

Semileptonic Form Factors for $B \rightarrow \pi \ell \nu$ decays

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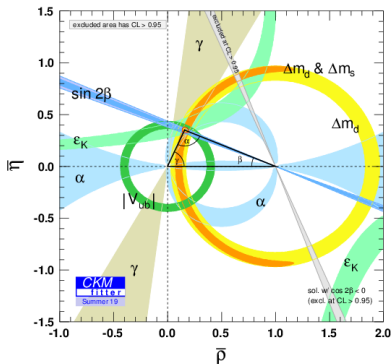
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Motivation

- Test unitarity of CKM matrix
- 2-3 σ discrepancy between exclusive ($B \rightarrow \pi \ell \nu$) and inclusive ($B \rightarrow X_u \ell \nu$)
- Lepton universality ratio predictions



CKMfitter Group (J. Charles et al.), Eur. Phys. J. C41,
1-131 (2005) [hep-ph/0406184], updated results and
plots available at: <http://ckmfitter.in2p3.fr>

- Differential $B \rightarrow \pi \ell \nu$ decay rate:

$$\underbrace{\frac{d\Gamma(B \rightarrow \pi \ell \nu)}{dq^2}}_{\text{Experiment}} = \underbrace{|V_{ub}|^2}_{\text{CKM}} \times \left(\kappa_1 \underbrace{|f_+(q^2)|^2}_{\text{non-pert.}} + \kappa_2 \underbrace{|f_0(q^2)|^2}_{\text{non-pert.}} \right)$$

$$q^\mu = p_B^\mu - p_\pi^\mu$$

κ — Known factors

- Requires a theoretical computation of the form factors:

$$\langle \pi(\vec{p}) | \mathcal{V}^\mu | B \rangle = 2f_+(q^2) \left(p_B^\mu - \frac{p_B \cdot q}{q^2} q^\mu \right) + f_0(q^2) \left(\frac{M_B^2 - M_\pi^2}{q^2} q^\mu \right)$$

$$\mathcal{V}^\mu = \bar{u} \gamma^\mu b$$

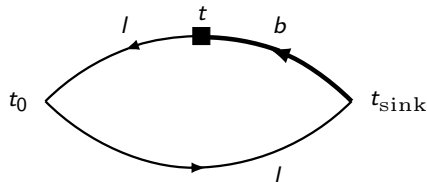
- RHQ Action for b quarks, Columbia interpretation

[Christ et al. PRD 76 (2007) 074505] [Lin and Christ PRD 76 (2007) 074506]

- Builds on original Fermilab action [El-Khadra et al. PRD 55 (1997) 3933]
- Related to Tsukuba interpretation [S. Aoki et al. PTP 109 (2003) 383]
- Clover action with anisotropic clover term
- Uses 3 parameters $(m_0 a, c_p, \zeta)$ that can be non-pertubatively tuned to remove $\mathcal{O}((m_0 a)^n)$, $\mathcal{O}(\vec{p}a)$, $\mathcal{O}((\vec{p}a)(m_0 a)^n)$ errors [PRD 86 (2012) 116003]
- Shamir DWF for l, s
- Use improvement terms to get one-loop improved **vector current**

	$L^3 \times T / a^4$	a^{-1} / GeV	m_π / MeV
C1	$24^3 \times 64$	1.78	340
C2	$24^3 \times 64$	1.78	430
M1	$32^3 \times 64$	2.38	300
M2	$32^3 \times 64$	2.38	360
M3	$32^3 \times 64$	2.38	410
F1S	$48^3 \times 96$	2.77	270
C0*	$48^3 \times 96$	1.73	139

- 2+1f ensembles: Degenerate light quark
- Sea quarks: Shamir domain-wall fermions
- **F1S** ensemble: new for this analysis
- *Physical pion ensemble C0 planned for future inclusion.



