Glue Imaging in Production of Vector Mesons $(J/\Psi, \phi \cdots)$

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2nd Precision Studies on QCD at EIC



An analogy to Fraunhofer Diffaction in Optics

[QCD at high energy, Kovchegov and Levin, 12]



- Treat the hadron target in DIS as a black disk. [Joseph von Fraunhofer, 1821]
- Similar pattern in optics ($\theta_i^{\min} \sim 1/(kR)$) and high energy QCD $t_i \sim \frac{1}{R^2}$.
- Two difference: 1. σ sensitive to gluon distribution; 2. Breakup of the target.
- This motivates us to use diffractive scattering to study gluon spatial distribution.



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Gluon GPDs and DVMP processes

For exclusive/diffractive processes with $V = J/\Psi, \phi \cdots$



 $\gamma^*(q) + p/A(p) \to V(q - \Delta) + p/A(p + \Delta)$

- The latter diagram is dominant at small-*x* (high energy) limit.
- Widely studied at small-x [Brodsky, Frankfurt, Gunion, Mueller, Strikman, 94; Kowalski, Teaney, 03; Kowalski, L. Motyka, Watt, 06, Watt, Kowalski, 08; Kowalski, Caldwell, 10; Berger, Stasto, 13; Rezaeian, Schmidt, 13]...
- Incoherent diffractive production for nucleon/nuclear targets [T. Lappi, H. Mantysaari, 11; Toll, Ullrich, 12; Lappi, Mantysaari, R. Venugopalan, 15; T. Lappi, Mantysaari, Schenke, 16]...; Review by [Mantysaari, 20]
- DVMP at NLO. [Boussarie, Grabovsky, Ivanov, Szymanowski, Wallon, 16]



Introduction Deeply Virtual Meson Production (DVMP) and Gluon GPDs

Wigner distribution

Wigner distributions [Ji, 03; Belitsky, Ji, Yuan, 2004] ingeniously encode all quantum information of how partons are distributed inside hadrons.



- Quasi-probability distribution; Not positive definite. [F. Yuan's talk]
- GPDs encode the parton spatial distributions.



The exact connection between dipole amplitude and Wigner distribution

[Hatta, Xiao, Yuan, 16] Def. of gluon Wigner distribution:

$$\begin{split} xW_g^T(x,\vec{q}_{\perp};\vec{b}_{\perp}) &= \int \frac{d\xi^- d^2\xi_{\perp}}{(2\pi)^3 P^+} \int \frac{d^2\Delta_{\perp}}{(2\pi)^2} e^{-ixP^+\xi^- -iq_{\perp}\cdot\xi_{\perp}} \\ &\times \quad \left\langle P + \frac{\Delta_{\perp}}{2} \left| F^{+i} \left(\vec{b}_{\perp} + \frac{\xi}{2} \right) F^{+i} \left(\vec{b}_{\perp} - \frac{\xi}{2} \right) \right| P - \frac{\Delta_{\perp}}{2} \right\rangle \,, \end{split}$$

Define GTMD [Meissner, A. Metz and M. Schlegel, 09]

$$xG(x,q_{\perp},\Delta_{\perp}) \equiv \int d^2 b_{\perp} e^{-i\Delta \cdot b_{\perp}} x W_g^T(x,\vec{q}_{\perp};\vec{b}_{\perp}).$$

With dipole-like gauge link, one finds

$$\begin{split} xG_{\rm DP}(x,q_{\perp},\Delta_{\perp}) &= \frac{2N_c}{\alpha_s} \int \frac{d^2 R_{\perp} d^2 R'_{\perp}}{(2\pi)^4} e^{iq_{\perp} \cdot (R_{\perp} - R'_{\perp}) + i\frac{\Delta_{\perp}}{2} \cdot (R_{\perp} + R'_{\perp})} \\ &\times \left(\nabla_{R_{\perp}} \cdot \nabla_{R'_{\perp}} \right) \frac{1}{N_c} \left\langle \operatorname{Tr} \left[U(R_{\perp}) U^{\dagger}(R'_{\perp}) \right] \right\rangle_x. \end{split}$$



Relation to TMDs and GPDs. Small-x tells us more information about them.

Scattering amplitude



- At small-*x*, the lifetime of $q\bar{q}$ fluctuation $\tau_f \sim \frac{q^+}{Q^2} + \frac{q^+}{M_V^2} \gg \tau_{int}$ with $\tau_{int} \sim R/\gamma$.
- **Factorize** the amplitude into three pieces in coordinate space:

$$\begin{split} \mathcal{A}(\Delta_{\perp}) &\sim \int d^2 r_{\perp} d^2 b_{\perp} e^{i b_{\perp} \cdot \Delta_{\perp}} \int_0^1 dz \Psi_{\gamma^*}(z, r_{\perp}) \Psi_V^*(z, r_{\perp}) \\ &\times \left\{ 1 - \frac{1}{N_c} \mathrm{Tr} \left[U(b_{\perp} + z r_{\perp}) U^{\dagger}(b_{\perp} - (1 - z) r_{\perp}) \right] \right\}. \end{split}$$

• I: $\gamma^* \to q\bar{q}$ splitting; II: $q\bar{q}$ scattering; III: CC of $V \to q\bar{q}$ splitting.



