

# Quarkonia and heavy-light mesons in a Bethe-Salpeter approach

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# Light-Cone Distribution Amplitudes

QCD factorization in B decays involves matrix elements which are convolution integrals:

$$\langle \pi^+ \pi^- | (\bar{u}b)_{V-A} (\bar{d}u)_{V-A} | \bar{B}_d \rangle \rightarrow \int_0^1 d\xi du dv \Phi_B(\xi) \Phi_\pi(u) \Phi_\pi(v) T(\xi, u, v; m_b)$$

The integrals are over a (hard) scattering kernel  $T(\xi, u, v, m)$  and light-cone distribution amplitudes (LCDA) expanded in Gegenbauer polynomials:

$$\left[ \begin{array}{l} \varphi_\pi(x) \\ \varphi_\pi^{\text{asy}}(x) = 6x(1-x) \end{array} \right]$$

- LCDA were poorly known for light mesons, in recent years improved determinations of the first two Gegenbauer moments of the pion and kaon, [RQCD Collaboration, Bali et al. \(2019\)](#).
- Next to nothing was known about heavy-light mesons, mostly models and asymptotic LCDA used.

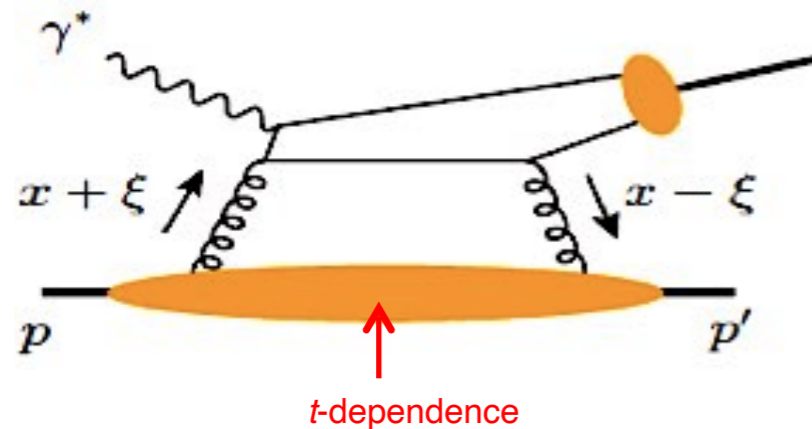
# Spatial imaging of glue in a nucleon/nucleus

Exclusive electroproduction: hard-scattering mechanism  $E_\gamma > 10$  GeV

J. C. Collins, L. Frankfurt, M. Strikman, Phys.Rev.D 56 (1997)

S.J. Brodsky, E. Chudakov, P. Hoyer, J.M. Laget, Phys. Lett. B (2001).

D.Y. Ivanov, A. Schäfer, L. Szymanowski and G. Krasnikov, Eur. Phys. J. C 34 (2004)



