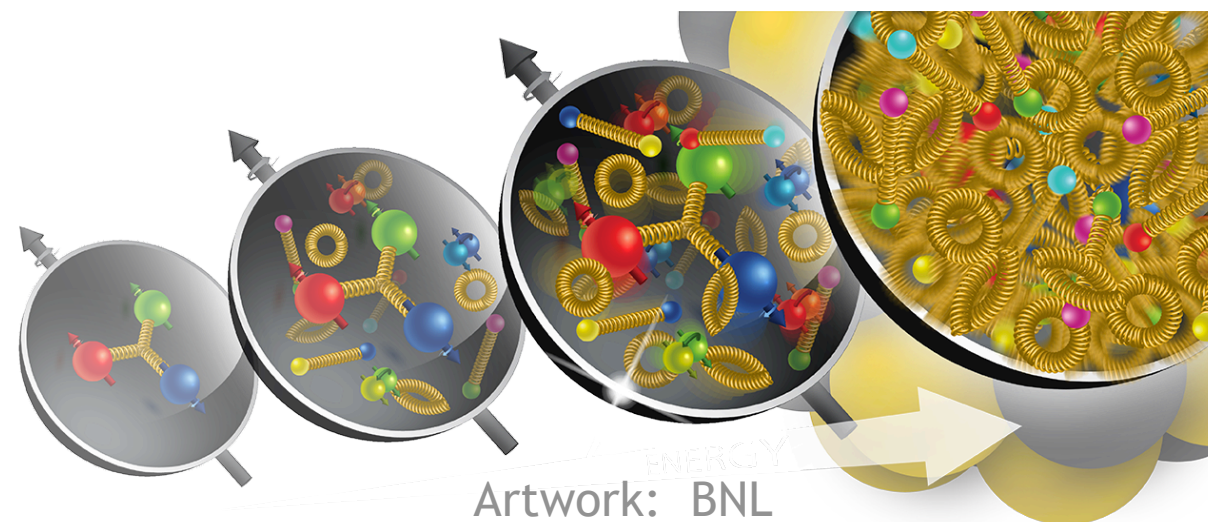


Gluon Saturation at the Electron-Ion Collider

RHIC/AGS Users Meeting
June 8th, 2021
Farid Salazar

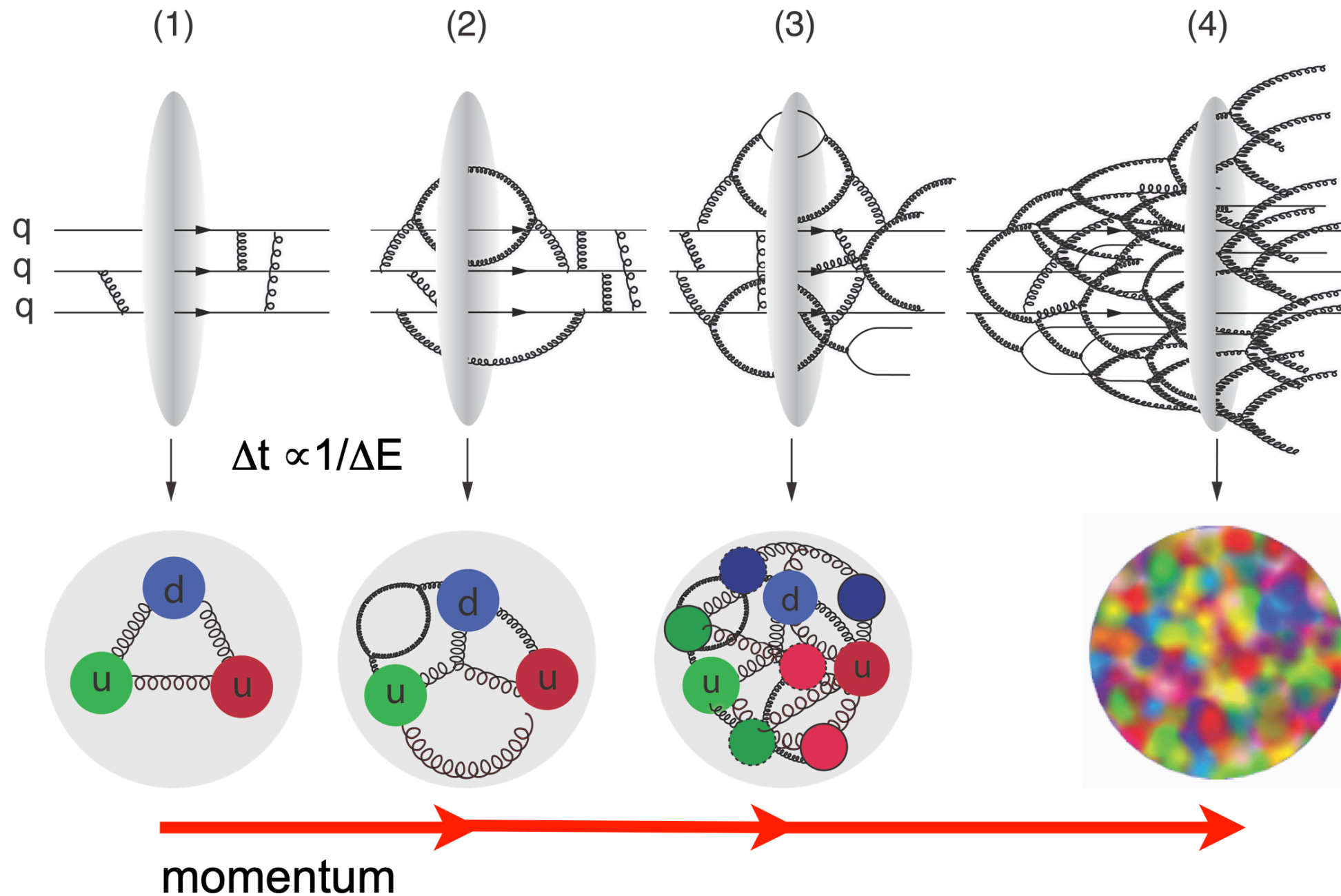
farid.salazarwong@stonybrook.edu



Outline

- Gluon saturation and the CGC
- Measurements at the EIC
 - Structure functions*
 - Diffraction processes*
 - Semi-inclusive measurements*
- Some recent theoretical developments
 - Matching CGC and TMD frameworks
 - Next-to-leading order computations

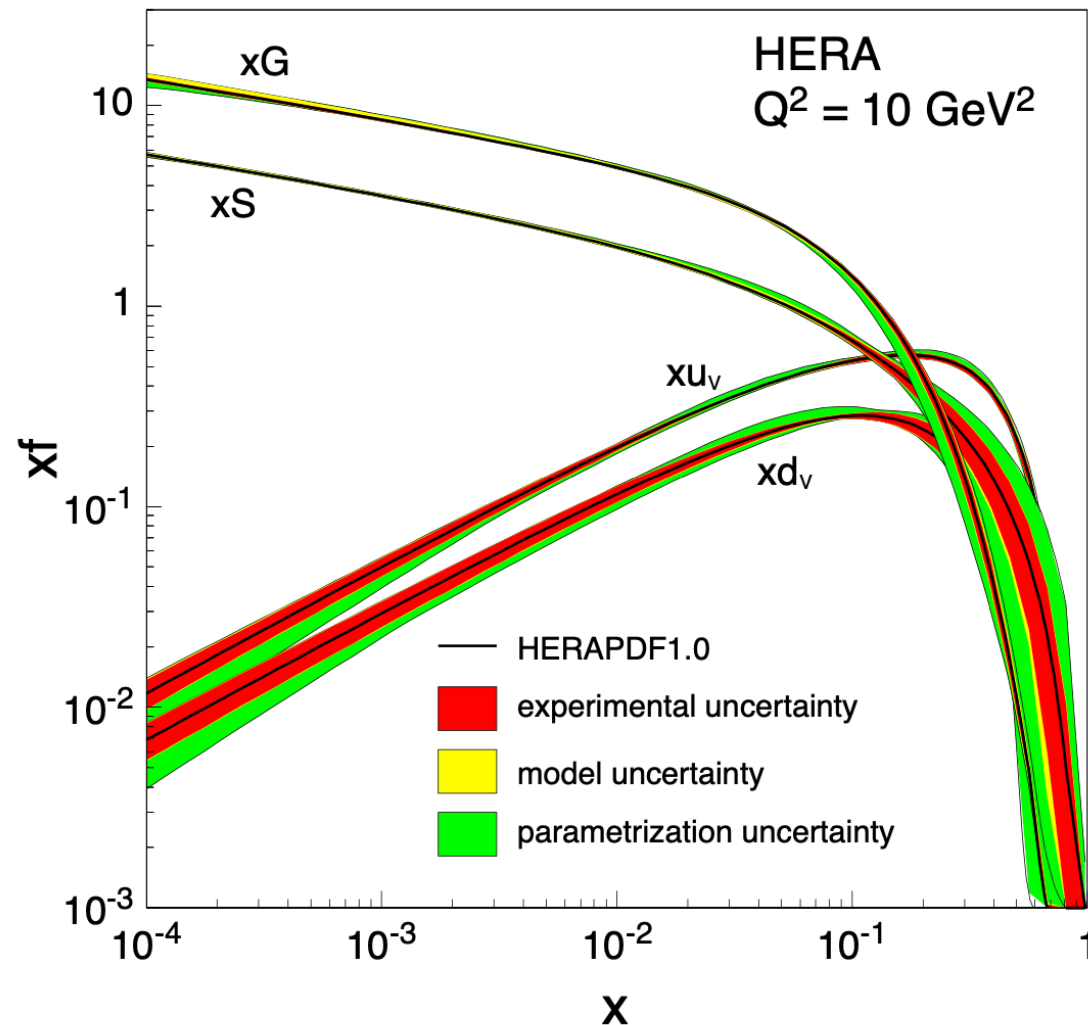
Gluon saturation



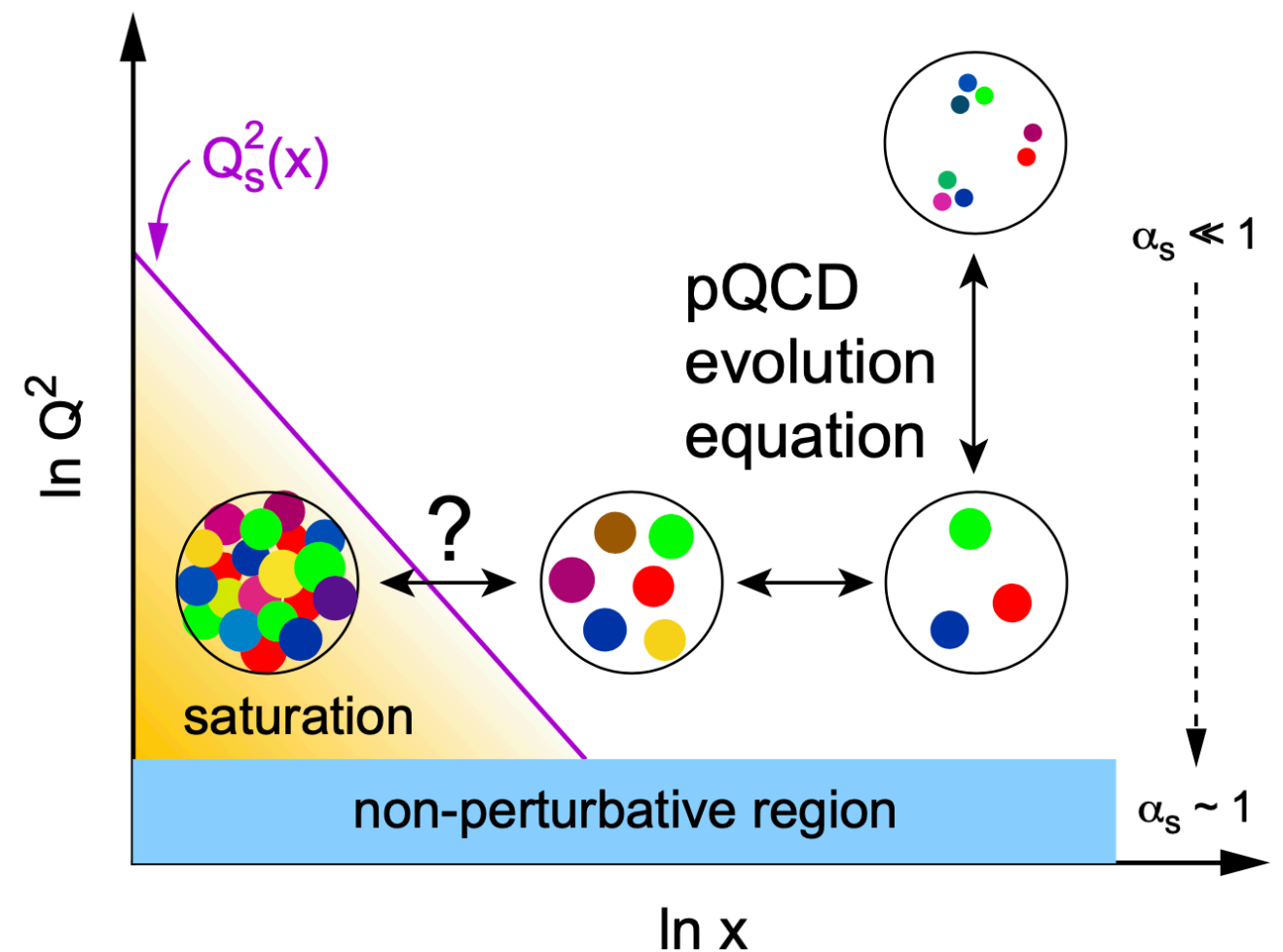
Artwork: T. Ullrich

Gluon Saturation

Rapid growth of gluon distributions at small x $\longrightarrow$ Violate unitarity



