# **Compton Polarimetry for RCS**



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### **RCS** properties

- RCS accelerates electron bunches from 0.4 to full beam energy (5-18 GeV)
- Bunch frequency  $\rightarrow$  2 Hz
- Bunch charge  $\rightarrow$  up to 28 nA
- Ramping time = 100 ms



### <u>Options</u>

- Compton polarimeter in/after RCS
  - Analyzing power depends on energy difficult to implement in RCS – could place after RCS in transfer line
  - Initial time estimates assumed same laser as ESR
     → measurement times much too long

- Møller just after RCS RCS
  - Analyzing power nearly constant
  - Requires spectrometer
  - Destructive difficult to fit in transfer line
  - Measurement times not short



## Hall A Møller in Transfer Line

First look at integrating Møller polarimeter starting with Hall A/JLab design

RCS extraction region very crowded

 $\rightarrow$  Extraction beam line shares tunnel with several beam lines

 $\rightarrow$  RCS starts at lower height, beam is kicked up to same elevation as ESR in the extraction line

→ Would like to avoid placing the polarimeter between height-adjusting dipoles (spin precession) Looking at placing the polarimeter before first dipole, or after last dipole

Beam



#### Polarimeter before vertical bend

Picture courtesy Bijan Bhandari



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$$\mathcal{L} = f_{coll} N_{\gamma} N_{e} \frac{\cos\left(\alpha_{c}/2\right)}{2\pi} \frac{1}{\sqrt{\sigma_{x,\gamma}^{2} + \sigma_{x,e}^{2}}} \frac{1}{\sqrt{\left(\sigma_{y,\gamma}^{2} + \sigma_{y,e}^{2}\right)\cos^{2}\left(\alpha_{c}/2\right) + \left(\sigma_{z,\gamma}^{2} + \sigma_{z,e}^{2}\right)\sin^{2}\left(\alpha_{c}/2\right)}}}$$

$$I = \frac{1}{N_{\gamma} = \langle \mathsf{P} \rangle / (\mathsf{f}_{\mathsf{laser}} \mathsf{E}_{\mathsf{laser}})}$$

Original rate estimates for RCS Compton (done for Yellow Report/CDR) assumed same laser as used for ESR:

 $\langle P \rangle = 5 \text{ W}, \text{ f}_{\text{laser}} = 25, 100 \text{ MHz} \rightarrow 5.4\text{E11}, 1.3\text{E11 photons/pulse}$ 

Assuming 28 nC electron bunches at 2 Hz, this gives about 2.9 Hz of backscattered photons (by design) Measurement time on the order of days for  $(dP/P)_{stat} = 1\%$ 

Should have assumed low duty cycle, higher peak power laser (similar to HERMES LPOL)  $\rightarrow$  Example system from RPMC lasers: Pulse energy = 30 mJ @ 2 Hz (<P>=60 mW)  $\rightarrow$  8.0E16 photons/pulse

For 15 ns pulses, backscattered photon rate is 240 kHz, "ideal" measurement times on the order of a few seconds



$$A_{\rm T} = \frac{2\pi r_o^2 a}{(d\sigma/d\rho)} \cos \phi \left[ \rho (1-a) \frac{\sqrt{4a\rho(1-\rho)}}{(1-\rho(1-a))} \right]$$

0.5  $E_{beam} = 5 \text{ GeV}$ 0.4  $E_{beam} = 10 \text{ GeV}$ φ=0 E<sub>beam</sub>=18 GeV ........ 0.3 0.2  $A_{\mathsf{T}}$ 0.1 0.0 -0.1-0.2 0.0 0.2 0.4 0.6 0.8 1.0  $\rho = E_{\gamma}/E_{\nu}^{\max}$ 

I contrast to longitudinal measurement, measurement of transverse polarization requires sensitivity to spatial dependence of asymmetry





### Simple rate, measurement time estimates

$$t^{-1} = \mathcal{L}\sigma \left(\frac{\Delta P}{P}\right)^2 P^2 A_{method}^2$$

Average analyzing power:  $A^2_{method} = \langle A \rangle^2 \rightarrow$  Average value of asymmetry over acceptance

Energy-weighted:  $A^2_{method} = \left(\frac{\langle EA \rangle}{\langle E \rangle}\right)^2 \rightarrow$  Energy deposited in detector for each helicity state

Differential:  $A^2_{method} = \langle A^2 \rangle \rightarrow$  Measurement of asymmetry bin-by-bin vs. energy, etc.

