

# IP8 Secondary Focus xL Coverage

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# Preliminaries

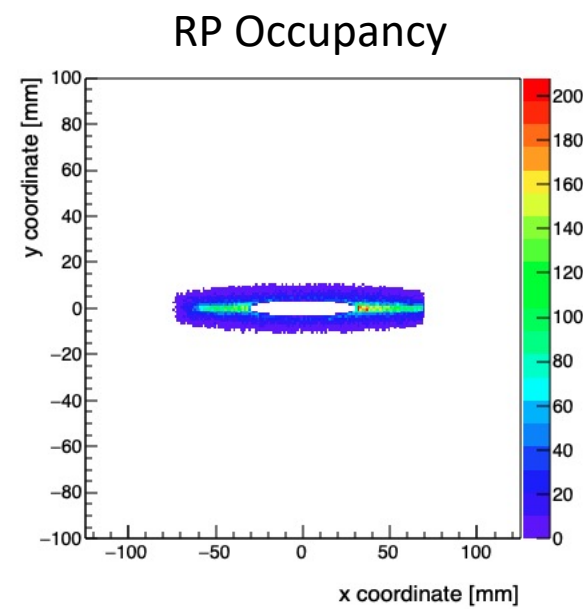
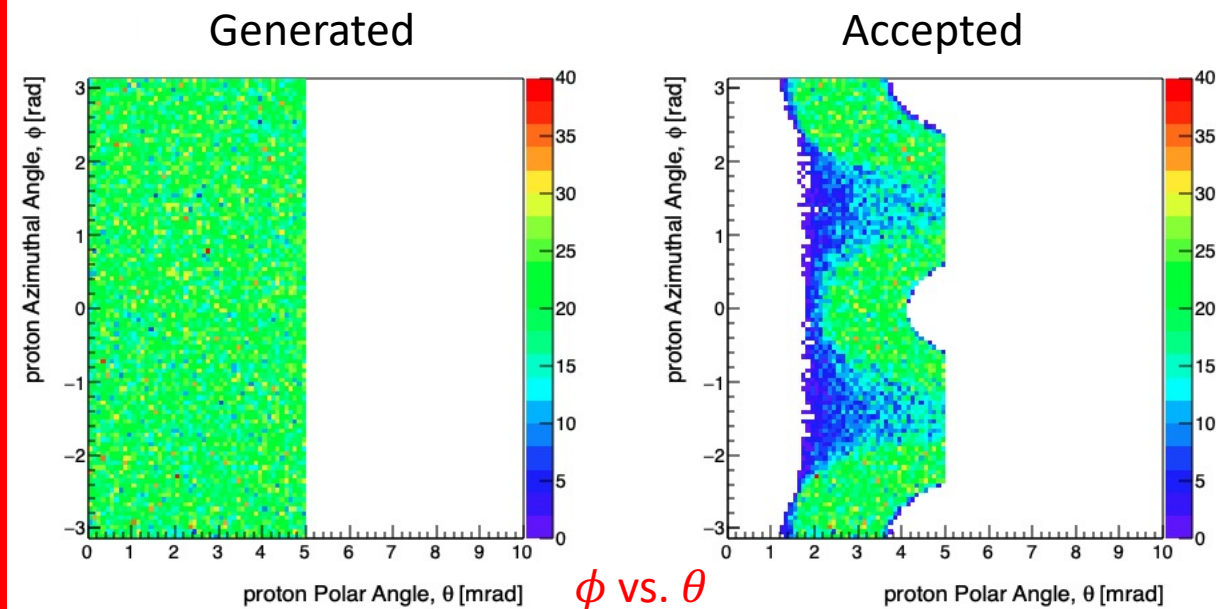
- Using the same IP8 lattice shown at the FF cross-collaboration meeting (<https://indico.bnl.gov/event/12068/>).
- There are some notable changes when using the correct IP8 optics numbers (compared to the IP6 ones).
  - Larger dispersion @ RPSF: Causes the  $10\sigma$  to be larger.
  - Now using the correct beta-functions at the RPSF – also causes the  $10\sigma$  to be larger (the beta-functions are a bit larger at the SF than at the IP, using IP6 numbers for comparison).
  - Beta functions at the drift RP are very different in x and y, with that difference getting smaller further downstream -> causes anisotropic acceptance in phi. Could re-think placement (in z) of drift RP (I have some results on this too).

