# Theoretical review on systematic uncertainties for CP violation searches

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#### **‡**Fermilab

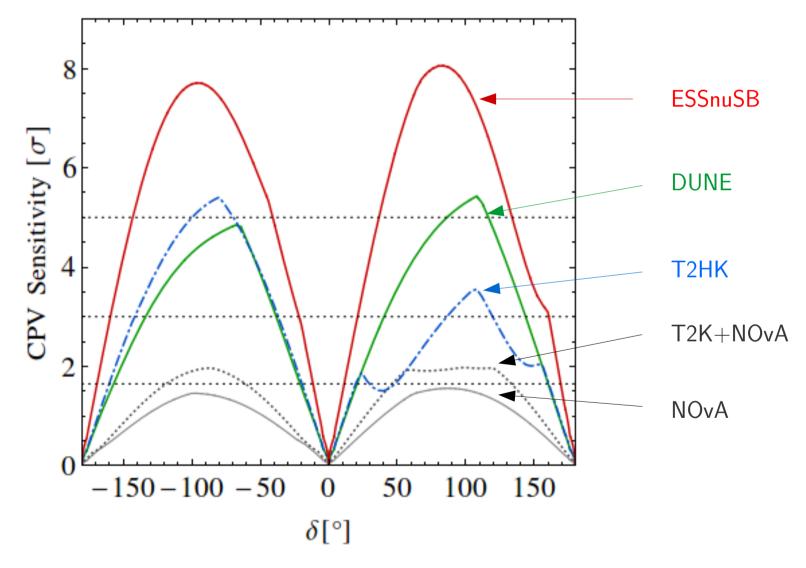


NNN'15 Stony Brook University, Oct 29<sup>th</sup>, 2015

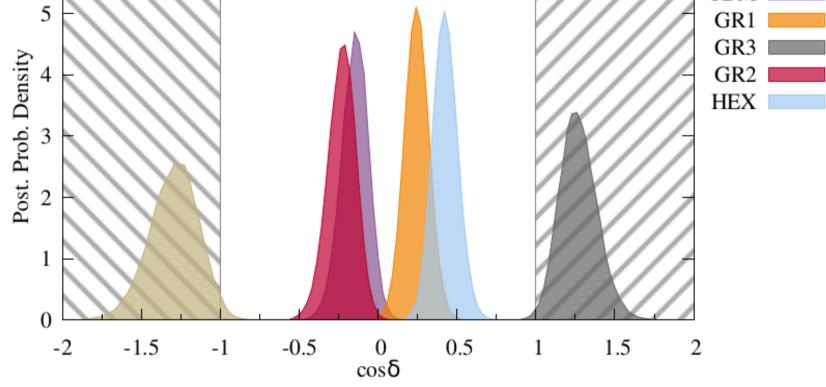
# Outline

- 1) Introduction
- 2) Normalization uncertainties
- 3) Shape uncertainties
- 4) Summary

#### Proposed experiments for CPV searches



# Why precision?



Ballett, King, Luhn, Pascoli, Schmidt, 1410.7573 [hep-ph] (see also, e.g., Girardi et al, 1410.8056, Meloni, 1308.4578)

(See talks in the UD2 afternoon session today)

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The golden channel in neutrino oscillations is:

$$\begin{split} p_{e\mu}^{\pm} &\equiv P(\overleftarrow{\nu_e}) \rightarrow \overleftarrow{\nu_{\mu}}) = \overbrace{s_{23}^2 \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta_{31} L}{2}\right)}_{2} + \overbrace{c_{23}^2 \sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta_{21} L}{2}\right)}_{2} \\ &+ \widetilde{J} \cos \left(\pm \delta - \frac{\Delta_{31} L}{2}\right) \sin \left(\frac{\Delta_{21} L}{2}\right) \sin \left(\frac{\Delta_{31} L}{2}\right), \end{split}$$

