

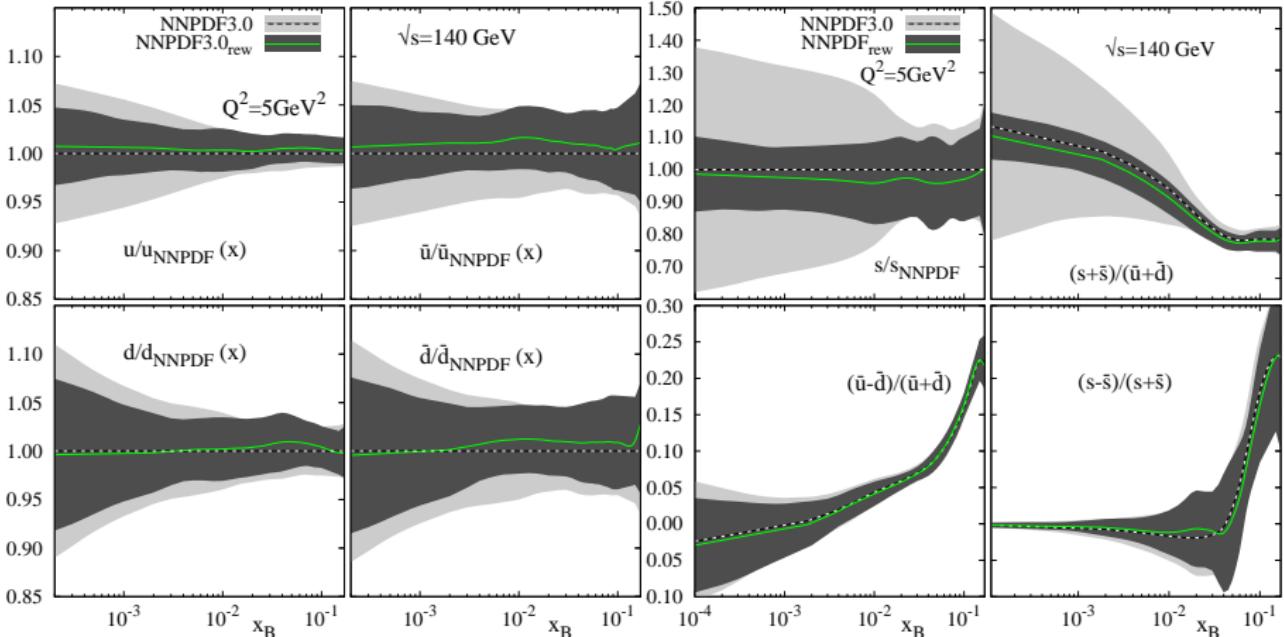
Challenges of QCD at the EIC

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Institute for Theoretical Physics
University of Regensburg

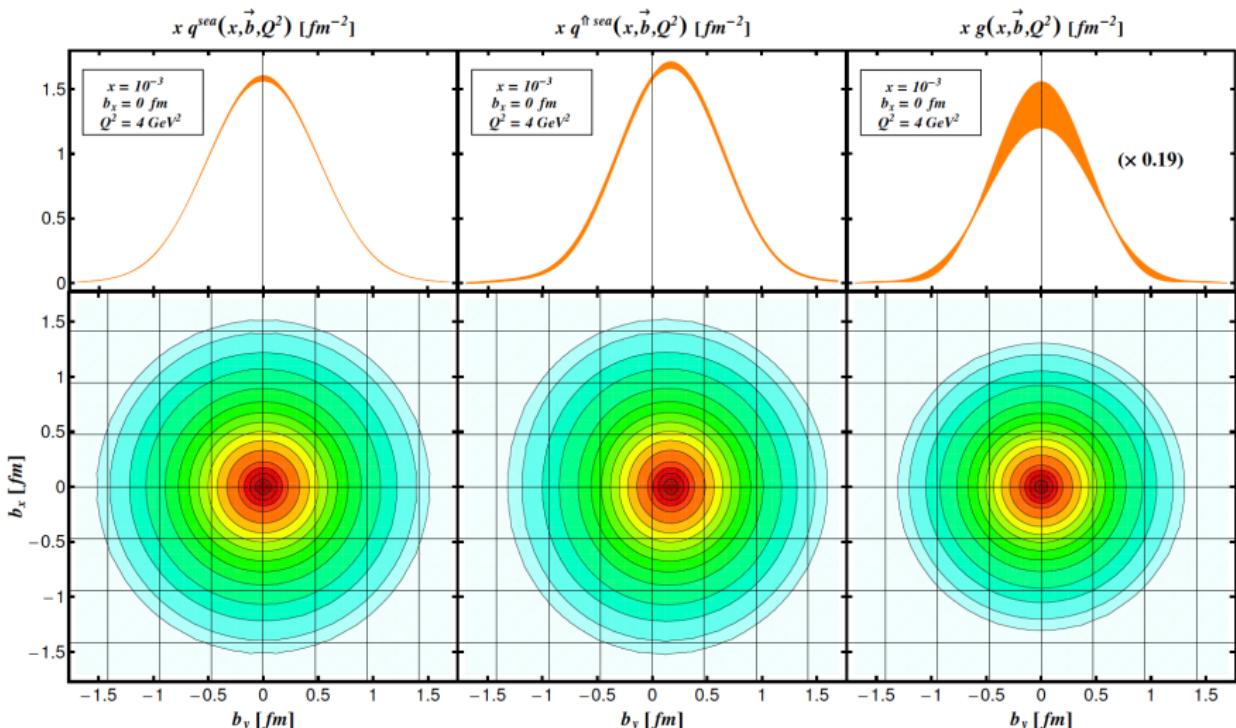
- The two EIC challenges: High precision and multitude of observables, a few topics from Yellow report 2103.05419
- Lattice QCD
- Power corrections, higher twist contributions
- models
- nucleon mass decomposition
- $e + A$: unifying hot and cold QCD

Impact plots from the Yellow Report I: PDFs and FFs

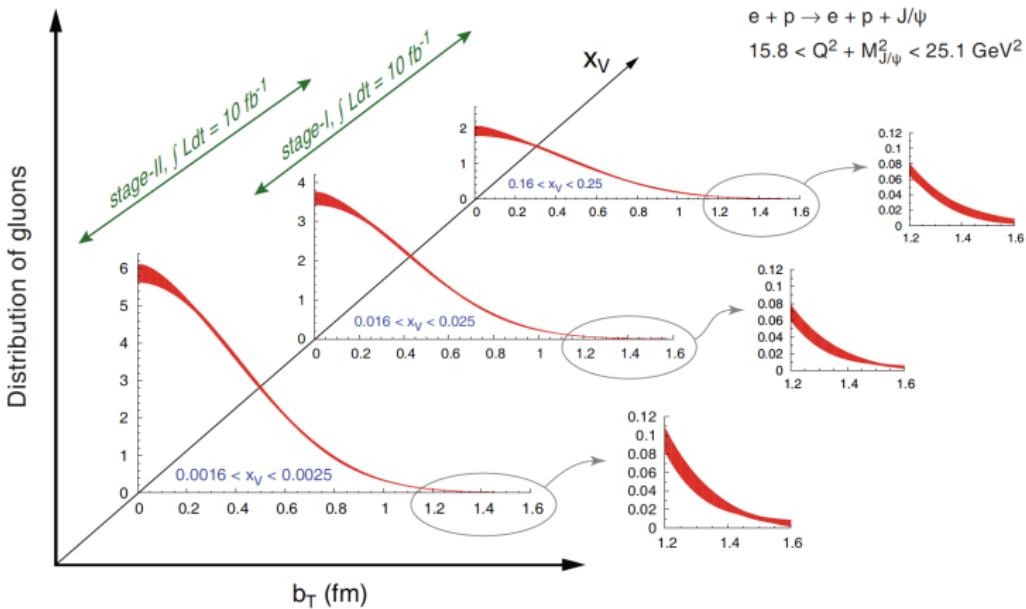


Results of a combined fit of PDFs and FFs for EIC pseudo data
Aschenauer, Borsa, Sassot, Van Hulsen, 1902.10663 by NNPDF

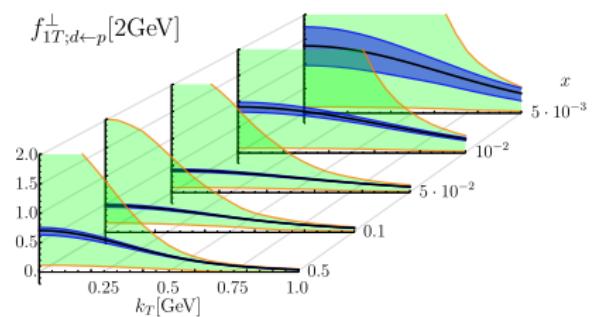
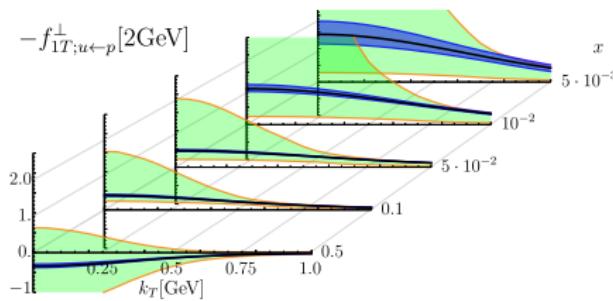
Impact plots from the Yellow Report II: quark GPDs



Impact plots from the Yellow Report III: Gluon GPDs



Impact plots from the Yellow Report IV: quark Sivers distributions



The big question: Can theory match this precision within one or two decades?

If not, you do not really know what you measure.

