

EIC Physics (mostly eA) that would benefit from lower energies and/or a second detector

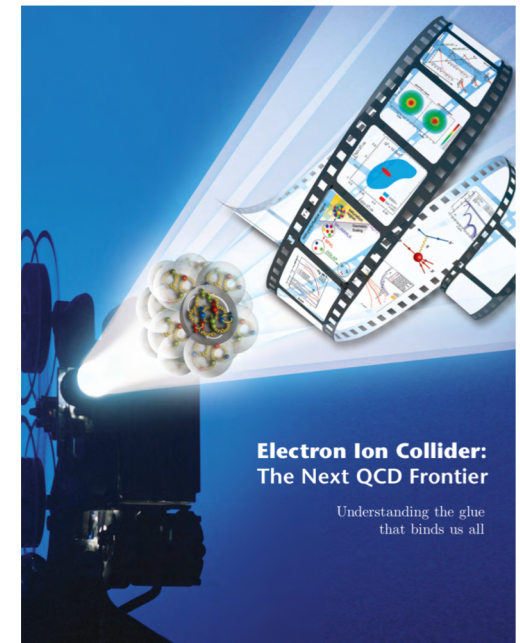
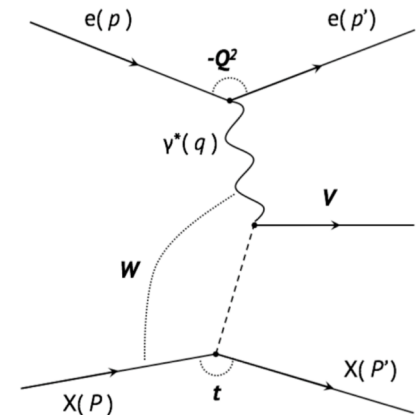
Spencer Klein, LBNL

Presented at the IR2@EIC Workshop

March 17-19, 2020

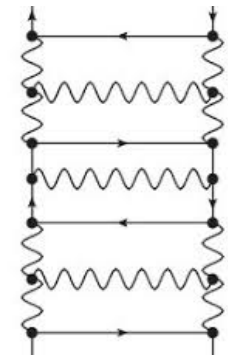
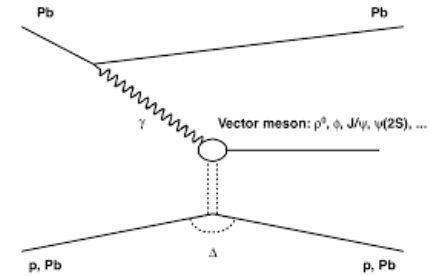
- Photo & electroproduction at an EIC
- Bjorken-x and rapidity
 - ◆ Near threshold photoproduction
 - ◆ Low-x production
- Charged photoproduction & the structure of exotics
- Backward photoproduction via baryon exchange & baryon stopping*
- Conclusions

*Work done in collaboration with Aaron Stanek & Sam Heppelman



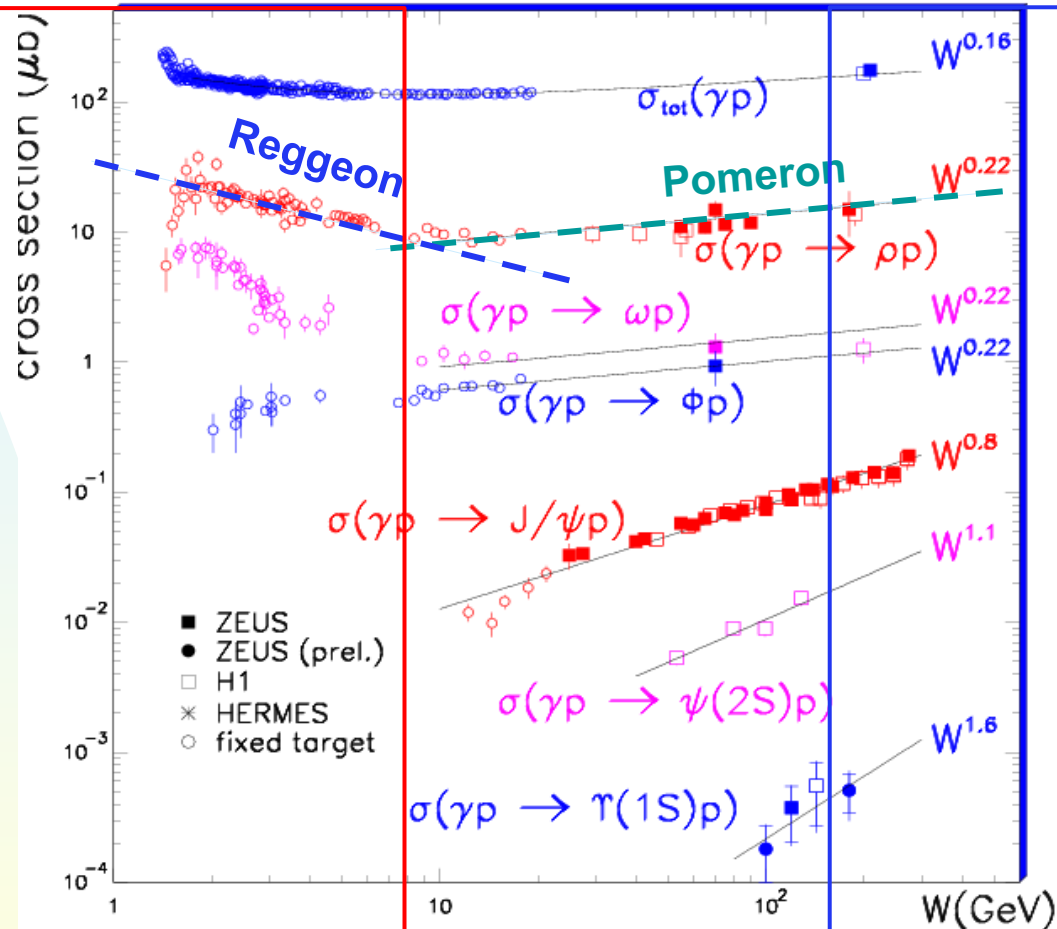
Photoproduction and Electroproduction

- Two models
- Gluon exchange
 - ◆ Lowest order is 2-gluon exchange
 - ◆ Higher orders is gluon ladder
- Pomeron + Reggeon exchange
 - ◆ Pomeron is gluon ladder
 - ✦ $J^{PC}=0^{++}$ explains vector meson dominance
 - Absorptive part of the potential
 - ◆ Reggeon involves quark exchange
 - ✦ meson trajectories (q-qbar)
 - ✦ Wide range of spin, parity and charge
 - ✦ Allows production of pentaquarks, other exotica etc.



Pomerons, Reggeons and kinematics

HERA + fixed target data



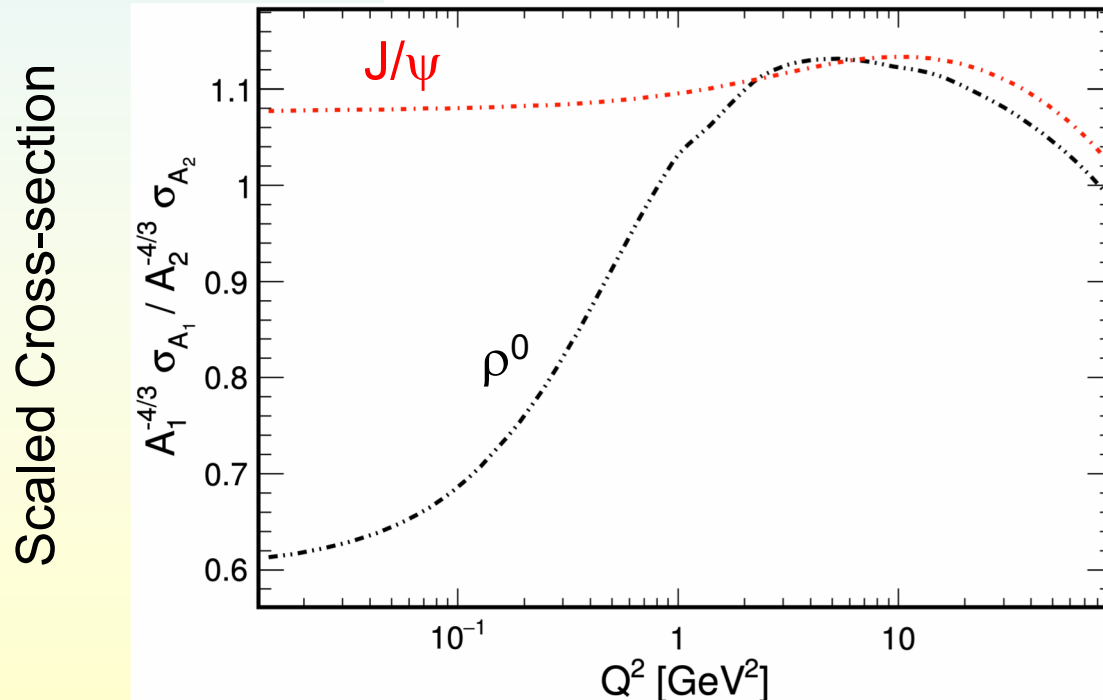
Ion-going direction
 Reggeon dominated
 Near threshold production
 Pentaquarks, other exotica

Rapidity:
 $y = \ln(W^2/m_p m_V)$

Electron-going direction
 Pomeron-dominated
 Low-x gluons
 saturation

Photoproduction vs. electroproduction

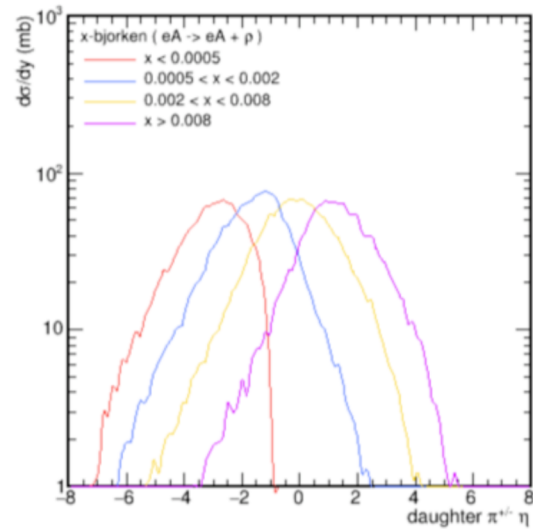
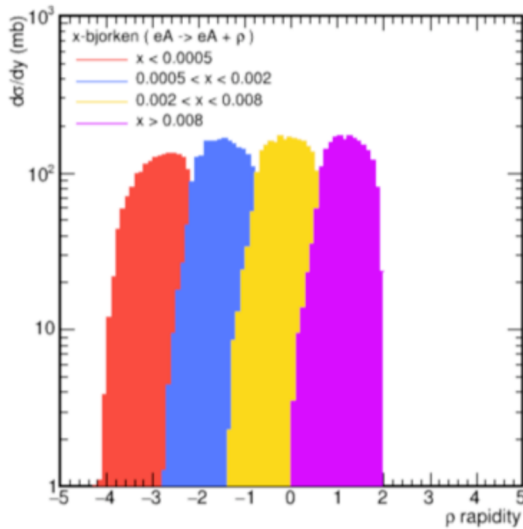
- Most EIC attention is on electroproduction ($Q^2 > 1 \text{ GeV}^2$)
- Photoproduction ($Q^2 < 1 \text{ GeV}^2$) is critical for studying shadowing, which should disappear at large Q^2
 - ◆ Good acceptance is needed for vector mesons at low p_T
- Shadowing is larger for lighter mesons
 - ◆ ρ/ϕ are experimentally important



