

Lattice Gauge Theory insights on Dark Matter

Enrico Rinaldi

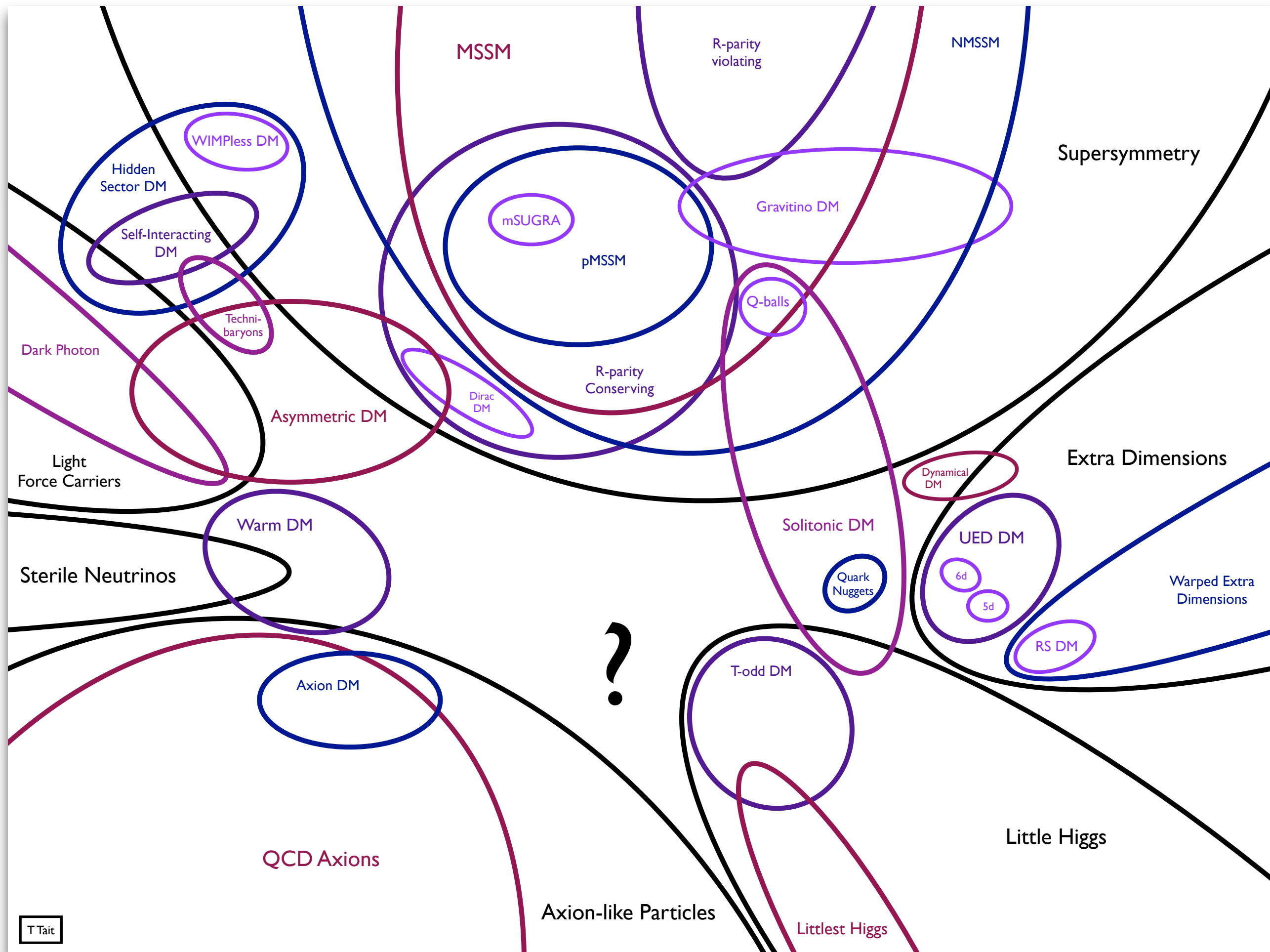


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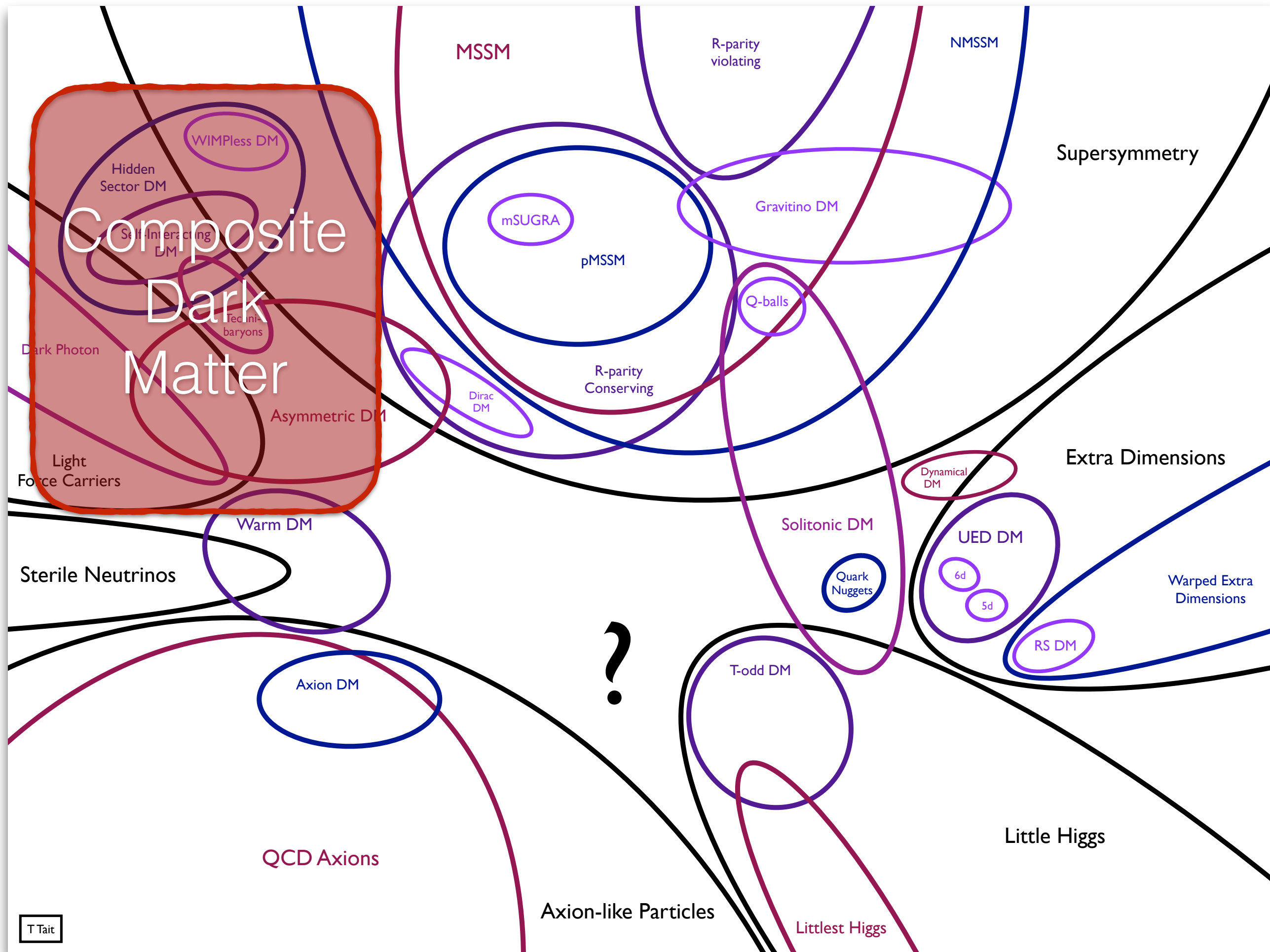
This research was performed under the auspices of the U.S. Department of Energy by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344 and supported by the LLNL LDRD "Illuminating the Dark Universe with PetaFlops Supercomputing" 13-ERD-023.

Computing support comes from the LLNL Institutional Computing Grand Challenge program.

What is Dark Matter?

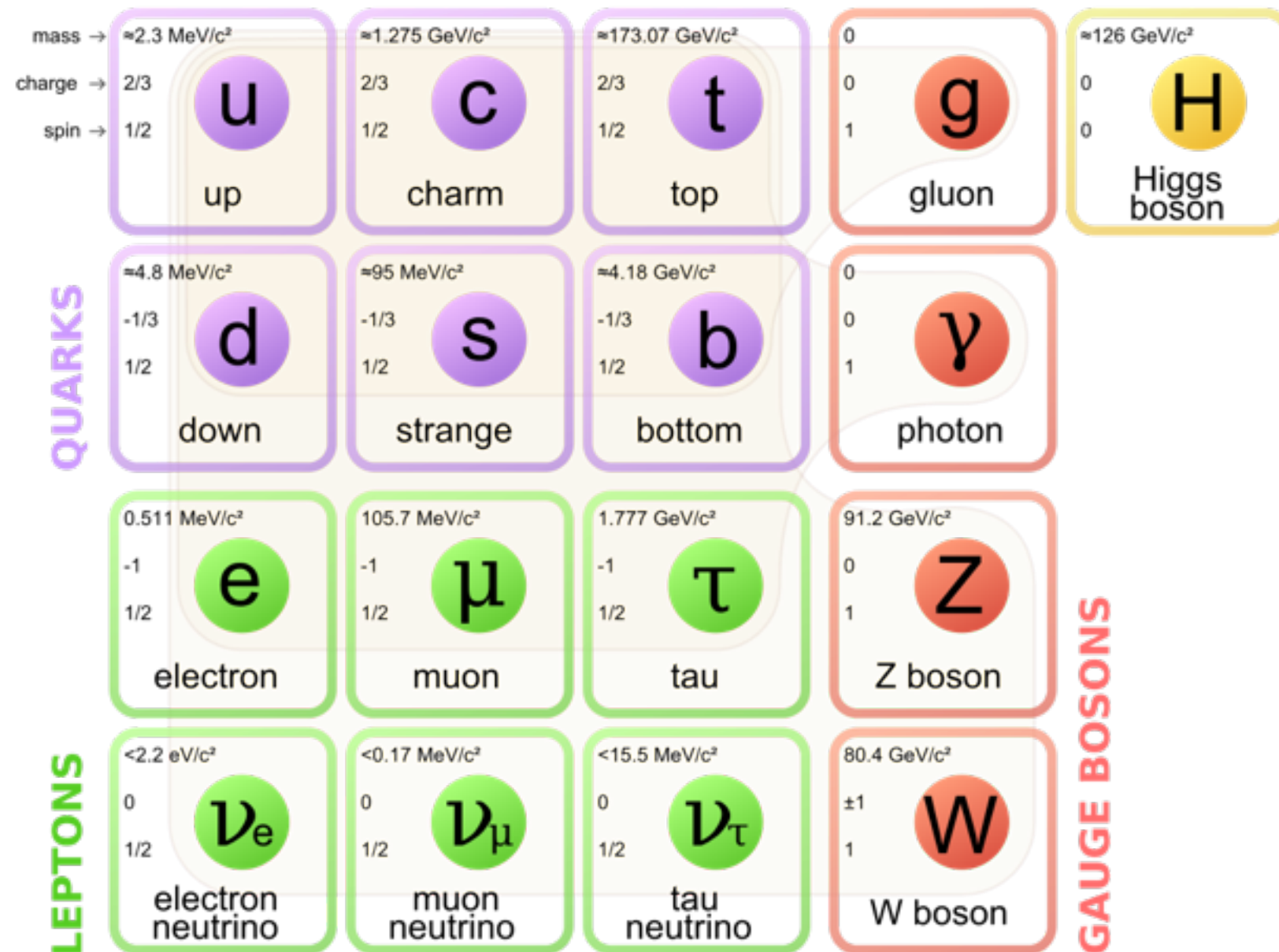


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- ★ Gravitational effects of DM show up in CMB, lensing and other large scale phenomena
 - ★ Direct Standard Model interactions are needed for production in the early Universe
 - ★ Direct detection and Collider experiments rely on SM interactions, but they are suppressed
 - ★ Strong exclusion bounds push theorists to explore a wider landscape of models for DM
 - ★ Problems with cosmological models can hint at strongly self-interacting dark matter



A very familiar picture

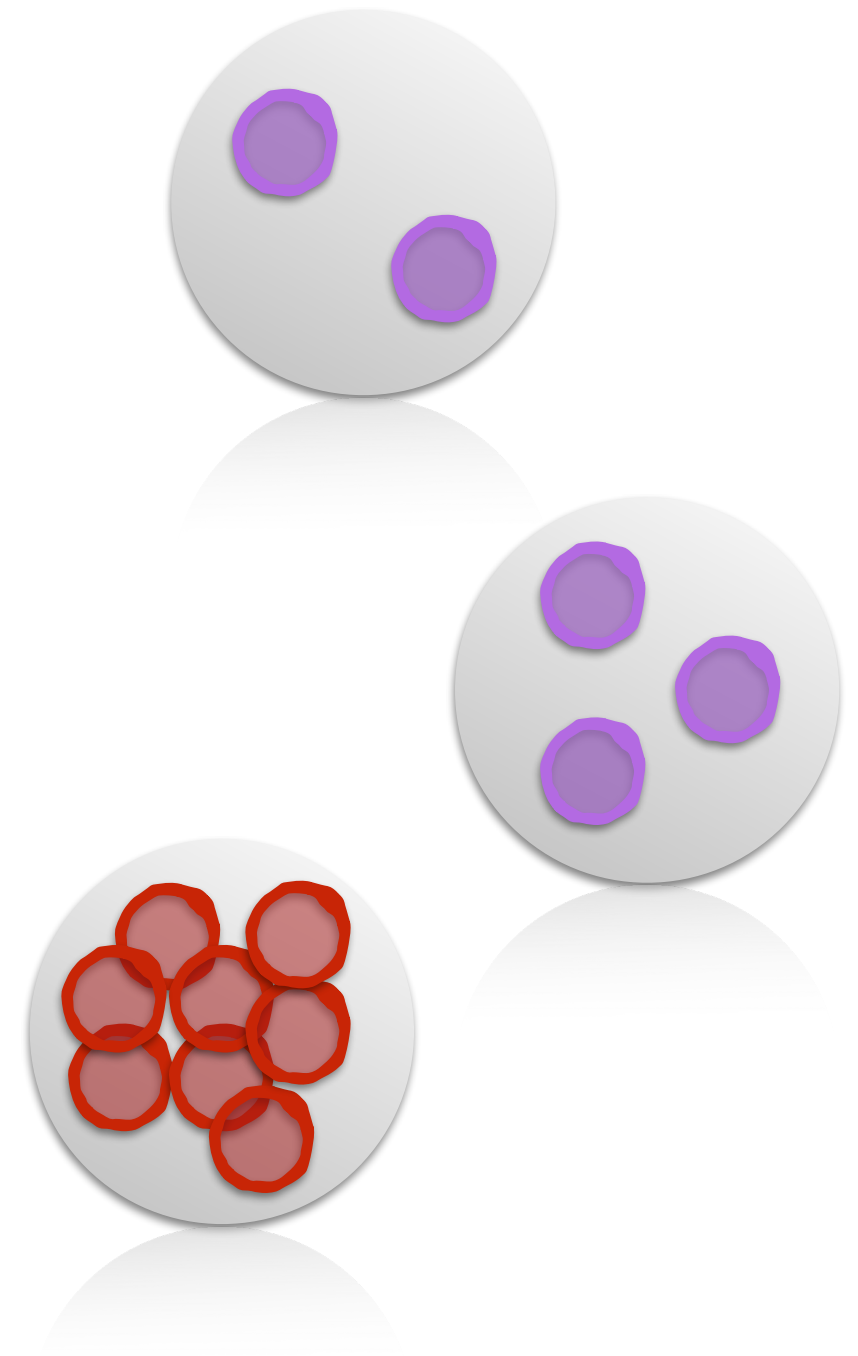
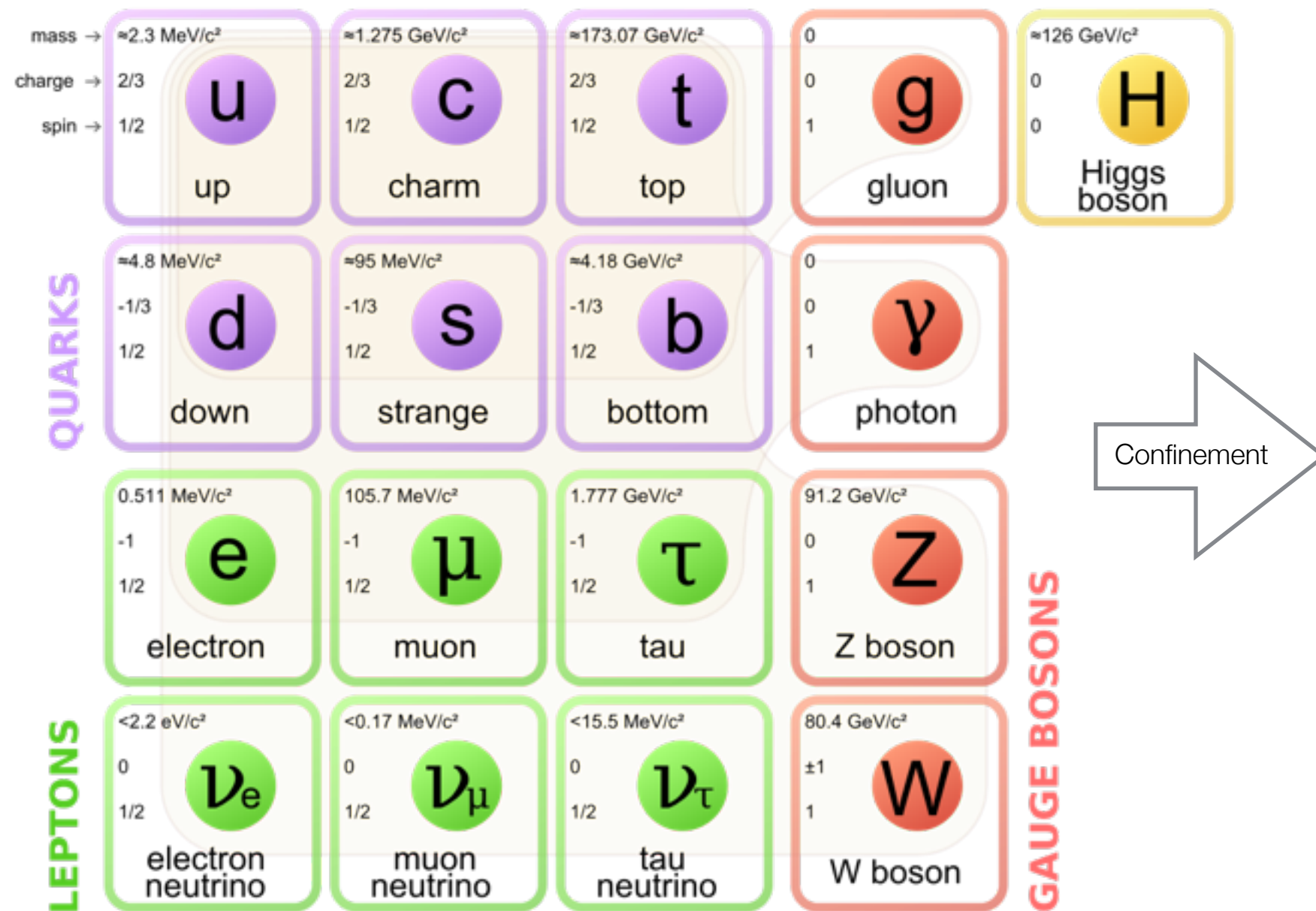
The Standard Model of particles



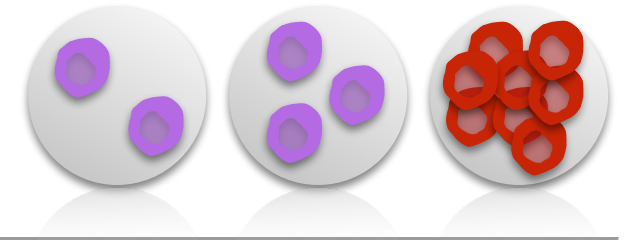
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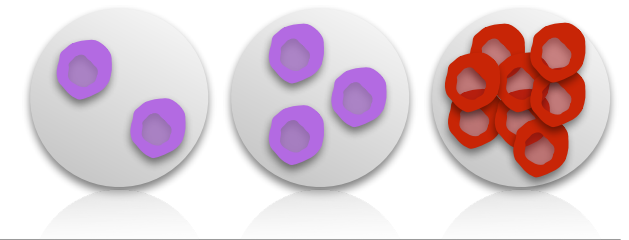
Mesons, Baryons and Glueballs



