



Unifying the URR PT approaches and covariance work in a consistent framework

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CSEWG

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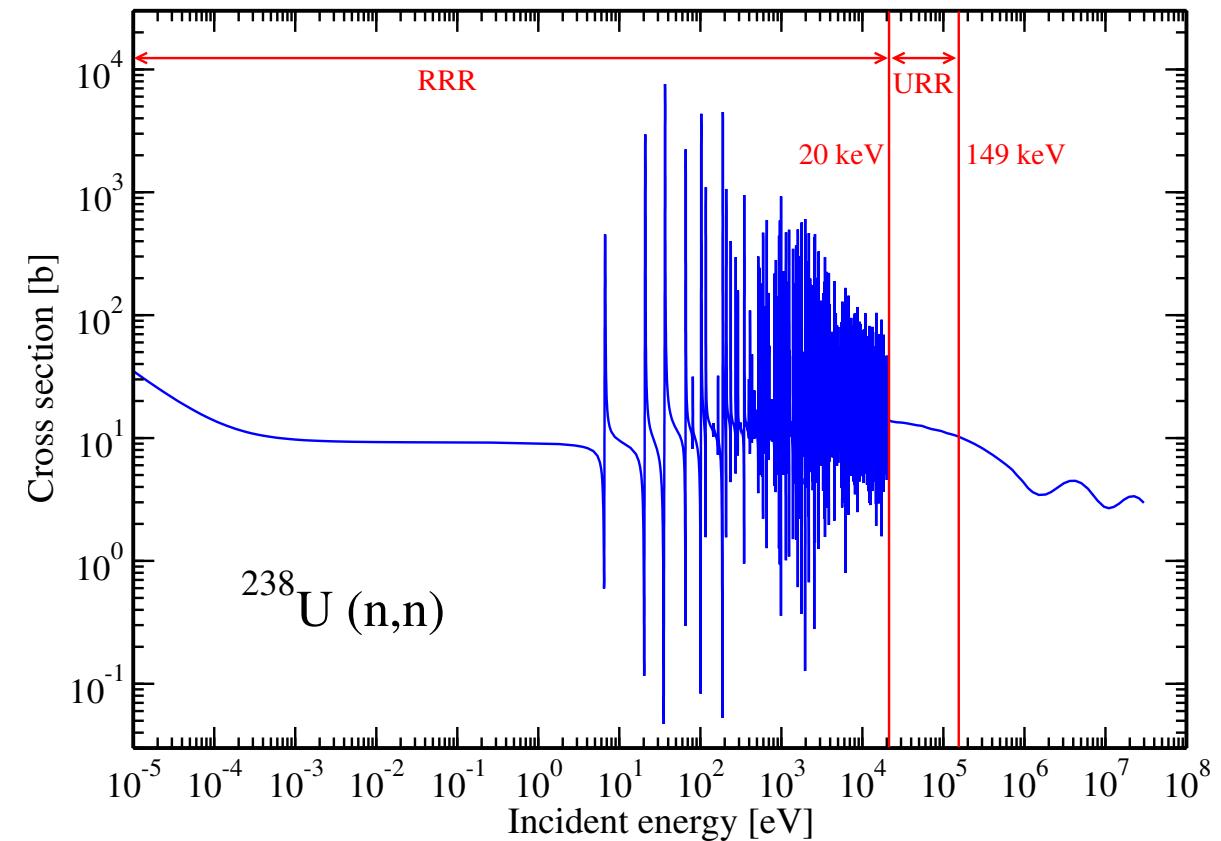


Unresolved resonance region

Thermalization of neutrons in ^{238}U by elastic scattering requires ~ 2200 collisions

Energy loss by elastic scattering in the URR amounts to ~ 240 collisions
 $\sim 11\%$ of collisions occur in the URR

This is something we should care about



Unresolved resonance region

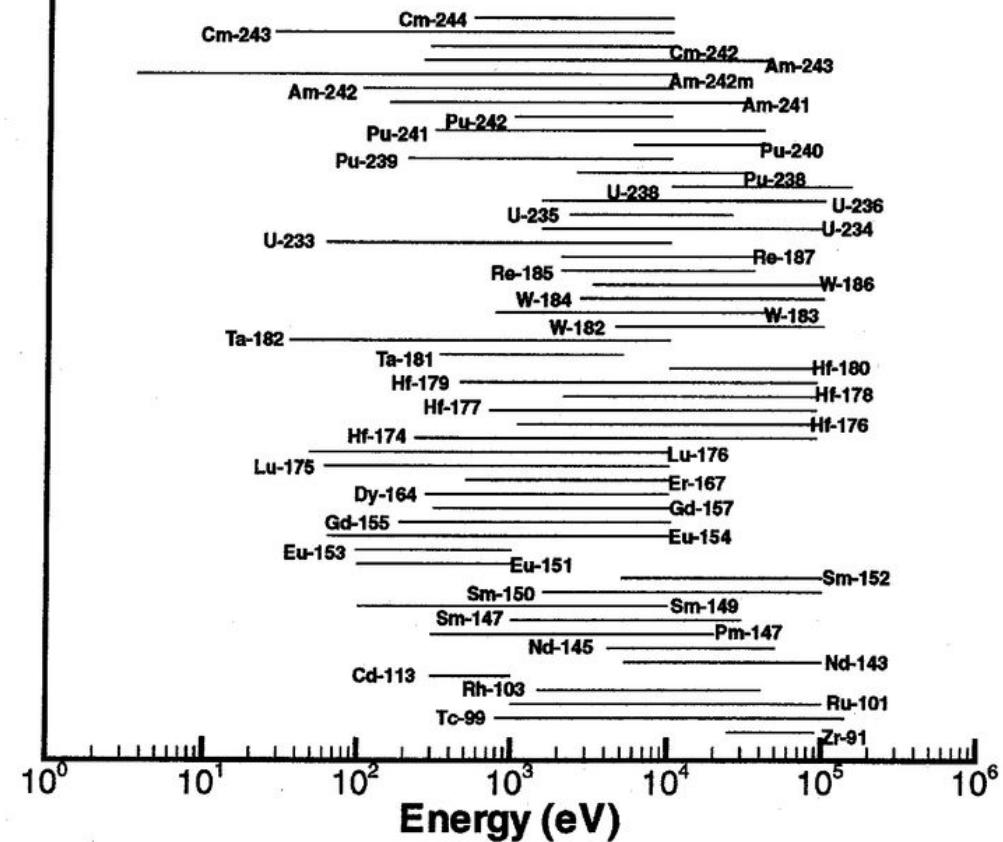
Resonances are too narrow to resolve experimentally

We cannot determine individual resonance parameters

We can only determine averages of parameters

Enough to determine the cross section PDF

URR Energy Ranges



Thomas Sutton, KAPL

Cross section probability distributions

The probability distribution is a measure of our ignorance of the cross section

If we know the cross section exactly, the PDF is a delta function

$$P(\tilde{\sigma}|E) = \delta(\tilde{\sigma} - \sigma(E))$$

$$P(\sigma|E, T = 0, \{x\})$$

If we want to express our ignorance of the URR parameters

$$P(\tilde{\sigma}|E, \bar{x}) = \int d\mathbf{x} P(\mathbf{x}|\bar{x}) \delta(\tilde{\sigma} - \sigma(E, \mathbf{x}))$$