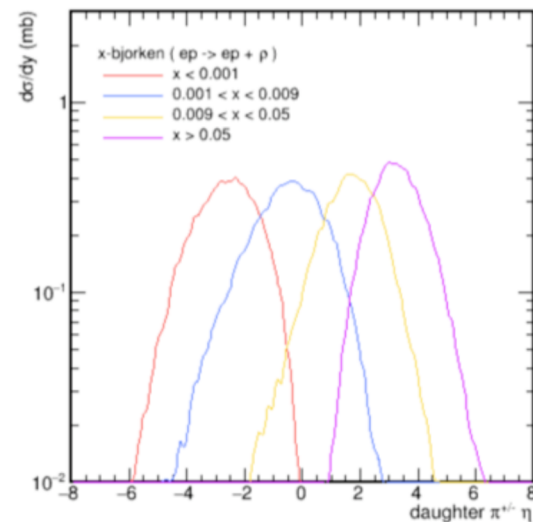
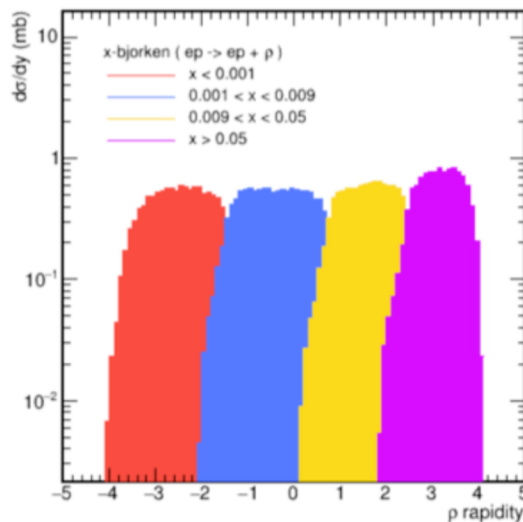
From eSTARlight. Similar plot by Mantysari and Venugopalan, Phys. Lett. B **781**, 664 (2018)

The ρ^0

- $10^{-4} < x < 1$ corresponds to $-4 < y < 4$
- Coverage up to rapidity $|y|$ requires coverage to $|\eta| > |y|+1$



eAu
18 GeV * 100 GeV

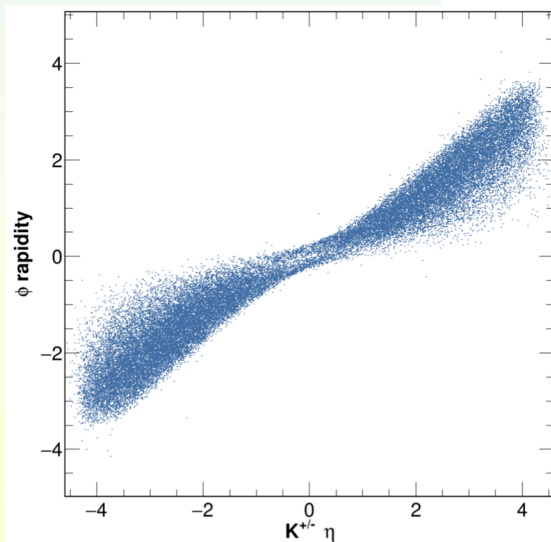


ep
18 GeV * 275 GeV

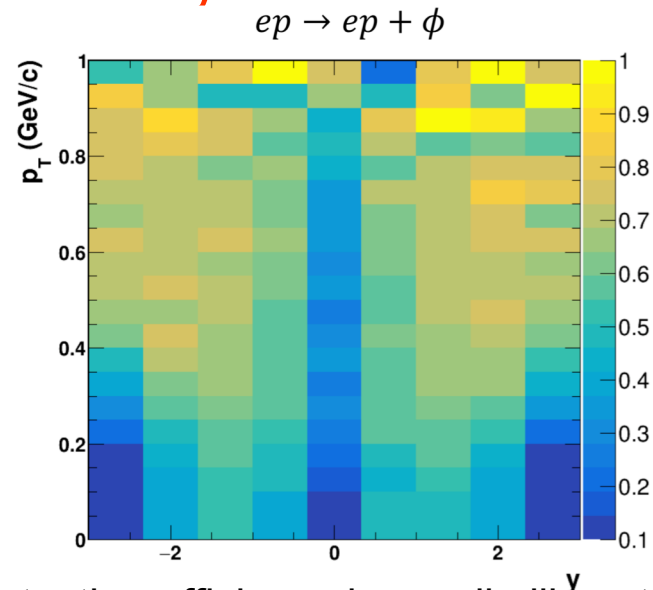
J. Arrington *et al.*,
arXiv:2102.08337

The ϕ

- $\Phi \rightarrow K^+K^-$ is the only viable decay mode
 - ◆ Kaon momentum is 135 MeV/c in ϕ frame
 - ✦ Relationship between kaon $\langle \eta \rangle$ and $\phi \langle y \rangle$ is nonlinear
 - Reduces detection efficiency at large $|y|$
- For photoproduction near $y=0$, kaons have $p=135$ MeV/c
 - ◆ Highly ionizing
 - ◆ Requires low B field and very low material to detect
- J/ψ & heavier mesons well handled by reference detector



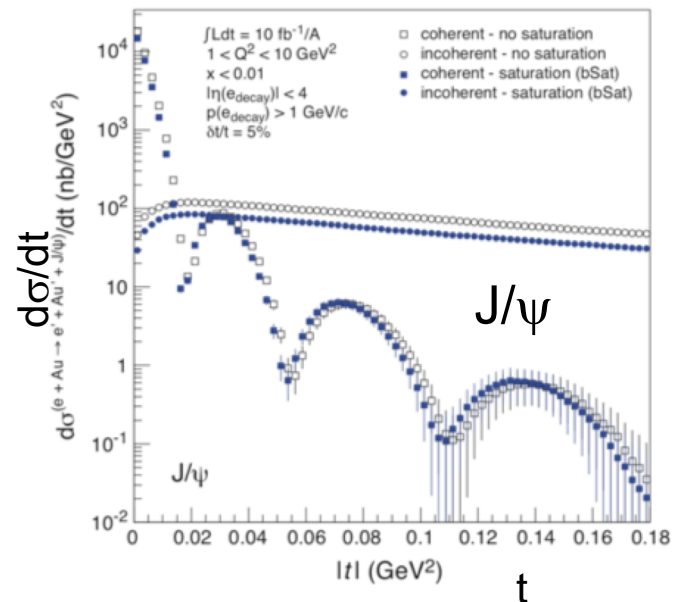
ϕ Rapidity vs. kaon pseudorapidity



ϕ detection efficiency in an all-silicon tracker

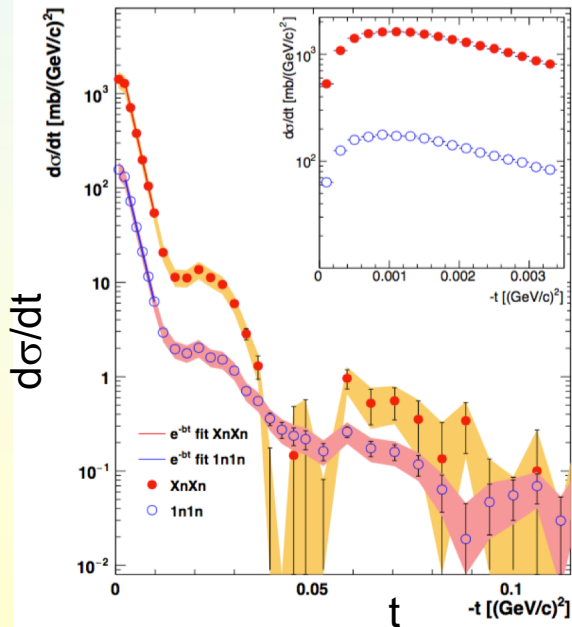
Separating Coherent and Incoherent production

- In the Good-Walker paradigm, coherent photoproduction is sensitive to the average nuclear shape, while the incoherent production is sensitive to event-by-event fluctuations
 - ◆ Need 500:1 coherent:incoherent separation at 2nd minimum
 - ✦ Neutron or proton emission is easy, but there are excitations that decay by emitting MeV (in the target frame) photons
 - ✦ Situation murky; requires good acceptance for far-forward $E < 100$ MeV
- Lead preferred over gold. It has no low-lying, long-lived states



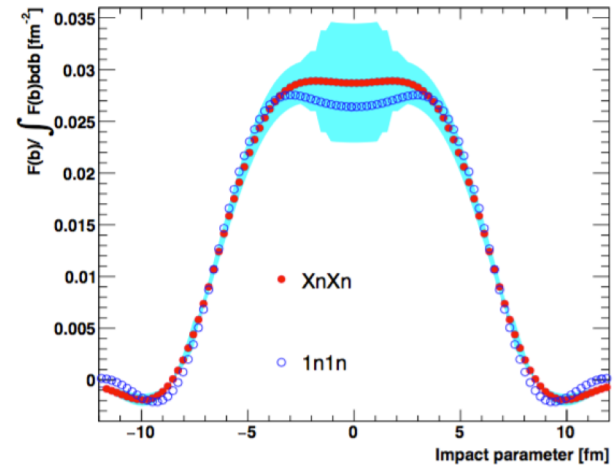
Observing diffractive dips

- $d\sigma/dt$ for coherent photoproduction probes the transverse gluon distribution in nuclei – like GPDs, but for nuclei.
- Requires good measurement of $t \sim p_{T,Pomeron}^2$
- $P_{T,VM} = P_{T,Pomeron} \oplus P_{T,photon} \oplus \text{Resolution}$
 - ◆ Need photon p_T to accurately determine Pomeron p_T
 - ✦ Observe scattered electron down to low Q^2
 - Limited by beam emittance; easier at higher k/E_e



