

# Exotic Spectroscopy (in Photoproduction) @ EIC

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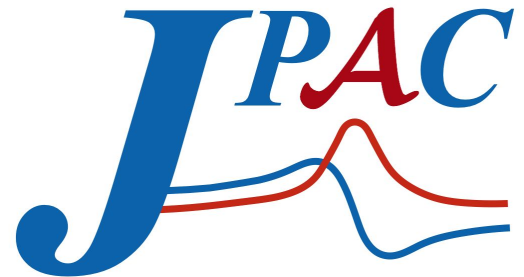
**[arXiv:2008.01001]**

**IR2 @ EIC**

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**INDIANA UNIVERSITY**  
BLOOMINGTON



# Exotic Hadrons

Plethora of quarkonium-like states observed since 2003 which do not fit into conventional  $q\bar{q}$  models.

**X(3872)** - large isospin violation

**Y(4260)** - anomalous coupling to open charm channels

**Z(3900)** - charged, charmonium-like state

Ambiguous interpretation of signals:

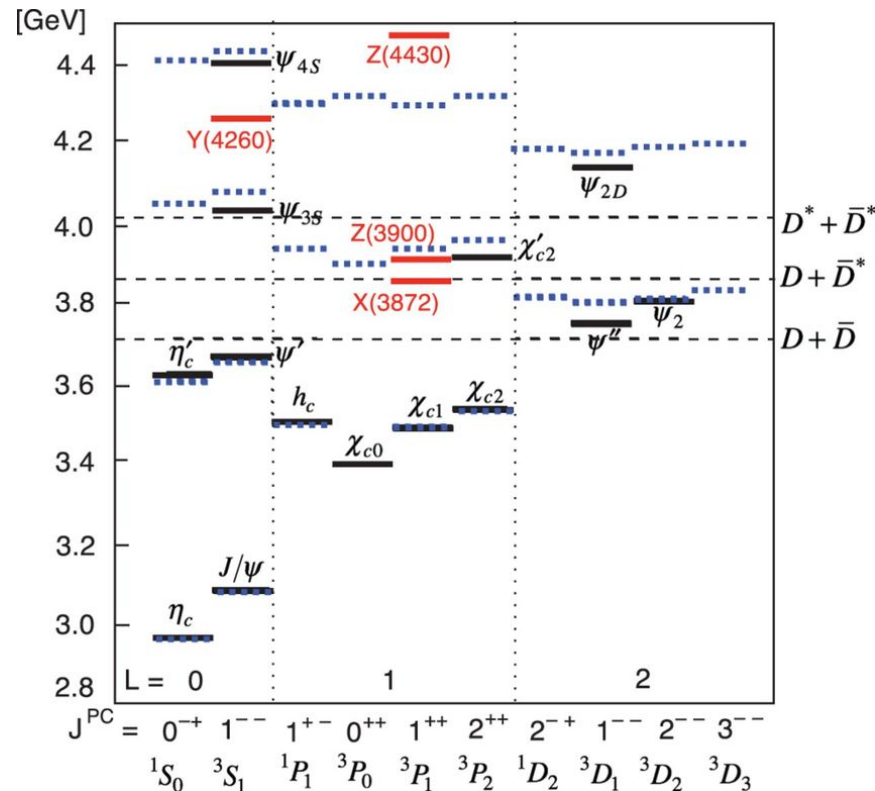
Multi-quark resonances, hadronic molecules, hadrocharmonia, **kinematic effect**, hybrid

For reviews of XYZs see e.g.:

A. Hosaka et al. [arXiv:1603.09229]

N. Brambilla et al. [arXiv:1907.07583]

F-K. Guo et al. [arXiv:1912.07030]

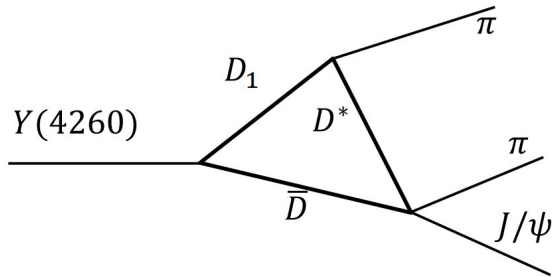


# Case of the $Z_c(3900)^+$

Charmonium-like state seen in  $\pi J/\psi$  and  $D^*D^*\bar{\text{bar}}$  and various production modes.