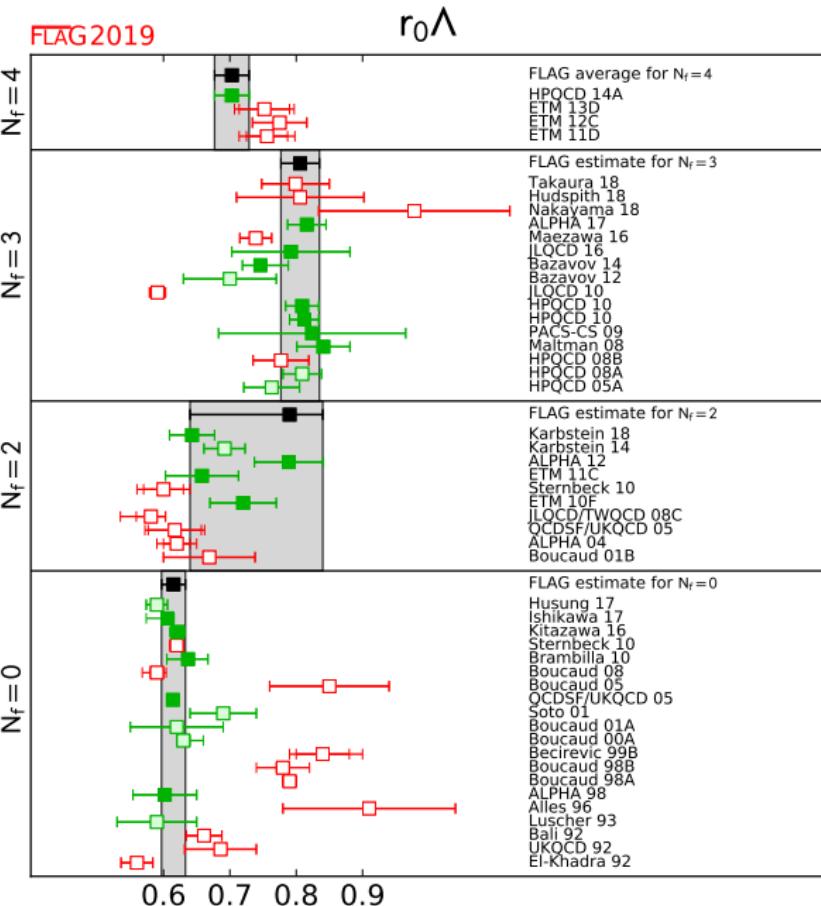
This is all about the control of systematic uncertainties

A paradigm: The FLAG Review

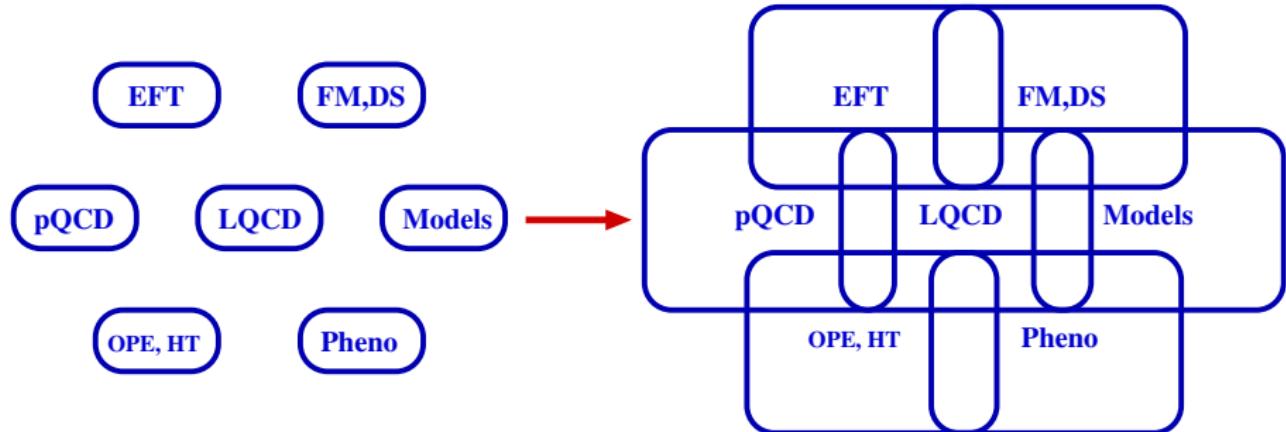
quality criteria:

- full control of systematic errors
- partial control of systematic errors
- no control of systematic errors

Sommer scale $r_0 = 0.472 \text{ fm}$; $r_0^2 \frac{V_{Q\bar{Q}}(r_0)}{dr_0} = 1.65$; $\Lambda = \Lambda_{QCD, \overline{MS}}$



The Changed Theory Landscape



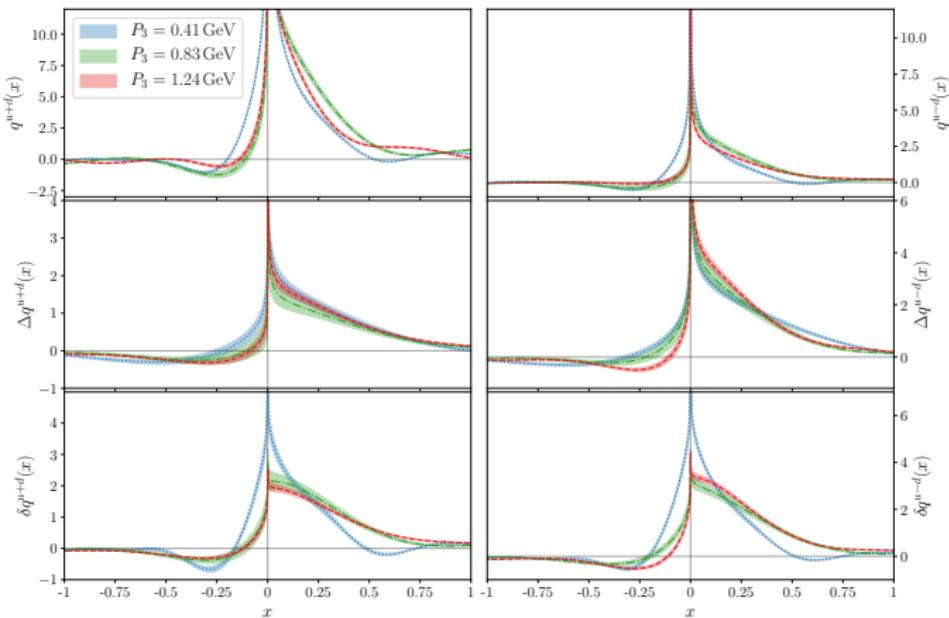
- Different theory subfields match in their regions of overlap
- The very broad EIC physics program needs consistent (!) input of all of them
- Only global collaboration can provide the needed qualified manpower and compute power

With novel lattice approaches formulated in recent years, lattice calculations can play a much more important role.

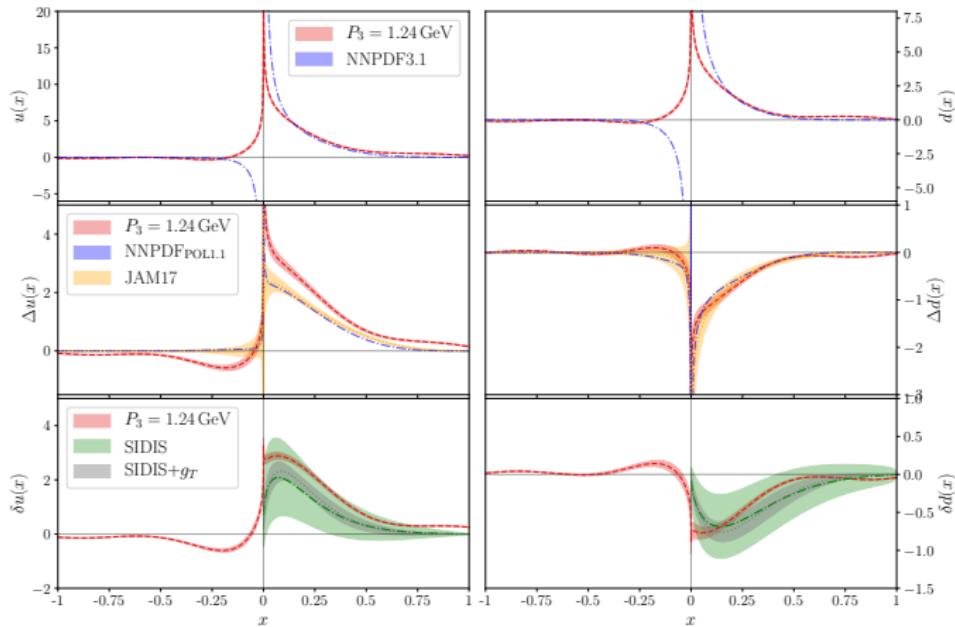
- X. Ji, PRL 110 (2013) 262002; arXiv:1305.1539
quasi-PDFs etc.
- Y.Q. Ma and J.W. Qiu PRD 98 (2018) 074021;
arXiv:1404.6860 lattice cross sections
- A. Radyushkin, PRD96 (2017) 034025; arXiv:1705.01488
pseudo-PDFs etc.
- U. Aglietti et al. Phys.Lett. B441 (1998) 371;
hep-ph/9806277
- V. Braun and D. Müller, EPJ C55 (2008) 349;
arXiv:0709.1348
- A.J. Chambers et al., PRL 118 (2017) 242001;
arXiv:1703.01153
- ...

I am a member of LPC (Y.B. Yang, X. Ji)