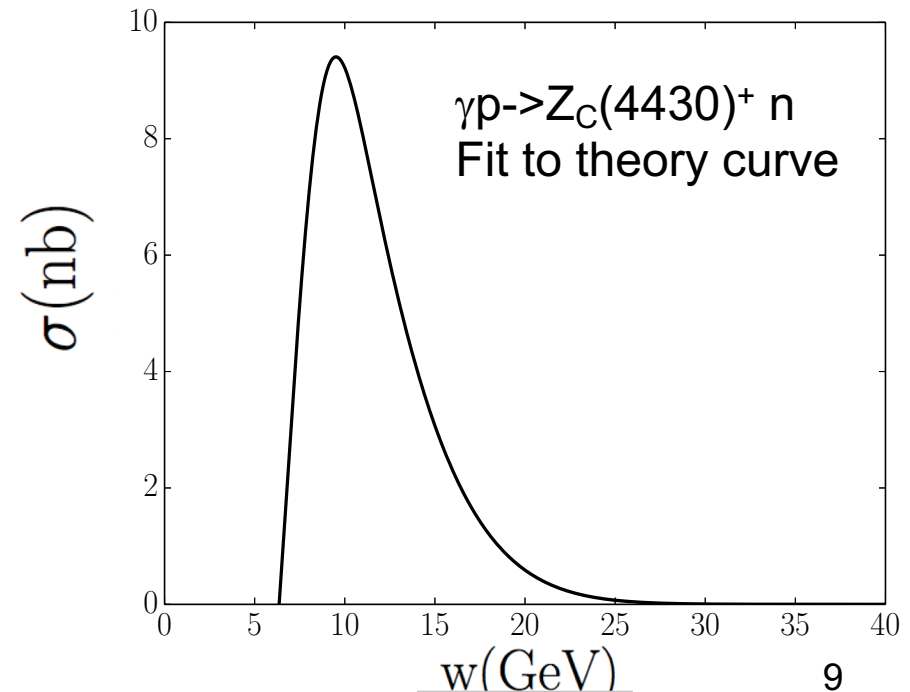
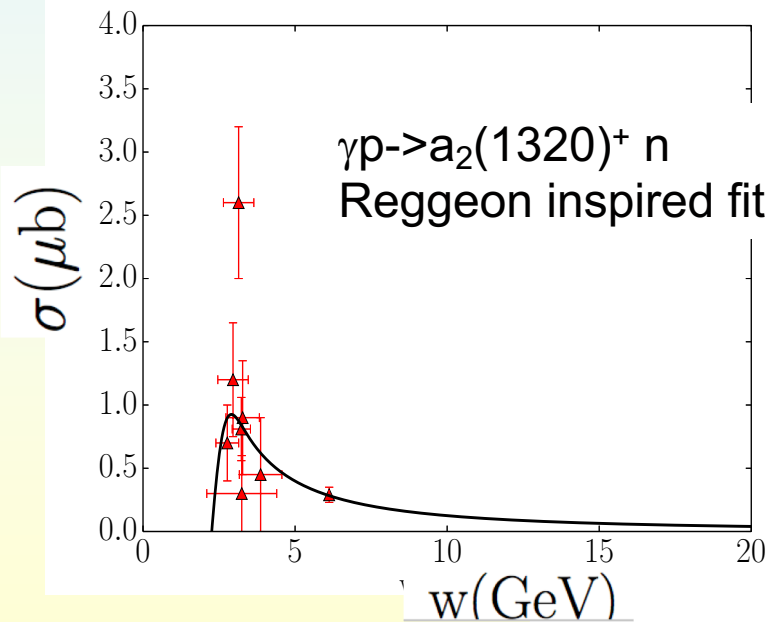
STAR ρ^0

$$F(b) \propto \frac{1}{2\pi} \int_0^\infty dp_{T,Pomeron} J_0(b p_{T,Pomeron}) \sqrt{\frac{d\sigma}{dt}}$$



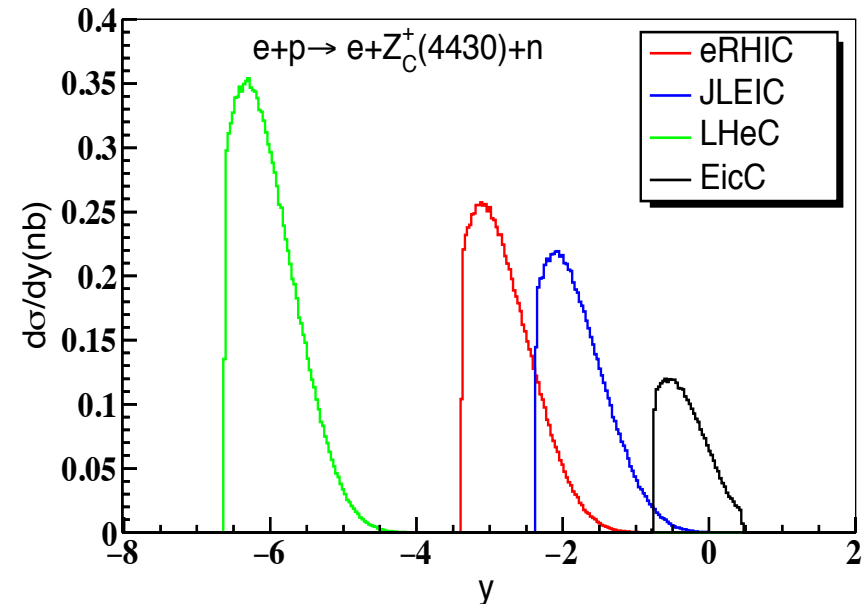
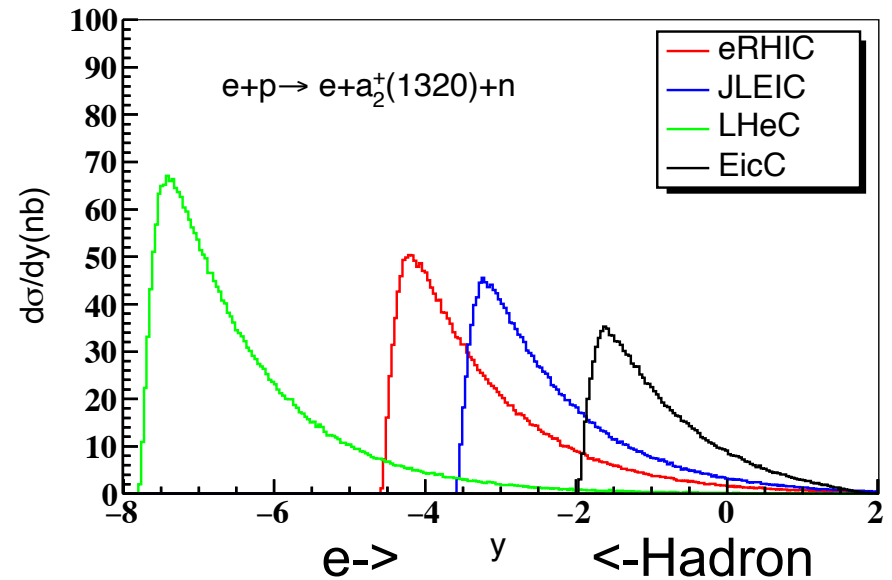
Reggeon exchange and forward production

- Examples: the $a_2^+(1320)$ standard candle and the exotic $Z_c^+(4430)$
- Use data/calculations of $\sigma(\gamma p \rightarrow X+n)$ as input to eSTARlight to predict $d\sigma/dy$ for the same process in EIC collisions/
 - Use the same Q^2 scaling as the ρ (for the a_2) and J/ψ (for the Z_c)



$a_2^+(1320)$ and $Z_c^+(4430)$ production in ep collisions at the EICs

- The $a_2^+(1320)$ is mainly at negative rapidity
 - ◆ $\sigma \sim 80$ nb at eRHIC
 - ✦ Copiously produced
- The $Z_c^+(4430)$ is heavier, and so somewhat more centrally produced.
- ◆ σ is 0.26 nb at eRHIC
- Both require good ion-going acceptance to be observable
- Both might be easier to observe at lower beam energies



Backward meson production

- Data from fixed-target experiments (including JLab), show that photoproduction can also occur in the backward production

 - ◆ Model via a baryon exchange trajectory

- Normally, photoproduction is maximal when t (momentum transfer from target) is small

 - ◆ $d\sigma/dt \sim \exp(-Bt)$

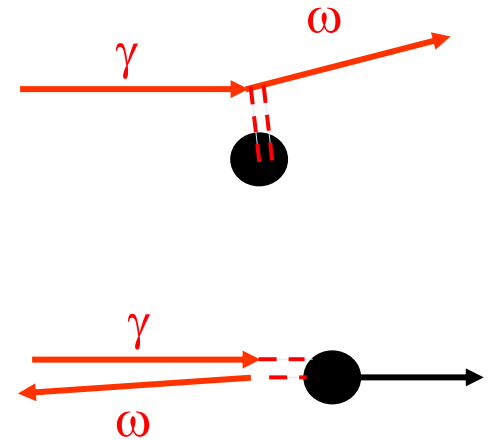
 - ✦ $B \sim \hbar/\text{target size}$

- In baryon exchange, in the CM frame, the meson scatters backward 180 degrees causing the baryon to recoil

 - ◆ In CM frame, baryon and photon/meson trade momentum

 - ◆ Mandelstam u is small, but t is large ($t > Q^2$)

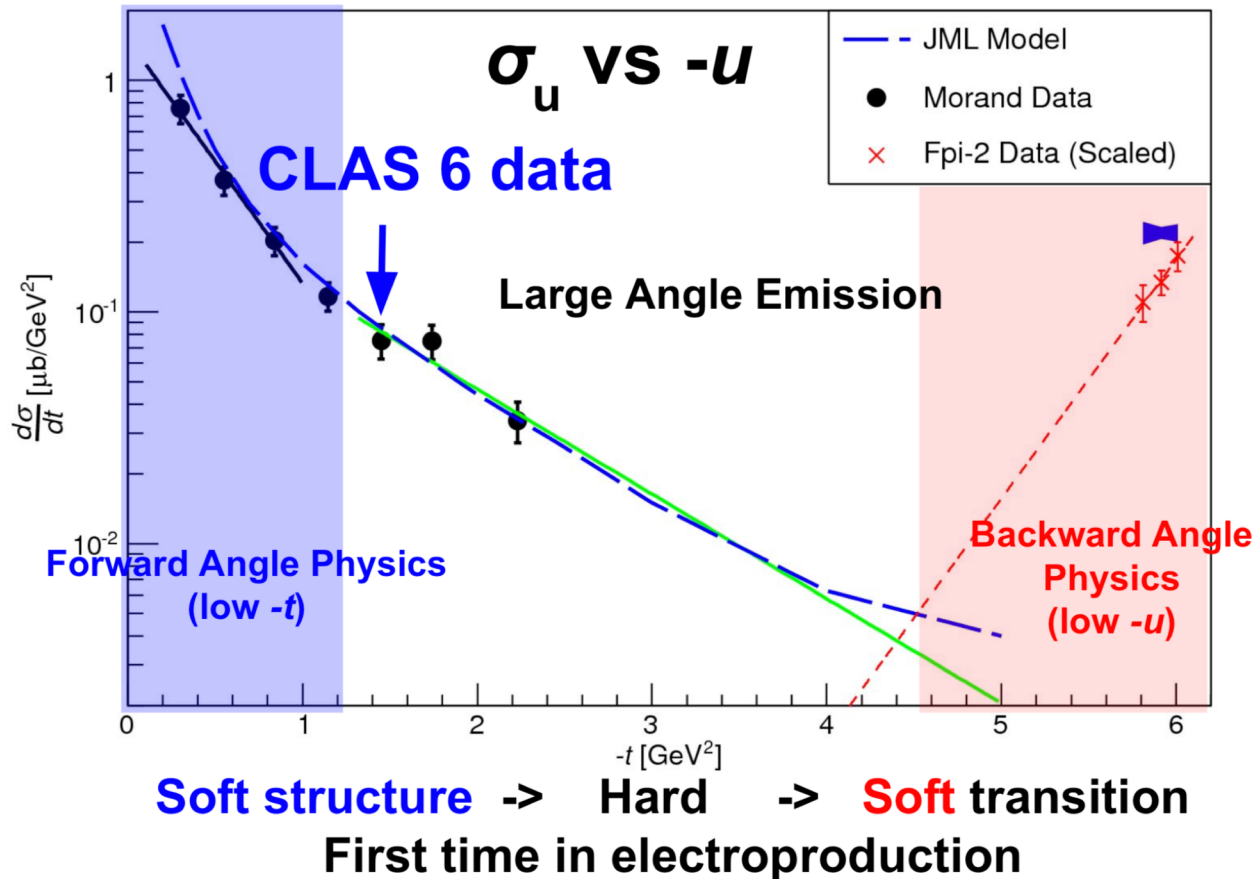
- How does an intact baryon recoil at high energies?
Similar to baryon stopping in RHI collisions