Composite Dark Matter

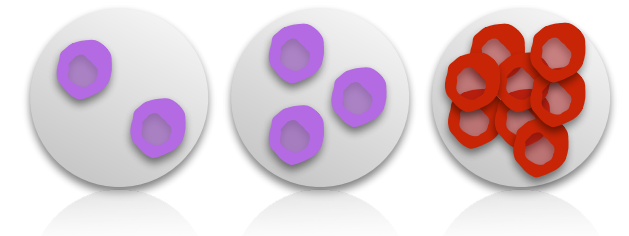


Composite Dark Matter



- ◆ Dark Matter is a **composite** object

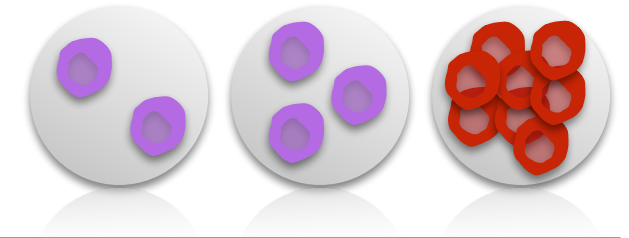
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e.g. **technibaryon** or
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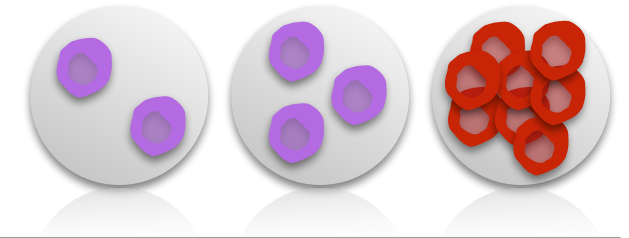
Composite Dark Matter



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- ◆ Interesting and complicated internal **structure**
- ◆ Properties dictated by **strong dynamics**
- ◆ Self-interactions are natural

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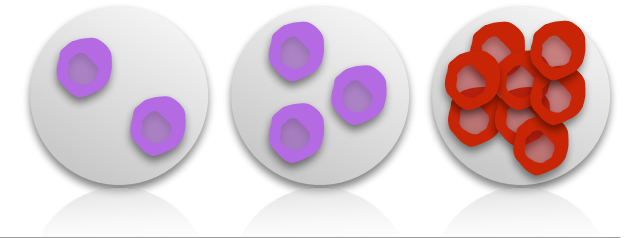


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Similar to **QCD**

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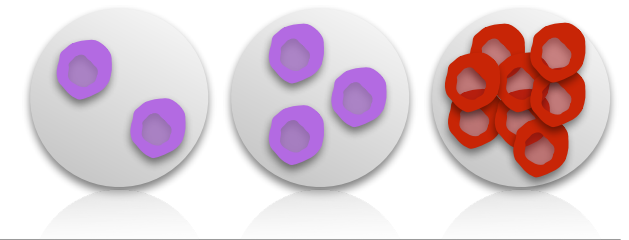


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- ♦ DM composite is **neutral and stable**
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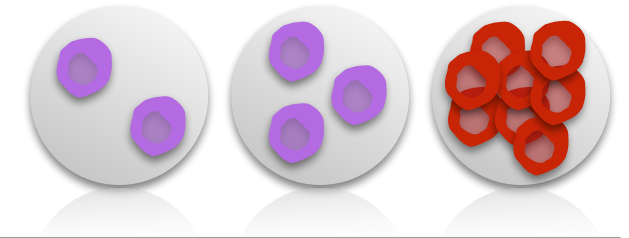
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Lattice Field Theory methods

- ♦ Properties dictated by **strong dynamics**

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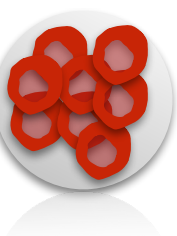
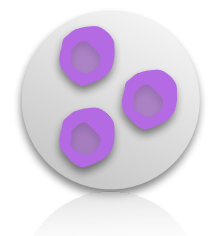
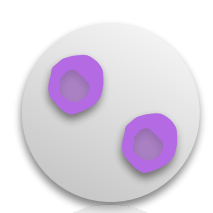
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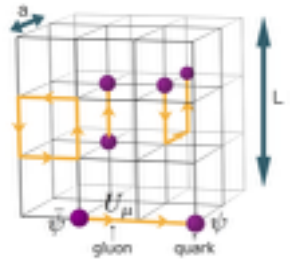
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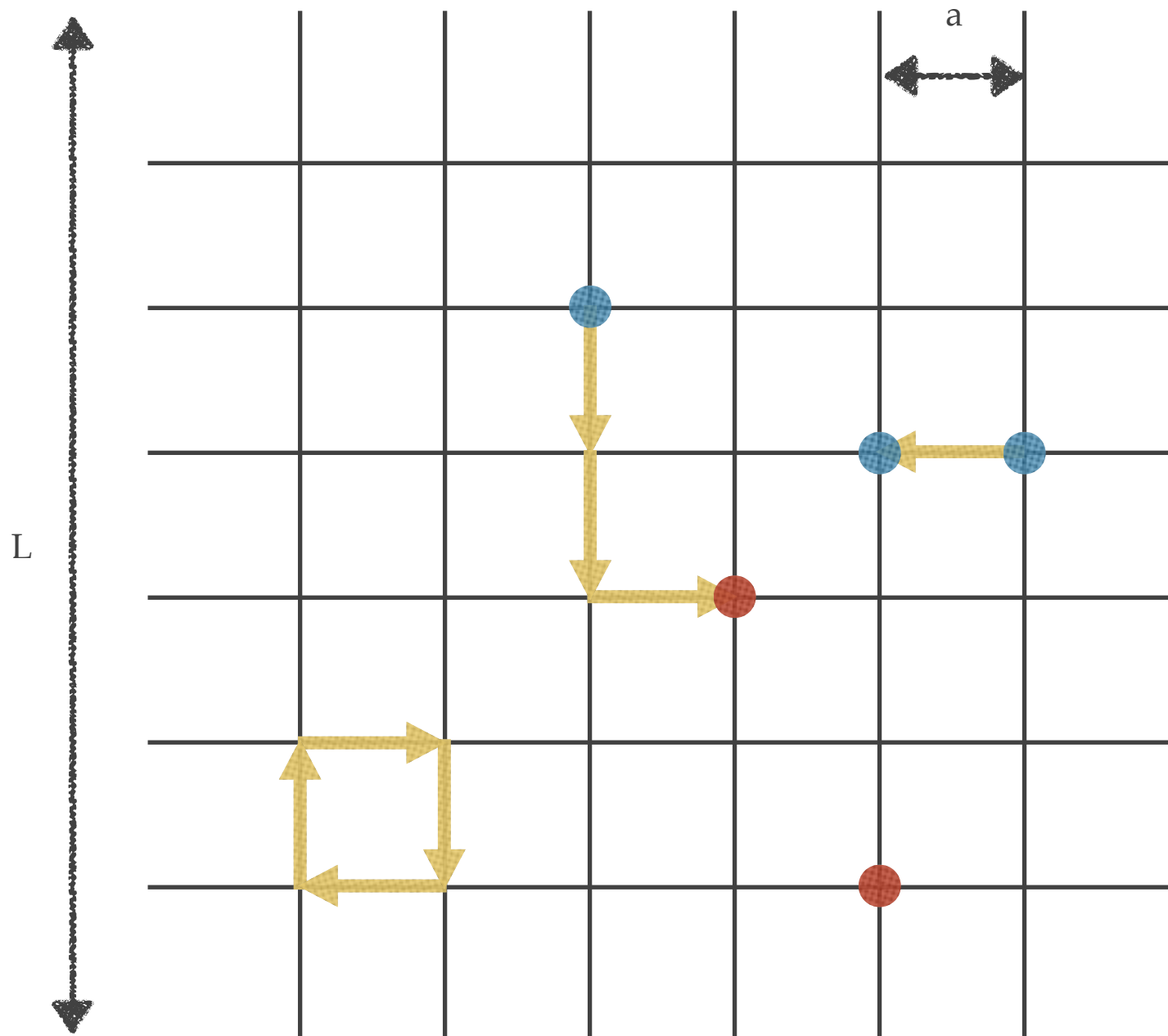
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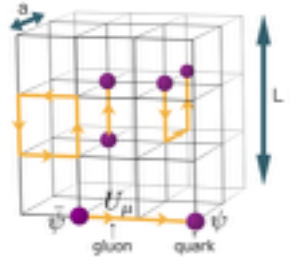
Self-interactions are included due to **strongly coupled** dynamics



Lattice Gauge Theory - basics



- Discretize space and time
 - lattice spacing “a”
 - lattice size “L”
- Keep all d.o.f. of the theory
 - not a model!
 - no simplifications
- Amenable to numerical methods
 - Monte Carlo sampling
 - use supercomputers
- Precisely quantifiable and improvable errors
 - Systematic
 - Statistical



Importance of lattice field theory simulations

- ◆ *lattice simulations are needed to solve the strong dynamics*
- ◆ naturally suited for models where dark fermion masses are comparable to the **confinement scale**
- ◆ **controllable** systematic errors and room for **improvement**
- ◆ Naive dimensional analysis and EFT approaches can miss important **non-perturbative** contributions
- ◆ NDA is **not precise enough** when confronting experimental results and might not work for certain situations: there are uncontrolled theoretical errors

Models for Composite Dark Matter

★ Pion-like (dark quark-antiquark)

- ◆ pNGB DM [*Hietanen et al.*, 1308.4130]
- ◆ Quirky DM [*Kribs et al.*, 0909.2034]
- ◆ Ectocolor DM [*Buckley&Neil*, 1209.6054]
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- ◆ Minimal SU(2)

★ Glueball-like (only gluons)

- ◆ SUNonia [*Boddy et al.*, 1402.3629]
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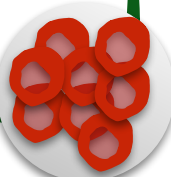
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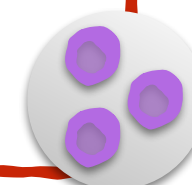
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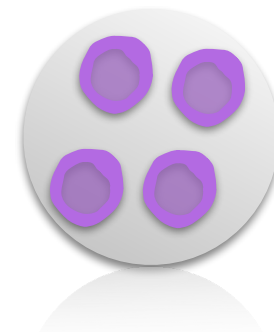


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“Stealth Dark Matter” Model

- ◆ **New strongly-coupled SU(4) gauge sector** “like” QCD with a **plethora of composite states** in the spectrum: all mass scales are technically natural for hadrons
- ◆ New **Dark fermions**: have **dark color** and also have **electroweak charges** ($W/Z, \gamma$)
- ◆ Dark fermions have **electroweak breaking masses** (Higgs) and **electroweak preserving masses** (not-Higgs)
- ◆ A global symmetry naturally stabilizes the **dark lightest baryonic** composite states (e.g. dark neutron)

“Stealth Dark Matter” model

- The field content of the model consists in *8 Weyl fermions*
- Dark fermions interact with the SM Higgs and obtain **current/chiral masses**
- Introduce **vector-like masses** for dark fermions that do not break EW symmetry
- Diagonalizing in the mass eigenbasis gives *4 Dirac fermions*
- Assume **custodial SU(2) symmetry** arising when **$u \leftrightarrow d$**

EW interactions




| Field | $SU(N)_D$ | $(SU(2)_L, Y)$ | Q |
|--|-------------------------|----------------------|--|
| $F_1 = \begin{pmatrix} F_1^u \\ F_1^d \end{pmatrix}$ | \mathbf{N} | $(\mathbf{2}, 0)$ | $\begin{pmatrix} +1/2 \\ -1/2 \end{pmatrix}$ |
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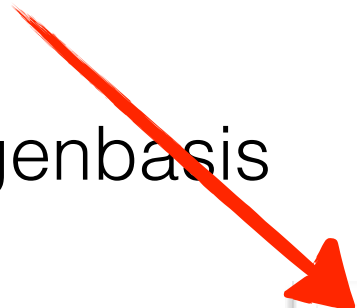
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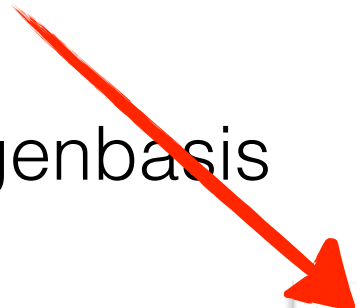
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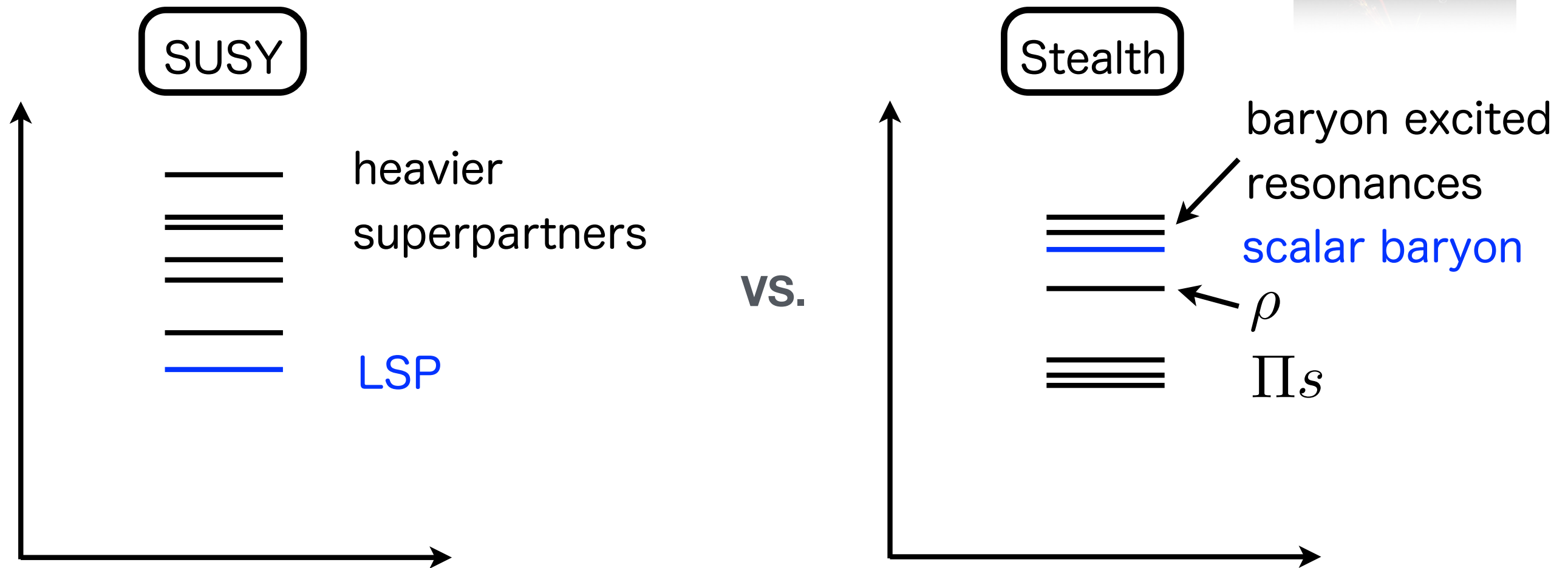
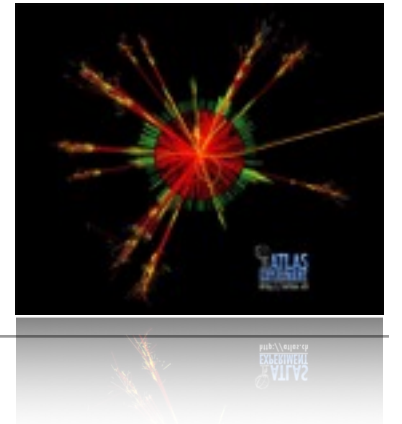


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Stealth DM at colliders

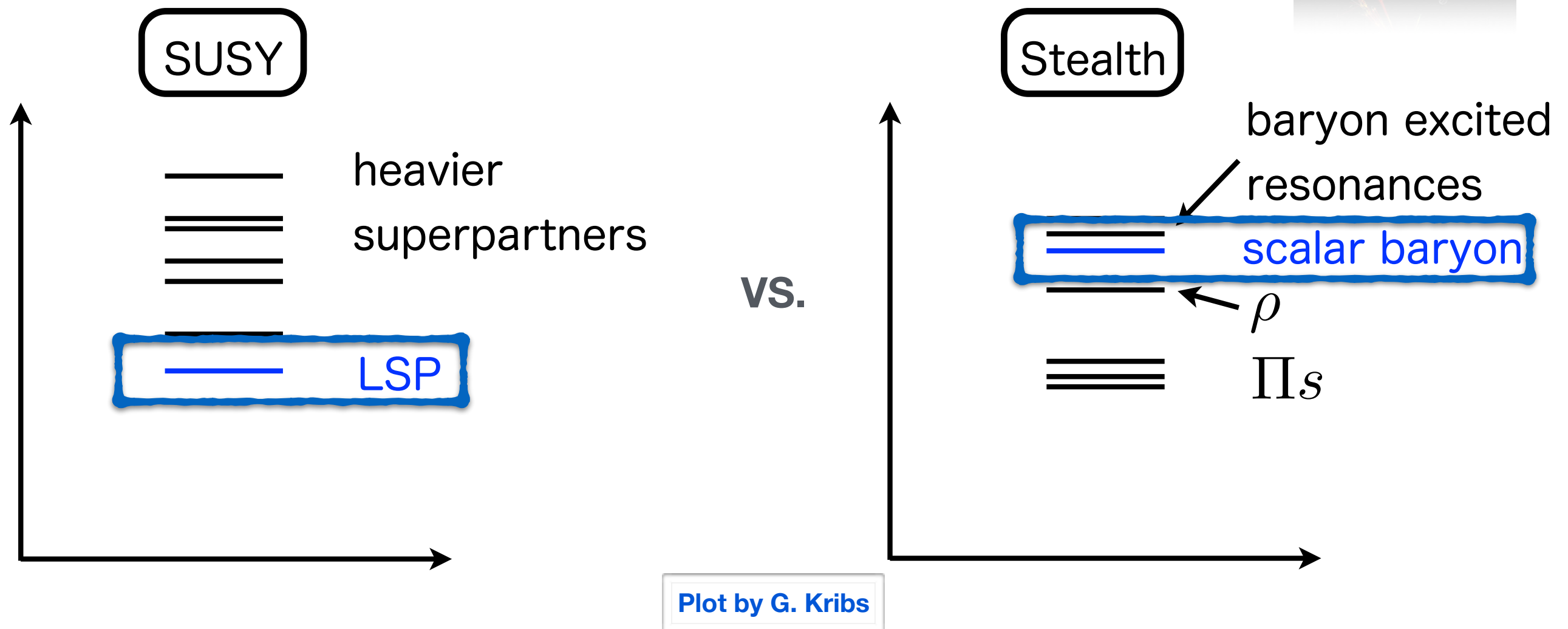
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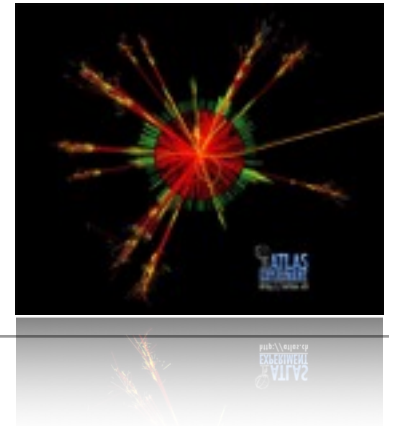
Plot by G. Kribs



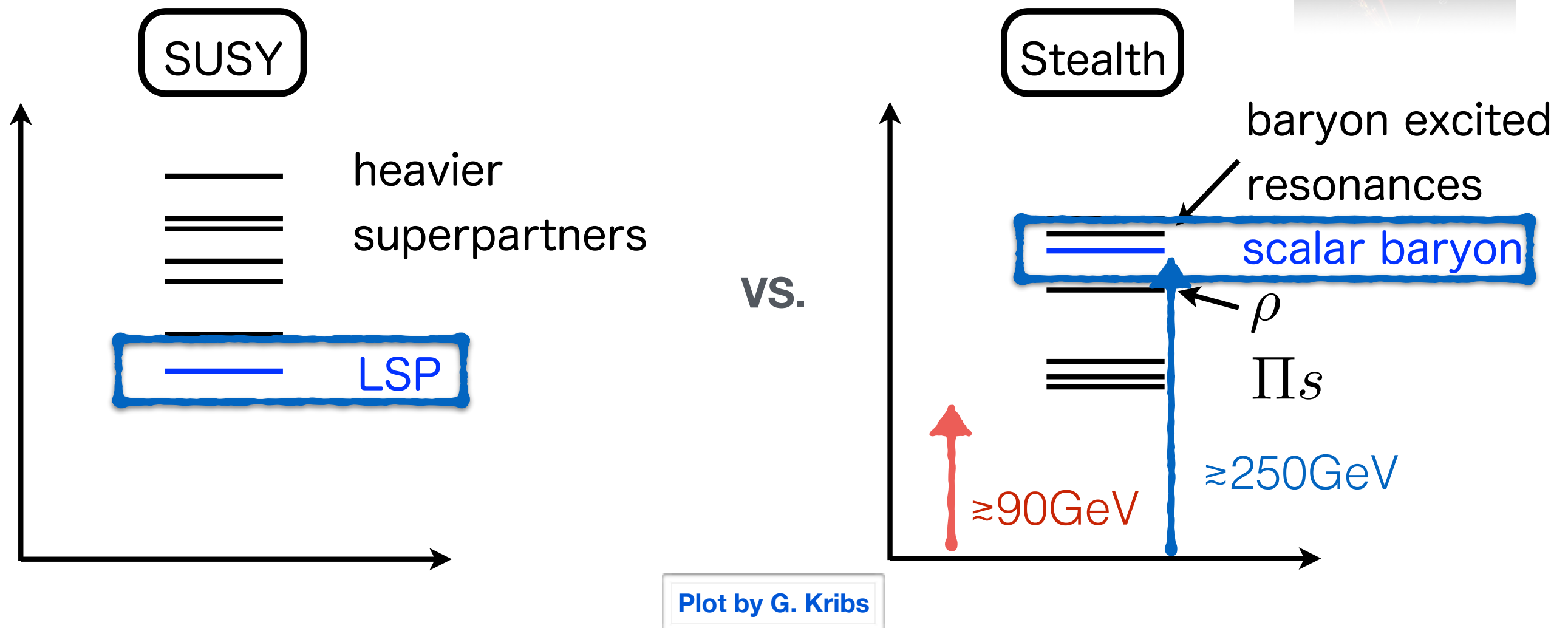
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Stealth DM at colliders