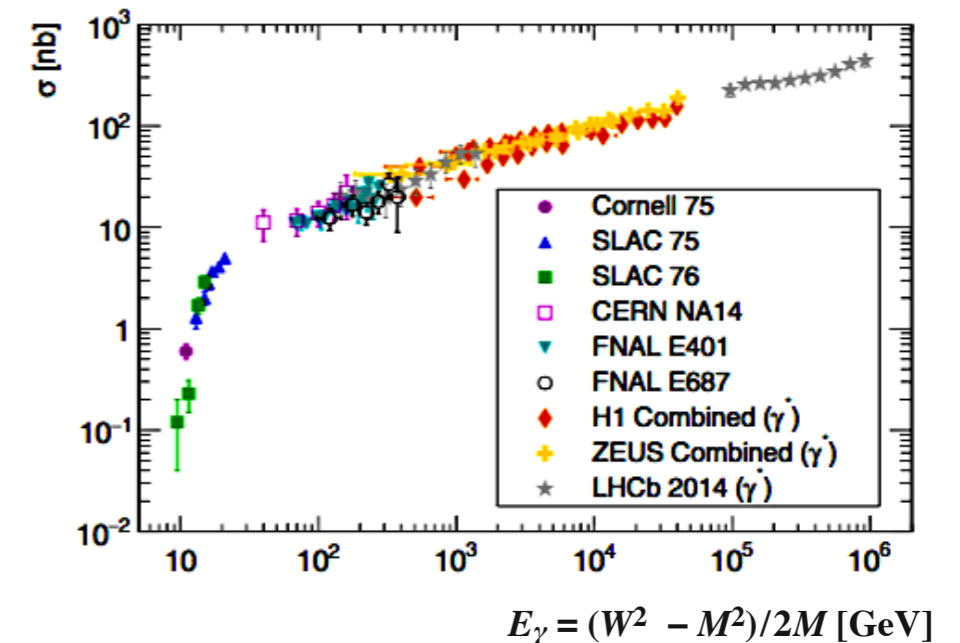
$J/\Psi, \Phi, \dots$

$$Q^2 = -(p_e^i - p_e^f)^2$$

$$W = \sqrt{2v \cdot M_p + M_p^2 - Q^2}$$

$$t = (p_p^i - p_p^f)^2 = (p_\gamma - p_{J/\Psi})^2$$

$$\xi = (P' - P) \cdot n/2$$



- Exclusive *charmonium* production: narrow quarkonium exchanges gluons with the nucleon's light quarks.
- Cross section is proportional to the square of the gluon density in the hadron.
- At large virtualities  $Q^2 \gg m_Q^2$ , electroproduction of heavy mesons is expressed in terms of *meson distribution amplitudes* in the factorization theorem (Collins, Frankfurt & Strikman 1997).

See also talks by Andreas Schäfer and Bowen Xiao

# Charged current production of light and charmed mesons

- Initial work by Marat Siddikov and Iván Schmidt for pion and rho production, Phys. Rev. D 99 (2019).
- Claim that their numerical estimates show that production could be studied in neutrino- or electron-induced processes at EIC.
- Recently extended to  $D_s$  and  $D_s^*$  mesons in collinear QCD where generalized gluon distributions factorize from perturbatively calculable coefficient functions.
- *Required input:* pseudoscalar and vector meson *distribution amplitudes*.

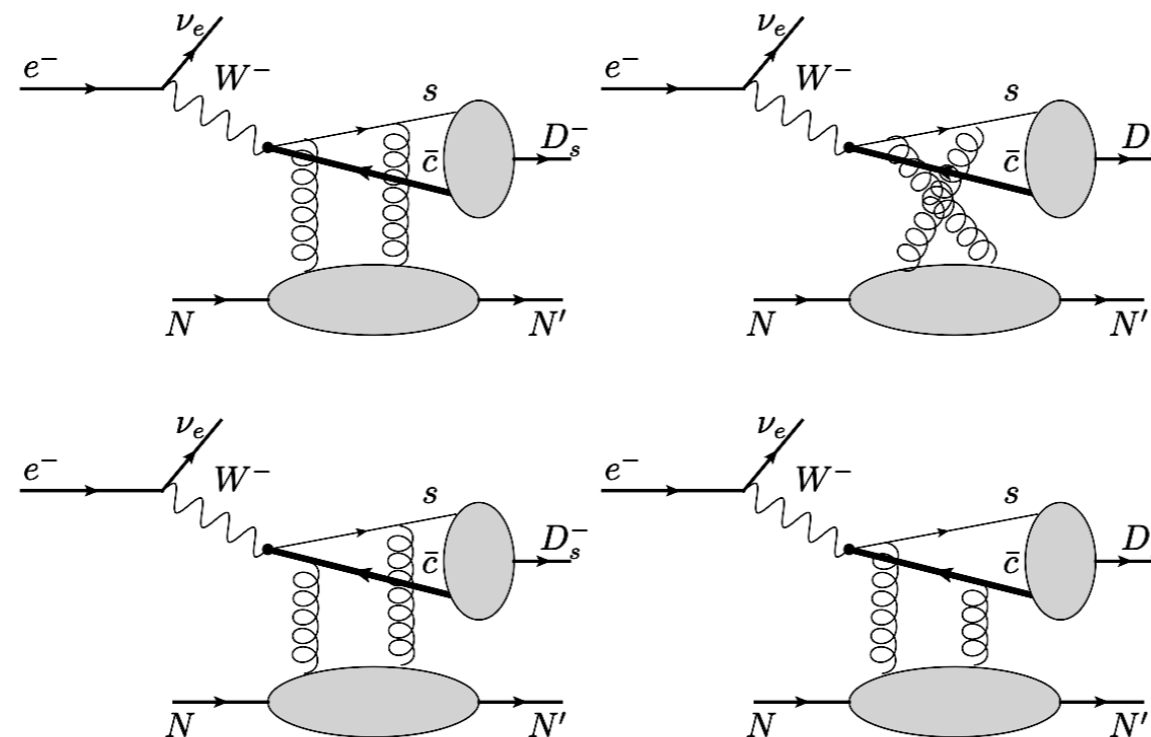


FIG. 1: Feynman diagrams for the factorized amplitude for the  $e^- + N \rightarrow \nu_e + D_s^- + N'$  process involving the gluon GPDs; the thick line represents the heavy anti-quark  $\bar{c}$ .