Non-linear effects in QCD at sufficiently small x (e.g. gluon recombination) tame the growth



Gluon saturation

Characterized by a scale Q_s^2 (saturation scale)

$$Q_s^2 \sim (A/x)^{1/3}$$

Gluon Saturation

The Color Glass Condensate

McLerran, Venugopalan.
[hep-ph/9309289](#), [hep-ph/9311205](#)

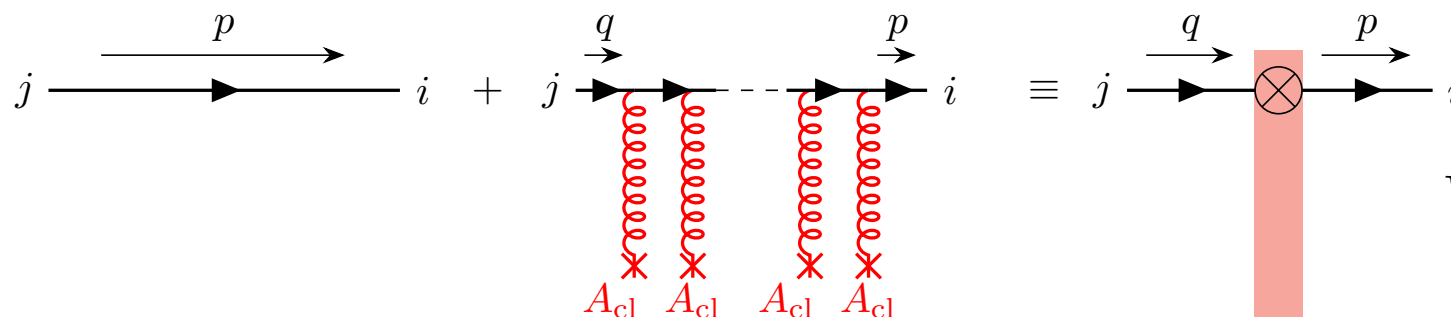
Large gluon densities at
small x

$$xG \sim \frac{1}{\alpha_s}$$

Characterize small x gluon as a classical
Yang-Mills field sourced by valence quarks

$$[D_\mu, F^{\mu\nu}] = J^\nu \longrightarrow A_{\text{cl}}^\mu$$

Computations in QCD +QED in the presence of color back-ground field A_{cl}^μ



Effective eikonal quark vertex

Light-like Wilson line

$$V_{ij}(\mathbf{x}) = P \exp \left\{ ig \int dx^- A_{cl}^{+,a}(\mathbf{x}, x^-) t^a \right\}$$

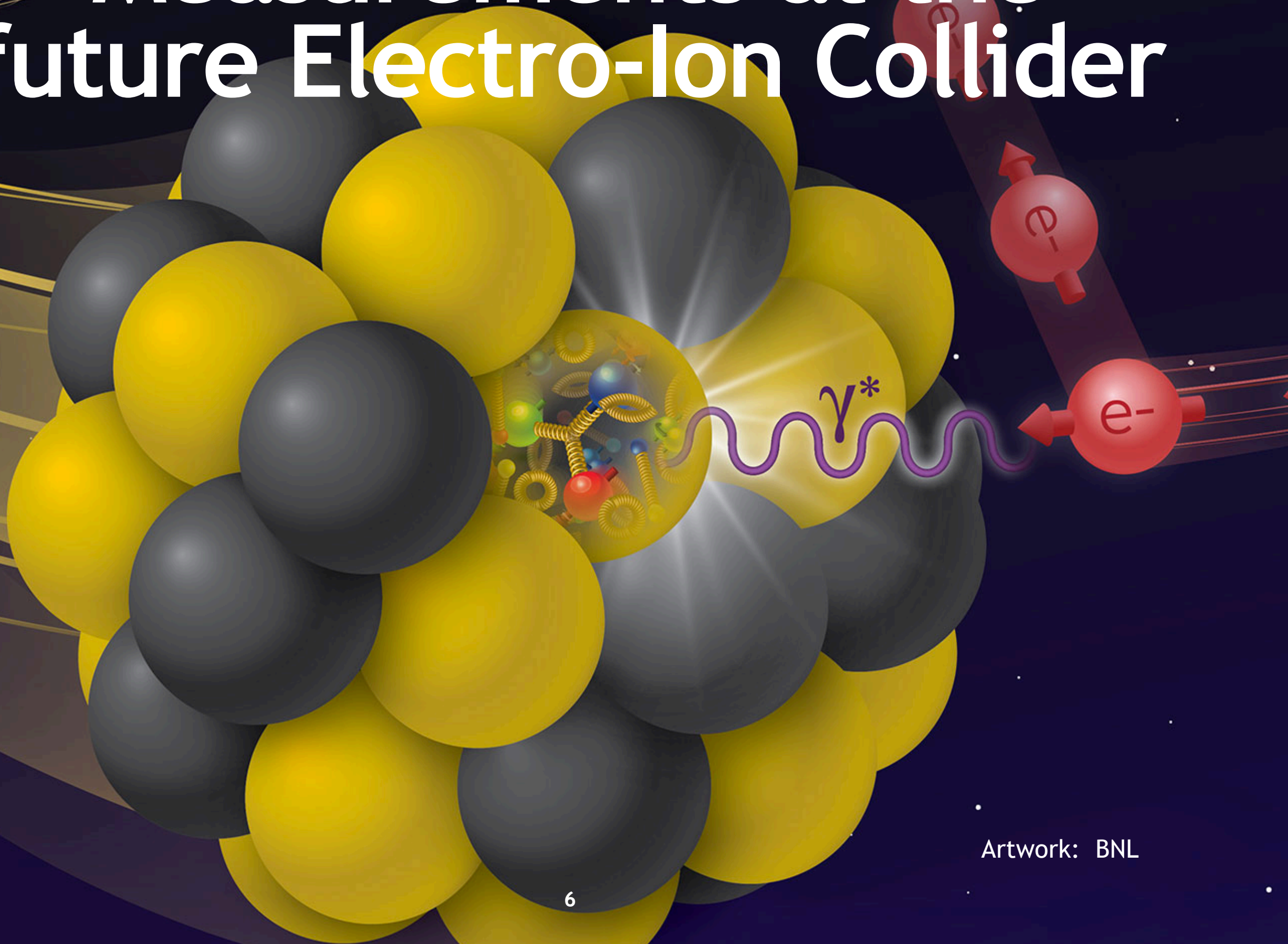
Observables depend on
correlators of Wilson Lines

$$\frac{1}{N_c} \langle \text{Tr}(V_{\mathbf{x}_{1\perp}} V_{\mathbf{x}_{2\perp}}^\dagger) \rangle_Y$$

$$\frac{1}{N_c} \langle \text{Tr}(V_{\mathbf{x}_{1\perp}} V_{\mathbf{x}_{2\perp}}^\dagger V_{\bar{\mathbf{x}}_{2\perp}} V_{\bar{\mathbf{x}}_{1\perp}}^\dagger) \rangle_Y$$

They obey RG equations (Balitsky-Kovchegov, JIMWLK)
and depend implicitly on saturation scale Q_s^2

Measurements at the future Electro-Ion Collider

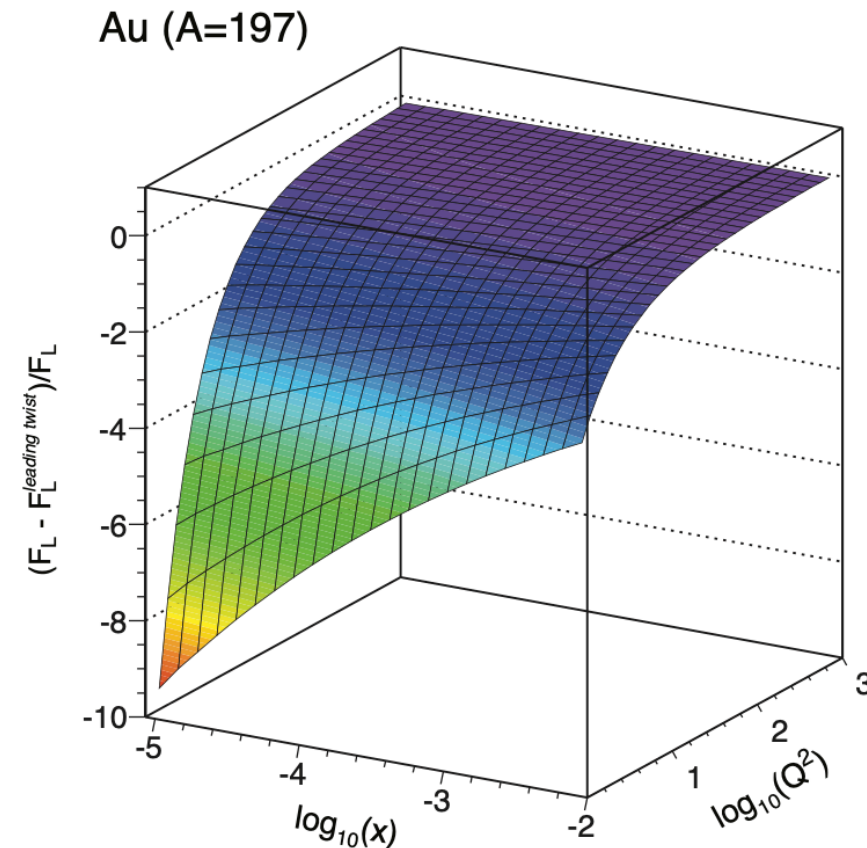
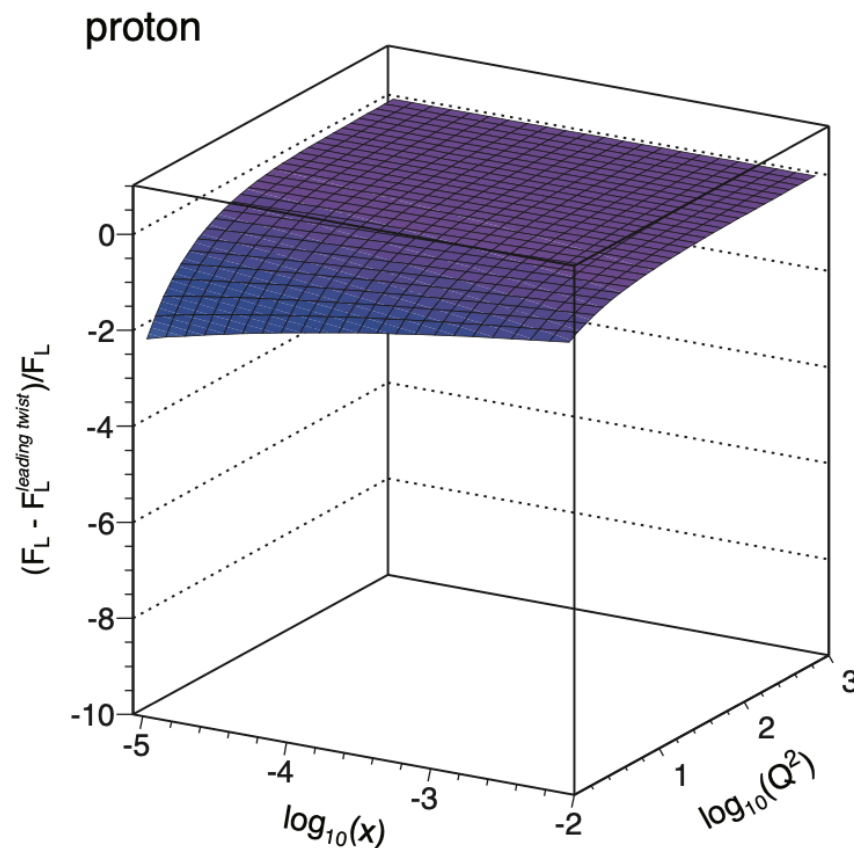


