



Nuclear modification of hard scattering processes in small systems at PHENIX

NIVEDITHA RAM

09/06/2022

PHENIX COLLABORATION MEETING 2022

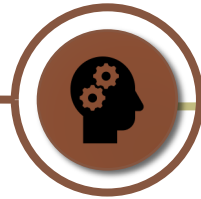
PhD Timeline and Projects



Aug, 2015

Design, construction and testing of the prototype-TPC

Prof. Tom Hemmick



Dec, 2018

Validation of the Glauber Model for centrality determination in d+Au collision system

Prof. Gabor David

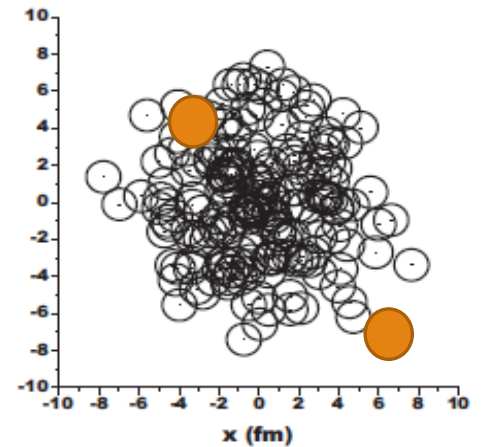
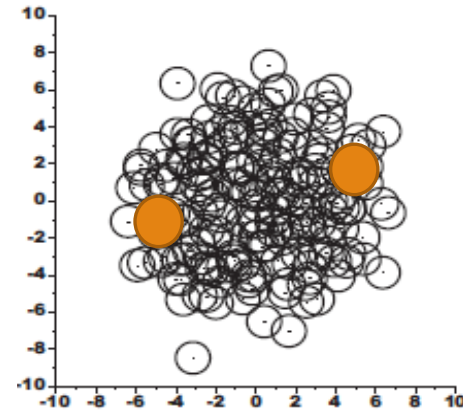
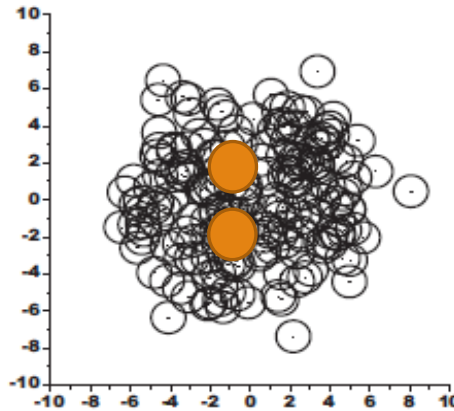
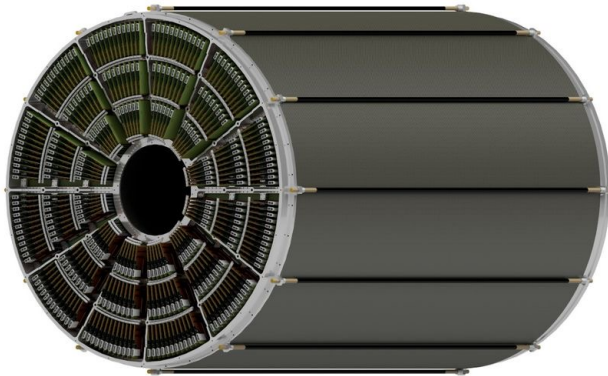


Aug, 2021

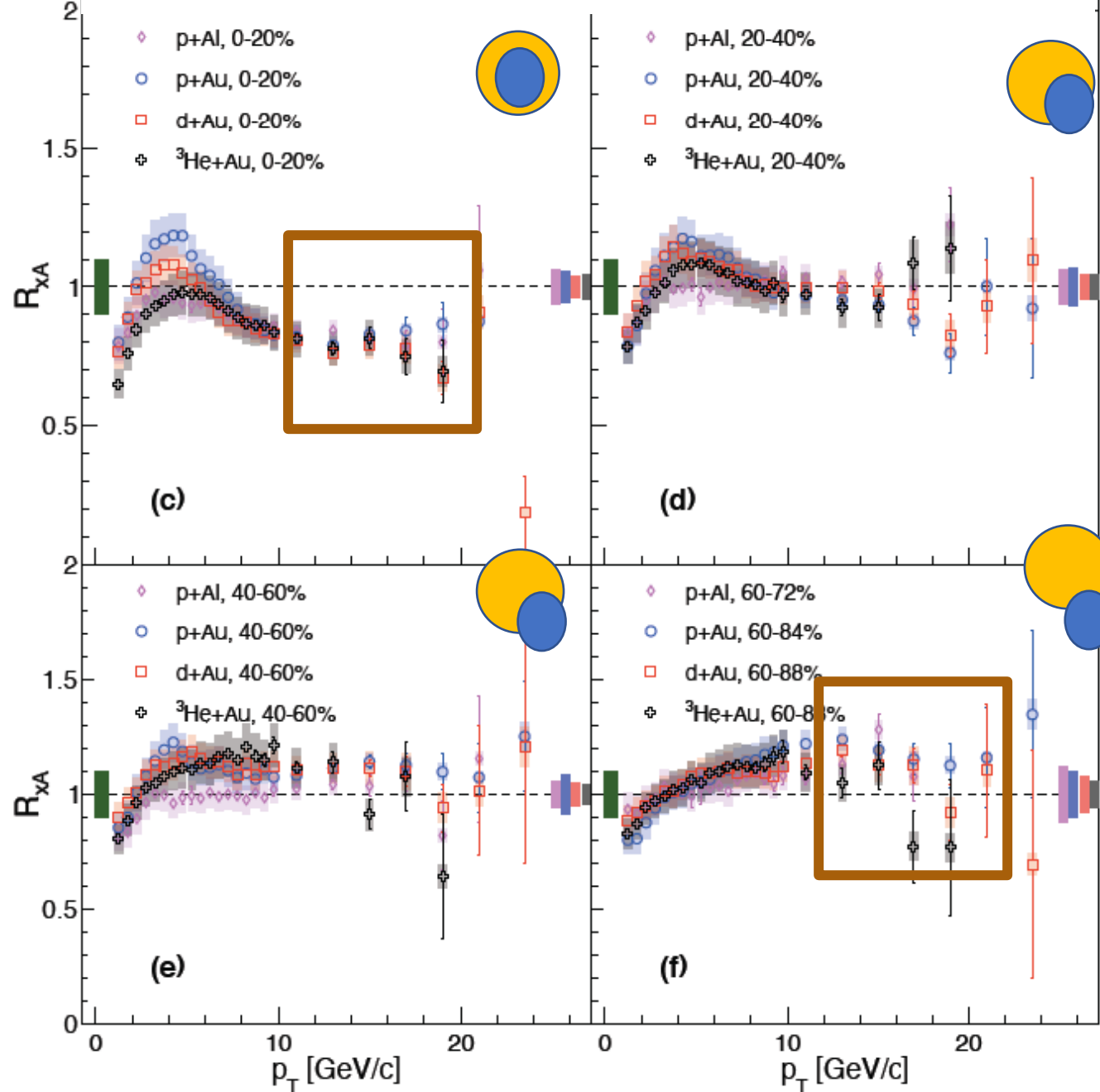


Dec, 2021

MA in physics



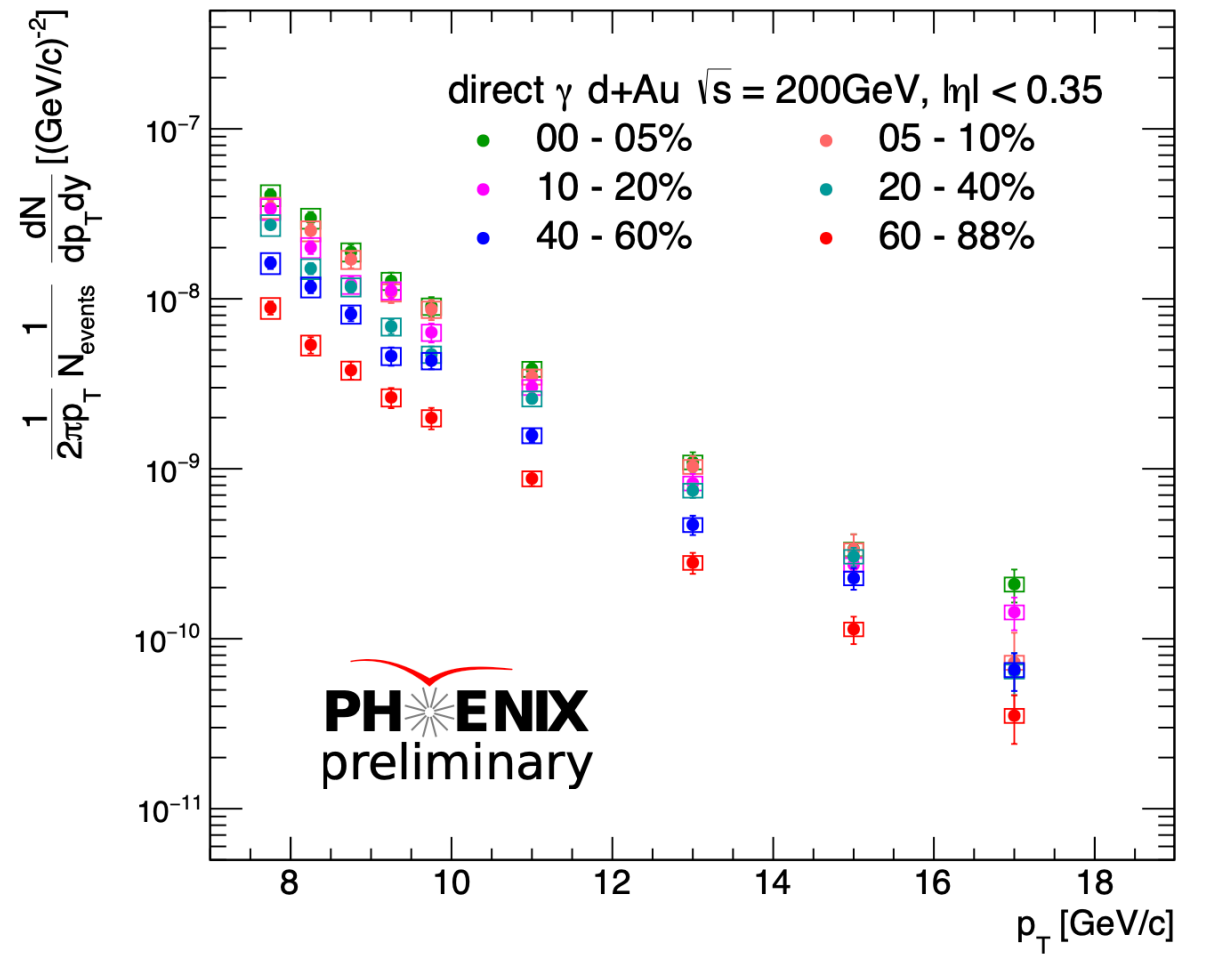
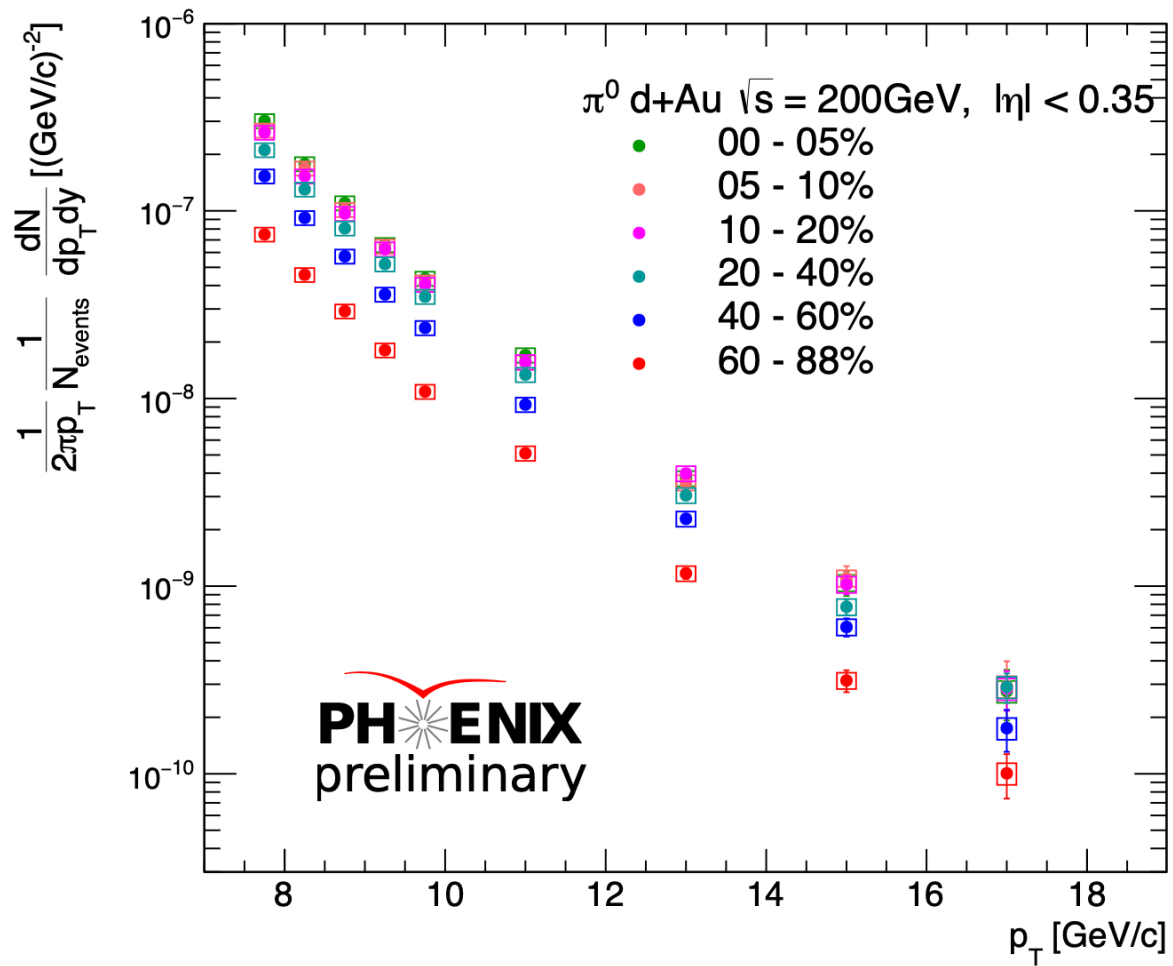
Puzzling behavior of R_{AA} in small systems



