

Chirality, Vorticity and Magnetic Field in Heavy Ion Collisions

UCLA Chirality Retreat

Perspectives on the STAR isobar and gold-gold data

Fuqiang Wang
Purdue University



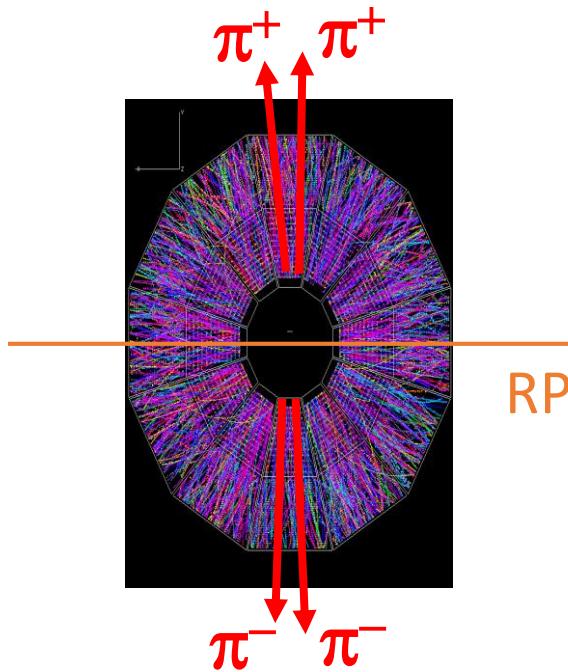
OUTLINE

1. Various methods to eliminate **flow-induced** backgrounds
2. Further background contamination from **nonflow** and
genuine 3-particle correlations
3. Hint of CME in Au+Au, and upper limit in isobar
4. Summary & outlook

THE $\Delta\gamma$ CORRELATOR

Voloshin, PRC 2004

Look for charge separation

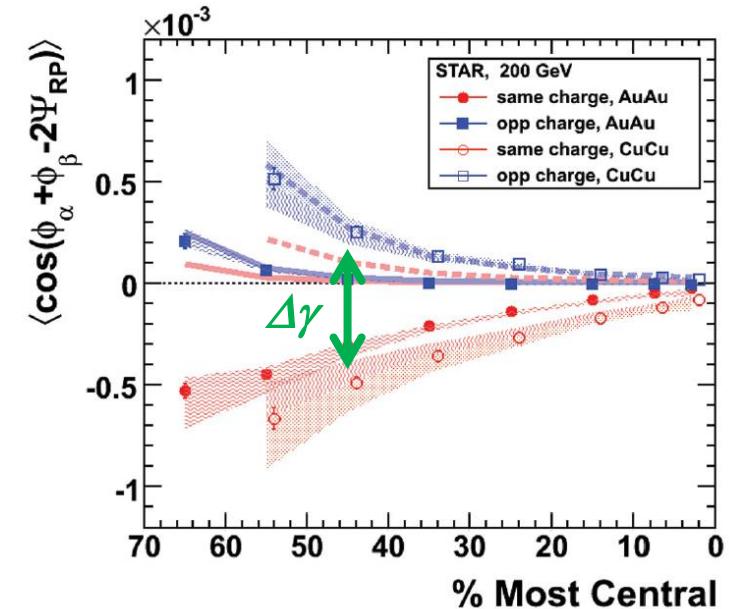


$$\gamma_{\alpha\beta} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{RP}) \rangle$$

$$\gamma_{+-,-+} > 0, \quad \gamma_{++,- -} < 0$$

$$\Delta\gamma = \gamma_{\text{opposite-sign}} - \gamma_{\text{same-sign}} > 0$$

STAR, PRL 2009, PRC 2010

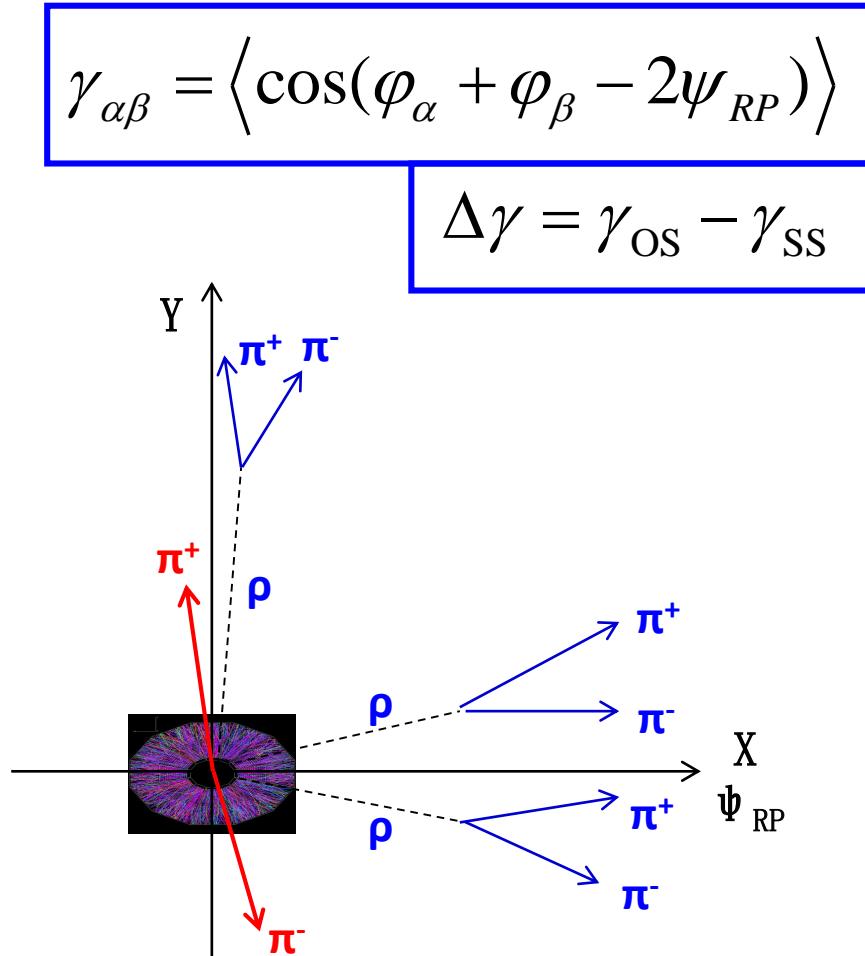


Significant signal

$$\Delta\gamma \sim 5 \times 10^{-4}$$

BACKGROUND IN $\Delta\gamma$ CORRELATOR

Voloshin 2004; FW 2009; Bzdak, Koch, Liao 2010; Pratt, Schlichting 2010; ...



$$dN_\pm / d\varphi \propto 1 + 2v_1 \cos \varphi^\pm + 2a_\pm \cdot \sin \varphi^\pm + 2v_2 \cos 2\varphi^\pm + \dots$$

$$\gamma_{\alpha\beta} = \left[\langle \cos(\varphi_\alpha - \psi_{RP}) \cos(\varphi_\beta - \psi_{RP}) \rangle - \langle \sin(\varphi_\alpha - \psi_{RP}) \sin(\varphi_\beta - \psi_{RP}) \rangle \right]$$

$$+ \left[\frac{N_{cluster}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{cluster}) \cos(2\varphi_{cluster} - 2\varphi_{RP}) \rangle \right]$$

$$= \left[\langle v_{1,\alpha} v_{1,\beta} \rangle - \langle a_\alpha a_\beta \rangle \right] + \frac{N_{cluster}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{cluster}) \rangle v_{2,cluster}$$

$$\Delta\gamma = 2 \langle a_1^2 \rangle + \frac{N_\rho}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_\rho) \rangle v_{2,\rho}$$

Flow-induced charge-dependent background:
nonflow coupled with flow

$$\Delta\gamma_{Bkg} \propto v_2 / N$$

SLIGHTLY MORE FORMALLY...

$$\Delta\gamma = 2\langle a_1^2 \rangle + \frac{N_\rho}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_\rho) \rangle v_{2,\rho}$$

$$\Delta C_3 = 2\langle a_1^2 \rangle v_{2,c \perp B} + \frac{N_{2p}}{N_\alpha N_\beta} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{2p}) \rangle v_{2,2p} v_{2,c} + \frac{N_{3p}}{N_\alpha N_\beta N_c} \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_c) \rangle$$

$$= 2\langle a_1^2 \rangle v_{2,c \perp B} + \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} v_{2,c} + \frac{C_{3p} N_{3p}}{2N^3}$$

$$\Delta\gamma = 2\langle a_1^2 \rangle \frac{v_{2,c \perp B}}{v_{2,c}^*} + \frac{C_{2p} N_{2p}}{N^2} \frac{v_{2,2p} v_{2,c}}{v_{2,c}^*} + \frac{C_{3p} N_{3p}}{2N^3 v_{2,c}^*}$$

