

# Finite Density QCD EoS and Net-Proton Cumulants in the BES Program

PRC 103 (2021) 3, 0314901

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# Outline

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1. Higher-order net-proton cumulants — why do we care?
2. Modeling the critical point and its effects (brief overview and comments)
3. Fourth order baryon number susceptibility in the presence of a QCD critical point  
(DM et al *Phys.Rev.C* 103 (2021) 3, 034901)
4. Conclusions and outlook

# Search for the Critical Point – Theory

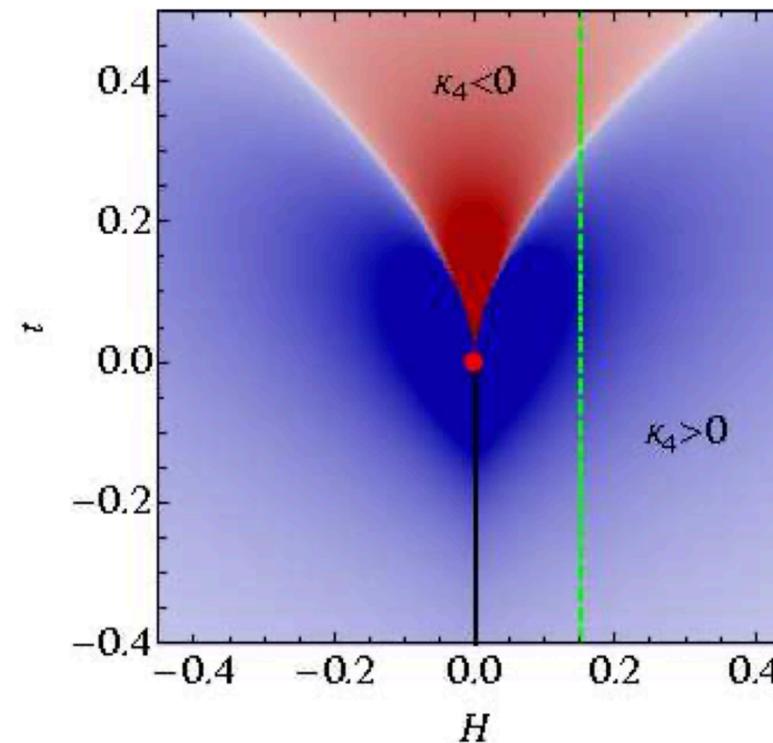
- ❖ Beam Energy Scan II at RHIC: looking for a critical point in the hadron-QGP transition.

M. A. Stephanov. PRL 107 (2011)

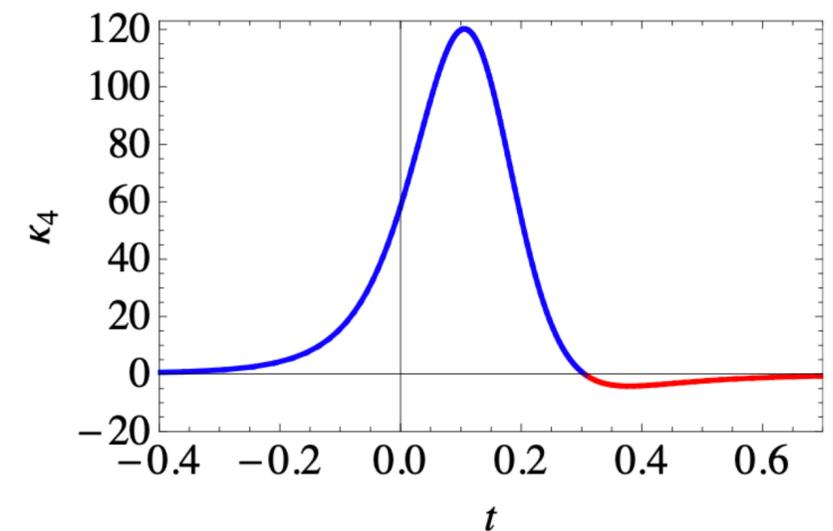
- QCD is in the same universality class as the 3D Ising model.
- Specific non-monotonic behavior of  $\chi_4^B$  as a function of  $\sqrt{s_{NN}}$ .

Candidate: **higher order cumulants**

1. Baryon number susceptibilities diverge at the critical point
2. Higher order  $\sim$  higher powers of correlation length



$$\chi_4^{\text{Ising}}(r, h) = \left( \frac{\partial^4 G}{\partial h^4} \right)_r = \left( \frac{\partial^3 M}{\partial h^3} \right)_r$$



# Search for the Critical Point – Experiment

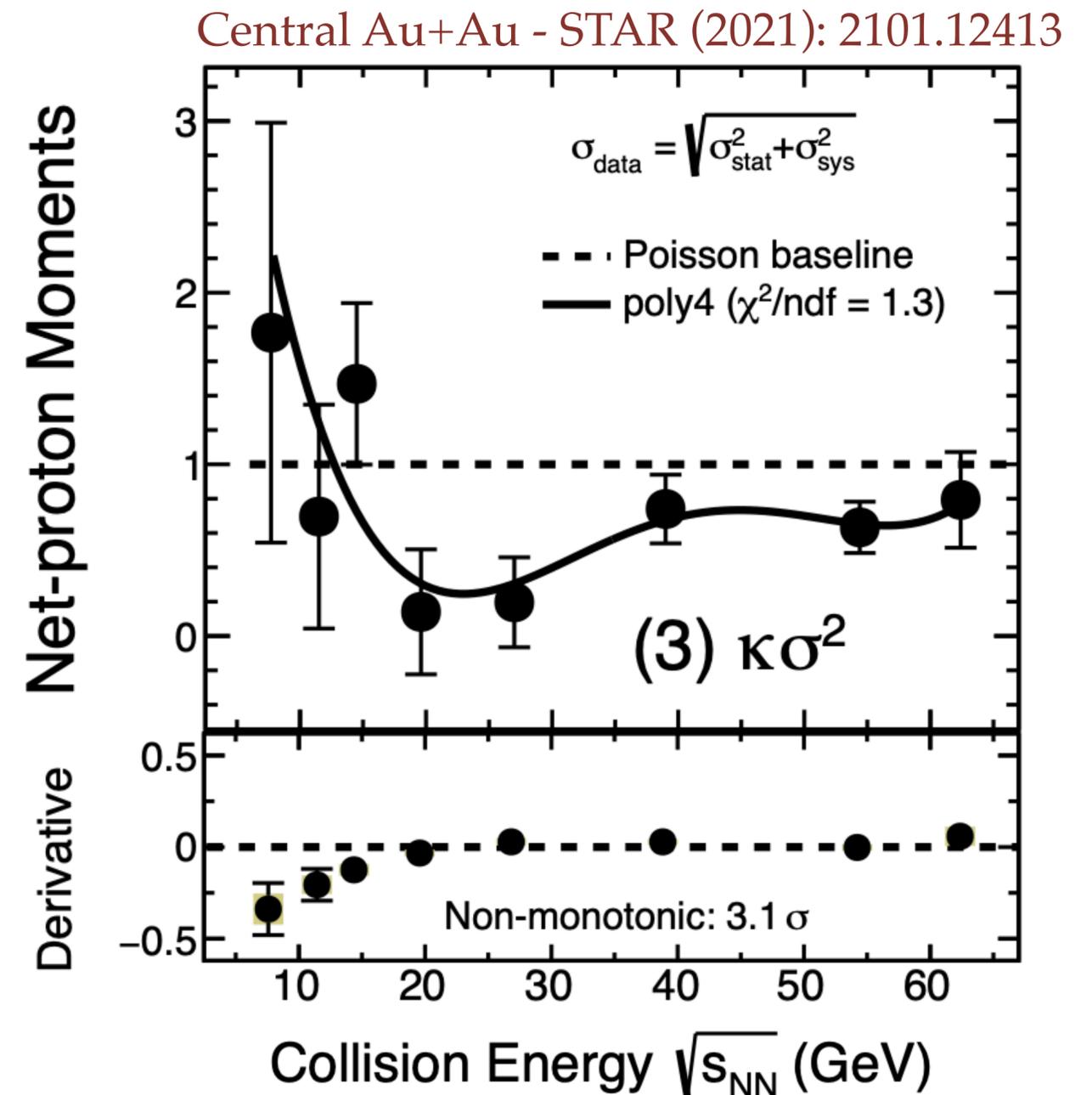
Experimental proxy for  $\chi_n^B \rightarrow$  net-proton moments

BES-I: Non-monotonic behavior of  $3.1\sigma$  significance, a.k.a “the dip”

BES-II (ongoing): Improved statistics  $\rightarrow$  smaller error bars

Is the dip followed by a diverging peak?

Theory: How do we interpret BES results?

