

Machine Learning for Collinear QCD Lecture 3

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Outline

A Layman's Recap

B Proton PDFs

C Nuclear PDFs

D Fragmentation functions

E Tutorial 2 — Fitting the gluon PDF

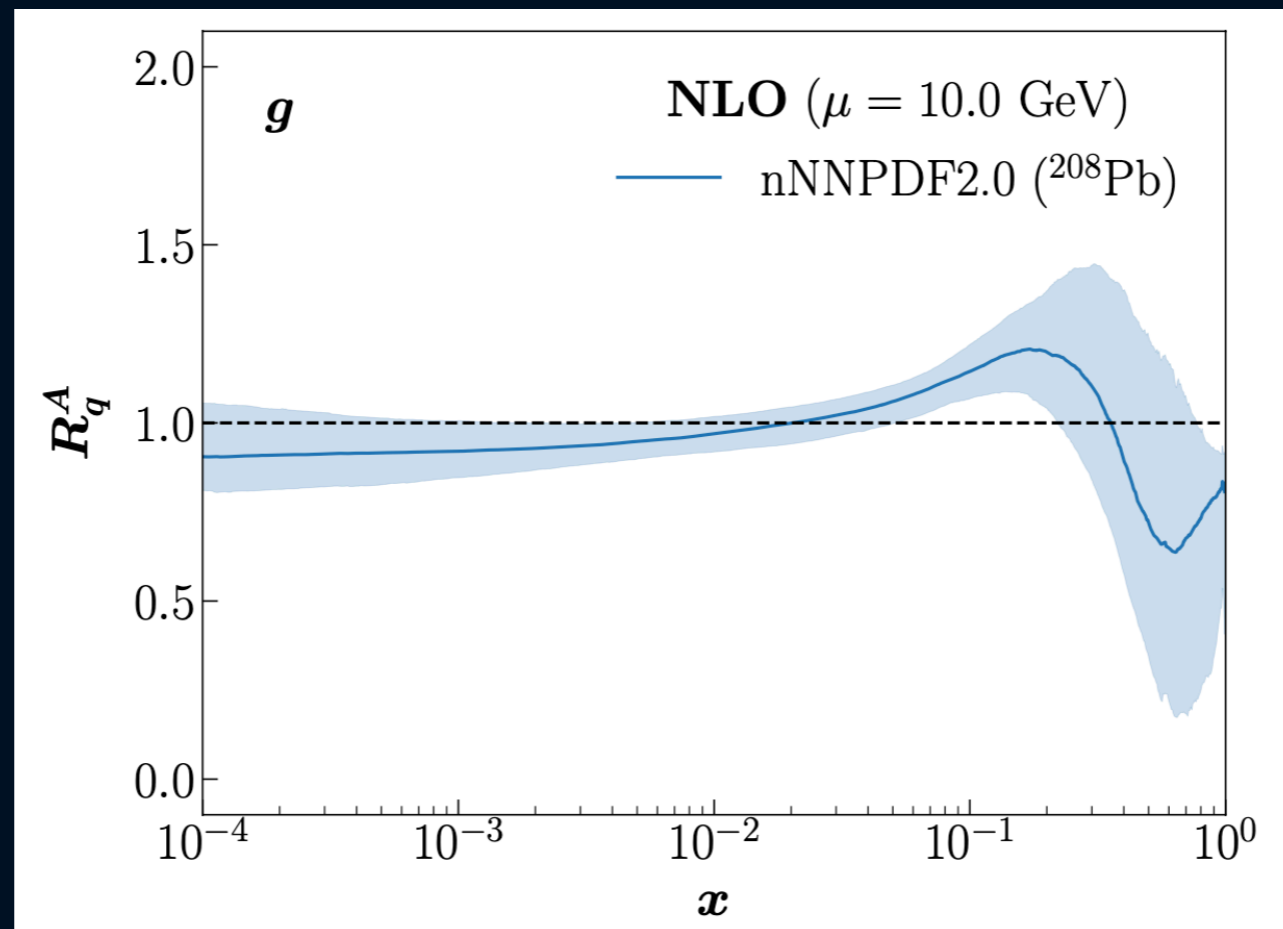
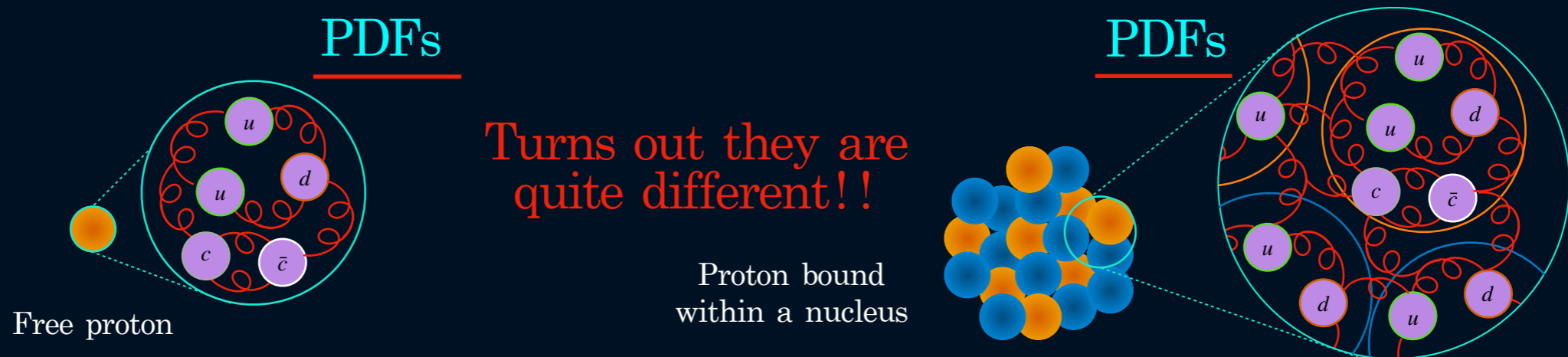
Substructure of nucleons and nuclei

To a certain approximation, the probability of finding quarks and gluons in a hadron carrying a momentum-fraction x of the hadron's momentum is encoded in non-perturbative parton distribution functions (PDFs).



