

Tau Neutrinos from GeV to EeV 2021 (NuTau2021)

Cross section measurement of ν_τ from GeV to TeV in Ar₄₀

Barbara Yaeggy
(9/30/2021)



Why cross sections are important?

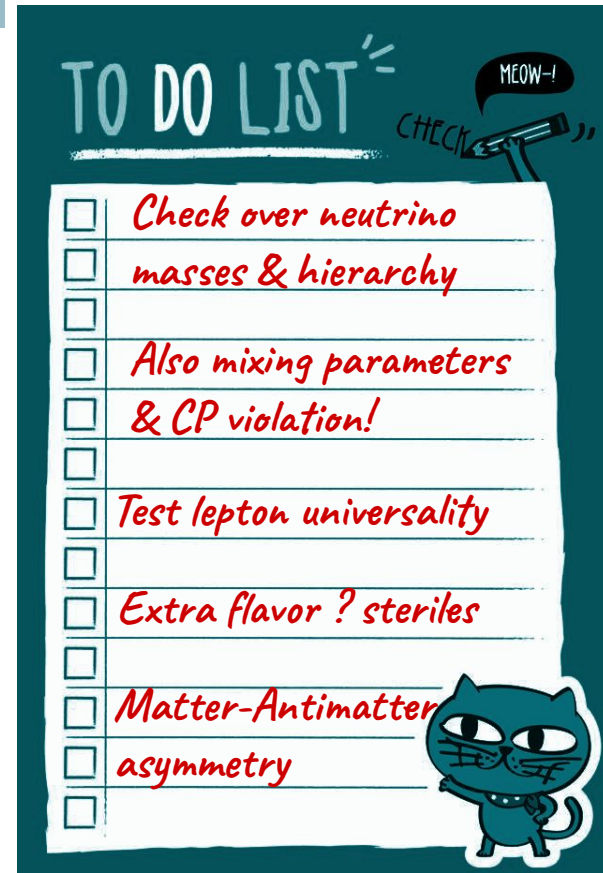
- **Neutrino interactions** (cross section) are the major contributor of systematic uncertainties in oscillation measurements (T2k, NOvA).
- E_ν & ν -nucleus interactions relies on **reconstruction techniques** either based on **kinematics** (T2K/HK) or **calorimetric methods** (DUNE/NOvA/SBN) and both requires reliable predictions from **interaction models**.
- Extraction of **oscillation parameter** is biased by the interaction model.

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta_{ij} \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4 E_\nu} \right)$$

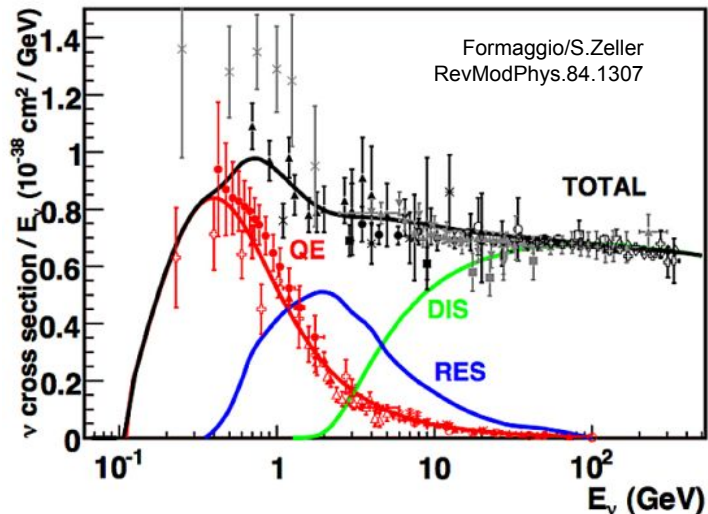
$$N^{\alpha \rightarrow \beta}(E_{\nu, \text{rec}}) \propto \sum_i \phi_\alpha(E_\nu) \times \sigma_\beta^i(E_\nu) \times P(\nu_\alpha \rightarrow \nu_\beta) \times \epsilon_\beta(E_\nu, E_{\nu, \text{rec}})$$

Nuclear and hadronic effects are energy dependent too!

Check talks from Laura Fields, Stephen Parke, Tyce deYoung, Julia Gehrlein



Neutrino Cross Sections at a Few GeV



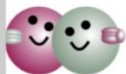
Nuclear models rely on different approximations, which are valid in specific kinematics and for specific process.

Can ν_τ give us access to a better understanding of some processes?

Short Range
Fermi Gas, Spectral Functions, Bodek-Ritchie rail to RFG, etc

Medium Range
Meson exchange current, transverse enhancement model

Long Range
Shell Model, RPA, EFT



by Cheryl Patrick

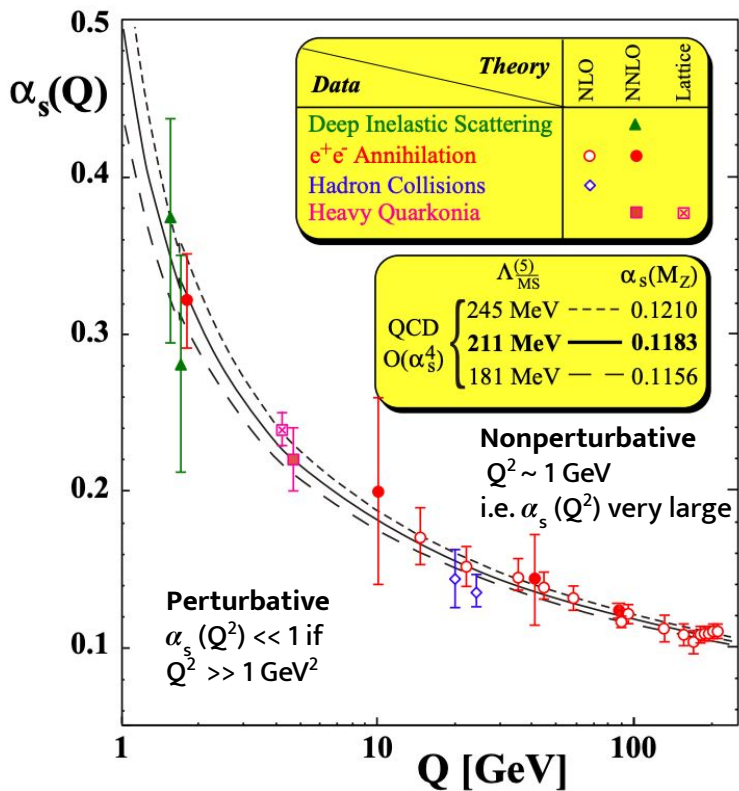
Nowadays, we are getting new accelerators and detectors!
→ precise interaction models and study phenomena.

Example: the oscillation of the ν_τ in long-baseline experiments with detectors with Ar40.

