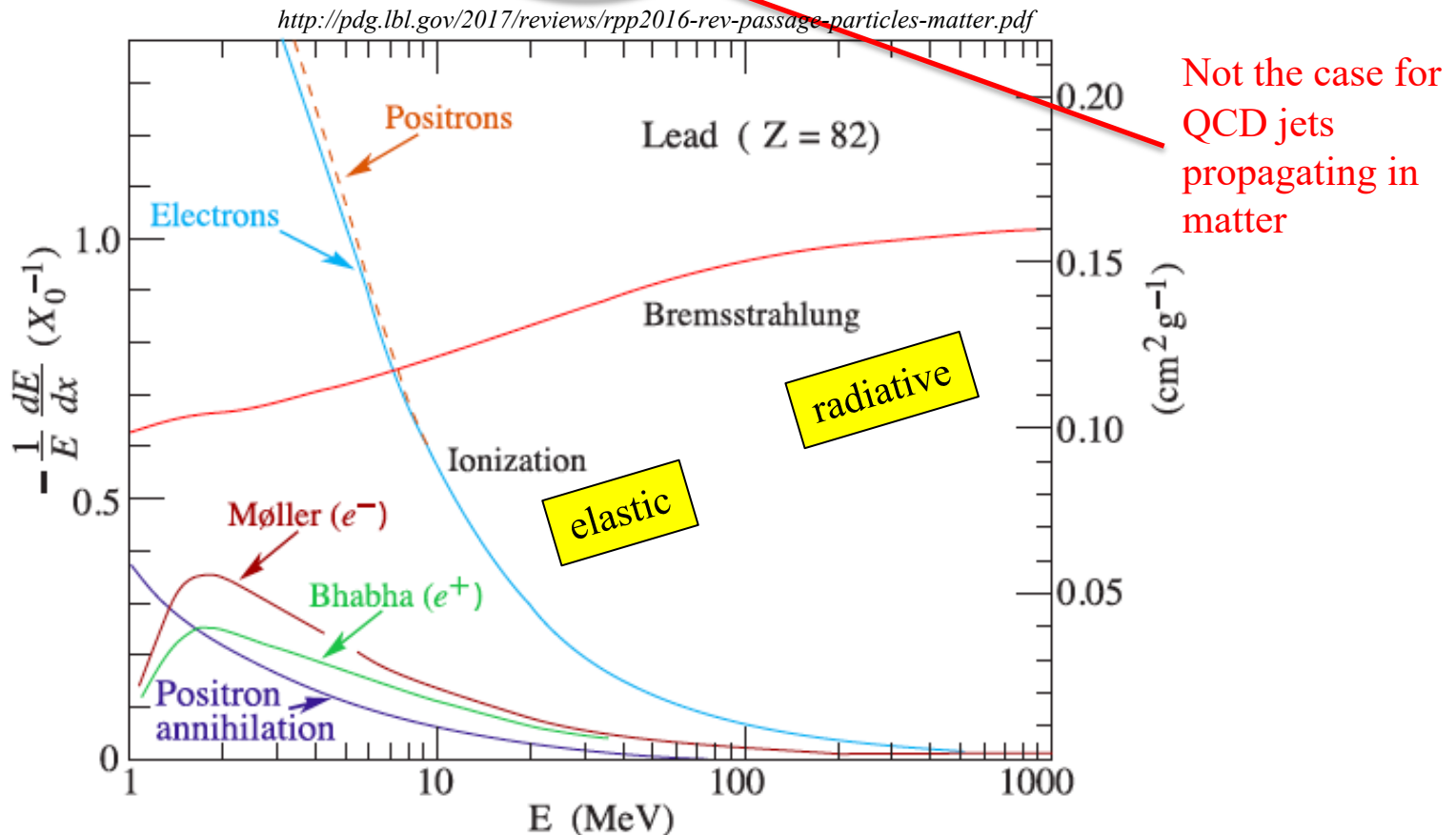
Fractional energy loss of an (on-shell) electron or positron in Lead



**Figure 33.11:** Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization

# Energy loss in QED

Fractional energy loss of an **(on-shell)** electron or positron in Lead

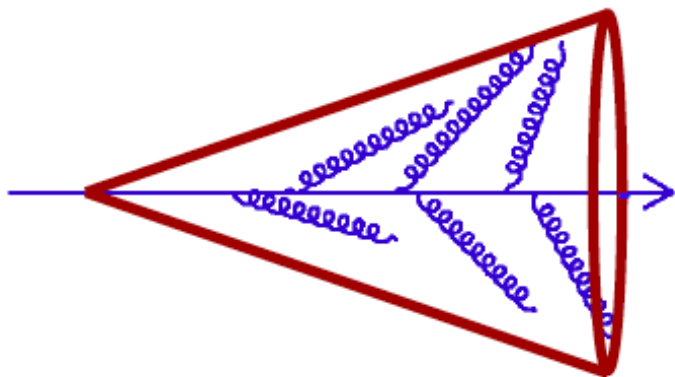


**Figure 33.11:** Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization

# Jet quenching in one slide

## Jet shower in-medium

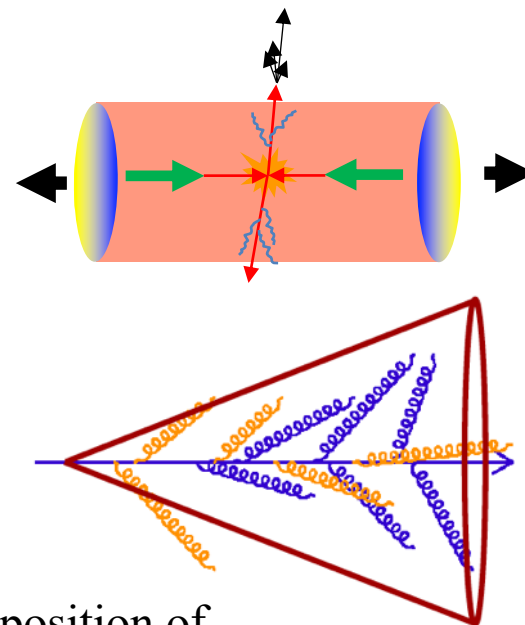
### Jet shower in vacuum



Evolution of highly virtual parton via gluon radiation

Quantum interference → angle-ordering

- hardest radiation is most collinear with jet axis
- Precise understanding in pQCD
- Accurately calculable with QCD-based Monte Carlo models



Superposition of

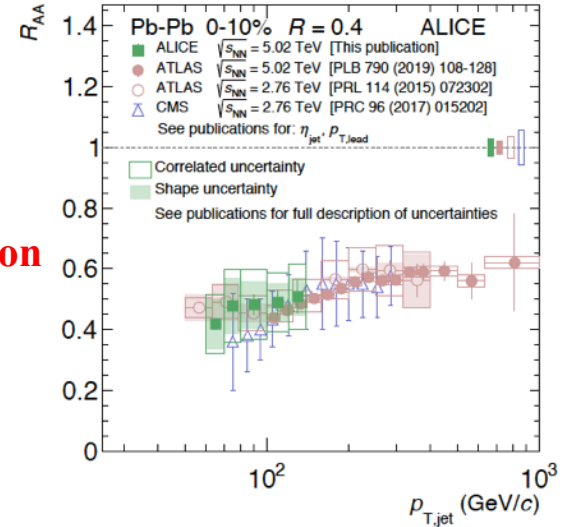
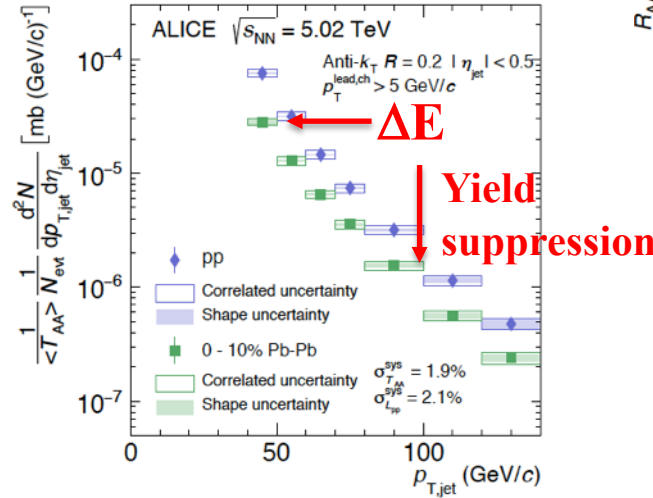
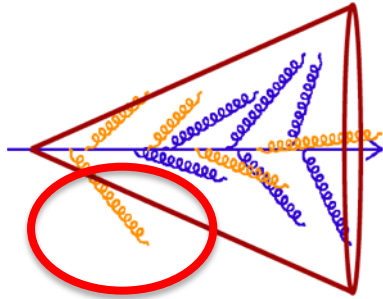
- vacuum shower
- medium-induced gluon emission

These processes happen simultaneously and interfere

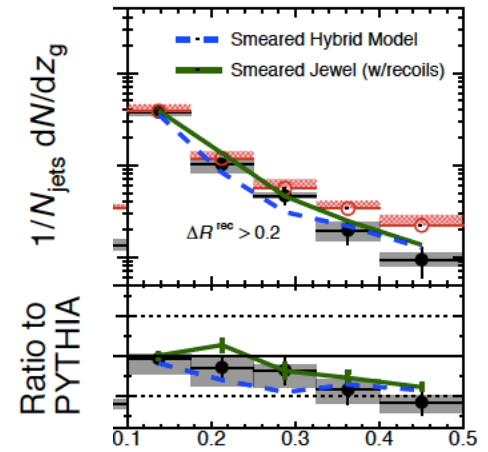
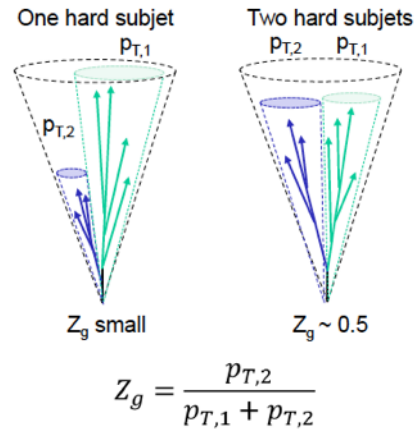
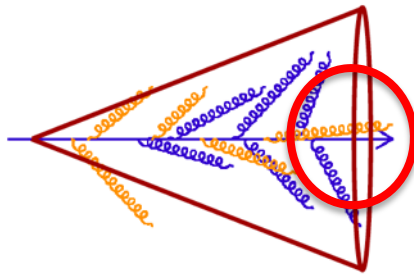
Angle-ordering is modified or destroyed

# Jet quenching: observable consequences I

## 1. Energy loss



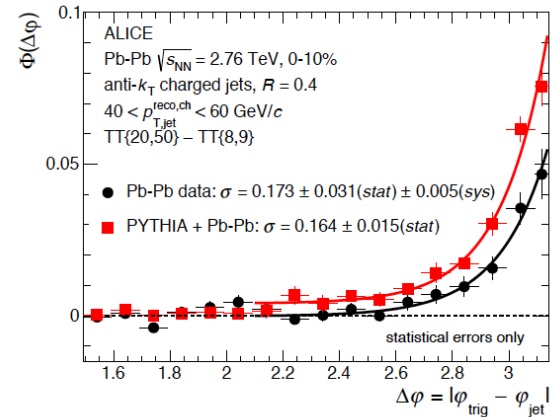
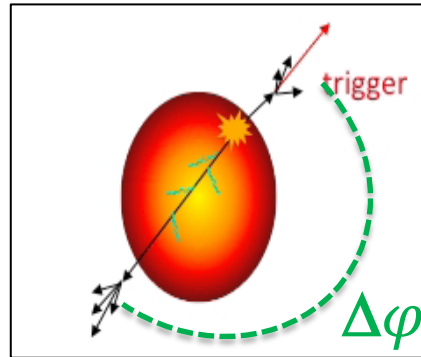
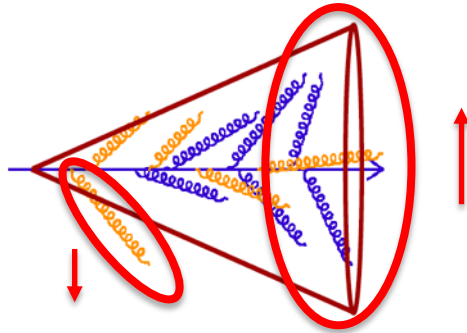
## 2. Modification of jet substructure



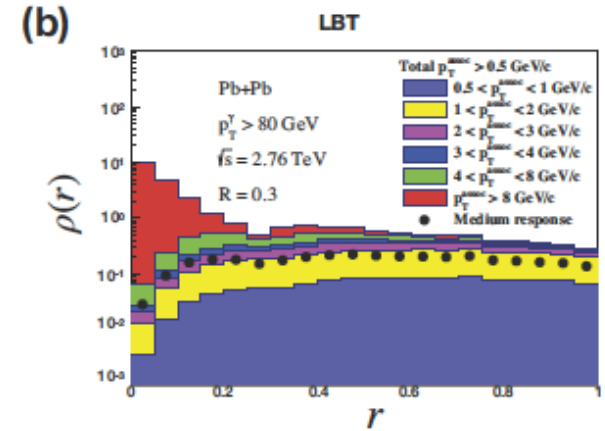
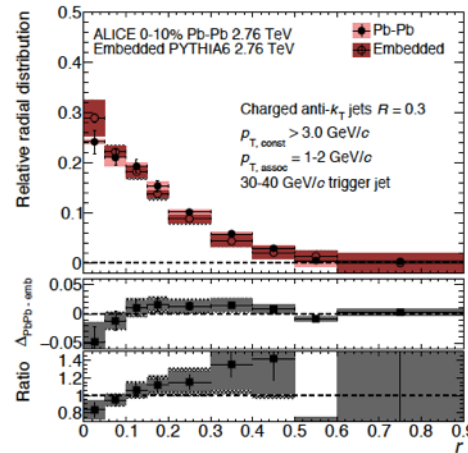
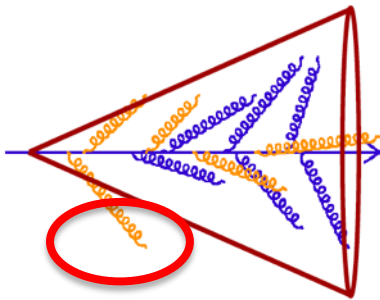


# Jet quenching: observable consequences II

## 3. Jet deflection



## 4. Recovery of large-angle radiation



# Jet quenching: observable consequences III

Four distinct manifestations of jet quenching:

- Jet energy loss
- Jet substructure modification
- Jet deflection
- Large-angle radiation

Different manifestations of same underlying physics

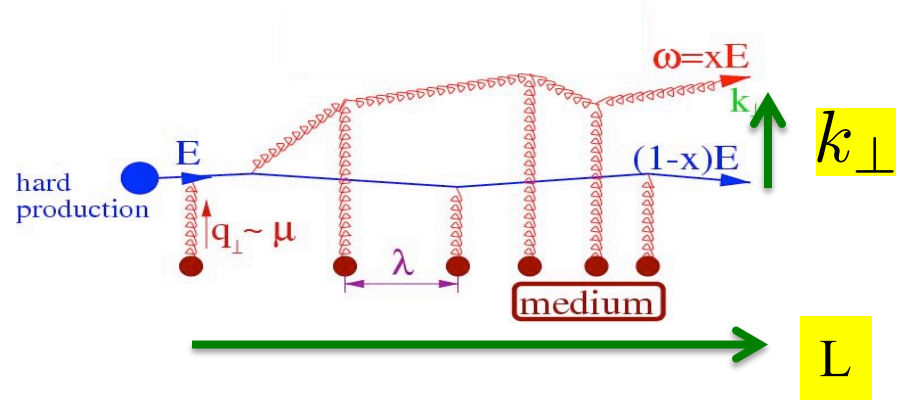
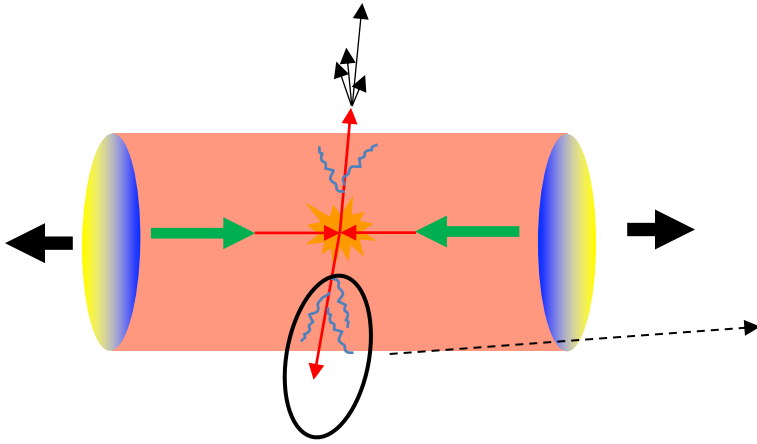
- All must occur if any of them does
- Probe different aspects of jet quenching
- Different experimental systematics as fn of kinematics and collision system
- Different theoretical sensitivity as fn of kinematics and collision system

This is an opportunity:

Measure the same physics multiple ways and require consistency

→ needs a theoretical framework...

