

Rare kaon $K \rightarrow \pi \ell^+ \ell^-$ decay amplitude

DWQ@25

Felix Erben

17 December 2021



THE UNIVERSITY
of EDINBURGH



European Research Council
Established by the European Commission

The RBC & UKQCD collaborations

[UC Berkeley/LBNL](#)

Aaron Meyer

[BNL and BNL/RBRC](#)

Yasumichi Aoki (KEK)
Peter Boyle (Edinburgh)
Taku Izubuchi
Yong-Chull Jang
Chulwoo Jung
Christopher Kelly
Meifeng Lin
Hiroshi Ohki
Shigemi Ohta (KEK)
Amarjit Soni

[CERN](#)

Andreas Jüttner (Southampton)

[Columbia University](#)

Norman Christ
Duo Guo
Yikai Huo
Yong-Chull Jang
Joseph Karpie
Bob Mawhinney
Ahmed Sheta
Bigeng Wang
Tianle Wang
Yidi Zhao

[University of Connecticut](#)

Tom Blum
Luchang Jin (RBRC)
Michael Riberdy
Masaaki Tomii

[Edinburgh University](#)

Matteo Di Carlo
Luigi Del Debbio
Felix Erben
Vera Gülpers
Tim Harris
Raoul Hodgson
Nelson Lachini
Michael Marshall
Fionn Ó hÓgáin
Antonin Portelli
James Richings
Azusa Yamaguchi
Andrew Z.N. Yong

[KEK](#)

Julien Frison

[University of Liverpool](#)

Nicolas Garron

[Michigan State University](#)

Dan Hoying

[Milano Bicocca](#)

Mattia Bruno

[Peking University](#)

Xu Feng

[University of Regensburg](#)

Davide Giusti
Christoph Lehner (BNL)

[University of Siegen](#)

Matthew Black
Oliver Witzel

[University of Southampton](#)

Nils Asmussen
Alessandro Barone
Jonathan Flynn
Ryan Hill
Rajnandini Mukherjee
Chris Sachrajda

[University of Southern Denmark](#)

Tobias Tsang

[Stony Brook University](#)

Jun-Sik Yoo
Sergey Syritsyn (RBRC)

$$K \rightarrow \pi \ell^+ \ell^-$$

- flavour-changing neutral current (FCNC)
- $s \rightarrow d$ transition, forbidden at tree level
- potential for new physics
- dominated by $K \rightarrow \pi \gamma^*$ virtual photon exchange

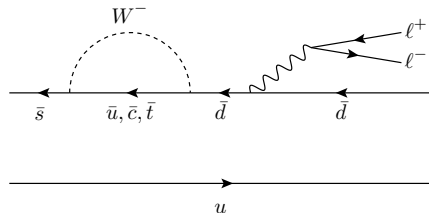
\Rightarrow dominated by long-distance effects

- measured at NA48, NA62 experiments
- NA62 measured 28,000 samples $K \rightarrow \pi \mu^+ \mu^-$ in 2016-2018 Run1

[Bician, PoS(ICHEP2020), 390(364)]

- more $K \rightarrow \pi \mu^+ \mu^-$ and $K \rightarrow \pi e^+ e^-$ in 2021-2024 Run2

[Lazeroni, CERN SPS, Mar. 2021]



Related talk by Raoul Hodgson, Thu 16 Dec, 12:00: Weak Hyperon Decays from Lattice QCD