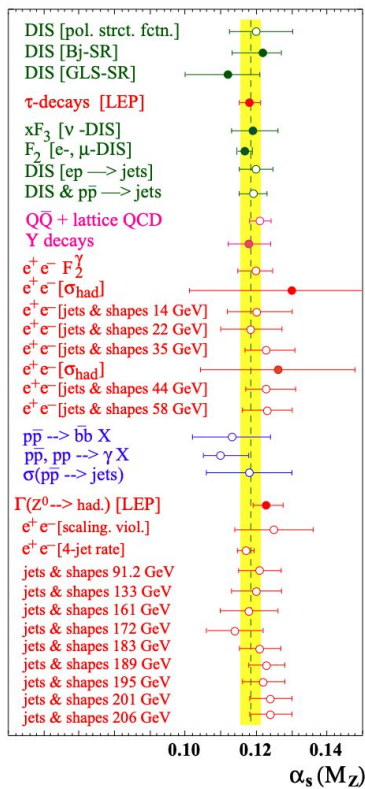
Given their resolution could provide higher statistics than previous experiments (OPERA/DONUT) but also:

More statistics → accurate insights for $\nu_\mu \rightarrow \nu_\tau$ oscillation and these results could be used as proof of the properties of neutrinos and their interactions.

Asymptotic freedom makes it possible to calculate the small distance interaction for quarks and gluons, assuming that they are free particles.



α_s 2002 by Siegfried Bethke (MPI of Physics, Munich, Germany)
[arXiv:hep-ex/0211012](https://arxiv.org/abs/hep-ex/0211012)



We can use **Deep ($Q^2 \gg M^2$) Inelastic ($W^2 \gg M^2$) Scattering** to probe the structure of hadrons.

The production of any particle can be determined by the cross section.

DIS experiments extract information from the lepton scattering cross sections to measure **Structure Functions** of the target, which are directly related to the nonperturbative Parton Distribution Functions, PDFs.

Bodek-Yang model aims for describing DIS cross section in all Q^2 regions
[arXiv:hep-ex/0308007](https://arxiv.org/abs/hep-ex/0308007)

Cross Section for NuTau: $\nu/\bar{\nu}_\tau(k) + N(p) \rightarrow \tau/\bar{\tau}(k')$

For the DIS the expression for the CC differential cross section is:

$$\frac{d^2\sigma^{\nu(\bar{\nu})}}{dxdy} = \frac{G_F^2 M E_\nu}{\pi(1 + Q^2/M_W^2)^2} \left((y^2 x + \frac{m_\tau^2 y}{2E_\nu M}) F_1 + \left[\left(1 - \frac{m_\tau^2}{4E_\nu^2}\right) - \left(1 + \frac{Mx}{2E_\nu}\right) \right] F_2 \right. \\ \left. \pm \left[xy \left(1 - \frac{y}{2}\right) - \frac{m_\tau^2 y}{4E_\nu M} \right] F_3 + \frac{m_\tau^2 (m_\tau^2 + Q^2)}{4E_\nu^2 M^2 x} F_4 - \frac{m_\tau^2}{E_\nu M} F_5 \right)$$

Suppressed for $\nu_{\mu,e}$

Albright and Jarlskog, in [Nucl. Phys. B 84, 467 \(1975\)](#), pointed out that there are two additional structure functions, F_4 and F_5 that contribute to the ν_τ XSec.

Structure Functions:

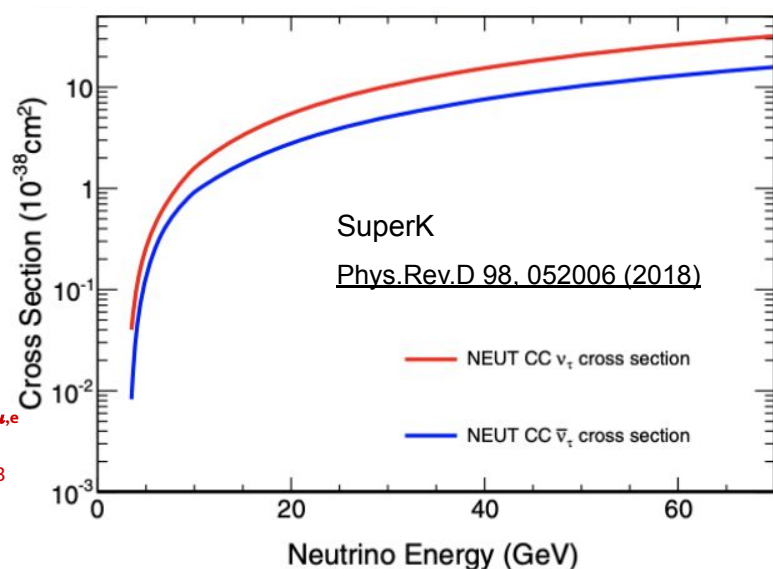
$2xF_1 = F_2$
 $-xF_3 = F_2$
 $xF_5 = F_2$
 $F_4 = 0$ also holds when the nucleon target is replaced by a lepton target.

The tau mass corrections to F_1 , F_2 and F_3 become negligible at high energies because the low-x rise of $q(x)$ is tempered by factors of x or y for these structure function.

$$W^2 = (p + q)^2 = M_N^2 + 2M_N\nu - Q^2,$$

$$Q^2 = -q^2 = -(k - k')^2,$$

$$x = Q^2/2p \cdot q = Q^2/2M_N\nu,$$



$$a = \frac{1 - m_\tau^2 \left(\frac{1}{2M_N E_\nu x} + \frac{1}{2E_\nu^2} \right)}{2 \left(1 + \frac{M_N x}{2E_\nu} \right)}$$

$$\frac{m_\tau^2}{2M_N(E_\nu - m_\tau)} \leq x \leq 1$$

$$b = \frac{\sqrt{\left(1 - \frac{m_\tau^2}{2M_N E_\nu x} \right)^2 - \frac{m_\tau^2}{E_\nu^2}}}{2 \left(1 + \frac{M_N x}{2E_\nu} \right)}$$

$$a - b \leq y \leq a + b$$



$$\frac{d^2\sigma^{\nu(\bar{\nu})}}{dx dy} = \frac{G_F^2 M_N E_\nu}{\pi(1 + Q^2/M_W^2)^2} \left\{ \left(y^2 x + \frac{m_\tau^2 y}{2E_\nu M_N} \right) F_1^{W^\pm} \right. \\ + \left[\left(1 - \frac{m_\tau^2}{4E_\nu^2} \right) - \left(1 + \frac{M_N x}{2E_\nu} \right) y \right] F_2^{W^\pm} \\ \pm \left[xy \left(1 - \frac{y}{2} \right) - \frac{m_\tau^2 y}{4E_\nu M_N} \right] F_3^{W^\pm} \\ \left. + \frac{m_\tau^2 (m_\tau^2 + Q^2)}{4E_\nu^2 M_N^2 x} F_4^{W^\pm} - \frac{m_\tau^2}{E_\nu M_N} F_5^{W^\pm} \right\}.$$