- Suppression in central collisions \rightarrow Formation of QGP droplets?
- Enhancement in peripheral collisions?

Is the determination of N_{coll} in different event classes biased

How do we study the centrality bias in experiments?



Invariant yield of π^0 and direct γ

No dependence on centrality

Dependence on centrality

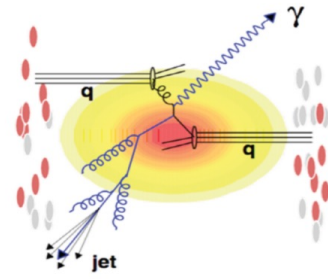
Invariant yield of direct γ

Dependence on centrality

Dependence on centrality

Invariant yield of π^0

Dependence on centrality
NO dependence on centrality



Hypothesis #1 : The centrality dependence of π^0 is from final state effect

Hypothesis #2 : The centrality dependence of π^0 is from bias in determination of N_{coll} in different centralities

Au+Au

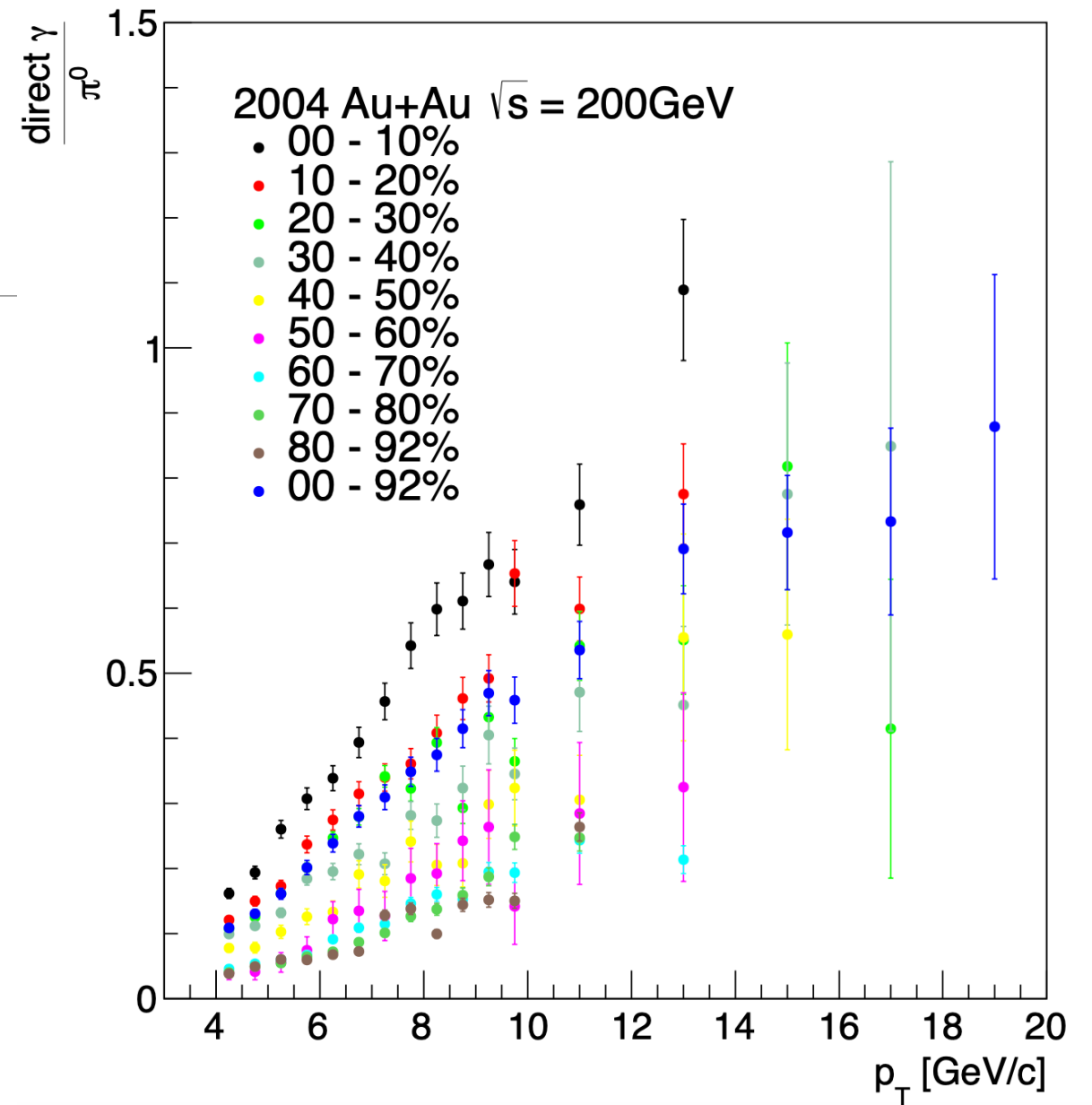
Ratio of direct γ over π^0

Clear centrality dependence



Hypothesis #1 :

The centrality dependence of π^0 is
from final state effect



d+Au

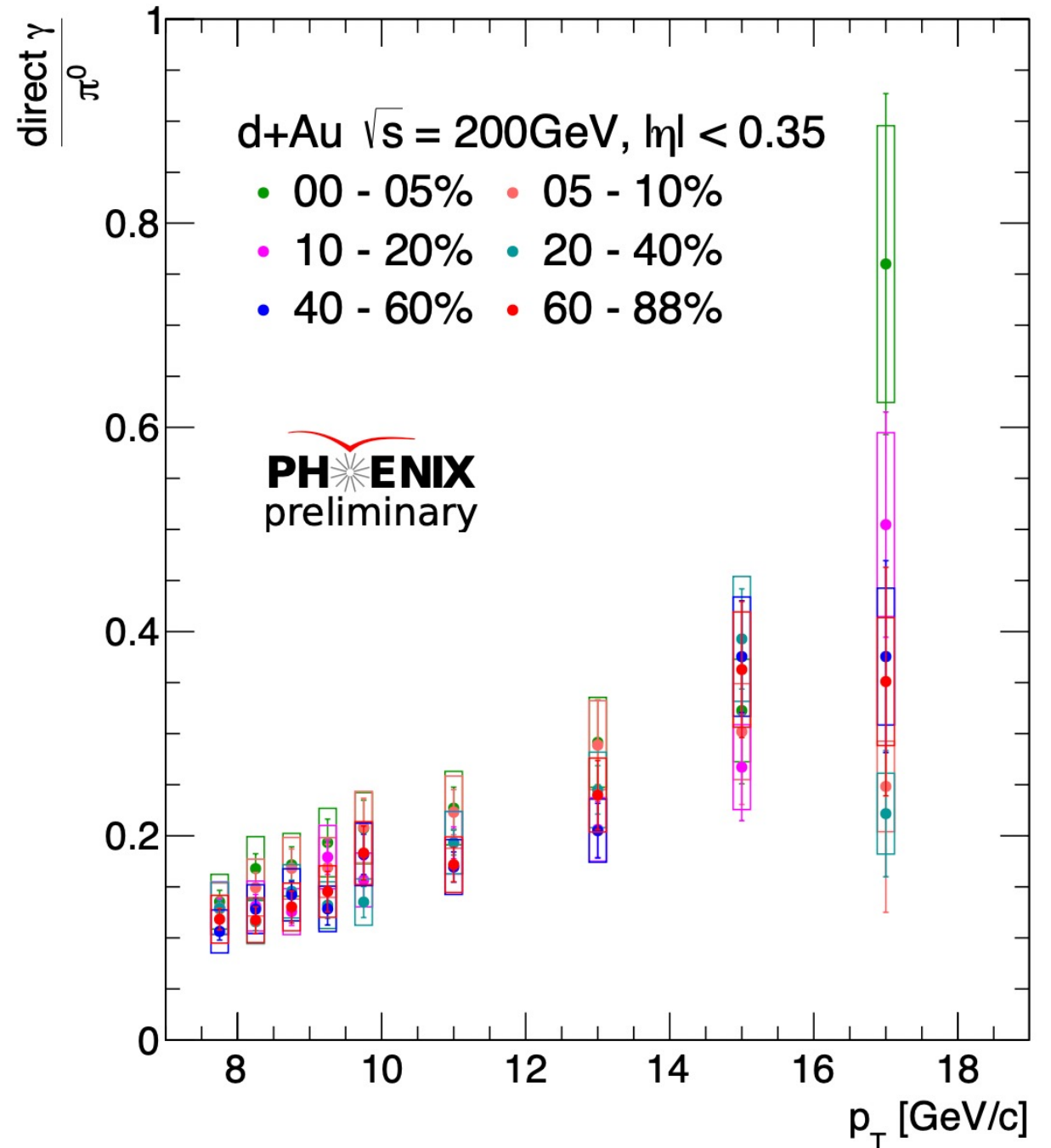
Ratio of direct γ over π^0

to first order, NO clear centrality
dependence



Hypothesis #2 :