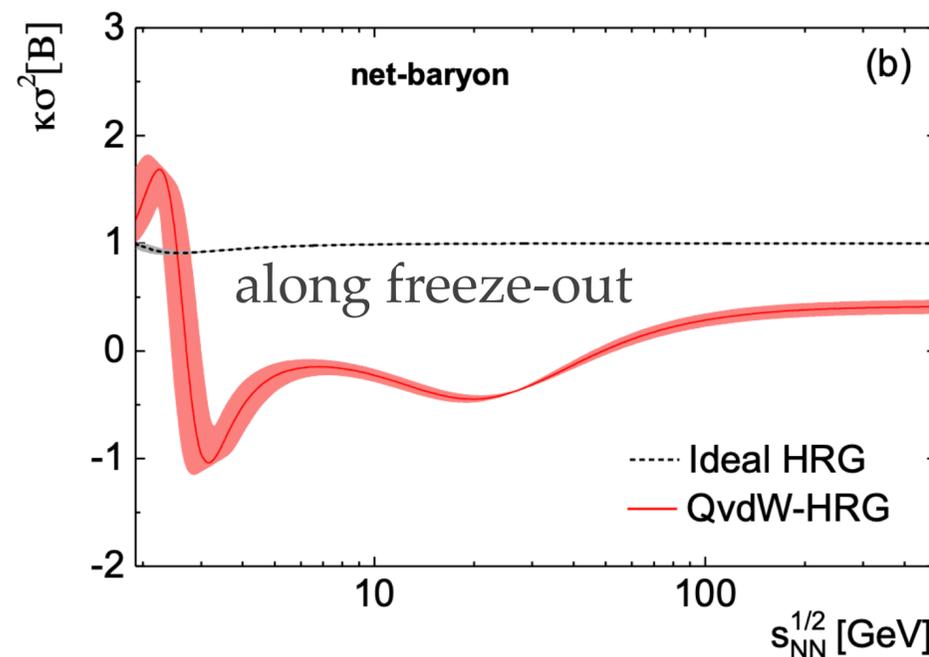
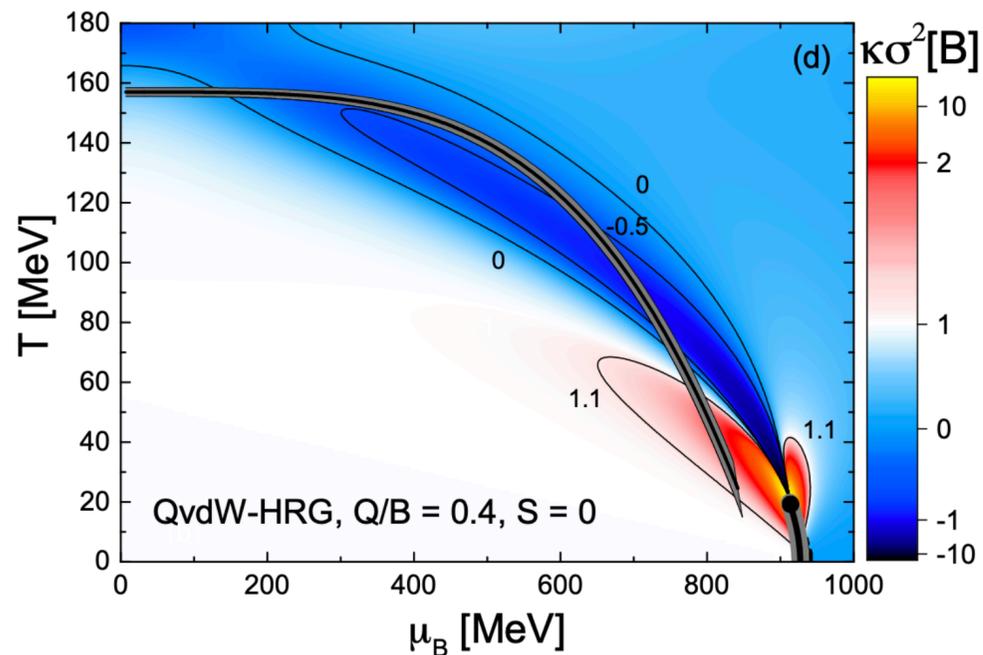


# Finite Density EoS for QCD with a CP

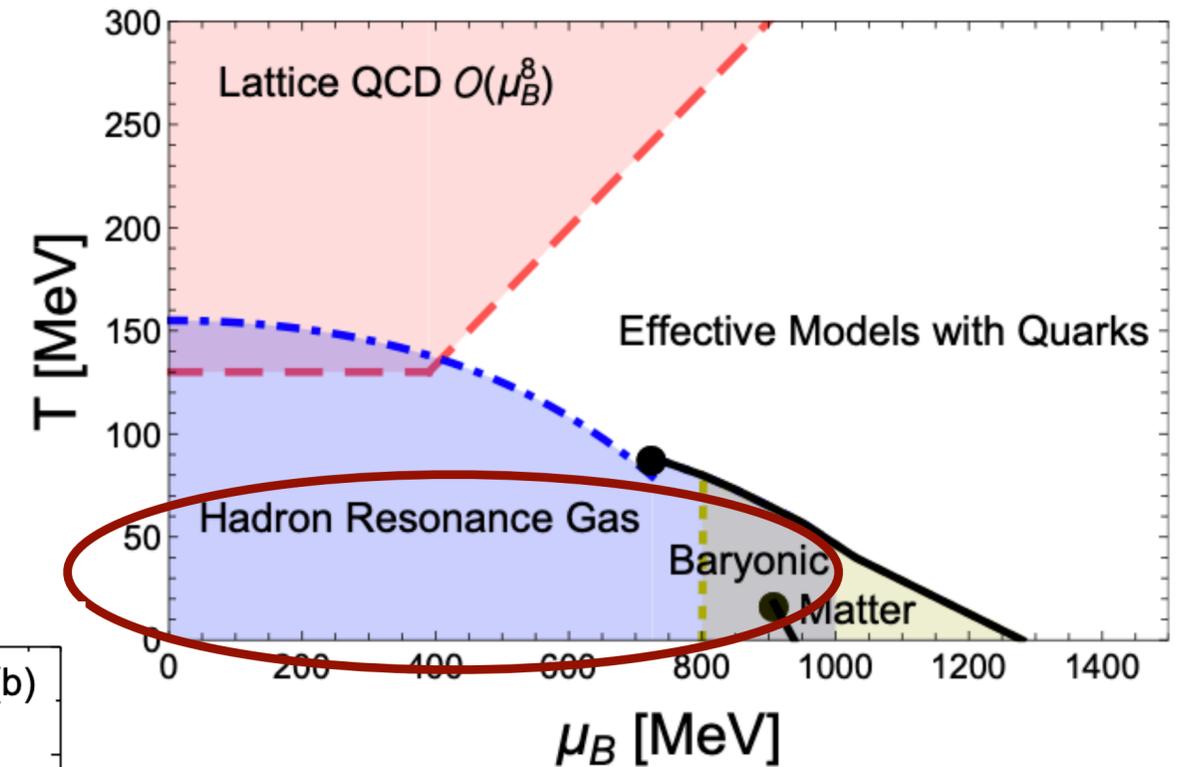
## I) HRG with vdW interactions + quantum statistics

- CP at high densities (liquid-gas)
- Impose charge conservation constraints
- Effect of global conservation laws

Talk: Vovchenko



V Dexheimer et al *J. Phys. G: Nucl. Part. Phys.* 48 (2021)



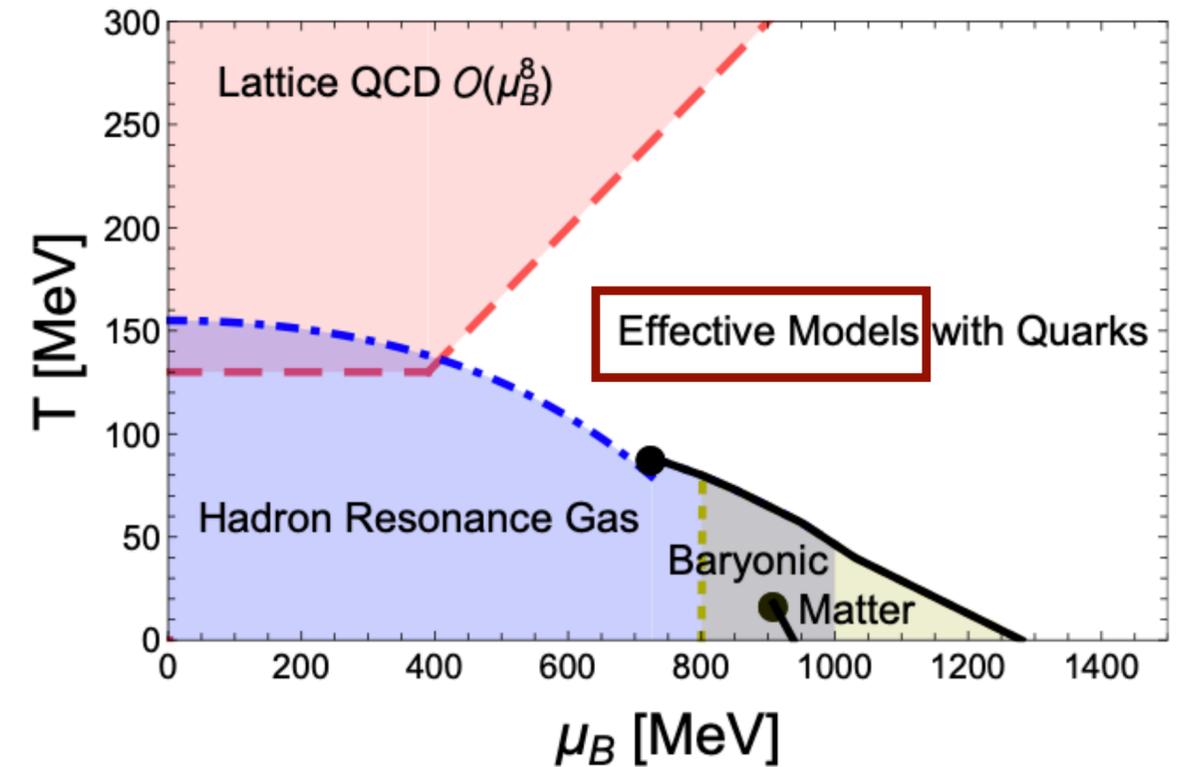
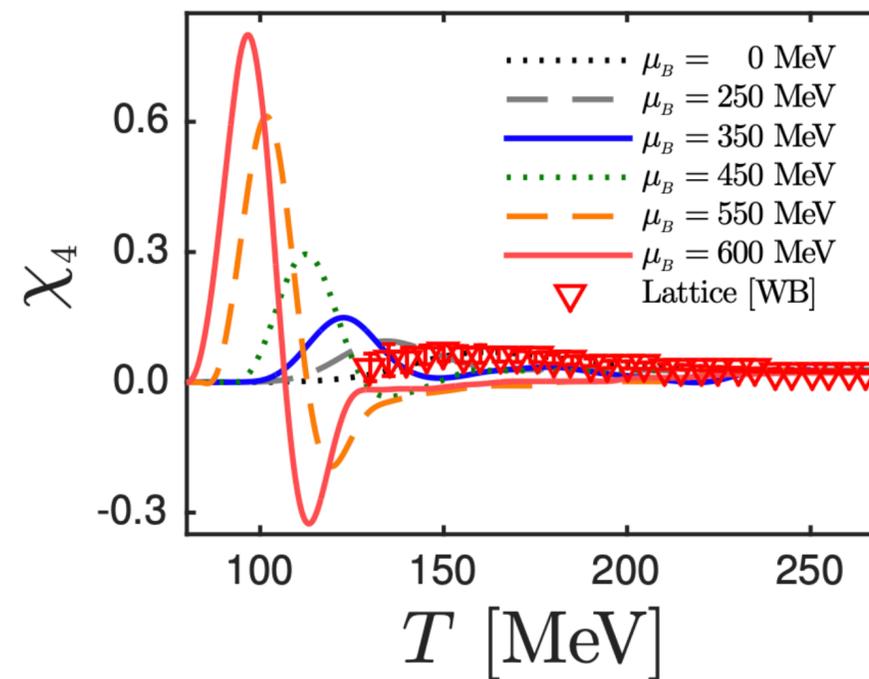
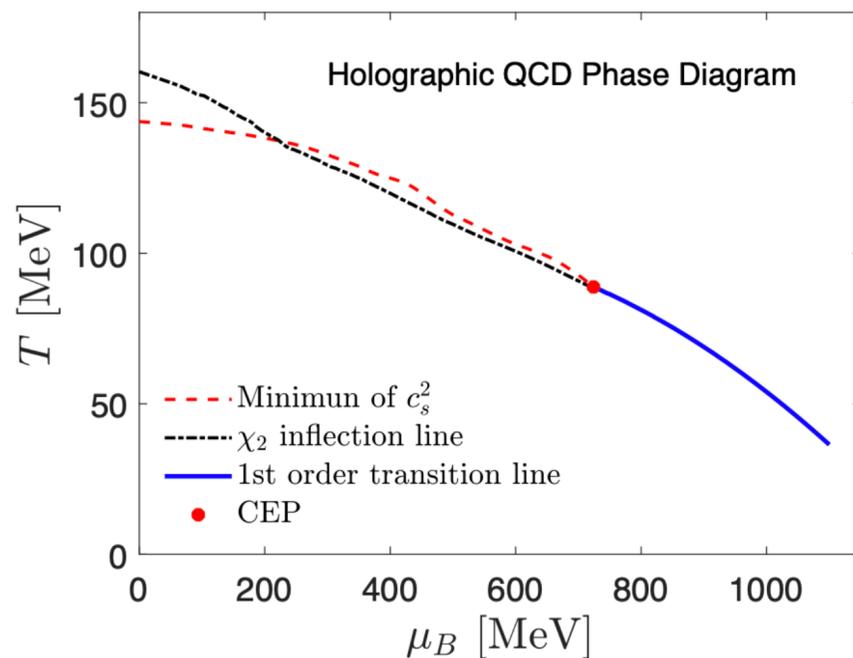
R. V. Poberezhnyuk, V. Vovchenko et al  
*Phys. Rev. C* 100, 054904