- ✦ Signatures are not dominated by missing energy: **DM is not the lightest particle!** The interactions are suppressed (form factors)
- ✦ Dark mesons production and decay give interesting signatures: **the model can be constrained by collider limits!**

Photon interactions

$$\langle \chi(p') | j_{\text{EM}}^\mu | \chi(p) \rangle = F(q^2) q^\mu$$

Expansion at low momentum through effective operators

◆ dimension 5 \Rightarrow magnetic dipole

◆ dimension 6 \Rightarrow charge radius

◆ dimension 7 \Rightarrow polarizability

$$\frac{(\bar{\chi} \sigma^{\mu\nu} \chi) F_{\mu\nu}}{\Lambda_{\text{dark}}}$$

$$\frac{(\bar{\chi} \chi) v_\mu \partial_\nu F^{\mu\nu}}{\Lambda_{\text{dark}}^2}$$

$$\frac{(\bar{\chi} \chi) F_{\mu\nu} F^{\mu\nu}}{\Lambda_{\text{dark}}^3}$$

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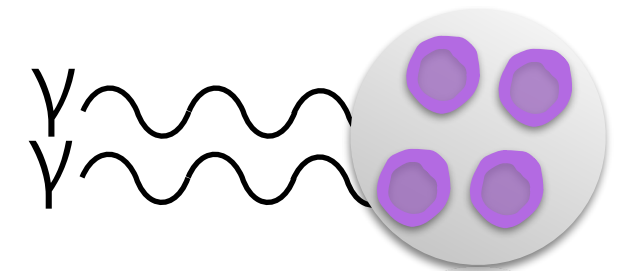
◆ dimension 6 \Rightarrow custodial

custodial SU(2)

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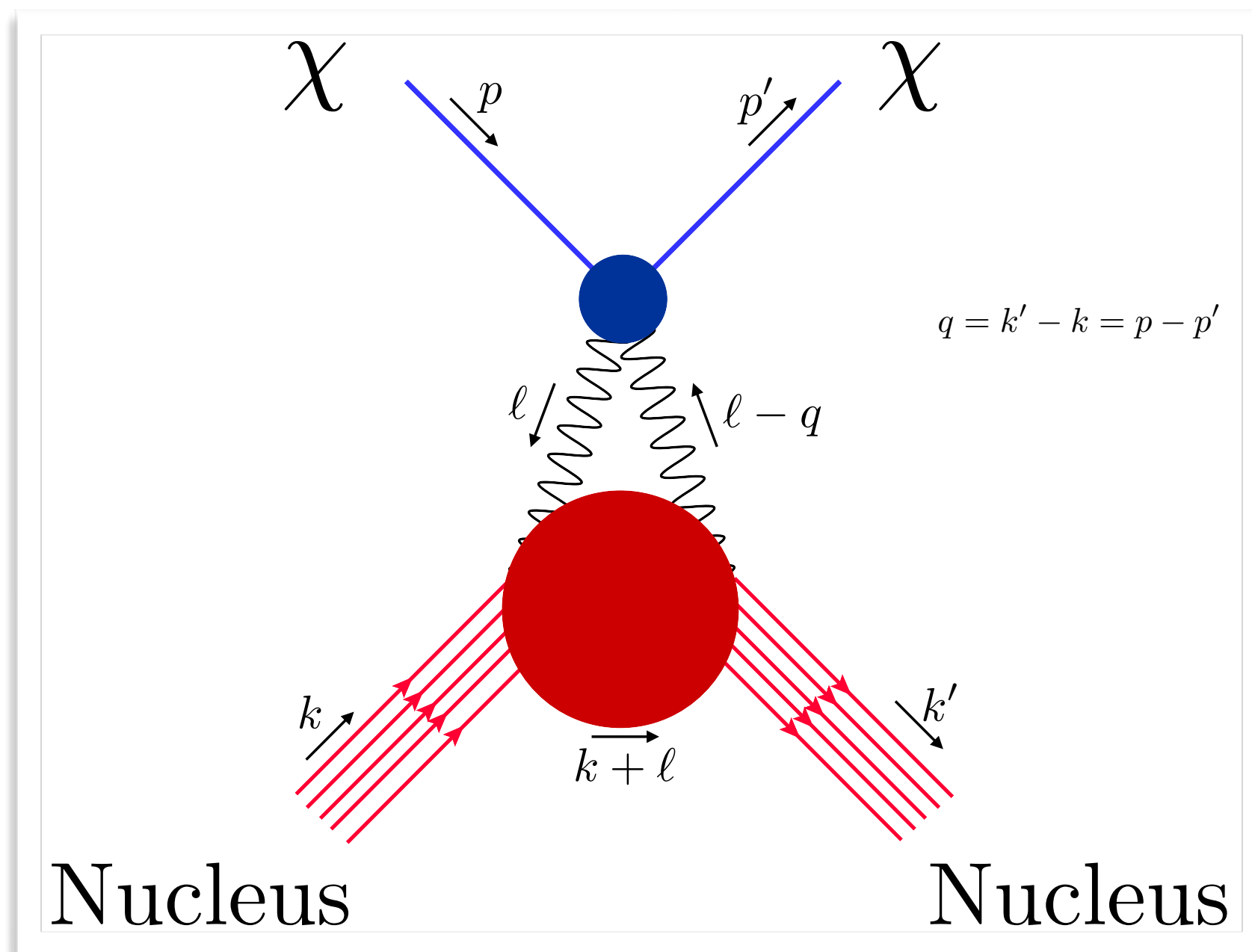
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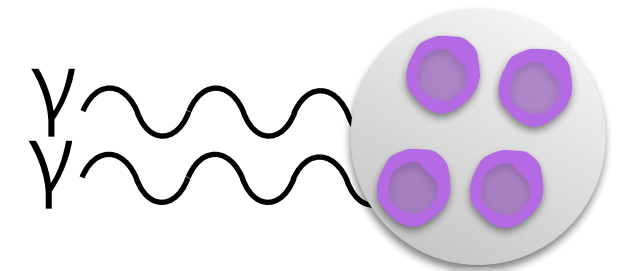
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Computing polarizability

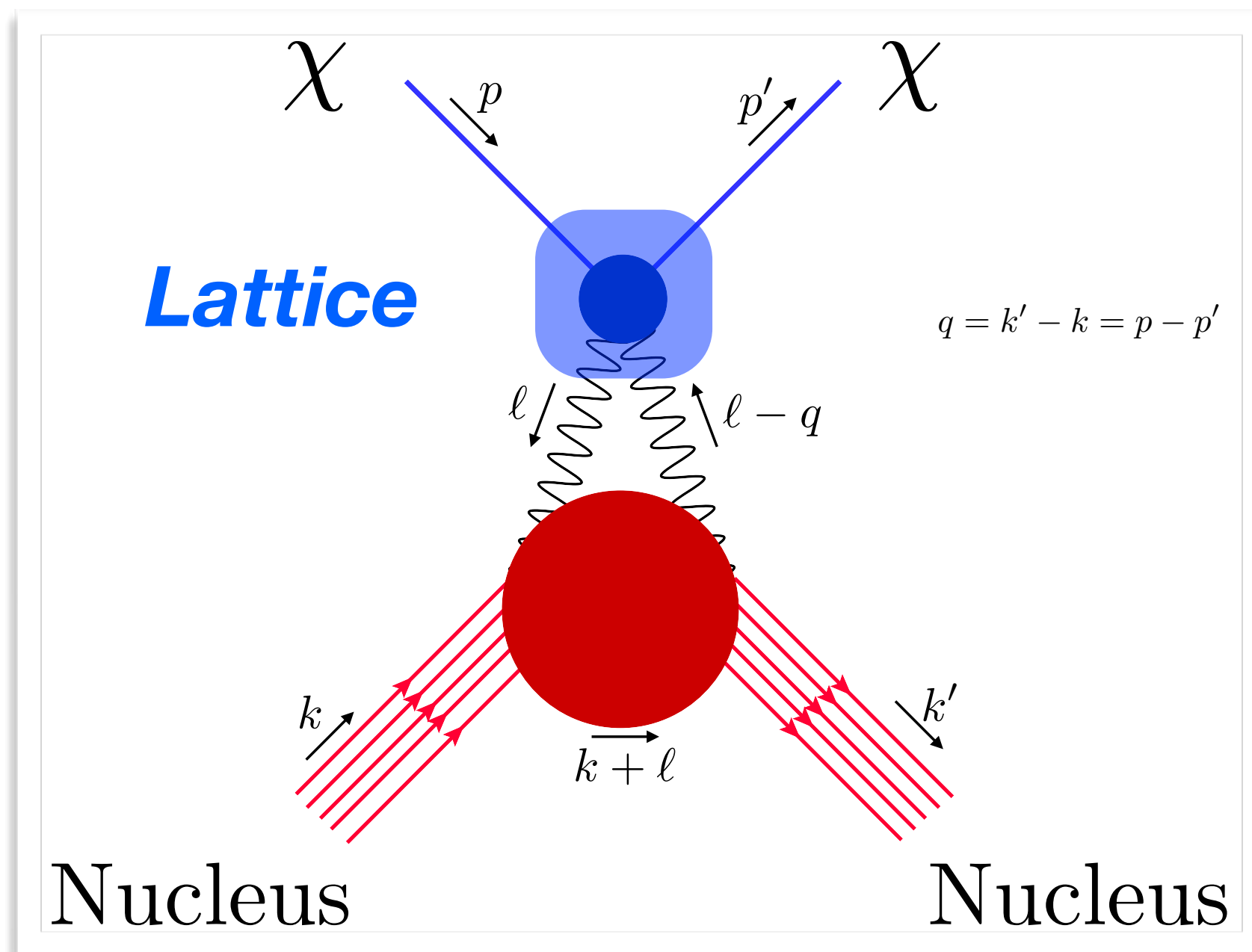
$$\frac{c_F e^2}{m_\chi^3} \chi^* \chi F^{\mu\alpha} F_\alpha^\nu v_\mu v_\nu$$

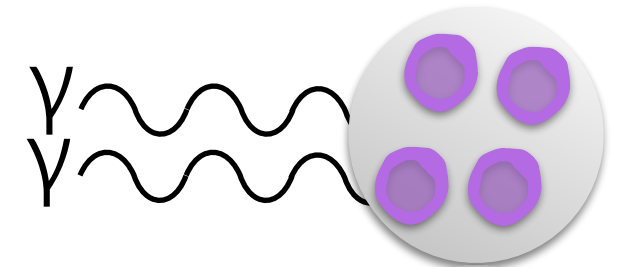




Computing polarizability

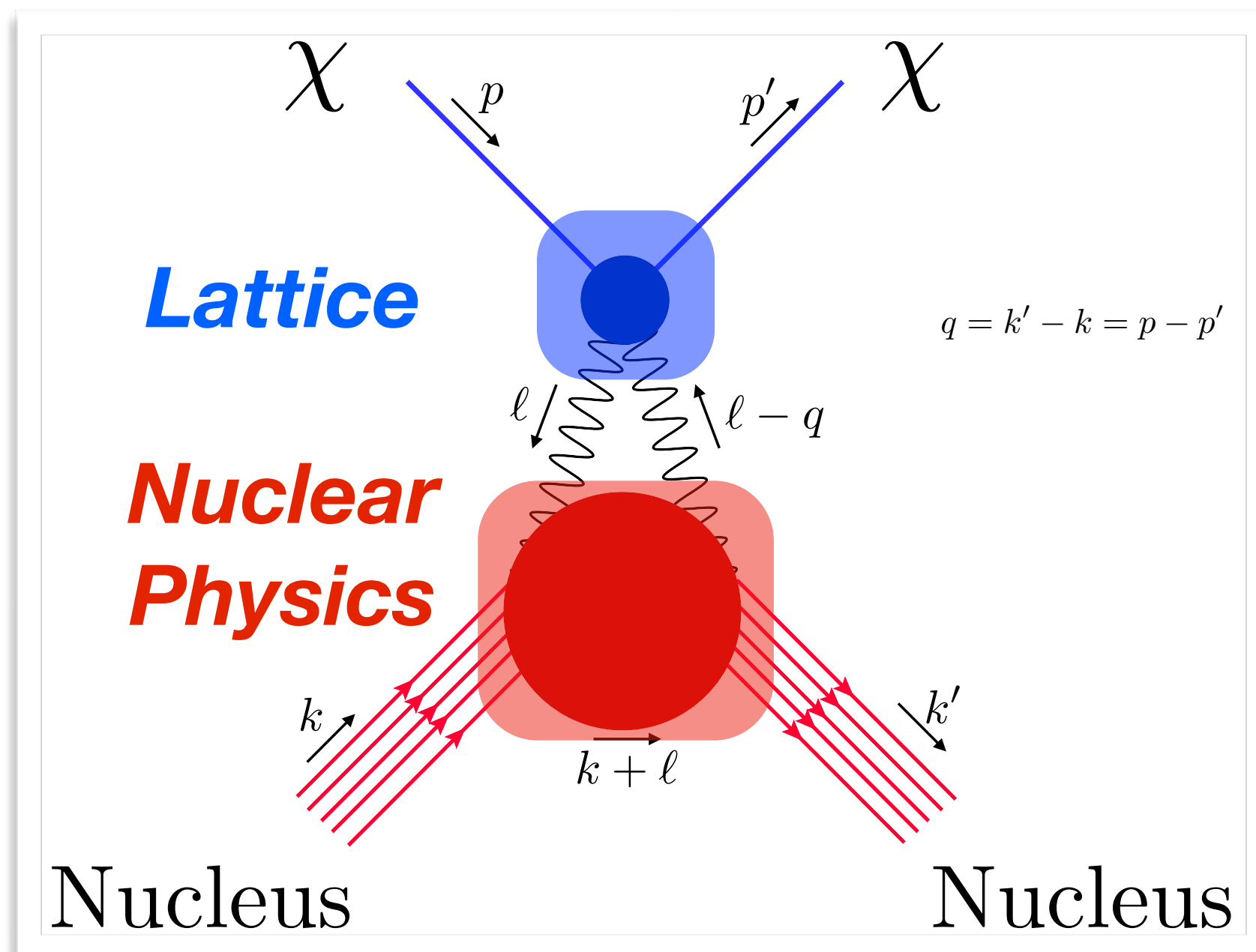
$$\frac{c_F e^2}{m_\chi^3} \chi^* \chi F^{\mu\alpha} F_\alpha^\nu v_\mu v_\nu$$



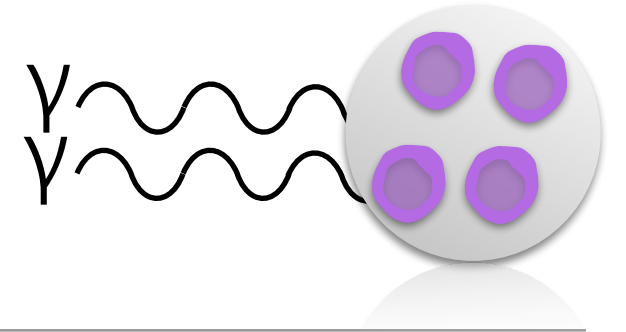


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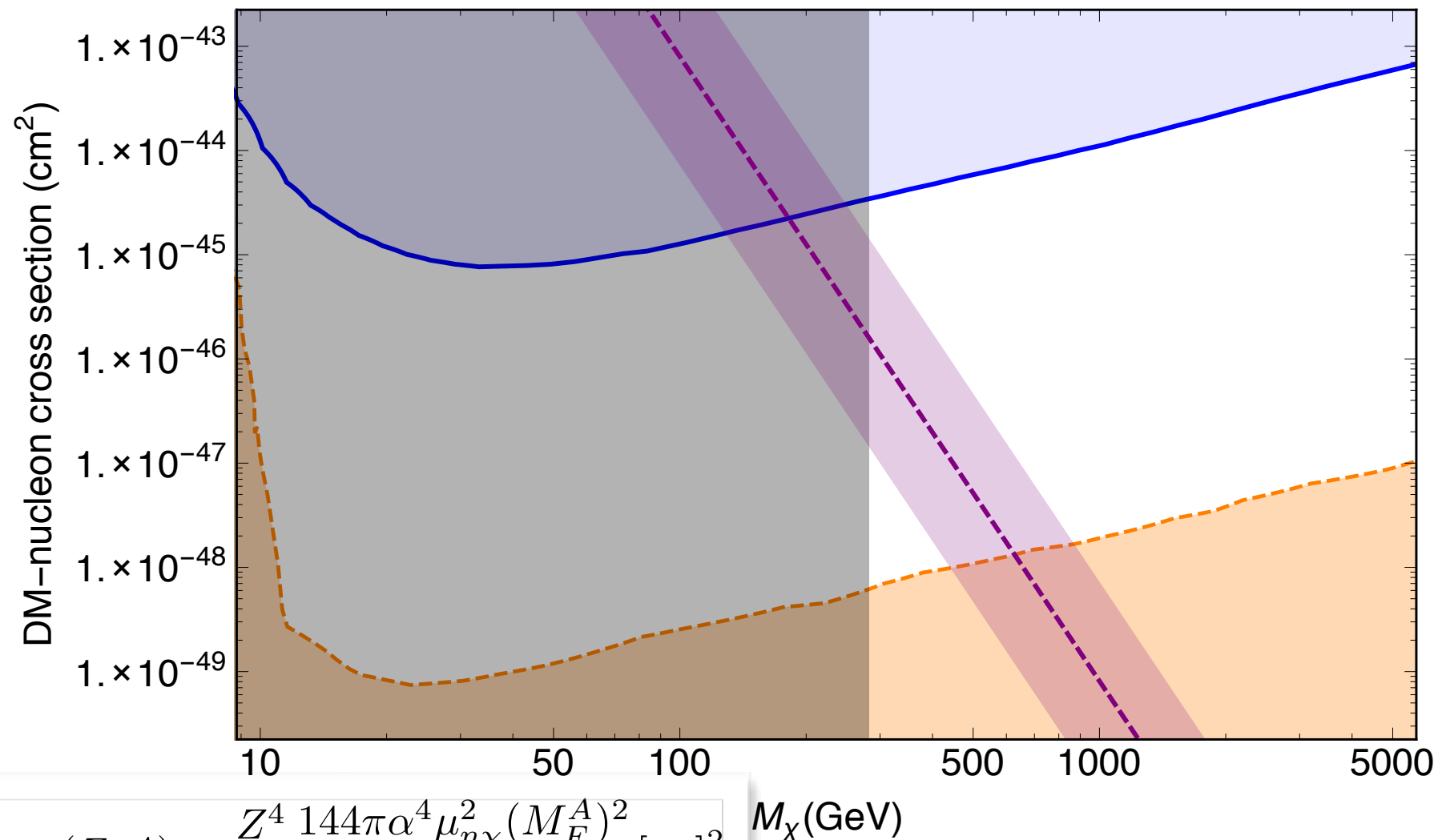
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Lowest bound from EM polarizability



Electric polarizability from lattice simulations with background fields

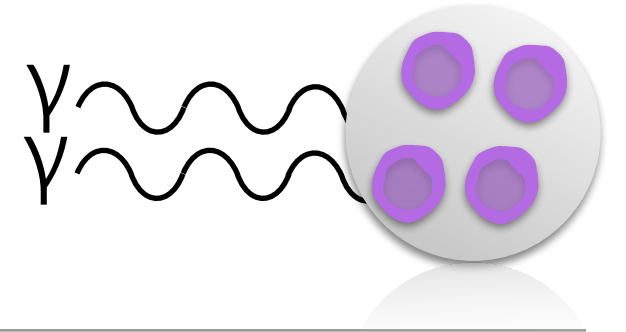


SU(4) $N_f=4$ Stealth DM

[LSD, 1503.04205]