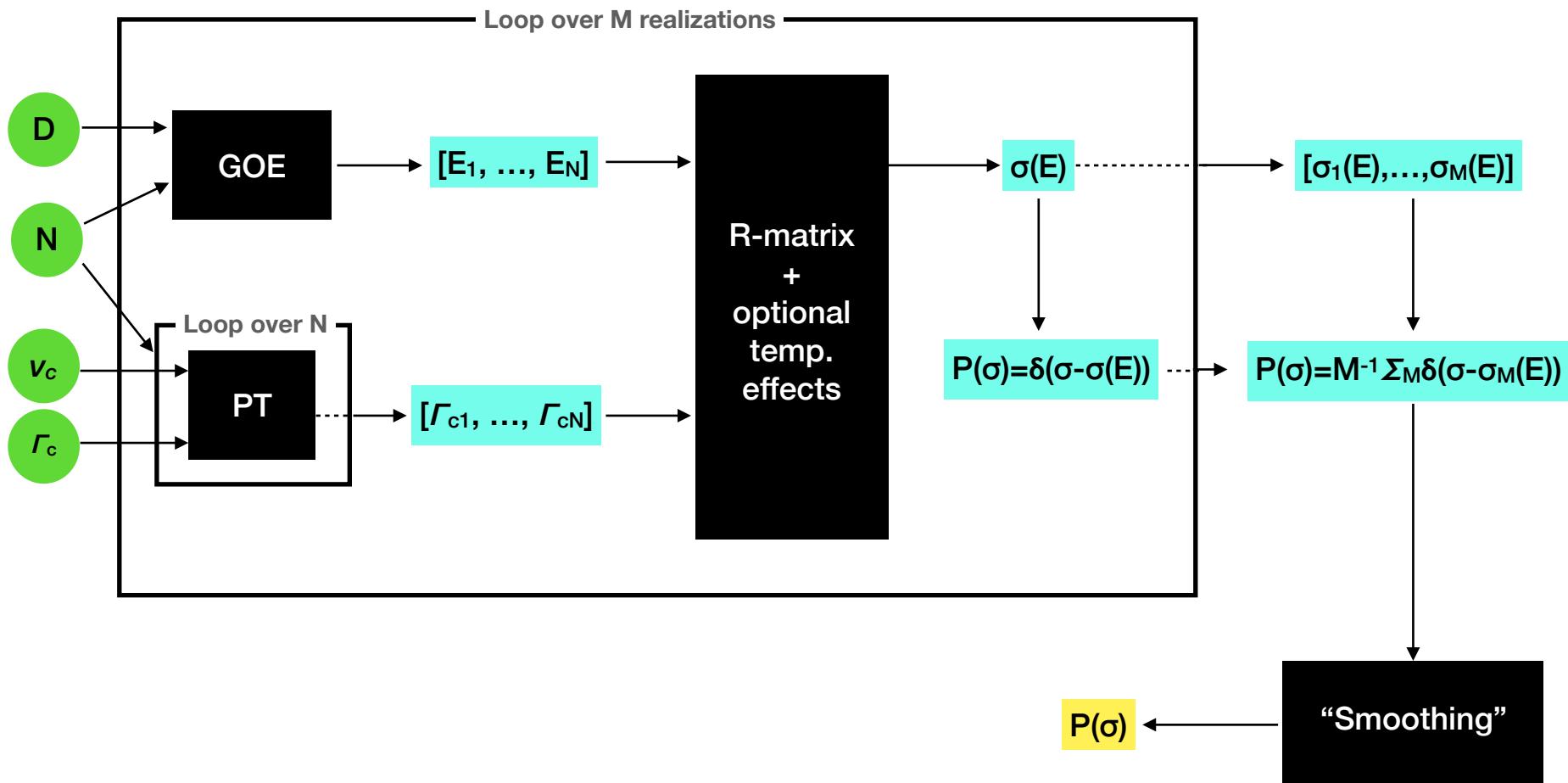
↑
Resonance
parameter
PDF

$$P(\{x\}|\bar{x})$$

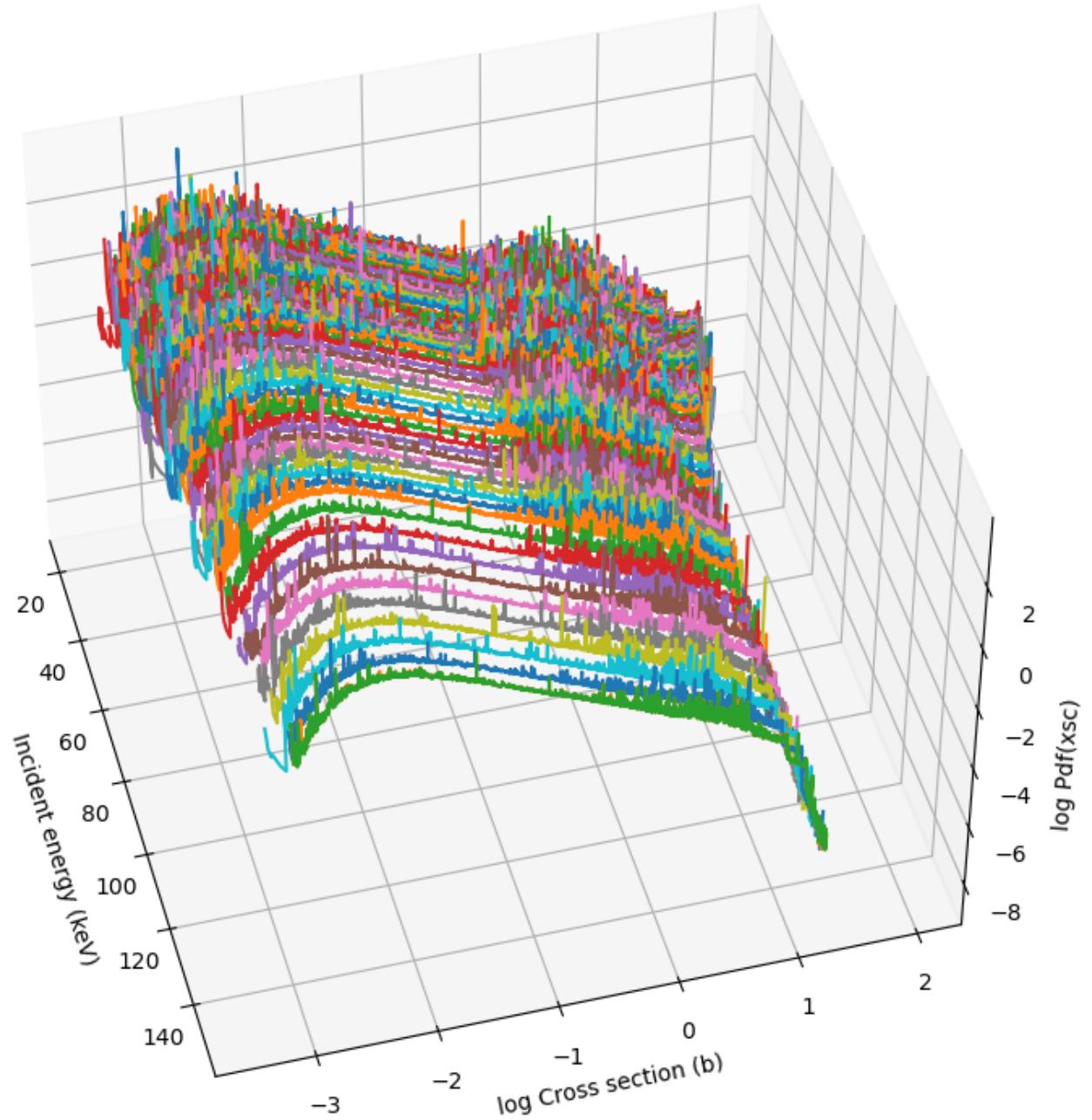
URR theory

$$P(\sigma|E, T = 0, \bar{x})$$

Resonance Ladders in FUDGE



U238 capture 0.0K



Including target temperature effects

Standard method

$$P_{\text{STND}}(\tilde{\sigma}|E, T) = \delta \left(\tilde{\sigma} - \int_{[V_t, V_r > 0]} dV_t P(V_t|T) \frac{V_r}{V} \sigma(E_r) \right)$$

Including URR parameter variations

$$P_{\text{STND}}(\tilde{\sigma}|E, T, \bar{x}) = \int dx P(x|\bar{x}) P_{\text{STND}}(\tilde{\sigma}|E, T)$$

Standard

1. Sample cross section
2. **Doppler broadening**
3. **Accumulate PDF**

$$P(\bar{\sigma}|E, T, \{x\})$$

$$P(\{x\}|\bar{x})$$

URR theory

$$P(\bar{\sigma}|E, T, \bar{x})$$

We propose an alternative method

Standard method

$$P_{\text{STND}}(\tilde{\sigma}|E, T) = \delta \left(\tilde{\sigma} - \int_{[V_t, V_r > 0]} dV_t P(V_t|T) \frac{V_r}{V} \sigma(E_r) \right)$$

Alternate method

$$P_{\text{ALT}}(\tilde{\sigma}|E, T) = \int_{[V_t, V_r > 0]} dV_t P(V_t|T) \delta \left(\tilde{\sigma} - \frac{V_r}{V} \sigma(E_r) \right)$$

Zero-temperature
PDF

Standard

1. Sample cross section
- 2. Doppler broadening**
- 3. Accumulate PDF**

Alternate

1. Sample cross section
- 2. Accumulate PDF**
- 3. Doppler broadening**

How are these PDF's connected?