CME 2p nonflow flow 3p nonflow

flow-induced bkgd

$$N \approx N_+ \approx N_-$$

$$C_{2p} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_{2p}) \rangle$$

$$C_{3p} = \langle \cos(\varphi_\alpha + \varphi_\beta - 2\varphi_c) \rangle_{3p}$$

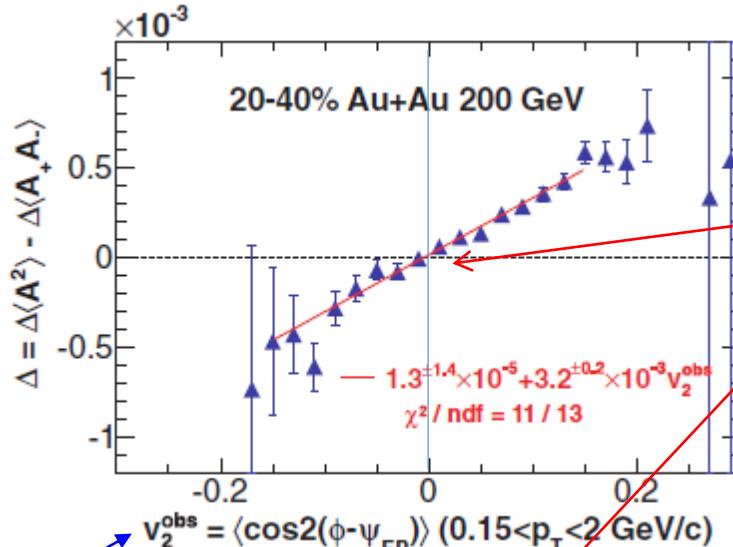
$v_{2,c \perp B}$: v_2 of c particle wrt direction $\perp B$

$v_{2,c}^*$: measured v_2 of c particle containing nonflow

STATISTICAL EVENT-SHAPE-ENGINEERING

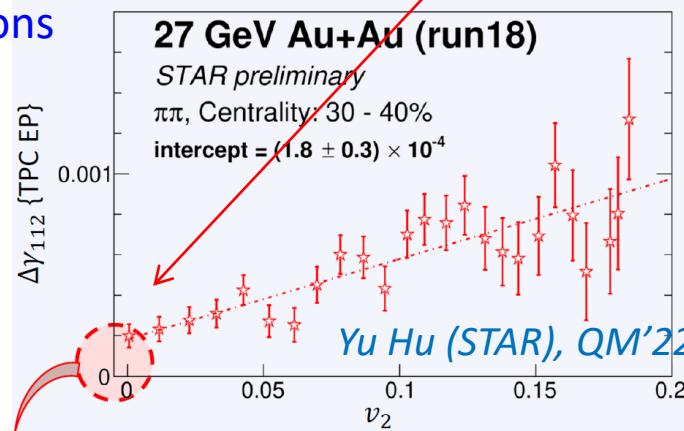
Here Δ is similar to $\cos(\alpha + \beta - 2\psi)$ correlator

STAR, PRC 89 (2014) 044908



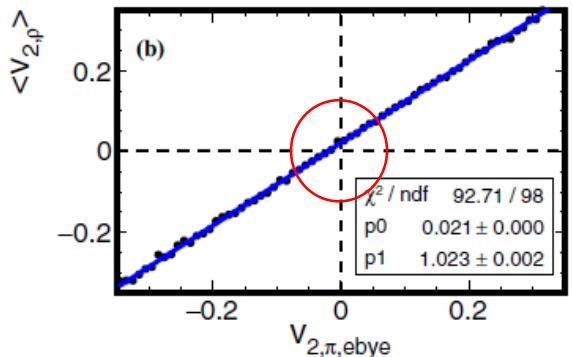
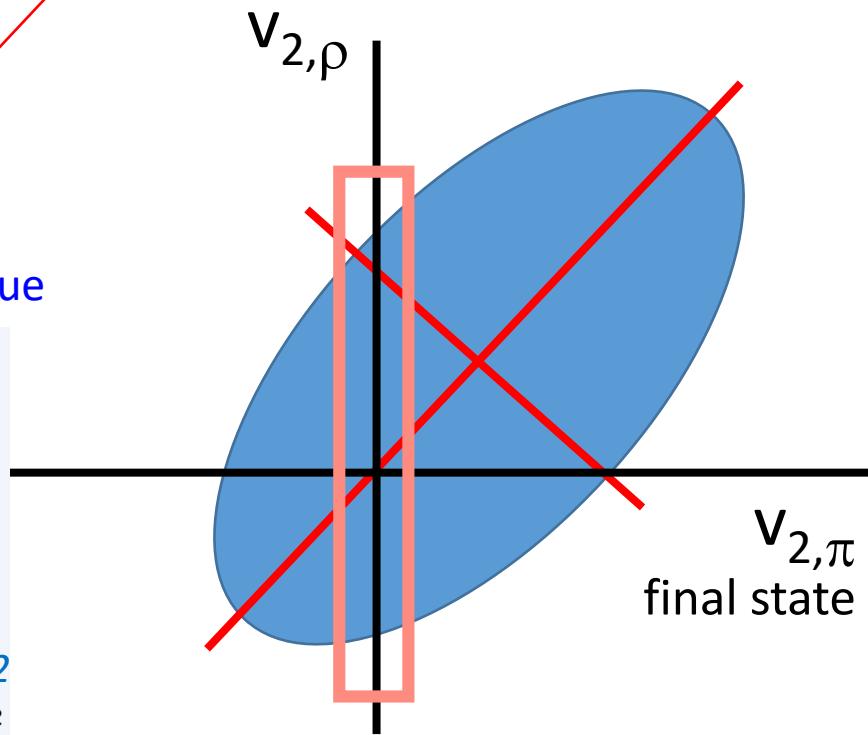
Primarily stat. fluctuations

Event-by-event v_2 technique



$$\Delta\gamma = 2 \left\langle a_1^2 \right\rangle \frac{v_{2,c \perp B}}{v_{2,c}^*} + \frac{C_{2p} N_{2p}}{N^2} \frac{v_{2,2p} v_{2,c}}{v_{2,c}^*} + \frac{C_{3p} N_{3p}}{2N^3 v_{2,c}^*}$$

Still has residual background,
because background $\sim v_{2,p}$ not $v_{2,\pi}$
FW, Jie Zhao, PRC 95 (2017) 051901(R)



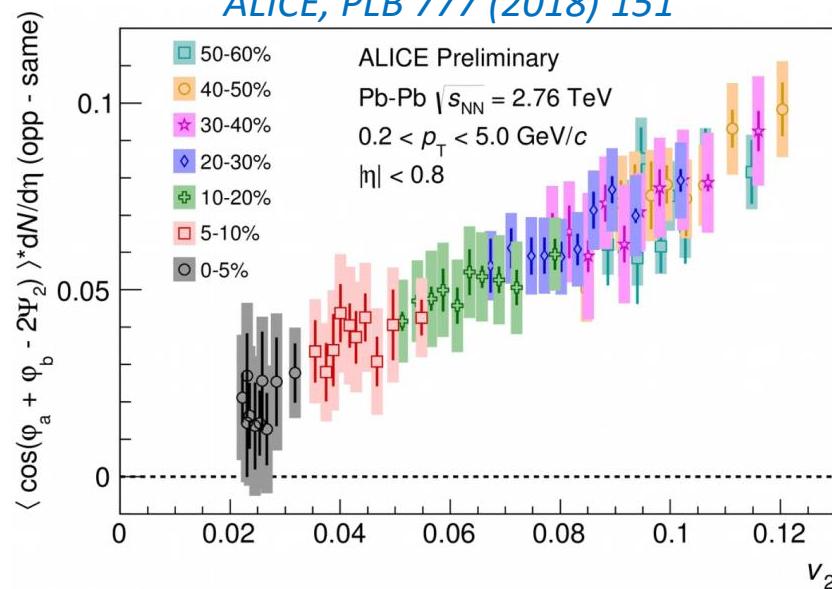
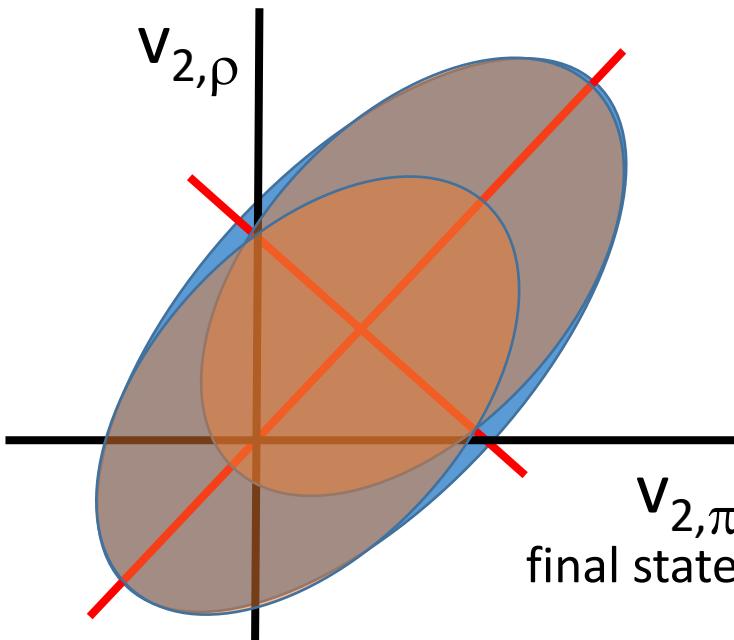
- Engineering on stat. fluctuations
- Background suppressed, but not totally eliminated
- LHC does not have this issue as v_2 selection in different phase space