Assuming 80% polarization,  $\langle P_{laser} \rangle = 6$  mW, 300  $\mu$ m beam spot size...,  $\rightarrow$  time for 1% measurement

E <sub>beam</sub>	A <sub>avg</sub>	T <sub>avg</sub>	<b>A</b> <sub>energy</sub>	T <sub>energy</sub>	A <sub>diff</sub>	T <sub>diff</sub>
5	4.51%	243 s	5.78%	148 s	5.48%	164 s
10	7.79%	92 s	10.15%	54 s	9.56%	61 s
18	11.29%	51 s	14.91%	29 s	13.96%	33 s



# **Toy Monte Carlo**

- Used toy Monte Carlo to look at asymmetries, energy/per bunch etc.
- Generated events for 240 bunches (i.e., 2 minutes of running)
  - For each bunch/crossing, generated number of backscattered photons based on simple luminosity estimates
  - -Assumed luminosity 10x smaller than previous slide  $\rightarrow$  6 mW average laser power
  - 12,000-15,800 backscattered photons/bunch crossing
- Assumed detector 25 m away from laser-beam collision point
  - No realistic simulation of detector just a point in space to look at distributions
- Event generator ignores hourglass effect, assumes all events generated at (0,0) → does incorporated finite beam sizes, but only at that point



LEP polarimeter measured difference in vertical position for h+ and h- induce by transverse analyzing power

At 5 GeV, average position difference for h+ and h- is

 $Y_{h+} - Y_{h-} = 0.49 \text{ mm}$ 

At 18 GeV, difference is similar magnitude:

 $Y_{h+} - Y_{h-} = 0.38 \text{ mm}$ 

While measurable, this quantity extremely sensitive to knowledge of absolute detector position and laserbeam collision point







Differential measurement of asymmetry vs. position at detector allows us to incorporate offsets in the fit

Example using Toy MC for counting-mode asymmetry vs. y assuming 0.1 mm segmentation (240 bunches)

→ Requires detector operated in integrating mode (~10,000 photons/bunch) with signal proportional to number of photons in each channel





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# Summary

- Compton polarimeter looks like a viable solution for RCS polarimetry
  - Significantly shorter measurement times
  - Requires less space along beamline
- Issues
  - Can only operate in integrating mode requires additional characterization of detector response compared to counting mode
  - Integrating mode requires more attention to detector noise, pedestal
  - No bunch-by-bunch measurement will need to average over several bunches
  - -Large total energy/bunch when integrated over all backscattered photons  $\rightarrow$  @ 18 GeV, avg. < $E_{\gamma}$ >=3.41 GeV \* 12,000 photons = 42,000 GeV! Is this an issue?
- To-do:
  - Identify possible locations in ESR transfer line main challenge will likely be finding space for long drift to photon detector
  - Investigate detectors for integrating mode operation (diamond? silicon?)
  - More detailed simulations incorporating detector response, noise



### **Analyzing Power and Measurement Times**

Measurement time depends on luminosity, analyzing power, and measurement technique

$$t^{-1} = \mathcal{L}\sigma \left(\frac{\Delta P}{P}\right)^2 A_{method}^2$$

Average analyzing power: 
$$A^2_{method} = \langle A \rangle^2$$

 $\rightarrow$  Average value of asymmetry over acceptance

Energy-weighted:

$$A_{method}^2 = \left(\frac{\langle EA \rangle}{\langle E \rangle}\right)^2$$

 $\rightarrow$  Energy deposited in detector for each helicity state

Differential:

$$A^2_{method} = \langle A^2 \rangle$$

 $\rightarrow$  Measurement of asymmetry bin-by-bin vs. energy, etc.

$$\langle A \rangle^2 < \left( \frac{\langle EA \rangle}{\langle E \rangle} \right)^2 < \langle A^2 \rangle$$