An “history” in a reactor simulation represents many “real-world histories” ($N \gg 10^{10}$)

Each “real-world history” is assumed to be statistically independent so the central limit theorem applies

$$P_{\text{ALT}}(\tilde{\sigma}|E, T) = \int_{[V_t, V_r > 0]} dV_t P(V_t|T) \delta \left(\tilde{\sigma} - \frac{V_r}{V} \sigma(E_r) \right)$$

stddev $\sim 1/\sqrt{N}$

$$P_{\text{STND}}(\tilde{\sigma}|E, T) = \delta \left(\tilde{\sigma} - \int_{[V_t, V_r > 0]} dV_t P(V_t|T) \frac{V_r}{V} \sigma(E_r) \right)$$

Our alternate approach is essentially an event-by-event PDF

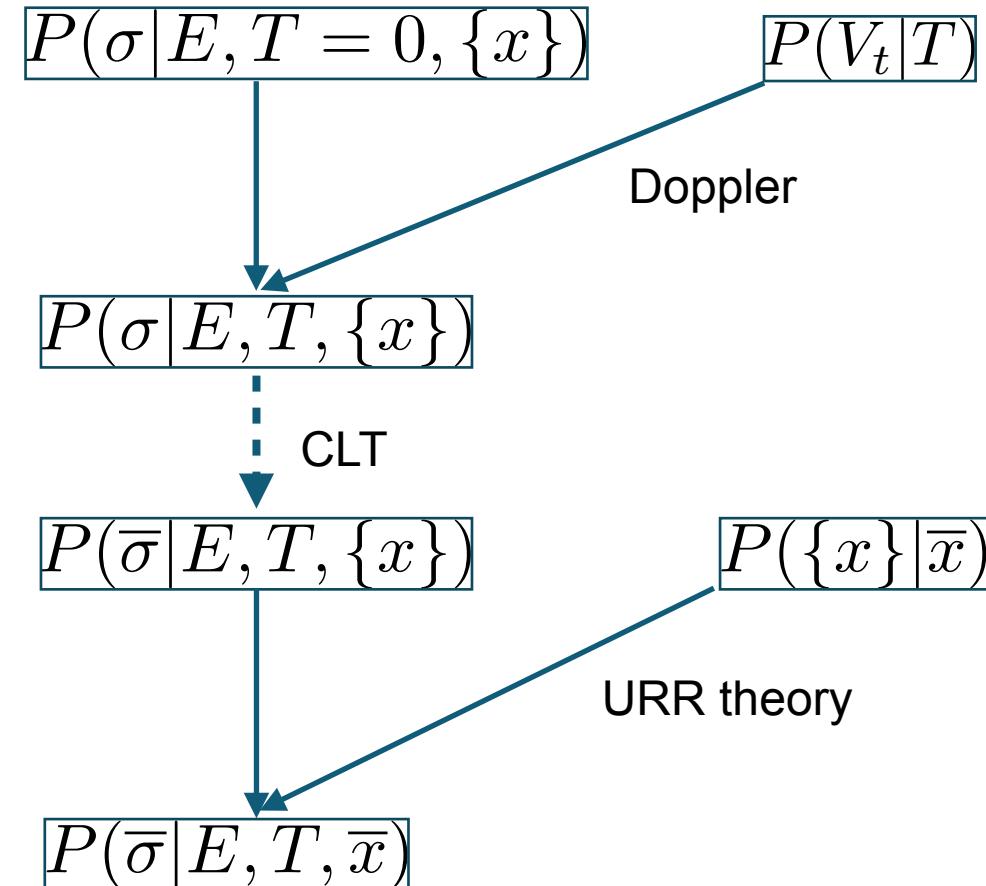
What this means for the standard approach

We've rephrased URR PT construction as composition of probabilities:

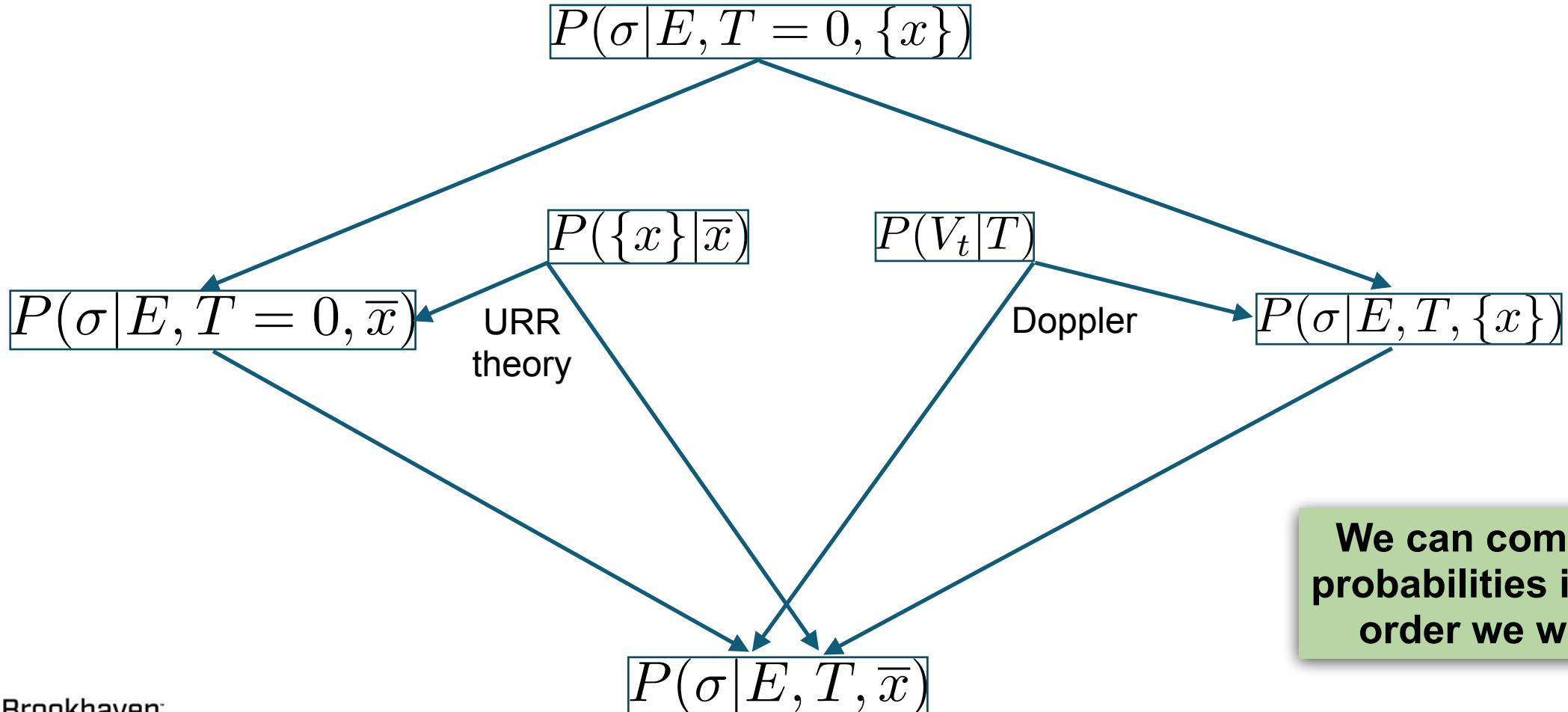
- target temp PDF
- URR parameters

Central Limit Theorem connects alternate(e-by-e) and standard URR PDFs

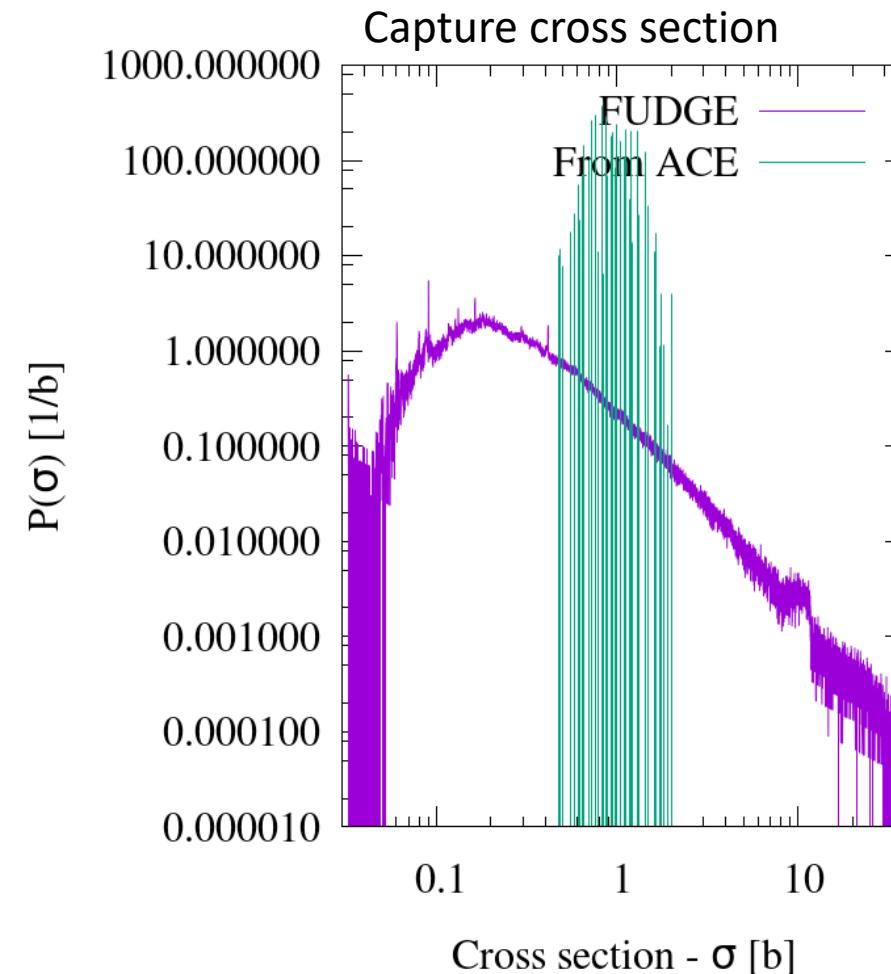
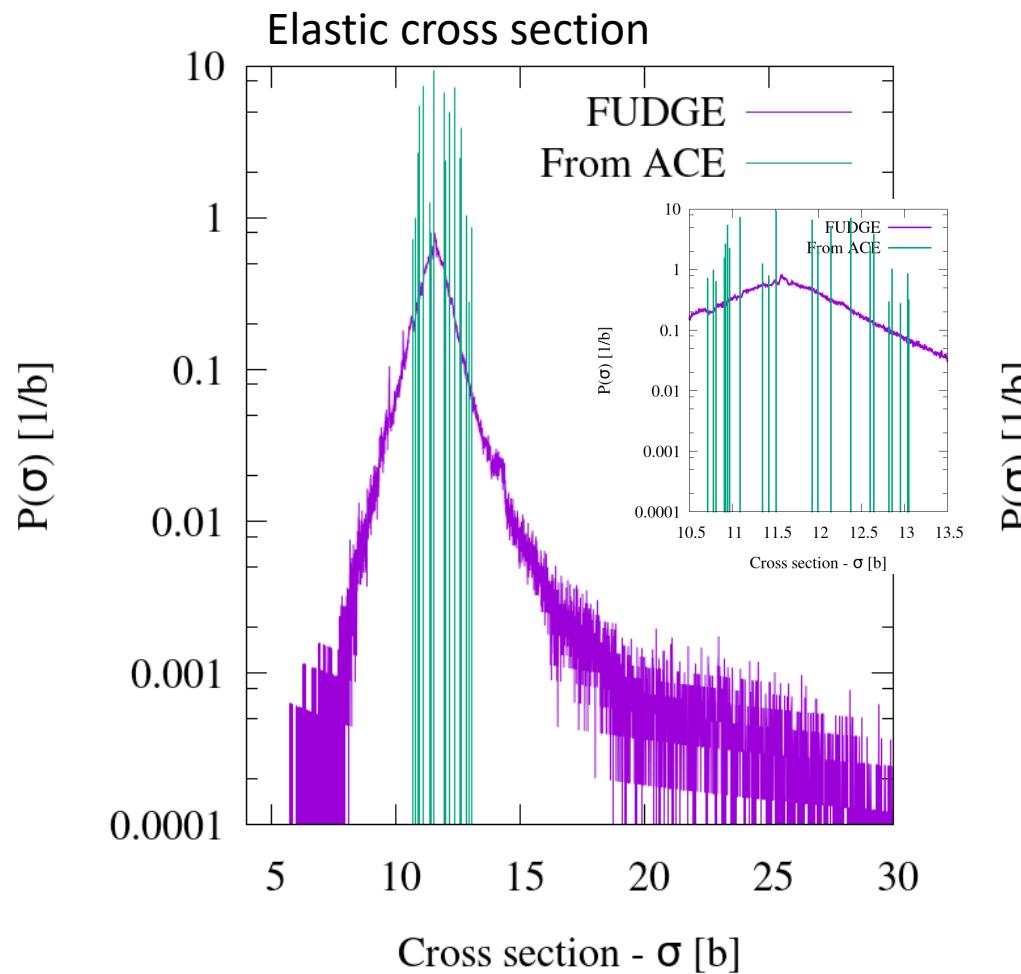
We can combine other sources of information, namely experiment!