Artwork: BNL

Structure functions at small-x

F_L in the saturation framework vs leading twist

$$F_L(x, Q^2) = \frac{4\pi^2\alpha_{em}}{Q^2} \sigma_L^{\gamma^* A}(x, Q^2)$$



Bartels, Golec-Biernat, Motyka. 0911.1935

F_L is affected by multiple gluon scatterings at all orders in twists ($1/Q^2$)

Enhanced for nuclear collisions and at smaller values of x

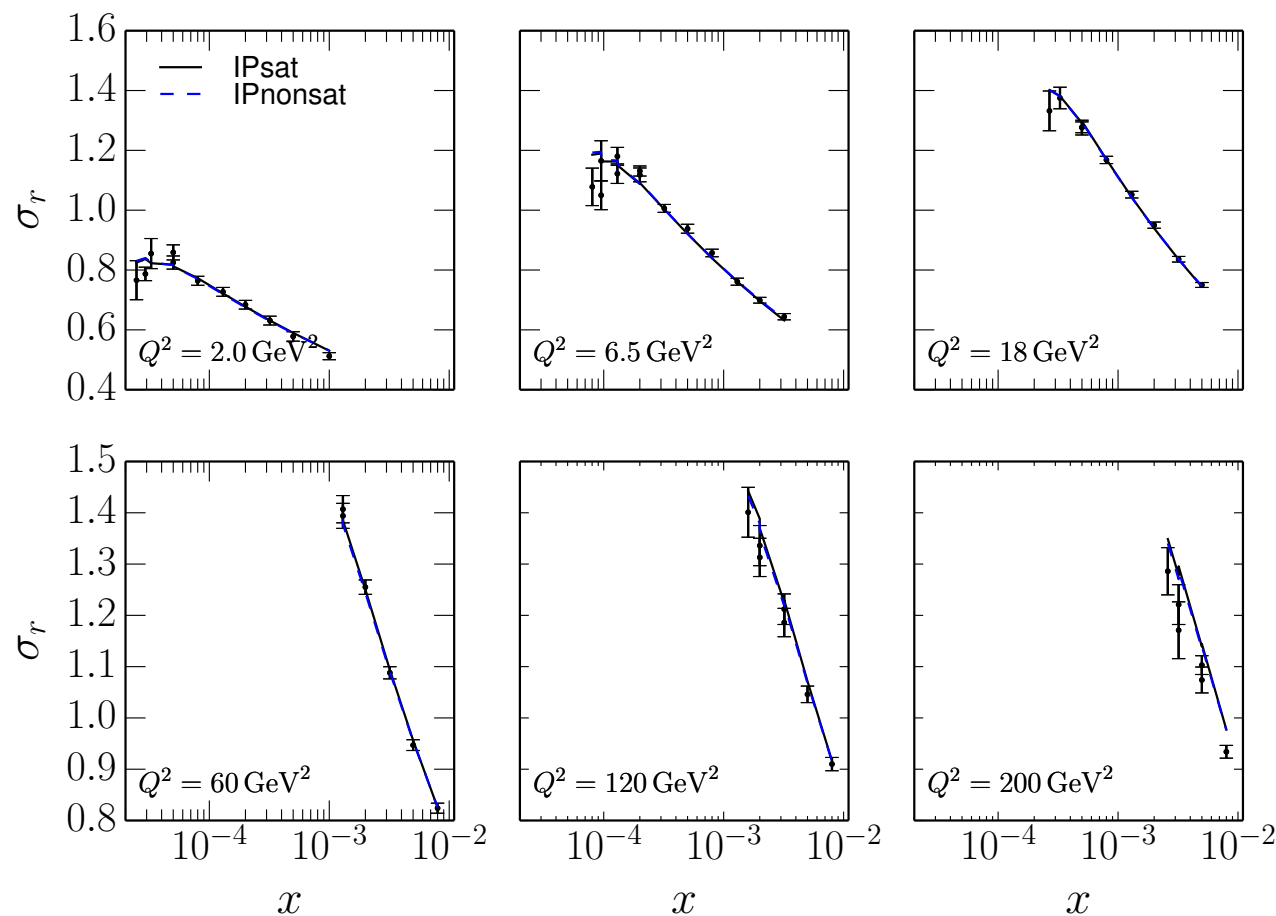
Will these differences manifest at colliders?

Structure functions at small-x

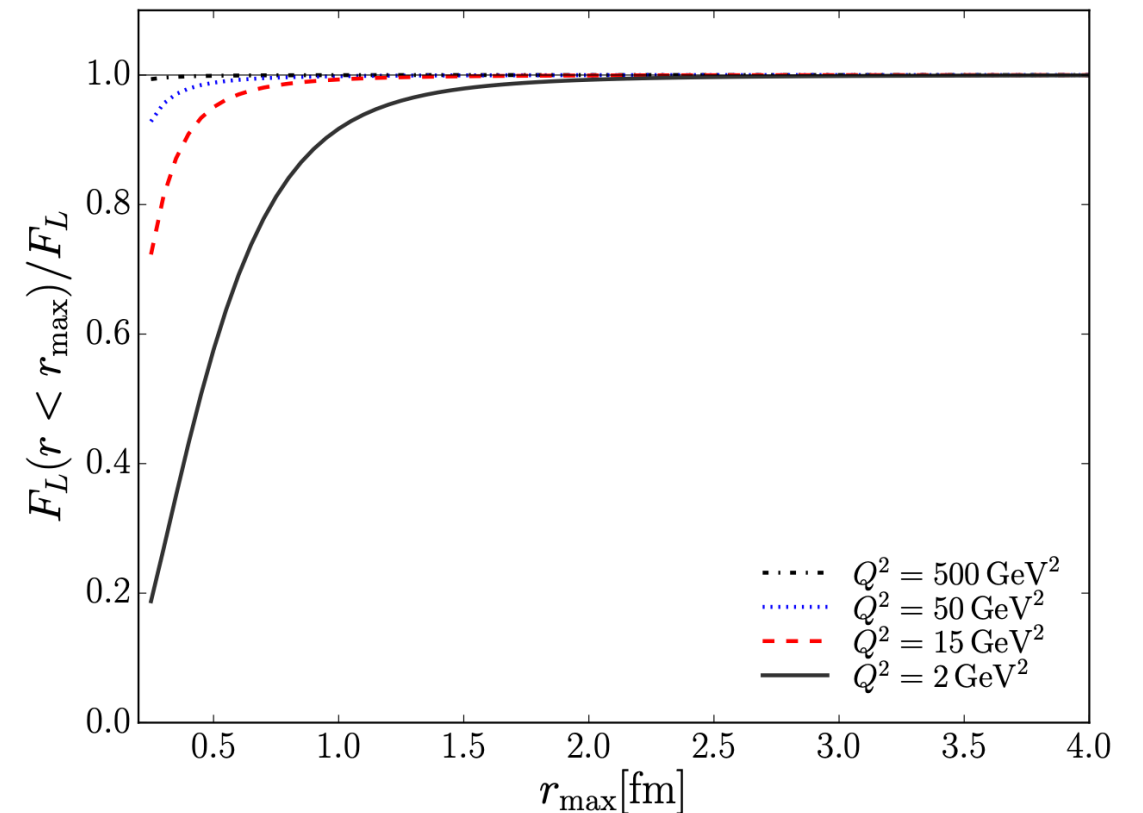
Reduced cross-section and F_L at HERA

$$\sigma_r(x, y, Q^2) = F_2(x, Q^2) - \frac{y^2}{1 + (1 - y)^2} F_L(x, Q^2)$$

IPsat & IPnonsat (linearized IPsat)



Dipole size contributions to F_L



At HERA saturation searches are hindered by non-perturbative contributions

Independent fits using IPSAT and IPnonsat describe data equally well.

No clear evidence of gluon saturation at HERA

Mäntysaari, Zurita. [1804.05311](#)

See also Mäntysaari, Schenke. [1806.06783](#)

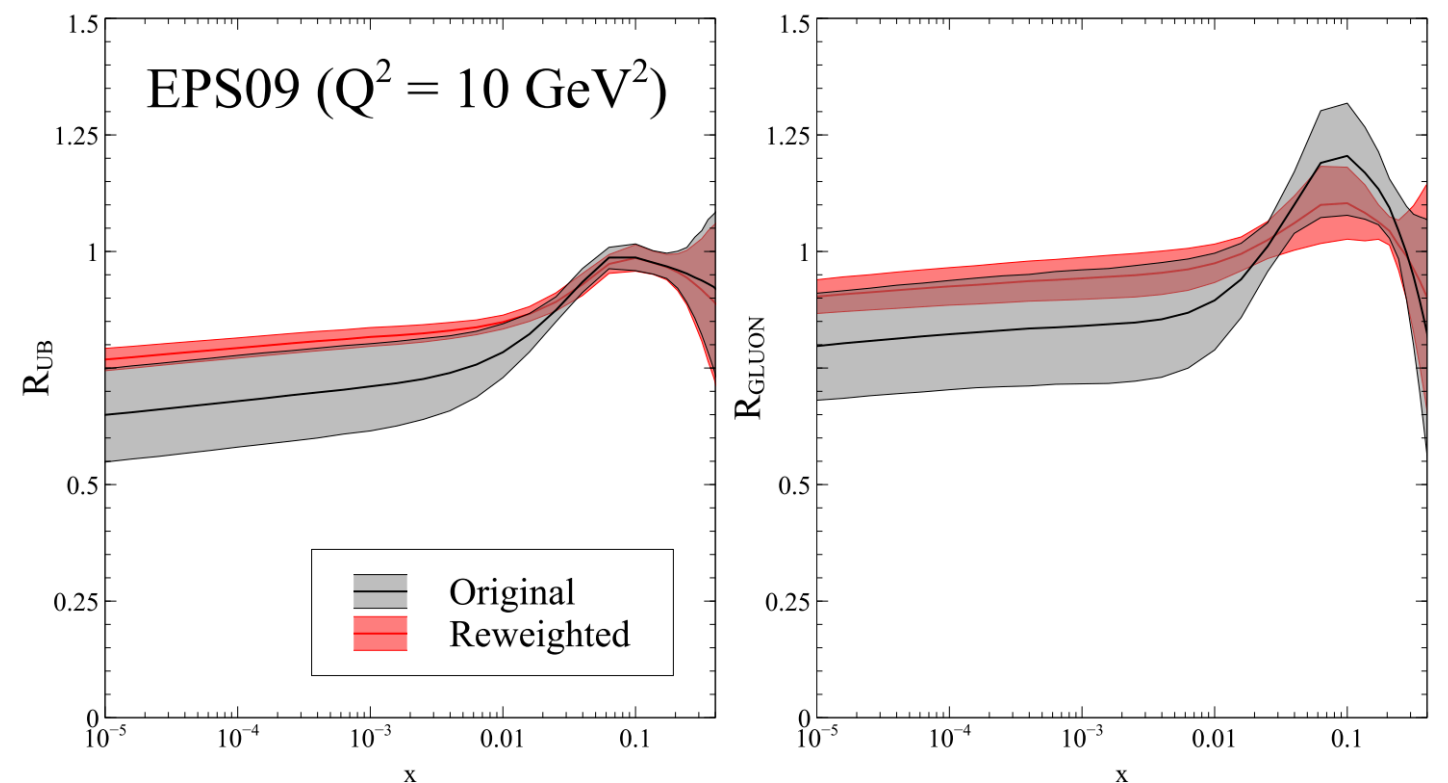
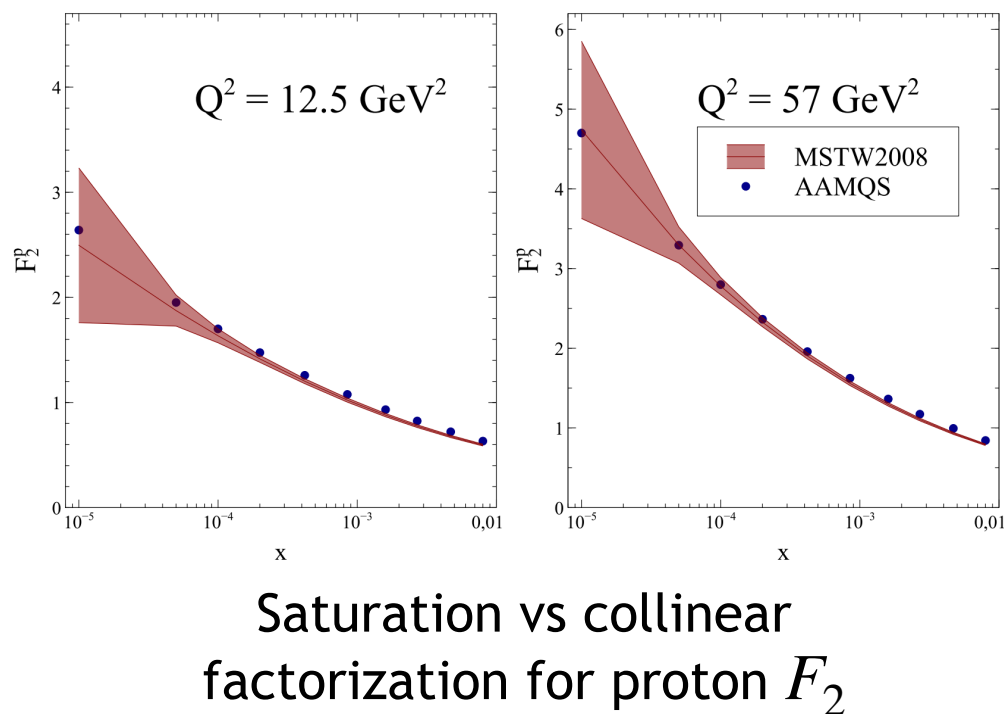
Structure functions at small-x