# Radiative energy loss in QCD



Thermal field theory:

$$C(\mathbf{q}) = \frac{g_s^2 m_D^2 T}{\mathbf{q}^2 (\mathbf{q}^2 + m_D^2)}$$

$$m_D^2 = 3g_s^2 T^2 / 2$$

$C(\mathbf{q}) =$  Scattering kernel

$\mathbf{q} =$  Momentum transfer

$T =$  Temperature

$m_D =$  Debye mass

} QGP properties

$$\hat{q} \equiv \frac{\langle k_{\perp}^2 \rangle}{L} \sim \frac{1}{L} \int d\mathbf{q}^2 \mathbf{q}^2 C(\mathbf{q})$$

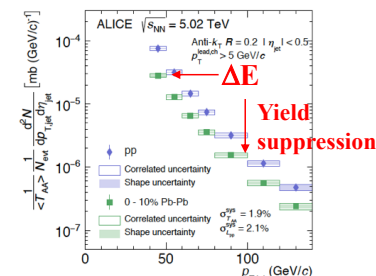
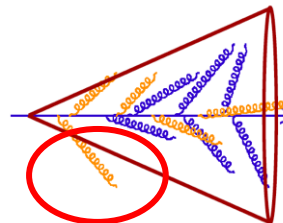
# Connecting qhat to measurements

Useful example: BDMPS

- multiple soft scattering approximation
- gives insight into parametric dependencies
- connection to more complete approaches must be checked

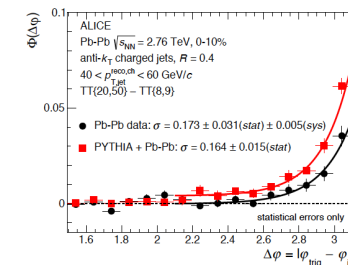
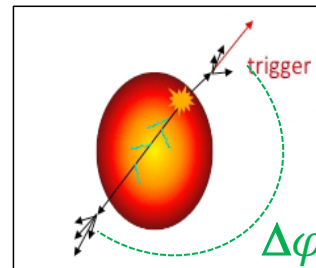
Medium-induced jet energy loss:

$$\Delta E_{med} \sim \alpha_s \hat{q} L^2$$



Medium-induced angular broadening:

$$\langle k_T^2 \rangle \sim \langle \Delta\phi^2 \rangle \sim \alpha_s \hat{q} L$$

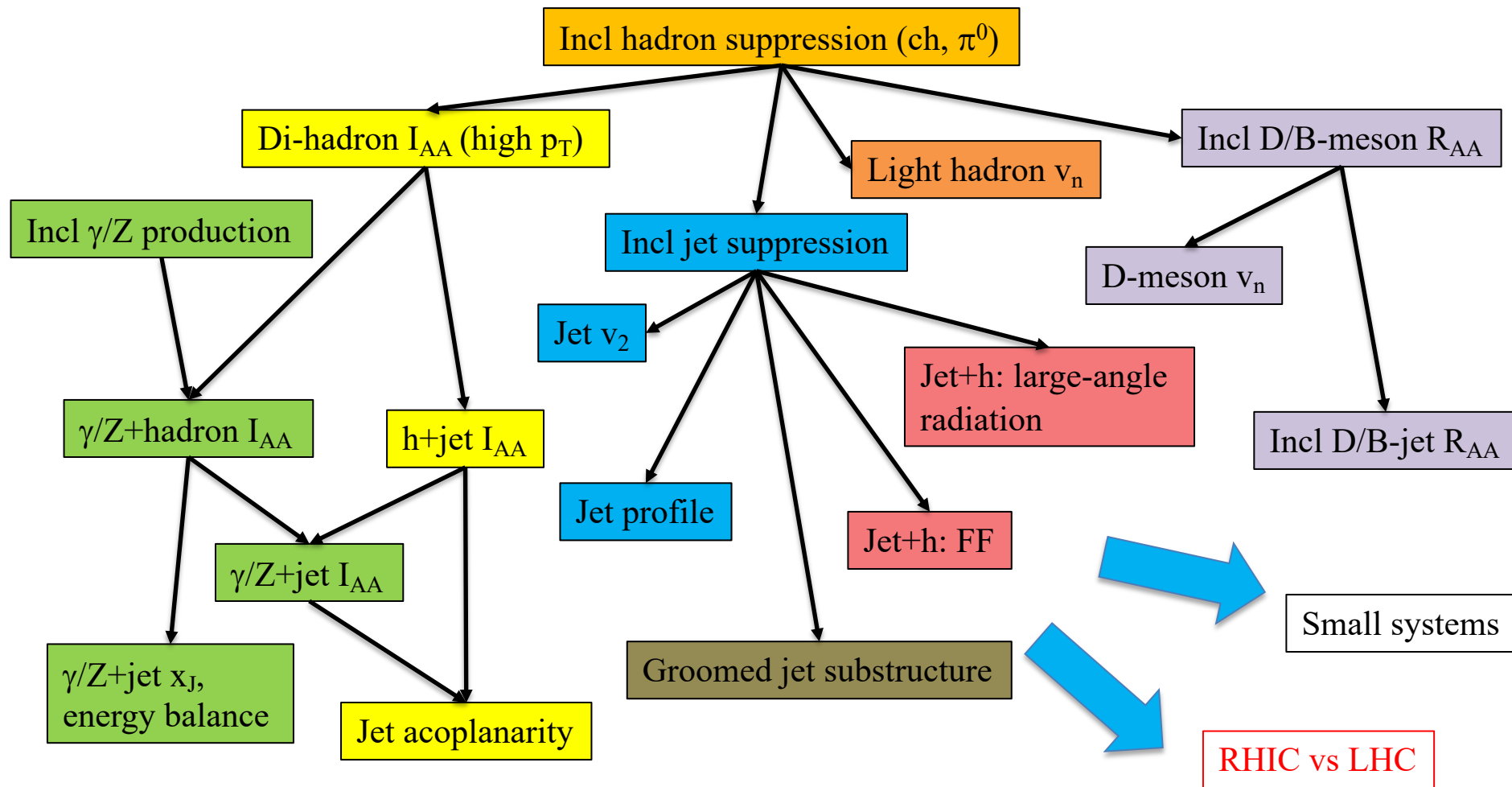


# Taxonomy of current jet quenching measurements

Map driven by experimental considerations:

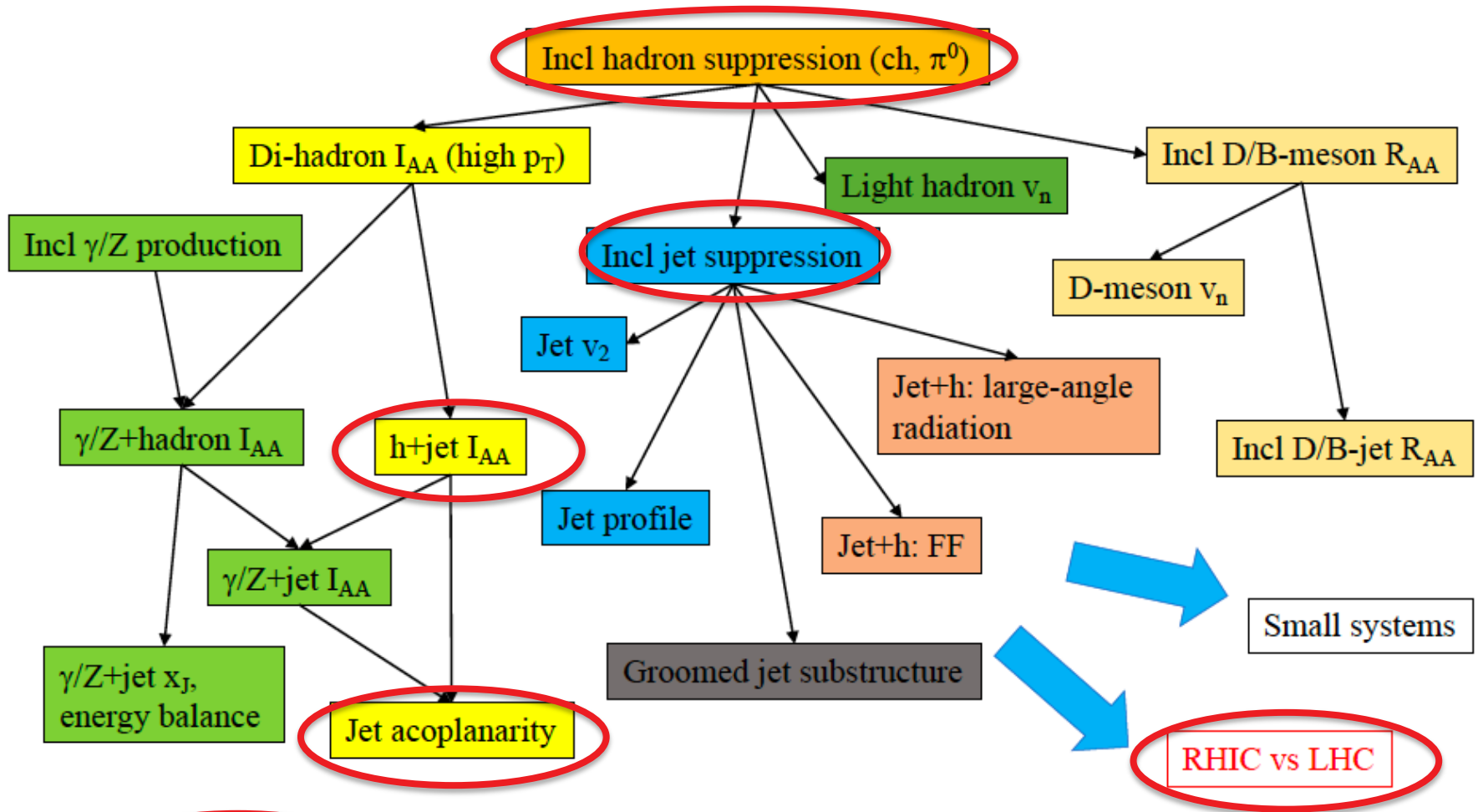
- arrows connect observables with just one thing changed

How do these map onto theory?



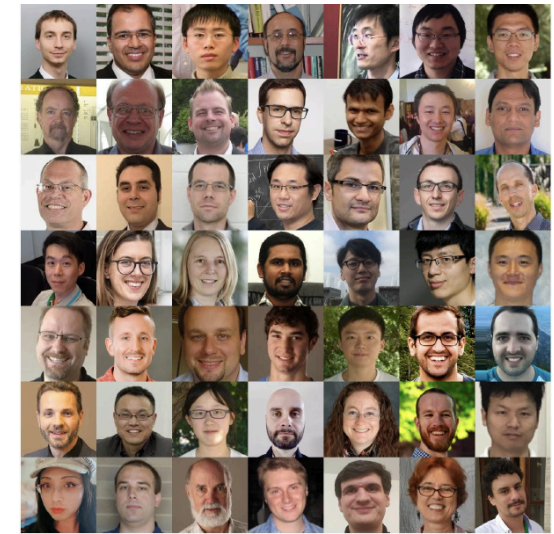
# Confusing! How to make sense of so many observables?

Go systematically: start with a few select measurements and build up the picture...

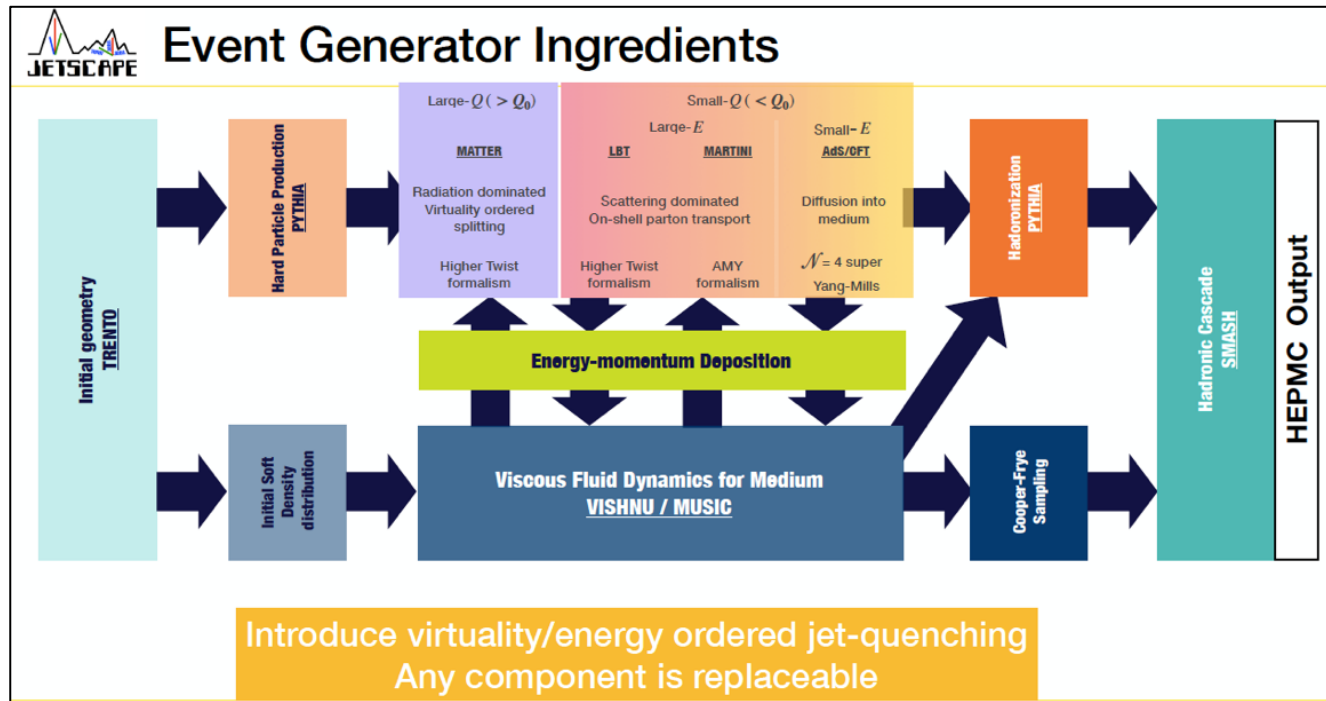


= my choices; other choices equally valid

# Connecting experiment and theory...

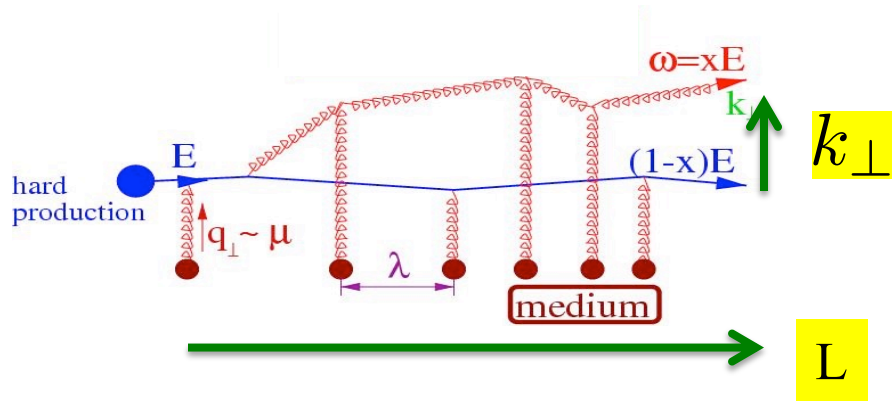
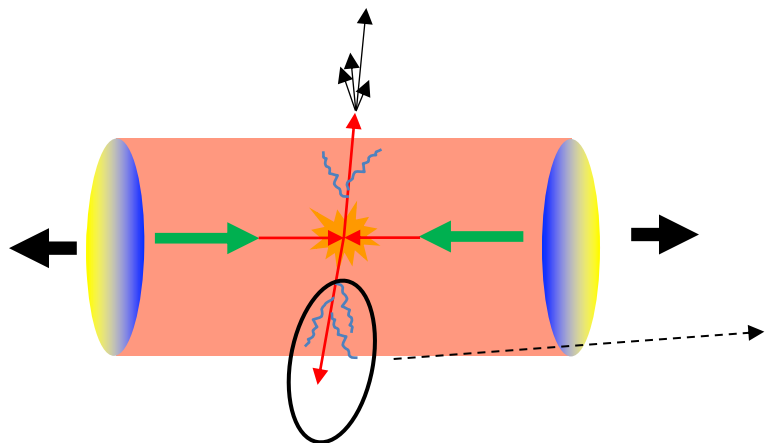


Modular framework: multi-stage jet quenching calculations  
**Parameter extraction via Bayesian Inference**  
 Goal: general tool for entire HI community





# JETSCAPE: measuring $\hat{q}$ using incl hadrons



Thermal field theory:

$$C(\mathbf{q}) = \frac{g_s^2 m_D^2 T}{\mathbf{q}^2 (\mathbf{q}^2 + m_D^2)}$$

$$m_D^2 = 3g_s^2 T^2 / 2$$

$C(\mathbf{q})$  = Scattering kernel

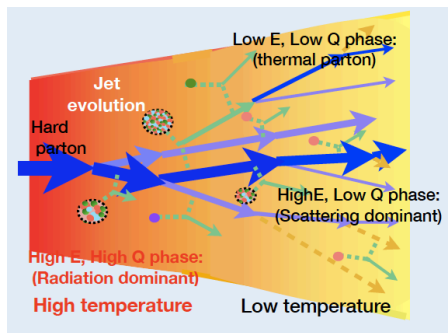
$\mathbf{q}$  = Momentum transfer

$T$  = Temperature

$m_D$  = Debye mass

$$\hat{q} \equiv \frac{\langle k_{\perp}^2 \rangle}{L} \sim \frac{1}{L} \int d\mathbf{q}^2 \mathbf{q}^2 C(\mathbf{q})$$

JETSCAPE  
parametrization



High jet virtuality  
 $Q \gg T$

Low jet virtuality  $Q \sim T$   
(sensitive to thermal medium)

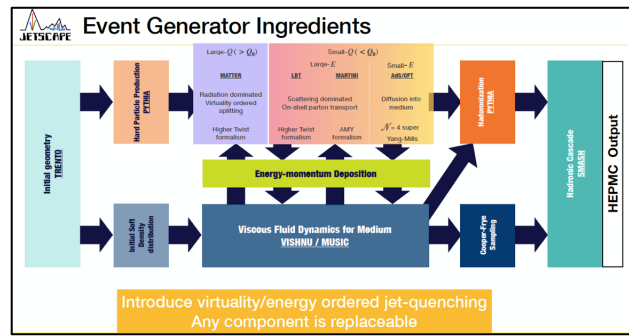
$$\frac{\hat{q}(E, T) |_{A,B,C,D}}{T^3} = 42 C_R \frac{\zeta(3)}{\pi} \left( \frac{4\pi}{9} \right)^2 \left\{ \frac{A \left[ \ln \left( \frac{E}{\Lambda} \right) - \ln(B) \right]}{\left[ \ln \left( \frac{E}{\Lambda} \right) \right]^2} \frac{C \left[ \ln \left( \frac{E}{T} \right) - \ln(D) \right]}{\left[ \ln \left( \frac{ET}{\Lambda^2} \right) \right]^2} \right\}$$

# Bayesian inference: inclusive hadrons

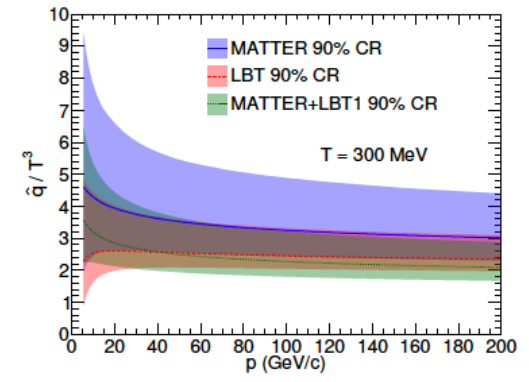
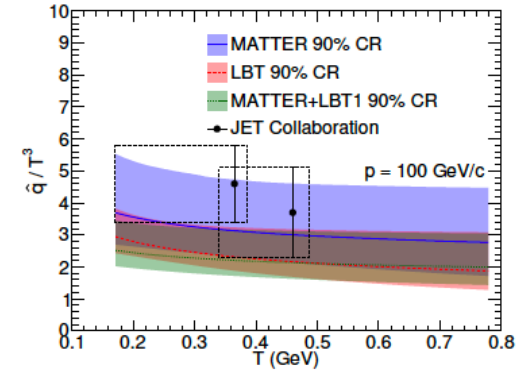
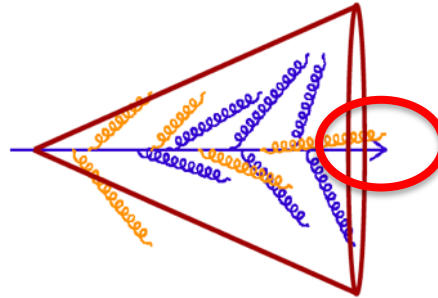
paper in preparation

$$\frac{\hat{q}(E, T) |_{A,B,C,D}}{T^3} = 42 C_R \frac{\zeta(3)}{\pi} \left( \frac{4\pi}{9} \right)^2 \left\{ \frac{A [\ln(\frac{E}{\Lambda}) - \ln(B)]}{[\ln(\frac{E}{\Lambda})]^2} + \frac{C [\ln(\frac{E}{T}) - \ln(D)]}{[\ln(\frac{ET}{\Lambda^2})]^2} \right\}$$