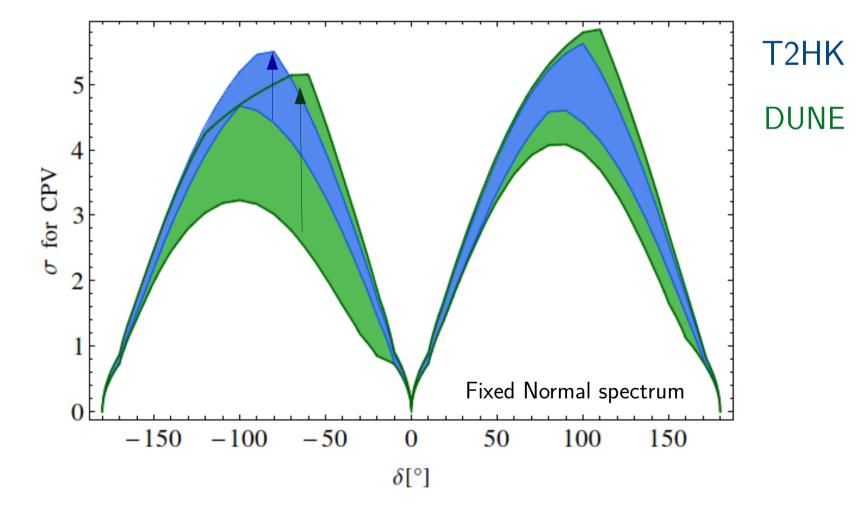
$$\begin{aligned} &\text{CP-violating} \\ &\text{interference} \end{aligned}$$

$$\Delta_{ij} = \frac{\Delta m_{ij}^2 L}{2E}$$

Cervera et al., hep-ph/0002108

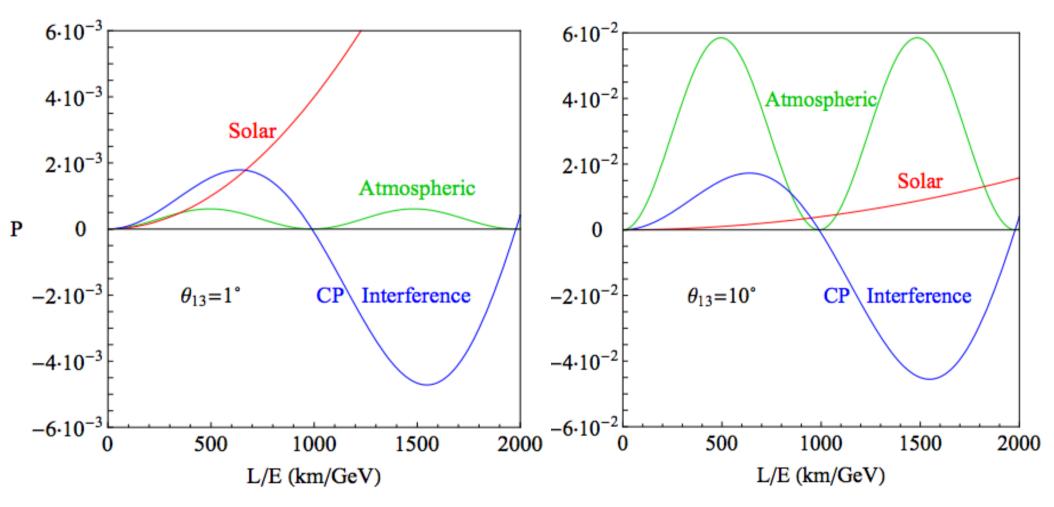
$$\begin{split} P_{e\mu}^{\pm} &\equiv P(\overleftarrow{\nu_e} \to \overleftarrow{\nu_{\mu}}) = s_{23}^2 \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta_{31} L}{2}\right) + c_{23}^2 \sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta_{21} L}{2}\right) \\ &+ \tilde{J} \cos \left(\pm \delta + \frac{\Delta_{31} L}{2}\right) \sin \left(\frac{\Delta_{21} L}{2}\right) \sin \left(\frac{\Delta_{31} L}{2}\right), \end{split}$$
$$\begin{aligned} &\quad For neutrinos, \\ there is non-trivial \\ energy dependence \\ too \end{split}$$
Wide band \\ beams exploit \\ this \end{aligned}

An example of relative importance of energy resolution:



#### Normalization uncertainties

#### Impact of systematics on CPV



Coloma and Fernandez-Martinez, 1110.4583 [hep-ph]

See also Marciano, hep-ph/0108181, Hagiwara et al, hep-ph/0607255 and Meregaglia, Rubbia, 0801.4035

P. Coloma - Systematics

#### Near/Far cancellation?

 $n_{\alpha \to \beta}(L, E) \sim \frac{1}{L^2} \epsilon_{\beta}(E) \times \sigma_{\beta}(E) \times \phi_{\alpha}(E) \times P_{\alpha\beta}(L, E)$ 

