



# Azimuthal angular correlation of jets from soft gluon radiation

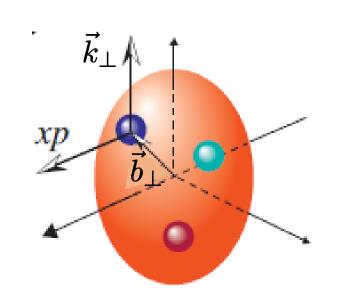
Yoshitaka Hatta BNL & RIKEN BNL

in collaboration with Bowen Xiao, Feng Yuan, Jian Zhou

2010.10774 (PRL) 2106.05307 (PRD)

sPHENIX RBRC workshop July 20-22, 2022

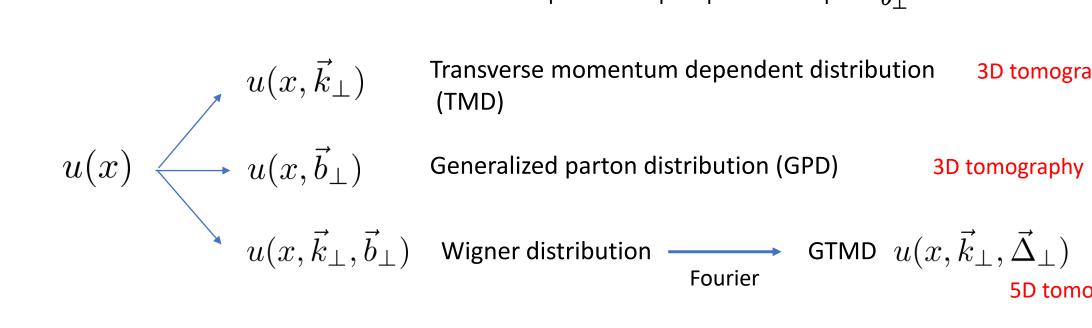
### Multi-dimensional tomography



$$u(x) = \int \frac{dz^{-}}{4\pi} \langle P|\bar{u}(0)\gamma^{+}u(z^{-})|P\rangle e^{ixP^{+}z^{-}}$$

Ordinary parton distribution functions (PDF) can be viewed as the 1D tomographic image of the nucleon

The nucleon is much more complicated! Partons also have transverse momentum  $ec{k}_{\perp}$ and are spread in impact parameter space  $\, ec{b}_{\,\, |} \,$ 

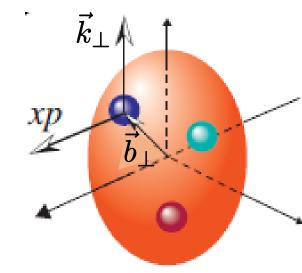


3D tomography

Wigner distribution 
$$\longrightarrow$$
 GTMD  $u(x, \vec{k}_\perp, \vec{\Delta}_\perp)$  SD tomography

## Angular correlations

Rich angular correlations between impact parameter  $\, ec{b}_{\perp} \,$  and transverse momentum  $\, ec{k}_{\perp} \,$ 



$$\begin{split} W(x,\vec{b}_{\perp},\vec{k}_{\perp}) \\ = W_0 + \cos 2(\phi_{b_{\perp}} - \phi_{k_{\perp}})W_2 + \cos(\phi_{b_{\perp}} - \phi_{k_{\perp}})W_O + \sin(\phi_{b_{\perp}} - \phi_{k_{\perp}})W_{OAM} \ + \cdots \\ \text{Elliptic Wigner} & \text{Odderon} & \text{Orbital angular momentum} \end{split}$$

Fourier transform → GTMD

$$W(x, \vec{\Delta}_{\perp}, \vec{k}_{\perp}) = W_0 + \cos 2(\phi_{\Delta_{\perp}} - \phi_{k_{\perp}})W_2 + \cdots$$

### Wigner distribution: Is it measurable?

#### In quantum optics, yes!

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## Measurement of the Wigner Distribution and the Density Matrix of a Light Mode Using Optical Homodyne Tomography: Application to Squeezed States and the Vacuum

D. T. Smithey, M. Beck, and M. G. Rayme Department of Physics and Chemical Physics Institute, University of Ore;

#### A. Faridani.

Department of Mathematics, Oregon State University, Corvalli (Received 16 November 1992)