Nuclear PDFs vs saturation framework

Saturation effects
enhanced in nuclei

$$Q_s^2 \sim A^{1/3}$$

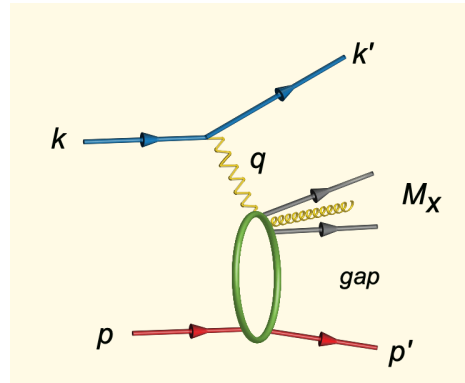
Less sensitivity to
non-perturbative effects



Sea quark and gluon distribution before
and after Bayesian re-weighting

At the EIC it might be possible to find tension between nuclear PDFs and saturation predictions

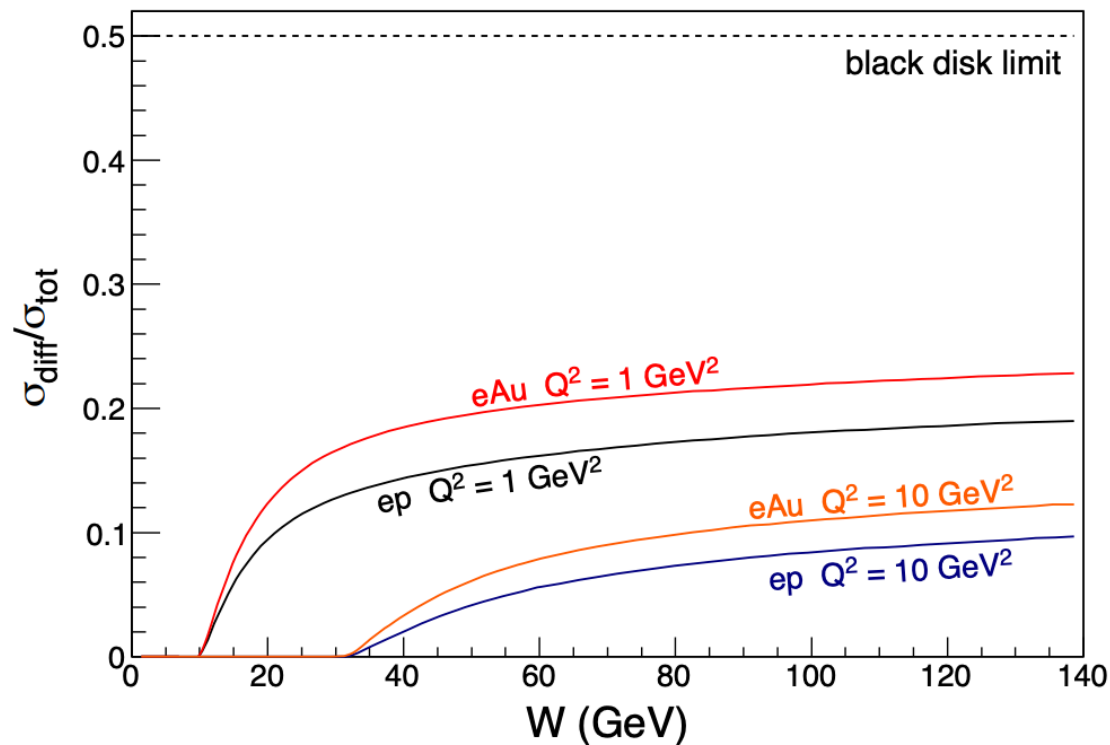
Diffractive processes



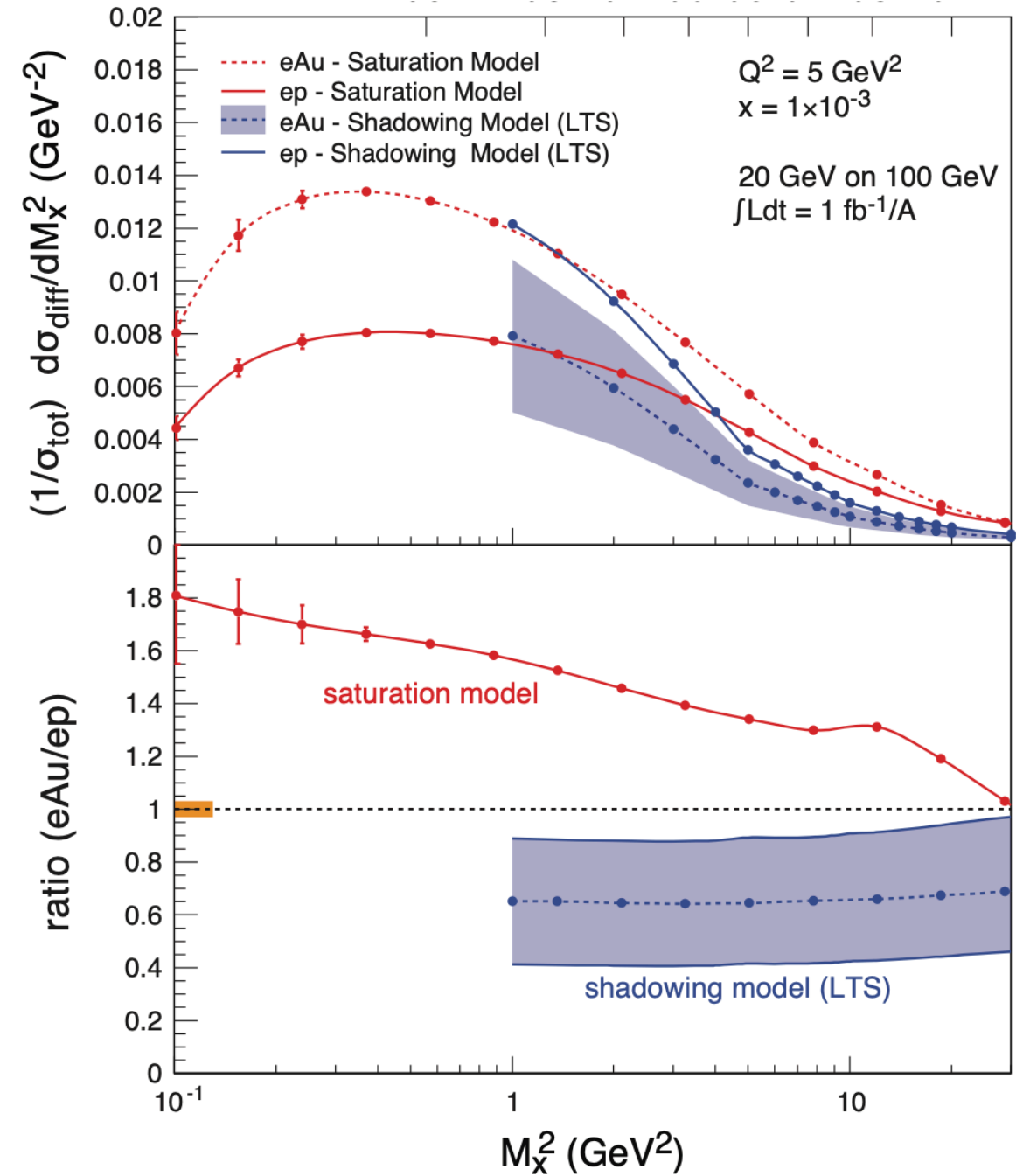
Diffractive events are characterized by rapidity gap

Neutral color exchange requires at least **two-gluons** (enhanced sensitivity to gluon sat)

Ratio of diffractive and total cross-section in ep and eAu collisions



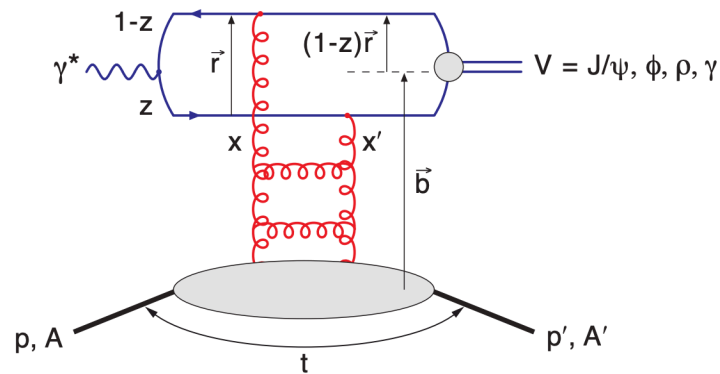
Diffractive events enhanced at lower Q^2 and have weak dependence on energy



Clear difference between saturation models and leading twist shadowing (LTS)