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At reactor experiments, the cancellation of systematics between near/far detectors is very effective:

$$\frac{n_{ee}^{FD}}{n_{ee}^{ND}} \sim \frac{L_{ND}^2}{L_{FD}^2} \frac{\epsilon_e \sigma_e \phi_e}{\epsilon_e \sigma_e \phi_e} P_{ee}$$

$$n_{\alpha \to \beta}(L, E) \sim \frac{1}{L^2} \epsilon_{\beta}(E) \times \sigma_{\beta}(E) \times \phi_{\alpha}(E) \times P_{\alpha\beta}(L, E)$$

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At Daya Bay this works extremely well!

$$\sin^2 2\theta_{13} = 0.084^{+0.005}_{-0.005}$$

(Seminar given by Xin Qian at Fermilab, on Oct 15)

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# Near/Far cancellation?

• For CP violation searches, we need an appearance experiment

• An ideal near detector can be used to predict some backgrounds:

$$n_{\nu_e}^{FD,bg} \sim n_{\nu_e}^{ND,bg} \frac{L_{ND}^2}{L_{FD}^2} \frac{V_{FD}}{V_{ND}}$$

Huber, Mezzetto and Schwetz, 0711.2950 [hep-ph] Coloma, Huber, Kopp, Winter, 1209.5973 [hep-ph]

# Near/Far cancellation?

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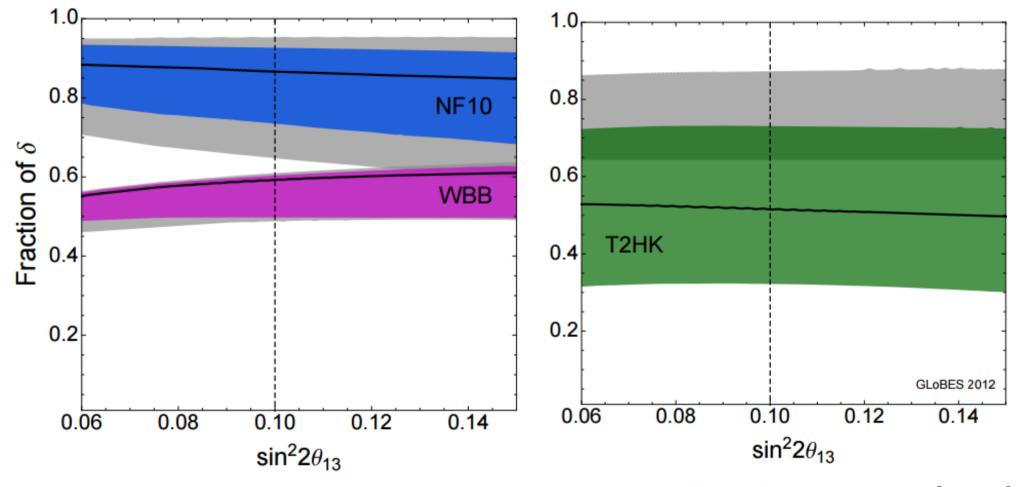
$$n_{\nu_e}^{FD,bg} \sim n_{\nu_e}^{ND,bg} \frac{L_{ND}^2}{L_{FD}^2} \frac{V_{FD}}{V_{ND}}$$

• However, a similar extrapolation for the signal will not work that well:

$$n_{\nu_e}^{FD,sig} \sim n_{\nu_{\mu}}^{ND,sig} \frac{L_{ND}^2}{L_{FD}^2} \frac{V_{FD}}{V_{ND}} \underbrace{\tilde{\sigma}_{\nu_e}}{\tilde{\sigma}_{\nu_{\mu}}} \times P(\nu_{\mu} \to \nu_e)$$

Huber, Mezzetto and Schwetz, 0711.2950 [hep-ph] Coloma, Huber, Kopp, Winter, 1209.5973 [hep-ph]

#### Impact of normalization uncertainties



Coloma, Huber, Kopp, Winter, 1209.5973 [hep-ph] (See also Huber, Mezzetto and Schwetz, 0711.2950 [hep-ph])

#### Possible ways out

• Correlations can help to reduce impact of systematics:

- the far detector can act as a "near detector"

• Possible ways to reduce the effect of <u>normalization uncertainties</u>:

 measure final flavor cross sections at a near detector (intrinsic contamination).