- theoretical description of form factor via two-loop low-energy expansion
[D'Ambrosio et al., JHEP 2019 10.1007/jhep02(2019)049], [D'Ambrosio et al., Physics Letters B 797, 134891 (2019)]
- lattice methodology for long-distance contribution first proposed in
[Isidori et al., Physics Letters B 633, 75 (2006)]
- additional details of strategy, including control over UV divergences
[Christ et al., Physical Review D 92, (2015)]
- exploratory work at $M_\pi = 430$ MeV
[Christ et al., Phys. Rev. D 94, 114516 (2016)]
- this work repeats the exploratory study at M_π^{phys}

$$K \rightarrow \pi \gamma^* \rightarrow \pi \ell^+ \ell^-$$

Long-distance part can be parametrised via form factor V_+ : [D'Ambrosio et al., hep-ph/9808289]

$$A_\mu = \frac{-iG_F}{(4\pi)^2} V_+(z) \left(q^2 (k+p)_\mu - (M_K^2 - M_\pi^2) q_\mu \right)$$

$$z = q^2/M_K^2, \quad q = k - p, \quad V_+ = a_+ + b_+ z + V^{\pi\pi}(z)$$

Predictions from fitting NA48/NA62 data:

$$K \rightarrow \pi e^+ e^- : \quad a_+ = -0.578(16), \quad b_+ = -0.779(66)$$

$$K \rightarrow \pi \mu^+ \mu^- : \quad a_+ = -0.592(15), \quad b_+ = -0.699(58)$$

theory, matching long-distance to short-distance from perturbation theory: [D'Ambrosio et al.,

1906.03046]

$$K \rightarrow \pi \ell^+ \ell^- : \quad a_+ = -1.59(8), \quad b_+ = -0.82(6)$$

Long-distance amplitude:

$$A_\mu(q^2) = \int d^4x \langle \pi(\vec{p}) | T[J_\mu(0) \mathcal{H}_W(x)] | K(\vec{k}) \rangle$$

J_μ : electromagnetic current, \mathcal{H}_W : $s \rightarrow d$ weak current

$$\mathcal{H}_W = \frac{G_F}{\sqrt{2}} V_{us}^* V_{ud} [C_1(Q_1^u - Q_1^c) + C_2(Q_2^u - Q_2^c)]$$

$$Q_1^q = (\bar{s}_i \gamma_\mu^L d_i)(\bar{q}_j \gamma_\mu^L q_j), \quad Q_2^q = (\bar{s}_i \gamma_\mu^L q_i)(\bar{q}_j \gamma_\mu^L d_j), \quad q \in \{u, c, t\}$$

LATTICE STRATEGY

on the lattice, four-point correlation function

$$\Gamma^{(4)}(t_J, t_H) = \langle \phi_\pi(t_\pi) T[J_\mu(t_J) H_W(t_H)] \phi_K^\dagger(t_K) \rangle$$

integrate t_H around fixed t_J :

$$I(T_a, T_b) = \int_{t_J - T_a}^{t_J + T_b} dt_H \tilde{\Gamma}^{(4)}(t_J, t_H)$$

with the reduced correlator

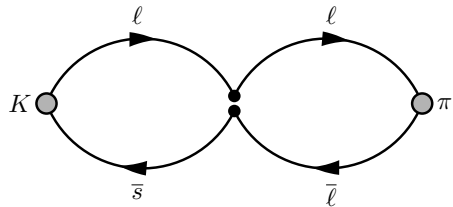
$$\tilde{\Gamma}^{(4)}(t_J, t_H) = \Gamma^{(4)}(t_J, t_H) / Z_{\pi K}$$

$Z_{\pi K}$ from 2pt correlators. Amplitude via:

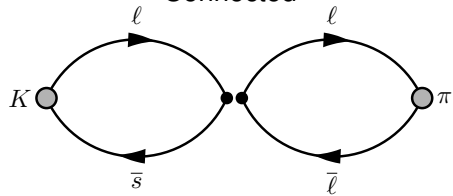
$$A_\mu(q^2) = \lim_{T_a, T_b \rightarrow \infty} I(T_a, T_b)$$

LATTICE STRATEGY

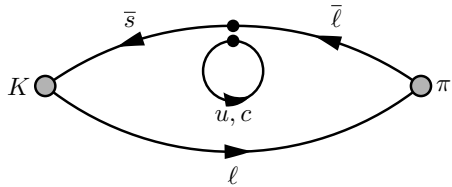
Types of diagrams, H_W 3pt function:



Connected

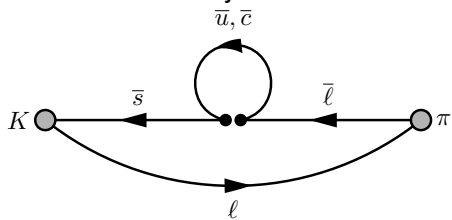


Wing



Eye

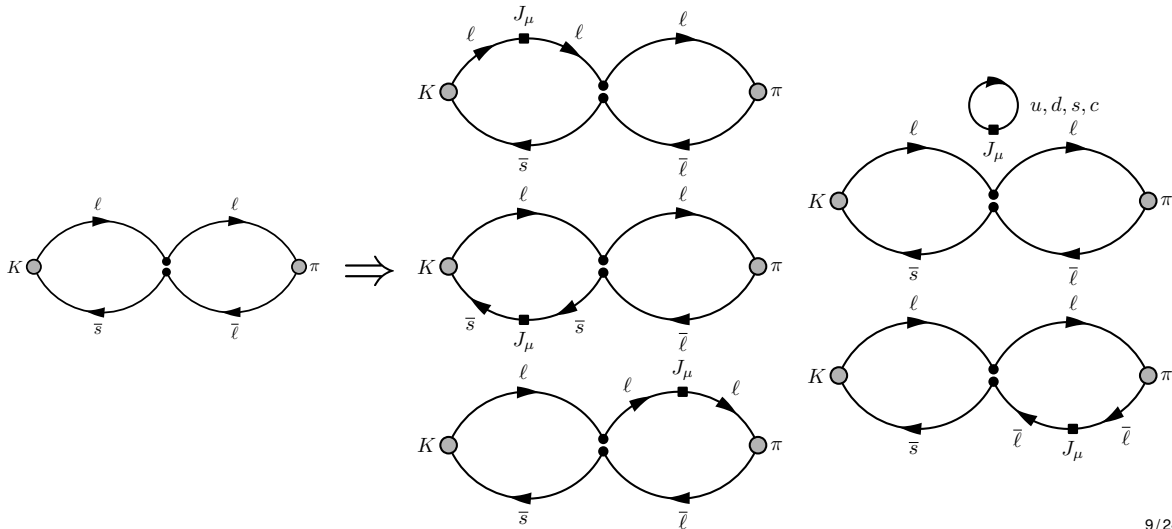
\bar{u}, \bar{c}



Saucer

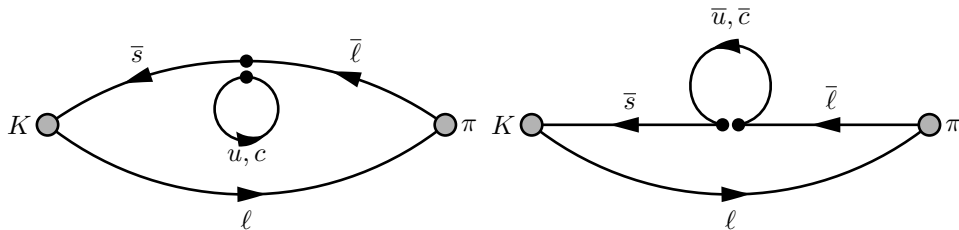
LATTICE STRATEGY

4pt function including J_μ , for the C diagram



LATTICE STRATEGY

- J_μ in loop of S, E diagrams
⇒ quadratic divergences as J_μ, H_W approach each other
- conserved current guarantees a logarithmic at most
- remaining divergence dealt with by GIM subtraction [Glashow et al. Phys. Rev. D 2, 1285 (1970)]
⇒ needs charm-quark loops
⇒ difference leads to cancellation of divergence