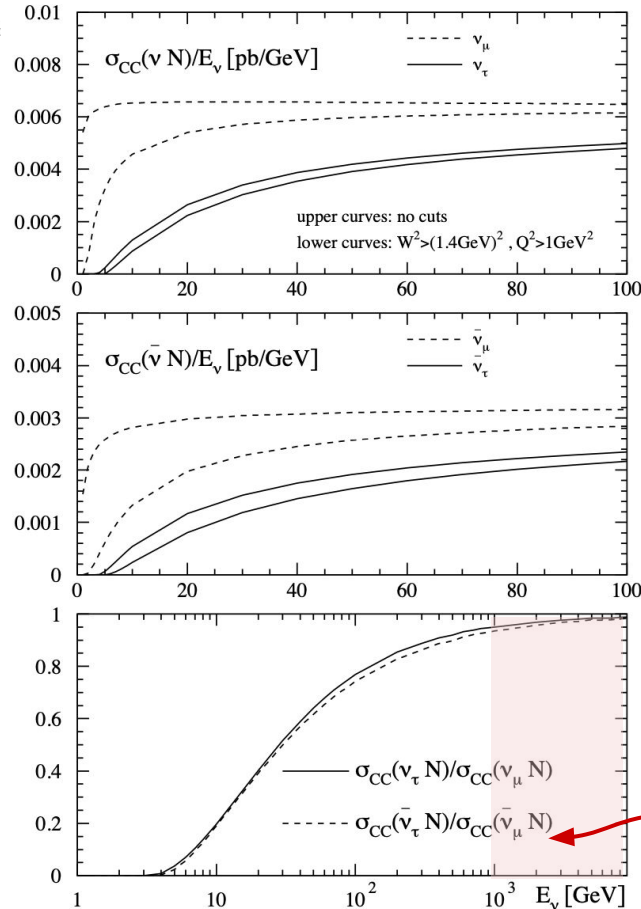
Reasons of the deficit in CC ν_τ :

- Reduce phase space: integration limits (x,y).
- F_5 contribution: independent of x and minus sign.

Also:

$$F_5 \sim F_1 \sim q(x, Q^2)$$

There is a small-x enhancement of its contribution to the cross section at high energies.



Cuts

With cuts, W_{\min} and Q^2 cuts agree to within 3% for $E_\nu > 20$ GeV

Ratios

The ratios ν_τ / ν_μ & $\bar{\nu}_\tau / \bar{\nu}_\mu$ are insensitive to DIS cuts

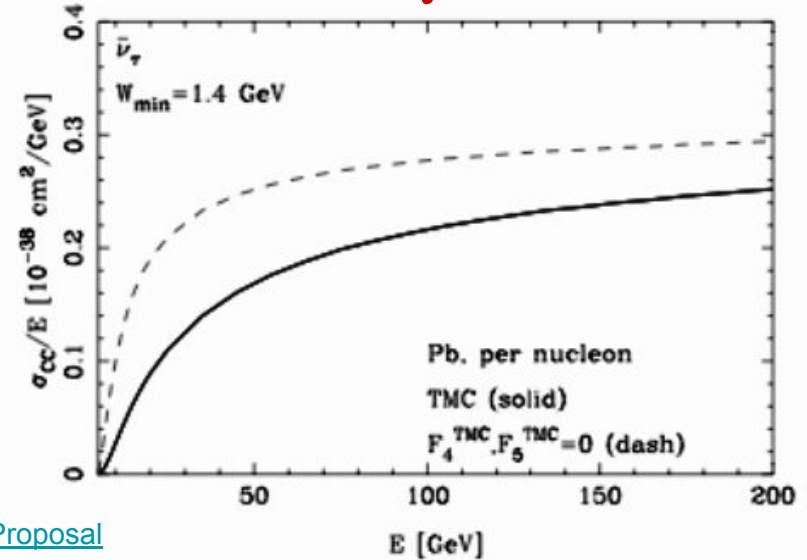
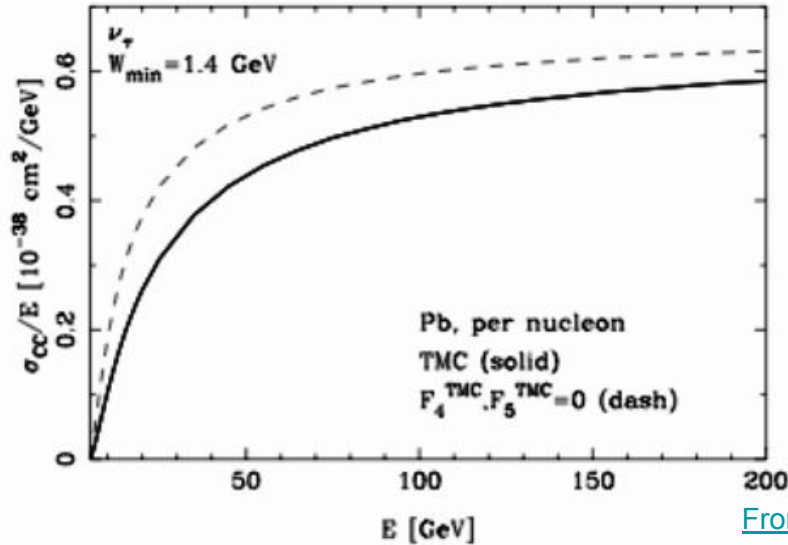
There is a deficit for CC- ν_τ Even at 10^3 GeV the ratio is below 1.

Notice the difference between the cross-sections in the $F_4 = F_5 = 0$ hypothesis and the SM prediction is larger for lower neutrino energies.

ν_τ on Pb

SM prediction (solid)
 $F_4 = F_5 = 0$ hypothesis (dashed)

$\bar{\nu}_\tau$ on Pb

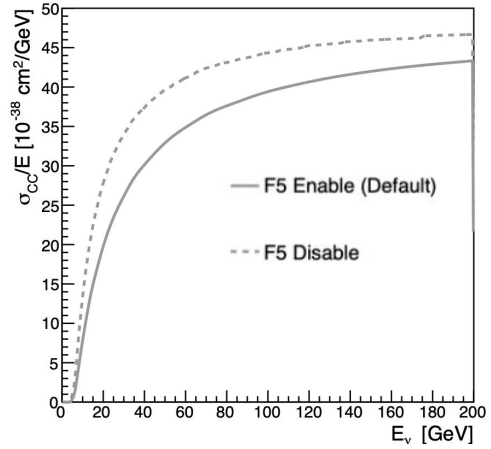


[From the SHIP Proposal](#)