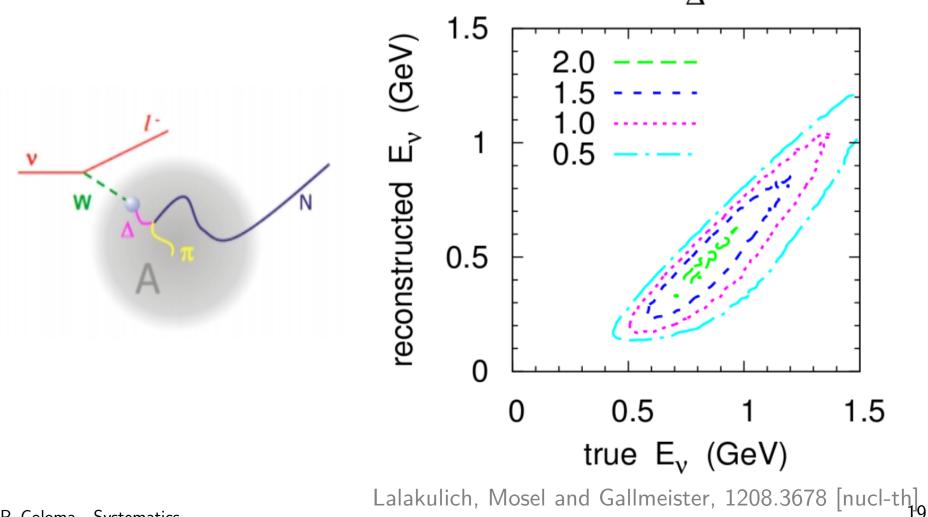
 put theoretical constraints on ratios between cross sections for different flavors

 Caveats: near/far flux extrapolation is tricky; near/far detectors may not be identical, etc

#### Shape uncertainties

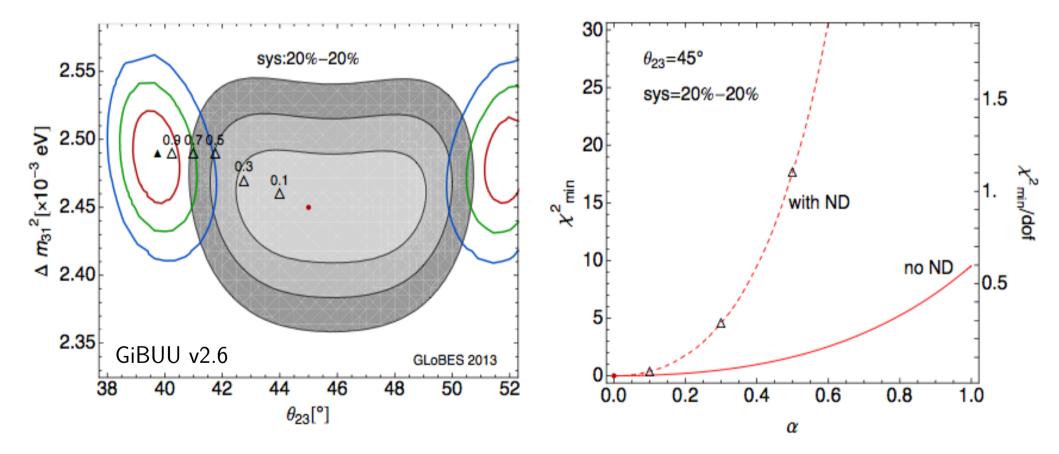
# Energy reconstruction effects

These effects can be parametrized as migration matrices from true to reconstructed energy:  $\Lambda$ 



