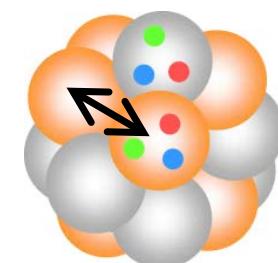
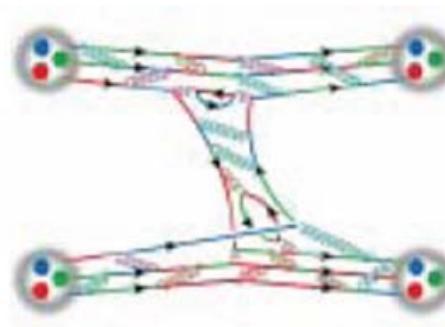
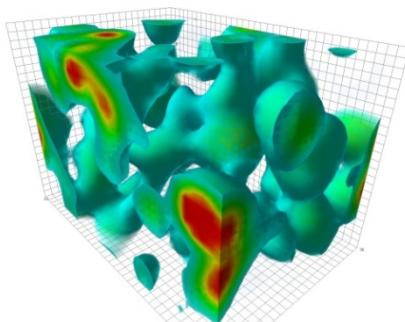


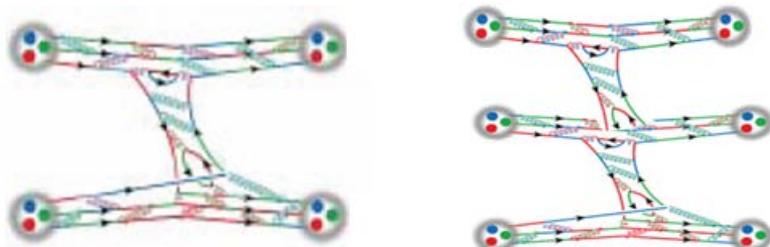
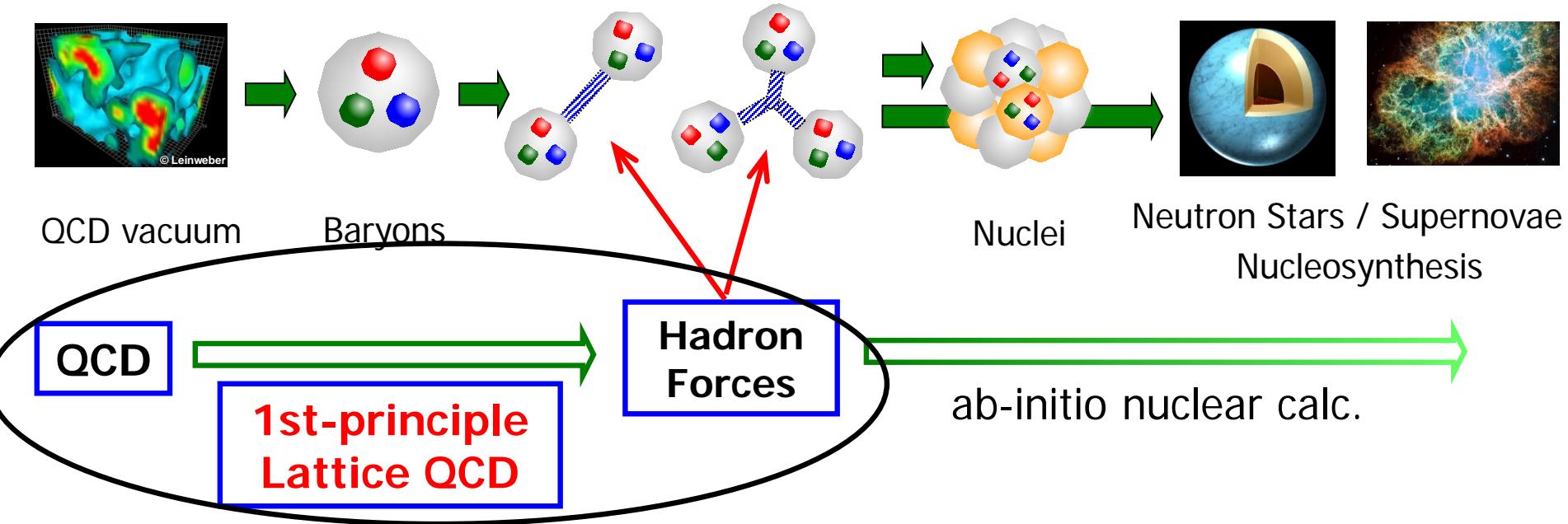
Hadron Interactions from Lattice QCD

Takumi Doi

(RIKEN Nishina Center / iTHEMS)

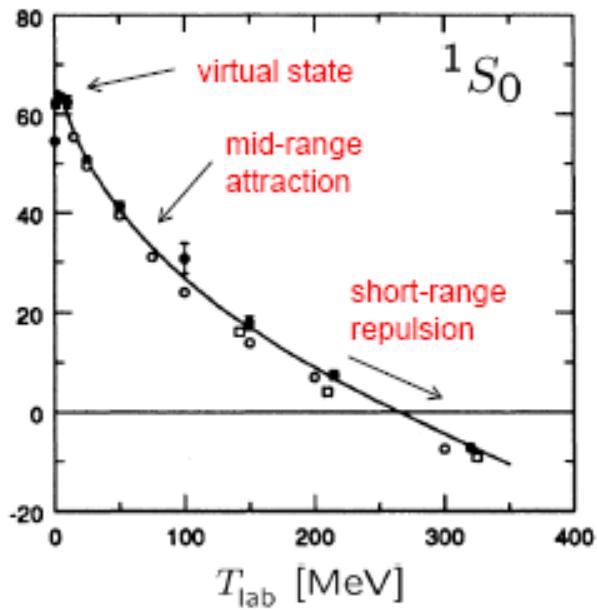


The Odyssey from Quarks to Universe

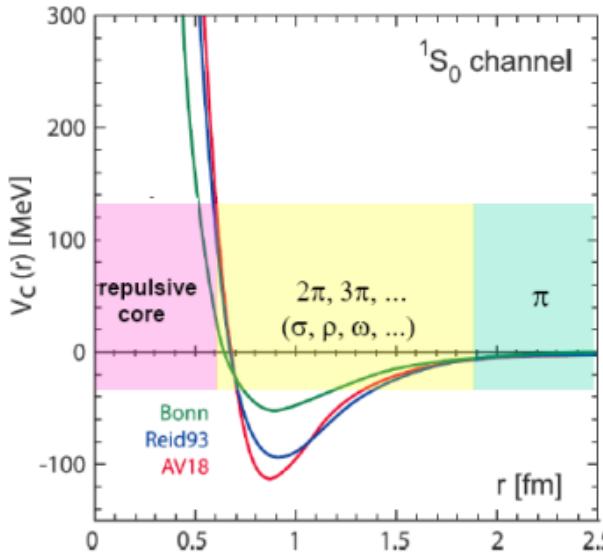
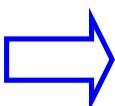


**Nuclear Physics
directly based on QCD**

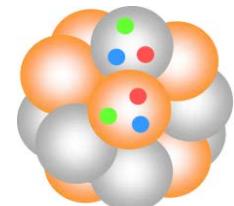
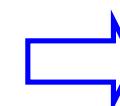
Nuclear Forces: Foundation of nuclear physics



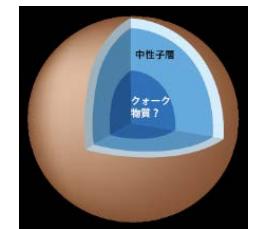
NN phase shifts
from experiments



Phenomenological
Nuclear Forces



Nuclei



Neutron Stars



Super Novae

Various
applications

- ***Nuclear Forces* play crucial roles**
 - Yet, no clear connection to QCD so far

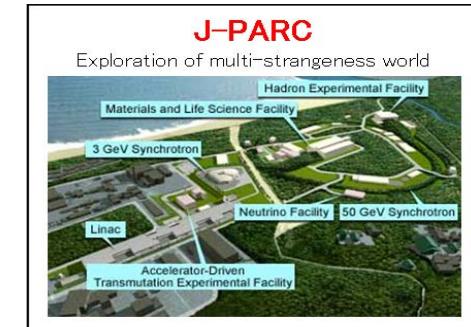
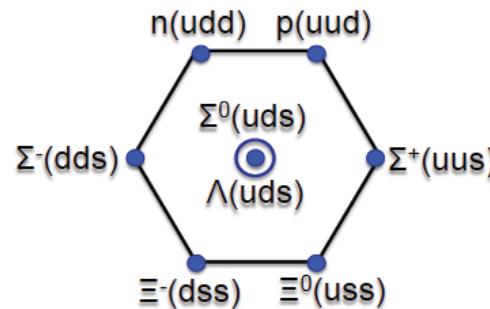
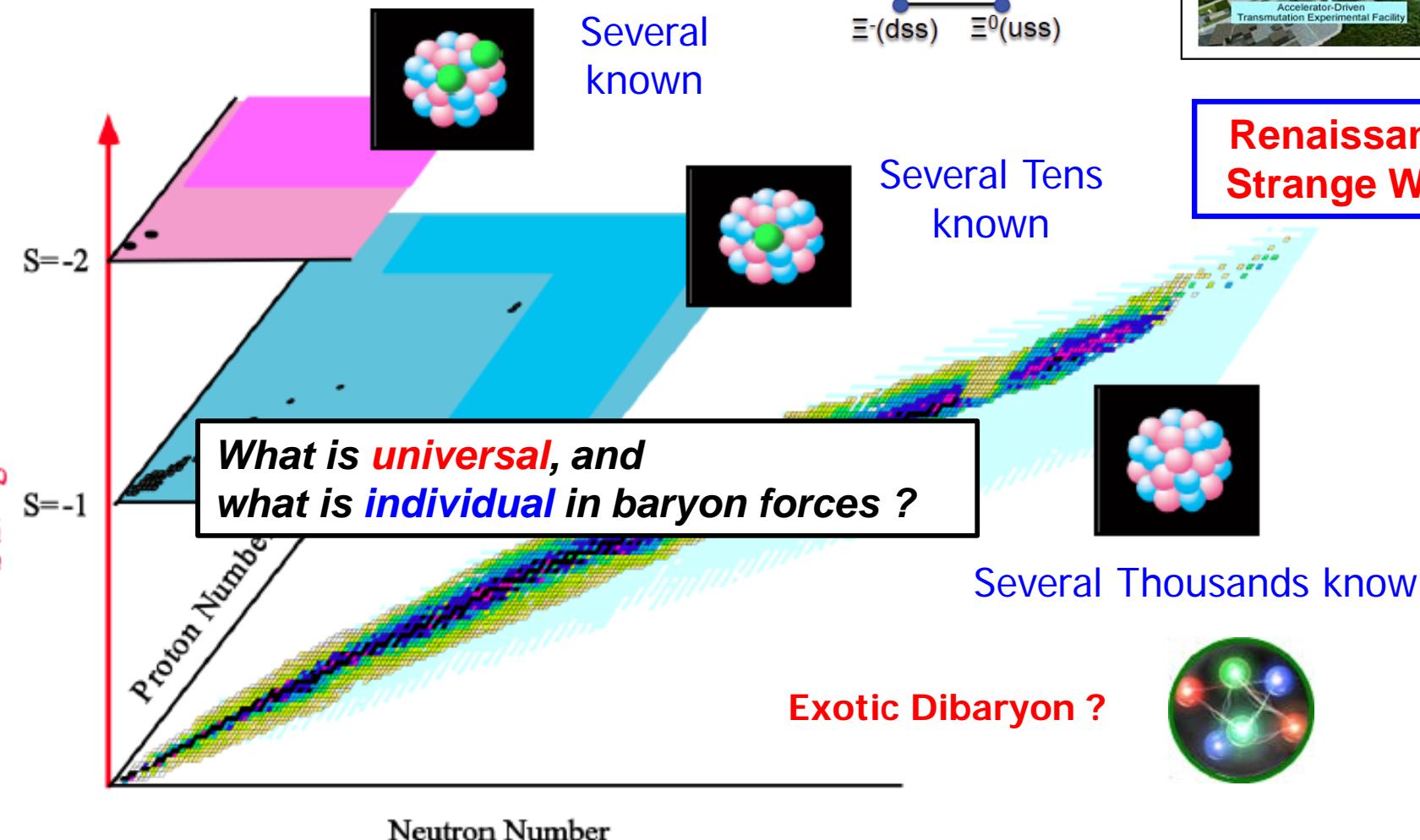
Phen. NN potentials: #params = 30~40

↔ QCD: #inputs = 6 : quark masses (m_u, m_d, m_s, m_c, m_b) & coupling α_s

Nuclear Forces → Baryon Forces (incl. Hyperons)

3D Nuclear Chart

Nucleons : u, d quarks
Hyperons : u, d, s quarks

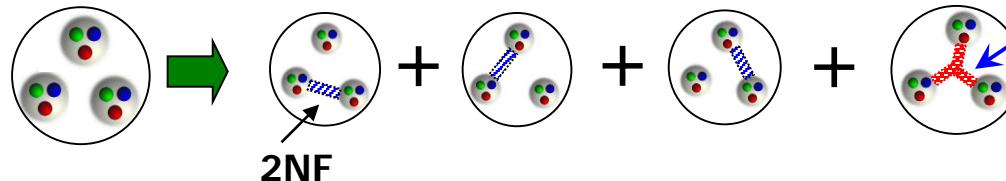


Renaissance in Strange World !



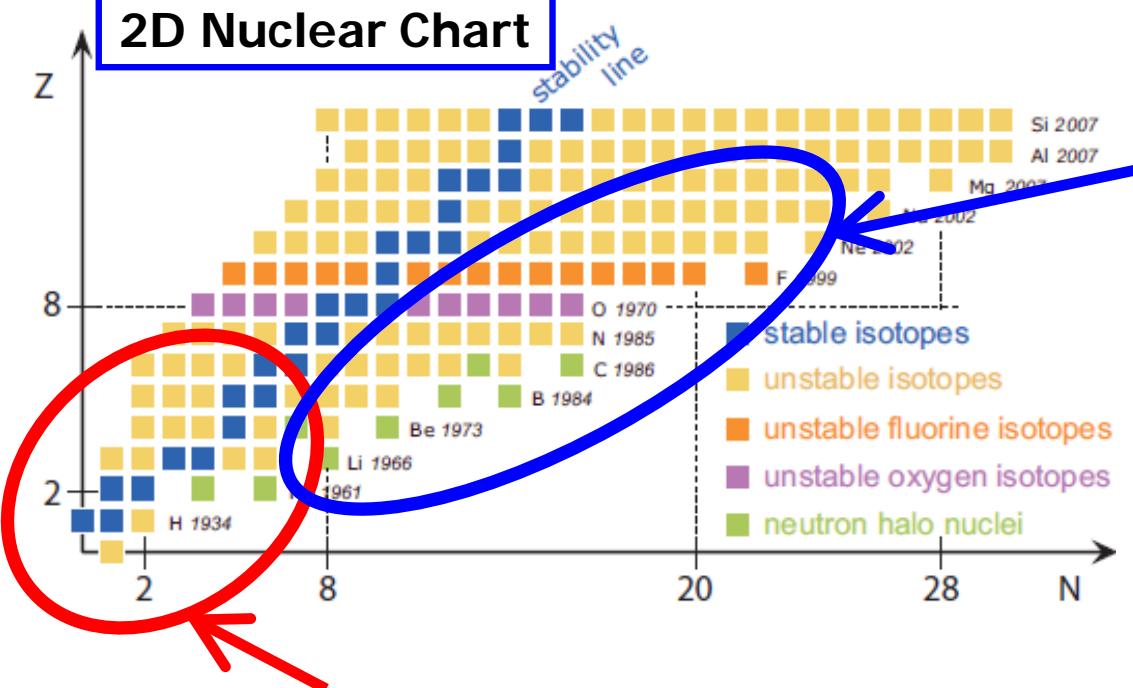
Nuclear Forces → Three-Nucleon Forces (3NF)

What is 3NF ?



3NF: Forces which cannot be explained by pair-wise 2NF

2D Nuclear Chart



**Paradigm Shift in
Unstable Nuclei**
(New Magic Numbers !)

← Important role of 3NF

T.Otsuka et al., PRL105(2010)032501

→ r-process Nucleosynthesis

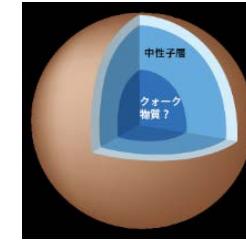


Precise ab initio calculations
show 3NF is indispensable

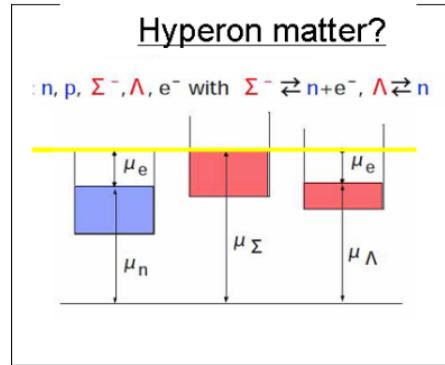
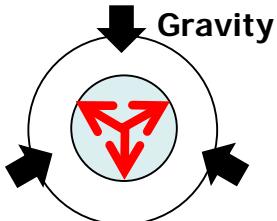
RIBF/FRIB

Dense Matter ← Interactions of YN, YY, + NNN, YNN,... are crucial

- Neutron Stars, Supernovae
↔ EoS of dense matter

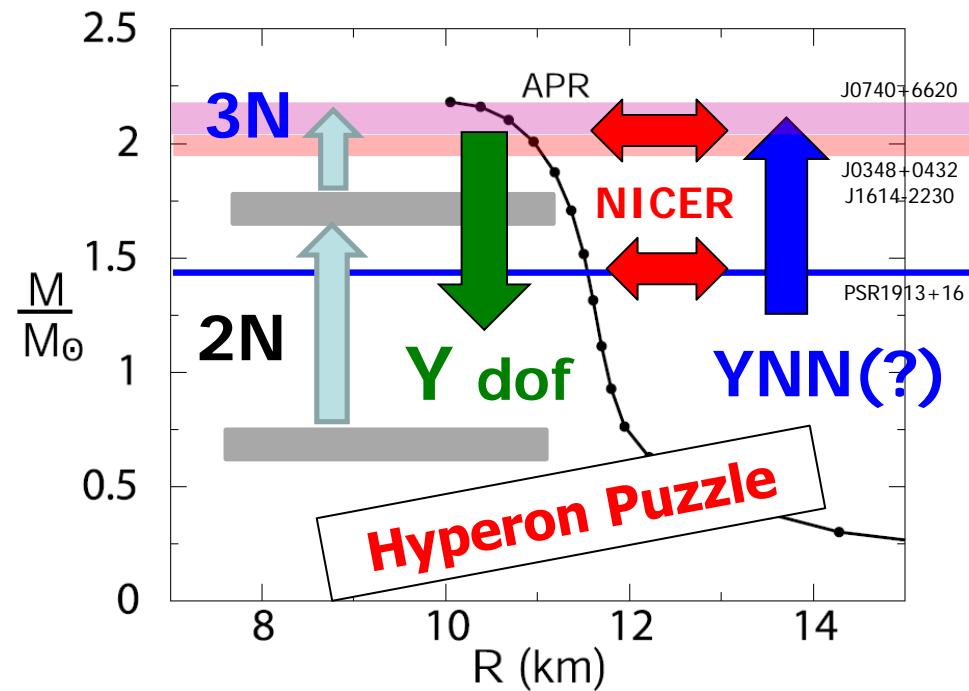


~10km
1-2 Msolar

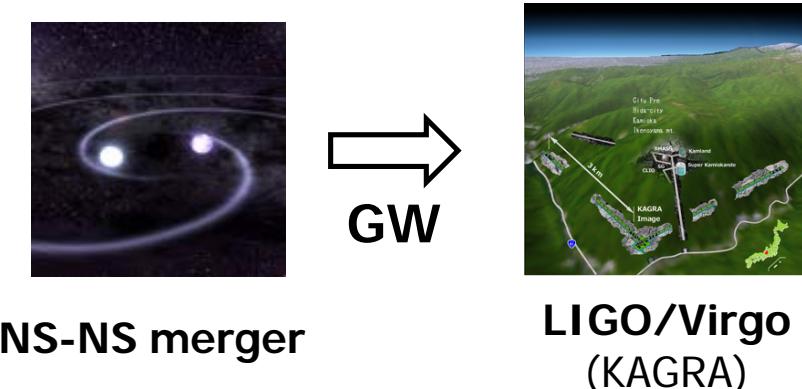


*How to sustain a neutron star
against gravitational collapse ?*

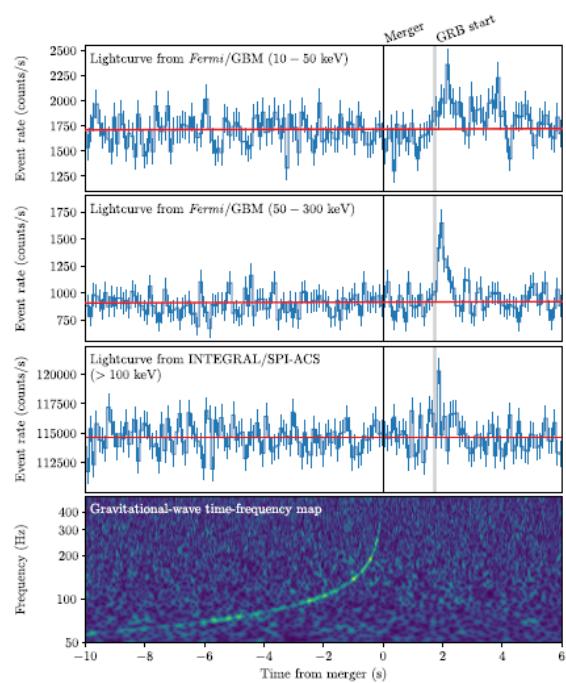
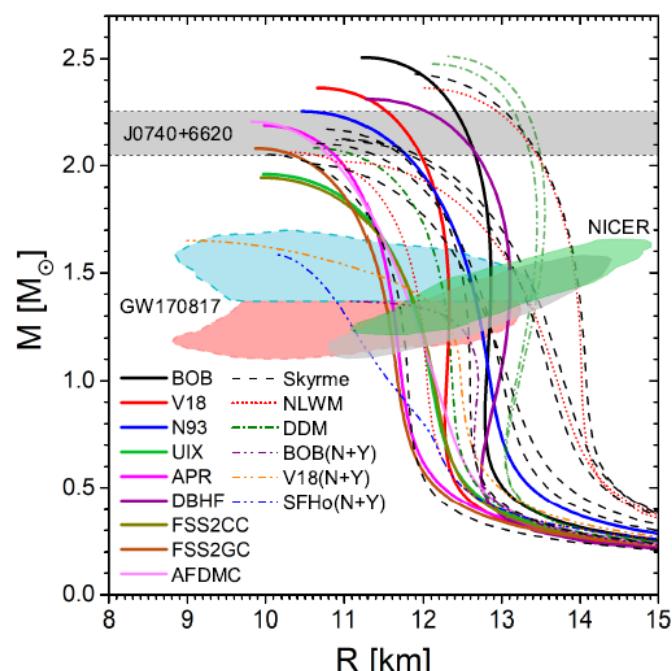
Quark matter ?