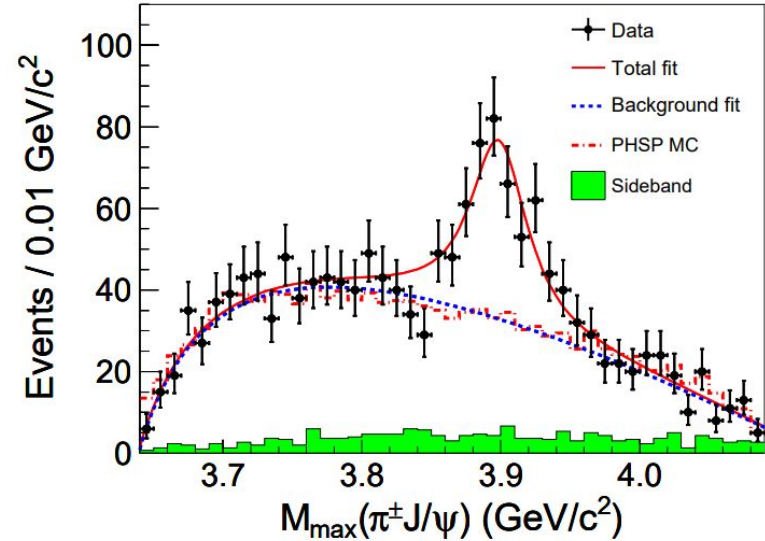
*but always with extra final state particles*

Charged nature unambiguously points to exotic nature but precise structure still unknown.



Guo, Liu, Sakai [arXiv:1912.07030]  
Szczepaniak [arXiv:1501.01691]

BESIII [Phys. Rev. Lett. 110, 252001 (2013)]



Triangle rescattering with final state pions unlikely to explain away Z-like signal but must still be accounted for in amplitude analysis and parameter extraction.

# Exclusive photoproduction

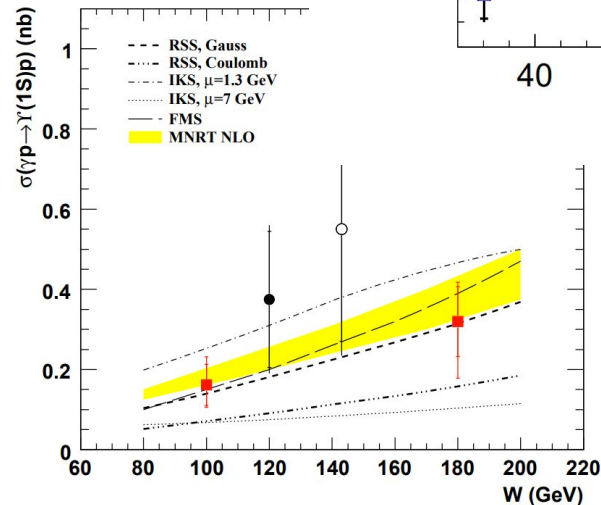
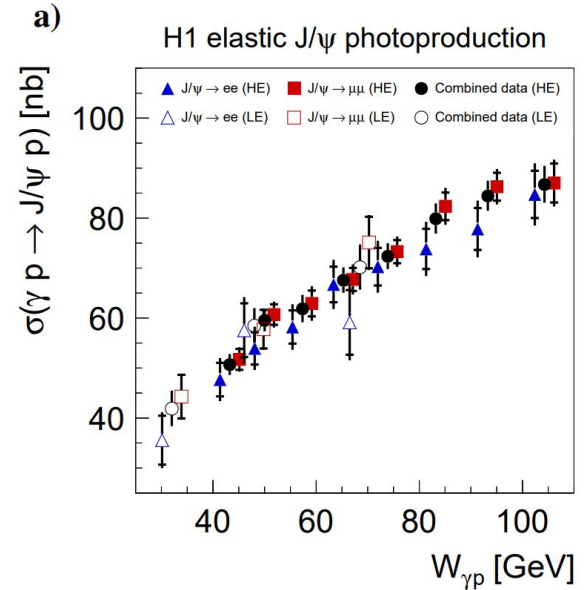
None of the XYZ's have been observed in photon-induced reactions

Current dedicated photoproduction facilities (e.g. GlueX@JLab) too low in energy reach.