DYNAMICAL EVENT-SHAPE-ENGINEERING

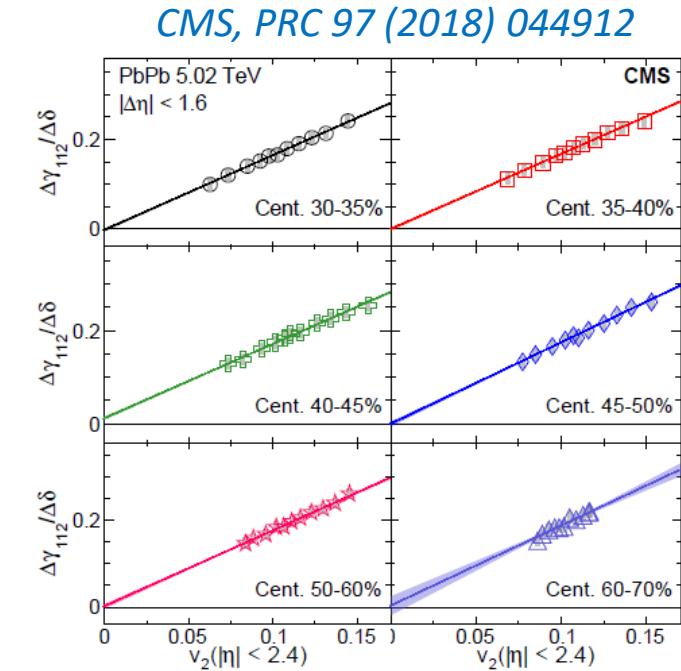
Schukraft, Timmins, Voloshin, PLB 719 (2013) 394

$$\Delta\gamma = 2\left\langle a_1^2 \right\rangle \frac{v_{2,c \perp B}}{v_{2,c}^*} + \frac{C_{2p} N_{2p}}{N^2} \frac{v_{2,2p} v_{2,c}}{v_{2,c}^*} + \frac{C_{3p} N_{3p}}{2N^3 v_{2,c}^*}$$

ESE displaced from
 α, β phase space



More sophisticated (model-dep.)
assumption of B dependence on v2
within a given centrality
Upper limit 26%



Assume CME does not depend
on v2 within a given centrality
Upper limit 7%

Promising way to extract possible CME signal. Will need to assess nonflow effects

THE R OBSERVABLE

Ajitanand, Lacey, et al., PRC 83 (2011) 011901
 Magdy, Lacey, et al., PRC 97 (2018) 061901(R)

$$\Delta S = \langle \sin(\varphi - \psi_2) \rangle_+ - \langle \sin(\varphi - \psi_2) \rangle_-$$

$$R(\Delta S) = \frac{N(\Delta S_{\text{real}})}{N(\Delta S_{\text{shuffled}})} \sqrt{\frac{N(\Delta S_{\text{real}}^\perp)}{N(\Delta S_{\text{shuffled}}^\perp)}}$$

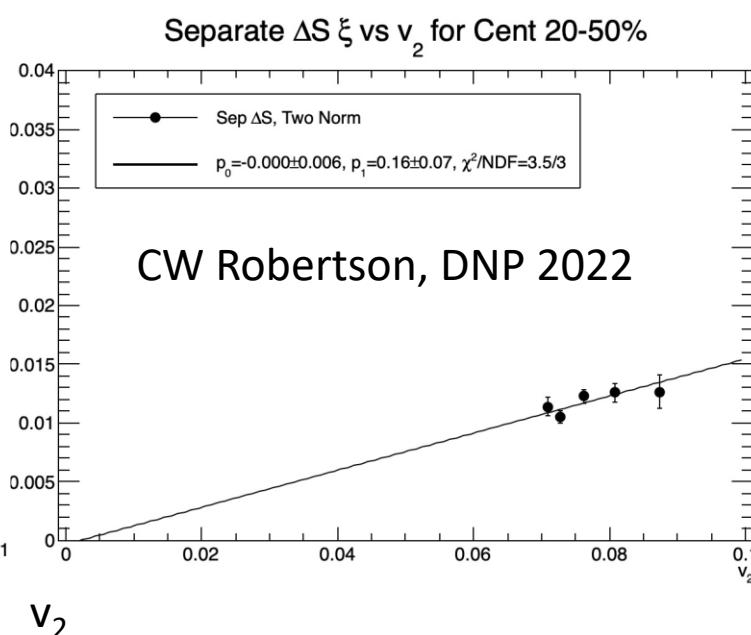
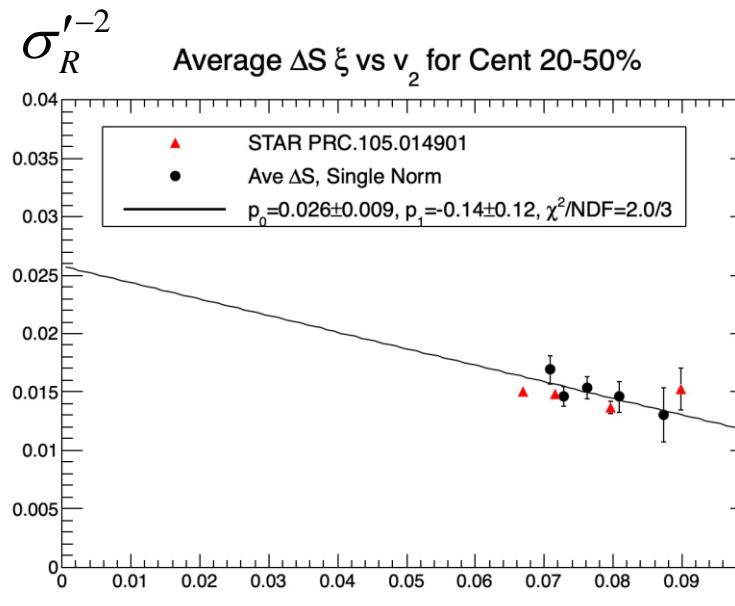
Width of $R(\Delta S)$ distribution reduces to variance $\sin^* \sin, \cos^* \cos \rightarrow$ equivalently the $\Delta\gamma$ variable

Normalized by shuffled width: $\sigma_R'^{-2} = N\Delta\gamma \propto v_2$

Choudhury et al., CPC 46 (2022) 014101

Earlier claims of CME; many issues identified.
 ESE studies indicate “v2-independence.”
Latest claim by Lacey et al., arXiv:2203.10029:
Scale by N, so isobar R compatible with CME.
 See Comment by FW, arXiv:2204.08450

Execution of R analysis: $\Delta S = (\Delta S_{\text{West-POI}}^{\text{East-EP}} + \Delta S_{\text{East-POI}}^{\text{West-EP}})/2$