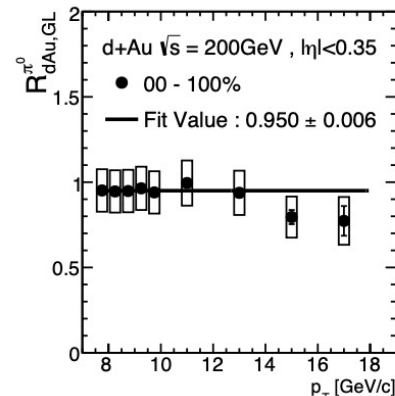
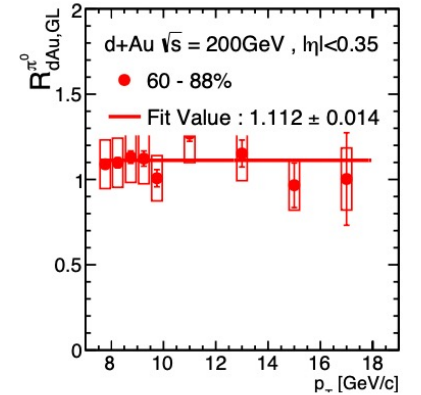
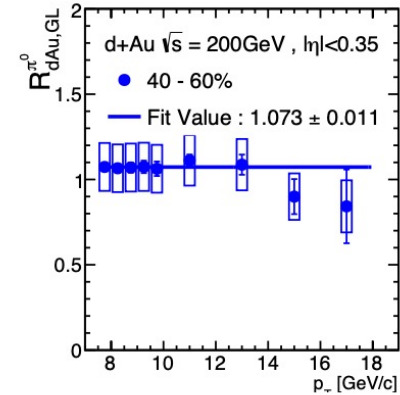
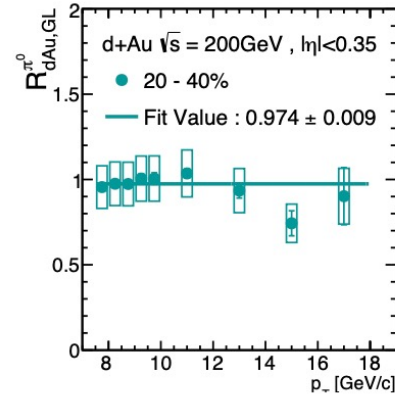
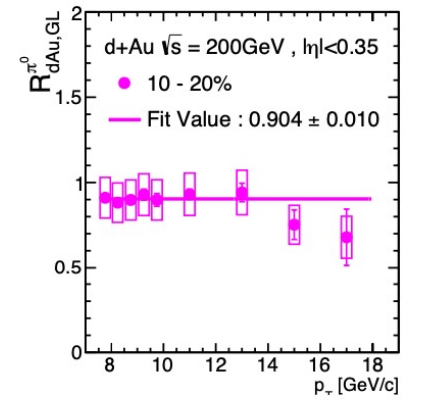
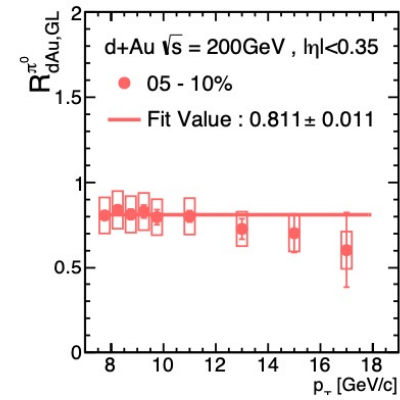
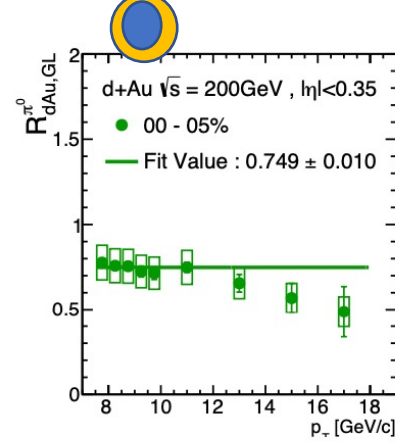
The centrality dependence of π^0 is
from bias in determination of N_{coll} in
different centralities



Nuclear Modification Factor of π^0 s

$$R_{AB,GL}(p_T) = \frac{\left(\frac{d^2N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll}^{GL} \rangle_{AB} * \left(\frac{d^2N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB)}{\langle N_{coll}^{GL} \rangle_{AB} * Y(pp)}$$

- ❖ There is a centrality dependence of π^0
- ❖ The most central events are suppressed (<1) and peripheral events are enhanced (>1)
- ❖ In the given p_T range, to first order R_{dAu} appears to be flat.

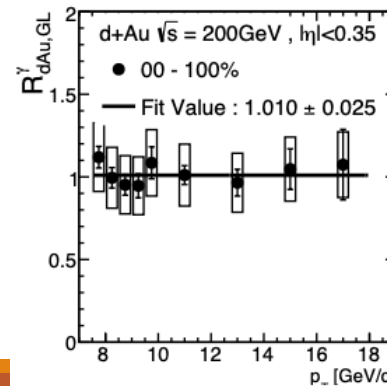
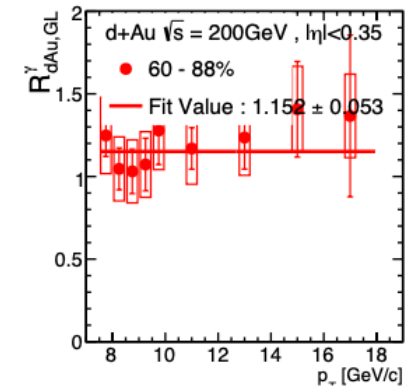
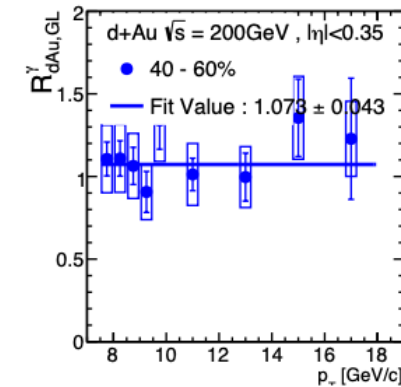
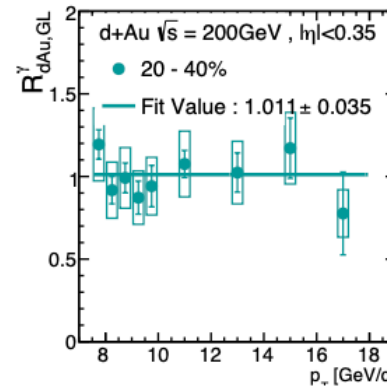
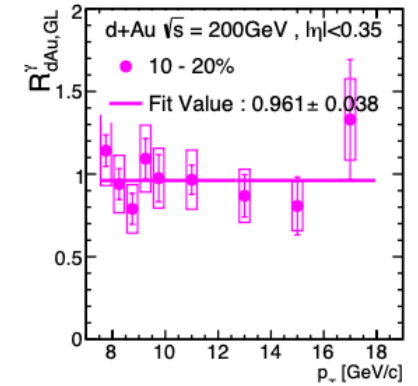
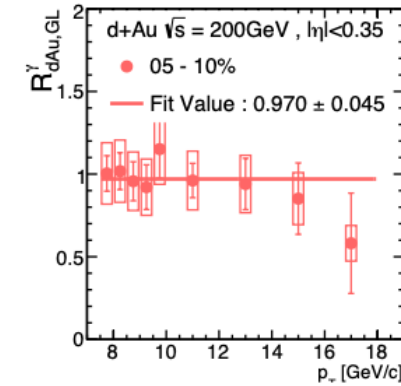
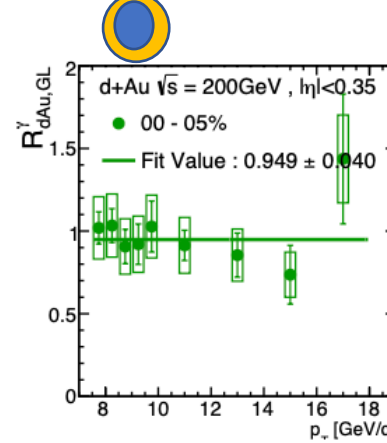