$\gamma p \rightarrow \omega p + \rho p$

- Electroproduction data from Clas 6 at Jlab
- Forward & backward interactions are soft; intermediate is hard

$$\gamma^* + p \rightarrow p + \omega, W = 2.47 \text{ GeV}, Q^2 = 2.35 \text{ GeV}^2$$



Plot from Bill Li (William & Mary).

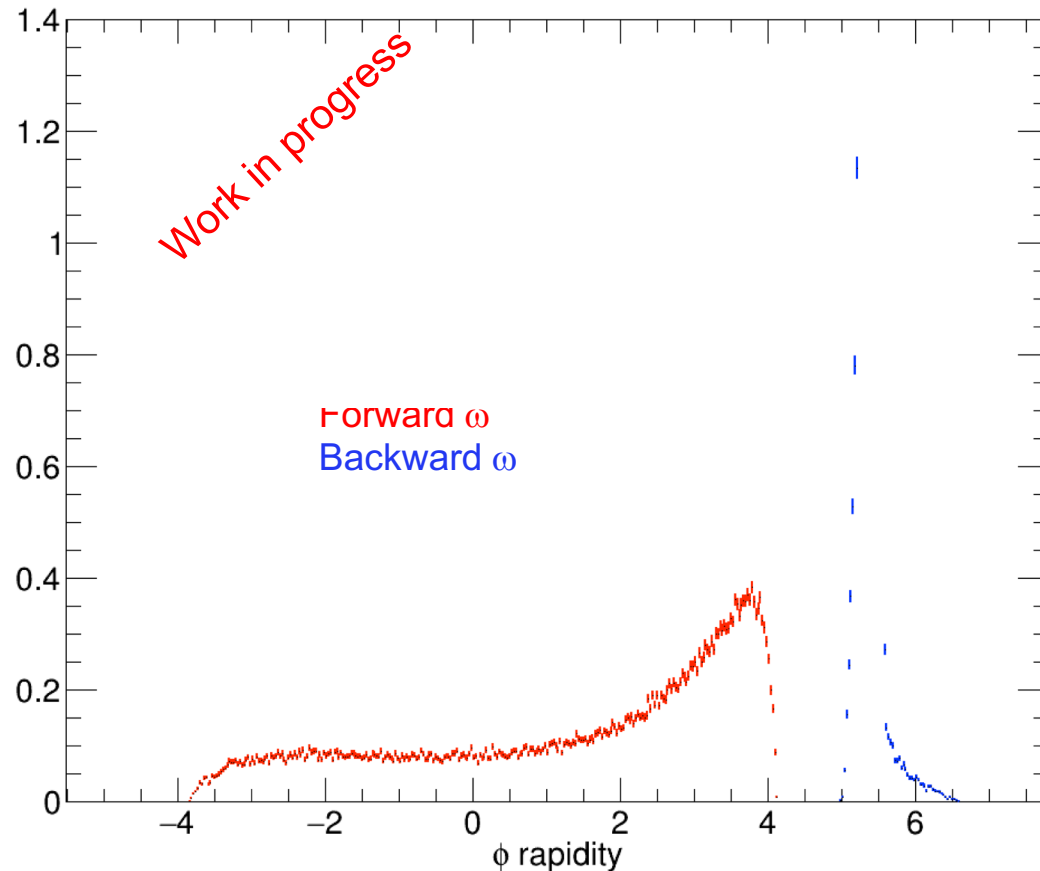
Parameterization of backward $\gamma p \rightarrow \omega p$

- ω is best studied backward photoproduction case
 - ◆ Fit to data from two experiments
- Assume same Reggeon-like form as forward production:
 - ◆ $d\sigma/dt|_{t=0} \sim A (s/1\text{GeV})^B$ embodies physics of reaction
 - ◆ $d\sigma/dt \sim \exp(-Ct)$ accounts for form factor (size) of target
 - ◆ Swap u for t , to match behavior of backward kinematics
- $d\sigma/du|_{u \sim 0} = A (s/1\text{GeV})^B$
 - ◆ $A = 4.4 \mu\text{b}/\text{GeV}^2$
 - ✦ $A = 180 \mu\text{b}/\text{GeV}^2$ for forward ω photoproduction
 - ◆ $B = -2.7$
 - ✦ $B = -1.92$ for forward ω photoproduction
- $d\sigma/du \sim \exp(-Cu)$, with $C = -21 \text{ GeV}^{-2}$
 - ◆ Similar slope as C in e^{Ct} term for forward $\gamma p \rightarrow \rho p$
- Rate is few % of the forward rate for $k \sim \text{GeV}$
 - ◆ Falls off a bit faster with increasing energy.
 - ◆ Cross-sections are large enough to be easily accessible.

EIC backward production kinematics

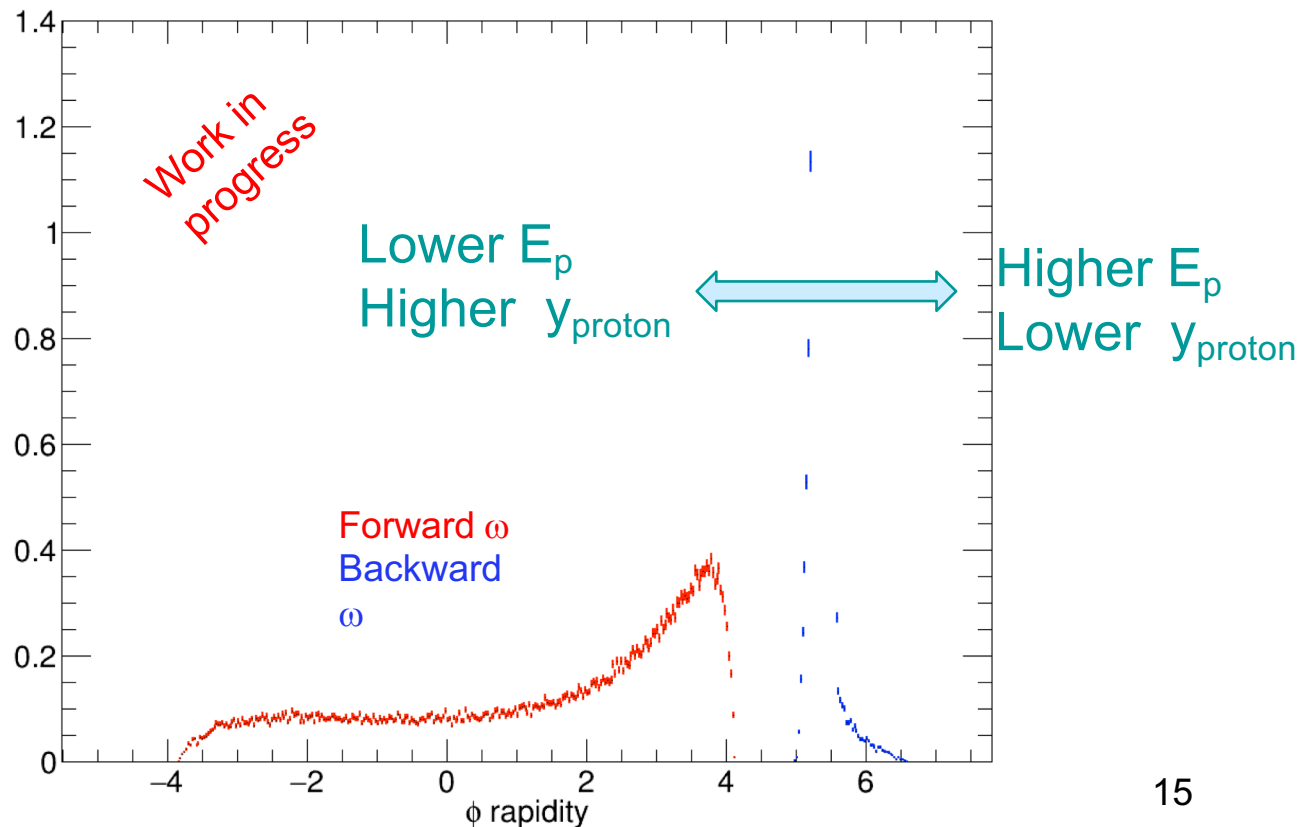
- An ω at near-beam rapidity, and a mid-rapidity proton
 - ◆ The proton is easily detectable
 - ◆ The forward vector meson looks tough.
 - ✦ Charged particle tracking problematic (?)
 - ✦ $\omega \rightarrow \pi^0 \gamma$ is a promising channel, since it is fully calorimetric

ep \rightarrow e p ω
18 GeV e on 100 GeV
eSTARlight simulation



Beam energy dependence of ω peak

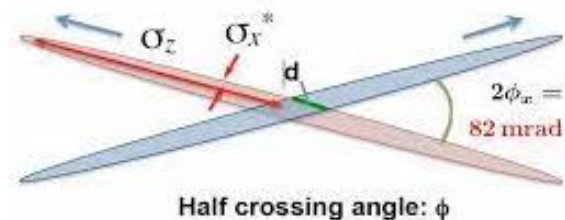
- For 275 GeV proton beams, y_ω rapidity ~ 6.5
- For 41 GeV protons beams
 - ◆ Proton rapidity = 0.0 \rightarrow typical ω rapidity is 4.6
 - ◆ Proton rapidity = 4.0 \rightarrow typical ω rapidity is 3.7
- Need to explore full phase space, but lower proton beam energies seem better



ep \rightarrow e p ω
18 GeV e on 100 GeV
eSTARlight simulation

How would a 2nd detector/low energy improve the rapidity/energy coverage?