- Quark propagator $(t_0 \xrightarrow{l} t)$
- Sequential propagator $(t_0 \xrightarrow{l} t_{\text{sink}} \xrightarrow{b} t)$
- Contract propagators at t
- Smeared point sources
- Project B meson to rest frame; momentum given to final state

- Start by finding the parallel and perpendicular form factors $f_{\parallel}$ and $f_{\perp}$
- Simpler to relate to lattice data in the rest frame of the B -meson:

$$f_{\parallel} = \frac{\langle \pi | \mathcal{V}^0 | B \rangle}{\sqrt{2M_B}} \qquad f_{\perp} p^i = \frac{\langle \pi | \mathcal{V}^i | B \rangle}{\sqrt{2M_B}}$$

- Neatly separates into spatial and temporal components
- Can form linear combinations to find f_0 and f_+ :

$$f_0(q^2) = \frac{\sqrt{2M_B}}{M_B^2 + E_{\pi}^2} \left[(M_B - E_{\pi}) f_{\parallel}(q^2) + (E_{\pi}^2 - M_{\pi}^2) f_{\perp}(q^2) \right]$$

$$f_+(q^2) = \frac{1}{\sqrt{2M_B}} \left[f_{\parallel}(q^2) + (M_B - E_{\pi}) f_{\perp}(q^2) \right]$$

Matrix Element Ratio

- Make plateau fits to the ratio of 3pt and 2pt functions:

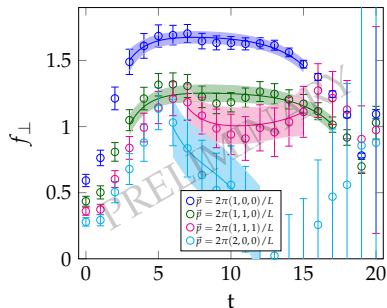
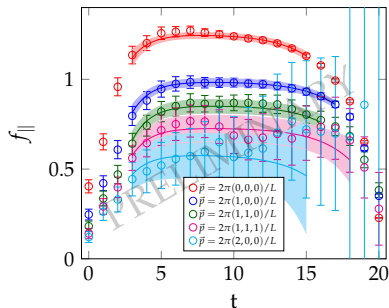
$$\langle \pi^{(0)} | V^\mu | B^{(0)} \rangle = \lim_{0 \ll t \ll t_{\text{sink}}} \frac{C_3^\mu(t)}{\sqrt{C_2^\pi(t) C_2^B(t_{\text{sink}} - t)}} \sqrt{\frac{4E_\pi^{(0)} E_B^{(0)}}{e^{-tE_\pi^{(0)}} e^{-(t_{\text{sink}} - t)E_B^{(0)}}}}$$

- Alternatively, include excited states fitting additional terms

$$\langle \pi^{(1)} | V^\mu | B^{(0)} \rangle \propto \sqrt{e^{-t(E_\pi^{(1)} - E_\pi^{(0)})}}$$

$$\langle \pi^{(0)} | V^\mu | B^{(1)} \rangle \propto \sqrt{e^{-(t_{\text{sink}} - t)(E_B^{(1)} - E_B^{(0)})}}$$

Excited state fits on the C1 ensemble:



$B \rightarrow \pi$ Chiral Continuum Fits

- Extrapolate to physical pion mass and zero lattice spacing simultaneously
- Use NLO hard-pion SU(2) HM χ PT [PRD 67 (2003) 054010]

$$f(M_\pi, E_\pi, a) = g(E_\pi, \Delta) \left(c_0 \left(1 + \frac{\delta f(M_\pi)}{(4\pi f_\pi)^2} \right) + c_1 \frac{M_\pi^2}{\Lambda^2} + c_2 \frac{E_\pi}{\Lambda} + c_3 \left(\frac{E_\pi}{\Lambda} \right)^2 + c_4 (a\Lambda)^2 \right)$$

where

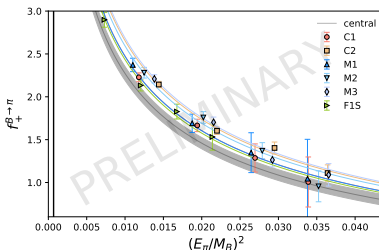
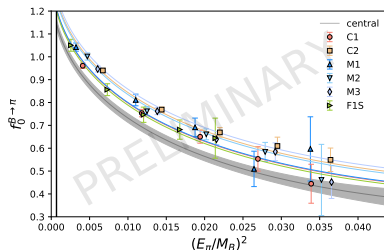
$$\Lambda = 1 \text{ GeV},$$