# 10 $\sigma$ Calculation

- We have normally calculated the transverse beam size as  $\sigma_{x,y} = \sqrt{\beta(z)_{x,y}\epsilon_{x,y}}$ , but this is only approximately true if the dispersion is very small, or if the beta-function times emittance is very large in comparison.
  - My code has the correct formula,  $\sigma_{x,y} = \sqrt{\beta(z)_{x,y}\epsilon_{x,y} + \left(D_{x,y} \frac{\Delta p}{p}\right)^2}$ .
  - The second term is very small compared to the first in the case of the normal RP, so  $\sigma_{x,y} = \sqrt{\beta(z)_{x,y}\epsilon_{x,y}}$  is perfectly fine to quote in that case.
  - However, at the RPSF, this second term is large enough to affect the total transverse beam size.
- Using the IP6 dispersion value for the standard RP and RPSF, which is smaller than the IP8 number caused the 10 $\sigma$  to be smaller in my previous study.
  - This is now fixed.

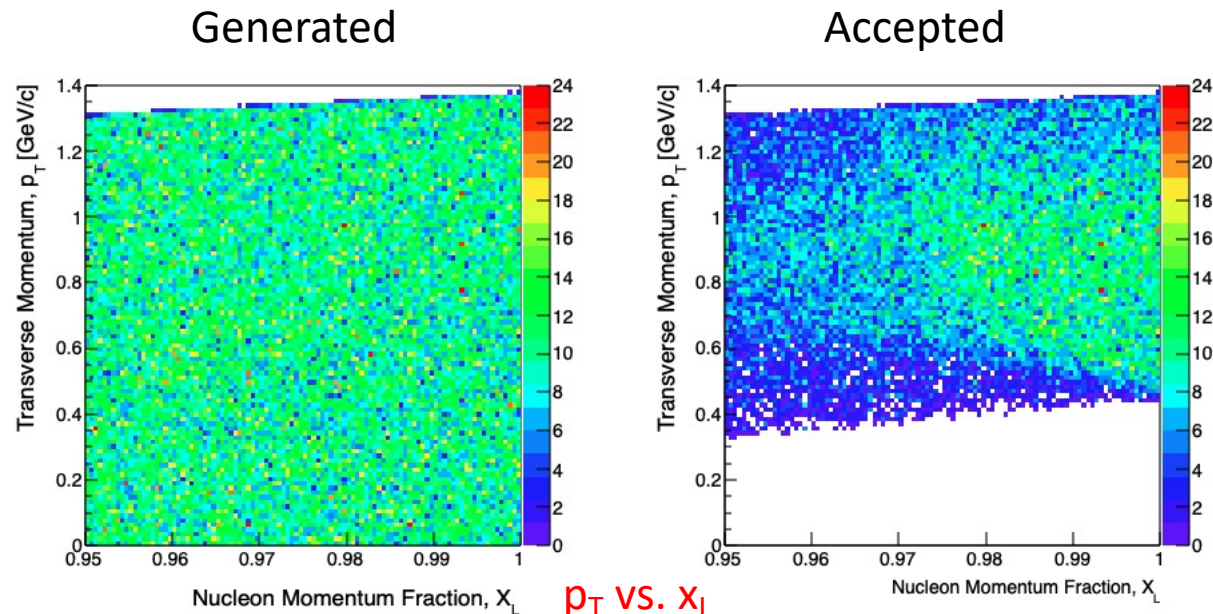
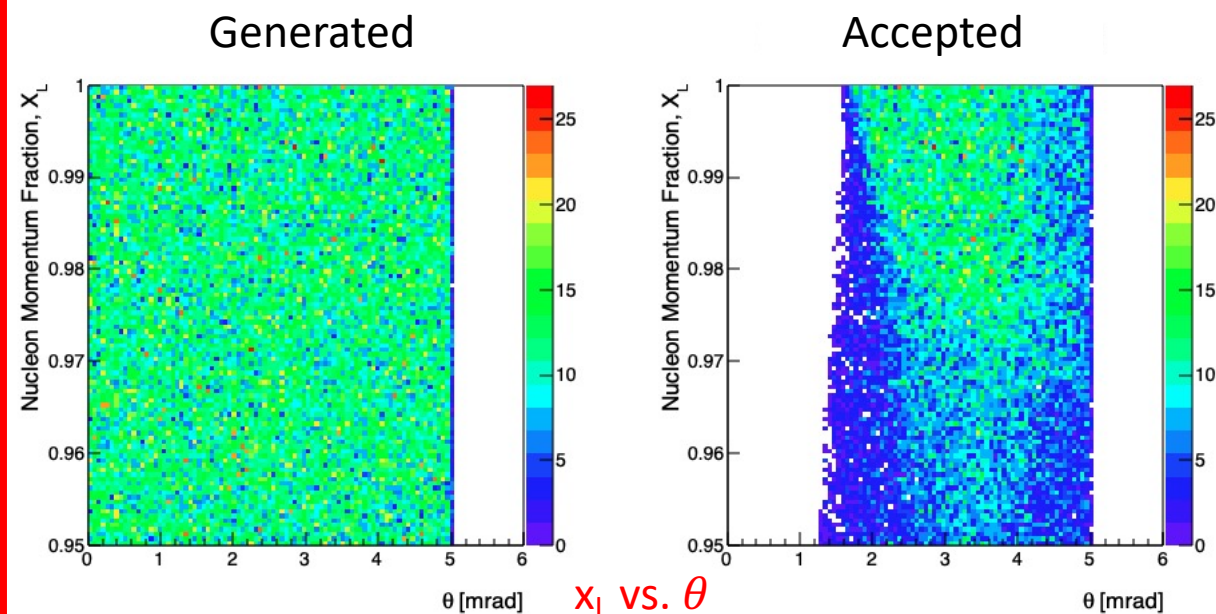
# Drift Roman Pots ( $z = 26\text{m}$ )