Ideal laboratory for spectroscopy

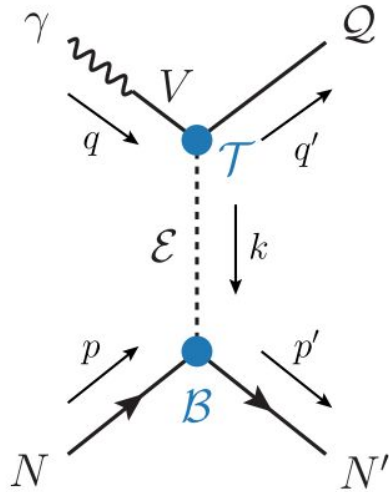
- Constrained kinematics  $\Rightarrow$  precise determination of production mechanism
- Direct production  $\Rightarrow$  eliminates contribution from triangle rescattering and FSI
- Phenomenology well understood
- Heavy quarkonium photoproduction studies at  $ep$  colliders demonstrated at HERA

**Higher luminosity and energy at EIC  
allow study of exotic hadrons!**



[arXiv:0903.4205]

# Exclusive photoproduction



## Recipe for an amplitude:

1. Identify relevant exchanges
2. Photon couplings fixed by observed decay widths and VMD
3. Bottom couplings from other reactions

## Necessary distinction between production near-threshold and high-energy

- In near threshold production lowest spin exchanges dominate
- At high-energies Regge physics kicks in

**JPAC [arXiv:2008.01001]**

See also e.g.

Galata [arXiv:1102.2070]

Lin, Liu, & Xu [ arXiv:1308.6345]

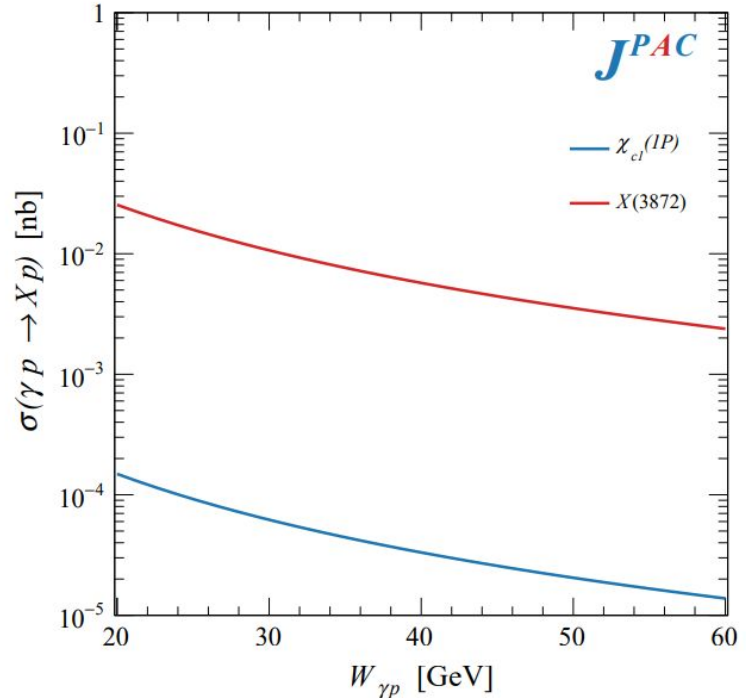
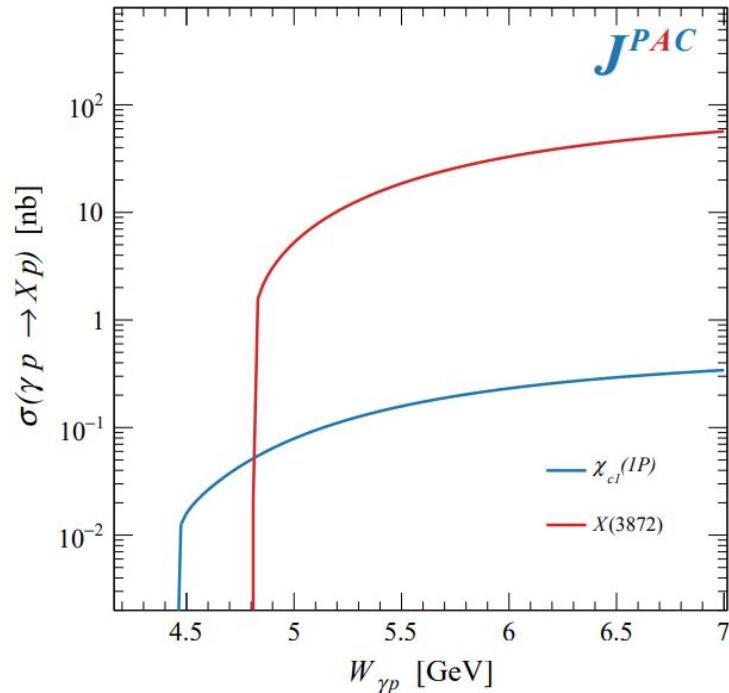
Wang et al. [arXiv:2009.05789]