#### Toy model

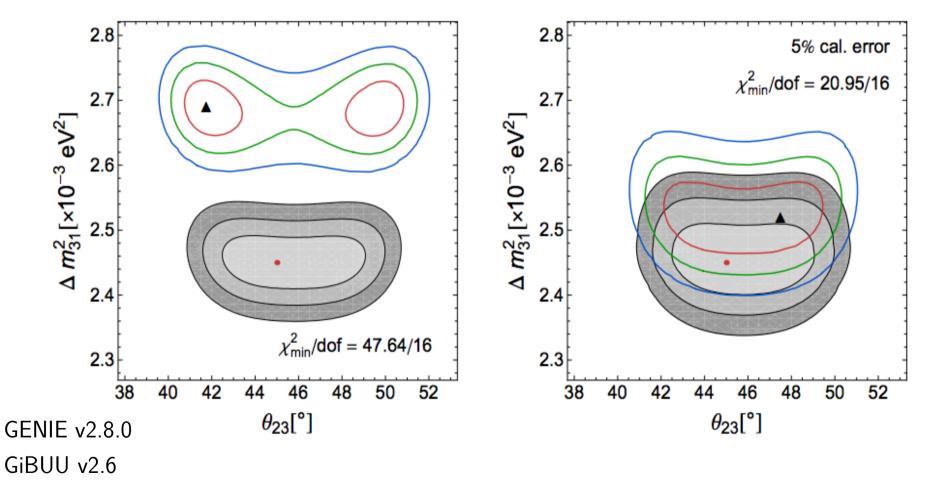
$$N_i^{\text{test}}(\alpha) = \alpha \times N_i^{QE} + (1 - \alpha) \times N_i^{QE-like}$$



Coloma and Huber, 1307.1243 [hep-ph]

#### Impact of nuclear model

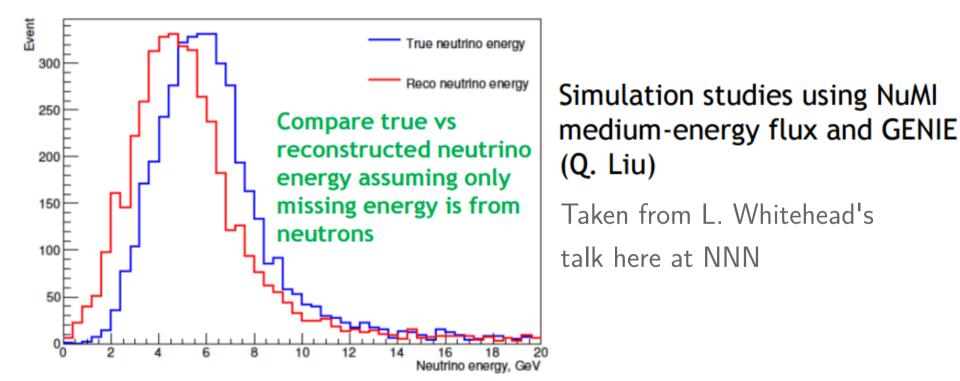
How large can these effects be?



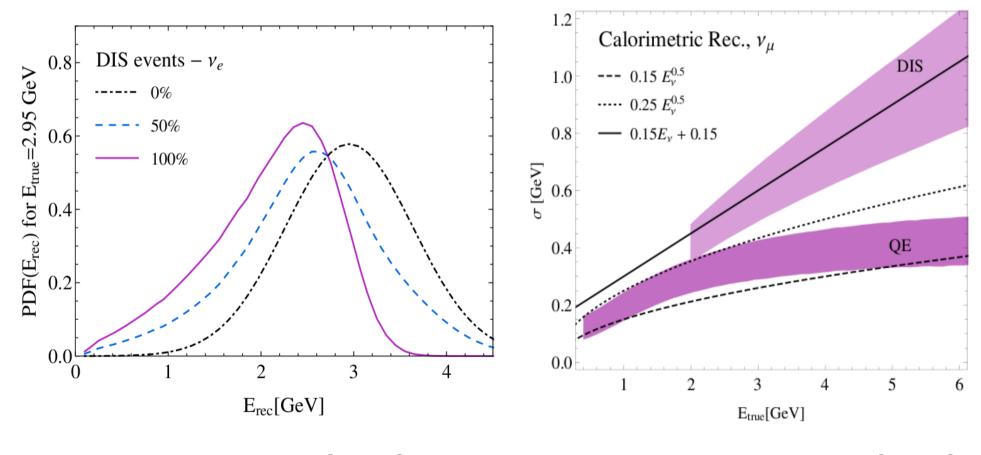
Coloma, Huber, Mariani and Jen, 1311.4506 [hep-ph]

# Does this improve with calorimetry?

Calorimetry relies on observing all particles produced in the interaction



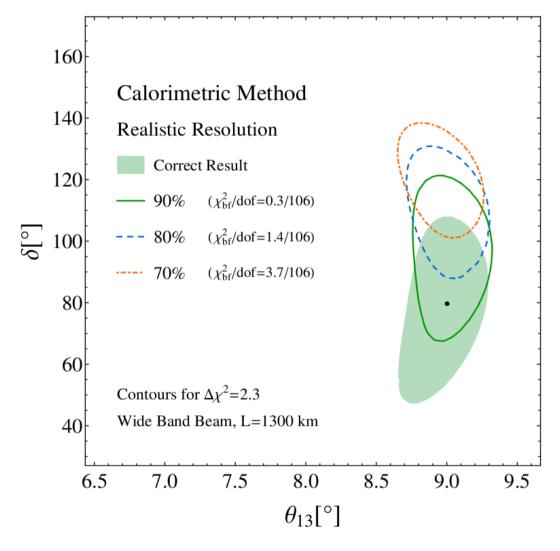
#### Does this improve with calorimetry?



