# Production and polarization of $J/\psi$ to $\mathcal{O}(\alpha_s^3)$ in the improved color evaporation model

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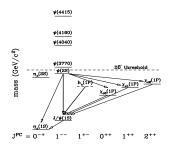
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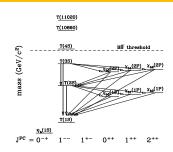


#### Overview

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  - Polarization
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- 2 ICEM Approach
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- Conclusion and Future

# Quarkonium Families





#### Quarkonia: bound states of $c\overline{c}$ or $b\overline{b}$

- ullet combination of two spin 1/2 particles with orbital angular momentum ullet different spin states  $^{2S+1}L_J$
- all color singlets  ${}^{2S+1}L_J{}^{[1]}$
- produced in hh,  $\gamma$ p,  $\gamma\gamma$ , and e<sup>+</sup>e<sup>-</sup>
- S states below the  $H\overline{H}$  (H=D,B) threshold decay electromagnetically into  $\ell^+\ell^-$

# Polarization and Angular Distribution

$$\begin{split} |\psi\rangle &= a_{-1} \, |J_z = -1\rangle + a_0 \, |J_z = 0\rangle + a_{+1} \, |J_z = +1\rangle, \qquad \qquad \sum |a_{J_z}|^2 = 1 \\ \lambda_{\vartheta} &= \frac{1 - 3|a_0|^2}{1 + |a_0|^2}, \qquad \qquad \lambda_{\varphi} = \frac{2\text{Re}[a_{+1}a_{-1}^*]}{1 + |a_0|^2}, \qquad \qquad \lambda_{\vartheta\varphi} = \frac{\sqrt{2}\text{Re}[a_0^*(a_+ - a_-)]}{1 + |a_0|^2} \end{split}$$

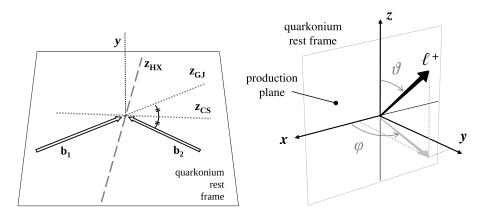
$$rac{d\sigma}{d\Omega} \propto rac{1}{3+\lambda_{artheta}} \Biggl[ 1 + \lambda_{artheta} \cos^2{artheta} + \lambda_{arphi} \sin^2{artheta} \cos(2arphi) + \lambda_{artheta arphi} \sin(2artheta) \cos{arphi} \Biggr]$$

- For a single elementary process, the polarized-to-total cross section can be calculated as  $a_{J_z}$ 's. Combinations of  $a_{J_z}$ 's gives different angular distributions.
- However, there is no combination that would give  $\lambda_{\vartheta} = \lambda_{\varphi} = \lambda_{\vartheta\varphi} = 0$ .
- An unpolarized production can only be described by a mixture of sub-processes or randomization modeling.



Pietro Faccioli, QWG 2010.

#### Polarization Measurement

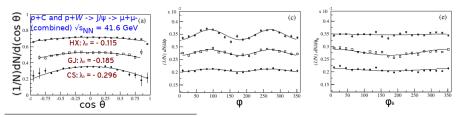


- There are three commonly used choices for the z-axis, namely  $z_{HX}$  (helicity),  $z_{CS}$  (Collins-Soper), and  $z_{GJ}$  (Gottfried-Jackson)
- $\vartheta$  is defined as the angle between the z-axis and the direction of travel for the  $\ell^+$  in the quarkonium rest frame

# **Extracting Polarization**

$$rac{d\sigma}{d\Omega} ~\propto ~ rac{1}{3+\lambda_{artheta}}[1+\lambda_{artheta}\cos^2artheta+\lambda_{arphi}\sin^2artheta\cos(2arphi)+\lambda_{arthetaarphi}\sin(2artheta)\cosarphi]$$

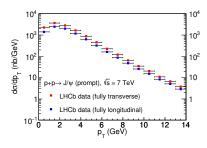
- Polarization parameters can be obtained by fitting the angular spectra as a function of  $\vartheta$  and  $\varphi$
- One can write  $\varphi_{\vartheta} = \varphi \frac{\pi}{2} \mp \frac{\pi}{4}$  for  $\cos \vartheta \lessgtr 0$ , then<sup>[1]</sup>
- ullet  $rac{d\sigma}{darphi_artheta} \propto 1 + rac{\sqrt{2}\lambda_{arthetaarphi}}{3+\lambda_artheta}\cosarphi_artheta$

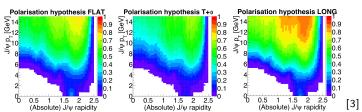


 $^{1}$ I. Abt et al. (HERA-B Collaboration), Eur. Phys. J. C 60, 517 (2009).