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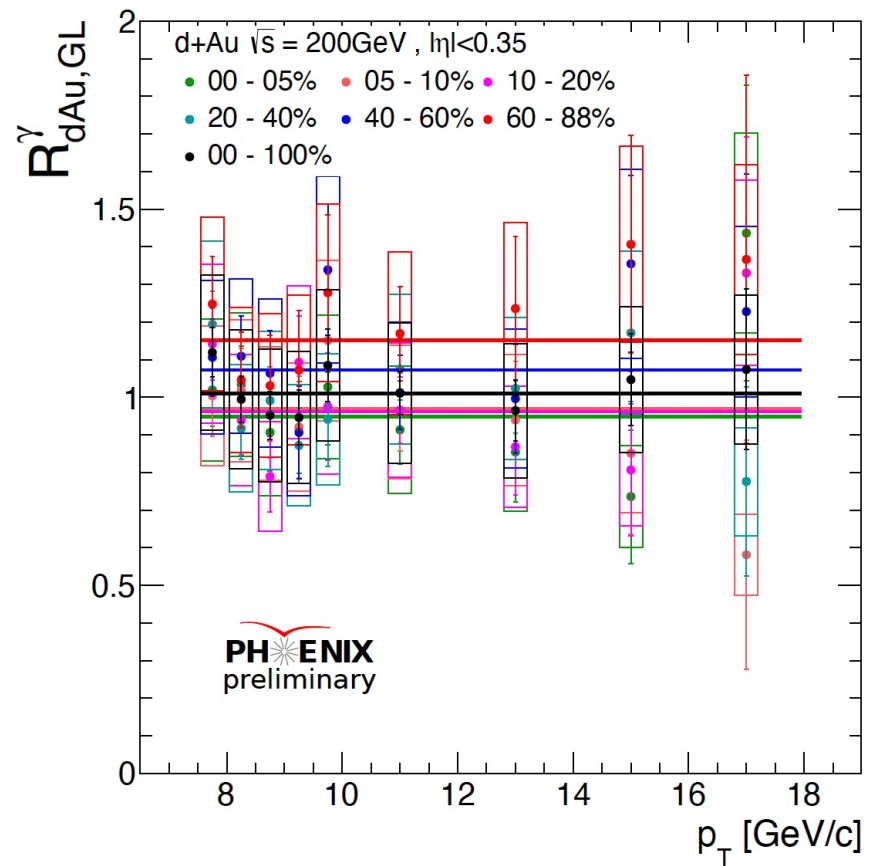
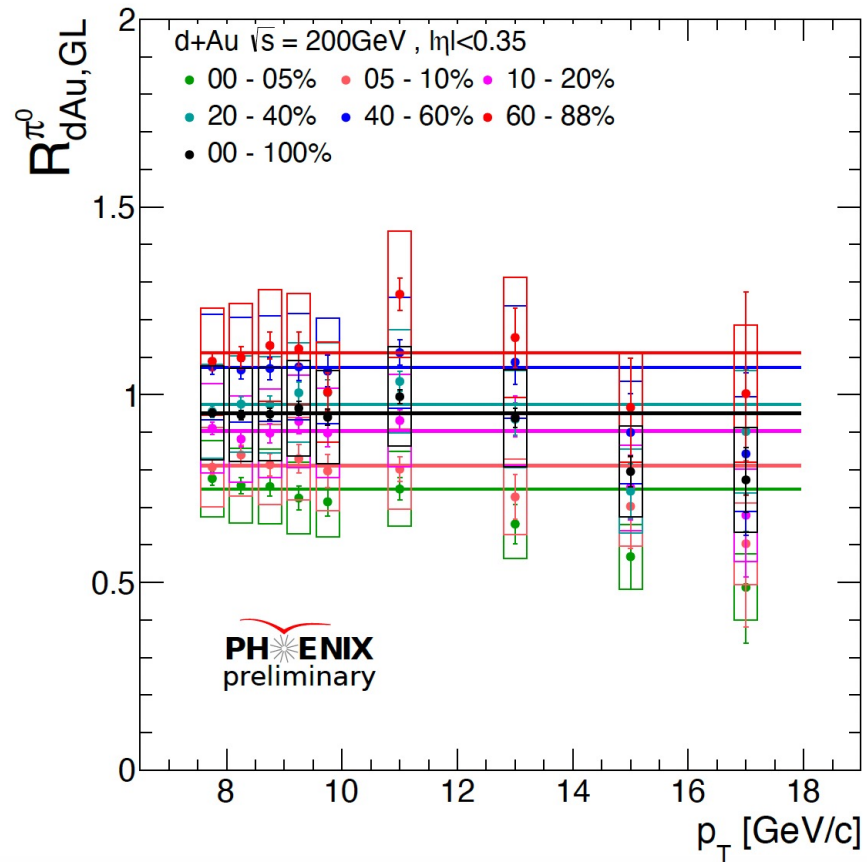
Nuclear Modification Factor of direct γ s

$$R_{AB,GL}(p_T) = \frac{\left(\frac{d^2N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll}^{GL} \rangle_{AB} * \left(\frac{d^2N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB)}{\langle N_{coll}^{GL} \rangle_{AB} * Y(pp)}$$

- ❖ There is a centrality dependence of direct γ s
- ❖ The most central events are suppressed (<1) and peripheral events are enhanced (>1)
- ❖ In the given p_T range, to first order R_{dAu} appears to be flat.



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The straight lines are fit to the datapoints at different centralities

- ❖ In central events the π^0 s are suppressed and no suppression observed in direct γ .
- ❖ In most peripheral events, the degree of enhancement of π^0 matches that of direct γ .

Nuclear Modification Factor of π^0 s and direct γ s

Experimentally determined N_{coll}

$$\langle N_{coll}^{exp} \rangle = \frac{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right)_{dAu}}{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right)_{pp}}$$

High p_T direct γ s are transparent to QGP and scales exactly with the number of binary collisions in an event sample.

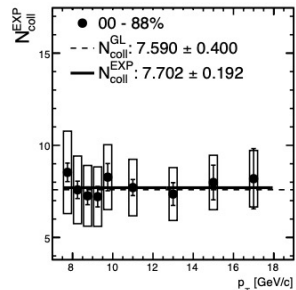
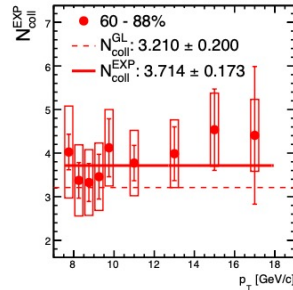
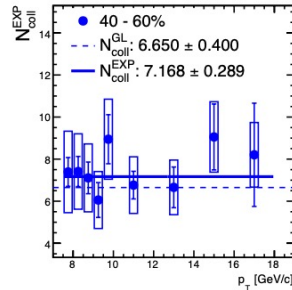
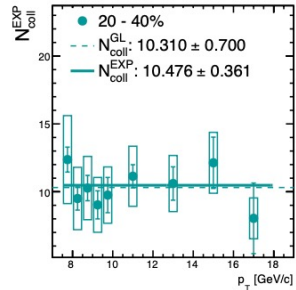
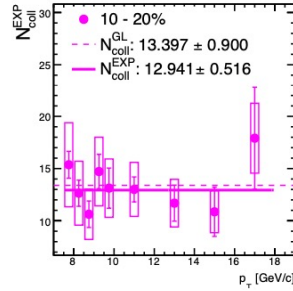
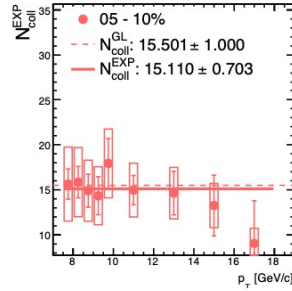
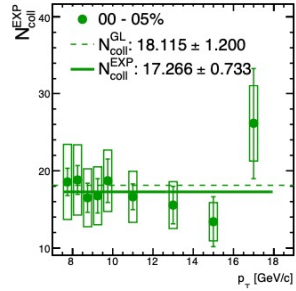
Better way to measure N_{coll} than standard Glauber model.

Can also be used to test the validity of any other modified Glauber models.

Note : on the right hand side, the numerator and the denominator are p_T dependent. The value on the left is obtained by dividing each p_T point and then doing a linear fit to get an average value.

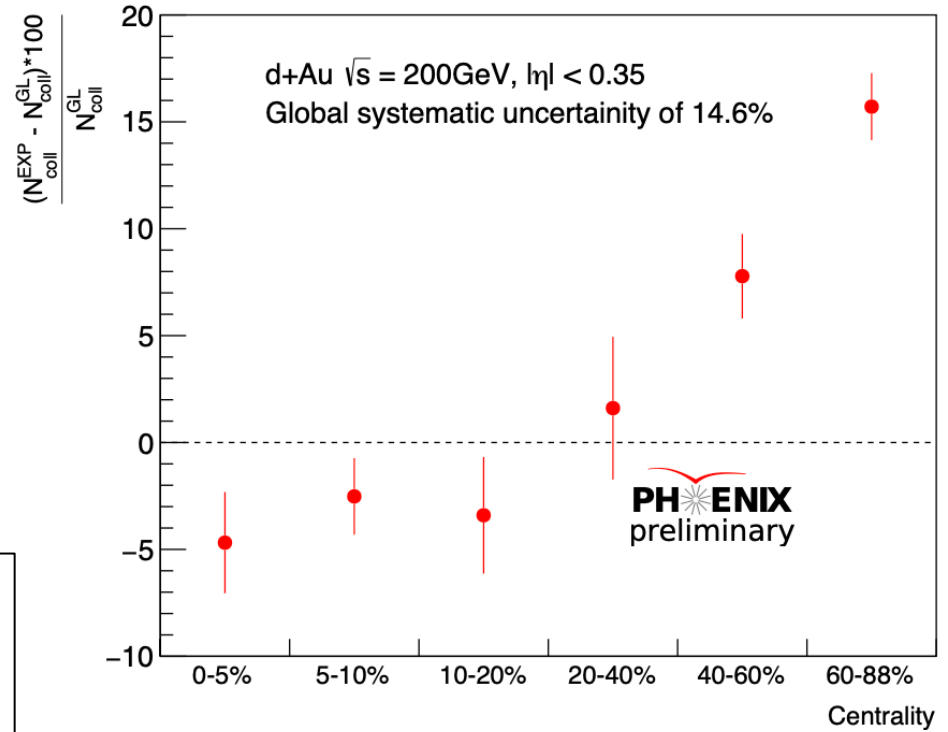
Deriving N_{coll} from direct γ s

Bias in N_{coll} is observed in peripheral d+Au collisions



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$$\langle N_{coll}^{exp} \rangle = \frac{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right) dAu}{\left(\frac{d^2 N_\gamma}{dp_T d\eta} \right) pp}$$



