# S-wave $\pi \pi \mathrm{I}=0$ and $\mathrm{I}=2$ scattering with physical pion mass 

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## Why shall we perform this calculation?

- $K \rightarrow \pi \pi$ calculation
- Need $E_{\pi \pi}$ and $A_{\pi \pi}$ to compute lattice matrix element $M_{l a t}$
- Need $\delta_{0}$ and $\frac{d \delta_{0}}{d k}$ to compute LL factor
- Understand $\pi \pi$ final state
- First calculation with physical pion mass and disconnected diagram
- A calculation on GPBC lattice
- 2015 results gives $\pi \pi$ energy which is $3 \sigma(7 \sigma$ with more statistics) higher than the phenomenological prediction(Pi-Pi puzzle)

$$
\begin{aligned}
\delta_{0} & =23.8(4.9)(1.2)^{\circ}(P R L, 2015) \\
& =19.1(2.5)^{\circ}(1386 \text { confs }) \\
& \approx 36^{\circ}(\text { Dispersion })
\end{aligned}
$$

- Possible reason: Excited state contamination for $\pi \pi$ state
- Solution: Introducing more operators


## Calculation details and techniques

Operators list:

- Original $\pi \pi$ interpolating operator ( $O_{a}$ or $\pi \pi(111,111)$ ): two single $\pi$ operator with momentum $( \pm 1, \pm 1, \pm 1) \frac{\pi}{L}[\pi(111)]$
- $\sigma$ operator $\left(O_{c}\right.$ or $\sigma$ ):
$\frac{i}{\sqrt{2}}(\bar{u} u+\bar{d} d)$. This operator has the same quantum number as $\pi \pi_{l=0}$
- Extra $\pi \pi$ operator ( $O_{b}$ or $\pi \pi(311,311)$ ):
two new $\pi$ operator with one of their component replaced by $\pm \frac{3 \pi}{L}[\pi(311)]$
General techniques:
- G-parity boundary condition
- All to all propagator
- Non-overlapping blocked bootstrap [cf. 1911.04582]


## Fitting strategies and results

$$
\begin{equation*}
C_{i j}\left(t_{\text {snnk }}, t=t_{\text {snk }}-t_{\text {src }}\right)=\left\langle O_{i}^{\dagger}\left(t_{\text {snk }}\right) O_{j}\left(t_{\text {src }}\right)\right\rangle-\langle 0| O_{i}\left(t_{\text {snk }}\right)|0\rangle\langle 0| O_{j}\left(t_{\text {src }}\right)|0\rangle \times \delta_{l, 0} \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
C_{i j}(t)=\sum_{x=1}^{N} A_{i x} A_{j x}\left(e^{-E_{x} t}+e^{-E_{x}(T-t)}\right)+B_{i j} \times \delta_{l, 2} \tag{2}
\end{equation*}
$$



- Non-negligible excited state contamination in $\pi \pi(111,111)$ operator
- $\sigma$ significantly suppresses both excited state error and statistical error
- $\pi \pi(311,311)$ improves statistical error


## Pi-Pi puzzle: Solution

- Normalized overlap matrix also supports this argument.

|  | state $_{0}$ | state $_{1}$ | state $_{2}$ |
| :---: | :---: | :---: | :---: |
| $\pi \pi(111,111)$ | $1.0(0.0)$ | $0.47(2)$ | $0.31(7)$ |
| $\sigma$ | $1.0(0.0)$ | $-0.83(3)$ | $-0.87(22)$ |
| $\pi \pi(311,311)$ | $0.053(9)$ | $-0.84(12)$ | $1.0(0.0)$ |

- We decide to include all operators

$$
\begin{aligned}
\delta_{0} & =23.8(4.9)(1.2)^{\circ}(\text { PRL, 2015, 1op }) \\
& =19.1(2.5)^{\circ}(1386 \text { confs, } 1 o p) \\
& \approx 36^{\circ}(\text { Dispersion }) \\
& =32.3(1.0)^{\circ}(741 \text { confs, } 3 o p)
\end{aligned}
$$

- Better control over excited state error
- Better statistical error with smaller number of confs


## Moving frame (non-zero total momentum)

- Two sets of $\pi$ operator ( $\pi(111)$ and $\pi(311)$ ) allow us to construct $\pi \pi$ operator with non-zero total momentum $\rightarrow$ Moving frame calculation
- Phase shift at multiple $\sqrt{s}$
- LL factor on lattice from finite difference
- Three different magnitudes of total momentum:

$$
P_{\text {tot }}=(2,0,0) \frac{\pi}{L},(2,2,0) \frac{\pi}{L},(2,2,2) \frac{\pi}{L}
$$