Diffractive processes

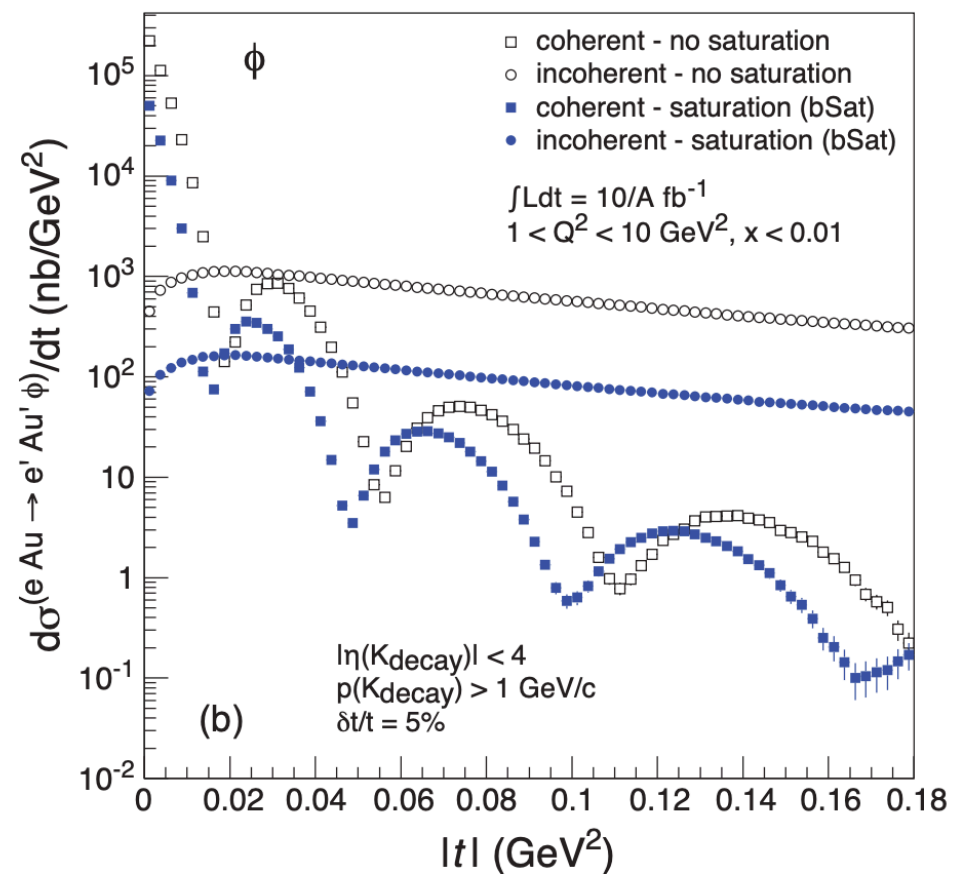
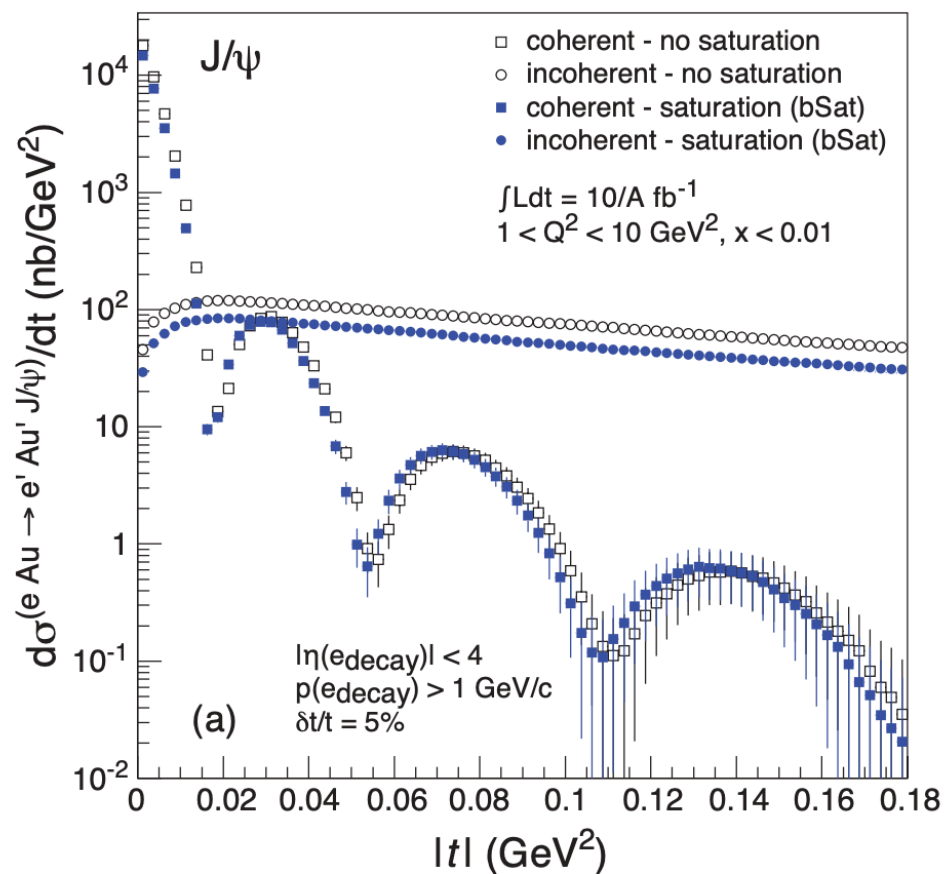
Exclusive nuclear electro-production of VM



*Sensitive to spatial distribution
(tomography)*

$$t = -\Delta_{\perp}^2$$

$$\Delta_{\perp} \leftrightarrow b_{\perp}$$

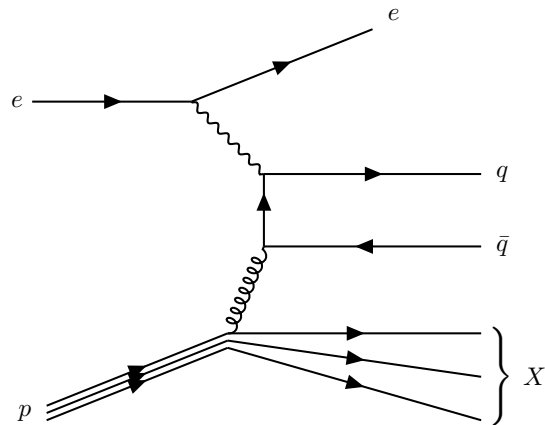


- Sartre event generator (bSat & bNonSat = linearized bSat)
- Large difference for ϕ less so for J/ψ

Semi-inclusive measurements

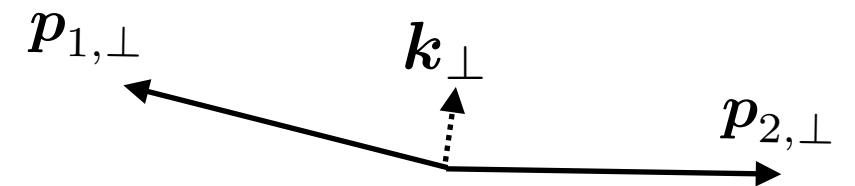
Back-to-back di-hadron suppression

Dominguez, Marquet, Xiao, Yuan. [1101.0715](#)



TMD factorization at
small- x Weizsäcker-Williams
 $xG_{WW}(x, k_{\perp})$ gluon TMD

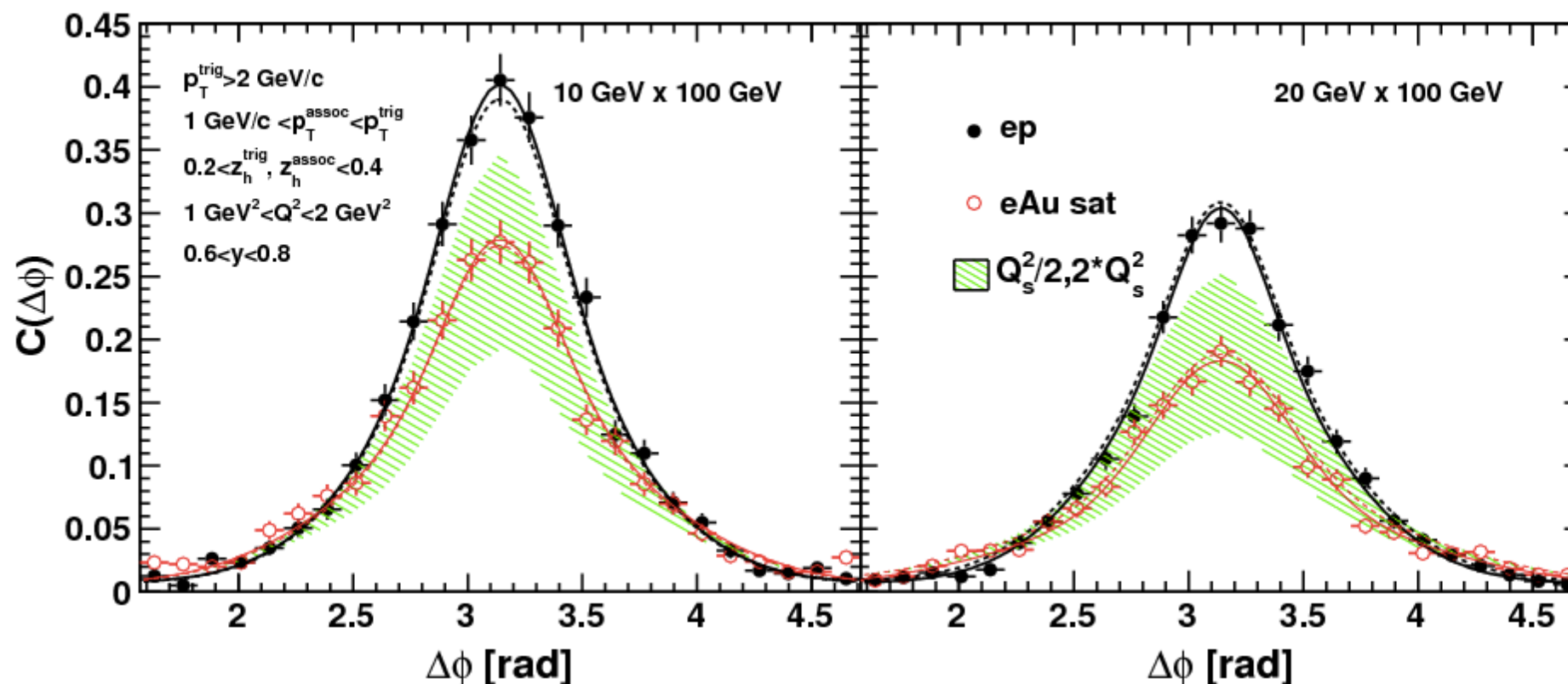
WW at small- x depends on
sat scale Q_s



$\Delta\phi$ angle between $p_{1\perp}$ and $p_{2\perp}$

Typical momentum transfer
from hadron/nucleus to

$$k_{\perp} \sim Q_s$$



Gluon saturation $\longrightarrow$ back-to-back
suppression

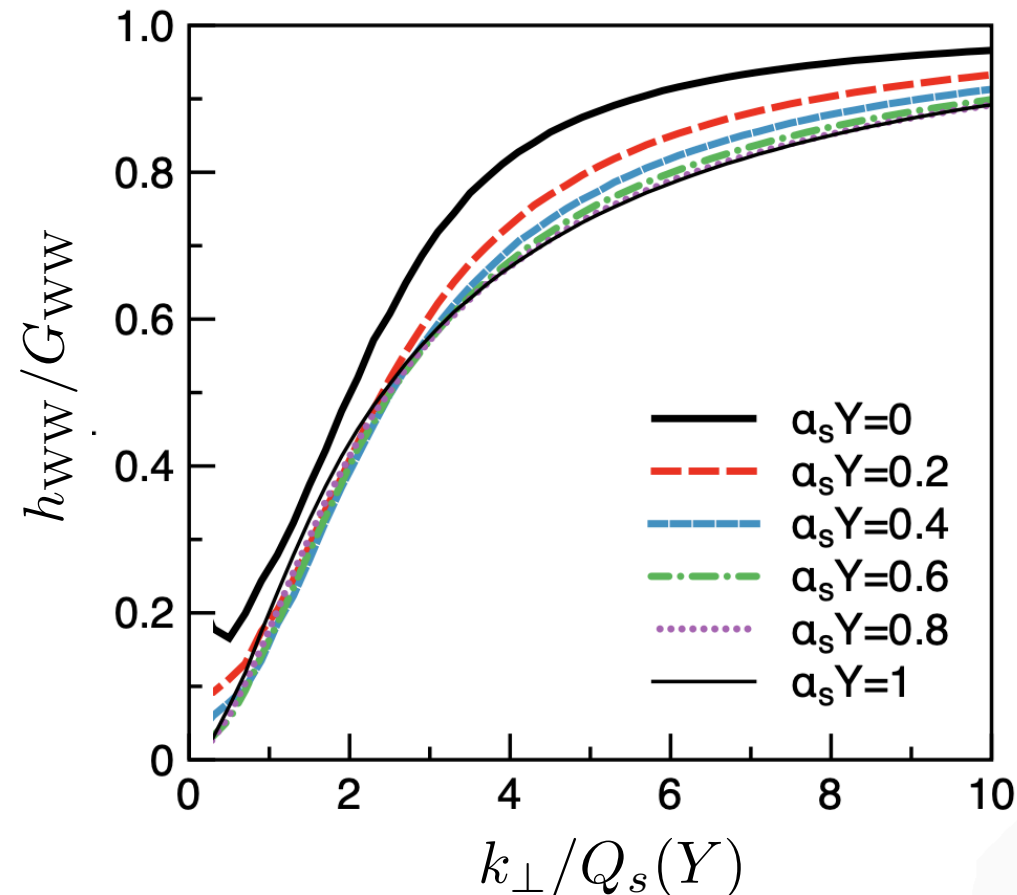
Zheng, Aschenauer, Lee, Xiao. [1403.2413](#)

Sudakov factor included!