Observation of NS-NS merger (GW170817)



Impact on EoS

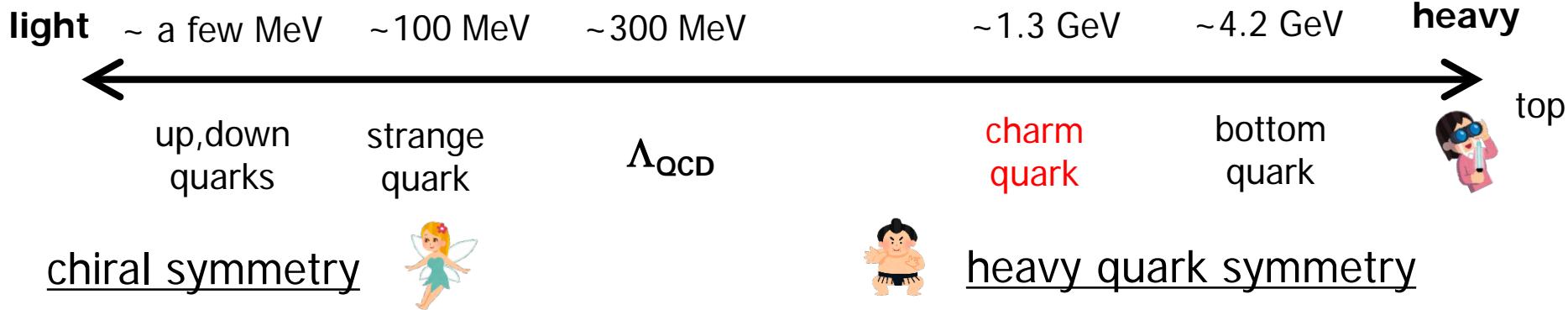


Gravitational Wave
& EM Wave observed

Nucleosynthesis of heavy
elements observed

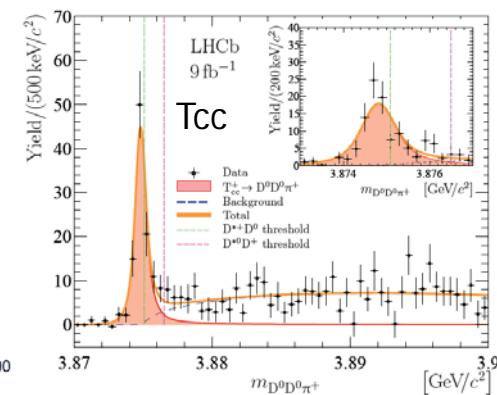
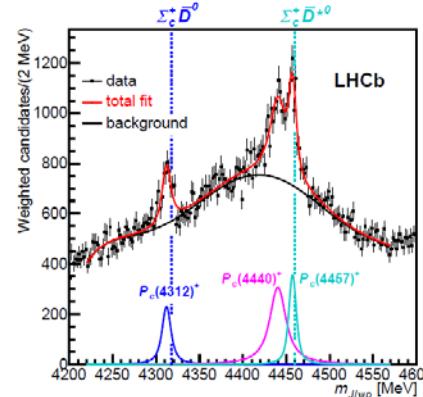
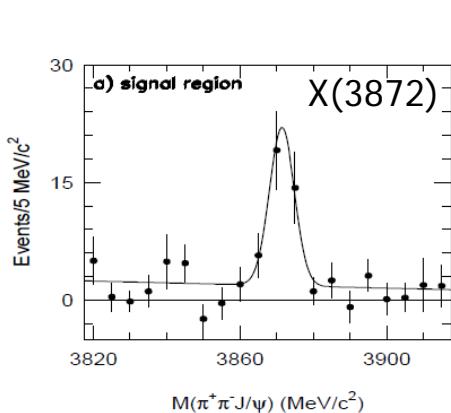
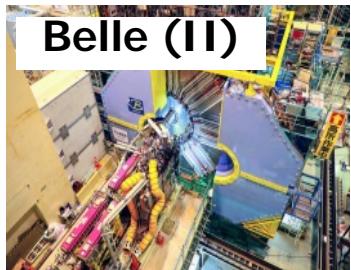
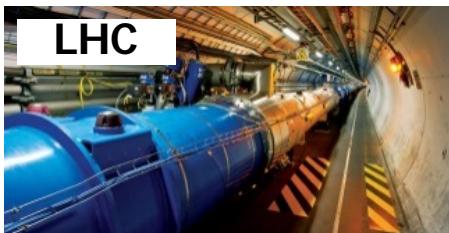
Much more will come!

Nuclear/Hyperon Forces → Charmed Forces



Heavy quarks: New doorway to the mysteries of QCD

Many new exotic particles being reported!

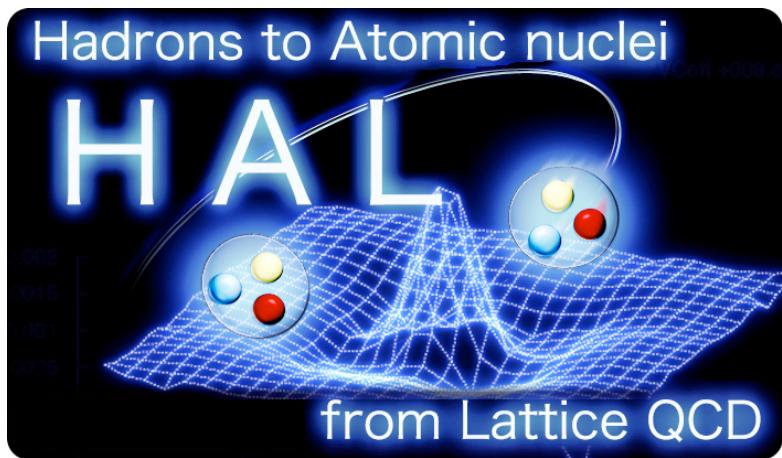


Hadron interactions crucial to understand these "signals" !

e.g., Zc(3900) from HAL LQCD → threshold cusp

Y. Ikeda et al. (HAL), PRL117(2016)242001

Hadrons to **A**tomic nuclei from **L**attice QCD (**HAL** QCD Collaboration)



Y. Akahoshi, S. Aoki,

K. Murakami, H. Nemura (YITP)

T. Aoyama (KEK)

T. Doi, T. Hatsuda, T. Miyamoto, T. Sugiura (RIKEN)

T. M. Doi, N. Ishii, K. Sasaki (Osaka Univ.)

F. Etminan (Univ. of Birjand)

Y. Ikeda (Kyushu Univ.)

T. Inoue (Nihon Univ.)

Y. Lyu (Peking Univ.)

H. Tong (Tianjin Normal Univ.)

+

E. Itou (RIKEN)

I. Kanamori (RIKEN)

K.-I. Ishikawa (Hiroshima Univ.)

「20XX年宇宙の旅」
from Quarks to Universe



- Outline

- Introduction
- Theoretical framework
 - Luscher's formula
 - HAL QCD method
 - Reliability test of LQCD methods
- (Results at heavy quark masses)
- Results near physical quark masses
- Summary / Prospects

Luscher's formula: Scatterings on the lattice

- Consider Schrodinger eq at asymptotic region

$$(\nabla^2 + k^2)\psi_k(r) = mV_k(r)\psi_k(r)$$

$$V_k(r) = 0 \text{ for } r > R$$



- (periodic) Boundary Condition in finite V
 → constraint on energies of the system
- Energy E and phase shift (at E) are related

$$E = 2\sqrt{m^2 + k^2} \quad (\text{QFT: } \psi_k(r) \rightarrow \text{NBS w.f.})$$

$$k \cot \delta_E = \frac{2}{\sqrt{\pi}L} Z_{00}(1; q^2), \quad q = \frac{kL}{2\pi}$$

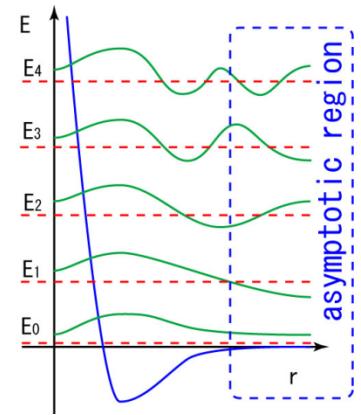
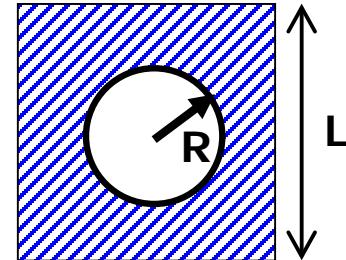
$$Z_{00}(s; q^2) = \frac{1}{\sqrt{4\pi}} \sum_{\mathbf{n} \in \mathbf{Z}^3} \frac{1}{(\mathbf{n}^2 - q^2)^s}$$

Large V expansion

$$\Delta E = E - 2m = -\frac{4\pi \mathbf{a}}{mL^3} \left[1 + c_1 \frac{a}{L} + c_2 \left(\frac{a}{L} \right)^2 + \mathcal{O}\left(\frac{1}{L^3}\right) \right]$$

\mathbf{a} : scattering length

c_1, c_2 : geometric constants



[HAL QCD method]

- “Potential” defined through phase shifts (S-matrix)
- Nambu-Bethe-Salpeter (NBS) wave function

$$\psi(\vec{r}) = \langle 0 | N(\vec{x} + \vec{r}) N(\vec{x}) | N(k) N(-k); W \rangle$$

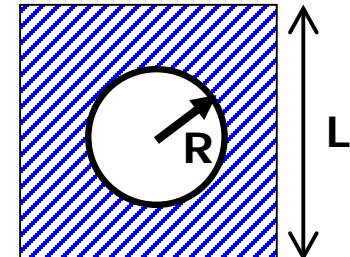
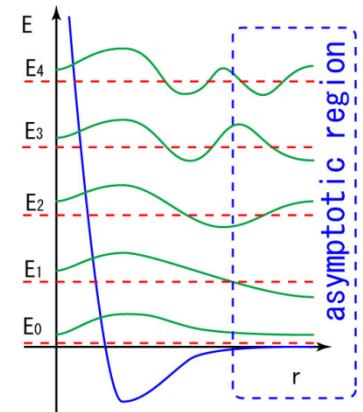
$$(\nabla^2 + k^2)\psi(\vec{r}) = 0, \quad r > R \quad W = 2\sqrt{m^2 + k^2}$$

– Wave function \leftrightarrow phase shifts

$$\psi(r) \simeq A \frac{\sin(kr - l\pi/2 + \delta(k))}{kr}$$

(below inelastic threshold)

Extended to multi-particle systems



M.Luscher, NPB354(1991)531

Ishizuka, PoS LAT2009 (2009) 119

C.-J.Lin et al., NPB619(2001)467

Aoki-Hatsuda-Ishii PTP123(2010)89

CP-PACS Coll., PRD71(2005)094504

S.Aoki et al., PRD88(2013)014036

“Potential” as a representation of S-matrix

- Consider the wave function at “interacting region”

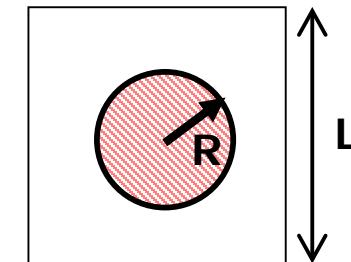
$$(\nabla^2 + k^2)\psi(\mathbf{r}) = m \int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \psi(\mathbf{r}'), \quad r < R$$

- $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: faithful to the phase shift by construction
 - $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: NOT an observable, but well defined
 - $\mathbf{U}(\mathbf{r}, \mathbf{r}')$: E-independent, while non-local in general
 - “Proof of Existence”: Explicit form can be given as

$$\mathbf{U}(\mathbf{r}, \mathbf{r}') = \frac{1}{m} \sum_{n,n'}^{n_{\text{th}}} (\nabla_{\mathbf{r}}^2 + k_n^2) \psi_n(\mathbf{r}) \mathcal{N}_{nn'}^{-1} \psi_{n'}^*(\mathbf{r}') \quad \mathcal{N}_{nn'} = \int d\mathbf{r} \psi_n^*(\mathbf{r}) \psi_{n'}(\mathbf{r})$$

- Non-locality → derivative expansion Okubo-Marshak(1958)

$$\mathbf{U}(\vec{r}, \vec{r}') = \begin{matrix} [V_c(r) + S_{12}V_T(r) + \vec{L} \cdot \vec{S} V_{LS}(r) + \mathcal{O}(\nabla^2)] \delta(\vec{r} - \vec{r}') \\ \text{LO} \qquad \text{LO} \qquad \text{NLO} \qquad \text{NNLO} \end{matrix}$$



Three-Body Forces

- Unitarity of S-matrix

S.Aoki et al. (HAL Coll.), PRD88(2013)014036
Gongyo-Aoki PTEP2018(2018)093B03

$$T_{[L]}(Q) = -\frac{2n^{3/2}}{mQ^{3n-5}} \boxed{e^{i\delta_{[L]}(Q)} \sin \delta_{[L]}(Q)}$$

c.f. R.B. Newton (1974) for n = 3

Similar formula to 2-body system

(w/ diagonalization matrix U which includes dynamics)

- NBS wave function

$$\psi_\alpha([x]) =_{in} \langle 0 | \phi([x]) | \alpha \rangle_{in} =_{in} \langle 0 | N(\vec{x}_1) N(\vec{x}_2) \cdots N(\vec{x}_n) | \alpha \rangle_{in}$$

$$\psi_{[L],[K]}(R, Q_A) \propto \sum_{[N]} U_{[L][N]}(Q_A) e^{i\delta_{[N]}(Q_A)} \boxed{\frac{\sin(Q_A R - \Delta_L + \delta_{[N]}(Q_A))}{(Q_A R)^{(D-1)/2}}} U_{[N][K]}^\dagger(Q_A)$$

Similar asymptotic behavior to 2-body system

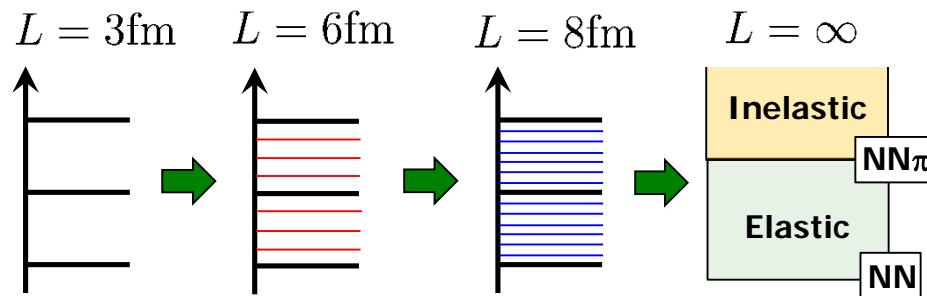
arbitrary n-body
(non-rela approx.)

c.f. Finite V spectrum, n=3 only, relativistic: Hansen, Sharpe, Briceno, ...

The Challenge in multi-baryons on the lattice

Existence of elastic scatt. states

- (almost) No Excitation Energy
- LQCD method based on G.S. saturation impossible



Signal/Noise issue

$$S/N \sim \exp[-\mathbf{A} \times (\mathbf{m_N} - \mathbf{3/2m_\pi}) \times \mathbf{t}]$$

Parisi('84), Lepage('89)

"Sign Problem"

$$L=8\text{fm} @ \text{physical point} \quad (E_1 - E_0) \simeq 25\text{MeV} \implies t > 10\text{fm}$$

$$S/N \sim 10^{-32}$$

Direct method (naïve plateau fitting at $t \sim 1\text{fm}$)

→ Does it really reliable?

Time-dependent HAL method

N.Ishii et al. (HAL QCD Coll.) PLB712(2012)437

E-indep of potential $U(r,r')$ \rightarrow (excited) scatt states share the same $U(r,r')$
They are not contaminations, but signals

Original (t-indep) HAL method

$$G_{NN}(\vec{r}, t) = \langle 0 | N(\vec{r}, t) N(\vec{0}, t) \overline{\mathcal{J}_{\text{src}}(t_0)} | 0 \rangle$$

$$R(\mathbf{r}, t) \equiv G_{NN}(\mathbf{r}, t)/G_N(t)^2 = \sum A_{W_i} \psi_{W_i}(\mathbf{r}) e^{-(W_i - 2m)t}$$

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{\psi_{W_0}(\mathbf{r}')} = (\underline{E_{W_0}} - H_0) \underline{\psi_{W_0}(\mathbf{r})}$$

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{\psi_{W_1}(\mathbf{r}')} = (\underline{E_{W_1}} - H_0) \underline{\psi_{W_1}(\mathbf{r})}$$

...

← Many states contribute

Partial solution of Sign Problem

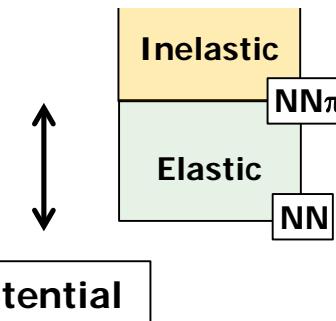
New t-dep HAL method

All equations can be combined as

$$\int d\mathbf{r}' \mathbf{U}(\mathbf{r}, \mathbf{r}') \underline{R(\mathbf{r}', t)} = \left(-\frac{\partial}{\partial t} + \frac{1}{4m} \frac{\partial^2}{\partial t^2} - H_0 \right) \underline{R(\mathbf{r}, t)}$$

~~G.S. saturation~~ \rightarrow "Elastic state" saturation

[Exponential Improvement]



Reliability test of LQCD methods

NN @ heavy quark masses

HAL method (HAL) : unbound

Direct method (NPL/CallLat/PACS-CS(Yamazaki et al.)): bound

Inconsistent!

What is the most plausible systematics ?

HAL QCD method

Energy-indep potential: “signal” from all elastic states

Non-locality of potential: derivative expansion could
introduce systematics

Direct method (= plateau fitting + Luscher’s formula)

Plateau fitting at $t \sim 1\text{fm}$ (much less than $1/(E_1 - E_0)$)

→ Excited states give “noises”

Examine the reliability of the HAL QCD method

Convergence of the derivative expansion of potential

LQCD data: $\Xi\Xi(^1S_0)$ @ $m\pi=0.51\text{GeV}$
wall source & smeared source
Same confs in Yamazaki et al.('12)

Higher Order Approximation (N²LO) (2)

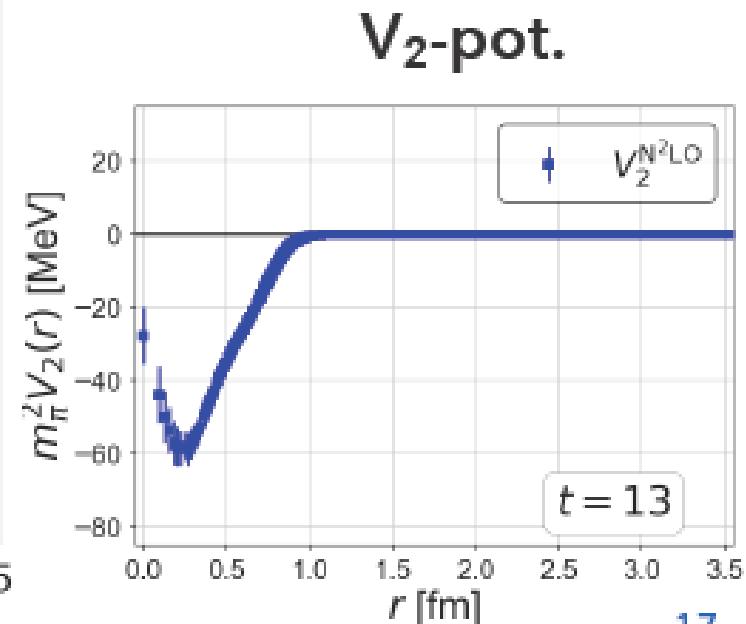
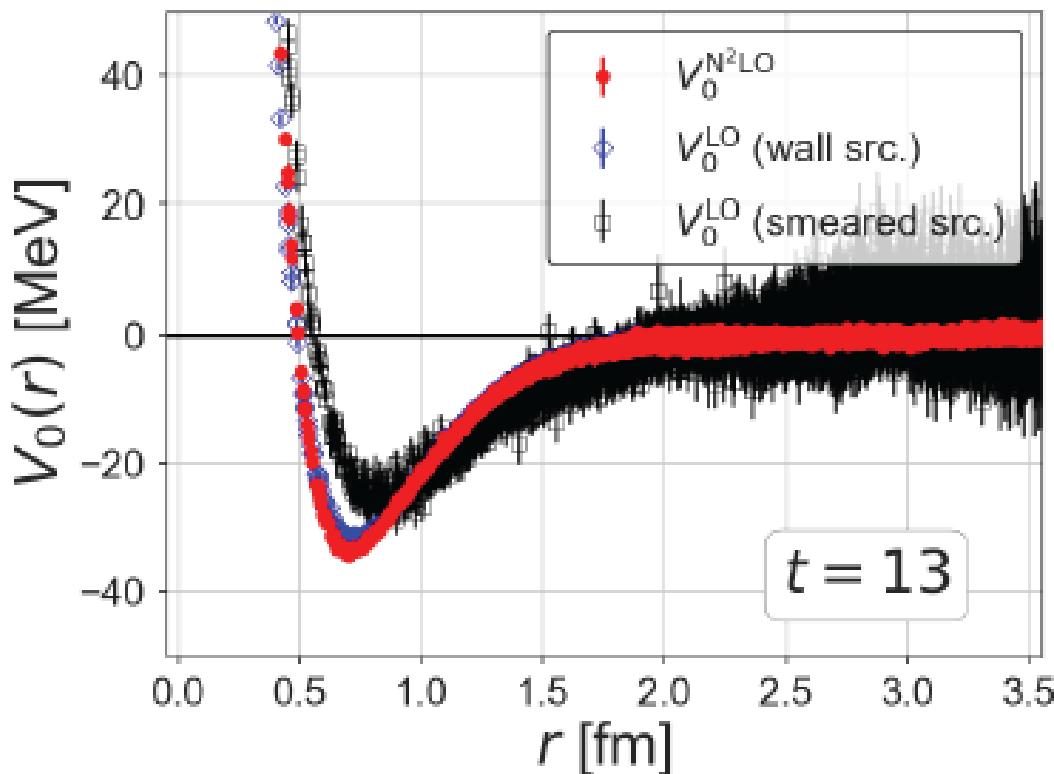
$$U(r, r') \simeq \left[V_0^{\text{N}^2\text{LO}}(r) + V_2^{\text{N}^2\text{LO}}(r) \nabla^2 \right] \delta(r - r')$$

wall src. \rightarrow small $V_2 \nabla^2$ correction

smeared src. \rightarrow large $V_2 \nabla^2$ correction

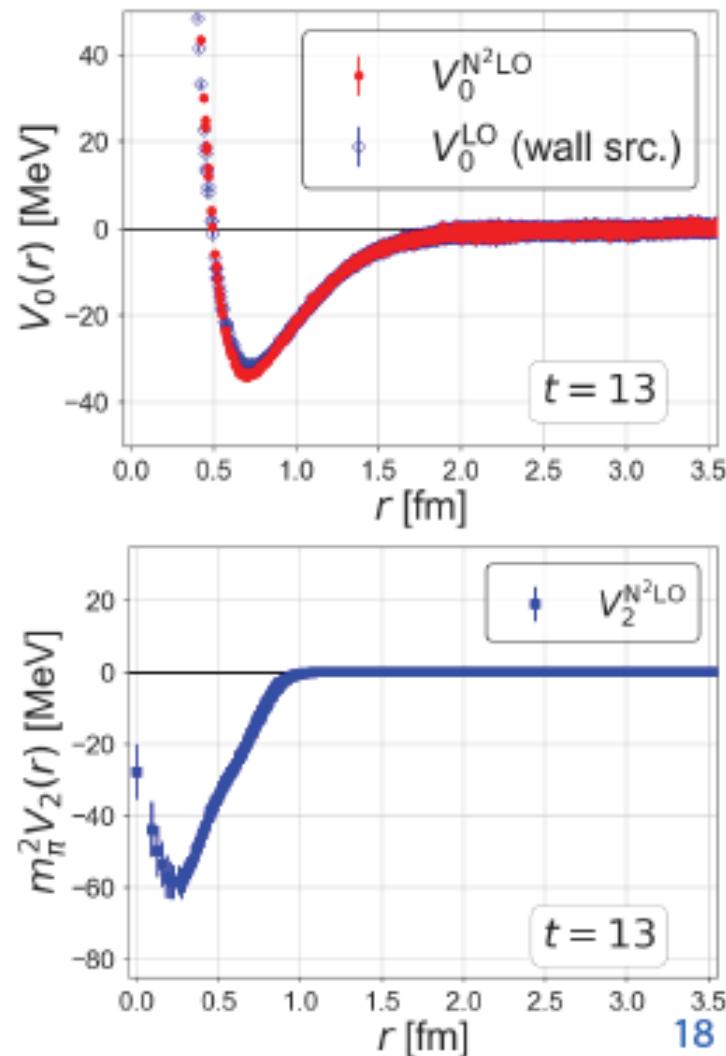
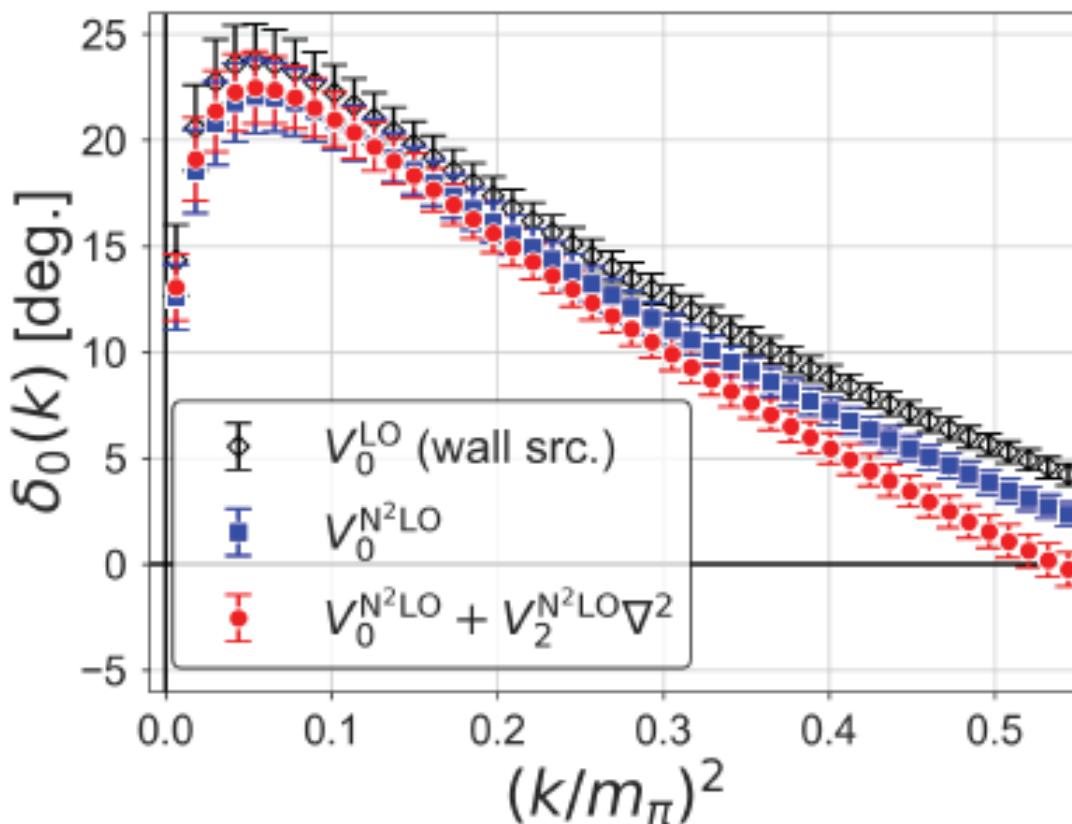
$\rightarrow V_2(r) \underline{\nabla^2 R^{\text{wall/smear}}(r)}$

dep. on shape of R



Phase Shift and Uncertainties in Velocity Expansion

- Wall src. LO approx. (standard of HAL QCD studies)
works well at low energy.
- **V_2 correction** at high energy



