

# About Radiative Corrections

11. 1. 2021

H. Spiesberger

PRISMA<sup>+</sup> Cluster of Excellence,  
Institut für Physik,

Johannes Gutenberg-Universität Mainz



# Overview

- About some general properties of radiative corrections
- Remarks about radiative corrections for DVCS
- To Do

- Emission of **real photons**  
experimentally often not distinguished from non-radiative processes:  
soft photons, collinear photons  
→ "radiative corrections"
- Virtual corrections: **loop diagrams**  
needed to cancel infrared divergences (Bloch-Nordsieck)

## Corrections for electron scattering

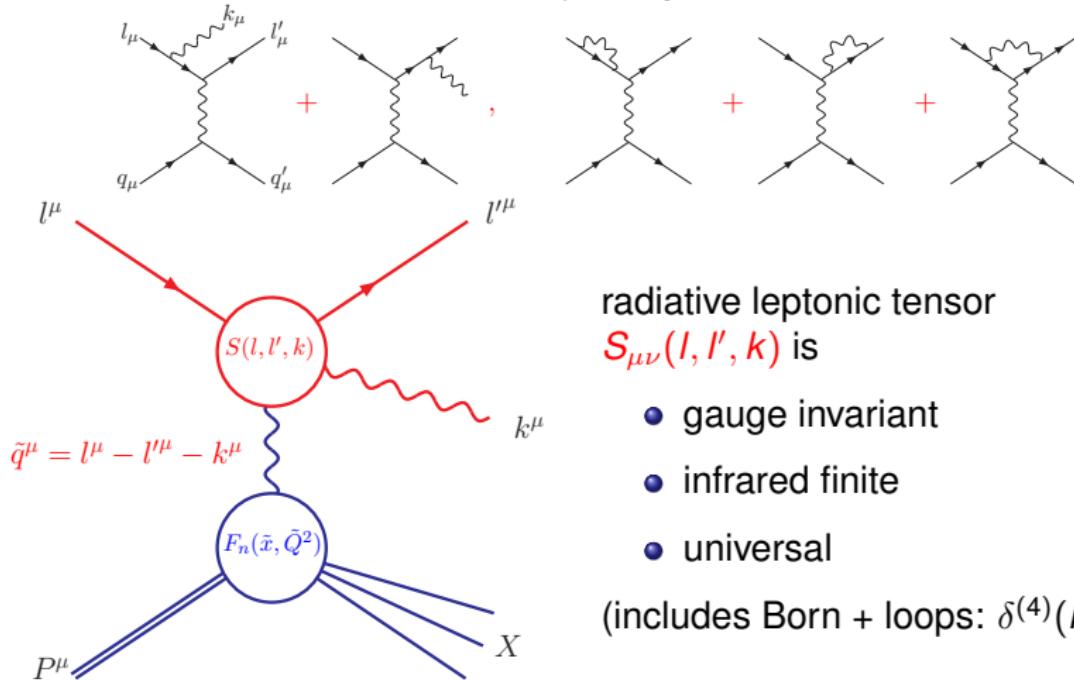
- Radiation from the electron  
model independent (universal)
- Radiation from the hadronic initial/final state  
parton model: radiation from quarks  
part of the nucleon structure
- Interference of electronic and hadronic radiation  
 $2\gamma$  exchange  
new structure
- vacuum polarization  
universal
- purely weak corrections

Note: for NC-scattering straightforward separation  
IR divergences: need to combine real and virtual radiation

# Leptonic radiation

Feynman diagrams for leptonic radiation at  $O(\alpha)$  (NC)

for eq scattering:



radiative leptonic tensor  
 $S_{\mu\nu}(l, l', k)$  is

- gauge invariant
- infrared finite
- universal

(includes Born + loops:  $\delta^{(4)}(k^\mu)$ )

## Radiator function

Observed cross section:

Convolution of true cross section  $\otimes$  radiator functions

$$d\sigma^{\text{obs}}(P, q) = \int \frac{d^3 k}{2k^0} \sum_n R_n(I, I', k) d\hat{\sigma}_n^{\text{true}}(P, q - k)$$