RBC/UKQCD $N_f = 2 + 1$ Möbius DWF ensemble [Blum et al. 1411.7017]

- $48^3 \times 96$
- $a^{-1} = 1.73 \text{ GeV}$
- $m_\pi L = 3.86$
- $m_\pi = 139 \text{ MeV}$
- $m_K = 499 \text{ MeV}$

measurements:

- 87 configurations, separated by 20 MC steps
- 6 time translations per configuration
- on this ensemble:

- $\vec{p}_K = (0, 0, 0), \vec{p}_\pi = \frac{2\pi}{L}(1, 0, 0)$

$$\Rightarrow z = \frac{q^2}{M_K^2} = 0.013(2)$$

SETUP OF COMPUTATION

light quarks:

- Coulomb-gauge fixed wall sources
- zMöbius action
 - approximation of Möbius using complex parameters
 - 5th dimension extent reduced from $L_s = 24$ to $L_s = 10$
 - all-mode-averaging style bias correction needed for Möbius action

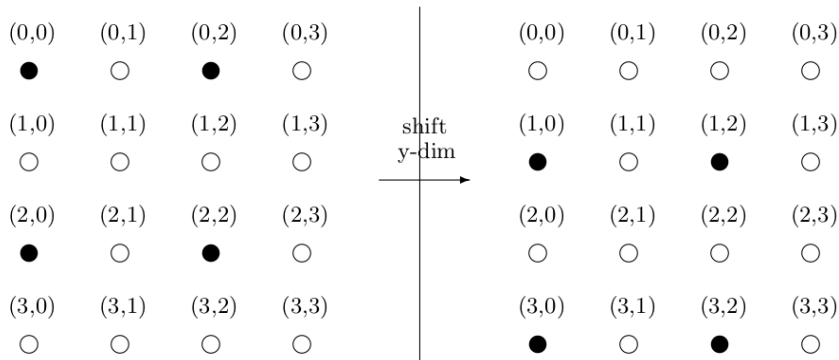
charm quarks:

- GIM demands same action as for l quark to work
- zMöbius action parameter choice makes simulation at m_c^{phys} impossible
 - simulation at unphysical $\alpha m_c \in \{0.25, 0.30, 0.35\}$

loops:

- spin-colour diluted sparse sources
- all-mode averaging
 - 1 hit exact ($10^{-8}, 10^{-10}, 10^{-12}, 10^{-14}$ for l, c_1, c_2, c_3)
 - 10 hits inexact (10^{-4} for all)

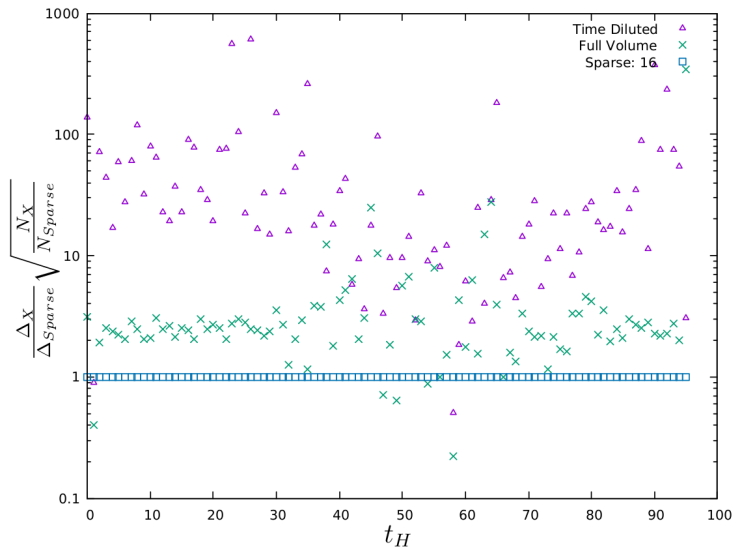
SPARSE SOURCES (LOOPS)



$$\kappa_{\text{sparse}}^{\mathbf{a}}(\mathbf{x}) = \begin{cases} \kappa_{\mathbb{Z}_2}(\mathbf{x}), & \mathbf{x} = \mathbf{a}, \mathbf{a} \in (\mathbb{Z}/2\mathbb{Z})^4 \\ 0, & \text{else} \end{cases}$$