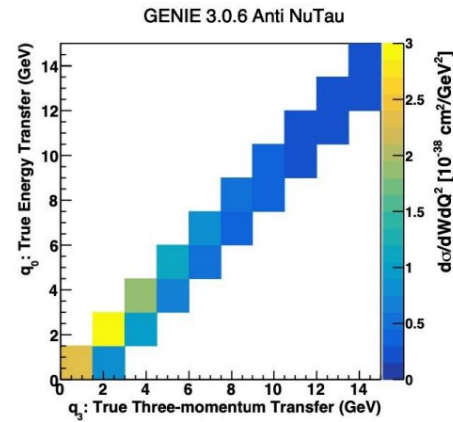
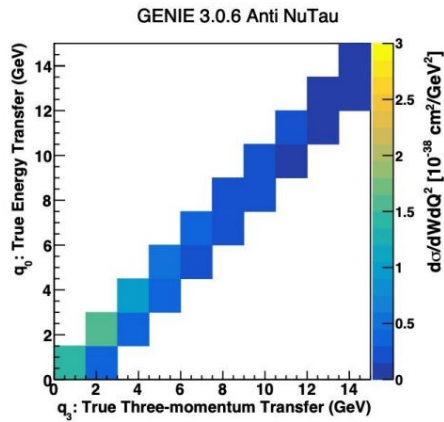
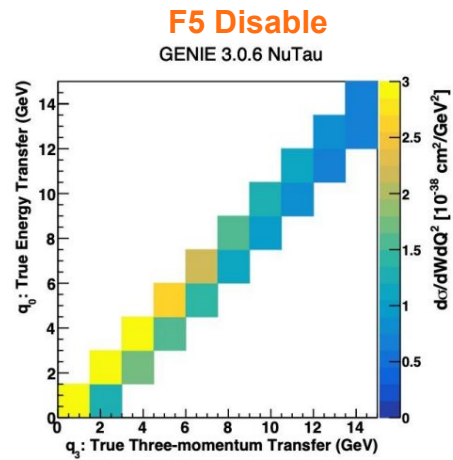
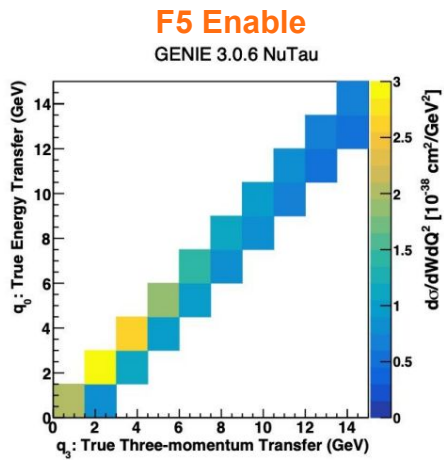
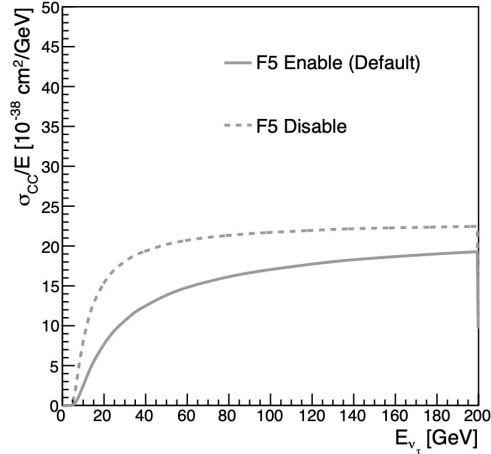
In short, we want to check ν_τ cross-sections in the DIS regime and due its relevance, we choose to conduct those studies in Ar_{40} .

CC - ν_τ TRUTH Level studies show that indeed, when DIS cuts are applied and $F_5 = 0$ we can extract new information from the lepton cross section.

**GENIE 3.0.6
CC-NuTau Cross
Section**

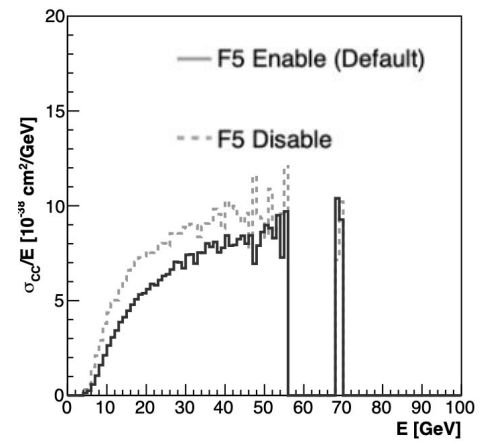


**GENIE 3.0.6
CC- Anti NuTau
Cross-Section**

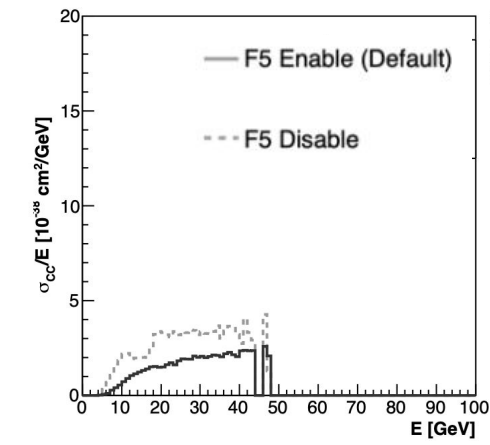


CC - ν_τ TRUTH Level studies show that indeed, when DIS cuts are applied and $F_5 = 0$ we can extract new information from the lepton cross section.