Examine the reliability of the Direct method

LQCD data: $\Xi\Xi(^1S_0)$ @ $m\pi=0.51\text{GeV}$

wall source & smeared source

Same confs in Yamazaki et al.('12)

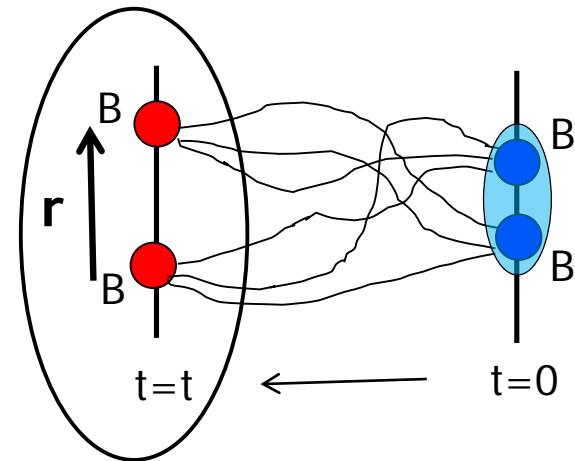
T. Iritani et al. (HAL Coll.) JHEP1610(2016)101

T. Iritani et al. (HAL Coll.) PRD96(2017)034521

Operator dependence in the direct method

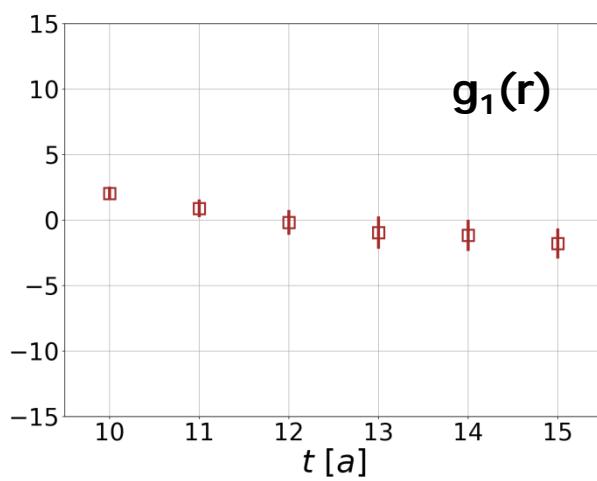
$$C_{2B}(t) = \langle 0 | T[\mathcal{J}_{\text{sink}}^{2B}(t) \overline{\mathcal{J}}_{\text{src}}^{2B}(0)] | 0 \rangle$$

Study sink op dep
w/ smeared src tuned in single-baryon

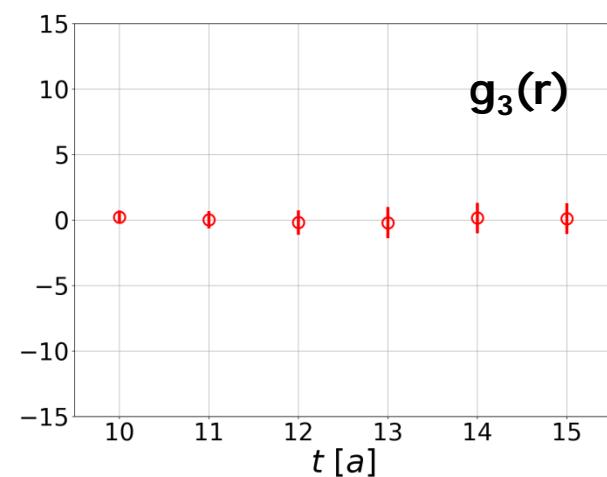
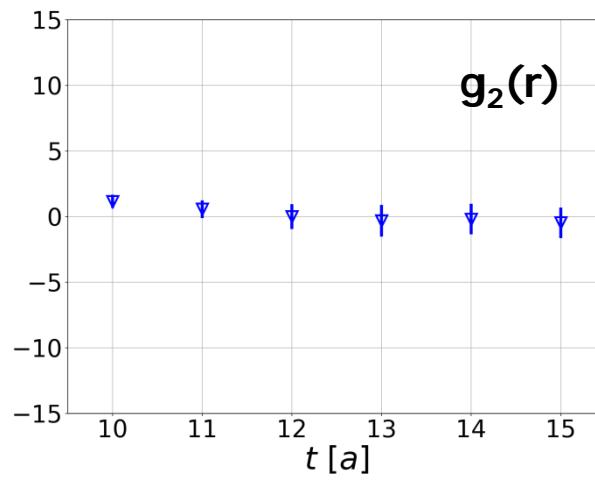


$$\mathcal{J}_{\text{sink}}^{2B} = \sum_{\vec{r}} g(r) \sum_{\vec{x}} B(\vec{r} + \vec{x}) B(\vec{x})$$

Usual direct method: $g(r)=1$ only



Effective Energy shift ΔE



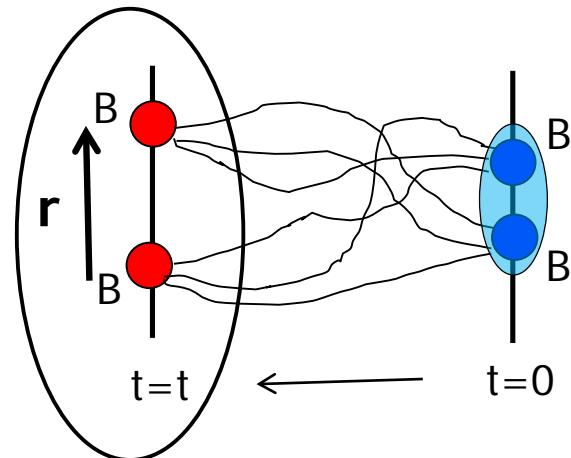
All plateaux “look” reliable

In reality, I shift data vertically “by hand”

Operator dependence in the direct method

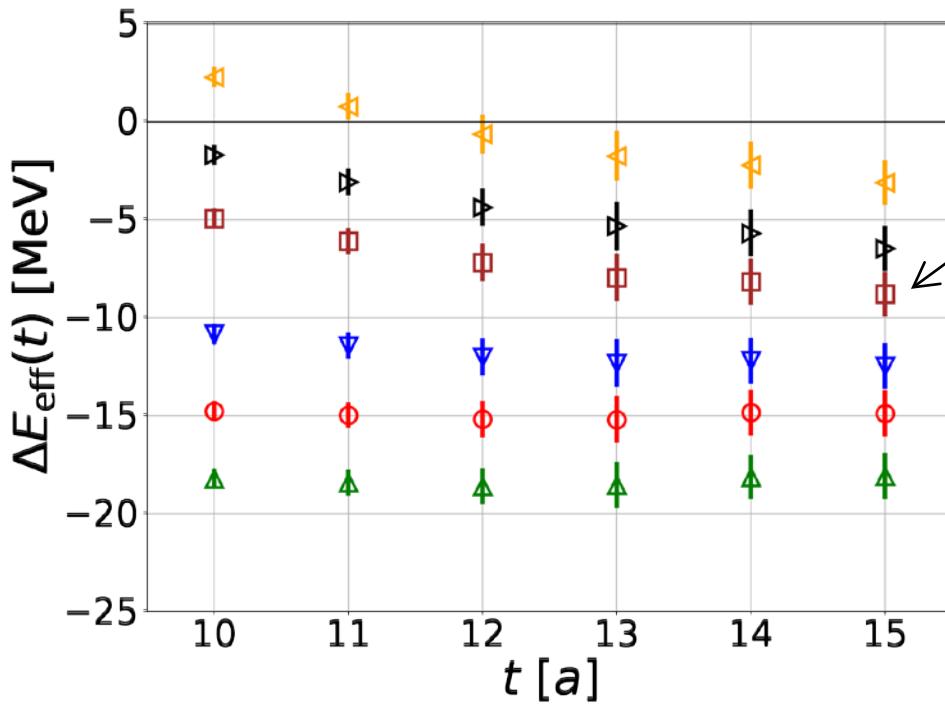
$$C_{2B}(t) = \langle 0 | T[\mathcal{J}_{\text{sink}}^{2B}(t) \overline{\mathcal{J}}_{\text{src}}^{2B}(0)] | 0 \rangle$$

Study sink op dep
w/ smeared src tuned in single-baryon



$$\mathcal{J}_{\text{sink}}^{2B} = \sum_{\vec{r}} g(r) \sum_{\vec{x}} B(\vec{r} + \vec{x}) B(\vec{x})$$

Usual direct method: $g(r)=1$ only



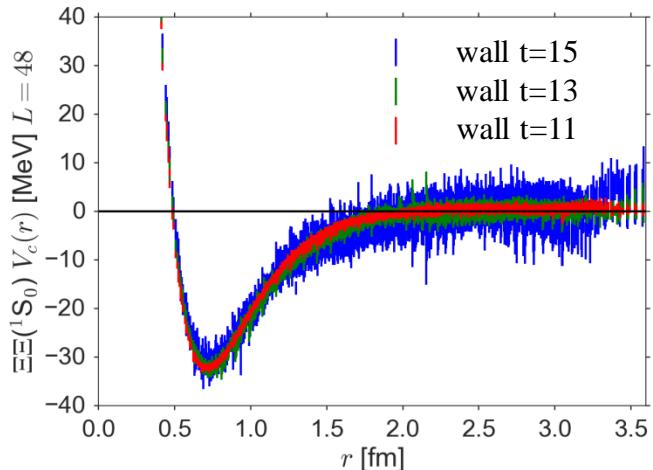
The results depend on
the choice of sink op
→ predictive power is lost!

Anatomy of the Direct method and the consistency between Luscher's formula and HAL method

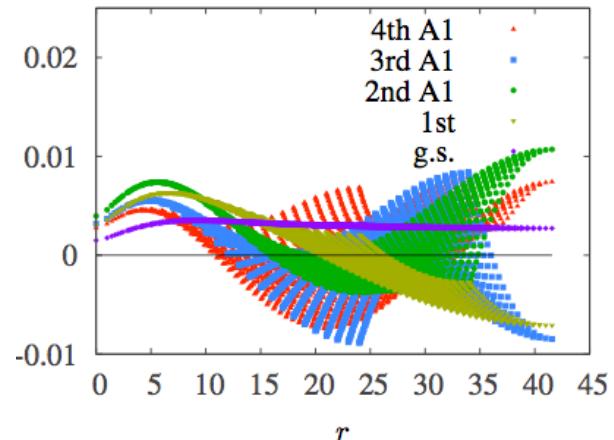
T. Iritani et al. (HAL) JHEP03(2019) 007

Understand the origin of “pseudo-plateaux”

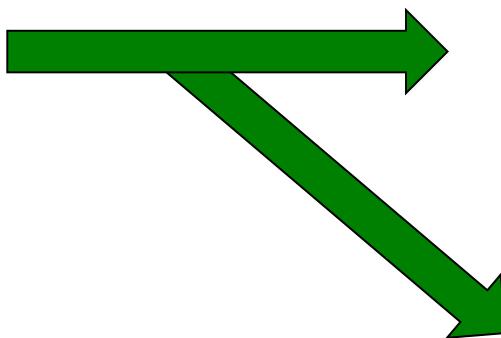
Potential



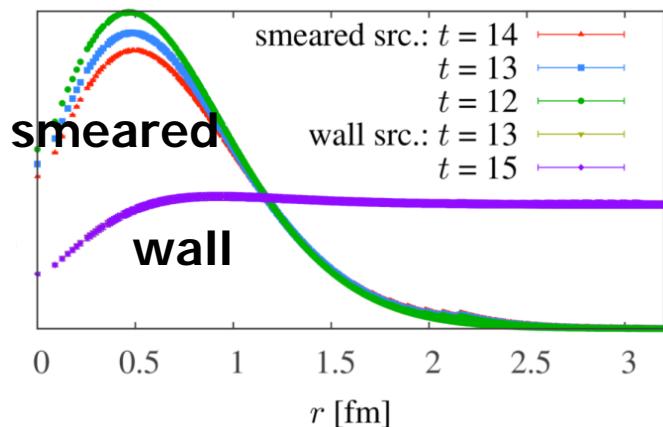
Eigen-wave functions



Solve Schrodinger eq.
in Finite V



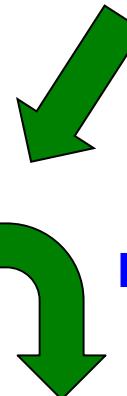
NBS correlator R(r,t)



Eigen-energies

n-th A1	ΔE_n [MeV]
0	-2.58(1)
1	52.49(2)
2	112.08(2)
3	169.78(2)
4	224.73(1)

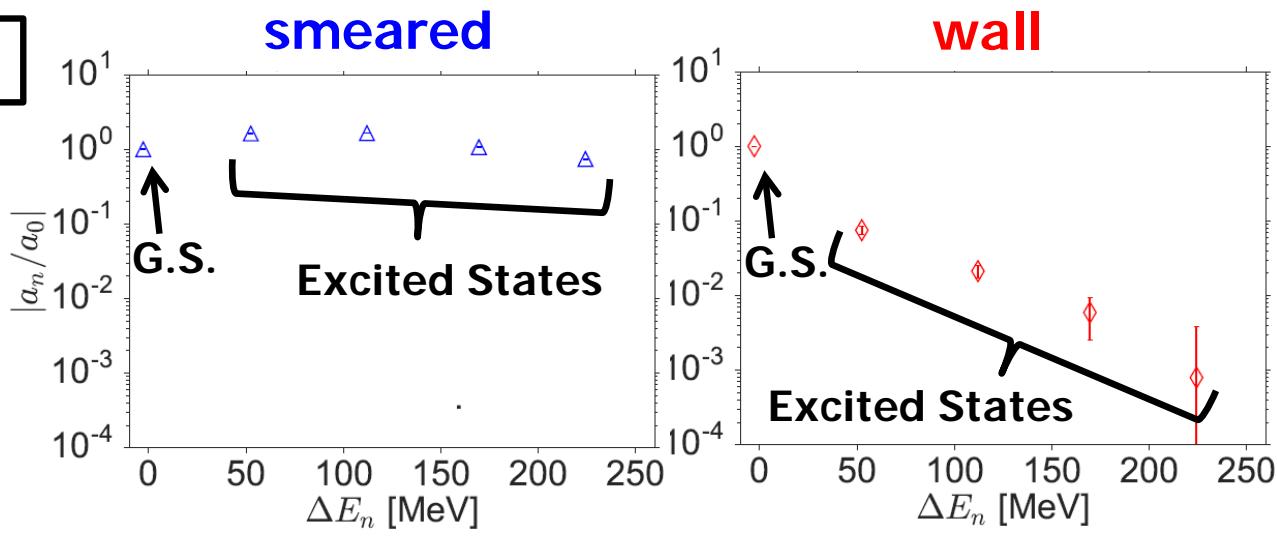
Decompose NBS correlator
to each eigenstates



Decompose NBS correlator
to each eigenstates

NBS correlator $R(r,t)$

Contribution from
each (excited) states
(@ $t=0$)



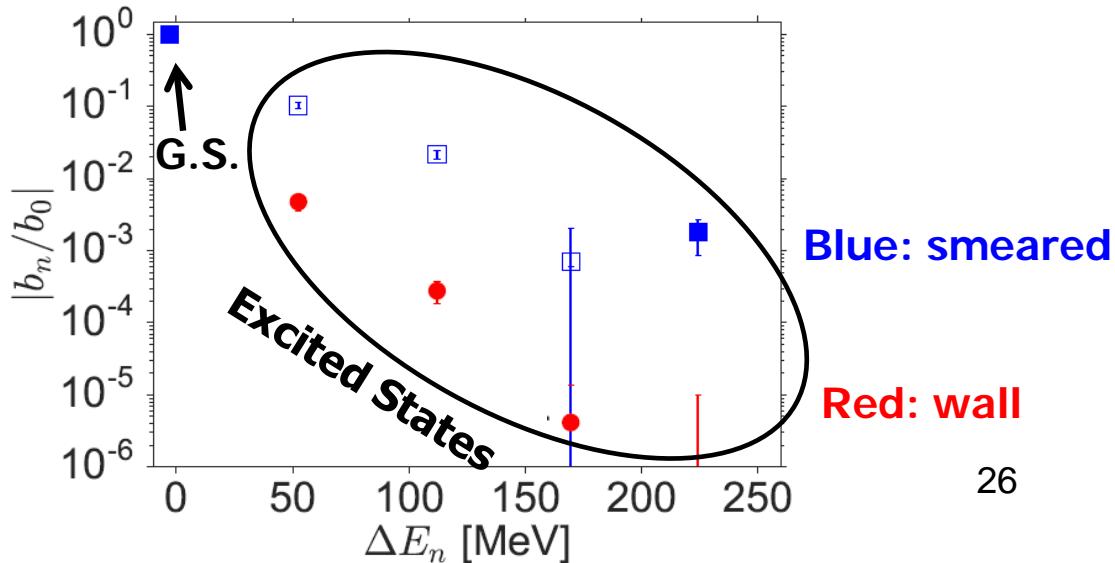
excited states NOT suppressed

excited states suppressed

Temporal-correlator $R(t) = \sum_r R(r,t)$

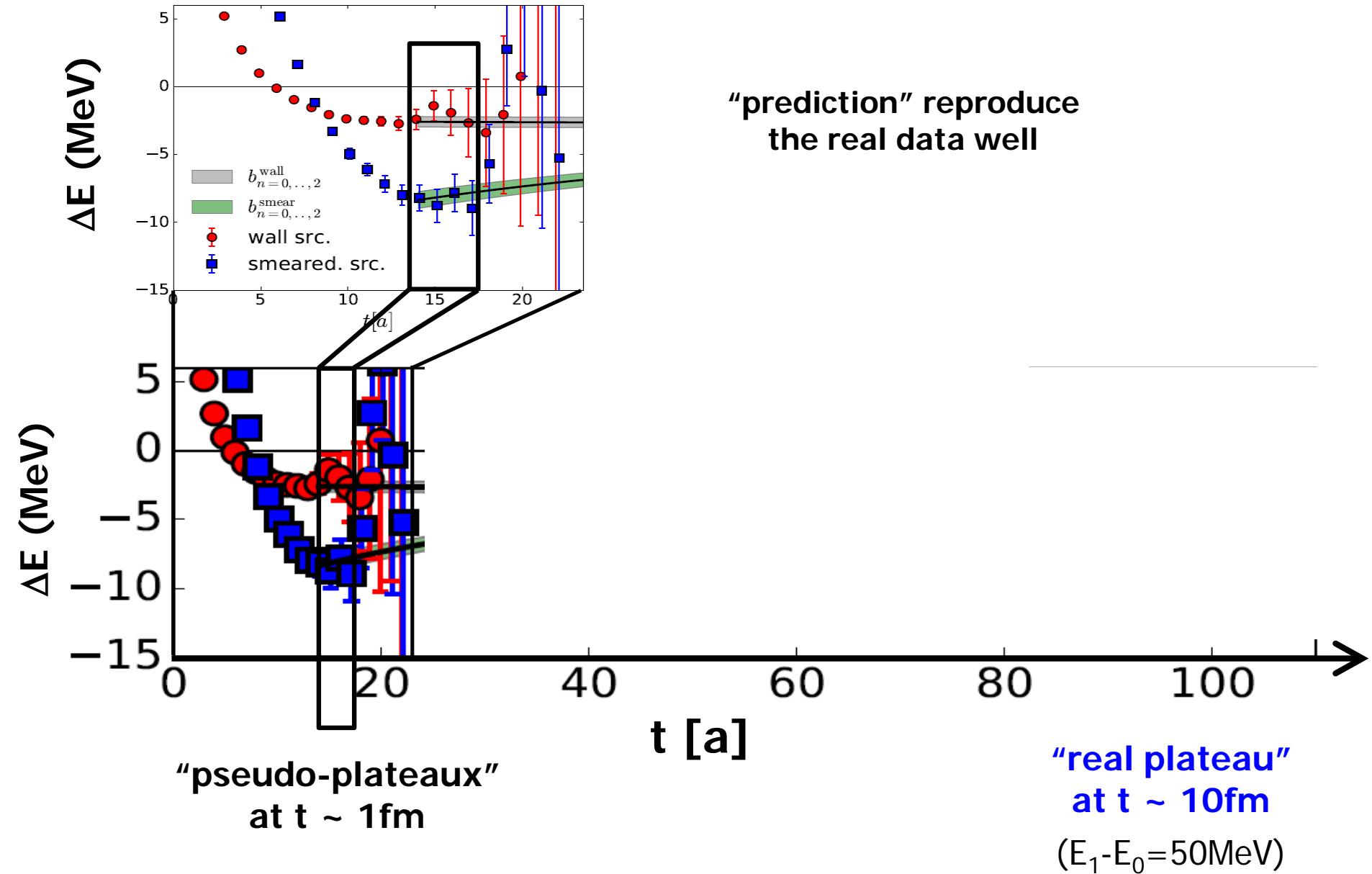
($R(t)$ w/ smeared has been
used in Direct method)

Contribution from
each (excited) states
(@ $t=0$)



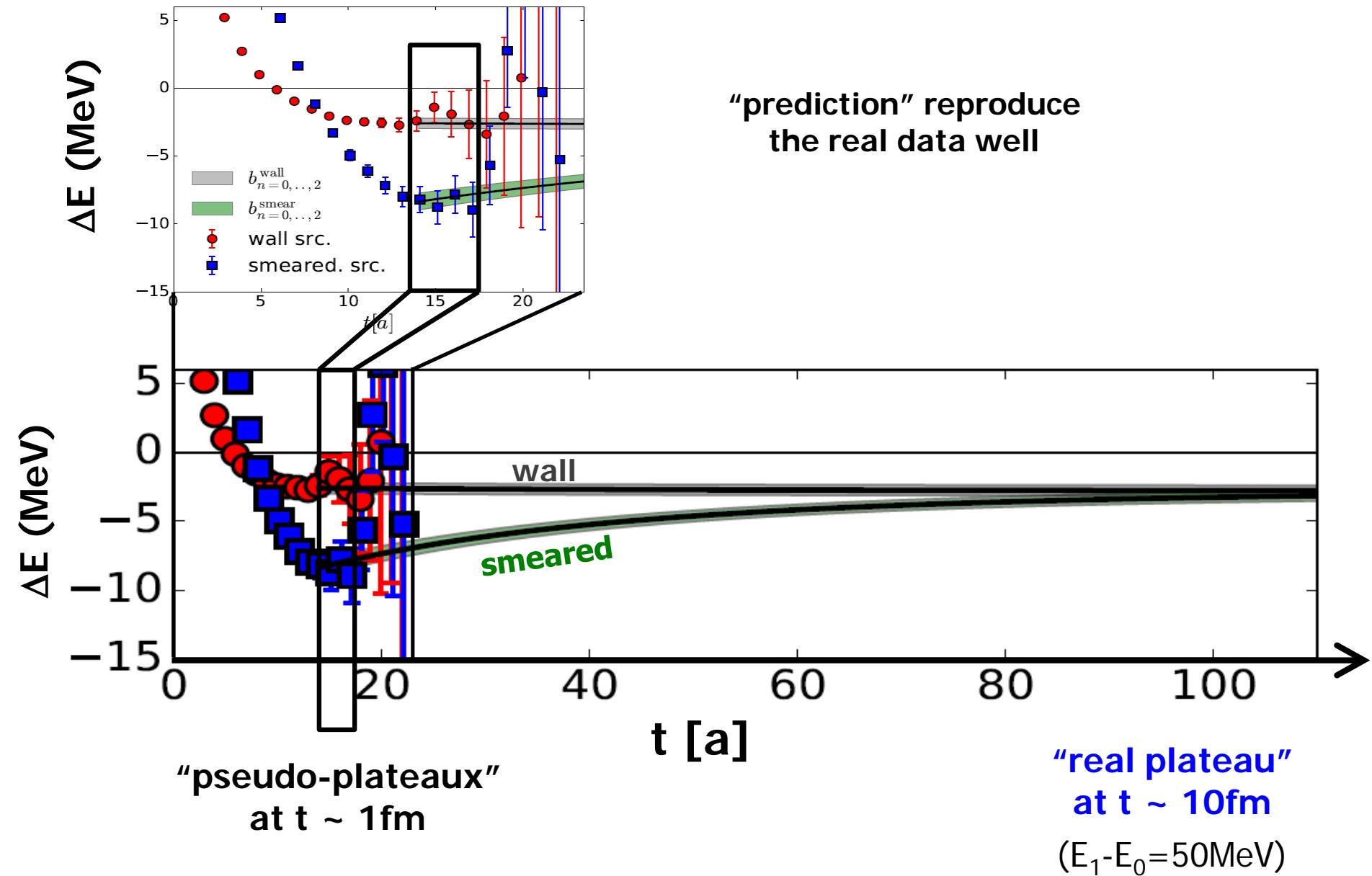
Understand the origin of “pseudo-plateaux”

