

RBC/UKQCD: status & near-term prospects for  
the muon  $g-2$  HVP contribution

Christoph Lehner  
(Uni Regensburg & Brookhaven National Laboratory)

December 14, 2021 – DWQ @ 25

## The RBC & UKQCD collaborations

### [UC Berkeley/LBNL](#)

Aaron Meyer

### [BNL and BNL/RBRC](#)

Yasumichi Aoki (KEK)

Peter Boyle (Edinburgh)

Taku Izubuchi

Yong-Chull Jang

Chulwoo Jung

Christopher Kelly

Meifeng Lin

Hiroshi Ohki

Shigemi Ohta (KEK)

Amarjit Soni

### [CERN](#)

Andreas Jüttner (Southampton)

### [Columbia University](#)

Norman Christ

Duo Guo

Yikai Huo

Yong-Chull Jang

Joseph Karpie

Bob Mawhinney

Ahmed Sheta

Bigeng Wang

Tianle Wang

Yidi Zhao

### [University of Connecticut](#)

Tom Blum

Luchang Jin (RBRC)

Michael Riberdy

Masaaki Tomii

### [Edinburgh University](#)

Matteo Di Carlo

Luigi Del Debbio

Felix Erben

Vera Gülpers

Tim Harris

Raoul Hodgson

Nelson Lachini

Michael Marshall

Fionn Ó hÓgáin

Antonin Portelli

James Richings

Azusa Yamaguchi

Andrew Z.N. Yong

### [KEK](#)

Julien Frison

### [University of Liverpool](#)

Nicolas Garron

### [Michigan State University](#)

Dan Hoying

### [Milano Bicocca](#)

Mattia Bruno

### [Peking University](#)

Xu Feng

### [University of Regensburg](#)

Davide Giusti

Christoph Lehner (BNL)

### [University of Siegen](#)

Matthew Black

Oliver Witzel

### [University of Southampton](#)

Nils Asmussen

Alessandro Barone

Jonathan Flynn

Ryan Hill

Rajnandini Mukherjee

Chris Sachrajda

### [University of Southern Denmark](#)

Tobias Tsang

### [Stony Brook University](#)

Jun-Sik Yoo

Sergey Syritsyn (RBRC)

## Calculation of the Hadronic Vacuum Polarization Contribution to the Muon Anomalous Magnetic Moment

T. Blum,<sup>1</sup> P. A. Boyle,<sup>2</sup> V. Gülpers,<sup>3</sup> T. Izubuchi,<sup>4,5</sup> L. Jin,<sup>1,5</sup> C. Jung,<sup>4</sup> A. Jüttner,<sup>3</sup> C. Lehner,<sup>4,\*</sup> A. Portelli,<sup>2</sup> and J. T. Tsang<sup>2</sup>

(RBC and UKQCD Collaborations)

<sup>1</sup>*Physics Department, University of Connecticut, Storrs, Connecticut 06269-3046, USA*

<sup>2</sup>*School of Physics and Astronomy, The University of Edinburgh, Edinburgh EH9 3FD, United Kingdom*

<sup>3</sup>*School of Physics and Astronomy, University of Southampton, Southampton SO17 1BJ, United Kingdom*

<sup>4</sup>*Physics Department, Brookhaven National Laboratory, Upton, New York 11973, USA*

<sup>5</sup>*RIKEN-BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973, USA*



(Received 25 January 2018; published 12 July 2018)

We present a first-principles lattice QCD + QED calculation at physical pion mass of the leading-order hadronic vacuum polarization contribution to the muon anomalous magnetic moment. The total contribution of up, down, strange, and charm quarks including QED and strong isospin breaking effects is  $a_{\mu}^{\text{HVP LO}} = 715.4(18.7) \times 10^{-10}$ . By supplementing lattice data for very short and long distances with  $R$ -ratio data, we significantly improve the precision to  $a_{\mu}^{\text{HVP LO}} = 692.5(2.7) \times 10^{-10}$ . This is the currently most precise determination of  $a_{\mu}^{\text{HVP LO}}$ .