# Bethe-Salpeter Equation for QCD Bound States

Diagrammatic representation of the Bethe-Salpeter equation. The left side shows a quark propagator with a self-energy correction (loop diagram) equal to the sum of a bare propagator and a loop diagram with a gluon exchange. The right side shows the transition kernel  $T$  (yellow circle) equal to the sum of a single gluon exchange  $K$  (blue rectangle) and a ladder diagram with two  $K$ 's and a  $T$  kernel.

Rainbow-ladder truncation (leading symmetry-preserving approximation)

Diagrammatic representation of the rainbow-ladder truncation. The top row shows the quark propagator  $S(p)$  as a sum of a bare propagator and diagrams with one, two, and three gluon exchanges. The bottom row shows the four-point kernel  $G^4(k, q, P)$  as a sum of diagrams with one, two, and three gluon exchanges.

- Quark propagators are obtained by solving the gap equation (DSE) for space-like momenta.
- In solving the BSE in Euclidean space, the propagators are functions of  $(k+P)^2$ ,  $P = (0,0,0,iM)$ .
- Extension to complex plane via Cauchy's integral theorem.

# Bethe-Salpeter Equation for QCD Bound States

$$\Gamma_M^{fg}(k, P) = \int^{\Lambda} \frac{d^4 q}{(2\pi)^4} K_{fg}(k, q; P) S_f(q_\eta) \Gamma_M^{fg}(q, P) S_g(q_{\bar{\eta}})$$

$K_{fg}(q, k; P)$  = Quark-antiquark scattering kernel

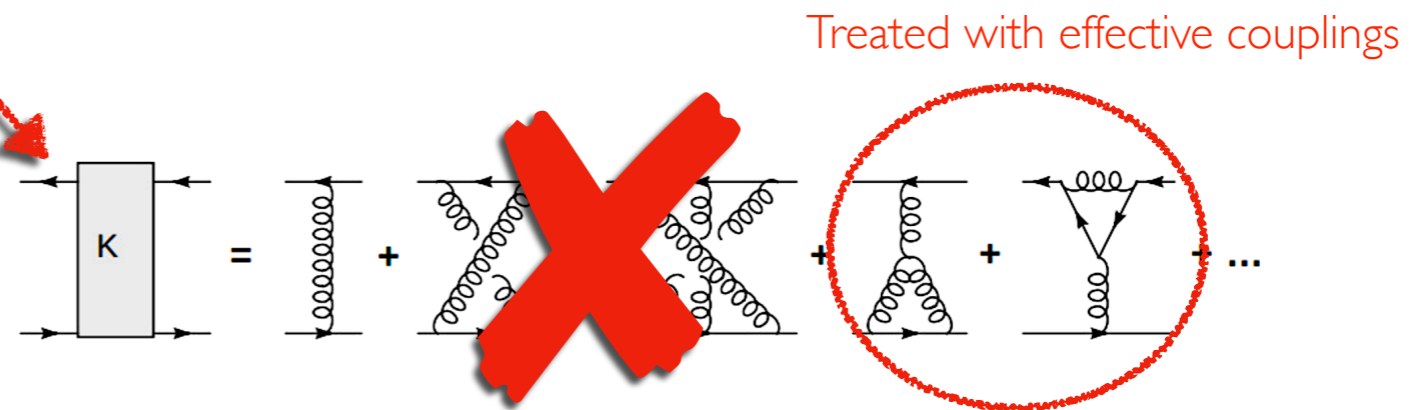
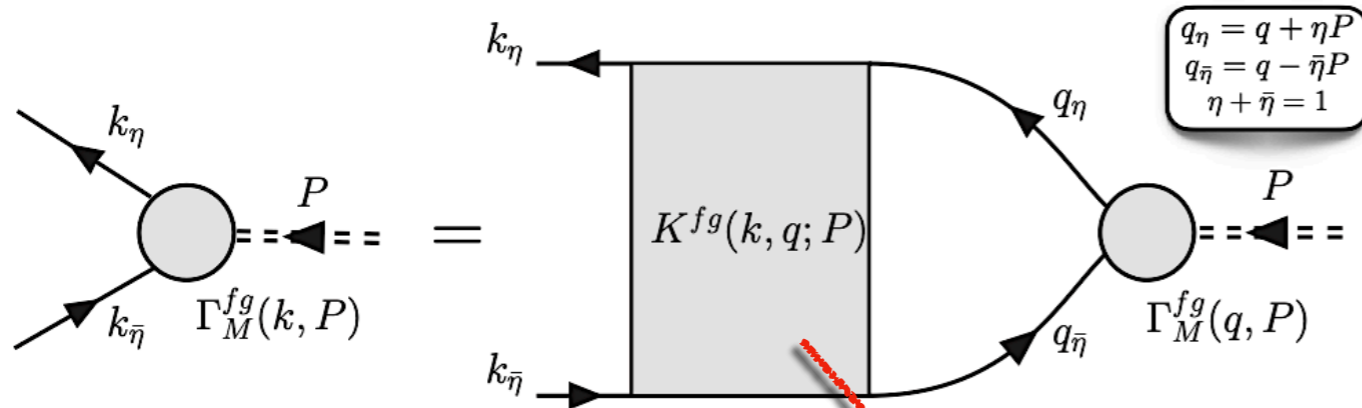
Poincaré invariant BSA solution for pseudoscalar mesons:

$S_f(q_\eta)$  = Dressed quark propagator

$$\Gamma_M^{fg}(k, P) = \gamma_5 \left[ i E_M^{fg}(k, P) + \gamma \cdot P F_M^{fg}(k, P) \right.$$

$\Gamma_M^{fg}(k, P)$  = Meson's Bethe-Salpeter Amplitude (BSA)

$$\left. + \gamma \cdot k k \cdot P G_M^{fg}(k, P) + \sigma_{\mu\nu} k_\mu P_\nu H_M^{fg}(k, P) \right]$$



# Pseudoscalar Meson Spectrum

Mesons/Observables	$m_M$	$m_M^{\text{exp.}}$	$\epsilon_r^m$ [%]	$f_M$	$f_M^{\text{exp./lQCD}}$	$\epsilon_r^f$ [%]
$\pi(u\bar{d})$	0.136	0.140	2.90	$0.094^{+0.001}_{-0.001}$	0.092(1)	2.17
$K(s\bar{u})$	0.494	0.494	0.0	$0.110^{+0.001}_{-0.001}$	0.110(2)	0.0
$D_u(c\bar{u})$	$1.867^{+0.008}_{-0.004}$	1.870	0.11	$0.144^{+0.001}_{-0.001}$	0.150(0.5)	4.00
$D_s(c\bar{s})$	$2.015^{+0.021}_{-0.018}$	1.968	2.39	$0.179^{+0.004}_{-0.003}$	0.177(0.4)	1.13
$\eta_c(c\bar{c})$	$3.012^{+0.003}_{-0.039}$	2.984	0.94	$0.270^{+0.002}_{-0.005}$	0.279(17)	3.23
$\eta_b(b\bar{b})$	$9.392^{+0.005}_{-0.004}$	9.398	0.06	$0.491^{+0.009}_{-0.009}$	0.472(4)	4.03

Mesons/Observables	$m_M$	$m_M^{\text{exp.}}$	$\epsilon_r^m$ [%]	$f_M$	$f_M^{\text{lQCD}}$	$\epsilon_r^f$ [%]
$B_u(b\bar{u})$	$5.277^{+0.008}_{-0.005}$	5.279	0.04	$0.132^{+0.004}_{-0.002}$	0.134(1)	4.35
$B_s(b\bar{s})$	$5.383^{+0.037}_{-0.039}$	5.367	0.30	$0.128^{+0.002}_{-0.003}$	0.162(1)	20.50
$B_c(b\bar{c})$	$6.282^{+0.020}_{-0.024}$	6.274	0.13	$0.280^{+0.005}_{-0.002}$	0.302(2)	7.28
$\eta_b(b\bar{b})$	$9.383^{+0.005}_{-0.004}$	9.398	0.16	$0.520^{+0.009}_{-0.009}$	0.472(4)	10.17