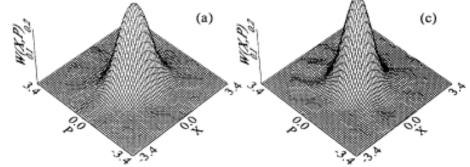
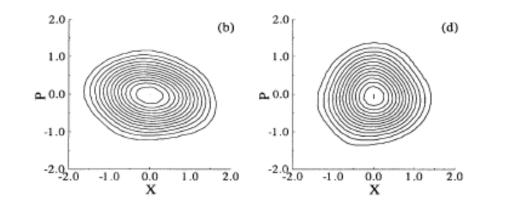
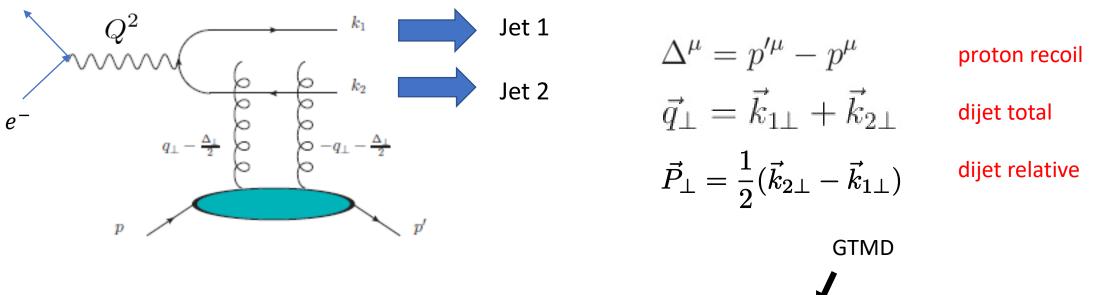


FIG. 1. Measured Wigner distributions for (a),(b) a squeezed state and (c),(d) a vacuum state, viewed in 3D and as contour plots, with equal numbers of constant-height contours. Squeezing of the noise distribution is clearly seen in (b).



#### Probing gluon Wigner in exclusive dijet production in DIS

YH, Xiao, Yuan (2016)



$$\frac{d\sigma}{dy_{1}dy_{2}d^{2}\vec{\Delta}_{\perp}d^{2}\vec{P}_{\perp}^{\perp}} \propto z(1-z)[z^{2}+(1-z)^{2}]\int d^{2}q_{\perp}d^{2}q'_{\perp}S(q_{\perp},\Delta_{\perp})S(q'_{\perp},\Delta_{\perp}) 
\times \left[\frac{\vec{P}_{\perp}}{P_{\perp}^{2}+\epsilon^{2}} - \frac{\vec{P}_{\perp}-\vec{q}_{\perp}}{(P_{\perp}-q_{\perp})^{2}+\epsilon^{2}}\right] \cdot \left[\frac{\vec{P}_{\perp}}{P_{\perp}^{2}+\epsilon^{2}} - \frac{\vec{P}_{\perp}-\vec{q}'_{\perp}}{(P_{\perp}-q'_{\perp})^{2}+\epsilon^{2}}\right] \quad \epsilon^{2} = z(1-z)Q^{2}$$

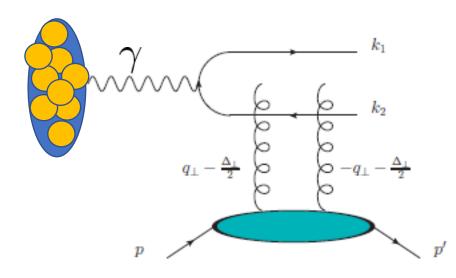
 $\sim d\sigma_0 + 2\cos 2(\phi_P - \phi_\Delta)d\tilde{\sigma}$ 

`elliptic flow', expected to be a few percent effect

## Exclusive dijet in UPC

UPC: where the heavy-ion and EIC communities can meet

Hagiwara, YH, Pasechnik, Tasevsky, Teryaev (2017)