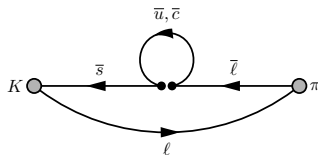
$$\mathbb{Z}/2\mathbb{Z} = \{0 \pmod{2}, 1 \pmod{2}\}$$

NOISE COMPARISON, SAUCER DIAGRAM

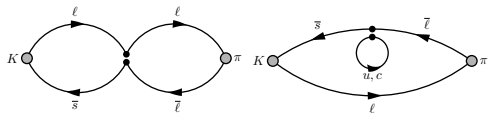


Cost comparison of 3 noise strategies

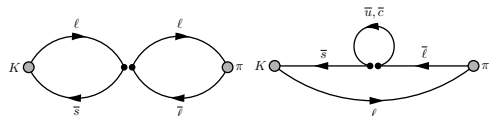
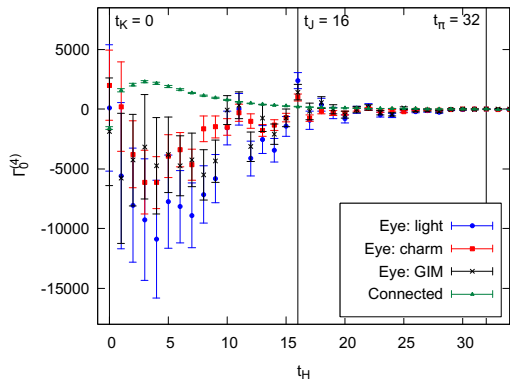
- full-volume Z_2 noise sources
- sparse Z_2 noise sources
- time-diluted A2A vectors



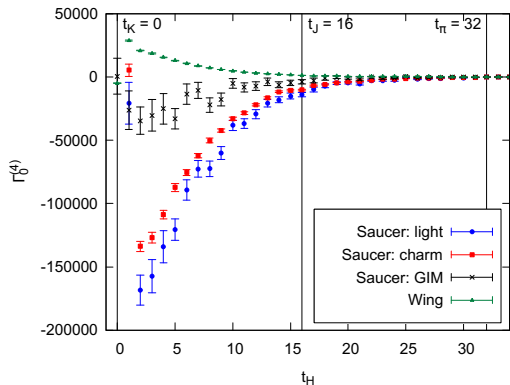
INDIVIDUAL RESULTS AT $m_c = 0.25$



Q_1



Q_2

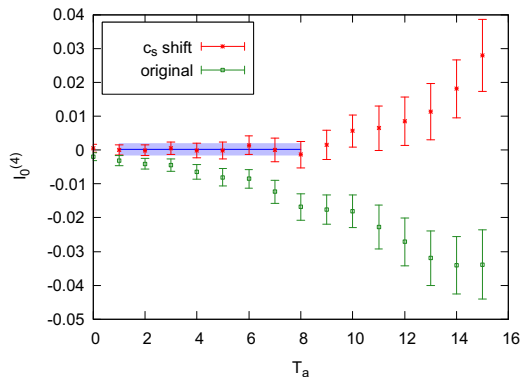


- any intermediate state with $E < E_K$ leads to exponentially growing contribution
 \Rightarrow needs to be subtracted [Christ et al., 1507.03094]
- Single-pion intermediate state:
- method 1: construct π state directly from correlation functions, subtract explicitly
- method 2: shift H_W by scalar density

$$H'_W = H_W - c_s \bar{s}d, \quad \langle \pi(\vec{k}) | H'_W | K(\vec{k}) \rangle = 0$$