Semi-inclusive measurements

Azimuthal anisotropies of dijet momentum imbalance

Ratio linearly polarized to unpolarized
WW gluon TMD in the CGC



Dumitru, Lappi, Skokov. [1508.04438](#)

Accessed in azimuthal asymmetry of
momentum imbalance in dijet pairs

Dominguez, Qiu, Xiao, Yuan. [1109.6293](#)

For effect of soft gluon final state radiation see
Hatta, Xiao, Yuan, Zhou. [2010.10774](#)

WW gluon TMD
(linearly polarized) $\rightarrow$ Dijet azimuthal
asymmetries

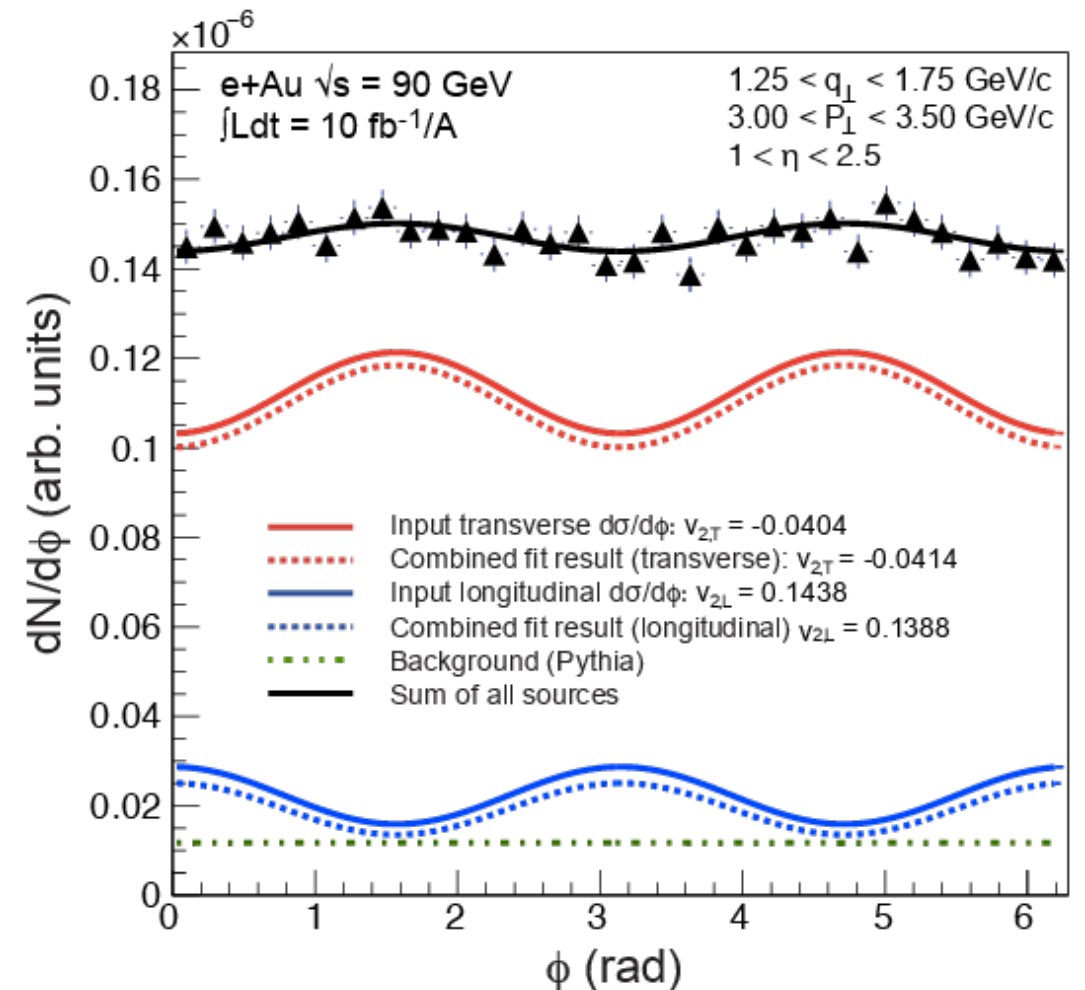
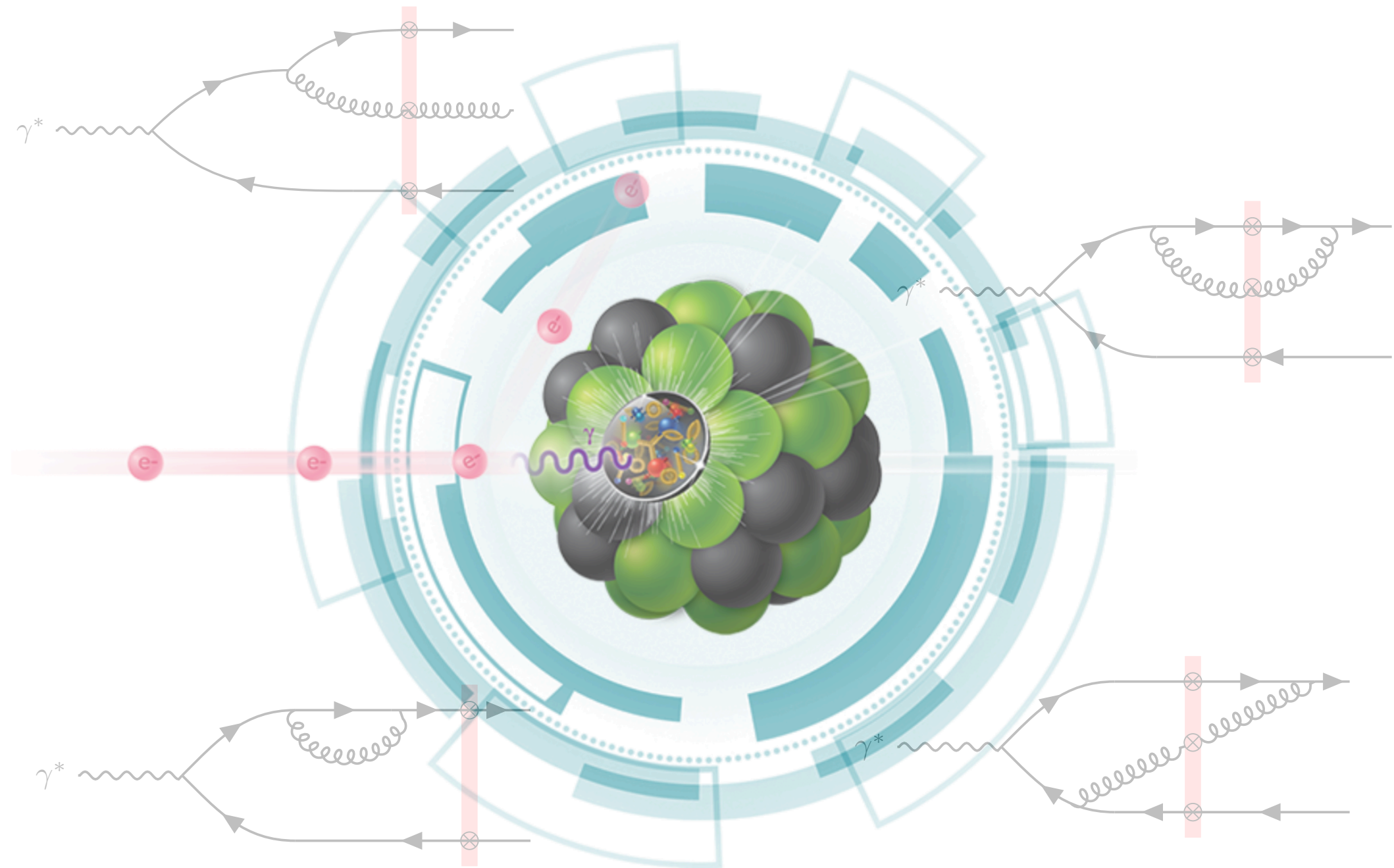


Diagram illustrating the geometry of dijet momentum imbalance. The vectors $p_{1,\perp}$ and $p_{2,\perp}$ are shown, along with the total transverse momentum $2P_{\perp}$ and the momentum imbalance vector $k_{\perp}$. The angle ϕ is defined as the angle between $P_{\perp}$ and $k_{\perp}$.

$$|v_2| \sim \frac{x h_{WW}}{x G_{WW}}$$

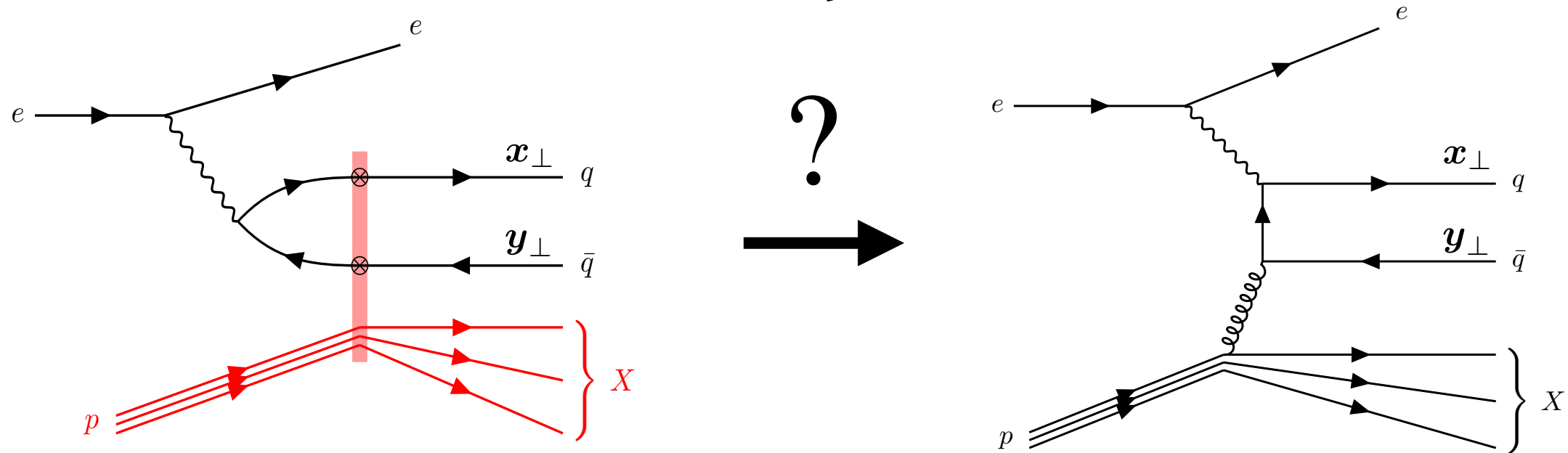
$$v_2 = \langle \cos(2\phi) \rangle$$

Recent theoretical developments



Matching TMD and CGC frameworks at small-x

From Wilson Lines to TMDs and beyond



Dominguez, Marquet, Xiao, Yuan. [1101.0715](#)

Small dipole size expansion amplitude:

$$1 - V(\mathbf{x}_\perp)V^\dagger(\mathbf{y}_\perp) = \underbrace{ig\mathbf{r}_\perp^i \tilde{\mathbf{A}}_\perp^i}_{\text{TMD}} + \mathcal{O}(r_\perp^2)$$

$$\mathbf{r}_\perp = \mathbf{x}_\perp - \mathbf{y}_\perp$$

Altinoluk, Boussarie. [1902.07930](#)