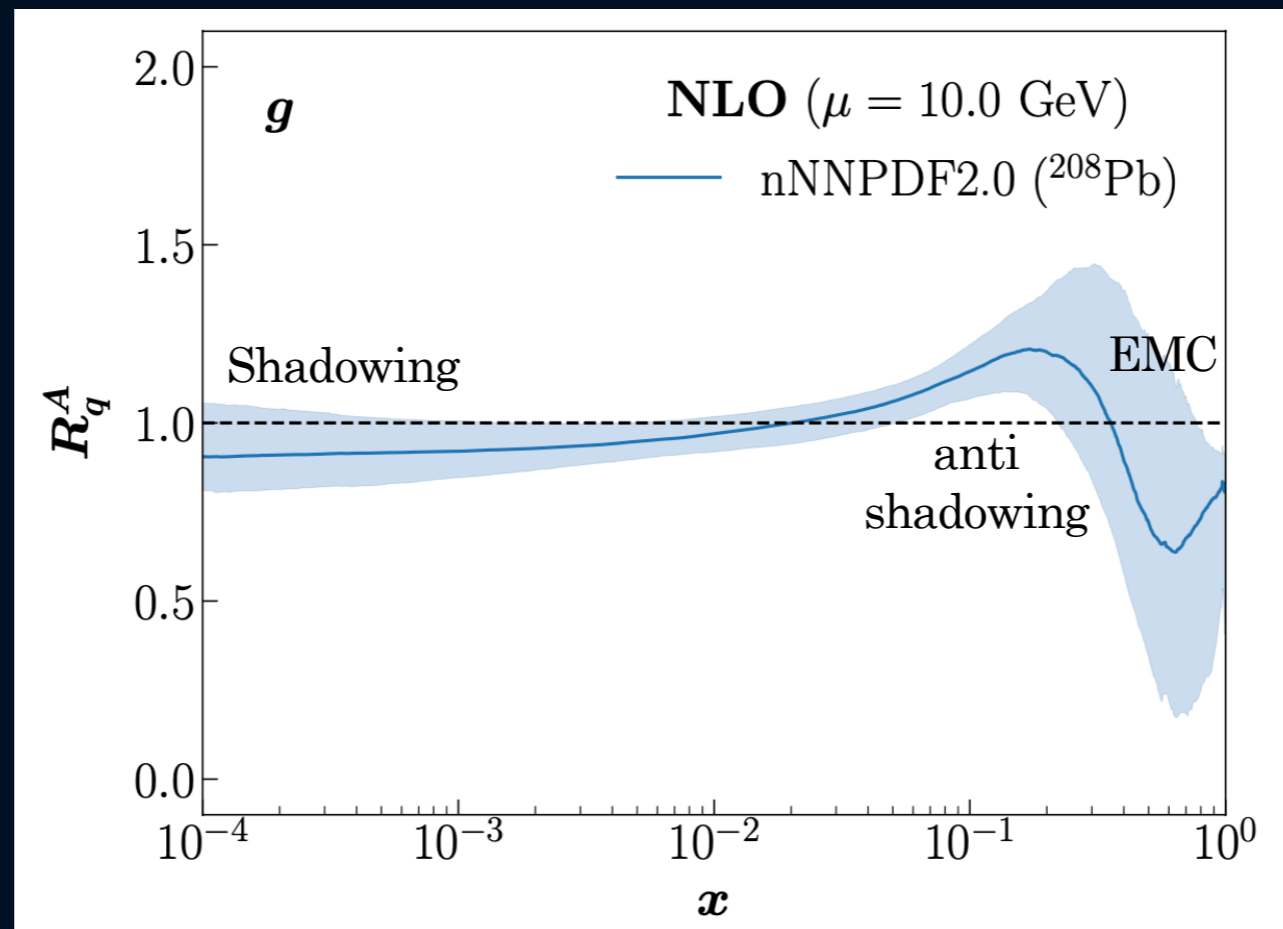
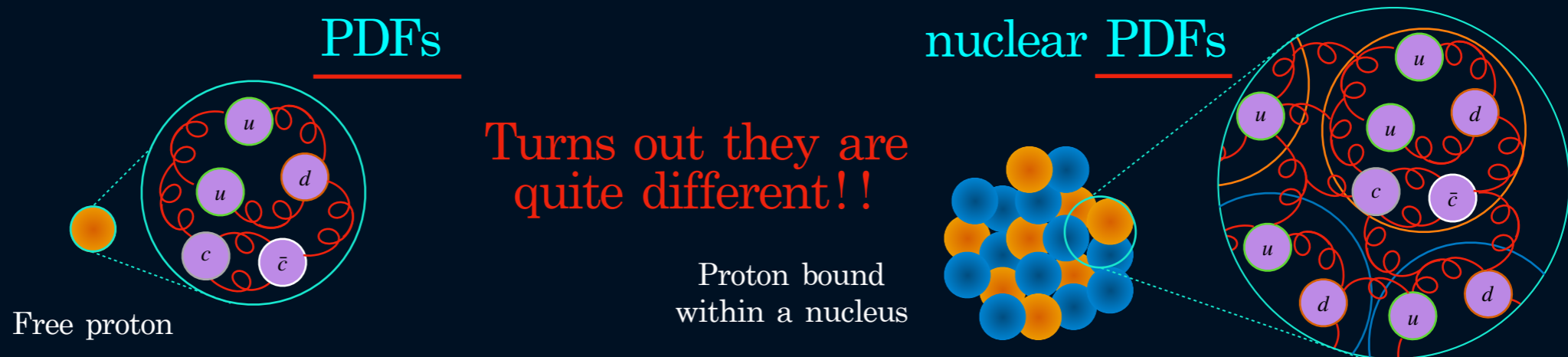
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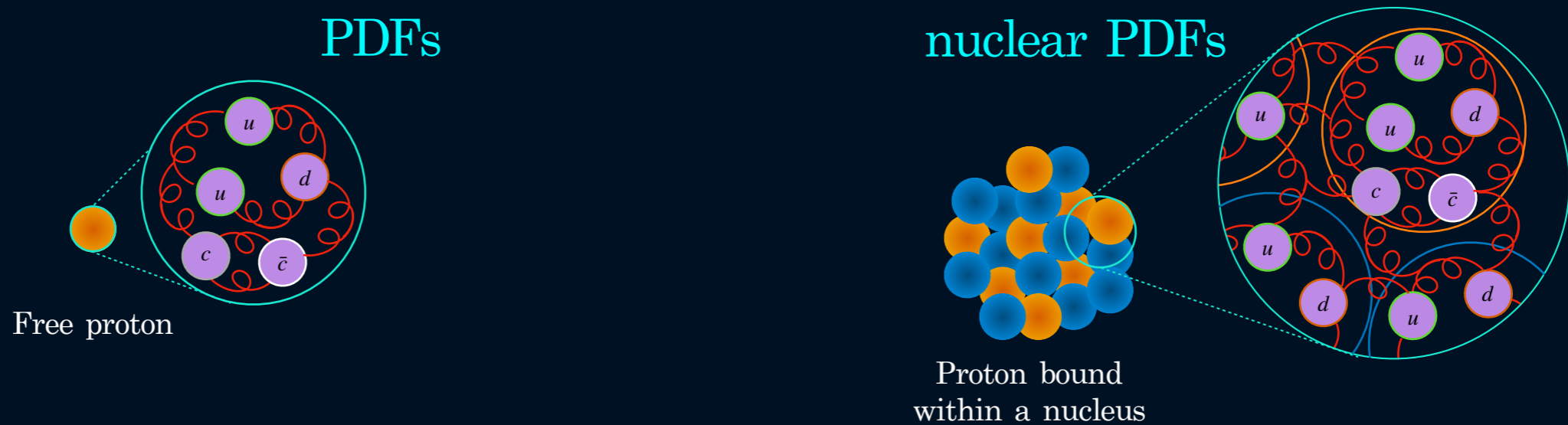
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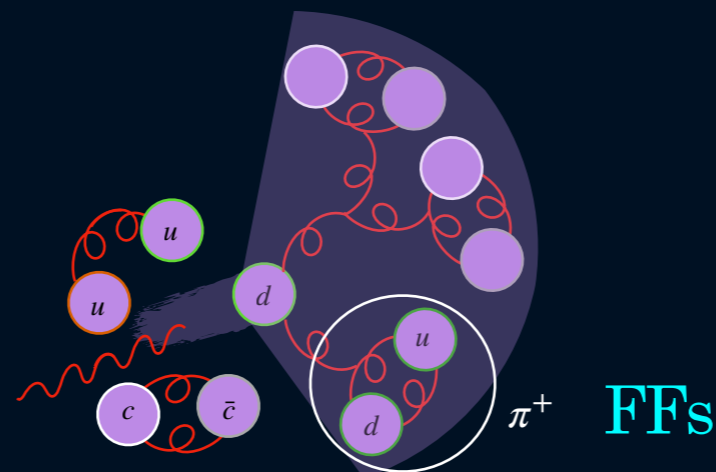


Substructure of nucleons and nuclei

To a certain approximation, the probability of finding quarks and gluons in a hadron carrying a momentum-fraction x of the hadron's momentum is encoded in non-perturbative parton distribution functions (PDFs).



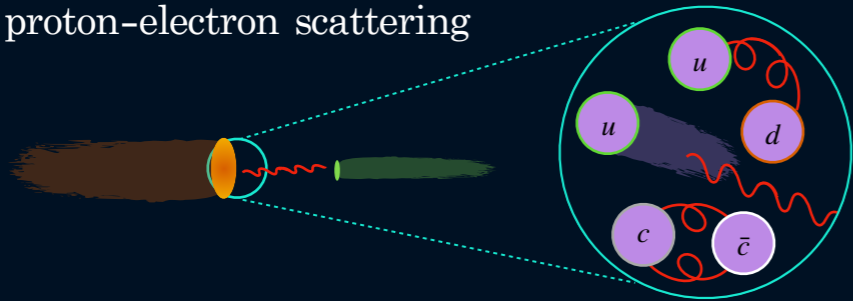
In addition, fragmentation functions (FFs) encodes the probability of producing a hadron from a quark fragmentation.



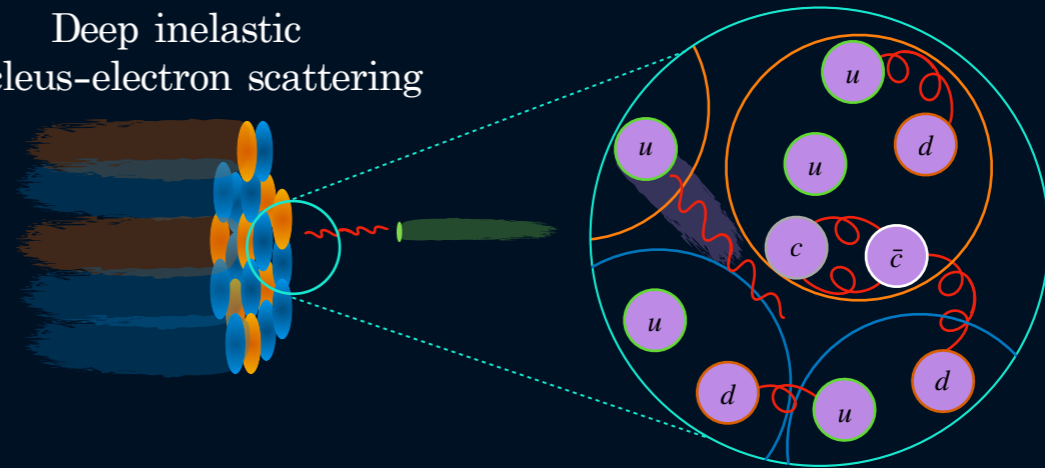
Substructure of nucleons and nuclei

We extract these **objects** from the scattering-probabilities off protons and nuclei that we measure in collider-**experiments**.