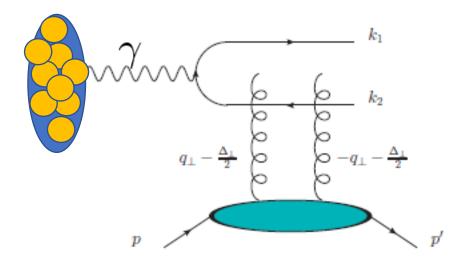
$$\frac{d\sigma^{pA}}{dy_1 dy_2 d^2 \vec{k}_{1\perp} d^2 \vec{k}_{2\perp}} = \omega \frac{dN}{d\omega} \frac{N_c \alpha_{em} (2\pi)^4}{P_\perp^2} \sum_f e_f^2 2z (1-z) (z^2 + (1-z)^2) \left( \mathbf{A^2} + 2\cos 2(\phi_P - \phi_\Delta) \mathbf{AB} \right)$$
 photon flux  $\propto Z^2$ 

When  $Q^2 \approx 0$  inversion can be done analytically.  $S(\vec{q}_\perp, \vec{\Delta}_\perp) = S_0(q_\perp, \Delta_\perp) + 2\cos 2(\phi_q - \phi_\Delta)\tilde{S}(q_\perp, \Delta_\perp)$ 

$$S_0(P_{\perp}, \Delta_{\perp}) = \frac{1}{P_{\perp}} \frac{\partial}{\partial P_{\perp}} A(P_{\perp}, \Delta_{\perp}). \qquad S_1(P_{\perp}, \Delta_{\perp}) = \frac{\partial B(P_{\perp}, \Delta_{\perp})}{\partial P_{\perp}^2} - \frac{2}{P_{\perp}^2} \int^{P_{\perp}^2} \frac{dP_{\perp}'^2}{P_{\perp}'^2} B(P_{\perp}, \Delta_{\perp})$$

## CMS dijet measurements

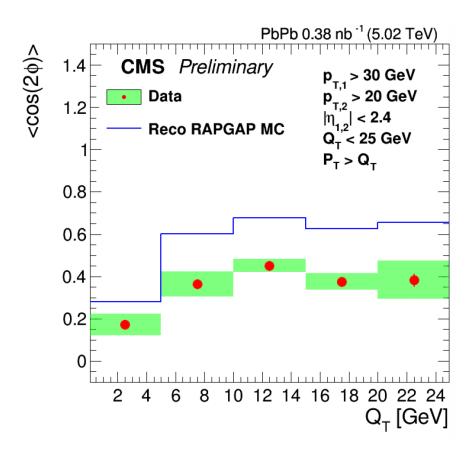
Partly motivated by our proposal, CMS measured dijet angular correlation in PbPb UPC



Use  $ec{q}_\perp=ec{k}_{1\perp}+ec{k}_{2\perp}$  as a proxy for  $ec{\Delta}_\perp$   $ec{\Delta}_\perp=-ec{q}_\perp \ \ ext{to leading order}$ 

Measured the  $\cos 2\phi$  correlation between  $q_{\perp}$  and  $P_{\perp}$ , instead of that between  $\vec{\Delta}_{\perp}$  and  $P_{\perp}$ 

#### **CMS-PAS-HIN-18-011**



Very large asymmetry!
Anything to do with the elliptic Wigner?

#### Dijet with soft gluons: general consideration

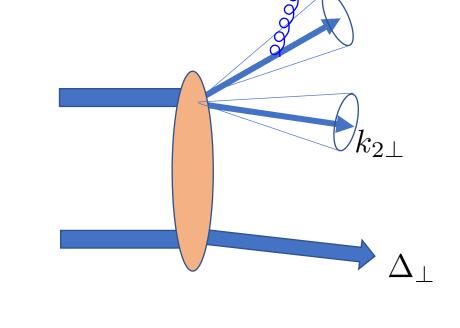
To leading order, dijet total momentum is equal to the proton recoil momentum

$$q_{\perp} = k_{1\perp} + k_{2\perp} = -\Delta_{\perp}$$

$$\frac{d\sigma}{d^2q_{\perp}} = \frac{d\sigma}{d^2\Delta_{\perp}}$$

With soft radiation, this becomes

$$\vec{q}_{\perp} = -\vec{\Delta}_{\perp} - \sum_{i} \vec{k}_{g\perp}^{i}$$



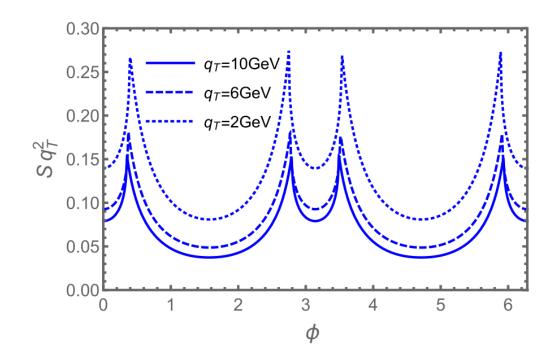
 $k_{g\perp}$  tend to be along jet directions  $o q_{\perp}$  tends to be along jet directions  $o \cos 2\phi$ 