## Shifted kinematics

observed momentum transfer:  $Q^2 = -q^2$ ,  $q^\mu = I^\mu - I'^\mu$ ,

→ true, shifted momentum transfer:  $\tilde{Q}^2 = -(q - k)^2$

→ correction factor: enhancement by  $Q^2/\tilde{Q}^2$  → radiative tail

→ expect strong dependence on experimental prescriptions for measuring kinematic variables

→ need full Monte-Carlo modelling

Can be extended to include higher-order effects: multi-photon emission, soft-photon exponentiation,  $e^+ e^-$ -pair creation

## Properties of leptonic radiation

with partial fractioning, write:  $R_n(l, l', k) = \frac{J}{k \cdot l} + \frac{F}{k \cdot l'} + \frac{C}{\tilde{Q}^2} + \dots$

- initial state radiation,  $k \cdot l$  small for  $\not{e}_{in}, \gamma \rightarrow 0$
- final state radiation,  $k \cdot l'$  small for  $\not{e}_{out}, \gamma \rightarrow 0$
- Compton peak,  $\tilde{Q}^2$  small for  $p_T(e_{out}) \simeq p_T(\gamma)$

ISR, FSR: narrow peaks, width  $\simeq \sqrt{m_l/E_l}$ : collinear or mass singularities

upon angular integration: large logarithm  $\propto \frac{\alpha}{\pi} \log \frac{Q^2}{m_e^2} \simeq 10\%$

Note: additional large logarithms from experimental cuts  $\propto \log \frac{\Delta E}{E_{max}}$

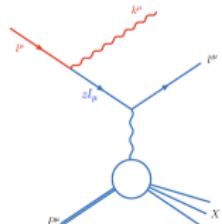
A technical remark about numerical integration:

Do partial fractioning, choose denominator as integration variable

## Collinear approximation (peaking approximation)

e.g., for initial-state radiation: assume  $k^\mu = (1 - z)l^\mu$   
→ Radiator function

$$R_{\text{ISR}} = \frac{\alpha}{2\pi} \frac{1 + z^2}{1 - z} \log \frac{Q^2}{m_e^2}$$



( $+\delta(1 - z)$  from loops →  $+$ -distribution  $1/(1 - z)_+$ )

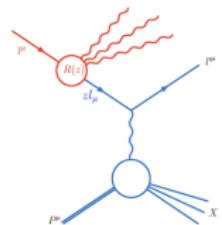
$$d\sigma_{\text{ISR}} = \int \frac{dz}{z} R_{\text{ISR}}(z) d\sigma_{\text{Born}}(l^\mu \rightarrow z l^\mu)$$

(similar for final-state radiation)

Can be extended to include multi-photon emission:

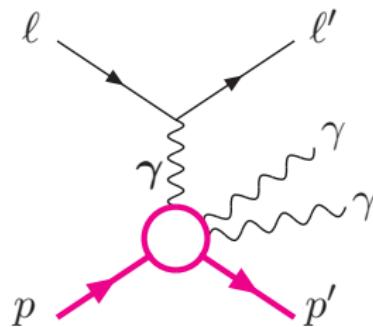
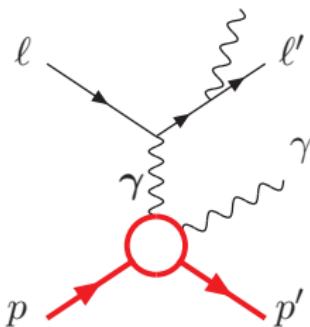
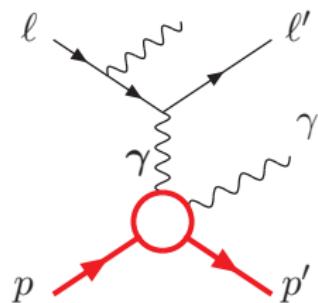
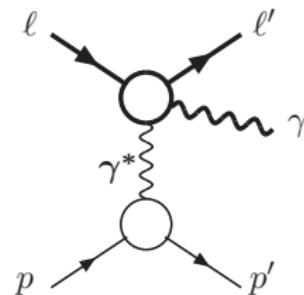
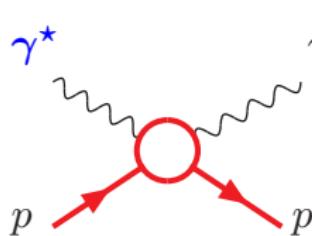
$$R_{\text{ISR}}^{(2)}(z) = \int_z^1 \frac{dz'}{z'} R_{\text{ISR}}^{(1)}(z') R_{\text{ISR}}^{(1)}(z/z') + \dots$$