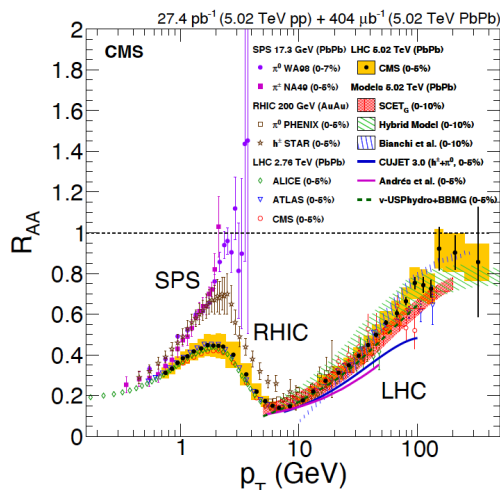
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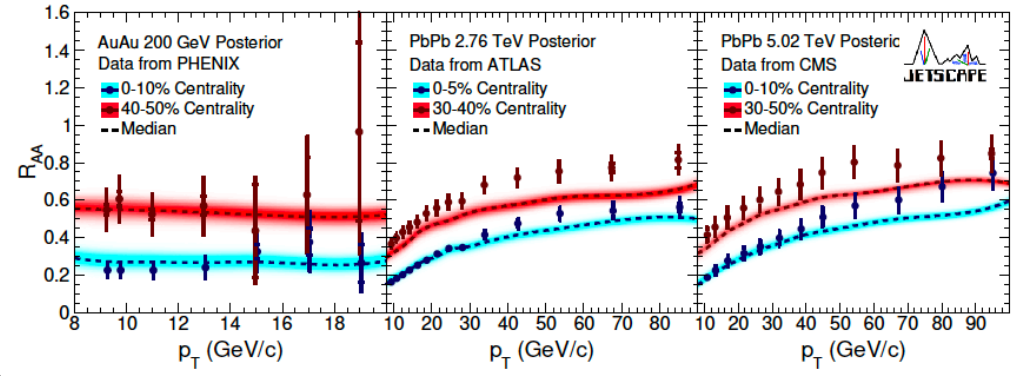
Rigorous quantitative  
determination of  $\hat{q}$



+



Posterior distributions



=



# JETSCAPE determination of $\hat{q}$ at using inclusive hadron $R_{AA}$ :

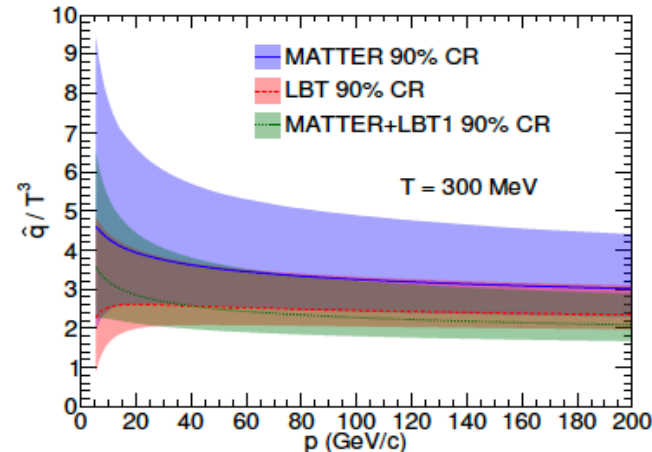
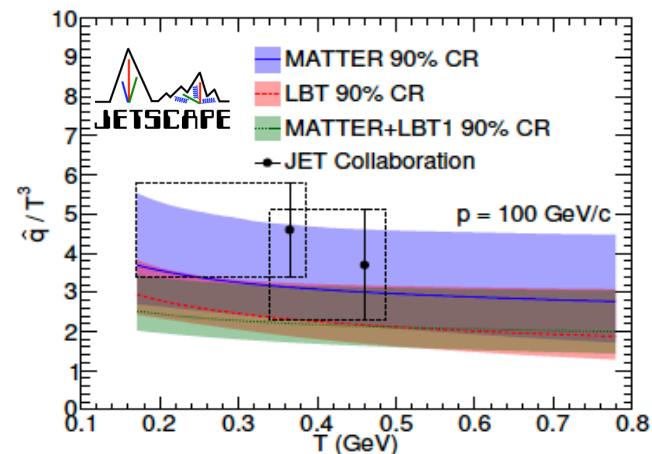
Current state-of-the-art quantitative analysis of jet quenching

Are we done?

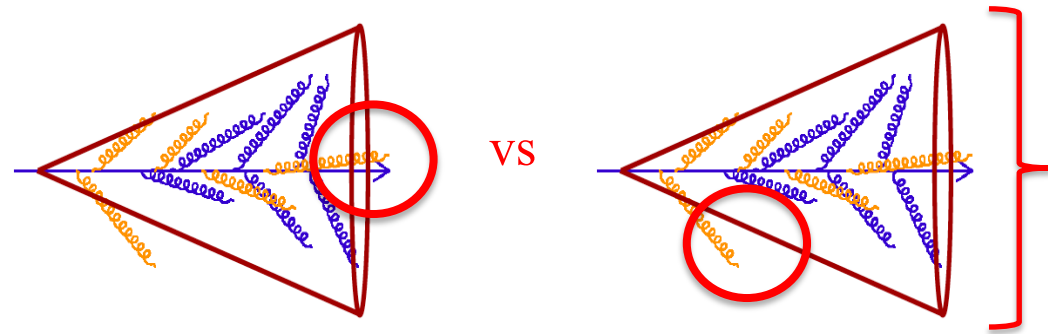
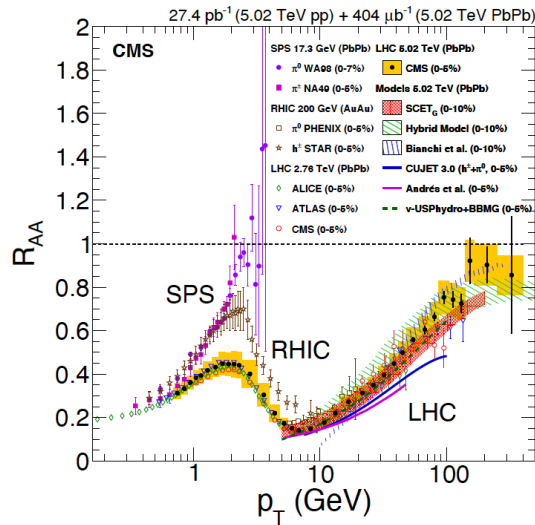
→ real progress, but hardly the complete story

Consider some next steps....

paper in preparation



# Inclusive hadron vs inclusive jet suppression

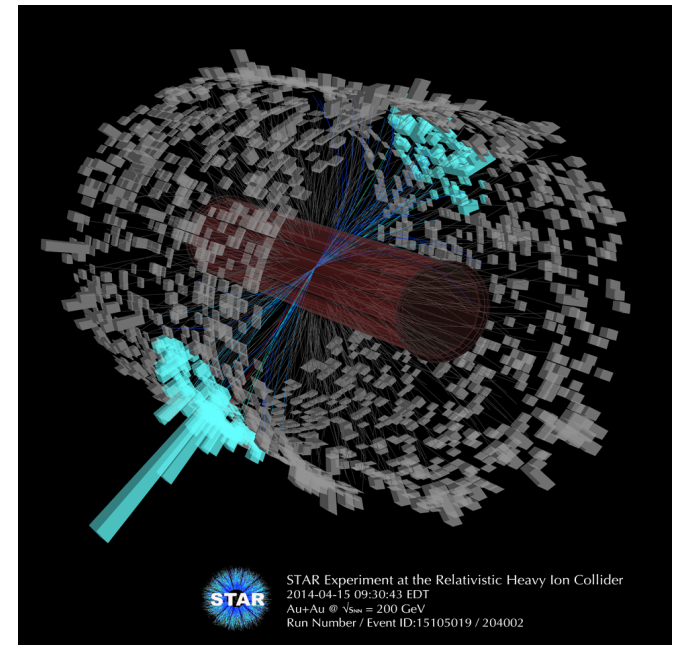


Inclusive hadron suppression driven by energy transport away from the hardest branch in the jet

- Insensitive to specific mechanisms of energy transport

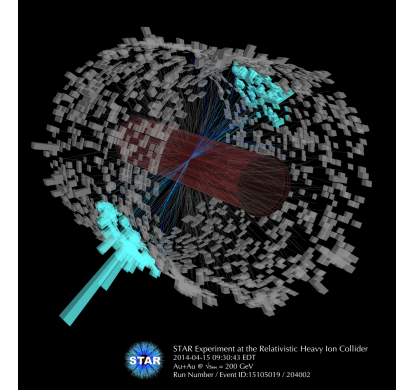
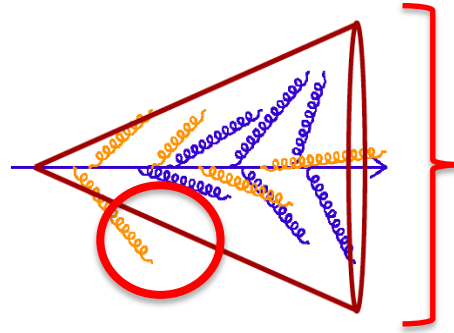
More comprehensive: reconstructed jets

- very challenging due to large backgrounds, especially at RHIC
- but problem has been solved

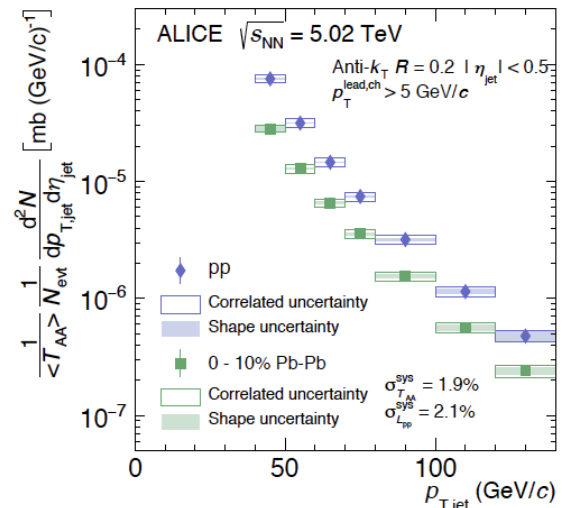
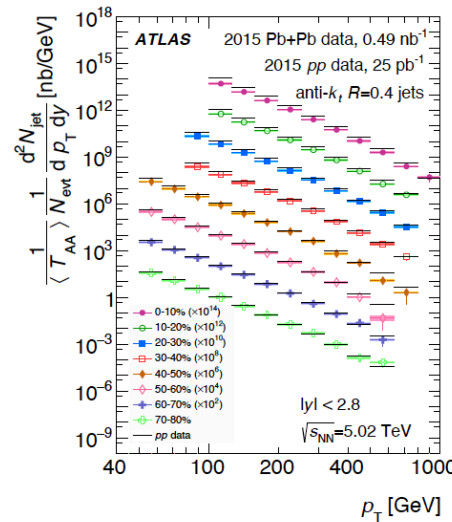
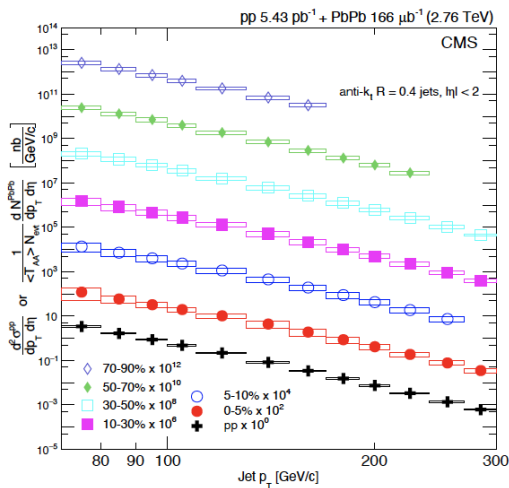


# Inclusive jets in A+A: spectra

RHIC



High-quality data over a vast kinematic range



**Measurement of inclusive charged-particle jet production in Au + Au collisions at  $\sqrt{s_{NN}} = 200$  GeV**

J. Adam,<sup>6</sup> L. Adamczyk,<sup>2</sup> J. R. Adams,<sup>39</sup> J. K. Adkins,<sup>30</sup> G. Agakishiev,<sup>28</sup> M. M. Aggarwal,<sup>41</sup> Z. Ahammed,<sup>61</sup> I. Alekseev,<sup>3,35</sup>  
 D. M. Anderson,<sup>55</sup> A. Aparin,<sup>28</sup> E. C. Aschenauer,<sup>6</sup> M. U. Ashraf,<sup>11</sup> F. G. Atetalla,<sup>29</sup> A. Attri,<sup>41</sup> G. S. Averichev,<sup>28</sup>  
 V. Bairathi,<sup>53</sup> K. Barish,<sup>10</sup> A. Behera,<sup>52</sup> R. Bellwied,<sup>20</sup> A. Bhasin,<sup>27</sup> J. Bielcik,<sup>14</sup> J. Bielcikova,<sup>38</sup> L. C. Bland,<sup>6</sup>  
 I. G. Bordyuzhin,<sup>3</sup> J. D. Brandenburg,<sup>49,6</sup> A. V. Brandin,<sup>35</sup> J. Butterworth,<sup>45</sup> H. Caines,<sup>64</sup> M. Calderón de la Barca Sánchez,<sup>8</sup>  
 D. Cebra,<sup>8</sup> I. Chakaberia,<sup>29,6</sup> P. Chaloupka,<sup>14</sup> B. K. Chan,<sup>9</sup> F.-H. Chang,<sup>37</sup> Z. Chang,<sup>6</sup> N. Chankova-Bunzarova,<sup>28</sup>  
 A. Chatterjee,<sup>11</sup> D. Chen,<sup>10</sup> J. H. Chen,<sup>18</sup> X. Chen,<sup>48</sup> Z. Chen,<sup>49</sup> J. Cheng,<sup>57</sup> M. Cherney,<sup>13</sup> M. Chevalier,<sup>10</sup> S. Choudhury,<sup>18</sup>

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 V. Bairathi,<sup>53</sup> K. Barish,<sup>10</sup> A. Behera,<sup>52</sup> R. Bellwied,<sup>20</sup> A. Bhasin,<sup>27</sup> J. Bielcik,<sup>14</sup> J. Bielcikova,<sup>38</sup> L. C. Bland,<sup>6</sup>  
 I. G. Bordyuzhin,<sup>3</sup> J. D. Brandenburg,<sup>49,6</sup> A. V. Brandin,<sup>35</sup> J. Butterworth,<sup>45</sup> H. Caines,<sup>64</sup> M. Calderón de la Barca Sánchez,<sup>8</sup>

D. A. Morozov,<sup>43</sup> M. Nagy,<sup>16</sup> J. D. Nam,<sup>54</sup> Md. Nasim,<sup>22</sup> K. Nayak,<sup>11</sup> D. Neff,<sup>9</sup> J. M. Nelson,<sup>7</sup> D. B. Nemes,<sup>64</sup> M. Nie,<sup>49</sup>  
 G. Nigmatkulov,<sup>35</sup> T. Niida,<sup>58</sup> L. V. Nogach,<sup>43</sup> T. Nonaka,<sup>58</sup> G. Odyniec,<sup>31</sup> A. Ogawa,<sup>6</sup> S. Oh,<sup>31</sup> V. A. Okorokov,<sup>35</sup> B. S. Page,<sup>6</sup>  
 R. Pak,<sup>6</sup> A. Pandav,<sup>36</sup> Y. Panebratsev,<sup>28</sup> B. Pawlik,<sup>40</sup> D. Pawlowska,<sup>62</sup> H. Pei,<sup>11</sup> C. Perkins,<sup>7</sup> L. Pinsky,<sup>20</sup> R. L. Pintér,<sup>16</sup>  
 J. Pluta,<sup>62</sup> J. Porter,<sup>31</sup> M. Posik,<sup>54</sup> N. K. Pruthi,<sup>41</sup> M. Przybycien,<sup>2</sup> J. Putschke,<sup>63</sup> H. Qiu,<sup>26</sup> A. Quintero,<sup>54</sup>  
 S. K. Radhakrishnan,<sup>29</sup> S. Ramachandran,<sup>30</sup> R. L. Ray,<sup>56</sup> R. Reed,<sup>32</sup> H. G. Ritter,<sup>31</sup> J. B. Roberts,<sup>45</sup> O. V. Rogachevskiy,<sup>28</sup>  
 J. L. Romero,<sup>8</sup> L. Ruan,<sup>6</sup> J. Rusnak,<sup>38</sup> N. R. Sahoo,<sup>49</sup> H. Sako,<sup>58</sup> S. Salur,<sup>46</sup> J. Sandweiss,<sup>64</sup> S. Sato,<sup>58</sup> W. B. Schmidke,<sup>6</sup>  
 N. Schmitz,<sup>33</sup> B. R. Schweid,<sup>52</sup> F. Seck,<sup>15</sup> J. Seger,<sup>13</sup> M. Sergeeva,<sup>9</sup> R. Seto,<sup>10</sup> P. Seyboth,<sup>33</sup> N. Shah,<sup>24</sup> E. Shalahiev,<sup>28</sup>  
 P. V. Shanmuganathan,<sup>6</sup> M. Shao,<sup>48</sup> F. Shen,<sup>49</sup> W. Q. Shen,<sup>50</sup> S. S. Shi,<sup>11</sup> Q. Y. Shou,<sup>50</sup> E. P. Sichtermann,<sup>31</sup> R. Sikora,<sup>2</sup>  
 M. Simko,<sup>38</sup> J. Singh,<sup>41</sup> S. Singha,<sup>26</sup> N. Smirnov,<sup>64</sup> W. Solyst,<sup>25</sup> P. Sorensen,<sup>6</sup> H. M. Spinka,<sup>4</sup> B. Srivastava,<sup>44</sup>  
 T. D. S. Stanislaus,<sup>60</sup> M. Stefaniak,<sup>62</sup> D. J. Stewart,<sup>64</sup> M. Strikhanov,<sup>35</sup> B. Stringfellow,<sup>44</sup> A. A. P. Suaide,<sup>47</sup> M. Sumera,<sup>38</sup>  
 B. Summa,<sup>42</sup> X. M. Sun,<sup>11</sup> X. Sun,<sup>12</sup> Y. Sun,<sup>48</sup> Y. Sun,<sup>21</sup> B. Surrow,<sup>54</sup> D. N. Svirida,<sup>3</sup> P. Szymanski,<sup>62</sup> A. H. Tang,<sup>6</sup> Z. Tang,<sup>48</sup>  
 A. Taranenko,<sup>35</sup> T. Tarnowsky,<sup>34</sup> J. H. Thomas,<sup>31</sup> A. R. Timmins,<sup>20</sup> D. Tlusty,<sup>13</sup> M. Tokarev,<sup>28</sup> C. A. Tomkiel,<sup>32</sup>  
 S. Trentalange,<sup>9</sup> R. E. Tribble,<sup>55</sup> P. Tribedy,<sup>6</sup> S. K. Tripathy,<sup>16</sup> O. D. Tsai,<sup>9</sup> Z. Tu,<sup>6</sup> T. Ullrich,<sup>6</sup> D. G. Underwood,<sup>4</sup> I. Upsal,<sup>49,6</sup>  
 G. Van Buren,<sup>6</sup> J. Vanek,<sup>38</sup> A. N. Vasiliev,<sup>43</sup> I. Vassiliev,<sup>17</sup> F. Videbak,<sup>6</sup> S. Vokal,<sup>28</sup> S. A. Voloshin,<sup>63</sup> F. Wang,<sup>44</sup> G. Wang,<sup>9</sup>  
 J. S. Wang,<sup>21</sup> P. Wang,<sup>48</sup> Y. Wang,<sup>11</sup> Y. Wang,<sup>57</sup> Z. Wang,<sup>49</sup> J. C. Webb,<sup>6</sup> P. C. Weidenkaff,<sup>19</sup> L. Wen,<sup>9</sup> G. D. Westfall,<sup>34</sup>  
 H. Wieman,<sup>31</sup> S. W. Wissink,<sup>25</sup> R. Witt,<sup>59</sup> Y. Wu,<sup>10</sup> Z. G. Xiao,<sup>57</sup> G. Xie,<sup>31</sup> W. Xie,<sup>44</sup> H. Xu,<sup>21</sup> N. Xu,<sup>31</sup> Q. H. Xu,<sup>49</sup> Y. F. Xu,<sup>50</sup>  
 Y. Xu,<sup>49</sup> Z. Xu,<sup>6</sup> Z. Xu,<sup>9</sup> C. Yang,<sup>49</sup> Q. Yang,<sup>49</sup> S. Yang,<sup>37</sup> Y. Yang,<sup>11</sup> Z. Yang,<sup>11</sup> Z. Ye,<sup>45</sup> Z. Ye,<sup>12</sup> L. Yi,<sup>49</sup> K. Yip,<sup>6</sup>  
 H. Zbroszczyk,<sup>62</sup> W. Zha,<sup>48</sup> D. Zhang,<sup>11</sup> S. Zhang,<sup>48</sup> S. Zhang,<sup>50</sup> X. P. Zhang,<sup>57</sup> Y. Zhang,<sup>48</sup> Y. Zhang,<sup>11</sup> Z. J. Zhang,<sup>37</sup>  
 Z. Zhang,<sup>6</sup> Z. Zhang,<sup>12</sup> J. Zhao,<sup>44</sup> C. Zhong,<sup>50</sup> C. Zhou,<sup>50</sup> X. Zhu,<sup>57</sup> Z. Zhu,<sup>49</sup> M. Zurek,<sup>31</sup> and M. Zyzak<sup>17</sup>