Pure lattice result and dispersive result with reduced  $\pi\pi$  dependence (window method)

Aaron Meyer (BNL → LBNL) & Mattia Bruno (BNL → CERN → Milano) joined since this 2018 paper

## Lattice QCD – Time-Moment Representation

Starting from the vector current  $J_\mu(x) = i \sum_f Q_f \bar{\Psi}_f(x) \gamma_\mu \Psi_f(x)$  we may write

$$a_\mu^{\text{HVP LO}} = \sum_{t=0}^{\infty} w_t C(t)$$

with

$$C(t) = \frac{1}{3} \sum_{\vec{x}} \sum_{j=0,1,2} \langle J_j(\vec{x}, t) J_j(0) \rangle$$

and  $w_t$  capturing the photon and muon part of the HVP diagrams (Bernecker-Meyer 2011).

The correlator  $C(t)$  is computed in lattice QCD+QED at physical pion mass with non-degenerate up and down quark masses including up, down, strange, and charm quark contributions. The missing bottom quark contributions are computed in pQCD.

## Window method (introduced in RBC/UKQCD 2018)

We also consider a window method. Following Meyer-Bernecker 2011 and smearing over  $t$  to define the continuum limit we write

$$a_\mu = a_\mu^{\text{SD}} + a_\mu^{\text{W}} + a_\mu^{\text{LD}}$$

with

$$a_\mu^{\text{SD}} = \sum_t C(t) w_t [1 - \Theta(t, t_0, \Delta)],$$

$$a_\mu^{\text{W}} = \sum_t C(t) w_t [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)],$$

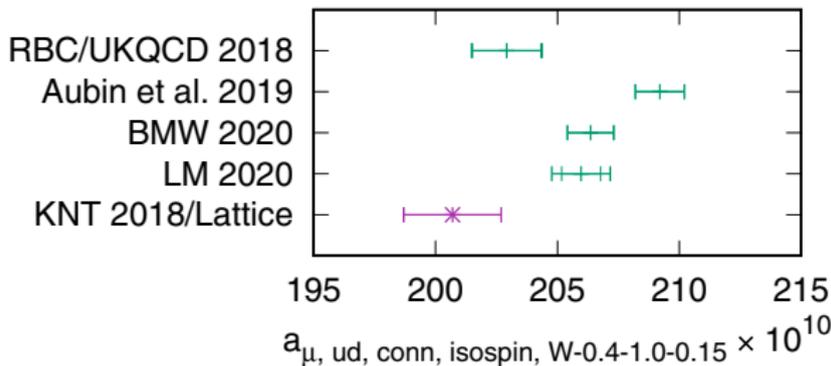
$$a_\mu^{\text{LD}} = \sum_t C(t) w_t \Theta(t, t_1, \Delta),$$

$$\Theta(t, t', \Delta) = [1 + \tanh [(t - t')/\Delta]] / 2.$$

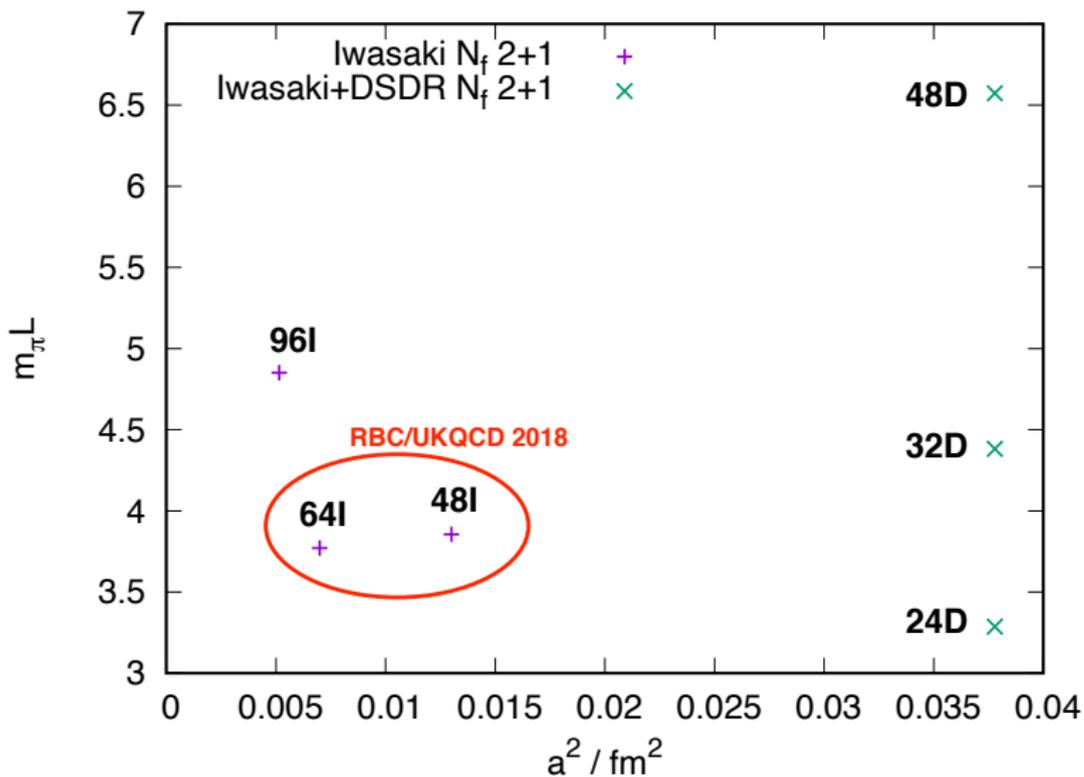
All contributions are well-defined individually and can be computed from lattice or R-ratio via  $C(t) = \frac{1}{12\pi^2} \int_0^\infty d(\sqrt{s}) R(s) s e^{-\sqrt{s}t}$  with  $R(s) = \frac{3s}{4\pi\alpha^2} \sigma(s, e^+ e^- \rightarrow \text{had})$ .