Solution of evolution equations like DGLAP  
Known at  $O(\alpha^2)$  (complete) and partially at  $O(\alpha^3)$



# Corrections for Virtual Compton Scattering

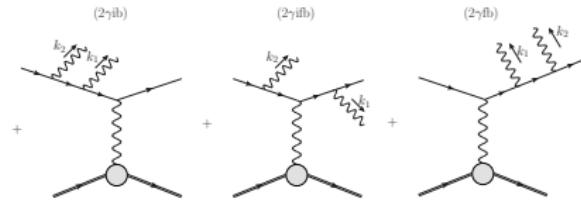
Leading order



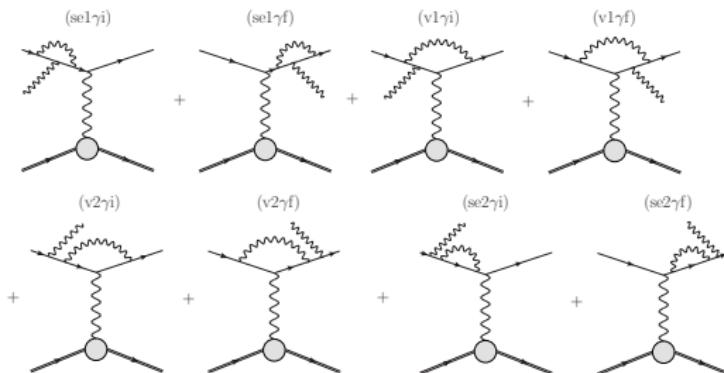
Leptonic corrections for VCS  
(+ loops)

VCS + 1 $\gamma$

# Corrections for VCS: leptonic $2\gamma$ radiation



## 2-photon radiation



## 1-loop corrected 1-photon radiation

VCS at LO,  $e(p_-) + p(p) \rightarrow e'(p'_-) + p'(p') + \gamma(k_1)$

$$\sigma_{\text{LO}} = \int_{x_{\min}}^{x_{\max}} dx \int_{Q_{\min}^2}^{Q_{\max}^2} dQ^2 \int_{t_{1,\min}}^{t_{1,\max}} dt_1 \int_{\phi_{\min}}^{\phi_{\max}} d\phi \hat{\sigma}_{\text{LO}}(x, Q^2, t_1, \phi)$$

$$x_{\min} = \text{input},$$

$$x_{\max} = \text{input},$$

$$Q_{\min}^2 = \text{input},$$

$$Q_{\max}^2 = \min(x s, \text{input}),$$

$$t_{1,\min} = \max(t_1^-, \text{input}),$$

$$t_{1,\max} = \min(t_1^+, \text{input}),$$

$$\phi_{\min} = 0,$$

$$\phi_{\max} = 2\pi.$$

$$t_1^\pm(r=0) = -\frac{Q^2}{1+4\tau\frac{1-x}{x}} \left( 1 + 2\tau \frac{1-x}{x^2} \left( 1 \pm \sqrt{1 + \frac{x^2}{\tau}} \right) \right).$$