(STAR Collaboration)

1 Abilene Christian University, Abilene, Texas 79600

# STAR heavy ion jet measurements: subsystems and datasets

Charged-particle jets (this paper):

- Time Projection Chamber (TPC)
- Vertex Position Detector (VPD)

Calorimetric jets (in progress)

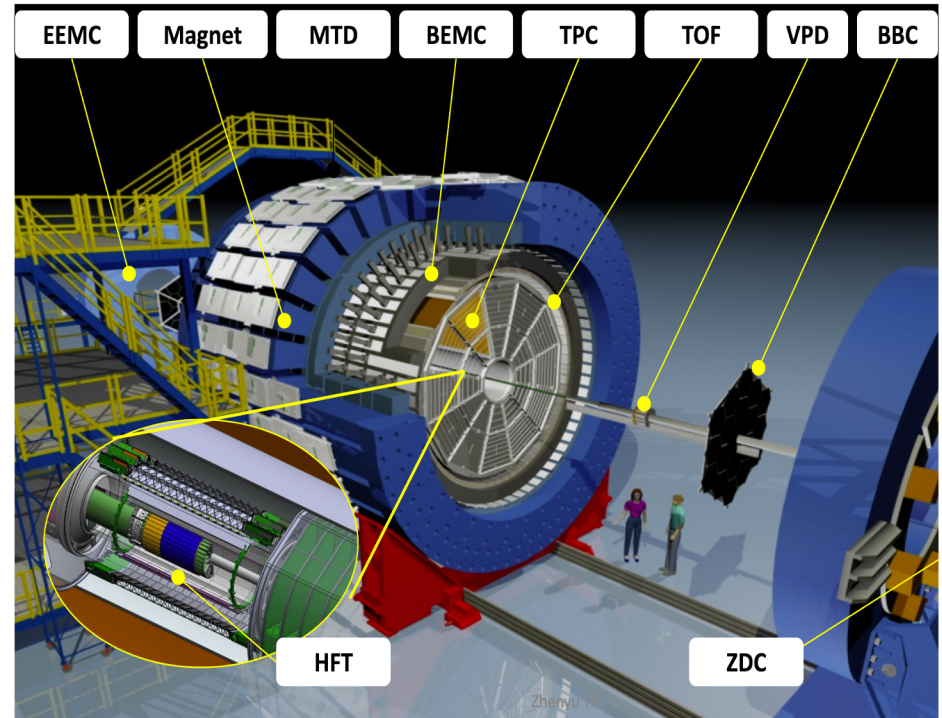
- + Barrel EM Calorimeter (BEMC)

Dataset: Au+Au,  $\sqrt{s_{NN}}=200$  GeV

- 2011 minimum bias,  $L_{int}=6 \mu\text{b}^{-1}$
- 2014 minimum bias; BEMC-triggered,  $L_{int}=5.2 \text{ nb}^{-1}$

Centrality selection:

- charged-track multiplicity,  $|\eta| < 0.5$
- central: 0-10%
- peripheral: 60-80%



# Charged-jet reference: 200 GeV pp collisions

Charged jets: cannot trigger, need MB pp

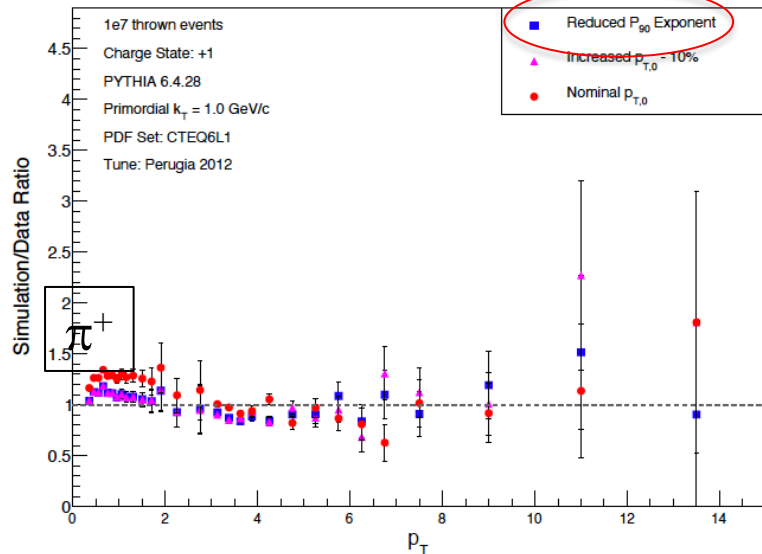
- But insufficient MB pp @ 200 GeV  
 → PYTHIA 6.428 Perugia 2012, STAR tune

PHYSICAL REVIEW D **100**, 052005 (2019)

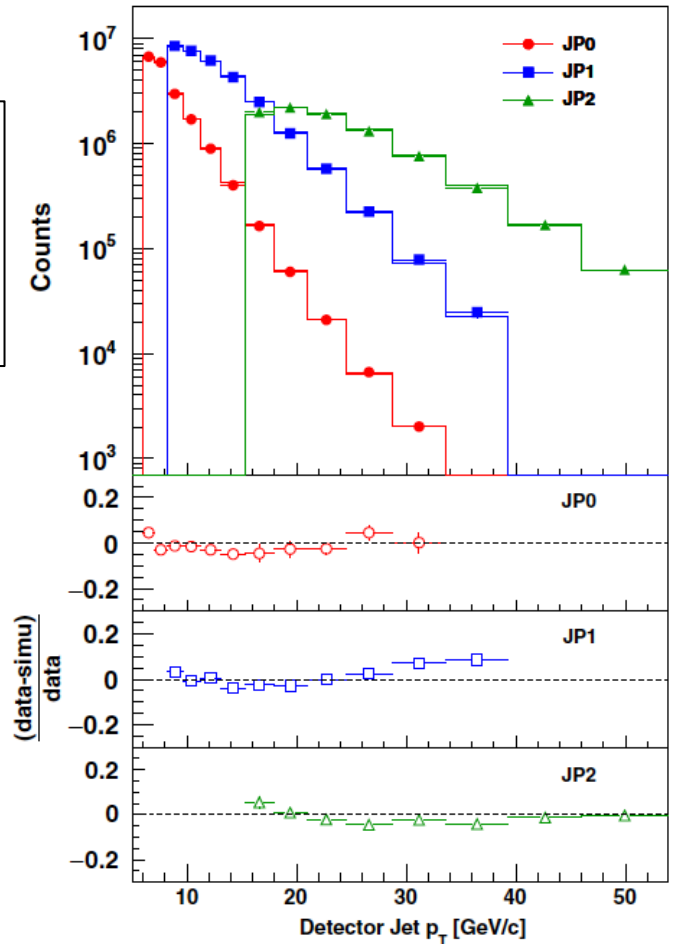
## Longitudinal double-spin asymmetry for inclusive jet and dijet production in pp collisions at $\sqrt{s}=510$ GeV

J. Adam,<sup>6</sup> L. Adamczyk,<sup>2</sup> J. R. Adams,<sup>39</sup> J. K. Adkins,<sup>30</sup> G. Agakishiev,<sup>28</sup> M. M. Aggarwal,<sup>40</sup> Z. Ahammed,<sup>60</sup> I. Alekseev,<sup>3,35</sup> D. M. Anderson,<sup>54</sup> R. Aoyama,<sup>57</sup> A. Aparin,<sup>28</sup> D. Arkhipkin,<sup>6</sup> E. C. Aschenauer,<sup>6</sup> M. U. Ashraf,<sup>56</sup>

*K. Adkins, PhD Thesis, U. Kentucky*



PYTHIA STAR tune vs STAR data:  
detector-level jets



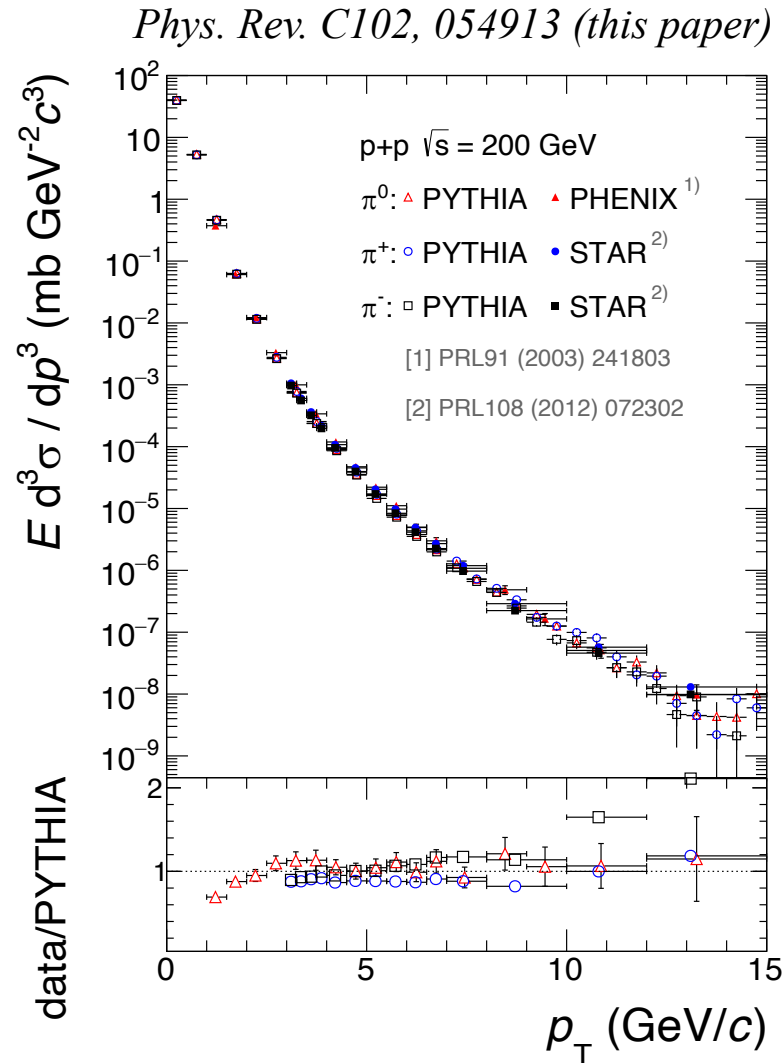
IR regularization scale for MPIs, tuned to LHC 7 TeV pp:

$$p_{T,0}^2(s) = p_{T,0}^2(s_{ref}) \left( \frac{s}{s_{ref}} \right)^{P_{90}}$$

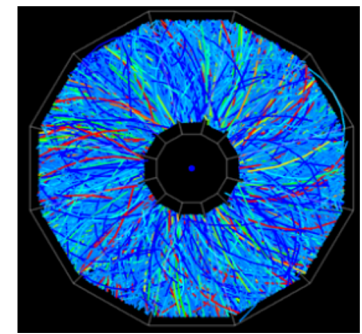


# PYTHIA STAR tune: additional check

Compare inclusive pion yield



# Jet reconstruction

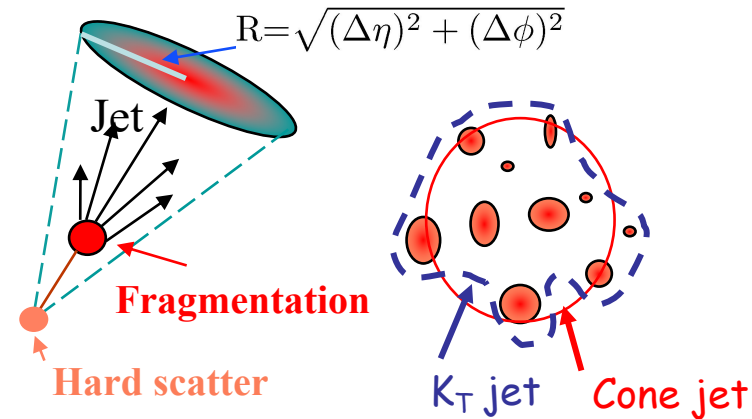


Charged jets:

- all ch. tracks  $|\eta| < 1$
- $0.2 < p_T < 30 \text{ GeV}/c$

Jet reconstruction:

- Anti- $k_T$ ,  $R=0.2, 0.3, 0.4$
- Recombination: boost-invariant  $p_T$  (3-vec)
- Jet centroid acceptance:  $|\eta| < 1-R$



This gives a population of jet candidates that is a combination of

- Jets from hard (**high  $Q^2$** ) processes with  $p_T$  smeared by complex uncorrelated event 👍
- Combinatorial “jets” from random combination of hadrons from soft (**low  $Q^2$** ) processes 👎



High:  $Q^2 > \sim (\text{few GeV})^2$ , somewhat arbitrary

- need an operational procedure to discriminate in measurement

# Analysis strategy: uncorrelated background suppression

G. De Barros et al., arXiv:1208.1518

Correction via unfolding is a linear transformation:

$$\mathbf{m} = \mathbf{R}\mathbf{t}$$

*measured* (green arrow pointing to  $\mathbf{m}$ )  
*truth* (red arrow pointing to  $\mathbf{t}$ )  
*Response matrix* (blue arrow pointing to  $\mathbf{R}$ )

$$R_{ij} = \text{Pr}(\text{measure } i | \text{truth is } j)$$

Regularized inversion:

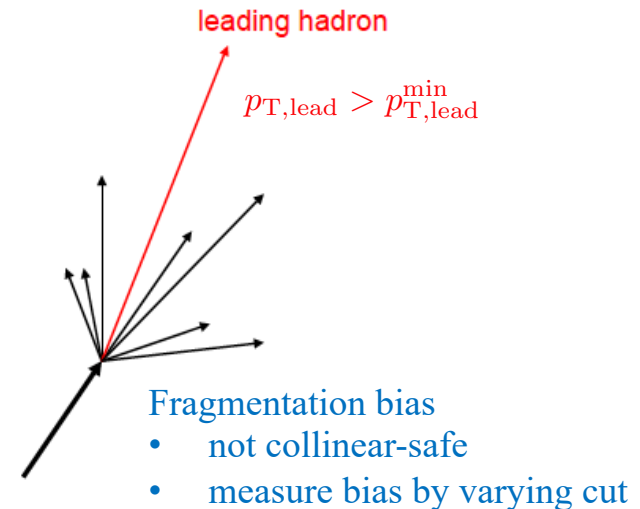
$$\mathbf{t}' = \widetilde{\mathbf{R}}^{-1}\mathbf{m}$$

Solution: bias signal jet population by requiring a hard leading hadron

If jet population contains significant non-jet background yield

- “Response” not meaningful
- Unfolding fails: doesn’t know where to put the counts

→ need to suppress non-jet bkgd prior to unfolding



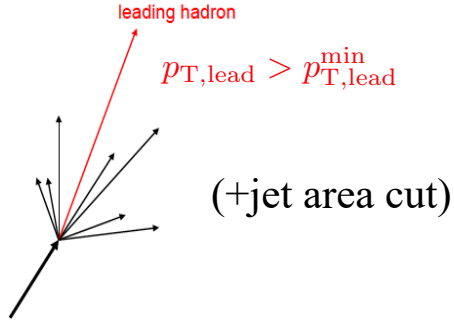
But: no cut on  $p_T^{\text{jet}}$

- unique to this analysis (for incl. jets)
- enables measurement to low  $p_T^{\text{jet}}$ , large R

# Cuts and corrections

Event-wise:

Yield correction (“vertical”)

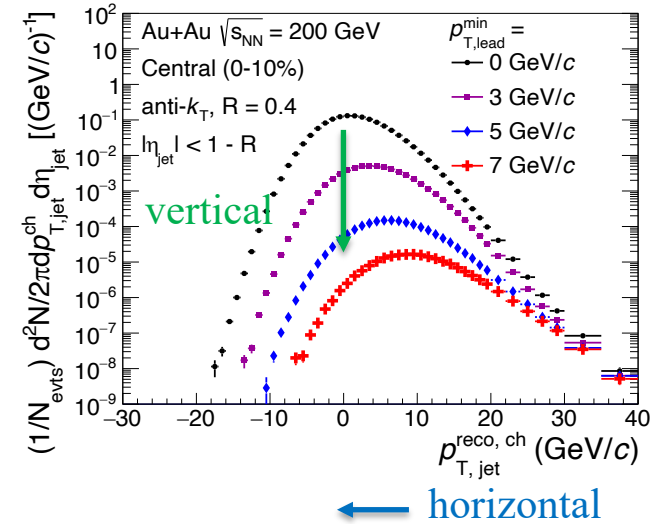


$p_T$ -shift for UE (“horizontal”)

Standard Fastjet procedure:

$$\rho = \text{median} \left\{ \frac{P_{T,jet}^{\text{raw},i}}{A_{jet}^i} \right\}$$

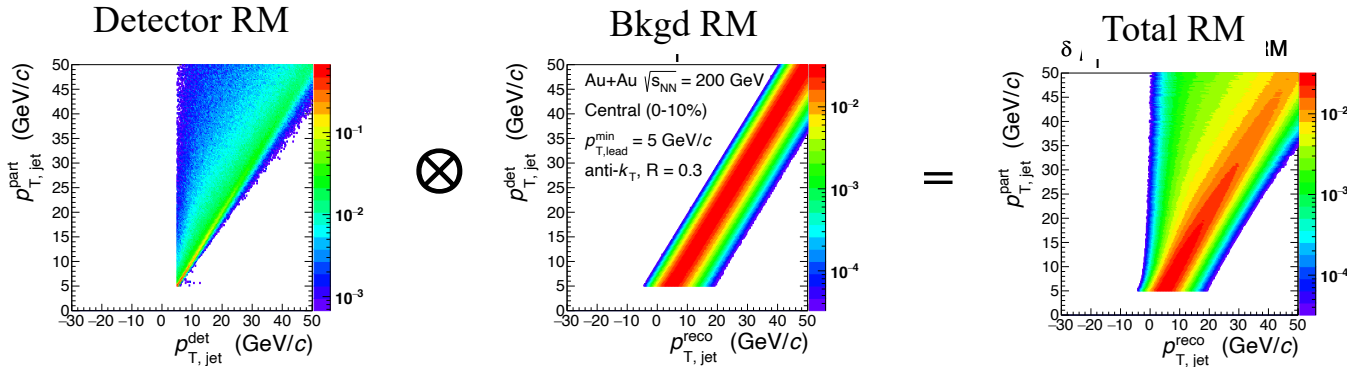
$$p_{T,jet}^{\text{reco},i} = p_{T,jet}^{\text{raw},i} - \rho A_{jet}^i$$



Ensemble-averaged distribution:

Unfolding  $t' = \widetilde{\mathbf{R}}^{-1} \mathbf{m}$

$\mathbf{R} =$  Detector effects  $\otimes$  Bkgd fluctuations



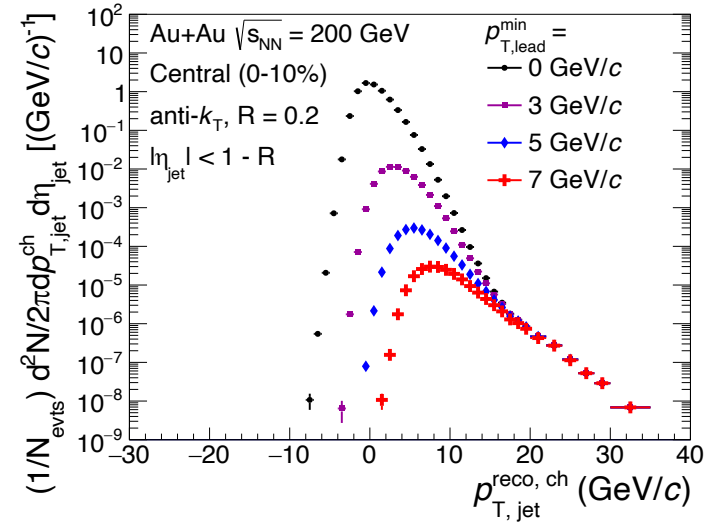
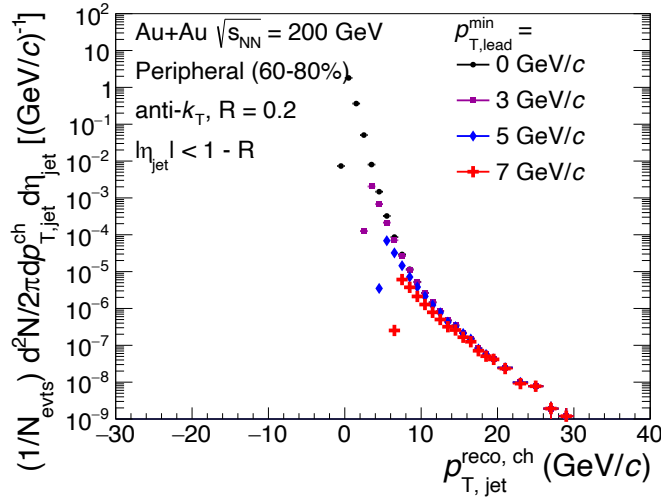
Syst. Uncert.  
 details in backup  
 slides

# Inclusive charged jets: raw data

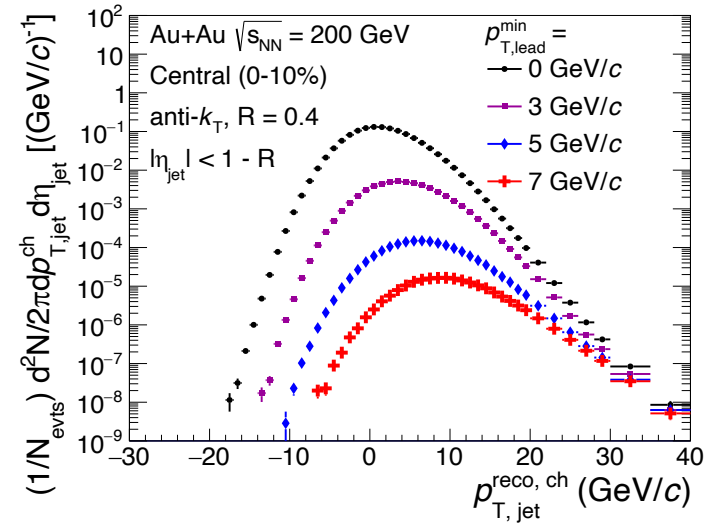
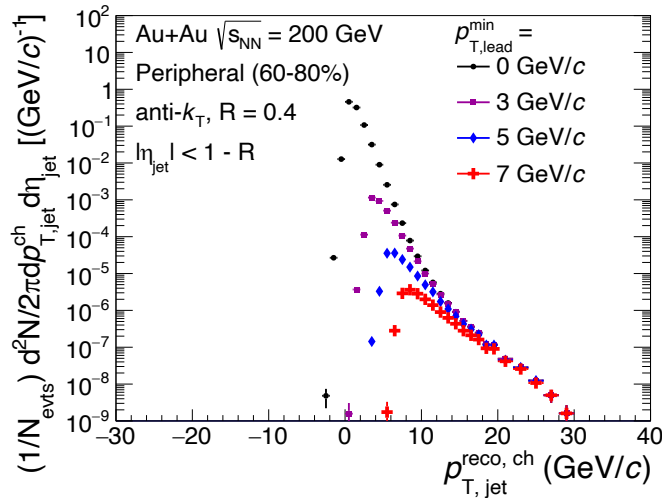
Peripheral Au+Au

Central Au+Au

R=0.2



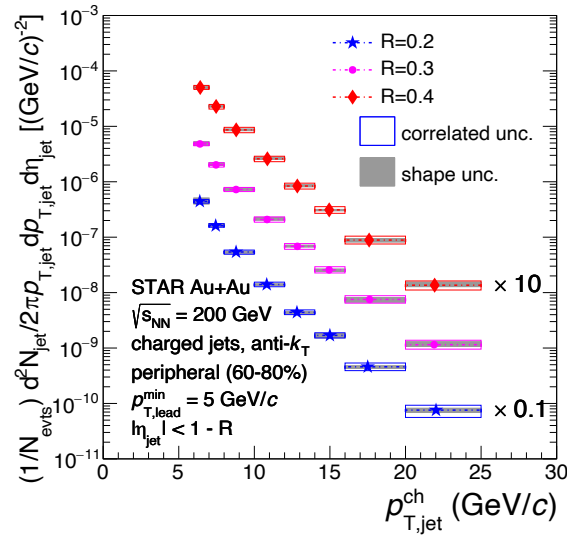
R=0.4



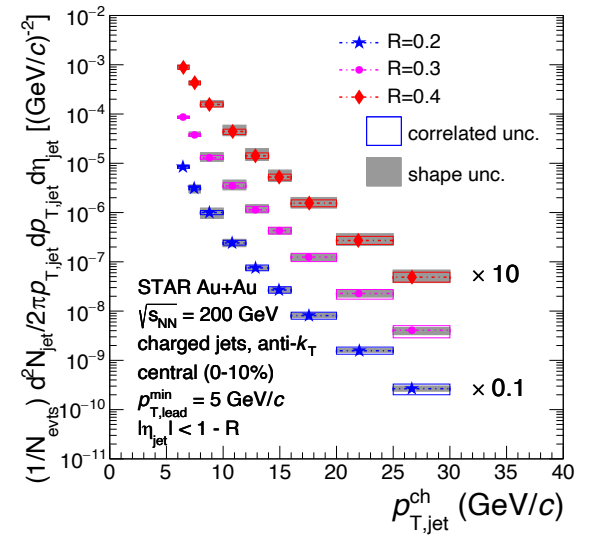
# Inclusive charged jets: corrected spectra

$$p_{T,\text{lead}}^{\text{min}} = 5 \text{ GeV}/c$$

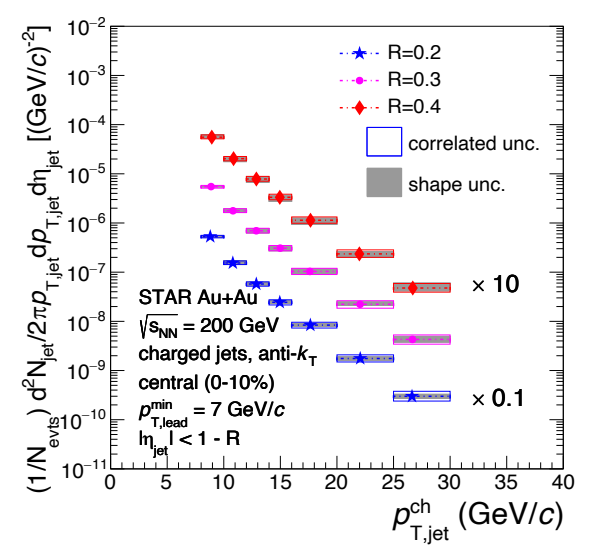
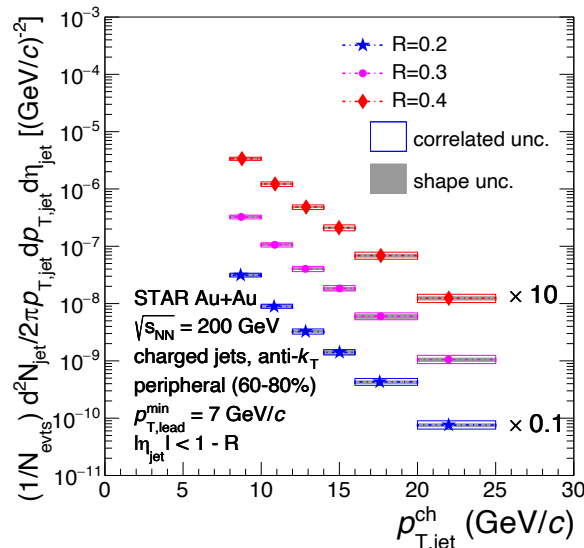
## Peripheral Au+Au



## Central Au+Au



$$p_{T,\text{lead}}^{\text{min}} = 7 \text{ GeV}/c$$



# Closure Test

## Full analysis on simulated data

- answer is known
- close the circle and check consistency

## Parametrized model

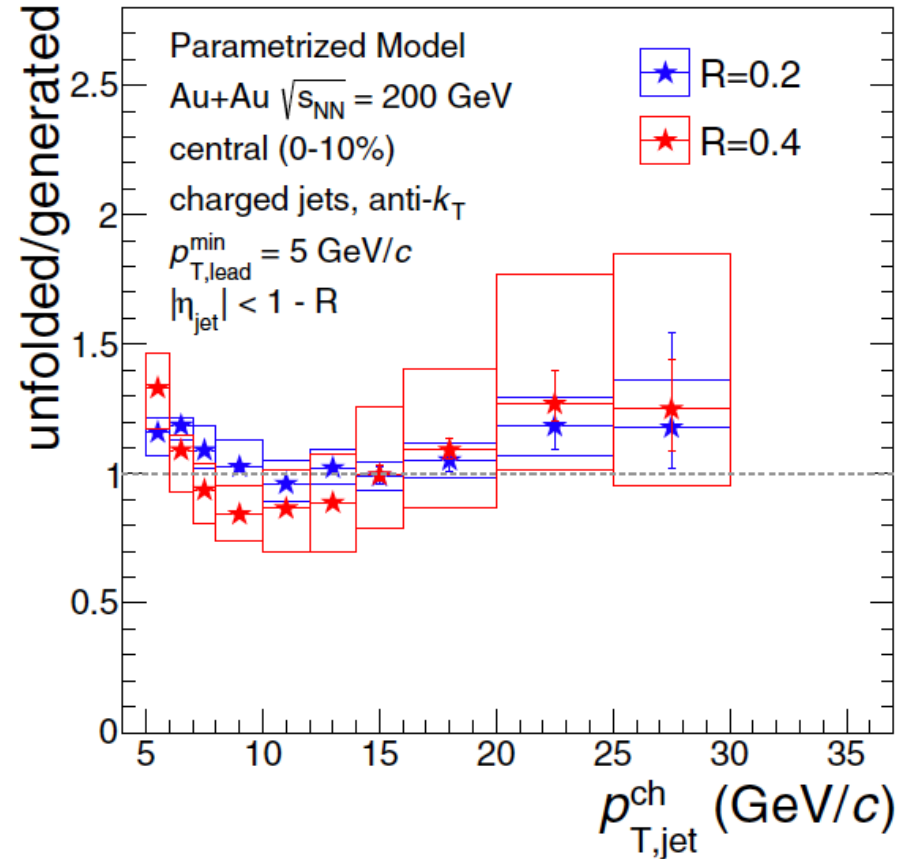
- “thermal” bkgd + PYTHIA jets + yield suppression
- good agreement with real data distributions (backup slides)

## Event generation

- similar statistical precision as real dataset

## Complete analysis chain

- including syst. uncert.



→ no evidence of bias beyond sys uncert band

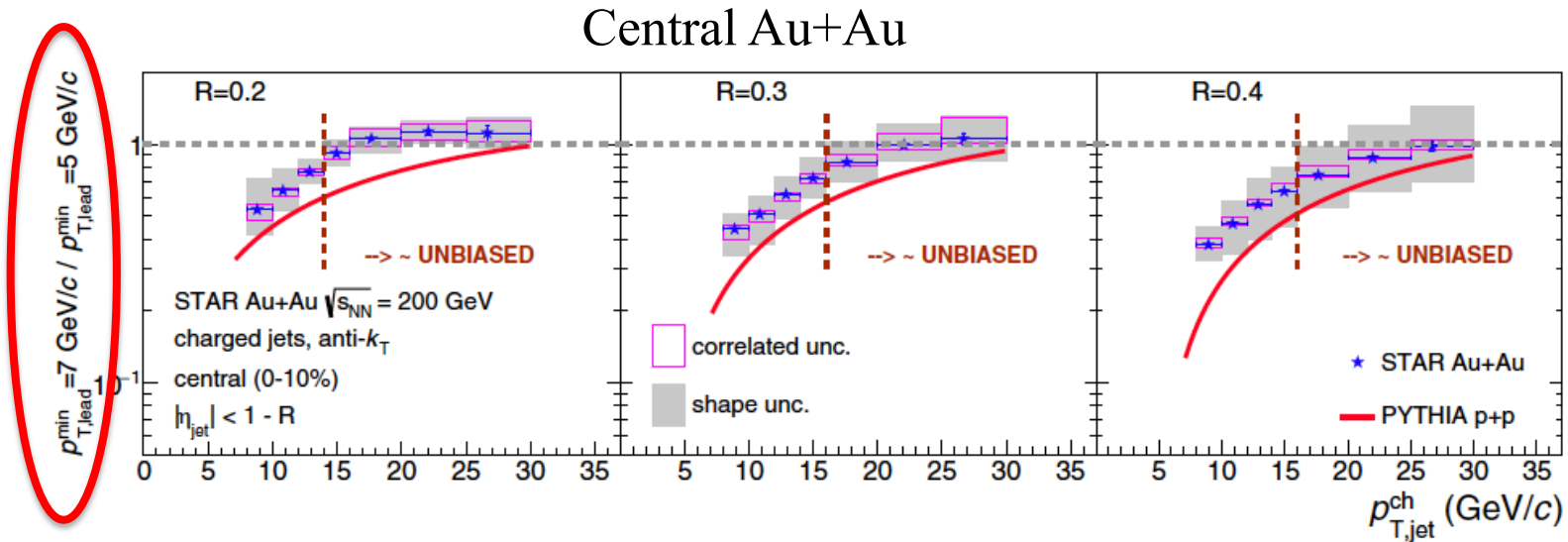


# Measure bias due to $p_{T,\text{lead}}^{\text{min}}$

Assertion: larger  $p_{T,\text{lead}}^{\text{min}} \rightarrow$  larger bias

Compare  $p_{T,\text{lead}}^{\text{min}} = 5$  and  $7 \text{ GeV}/c$

- ratio  $\sim$  unity within uncert.  $\rightarrow$  bias is negligible

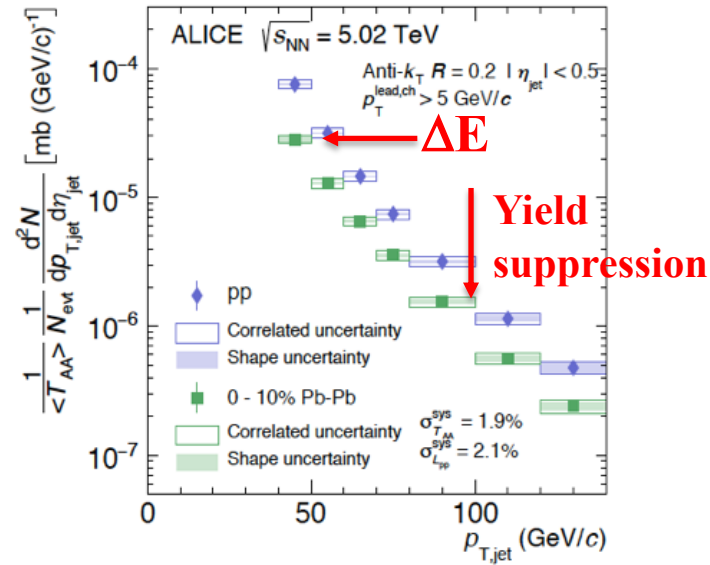


Curious fact: bias is smaller in central Au+Au than in PYTHIA p+p....?

- non-trivial fragmentation+quenching physics
- explore with next-generation calorimetric measurement, TBD

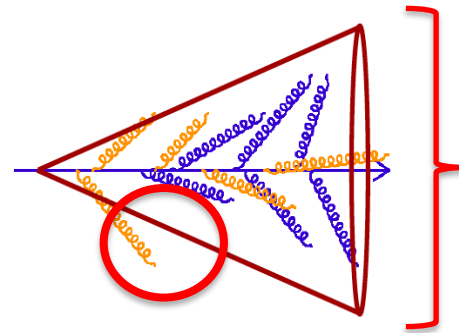


# Measuring jet energy loss



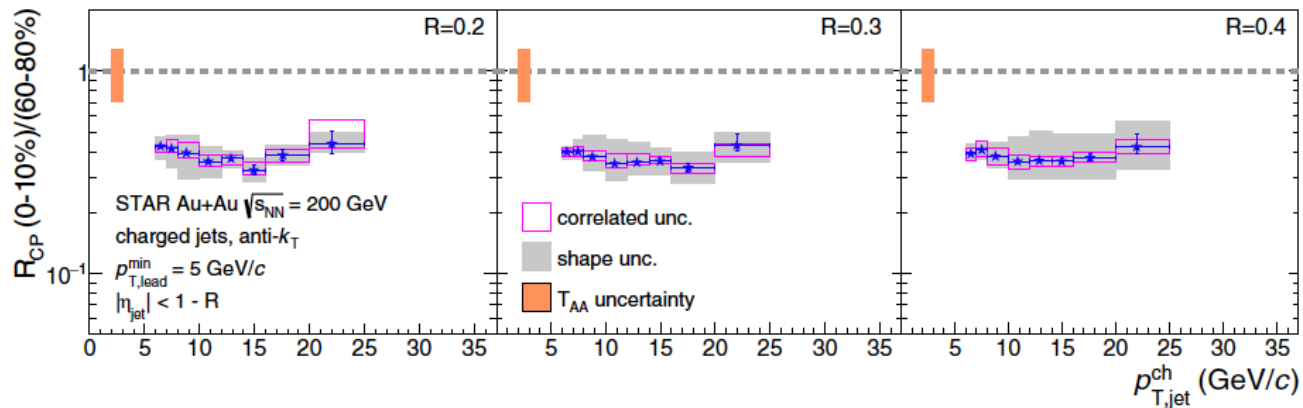
$$R_{AA} = \frac{\text{Rate in central A + A}}{\text{Rate in p + p} \otimes \text{geometry}}$$

$$R_{CP} = \frac{\text{Rate in central AA}}{\text{Rate in periph AA} \otimes \text{geometry}}$$