# Finite Density EoS for QCD with a CP

V Dexheimer *et al* 2021 *J. Phys. G: Nucl. Part. Phys.* **48** 073001

## II) Black hole holography

- CP at high densities:  $T_{CEP} = 89 \text{ MeV}$ ,  $\mu_B^{CEP} = 724 \text{ MeV}$
- Large coverage of phase diagram
- Agrees with LQCD



J. Grefa *et al.* 2102.12042

R. Critelli *et al* *Phys. Rev. D* **96**, 096026

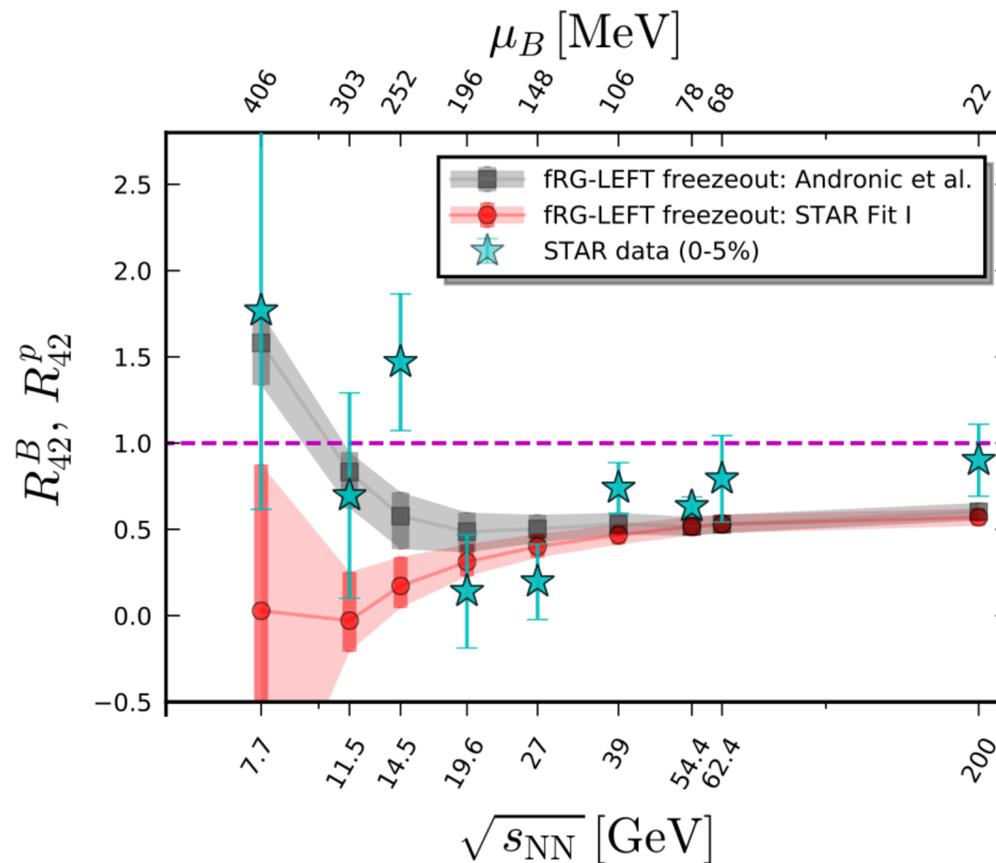
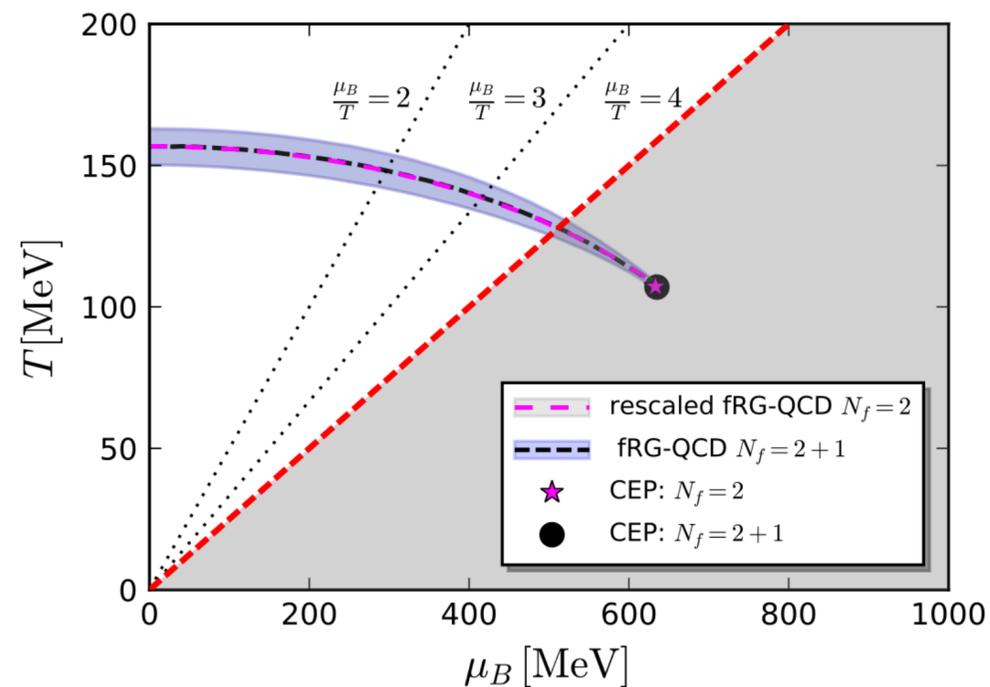
# Finite Density EoS for QCD with a CP