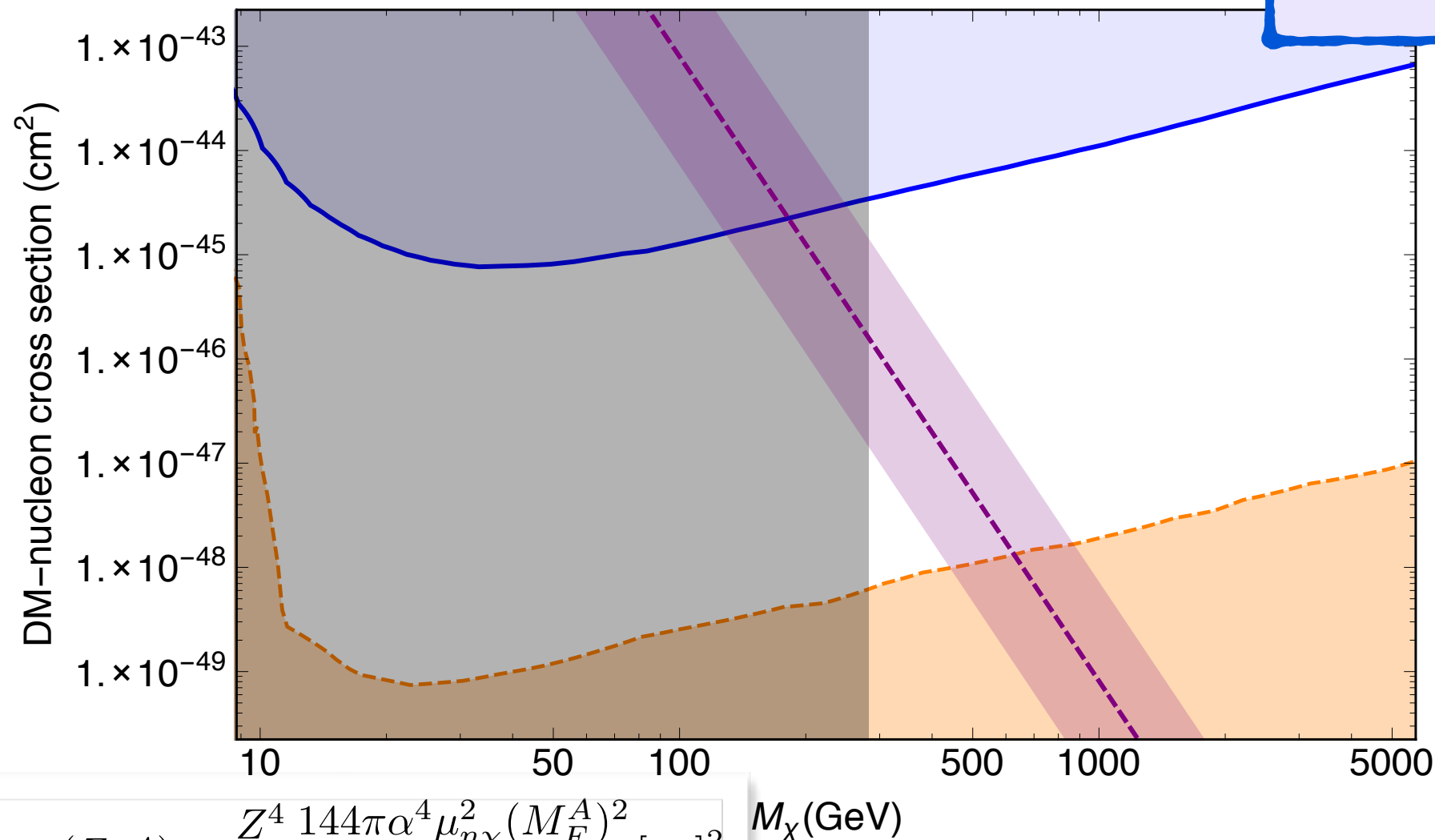
$$\sigma_{\text{nucleon}}(Z, A) = \frac{Z^4}{A^2} \frac{144\pi\alpha^4 \mu_{n\chi}^2 (M_F^A)^2}{m_\chi^6 R^2} [c_F]^2$$

Lowest bound from EM polarizability



Electric polarizability from lattice simulations with background fields

LUX exclusion bound for spin-independent cross section

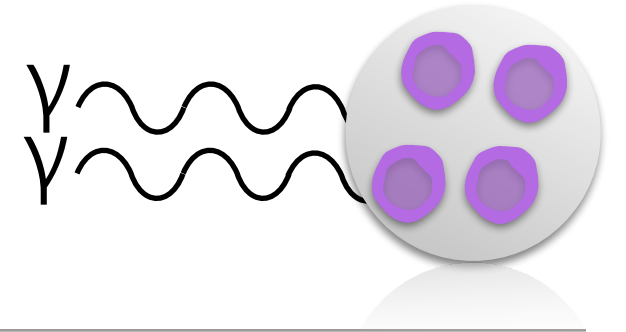


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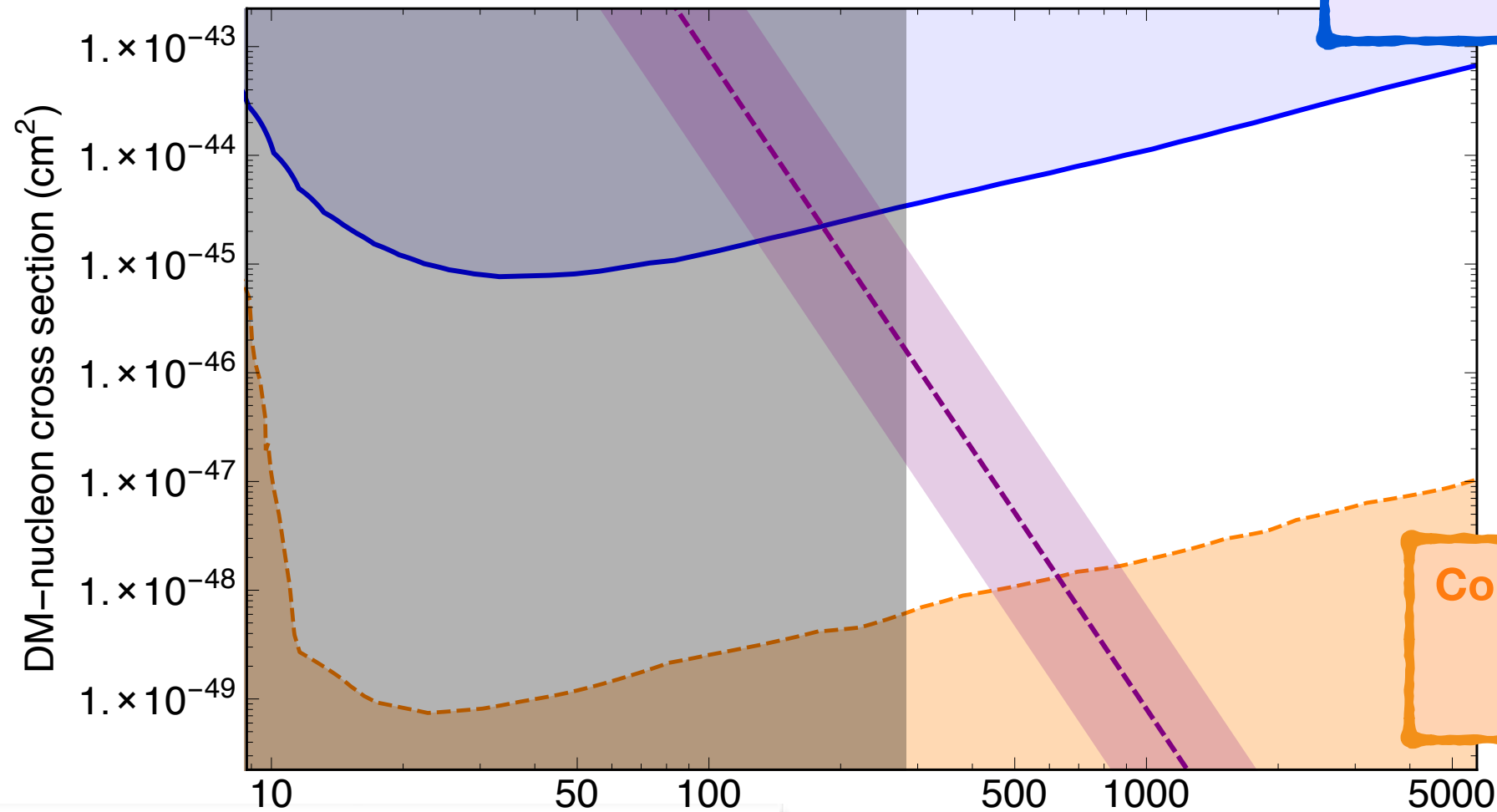
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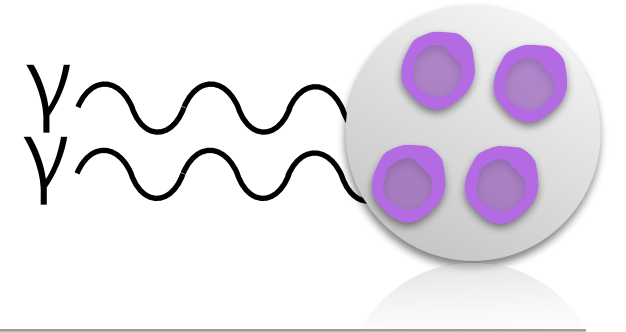
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Coherent neutrino scattering background

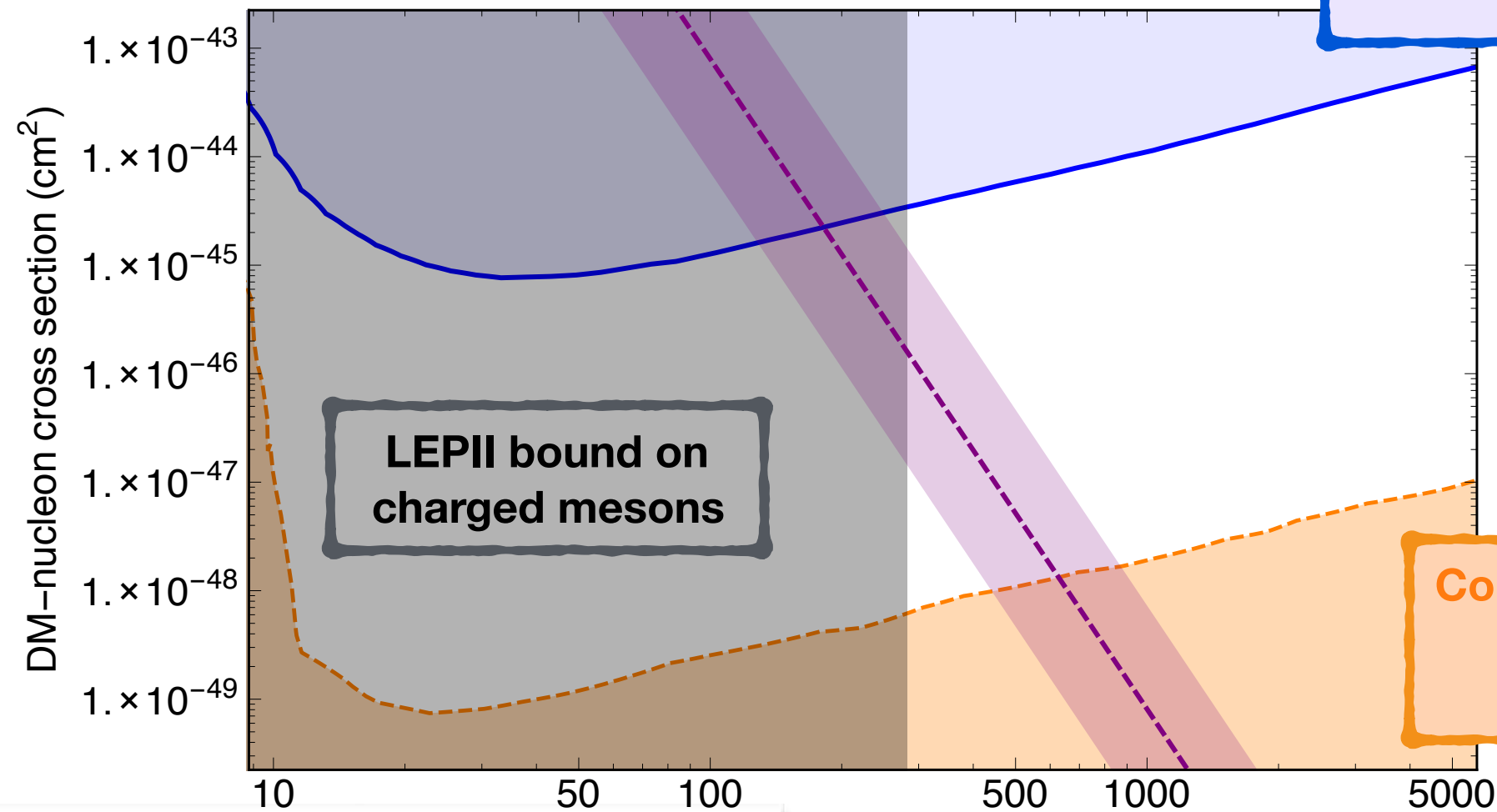
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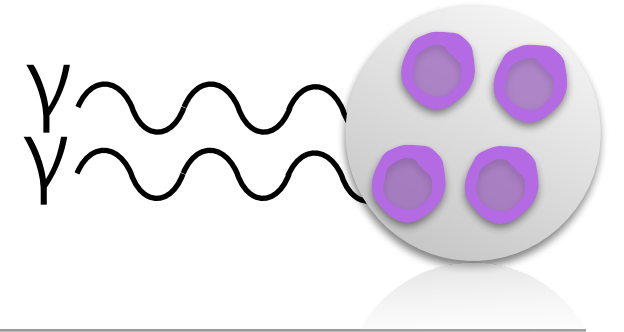
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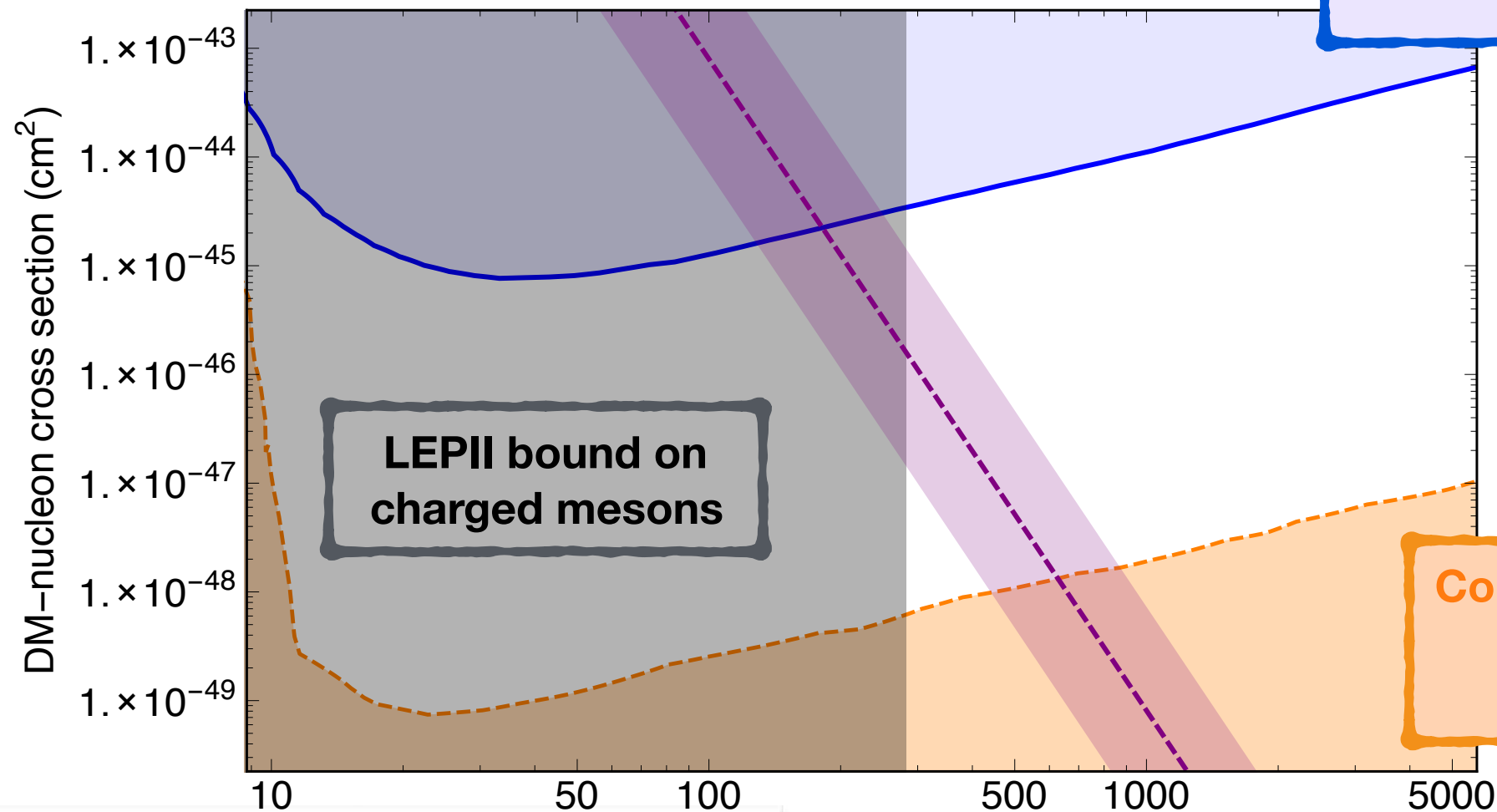
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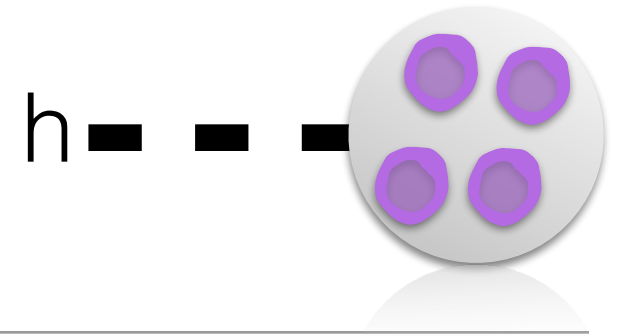
lowest allowed direct detection cross-section for composite dark matter theories with EW charged constituents

Concluding remarks

- ★ QCD ideas and lattice QCD techniques can be borrowed when exploring the DM landscape (BSM)
- ★ Composite dark matter is a viable interesting possibility with rich phenomenology
- ★ Lattice methods can help in calculating direct detection cross sections, production rates at colliders, and self-interaction cross sections of phenomenological relevance.
- ★ Dark matter constituents can carry electroweak charges and still the stable composites are currently undetectable. Stealth cross section.

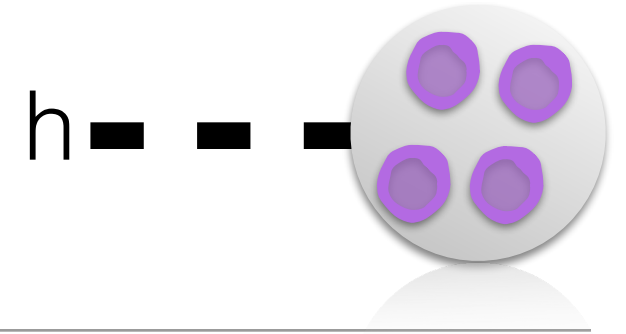
extra

Computing Higgs exchange



- ◆ Need to **non-perturbatively** evaluate the dark **σ -term**

$$\mathcal{M}_a = \frac{y_f y_q}{2m_h^2} \sum_f \langle B | \bar{f} f | B \rangle \sum_q \langle a | \bar{q} q | a \rangle$$

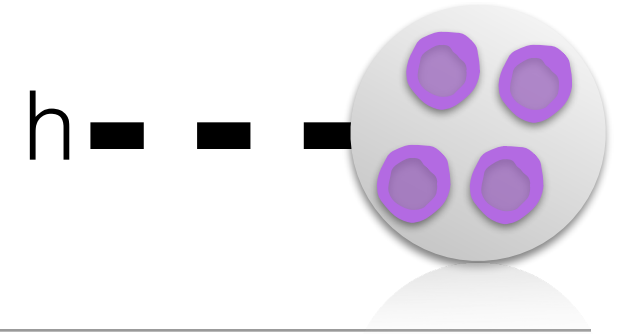


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2. dark baryon scalar form factor: need lattice input for generic DM models!
3. nucleon scalar form factor: ChPT and lattice input



Computing Higgs exchange

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- ◆ **Effective Higgs coupling** non-trivial with mixed chiral and vector-like masses

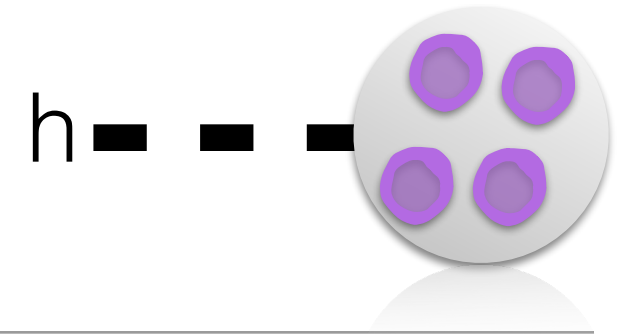
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$$y_f \langle B | \bar{f} f | B \rangle = \frac{m_B}{v} \sum_f \frac{v}{m_f} \left. \frac{\partial m_f(h)}{\partial h} \right|_{h=v} f_f^{(B)}$$

$$m_f(h) = m + \frac{y_f h}{\sqrt{2}}$$

$$\alpha \equiv \frac{v}{m_f} \left. \frac{\partial m_f(h)}{\partial h} \right|_{h=v} = \frac{y v}{\sqrt{2} m + y v}$$



Computing Higgs exchange

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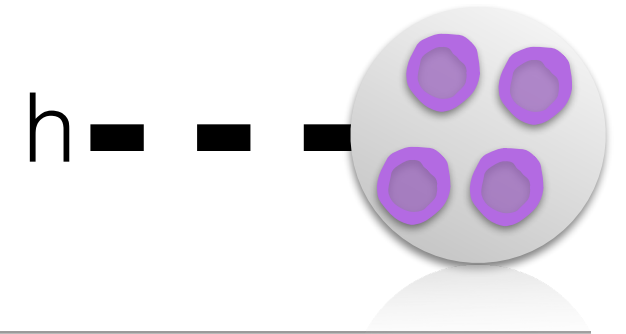
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- ◆ **Model-dependent** answer for the cross-section
- ◆ **Lattice input** is necessary: compute mass and form factor

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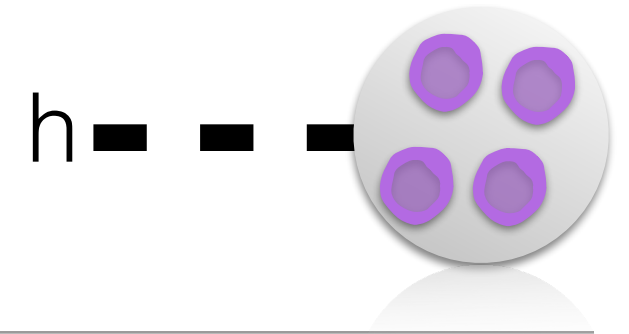
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Lattice!