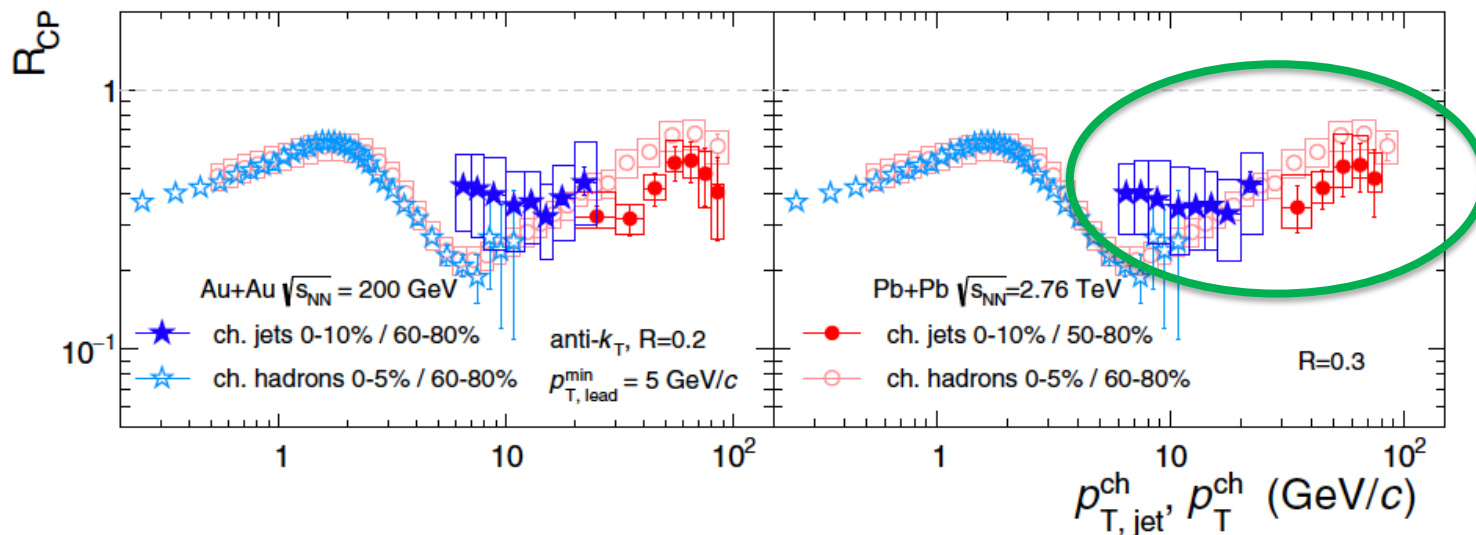


# Jet energy loss: $R_{CP}$

$p_{T,lead}^{min\ bias} \sim$  cancels in ratio, show full  $p_T$  range



Compare to charged hadrons & LHC charged jets



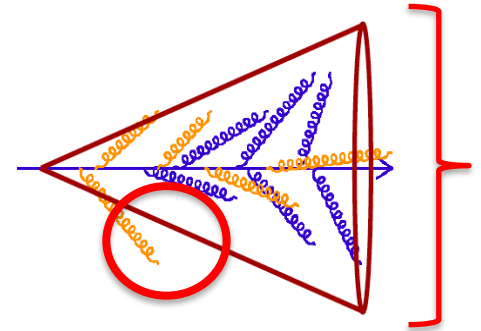
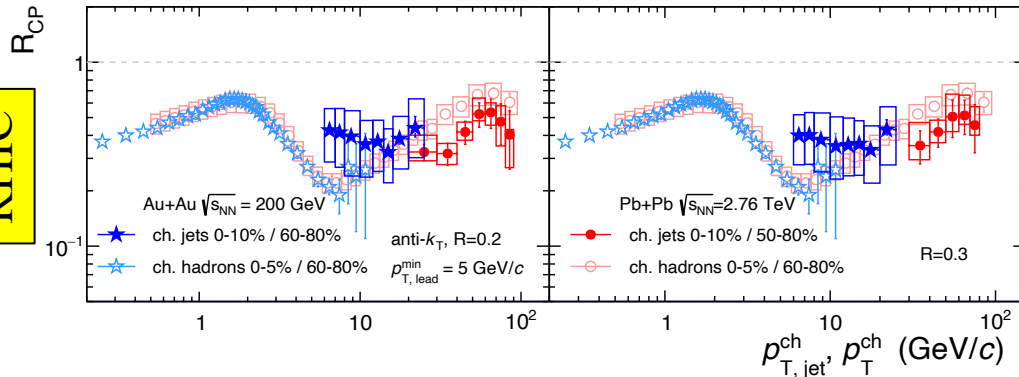
STAR  
 ALICE

Similar suppression w/ different spectrum shapes...??

# Jet suppression RHIC vs LHC: additional comparisons

## Charged jet $R_{CP}$

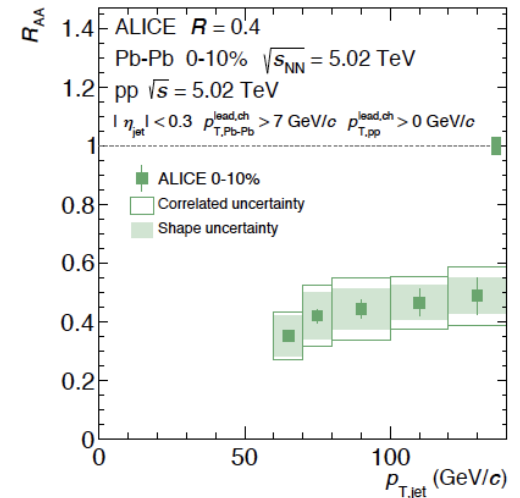
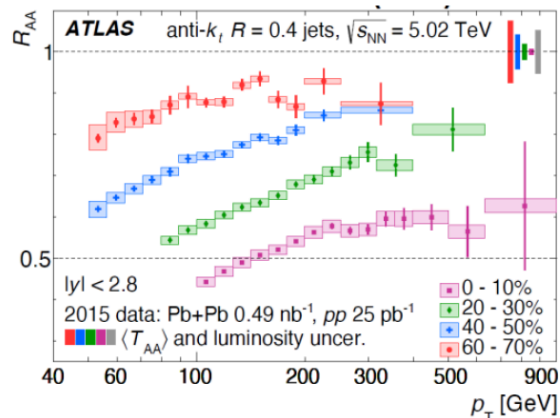
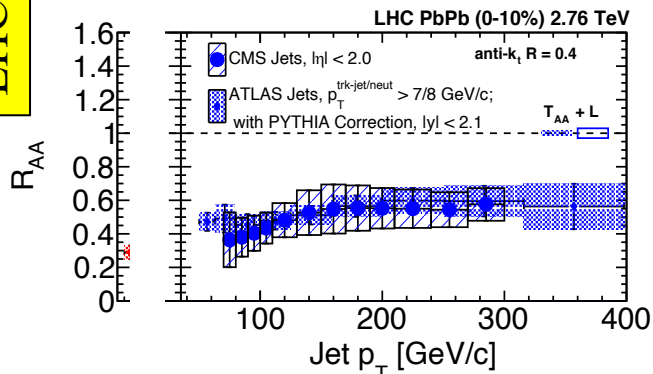
RHIC



- Strong jet yield suppression
- Suppression  $\sim$  similar magnitude at RHIC and LHC...?

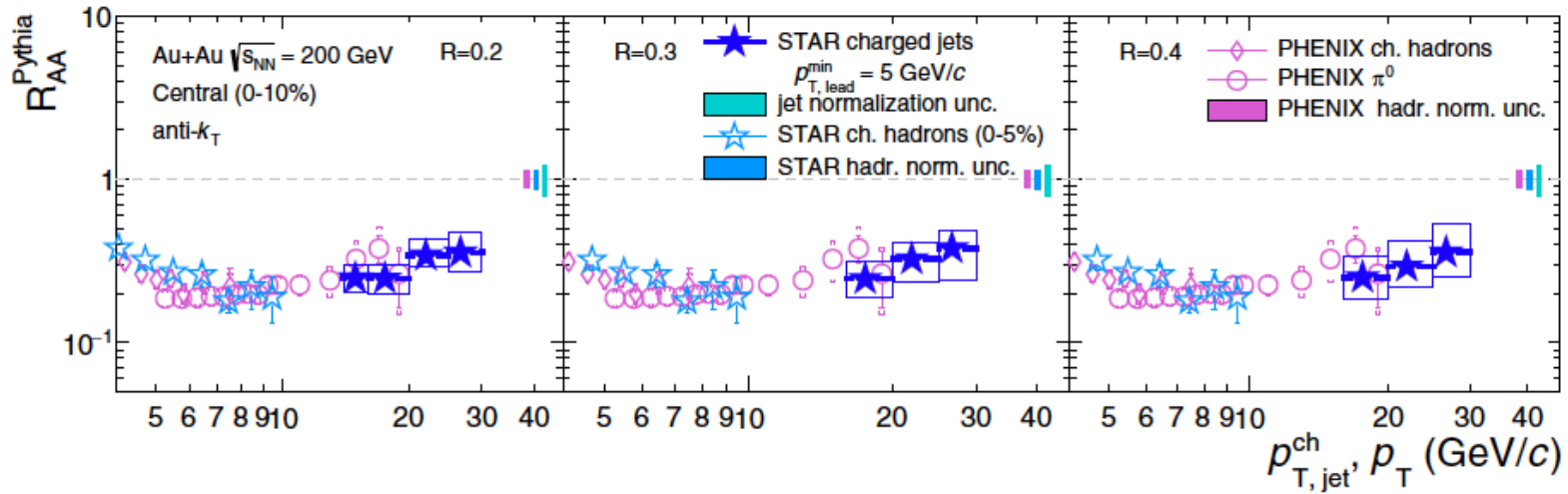
## Calorimetric jet $R_{AA}$

LHC

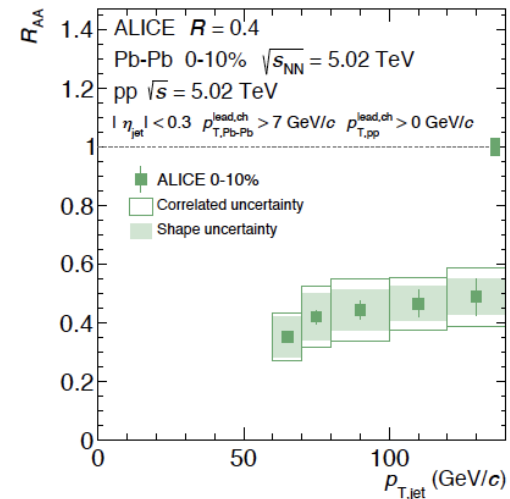


# Jet energy loss: $R_{AA}$

pp reference: PYTHIA STAR tune



Similar picture to  $R_{CP}$



# Inclusive jet $R_{AA}$ : comparison to models

Diverse jet quenching calculations based on pQCD  
+ various approximations for jet+medium interaction

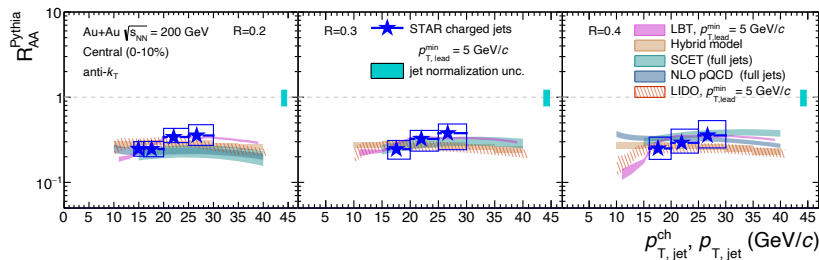
Current models work well over a wide range

Data relatively featureless, do not discriminate

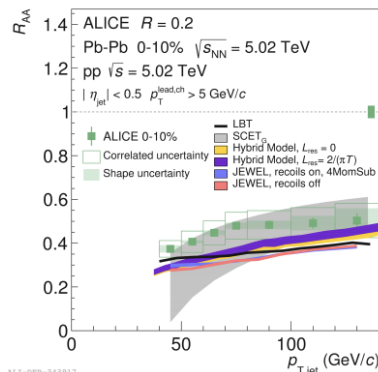
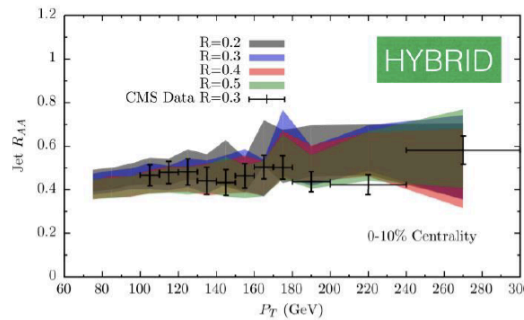
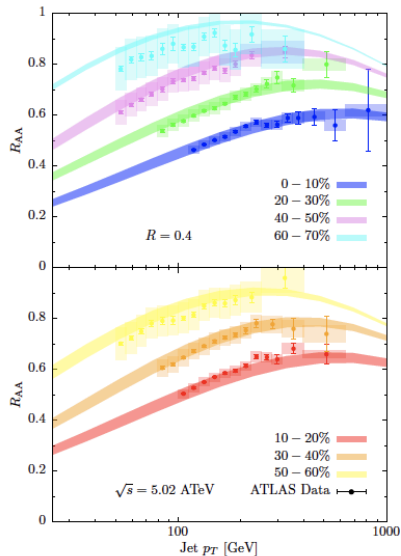
How to make progress?

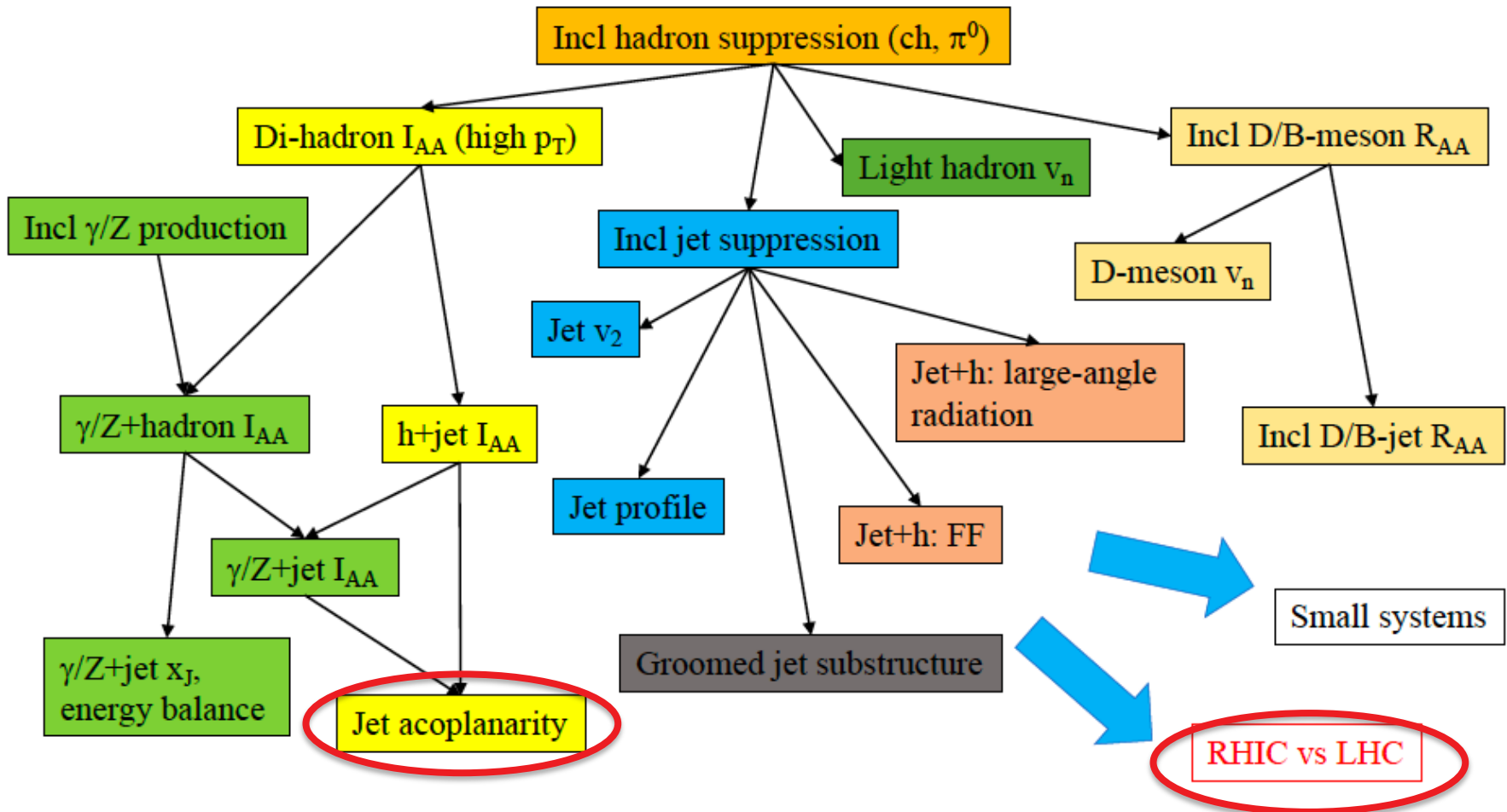
1. JETSCAPE: go beyond current formulation of qhat to capture full dynamics of jet-medium interaction  
→ global fits to hadron&jet data

2. Other observables with orthogonal parametric dependencies



arXiv:2101.01742

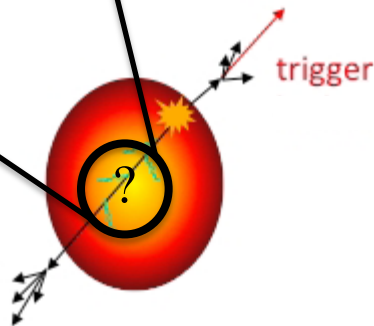
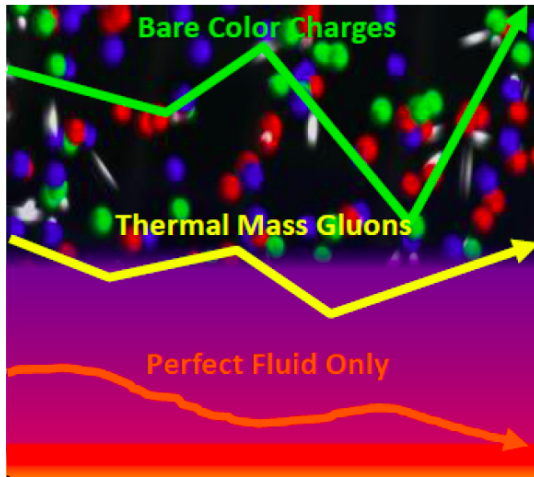




# Jet acoplanarity: in-medium hard scattering

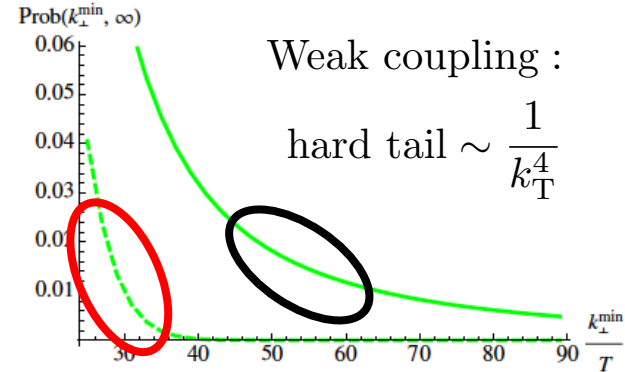
("Rutherford experiment")

Discrete scattering centers or effectively continuous medium?



*d'Eramo et al., JHEP 1305 (2013) 031*

Distribution of momentum transfer  $k_T$



Strong coupling:  
Gaussian distribution

What are the quasi-particles?

- high  $Q^2$ : bare q and g
- low-ish  $Q^2$ :
  - thermal-mass glue
  - magnetic monopoles
  - ...?