We are now ready to “predict” the behavior of $m(\text{eff})$ of ΔE at any “t”



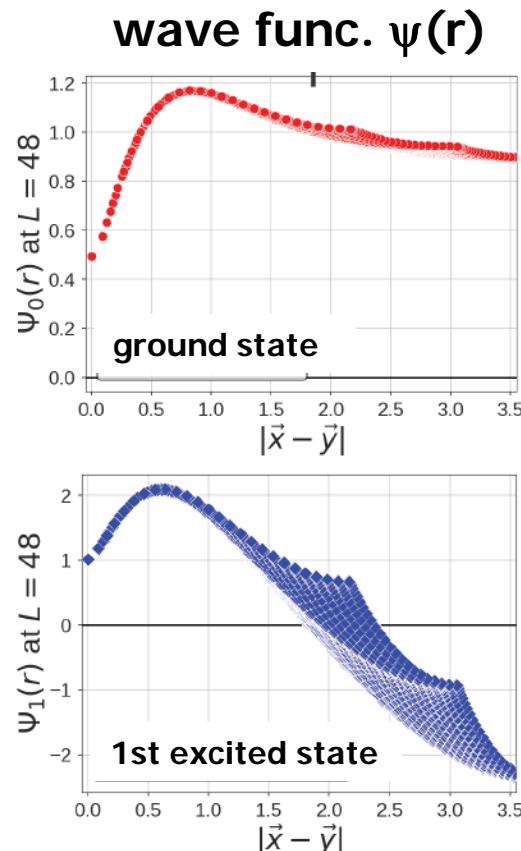
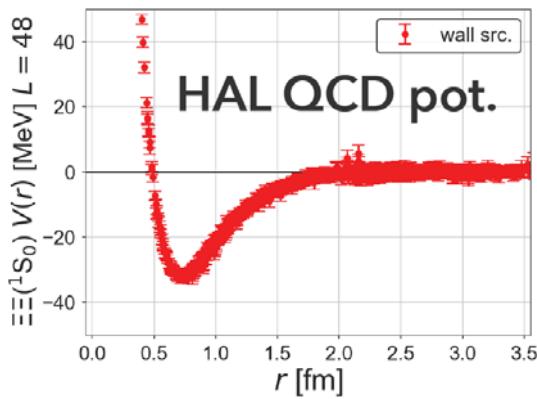
Understand the origin of “pseudo-plateaux”

We are now ready to “predict” the behavior of $m(\text{eff})$ of ΔE at any “t”



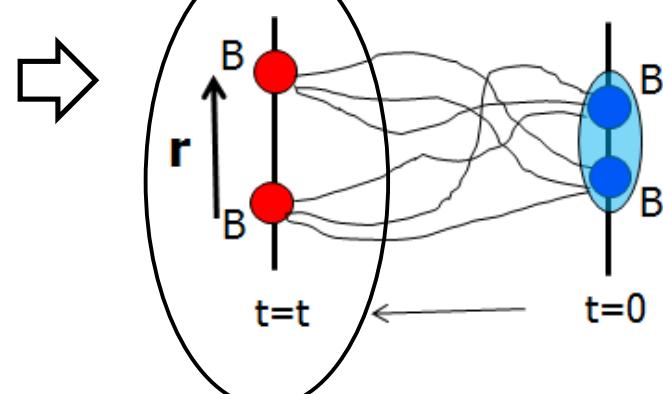
Operator optimized for 2-body system by HAL

- HAL method → HAL pot → 2-body wave func. @ finite V
- 2-body wave func. → optimized operator
 - Applicable for sink and/or src op : Here we apply for sink op
- While utilizing info by HAL, formulation is Luscher's method



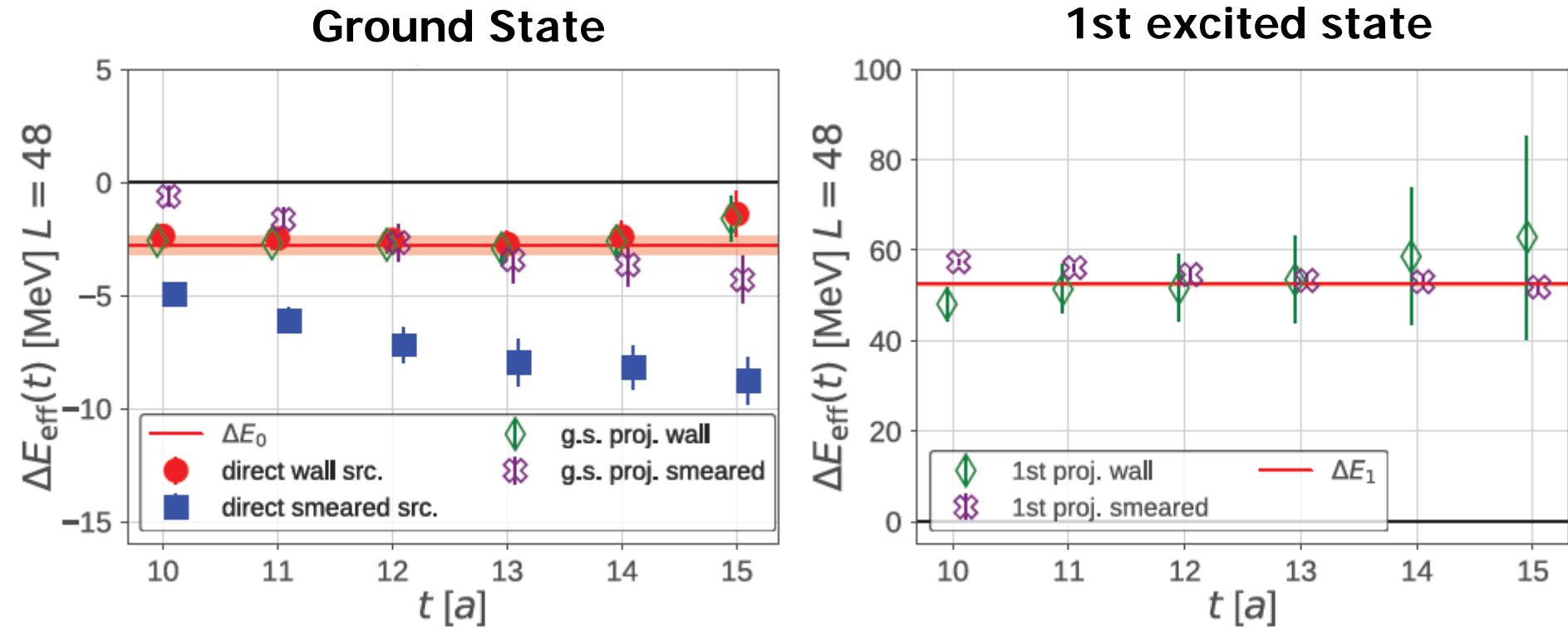
HAL-optimized sink op

$$\mathcal{J}_{\text{sink}}^{2B} = \sum_{\vec{r}} \psi^\dagger(\vec{r}) \sum_{\vec{x}} B(\vec{r} + \vec{x}) B(\vec{x})$$



Effective energy shift ΔE from “HAL-optimized op”

HAL-optimized sink op → projected to each state → “True” plateaux



HAL QCD pot = Luscher's formula w/ proper projection
≠ Direct method w/ naïve plateau fitting

Luscher's formula requires state-projection (a la HAL) or variational calc

Reliability test of LQCD methods

NN @ heavy quark masses

HAL method (HAL) : unbound

Direct method (NPL/CalLat/PACS-CS(Yamazaki et al.)) : bound

Inconsistent!

----- T. Iritani et al. (HAL QCD Coll.) JHEP10(2016)101, PRD96(2017)034521,
PRD99(2019)014514, JHEP03(2019)007 -----

Semi-improved calc w/ Luscher's formula (Mainz2019) : unbound

Variational calc w/ Luscher's formula (CalLat2020) : unbound

Variational calc w/ Luscher's formula (NPL2021) : (unbound)

Issue was essentially settled

- Outline

- Introduction
- Theoretical framework
 - Luscher's formula
 - HAL QCD method
 - Reliability test of LQCD methods
- (Results at heavy quark masses)
- Results near physical quark masses
- Summary / Prospects

- Baryon Forces from LQCD Ishii-Aoki-Hatsuda (2007)
- Exponentially better S/N Ishii et al. (2012)
- Coupled channel systems Aoki et al. (2011,13)

[Theory] = HAL QCD method

Baryon Interactions near the Physical Point

[Hardware]

= K-computer [10PFlops]

- + FX100 [1PFlops] @ RIKEN
- + HA-PACS [1PFlops] @ Tsukuba
- HPCI Field 5 / Post K Priority Issue 9



[Software]

= Unified Contraction Algorithm

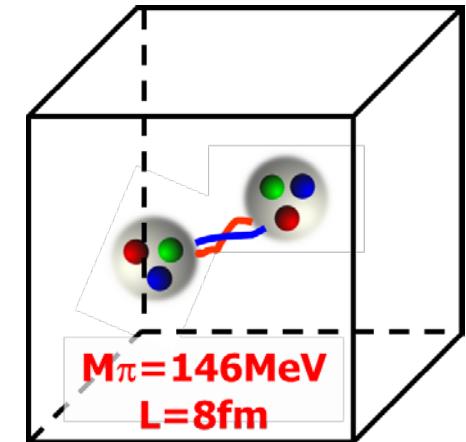
- Exponential speedup Doi-Endres (2013)

$^3\text{H}/^3\text{He}$:	$\times 192$
^4He	:	$\times 20736$
^8Be	:	$\times 10^{11}$

Lattice QCD Setup

- **Nf = 2 + 1 gauge configs**
 - clover fermion + Iwasaki gauge w/ stout smearing
 - $V=(8.1\text{fm})^4$, $a=0.085\text{fm}$ ($1/a = 2.3 \text{ GeV}$)
 - **$m(\pi) \approx 146 \text{ MeV}, m(K) \approx 525 \text{ MeV}$**
 - #traj ≈ 2000 generated

PACS Coll., PoS LAT2015, 075

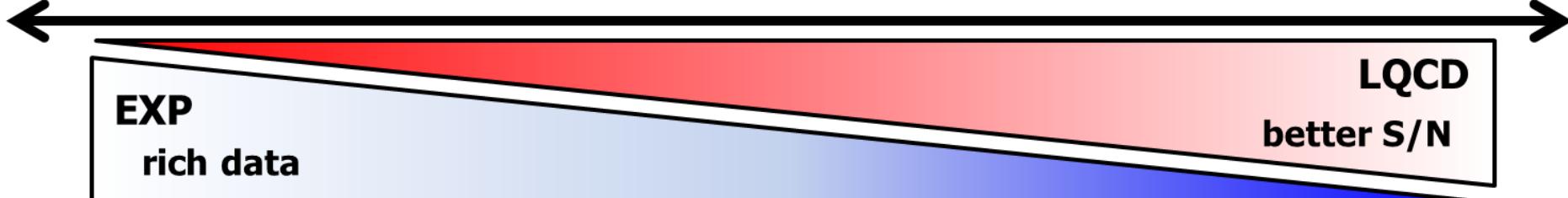


• Measurement

- All of NN/YN/YY for central/tensor forces in $P=(+)$ (S, D-waves)

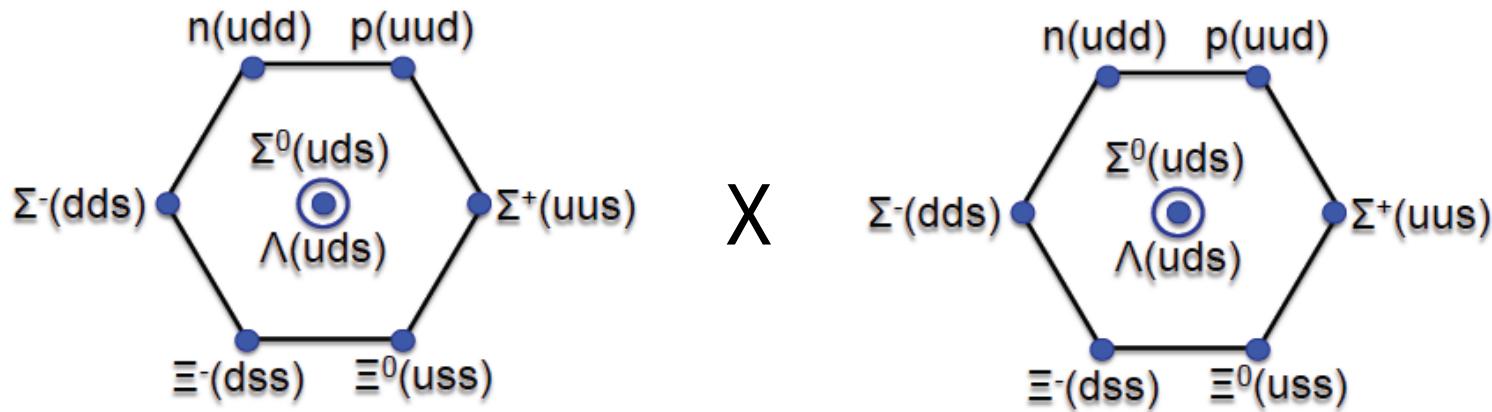
Predictions for Hyperon forces

S=0	S=-1	S=-2	S=-3	S=-4	S=-5	S=-6
NN	NΛ, NΣ	ΛΛ, ΛΣ, ΣΣ, NΞ	ΛΞ, ΣΞ, NΩ	ΞΞ	ΞΩ	ΩΩ



Birds-eye View

classification w/ flavor SU(3)-irrep base



$$8 \times 8 = \underbrace{27}_{\text{symmetric}} + \underbrace{8s + 1}_{\text{anti-symmetric}} + \underbrace{10^* + 10 + 8a}_{\text{NN channel}}$$

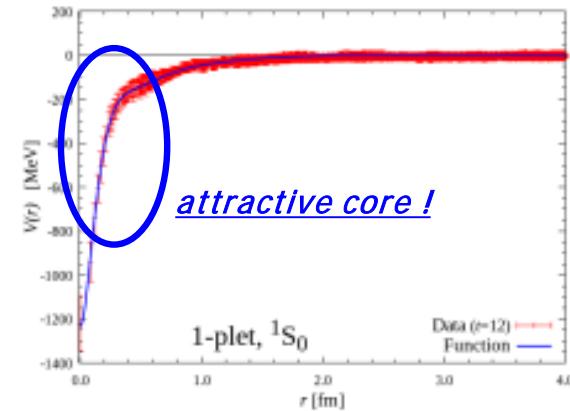
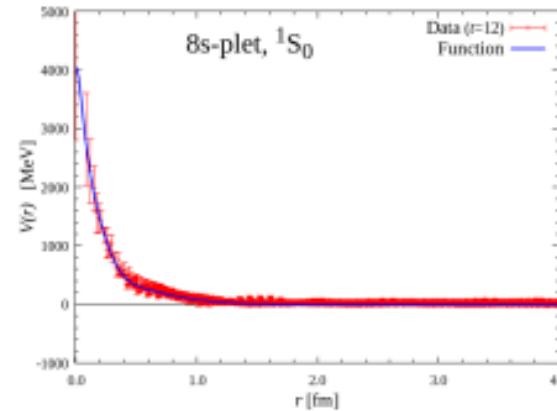
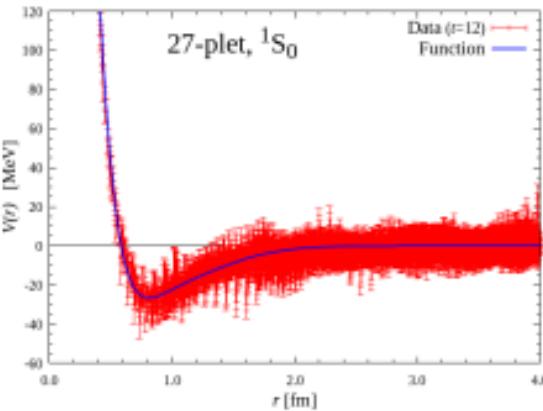
c.f. Exact SU(3) limit LQCD calc @ heavy masses

Diagonal Potentials in SU(3)f-irrep base in S=-2

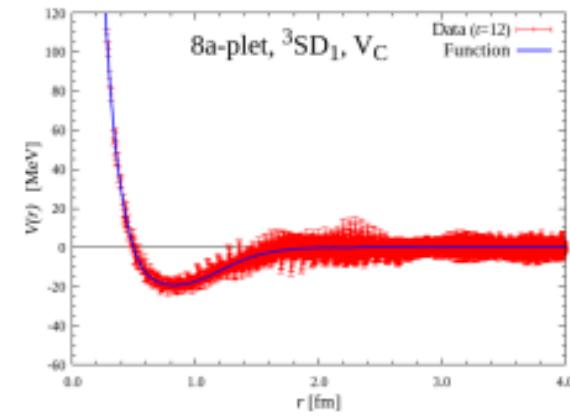
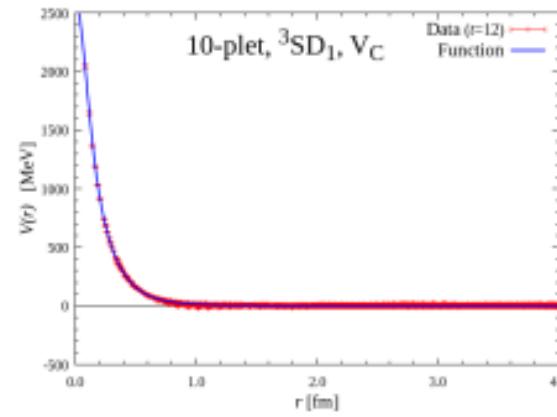
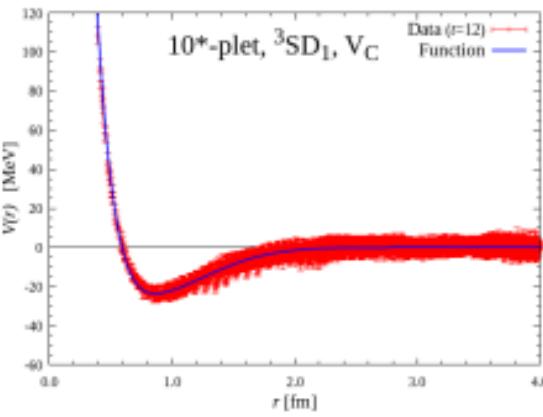
(only) S= -2 can access all irreps
off-diag pot relatively small

T.Inoue (HAL), AIP Conf. Proc. 2130 (2019) 020002

$1S_0$



$3S_1 - 3D_1$



27,10*:
NN-type

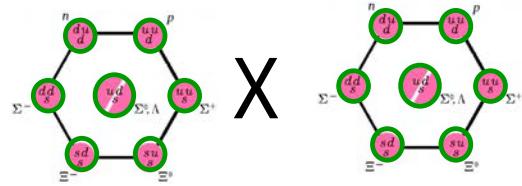
8s,10:
strong repulsive core

1s: deep attractive pocket
8a: weak repulsive core

Quark Pauli repulsion + OGE for short range

M.Oka et al., NPA464(1987)700

Candidates of di-baryons



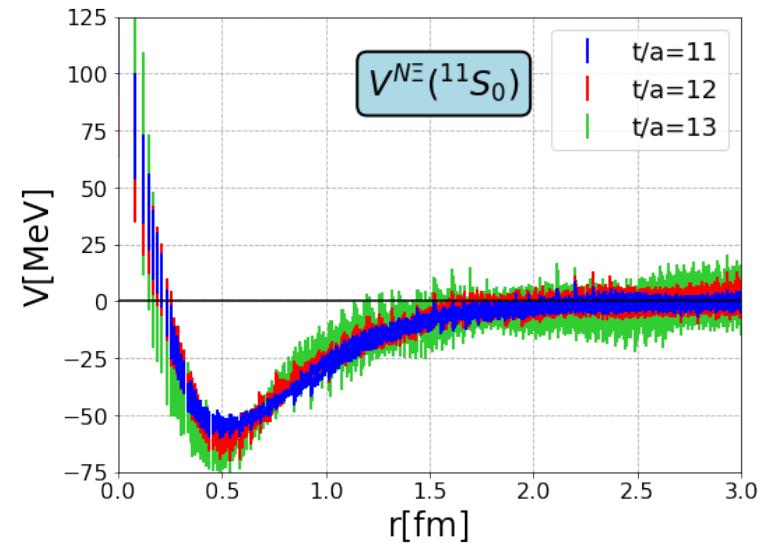
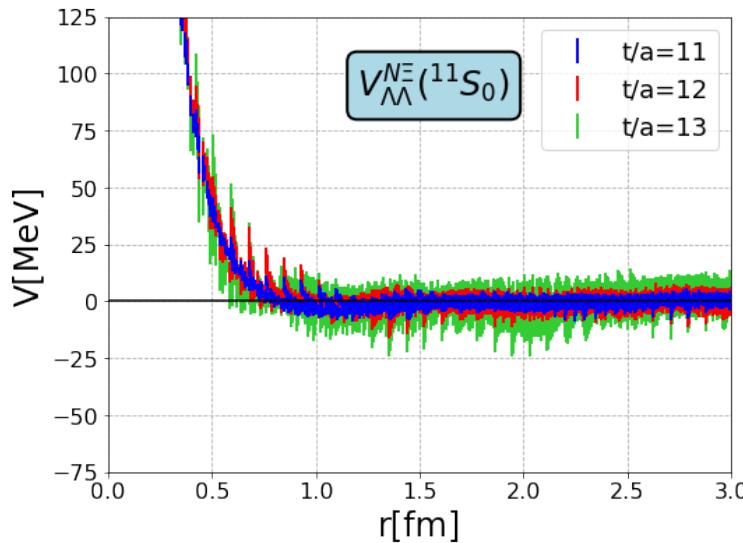
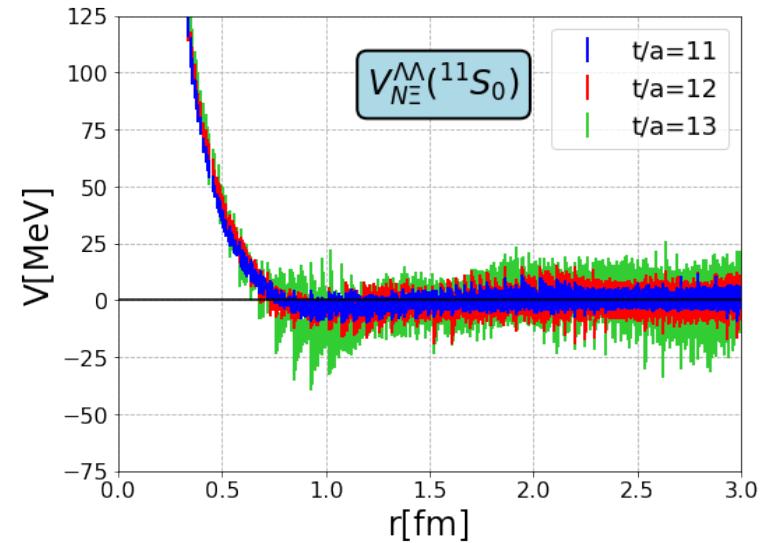
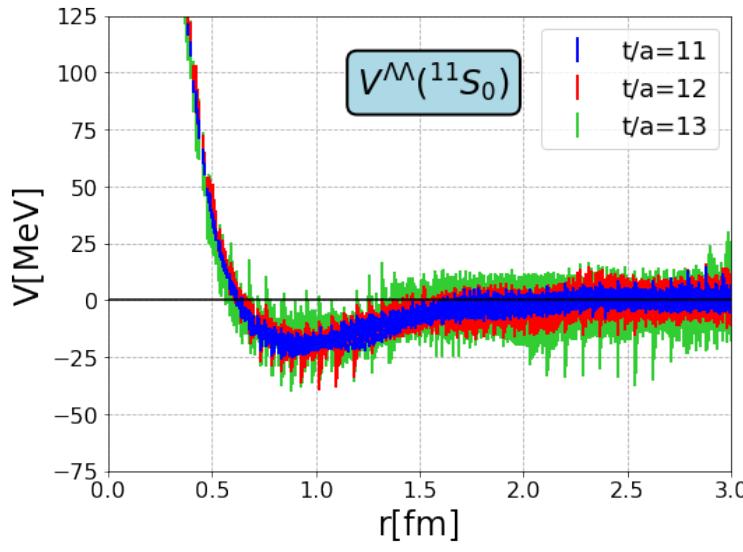
$$8 \times 8 = 27 + 8s + 1 + 10^* + 10 + 8a$$

dineutron, $\Xi\Xi$ etc. H-dibaryon Deuteron
 $(J=0)$ $(J=0)$ $(J=1)$

There may also exist $S = -2$ hypernuclei
relevant to these strong attractions

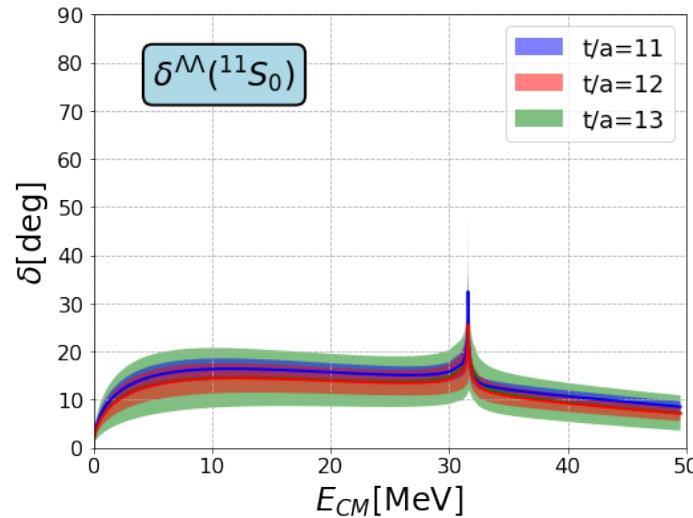
→ Detailed study w/ SU(3) breaking effects
(particle-base)