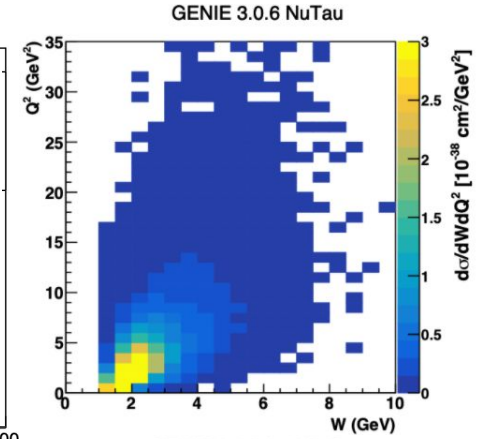
GENIE 3.0.6 NuTau



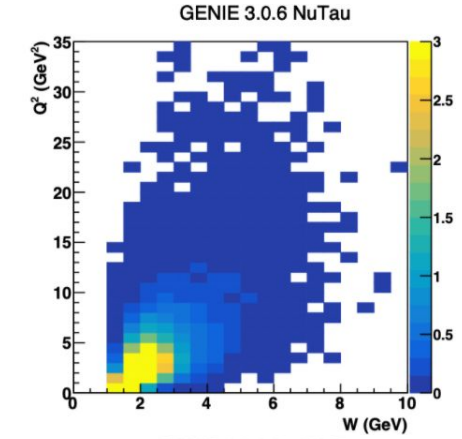
GENIE 3.0.6 Anti NuTau



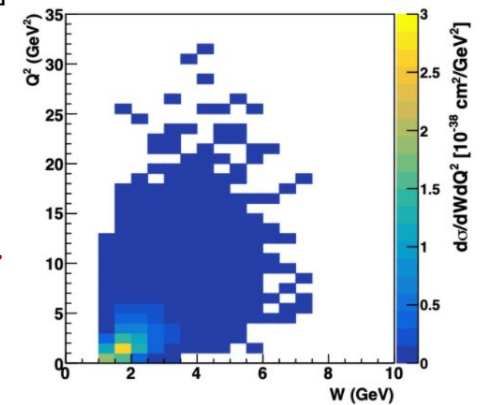
F5 Enable



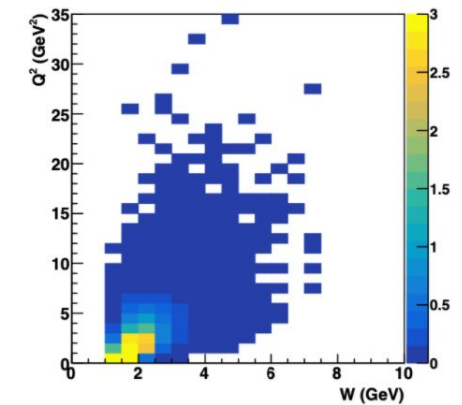
F5 Disable



GENIE 3.0.6 Anti NuTau

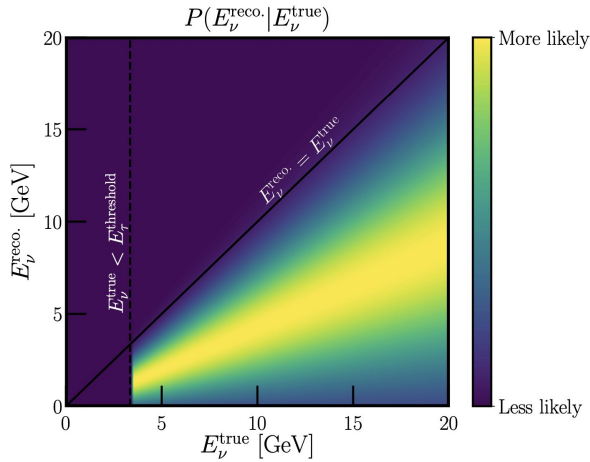


GENIE 3.0.6 Anti NuTau



*Can we reconstruct kinematic variables?
Can we have a minimum resolution to get a cross section?*

Tau Neutrino Interactions



A.Gouvea, K. Kelly, G.Stenico, P.Pasquini
[PhysRevD.100.016004](https://arxiv.org/abs/1001.01604)

P.Machado, J.Turner, H.r Schulz
[arXiv:2007.00015](https://arxiv.org/abs/2007.00015)

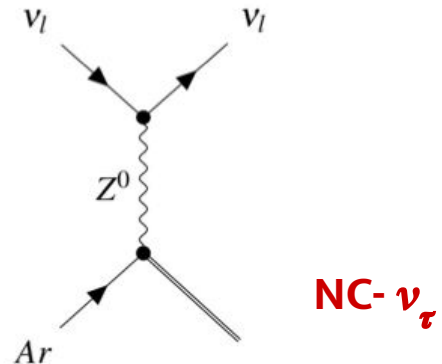
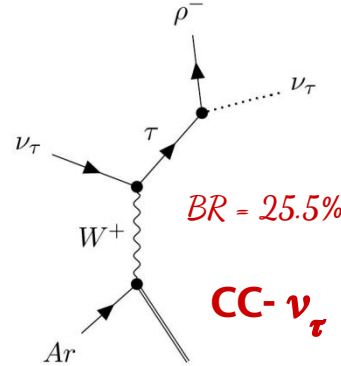
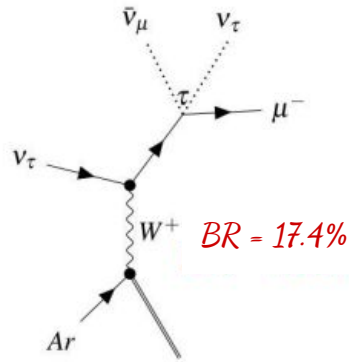
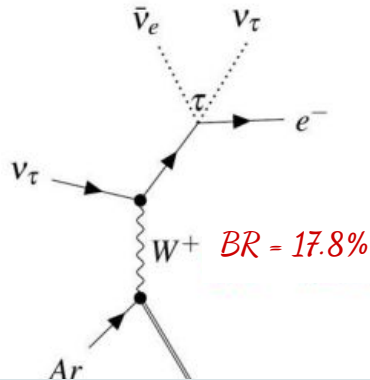
Decay mode	Branching ratio
Leptonic	35.2%
$e^- \bar{\nu}_e \nu_\tau$	17.8%
$\mu^- \bar{\nu}_\mu \nu_\tau$	17.4%
Hadronic	64.8%
$\pi^- \pi^0 \nu_\tau$	25.5%
$\pi^- \nu_\tau$	10.8%
$\pi^- \pi^0 \pi^0 \nu_\tau$	9.3%
$\pi^- \pi^- \pi^+ \nu_\tau$	9.0%
$\pi^- \pi^- \pi^+ \pi^0 \nu_\tau$	4.5%
other	5.7%

τ is heavy, $\sim 1.777 \text{ GeV}$
 τ energy threshold $\sim 3.5 \text{ GeV}$

Low statistics

$\tau_{\text{life}} \sim 2.9 \times 10^{-13} \text{ sec}$

Challenge: ν_τ reconstruction and the background rejection from NC.



Rho Channel: $\tau \rightarrow \nu_\tau \rho^- \rightarrow \pi^- \pi^0$

Background & Signal Separation

Based on the BDT work by **Miriana Rajaoalisoa**, see also **Thomas Kosc** talk

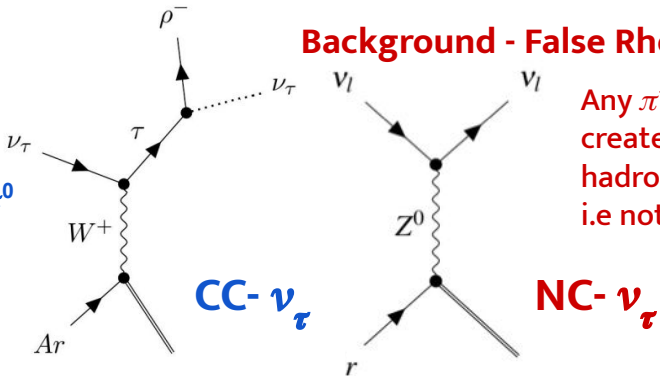
- Select Rho candidates
- Use a **BDT (Boosted Decision Tree)** classifier to separate the NC- ν_τ from CC- ν_τ
- BDT trained on transverse kinematic variables.
- BDT (Pre-selection): Check for rho decay products identification.