protons (275 GeV settings)  
 $0 < \theta < 5 \text{ mrad}$   
 $0.95 < x_L < 1.0$



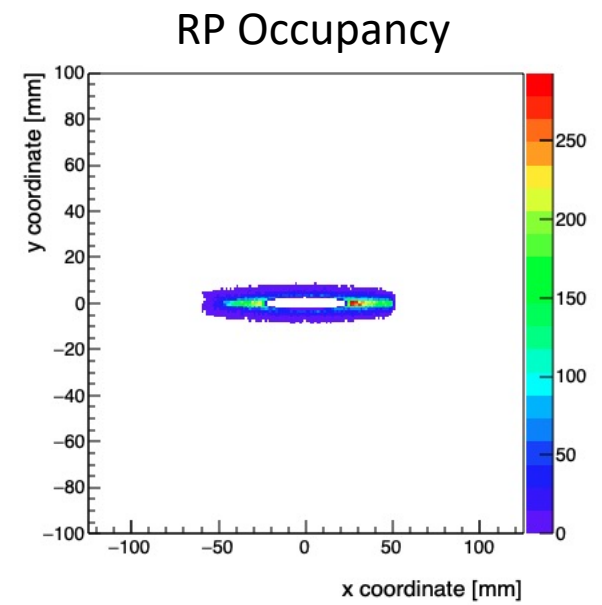
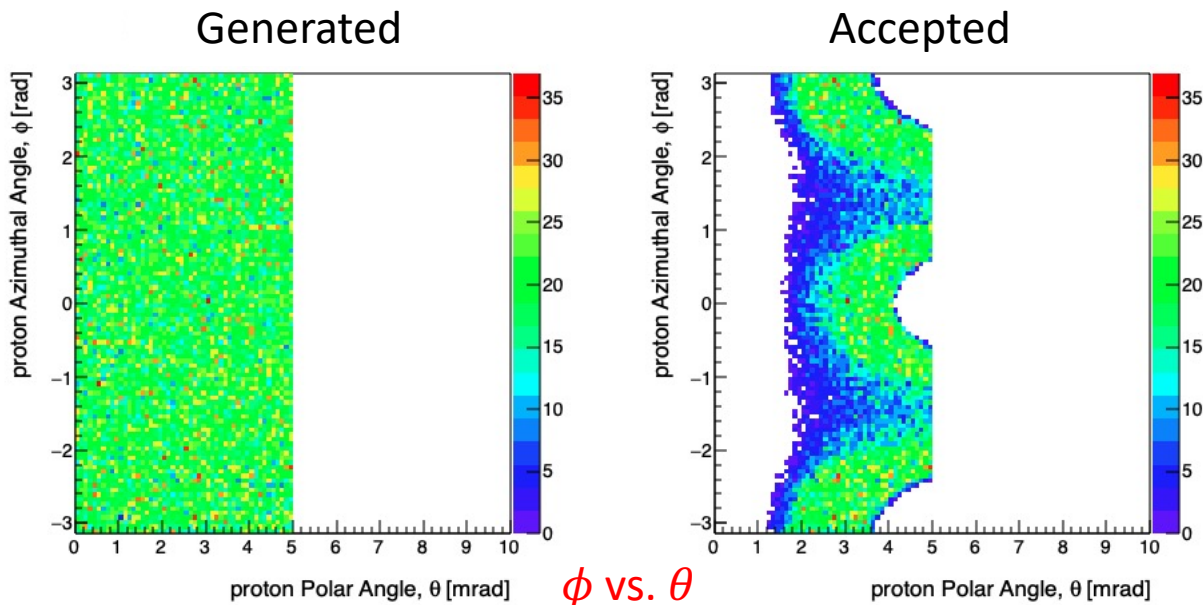
$$\beta_x (z = 25.7\text{m}) = 548 \text{ m}$$
$$\beta_y (z = 25.7\text{m}) = 87 \text{ m}$$
$$D_x = 0.126$$