## III) Functional methods

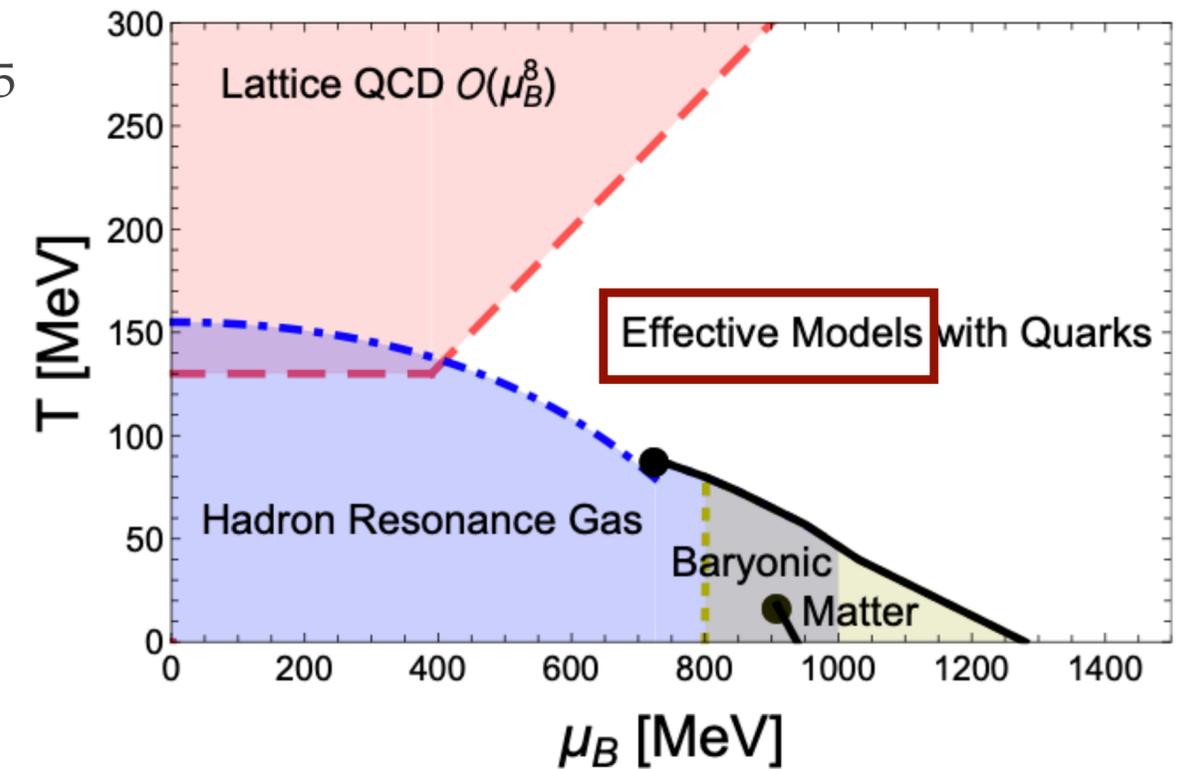
$$N_f = 2+1$$

F. Gao, JM. Pawłowski. 2010.13705

- CP at high densities:  $T_{CEP} = 109 \text{ MeV}, \mu_B^{CEP} = 610 \text{ MeV}$
- Large coverage of phase diagram
- Agrees with LQCD



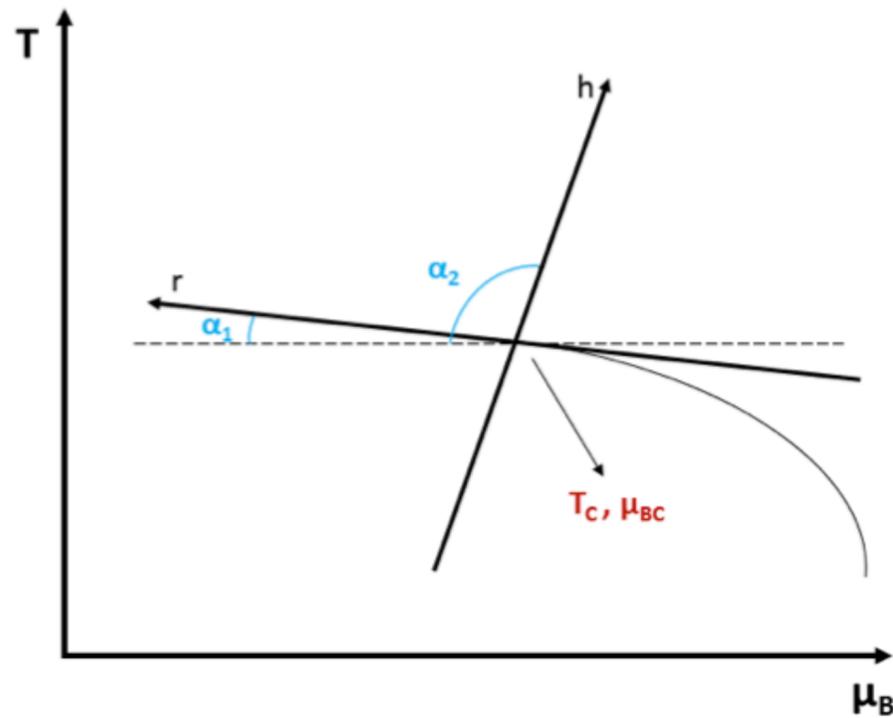
V Dexheimer et al 2021 J. Phys. G: Nucl. Part. Phys. 48 073001



W. Fu et al 2101.06035

# Finite Density EoS for QCD with a CP

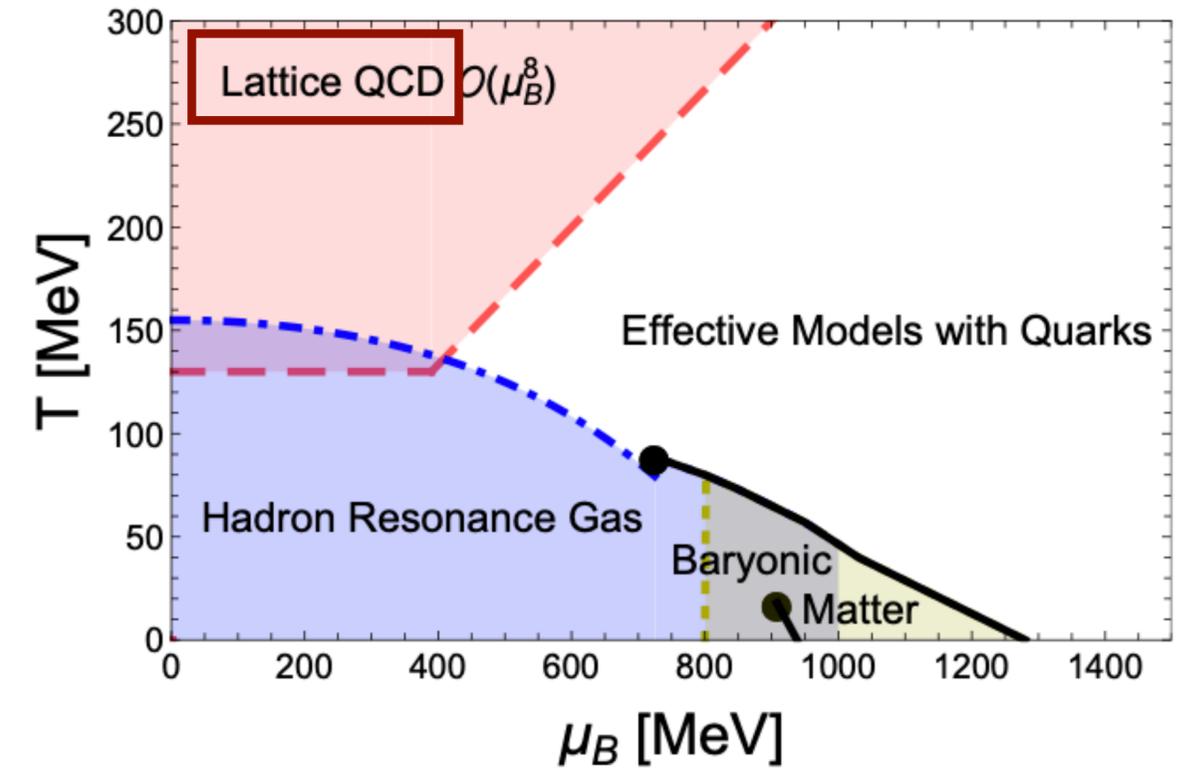
## IV) CP from 3D Ising Universality class (BEST EoS)



Up to  $\mathcal{O}(\mu_B^4)$ :

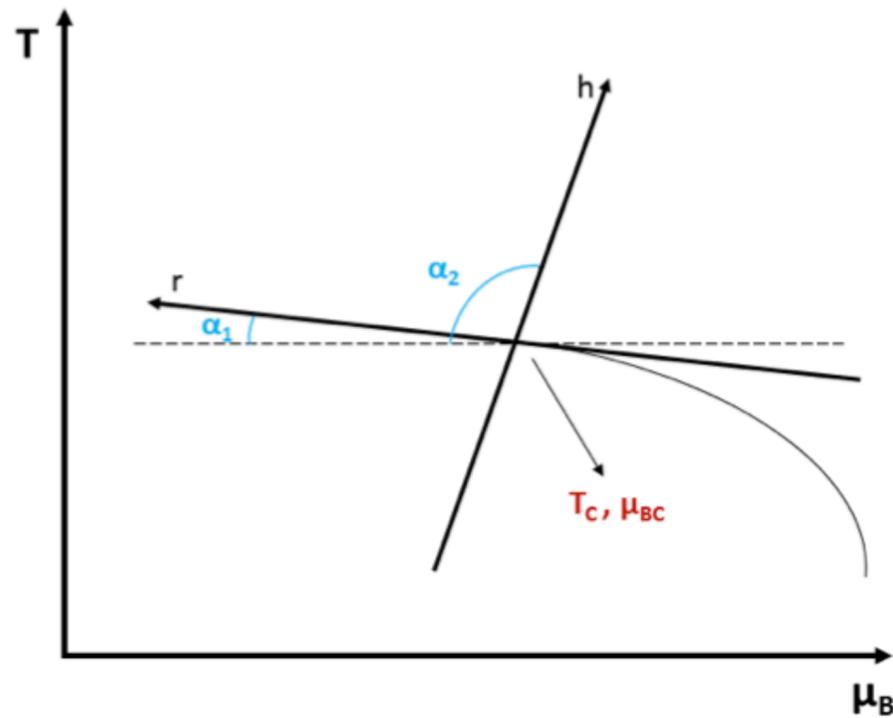
P. Parotto, DM, et al Phys. Rev C (2020)