Deep inelastic proton-electron scattering



Deep inelastic Nucleus-electron scattering

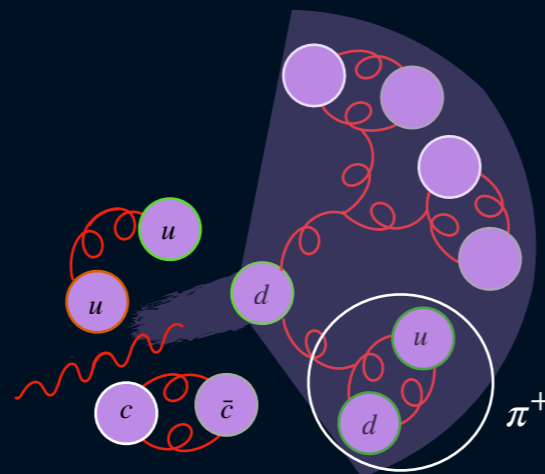


$$\text{Scattering probability} = \hat{\sigma} \otimes f^{(p)}$$

Perturbative
Non-perturbative PDFs

$$\text{Scattering probability} = \hat{\sigma} \otimes f^{(A)}$$

Perturbative
Non-perturbative nuclear PDFs



$$\text{Scattering probability} = \hat{\sigma} \otimes f^{(p)} \otimes D^{(\pi^\pm)}$$

Perturbative
Non-perturbative

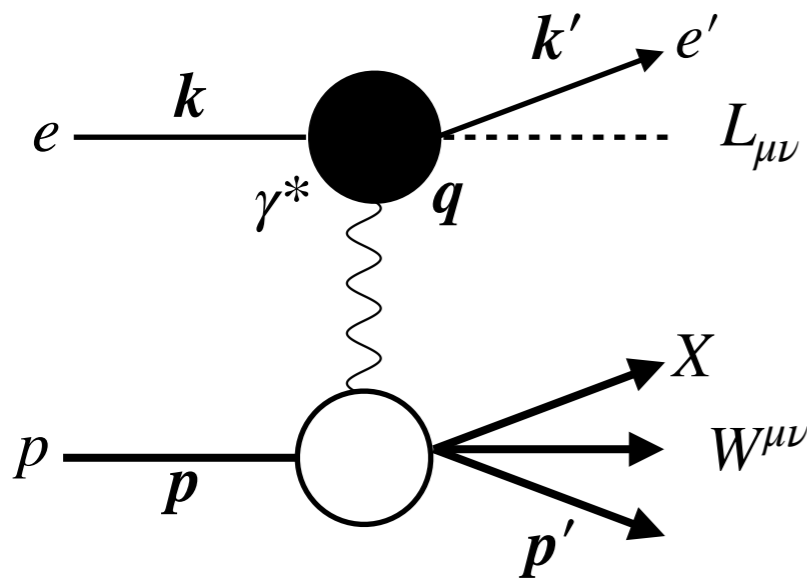
Proton PDFs

To calculate a cross section involving a hadron in renormalizable perturbation series:

$$\mu \sim Q \gg 0 \implies \alpha_s(\mu^2) \ll 1 \implies O(Q, m, \mu) = O^{(0)} + O^{(1)} + O^{(2)} + \dots$$

$$\mu \text{ appears in } O(Q, m, \mu) \text{ as } \begin{cases} \frac{Q}{\mu} \sim 1 \implies \text{short-distance} \\ \frac{\mu}{m} \gg 1 \implies \text{long-distance} \end{cases}$$

Deep inelastic scattering



Collinear factorization theorems allows us to separate them systematically.

$$W^{\mu\nu} = \frac{1}{4\pi} \int d^4y e^{iq \cdot y} \sum_X \langle p | j^\mu(y) | X \rangle \langle X | j^\nu(0) | p \rangle$$

$$\stackrel{\text{Bjorken limit}}{\equiv} \sum_a \int_x^1 \frac{d\xi}{\xi} \underbrace{f_{a/p}(\xi, \mu)}_{\text{PDF}} \underbrace{H_a^{\mu\nu}(q^\mu, \xi p^\mu, \mu, \alpha_s(\mu))}_{\text{Hard scattering coefficient}} + \text{H.T.}$$

$$\frac{d^2 \sigma^{lp \rightarrow lX}}{dx dQ^2} \propto L_{\mu\nu} W^{\mu\nu} \propto \sum_a \int_x^1 d\xi f_{a/H}(\xi, Q^2) \frac{d^2 \hat{\sigma}}{dx dQ^2}(x/\xi, Q^2) = \hat{\sigma} \otimes f_{a/H}$$

Nuclear PDFs

EMC Effect: $\frac{F_2^A}{F_2^D} \neq \frac{2}{A} \cdot C \otimes \frac{Zf^p + (A-Z)f^n}{f^p + f^n}$ Nucleus is **not** an ensemble of Z free protons and $(A-Z)$ free neutrons.

Four equally possible scenarios:

- a) the **fundamental interactions are the same** but **PDFs in nucleus are different**.
- b) The fundamental interactions are different in the medium but PDFs are the same.
- c) Both (a) and (b).
- d) The factorization picture is no longer valid.

Working with the assertion (a): $\tilde{W}^{\mu\nu} = \frac{1}{4\pi} \int d^4y e^{iq \cdot y} \sum_X \langle \overset{\downarrow}{A} | j^\mu(y) | X \rangle \langle X | j^\nu(0) | \overset{\downarrow}{A} \rangle$ Replace proton with nuclear states

$\underset{\text{Bjorken limit}}{\equiv} \sum_a \int_x^1 \frac{d\xi}{\xi} \underbrace{f_{a/A}(\xi, \mu)}_{\text{Nuclear PDF}} \underbrace{H_a^{\mu\nu}(q^\mu, \xi p^\mu, \mu, \alpha_s(\mu))}_{\text{Hard scattering coefficient}} + \text{H.T.}$

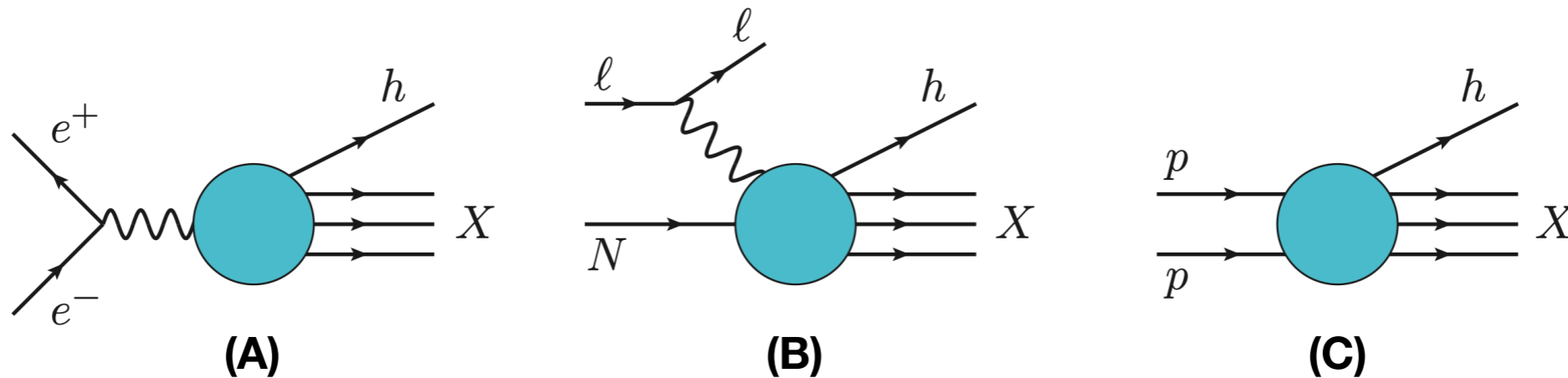
Physical quantity to be fitted

$f_{a/A} \overset{\downarrow}{=} Z \cdot \underset{\downarrow}{\tilde{f}_{alp}} + (A-Z) \cdot \underset{\downarrow}{\tilde{f}_{aln}} \implies \frac{d^2\sigma^{lA \rightarrow lX}}{dx dQ^2} \propto L_{\mu\nu} \tilde{W}^{\mu\nu} \propto Z \cdot \tilde{F}_2^{p/A} + (A-Z) \cdot \tilde{F}_2^{n/A} = \boxed{\hat{\sigma} \otimes f_{a/A}}$

Effective quantities

Fragmentation Functions

In measurements where a hadron is identified...