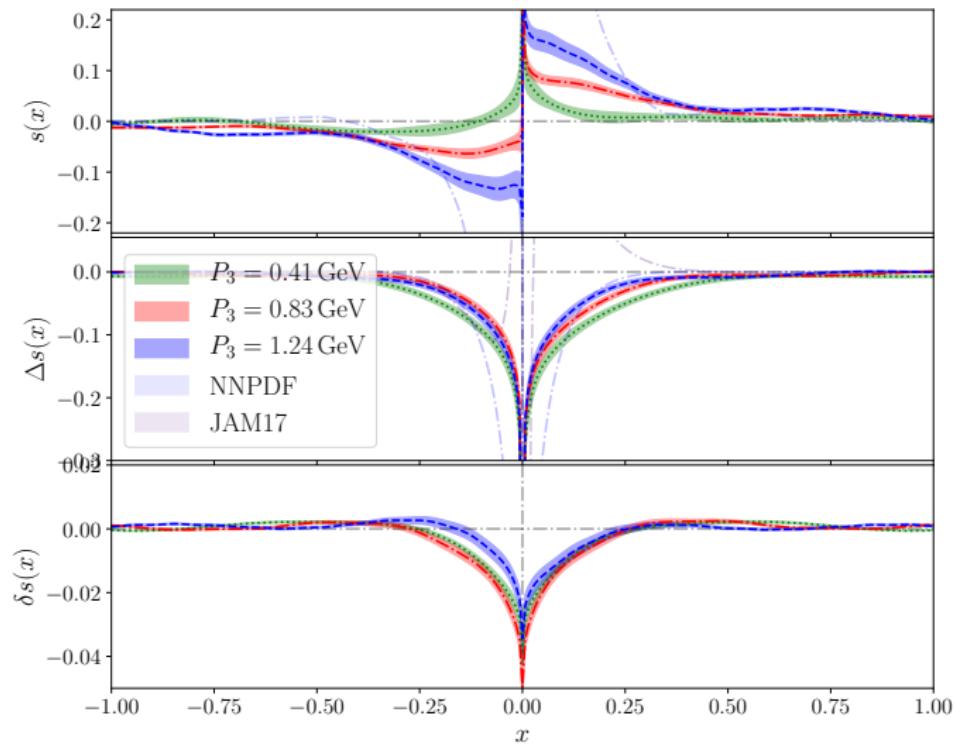
some new results of ETMC 2106.16065



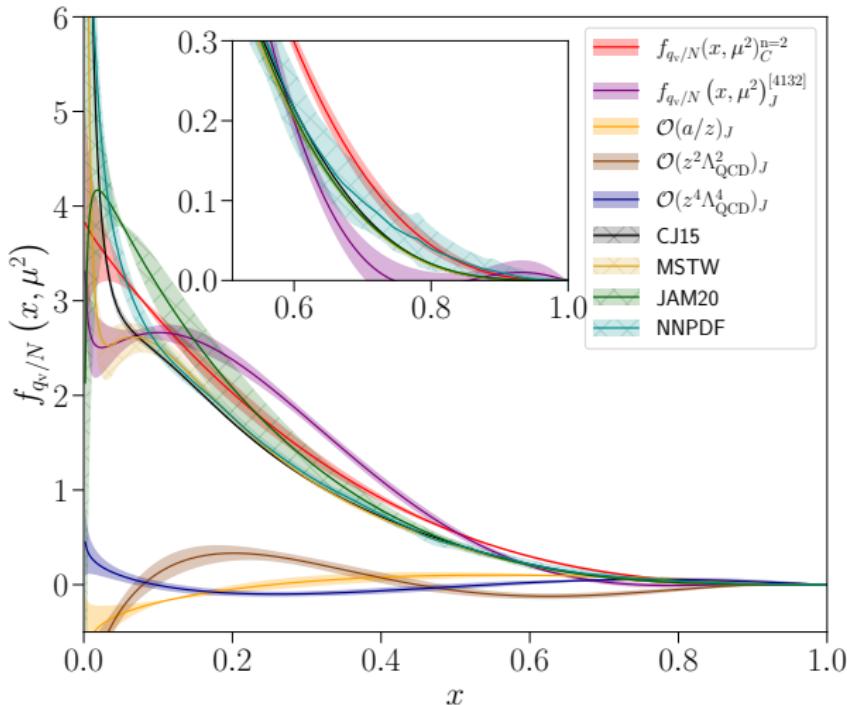
light quark PDFs, quasi-PDFs



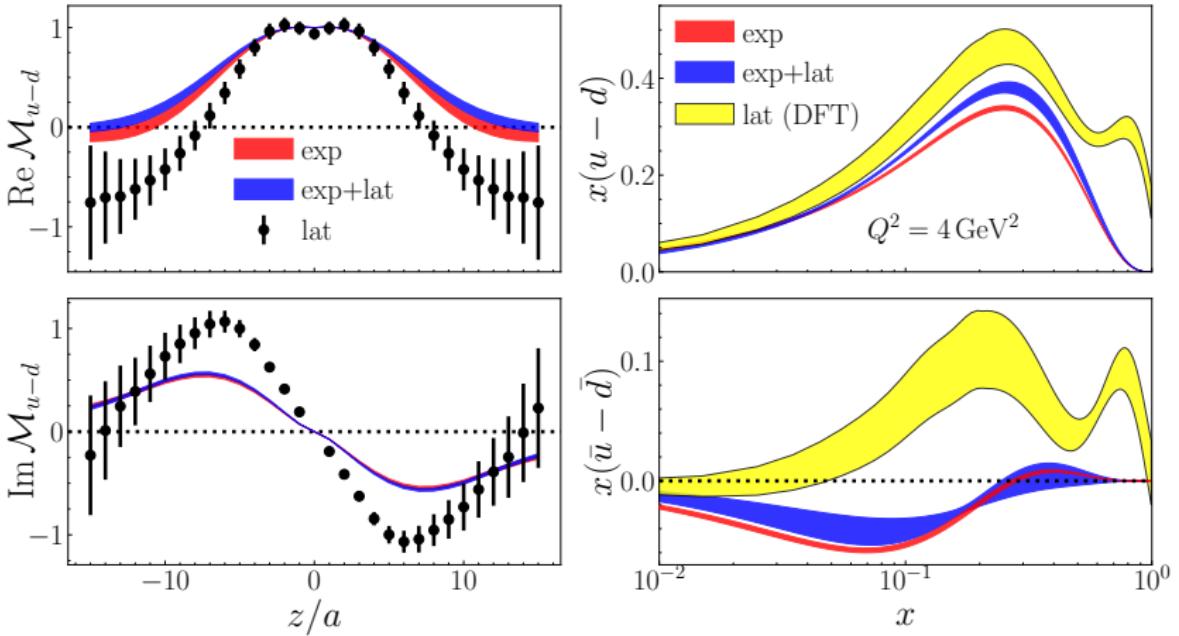
the light quark PDFs



the strange quark PDFs

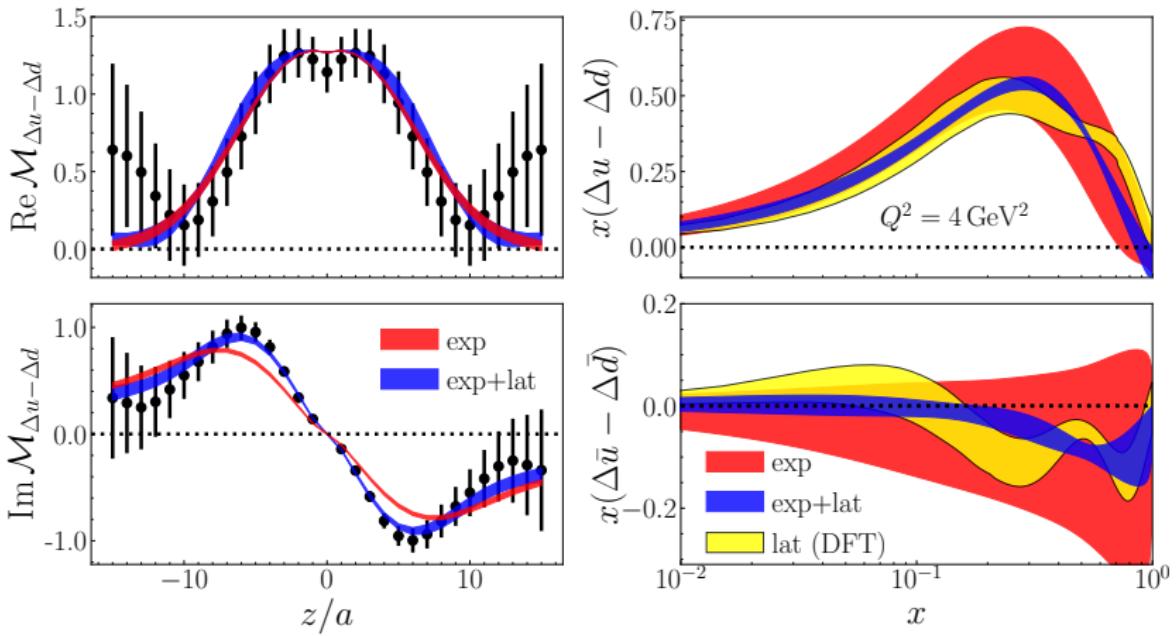


leading twist PDFs plus corrections
Ioffe-Time pseudo Distributions (pITD)



The future will see ever more Lattice augmented PDF fits
 Real and imaginary parts of the isovector unpolarized matrix element M_{u-d} ; quasi-PDFs; the resulting PDFs

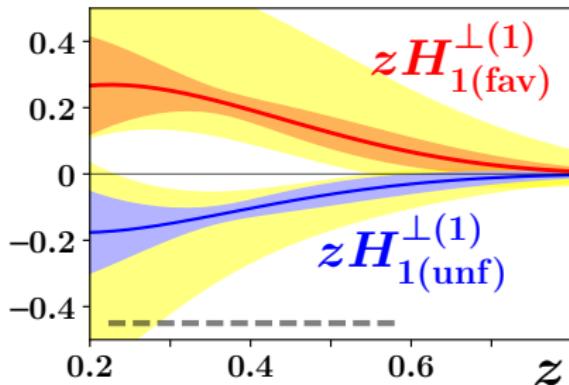
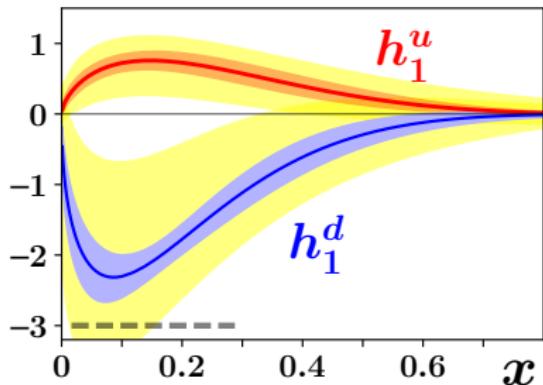
Bringewatt et al. arXiv:2010.00548



Lattice augmented PDF fits

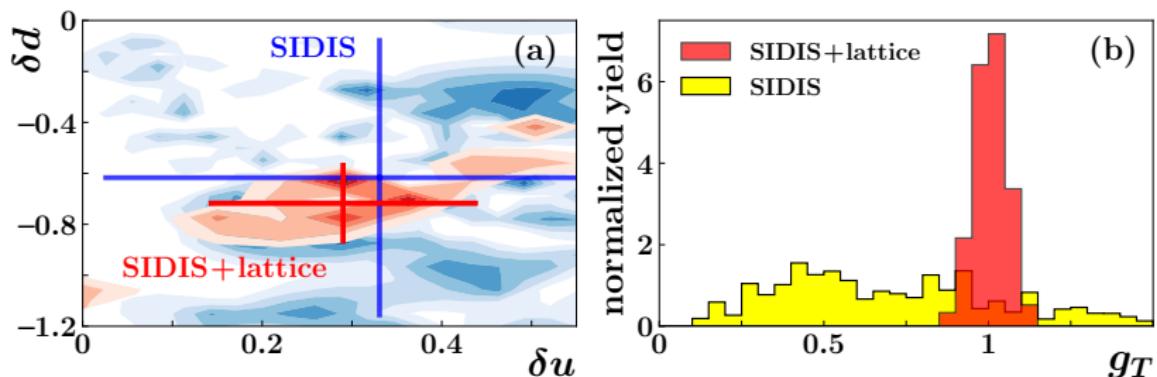
Real and imaginary parts of the isovector polarized matrix element M_{u-d} ; quasi-PDFs; the resulting PDFs

Bringewatt et al. arXiv:2010.00548



Lattice augmented transversity fits from JAM
Lin et al. arXiv:1710.09858

Transversity PDFs $h_1^{u,d}$ and favored $zH_{1(fav)}^{\perp(1)}$ and un-favored $zH_{1(unfav)}^{\perp(1)}$ Collins FFs



Lattice augmented transversity from fits from JAM
 Lin et al. arXiv:1710.09858

Contour plot of δu and δd for the SIDIS only (blue) and SIDIS+lattice(red);
 Normalized yields for the isovector tensor charge g_T

The χ QCD collaboration has very carefully investigated the continuum limit for quasi-PDF operators 2012.05448

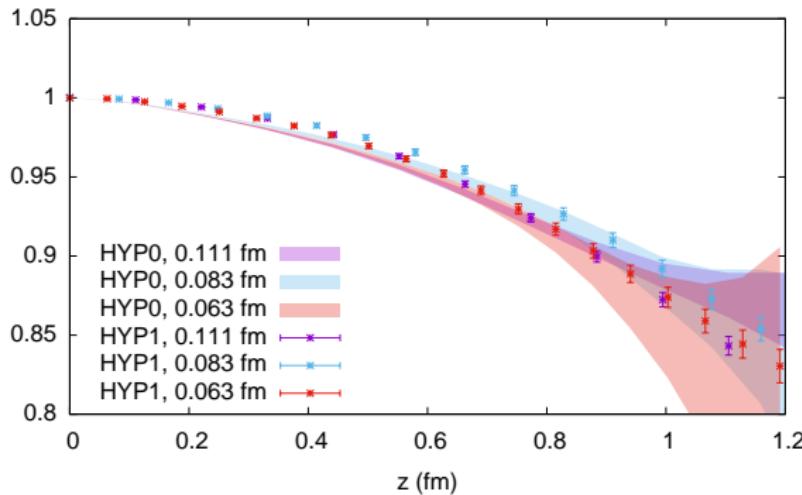
$$O_\Gamma(z) = \bar{\psi}(0)\Gamma U(0, z)\psi(z)$$