Ankowski et al, 1507.08561 [hep-ph]

Ankowski et al, 1507.08560 [hep-ph]

## Does this improve with calorimetry?



Ankowski et al, 1507.08561 [hep-ph]

# Summary (I/II)

- The most relevant systematics for appearance experiments are those related to cross sections. Challenges:
  - Unavailability of final flavor at the near detector
  - near-far detector extrapolation
- Systematic effects may be kept under control under several assumptions:
  - no flux shape uncertainties
  - no cross section shape uncertainties
  - disappearance data can be used to reduce uncertainties in appearance Challenging!

# Summary (II/II)

- Shape uncertainties are more dangerous. We find:
  - Large impact on the determination of the disappearance parameters for kinematic reconstruction due to non-QE contamination of the QE sample
  - Significant bias for deltaCP from missing energy in calorimetric detectors, if not accurately calibrated
- Failure to include nuclear effects properly may induce significant bias on the oscillation parameters

#### Thanks!

### Matrix generation details

Detection thresholds: 20 MeV for mesons and 40 MeV for protons Detection efficiencies: 60% for pi0, 80% for other mesons, 50% for protons Resolution details assumed:

$$\sigma(|\mathbf{k}_{\mu}|) = 0.02|\mathbf{k}_{\mu}| \quad \text{and} \quad \sigma(\theta) = 0.7^{\circ},$$
  

$$\sigma(E_{e}) = 0.10E_{e} \quad \text{and} \quad \sigma(\theta) = 2.8^{\circ}$$
  

$$\frac{\sigma(E_{\pi^{0}})}{E_{\pi^{0}}} = \max\left\{\frac{a_{\pi^{0}}}{\sqrt{E_{\pi^{0}}}}, \frac{b_{\pi^{0}}}{E_{\pi^{0}}}\right\} \quad \frac{\sigma(E_{h})}{E_{h}} = \max\left\{\frac{a_{h}}{\sqrt{E_{h}}}, b_{h}\right\}$$
  

$$a_{\pi^{0}} = 0.107 \text{ and } b_{\pi^{0}} = 0.02, \qquad a_{h} = 0.145 \text{ and } b_{h} = 0.067.$$

(Neutrons assumed to exit undetected)

#### Matrix generation details

Calorimetric reconstruction:

knocked-out

Energy of mesons produced

Kinematic reconstruction:

$$E_{\nu}^{\text{kin}} = \frac{2(nM - \epsilon_n)E_{\ell} + W^2 - (nM - \epsilon_n)^2 - m_{\ell}^2}{2(nM - \epsilon_n - E_{\ell} + |\mathbf{k}_{\ell}| \cos \theta)}$$
Single nucleon  
Number of nucleons knocked Invariant hadronic mass squared.  
out of the nucleus (For single-nucleon knock-out, W<sup>2</sup> = M<sup>2</sup>)

	SB			BB			NF		
Systematics	Opt.	Def.	Cons.	Opt.	Def.	Cons.	Opt.	Def.	Cons.
Fiducial volume ND	0.2%	0.5%	1%	0.2%	0.5%	1%	0.2%	0.5%	1%
Fiducial volume FD	1%	2.5%	5%	1%	2.5%	5%	1%	2.5%	5%
(incl. near-far extrap.)									
Flux error signal $\nu$	5%	7.5%	10%	1%	2%	2.5%	0.1%	0.5%	1%
Flux error background $\nu$	10%	15%	20%	correlated			correlated		
Flux error signal $\bar{\nu}$	10%	15%	20%	1%	2%	2.5%	0.1%	0.5%	1%
Flux error background $\bar{\nu}$	20%	30%	40%	correlated			correlated		
Background uncertainty	5%	7.5%	10%	5%	7.5%	10%	10%	15%	20%
Cross secs $\times$ eff. QE <sup>†</sup>	10%	15%	20%	10%	15%	20%	10%	15%	20%
Cross secs $\times$ eff. RES <sup>†</sup>	10%	15%	20%	10%	15%	20%	10%	15%	20%
Cross secs $\times$ eff. DIS <sup>†</sup>	5%	7.5%	10%	5%	7.5%	10%	5%	7.5%	10%
Effec. ratio $\nu_e/\nu_\mu \ QE^*$	3.5%	11%	-	3.5%	11%	-	_	_	_
Effec. ratio $\nu_e/\nu_\mu$ RES <sup>*</sup>	2.7%	5.4%	-	2.7%	5.4%	-	-	_	-
Effec. ratio $\nu_e/\nu_\mu$ DIS <sup>*</sup>	2.5%	5.1%	-	2.5%	5.1%	-	-	_	_
Matter density	1%	2%	5%	1%	2%	5%	1%	2%	5%