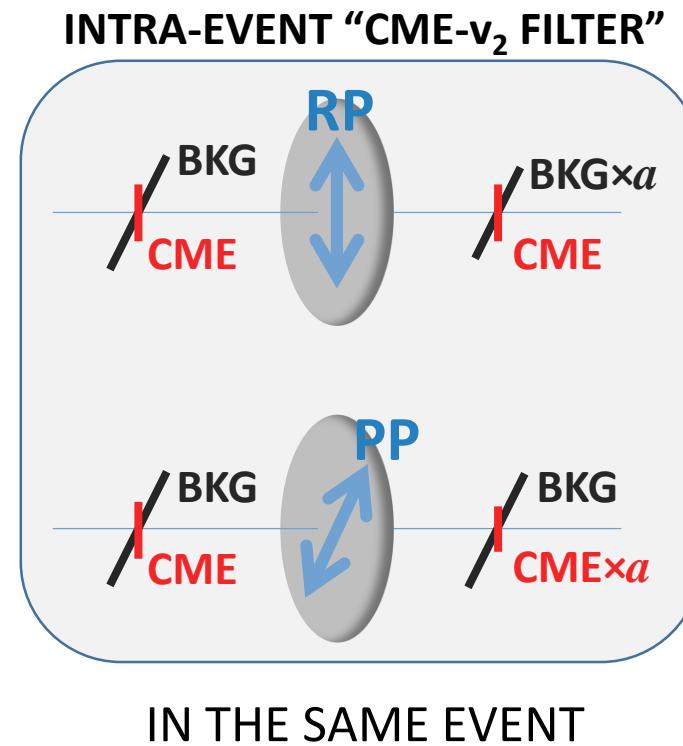
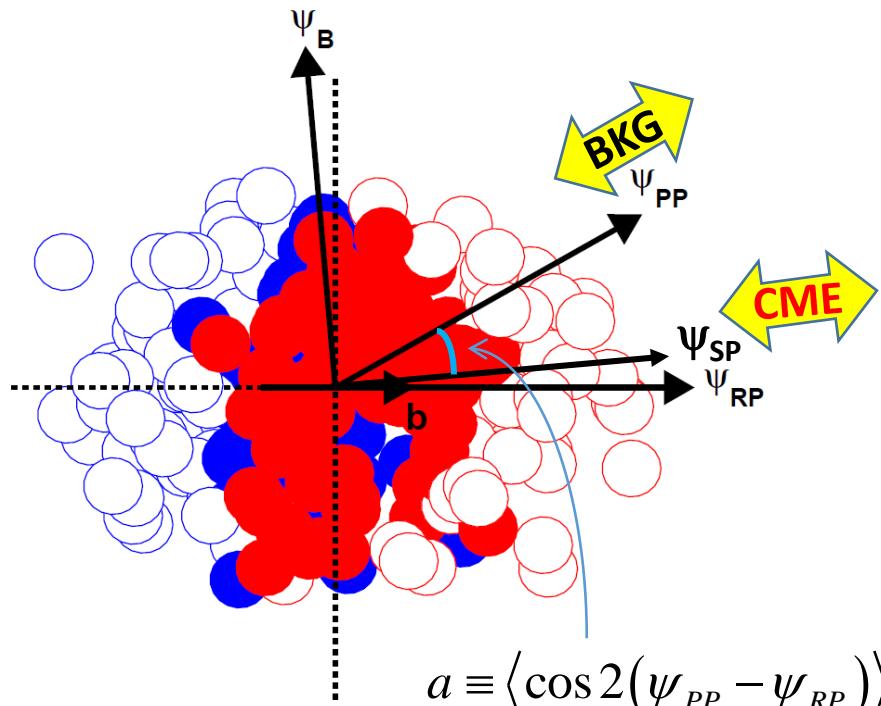
Issue identified: Reuse of data in average causes autocorrelation, yielding non-zero intercept. Feng et al, PRC 103 (2021) 034912

Positive intercept not present in analysis with separate ΔS

SPECTATOR & PARTICIPANT PLANES

H.-j. Xu, FW, et al., CPC 42 (2018) 084103, arXiv:1710.07265
 S.A. Voloshin, PRC 98 (2018) 054911, arXiv:1805.05300

$$\Delta\gamma = 2\langle a_1^2 \rangle \frac{v_{2,c \perp B}}{v_{2,c}^*} + \frac{C_{2p} N_{2p}}{N^2} \frac{v_{2,2p} v_{2,c}}{v_{2,c}^*} + \frac{C_{3p} N_{3p}}{2N^3 v_{2,c}^*}$$



$$\Delta\gamma_{\{SP\}} = \frac{\Delta\gamma_{CME\{PP\}}}{a} + a\Delta\gamma_{Bkg\{PP\}}$$

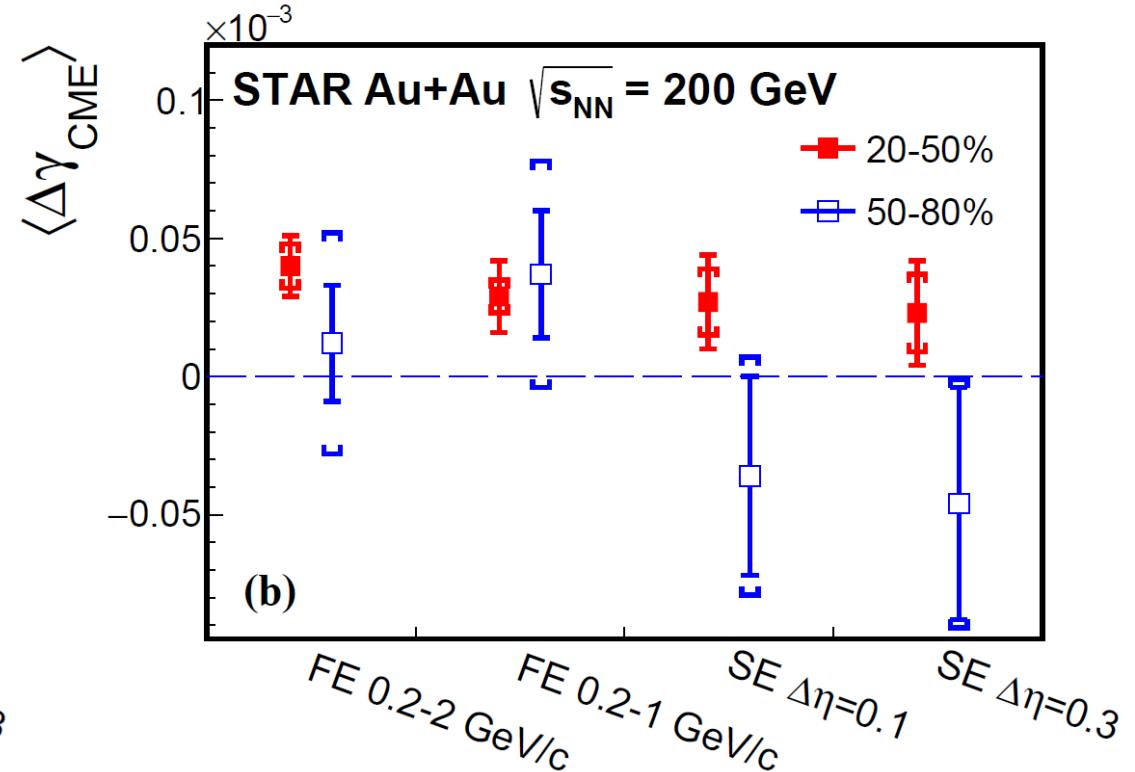
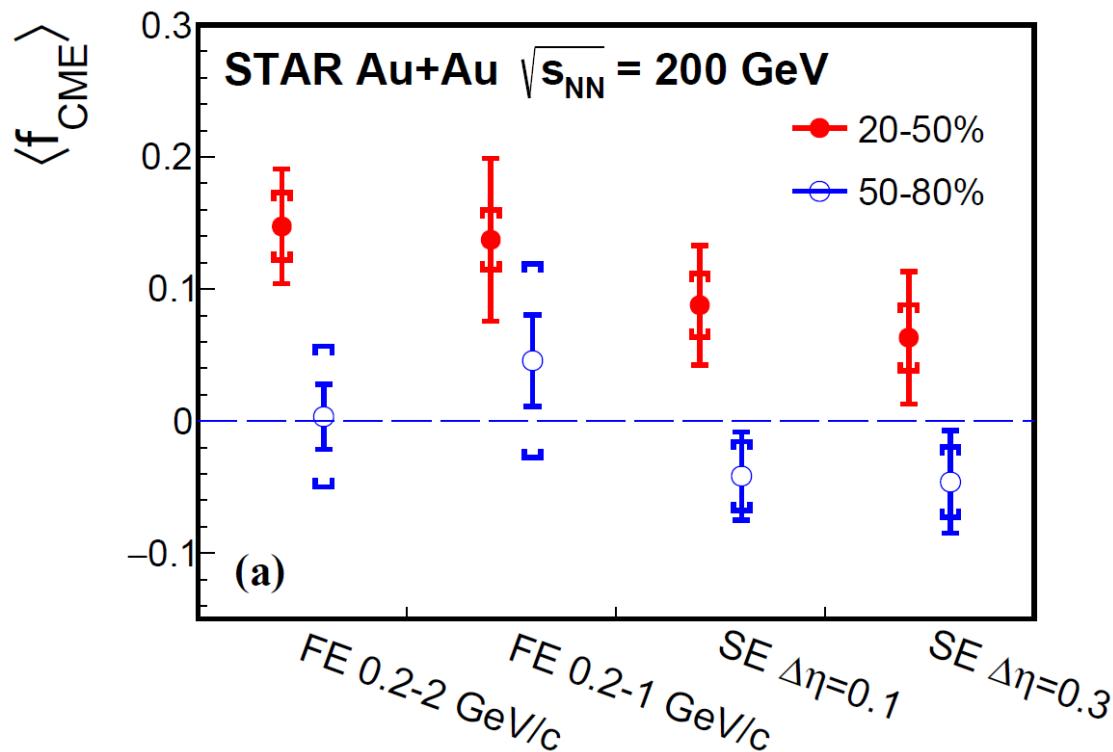
$$\Delta\gamma_{\{PP\}} = \Delta\gamma_{CME\{PP\}} + \Delta\gamma_{Bkg\{PP\}}$$