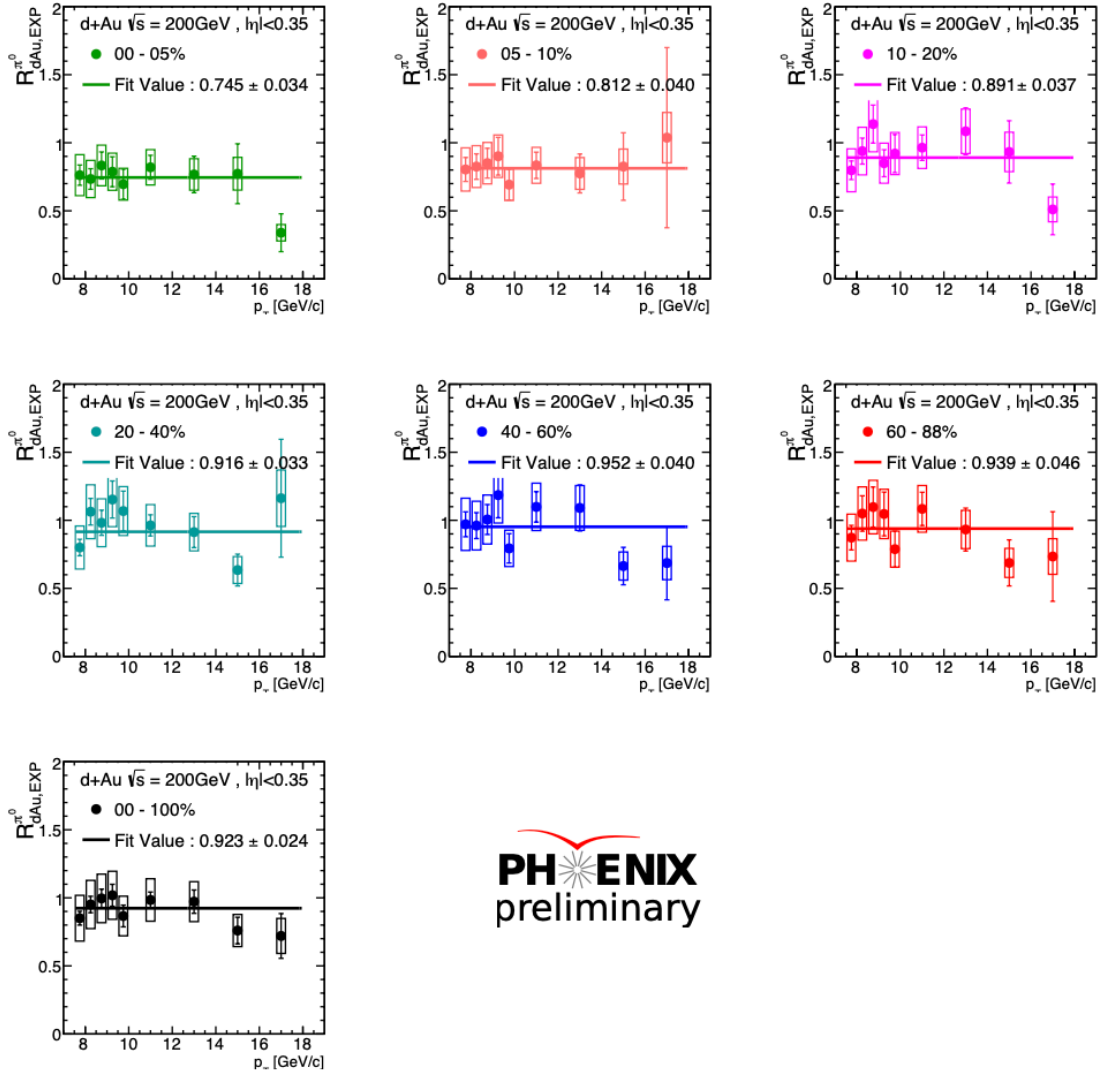
Deriving unbiased $R_{dAu}^{\pi^0}$

$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

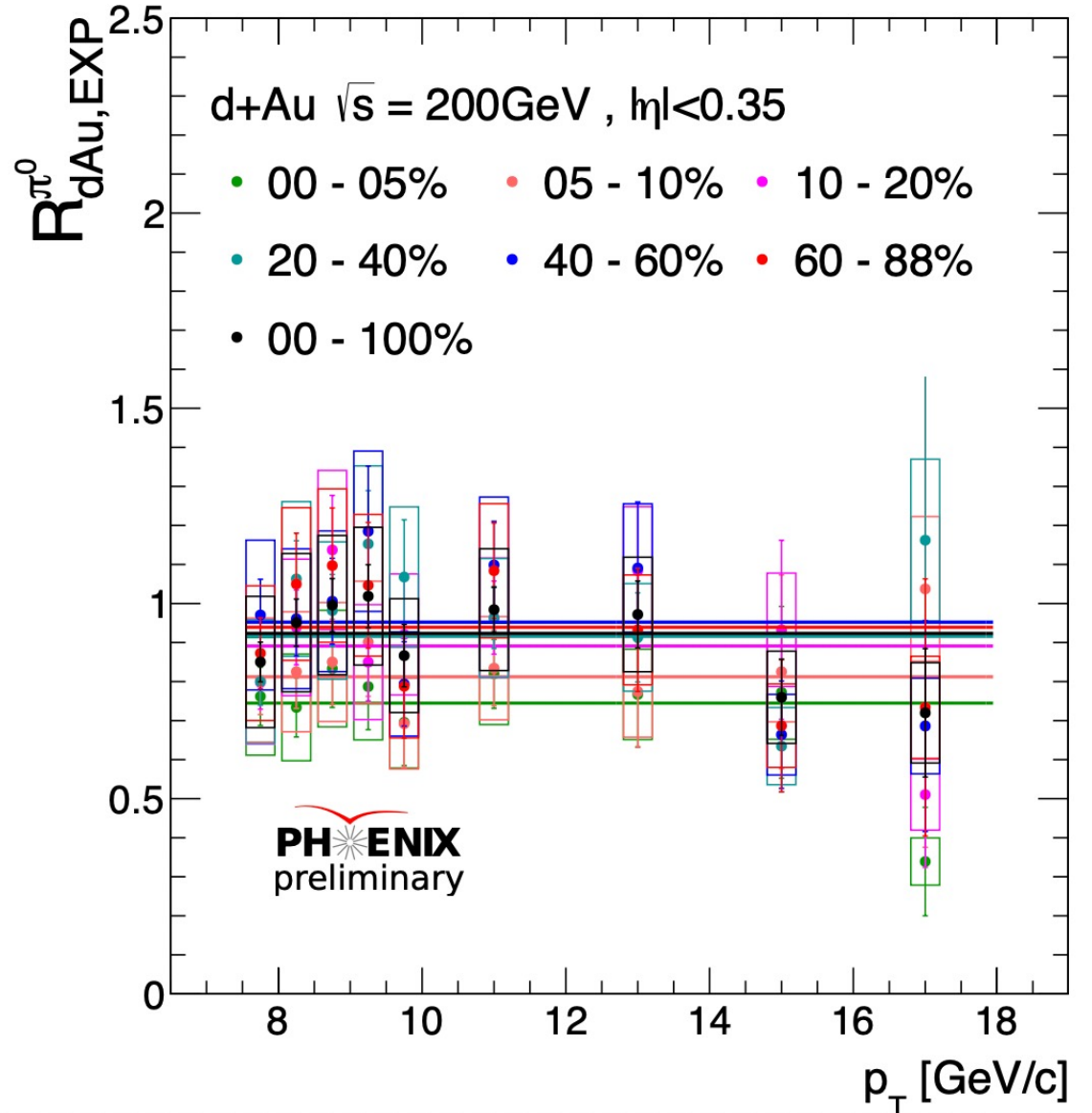
Using experimentally derived Ncoll = Renormalising $R_{dAu,GL}^{\pi^0}$ with $R_{dAu,GL}^{\gamma}$

Unbiased $R_{dAu}^{\pi^0}$

$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

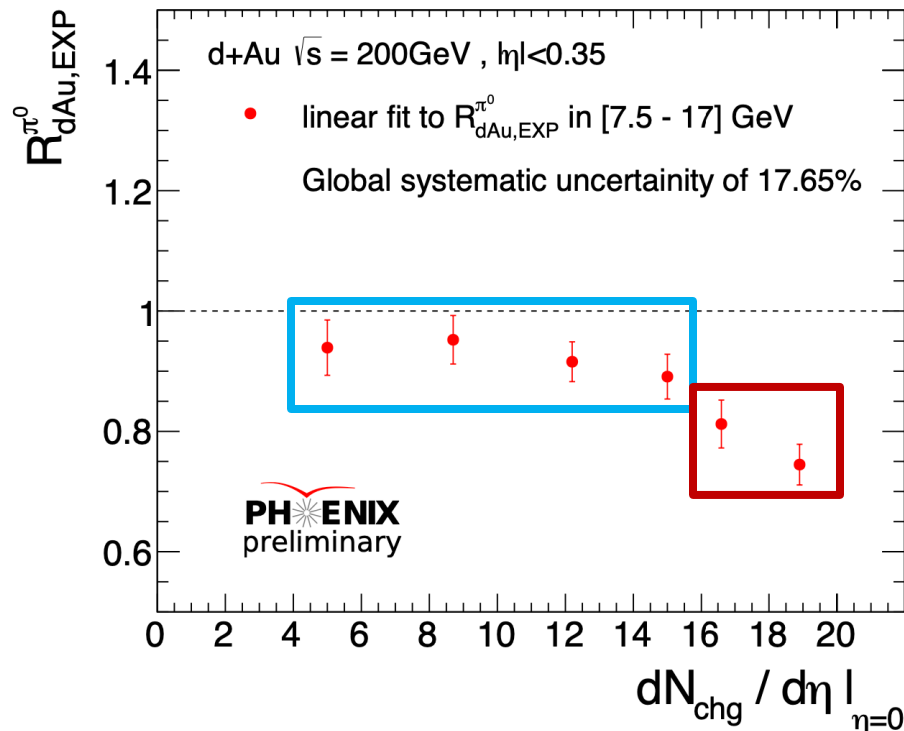


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$$R_{dAu,EXP}^{\pi^0} = \frac{R_{dAu,GL}^{\pi^0}}{R_{dAu,GL}^{\gamma}} = \frac{(Y_{dAu}^{\pi^0}/Y_{pp}^{\pi^0})}{(Y_{dAu}^{\gamma}/Y_{pp}^{\gamma})} = \frac{(Y_{dAu}^{\pi^0}/Y_{dAu}^{\gamma})}{(Y_{pp}^{\pi^0}/Y_{pp}^{\gamma})}$$

First evidence of suppression in $R_{dAu,EXP}^{\pi^0}$



- ❖ After correcting for the bias, the $R_{dAu,EXP}^{\pi^0}$
 - ❖ Does not show any enhancement in peripheral collisions
 - ❖ Does show suppression in central collisions
- ❖ There is no identified source of systematic error which has centrality dependence.
- ❖ Detailed study including,
 - ❖ ongoing analysis of p+Au, $^3\text{He}+\text{Au}$
 - ❖ initial state effects on the production mechanism of π^0 s and direct γ s

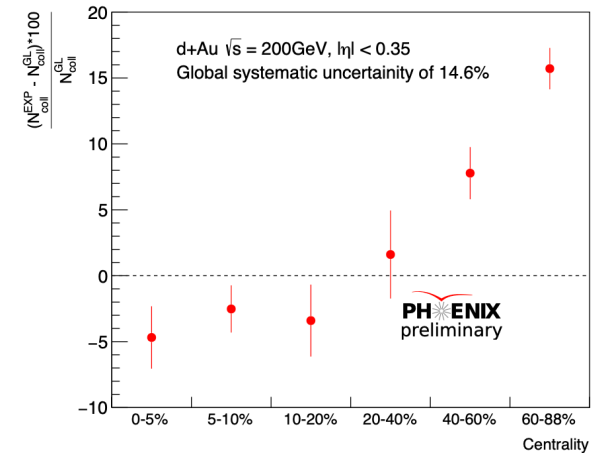
is necessary to understand whether this observed suppression is an initial or final state effect.



SUMMARY

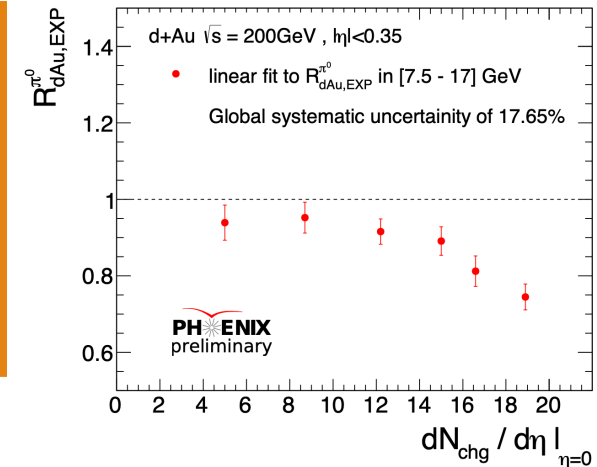
Bias in centrality determination in small system collisions

- ❖ Determination of N_{coll} by using Glauber Model is biased. Specifically for peripheral events



Evidence of $R_{dAu}^{\pi^0}$ suppression in most central

- ❖ Normalized nuclear modification factor of π^0 s using direct γ s
- ❖ Most central events show suppression of $\sim 15\%$