Coloma, Huber, Kopp, Winter, 1209.5973 [hep-ph]

# Toy model

 Neglecting all FSI and multinucleon contributions, we can compute the number of events as:

$$N_i^{QE} = \sigma_{QE}(E_i)\phi(E_i)P_{\mu\mu}(E_i)$$

 However, in practice we will observe a different distribution at the detector, given by:

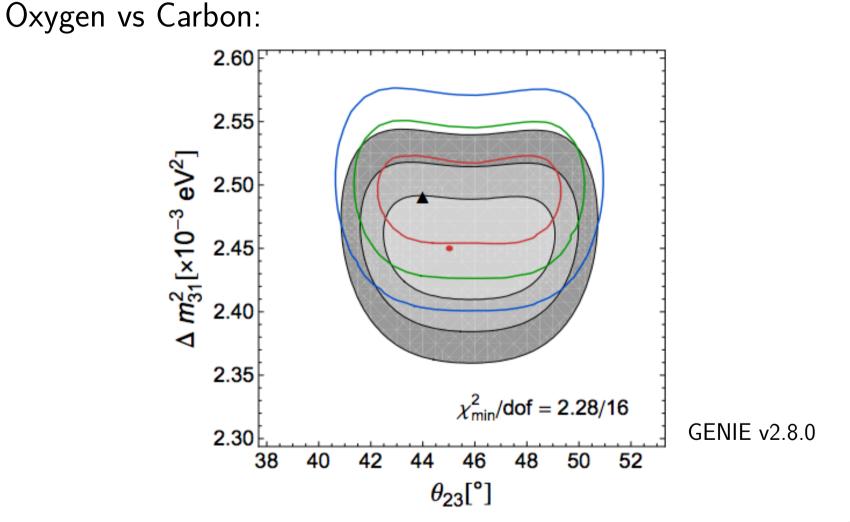
$$N_i^{QE-like} = \sum_j M_{ij}^{QE} N_j^{QE} + \sum_{non-QE} \sum_j M_{ij}^{non-QE} N_j^{non-QE}$$

An intermediate situation would most likely take place:

$$N_i^{test}(\alpha) = \alpha N_i^{QE} + (1 - \alpha) N_i^{QE-like}$$

Coloma and Huber, 1307.1243 [hep-ph]

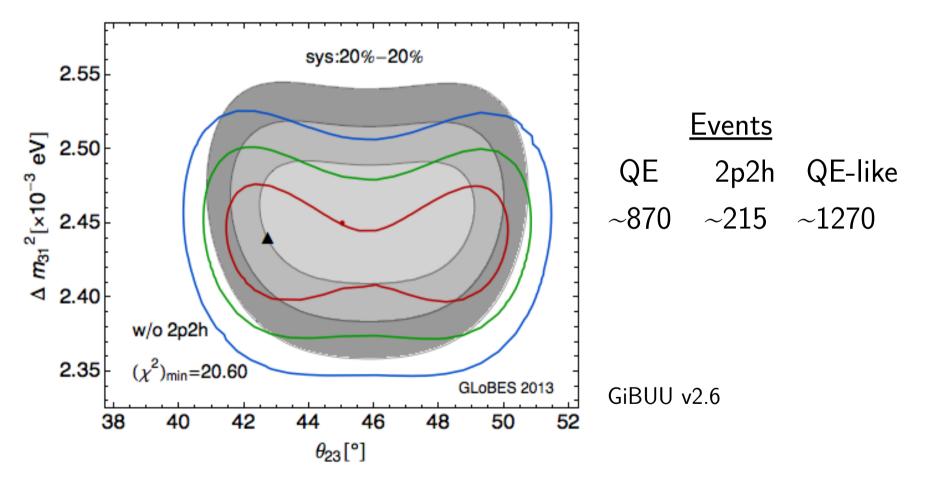
#### Impact of target nucleus



Coloma, Huber, Mariani and Jen, 1311.4506 [hep-ph]