A. SIA: single-inclusive hadron production in electron-positron annihilation, $e^+ + e^- \rightarrow h + X$.

$$\sigma^{e^+ + e^- \rightarrow h + X} = \hat{\sigma} \otimes D^h$$

B. SIDIS: semi-inclusive deep-inelastic lepton-nucleon scattering, $\ell + N \rightarrow \ell + h + X$

$$\sigma^{\ell + p \rightarrow \ell + h + X} = \hat{\sigma} \otimes f^p \otimes D^h$$

C. Single-inclusive hadron production in proton-proton collisions, $p + p(\bar{p}) \rightarrow h + X$

$$\sigma^{p + p(\bar{p}) \rightarrow h + X} = \hat{\sigma} \otimes f^p \otimes f^p \otimes D^h$$

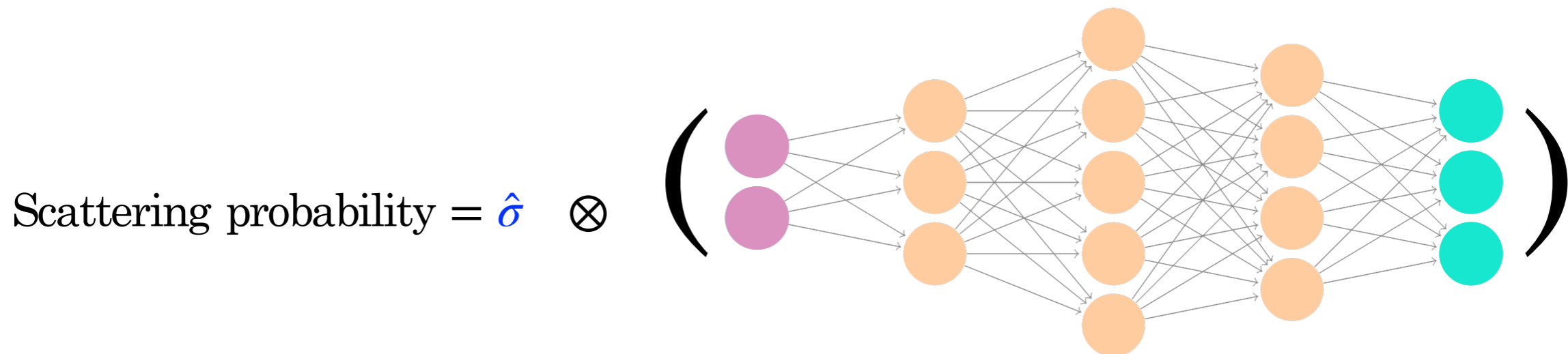
Neural Networks parameterisation

Due to DGLAP being **perturbative**, we only need to extract the renormalised PDFs or FFs at one arbitrary scale μ

$$t(\theta) = \boxed{H^{\text{hard}}(Q) \otimes \Gamma^{\text{DGLAP}}(Q, \mu)} \otimes [xf(\mu, \theta) || zD(\mu, \theta)]$$

$[xf(\mu, \theta) || zD(\mu, \theta)] \rightarrow$ What kind of Parameterisation?

Since PDFs and FFs functional form is not known from first principles, we parameterise them using feed-forward neural networks (NNs) since NNs can approximate any continuous function within the data range

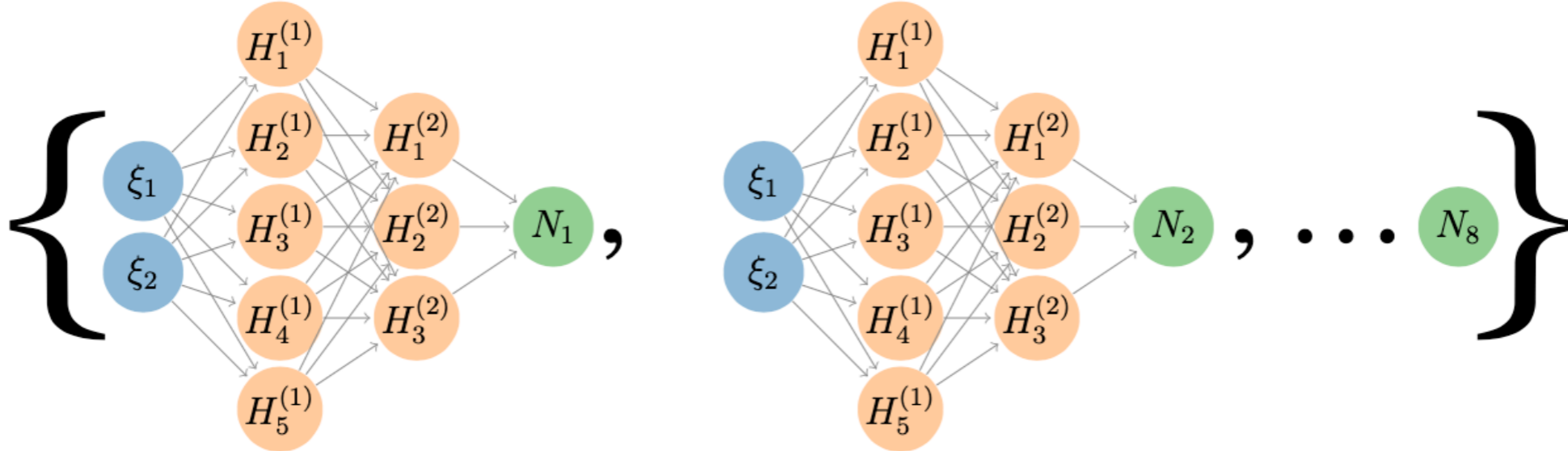


NNPDF3.1 — proton PDFs

Parameterisation

$\{g \quad \Sigma \quad T_3 \quad T_8 \quad V \quad V_3 \quad V_8 \quad (c^+)\}$

NNPDF3.1 — arXiv:1706.00428
NNPDF4.0 — arXiv:2109.02653



$$f_i(x, \mu_0) = A_i \hat{f}_i(x, \mu_0), \quad \hat{f}_i(x, \mu_0) = x^{\alpha_i} (1-x)^{\beta_i} N_i(x)$$

$$A_g = \frac{1 - \int_0^1 dx x \Sigma(x, \mu_0)}{\int_0^1 dx x \hat{g}(x, \mu_0)}, \quad A_\Sigma = A_{T_3} = A_{T_8} = 1$$

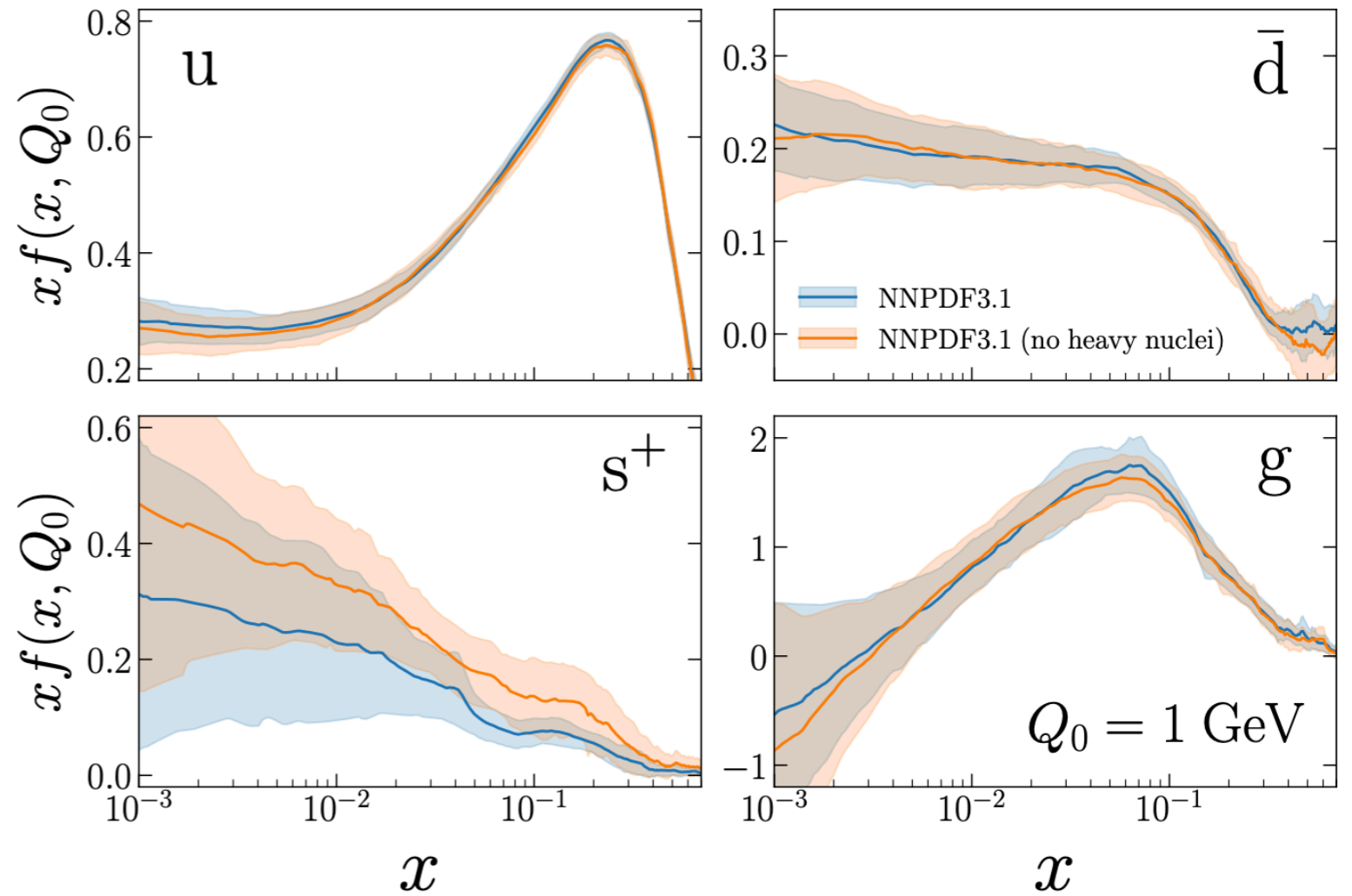
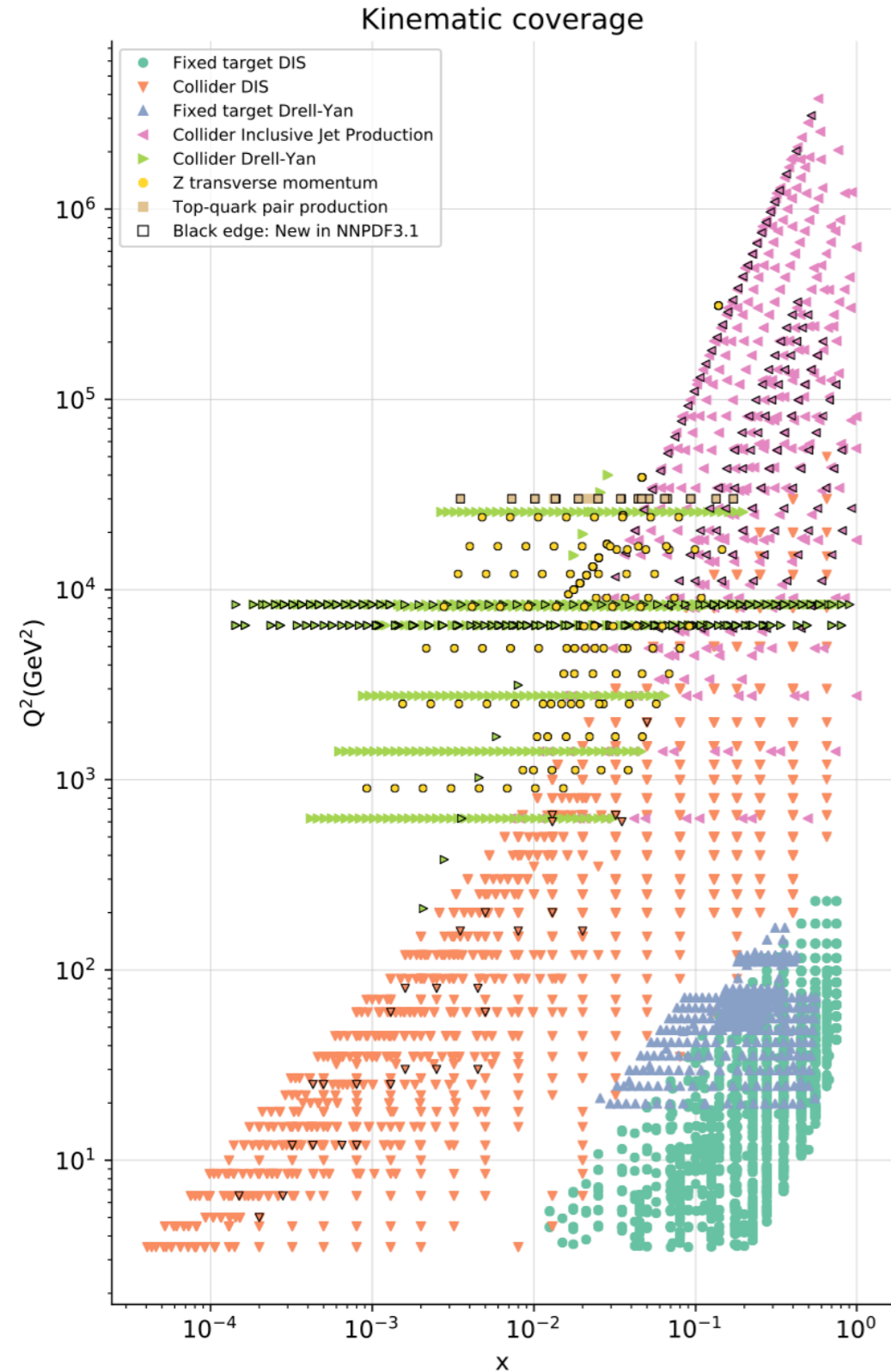
$$A_V = \frac{3}{\int_0^1 dx \hat{V}(x, \mu_0)}, \quad A_{V_3} = \frac{1}{\int_0^1 dx \hat{V}_3(x, \mu_0)}, \quad A_{V_8} = \frac{3}{\int_0^1 dx \hat{V}_8}$$