Production cross-sections by meson exchanges fall at high energies as a consequence of unitarity (Reggeization).

**Lower  $\gamma p$  invariant mass  $\Rightarrow$  higher cross-section for X and Z**

# X(3872)

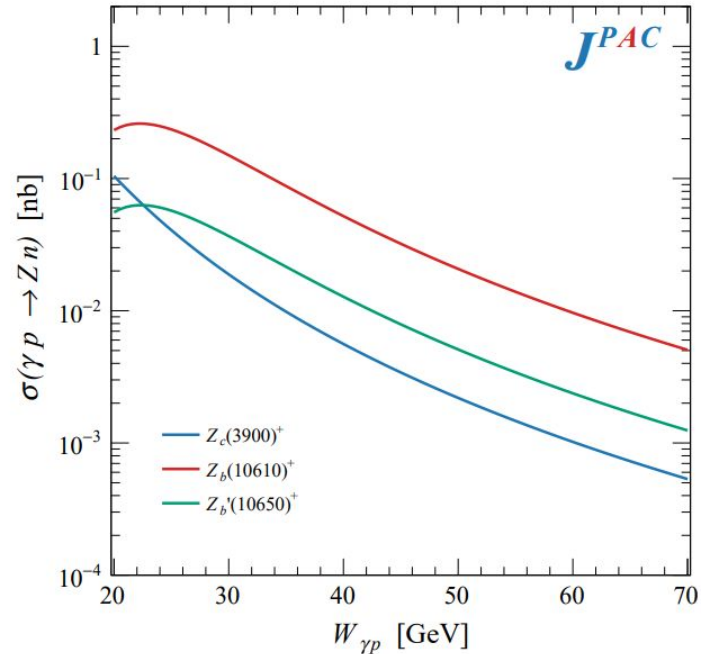
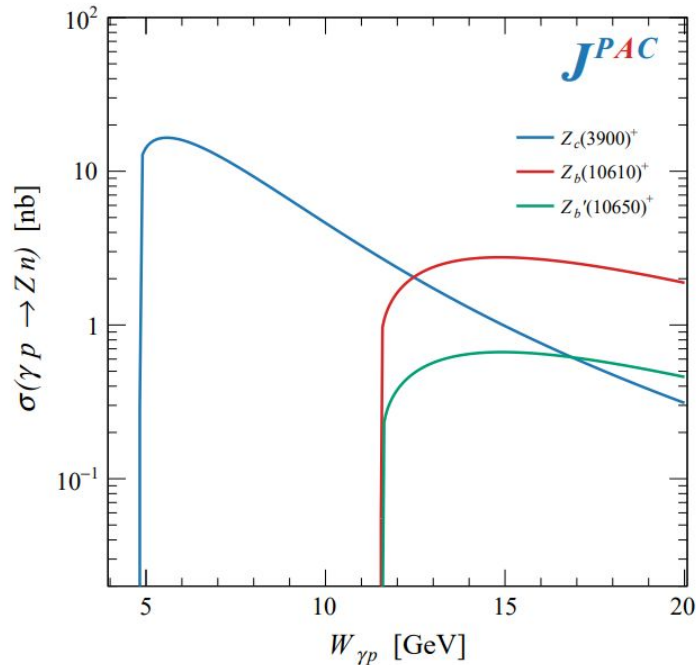
Effect very dramatic for vector exchanges, X(3872) production  
(note: interpolation between high and low energy not really straightforward)