$a_\mu^{\text{W}}$  has small statistical and systematic errors on lattice!

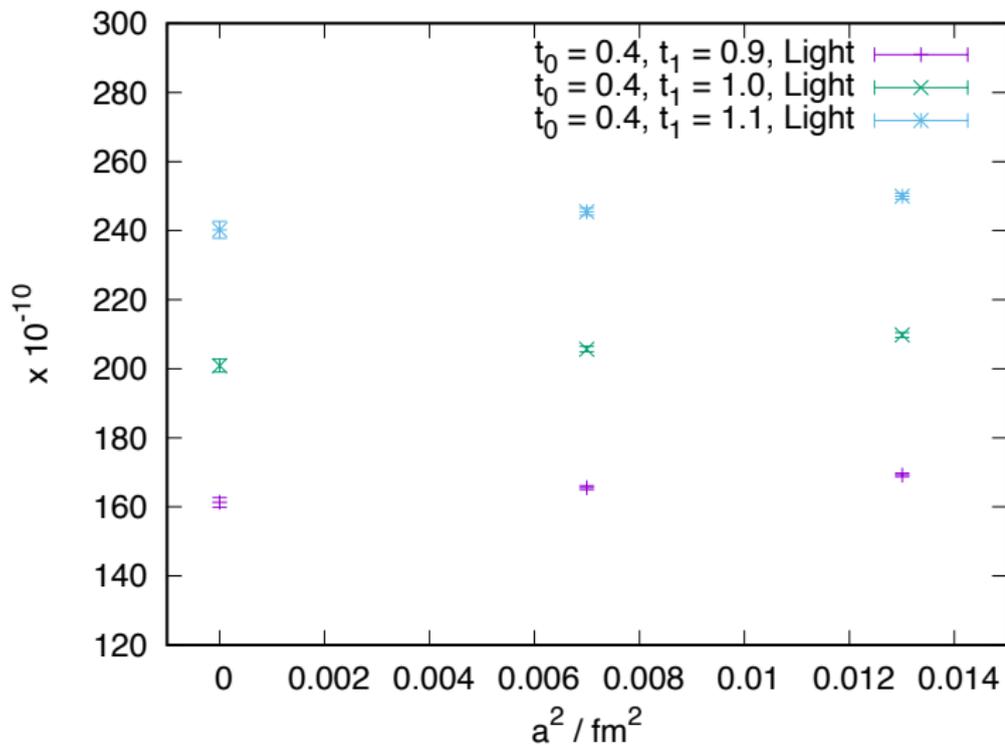
- ▶ In last few years, we reported on our progress for the complete result (improved bounding method,  $(\pi\pi)_{I=1}$  phase shift study, improvements for disconnected/QED/SIB diagrams), this talk focuses entirely on progress on the Euclidean time window (RBC/UKQCD 2018) in the isospin symmetric limit with  $t_0 = 0.4$  fm,  $t_1 = 1.0$  fm,  $\Delta = 0.15$  fm.
- ▶ This quantity promises reduced systematic lattice uncertainties, however, currently exhibits tensions between different lattice and R-ratio results:



Our extended list of ensembles, all with  $m_\pi = 135 \pm 5$  MeV:

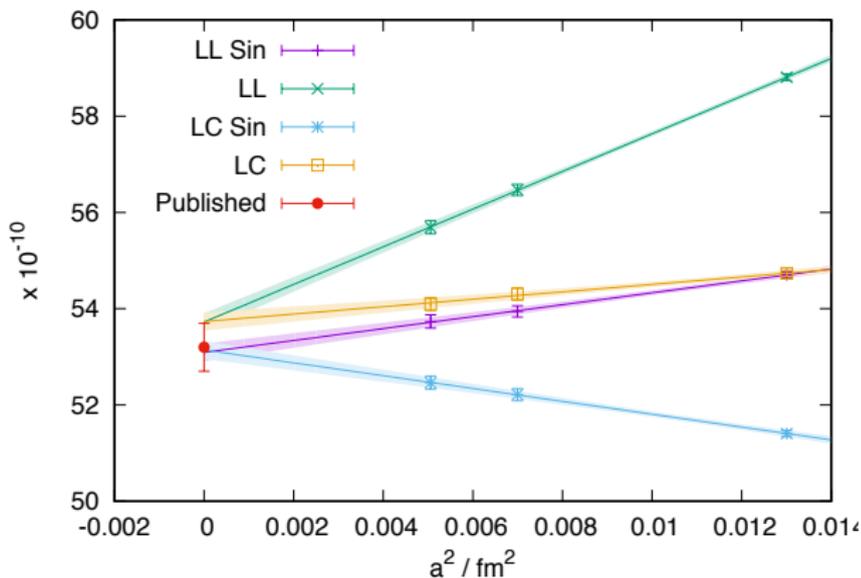


Previous continuum extrapolation mild (plot from [RBC/UKQCD 18](#)), will now check with even finer lattice spacing:



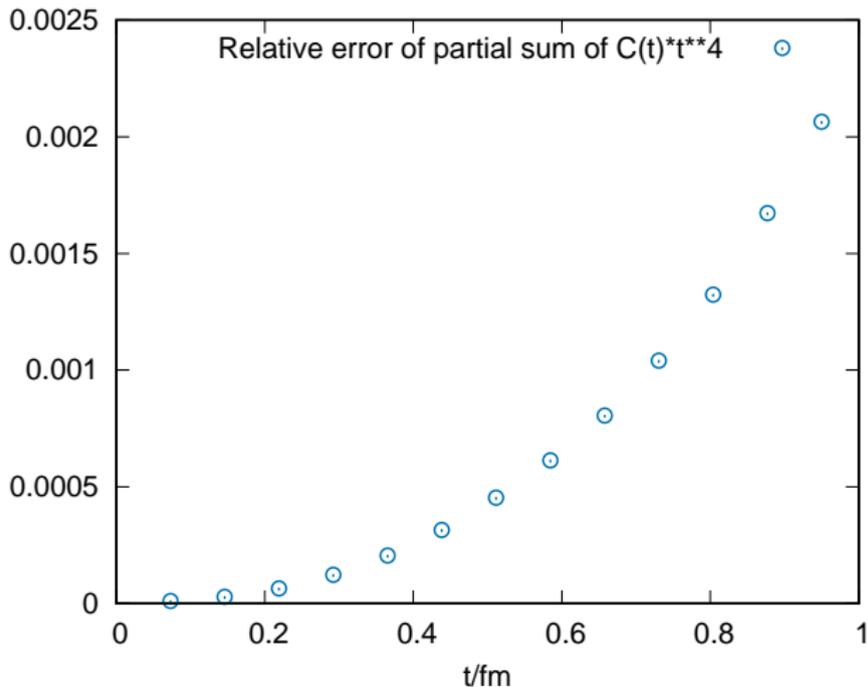
Previously already added  $a^{-1} = 2.77$  GeV ensemble for strange quark:

- ▶ Third lattice spacing for strange data ( $a^{-1} = 2.77$  GeV with  $m_\pi = 234$  MeV with sea light-quark mass corrected from global fit):