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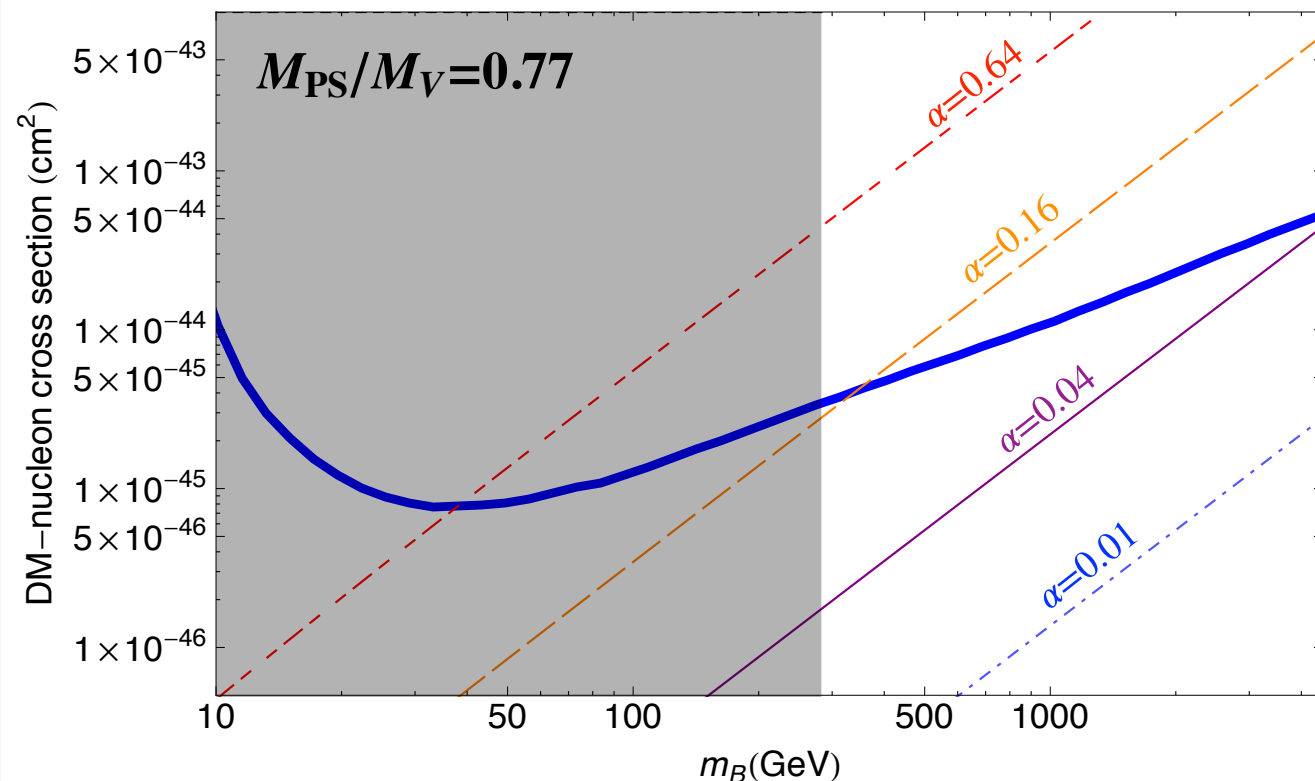
Bounds from Higgs exchange



- ◆ Lattice results for the cross-section are compared to **experimental** bounds
- ◆ Coupling space in specific models can be vastly constrained

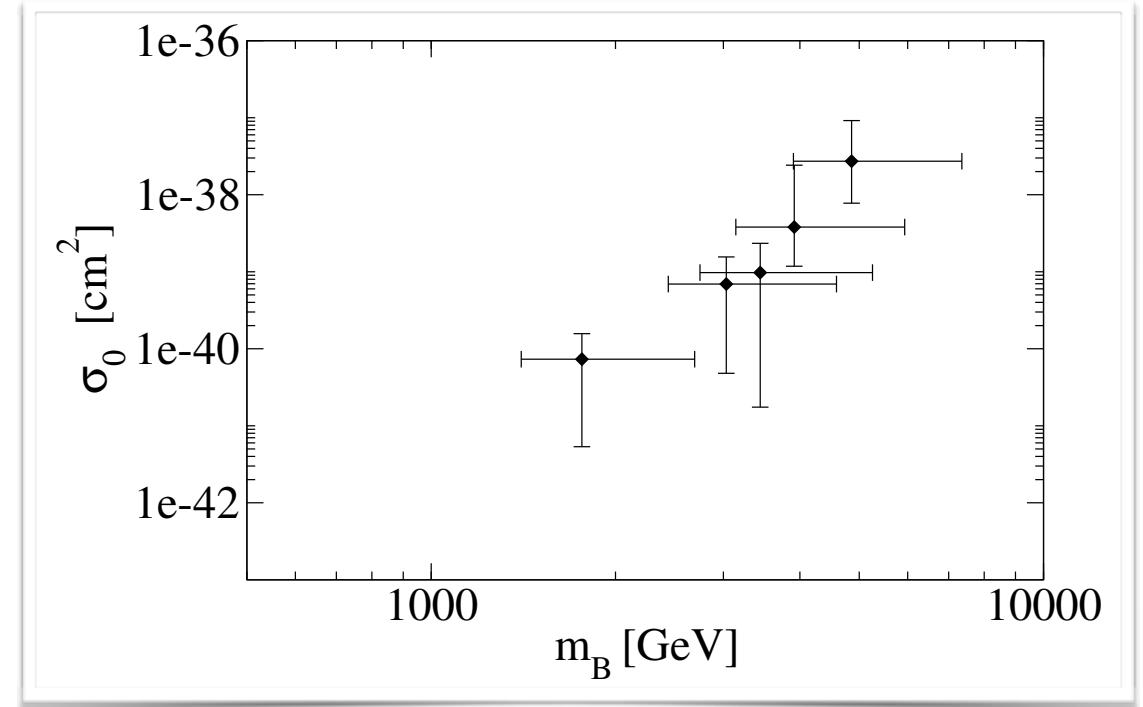
SU(4) $N_f=4$ Stealth DM

[LSD, 1402.6656-1503.04203]



SU(3) $N_f=8$ “technibaryon”

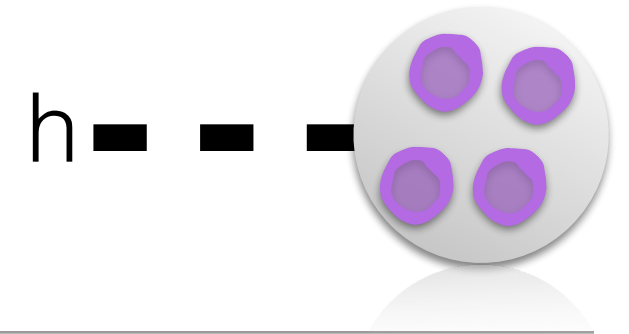
[LatKML, 1510.07373]



- ◆ Some candidates can be excluded as *dominant sources of dark matter
- ◆ There is lattice evidence for universality of dark scalar form factors

[DeGrand et al., 1501.05665]

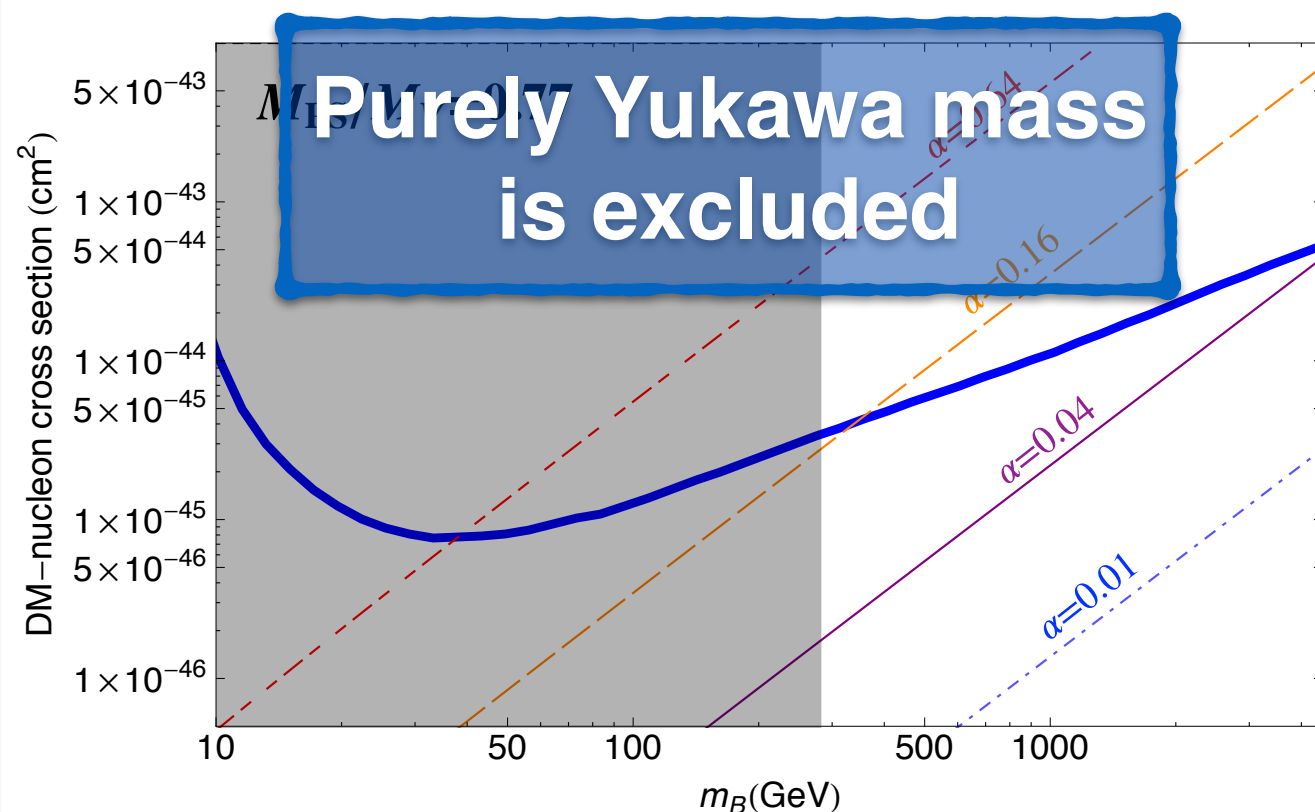
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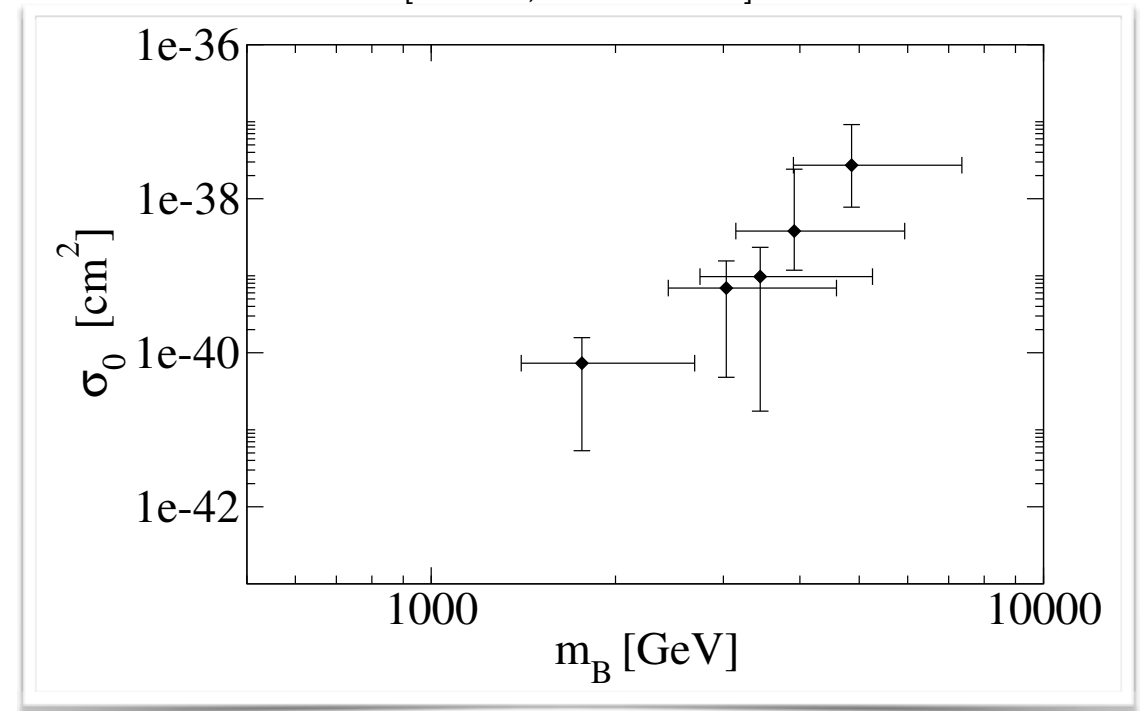
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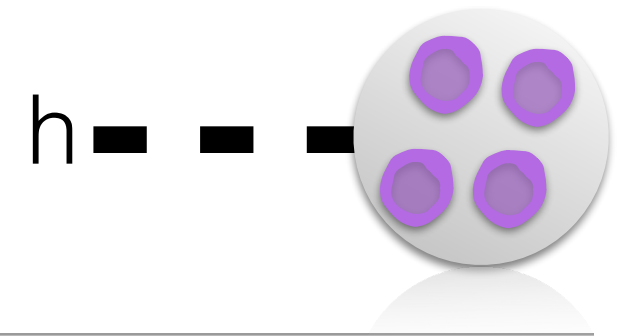
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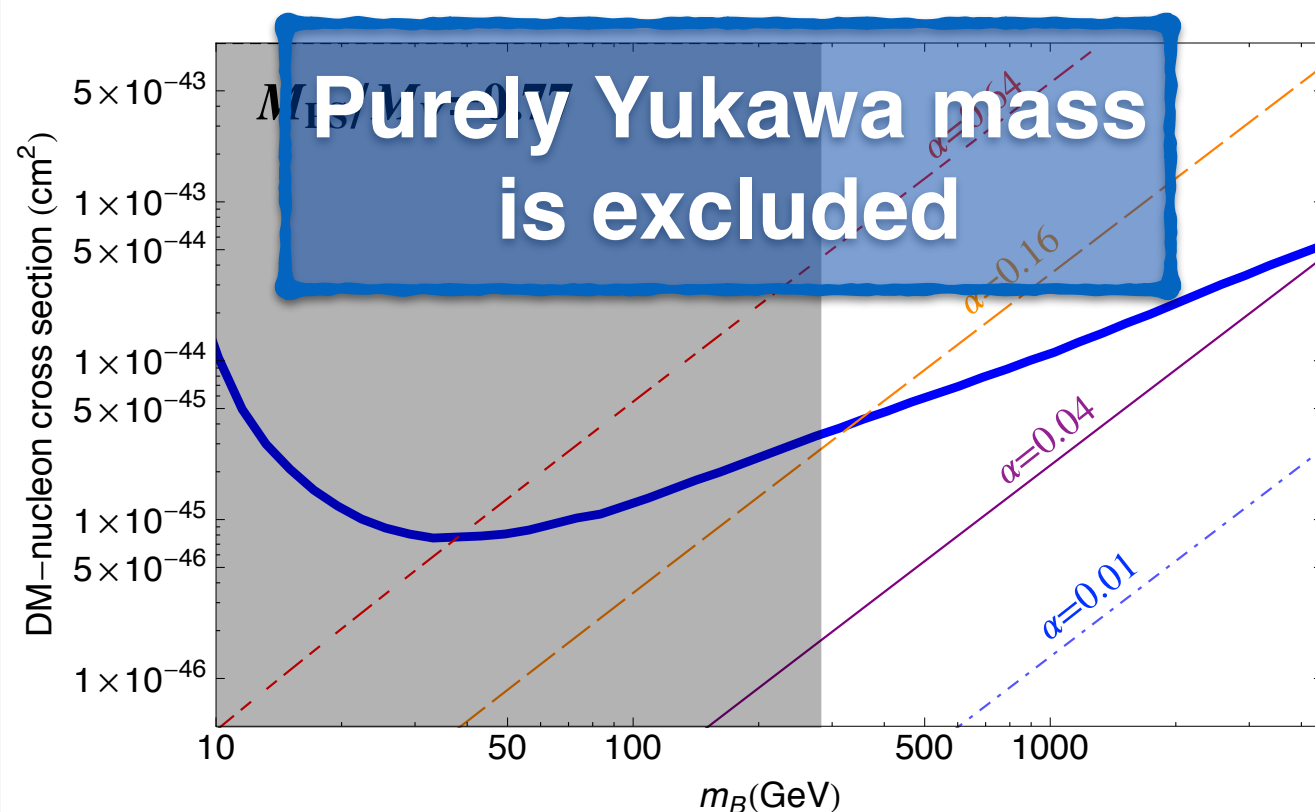
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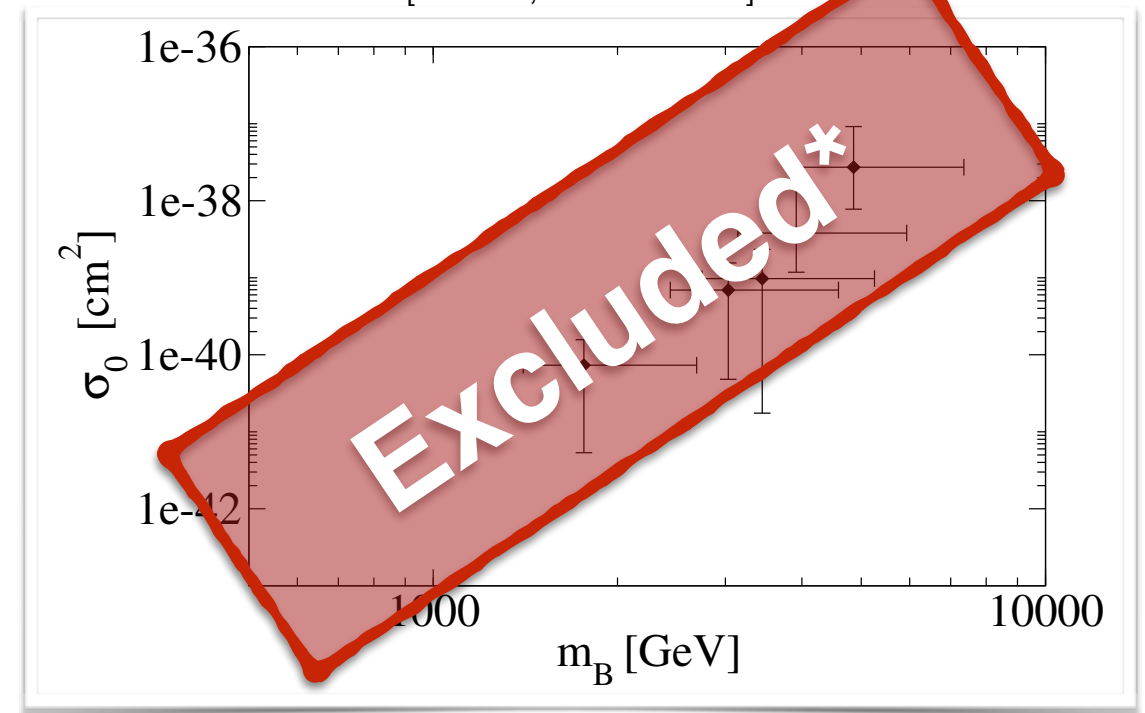
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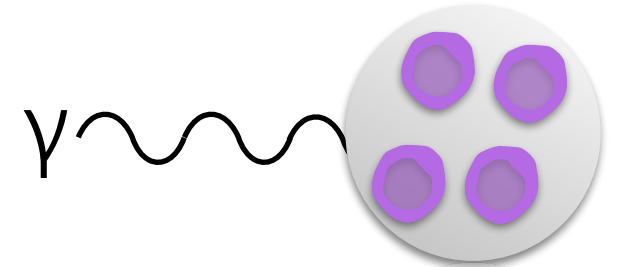
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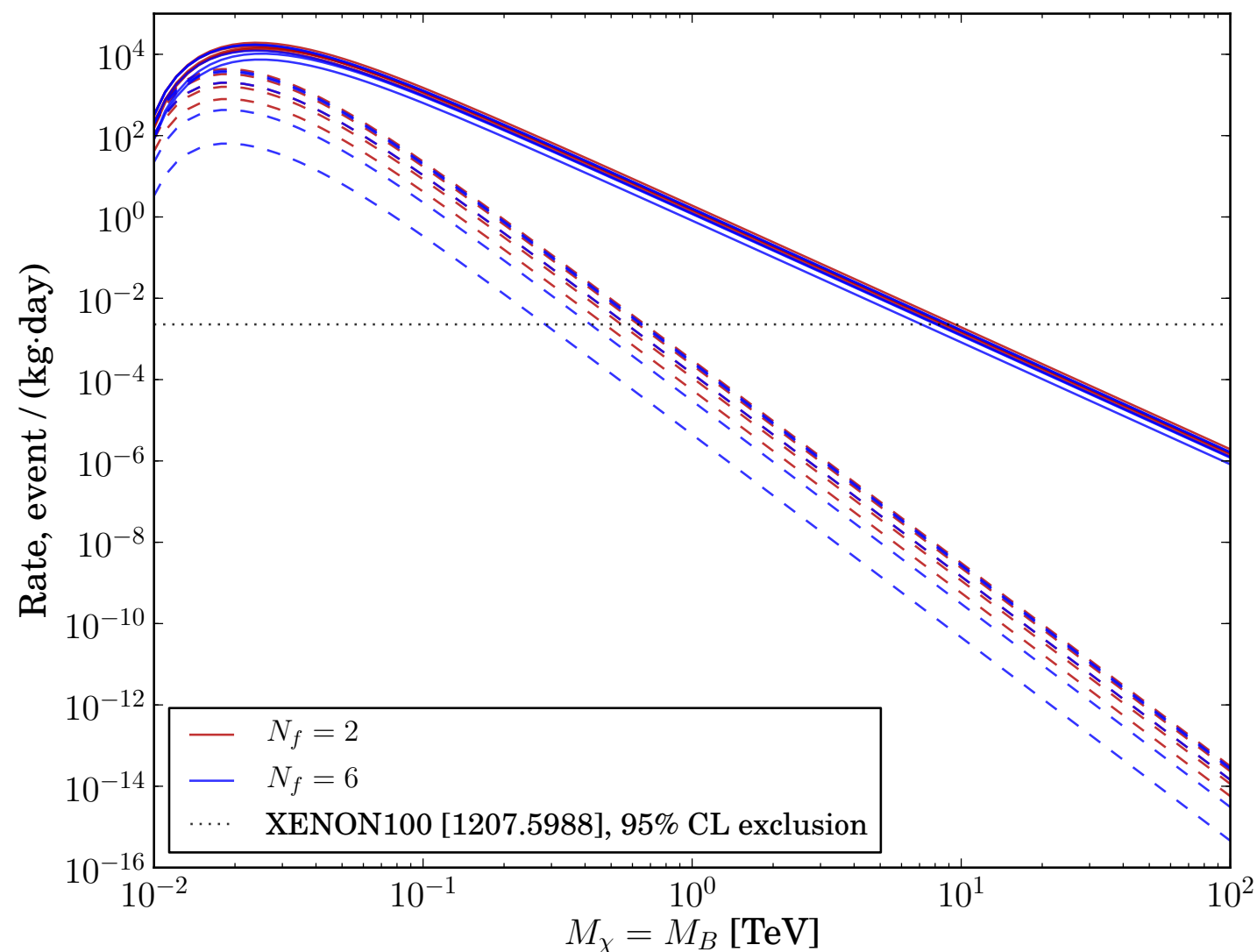
Bounds from EM moments