# Vector Meson Spectrum

Mesons/Properties	$m_M$	$m_M^{\text{exp}}$	$\epsilon_r^m \%$	$f_M$	$f_M^{\text{exp}/\text{IQCD}}$	$\epsilon_r^f \%$
$\rho(u\bar{d})$	0.730	0.775	5.810	0.145	0.153	5.229
$\phi(s\bar{s})$	1.070	1.019	5.197	0.187	0.168	11.309
$K^*(s\bar{u})$	0.942	0.896	5.134	0.177	0.159	11.321
$J/\Psi(c\bar{c})$	3.124	3.097	0.872	0.277	0.294	5.782
$\Upsilon(b\bar{b})$	9.411	9.460	0.518	0.594	0.505	17.624

Mesons/Properties	$m_M$	$m_M^{\text{exp}}$	$\epsilon_r^m \%$	$f_M$	$f_M^{\text{exp}/\text{IQCD}}$	$\epsilon_r^f \%$
$D_u^*(c\bar{u})$	2.021	2.009	0.597	0.165	0.158	4.430
$D_s^*(c\bar{s})$	2.169	2.112	2.699	0.205	0.190	7.895

F. Serna, R. Correa da Silveira, **B.E.**, in preparation



# Light-Cone Distribution Amplitudes

- We do not obtain the LCDA directly from the light-front wave functions:

$$f_M \phi_M(x, \mu) = \int \frac{d^2 k_\perp}{16\pi^2} \psi_M(x, k_\perp)$$

- Instead, we compute the Mellin moments:

$$\langle x^m \rangle = \int_0^1 dx x^m \phi_M(x, \mu) \quad \langle x^0 \rangle = \int_0^1 dx \phi_M(x, \mu) = 1$$

$$f_M \phi_M(x, \mu) = Z_2 \text{tr}_{\text{CD}} \int \frac{d^4 k}{(2\pi)^4} \delta(n \cdot k - x n \cdot P) \gamma_5 \gamma \cdot n S(k) \Gamma_\pi(k; P) S(k - P)$$

$$f_M (n \cdot P)^{m+1} \int_0^1 dx x^m \phi_M(x, \mu) = Z_2 \text{tr}_{\text{CD}} \int \frac{d^4 k}{(2\pi)^4} (n \cdot k)^m \gamma_5 \gamma \cdot n \chi_\pi(k; P)$$