$$f_{CME} = \frac{\Delta\gamma_{CME\{PP\}}}{\Delta\gamma_{\{PP\}}} = \frac{A/a - 1}{1/a^2 - 1}$$

$$A = \Delta\gamma_{\{SP\}} / \Delta\gamma_{\{PP\}}, \quad a = v_2\{SP\} / v_2\{PP\}$$

Au+Au Collisions at 200 GeV (2.4B MB)

STAR, PRL 128 (2022) 092301, arXiv:2106.09243



- Peripheral 50-80% collisions: consistent-with-zero signal with relatively large errors
- Mid-central 20-50% collisions: indication of finite CME signal with $1-3\sigma$ significance
- How much is there remaining nonflow contamination?

REMAINING NONFLOW EFFECTS

Feng, FW, et al., PRC 105 (2022) 024913, arXiv:2106.15595

$$f_{\text{CME}} = \frac{\Delta\gamma_{\text{CME}}\{\text{PP}\}}{\Delta\gamma\{\text{PP}\}} = \frac{A/a - 1}{1/a^2 - 1}$$

$$\frac{A}{a} = \frac{\Delta\gamma\{\text{SP}\}/\Delta\gamma\{\text{PP}\}}{v_2\{\text{SP}\}/v_2^*\{\text{PP}\}} = \frac{\Delta C_3\{\text{SP}\}}{v_2^2\{\text{SP}\}} \cdot \frac{v_2^{*2}\{\text{PP}\}}{C_3\{\text{PP}\}} = \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3/\epsilon_2}{N v_2^2\{\text{PP}\}}}$$

$$\Delta C_3 = 2 \langle a_1^2 \rangle v_{2,c \perp B} + \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} v_{2,c} + \frac{C_{3p} N_{3p}}{2N^3}$$

Nonflow in $\Delta\gamma \rightarrow$ negative f_{CME}

$$\Delta C_3^{\text{Bkg}}\{\text{SP}\} = \frac{C_{2p} N_{2p}}{N^2} v_{2,2p}\{\text{SP}\} v_{2,c}\{\text{SP}\}$$

$$\Delta C_3^{\text{Bkg}}\{\text{EP}\} = \frac{C_{2p} N_{2p}}{N^2} v_{2,2p}\{\text{EP}\} v_{2,c}\{\text{EP}\} + \frac{C_{3p} N_{3p}}{2N^3}$$

$$\begin{aligned} \epsilon_2 &= \frac{C_{2p} N_{2p}}{N} \cdot \frac{v_{2,2p}}{v_2} \\ &\approx N \Delta\gamma / v_2 \end{aligned}$$

$$\epsilon_3 = \frac{C_{3p} N_{3p}}{2N}$$

Nonflow in $v_2 \rightarrow$ positive f_{CME}

$$v_2^*\{\text{PP}\} = \sqrt{v_2^2\{\text{PP}\} + v_{2,\text{nf}}^2}$$

$$\epsilon_{\text{nf}} \equiv v_{2,\text{nf}}^2 / v_2^2$$

$$f_{\text{CME}}^* \approx \left(\epsilon_{\text{nf}} - \frac{\epsilon_3/\epsilon_2}{N v_2^2\{\text{EP}\}} \right) \Bigg/ \left(\frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$

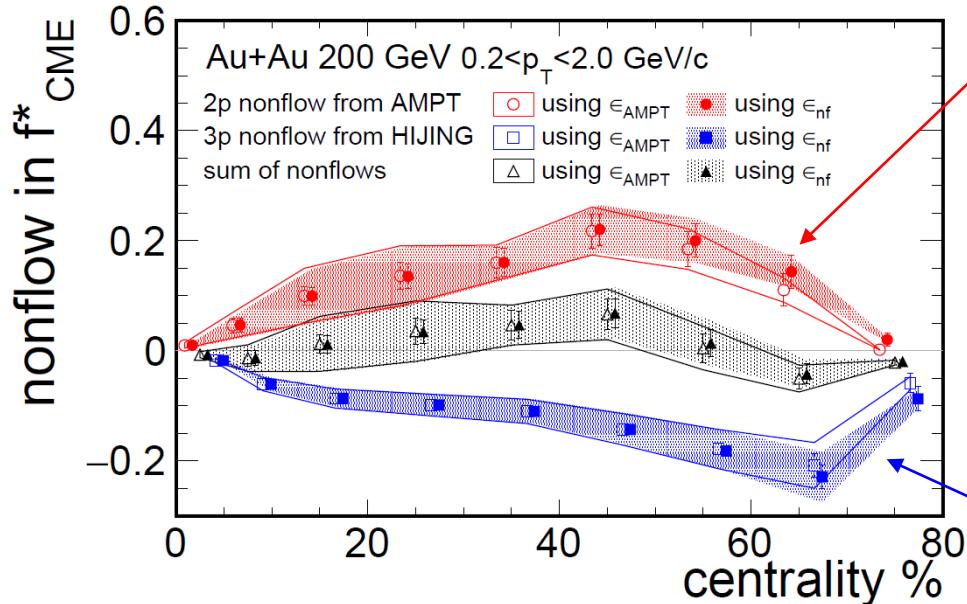
$$f_{\text{CME}}^* = \left(\frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3/\epsilon_2}{N v_2^2\{\text{EP}\}}} - 1 \right) \Bigg/ \left(\frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right) = \left(\frac{1 + \epsilon_{\text{nf}}}{1 + \frac{(1 + \epsilon_{\text{nf}})\epsilon_3/\epsilon_2}{N v_2^{*2}\{\text{EP}\}}} - 1 \right) \Bigg/ \left(\frac{1}{a^{*2}} - 1 \right)$$

MODEL ESTIMATES OF NONFLOW

Feng, FW, et al., PRC 105 (2022) 024913, arXiv:2106.15595

$$\frac{A}{a} = \frac{\Delta\gamma\{\text{SP}\}/\Delta\gamma\{\text{PP}\}}{v_2\{\text{SP}\}/v_2^*\{\text{PP}\}} = \frac{C_3\{\text{SP}\}}{v_2^2\{\text{SP}\}} \cdot \frac{v_2^{*2}\{\text{PP}\}}{C_3\{\text{PP}\}} = \frac{1 + \epsilon_{\text{nf}}}{1 + \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{PP}\}}}$$

$$f_{\text{CME}}^* \approx \left(\epsilon_{\text{nf}} - \frac{\epsilon_3 / \epsilon_2}{Nv_2^2\{\text{EP}\}} \right) \Big/ \left(\frac{1 + \epsilon_{\text{nf}}}{a^2} - 1 \right)$$