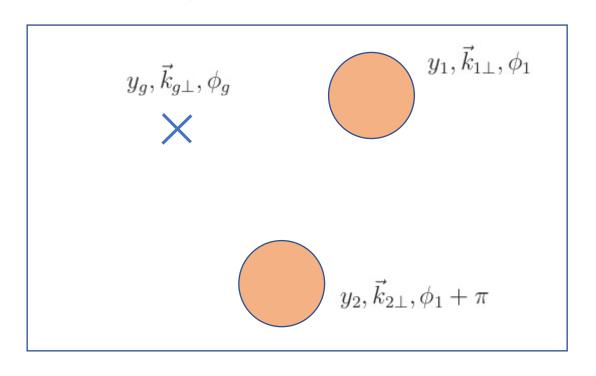
$$rac{d\sigma}{d^2P_\perp d^2q_\perp} = \int d^2q_\perp' rac{d\sigma_0}{d^2P_\perp d^2q_\perp'} S_J(q_\perp - q_\perp'),$$
 measured primordial Soft factor

#### Angular distribution of soft gluons around a dijet system

To one-loop,

$$S_J(q_{\perp}) = \delta(q_{\perp}) + \frac{\alpha_s}{2\pi^2} \int dy_g \left(\frac{k_1 \cdot k_2}{k_1 \cdot k_g k_2 \cdot k_g}\right)_{\vec{q}_{\perp} = -\vec{k}_{g\perp}}$$





Carefully integrate over the soft gluon phase space avoiding jet cones

 $\rightarrow$  all even harmonics  $\cos 2n(\phi_{q_{\perp}} - \phi_{P_{\perp}})$ 

## One-loop analytical result

$$g^{2} \int \frac{d^{3}k_{g}}{(2\pi)^{3} 2E_{k_{g}}} \delta^{(2)}(q_{\perp} + k_{g\perp}) C_{F} S_{g}(k_{1}, k_{2})$$

$$= \frac{C_{F} \alpha_{s}}{\pi^{2} q_{\perp}^{2}} [c_{0}^{\text{diff}}(q_{\perp}^{2}) + 2\cos(2\phi) c_{2}^{\text{diff}}(q_{\perp}^{2}) + \cdots].$$

$$c_n(q_\perp) = c_n(0) + \mathcal{O}(q_\perp^a/P_\perp^a)$$
 power corrections

$$\begin{split} c_0^{\text{diff}}(0) &= \ln \frac{a_0}{R^2}, \qquad a_0 = 2 + 2 \cosh(\Delta y_{12}) \\ c_2^{\text{diff}}(0) &= \ln \frac{a_2}{R^2}. \qquad \ln a_2 = \Delta y_{12} \sinh \Delta y_{12} - \cosh \Delta y_{12} \ln \left[ 2(1 + \cosh \Delta y_{12}) \right] \end{split}$$

$$\frac{d\sigma}{d\phi} \sim \alpha_s \cos 2\phi \int \frac{dq_\perp^2}{q_\perp^2} = \infty$$

Angular dependent cross section divergent at fixed order.