Mesonic and Baryonic EM form factors
directly from lattice simulations

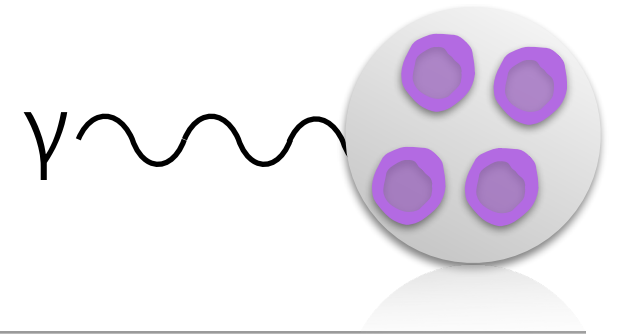
SU(3) $N_f=2,6$ dark fermionic baryon

[LSD, 1301.1693]



- ★ baryon similar to QCD neutron
- ★ dark quarks with $Q=Y$
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- ★ results independent of N_f

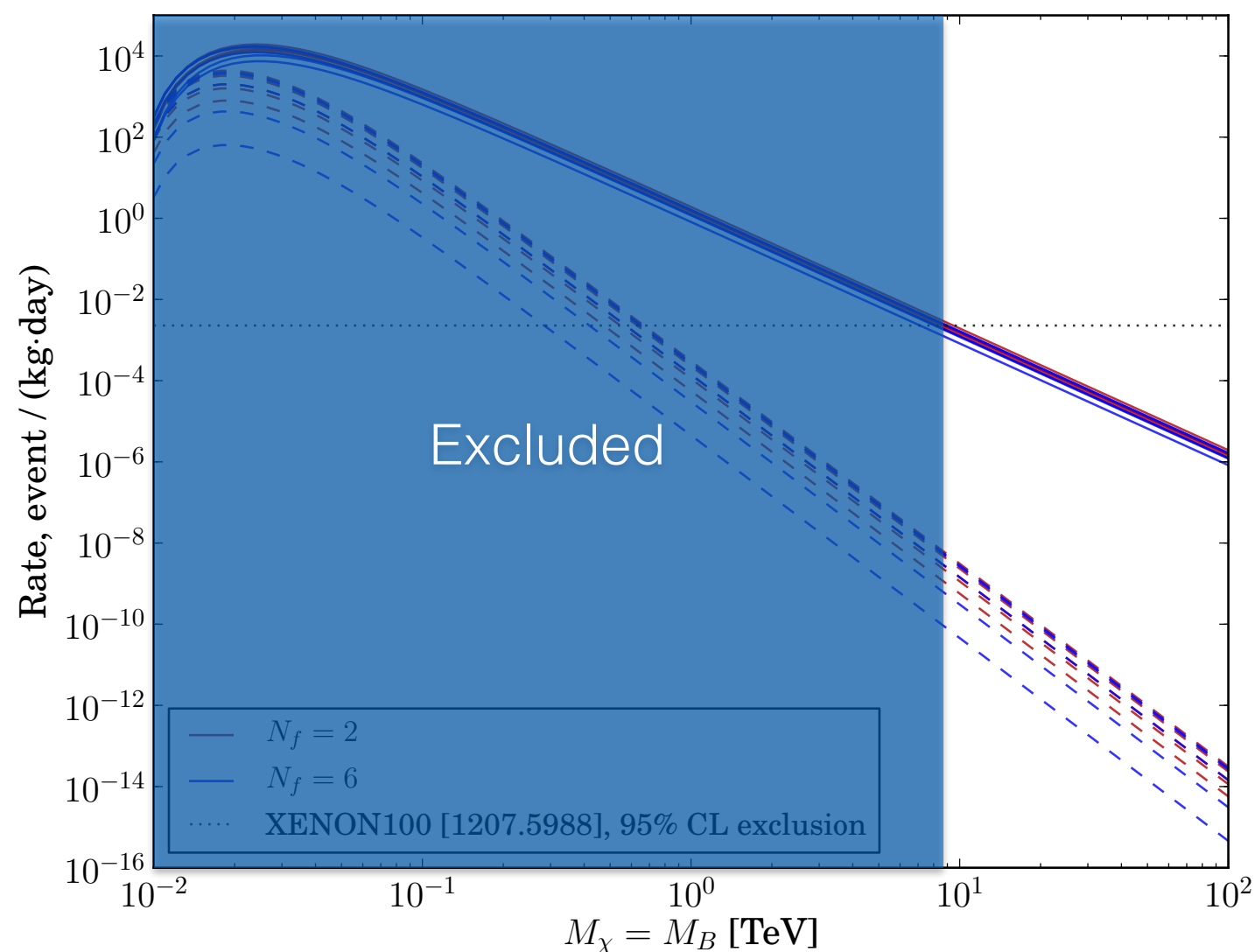
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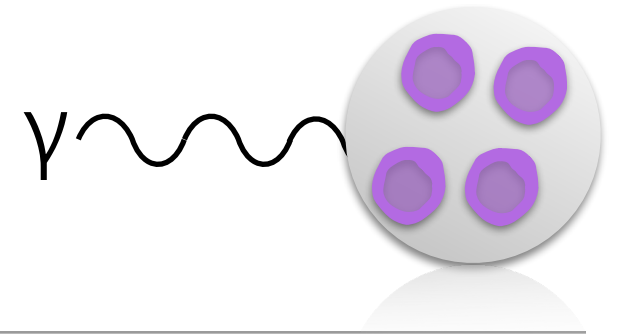
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$M_B > \sim 10 \text{ TeV}$

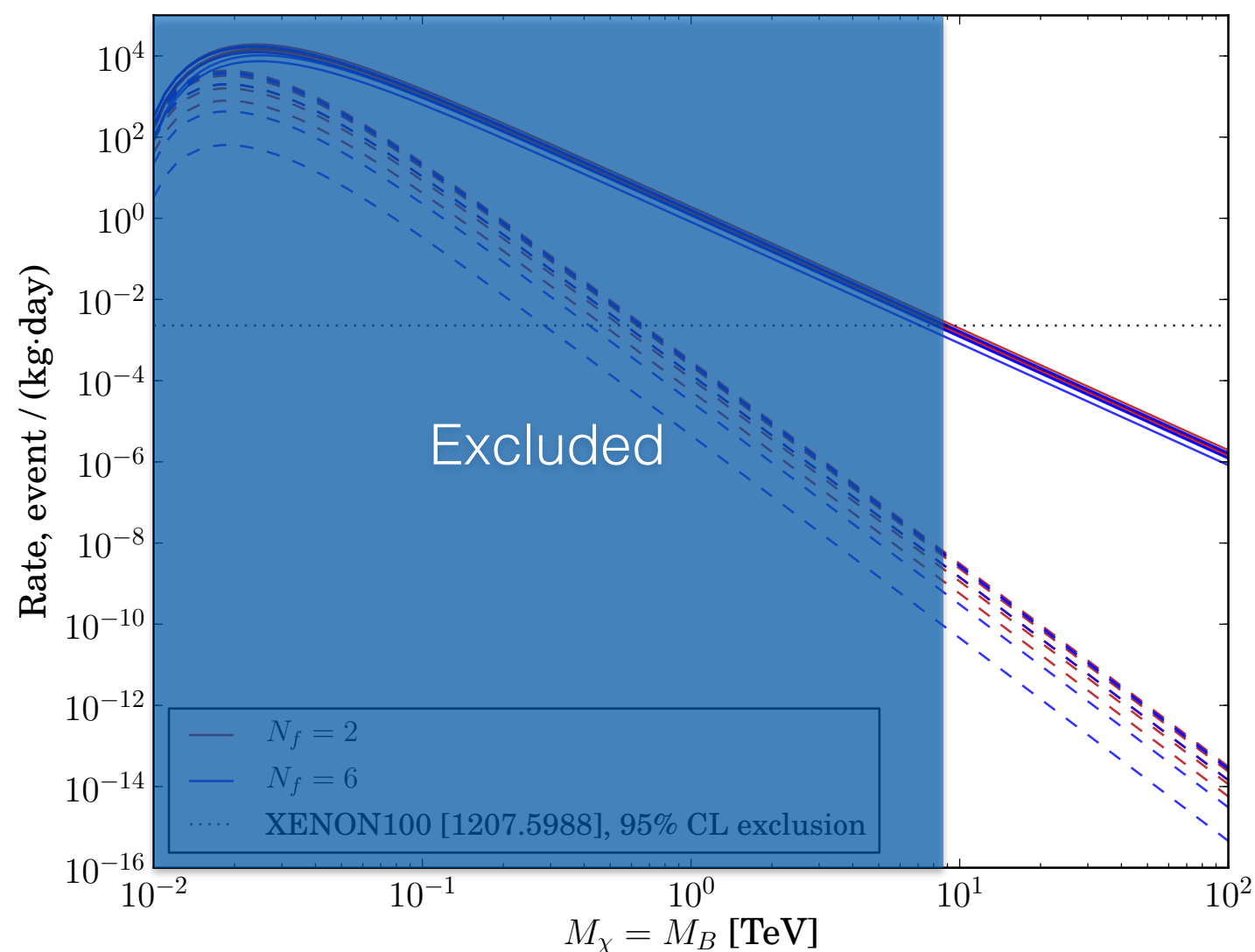
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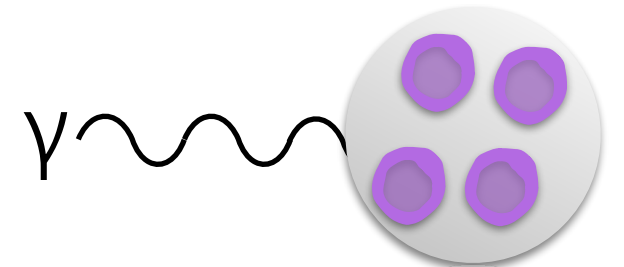


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pushed to $\sim 100 \text{ TeV}$
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Bounds from EM moments

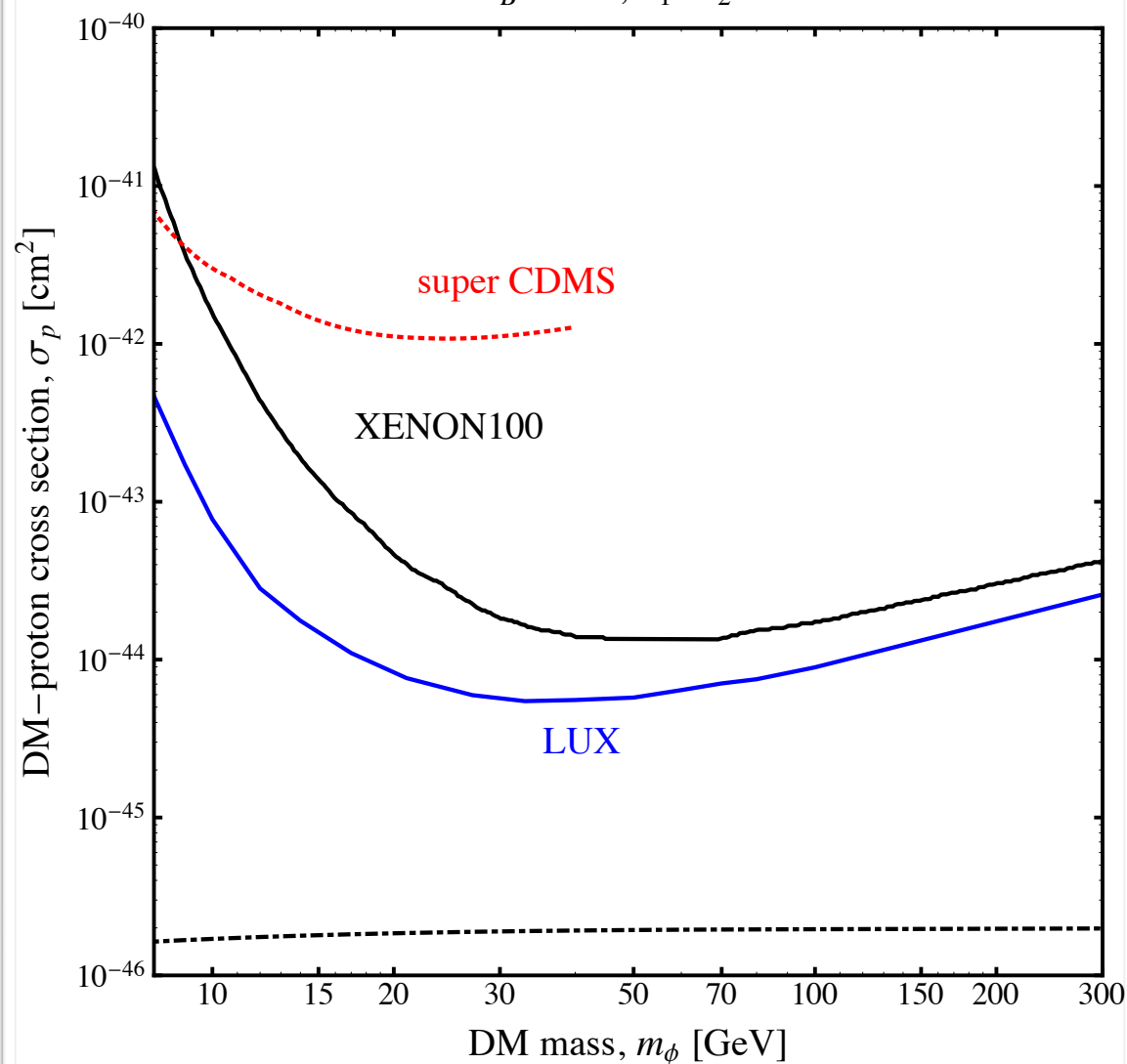


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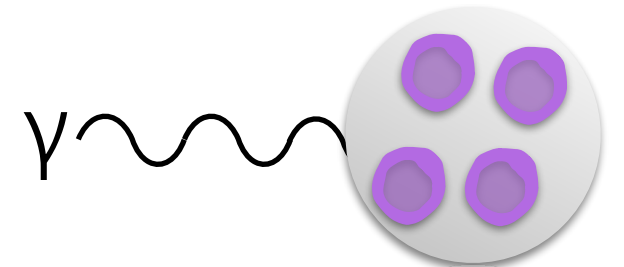
[Hietanen et al., 1308.4130]

$d_B = -0.1, d_1 + d_2 = 1$



- ★ dm is “mesonic” pNGB
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- ★ use VMD with lattice ρ mass
- ★ scale set by $F_\pi = 256$ GeV
- ★ depends on isospin breaking d_B
- ★ also couples to Higgs ($d_1 + d_2$)