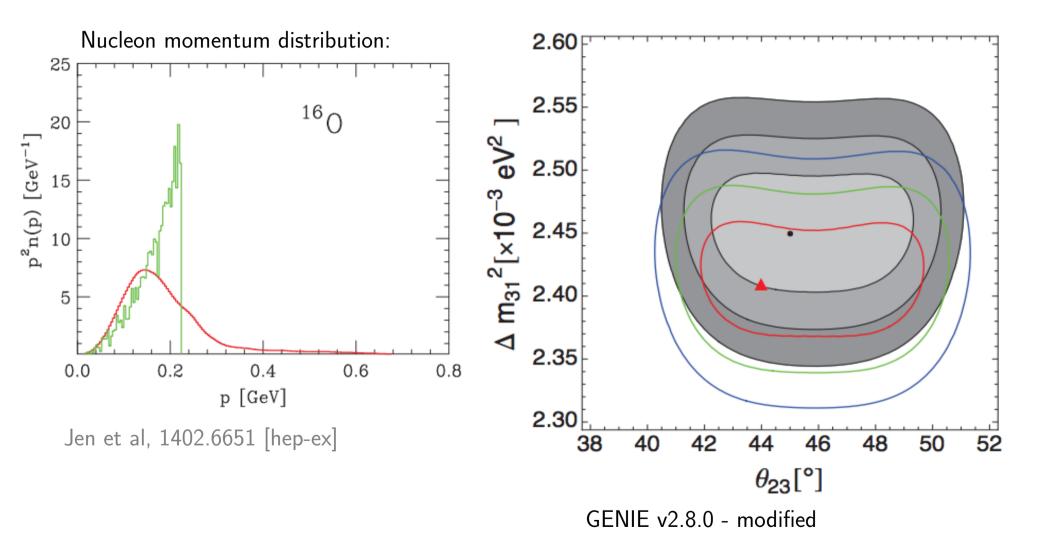
#### Impact of 2p2h

Even if we get all contributions right except 2p2h...



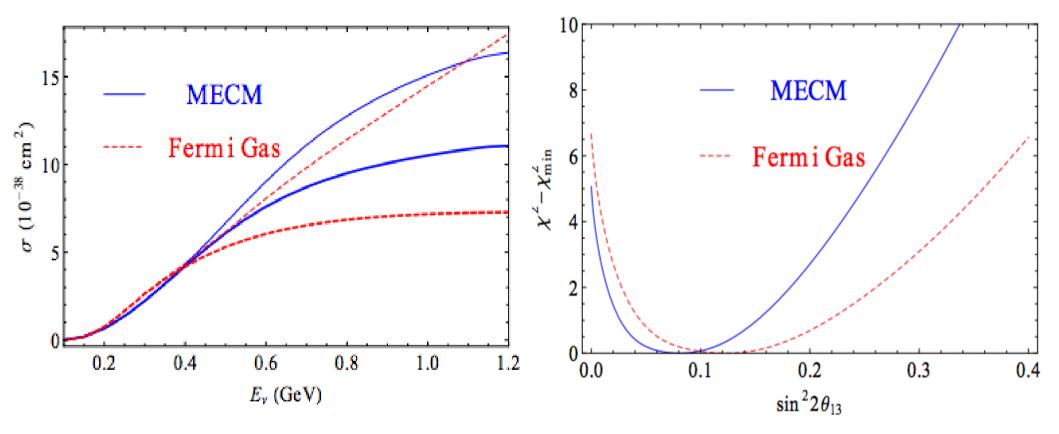
Coloma and Huber, 1307.1243 [hep-ph]

#### Other factors: RFGM vs SF



#### Cross section models

Impact on an analysis which reproduces T2K results in 1106.2822 [hep-ex]



Martini, Meloni, 1203.3335 [hep-ph] MECM = model from Martini, Ericson, Chanfray, Marteau, 0910.2622 [nucl-th]