#### Resummation

Catani, Grazzini, Torre, (2014), Catani, Grazzini, Sargsyan (2017) YH, Yuan, Xiao, Zhou (2020~)

Fourier transform  $q_{\perp} \rightarrow b_{\perp}$ 

part

$$\begin{array}{ll} \text{Angular-independent} & \int \frac{d\vec{q}_{\perp}}{(2\pi)^2} e^{i\vec{b}_{\perp} \cdot \vec{q}_{\perp}} \left[ \frac{1}{q_{\perp}^2} \right]_{+} = -\frac{1}{4\pi} \ln \frac{b_{\perp}^2 P_{\perp}^2}{c_0^2} & \xrightarrow{\text{IR divergent}} \\ \text{part} & \xrightarrow{\text{regularized by the plus prescription}} \end{array}$$



exponentiation 
$$\exp\left(-\frac{2C_Fc_0}{\pi}\int_{\mu_b}^{P_\perp}\frac{d\mu}{\mu}\alpha_s(\mu)\right)$$

Angular-dependent part

$$\int \frac{d\vec{q}_{\perp}}{(2\pi)^2} e^{i\vec{b}_{\perp} \cdot \vec{q}_{\perp}} \frac{\cos 2\phi_{q_{\perp}}}{q_{\perp}^2} = -\frac{\cos 2\phi_{b_{\perp}}}{2\pi} \int_0^{\infty} \frac{dq_{\perp}}{q_{\perp}} J_2(b_{\perp}q_{\perp})$$

 $=rac{1}{2}$  No plus-prescription, but finite!

At higher orders, there are logarithms in the b-space even for the angular-dependent part.

→ Can be resummed by the same Sudakov factor

Fourier transform back to q-space

$$\frac{d\sigma}{dq_{\perp}^2 d\phi} = \cos n\phi \int_0^{\infty} db_{\perp} b_{\perp} J_n(b_{\perp} q_{\perp}) e^{-Sudakov} \dots$$

$$\langle \cos n\phi \rangle \sim \frac{1}{q_{\perp}^2} \longleftrightarrow \langle \cos(n\phi) \rangle \propto q_{\perp}^n$$

before resummation

after resummation

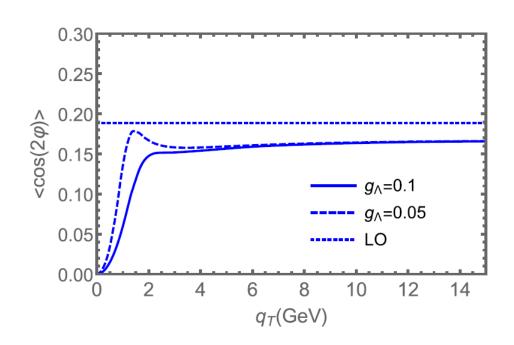
Azimuthal asymmetries in QCD hard scattering: infrared safe but divergent

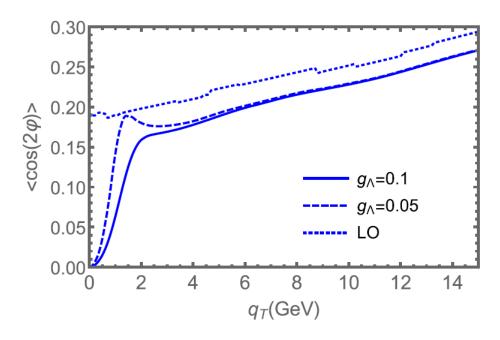
Stefano Catani,<sup>a</sup> Massimiliano Grazzini<sup>b</sup> and Hayk Sargsyan<sup>b</sup>

<sup>a</sup>INFN — Sezione di Firenze and Dipartimento di Fisica e Astronomia, Università di Firenze, I-50019 Sesto Fiorentino, Florence, Italy IEP06 (2017) 017