V Dexheimer et al 2021 J. Phys. G: Nucl. Part. Phys. **48** 073001



# Parameterized EoS with CP from 3D Ising

## IV) CP from 3D Ising Universality class (BEST EoS)



Up to  $\mathcal{O}(\mu_B^4)$ :

P. Parotto, DM, et al Phys. Rev C (2020)

### Critical Contribution

1. Define parameterization of the 3D Ising Model in the vicinity of the critical point:

$$M = M_0 R^\beta \theta$$

$$h = h_0 R^{\beta\delta} \tilde{h}(\theta)$$

$$r = R(1 - \theta^2)$$

$$M_0 \simeq 0.605 \quad h_0 \simeq 0.364$$

$$\tilde{h}(\theta) = \theta(1 + a\theta^2 + b\theta^4)$$

$$\beta \simeq 0.326, \quad \delta \simeq 4.80$$

C. Nonaka, M. Asakawa, Phys. Rev C (2005)

2. Map the 3D-Ising phase diagram to QCD variables:

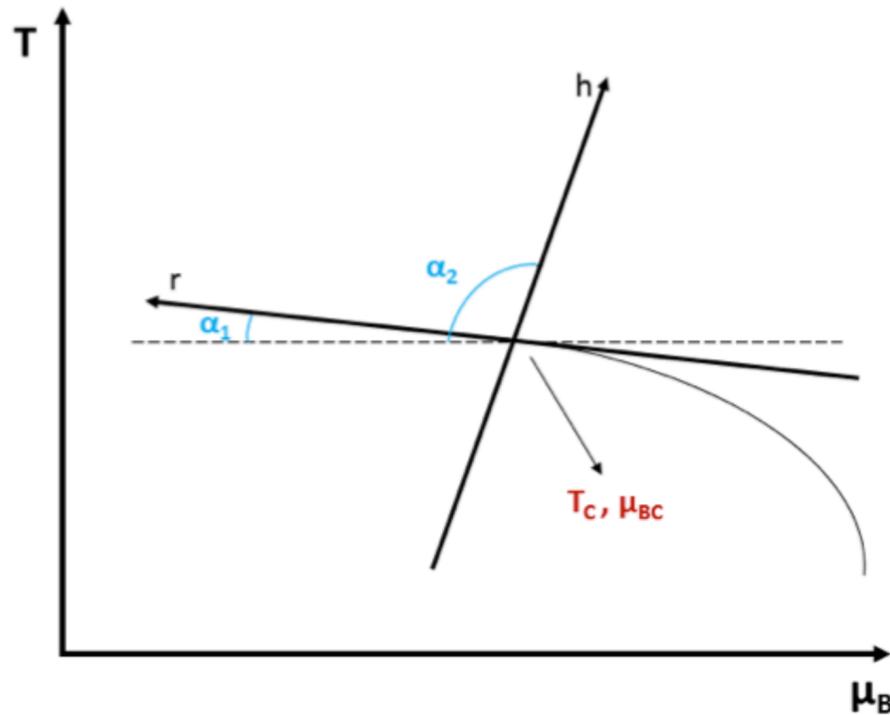
$$(\mathbf{r}, \mathbf{h}) \longleftrightarrow (\mathbf{T}, \mu_B) : \frac{T - T_C}{T_C} = \mathbf{w} (r \rho \sin \alpha_1 + h \sin \alpha_2)$$

$$\frac{\mu_B - \mu_{BC}}{T_C} = \mathbf{w} (-r \rho \cos \alpha_1 - h \cos \alpha_2)$$

# Parameterized EoS with CP from 3D Ising

## IV) CP from 3D Ising Universality class (BEST EoS)

### Non-Ising Contribution



Up to  $\mathcal{O}(\mu_B^4)$ :

P. Parotto, DM, et al Phys. Rev C (2020)

3. Impose matching to lattice QCD at  $\mu_B = 0$ :

$$T^4 c_n^{\text{LAT}}(T) = T^4 c_n^{\text{Non-Ising}}(T) + c_n^{\text{Ising}}(T)$$

4. Reconstruct the full pressure:

$$P(T, \mu_B) = T^4 \sum_n c_n^{\text{Non-Ising}}(T) \left( \frac{\mu_B}{T} \right)^n + P_{\text{crit}}^{\text{QCD}}(T, \mu_B)$$

- Reduce number of free parameters by imposing:

$$T = T_0 + \kappa T_0 \left( \frac{\mu_B}{T_0} \right)^2 + \mathcal{O}(\mu_B^4), \quad \alpha_1 = \tan^{-1} \left( 2 \frac{\kappa}{T_0} \mu_{BC} \right)$$

# Size and shape of Ising Contribution

❖ Dependence on mapping parameters?

We can estimate the size of the critical region along the crossover,  $h=0$ :

$$\chi_4^{\text{Ising}} \sim AG_{\mu\mu\mu\mu}(r,0) \sim AG_{hhhh}(r,0)h_\mu^4$$

Along the critical chemical potential line,  $r \sim 0$ :

$$\chi_4^{\text{Ising}} \sim AG_{\mu\mu\mu\mu}(0,h) \sim AG_{hhhh}(0,h)h_\mu^4$$

Comparing to the regular contribution  $\chi_4^{\text{reg}} \sim 1$ , we find

$$\Delta\mu_B \sim T_C \rho w \cos \alpha_1 \left( \frac{A^{1/4}}{T_C} \frac{\sin \alpha_1}{w T_C \sin \alpha_1 - \alpha_2} \right)^{\frac{4}{\beta(3\delta-1)}}$$

$$\Delta T \sim T_C \left( \frac{A}{T_C^4} \right)^{\frac{\delta}{3\delta-1}} \frac{\sin \alpha_1}{\cos \alpha_1} \left( \frac{\sin \alpha_1}{w \sin \alpha_1 - \alpha_2} \right)^{\frac{\delta+1}{3\delta-1}}$$

**Weak dependence on  $w$ , while smaller  $\alpha_1 - \alpha_2$  yields a larger critical region for the same  $w, \rho$ :**

$$\Delta\mu_B \sim w^{1/7}, \quad \Delta T \sim w^{-3/7}$$

$$\Delta\mu_B \sim \sin \alpha_1 - \alpha_2^{-6/7}, \quad \Delta T \sim \sin \alpha_1 - \alpha_2^{-3/7}$$

# Parameter choices

- ❖ Two sets of parameters

	$\mu_{BC}$	$T_C$	$\alpha_1$	$\alpha_2 - \alpha_1$	$w$	$\rho$
<b>I.</b>	420 MeV	138 MeV	$4.6^\circ$	$90^\circ$	0.5, 1, 2	0.5, 1, 2
<b>II.</b>	420 MeV	138 MeV	$4.6^\circ$	$-3^\circ$	0.5, 1, 2	0.5, 1, 2

Common choice in literature. Orthogonal axes.

Motivated by M.S. Pradeep, M. Stephanov,  
Phys. Rev. D 100 (2019)

- ❖ Large  $\mu_{BC}$  to allow for maximum freedom in parameter choice within the range of the Taylor expansion.
- ❖  $T_C$  and  $\alpha_1$  not free  $\rightarrow$  follow from choice of  $\mu_{BC}$ .
- ❖ Compute critical contribution to  $\chi_4^B$  for each choice of parameters.