$\Lambda\Lambda$, $N\Xi$ (effective) 2x2 coupled channel analysis



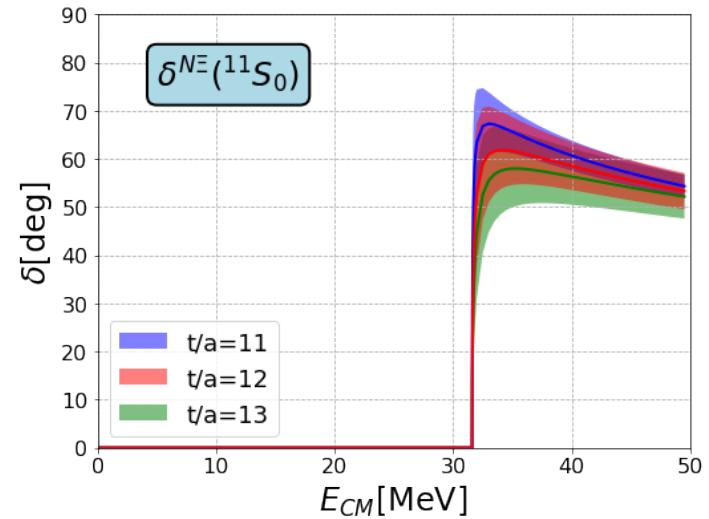
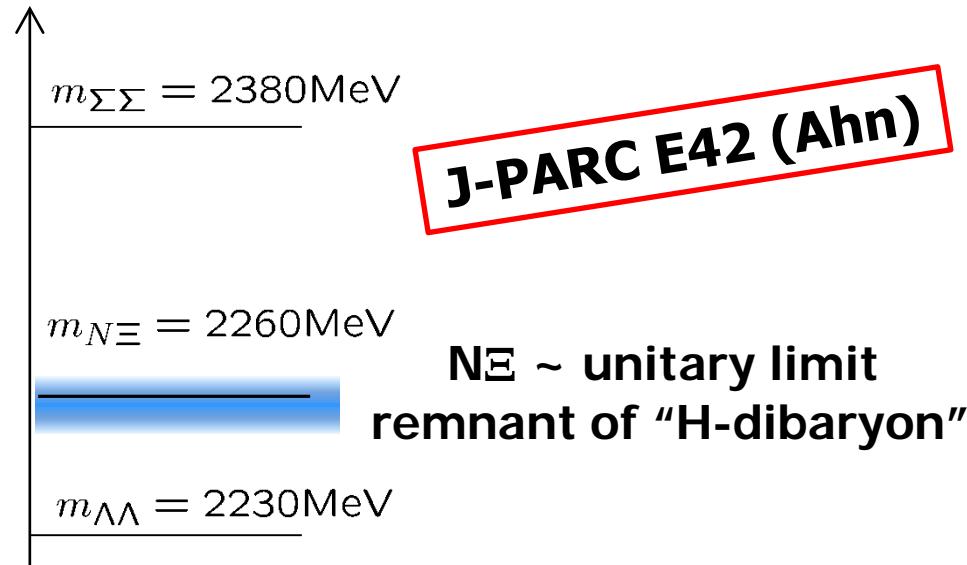
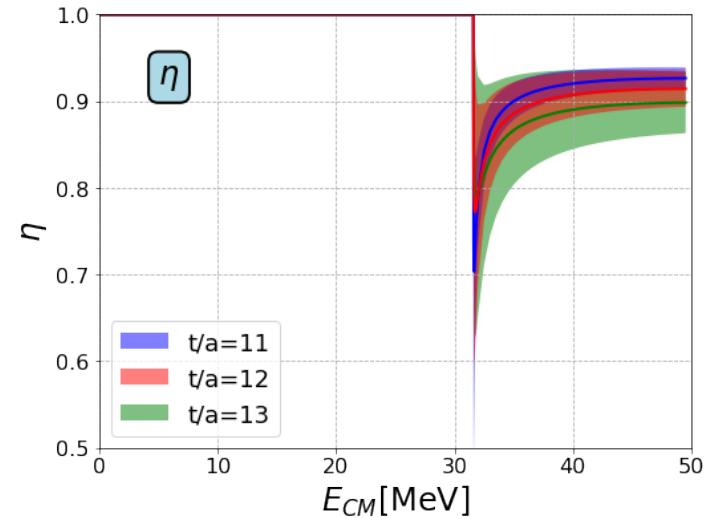
$N\Xi$ ($^{11}S_0$) channel is attractive
 $N\Xi-\Lambda\Lambda$ coupling is small

$\Lambda\Lambda$, $N\Xi$ 2x2 coupled channel analysis



$$a_0 = -0.81(23)(+0.00/-0.13) \text{ [fm]}$$

$$r_{\text{eff}} = 5.47(78)(+0.09/-0.55) \text{ [fm]}$$



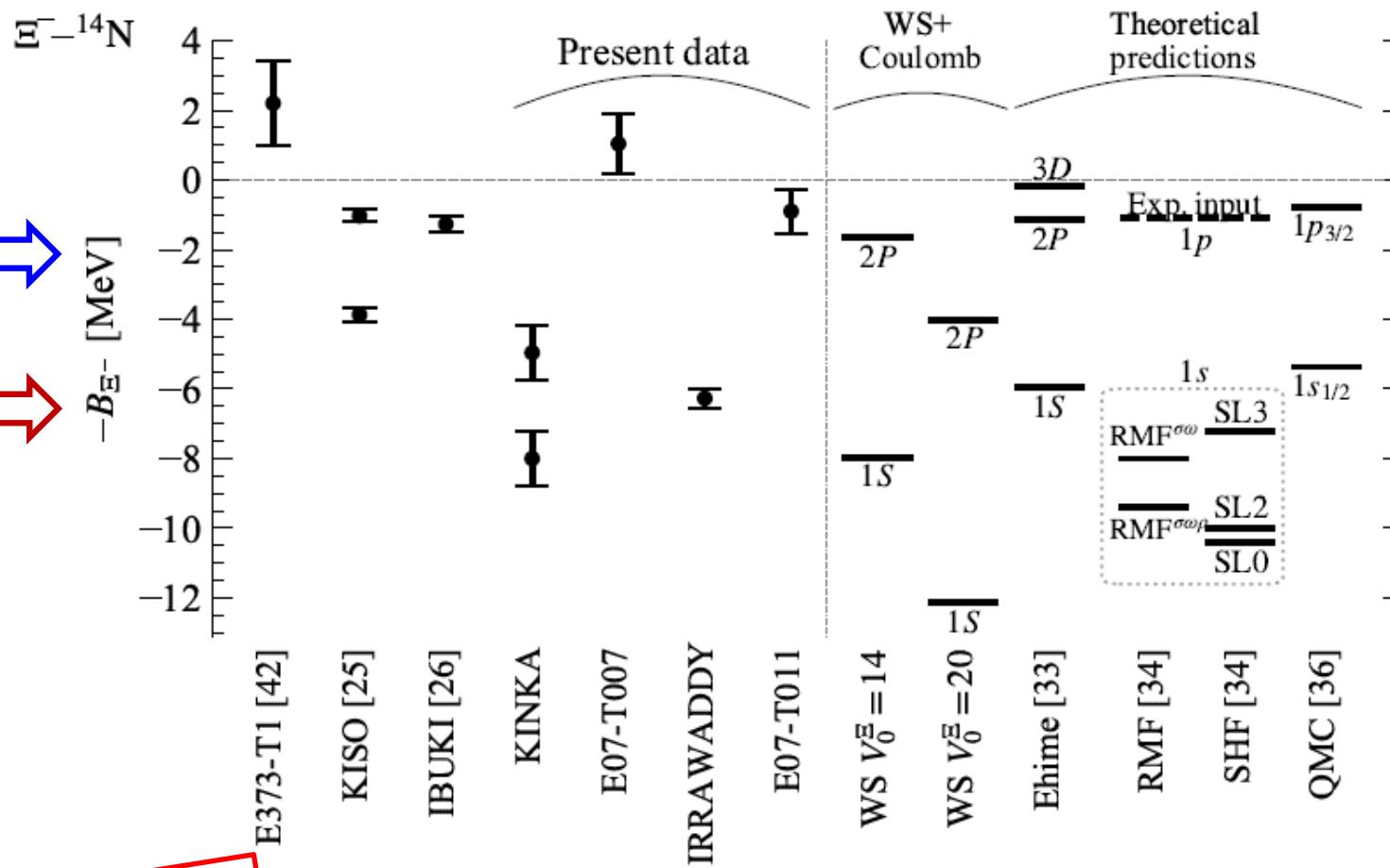
(N.B. $N\Xi = 1 \text{ rep } 50\%, 27 \text{ rep } 30\% \text{ in SU}(3)$)

Recent experimental progress on Ξ -Hypernuclei

Excited state?

Ground state?

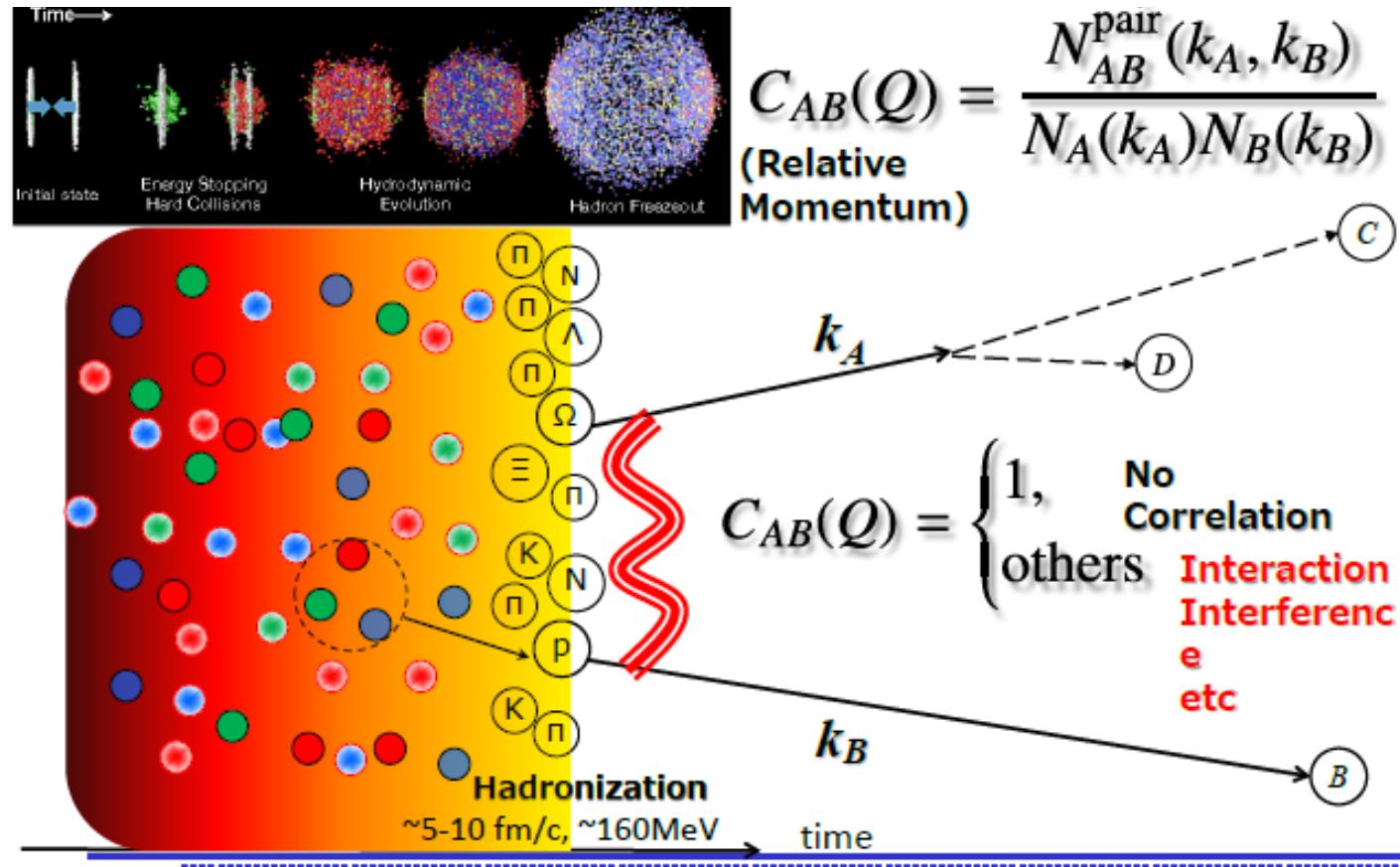
J-PARC E07, E70,
more in HEF-EX



Attractive $N\Xi$ -int well established
Small $N\Xi-\Lambda\Lambda$ coupling indicated

M. Yoshimoto et al.,
PTEP2021, 073D02

Baryon-Baryon correlation in HIC



BB-correlation

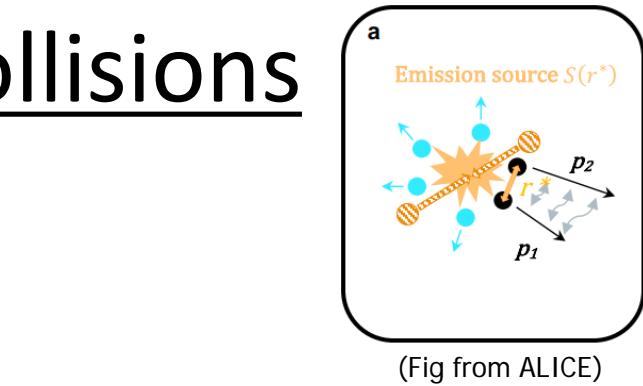
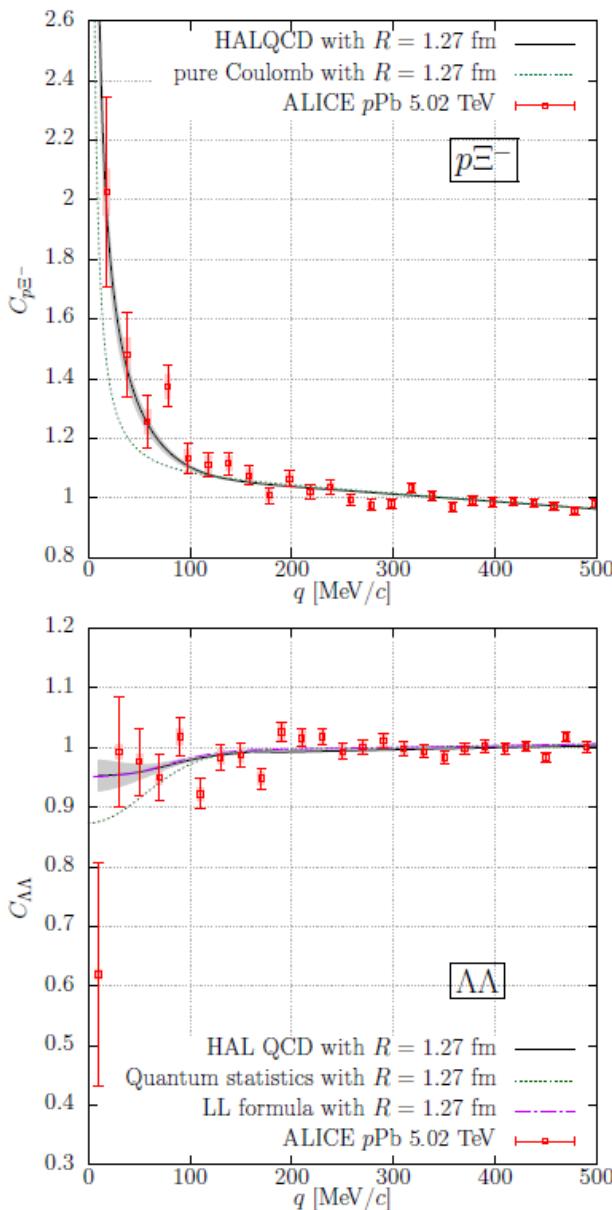
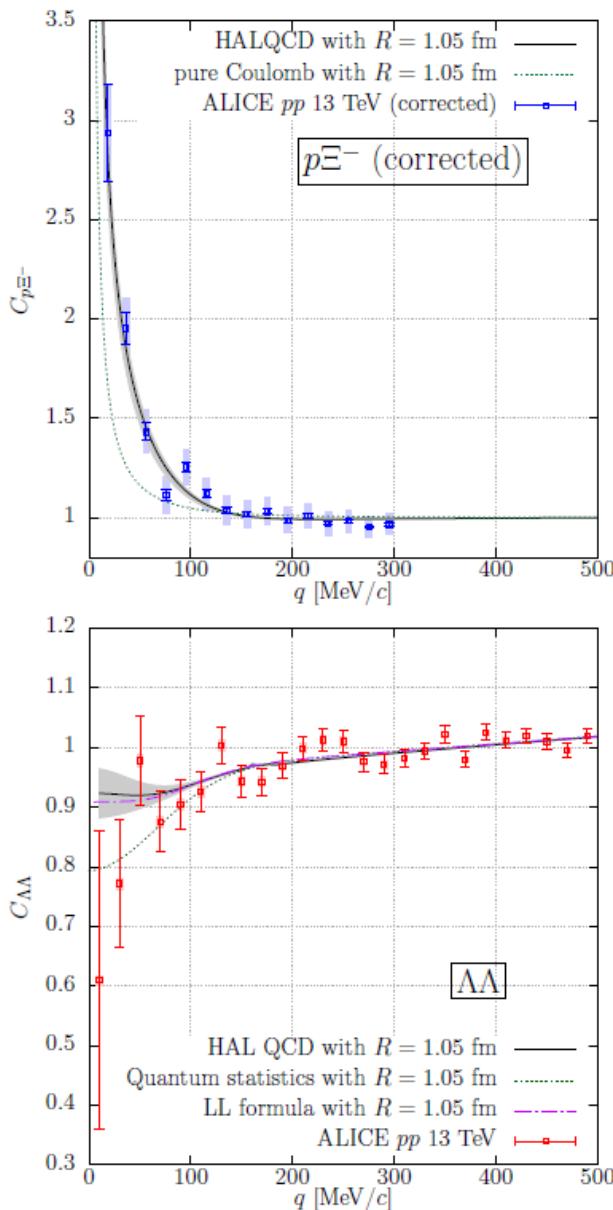
↔ BB- (final state) interaction

Ratio of correlation between small/large source size is useful to mask Coulomb effect

Fig. from K. Morita

K. Morita et al., PRC94(2016)031901
K. Morita et al., PRC102(2020)015201

Femtoscopy from nucleus collisions



← **$p\Xi^-$**

LQCD prediction confirmed
by experiment!

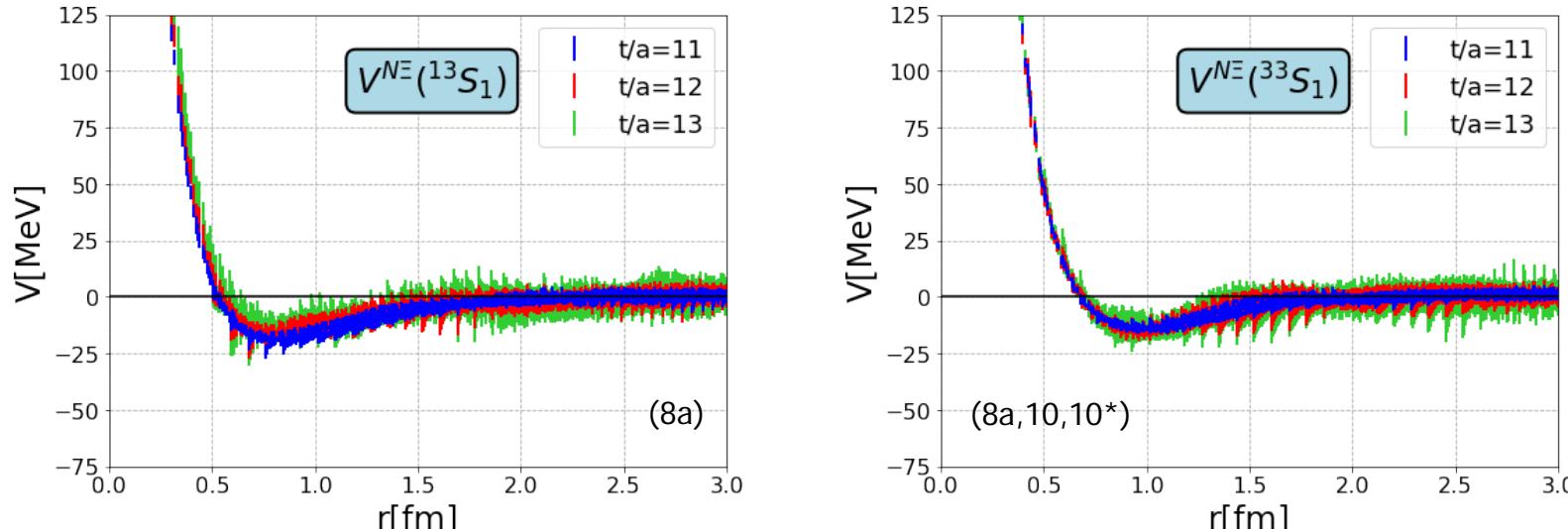
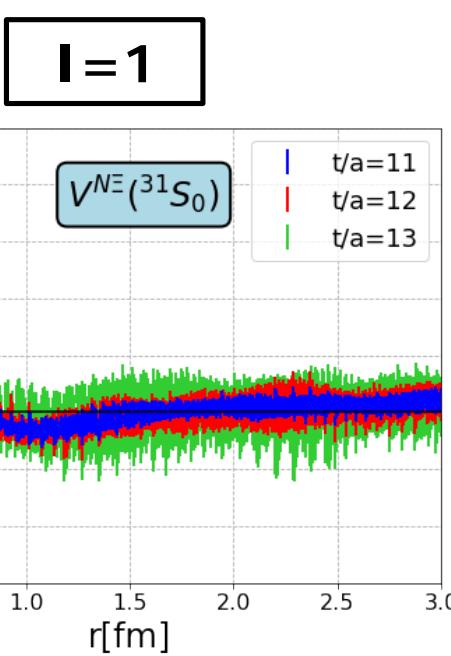
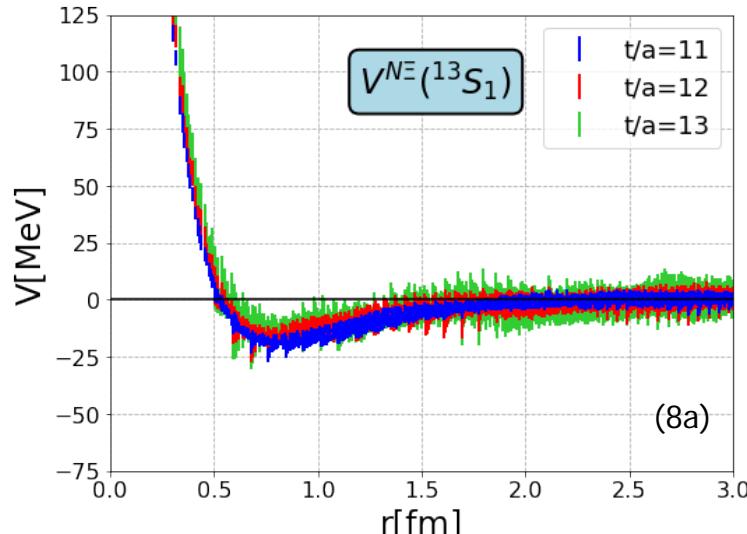
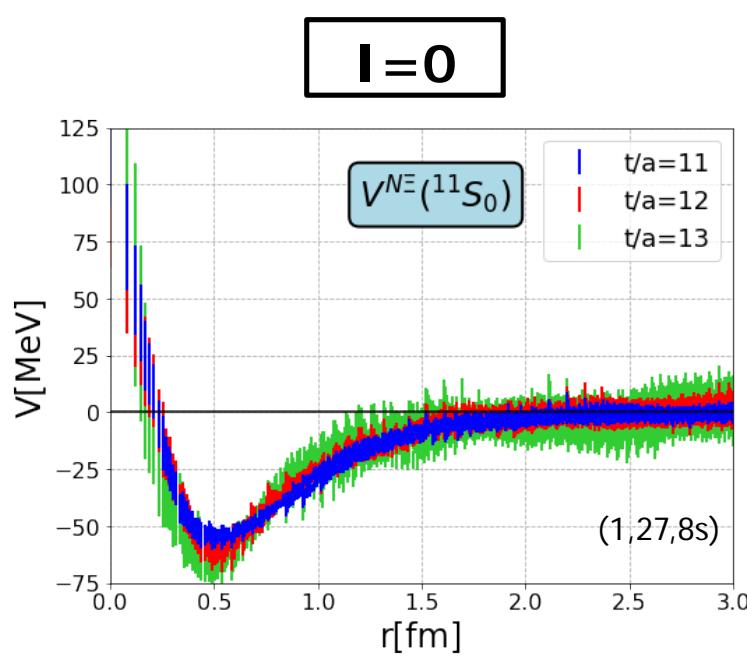
← **$\Lambda\Lambda$**