The experimental nuclear physics group, Stony Brook, Aug 2021

THANK YOU

For all the important lessons in physics and beyond

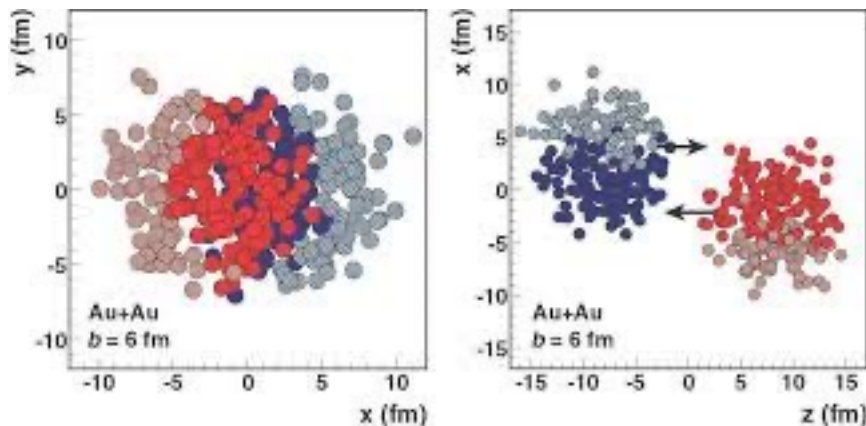
Backup

Definition

Nuclear Modification Factor

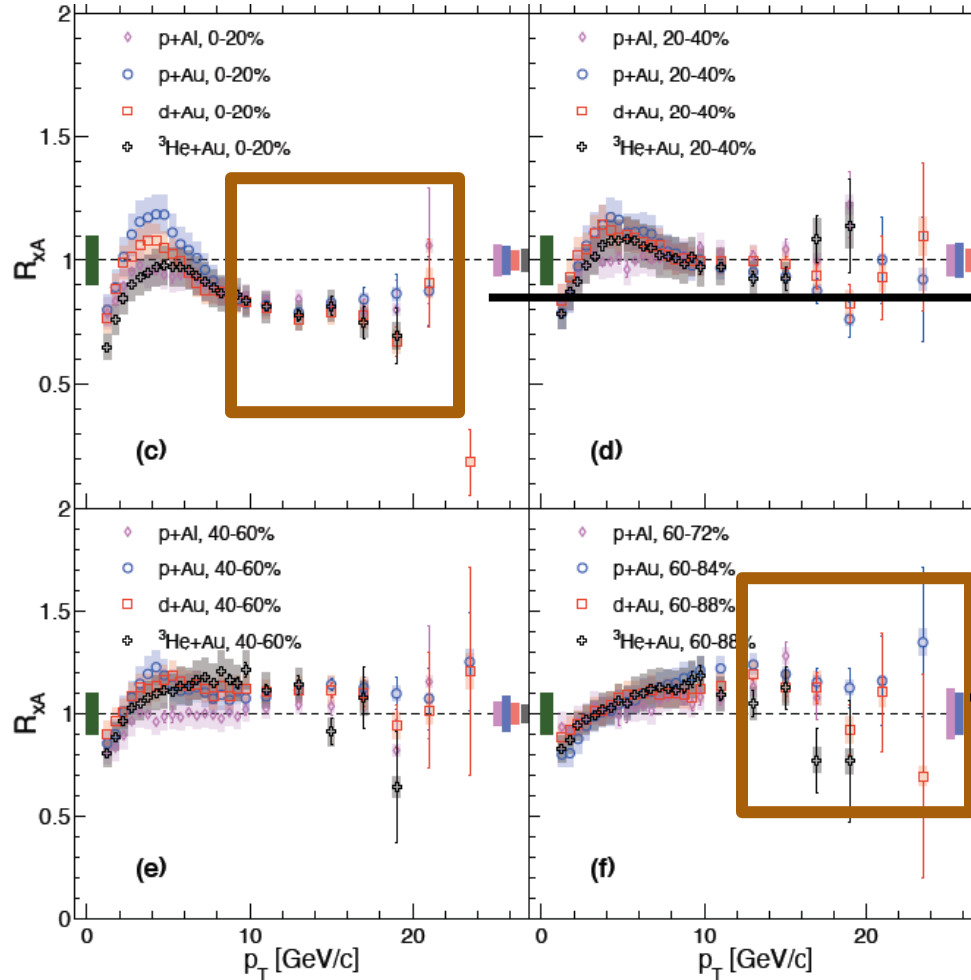
$$R_{AB}(p_T) = \frac{\left(\frac{d^2 N}{dp_T d\eta}\right)_{AB}}{\langle N_{coll} \rangle_{AB} * \left(\frac{d^2 N}{dp_T d\eta}\right)_{pp}} = \frac{Y(AB) \text{ d+Au}}{\langle N_{coll} \rangle_{AB} * Y(pp) \text{ p+p}}$$

How is heavy ion collision different from a scaled p+p collision



Initial motivation

Using direct photons at high p_T to measure the number of binary collisions (N_{coll}) in a system

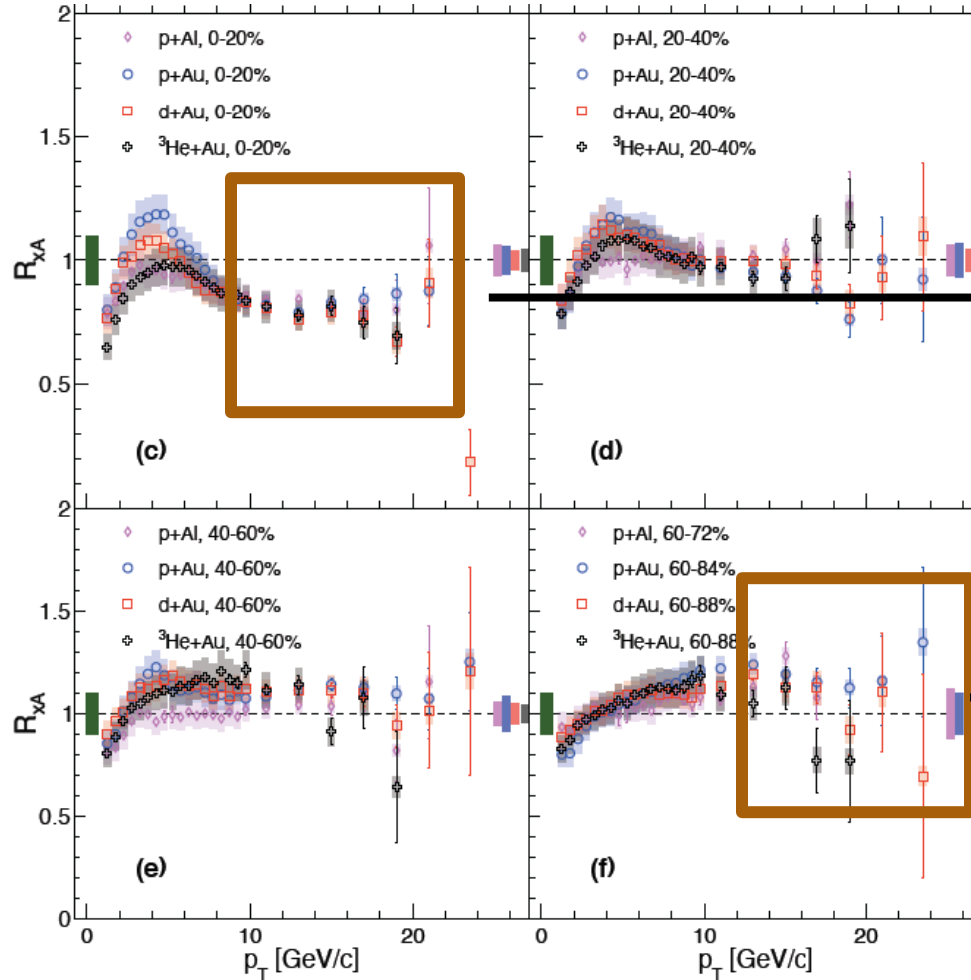


Observed suppression in central collisions remains even after correcting for bias in N_{coll} determination

Observed enhancement in peripheral collisions is an artifact of bias in the N_{coll} determination using Glauber model for small systems

Initial motivation

Using direct photons at high p_T to measure the number of binary collisions (N_{coll}) in a system



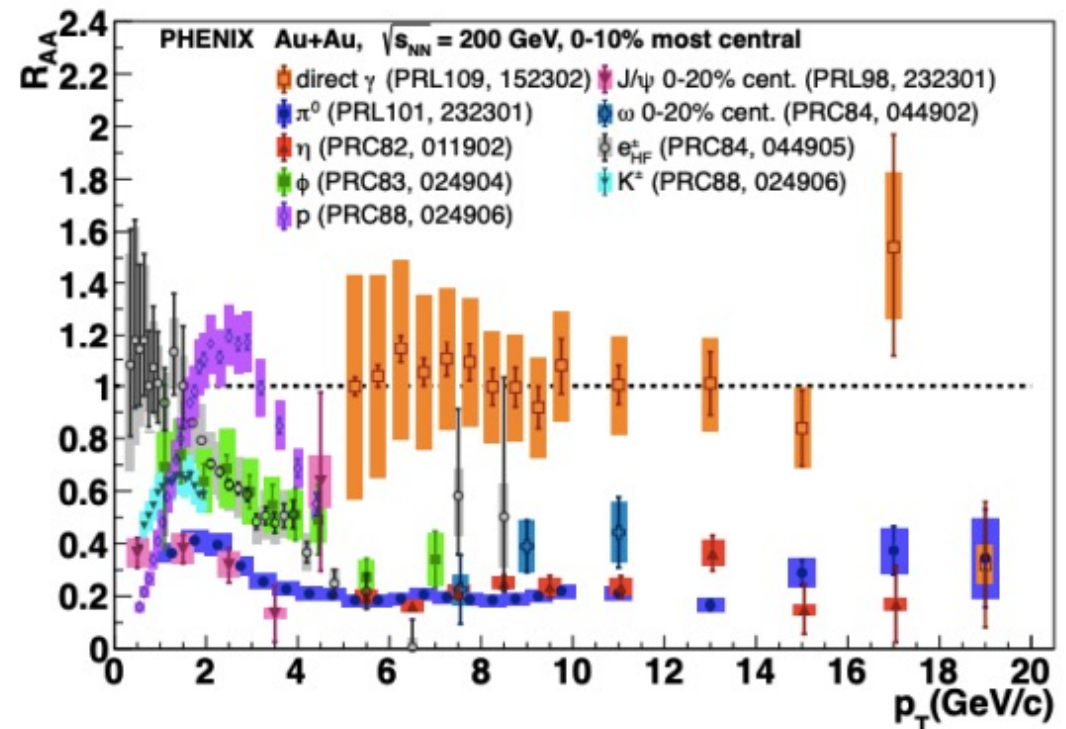
Observed suppression in central collisions remains even after correcting for bias in N_{coll} determination

Observed enhancement in peripheral collisions is an artifact of bias in the N_{coll} determination using Glauber model for small systems

$R_{AB}(p_T) < 1$ is a signature of QGP

1) $\pi^0, \eta, \phi, J/\psi, \omega$ interact with the QGP $\rightarrow R_{AB}(p_T)$ is suppressed