$$\delta f(M_\pi) = -\frac{3}{4}(3g_b^2 + 1) \left(M_\pi^2 \log \left(\frac{M_\pi^2}{\Lambda^2} \right) + \text{FVE} \right)$$

$$g_+(E_\pi, \Delta_+) = \frac{\Lambda}{E_\pi + \Delta_+}; \Delta_+ = M_{B^*} - M_B \approx 45.2 \text{ MeV [Pole term]}$$

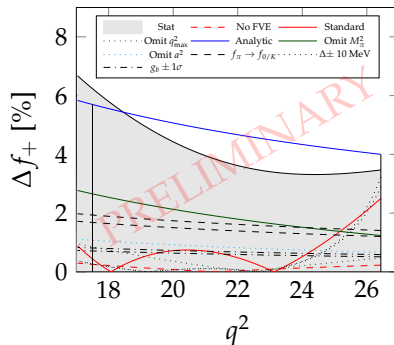
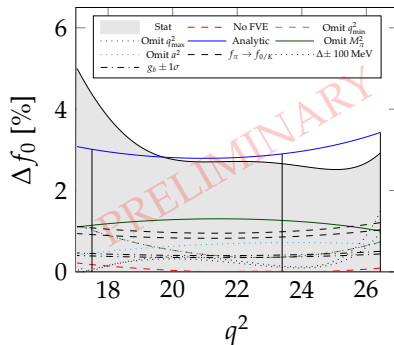
$$g_0(E_\pi, \Delta_0) = \frac{\Lambda}{E_\pi + \Delta_0}; \Delta_+ = M_{B^*(0^+)} - M_B \approx 305 \text{ MeV [Pole term]}$$

$B \rightarrow \pi$ Chiral Continuum Fits



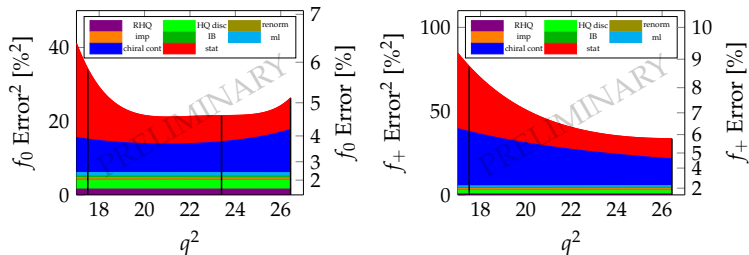
- Continuum form factor given by $f(M_\pi^{\text{phys}}, E_\pi, a = 0)$
- Need to assess the systematic error on the continuum results
- Include variations on the continuum fit ansatz as a systematic

Fit Variation Systematic



- Assess percentage deviations of possible alternatives the preferred fit
- Take largest deviation at reference q^2 values, which are then used in synthetic data points for the q^2 -extrapolation

Total Error Budgets



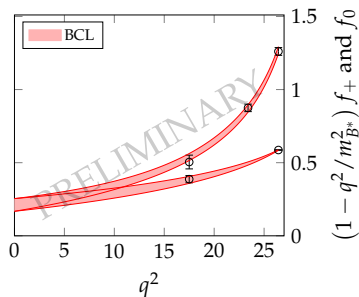
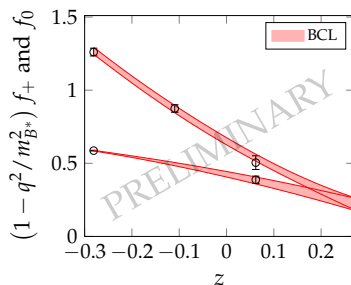
- Substantially reduced statistical errors over prior 2015 analysis
- Combine systematics at reference q^2 values with statistics and pass on to the q^2 extrapolation

z-expansions

- z-expansions given the kinematic constraint that $f_0 = f_+$ at $q^2 = 0$
- Can be compared with experimental data to determine $|V_{ub}|$, lepton-universality ratios
- z-expansions shown are BCL fits, also look at BGL