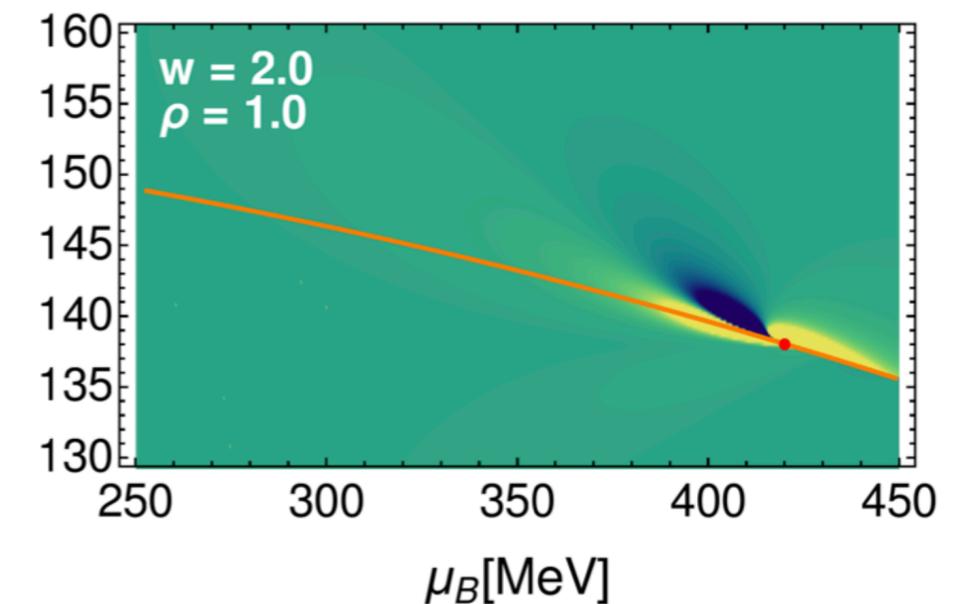
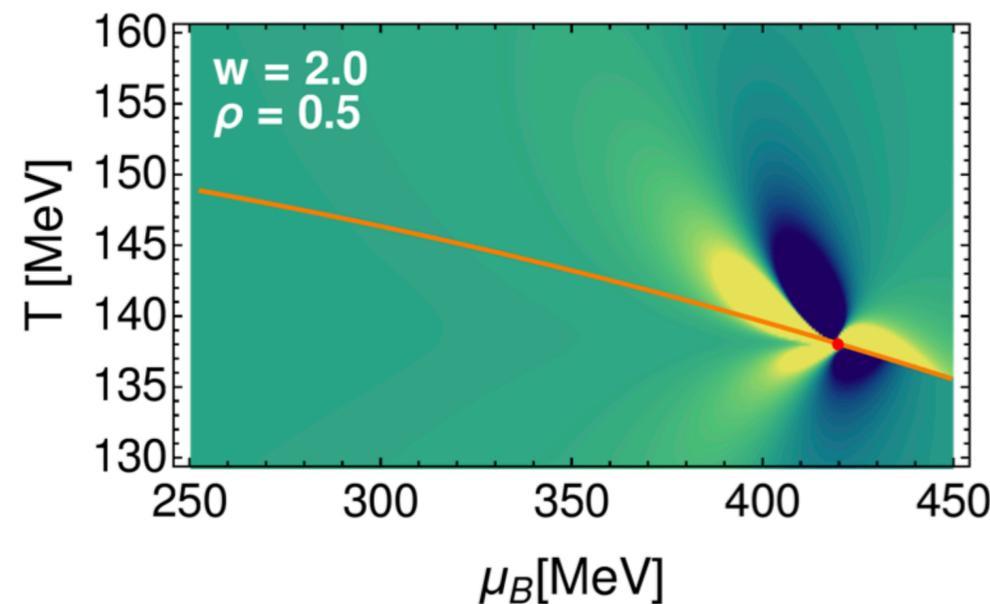
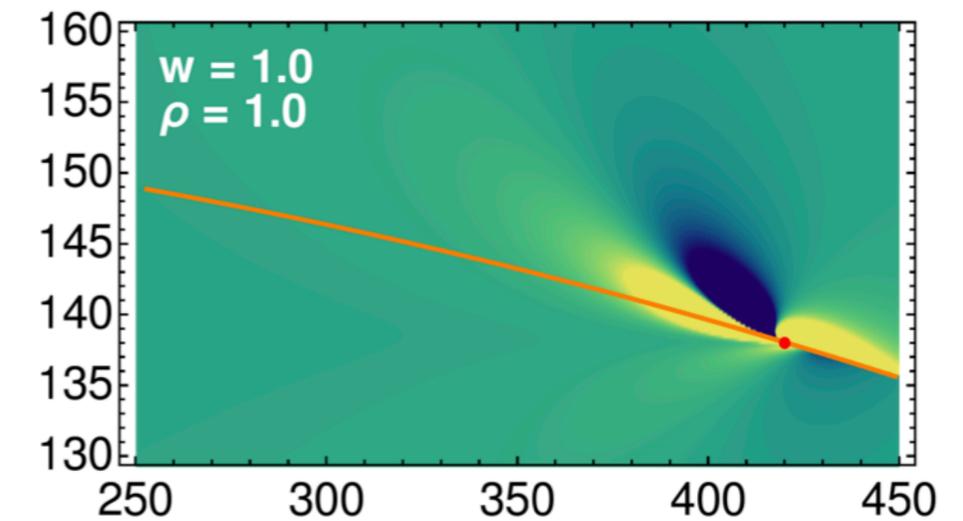
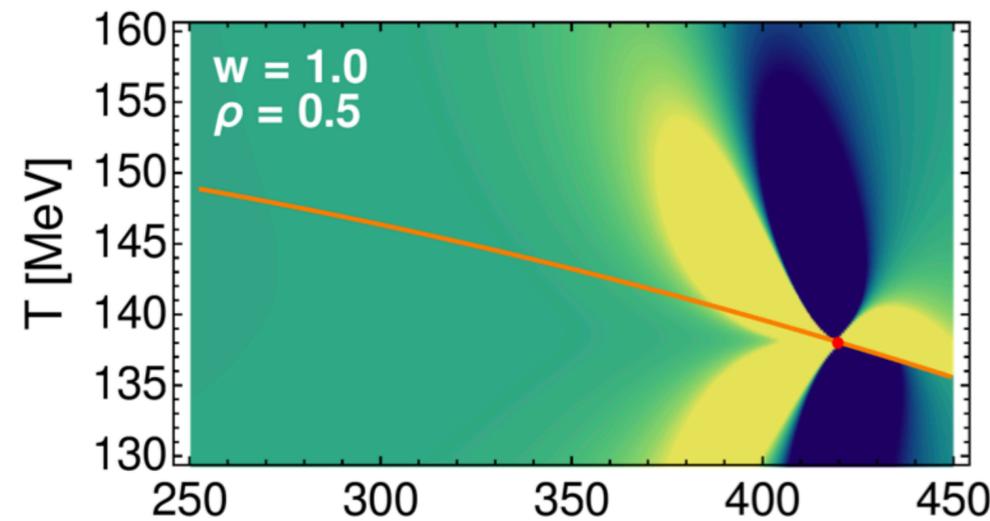
# Critical contribution to $\chi_4^B$

$$(\alpha_2 - \alpha_1 = 90^\circ)$$

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green  $\rightarrow$  yellow: positive values, blue:  
negative.  
orange curve: parameterized transition  
line.

- ❖ Smaller  $w$  yields larger critical region.
- ❖  $\rho$  stretches critical region along  $\mu_B$
- ❖ Approach to critical point characterized by peak rather than dip, except for immediate vicinity.



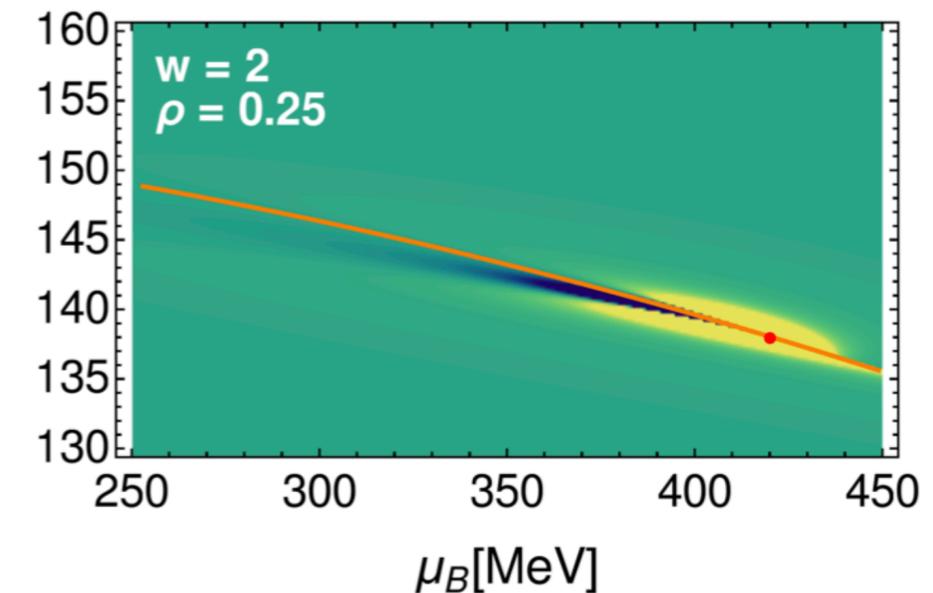
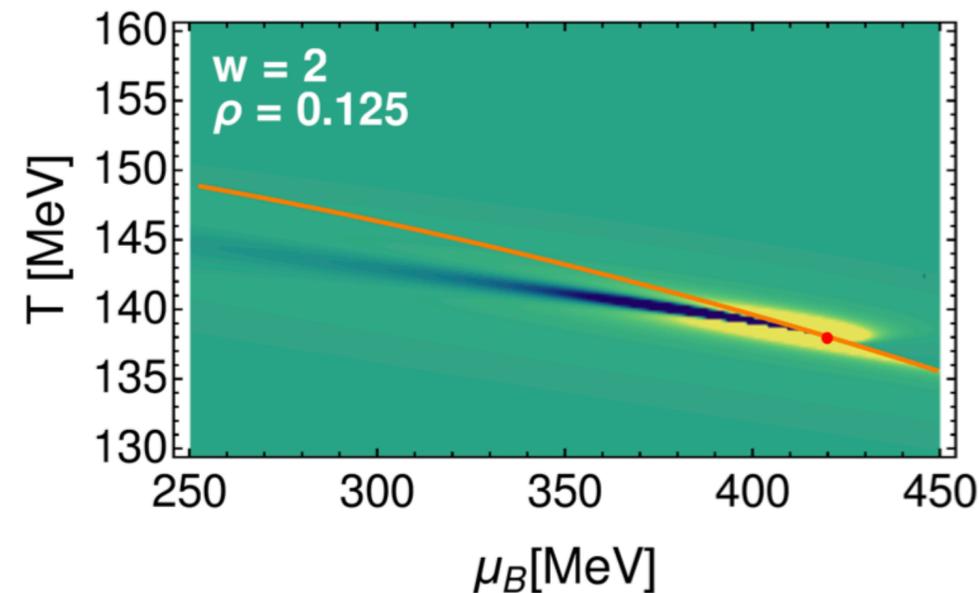
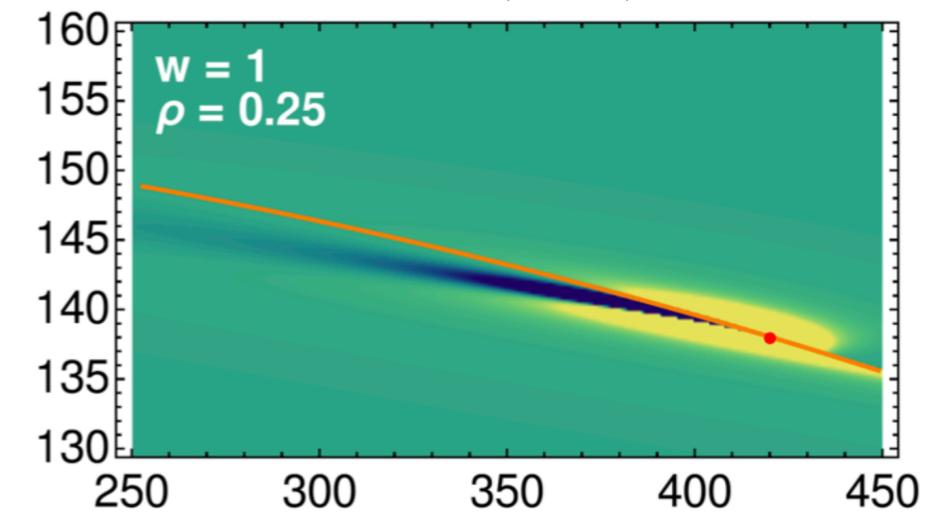
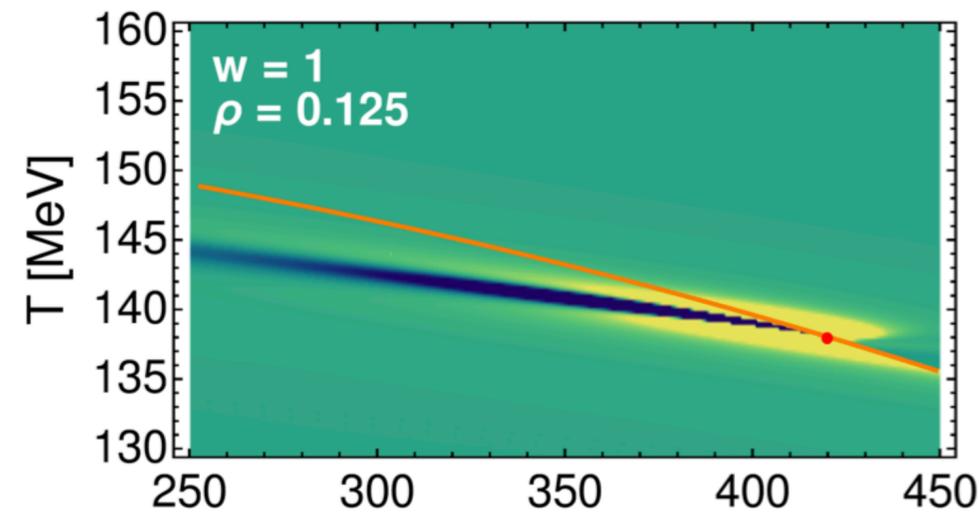
# Critical contribution to $\chi_4^B$