- For each total momentum, three $\pi \pi$ operators:
- $\pi \pi(111,111)$
- $\pi \pi(311,111)$ (Not present in stationary frame calculation)
- $\pi \pi(311,311)$
- Only compute extra contractions and no need to generate operators in quark level
- Use symmetry to reduce the number of contractions from 7848 to 1037 by removing statistically redundant contractions


## Moving frame

- Moving frame calculation is more vulnerable to excited state contamination error due to the denser spectrum of states

- As total momentum increases ( $\sqrt{s}$ decreases), the interaction between two pions decreases, making the inter-coupling between different states and operators decreases (will see that later)
$\rightarrow$
Multiple operators become less powerful


## Moving frame




- Effect of additional operators is less significant
- No improvement in statistical error


## Excited state error

- Simply look at plateau region is less reliable due to rapid increase of error as a function of $t_{\text {min }}$ (fake plateau)
- Including additional state in fitting function makes the result much noisier, even if we freeze the energy of that state
- Different patterns in data: The overlap matrix in moving frames are highly diagonal, while most elements have similar size in stationary frame $I=0$

| $P_{C M}=(2,2,2)$ | state $_{0}$ | state $_{1}$ | state $_{2}$ |
| :---: | :---: | :---: | :---: |
| $\pi \pi(111,111)$ | $1.0(0.0)$ | $-0.07(1)$ | $-0.035(8)$ |
| $\pi \pi(111,311)$ | $-0.013(6)$ | $1.0(0.0)$ | $-0.19(5)$ |
| $\pi \pi(311,311)$ | $-0.015(2)$ | $0.05(2)$ | $1.0(0.0)$ |


| $P_{C M}=(0,0,0)$ | state $_{0}$ | state $_{1}$ | state $_{2}$ |
| :---: | :---: | :---: | :---: |
| $\pi \pi(111,111)$ | $1.0(0.0)$ | $0.47(2)$ | $0.31(7)$ |
| $\sigma$ | $1.0(0.0)$ | $-0.83(3)$ | $-0.87(22)$ |
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- Use different error estimation methods on data with different patterns


## Excited state error

- Include one extra state in the optimal fit
- Energy is obtained from dispersive model
- Nearly diagonal overlap matrix $\Rightarrow$ obtain overlap factors from single operator fit
- Operators couple to all states strongly $\Rightarrow$ obtain overlap factors from multiple operator fit with smaller $t_{\text {min }}$
- Calculate the maximum energy difference between optimal fit and extra state fit


## Final result

| $P_{\text {tot }}$ | I | $\sqrt{s}(\mathrm{MeV})$ | $\delta$ | $\Delta \delta_{\text {dis }}$ | $\Delta \delta_{\mathrm{FV}}$ | $\Delta \delta_{\text {unphy }}$ | $\Delta \delta_{\text {exc }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $(0,0,0) \frac{\pi}{L}$ | 0 | 471.0 | $32.3(\mathbf{1 . 0})(1.4)$ | 0.64 | 0.32 | 0.83 | $\mathbf{0 . 9 0}$ |
| $(2,0,0) \frac{L}{L}$ | 0 | 435.1 | $24.0(\mathbf{3 . 4})(7.6)$ | 0.46 | 0.23 | 0.71 | $\mathbf{7 . 6}$ |
| $(2,2,0) \frac{\pi}{L}$ | 0 | 365.6 | $18.0(\mathbf{4 . 5})(4.9)$ | 0.36 | 0.18 | 0.47 | $\mathbf{4 . 9}$ |
| $(0,0,0) \frac{\pi}{L}$ | 2 | 565.4 | $-10.98(22)(44)$ | 0.20 | 0.10 | 0.18 | $\mathbf{0 . 3 4}$ |
| $(2,0,0) \frac{\pi}{L}$ | 2 | 479.1 | $-7.96(23)(29)$ | 0.16 | 0.08 | 0.03 | $\mathbf{0 . 2 3}$ |
| $(2,2,0) \frac{\pi}{L}$ | 2 | 386.7 | $-4.48(40)(77)$ | 0.09 | 0.04 | 0.06 | $\mathbf{0 . 7 6}$ |
| $(2,2,2) \frac{\pi}{L}$ | 2 | 271.5 | $-0.32(20)(63)$ | 0.01 | 0.00 | 0.02 | $\mathbf{0 . 6 3}$ |

- Large statistical error with moving $I=0$
- Small excited state error with stationary $I=0$ and all $I=2$ cases
- Huge excited state error with moving $I=0$
$\sigma$ operator might be critical


## Final plot



## Conclusion and future steps

Conclusions:

- Understand and solve Pi-Pi puzzle
- Calculate $\pi \pi$ scattering with $I=0$ and $I=2$ at seven different energies
- Estimate excited state error carefully

Future steps:

- Add moving $\sigma$ operator in moving frame $I=0$ calculation
- Using ensembles with different lattice spacing to extrapolate to the continuum limit
- Combine this result with other lattice results to improve the dispersive prediction