The bare $O_\Gamma(z)$ in lattice regularisation at 1-loop level contains an $1/a$ term

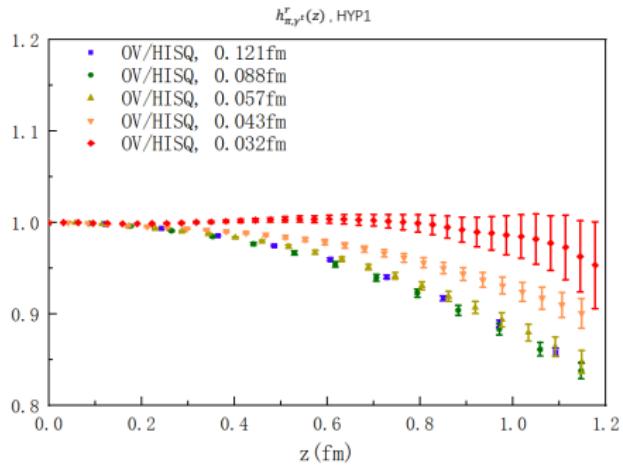
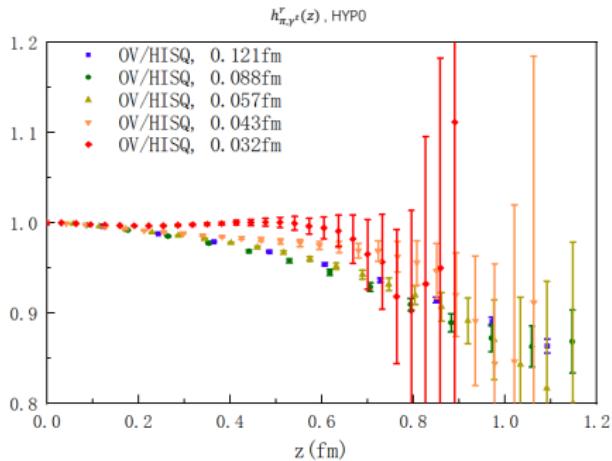
$$O_\Gamma(z) = \Gamma\left(1 + g^2(\gamma \log(p^2 a^2) - m_{-1} \frac{z}{a}) + \dots\right)$$

$$h_{\pi,\Gamma}(z) = \langle \pi | O_\Gamma(z) | \pi \rangle$$

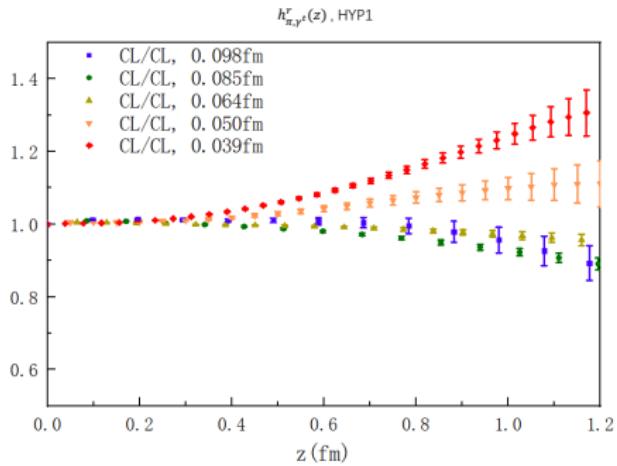
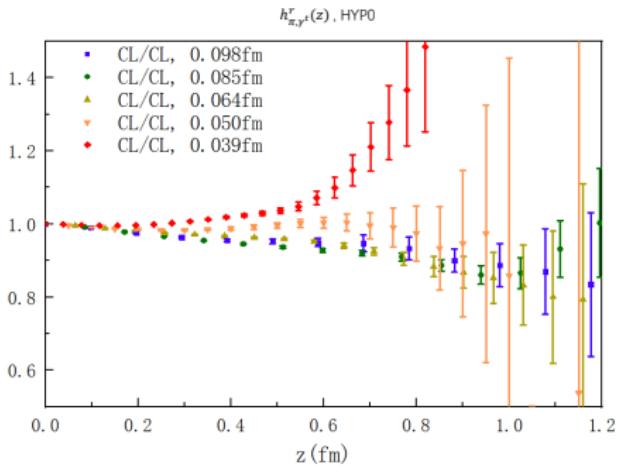
$$h_{\pi,\Gamma}(z) = \frac{\langle O_\pi(T/2) \sum_x (O_\Gamma(z; (\vec{x}, t)) + O_\Gamma(z; (\vec{x}, T-t))) O_\pi^\dagger(0) \rangle}{\langle O_\pi(T/2) O_\pi^\dagger(0) \rangle}$$



the RI/MOM renormalized pion matrix element for overlap valence on DWF sea

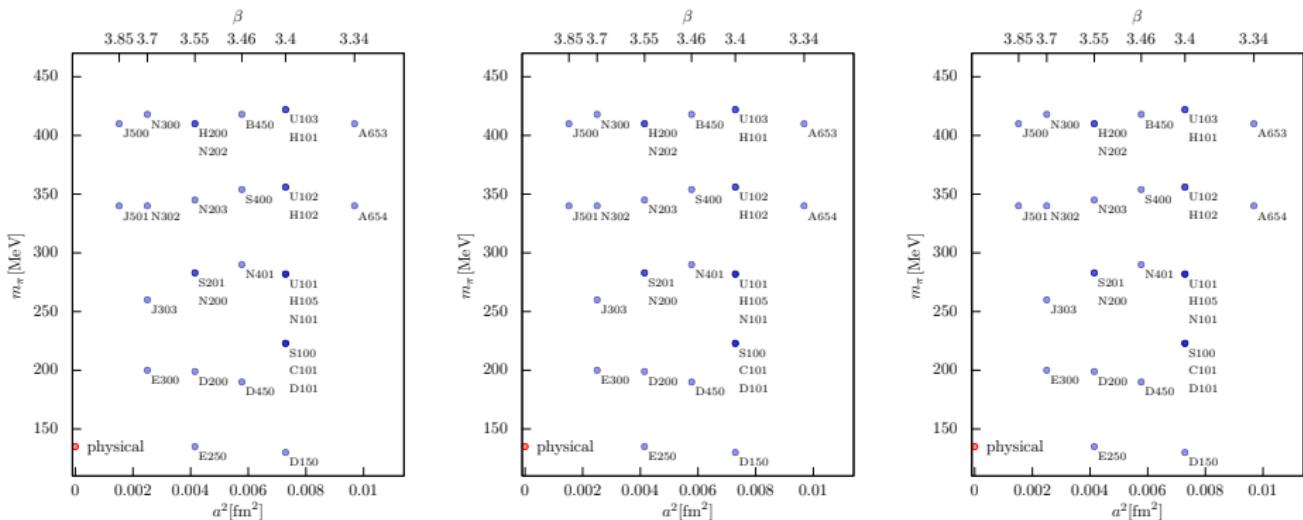


the RI/MOM renormalized pion matrix element for overlap valence on HISQ sea



the RI/MOM renormalized pion matrix element for clover fermions (CLS)

Bottom line: Taking the continuum limit is non-trivial

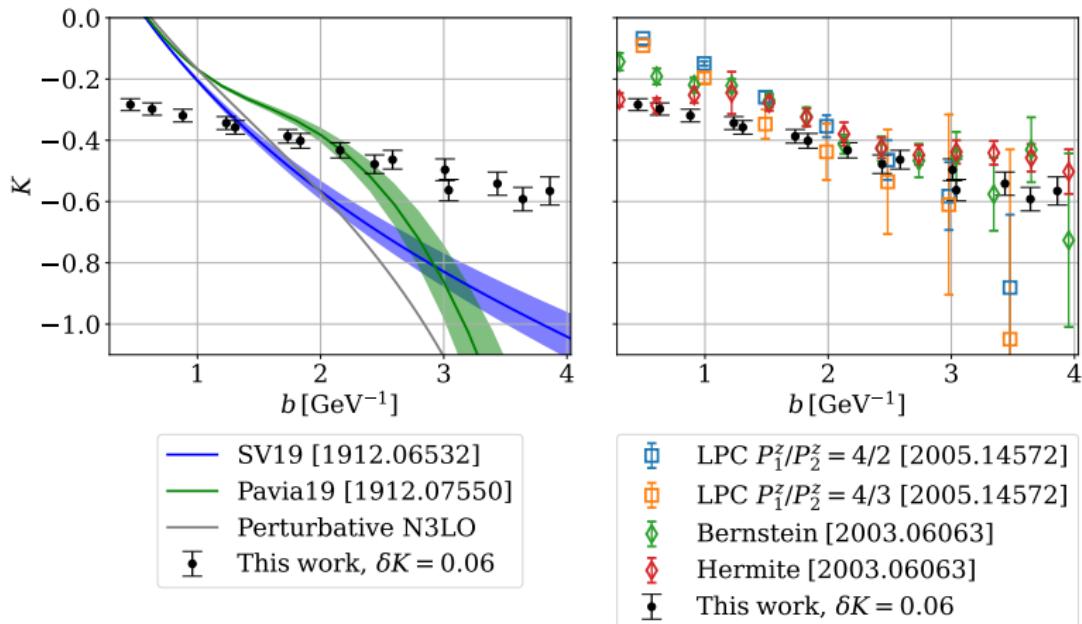


The CLS ensembles

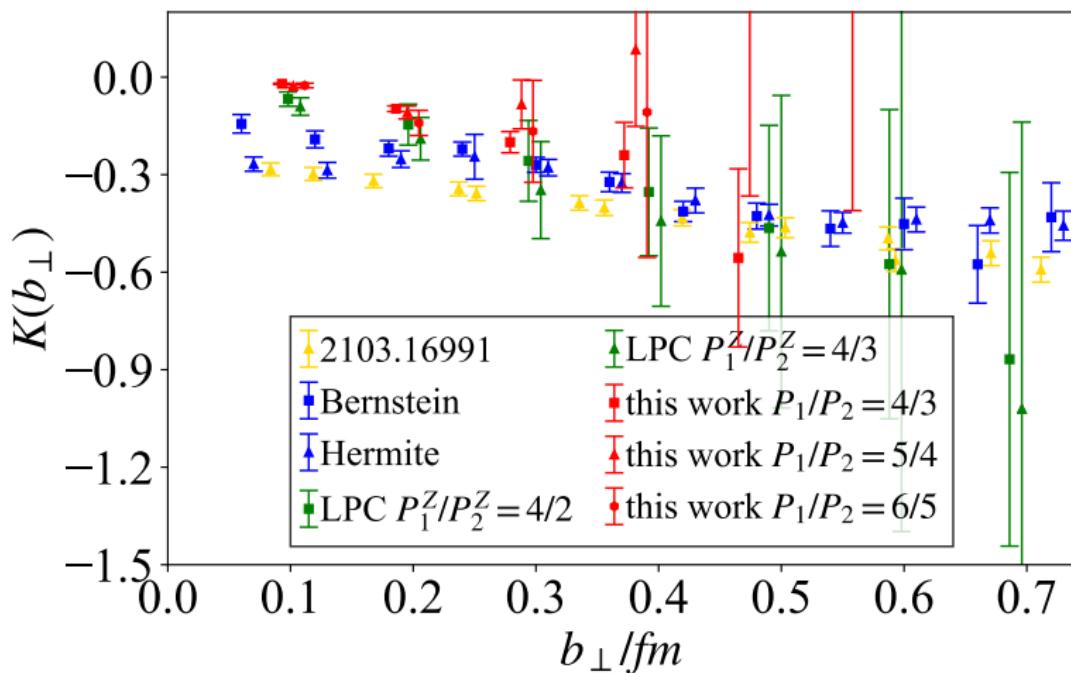
the Collins Soper kernel

A prime example for the speed of progress

- MIT-Fermilab-BNL group P. Shanahan, M. Wagman, Y. Zhao 2003.06063 quasi-PDFs
- LPC: Q.-A. Zhang et al. 2005.14572 LaMET, soft function
- TMDC-RQCD, M. Schlemmer et al. 2103.16991 moment method 2002.07527
- Beijing-ETMC Y. Li et al. 2106.13027 LaMET, soft function