- we compared both, settled for method 2
- Other states ($\pi\pi$, $\pi\pi\pi$) are expected to be insignificant until calculations reach sub-percent precision

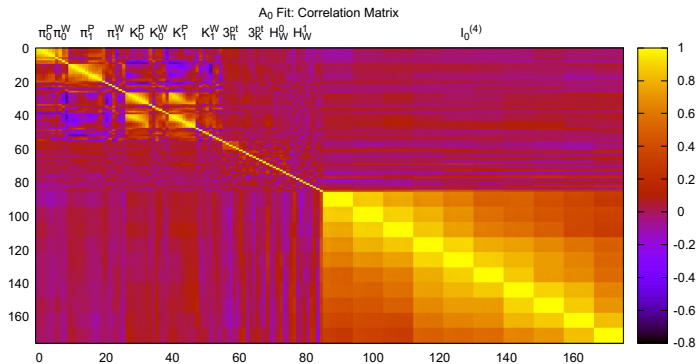
T_a DEPENDENCY AT $m_c = 0.25$



$$I(T_a, T_b) = \int_{t_J - T_a}^{t_J + T_b} dt_H \Gamma^{(4)}(t_J, t_H)$$

- $T_b = 8$ fixed
- green data has exponentially growing terms from intermediate π states
- red: π states manually removed
- fit result: $A_0 = 0.00022(172)$

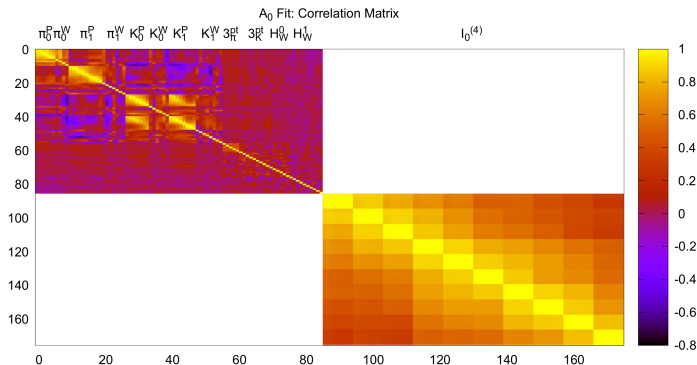
COMBINED FIT



A_0 via combined fit to

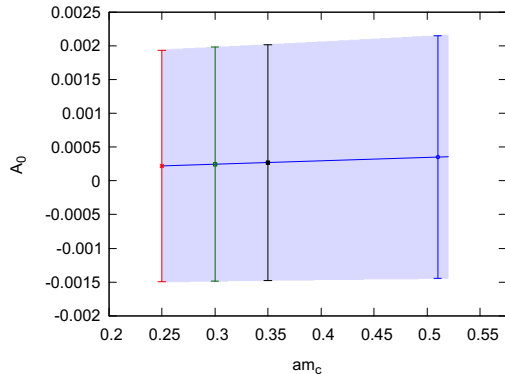
- 2pt functions for π , K
 - point (P) and wall (W) sinks
 - zero momentum or $\frac{2\pi}{L}(1, 0, 0)$
- vector currents 3_π^{pt} and (3_K^{pt})
- weak current without (H_W^0) and with momentum (H_W^1)
- 4pt correlator $I_0^{(4)}(T_a, T_b)$

COMBINED FIT



- Highly correlated
- we treat $I_0^{(4)}(T_a, T_b)$ as de-correlated from all other fit variables

AMPLITUDE



- all 3 charm masses m_c give compatible results
- ... but also compatible with 0
- extrapolated result $A_0 = 0.00035(180)$
- due to high statistical error, we ignored systematic effects

COMPARISON TO EARLIER WORK (Q1 DIAGRAMS)