Y. Kamiya et al.,
PRC105(2022)014915

See also ALICE Coll., PLB797(2019)134822,
PRL123(2019)112002, Nature 588(2020)232

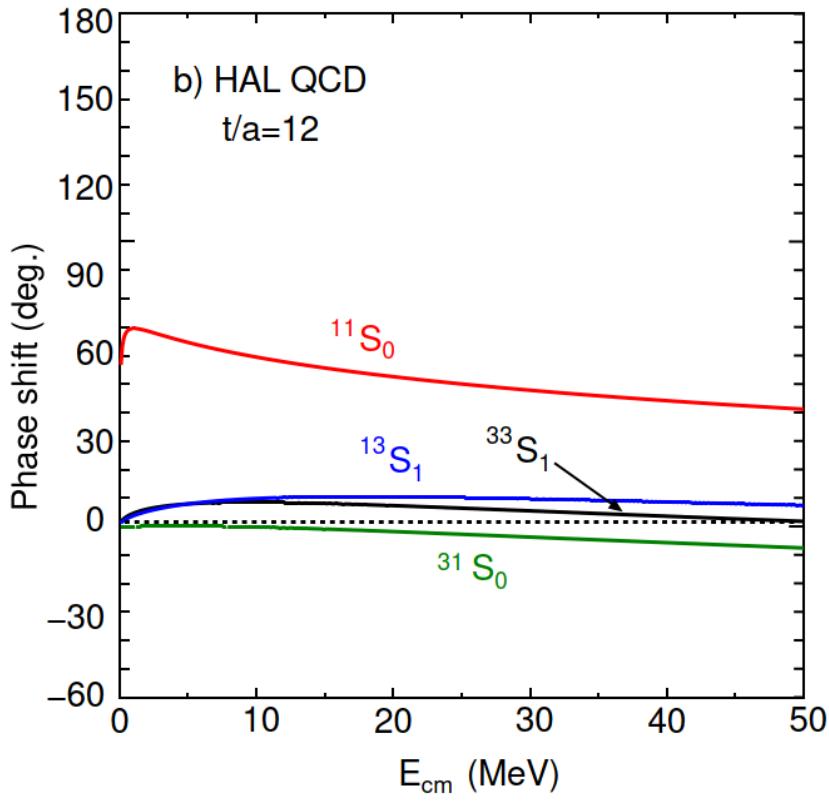
Spin-Isospin dependence of $N\Xi$ potentials

S=0



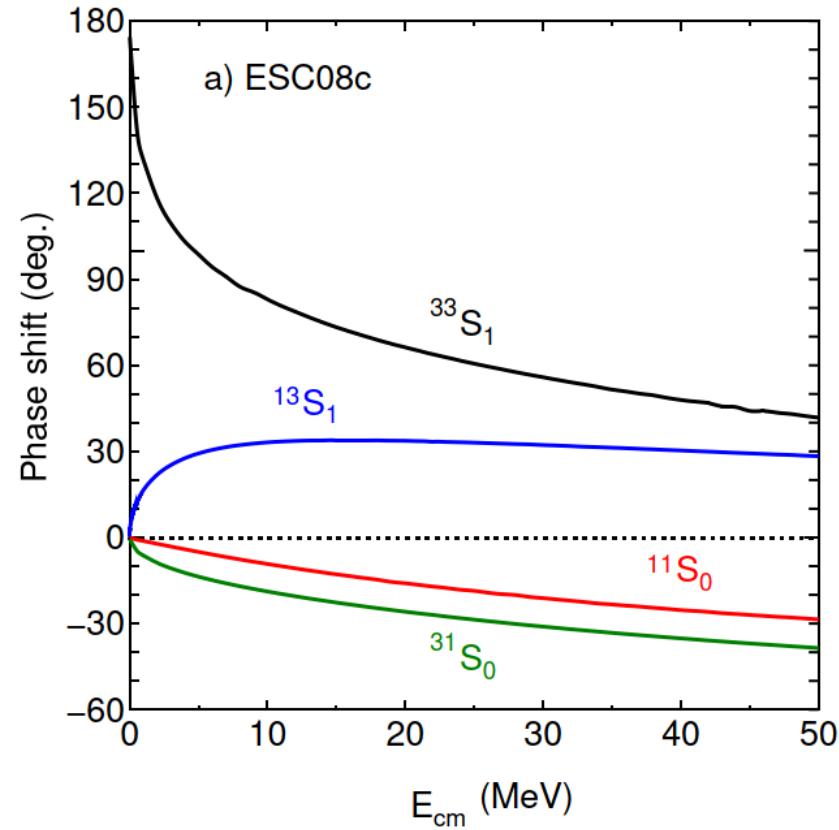
Prediction of NE scattering phase shifts

E. Hiyama et al., PRL124(2020)092501



[HAL]

- I=0, S=0: **attractive**
- I=0, S=1: weakly attractive
- I=1, S=0: weakly repulsive
- I=1, S=1: **weakly attractive**

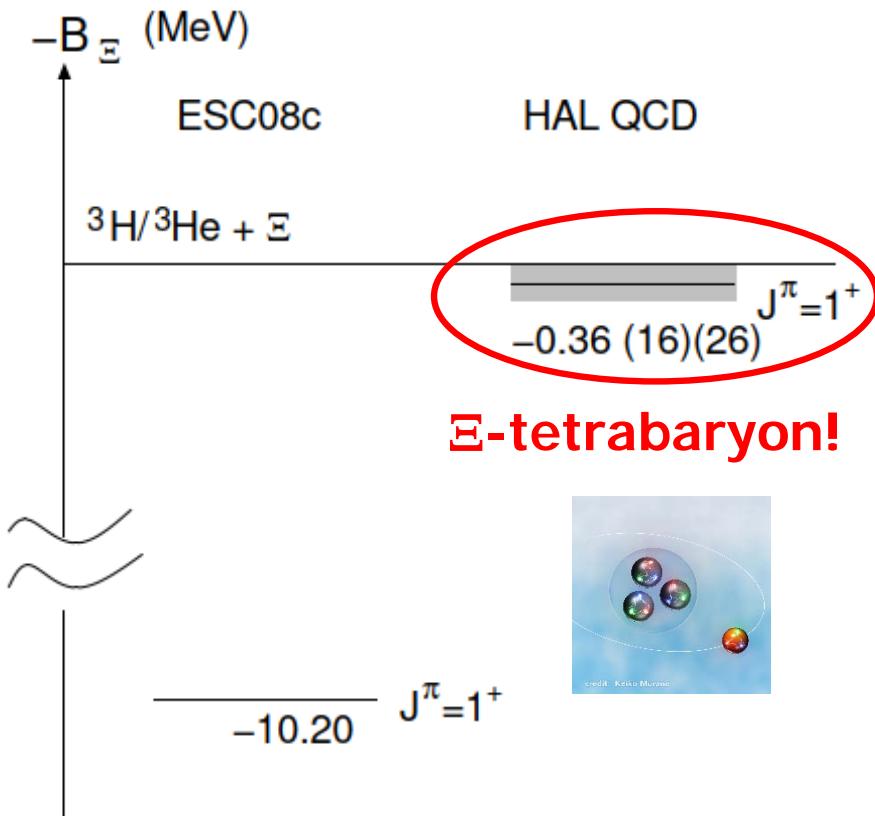


[Nijmegen (ESC08c)]

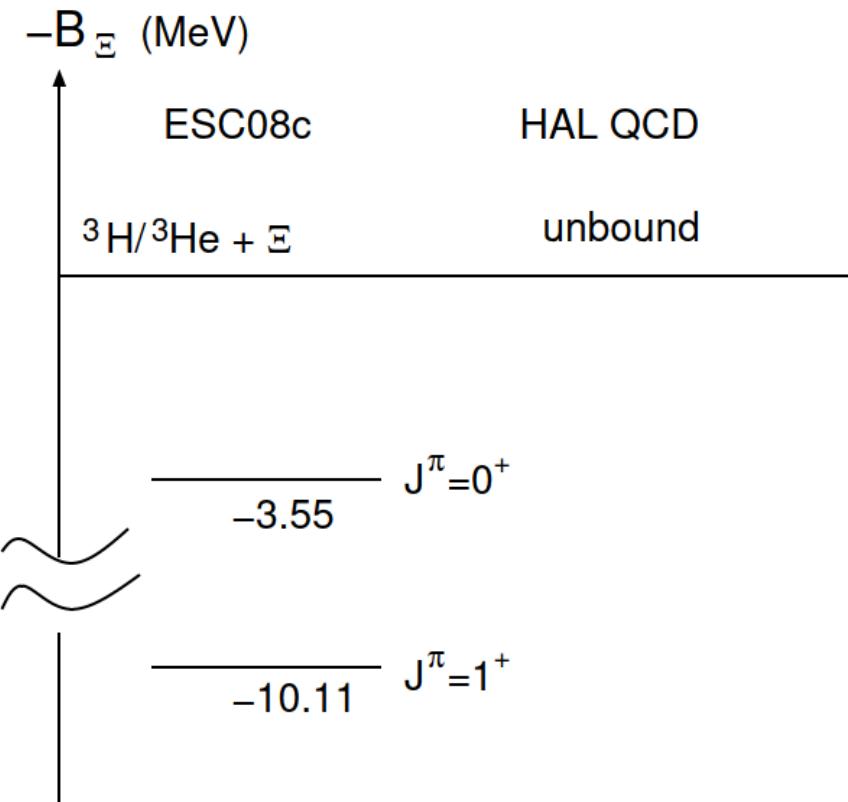
- I=0, S=0: **repulsive**
- I=0, S=1: attractive
- I=1, S=0: repulsive
- I=1, S=1: **strongly attractive**

S= -2 Light Hypernuclei from LQCD

a) NNN Ξ (T=0)



b) NNN Ξ (T=1)

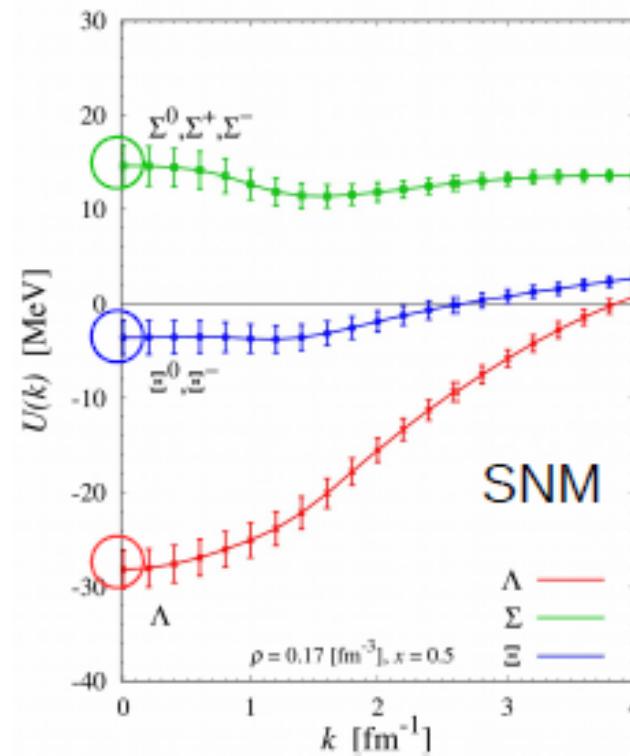
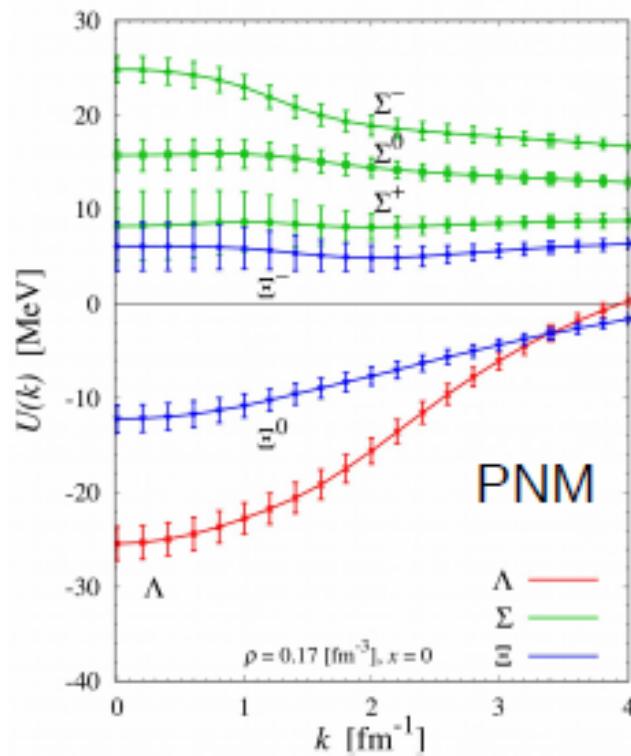


Guiding experiments

HIC@LHC, J-PARC, ...

E. Hiyama et al., PRL124(2020)092501

“Super-super heavy nuclei”: Dense matter from LQCD Hyperon single-particle potential



@ $\rho = 0.17 \text{ fm}^{-3}$

Preliminary

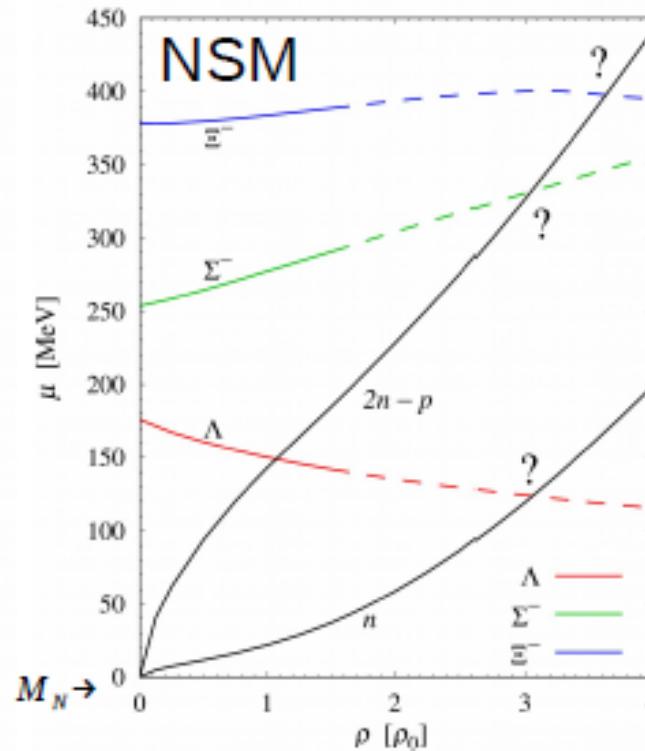
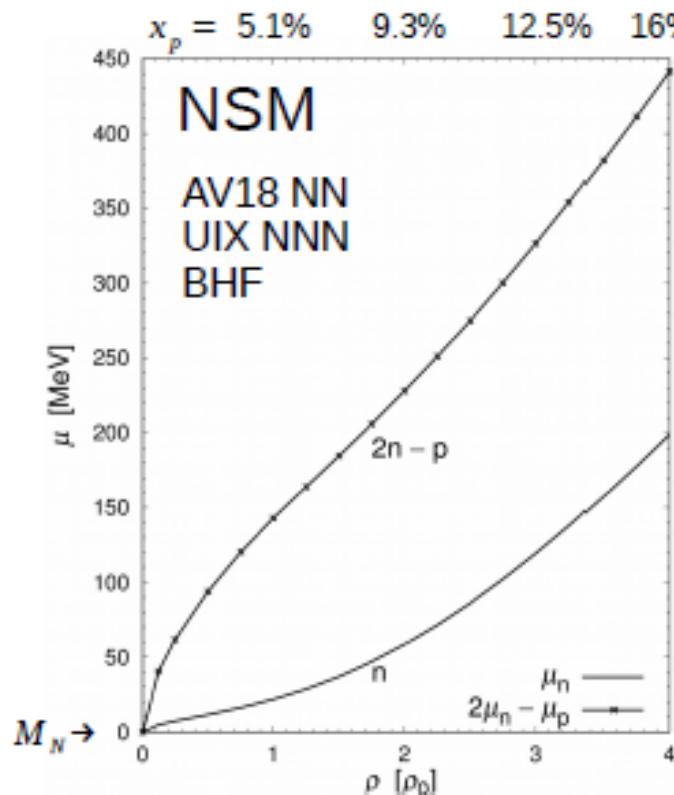
T. Inoue (HAL Coll.)
PoS INPC2016, 277

- obtained by using YN, YY S-wave forces from QCD.
 - Results are compatible with experimental suggestion.

$$U_{\Lambda}^{\text{Exp}}(0) \simeq -30, \quad U_{\Xi}(0)^{\text{Exp}} \simeq -10?, \quad U_{\Sigma}^{\text{Exp}}(0) \geq +20? \quad [\text{MeV}]$$

attraction attraction small repulsion

Hyperon onset in NSM (just for fun)



Preliminary

- Result indicate Λ , Σ^- , Ξ^- appear around $\rho = 3.0 - 4.0 \rho_0$
- However,
 - $YN^{L=1,2,\dots}$ and YNN force could be important
 - We may need to compare with more SBF

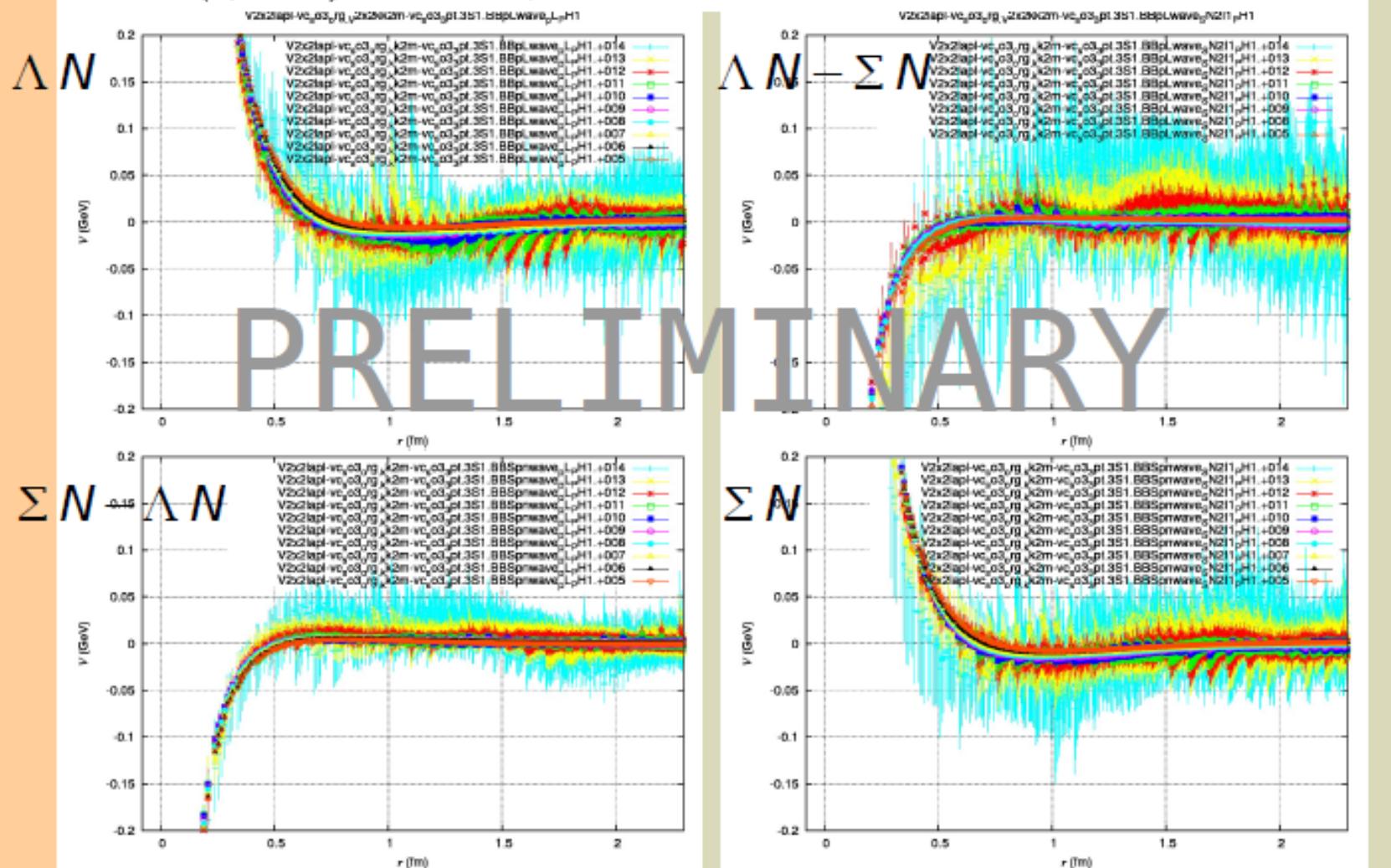
[Challenges]
Precision for $|S| \leq 1$
3-baryon forces
P-wave/LS forces

$\Lambda N - \Sigma N$ Vc potential in $^3S_1 - ^3D_1$

H. Nemura
(Lat2021)

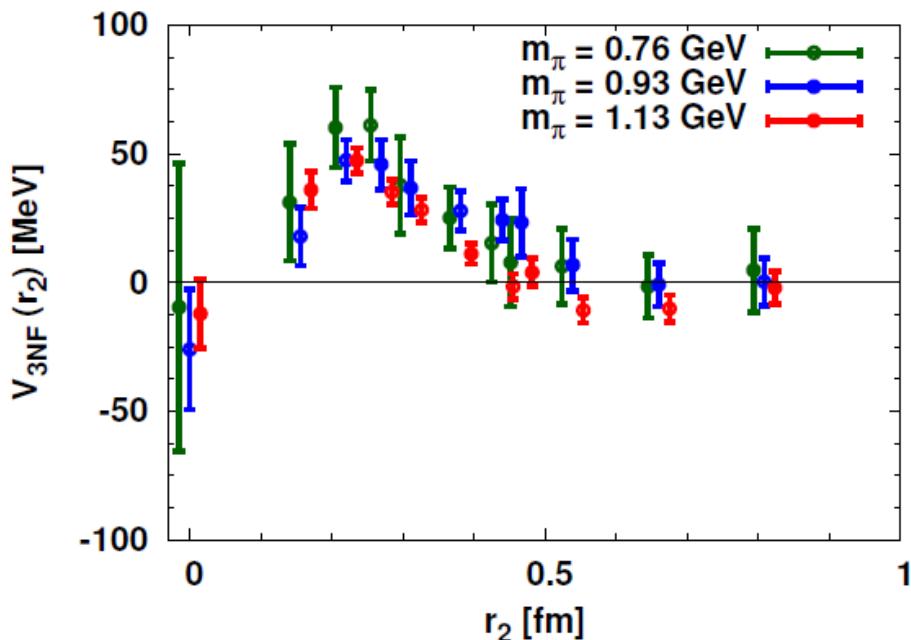
Very preliminary result of LN potential at the physical point

$$\left(\frac{\nabla^2}{2\mu} - \frac{\partial}{\partial t} \right) R(\vec{r}, t) = \int d^3r' U(\vec{r}, \vec{r}') R(\vec{r}', t) + O(k^4) = V_{\text{LO}}(\vec{r}) R(\vec{r}, t) + \dots (8)$$



3N-forces (3NF)

Nf=2, $m\pi=0.76-1.1$ GeV

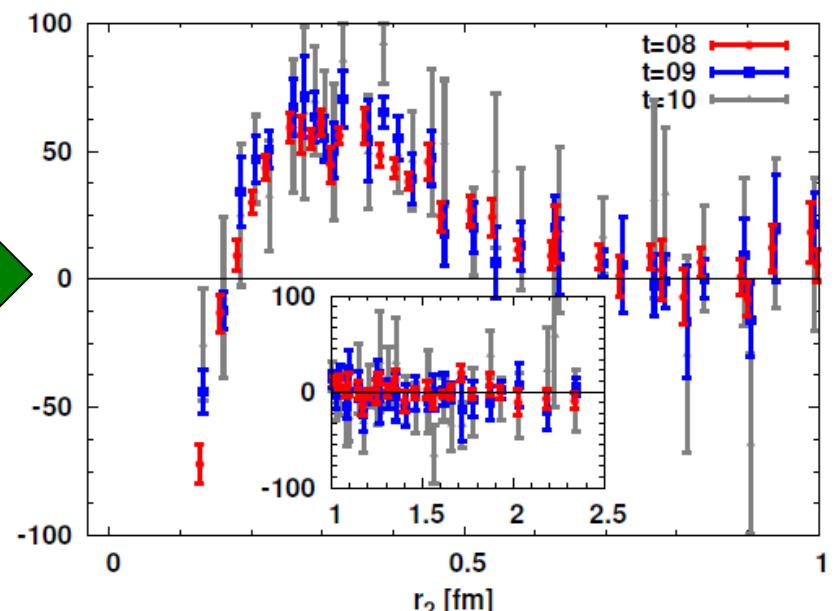


Triton channel

r_2

r_2

Nf=2+1, $m\pi=0.51$ GeV

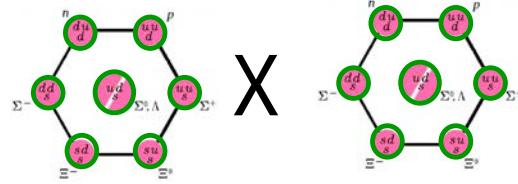


Magnitude of 3NF is similar for all masses

Range of 3NF tend to be enlarged for $m(\pi)=0.5\text{GeV}$

Next challenge: **Calc of P-wave 2BF** : better subtraction of 2BF in 3-body systems
YNN (w/o or w/ P-wave 2BF) : gauge conf generation on Fugaku

Candidates of di-baryons



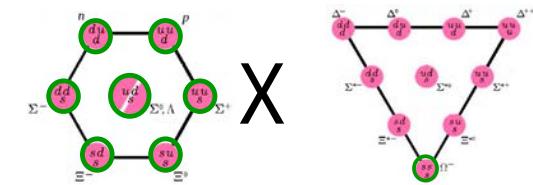
X

$$8 \times 8 = 27 + 8s + 1 + 10^* + 10 + 8s$$

dineutron, $\Xi\Xi$ etc.