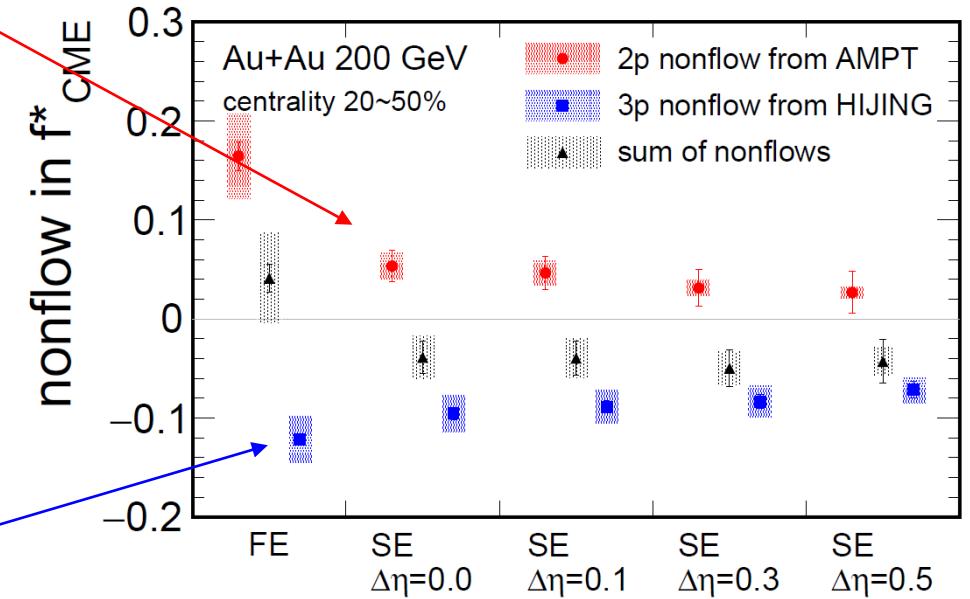


AMPT 2p estimates

v_2 nonflow data analysis ongoing

Limited-scope data comparison ongoing

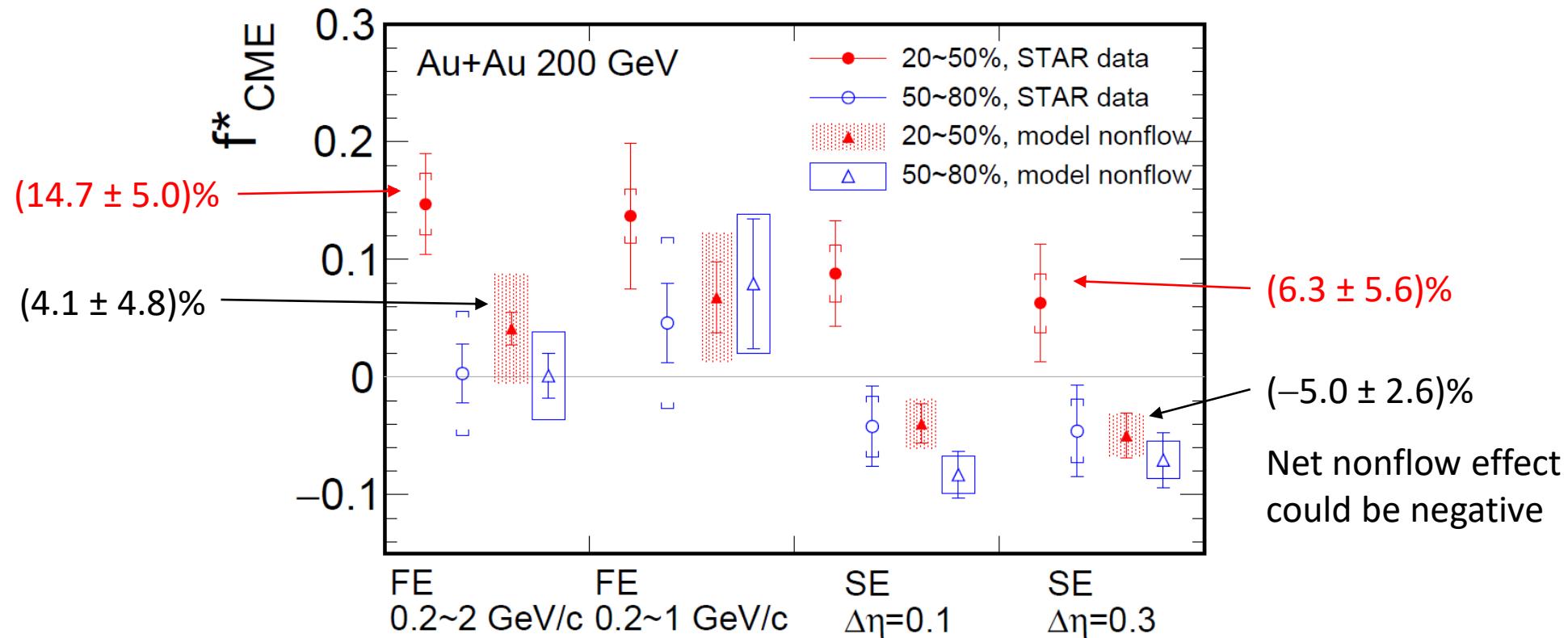
HIJING 3p estimates



IMPLICATIONS TO Au+Au DATA

STAR, PRL 128 (2022) 092301, arXiv:2106.09243

Feng, FW, et al., PRC 105 (2022) 024913, arXiv:2106.15595



There may indeed be hint of CME in the Au+Au data

BACK TO ISOBAR DATA

$$\Delta C_3 = 2 \left\langle a_1^2 \right\rangle v_{2,c \perp B} + \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} v_{2,c} + \frac{C_{3p} N_{3p}}{2N^3}$$

$$\Delta C_3^{\text{Ru}} \{ \text{EP} \} = \left(\frac{C_{2p} N_{2p}}{N^2} v_{2,2p} \{ \text{EP} \} v_{2,c} \{ \text{EP} \} \right)^{\text{Ru}} + \left(\frac{C_{3p} N_{3p}}{2N^3} \right)^{\text{Ru}}$$

$$\Delta C_3^{\text{Zr}} \{ \text{EP} \} = \left(\frac{C_{2p} N_{2p}}{N^2} v_{2,2p} \{ \text{EP} \} v_{2,c} \{ \text{EP} \} \right)^{\text{Zr}} + \left(\frac{C_{3p} N_{3p}}{2N^3} \right)^{\text{Zr}}$$

$\Delta C_3^{\text{Bkg}} \{ \text{SP} \} = \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} \{ \text{SP} \} v_{2,c} \{ \text{SP} \}$ $\Delta C_3^{\text{Bkg}} \{ \text{EP} \} = \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} \{ \text{EP} \} v_{2,c} \{ \text{EP} \} + \frac{C_{3p} N_{3p}}{2N^3}$

$$\frac{\left(N \Delta \gamma / v_2^* \right)^{\text{Ru}}}{\left(N \Delta \gamma / v_2^* \right)^{\text{Zr}}} = \frac{\left(N C_3 / v_2^{*2} \right)^{\text{Ru}}}{\left(N C_3 / v_2^{*2} \right)^{\text{Zr}}} = \frac{\left(C_{2p} \frac{N_{2p}}{N} \frac{v_{2,2p}}{v_2} \right)^{\text{Ru}}}{\left(C_{2p} \frac{N_{2p}}{N} \frac{v_{2,2p}}{v_2} \right)^{\text{Zr}}} \cdot \frac{\left(1 + \epsilon_{\text{nf}} \right)^{\text{Zr}}}{\left(1 + \epsilon_{\text{nf}} \right)^{\text{Ru}}} \cdot \frac{\left(1 + \frac{\epsilon_3 / \epsilon_2}{N v_2^2} \right)^{\text{Ru}}}{\left(1 + \frac{\epsilon_3 / \epsilon_2}{N v_2^2} \right)^{\text{Zr}}}$$

- Depending on the relative Ru+Ru/Zr+Zr difference of various nonflow effects, the baseline can be above, equal, or below unity
- Final isobar conclusion requires detailed nonflow studies