True Rho:

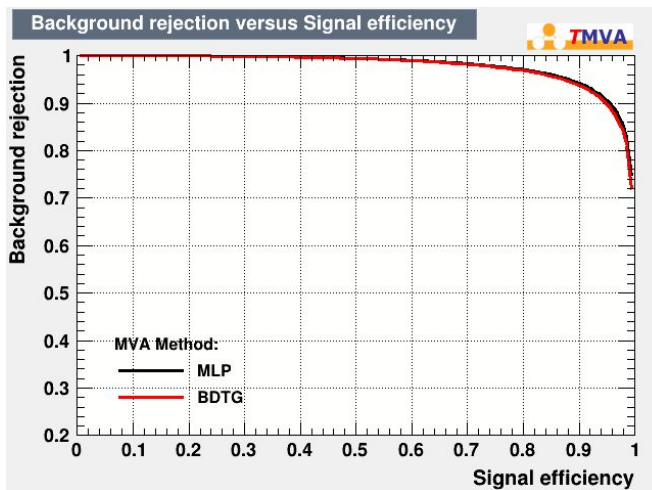
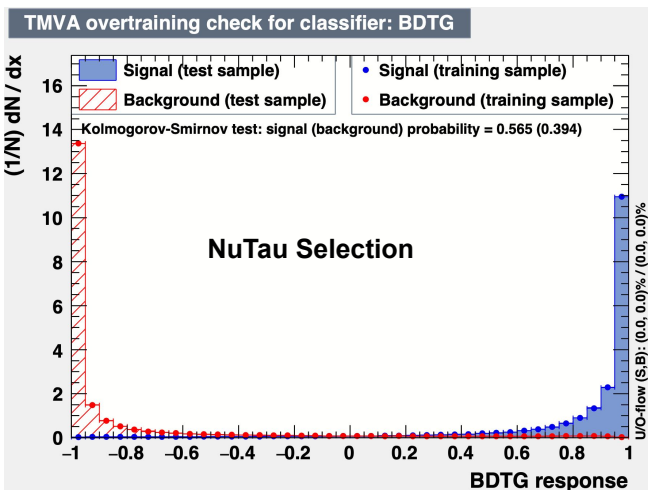
Signal CC: $\tau \rightarrow \nu_\tau \rho^-$

Signal NC: $\rho^- \rightarrow \pi^- \pi^0$

BR = 25.5%



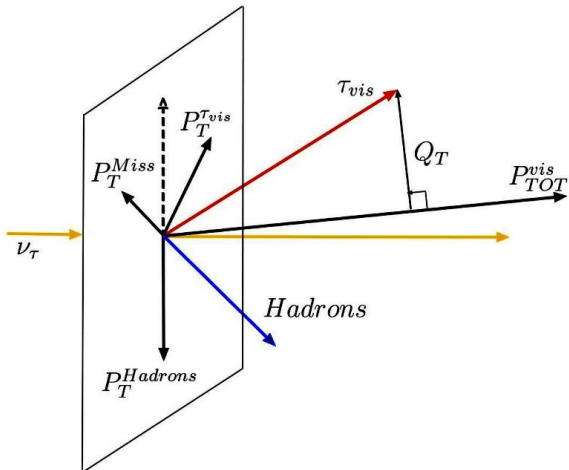
Any $\pi^- \pi^0$ couple created out of the hadronic τ^- decay, i.e. not a True Rho.



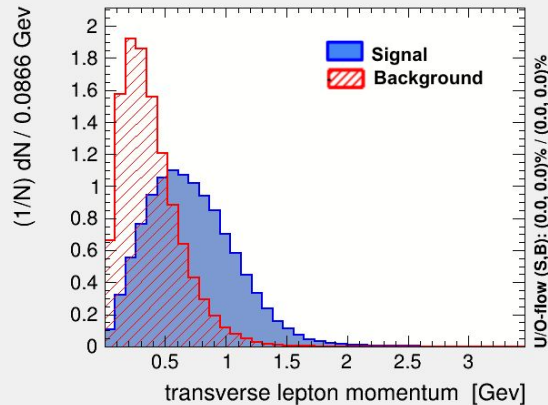
MVA: Multivariate analysis

MLP: Multi Layered Perceptrons

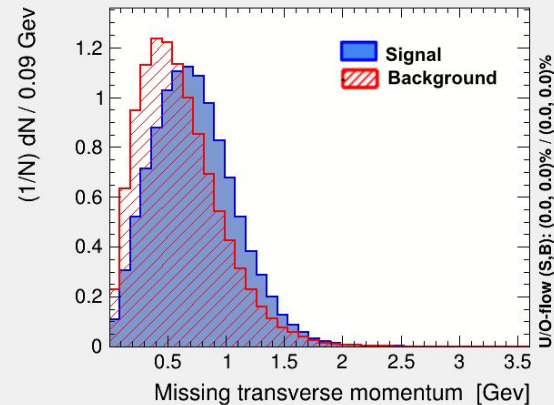
Transverse Kinematic Variables - BDT trained with non-smearred



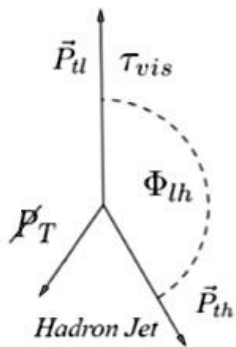
Input variable: transverse lepton momentum



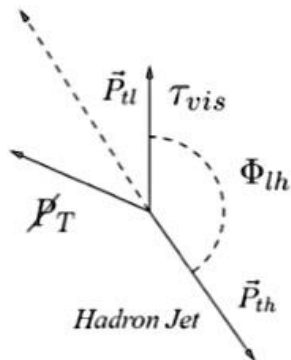
Input variable: Missing transverse momentum