H-dibaryon ($J=0$)

Deuteron ($J=1$)

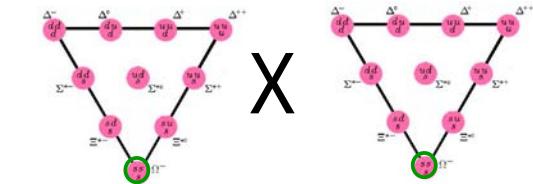


X

$$8 \times 10 = 35 + 8 + 10 + 27$$

$N\Omega$ ($J=2$)

Goldman et al. ('87)
Oka ('88)



X

$$10 \times 10 = 28 + 27 + 10^* + 35$$

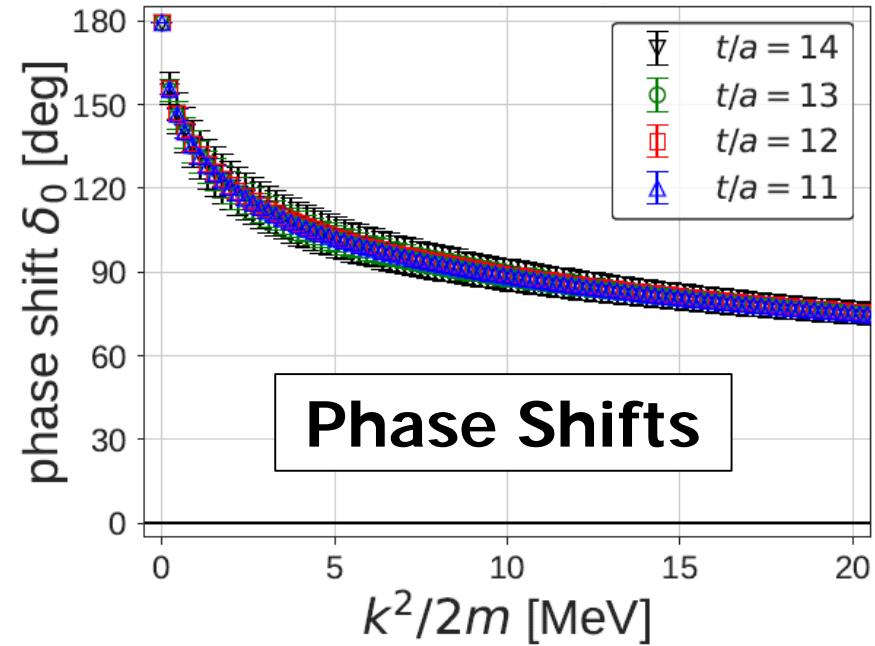
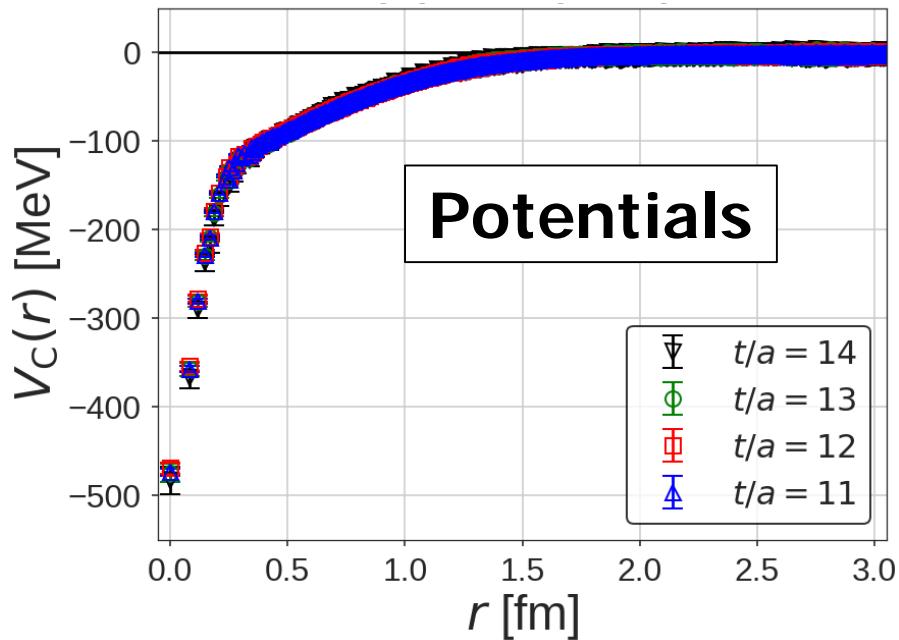
$\Omega\Omega$ ($J=0$)

$\Delta\Delta$ ($J=3$)

Zhang et al. ('97)

Dyson-Xuong ('64)
Kamae-Fujita ('77)
Oka-Yazaki ('80)

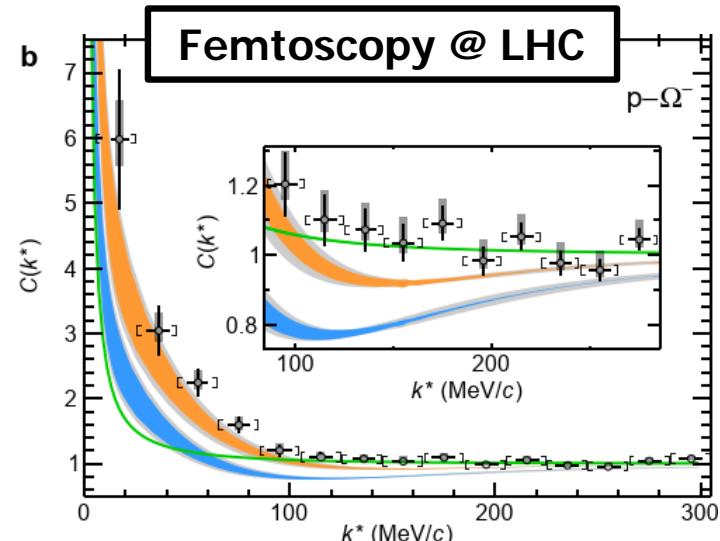
$\text{N}\Omega$ system (${}^5\text{S}_2$)



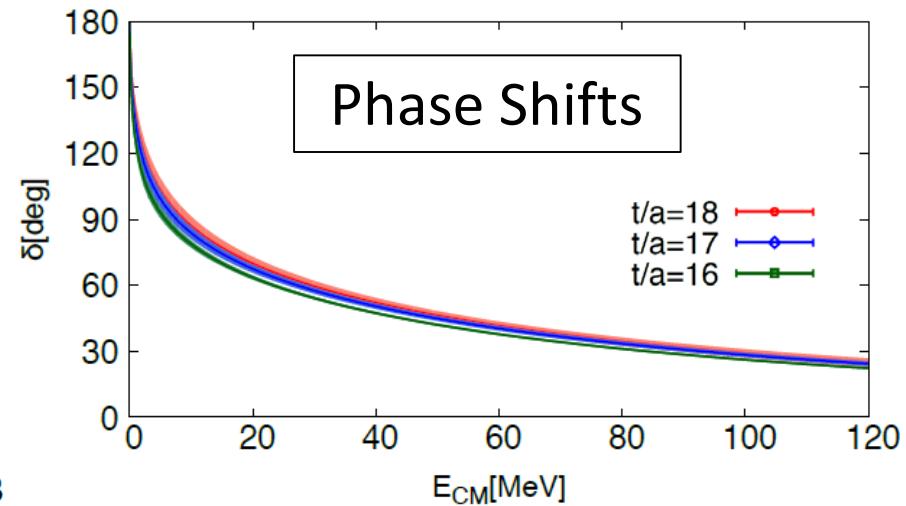
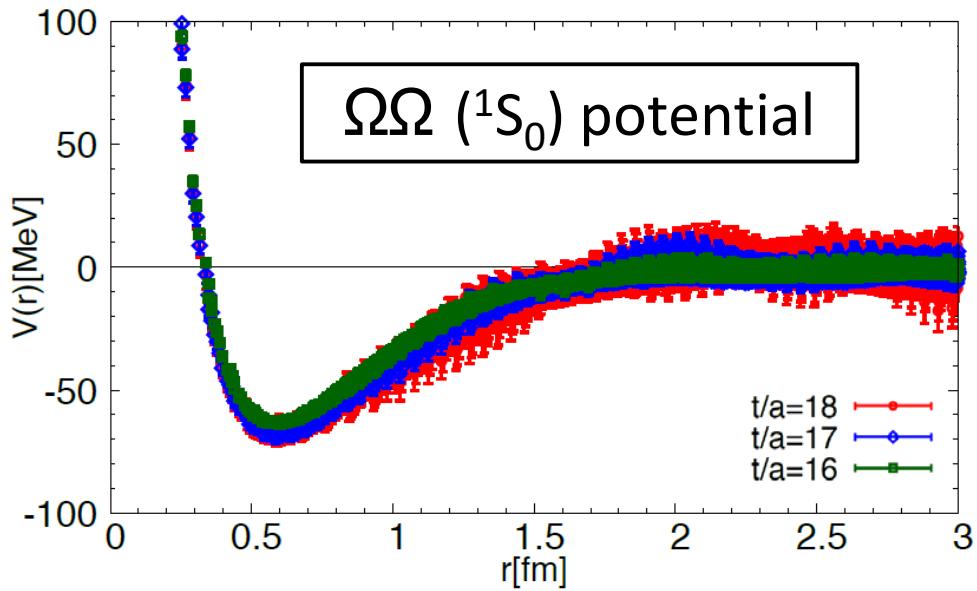
(Quasi) Bound state
[~ Unitary limit]

$$B_{N\Omega} = 1.54(0.30)(^{+0.04}_{-0.10}) \text{ MeV}$$

$$B_{p\Omega^-} = 2.46(0.34)(^{+0.04}_{-0.11}) \text{ MeV}$$

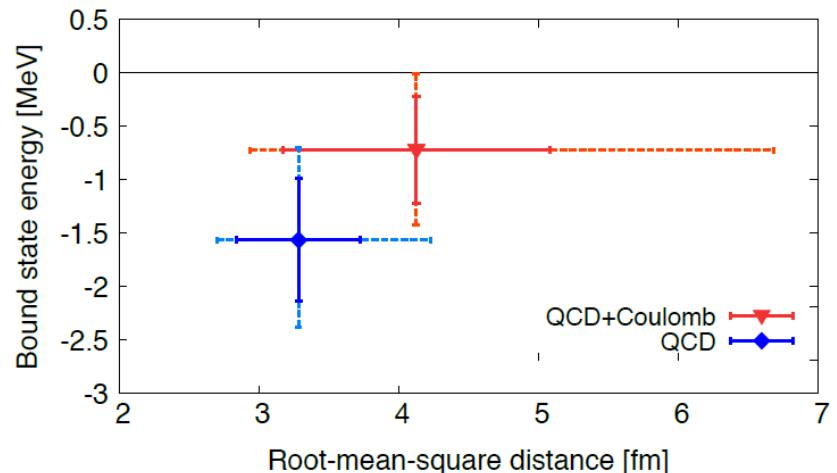


"Most Strange" Dibaryon : $\Omega\Omega$

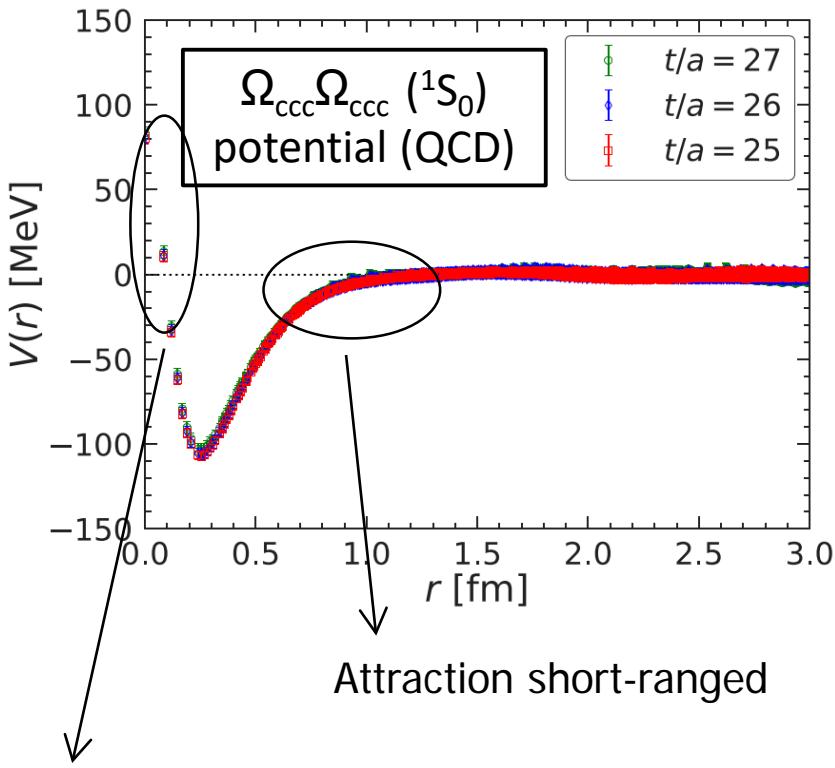


"di-Omega"
[~ Unitary limit]

could be searched in LHC RUN3

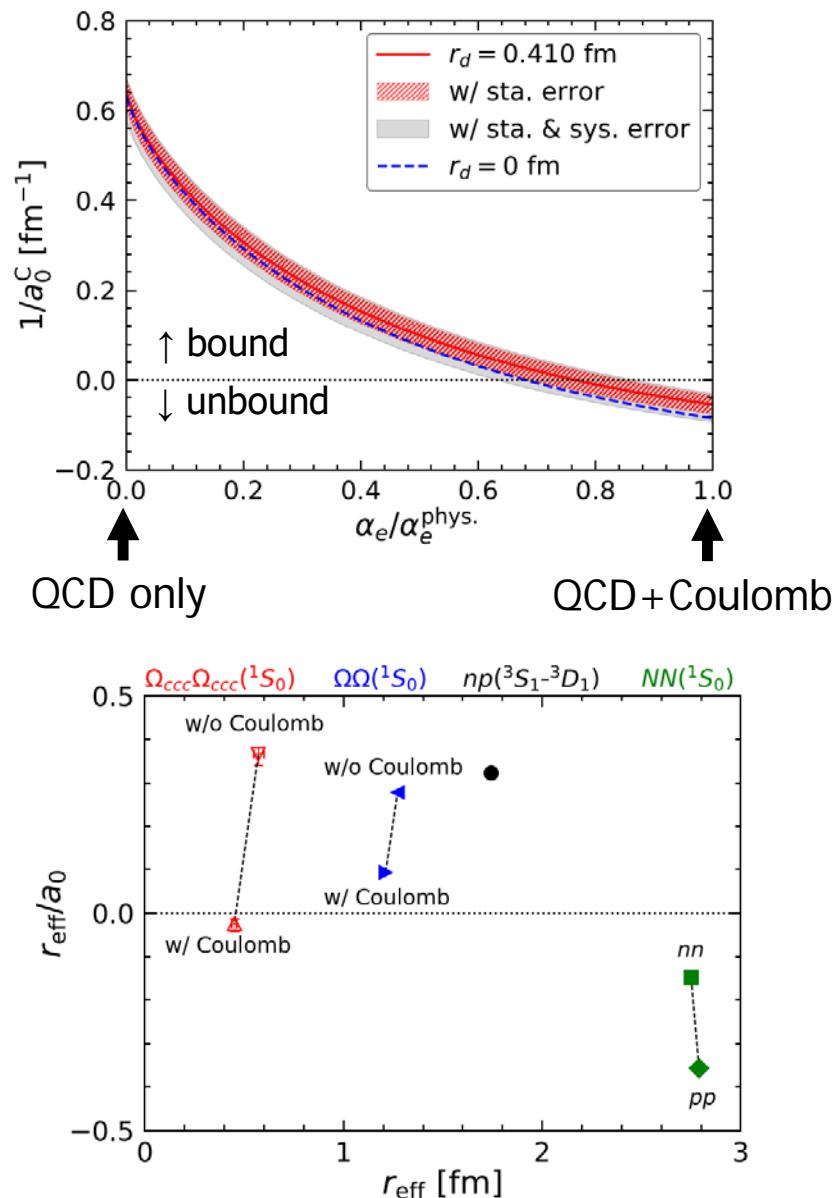


"Most Charming" Dibaryon: $\Omega_{ccc}\Omega_{ccc}$



Color-magnetic int suppressed by $(m_s/m_c)^2 \sim 0.1$

Dibaryon closest to the unitarity



Hadron Interactions from Lattice QCD

Where are we going?

Challenges, new development and future prospects

w/ new supercomputer “Fugaku”

New calc on Fugaku: physical point simulation!

Fugaku (富岳) : 440 PFlops

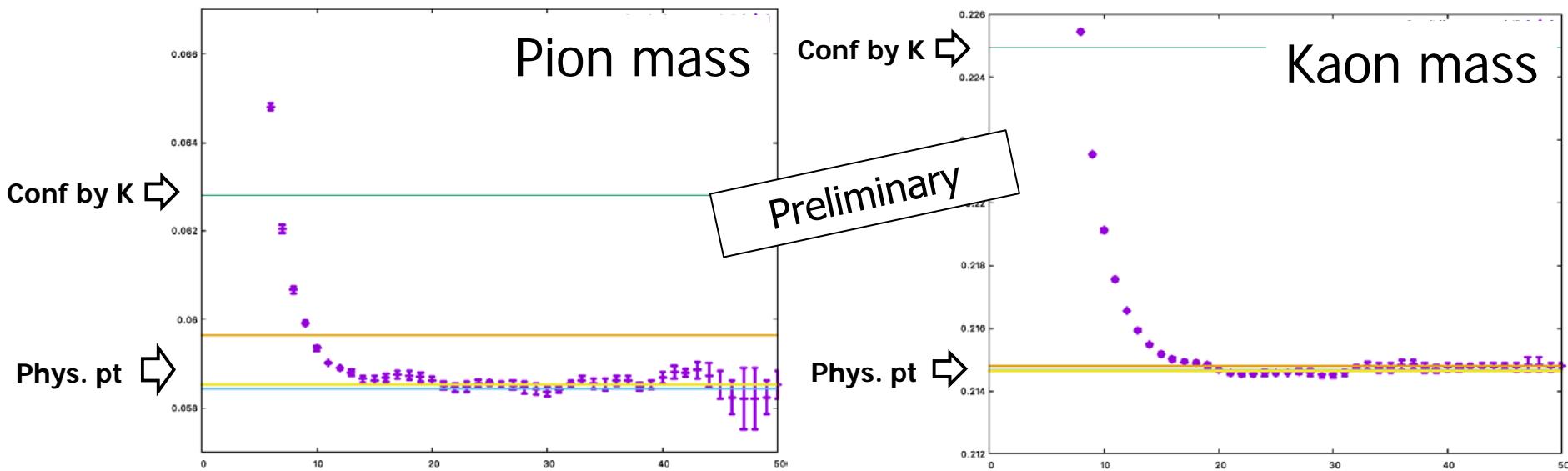
Successor of K-computer
Developed since 2014-,
Full operation since 2021-



Codesign of hardware/software
(LQCD was one of 9 targets)

Fastest in the world! (2020-)

Our Efficiency = ~17% (w/ naïve double prec count)



[E. Itou+]

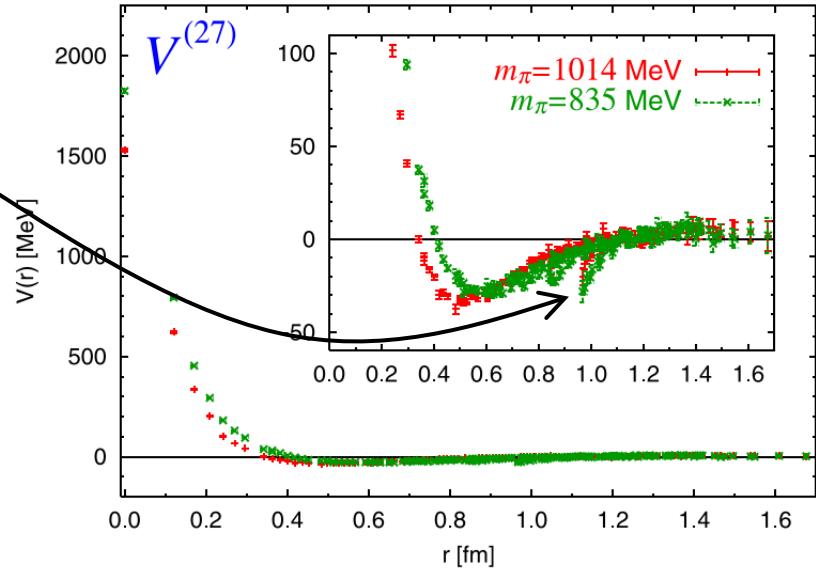
We are on the physical point!

S/N improvement by partial wave decomposition

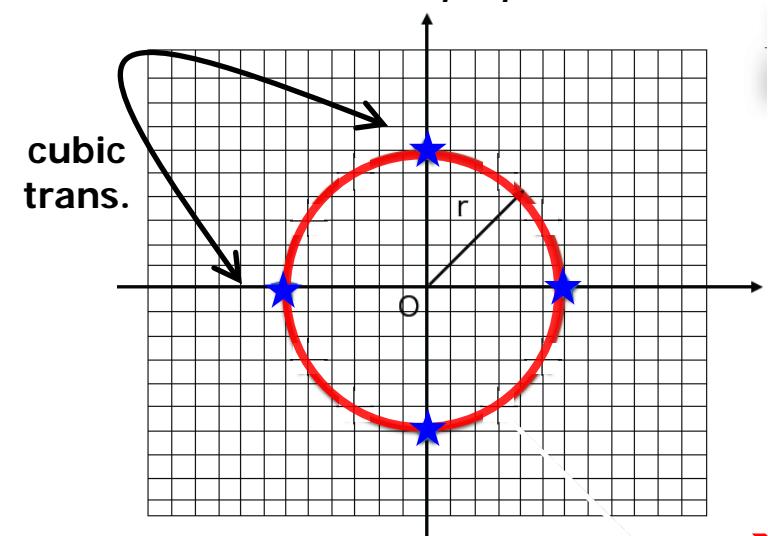
T. Miyamoto et al. (HAL Coll.), PRD101(2020)074514

- Cubic group irrep on lattice
→ sys err by mixing of partial waves
- Partial wave decomposition
can cure the problem

C. W. Misner, Class. Quantum Grav. 21 (2004) S243-S247

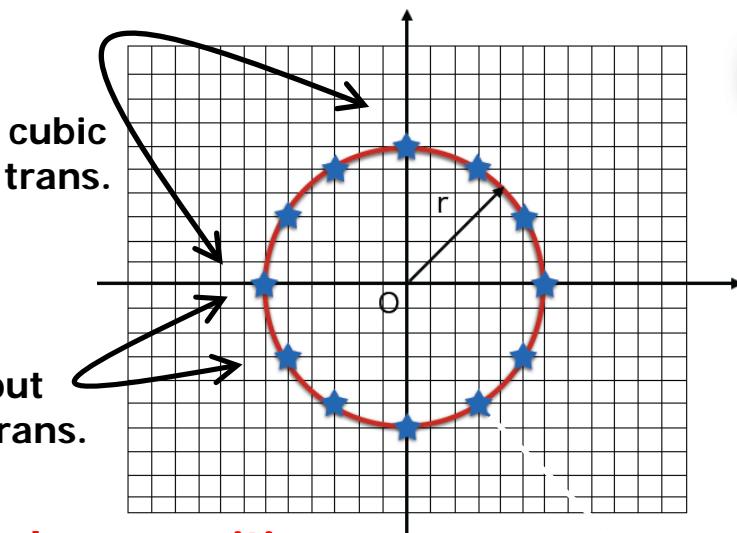


A_1^+ -irrep in cubic group
 $\longleftrightarrow L=0, 4, \dots$



We have
more info !