What this means for the alternate (event-by-event) approach



FUDGE pdf vs ACE probability tables

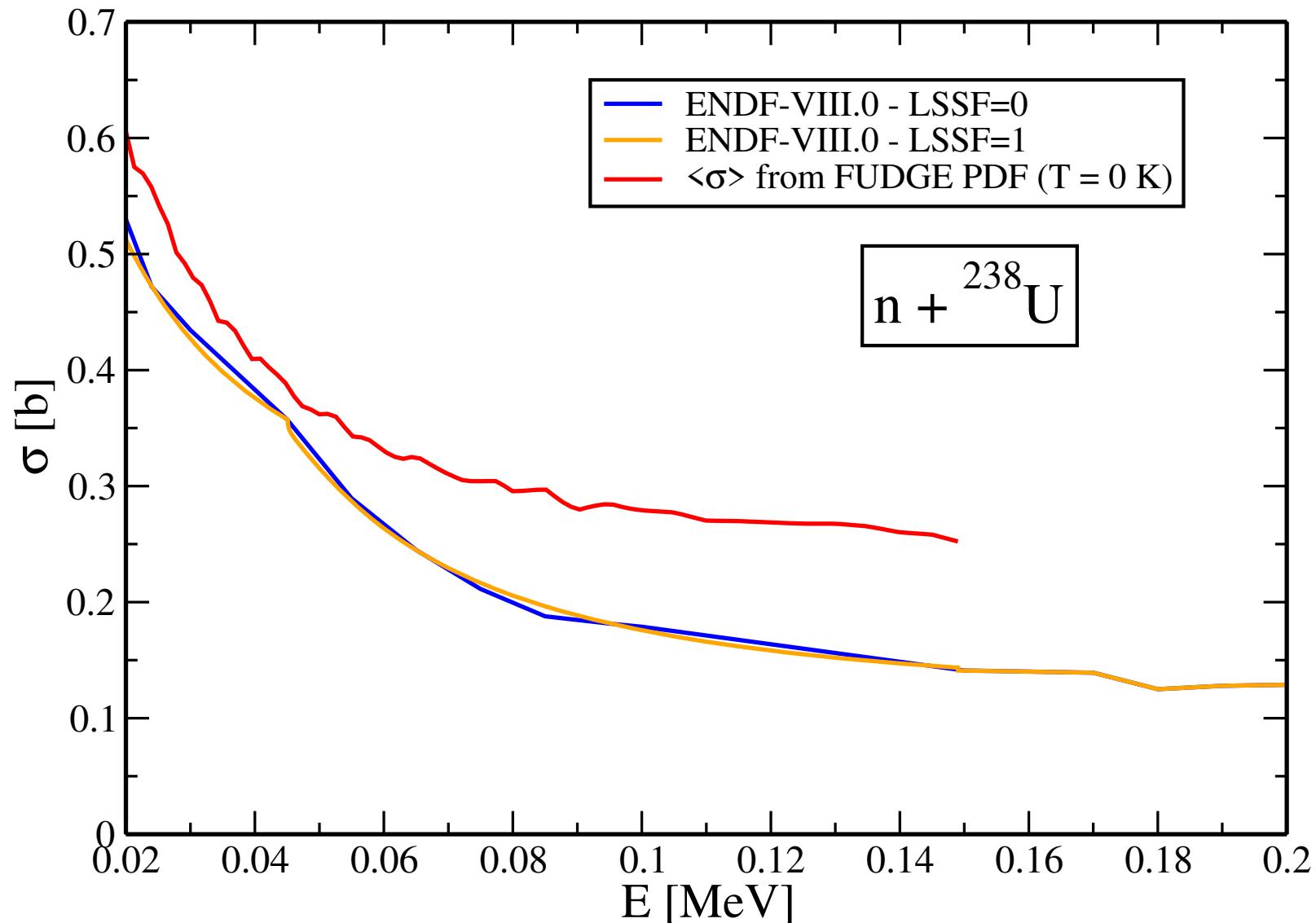


**Without CLT,
alternate
(e-by-e) PDF
is very broad
and capture
peak in
wrong place
(mean value
still same)**

Average cross section

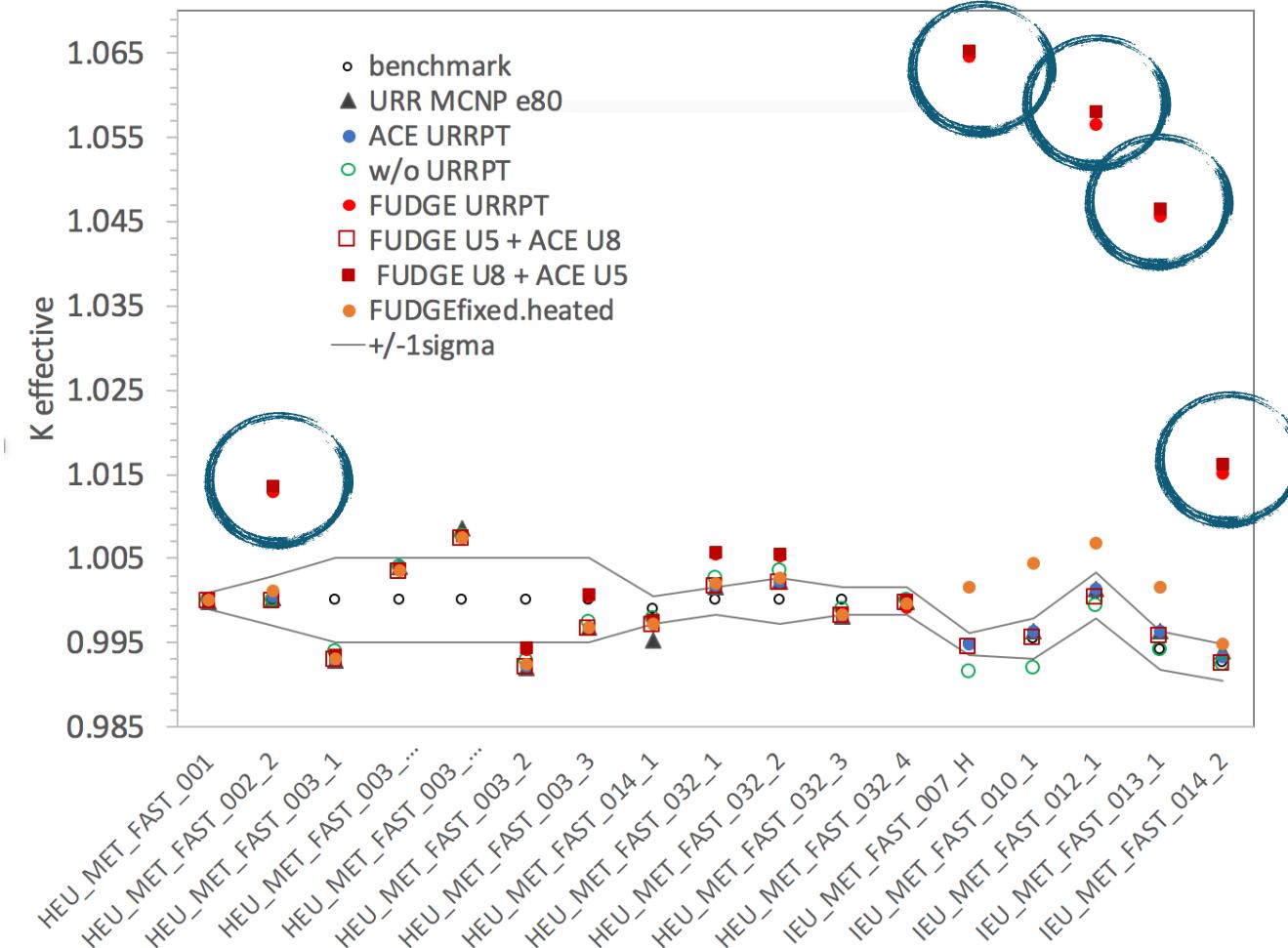
- The e-by-e PDF leads to an average cross section that is in disagreement with the one extracted from ENDF

$$\langle \sigma \rangle = \int d\sigma \sigma P(\sigma)$$



Testing probability tables in V&V suite

Poor initial results for ZPRs and BigTen now much improved

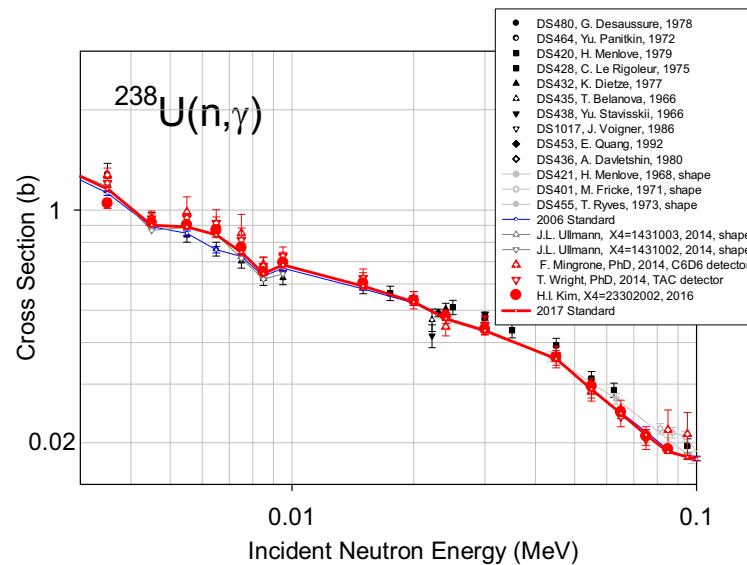


Because alternate (e-by-e) PDF is very broad & peaked in wrong place, it performs poorly in benchmarks

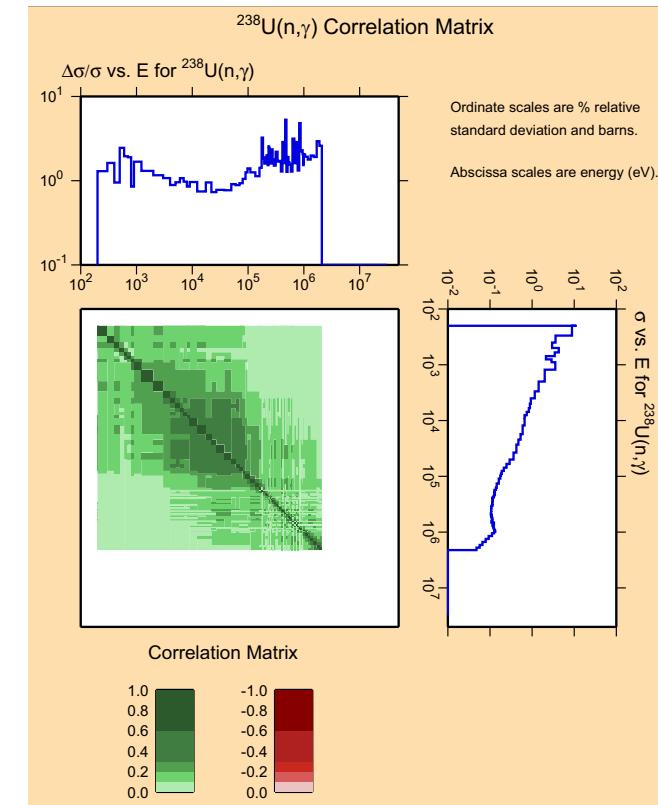
Combining theory and experiment

We can combine theory and experiment at **zero temperature**

$$P_{\text{th-exp}}(\tilde{\sigma}|E) = P_{\text{th}}(\tilde{\sigma}|E) P_{\text{exp}}(\tilde{\sigma}|E)$$