# Importance of Polarization

- Polarization predictions are strong tests of production models
- Detector acceptance depends on polarization hypothesis
- Understanding polarization helps narrow systematic uncertainties





<sup>2</sup>R. Aaij *et al.* (LHCb Collaboration), Eur. Phys. J. C **71**, 1645 (2011). <sup>3</sup>G. Aad *et al.* (ATLAS Collaboration), Nucl. Phys. B **850**, 387 (2011). [2]

# Quarkonium Polarization Puzzle

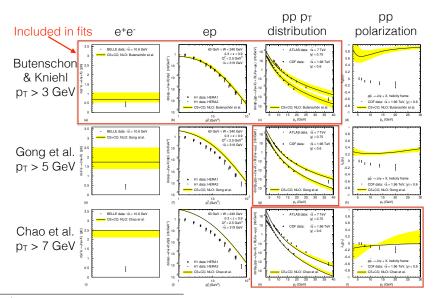
#### Quarkonium Polarization Puzzle

- mechanism of producing quarkonium has not yet been understood
- non-relativistic QCD (NRQCD), a common method to calculate quarkonium production, has difficulties describing yield and polarization simultaneously with a low- $p_T$  cut

# Non Relativistic QCD (NRQCD) [Bodwin, Braaten, Lepage 95]

- ullet e.g. for  $J/\psi$ ,  $\sigma_{J/\psi}=\sum_{\pmb{n}}\sigma_{c\overline{c}[\pmb{n}]}\langle\mathcal{O}^{J/\psi}[\pmb{n}]\rangle$
- both color singlet term  $n = {}^3S_1^{[1]}$  and color octet terms  ${}^1S_0^{[8]}$ ,  ${}^3S_1^{[8]}$ , and  ${}^3P_J^{[8]}$  contributes to the production
- mixing of Long Distance Matrix Elements (LDMEs =  $\langle \mathcal{O}^{J/\psi}[n] \rangle$ ) are determined by fitting to data, usually  $p_T$  distributions above some  $p_T$  cut

### Polarization Puzzle<sup>[4]</sup>



<sup>&</sup>lt;sup>4</sup>N. Brambilla *et al.*, Eur. Phys. J. C **74**, 2981 (2014)

# The Improved Color Evaporation Model (ICEM)

[Ma, Vogt (PRD **94**, 114029 (2016).)]

$$\sigma = F_{\mathcal{Q}} \sum_{i,j} \int_{M_{\psi}}^{2m_H} dM \int dx_i dx_j f_i(x_i, \mu_F) f_j(x_j, \mu_F) d\hat{\sigma}_{ij \to c\bar{c} + X}(p_{c\bar{c}}, \mu_R)|_{p_{c\bar{c}} = \frac{M}{M_{\psi}} p_{\psi}},$$

where  $M_{\psi}$  is the mass of the charmonium state,  $\psi$ .

- all Quarkonium states are treated like  $Q\overline{Q}$  (Q=c,b) below  $H\overline{H}$  (H=D,B) threshold
- ullet all diagrams for  $Q\overline{Q}$  production included, independent of color
- ullet able to describe relative production of  $\psi(2{\sf S})$  to  $J/\psi$
- ullet fewer parameters than NRQCD (one  $F_{\mathcal{Q}}$  for each Quarkonium state)
- distinction between the momentum of the  $c\bar{c}$  pair and that of charmonium so that the  $p_T$  spectra will be softer and thus may explain the high  $p_T$  data better
- $F_{\mathcal{Q}}$  is fixed by comparison of NLO calculation of  $\sigma_{\mathcal{Q}}^{CEM}$  to  $\sqrt{s}$  for  $J/\psi$  and  $\Upsilon$ ,  $\sigma(x_F>0)$  and  $Bd\sigma/dy|_{y=0}$  for  $J/\psi$ ,  $Bd\sigma/dy|_{y=0}$  for  $\Upsilon$

# Collinear Polarized ICEM at $\mathcal{O}(\alpha_s^3)^{[5]}$

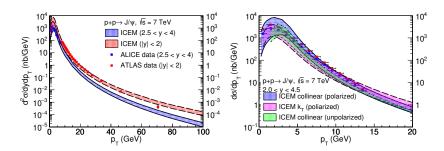
#### Production distribution

$$\frac{d^2\sigma}{d\rho_T dy} = F_{\mathcal{Q}} \sum_{i,j=\{q,\bar{q},g\}} \int_{M_{\mathcal{Q}}}^{2m_H} dM_{\psi} \int d\hat{s} dx_1 dx_2 f_{i/p}(x_1,\mu^2) f_{j/p}(x_2,\mu^2) d\hat{\sigma}_{ij\to c\bar{c}+X} \; , \label{eq:delta-dy}$$

- We consider all 16 diagrams from gg $\rightarrow$  c $\bar{c}$ g, 5(+5) from gq( $\bar{q}$ ) $\rightarrow$  c $\bar{c}$  q( $\bar{q}$ ), and 5 from q $\bar{q}$  $\rightarrow$  c $\bar{c}$ g with the projection operator applied at the diagram level.
- The  $c\bar{c}$  produced are the proto- $J/\psi$  before hardonization.
- We used the CT14 PDFs in our calculations.
- $k_T$ -smearing is applied to the initial state partons to provide better description at low  $p_T$
- First  $p_T$ -dependent polarization results using collinear factorization
- 1.18 <  $m_c$  < 1.36 GeV,  $\mu_F/m_T = 2.1^{+2.55}_{-0.85}$ ,  $\mu_R/m_T = 1.6^{+0.11}_{-0.12}$
- same set of variations used in MV (2016) and NVF [PRC **87**, 014908 (2013)]

<sup>5</sup>V. Cheung and R. Vogt, PRD (accepted).