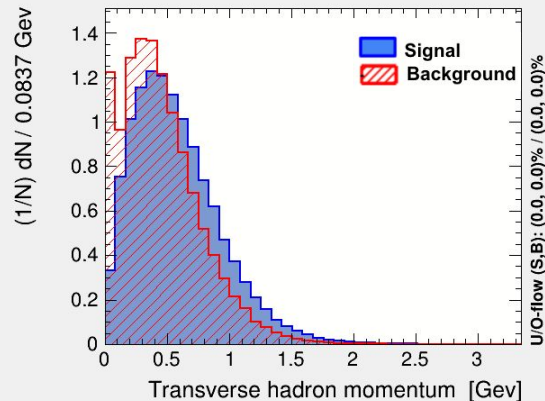
Background
Interaction
Products



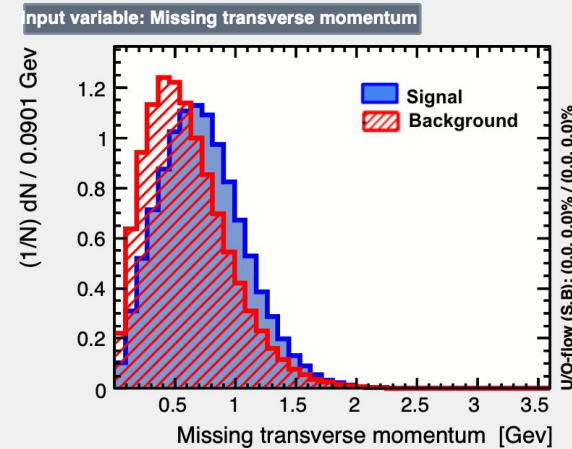
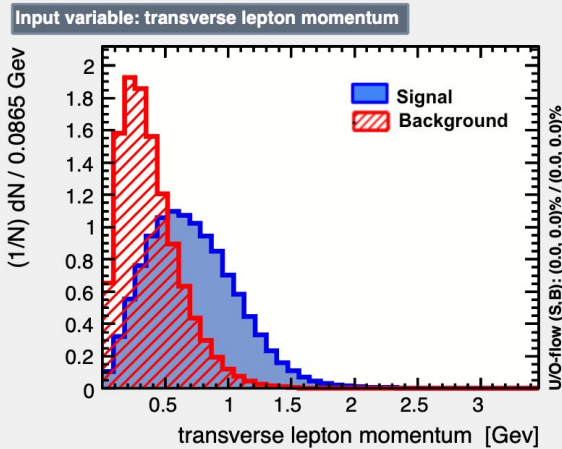
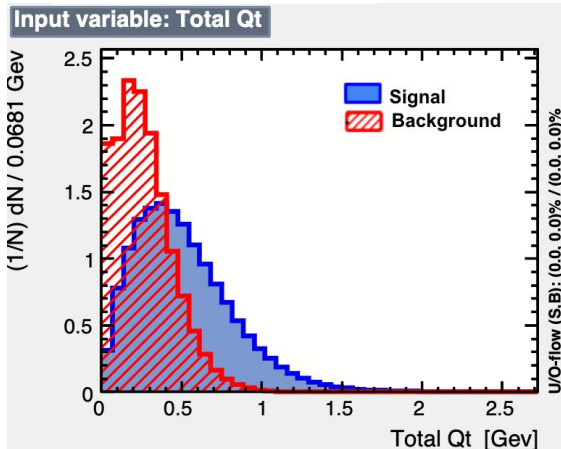
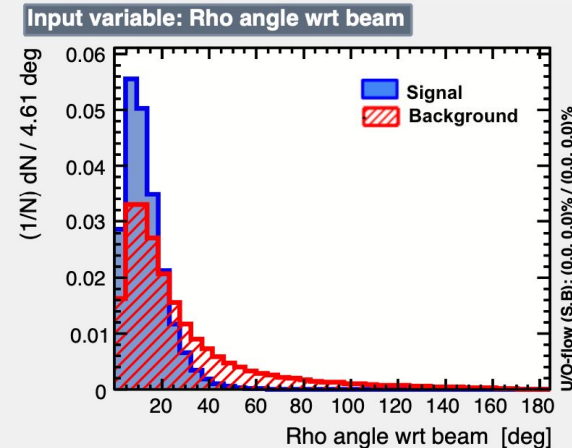
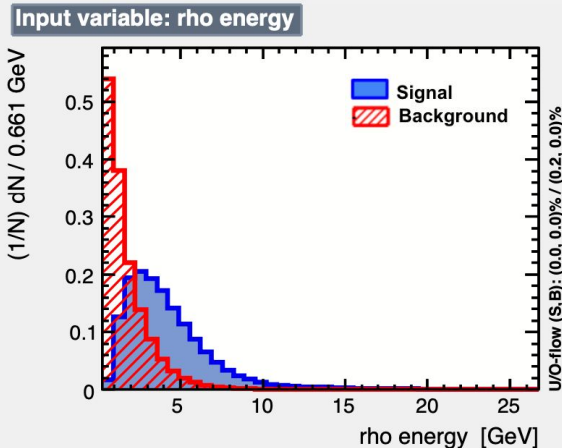
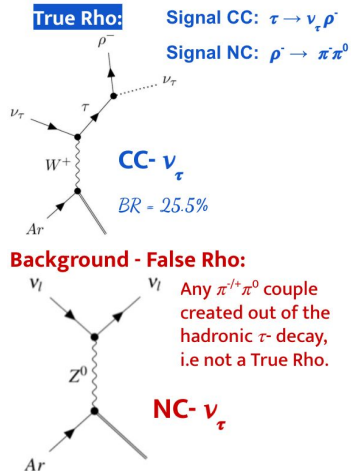
NuTau
Interaction
Products



Input variable: Transverse hadron momentum



Rho Selection - BDT trained with non-smearred



COMMENTS

In general:

- NuTau physics is challenging but everything that we could learn from it would be productive & worthwhile.

In specific:

- Signal and background separation over the transverse kinematic variables is acceptable.
- There is room to test other reconstruction tools different than the BDT score.
- **Coming soon: kinematic variables reconstruction.**

Thank you!



Milky Way and Volcan de Fuego , Guatemala

BACKUP

▶ Leptonic channels

- Missing transverse momentum P_T^{miss} .
- Angle between the transverse hadronic system momentum and transverse lepton (tau visible product) momentum φ_{lh} .
- Angle between the missing transverse momentum and transverse lepton (tau visible product) momentum φ_{ml} .
- Lepton transverse momentum P_T^l .
- Hadronic system transverse momentum P_T^h .
- Lepton (tau visible product) energy E_l .
- Total visible energy E_{Total}^{vis} .
- Transverse mass M_T .
- Angle between the neutrino direction and missing transverse momentum $\theta_{\nu m}$.
- Angle between the neutrino direction and the hadronic system $\theta_{\nu h}$.
- Angle between the neutrino direction and the lepton (tau visible product) $\theta_{\nu l}$.
- Ratio between the hadronic system and the total visible energy.
- Component of the lepton (tau visible product) momentum perpendicular to the total visible momentum.
- Ratio of the transverse lepton (tau visible product) momentum and the missing transverse momentum.