## Corrections for VCS: Radiative process: $e(p_-) + p(p) \rightarrow e'(p'_-) + p'(p') + \gamma(k_1) + \gamma(k_2)$

$$\sigma_{\text{rad}} = \int_{x_{\min}}^{x_{\max}} dx \int_{Q_{\min}^2}^{Q_{\max}^2} dQ^2 \int_0^{R_{\max}^2} dR^2 \int_{t_{1,\min}}^{t_{1,\max}} dt_1 \int_{\phi_{\min}}^{\phi_{\max}} d\phi \\ \times \int_{c_{1,\min}}^{c_{1,\max}} d\cos\theta_1 \int_{\phi_{1,\min}}^{\phi_{1,\max}} d\phi_1 \hat{\sigma}_{\text{rad}}(x, Q^2, t_1, \phi; R^2, \theta_1, \phi_1)$$

$x_{\min}$  = input ,

$x_{\max}$  = input ,

$Q_{\min}^2$  = input ,

$Q_{\max}^2$  = min (xs, input) ,

$$R_{\max}^2 = Q^2 \left( A_r + \sqrt{A_r^2 - \left( \frac{1-x}{x} \right)^2} \right) ,$$

$$A_r = \frac{1}{x\tau_x} + \frac{1-x}{x} \cdot \frac{\tau_x+x}{\tau_x+2x} ,$$

$t_{1,\min}$  = max ( $t_1^-$ , input) ,

$t_{1,\max}$  = min ( $t_1^+$ , input) ,

$\phi_{\min}$  = 0 ,

$\phi_{\max}$  =  $2\pi$  ,

$\cos\theta_{2,\min}$  = -1 ,

$\cos\theta_{2,\max}$  = +1 ,

$\phi_{2,\min}$  = 0 ,

$\phi_{2,\max}$  =  $2\pi$  .

$$t_1^\pm = -\frac{Q^2}{1+4\tau\frac{1-x}{x}} \left( 1 - r + \frac{2\tau}{x} \left( \frac{1-x}{x} - r \right) \pm \sqrt{B_t} \right) , \quad B_t = \frac{4\tau}{x^2} (\tau + x^2) \left( r^2 - 2rA_r + \left( \frac{1-x}{x} \right)^2 \right) .$$

## Corrections for VCS: Singularities from leptonic radiation

$$e(p_-) + p(p) \rightarrow e'(p'_-) + p'(p') + \gamma(k_1) + \gamma(k_2)$$

Propgator denominators of the diagrams for radiation

$$D_5 = 2p'_- \cdot k_1$$

$$D_6 = 2p'_- \cdot k_2$$

$$D_{56} = D_5 + D_6 + 2k_1 \cdot k_2$$

$$E_5 = -2p'_- \cdot k_1$$

$$E_6 = -2p'_- \cdot k_2$$

$$E_{56} = E_5 + E_6 + 2k_1 \cdot k_2$$

Partial fractioning

Subtractions to remove infrared and collinear divergences

e.g. following dipole subtraction scheme of S. Dittmaier

## Existing Monte Carlo programs

**HERACLES** and **DJANGOH**: for **inclusive DIS**

QCD-based event generation, valid at large  $Q^2$ : parton model

- Complete QED and electroweak corrections at  $O(\alpha)$  for NC and CC scattering, polarized lepton, polarized nucleon
- Interface to **LEPTO**, **JETSET**. Jets, parton showers, hadronic final state.  
**SOPHIA** for low-mass hadronic final states

Fortran code. Well tested, but less efficient for exclusive processes

**POLARES**: second-order corrections for elastic  $e p$

Includes  $O(\alpha)$  corrections for  $e p \rightarrow e p \gamma$

C++ code, to be published (with R. Bucoveanu)

# To Do

- Define objects to be implemented in the **PARTONS** framework
  - non-radiative and radiative processes, kinematics
- First step: **leading log approximation**
  - Use radiator functions
- Beyond the collinear approximation
  - Implement radiative leptonic tensor
    - Work out decomposition in terms of structure functions
- Radiative corrections for DVCS
  - Finish project with M. Hentschinski and M. Stratmann (?)
    - Use POLARES for cross-check