**Bethe-Salpeter wave function**

L. Chang, I.C. Cloët, J.J. Cobos-Martínez, C.D. Roberts, S.M. Schmidt, P. C. Tandy, Phys. Rev. Lett. 110 (2013)

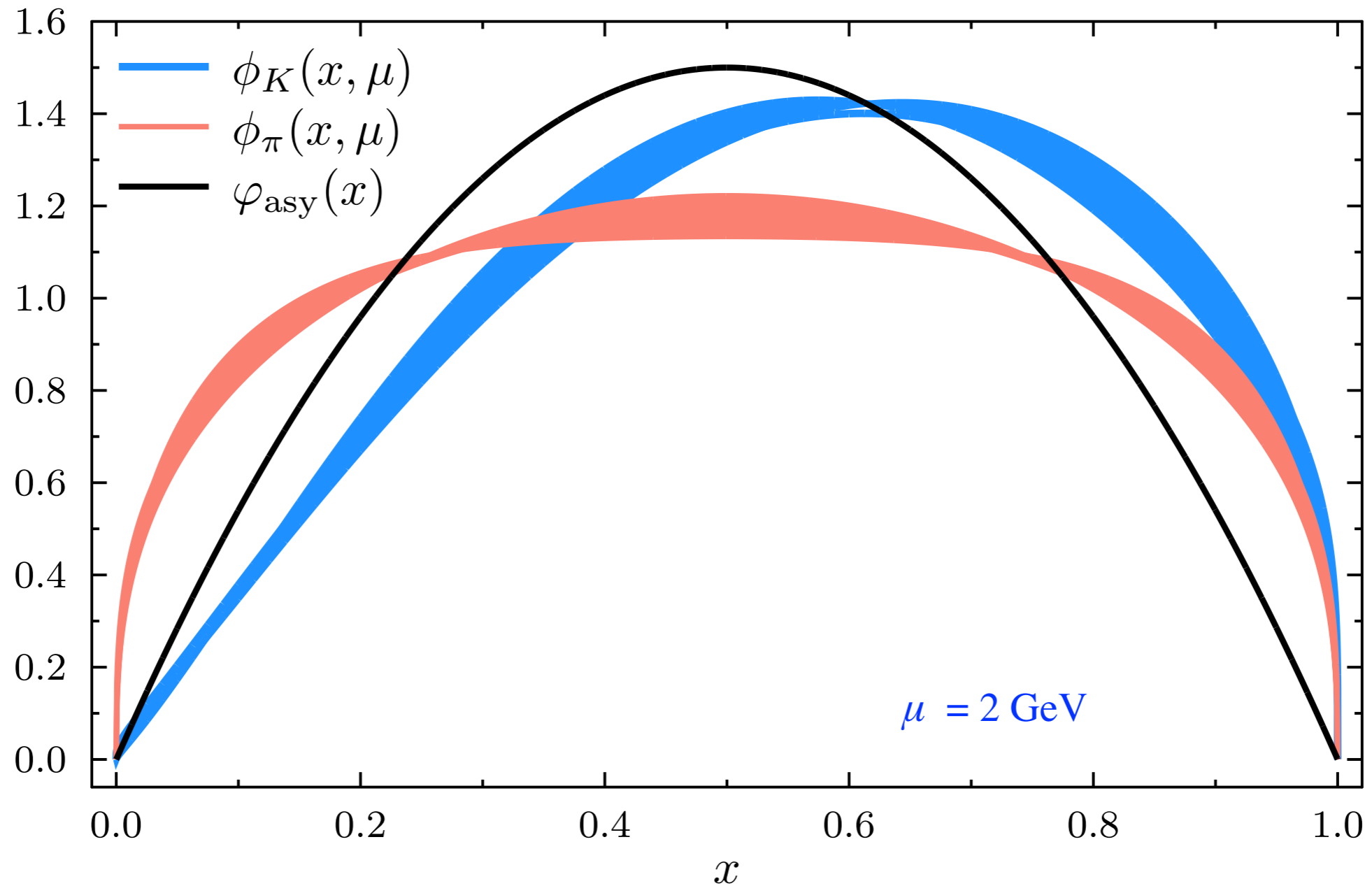
D. Binosi, L. Chang, M. Ding, F. Gao, I. Papavassiliou and C. D. Roberts, Phys. Lett. B 790 (2019)

- We reconstruct the distributions amplitudes by means of a Gegenbauer-like expansion or other functional forms:

$$\phi_\pi^{\text{rec.}}(x, \mu) = \mathcal{N}(\alpha) [x\bar{x}]^{\alpha-1/2} [1 + a_2 C_2^\alpha(2x - 1)]$$

$$\phi_H^{\text{rec.}}(x, \mu) = \mathcal{N}(\alpha, \beta) 4x\bar{x} e^{4\alpha x\bar{x} + \beta(x - \bar{x})}$$