$$(\alpha_2 - \alpha_1 = -3^\circ)$$

green  $\rightarrow$  yellow: positive values, blue:  
negative.  
orange curve: parameterized transition  
line.

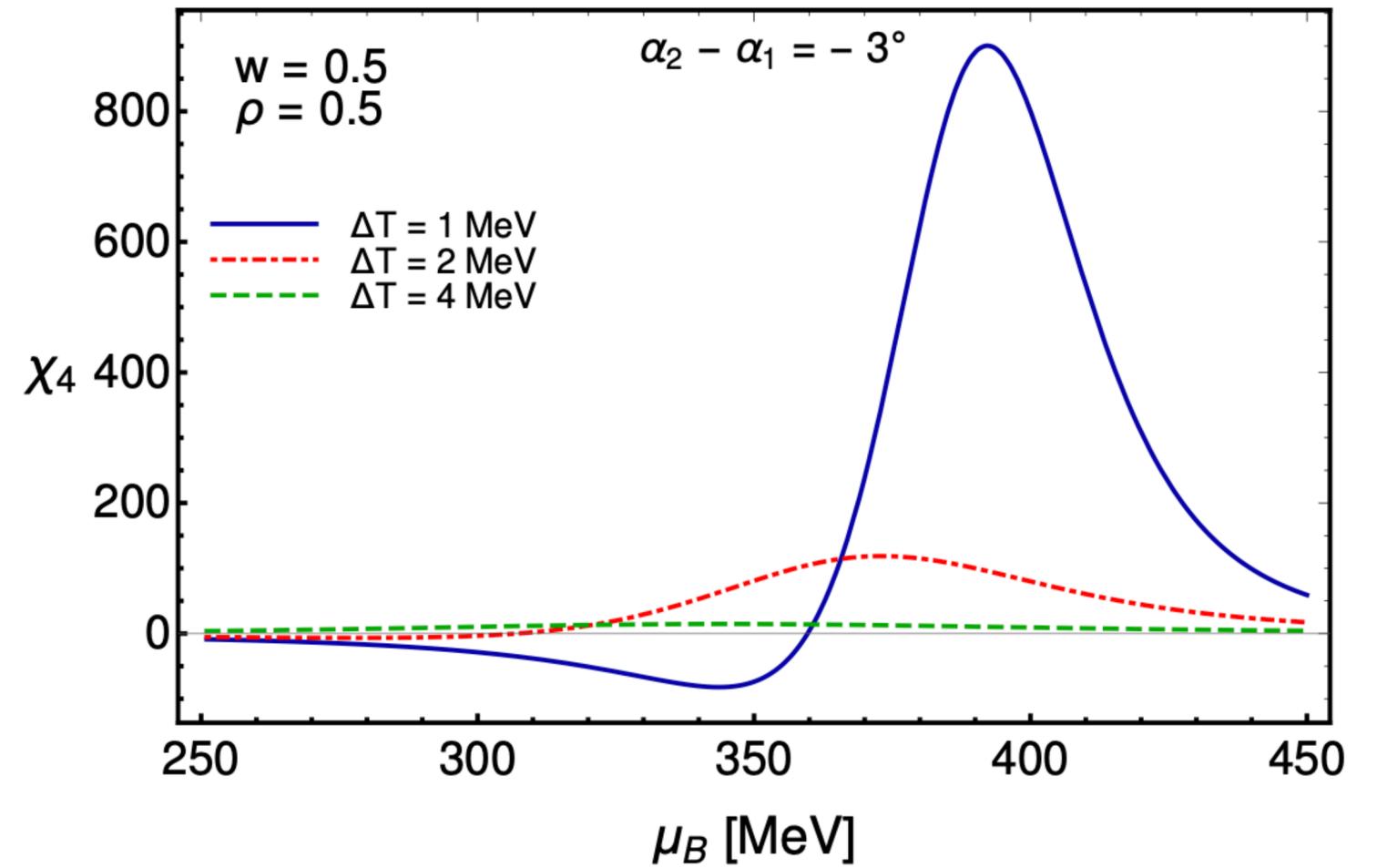
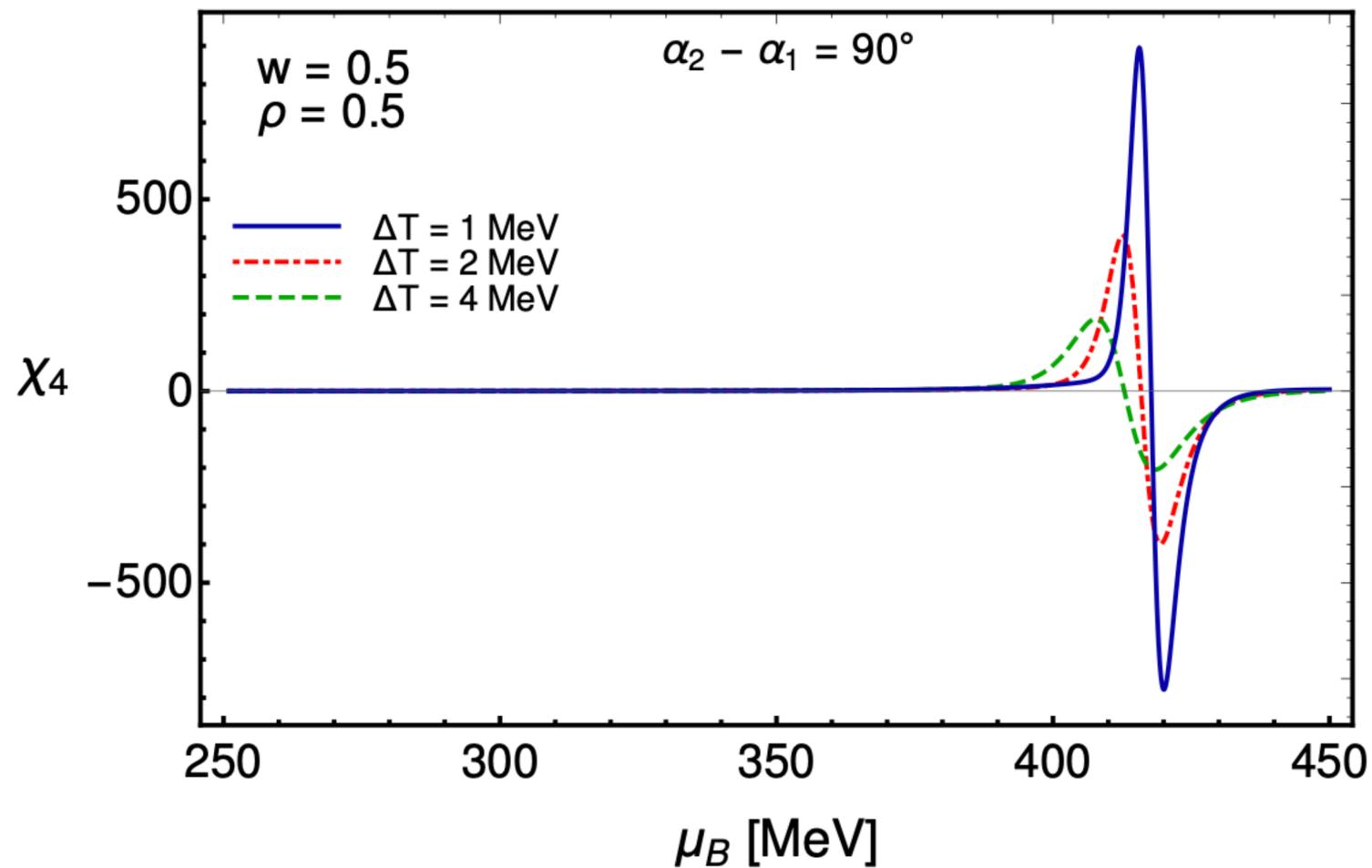
- ❖ **Negative angle difference bends negative lobe downwards.**
- ❖ **Agrees with leading singularity prediction.**
- ❖ **Note: this choice is not thermodynamically stable under current mapping.**

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# Peak sign of CP, not the dip

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**Dip only appears for one choice\*, close to the vicinity of the transition line.**

\*shown in this work

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# Why the difference?