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Relation between gluon GPDs and dipole scattering amplitude



[Hatta, Xiao, Yuan, 17] gluon GPDs with asymmetric dipole for arbitrary z

• Dipole scattering amplitude $S_x = \left\langle \frac{1}{N_c} \operatorname{Tr} \left[U(x_{1\perp}) U^{\dagger}(x_{2\perp}) \right] \right\rangle_x$ defines GTMD:

$$\mathcal{F}_{x}(\tilde{q}_{\perp},\Delta_{\perp}) = \frac{1}{(2\pi)^{4}} \int d^{2}x_{1\perp} d^{2}x_{2\perp} e^{ik_{1\perp} \cdot x_{1\perp} - ik_{2\perp} \cdot x_{2\perp}} S_{x}(x_{1\perp},x_{2\perp})$$

with $k_{1\perp} \equiv \tilde{q}_{\perp} + z \Delta_{\perp}$ and $k_{2\perp} \equiv \tilde{q}_{\perp} - (1-z)\Delta_{\perp}$.

Symmetric dipole definition with shifts of coordinate and momentum

$$F_x(q_{\perp}, \Delta_{\perp}) = \int \frac{d^2 r_{\perp} d^2 b_{\perp}}{(2\pi)^4} e^{ib_{\perp} \cdot \Delta_{\perp} + ir_{\perp} \cdot q_{\perp}} S_x \left(b_{\perp} + \frac{r_{\perp}}{2}, b_{\perp} - \frac{r_{\perp}}{2} \right)$$

Dipole scattering amplitude

Several properties of the scattering amplitude:

- $S_x = \left\langle \frac{1}{N_c} \operatorname{Tr} \left[U(b_{\perp} + \frac{r_{\perp}}{2}) U^{\dagger}(b_{\perp} \frac{r_{\perp}}{2}) \right] \right\rangle_x$
- Due to the Pomeron exchange, S_x is predominantly real in small-x.
- $S_x = S_x^* \Rightarrow S_x$ can only depend on b_{\perp}^2 , r_{\perp}^2 and $(r_{\perp} \cdot b_{\perp})^{2n}$.
- Double Fourier transform:

$$F_x(q_{\perp},\Delta_{\perp}) = \int \frac{d^2 r_{\perp} d^2 b_{\perp}}{(2\pi)^4} e^{ib_{\perp}\cdot\Delta_{\perp} + ir_{\perp}\cdot q_{\perp}} S_x\left(b_{\perp} + \frac{r_{\perp}}{2}, b_{\perp} - \frac{r_{\perp}}{2}\right)$$

Angular dependence after double Fourier transform

 $F_x(q_{\perp}, \Delta_{\perp}) = F_0(|q_{\perp}|, |\Delta_{\perp}|) + 2\cos 2(\phi_{q_{\perp}} - \phi_{\Delta_{\perp}})F_{\epsilon}(|q_{\perp}|, |\Delta_{\perp}|) + \cdots$

Non-trivial angular correlation between Δ_{\perp} and q_{\perp} due to small-*x* dynamics.



Explicit expressions for gluon GPDs



Comparing coefficients gives [Hatta, Xiao, Yuan, 17]

Helicity conserved:
$$xH_g(x, \Delta_{\perp}) = \frac{2N_c}{\alpha_s} \int d^2q_{\perp}q_{\perp}^2F_0$$

Helicity flipping: $xE_{Tg}(x, \Delta_{\perp}) = \frac{4N_c M^2}{\alpha_s \Delta_{\perp}^2} \int d^2 q_{\perp} q_{\perp}^2 F_{\epsilon}$



Probing gluon GPD at small-x

DVCS[Hatta, Xiao, Yuan, 17] and DVMP [Mantysaari, Roy, Salazar, Schenke, 20]



$$\frac{d\sigma_{TT}}{dx_B dQ^2 d^2 \Delta_{\perp}} = \frac{\alpha_{em}^3}{\pi x_{Bj} Q^2} \left\{ \left(1 - y + \frac{y^2}{2}\right) \left(\mathcal{A}_0^2 + \mathcal{A}_2^2\right) + (1 - y) 2\mathcal{A}_0 \mathcal{A}_2 \cos(2\phi_{\Delta l}) \right\}$$

- A_0 : helicity conserved amplitude; A_2 : helicity-flip amplitude
- Use the lepton plane as a reference, one can measure angular correlations.
- $\cos 2\phi_{\Delta l}$ correlation is sensitive to the helicity-flip gluon GPD xE_{Tg} .

Helicity conserved:
$$xH_g(x, \Delta_{\perp}) = \frac{2N_c}{\alpha_s} \int d^2q_{\perp}q_{\perp}^2F_0$$

Helicity flipping: $xE_{Tg}(x, \Delta_{\perp}) = \frac{4N_c M^2}{\alpha_s \Delta_{\perp}^2} \int d^2 q_{\perp} q_{\perp}^2 F_{\epsilon}$



Dipole model and HERA data



[Kowalski, Teaney, 03; Kowalski, Motyka, Watt, 06]



Dipole model and HERA data



[Kowalski, Motyka, Watt, 06]



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Diffractive vector meson production



 Incoherent DVMP is sensitive to the proton fluctuating shape. (Variance) [Mantysaari, Schenke, 16; Mantysaari, Roy, Salazar, Schenke, 20]

Good-Walker: measure of fluctuation

 $\frac{d\sigma_{\rm incoh}}{d\hat{t}} \sim \langle |\mathcal{A}|^2 \rangle - |\langle \mathcal{A} \rangle|^2$



Spatial Imaging at EIC



Ultimate goal: spatial distributions (via FT). [EIC white paper, 1212.1701]



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Introduction
Deeply Virtual Meson Production (DVMP) and Gluon GPDs

Summary



- Gluon spatial imaging through DVMP.
- Probing non-trivial angular correlation and helicity-flip GPD.
- Study of fluctuations via incoherent diffractive productions.
- Another interesting topic: Heavy quarkonium near threshold and proton mass.
 [D. Kharzeev, 96; ...; Hatta, Yang, 18; Kou, Wang, Chen, 21; Guo, Ji, Liu, 21; Sun, Tong, Yuan, 21;...] Small-x picture no longer applies. Paradigm shift!



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