- Low-energy running shifts the forward region toward mid-rapidity.
 - ◆ Near-threshold production, pentaquarks etc. become more central
- This does not work on the low-x side
 - ◆ Good rapidity coverage is needed to exploit the full EIC energy
- Another partial solution: instrument the IR above & below the plane where the beams diverge



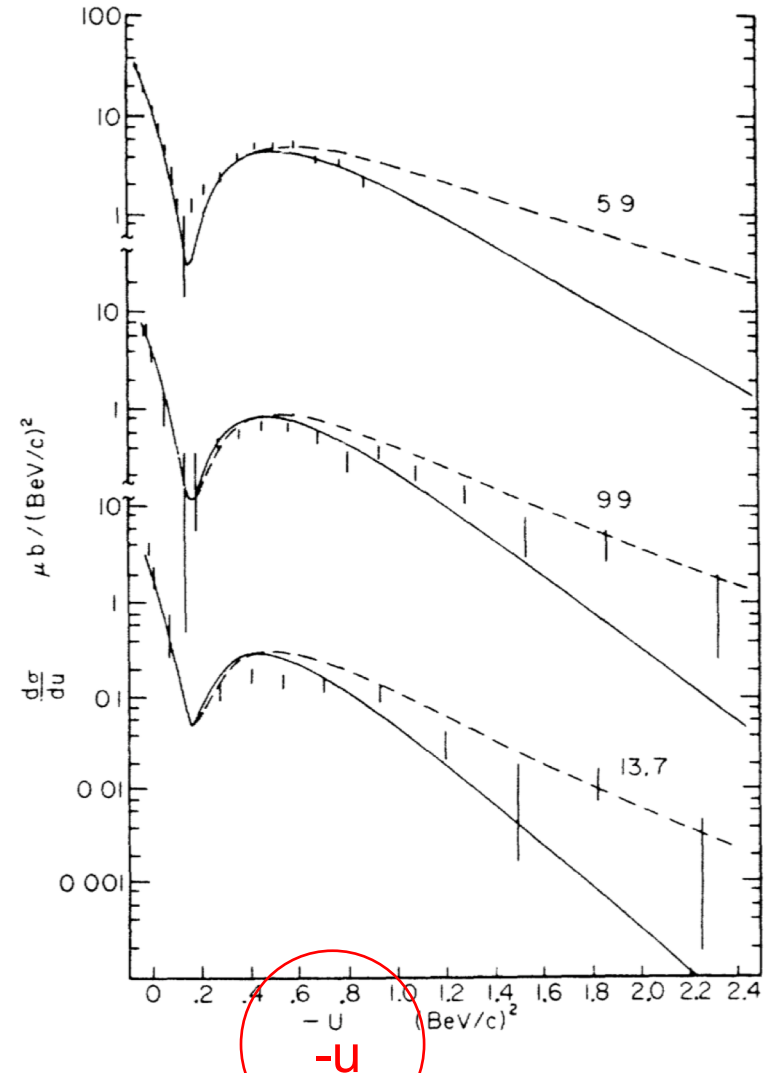
Conclusions

- Vector meson and other exclusive/semi-exclusive production reactions can challenge proposed EIC designs.
- Very wide pseudorapidity coverage is required to study vector meson production over the full range of Bjorken- x .
- Near-threshold production and Reggeon-exchange production, including exotica requires good acceptance in the ion-going direction.
 - ◆ Running at a reduced ion beam energy will shift this production toward mid-rapidity.
- Excellent far-forward ion-going detectors are required to separate coherent and incoherent photoproduction, and to study.
- Backward production reactions lead to mid-rapidity baryons and far-forward mesons. The later are a detector challenge, requiring more study.

Backup

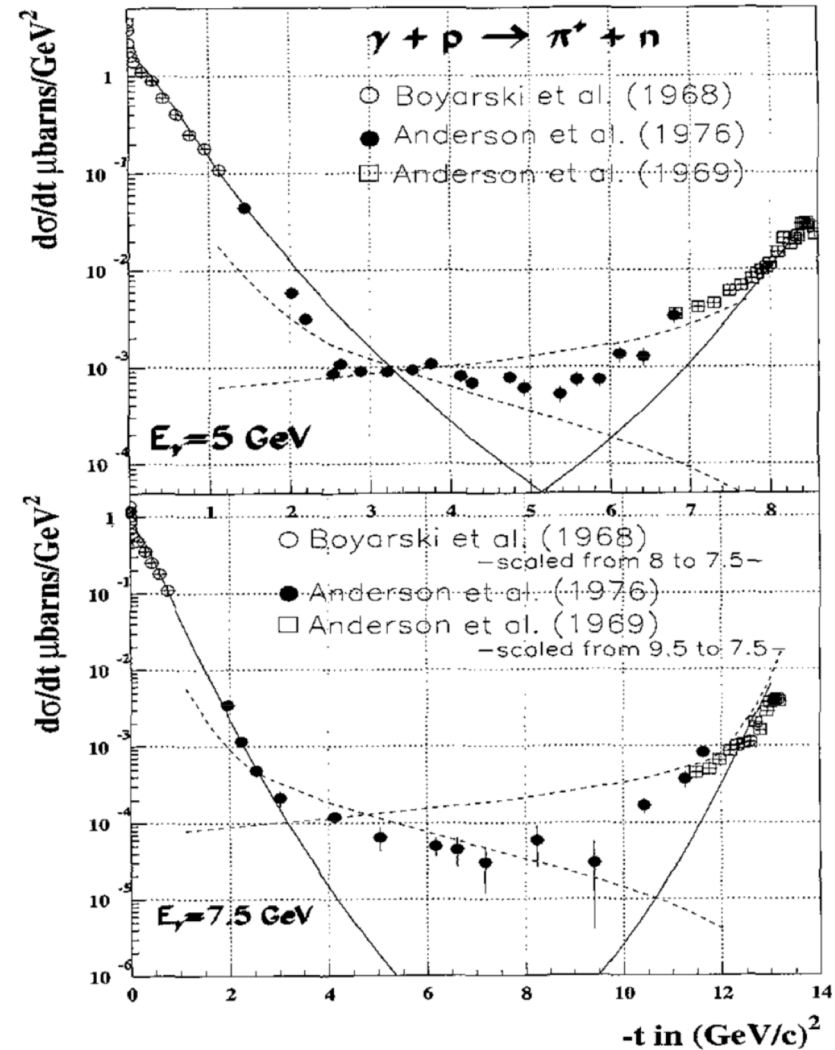
$\pi^+p \rightarrow \pi^+p$ elastic scattering

- $5.9 \text{ GeV} < E_\pi < 13.7 \text{ GeV}$
 - ◆ Above the resonance region
- Clear peak near $u=0$
 - ◆ Elastic scattering in the backward direction
- Diffractive minima visible in u -spectrum
 - ◆ Looks a lot like a form factor



$\gamma p \rightarrow \pi^+ n$

- Data from multiple experiments
- Data exists for $4 \text{ GeV} < E_\gamma < 16 \text{ GeV}$
- ◆ Again, above the resonance region



Backward ω data for fit

- The ω is one of the better studied mesons for backward production. There is more data available than for the ρ .
- ◆ Reasonable lever arm for photon energy.

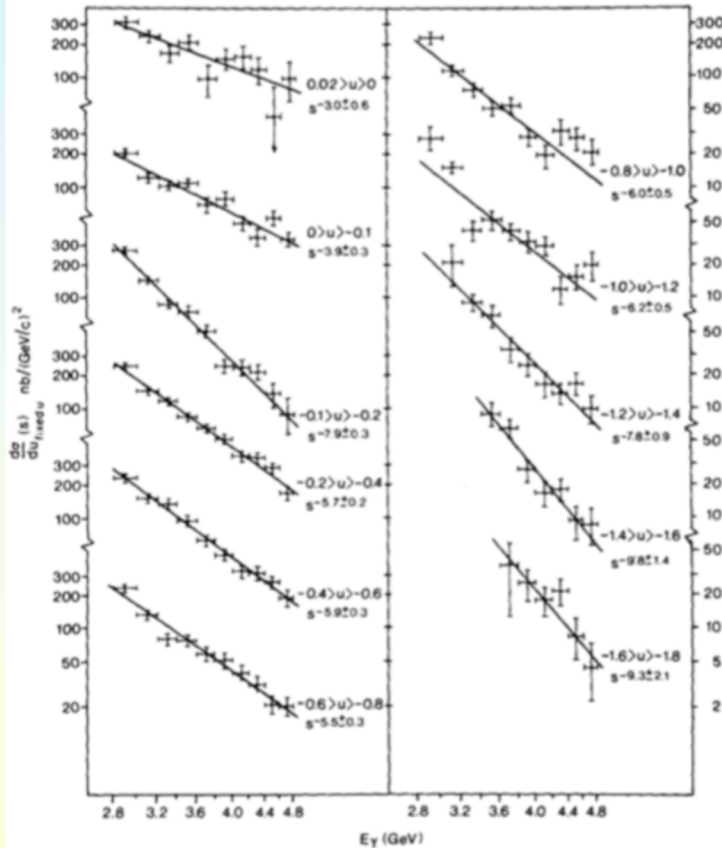


TABLE I. The compiled data. Errors on original data were around 25% of the listed value. Error due to transcription from figure is estimated to be less than 5%.