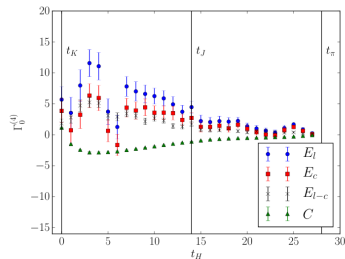
First exploratory calculation of the long-distance contributions to

the rare kaon decays $K \rightarrow \pi \ell^+ \ell^-$

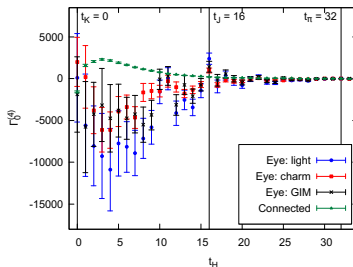
Norman H. Christ,¹ Xu Feng,^{1,2,3,4} Andreas Jüttner,⁵

Andrew Lawson,⁵ Antonin Portelli,^{5,6} and Christopher T. Sachrajda⁵

$M_\pi = 430\text{MeV}$



this work
 $M_\pi = 139\text{MeV}$



- at M_π^{phys} , u, c loops statistically decorrelate
- even at light m_c !
- GIM subtraction noisier than at unphysical M_π

COMPARISON TO EARLIER WORK (Q2 DIAGRAMS)

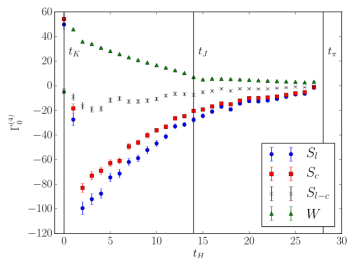
First exploratory calculation of the long-distance contributions to

the rare kaon decays $K \rightarrow \pi \ell^+ \ell^-$

Norman H. Christ,¹ Xu Feng,^{1,2,3,4} Andreas Jüttner,⁵

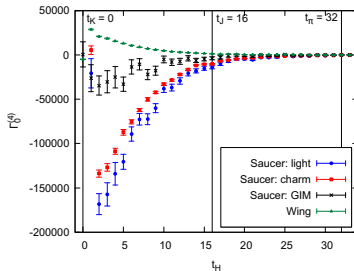
Andrew Lawson,⁵ Antonin Portelli,^{5,6} and Christopher T. Sachrajda⁵

$M_\pi = 430\text{MeV}$



this work

$M_\pi = 139\text{MeV}$



- at M_π^{phys} , u, c loops statistically decorrelate
- even at light m_c !
- GIM subtraction noisier than at unphysical M_π

CONCLUSIONS

- first lattice QCD calculation of $K \rightarrow \pi \ell^+ \ell^-$ and m_{π}^{phys}
- form factor of the long-distance amplitude: $V(0.013(2)) = -0.87(4.44)$
 $\Rightarrow \sim 10$ times experimental error
 - $V_{\mu}^{\text{exp}}(0) = a_{\mu}^{\text{exp}} = -0.578(16)$
 - $V_e^{\text{exp}}(0) = a_e^{\text{exp}} = -0.592(15)$
- Naive $1/\sqrt{N}$ scaling suggests factor ~ 100 in statistics to match experimental error, prohibitively expensive
- manuscript of this study to be submitted for publication within next month

OUTLOOK

- To go forward, new methodology is needed.
 - main source of noise / error comes from loop-type diagrams (E, S)
- ⇒ improvement in estimation of disconnected diagrams is crucial
- Further potential improvement: summed method for $I^{(4)}$
(see Raoul Hodgson's talk):

$$\int_{t_J - T_a}^{t_J + T_b} dt_H \rightarrow \int_0^T dt_J \int_0^T dt_H$$

- in 3-flavour QCD, c loop would not have to be computed
⇒ requires new renormalisation procedure
- on RBC/UKQCD $64^3 \times 96$ M_π^{phys} ensemble ($a^{-1} = 2.38\text{GeV}$), physical m_c could be simulated