## Towards explaining the CMS data

UPC at the LHC,  $P_{\perp}=35 {\rm GeV}$  R=0.4



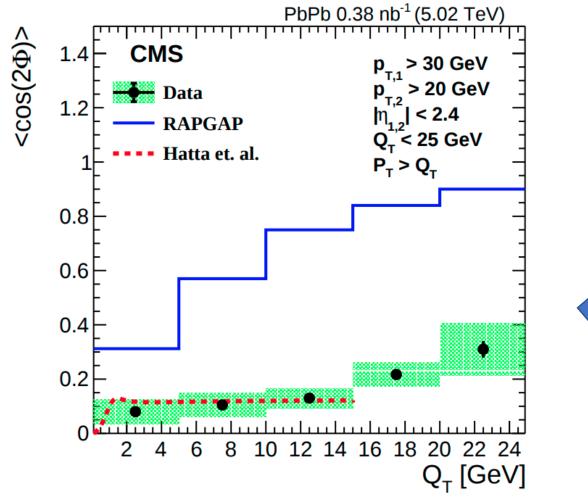


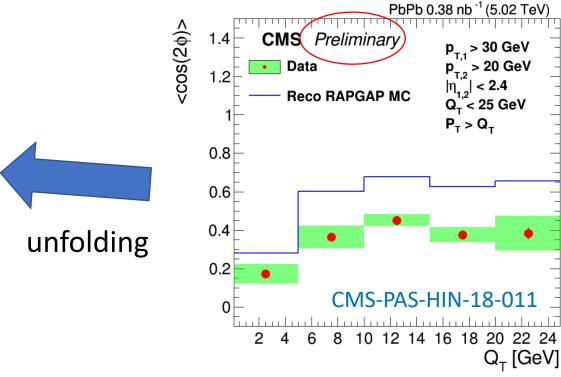
Part of power corrections included → monotonically rising behavior

#### Lesson learned:

$$ec{q}_{\perp} = ec{k}_{1\perp} + ec{k}_{2\perp}$$
 cannot be a proxy for  $ec{\Delta}_{\perp}$  .

## CMS-PAS-HIN-18-011 2205.00045 (updated plots)





# Dijet in UPC pA/AA at sPHENIX

Take 
$$\sqrt{s_{NN}} = 200 \text{ GeV}$$

Typical WW photon energy in UPC

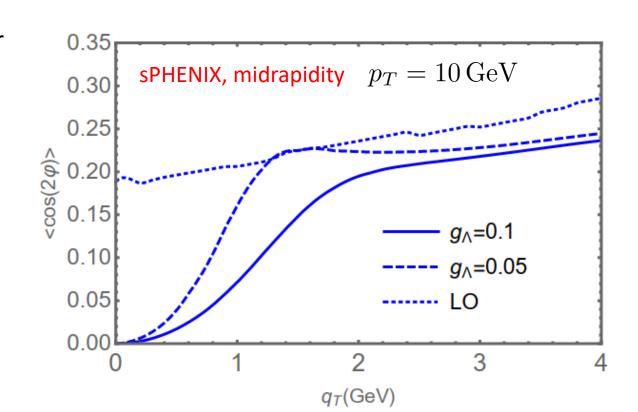
$$\omega = P_{\perp} e^y \sim 2 \text{GeV}$$

Choose y=-1 and/or go to larger photon energy  $\,\omega\,$  to make  $\,P_{\perp}$  larger.

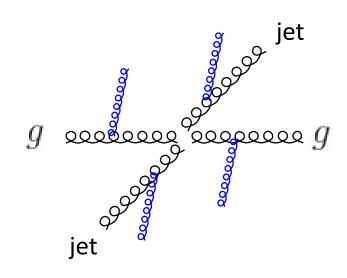
Stronger sensitivity to nonperturbative Sudakov factor

$$e^{-S_{sud}^{pert}(b_{\perp})-2g_{\Lambda}b_{\perp}^2}$$

Can the recoiling proton momentum be measured in pA?



# Inclusive dijet in pp at sPHENIX



$$\left. \frac{d^6 \sigma}{d\Omega} \right|_{\text{soft}} = \sigma_0^{gg} x_g f_g(x_g) \frac{\alpha_s}{2\pi^2} \frac{C_A}{q_\perp^2} \left[ 2 \ln \frac{P_\perp^2}{q_\perp^2} + 2(c_0^{gg} + c_2^{gg} 2 \cos(2\phi) + \cdots) \right]$$