$$t(x, Q; \theta) = \sum_i H_i^{\text{hard}}(Q) \otimes \sum_j \Gamma_{ij}^{\text{DGLAP}}(Q, \mu) \otimes f_j(\mu; \theta)$$

Loss function

$$\chi^2 = \chi_{\text{exp}}^2 + \lambda_{\text{pos}} \sum_{l=1}^{N_{\text{pos}}} \sum_{j=1}^{N_A} \sum_{i_l=1}^{N_{\text{dat}}^{(l)}} \max\left(0, -\mathcal{F}_{i_l}^{(l)}(A_j)\right)$$

NNPDF3.1 — proton PDFs

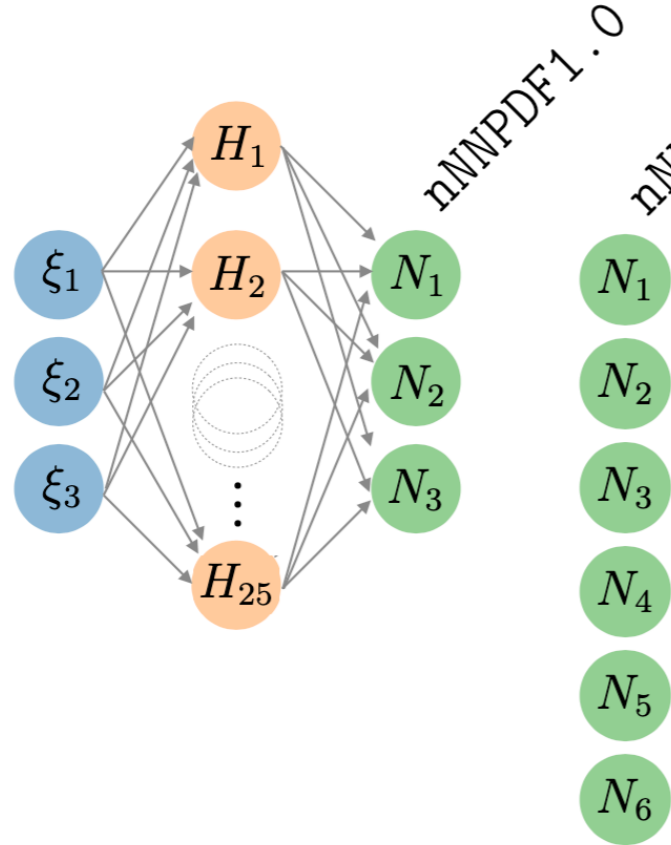


nNNPDF2.0 — nuclear PDFs

Parameterisation

$\{g \quad \Sigma \quad T_3 \quad T_8 \quad V \quad V_3 \quad V_8\}$

nNNPDF1.0 — arXiv:1904.00018
 nNNPDF2.0 — arXiv:2006.14629
 nNNPDF3.0 — arXiv:2201.12363