$$\sigma_x @ \text{RP} = 3.10607 \text{ mm}$$
$$\sigma_y @ \text{RP} = 0.372664 \text{ mm}$$



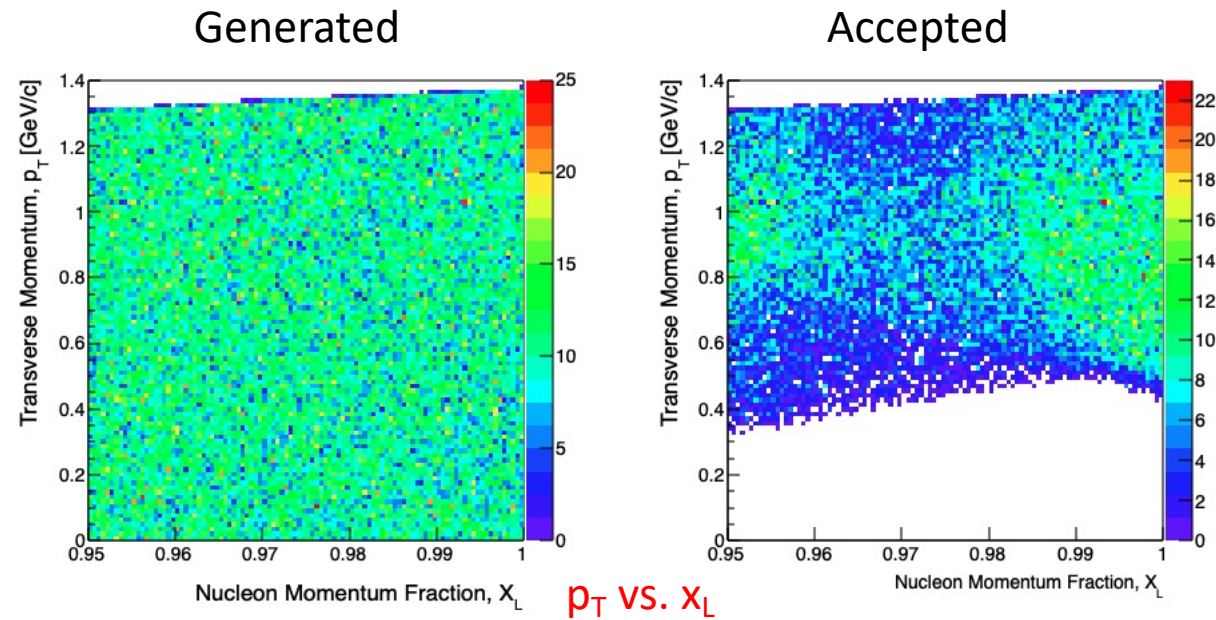
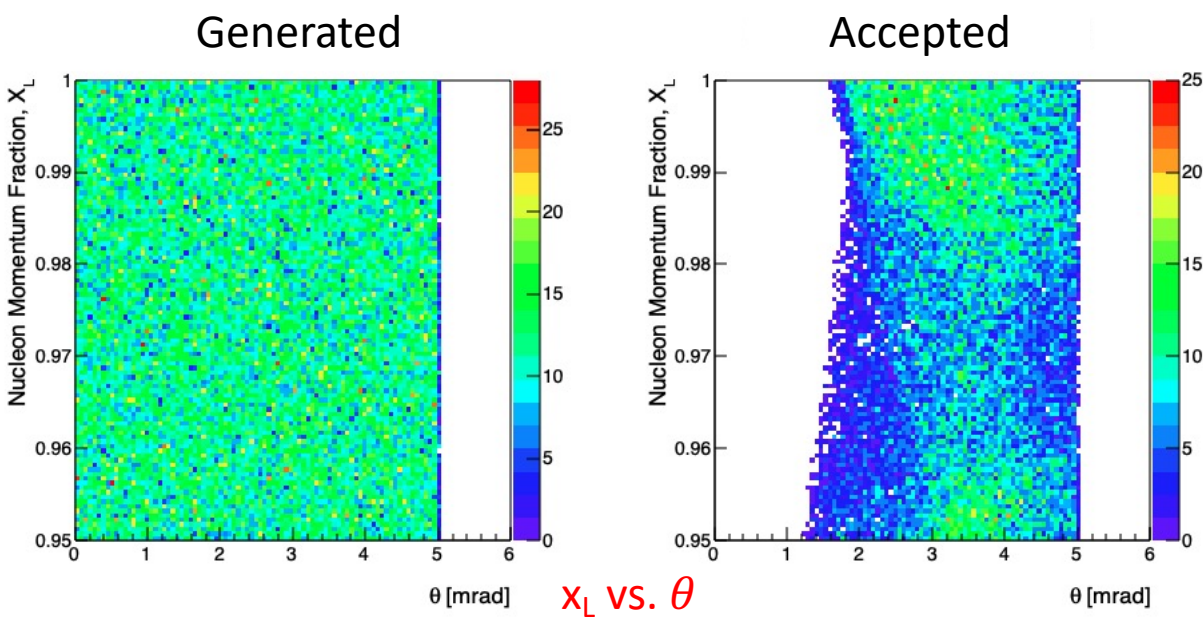
# Drift Roman Pots ( $z = 30\text{m}$ )

protons (275 GeV settings)  
 $0 < \theta < 5 \text{ mrad}$   
 $0.95 < x_L < 1.0$



$$\beta_x (z = 29.7\text{m}) = 338 \text{ m}$$