E_γ GeV	$d\sigma/du(u \approx 0)$ nb/GeV ²	Source
2.9	200	Sibirtsev et al. ⁶ Figure 1
3.0	300	Cliff et al. ⁴ Figure 3
3.0	200	Sibirtsev et al. ⁶ Figure 7
3.2	240	Cliff et al. ⁴ Figure 3
3.3	110	Sibirtsev et al. ⁶ Figure 7
3.5	170	Cliff et al. ⁴ Figure 2
3.5	170	Sibirtsev et al. ⁶ Figure 1
3.5	100	Sibirtsev et al. ⁶ Figure 7
3.6	210	Cliff et al. ⁴ Figure 3
3.6	100	Sibirtsev et al. ⁶ Figure 7
3.8	90	Cliff et al. ⁴ Figure 3
3.9	60	Sibirtsev et al. ⁶ Figure 7
4.0	150	Cliff et al. ⁴ Figure 3
4.1	70	Sibirtsev et al. ⁶ Figure 7
4.2	160	Cliff et al. ⁴ Figure 3
4.3	40	Sibirtsev et al. ⁶ Figure 7
4.4	120	Cliff et al. ⁴ Figure 3
4.4	30	Sibirtsev et al. ⁶ Figure 7
4.5	50	Sibirtsev et al. ⁶ Figure 7
4.6	30	Sibirtsev et al. ⁶ Figure 7
4.7	75	Cliff et al. ⁴ Figure 2
4.7	80	Sibirtsev et al. ⁶ Figure 1
4.8	100	Cliff et al. ⁴ Figure 3

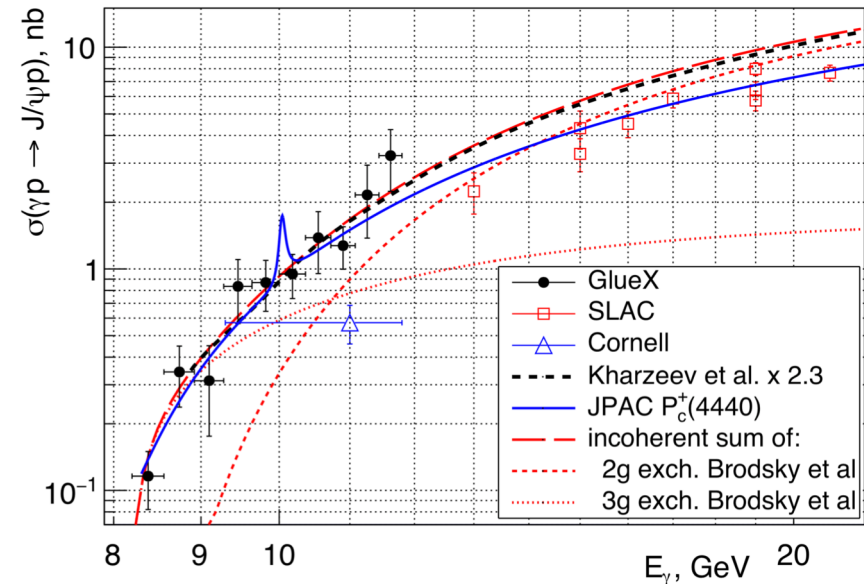
⁴R. Clift *et al.*, Physics Letters **72B**, 144 (1977).

⁵B.-G. Yu and K.-J. Kong, Physical Review D **99** (2019).

⁶R. Sibirtsev *et al.*, arXiv:nucl-th/0202083v1 (2002).

Near threshold quarkonium production

- Near-threshold quarkonium production is sensitive to new mechanisms (i. e. 3-gluon exchange)
 - ◆ GlueX data favors a mix gluon exchange for J/ψ
- Sensitive to near-threshold
 - ◆ $P_C^+(4440) \Rightarrow J/\psi p$
 - ◆ Pentaquark candidate
- EIC will study ψ' , Y states and probe the Q^2 dependence of multiple resonances
- For nuclei, near-threshold or sub-threshold production is sensitive to short-range nuclear correlations.
- Requires good acceptance in the ion-going direction



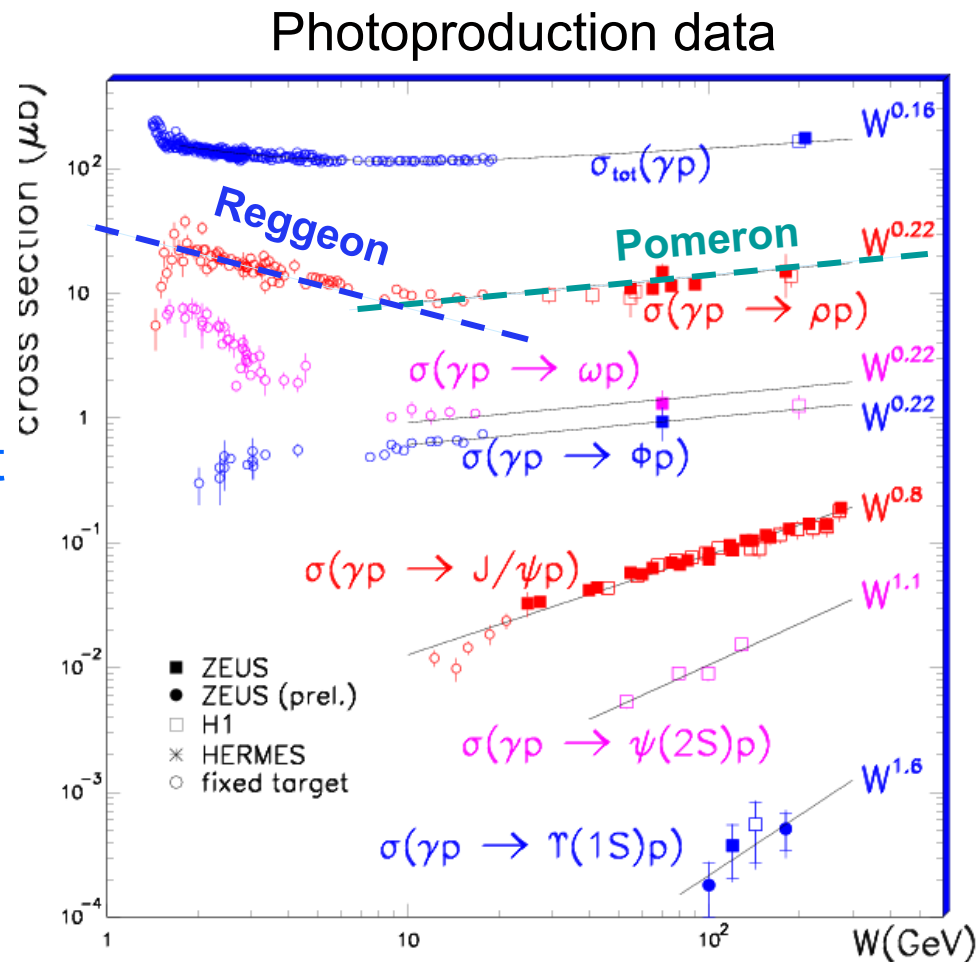
Photoproduction & electroproduction in eSTARlight

- Convolution of photon flux from electron with $\sigma(\gamma p \rightarrow V p)$
 - ◆ Both depend on Q^2
- Weizsacker-Williams photon flux (with non-zero Q^2)
- VM cross-sections parameterized from HERA data/theory....
 - ◆ Reggeon and Pomeron exchange
 - ✦ Q^2 dependence via a power law from HERA data
- Other cross-sections from theory predictions
- Nuclear targets included with a Glauber calculation
- Vector mesons retain the photon spin
 - ◆ For $Q^2 \sim 0$, transversely polarized
 - ◆ As Q^2 rises, longitudinal polarization enters
 - ◆ Spin-matrix elements quantified with HERA data
- Embodied in eSTARlight code, available at:
<https://github.com/eic/estarlight/>

n.b. new
location

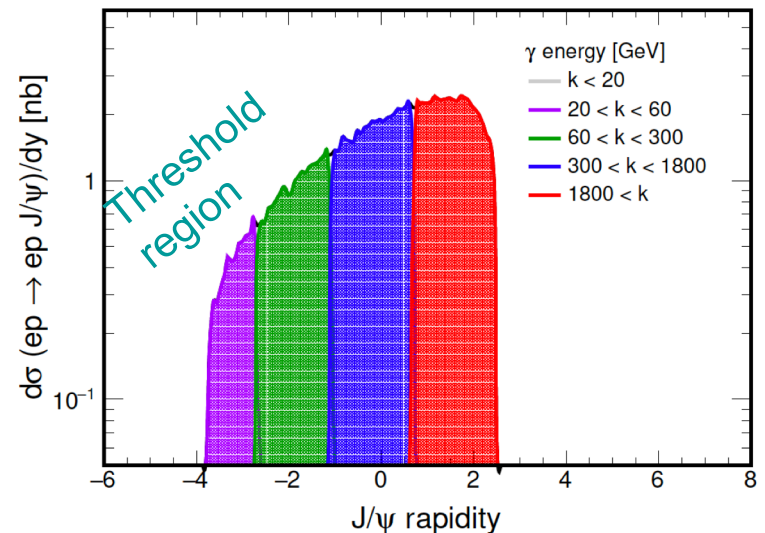
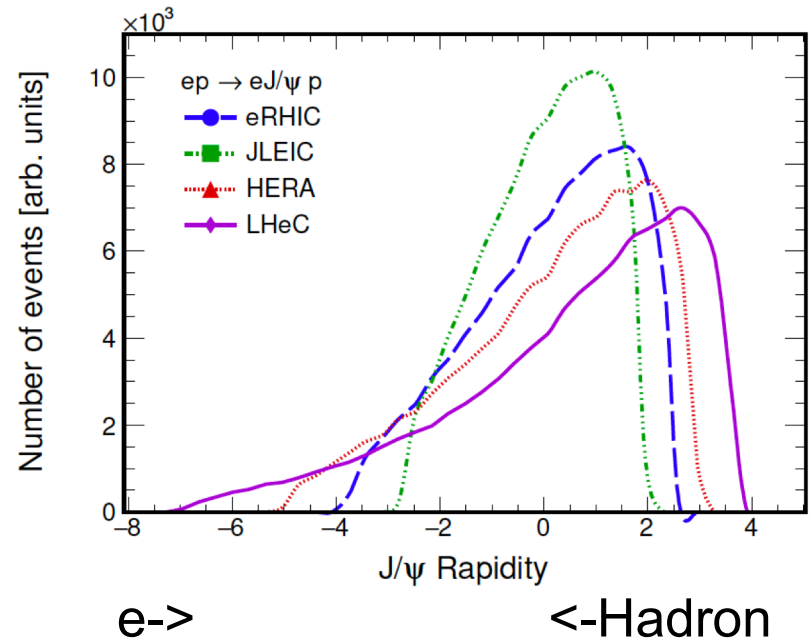
Pomerons and Reggeons in photoproduction

- HERA photoproduction cross-sections well fit by
- $\sigma(W) = XW^\epsilon + YW^{-\eta}$
 - ◆ $W = \gamma p$ CM energy
- XW^ϵ : Pomeron (gluons)
 - ◆ $\epsilon \sim > 0.2$ – meson dependent
 - ◆ $J^{PC} = 0^{++}$
- $YW^{-\eta}$: ‘Reggeon’ ($\sim\sim qq\bar{q}$)
 - ◆ $\eta \sim\sim 1.5$
 - ◆ Summed light-quark meson trajectories
 - ◆ \sim valence quarks
 - ◆ Zero for ϕ , J/ψ , etc.
 - ◆ Range of spin/parity
- Q^2 dependence – power law