Estimated higher twist uncertainty $\delta K = 0.06$



Four groups using three different lattice approaches get compatible results (most reliable way to estimate systematic uncertainties)

higher twist, power corrections

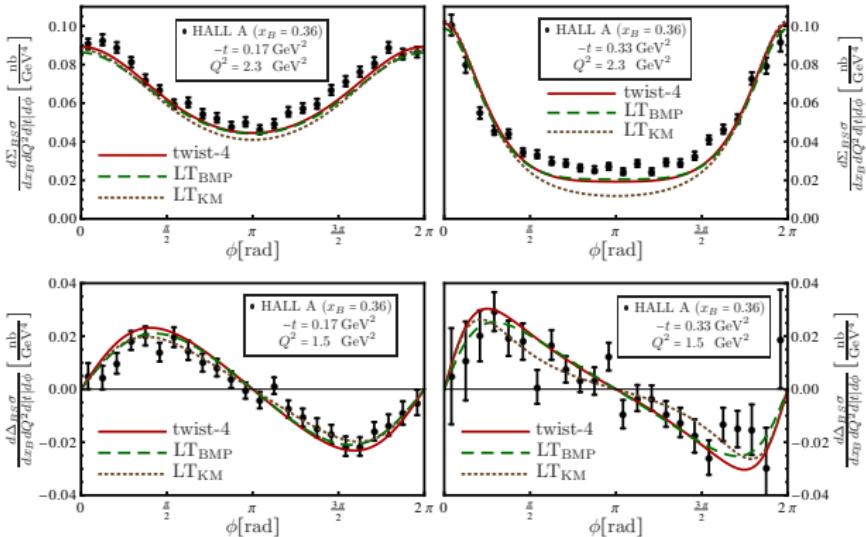
The task:

- disentangle all higher-twist distribution and fragmentation functions
- calculate all kinematic power corrections, e.g. $\sim \frac{m^2}{Q^2}$

One example: DVCS Braun, Manashov, Müller, Pirnay;
1401.7621

$$\begin{aligned}\frac{d\Sigma_{BS}\sigma}{dx_B dQ^2 d|t| d\phi} &:= \frac{1}{2} \left[\frac{d\sigma^\rightarrow}{dx_B dQ^2 d|t| d\phi} + \frac{d\sigma^\leftarrow}{dx_B dQ^2 d|t| d\phi} \right] \\ \frac{d\Delta_{BS}\sigma}{dx_B dQ^2 d|t| d\phi} &:= \frac{1}{2} \left[\frac{d\sigma^\rightarrow}{dx_B dQ^2 d|t| d\phi} - \frac{d\sigma^\leftarrow}{dx_B dQ^2 d|t| d\phi} \right]\end{aligned}$$

and all kinematic corrections $\sim \frac{m^2}{Q^2}$ and $\sim \frac{t}{Q^2}$



The difference between LT_{KM} and LT_{BMP} illustrates the ambiguity in defining leading twist for DVCS.

Power corrections from a renormalon model for quasi-PDFs

Braun, Vladimirov, Zhang, 1810.00048

- power corrections for quasi-PDFs have the generic behavior

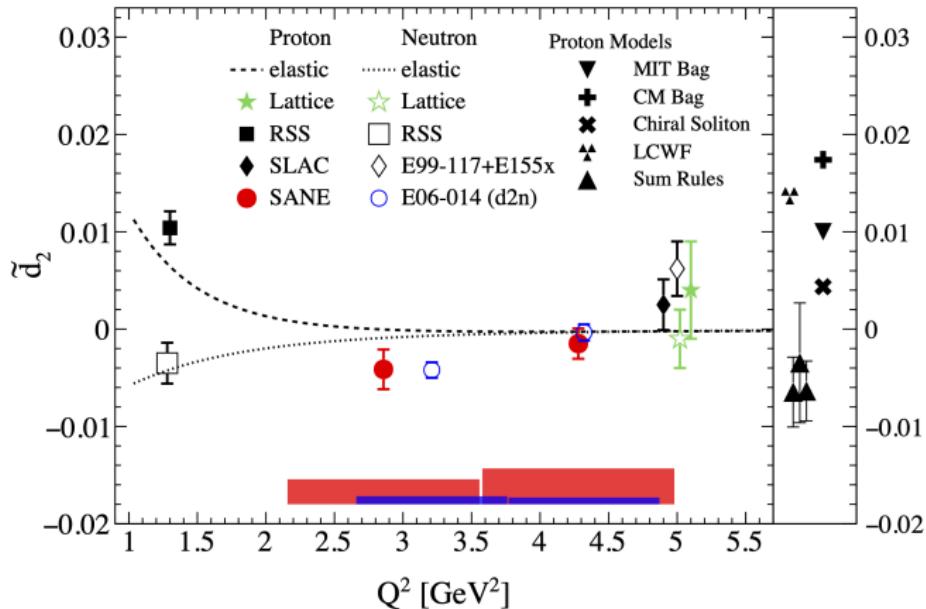
$$\mathcal{O} = q(x) \left\{ 1 + \mathcal{O} \left(\frac{\Lambda^2}{p^2} \frac{1}{x^2(1-x)} \right) \right\}$$

- power corrections for pseudo-PDFs have the generic behavior

$$\mathcal{O} = q(x) \left\{ 1 + \mathcal{O} \left(z^2 \Lambda^2 (1-x) \right) \right\}$$

bottom line: Interesting results, but you only know the size of power and higher-twist corrections after calculating them, which is possible on the lattice but extremely hard

the present status of \tilde{d}_2 SANE 1805.08835



Our 2005 lattice results look good, but they included no systematic errors.

Experiments can only measure d_2^p and d_2^n . LQCD gives also d_2^u and d_2^d which can be translated into the color force acting on a quark according to Burkhardt arXiv:0810.3589 (factor of 4 wrong)

$$\int_0^1 dx \, x^2 g_1^p(x, Q^2) = \frac{a_2^p}{2} + \mathcal{O}\left(\frac{m_N^2}{Q^2}\right)$$

$$\int_0^1 dx \, x^2 g_2^p(x, Q^2) = \frac{d_2^p - a_2^p}{3} + \mathcal{O}\left(\frac{m_N^2}{Q^2}\right)$$

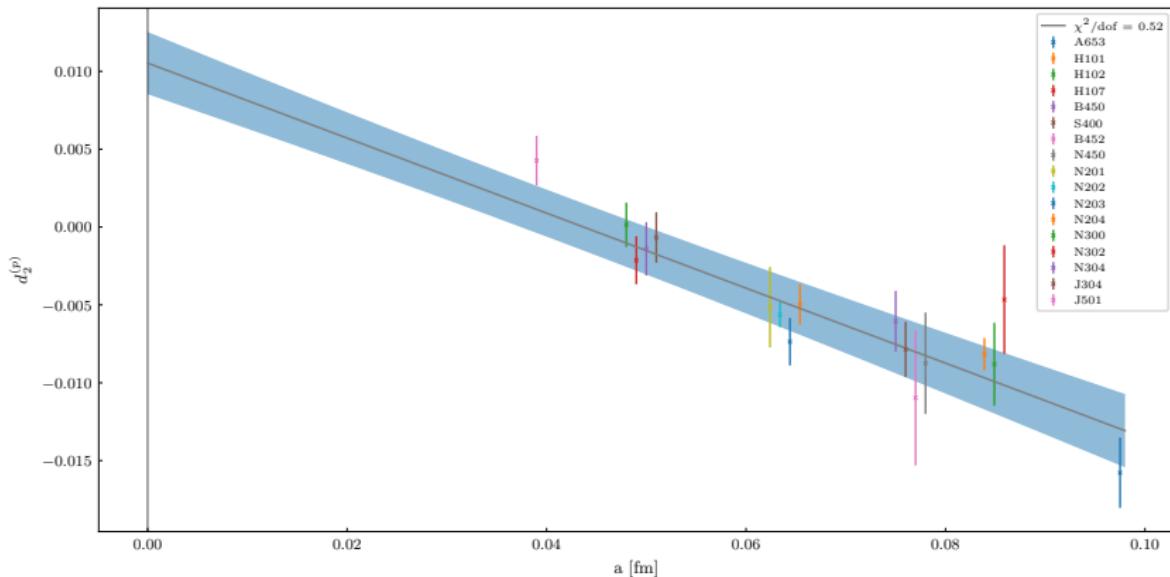
$$d_2^q = -\frac{1}{2M(P^+)^2} \langle P_N, S_T | \bar{q}(0)\gamma^+ g G^{+y}(0) q(0) | P_N, S_T \rangle$$