# $Z_c(3900)^+$

Spin-0, pion exchange  $\Rightarrow$  cross-sections for Z states fall less dramatically from threshold production.

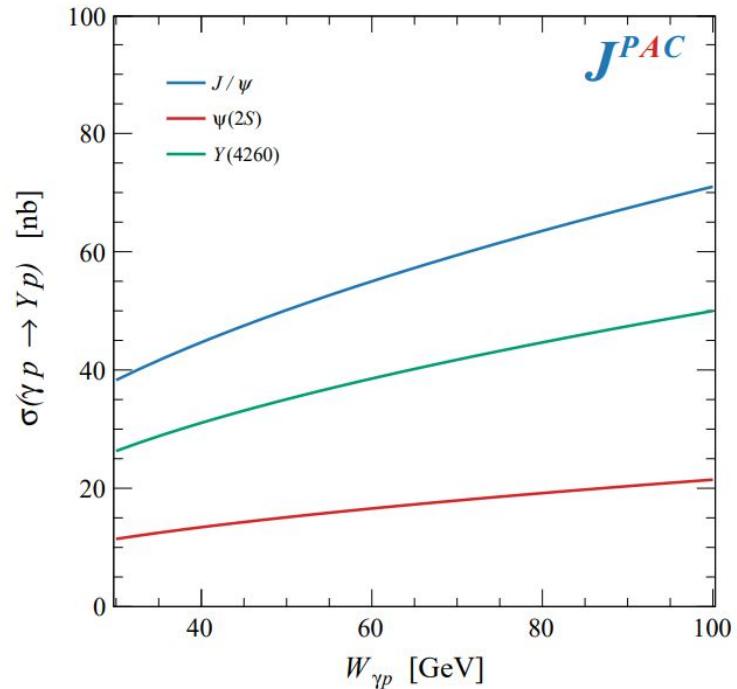
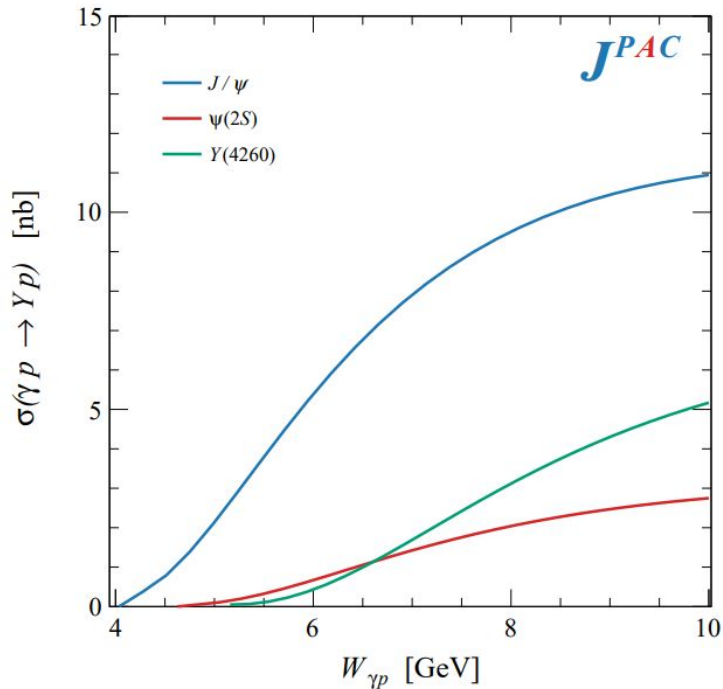
***Spectroscopy of Z's possible with high-luminosity at low CM energies!***



# Vector charmonia and Y(4260)

Unlike X and Z, vector states produced via diffractive Pomeron exchange.