EIC photoproduction kinematics

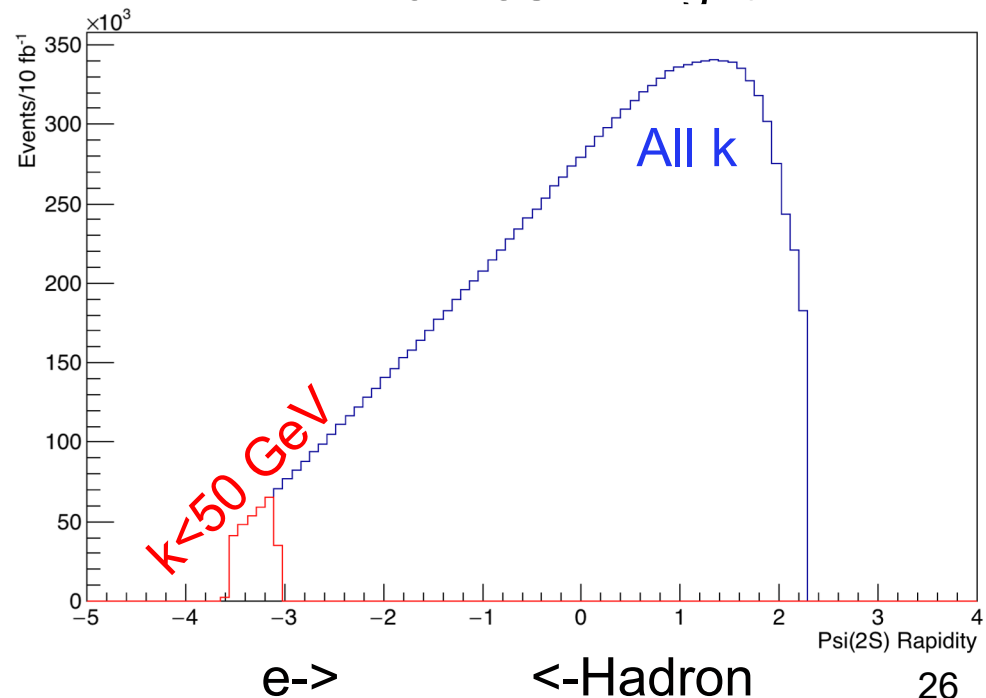
- Maps photon energy onto rapidity
- $k = \frac{M}{2} \exp(y)$
- $y = \ln(2k/M)$
- Reggeon activity strongest at low photon energies
 - ◆ Requires good acceptance in the hadron-going direction
- Highest photon energies correspond to electron-going direction
 - ◆ Need good e-going acceptance



$\Psi(2S)$ & Y photoproduction at eRHIC

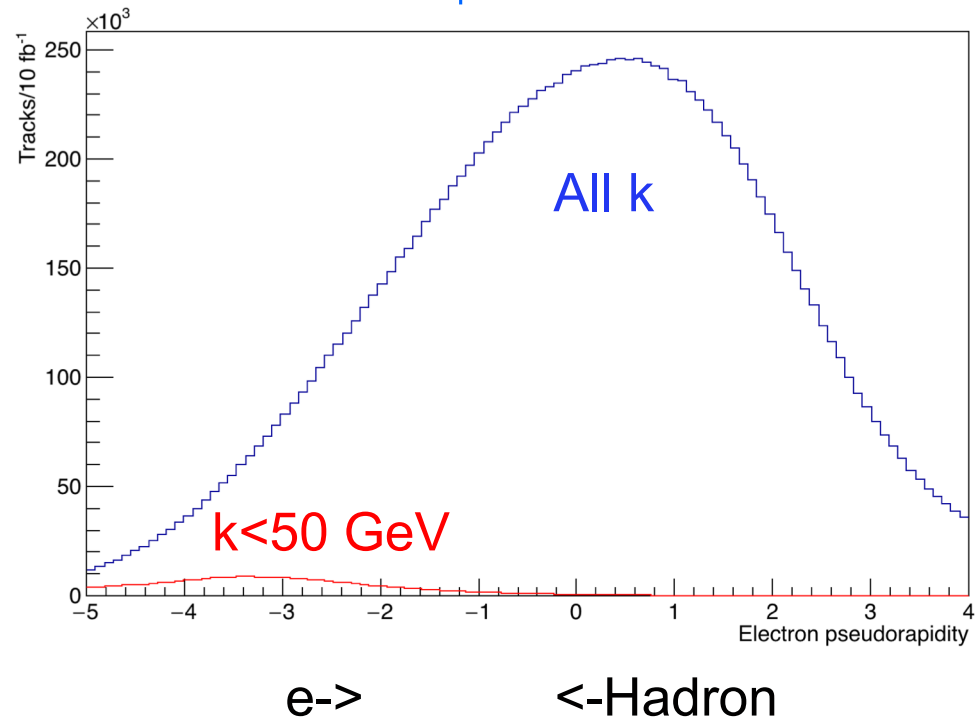
- 18 GeV e^- on 275 GeV protons
- $\Psi(2S)$: $\sigma=1.4$ nb (1/6 of $\sigma(J/\psi)$)
 - ◆ 14 million events in 10 fb^{-1}
- 300,000 events with photon energy <50 GeV (target frame)
 - ◆ $\Psi(2s)$ threshold region is $3.5 < y < 3.0$ for this configuration
 - ◆ $\sim 2,800$ each $\Psi(2S) \rightarrow ee, \mu\mu$
- $\sigma(Y(1S))=0.01\sigma(\psi')$
 - ◆ 140,000 events/ 10fb^{-1}
 - ◆ $\sim 3,000$ each to $ee, \mu\mu$
 - ◆ $\sim 3,000$ near-threshold events
 - ◆ ~ 75 each to $ee, \mu\mu$
 - ◆ More central than ψ'

From eSTARlight



$\Psi(2S) \rightarrow ee$ lepton pseudorapidities

- Lepton pseudorapidity depends on $\Psi(2S)$ rapidity, p_T and polarization (which depends on Q^2)
- Leptons from most near-threshold ($k < 50$ GeV target frame) $\Psi(2S) \rightarrow ll$ decays have $-5 < y < -2$
 - ◆ Good acceptance required in hadron-going direction
 - ◆ N.b. $\text{Br}(\Psi(2S) \rightarrow ee \text{ or } \mu\mu)$ is 0.7%. Plus $J/\Psi \pi^+ \pi^-$
- Rates for $\Psi(1S)$ smaller - usable.
- Higher Ψ states accessible



Expected event rate for vector mesons, $a_2^+(1320)$ and Z_c^+

■ Total cross sections and expected events for vector mesons and two charged particles in ep collisions

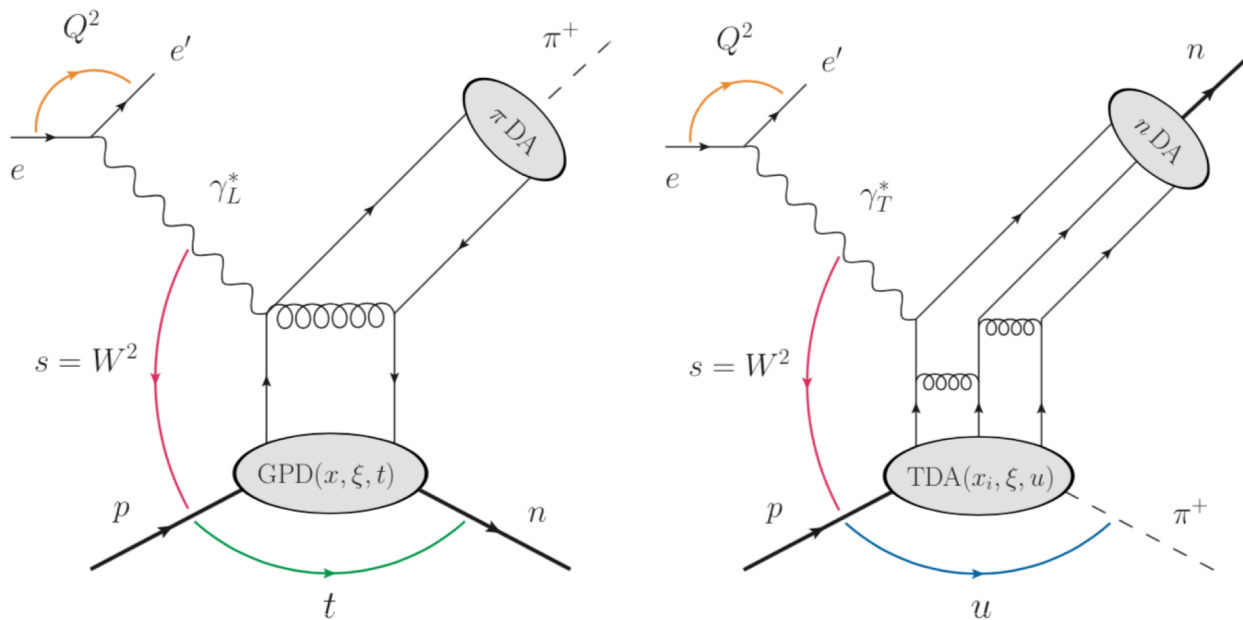
◆ 10 fb⁻¹ integrated luminosity

	Events ($0 < Q^2 < 1.0\text{GeV}^2$)				Events ($Q^2 > 1.0\text{GeV}^2$)			
	ρ	ϕ	J/ψ	ψ'	ρ	ϕ	J/ψ	ψ'
eRHIC -ep	50 giga	2.3 giga	85 mega	14 mega	140 mega	17 mega	5.7 mega	1.2 mega
eRHIC -eA	44 giga	2.8 mega	100 mega	16 mega	37 mega	5.6 mega	3.9 mega	960 kilo
JLEIC -ep	37 giga	1.6 giga	39 mega	6.0 mega	100.0 mega	12.0 mega	2.7 mega	550 kilo
JLEIC -eA	28 giga	1.6 giga	28 mega	3.9 mega	22 mega	3.2 mega	1.2 mega	250 kilo
LHeC -ep	100 giga	5.6 giga	470 mega	78 mega	260 mega	37 mega	29 mega	6.3 mega
LHeC -eA	110 giga	8.2 giga	720 mega	140 mega	100 mega	16 mega	27 mega	7.2 mega

	Events ($0 < Q^2 < 1.0\text{GeV}^2$)				Events ($1.0\text{GeV}^2 < Q^2 < 5.0\text{GeV}^2$)			
	eRHIC	JLEIC	LHeC	EicC	eRHIC	JLEIC	LHeC	EicC
$a_2^+(1320)$	0.79 giga	0.69 giga	1.06 giga	0.47 giga	5.1 mega	5.0 mega	5.2 mega	4.0 mega
$Z_c^+(4430)$	2.6 mega	2.2 mega	3.6 mega	0.94 mega	0.12 mega	0.12 mega	0.12 mega	68.0 kilo