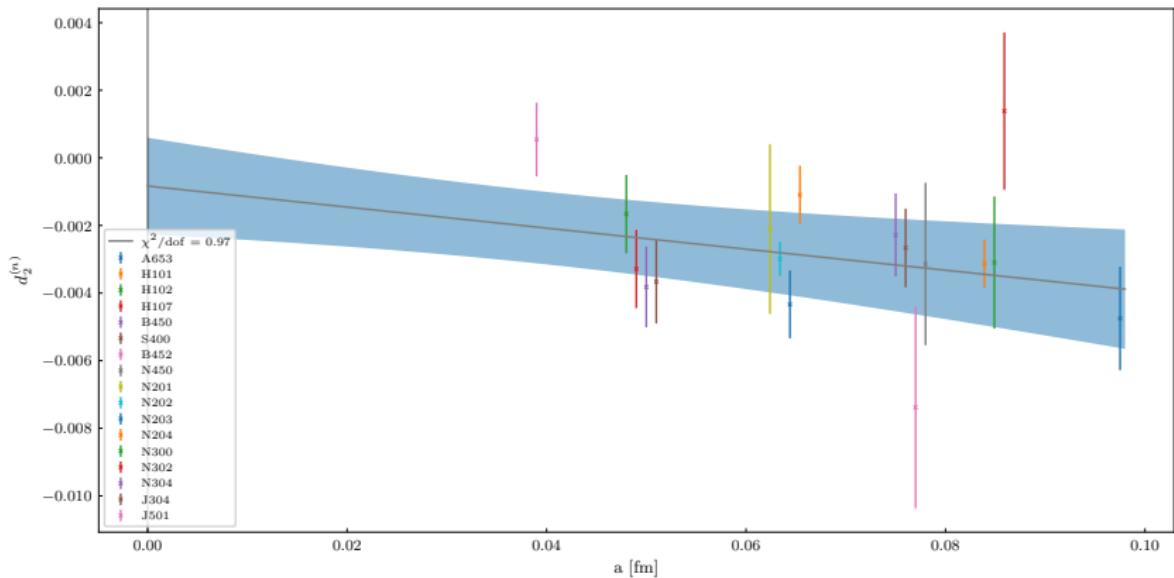
$$F^u = -25 \pm 30 \text{ MeV/fm} = -0.04 \pm 0.05 (\Lambda_{QCD})^2$$

$$F^d = -14 \pm 13 \text{ MeV/fm} = 0.02 \pm 0.02 (\Lambda_{QCD})^2$$

Can this be true?

d_2^p

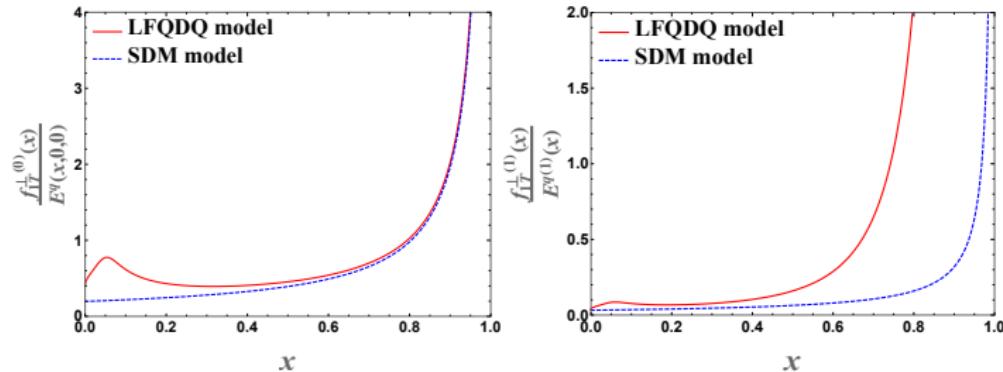


d_2^n 

models

It will take quite some time before LQCD can match the experimental precision. Therefore, models will continue to be important.

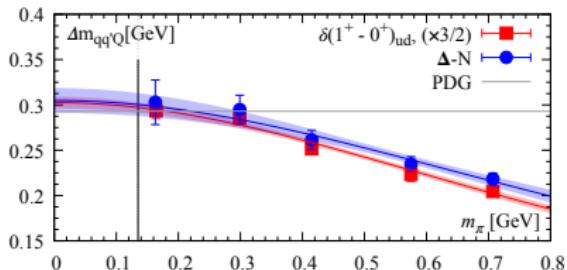
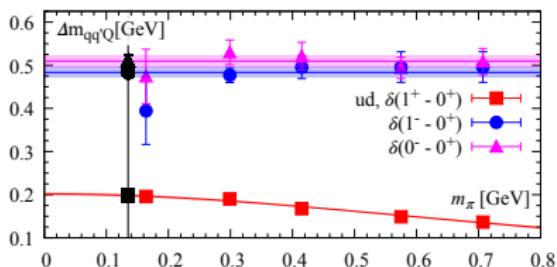
For example, one can relate GPDs to TMDs in quark diquark models, e.g. B. Pasquini et al. 1907.06960 B.Gujar et al. 2107.02216



Ratio of moments of the up quark Sivers function to the up quark GPD E^u

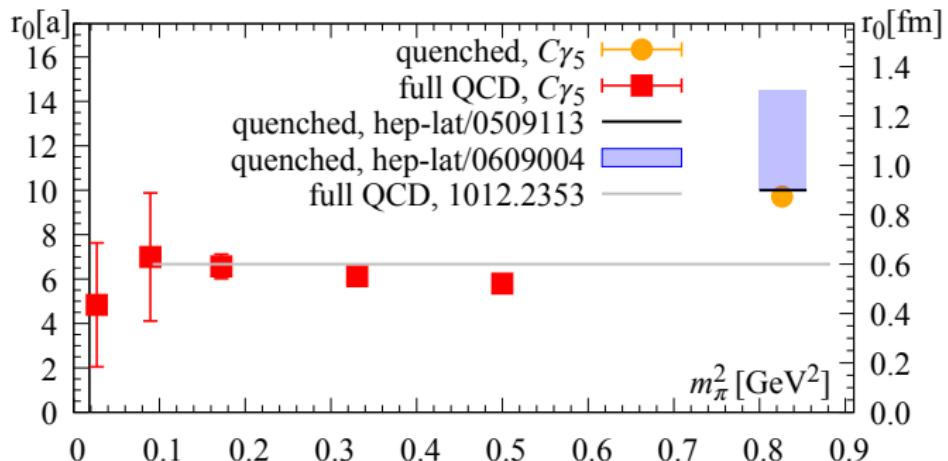
Diquarks from lattice QCD Francis, de Forcrand, Lewis,
Maltmann 2106.09080 2+1 Clover fermions down to
 $m_\pi = 164$ MeV

A LQCD analysis of the $Q(q^c C \Gamma q)$ system.



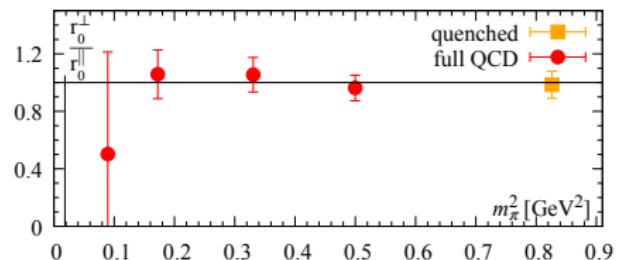
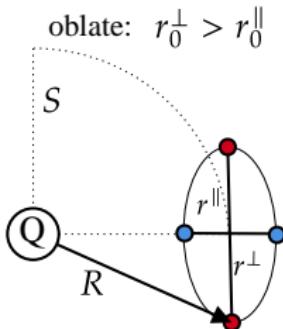
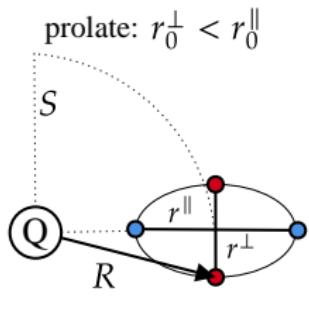
the 'good' diquark 0^+ dominates

One can study the geometry of the diquark

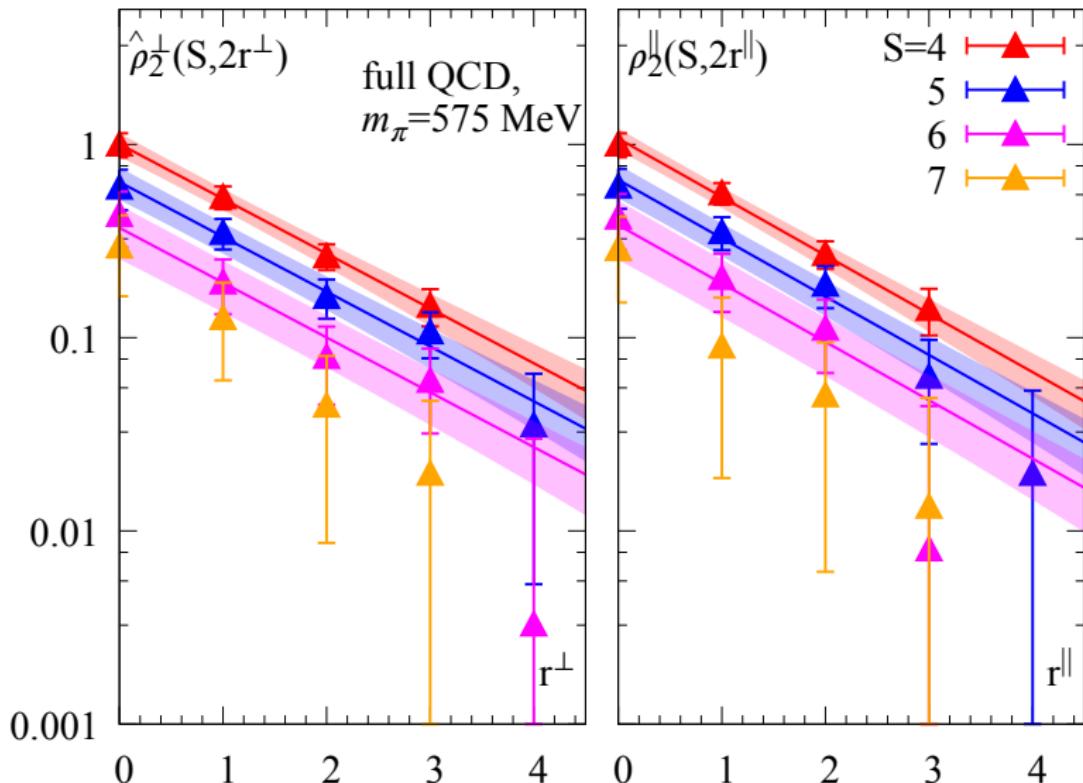


the diquark radius is roughly 0.6 fm

$$C_\Gamma(\vec{x}_1, \vec{x}_2, t) = \langle \mathcal{O}_\Gamma(\vec{0}, 2t) \rho(\vec{x}_1, t) \rho(\vec{x}_2, t) \rangle \langle \mathcal{O}_\Gamma(\vec{0}, 0)$$