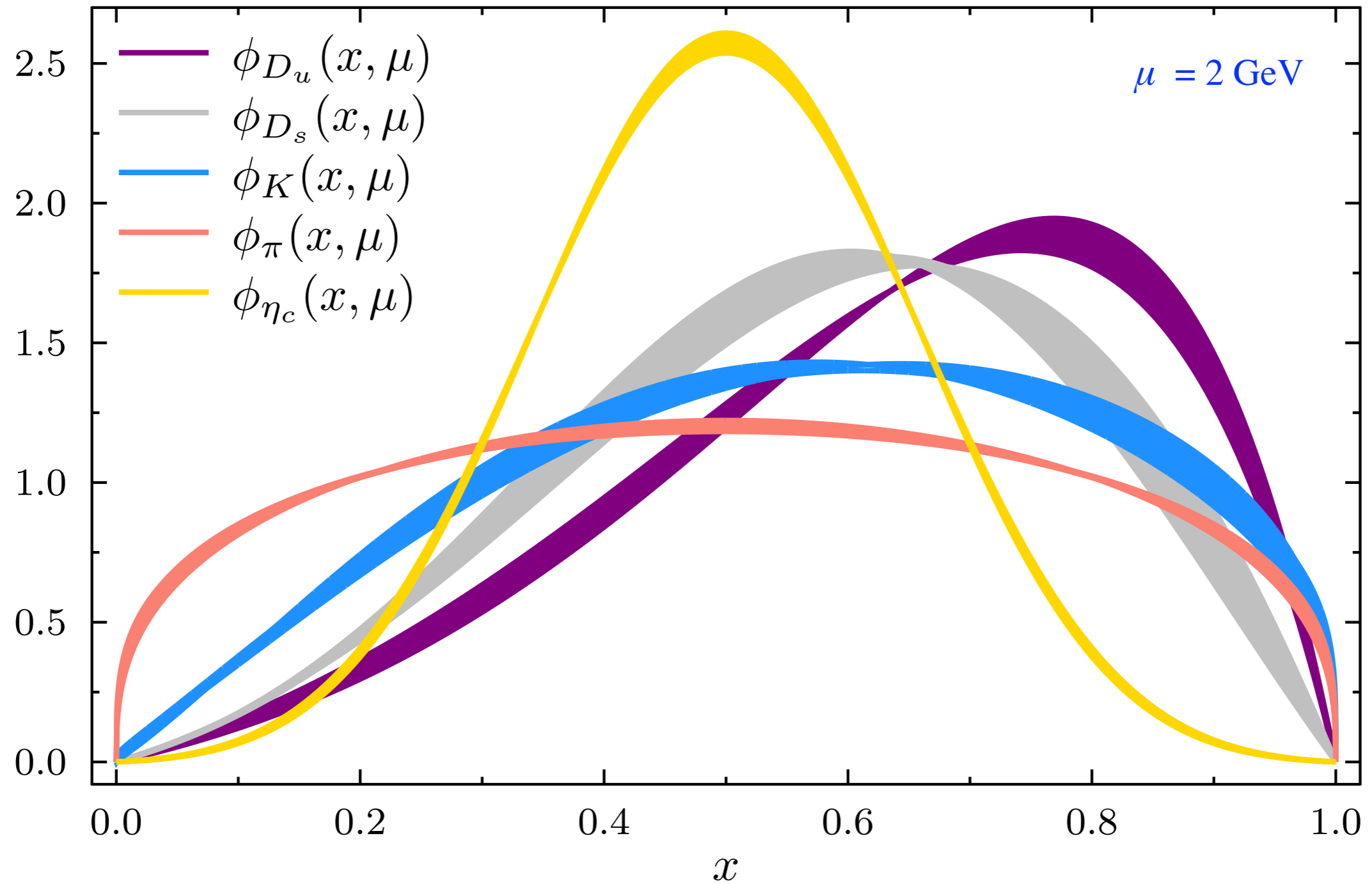
# Light-Cone Distribution Amplitudes



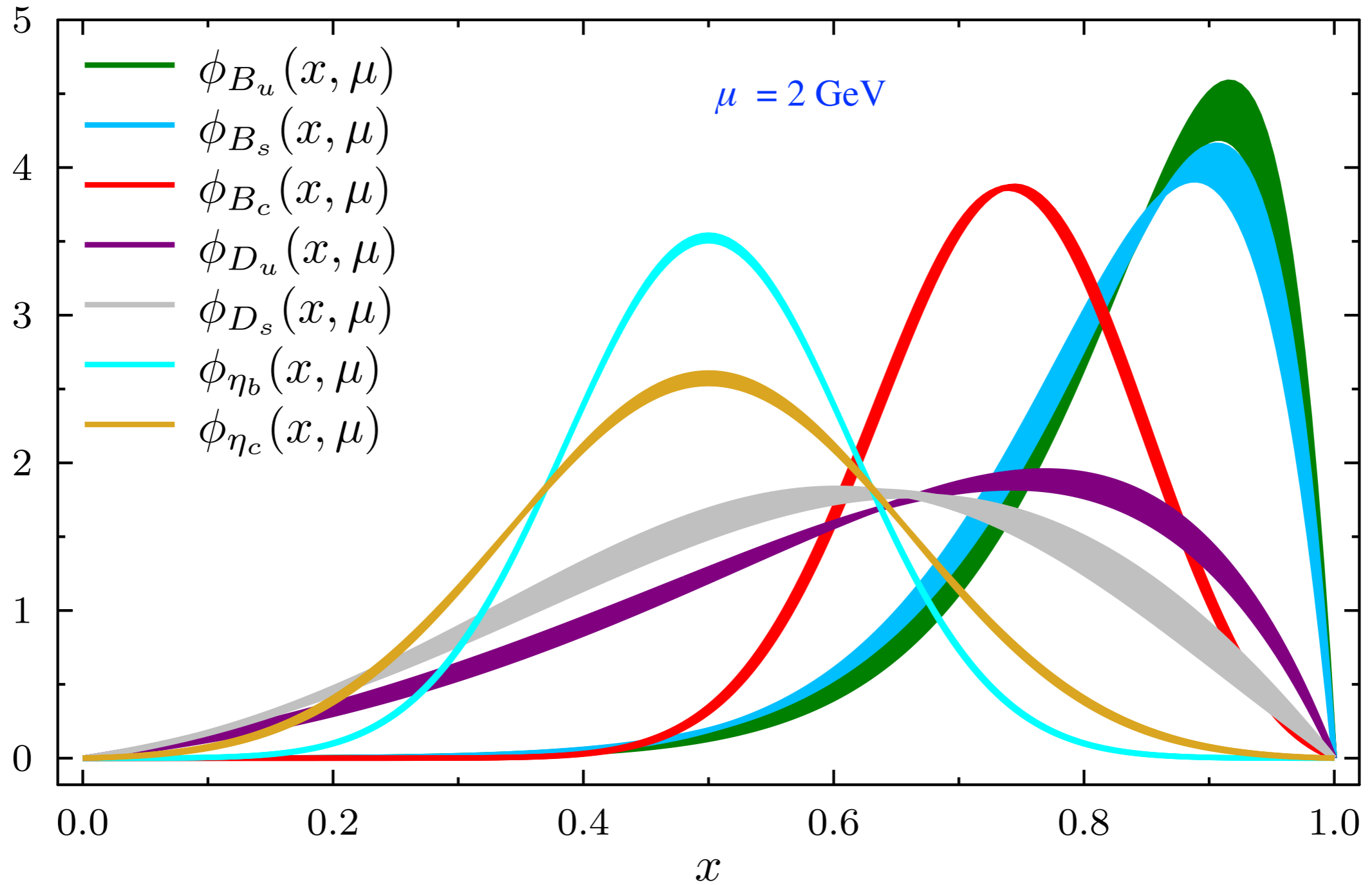
L. Chang, I.C. Cloët, J.J. Cobos Martínez, C.D. Roberts, S.M. Schmidt, P. C. Tandy, Phys. Rev. Lett. 110 (2013)

F. Serna, R. Correa da Silveira, J.J. Cobos Martínez, B.E., E. Rojas, Eur. Phys. J. C 80 (2020)

# Light-Cone Distribution Amplitudes



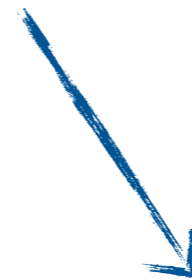
# Light-Cone Distribution Amplitudes



Only valid for full QCD !

Application to HQET requires the use of heavy-quark expansion of charm and bottom propagator in the Bethe-Salpeter equation.

$$\langle 0 | \bar{u}(zn) \mathcal{W}[z, 0] \gamma \cdot n \gamma_5 h_v(0) | \bar{B}(v) \rangle$$



$$S_Q(p) = \frac{\gamma \cdot p + m_Q}{p^2 - m_Q^2} \xrightarrow{m_Q \rightarrow \infty} \frac{1 + \gamma \cdot v}{2v \cdot k} + \mathcal{O}\left(\frac{k}{m_Q}\right),$$

$$p_\mu = m_Q v_\mu + k_\mu$$

$$k \sim \Lambda_{\text{QCD}}$$

# Conclusions & Progress

- Much progress was made from QCD based modeling toward nonperturbative numerical solutions of quark propagators and quark-antiquark bound states for flavored mesons satisfying chiral symmetry and Poincaré covariance.
- Good reproduction of charmonium and bottomonium as well as D and B meson mass spectrum and their weak decay constants.
- Improvements in Bethe-Salpeter kernels beyond ladder truncation underway ...  
⇒ needed for scalar and axialvector channels and their higher radially excited states, as well as better control of quark correlation functions on complex plane.
- LCDA of vector quarkonia and  $D$  and  $D_s$  mesons underway.
- Calculation of scalar  $D$  and  $D_s$  mesons & excitations (PANDA-FAIR, Jlab, EIC ?)