$$\beta_y (z = 29.7\text{m}) = 55 \text{ m}$$
$$D_x = 0.204$$

$$\sigma_x @ \text{RP} = 2.44373 \text{ mm}$$
$$\sigma_y @ \text{RP} = 0.297106 \text{ mm}$$



# Roman Pots @ SF

$\sigma_x$  @ RPSF = 0.328283 mm.

$\sigma_y$  @ RPSF = 0.085217 mm.

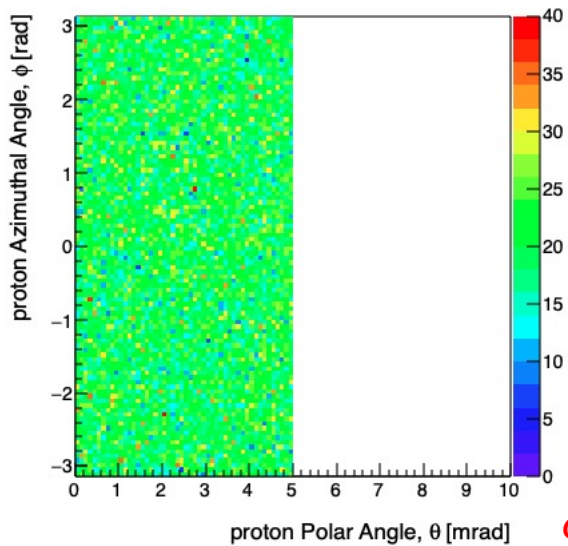
protons (275 GeV settings)