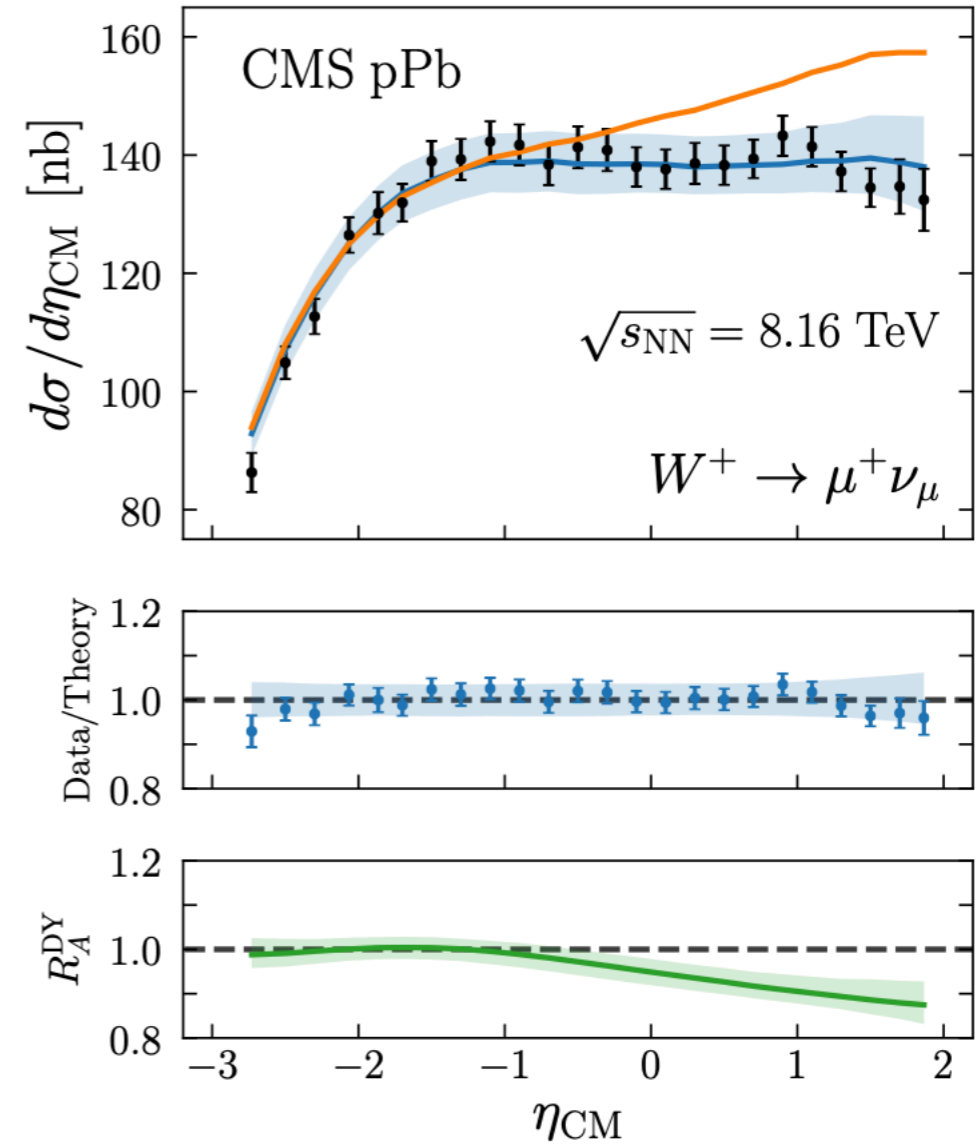
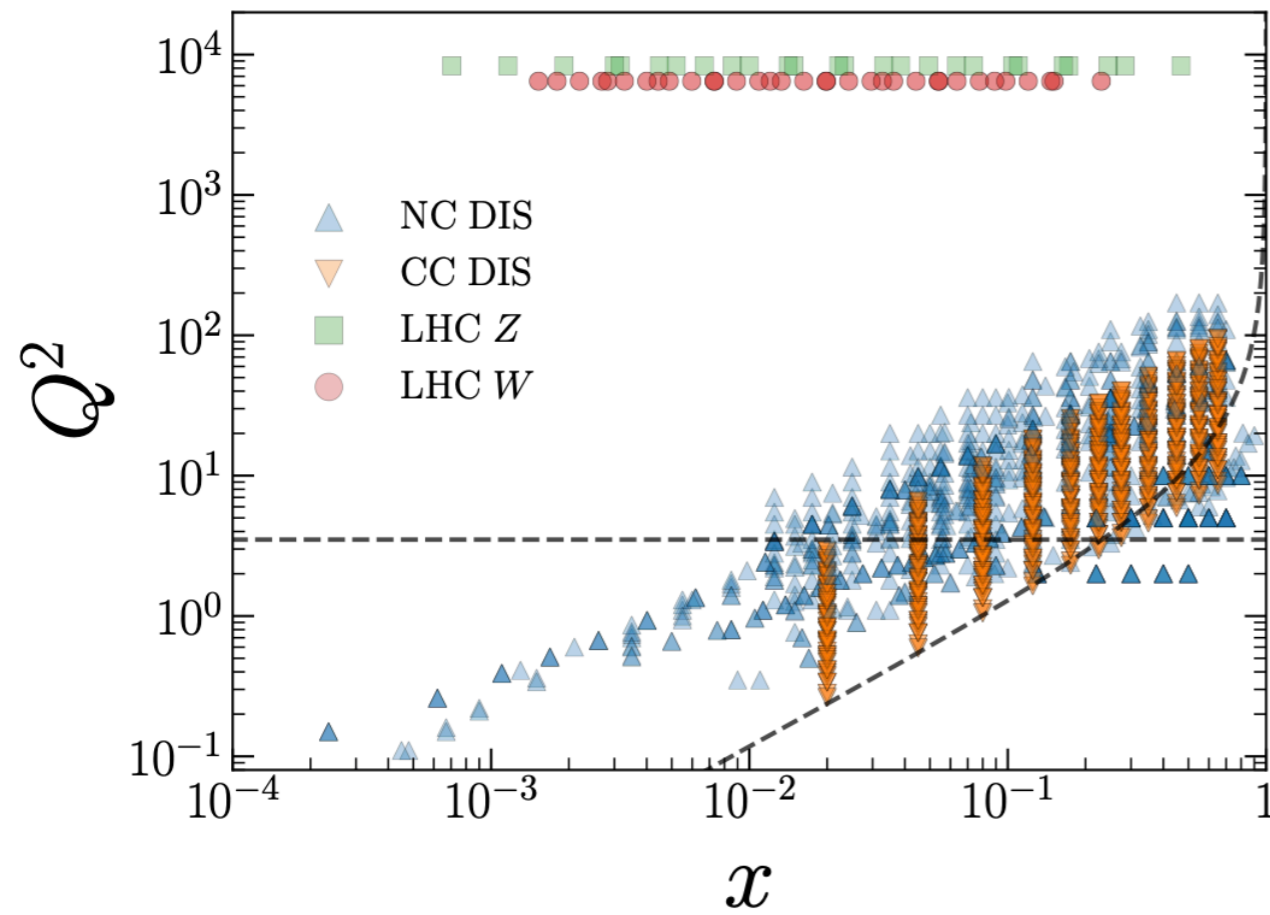
$$\begin{aligned}
 x\Sigma^{(p/A)}(x, Q_0) &= x^{\alpha_\Sigma} (1-x)^{\beta_\Sigma} N_1(x, A), \\
 xT_3^{(p/A)}(x, Q_0) &= x^{\alpha_{T_3}} (1-x)^{\beta_{T_3}} N_2(x, A), \\
 xT_8^{(p/A)}(x, Q_0) &= x^{\alpha_{T_8}} (1-x)^{\beta_{T_8}} N_3(x, A), \\
 xV^{(p/A)}(x, Q_0) &= B_V x^{\alpha_V} (1-x)^{\beta_V} N_4(x, A), \\
 xV_3^{(p/A)}(x, Q_0) &= B_{V_3} x^{\alpha_{V_3}} (1-x)^{\beta_{V_3}} N_5(x, A), \\
 xg^{(p/A)}(x, Q_0) &= B_g x^{\alpha_g} (1-x)^{\beta_g} N_6(x, A).
 \end{aligned}$$

$$t(x, Q; \theta) = \sum_i H_i^{\text{hard}}(Q) \otimes \sum_j \Gamma_{ij}^{\text{DGLAP}}(Q, \mu) \otimes f_j^A(\mu; \theta)$$

Loss function

$$\begin{aligned}
 \chi^2 &= \chi_{\text{exp}}^2 + \lambda_{\text{BC}} \sum_f \sum_{i=1}^{N_x} \left(q_f^{(p/A)}(x_i, Q_0, A=1) - q_f^{(p)}(x_i, Q_0) \right)^2 \\
 &+ \lambda_{\text{pos}} \sum_{l=1}^{N_{\text{pos}}} \sum_{j=1}^{N_A} \sum_{i_l=1}^{N_{\text{dat}}^{(l)}} \max\left(0, -\mathcal{F}_{i_l}^{(l)}(A_j)\right),
 \end{aligned}$$

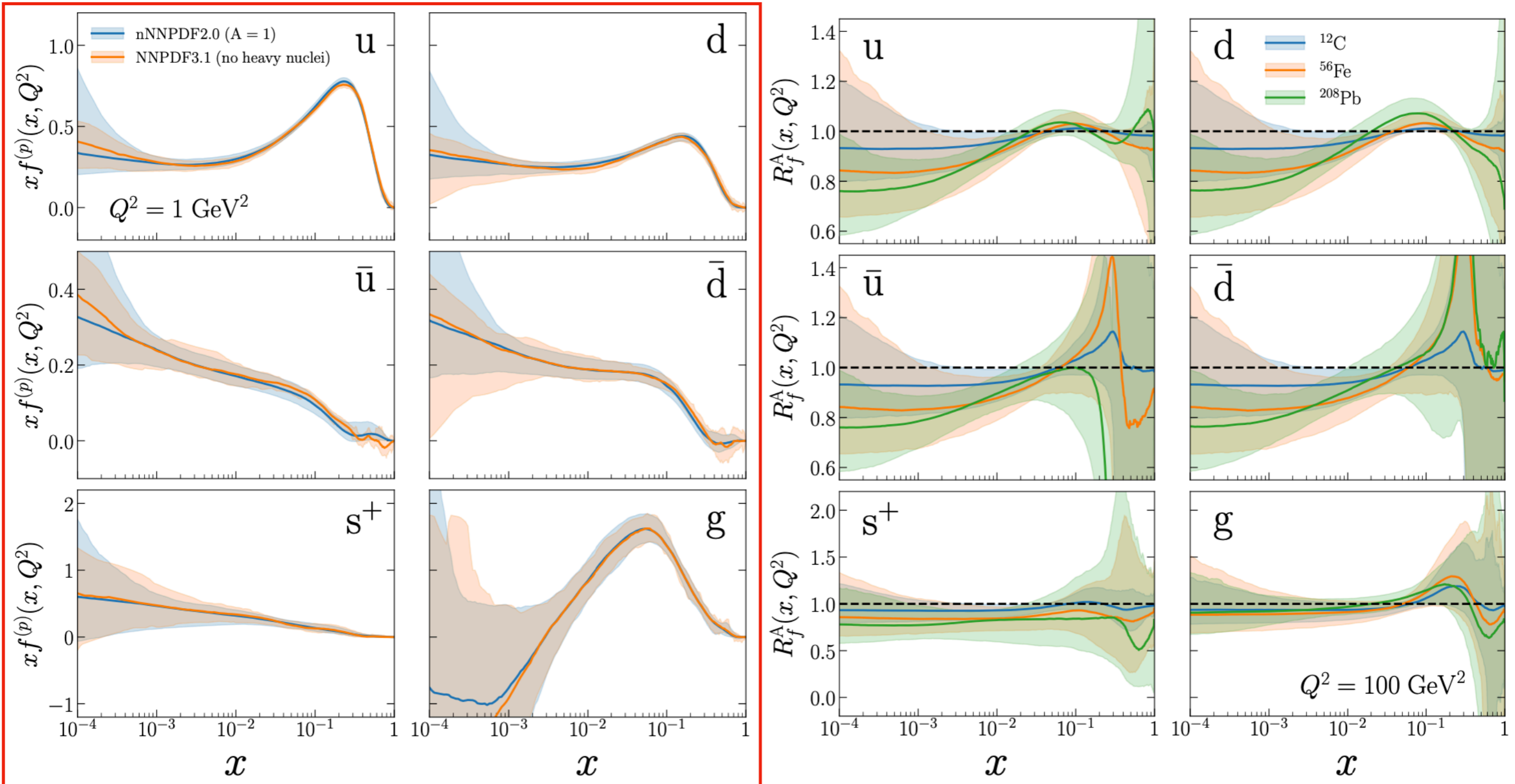
nNNPDF2.0 — nuclear PDFs



nNNPDF2.0 — nuclear PDFs

$$\chi^2 = \chi_{\text{exp}}^2 + \lambda_{\text{BC}} \sum_f \sum_{i=1}^{N_x} \left(q_f^{(p/A)}(x_i, Q_0, A=1) - q_f^{(p)}(x_i, Q_0) \right)^2$$