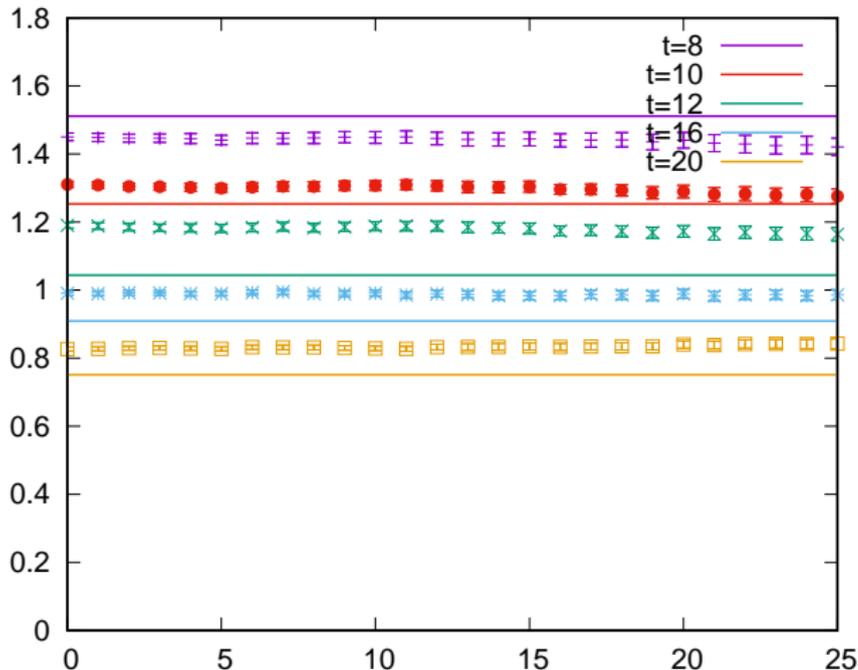
- ▶ For light quark use new 96l ensemble at physical pion mass. Data generated on Perlmutter and Summit in USA and Booster in Germany ( $a^{-1} = 2.77$  GeV with  $m_\pi = 139$  MeV)

We are still blinded, show only statistical relative error for new 96l:



Now full statistics (2x compared to what I showed at Lattice 21); 1 per mille statistical precision on finest lattice crucial for high-precision window tests

Master field error estimate versus simple Jackknife on 33 configurations on 96l:



y-axis: blinded noise, x-axis: distance cut for MF estimate

## Other improvements relevant for window:

- ▶ 4x statistics on 48l and 64l
- ▶ Explicit calculation of parametric derivatives at physical point
- ▶ Concluding study of missing charm determinant ( $2+1 \rightarrow 2+1+1$ )
- ▶ All of the above based on master-field-type study, see next slides

## Master-field calculation of gradients

For a local observable

$$O = \frac{1}{V} \sum_y O_y \quad (1)$$

we can define the truncated master-field covariance

$$\text{Cov}_R(O, A) \equiv \frac{1}{V} \sum_{x,y,|y|\leq R} (\langle O_x A_{x+y} \rangle_\beta - \langle O_x \rangle_\beta \langle A_{x+y} \rangle_\beta) \quad (2)$$

such that, e.g., the  $\beta$ -derivative of  $O$  is given by

$$\frac{\langle O \rangle_{\beta+\varepsilon} - \langle O \rangle_\beta}{\varepsilon} = 6 \lim_{R \rightarrow \infty} \text{Cov}_R(O, A). \quad (3)$$

In practice use exponential approach to plateau for  $R \rightarrow \infty$ .

We isolate the dependence on sea-quark mass  $m$  of an observable  $O$  by studying

$$\langle O \rangle_m \equiv \frac{\int \det(D(m)) O P}{\int \det(D(m)) P} \quad (4)$$

with Dirac matrix  $D(m)$  and residual weight  $P$ . Can show that

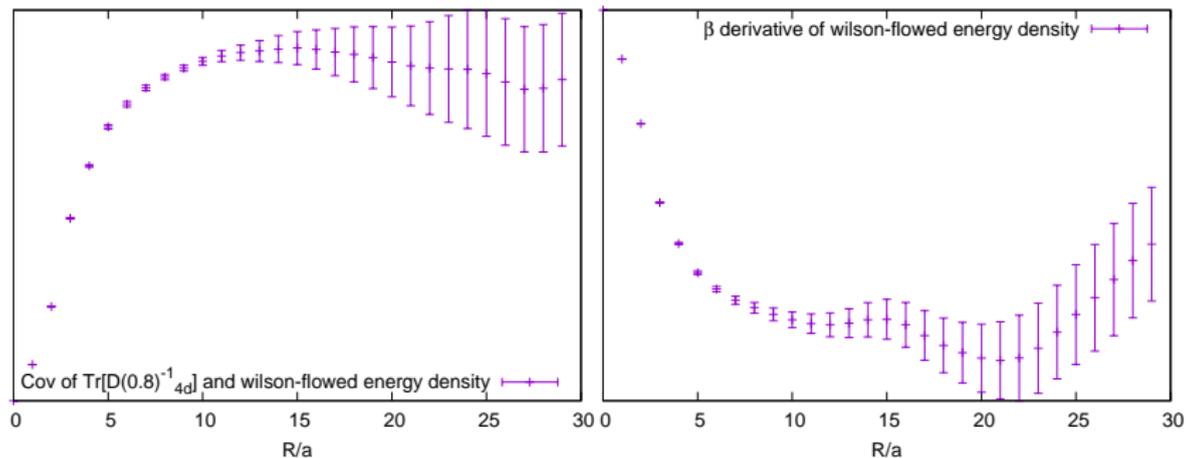
$$\frac{\langle O \rangle_{m+\varepsilon} - \langle O \rangle_m}{\varepsilon} = \text{Cov}(O, \text{Tr}[D_{4d}^{-1}(m)]) + \mathcal{O}(\varepsilon). \quad (5)$$