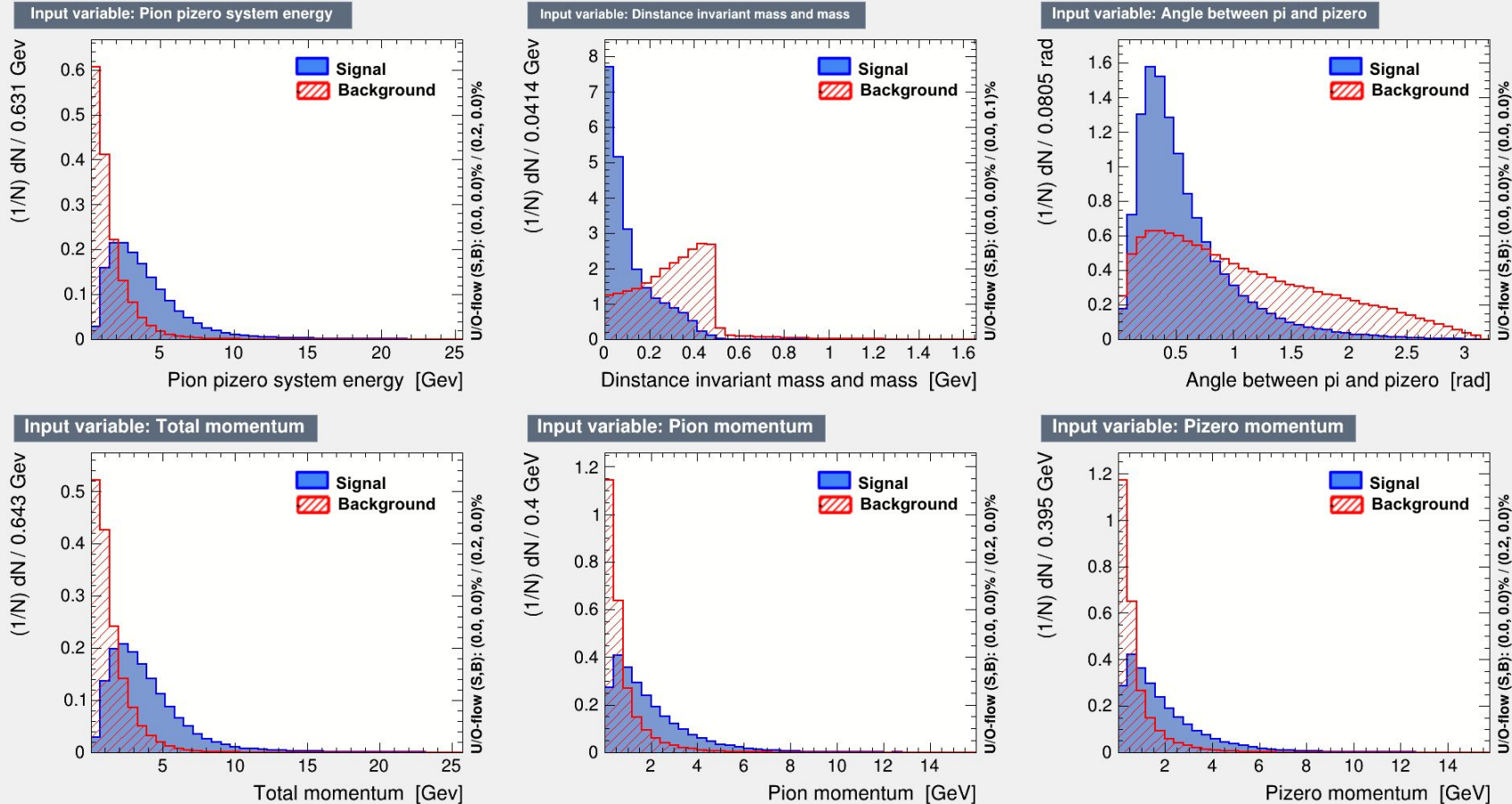
► Rho classification

- Energy of the π^0/π^\pm couple.
- Energy of the π^\pm .
- Invariant mass of the pion.
- Invariant mass of the π^0 .
- Pion energy sharing $r_\pi^K = \frac{E_\pi^K}{E_\pi^K + E_{\pi^0}^K}$.
- Distance invariant mass and mass $d = \sqrt{(M_{\pi^0}^{(inv)} - m_{\pi^0})^2 + (M_\rho^{(inv)} - m_\rho)^2}$.
- Angle between the π^\pm and π^0 .
- π^\pm momentum.
- π^0 momentum.

► Rho channel vs NC classification

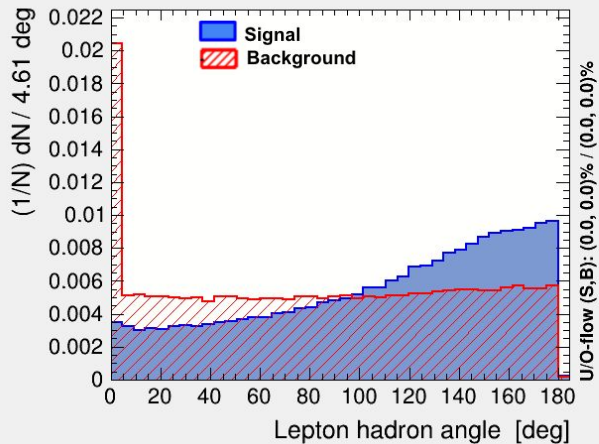
- All previous variables including those used for the leptonic channels.
- Minimum angle between the tau visible daughter and the hadronic system.
- Ratio between the transverse size of the hadronic system and the full event.

Rho Selection - BDT trained with non-smearred

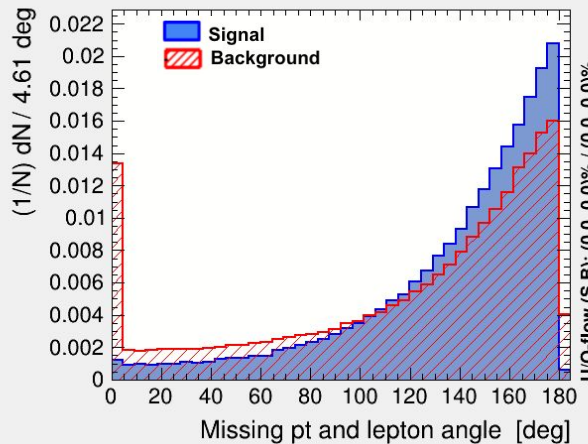


Transverse Kinematic Variables - BDT trained with non-smearred

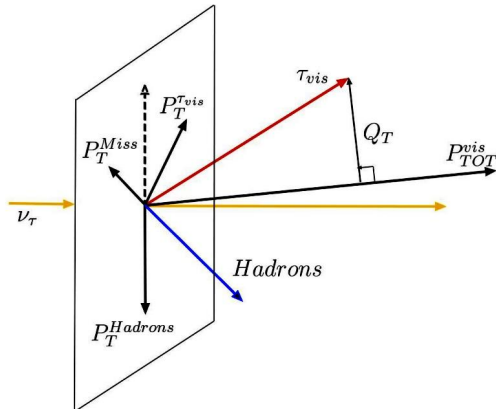
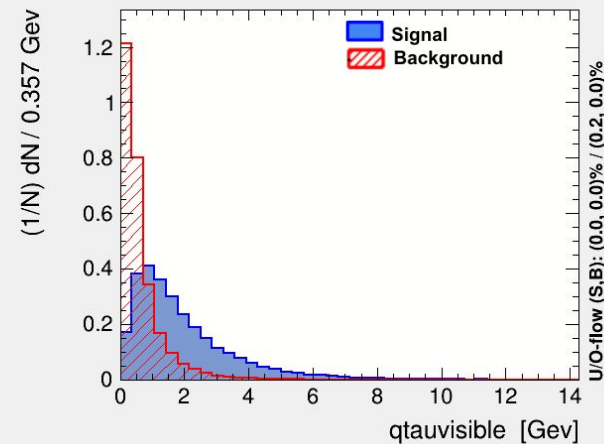
Input variable: Lepton hadron angle



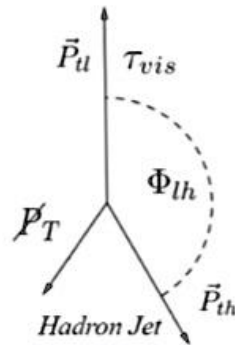
Input variable: Missing pt and lepton angle



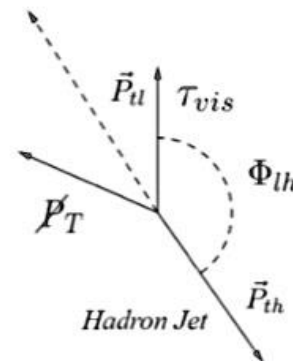
Input variable: qtauvisible



**Signal
interaction
products**



**Background
interaction
products**



Transverse Kinematic Variables - BDT trained with non-smear