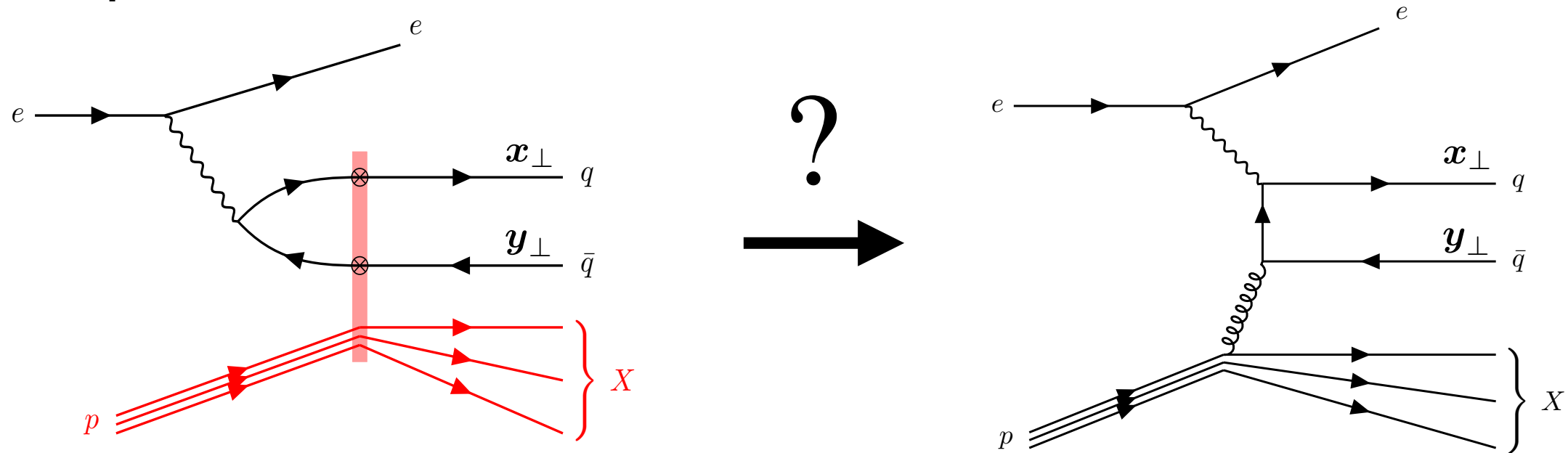
Boussarie, Mehtar-Tani. [2001.06449](#)

A gauge invariant expansion:

$$\mathbb{1} - V(\mathbf{x}_\perp)V^\dagger(\mathbf{y}_\perp) = \underbrace{ig \int_{\mathbf{y}_\perp}^{\mathbf{x}_\perp} d\mathbf{z}_\perp^i \tilde{\mathbf{A}}_\perp^i(\mathbf{z}_\perp)}_{\text{ITMD}} + \underbrace{g^2 \int_{\mathbf{y}_\perp}^{\mathbf{x}_\perp} d\mathbf{z}_{1\perp}^i \int_{\mathbf{y}_\perp}^{\mathbf{z}_{1\perp}} d\mathbf{z}_{2\perp}^j \tilde{\mathbf{A}}_\perp^i(\mathbf{z}_{1\perp}) V(\mathbf{z}_{1\perp}) V^\dagger(\mathbf{z}_{2\perp}) \tilde{\mathbf{A}}_\perp^j(\mathbf{z}_{2\perp})}_{\text{genuine higher twists (g.h.t.)}}$$

Matching TMD and CGC frameworks at small-x

CGC, improved TMD and TMD frameworks



Mäntysaari, Mueller, FS, Schenke. [1912.05586](#)

Boussarie, Mäntysaari, FS, Schenke.

[2106.XXXXX](#)

$$d\sigma_{\text{CGC}} = \underbrace{d\sigma_{\text{TMD}}}_{d\sigma_{\text{ITMD}}} + \overbrace{\mathcal{O}\left(\frac{k_{\perp}}{Q_{\perp}}\right)}^{\text{kinematic}} + \overbrace{\mathcal{O}\left(\frac{Q_s}{Q_{\perp}}\right)}^{\text{genuine}}$$

TMD valid $k_{\perp}, Q_s \ll Q_{\perp}$

back-to-back hadrons/jets
and transverse momenta larger than
sat scale

Improved TMD valid $Q_s \ll Q_{\perp}$

transverse momenta larger
than sat scale

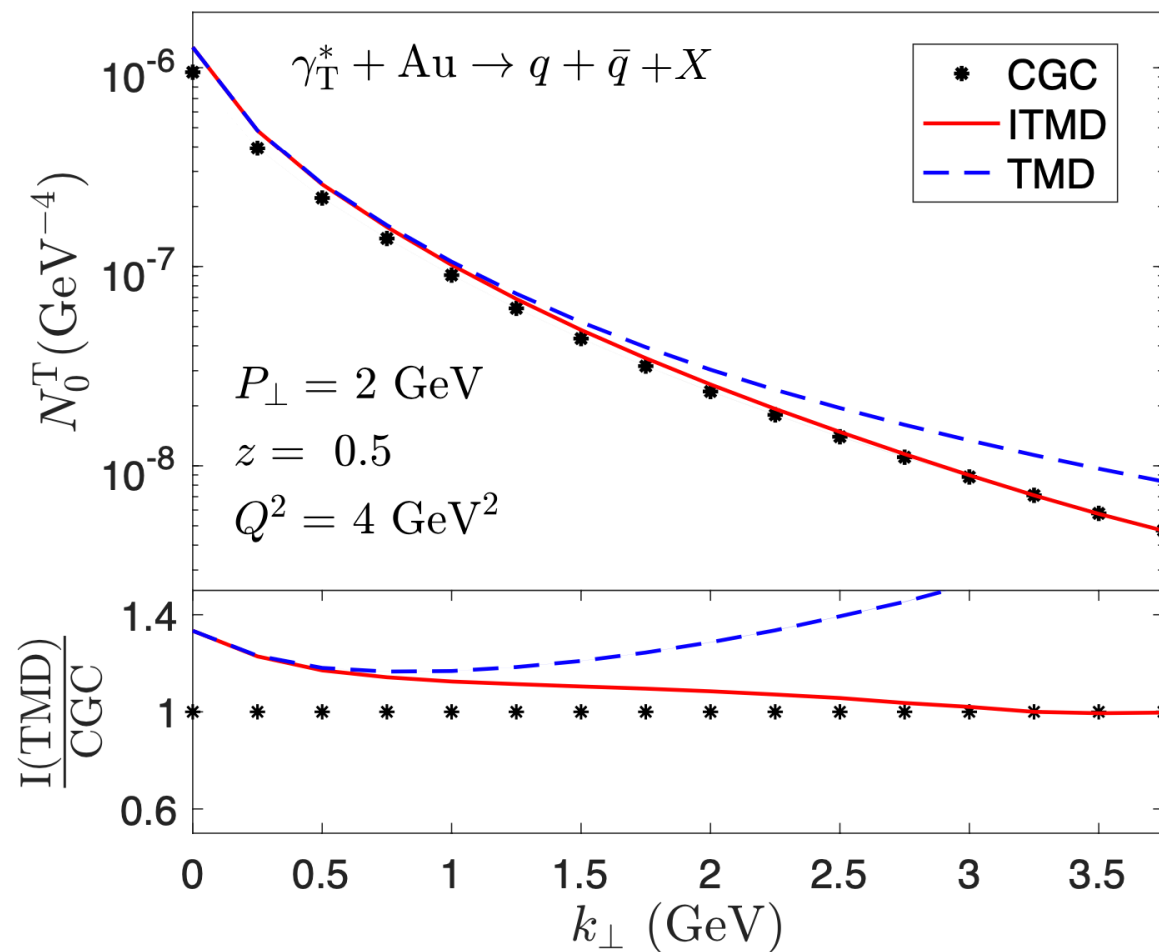
No need for back-to-back!

Matching TMD and CGC frameworks at small-x

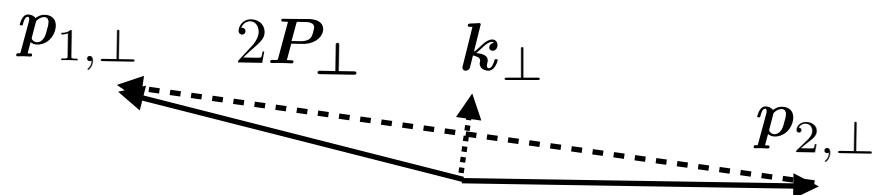
Kinematic and genuine higher saturation effects at the EIC

Boussarie, Mäntysaari, FS, Schenke. 2106.XXXXX

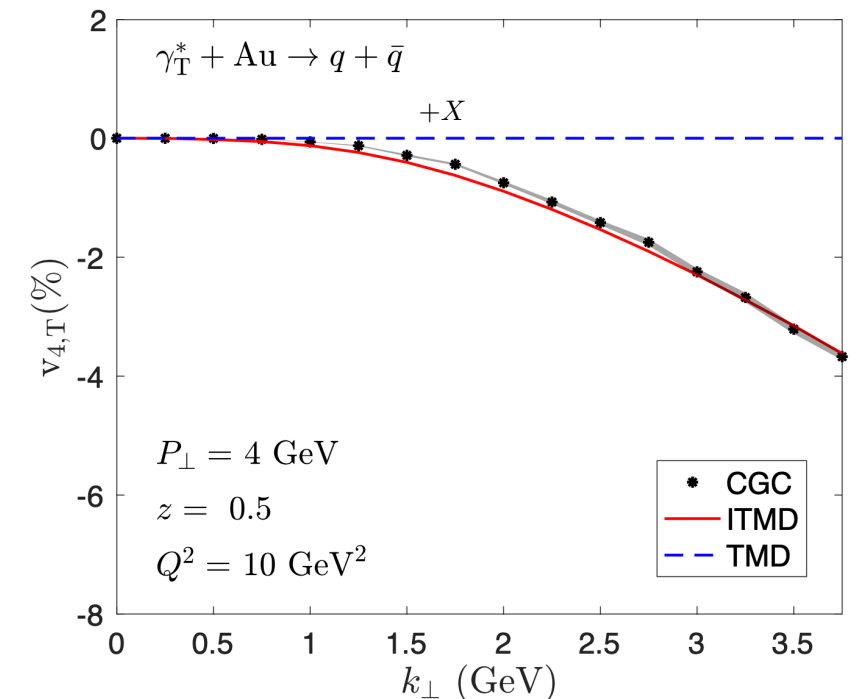
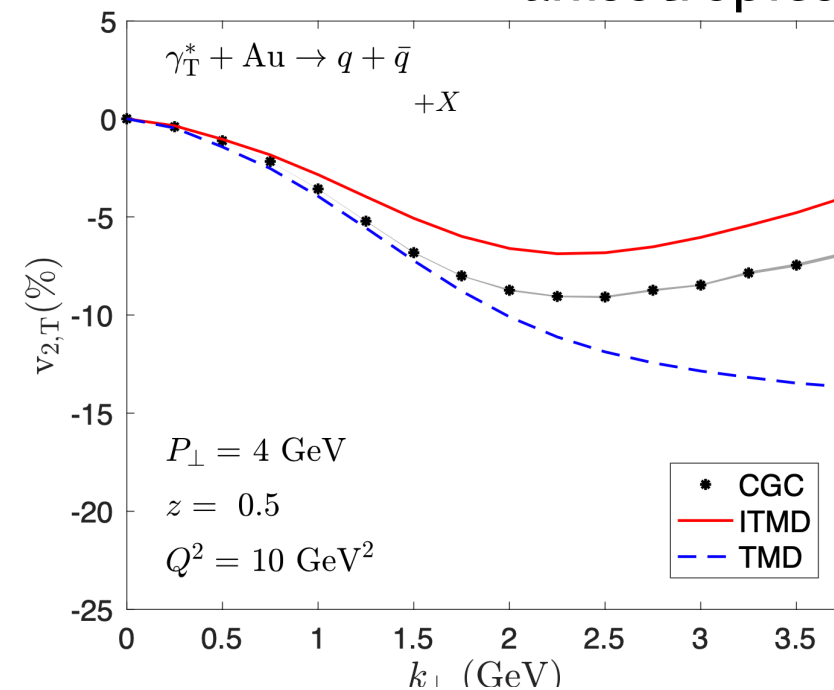
Differential yield