Bounds from EM moments

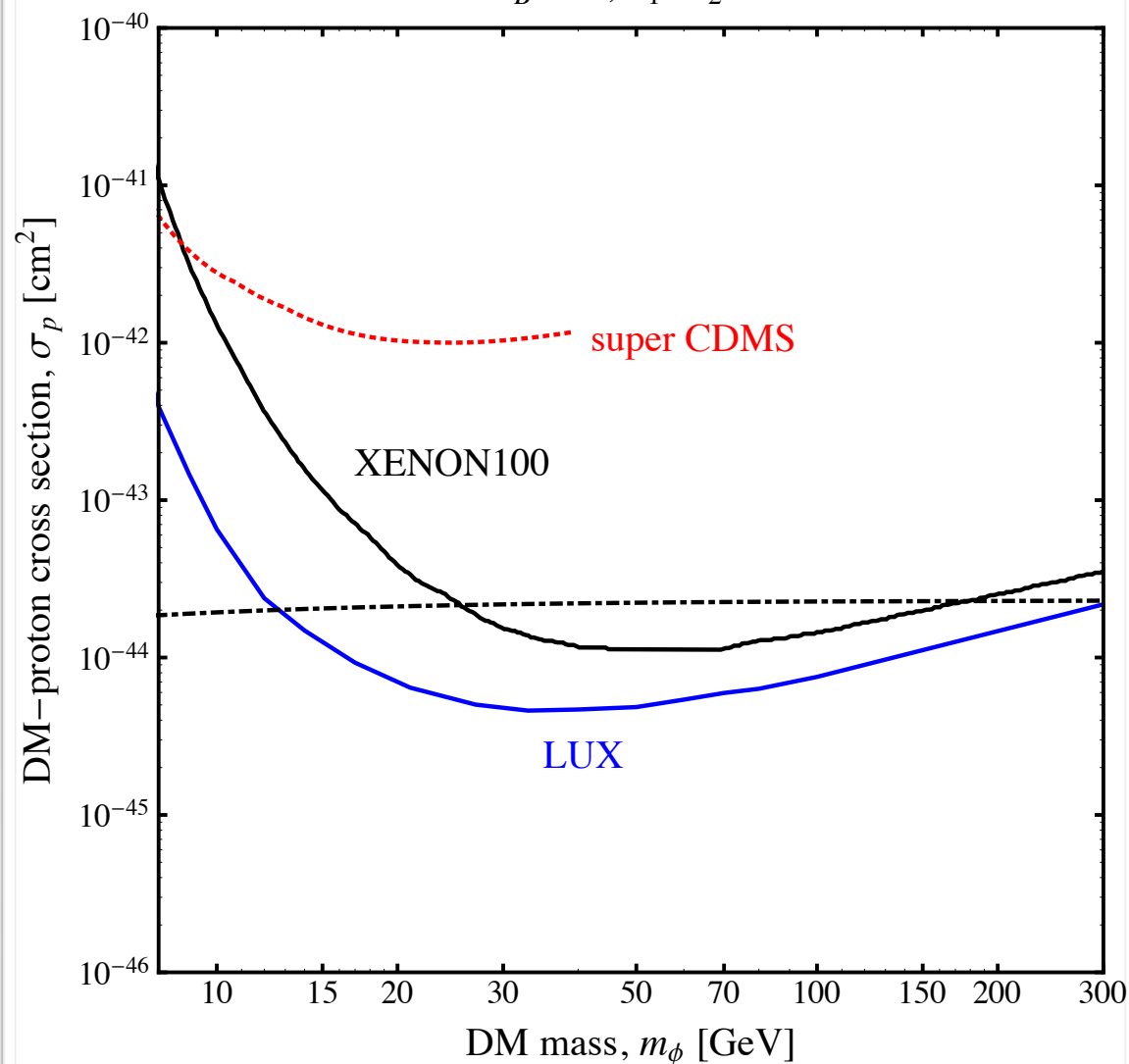


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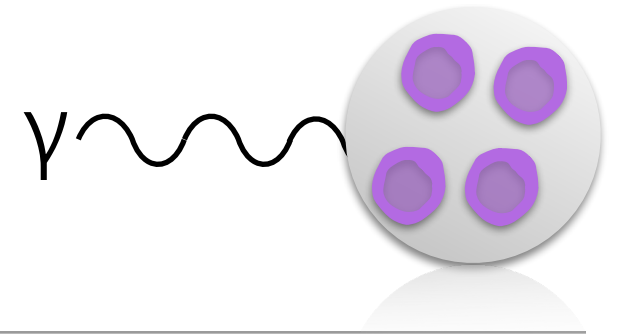
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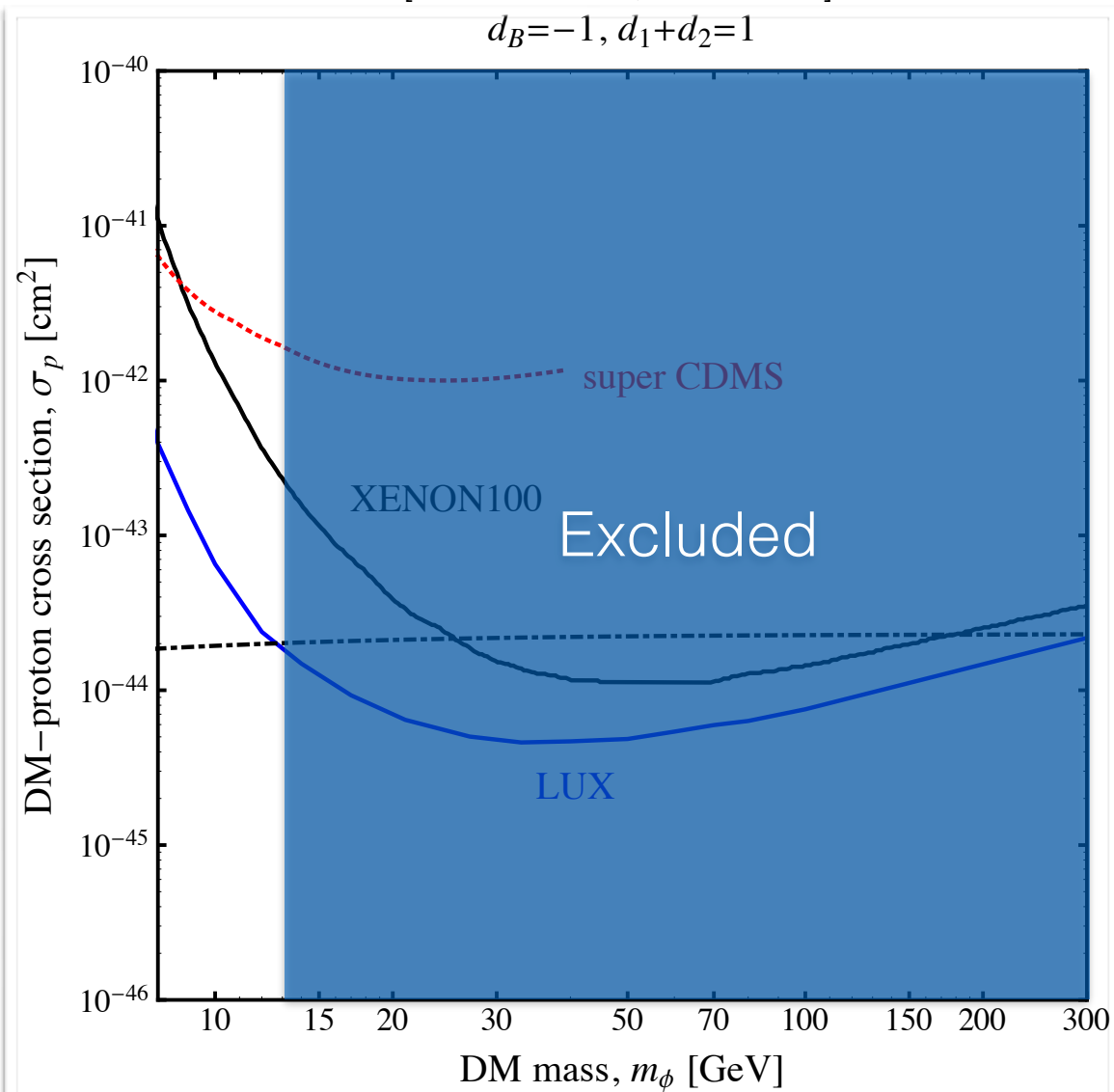


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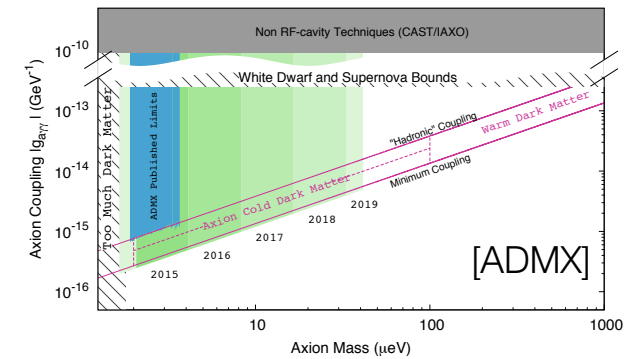
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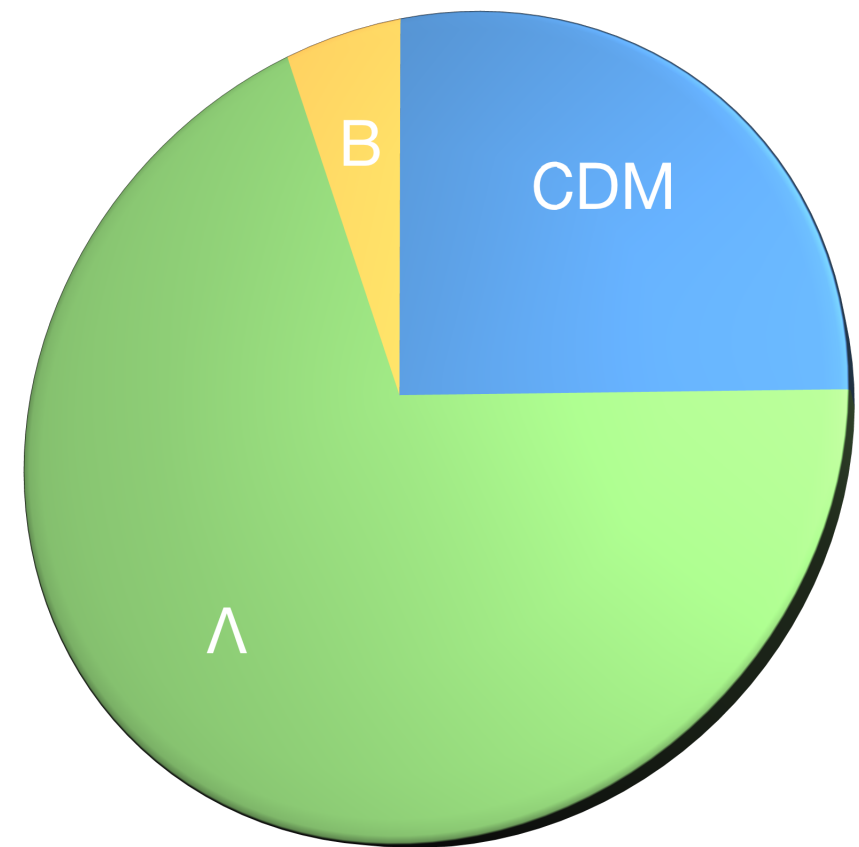
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$M_B \sim < 13$ GeV
depends on d_B

Axion dark matter

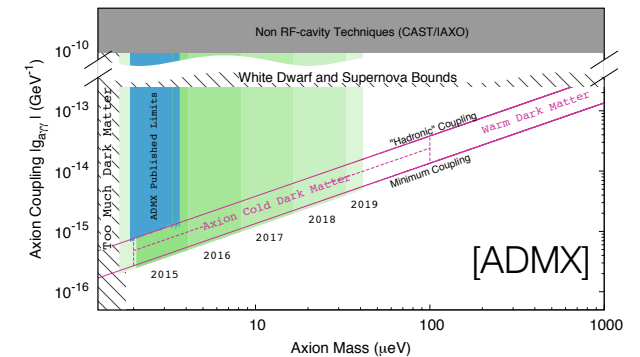


- Axions were originally proposed to deal with the Strong-CP problem
 - They also form a plausible DM candidate
 - The axion energy density requires non-perturbative QCD input
- Being sought in ADMX (LLNL, UW) & CAST-IAXO (CERN) with large discovery potential in the next few years
- Requiring $\Omega_a \leq \Omega_{\text{CDM}}$ yields a lower bound on the axion mass today



$\Omega_{\text{tot}} = 1.000(7)$
PDG 2014

Axion dark matter



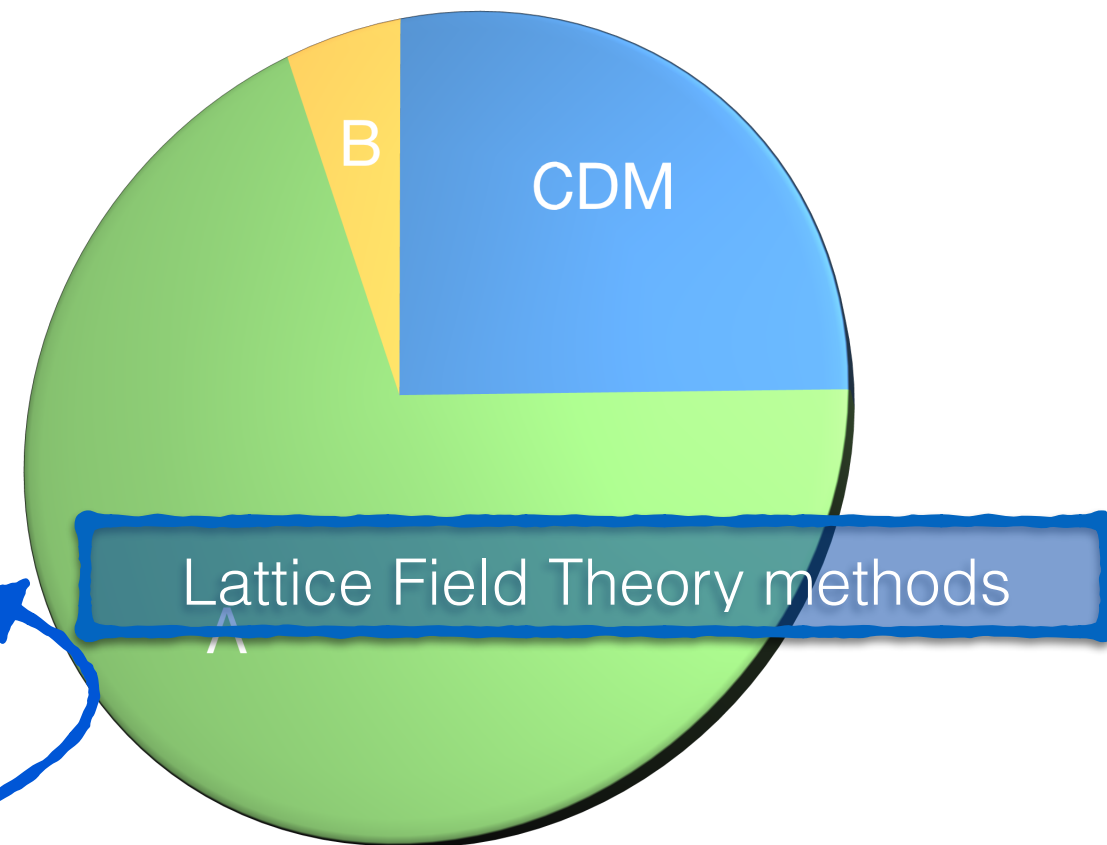
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PDG 2014

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Constraints from lattice simulations

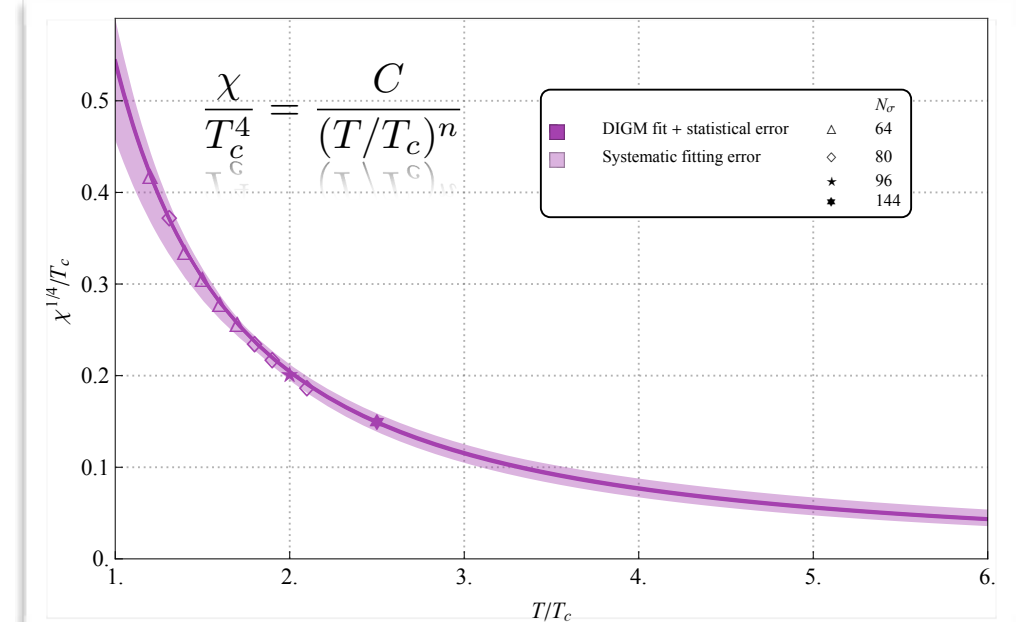
Non-perturbative calculation of QCD topology at finite temperature

- Pure gauge SU(3) topological susceptibility
 ➔ compatible with model predictions, but
large non-perturbative effects

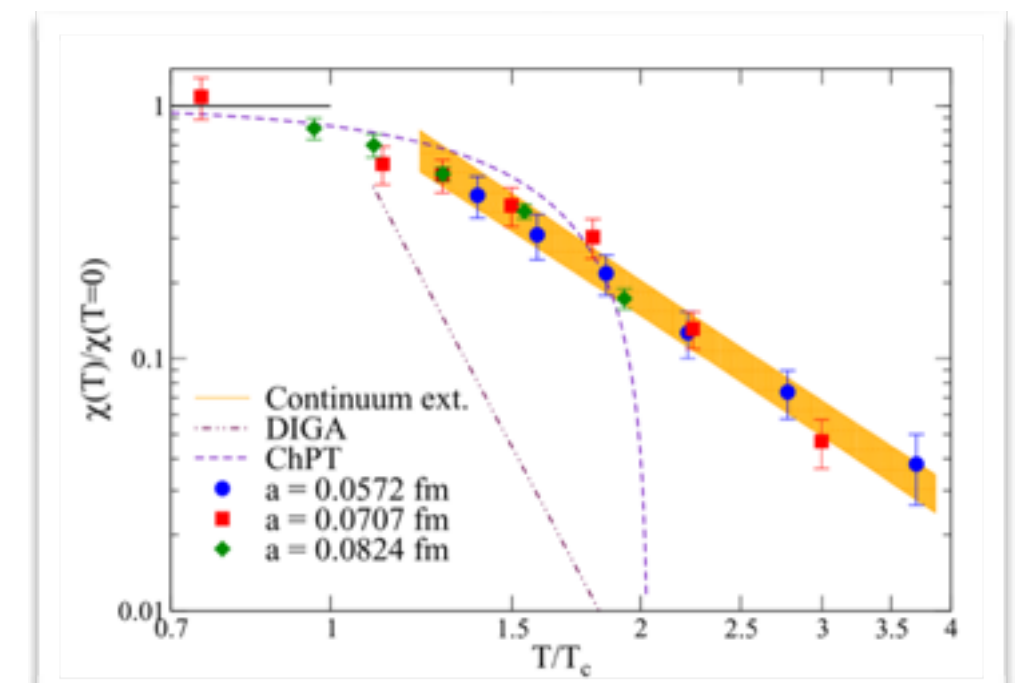
[Kitano&Yamada, 1506.00370][Borsanyi et al., 1508.06917][Frison et al., 1606.07175]

- is QCD topological susceptibility at high-T
well described by models? ➔ light
 fermions importantly affect the vacuum

[Trunin et al., 1510.02265][Petreczky et al., 1606.03145][Borsanyi et al., 1606.07494]



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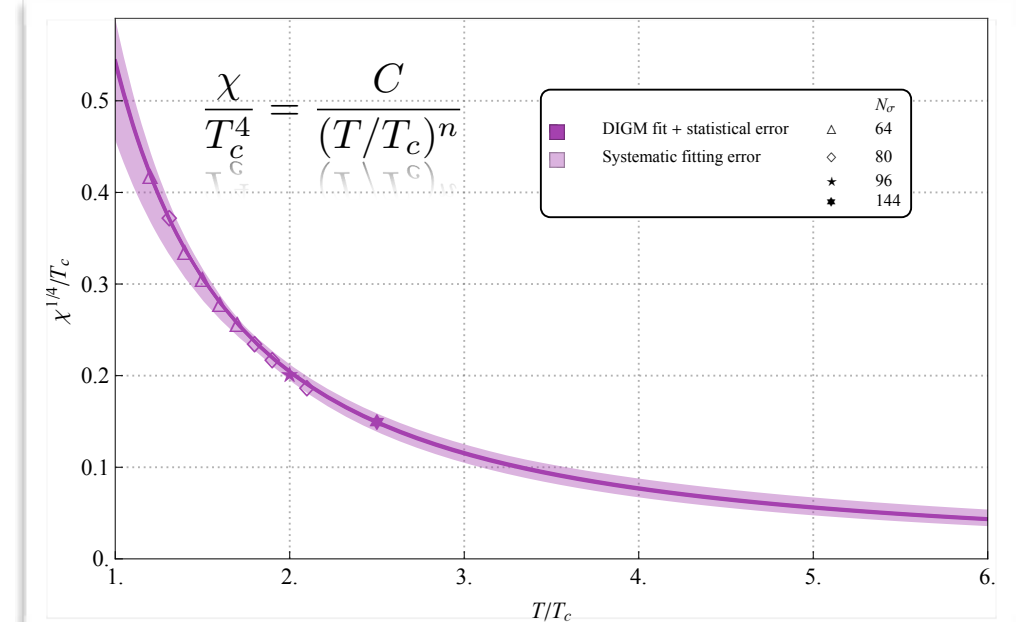
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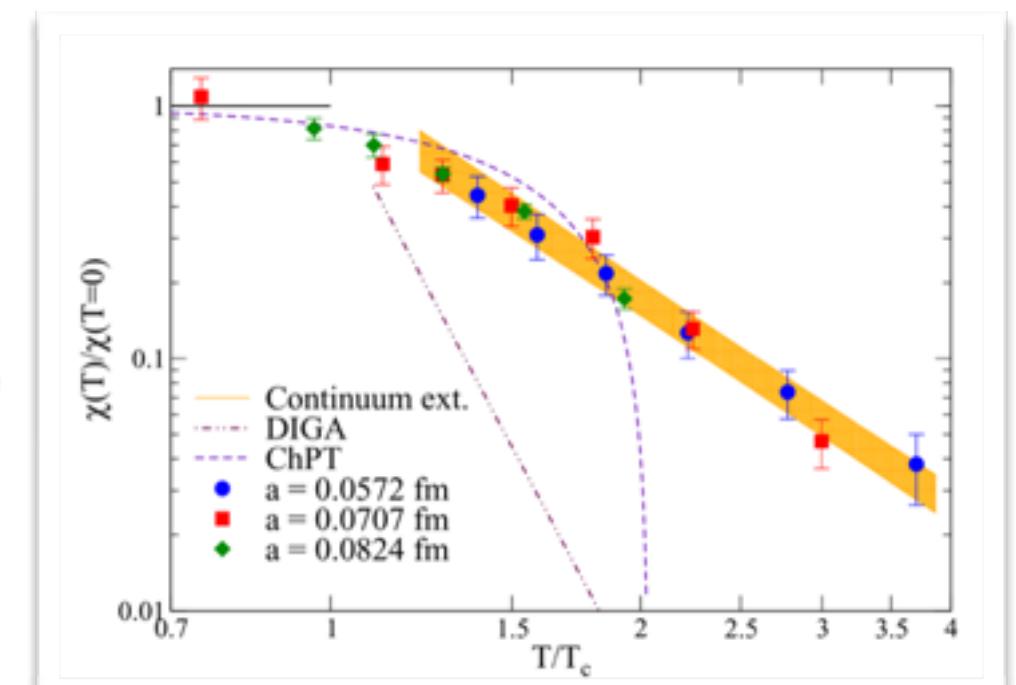
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Great effort to control all systematic lattice effects in order to impact experiments.
 This research has started only 1 year ago!