$0 < \theta < 5$  mrad

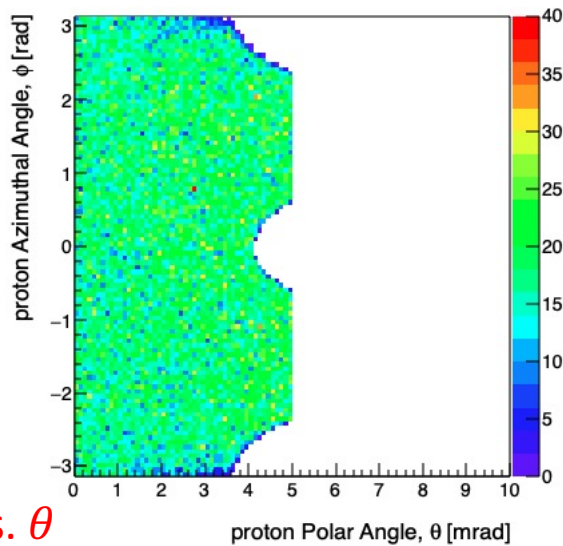
$0.95 < x_L < 1.0$

RPSF Occupancy

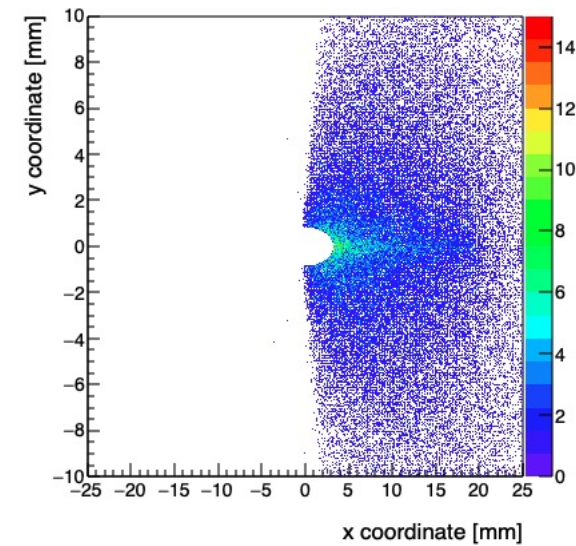
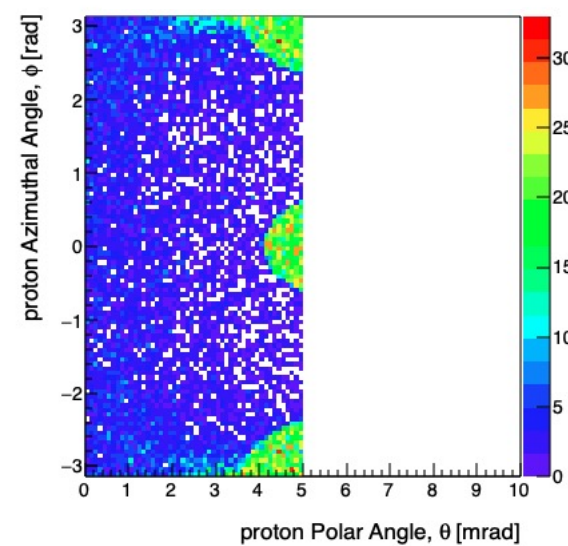
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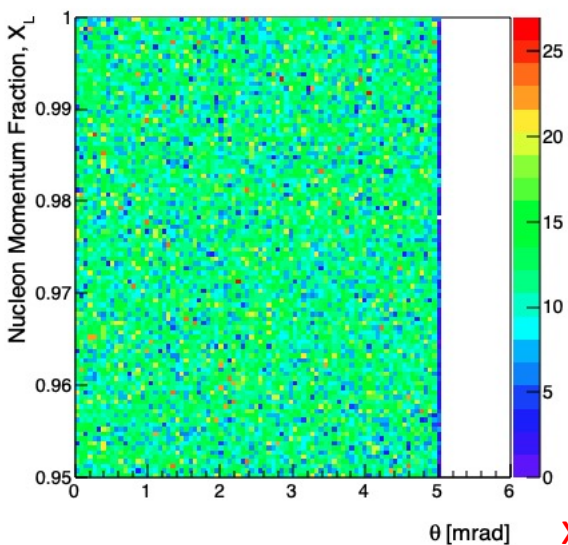