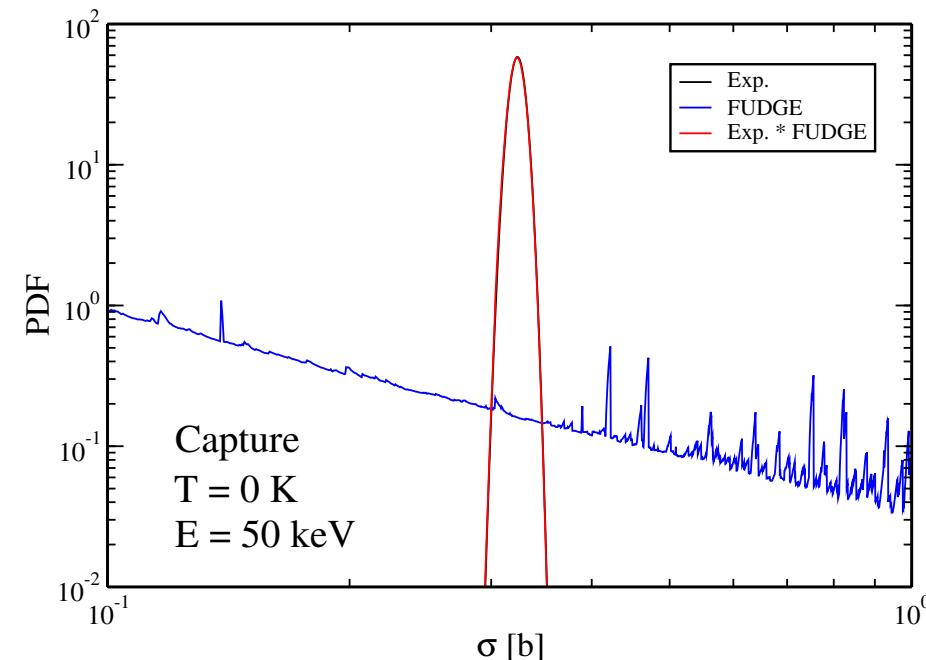
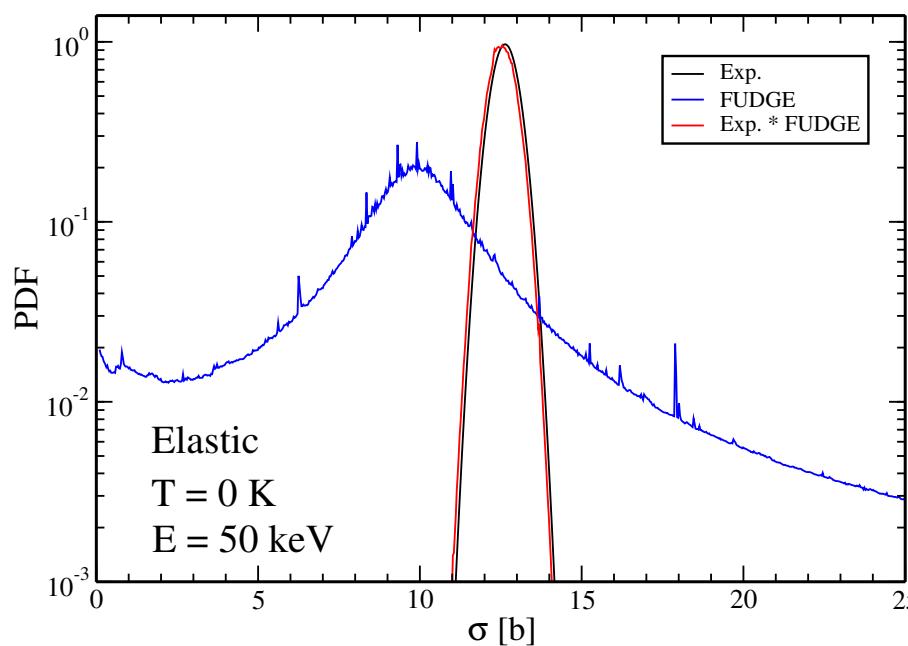
Evaluation of the neutron data standards,
Nuclear Data Sheets 148 (2018) 143-188



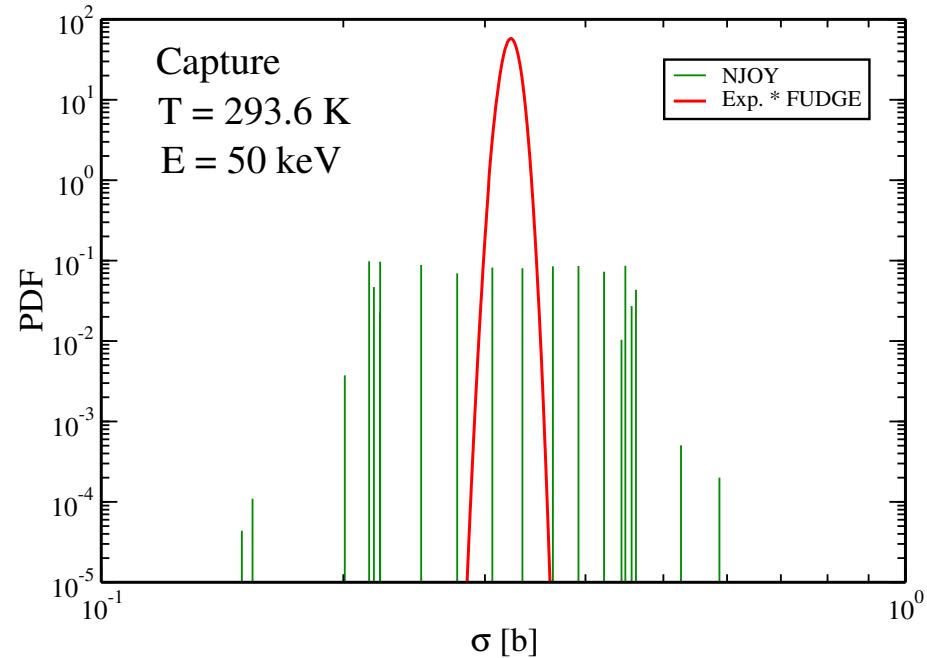
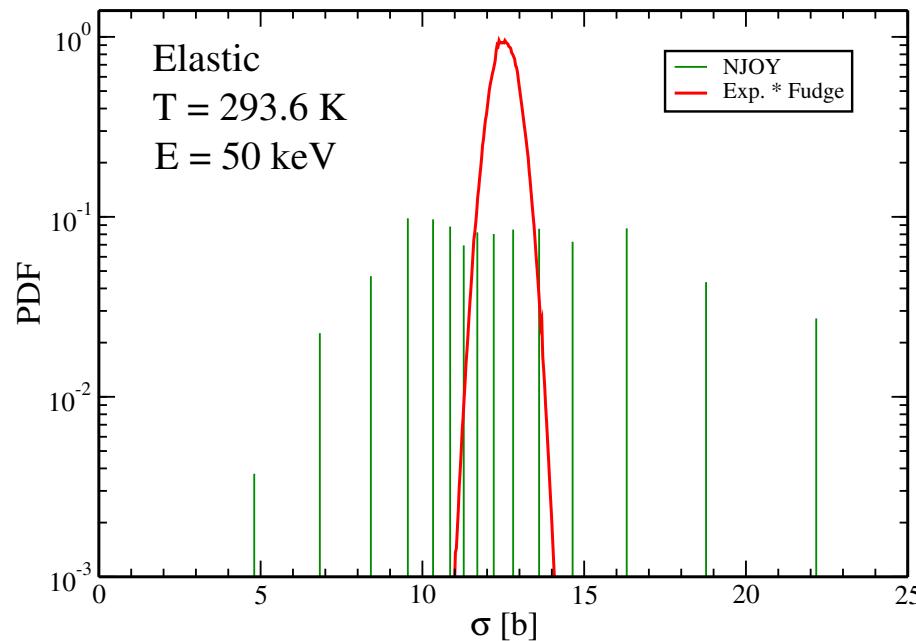
PDF at 0 K for n- ^{238}U

The experimental PDF is a strong constraint for the theoretical one

The combined PDF (normalized) is very similar to the experimental one



Heating the PDF for n- ^{238}U



The variance problem seems to be solved!
(we need to run some benchmarks!)

We can combine experiment with either PDF variant!

Summary & outlook

A Bayesian approach is more advisable for the PDF calculation in the URR

- Different probabilities are properly combined together
- Including experiment is straightforward
- Enables event-by-event treatment in URR
- Traditional UQ method and traditional self-shielding method should be interchangeable

Future work will be devoted to

- Speed up numerical calculations of PDFs at higher temperatures
- Fit the zero-temperature PDFs with a “universal” function

What if...