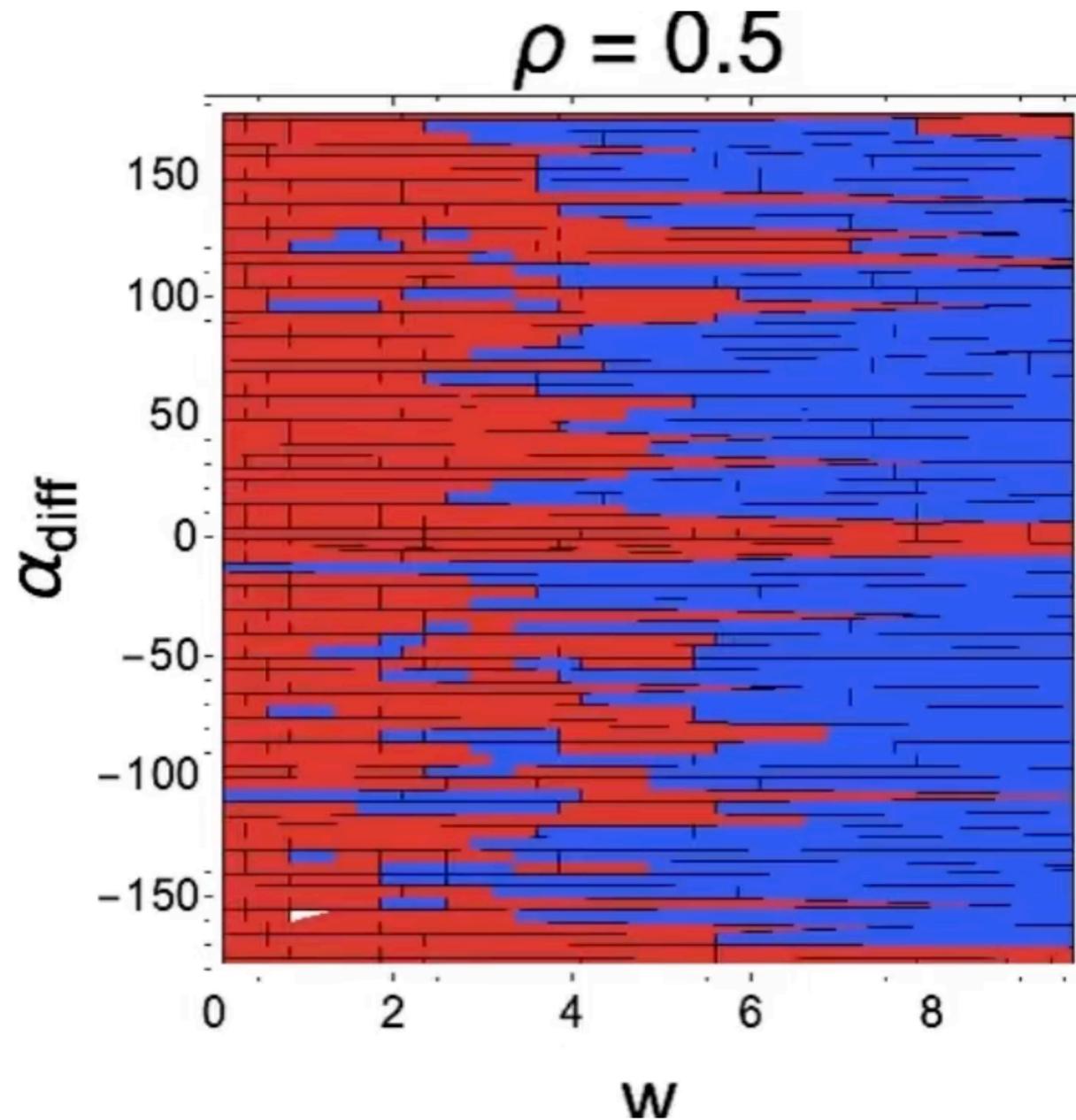
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- ❖ Ising to QCD map introduces mixing between  $r$ ,  $h$ :

$$\partial_{\mu_B} = \frac{1}{w\rho T_C \sin \alpha_1 - \alpha_2} (\sin \alpha_1 \partial_h + \sin \alpha_2 \partial_r)$$

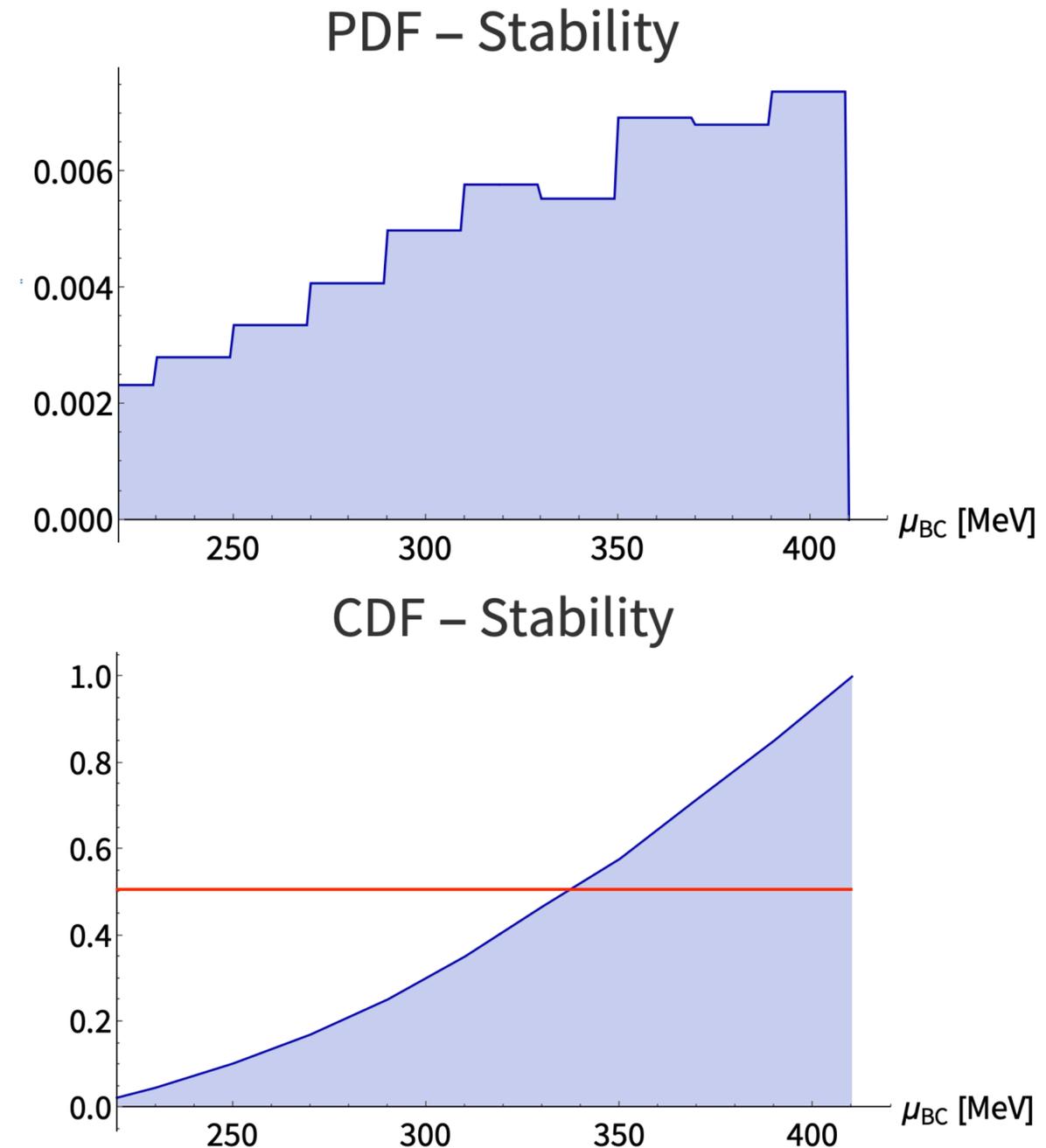
- $h$  has larger scaling dimension: dominant contribution close to the critical point.
- Since  $\alpha_1$  is small, when is  $\alpha_2$  not small  $\rightarrow$   $h$  contribution becomes suppressed (orthogonal case).
- Taking most divergent terms corresponds to  $\partial_{\mu_B} \sim \partial_h$
- **but** subleading terms may dominate if leading term is sufficiently suppressed.

# Parameter space scan (preliminary)



- Introduction of free-parameters with Ising  $\rightarrow$  QCD map allows for thermodynamically unstable realizations.
- ML classifier assisted classification of entire parameter space.

Preliminary results show:  
 i) Larger  $\mu_{BC}$  is preferred  
 ii) Hard upper limit on  $\rho$   
 iii) Small angles strongly disfavored



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# Conclusions, considerations, outlook

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- ❖ Investigated behavior of  $\chi_4^B$  in the presence of the critical point in the 3D Ising model universality class.
- ❖  $\chi_4^B$  can be affected by sub-leading terms.
- ❖ Diverging peak is a robust signature of the critical point.
- ❖ Dip only present in very few, thermodynamically unstable cases under current mapping.
- ❖ Current study: **equilibrium** properties of QCD EoS. HIC are dynamical systems, **need EbE relativistic viscous BSQ hydro evolution + critical fluctuations + hadronic transport.**
- ❖ Temperature difference between hadronization and freeze-out?
- ❖ Machine-learning-assisted study coming soon!