[PRD 79 (2009) 013008]

[PRL 74 (1995) 4603]



- Test lepton flavour universality with the R ratio

$$R(\pi) = \frac{\int_{q^2=m_\tau^2}^{q^2=q_{\max}^2} dq^2 \frac{d\Gamma(B \rightarrow \pi \tau \bar{\nu}_\tau)}{dq^2}}{\int_{q^2=m_\mu^2}^{q^2=q_{\max}^2} dq^2 \frac{d\Gamma(B \rightarrow \pi \mu \bar{\nu}_\mu)}{dq^2}}$$

- Also consider the alternative R -ratio proposed by Isidori and Sumensari [Isidori & Sumensari, Eur. Phys. J. C 80 (2020)]

$$R^{\text{new}}(\pi) = \frac{\int_{q^2=q_{\min}^2}^{q^2=q_{\max}^2} dq^2 \frac{d\Gamma(B \rightarrow \pi \tau \bar{\nu}_\tau)}{dq^2}}{\int_{q^2=q_{\min}^2}^{q^2=q_{\max}^2} \frac{\omega_\tau(q^2)}{\omega_\mu(q^2)} dq^2 \frac{d\Gamma(B \rightarrow \pi \mu \bar{\nu}_\mu)}{dq^2}}$$

- Common integration range

$$q_{\min}^2 \geq m_\tau^2$$

[Freytsis *et al.* Phys. Rev. D 92, 054018 (2015)]

[Bernlochner & Ligeti, Phys. Rev. D 95, 014022 (2017)]

[Soni, PoS ICHEP2020, 436 (2021)]

$$\frac{d\Gamma(B \rightarrow \pi \ell \bar{\nu})}{dq^2} = \phi(q^2) \omega_\ell(q^2) \left[F_V^2(q^2) + (F_S^\ell(q^2))^2 \right]$$

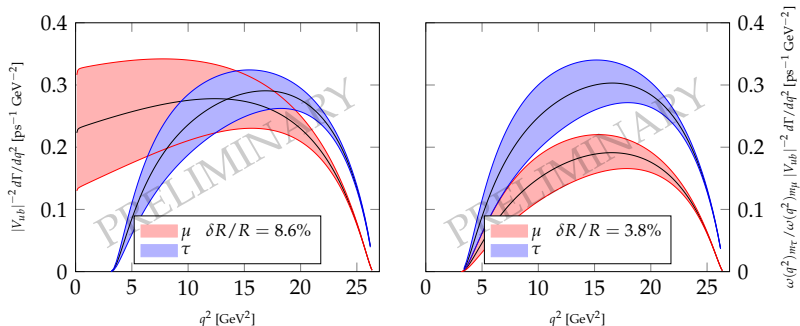
$$\phi(q^2) = \eta_{\text{EW}} \frac{G_F^2 |V_{ub}|^2}{24\pi^3} |\vec{k}|$$

$$\omega_\ell(q^2) = \left(1 - \frac{m_\ell^2}{q^2} \right)^2 \left(1 + \frac{m_\ell^2}{2q^2} \right)$$

$$F_V^2 = \vec{k}^2 |f_+(q^2)|^2$$

- Equal weights $\omega_\ell(q^2)$ in the integrand

$$(F_S^\ell)^2 = \frac{3}{4} \frac{m_\ell^2}{m_\ell^2 + 2q^2} \frac{(m_B^2 - m_\pi^2)^2}{m_\pi^2} |f_0(q^2)|^2$$



- “Standard” (left) and “New” (right) integrands for the R -ratio
- New ratio definition exhibits a reduction in relative uncertainty

- Updates to RBC-UKQCD 2015 results in the pipeline
 - More precisely determined lattice spacing and inclusion of third lattice spacing via F1S ensemble reduces errors
 - More details available in conference proceedings:
[arXiv:1912.09946](#), [arXiv:2012.04323](#)
 - Next step: Internal cross-checks
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Related talks:

- $B_s \rightarrow K\ell\nu, B_s \rightarrow D_s\ell\nu$: Jonathan Flynn @ 11:00 EST today