We didn't need to compute resonance ladders?

How can we fit the PDFs?

We are focusing on neutron-induced reactions

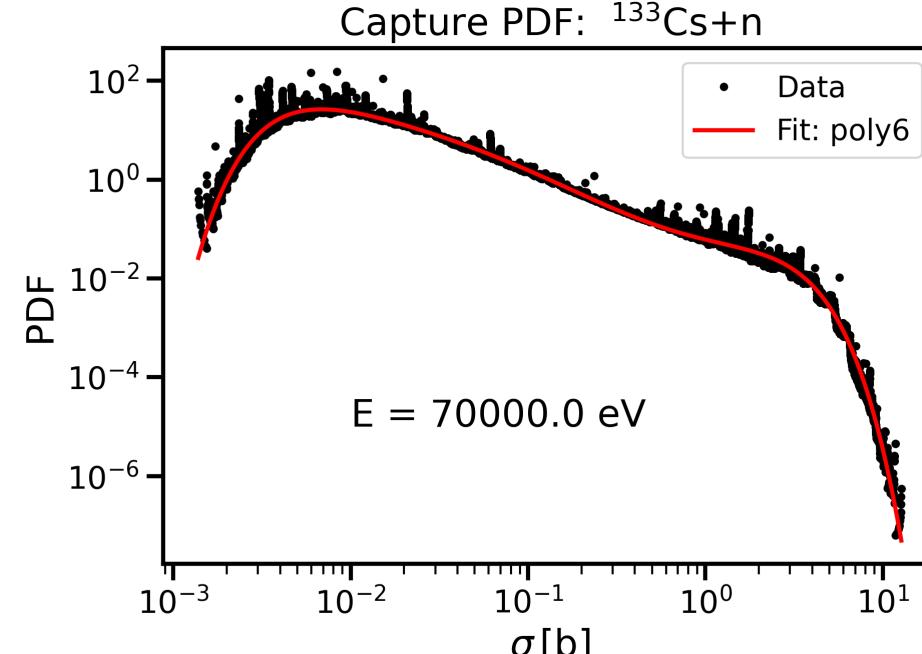
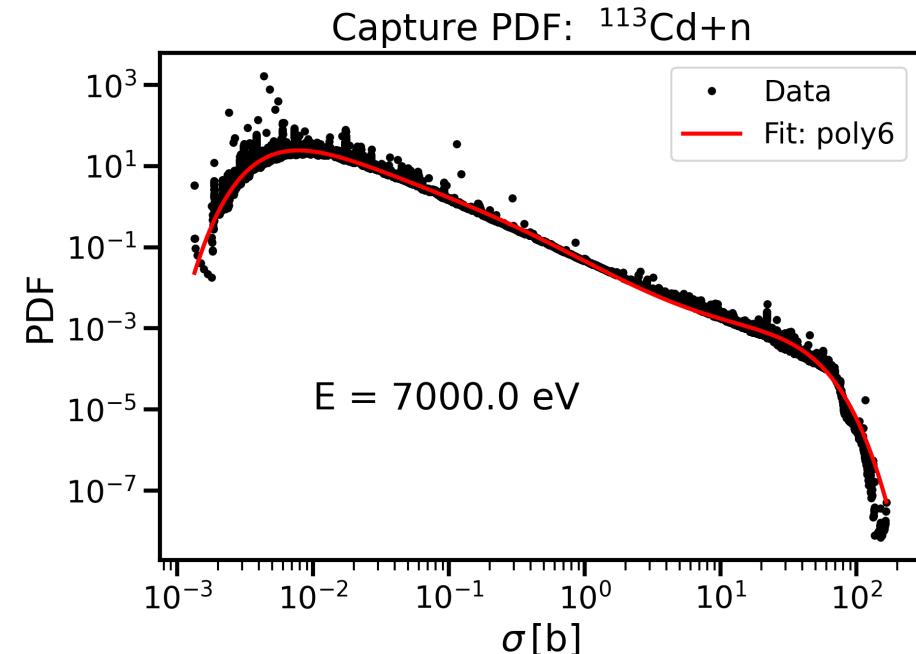
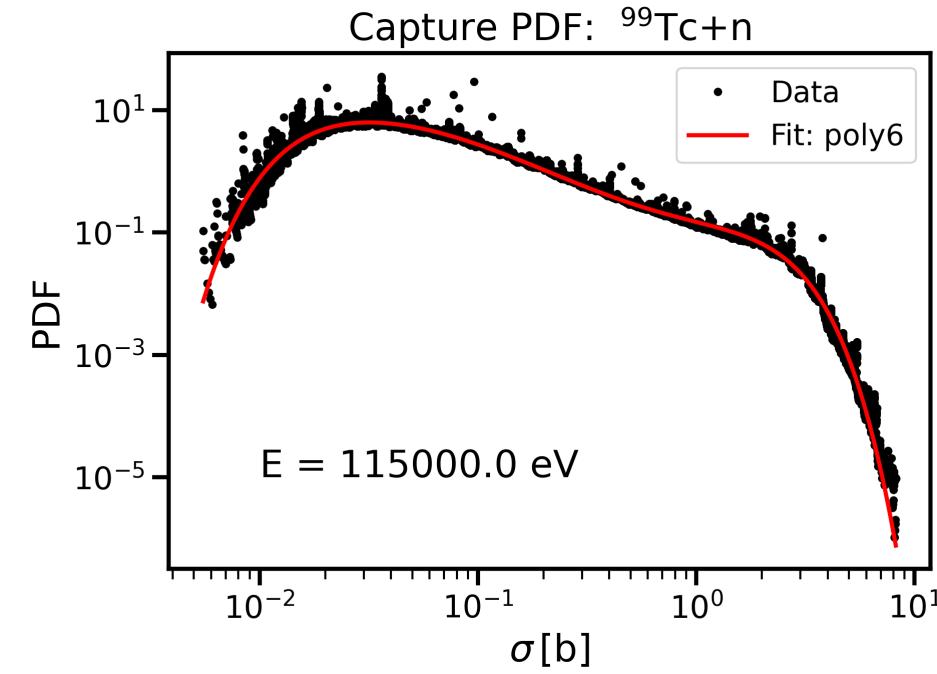
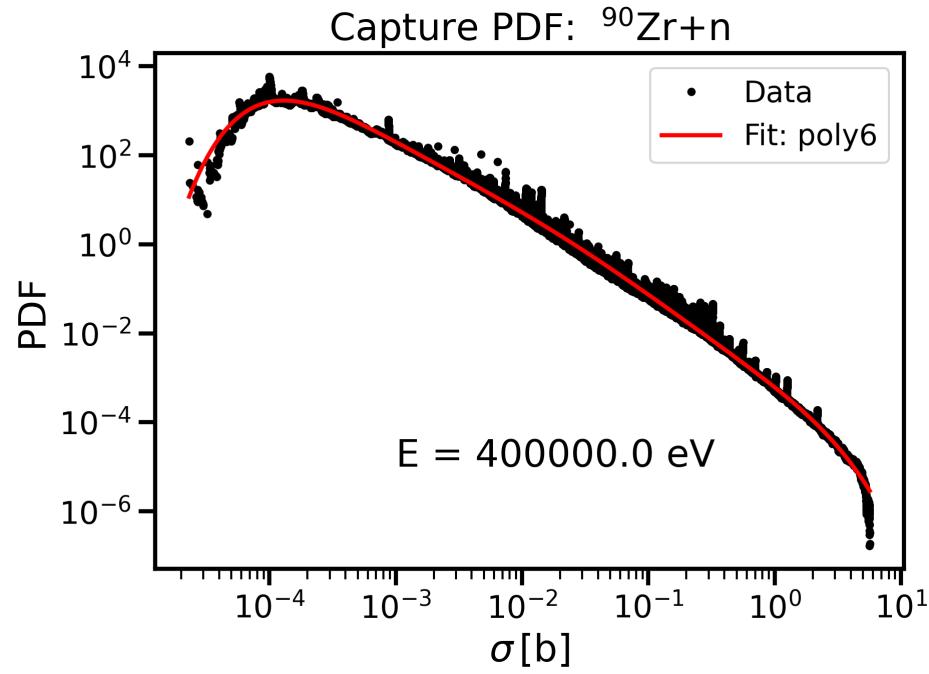
- Many target nuclei