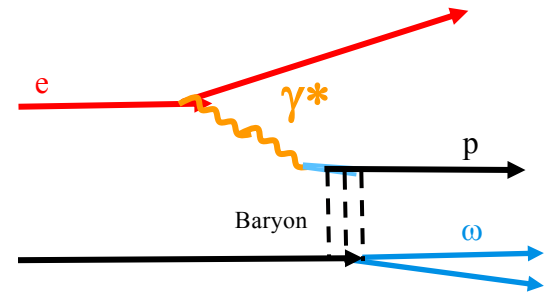
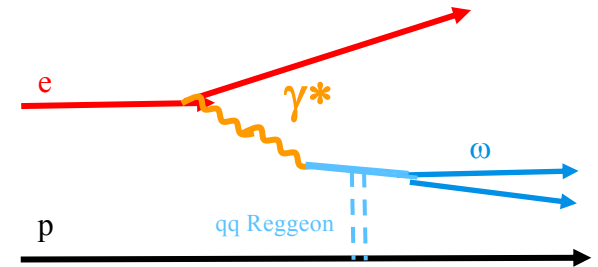
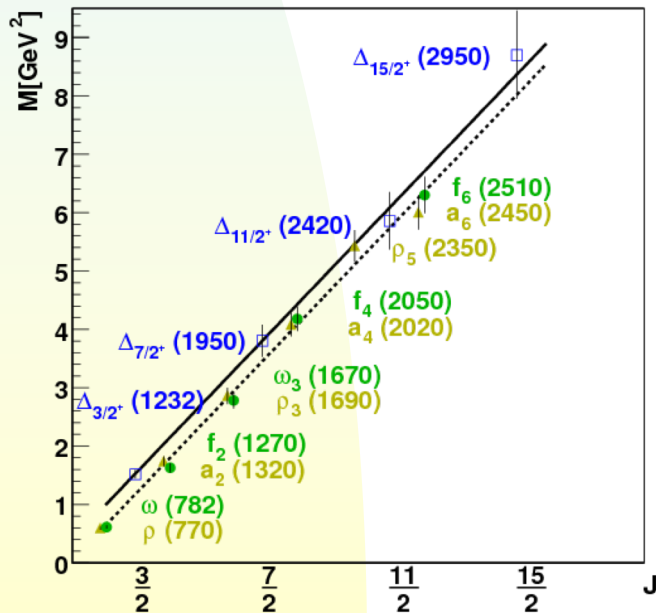
Theoretical approach - I

- GPD-like model, with Transition Distribution Amplitude quantifying baryon trajectories.



Theoretical approach II – Baryon trajectories

- For baryonic Regge trajectory
 - ◆ $\sigma(W) = XW^\epsilon + YW^{-\eta}$
 - ◆ Replace t with u , and much familiar behavior is restored.
 - ✦ Similar to meson trajectories
- Key trajectories: N , Δ ,
 - ◆ Λ/Σ for strangeness (not today)

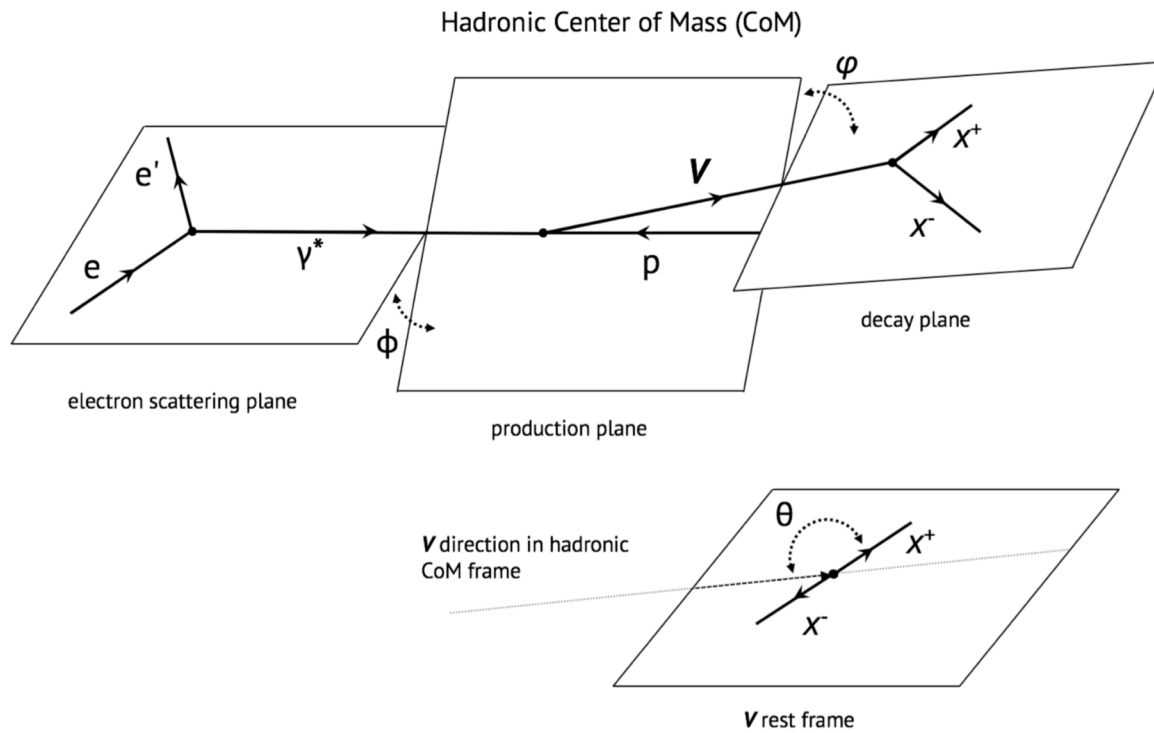


Implications for baryon stopping

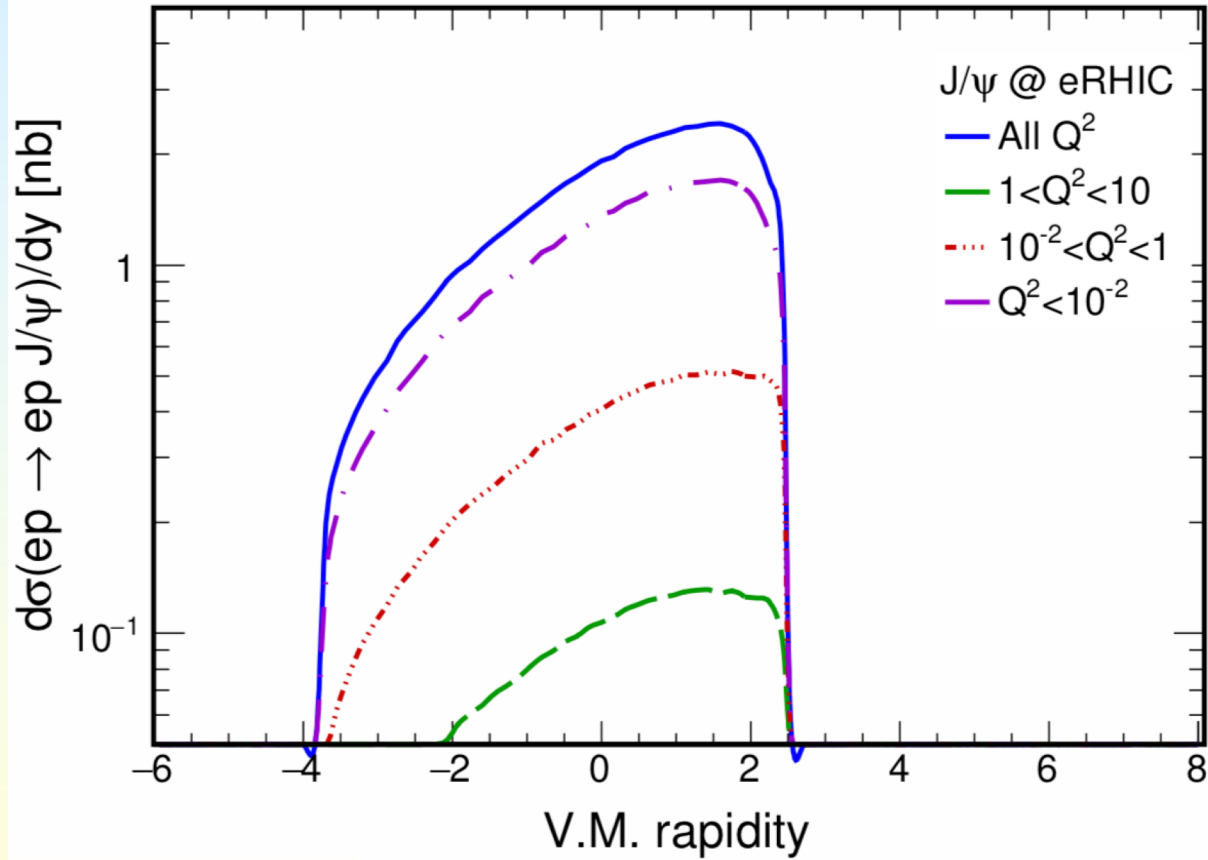
- Conventional wisdom: Regge phenomenology only matters at low energy
 - ◆ But... the relevant energy is the dipole-baryon CM energy.
 - ◆ soft dipole \rightarrow small CM energy.
 - ✦ Low-energy UPC photon
 - ✦ A soft virtual π
 - ✦ A low- x q - q bar dipole
 - ✦ Other configuration within an incident nucleus
- The baryon recoils but remains intact
 - ◆ Transport over multiple units in rapidity.
 - ✦ Like baryon stopping.
 - ◆ Phenomenology is very reminiscent of the baryon junction model.
 - ✦ Are there connections?

Backup

Angular definitions



Rapidity vs. Q^2



$\gamma p \rightarrow \omega p$ data

TABLE I. The compiled data. Errors on original data were around 25% of the listed value. Error due to transcription from figure is estimated to be less than 5%.

E_γ <i>GeV</i>	$d\sigma/du(u \approx 0)$ <i>nb/GeV²</i>	Source
2.9	200	Sibirtsev et al. ⁶ Figure 1
3.0	300	Clift et al. ⁴ Figure 3
3.0	200	Sibirtsev et al. ⁶ Figure 7
3.2	240	Clift et al. ⁴ Figure 3
3.3	110	Sibirtsev et al. ⁶ Figure 7
3.5	170	Clift et al. ⁴ Figure 2
3.5	170	Sibirtsev et al. ⁶ Figure 1
3.5	100	Sibirtsev et al. ⁶ Figure 7
3.6	210	Clift et al. ⁴ Figure 3
3.6	100	Sibirtsev et al. ⁶ Figure 7
3.8	90	Clift et al. ⁴ Figure 3
3.9	60	Sibirtsev et al. ⁶ Figure 7
4.0	150	Clift et al. ⁴ Figure 3
4.1	70	Sibirtsev et al. ⁶ Figure 7
4.2	160	Clift et al. ⁴ Figure 3
4.3	40	Sibirtsev et al. ⁶ Figure 7
4.4	120	Clift et al. ⁴ Figure 3
4.4	30	Sibirtsev et al. ⁶ Figure 7
4.5	50	Sibirtsev et al. ⁶ Figure 7
4.6	30	Sibirtsev et al. ⁶ Figure 7
4.7	75	Clift et al. ⁴ Figure 2
4.7	80	Sibirtsev et al. ⁶ Figure 1
4.8	100	Clift et al. ⁴ Figure 3

⁴R. Clift *et al.*, Physics Letters **72B**, 144 (1977).

⁵B.-G. Yu and K.-J. Kong, Physical Review D **99** (2019).

⁶R. Sibirtsev *et al.*, arXiv:nucl-th/0202083v1 (2002).