REMAINING NONFLOW EFFECTS

FENG, Yicheng (STAR): QM'2022, SQM'2022

Feng, FW, et al., PRC 105 (2022) 024913, arXiv:2106.15595

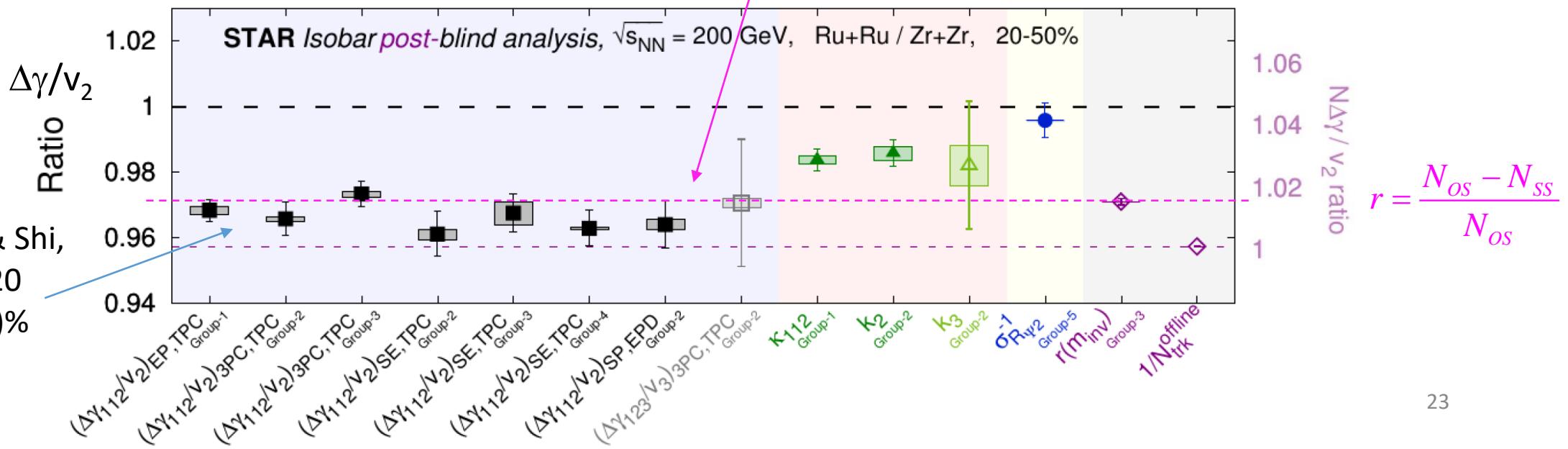
$$C_3 = \frac{C_{2p} N_{2p}}{N^2} v_{2,2p} v_2 + \frac{C_{3p} N_{3p}}{2N^3}; \quad C_{2p} \equiv \langle \cos(\alpha + \beta - 2\phi_{2p}) \rangle; \quad C_{3p} \equiv \langle \cos(\alpha + \beta - 2c) \rangle_{3p}$$

$$\begin{aligned} \varepsilon_2 &\equiv \frac{C_{2p} N_{2p}}{N} \cdot \frac{v_{2,2p}}{v_2} \\ \varepsilon_3 &\equiv \frac{C_{3p} N_{3p}}{2N} \end{aligned}$$

$$\begin{aligned} v_2^{*2} &= v_2^2 + v_{2,nf}^2 \\ \varepsilon_{nf} &\equiv v_{2,nf}^2 / v_2^2 \end{aligned}$$

$$\begin{aligned} N &\approx N_+ \approx N_- \\ \Delta X &\equiv X^{\text{Ru}} - X^{\text{Zr}} \end{aligned}$$

$$\frac{(N\Delta\gamma/v_2^*)^{\text{Ru}}}{(N\Delta\gamma/v_2^*)^{\text{Zr}}} \equiv \frac{(NC_3/v_2^{*2})^{\text{Ru}}}{(NC_3/v_2^{*2})^{\text{Zr}}} \approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} \cdot \frac{(1+\varepsilon_{nf})^{\text{Zr}}}{(1+\varepsilon_{nf})^{\text{Ru}}} \cdot \frac{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Ru}}}{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Zr}}} \approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} - \frac{\Delta\varepsilon_{nf}}{1+\varepsilon_{nf}} + \frac{\frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}}{1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}} \left(\frac{\Delta\varepsilon_3}{\varepsilon_3} - \frac{\Delta\varepsilon_2}{\varepsilon_2} - \frac{\Delta N}{N} - \frac{\Delta v_2^2}{v_2^2} \right)$$



REMAINING NONFLOW EFFECTS

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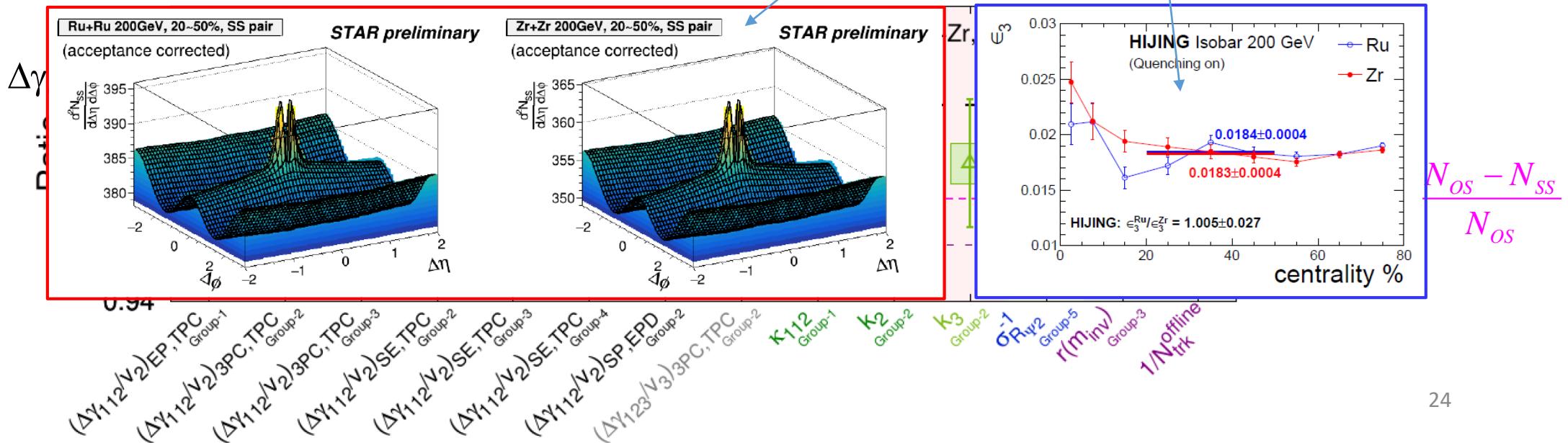
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$$\frac{\left(N\Delta\gamma/v_2^*\right)^{\text{Ru}}}{\left(N\Delta\gamma/v_2^*\right)^{\text{Zr}}} \equiv \frac{\left(NC_3/v_2^{*2}\right)^{\text{Ru}}}{\left(NC_3/v_2^{*2}\right)^{\text{Zr}}} \approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} \cdot \frac{(1+\varepsilon_{nf})^{\text{Zr}}}{(1+\varepsilon_{nf})^{\text{Ru}}} \cdot \frac{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Ru}}}{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Zr}}} \approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} - \frac{\Delta\varepsilon_{nf}}{1+\varepsilon_{nf}} + \frac{\frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}}{1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}} \left(\frac{\Delta\varepsilon_3}{\varepsilon_3} - \frac{\Delta\varepsilon_2}{\varepsilon_2} - \frac{\Delta N}{N} - \frac{\Delta v_2^2}{v_2^2} \right)$$



REMAINING NONFLOW EFFECTS

FENG, Yicheng (STAR): QM'2022, SQM'2022

Feng, FW, et al., PRC 105 (2022) 024913, arXiv:2106.15595

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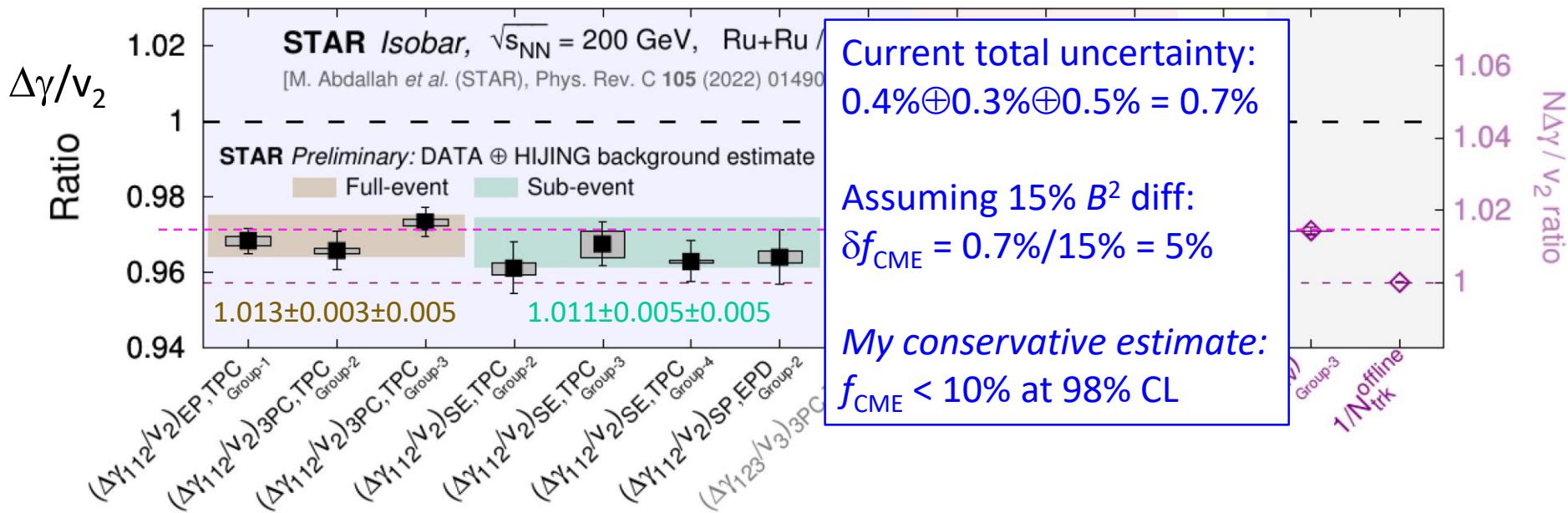
$$\begin{aligned}v_2^{*2} &= v_2^2 + v_{2,nf}^2 \\ \varepsilon_{nf} &\equiv v_{2,nf}^2 / v_2^2\end{aligned}$$