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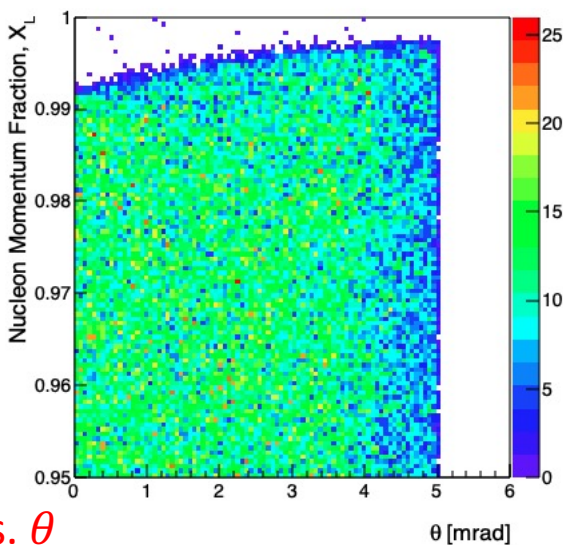


$\phi$  vs.  $\theta$

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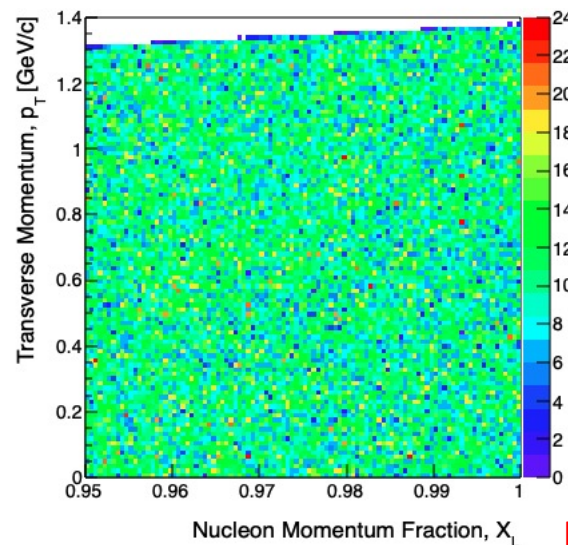


Accepted

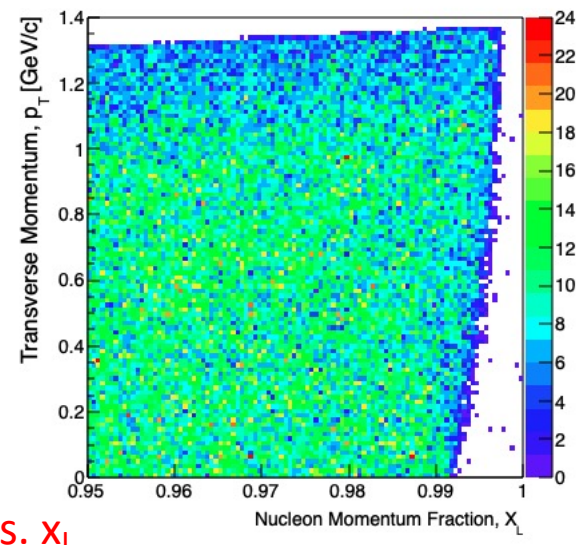


$x_L$  vs.  $\theta$

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Accepted



$p_T$  vs.  $x_L$

# Summary

- The secondary focus provides a significant improvement in coverage of  $x_L$  for protons.
- Finding a “sweet spot” for the RP drift detectors will require optimization.
- It would still be wise to have both detectors working together to maximize total phase space coverage in  $x_L$  and  $p_T$ .