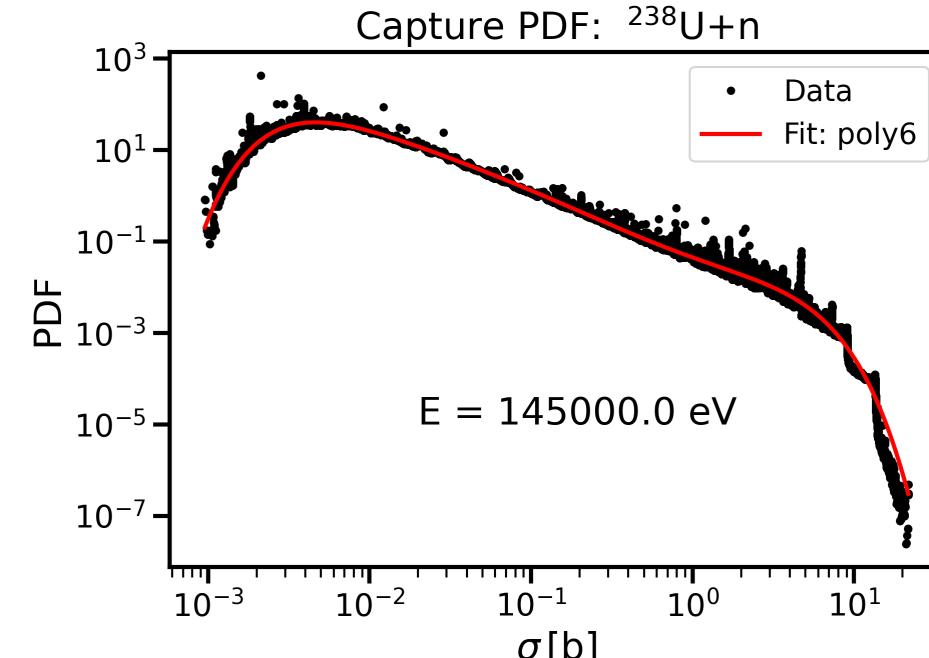
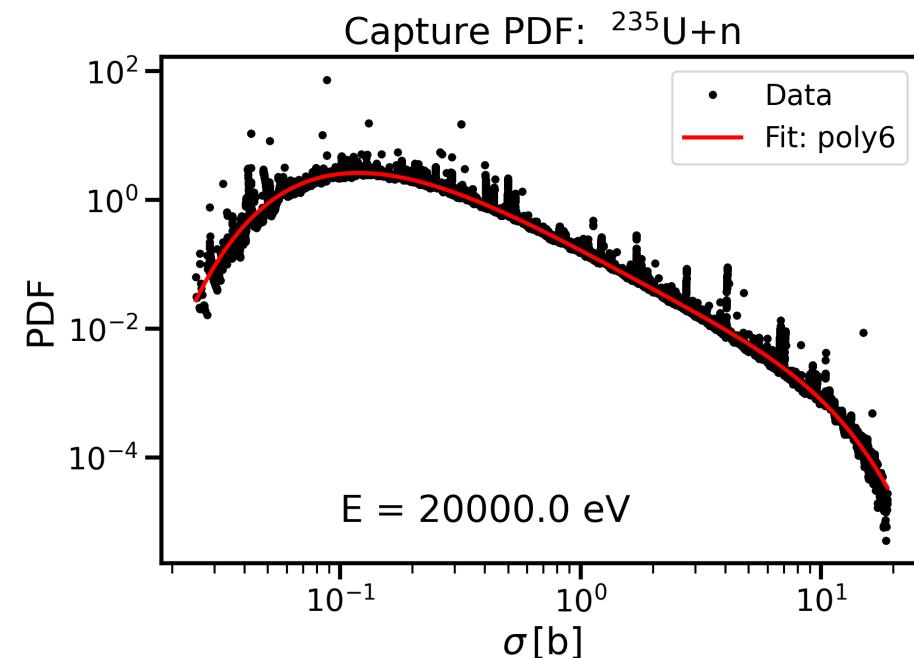
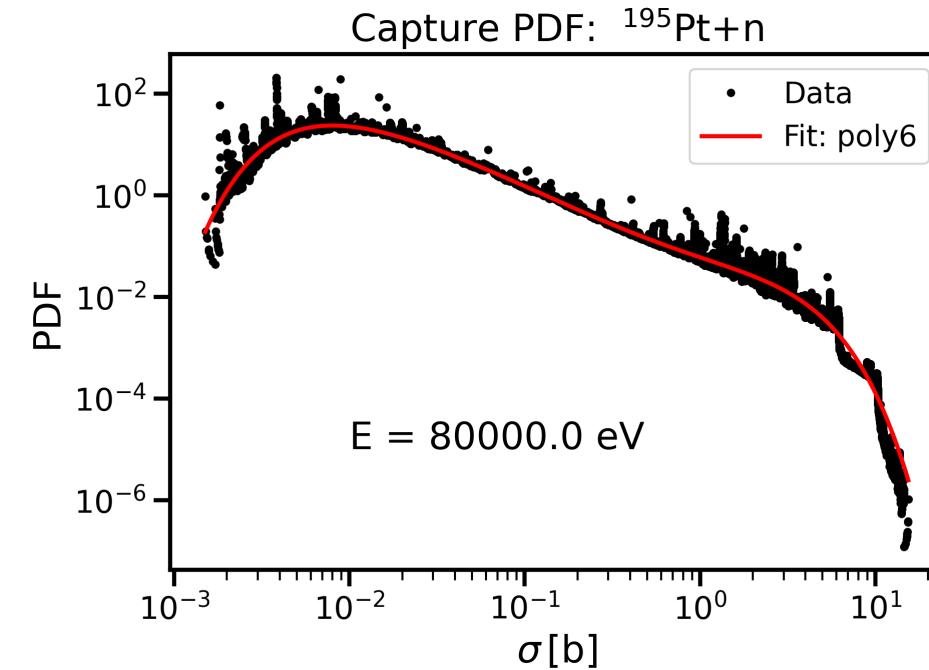
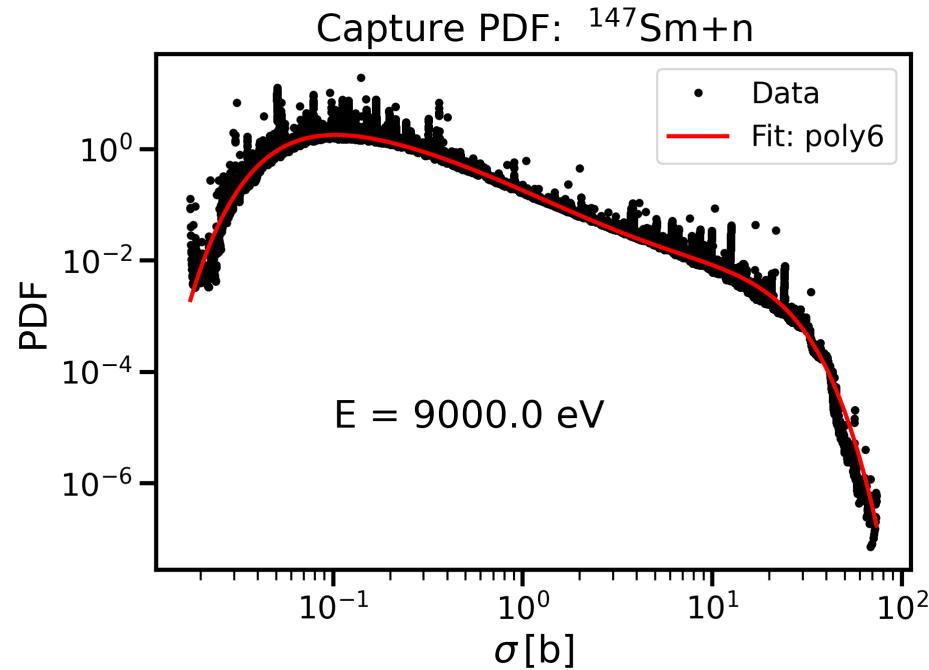
- Different reaction channels (elastic, capture, ...)

We are looking for a “universal” PDF for each reaction channel

We used standard fitting procedures and also symbolic regression techniques
to find a proper function

$$P(\sigma, E) = \exp \left\{ \sum_{n=0}^6 p_n(E) (\ln \sigma)^n \right\}$$





Next step

The idea seems to work, but for a full parametrization we need to fit the coefficients $p_n(E)$ as functions of energy
... or store them in a database!!!