$$+ \lambda_{\text{pos}} \sum_{l=1}^{N_{\text{pos}}} \sum_{j=1}^{N_A} \sum_{i_l=1}^{N_{\text{dat}}^{(l)}} \max\left(0, -\mathcal{F}_{i_l}^{(l)}(A_j)\right),$$



MAPFF1.0 — Fragmentation Functions

Parameterisation

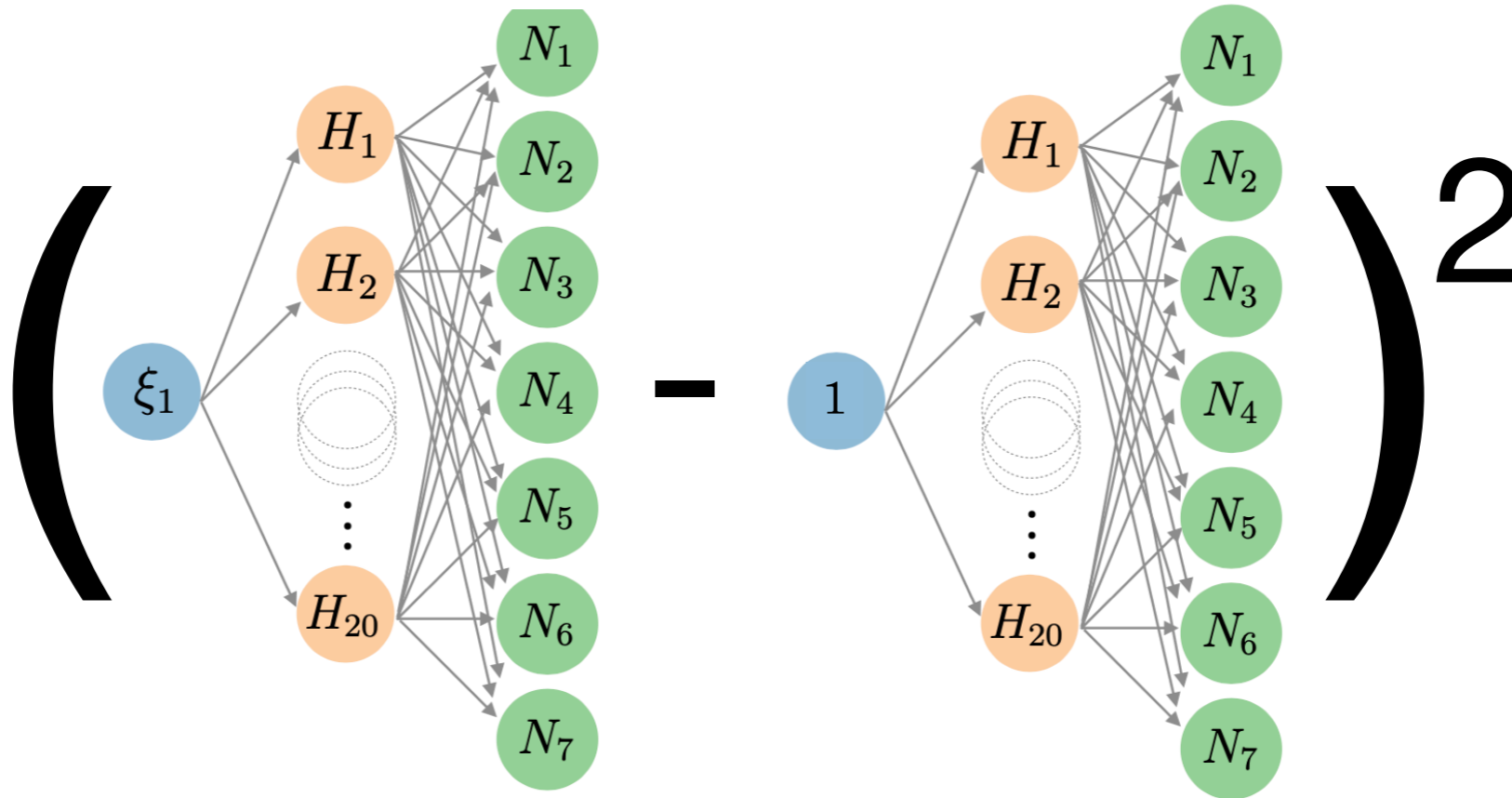
MAPFF1.0 — arXiv:2105.08725

7 flavours assuming a partially symmetric sea

$$\pi^+: \{u, \bar{d}, d = \bar{u}, s^+, c^+, b^+, g\}$$

$$K^+: \{u, \bar{s}, s = \bar{u}, d^+, c^+, b^+, g\}$$

$$zD_i^{\pi^+}(z, \mu_0 = 5 \text{ GeV}) = \left(N_i(z; \boldsymbol{\theta}) - N_i(1; \boldsymbol{\theta}) \right)^2$$