# Jet acoplanarity: in-medium soft deflection

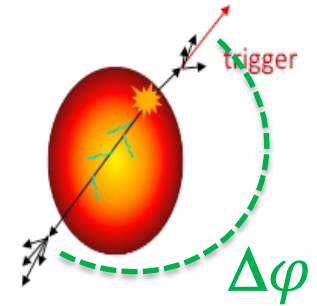
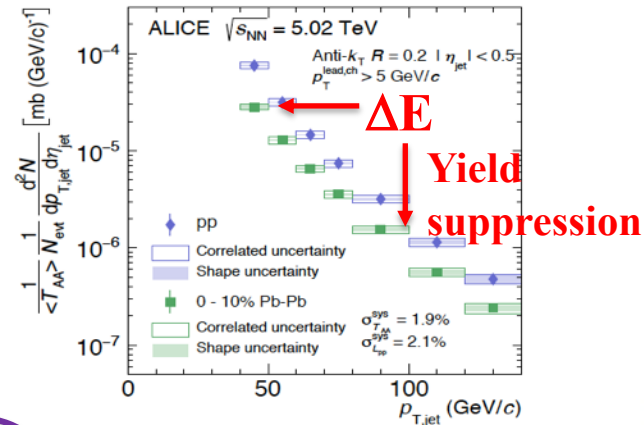
For intuition use BDMPS theory: multiple soft scattering approximation

Medium-induced jet energy loss:

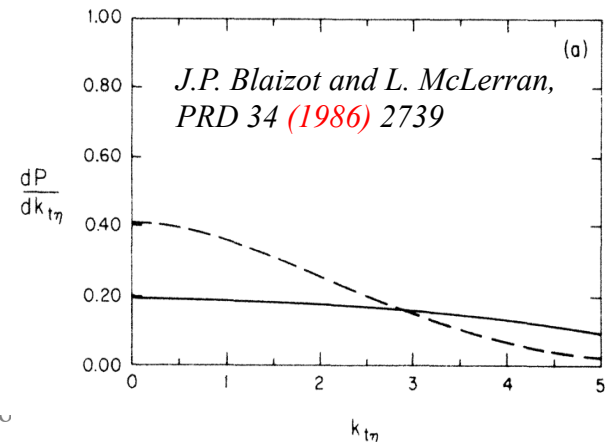
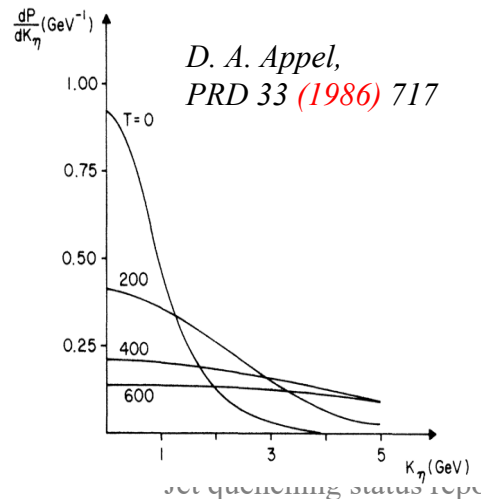
$$\Delta E_{med} \sim \alpha_s \hat{q} L^2$$

Medium-induced angular broadening:

$$\langle k_T^2 \rangle \sim \langle \Delta\varphi^2 \rangle \sim \alpha_s \hat{q} L$$



Different parametric dependencies → better model discrimination?

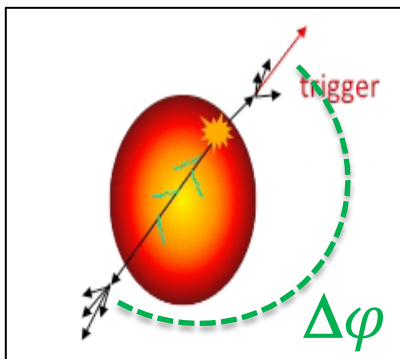


Side note: using jet scattering to measure the QGP is an old idea but experimentally very challenging

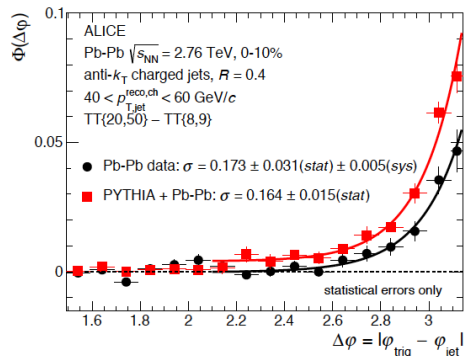
- techniques now in place



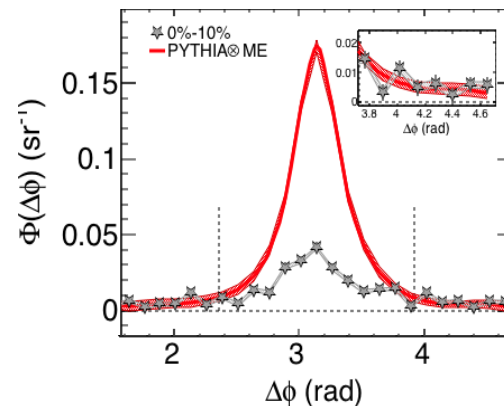
# Jet acoplanarity: data



ALICE, JHEP 09 (2015) 170

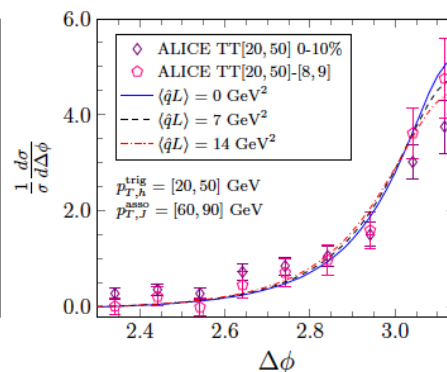
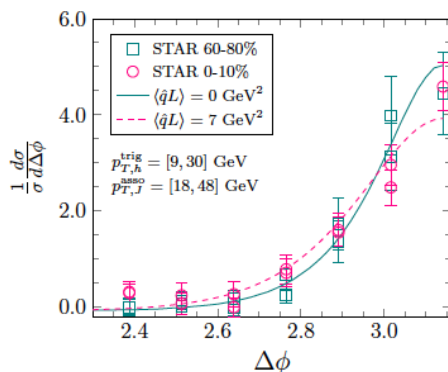
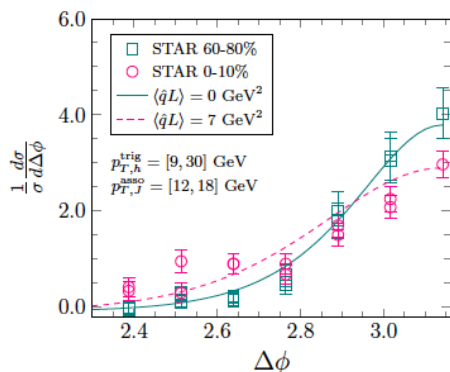


STAR, Phys Rev C 96 (2017) 024905



Significant background: Initial-state (Sudakov) radiation

L. Chen et al., Phys.Lett.B 773 (2017) 672

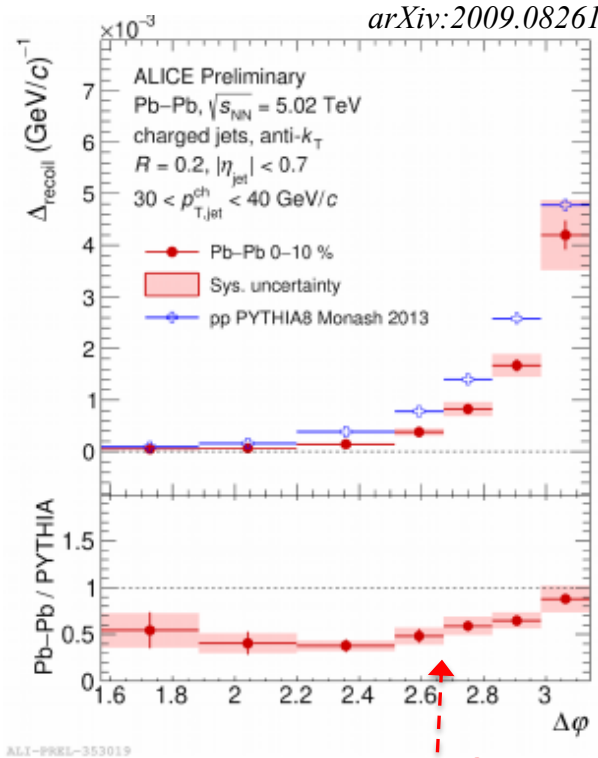
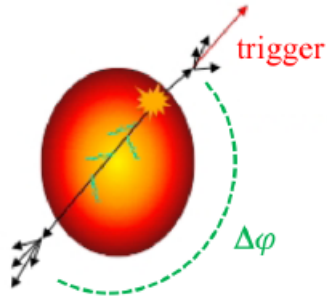


First- generation ALICE+STAR measurements:

no medium-induced acoplanarity observed above background

Second-generation measurements with greater precision in progress....

# Jet acoplanarity: ALICE Run 2



Work in progress:

- larger  $R$ , lower  $p_T^{\text{jet}}$
- measured p+p reference

Narrowing due to radiative corrections...?

arXiv:2003.10182

## Radiative $p_{\perp}$ -broadening of fast partons in an expanding quark-gluon plasma

B.G. Zakharov<sup>1</sup>

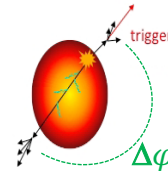
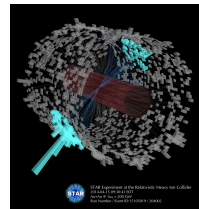
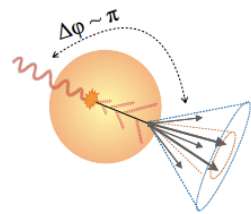
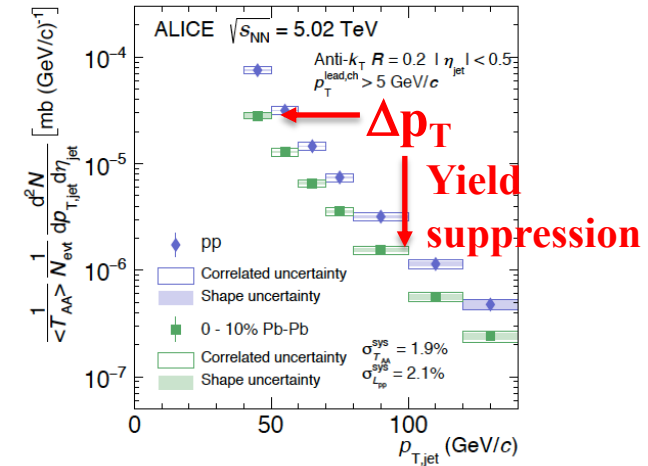
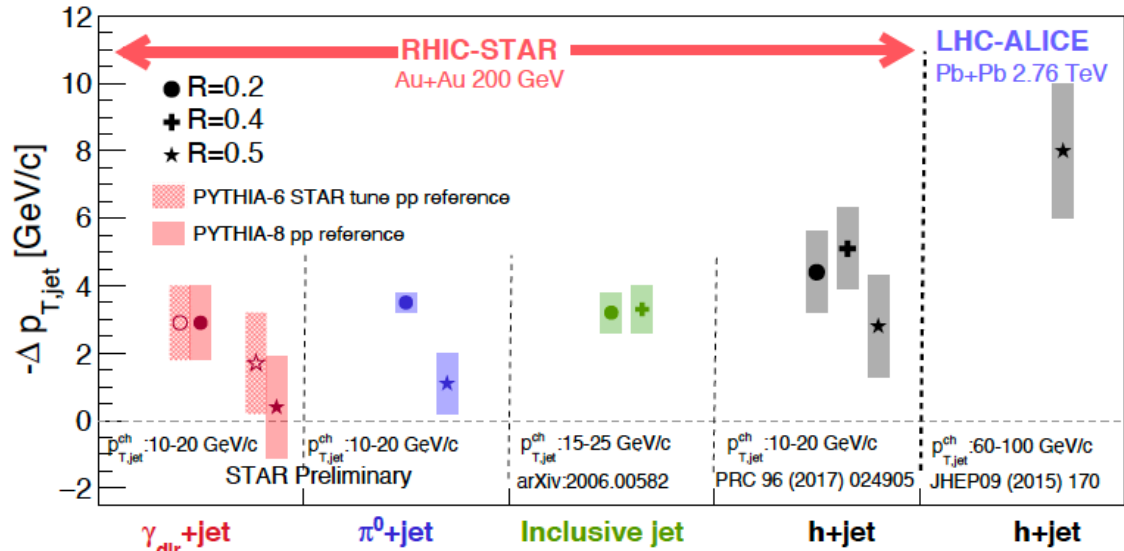
<sup>1</sup>L.D. Landau Institute for Theoretical Physics, GSP-1, 117940, Kosygina Str. 2, 117334 Moscow, Russia

(Dated: March 24, 2020)

We study contribution of radiative processes to  $p_{\perp}$ -broadening of fast partons in an expanding quark-gluon plasma. It is shown that the radiative correction to  $\langle p_{\perp}^2 \rangle$  for the QGP produced in AA-collisions at RHIC and LHC may be negative, and comparable in absolute value with the non-radiative contribution. We have found that the QGP expansion enhances the radiative suppression of  $p_{\perp}$ -broadening as compared to the static medium.

# Phenomenology: in-medium energy loss measured via jet spectrum shift

Inclusive jet and X+jet measurements



RHIC: energy loss similar for different probes

- possible R-dependence

LHC: energy loss larger than RHIC



Confrontation with theory calculations TBD

# Jet quenching: Outlook

## LHC

- Run 3 starts early 2022; factor  $\sim 10$  luminosity increase
- ALICE: essentially a new detector with vastly improved capabilities
- ATLAS/CMS moderate improvements (major upgrades  $\sim 2025$  for Run 4)
- Through Run 4 (2029): Pb+Pb @  $10 \text{ nb}^{-1}$

## RHIC

- New detector focused on jet physics: sPHENIX
- Upgraded STAR
- Through 2025: STAR Au+Au @  $110 \text{ nb}^{-1}$ ; sPHENIX Au+Au @  $23 \text{ nb}^{-1}$

→ At both facilities: factor  $\sim 10$  increase in data, much improved instrumentation

But experimental advances alone are not sufficient for quantitative understanding of jet quenching and the QGP

Theory and modelling:

- Conceptual and calculational advances in modelling of in-medium jet modification
  - Rigorous-large scale global fits to a wide range of judiciously chosen jet and hadron data
- Bayesian inference using JETSCAPE

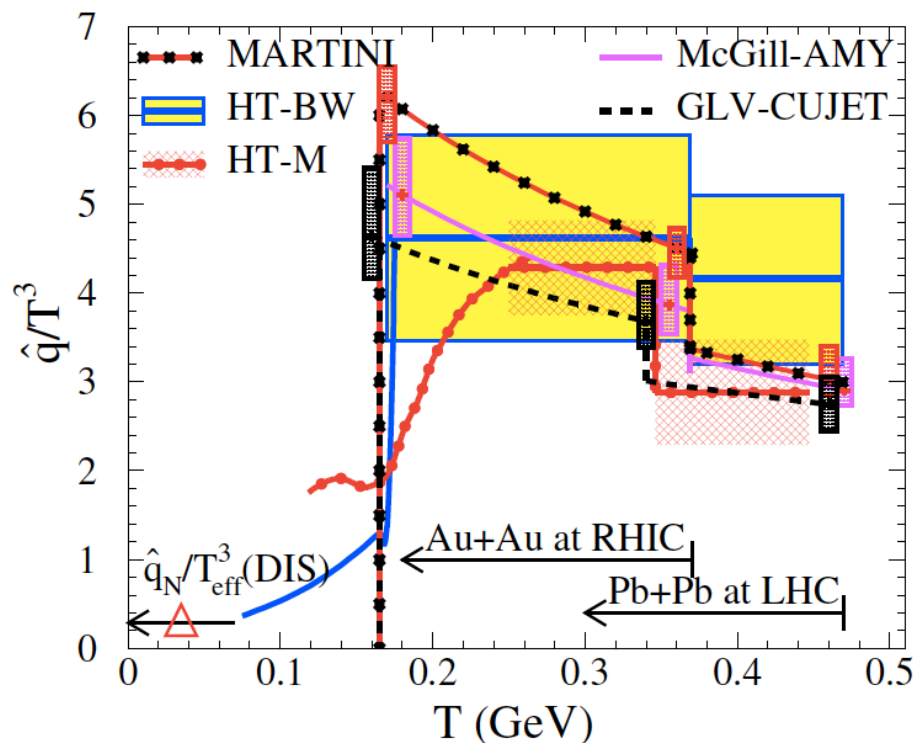
Jet quenching was discovered 20 years ago; still compelling, not yet solved...

# Extra slides

# Measuring $\hat{q}$ : inclusive hadron suppression

JET Collaboration

Phys.Rev. C90 (2014) 1, 014909



Fit pQCD-based models to **single-hadron suppression** data at RHIC and LHC

For a 10 GeV light quark at time 0.6 fm/c:

$$\text{RHIC} : \hat{q} \approx 1.2 \pm 0.3 \text{ GeV}^2/\text{fm}$$

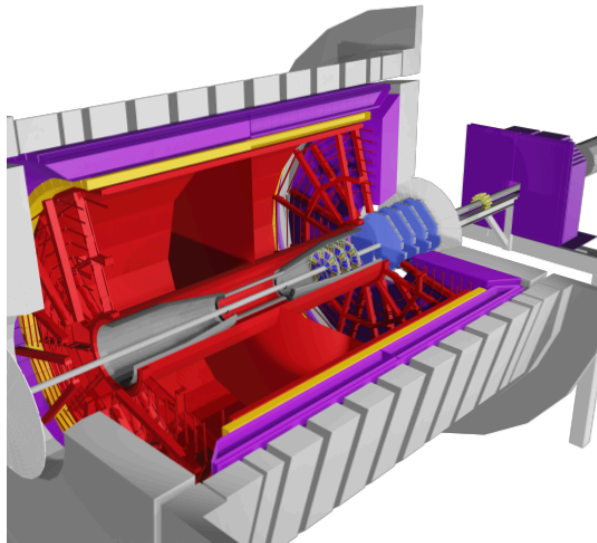
$$\text{LHC} : \hat{q} \approx 1.9 \pm 0.7 \text{ GeV}^2/\text{fm}$$