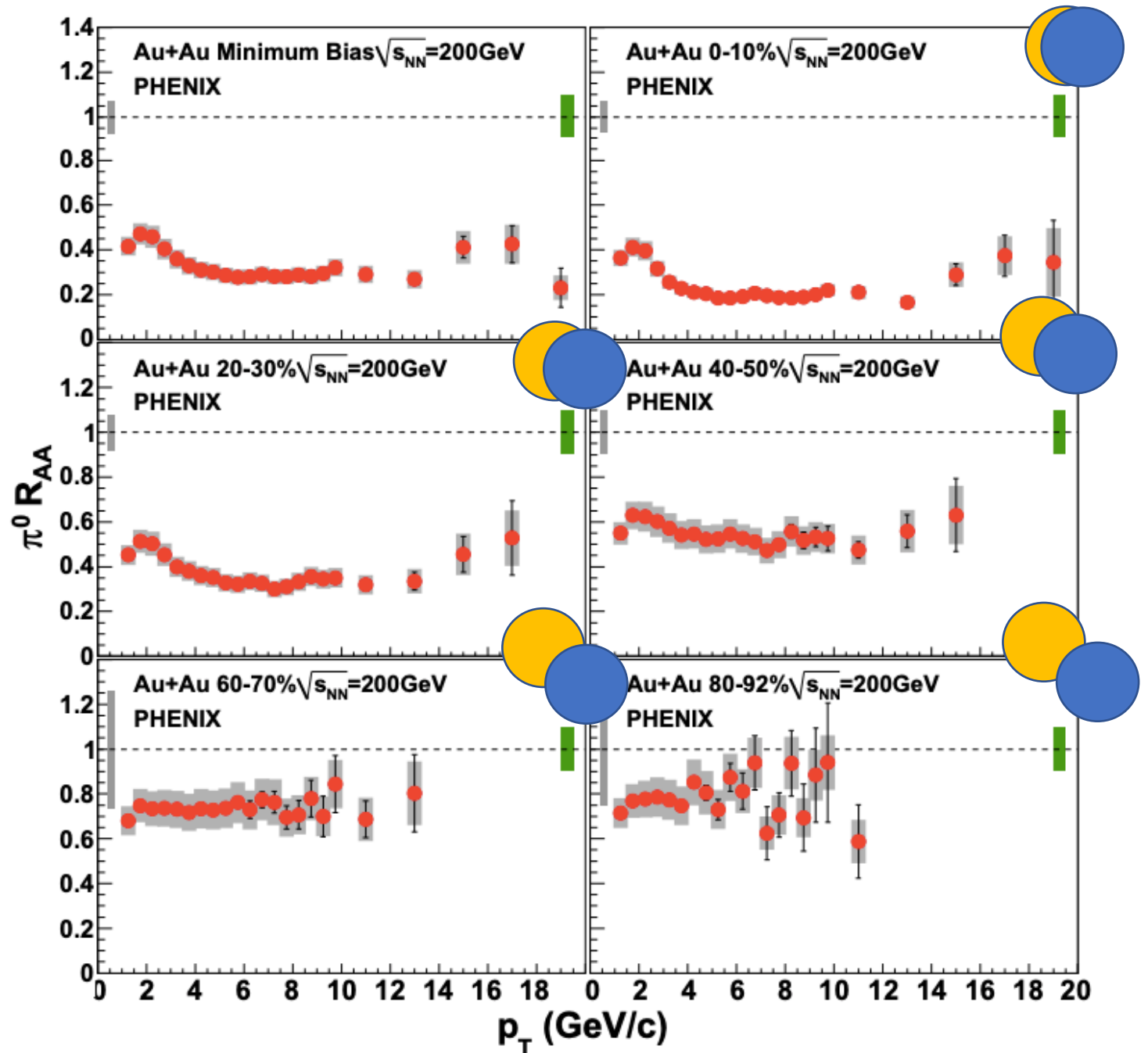
2) Direct photon is transparent to the QGP $\rightarrow R_{AB}(p_T)$ is unity



Centrality binned R_{AA} of π^0

in Au+Au collisions

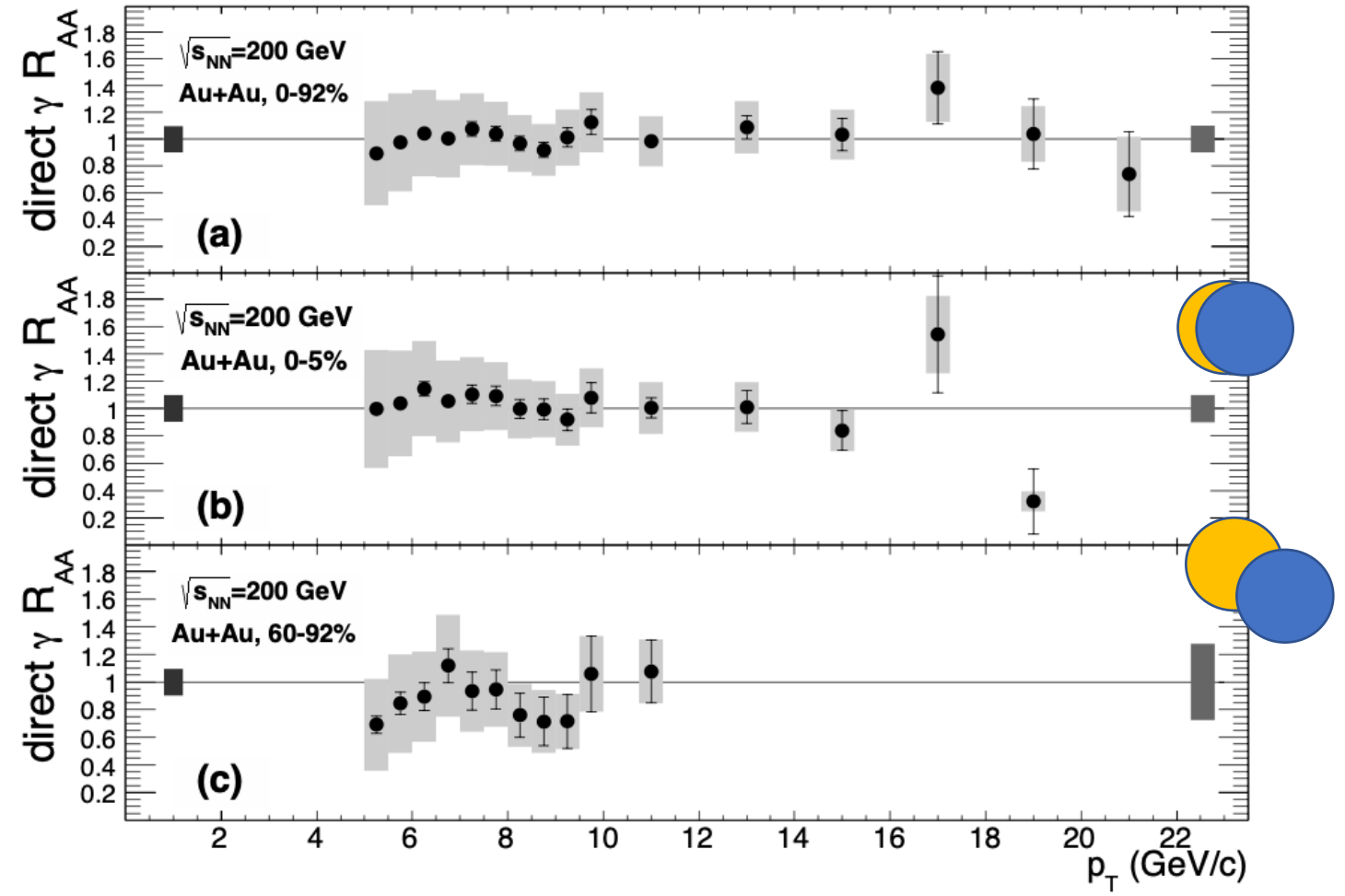
- Most central collisions show the most suppression.
- Vanishing suppression at peripheral collisions
- Intuitive trend in centrality for a system with QGP creation



Centrality binned R_{AA} of direct γ in Au+Au collisions

- Unity at all centralities.
- Direct photons are transparent to QGP
- Using high p_T direct photons, we obtain :

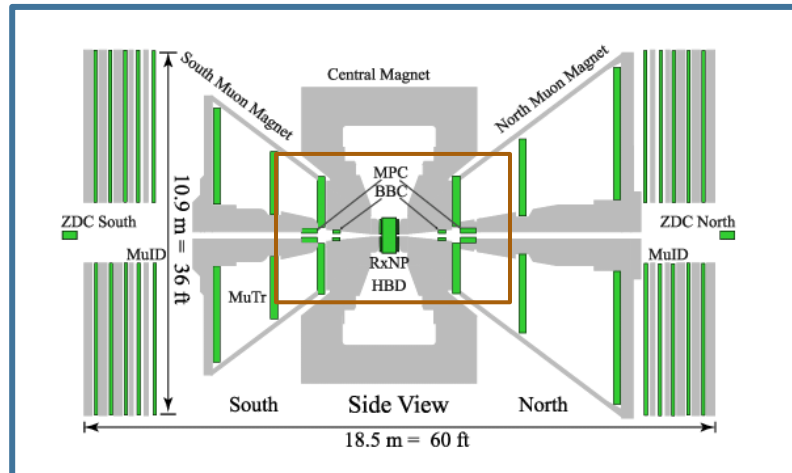
$$N_{coll}^{exp} = \frac{Yield_{AA}^{\gamma}}{Yield_{pp}^{\gamma}}$$



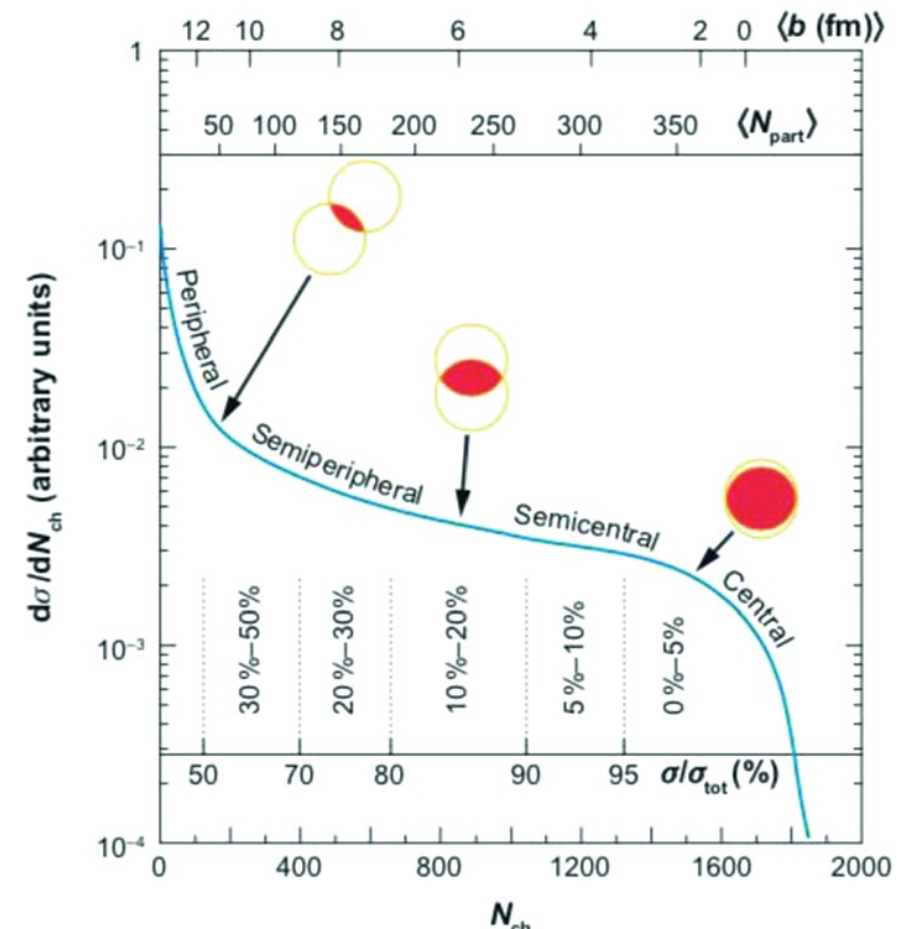
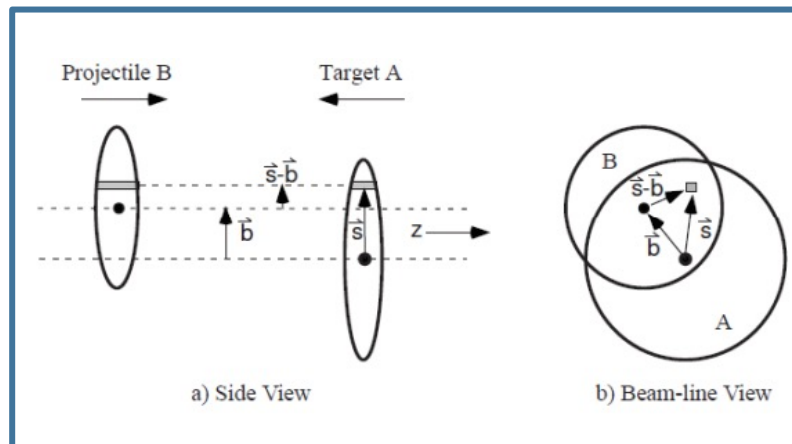
Measurement of average number of binary collisions from bulk observables