Finally, for DWF an additional flavor enters as

$$\det(D(m)D^{-1}(1)) \quad (6)$$

such that for  $m = 1$  the factor is trivial and we can view adding an additional flavor as changing the sea-quark mass down from  $m = 1$  to the target value.

## Example for wilson-flowed energy density (96l, $t_0 \approx 2$ )



Computed in similar way also derivatives of, e.g., VV and PP correlators.

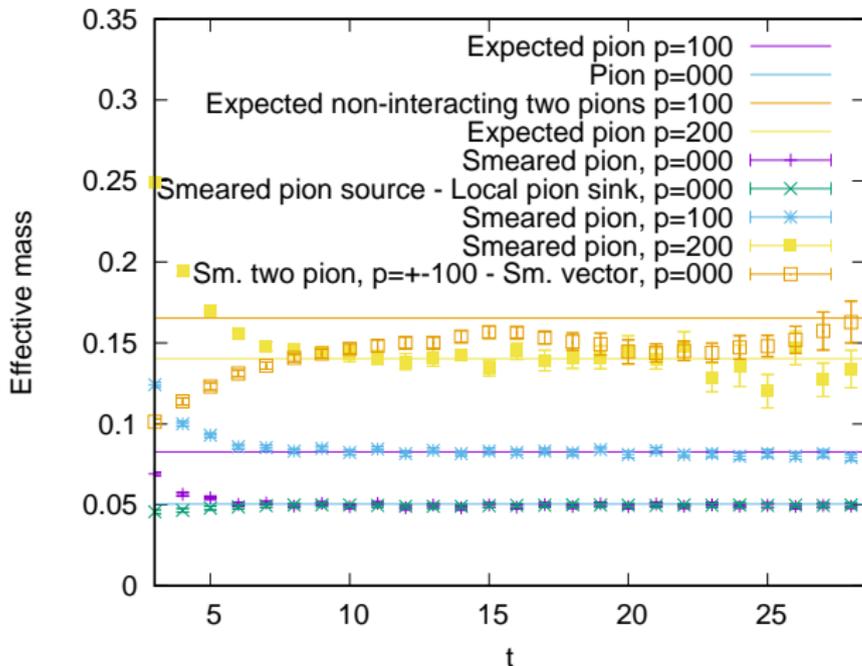
## Blinding

- ▶ 3 analysis groups for ensemble parameters (not blinded)
- ▶ 6 analysis groups for vector-vector correlators (blinded, to avoid bias towards other lattice/R-ratio results)
- ▶ Blinded vector correlator  $C_b(t)$  relates to true correlator  $C_0(t)$  by

$$C_b(t) = (b_0 + b_1 a^2 + b_2 a^4) C_0(t) \quad (7)$$

with appropriate random  $b_0, b_1, b_2$ , different for each analysis group. This prevents complete unblinding based on previously shared data on coarser ensembles.

Preview of distillation data on finest 96l ensemble for improved bounding method:



Early science time on Perlmutter was/is crucial! Above is master field study with one configuration and limited measurements on it.

## Conclusions and Outlook

- ▶ 4x statistics on 48l and 64l
- ▶ Third lattice spacing (96l) at  $a^{-1} \approx 2.7$  GeV
- ▶ 2+1  $\rightarrow$  2+1+1 and parametric gradients using master-field approach
- ▶ Still blinded, publish window results first
- ▶ For complete HVP analysis data set almost complete as well (still finishing distillation data on 96l, a lot of new data also on QCD+QED)