**Diffractive cross-section rise with energy, Y states instead benefit from high CM.**





# jpgacPhoto

Code implementations of all amplitudes mentioned here (and more) freely available for download and use at [github.com/dwinney/jpacPhoto](https://github.com/dwinney/jpacPhoto)

- Probability distribution (  $\Sigma_{\lambda} |A|^2$  )
- Differential cross section (  $d\sigma / dt$  )
- Integrated total cross section (  $\sigma$  )
- Polarization asymmetries (  $A_{LL}$  and  $K_{LL}$  )
- Spin density matrix elements (  $\rho^{\alpha_{\lambda}, \lambda'}$  )
- Integrated beam asymmetry (  $\Sigma_{4\pi}$  )
- Beam asymmetry in the y-direction (  $\Sigma_y$  )
- Parity asymmetry (  $P_{\sigma}$  )

Available amplitudes, so far, include:

- Baryon resonance (s-channel)
- Pomeron exchange (t-channel)
- (fixed-spin and reggeized) Charged pseudo-scalar meson exchange (t-channel)
- (fixed-spin and reggeized) Vector meson exchange (t-channel)
- Primakoff effect off nuclear target (t-channel)
- (fixed-spin) Dirac fermion exchange (u-channel)
- (fixed-spin) Rarita-Schwinger fermion exchange (u-channel)

Can incorporate production and subsequent decays in MC generators, **see Derek's talk up next!**

# Summary

The EIC offers a unique opportunity to study exotics in photoproduction to complement measurements at existing facilities.

In particular production of  $Z_c$  and  $Z_b$  states hugely benefits from a high-luminosity, low energy program (e.g. at IR2). Production cross-sections comparable to existing  $e^+e^-$  machines.

**Thank you!**

# Backup Slides

# Propagators

At low energies (near threshold) we expect the partial wave sum of the full amplitude to be  $\propto (p(s)q(s))^j$  such that only the lowest  $j$  contributes. Thus, we consider fixed-spin, *Feynman propagators* which contain full energy dependence at low energies.

Easily written in terms of Feynman rules:

$$\mathcal{P}^0 = \frac{1}{t - m_{\mathcal{E}}^2}, \quad \mathcal{P}_{\alpha,\beta}^1 = \frac{g_{\alpha\beta} - k_{\alpha}k_{\beta}/m_{\mathcal{E}}^2}{t - m_{\mathcal{E}}^2}$$

Equivalently contracting all Lorentz structures evaluated in the  $t$ -channel CM frame we may match this to a helicity amplitude proportional to

$$\frac{d_{\mu\mu'}^j(\theta_t)}{t - m_{\mathcal{E}}^2} = \frac{\text{polynomial of order } j \text{ in } s}{t - m_{\mathcal{E}}^2}$$

# Reggeization

At high energies, the above will like  $s^j$  which for  $j \geq 1$  exceeds unitarity bounds. Therefore we restrict fixed-spin exchanges heuristically to a few GeV above threshold.

Beyond that, we must consider the re-summed (Reggeized) tower of arbitrary spin with the replacement

$$\mathcal{N}_{\mu\mu'}^j \left( \frac{4p(t)q(t)}{s_0} \right)^{j-M} \frac{d_{\mu\mu'}^j(\theta_t)}{\xi_{\mu\mu'}^{(t)}(s,t)} \frac{1}{t - m_{\mathcal{E}}^2} \rightarrow -\alpha' \Gamma(j - \alpha(t)) \left[ \frac{1 + \tau e^{-i\pi\alpha(t)}}{2} \right] \left( \frac{s}{s_0} \right)^{\alpha(t)-M}$$

We use usual real, linear parameterizations for the  $\rho - \omega$  and  $\pi$  Regge trajectories:

$$\alpha_{\rho}(t) = 1 + 0.9(t - m_{\rho}^2) \quad \alpha_{\pi}(t) = 0.7(t - m_{\pi}^2) .$$

Note an intercept  $\alpha_0 = \alpha(t=0) < 1$  forces the Reggeized amplitude to decrease at high energies.