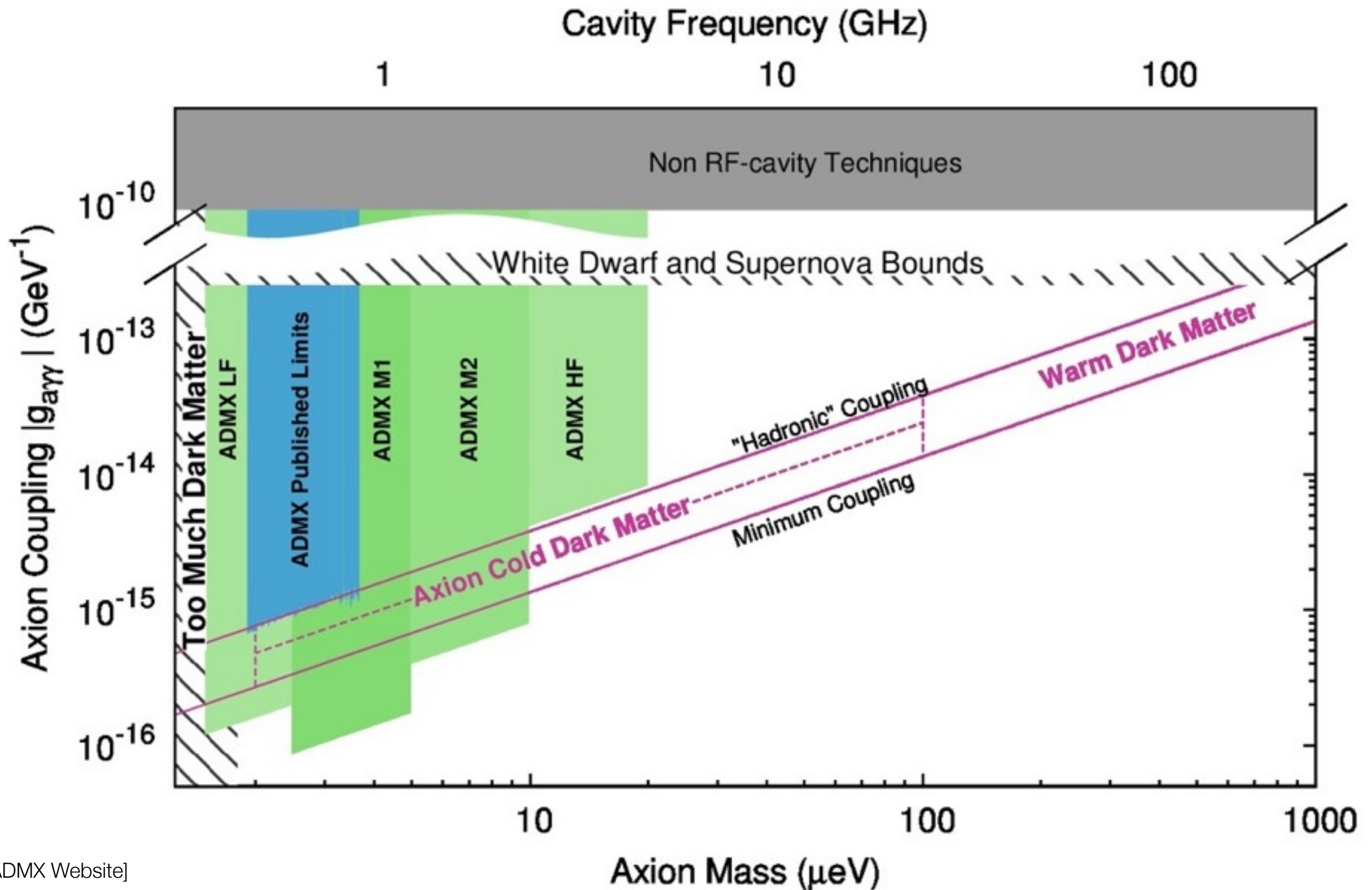
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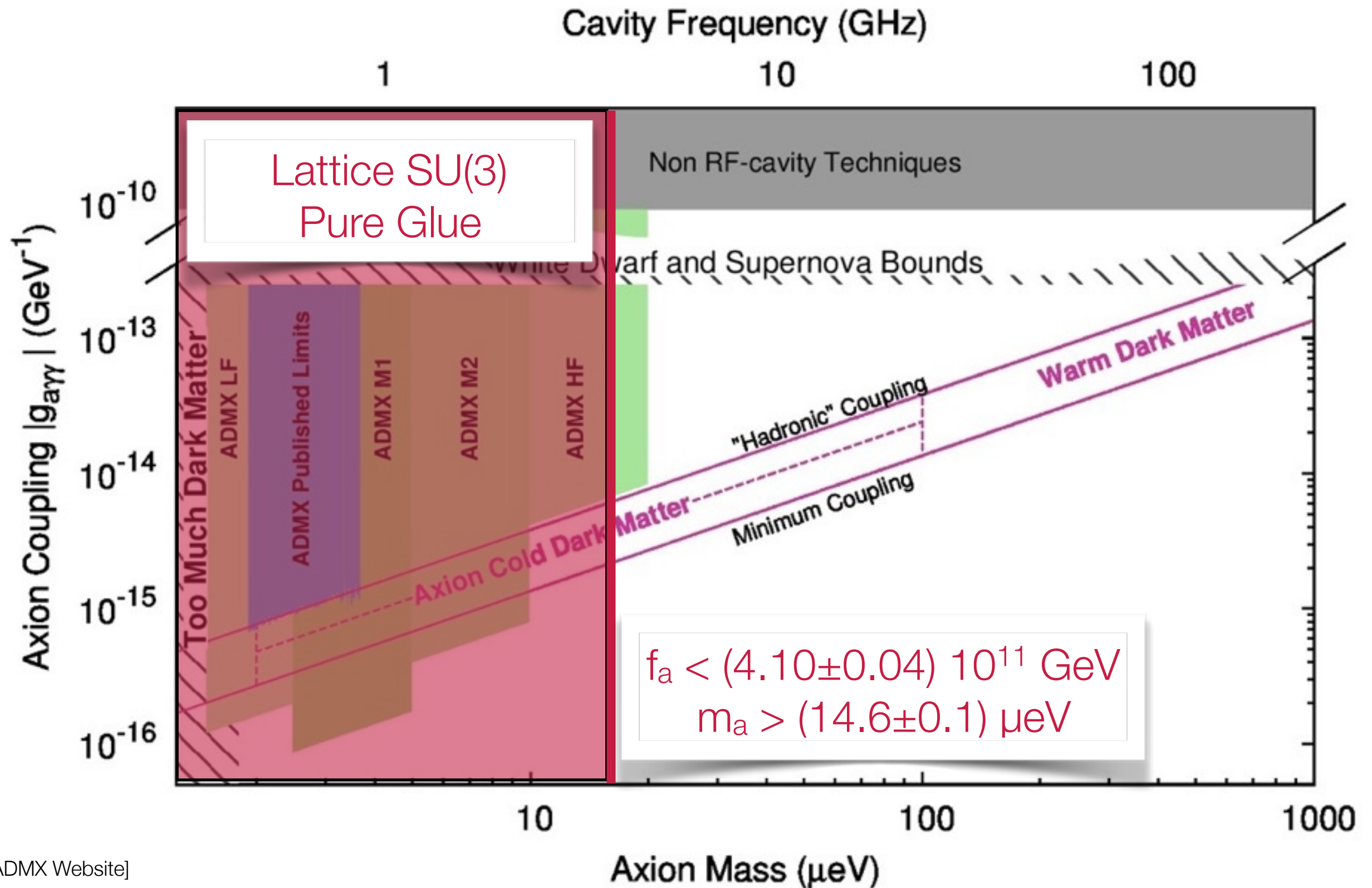
$$m_a^2 f_a^2 = \left. \frac{\partial^2 F}{\partial \theta^2} \right|_{\theta=0}$$

Axion mass lower bound



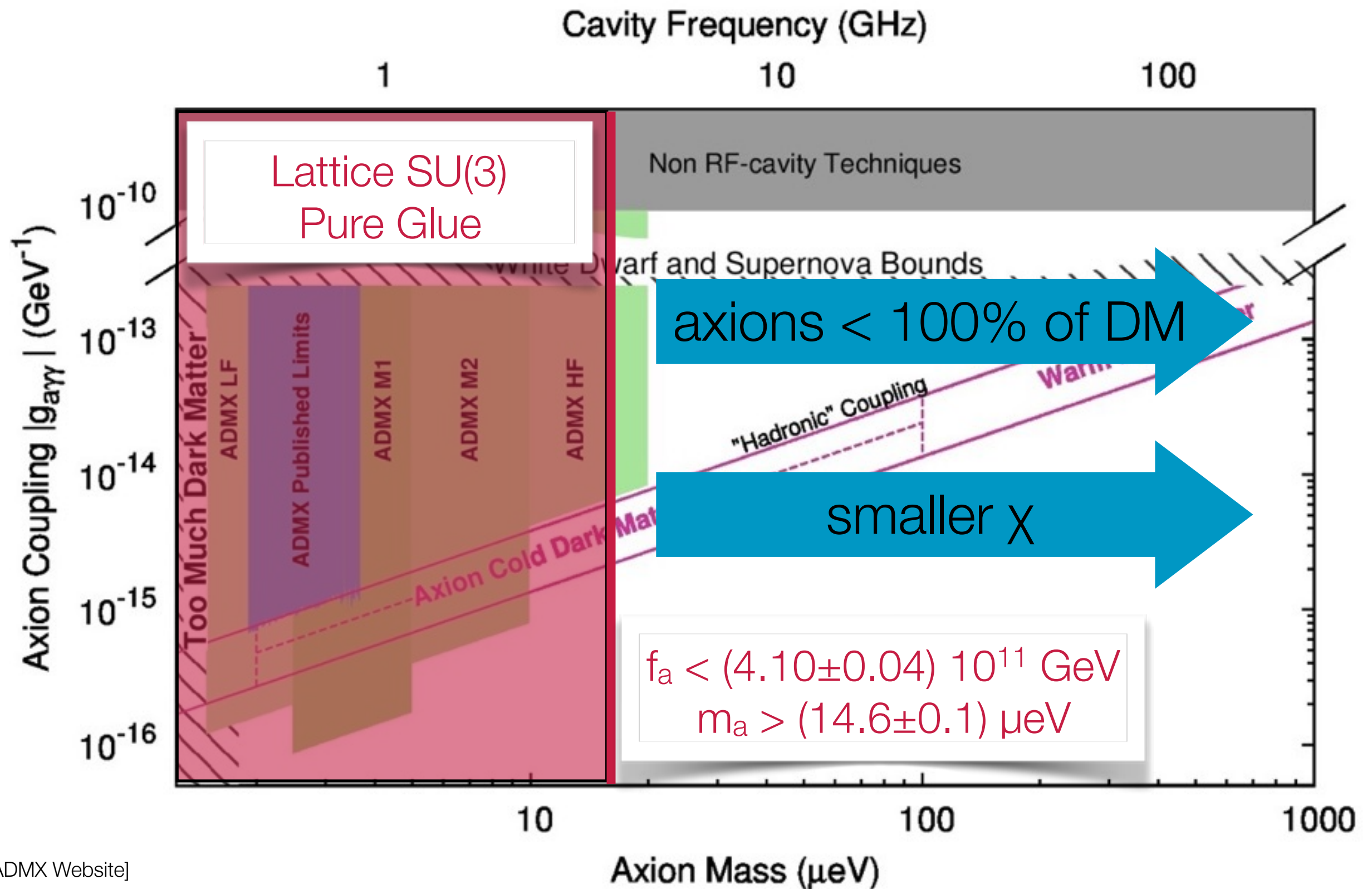
$$m_a^2 f_a^2 = \left. \frac{\partial^2 F}{\partial \theta^2} \right|_{\theta=0}$$

Axion mass lower bound



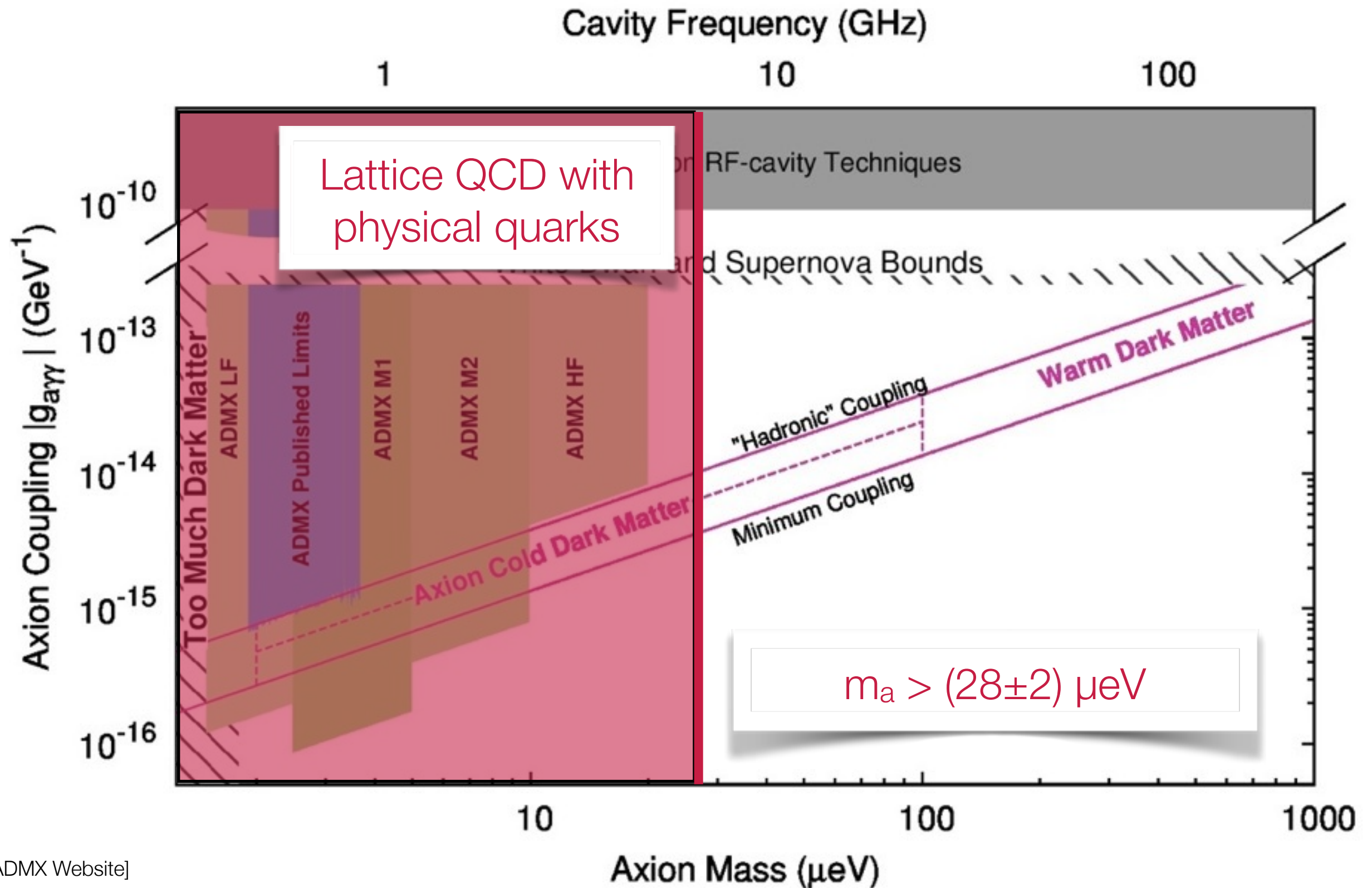
$$m_a^2 f_a^2 = \left. \frac{\partial^2 F}{\partial \theta^2} \right|_{\theta=0}$$

Axion mass lower bound



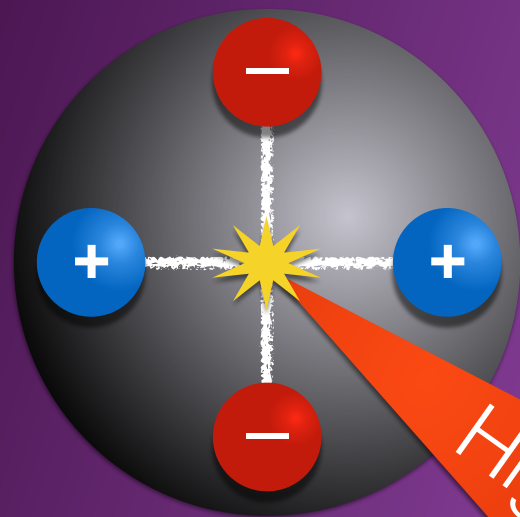
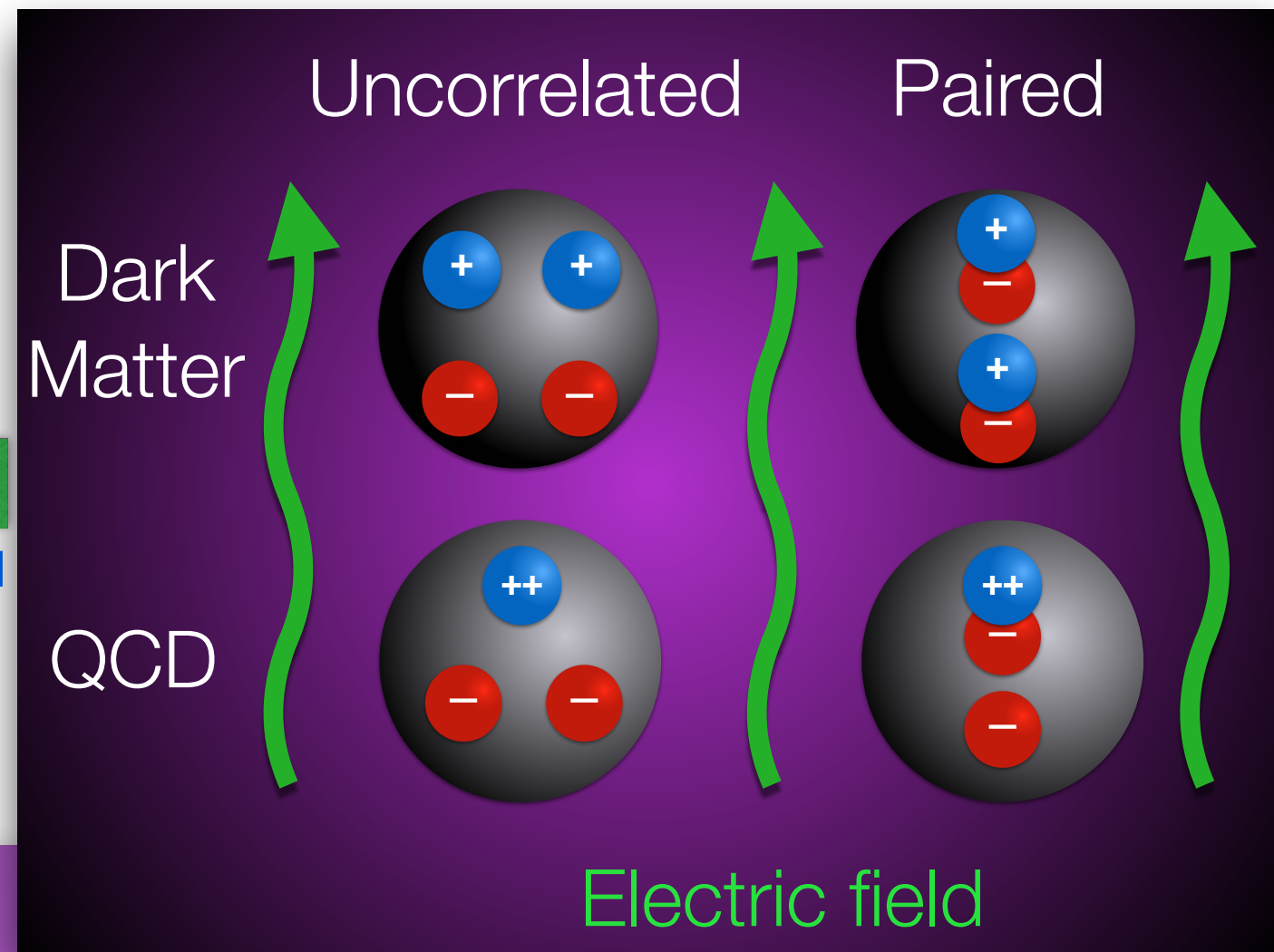
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Axion mass lower bound



PRL Editors' Suggestion: Polarizability

[LSD collab., Phys. Rev. Lett. 115 (2015) 171803]



PRD Editors' Suggestion: Higgs exchange

[LSD collab., Phys. Rev. D92 (2015) 075030]