Reasonable and improvable precision

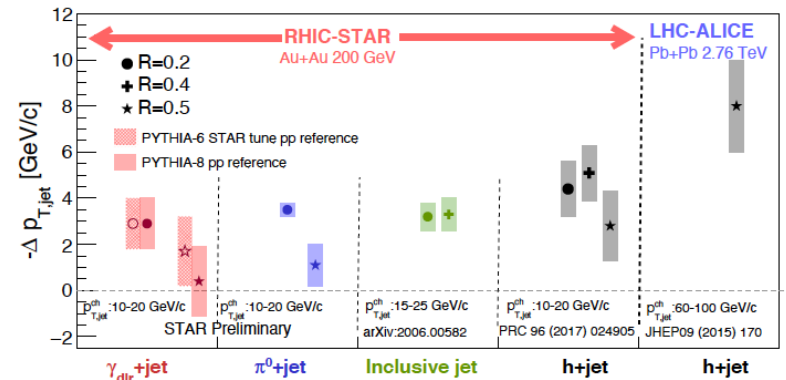
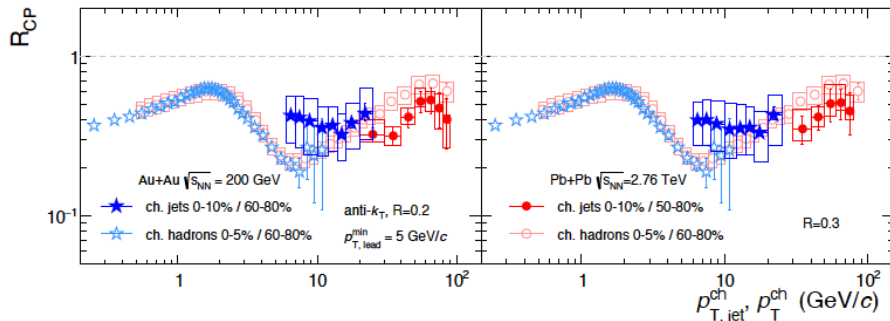
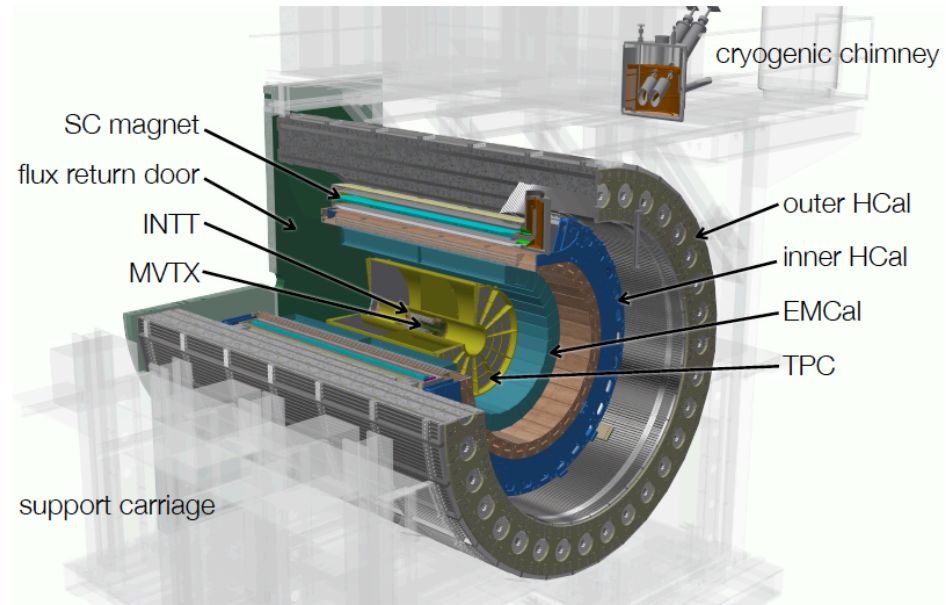
$$\text{Cold matter (e+A at HERA): } \hat{q} \approx 0.02 \text{ GeV}^2/\text{fm}$$

# RHIC & LHC: the present

STAR



sPHENIX (under construction)



# RHIC: the future

Beam Use Request to RHIC PAC, Sept 2020

## STAR

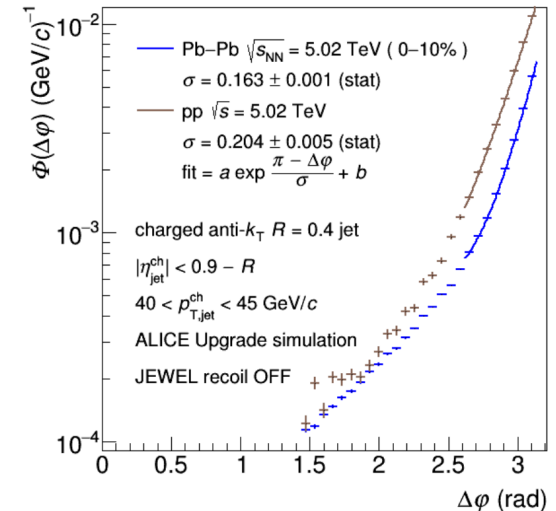
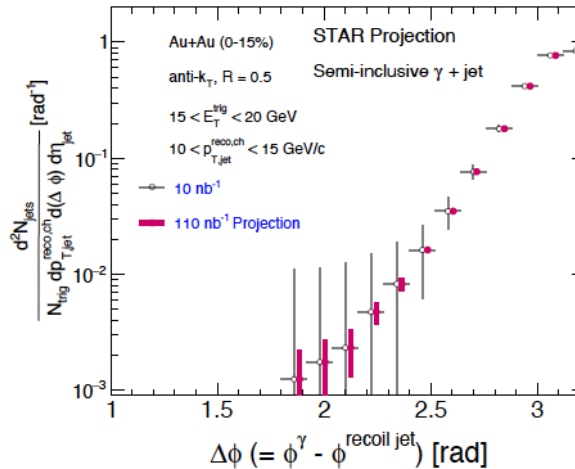
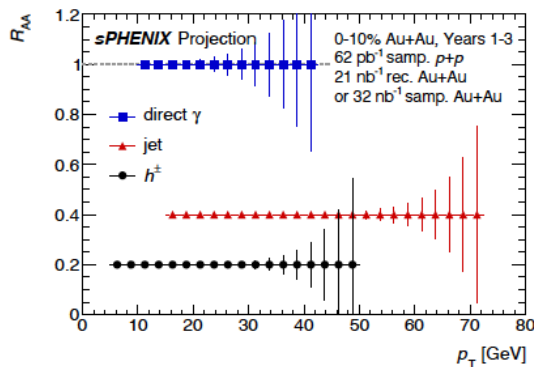
year	minimum bias [ $\times 10^9$ events]	high- $p_T$ int. luminosity [ $\text{nb}^{-1}$ ]		
		all vz	$ vz  < 70\text{cm}$	$ vz  < 30\text{cm}$
2014	2	26.5	19.1	15.7
2016				
2023	10	43	38	32
2025	10	58	52	43

## sPHENIX

Year	Species	$\sqrt{s_{NN}}$ [GeV]	Cryo Weeks	Physics Weeks	Rec. Lum. $ z  < 10\text{ cm}$	Samp. Lum. $ z  < 10\text{ cm}$
2023	Au+Au	200	24 (28)	9 (13)	3.7 (5.7) $\text{nb}^{-1}$	4.5 (6.9) $\text{nb}^{-1}$
2024	$p^\uparrow p^\uparrow$	200	24 (28)	12 (16)	0.3 (0.4) $\text{pb}^{-1}$ [5 kHz] 4.5 (6.2) $\text{pb}^{-1}$ [10%-str]	45 (62) $\text{pb}^{-1}$
2024	$p^\uparrow + \text{Au}$	200	-	5	0.003 $\text{pb}^{-1}$ [5 kHz] 0.01 $\text{pb}^{-1}$ [10%-str]	0.11 $\text{pb}^{-1}$
2025	Au+Au	200	24 (28)	20.5 (24.5)	13 (15) $\text{nb}^{-1}$	21 (25) $\text{nb}^{-1}$

Au+Au total int lumi through 2025:

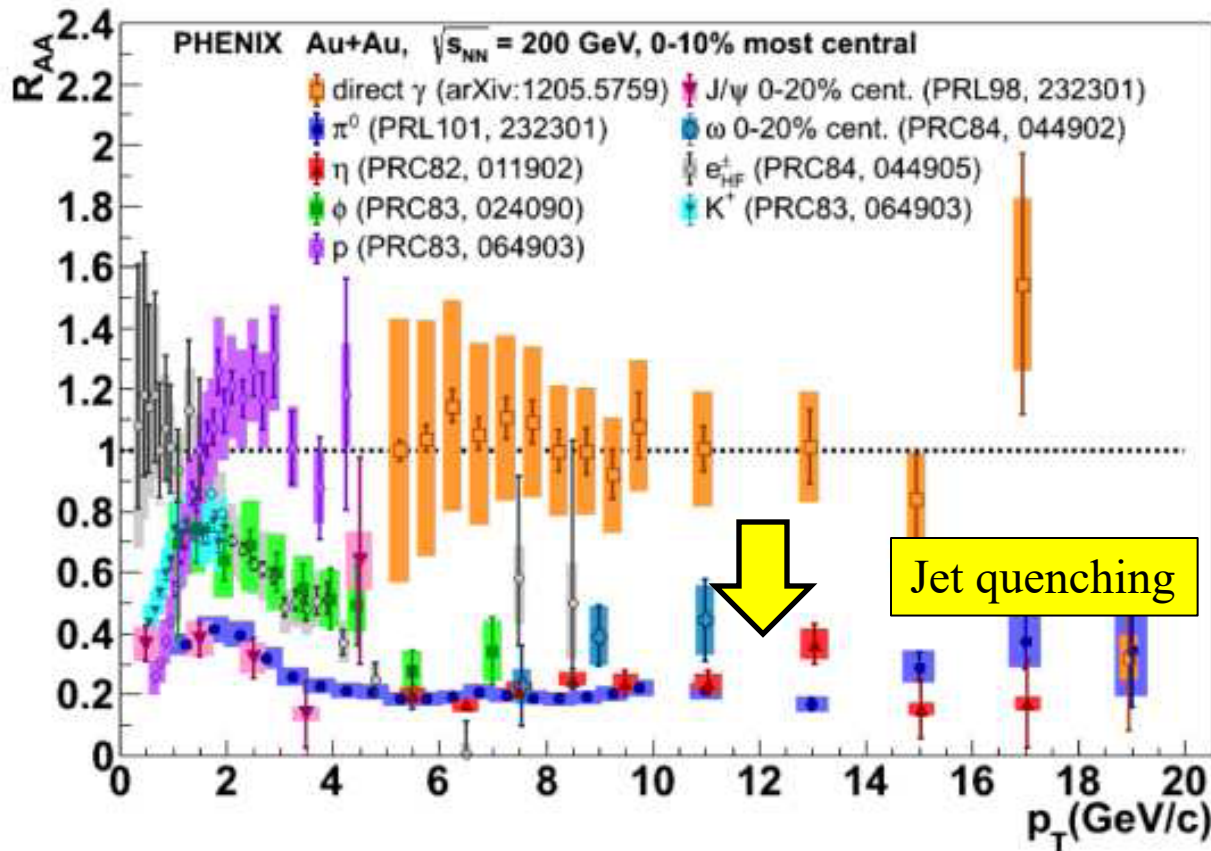
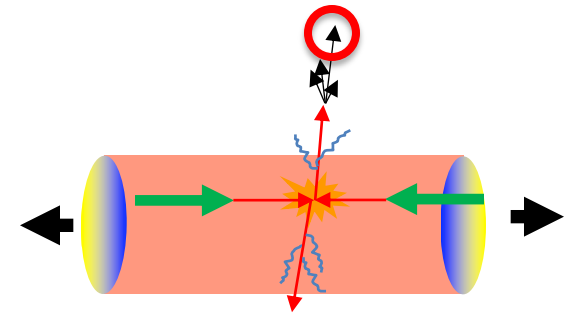
- STAR: 110  $\text{nb}^{-1}$
- sPHENIX: 23  $\text{nb}^{-1}$





# Jet quenching via high $p_T$ hadrons

$$R_{AA} = \frac{\text{Observed rate in AA}}{\text{Expected rate from } pp \otimes \text{ geometry}}$$

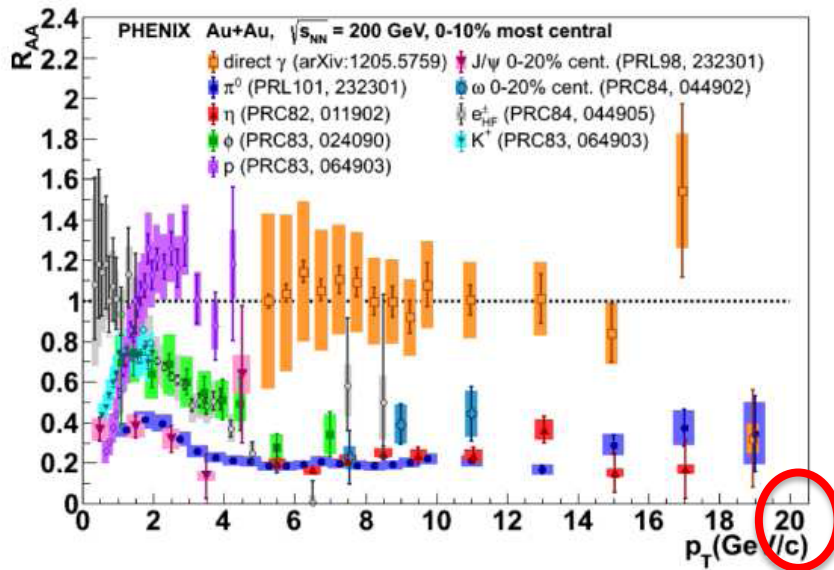


Photons (color-neutral)

Jet fragments (color-charged)

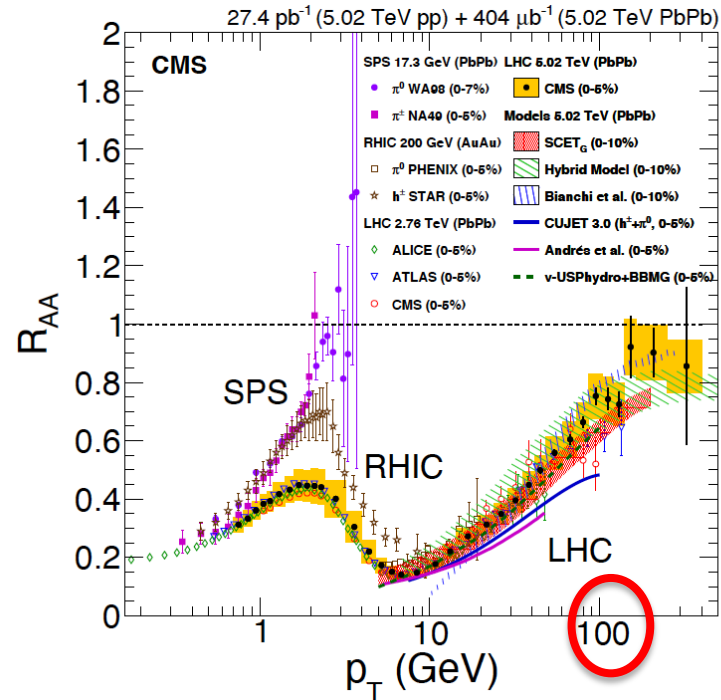
# Inclusive hadron suppression: RHIC vs LHC

RHIC



LHC

*JHEP 04 (2017) 039*



RHIC/LHC: Qualitatively similar, quantitatively different

- interplay between energy loss ( $\sim$ matter density) and spectrum shape

TABLE I. Components of the systematic uncertainty (%) for jets with  $R = 0.2, 0.3,$  and  $0.4$  in central and peripheral Au+Au collisions. See text for details.

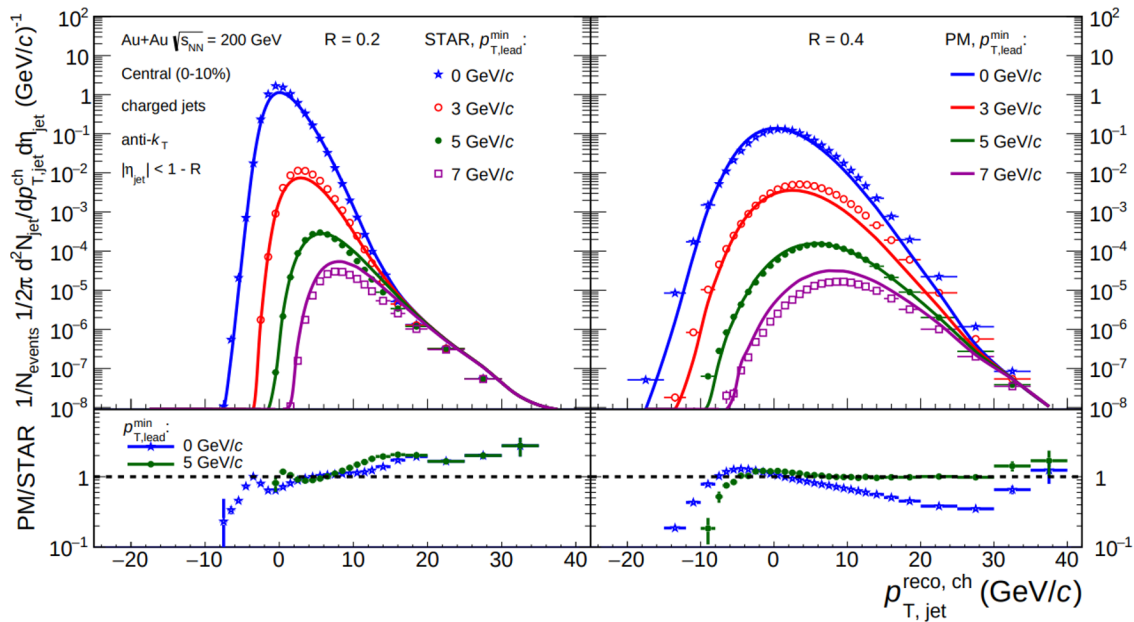
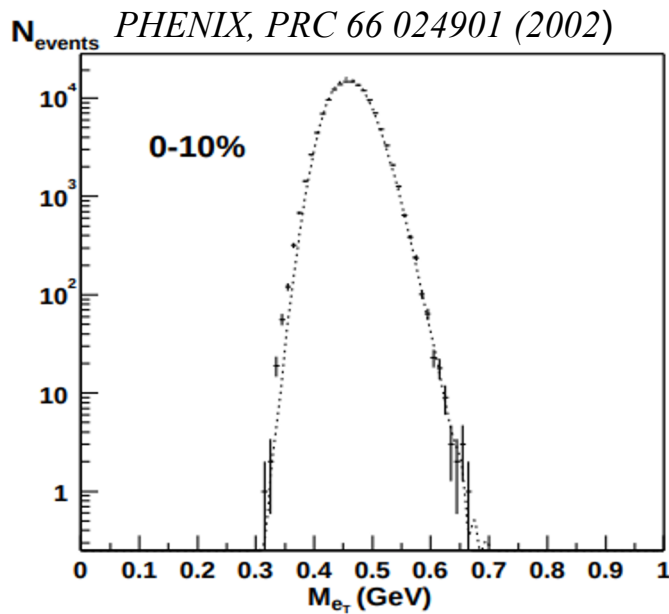
		Central Au+Au collisions, $\sqrt{s_{NN}}= 200$ GeV						Peripheral Au+Au collisions, $\sqrt{s_{NN}}= 200$ GeV					
$R$		0.2		0.3		0.4		0.2		0.3		0.4	
$p_{T,\text{jet}}^{\text{ch}}$ (GeV/ $c$ )		[14,16]	[20,25]	[14,16]	[20,25]	[14,16]	[20,25]	[14,16]	[18,20]	[14,16]	[18,20]	[14,16]	[18,20]
<b>Correlated</b>	Tracking efficiency	+15 -12	+16 -10	+16 -13	+12 -22	+14 -11	+18 -12	+6 -8	+10 -12	+12 -11	+14 -12	+13 -12	+16 -12
	Fragmentation for $R_{\text{det}}$	+1 -3	+3 -1	+3 -1	+4 -5	+4 -1	+12 -2	+0 -5	+0 -5	+0 -1	+2 -2	+2 -1	+3 -1
	$\delta p_T$	+8 -3	+16 -1	+10 -2	+17 -2	+7 -5	+14 -3	+10 -1	+15 -1	+9 -1	+11 -1	+8 -1	+11 -1
	$\rho$	+1 -1	+1 -1	+1 -0	+0 -1	+1 -1	+1 -1	+1 -3	+4 -1	+1 -3	+2 -4	+1 -3	+1 -4
	<b>Total correlated</b>	+17 -13	+24 -10	+19 -13	+21 -23	+17 -11	+26 -13	+12 -10	+18 -14	+15 -11	+18 -13	+15 -12	+20 -13
	<b>Shape</b>	+17 -14	+12 -10	+24 -19	+25 -18	+46 -29	+51 -31	+14 -11	+8 -7	+8 -6	+17 -12	+4 -3	+11 -9

# Background Description - Parametrized Model

Closure test utilizes simple model for background: Boltzmann-distributed independent emission with hard jet fragmentation based on PYTHIA p+p calculation

E-by-e ET fluctuations well-described by Boltzmann indep. emission

Also works well for jet measurements



Picture consistent with good description of jet background by Mixed Events (STAR h+jet)

Heavy-ion jet measurement background strongly dominated by statistical phase space  
 Contrary to conventional wisdom: the problem is simple!

# Jet broadening: $R=0.2/R=0.4$