$$\langle N_{coll} \rangle_{AB} * Y(pp)$$

Number of charged particle from BBC
 $3.1 < |\eta| < 4$



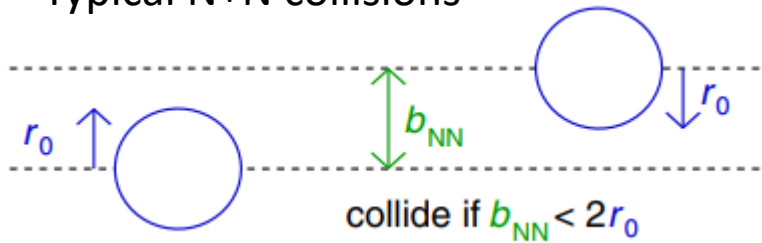
Standard Glauber model gives mapping of charged particle in forward region to number of binary collisions of the event.



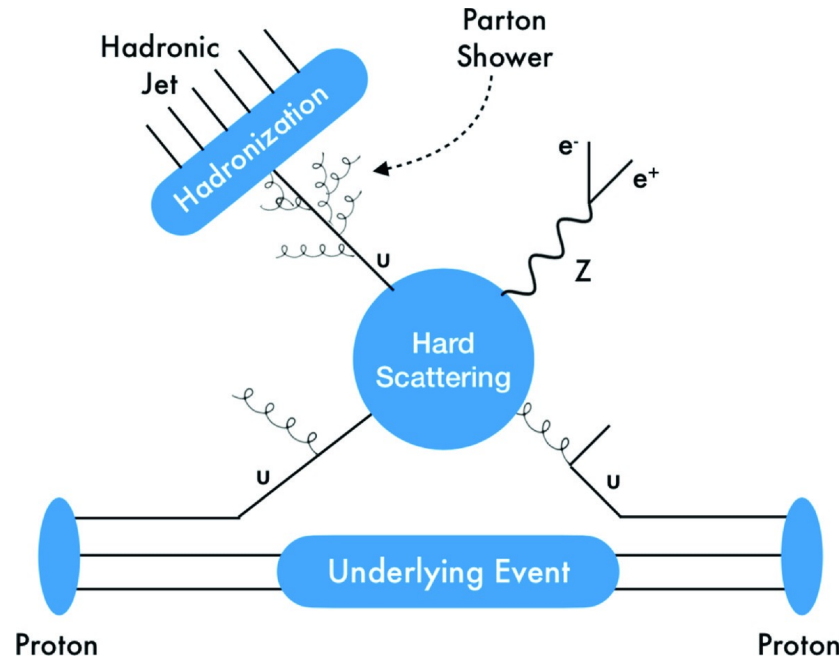
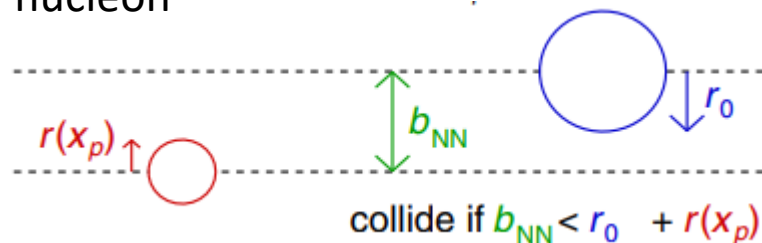
Origin of the bias...?

high-x (effective) size fluctuations

Typical N+N collisions



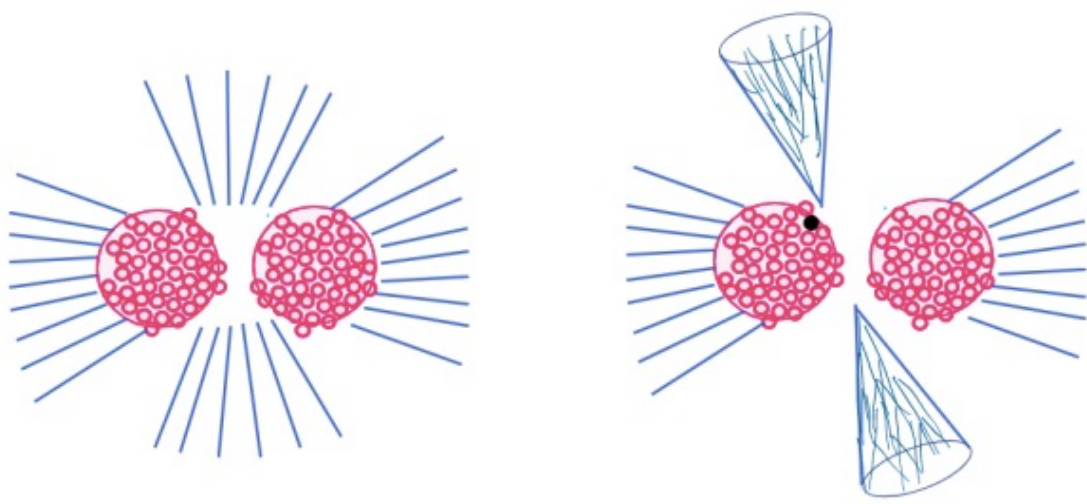
N+N collisions with large- x_p projectile nucleon



The high-X parton creates the hard scattering event. But the underlying event is severely depleted.

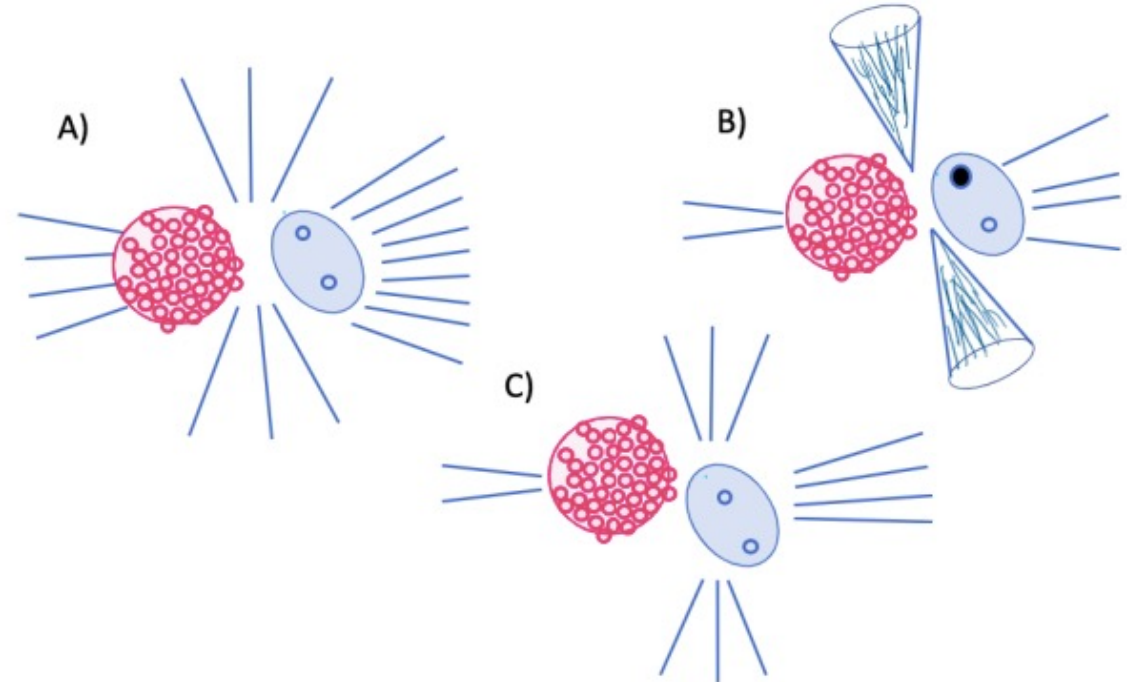
This can be thought of as
 a) energy conservation or
 b) change in the cross-section of the nuclei due to the presence of high-X parton.

Phys. Rev. C 94 (2016), 024915



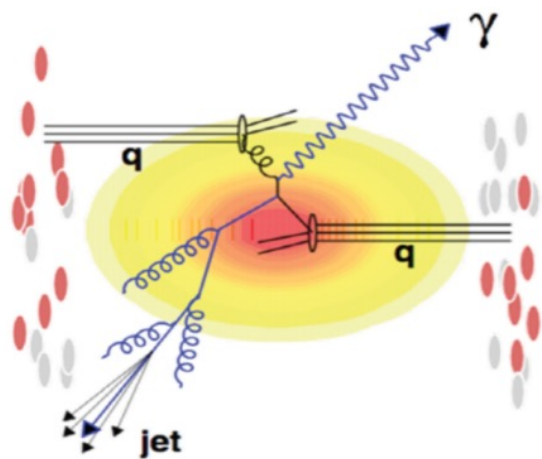
In a heavy-ion collision, the presence of one high- X parton nuclei, creates the jets, but the average underlying event isn't affected as there are several other partons for interactions.

In a d+Au collision, the presence of one high- X parton depletes the underlying event and there are not enough other interactions to compensate for this. Thus a central d+Au event will look like a peripheral d+Au event. This is a p_T (or x) dependent change. The bin-shift is larger at higher momentum.



This shrinking nucleon model has a prediction for $R_{dAu}(x)$ and thus we can compare it to our data.

$R_{dAu}^{\pi^0}$ and R_{dAu}^{γ} as a function of parton momentum x



$$x_p = 2p_T^{\text{jet}} / \sqrt{s_{NN}} \approx 2p_T^{\pi^0} / (0.75 * \sqrt{s_{NN}})$$

$$x_p = 2p_T^{\text{direct}\gamma} / \sqrt{s_{NN}}$$

Does our data fit the expectation from the “shrinking nucleon” picture?