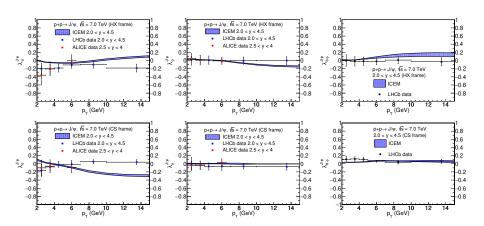
# Collinear ICEM Unpolarized Cross Sections<sup>[5]</sup>



- $k_T$ -smearing gives a small kick  $< k_T^2 > \sim 1 \text{ GeV}^2$  to the inital state parton.
- The uncertainty band<sup>[5]</sup> is constructed by varying the charm quark mass, factorization scale, and renormalization scale.
- We find agreement with the  $p_T$ -distribution measured by the LHCb<sup>[6]</sup>.
- We also find agreement with the unpolarized ICEM calculations [MV (2016)].

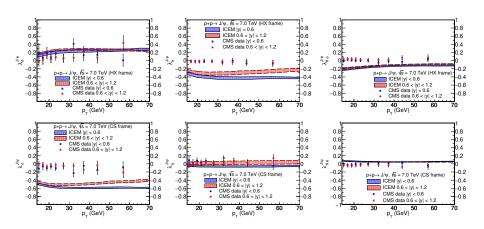
<sup>6</sup>R. Aaij et al. (LHCb Collaboration), Eur. Phys. J. C 73, 2631 (2013).

# Polarization Parameters in Collinear ICEM<sup>[5]</sup>



- We find agreement with LHCb data<sup>[6]</sup> at small and moderate  $p_T$ .
- Difference between the prediction and experimental results in high  $p_T$  is frame dependent.

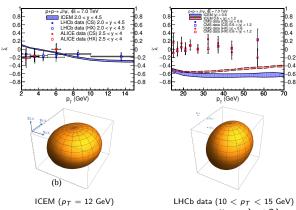
# Polarization Parameters in Collinear ICEM<sup>[5]</sup>



- We find partial agreement with CMS data<sup>[7]</sup> at high  $p_T$
- Agreement is frame dependent.

<sup>&</sup>lt;sup>7</sup>S. Chatrchyan et al. (CMS Collaboration), Phys. Lett. B **727**, 381-402 (2013).

# Invariant Polarization Parameter in Collinear ICEM<sup>[5]</sup>



- ullet The frame-invariant polarization parameter  $\tilde{\lambda}=\frac{\lambda_{\vartheta}+3\lambda_{\varphi}}{1-\lambda_{\varphi}}$
- Comparing the frame-invariant polarization paremeter removes frame-induced kinematic dependencies
- We find agreement with the invariant polarization at LHCb<sup>[6]</sup>, but discrepancy between high  $p_T$  data at CMS<sup>[7]</sup>.

#### Conclusion and Future

#### (I)CEM

- Less rigorous
- Fewer fit parameters
- Applied extensively to only hadroproduction (so far)

#### **NRQCD**

- More rigorous
- More fit parameters
- Applied to all collision systems

### In this talk, I

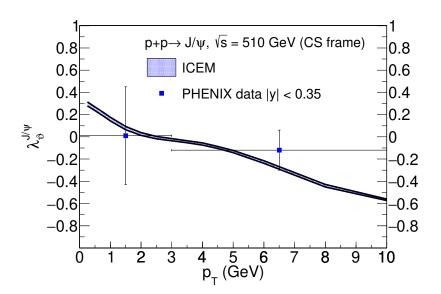
- outlined the quarkonium polarization puzzle
- showed the latest attempt to solve the polarization puzzle in the ICEM

#### In the future, we

- anticpate the feed down from P states can explain the discrepancies in high  $p_T$ .
- will move from hadroproduction to other collision systems.

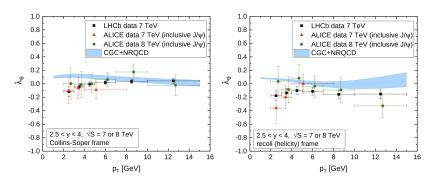
# Backup Slides

# Comparison to PHENIX Results



# CGC+NRQCD<sup>[8]</sup>

- is a solution to the polarization puzzle where gluon distribution is calculated using CGC and the conversion of  $Q\bar{Q}$  is described by NRQCD formulation
- ullet able to describe all polarization parameters for  $p_{\mathcal{T}} < 15$  GeV



<sup>&</sup>lt;sup>8</sup>Y. Q. Ma, T. Stebel, R. Venugopalan, JHEP12 (2018) 057.