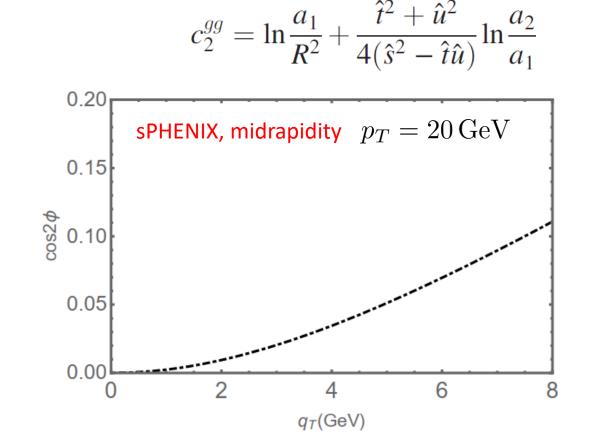
Four colored partons

Double log resummation

→ universal, diagonal in color

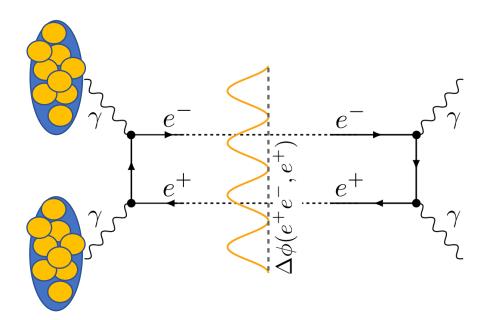
Single log resummation

→ should be done in a matrix form in color space.



## Ultraperipheral AA collisions (UPC) at RHIC and LHC

A large nucleus—copious source of linearly polarized photons



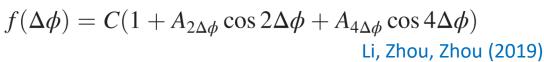
Dielectron production at midrapidity at low transverse momentum in peripheral and semi-peripheral Pb–Pb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV

**ALICE Collaboration** 

PHYSICAL REVIEW LETTERS 121, 212301 (2018)

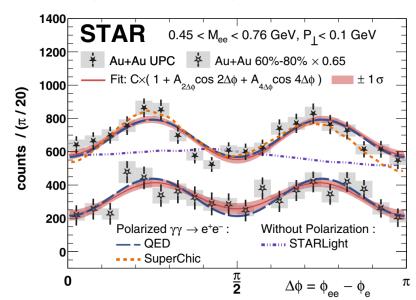
Observation of Centrality-Dependent Acoplanarity for Muon Pairs Produced via Two-Photon Scattering in Pb + Pb Collisions at  $\sqrt{s_{NN}}$  = 5.02 TeV with the ATLAS Detector

M. Aaboud *et al.*\*
(ATLAS Collaboration)

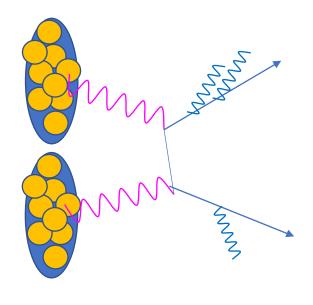


PHYSICAL REVIEW LETTERS 127, 052302 (2021)

Measurement of  $e^+e^-$  Momentum and Angular Distributions from Linearly Polarized Photon Collisions (STAR Collaboration)



## Soft photon resummation in UPC

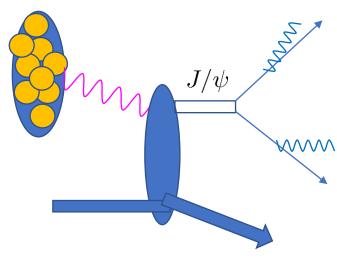


Linearly polarized photon  $\rightarrow \cos 2\phi$  correlation in lepton pairs Background from soft photon radiation

Large anisotropy 
$$\frac{c_2}{c_0}=rac{\lnrac{k_\perp}{2m}}{\lnrac{2k_\perp}{m}}\sim\mathcal{O}(1)$$
  $R\leftrightarrow m/k_\perp$ 

Resummation Klein, Mueller, Xiao, Yuan (2020); YH, Xiao, Yuan, Zhou (2021)

RHIC



 $Q_0 = 40 \text{ MeV} \sim 1/R_A$ 0.0
0
2
4
6
8
10  $q_T/Q_0$ 

linearly polarized photon distribution dominates when

$$q_{\perp} < 100 \, {\rm MeV}$$

At larger momentum, the final state radiation dominates.

Brandenburg et al, 2207.02478

#### Conclusions

- Multi-dimensional tomography (TMD, GPD, Wigner) important theme at the EIC, preview at sPHENIX
- New distributions often probed in jet angular correlations.
- Things took an interesting turn after the CMS measurement. Soft gluon radiations can diminish the signal. Interesting in its own right.
- Dijet total momentum  $\vec{q}_\perp = \vec{k}_{1\perp} + \vec{k}_{2\perp}$  very sensitive to higher order corrections, cannot be a proxy for  $\vec{\Delta}_\perp$