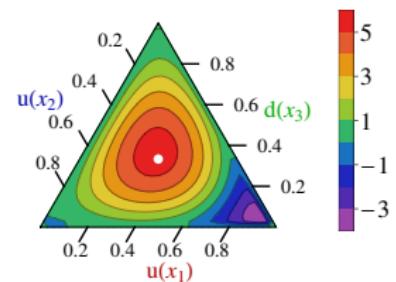
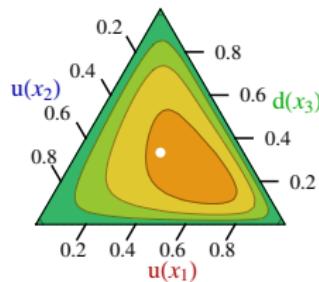
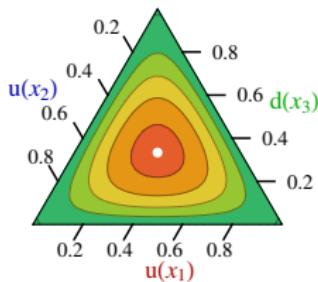


diquarks are roughly spherical



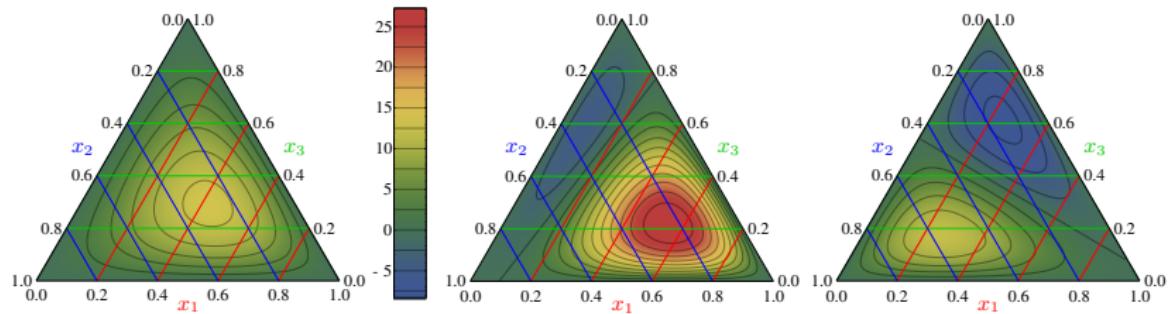
another example: Hadron distribution amplitudes

C. Mezrag, J. Segovia, L. Chang, C.D. Roberts; 1711.09101;
Fadeev-type 3 particle wave function



asymptotic, nucleon and roper wave function

V.Braun et al. 1403.4189 ; Lattice QCD



nucleon, $N^*(1650)$ and $N^*(1535)$

nucleon mass decomposition

The scale anomaly of QCD is of fundamental importance for the understanding of QFT.

Q1: Does it contribute to m_N ?

Ji hep-ph/9410274; Lorcé 1706.05853; Metz, Pasquini, Rodini 2006.11171; ...

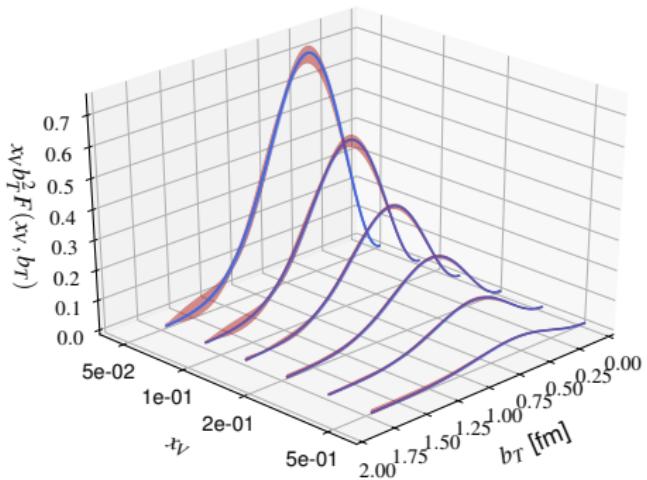
Q2: If so, can it be isolated from experiment or Lattice QCD?

Kharzeev, Satz, Syamtomov, Zinovjev hep-ph/9901375, near threshold quarkonium photoproduction; He, Sub, Yang 2101.04942 LQCD

The answers to both questions are still open

Experimentally, determining the gluon GPD from exclusive Υ production looks very promising

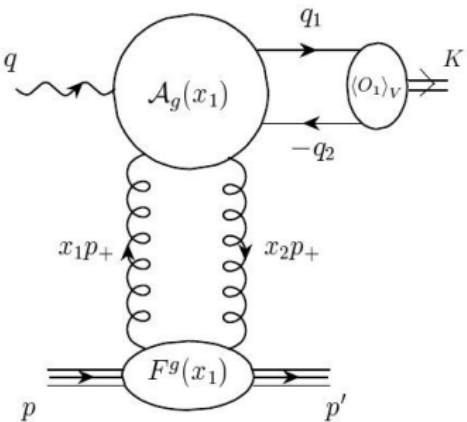
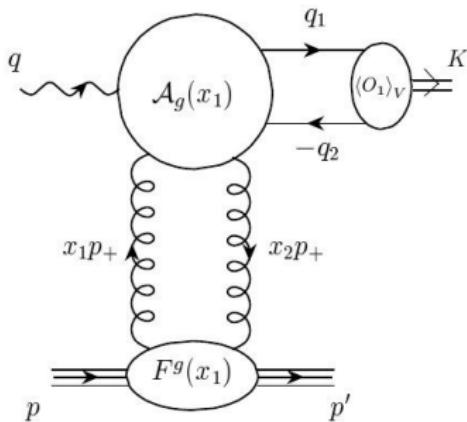
Yellow Report; Joosten and Meziani 1802.02616



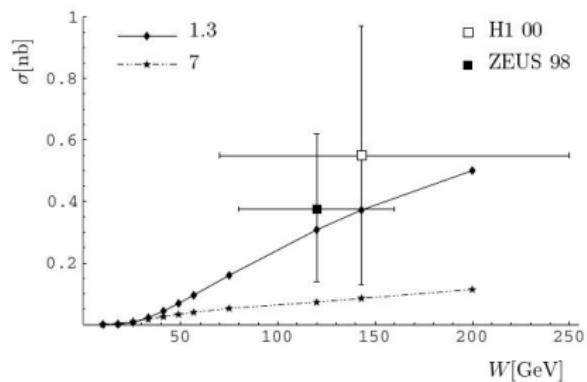
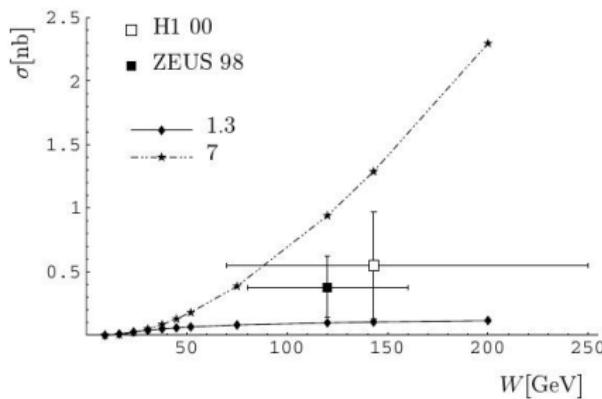
The problem is the control of systematic theoretical uncertainties:

A full NLO calculation of exclusive J/ψ and Υ photoproduction

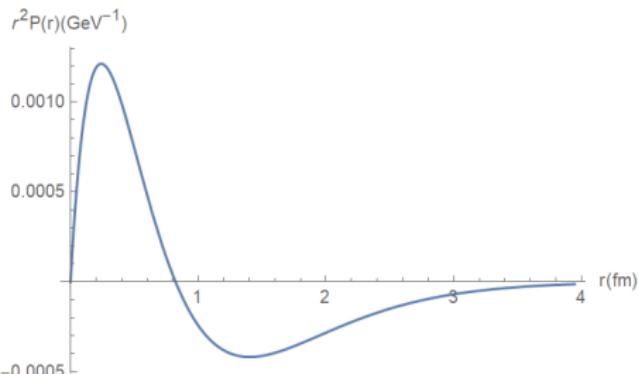
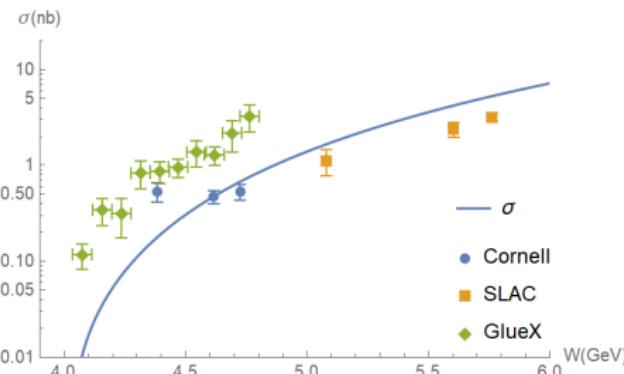
Ivanov, Schäfer, Szymanowski and Krasnikov, hep-ph/0401131



LO (left) and NLO (right) results for Y photoproduction



Guo, Ji, and Liu 2103.11506 argue that factorization should stay valid in the limit of threshold production and use lattice input for the gravitational form factors from Shanahan and Detmold 1810.04626 to fit J/ψ GlueX data and extract the gluon pressure inside a proton.



The J/ψ photoproduction cross-section; prediction using the gravitational formfactor from the lattice (line), and the gluon contribution to the pressure inside a proton

and there is much more:

Hatta and Yang 1808.02163

Sun, Sun, and Zhou 2012.09443

Wang, Kou, Xie, and Chen 2102.01610

Zahed 2102.08191

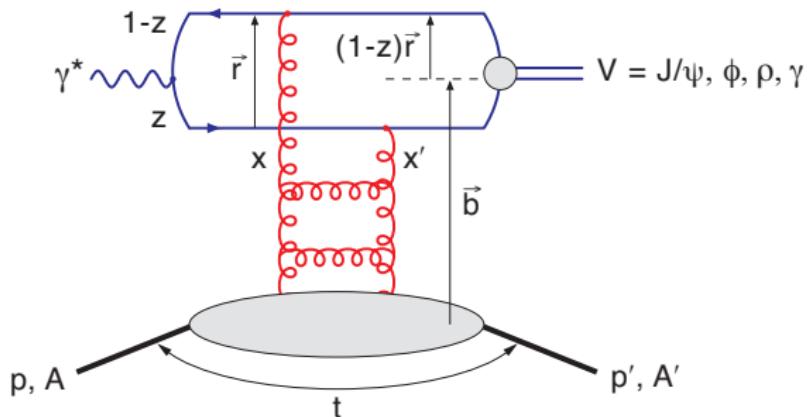
Kou, Wang and Chen 2103.1001

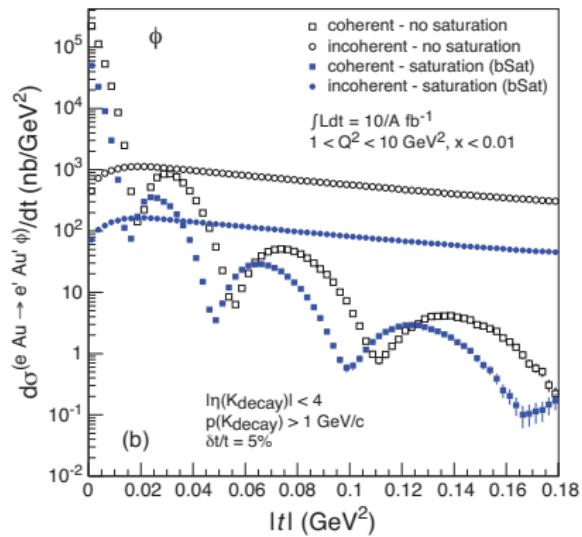
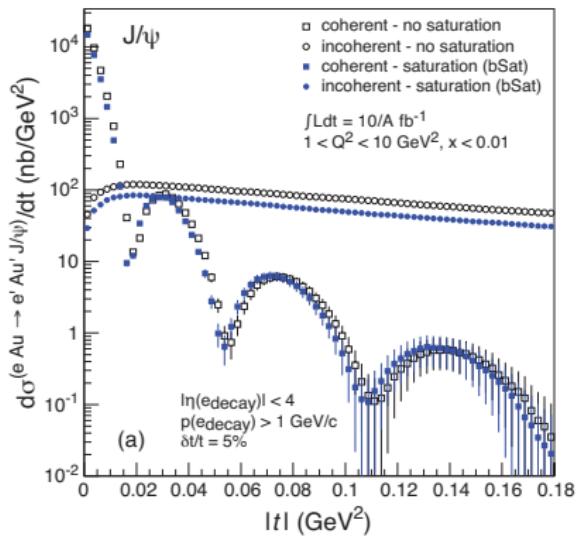
Sun, Tong and Yuan 2103.12047