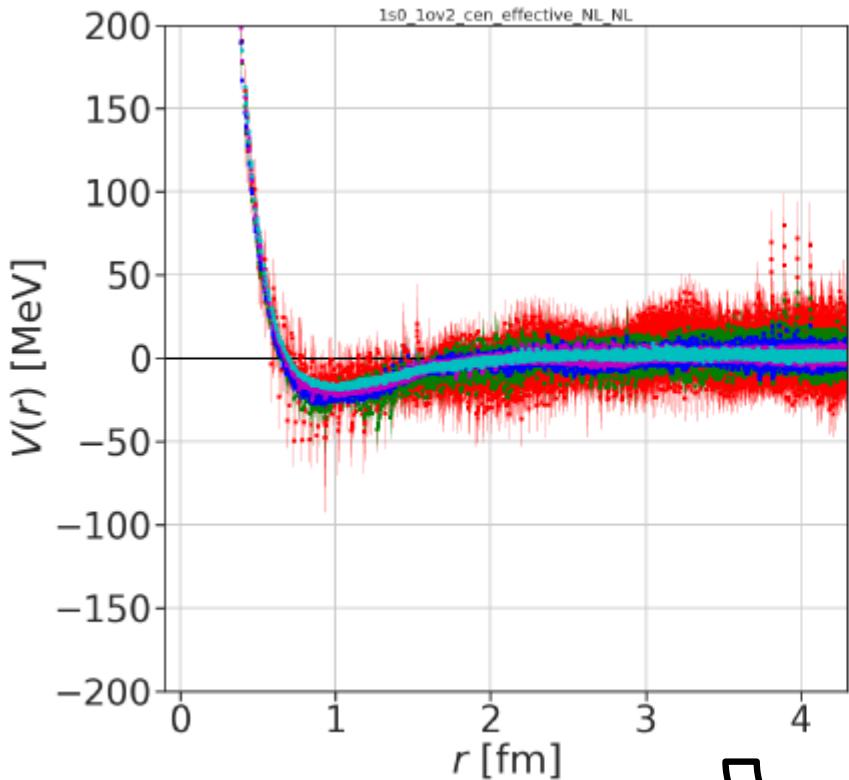
Same r but
not cubic trans.



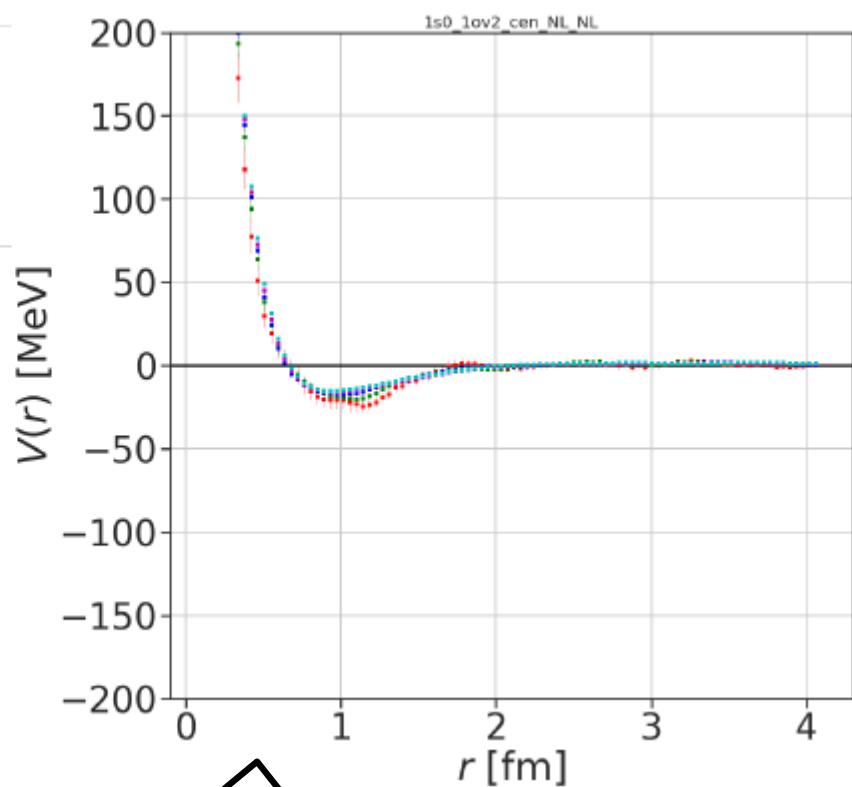
→ Partial wave decomposition

Effective $N\Lambda$ central potential in $l=1/2$, 1S_0 channel

w/o partial wave decomp.



w/ partial wave decomp.



**$N\Lambda$ scatt exp
@ J-PARC**

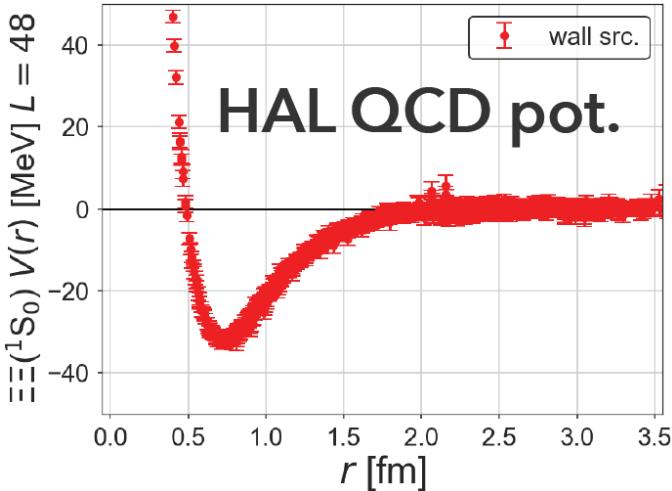
Significant Improvement!

[T. M. Doi]

(N.B. improvement in phase shifts would be much milder)

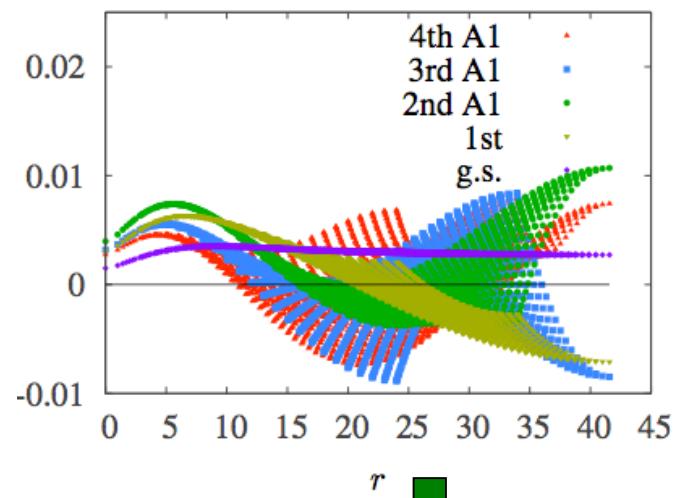
- t=12
- t=11
- t=10
- t=9
- t=8

New method to examine sys err in HAL potential

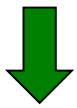


Solve
Schrodinger eq.
in Finite V

Eigen wave functions



Solve in FV



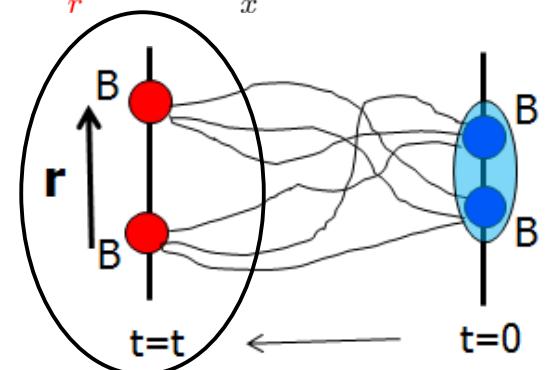
Eigen energies
= FV spectrum

n-th A1	ΔE_n [MeV]
0	-2.58(1)
1	52.49(2)
2	112.08(2)
3	169.78(2)
4	224.73(1)

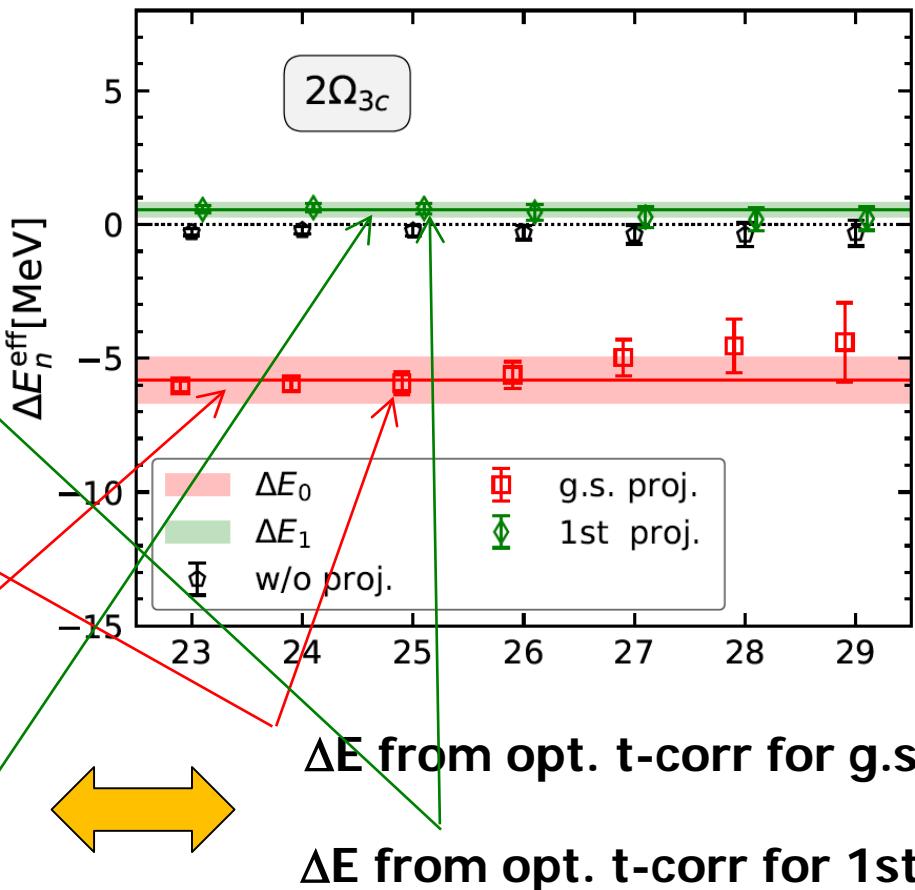
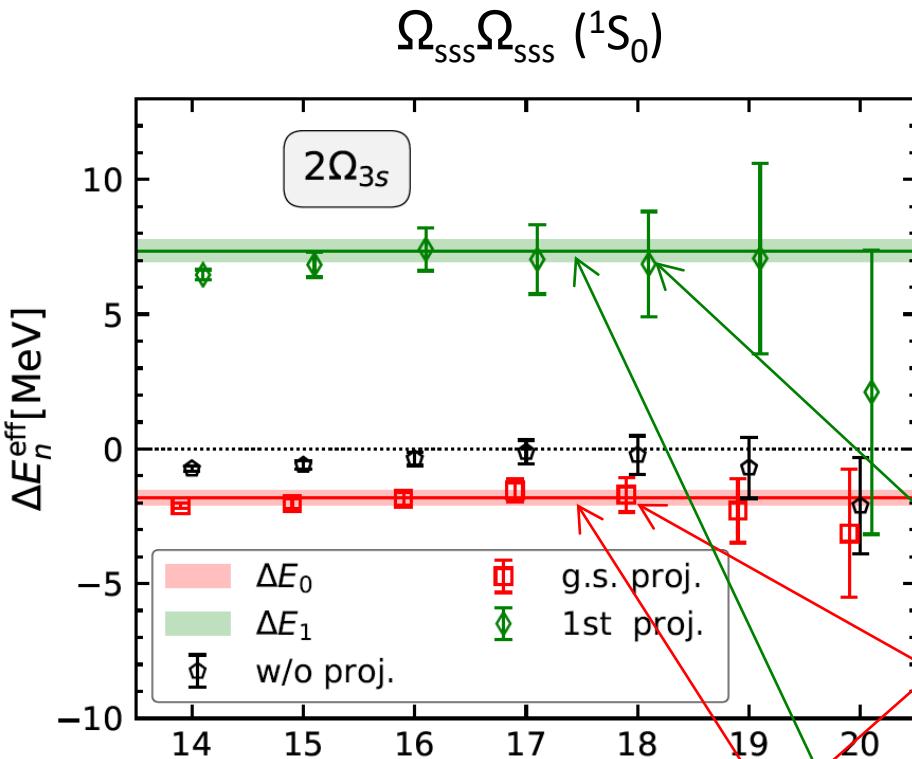
Consistency
Check!

FV spectrum
from
Optimized
temporal corr

$$\mathcal{J}_{\text{sink}}^{2B} = \sum_{\vec{r}} \psi^\dagger(\vec{r}) \sum_{\vec{x}} B(\vec{r} + \vec{x}) B(\vec{x})$$



Example: Comparison of FV spectrum for di-Omegas



ΔE of g.s. from HAL pot

ΔE of 1st from HAL pot

ΔE from opt. t-corr for g.s.

ΔE from opt. t-corr for 1st.

Good agreement!

Sys err in HAL pot under control

$\Omega_{sss}\Omega_{sss}$: g.s. dominant

$\Omega_{ccc}\Omega_{ccc}$: 1st. dominant

Y. Lyu et al., arXiv:2201.02782

Toward P-wave int by all-to-all method

LapH method

M. Peardon et al., PRD80(2009)054506

$$\mathcal{S}^{ab}(x, y) = \sum_{l=0}^{N_l} \omega_l v_l^a(x) v_l^{*b}(y)$$

$$\begin{aligned} \Delta^{ab}(x, y; U) &= \sum_{k=1}^3 \left\{ U_k^{ab}(x) \delta(y, x + \hat{k}) + U_k^{ba}(y)^* \delta(y, x - \hat{k}) - 2\delta(x, y) \delta^{ab} \right\} \\ &= (\mathbf{v}_1 \quad \mathbf{v}_2 \quad \cdots \quad \mathbf{v}_N) \begin{pmatrix} \lambda_1 & & & \mathbf{0} \\ & \lambda_2 & & \\ & & \ddots & \\ \mathbf{0} & & & \lambda_N \end{pmatrix} \begin{pmatrix} \mathbf{v}_1^\dagger \\ \mathbf{v}_2^\dagger \\ \vdots \\ \mathbf{v}_N^\dagger \end{pmatrix} \end{aligned}$$

Approximate all-to-all prop by N_l low-modes in gauge covariant Laplacian

New improvement: Free LapH method

T. Sugiura et al., PoS LAT2021,565

gauge cov. Laplacian \rightarrow free Laplacian

$$\Delta(x, y) = \sum_{k=1}^3 \left\{ \delta(y, x + \hat{k}) + \delta(y, x - \hat{k}) - 2\delta(x, y) \right\}$$

Comput. Cost can be reduced from $O(N_c^4 \times N_l^4) \rightarrow N_c! N_c O(N_l^3)$

\rightarrow Typically $O(100)$ speedup!

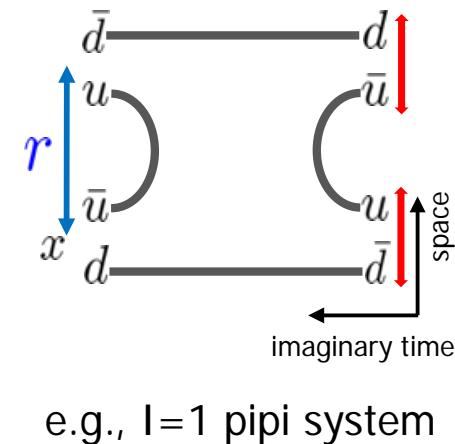
Explicit calc for NN in progress

Resonances, Exotics

Challenge: calc of quark-annihilation diagram

↔ all-to-all propagator required

$\propto O(L^4)$ cost

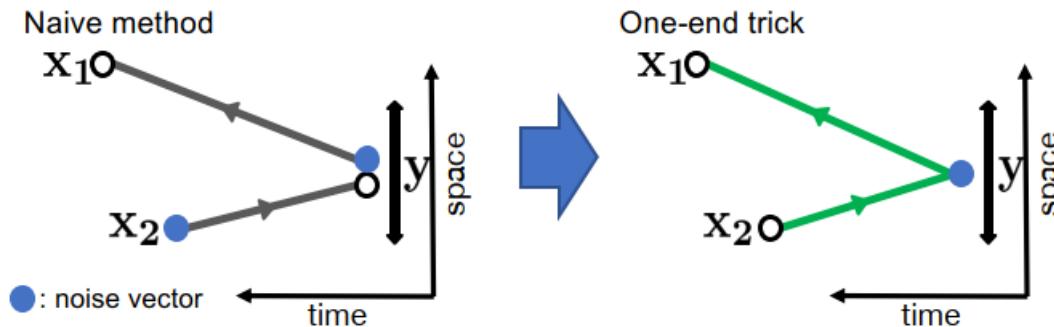


e.g., I=1 pipi system

Hybrid method (low-modes + stochastic estimate) is noisy

Y. Akahoshi et al., (HAL Coll.), ('19, '20)

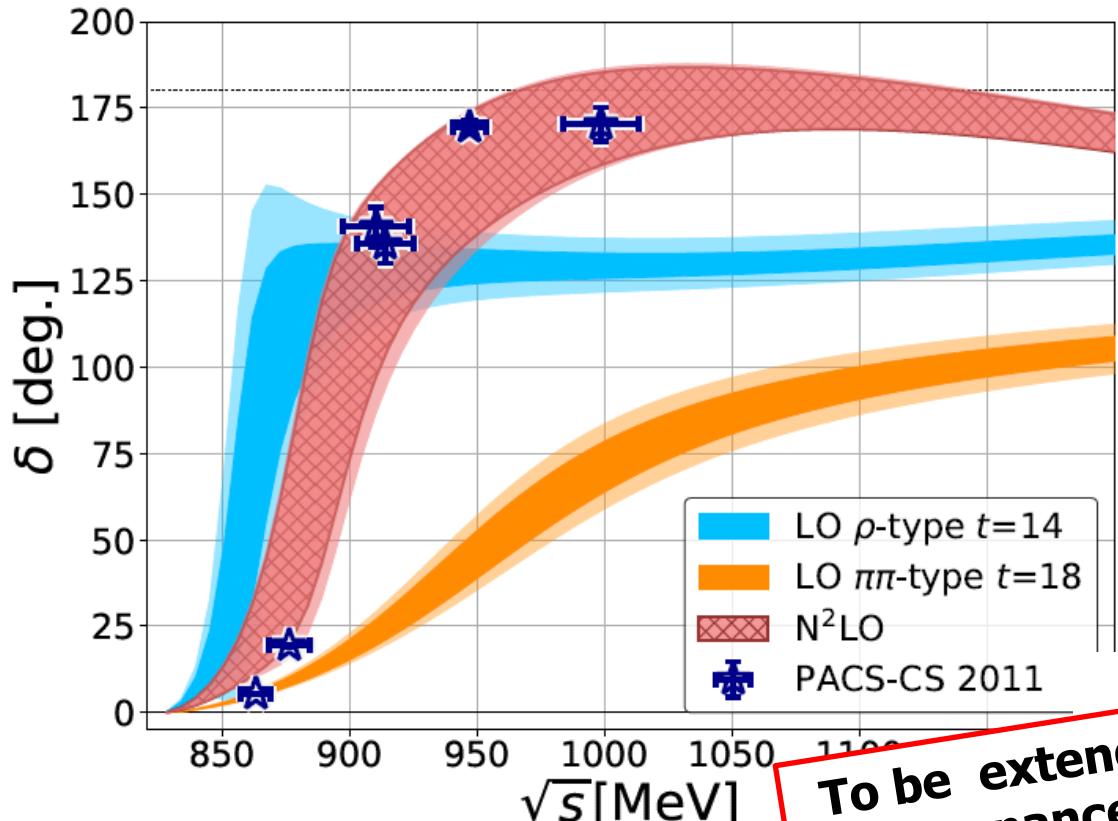
New method w/ one-end trick



McNeile-Michael (2006)

of noise vectors: 2 (naive method) → 1 (one-end trick) → $x_1/10$ stat error!

Study of $|l|=1$ pipi system: phase shifts



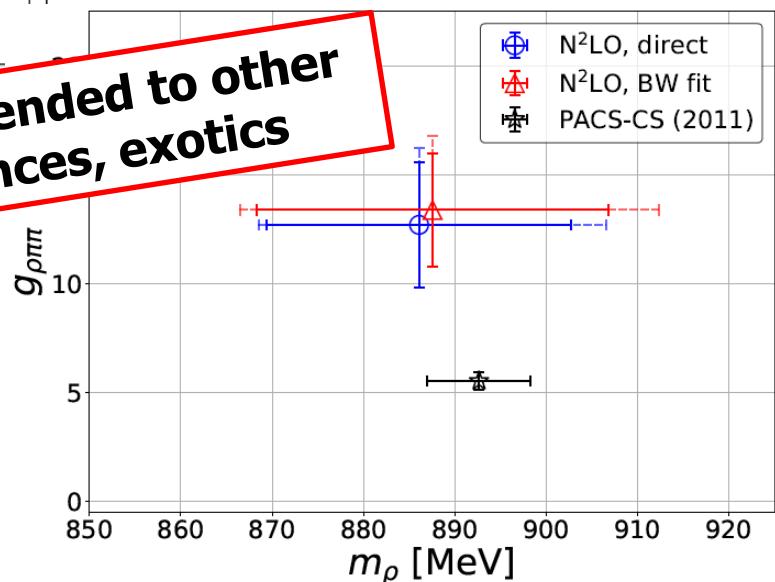
$$m(\text{pi}) \sim = 0.40 \text{GeV}$$

$$m(\text{rho}) \sim = 0.89 \text{GeV}$$

mass & $g(\rho\pi\pi)$

Rho-resonance from N2LO analysis
mass: consistent w/ FV method
width: improvement of N2LO necessary

To be extended to other resonances, exotics



2019

2020

2021

2022

2030s

From “near the physical point”
To “on the physical point”

(uncertainty from quark mass dep
~ B.E. of dibaryon, hypernuclei)

Config generation

Dibaryons

Hyperon forces

Charmed forces

K comp \leftrightarrow ~ 8yr
(K: 5.5yr + HOKUSAI: 2.5yr)



11 PFlops

Fugaku \leftrightarrow ~ 2yr
(S/D-waves)



440 PFlops

P-wave
LS-forces

3-body forces

Exotics, Resonances

2019

2020

2021

2022

2030s



J-PARC

YN, YY, (YNN)
Exotic hyper-nuclei



**LIGO
KAGRA**

NICER

GW in NS merger
NS radius

→ EoS



RIBF

3NF ($I=1/2, 3/2$)
r-process in NS merger

FRIB (2022-)

FAIR (2025(?) -)



LHC/RHIC

Baryon correlation
Exotic hadrons

LHC RUN3 (2022-24)



Belle II

Exotic hadrons



K-computer



Fugaku

Summary

- Renaissance in particle/nuclear/astro-physics
 - **Observations** of neutron stars (LIGO-Virgo-KAGRA, NICER, ...)
 - **Experiments** of hadrons/nuclei → J-PARC, LHC, Belle II, ...
 - **Theory** by LQCD calc of hadron interactions
- The 1st LQCD for Baryon Interactions near the phys. point
 - Central/Tensor forces for NN/YN/YY in $P=(+)$ channel
 - Dibaryons, Applications to Hypernuclei, EoS
- Prospects
 - **Baryon forces on the physical point by Fugaku supercomputer**
 - Dibaryons, Hypernuclei, charmed systems
 - Partial wave decomposition, all-to-all methods
 - Future: $P=(-)$ channel, LS-forces, 3-baryon forces, etc., & EoS
 - Resonances & Exotics