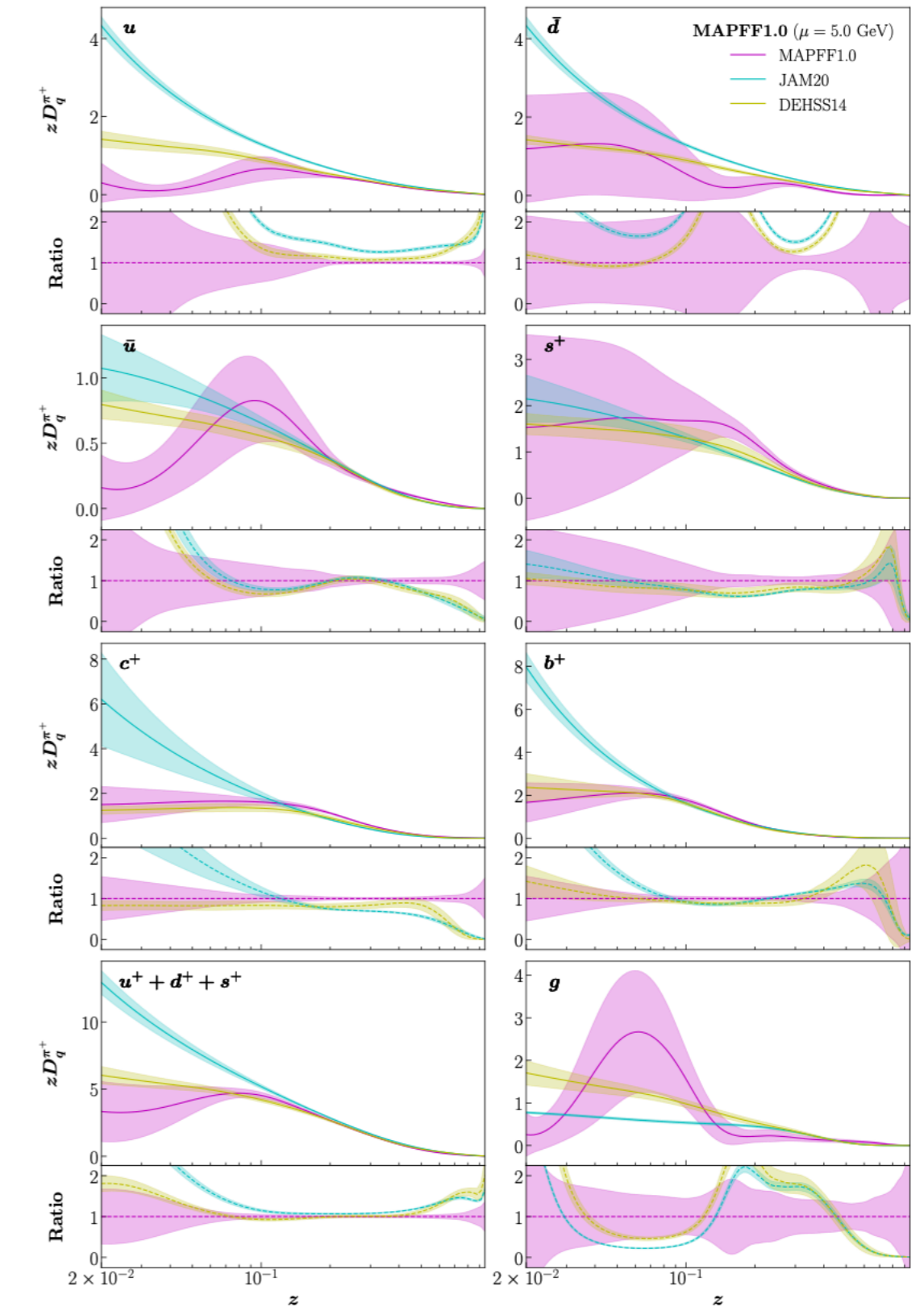
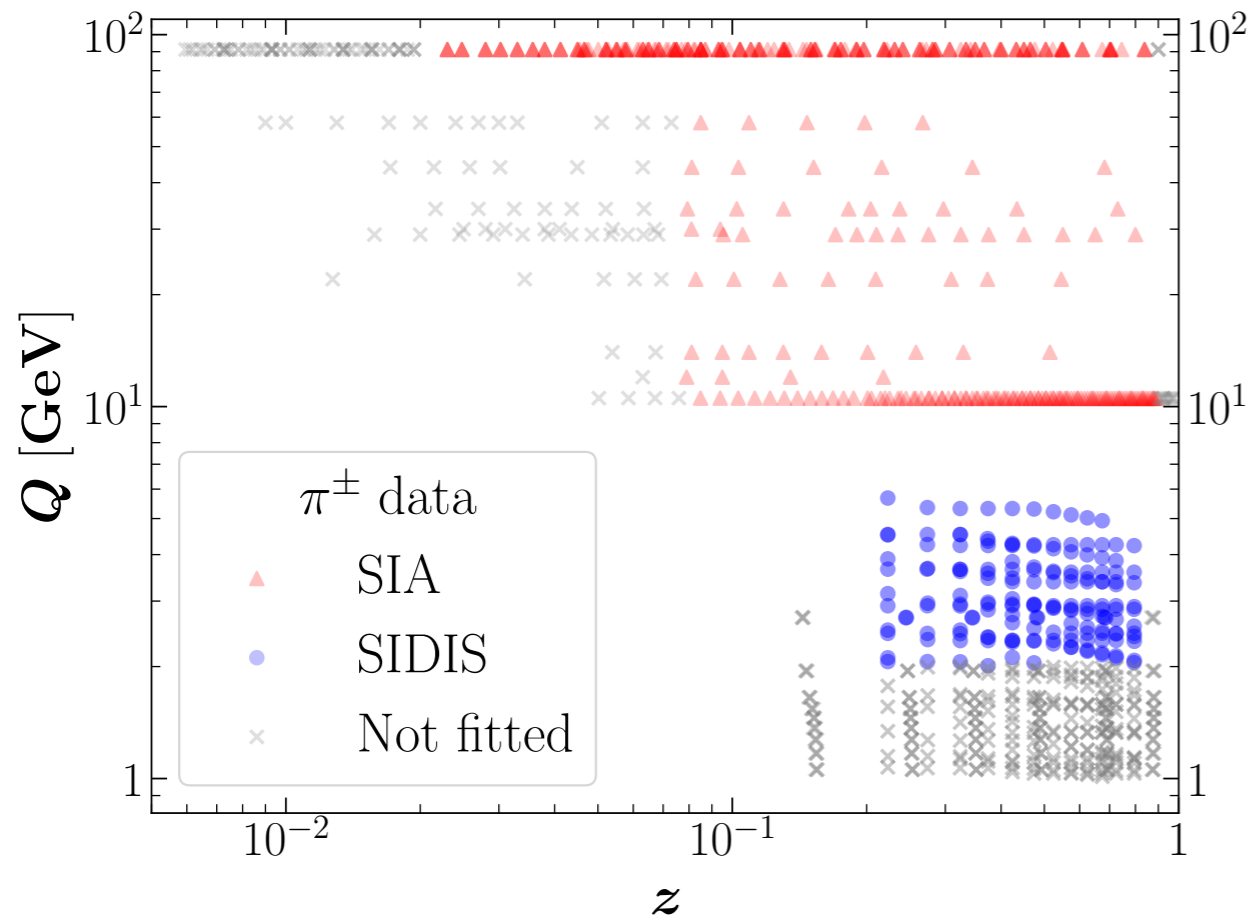


$$F_i(x, z, Q) = x \sum_{q\bar{q}} e_q^2 \left\{ [C_{i,qq}(x, z, Q) \otimes f_q(x, Q) + C_{i,qg}(x, z, Q) \otimes f_g(x, Q)] \otimes D_q^{\pi^\pm}(z, Q) \right. \\ \left. + [C_{i,gq}(x, z, Q) \otimes f_q(x, Q)] \otimes D_g^{\pi^\pm}(z, Q) \right\}, \quad i = 2, L.$$

Loss function

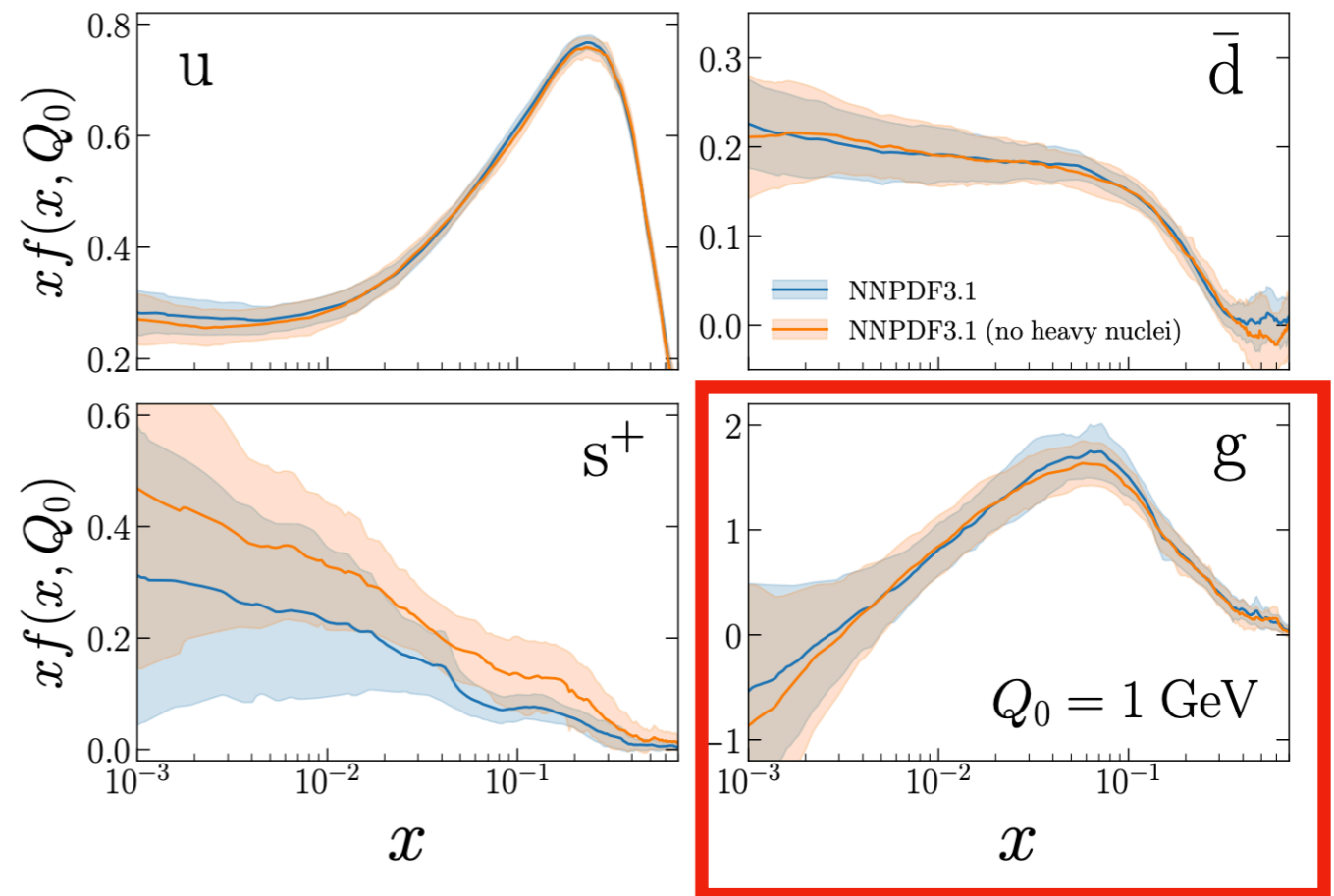
$$\chi^2 = \chi_{\text{exp}}^2$$

MAPFF1.0 — Fragmentation Functions



Tutorial 2 — Fitting the gluon PDF

Time for
Tutorial 2!



Tutorial 2: Non-linear regression with neural networks

Author and references

Rabah Abdul Khalek - khalek@jlab.org (Postdoc at JLab | Theory division)

https://github.com/rabah-khalek/TF_tutorials

Learning Goals

This notebook will serve as an introduction to non-linear regression using neural networks. We will learn how to build and train neural networks with TensorFlow, a powerful machine learning library.