$$\begin{aligned}N &\approx N_+ \approx N_- \\ \Delta X &\equiv X^{\text{Ru}} - X^{\text{Zr}}\end{aligned}$$

$$\frac{(N\Delta\gamma/v_2^*)^{\text{Ru}}}{(N\Delta\gamma/v_2^*)^{\text{Zr}}} \equiv \frac{(NC_3/v_2^{*2})^{\text{Ru}}}{(NC_3/v_2^{*2})^{\text{Zr}}} \approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} \cdot \frac{(1+\varepsilon_{nf})^{\text{Zr}}}{(1+\varepsilon_{nf})^{\text{Ru}}} \cdot \frac{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Ru}}}{\left(1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}\right)^{\text{Zr}}}$$

$$\approx \frac{\varepsilon_2^{\text{Ru}}}{\varepsilon_2^{\text{Zr}}} - \frac{\Delta\varepsilon_{nf}}{1+\varepsilon_{nf}} + \frac{\frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}}{1 + \frac{\varepsilon_3/\varepsilon_2}{Nv_2^2}} \left(\frac{\Delta\varepsilon_3}{\varepsilon_3} - \frac{\Delta\varepsilon_2}{\varepsilon_2} - \frac{\Delta N}{N} - \frac{\Delta v_2^2}{v_2^2} \right)$$

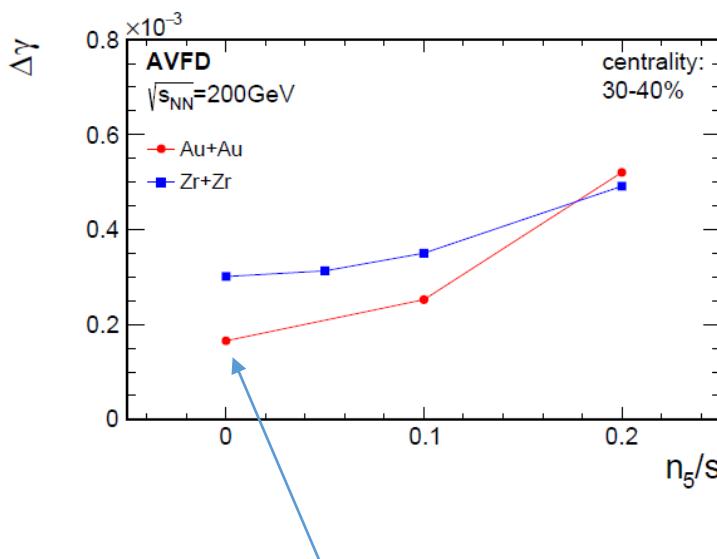
Good cancelation of v2 nonflow and 3p correlation



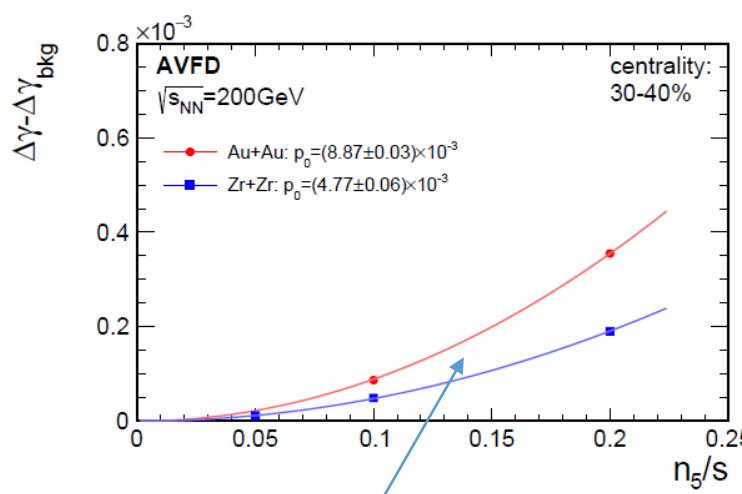
ISOBAR SIGNAL MAY BE SMALL

Feng, Lin, Zhao & FW, PLB 820 (2021) 136549, arXiv:2103.10378

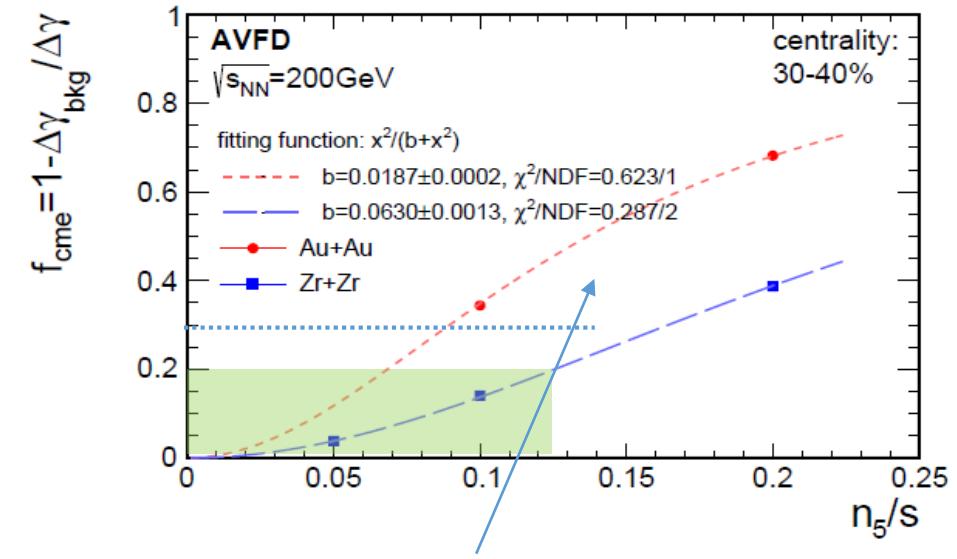
AVFD simulations in preparation for isobar blind analysis



Background $\propto 1/N$
isobar/AuAu ~ 2



Mag. field $B \sim A/A^{2/3} \sim A^{1/3}$
 $\Delta\gamma_{\text{CME}} \sim B^2 \sim A^{2/3}$
Signal: AuAu/isobar ~ 1.5



Could be x3 reduction in f_{CME} at the same n_5/s
If AuAu $f_{\text{CME}}=10\%$, then isobar 3% (1 σ effect)
 $R_{\text{U}/\text{Zr}} = 1 + 15\% * 3\% = 1.005 (\pm 0.004)$

Caveats: Axial charge densities and sphaleron transition probabilities could be different between Au+Au and isobar, e.g. AVFD-glasma μ_5/s : isobar/AuAu ~ 1.5

SUMMARY, NEAR FUTURE, OUTLOOK

- CME is a very important physics. Large efforts to eliminate backgrounds.
 - Flow-induced background
 - **Nonflow contamination in v_2 measurements**
 - **Genuine 3-particle correlation background**
 - Invariant mass dependence
 - Event-shape engineering
 - Isobar collisions
 - **Spectator/participant plane comparison**
- Isobar signal **consistent with zero** (not contradictory to Au+Au);
Stringent upper limit on the way.
- Au+Au data (2.4 B MB events) **hint at finite CME signal** with small significance;
Rigorous nonflow data analysis on the way. Expect 20B Au+Au from 2023+25 runs.