CGC shows further suppression relative to TMD at back-to-back limit



Momentum imbalance azimuthal anisotropies

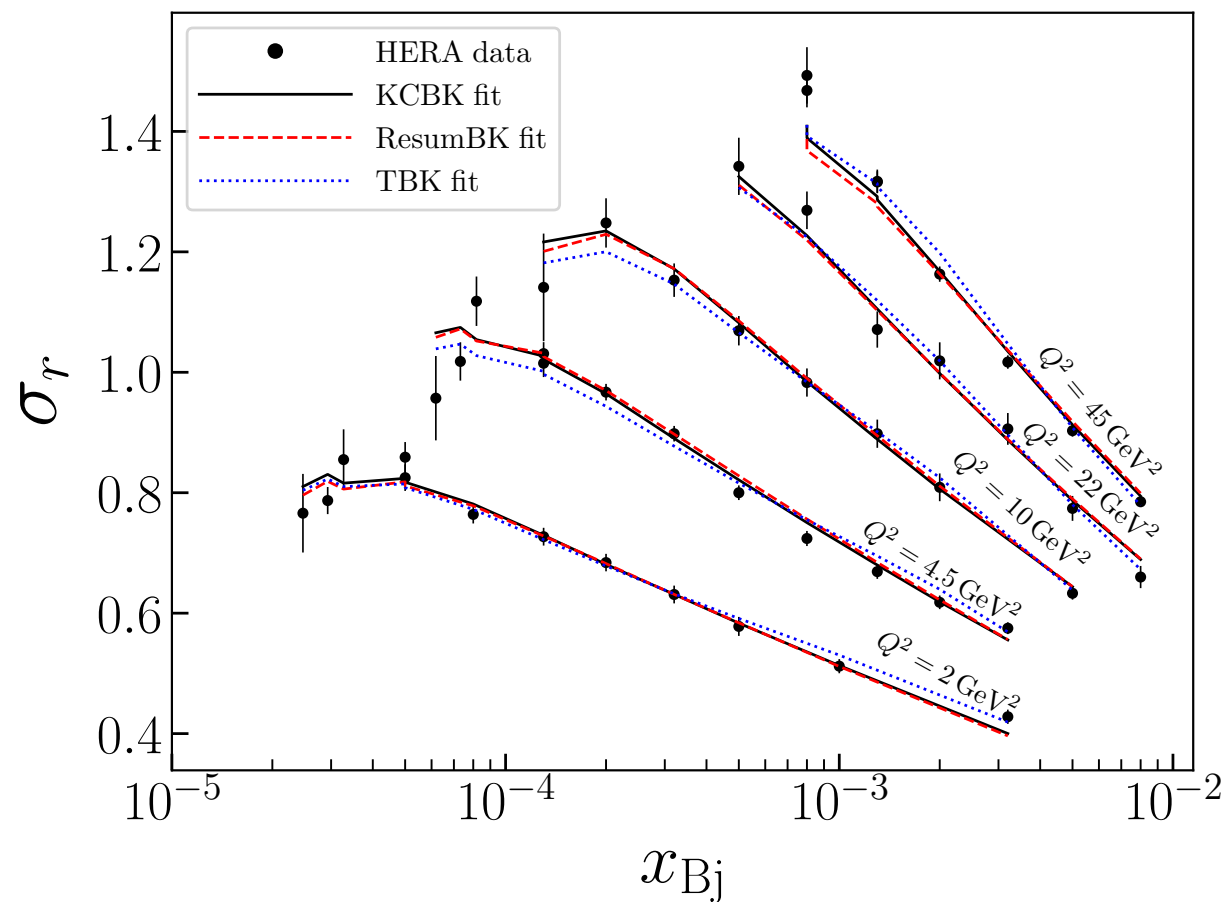


Anisotropies modified in ITMD and CGC

Next-to-leading order computations

Recent advances of NLO computations

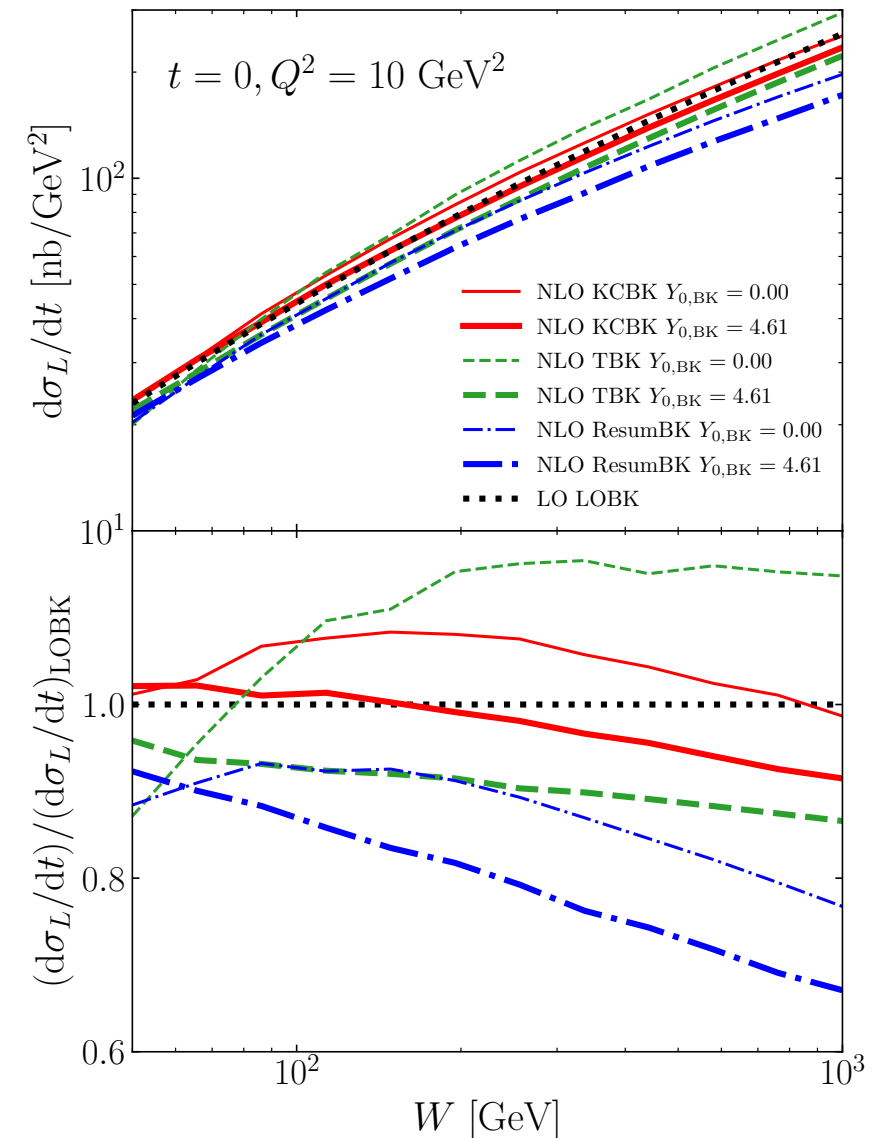
Reduced DIS cross-section fits with
NLO BK + NLO impact factor*



Beuf, Lappi, Hänninen, Mäntysaari. [2007.01645](#)

**Should be extended to nuclear DIS
for EIC predictions!**

J/Ψ electron production
(only L polarization) at NLO



Mäntysaari, Penttala. [2104.02349](#)

**Comparison with HERA data
once T pol is theoretically available**

* only light-quark contribution

Next-to-leading order computations

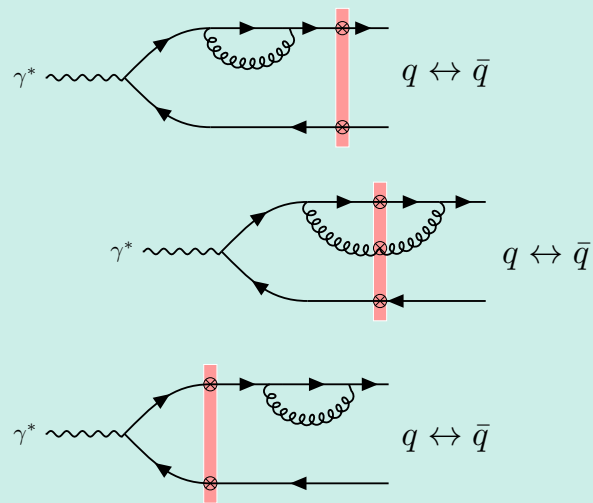
Dijet production at NLO

Caucal, FS, Venugopalan. 2107.XXXXX

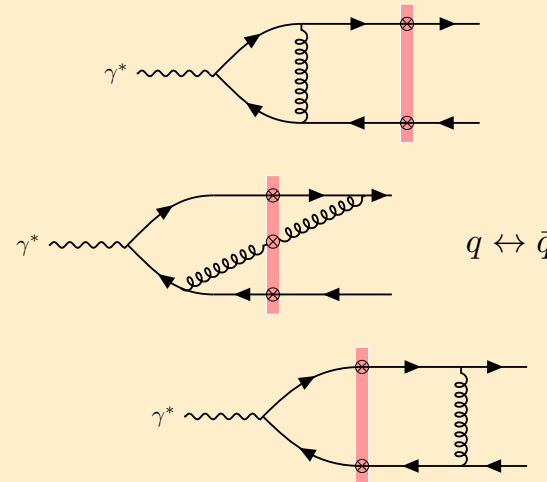
See also Roy, Venugopalan.
1911.04530 for photon+dijet

Computation of real and virtual contributions to dijet production

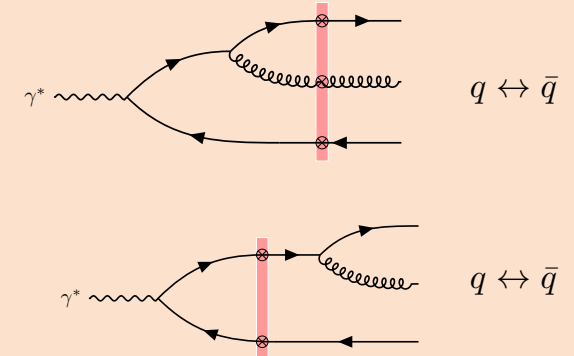
Self-energy contributions



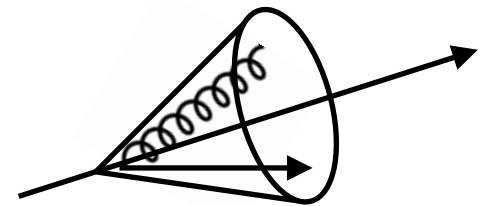
Vertex contributions



Real emission



Small-cone condition:



- Cancellation of divergences: UV, soft and collinear
- Factorization of small-x: Leading Log JIMWLK evolution
- Numerically tractable expressions for impact factor

$$\mathcal{C}_{qg,\perp}^2 \leq \mathcal{C}_{qg,\perp}^2|_{\max} = R^2 p_j^2 \min \left(\frac{z_g^2}{z_j^2}, \frac{(z_j - z_g)^2}{z_j^2} \right)$$

for proper jet definition

Numerical results will be available
in the near future!

Summary/Outlook

- Discovery of gluon saturation is one of the major goals of the future EIC
- The Color Glass Condensate is an EFT for this saturated regime that has been applied to calculate a variety of processes
- A wide variety of observables are available to search for gluon saturation:
structure functions, diffractive processes and semi-inclusive measurements
- Recent advances in higher order computations and connections to TMD framework will allow to extend the robustness of theoretical predictions for the EIC

Other measurements at the EIC

- Exclusive dijet measurements and Wigner distribution
Hatta, Xiao, Yuan. [1601.01585](#)
Mäntysaari, Mueller, Schenke. [1902.05087](#)
- Incoherent diffraction and fluctuations
Mäntysaari, Schenke. [1603.04349](#)
Mäntysaari, Roy, FS, Schenke. [2011.02464](#)
- Single hadron/jet semi-inclusive production R_{eA} ratios
Iancu, Mueller, Triantofyllopoulos, Wei. [2012.08562](#)
- Inclusive quarkonia and WW gluon TMD
Bacchetta, Boer, Pisano, Taelis. [1809.02056](#)
- Hadron+photon correlations
Kolbé, Roy, FS, Schenke, Venugopalan. [2008.04372](#)