K.-F. Liu 2103.15768

Freese and Miller 2104.03213

Saturation is a key element of theoretical Heavy-ion physics.
At EIC it can be probed, e.g., in diffractive meson production
Toll and Ullrich, arXiv:1211.3048

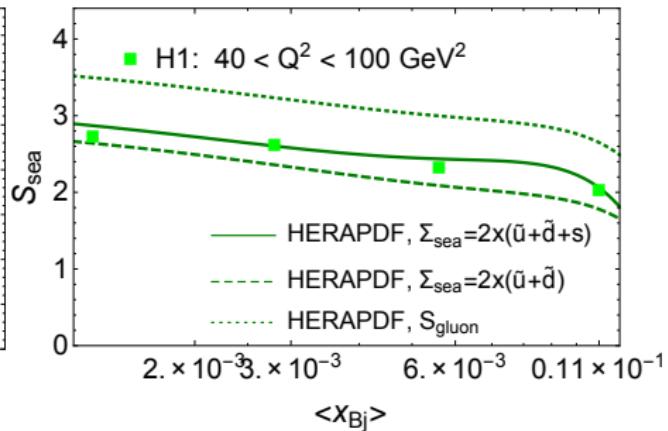
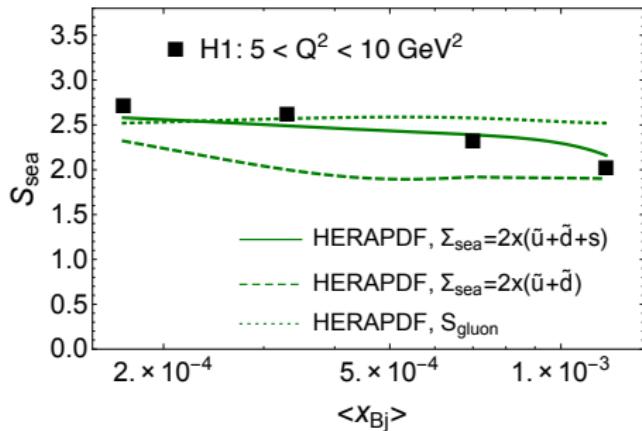




diffraction J/ψ and Φ production

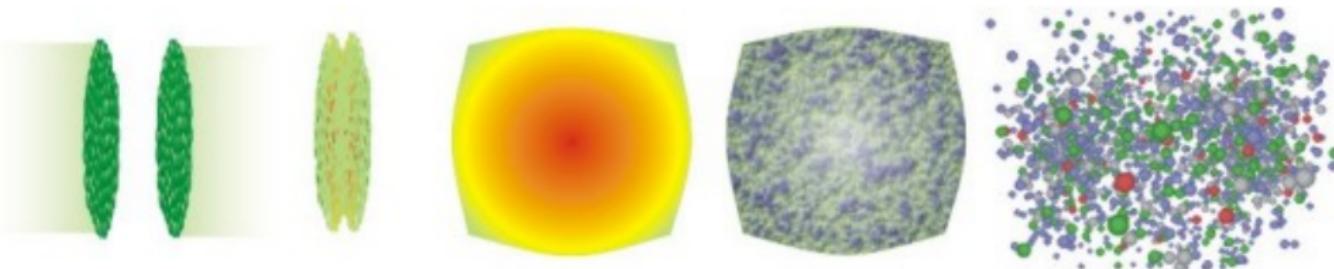
An information theoretical interpretation of DIS

Kharzeev and Levin arXiv:2102.09773

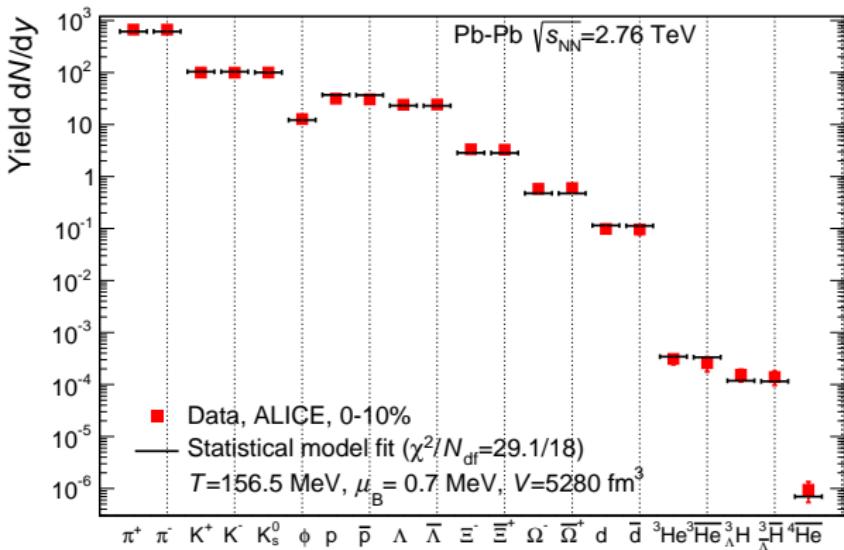


The entropy of partons resolved by DIS

QED and QCD, i.e. heavy-ion collisions and e+A are time reversal invariant no information loss, no entropy production, an entangled state

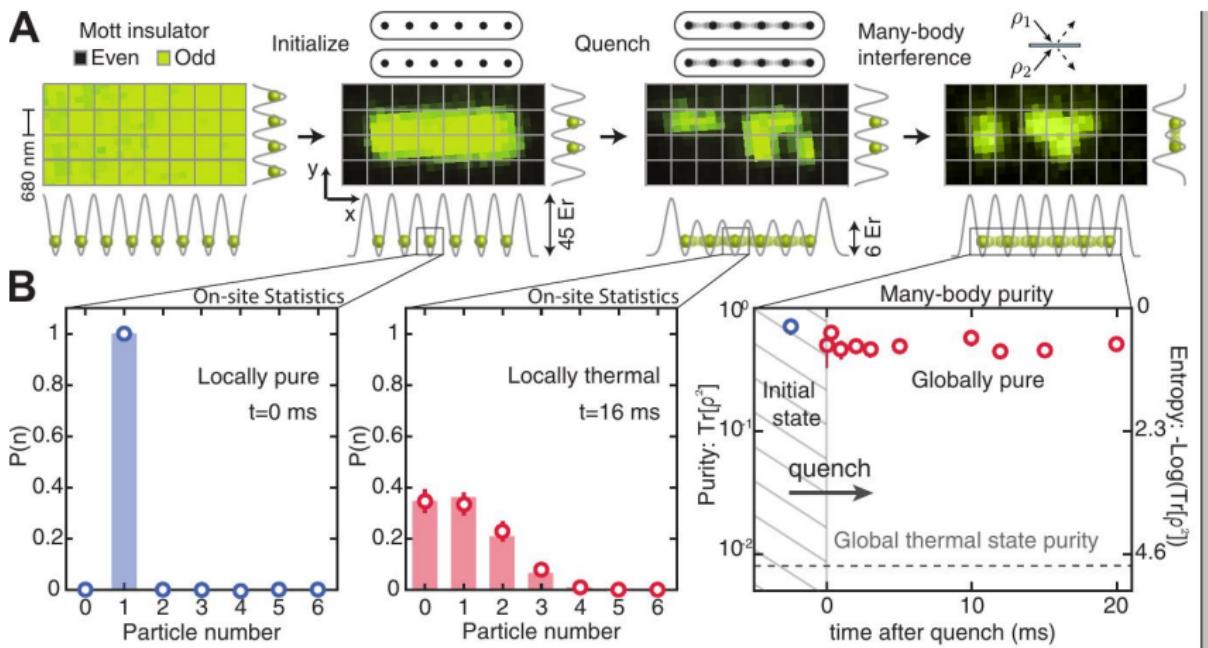


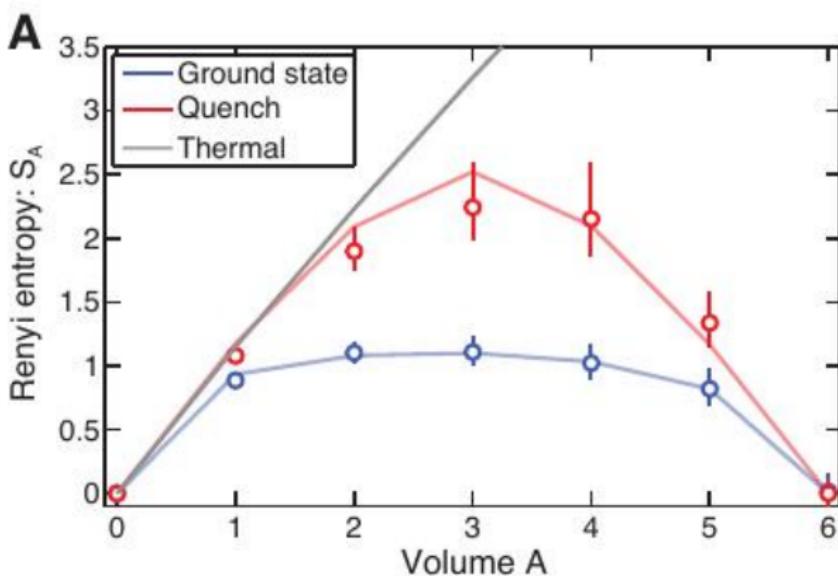
rapid thermalization and vorticity relaxation could be related to entanglement; EPR: Naive causality does not apply for entangled states; decoherence depends on the observable



But: $R(\text{rms}, {}^3\text{H}) = 10.6 \text{ fm} \sim 2R_{\text{Pb}}$; $E_B \ll 156 \text{ MeV}$ the yield should be suppressed

The experiment 1603.04409 “Quantum thermalization through entanglement in an isolated many-body system”





The message: Entanglement vs. statistics depends on the observable for all isolated quantum systems: $e+p$, $e+A$, $p+A$ and $A+A$
 a new field of research for the EIC ??

Conclusions

- At EIC accuracy the fine-print of theory calculations becomes crucial: LO vs. NLO vs. NNLO, w/o resummation, higher twist contributions, continuum extrapolation for Lattice calculations, solid all order proof for factorization or just plausibility arguments, ... there is no longer room for hand-waving arguments.
- On the other hand, the theory landscape has qualitatively changed over the last decade such that we can meet this challenge but only if all qualified manpower is combined. Even closer international collaboration is the key.
- Lattice QCD promises to play an ever more important role.