Luminosity Spectrometer

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Basic idea of ZEUS spectrometer



New work-in-progress design of ePIC spectrometer







DD4HEP Design

Exit Window

- Taken to be Aluminum, 1 cm thick
- Thickness and composition needs to be accurately known in order to account for the lost photons

Collimator

- Steel block to remove stray particles that may damage other components further downstream.
- Opening size defines phase space of measured photons: taken to be $5\sigma \sim \pm 2.4$ cm.



DD4HEP Design

Sweeper Magnet

- Taken to be the same as the ZEUS design
- Dimensions extracted from a CAD drawing.
 - 10 cm opening ($\sim \pm 7\sigma$ photon beam).

 - 0.5 T horizontal B field (electrons go up and down) $p_T = 0.3 BdL = 0.3 (\frac{\text{GeV}}{\text{T m}}) * 0.5 \text{T} * 0.78 \text{m} = 0.117 \text{GeV}$
- Still need to ensure fringe field < 10 Gauss @ electron beam.

He gas chamber

- $\pm 5\sigma$ dimensions
- Minimizes multiple scattering and unwanted conversions



CAD drawing obtained from Yulia Furletova

Implemented in DD4HEP by Justin Chan

DD4HEP Design

Converter

- Taken to be Aluminum, 1 mm thin
- Thickness and composition needs to be accurately known in order to account for the conversions

Spectrometer Magnet

- Taken to be the same as the Sweeper Magnet: 0.5 T horizontal B field.
- Placed just after the converter.
- 12 cm opening also provides ~±6sigma clearance for direct-Y beam.

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$$p_T = 0.3 BdL = 0.3 (\frac{\text{GeV}}{\text{T m}}) * 0.5 \text{T} * 0.78 \text{m} = 0.117 \text{GeV}$$

DAHEP Design

1 Tracker

Station directly in front of CAL



- 3 stations with a top & bottom plane each.
- $\pm 5\sigma$ gap between top and bottom for direct photon passage.
- 20 cm x 20 cm transverse dimensions
- 0.3 mm SiliconOxide (sensor) + 0.14 mm Copper (ASIC+cooling approximation): ~1% X₀

Calorimeter

- Taken to be PbWO4 with same module dimensions as scattered electron CAL
 - 2 cm x 2 cm x 20 cm.
- ±5σ gap between top and bottom for direct photon passage.
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CAL acceptance



- Shape of acceptance governed by geometry of spectrometer arm
- Current design with nominal B=0.5 T favors high end of spectrum
- One can "slide" the acceptance to the left by decreasing B

eA (very high Bremstrahlung)- Use a large B

ep - Use a smaller B

Expected Rates of electrons at spectrometer CALs

Bethe-Heitler formula for unpolarized ep Bremstrahlung



$$\frac{d\sigma}{dE_{\gamma}} = 4\alpha Z^2 r_e^2 \frac{E'_e}{E_{\gamma} E_e} \left(\frac{E_e}{E'_e} + \frac{E'_e}{E_e} - \frac{2}{3}\right) \left(\ln\frac{4E_p E_e E'_e}{m_p m_e E_{\gamma}} - \frac{1}{2}\right)$$

- Bremstrahlung σ is much larger for eAu than ep, but the bunch luminosity will be lower for eAu.
- These rates depend also on the design (acceptance) of the spectrometer CALs as well as the converter thickness (1 mm Al).
- Pileup greatly suppressed with low conversion rate!

See Bill Schmidke's talk for studies assuming 10 mm converter

Calculation of Polarized Bremstrahlung



- Bethe & Heitler first derived the cross section for bremstrahlung in 1934, <u>but for unpolarized projectile</u> <u>and targets.</u>
- Steps to calculate the polarized cross section outlined in Lifshitz QED textbook (section 93).

Before calculating the polarized cross section, a re-derivation of the unpolarized double-differential cross section has been done:

1) Write out Feynman amplitudes, square them, and work through the Dirac algebra.

2) Integrate over scattered electron angles and photon's azimuthal angle.

3) Take ultra-relativistic and small-angle limit (p ~ ϵ , θ <~ m/ ϵ), δ = θ ϵ /m.

Lifshitz Eq 93.16

$$\frac{d^2\sigma}{d\omega d\delta} = 8Z^2 \alpha r_e^2 \frac{1}{\omega} \frac{\varepsilon'}{\varepsilon} \frac{\delta}{(1+\delta^2)^2} \left(\begin{bmatrix} \frac{\varepsilon}{\varepsilon'} + \frac{\varepsilon'}{\varepsilon} - \frac{4\delta^2}{(1+\delta^2)^2} \end{bmatrix} \ln \frac{2\varepsilon\varepsilon'}{m_e\omega} - \frac{1}{2} \begin{bmatrix} \frac{\varepsilon}{\varepsilon'} + \frac{\varepsilon'}{\varepsilon} + 2 - \frac{16\delta^2}{(1+\delta^2)^2} \end{bmatrix} \right)$$
Unpolarized spinor simplification
$$\frac{1}{2} \sum_{spins} u(p)\bar{u}(p) = \frac{1}{2}(\gamma p + m) \qquad u(p)\bar{u}(p) \rightarrow \frac{1}{2}(\gamma p + m)(1-\gamma^5(\gamma s)) \qquad \begin{array}{c} \text{Contains} \\ \text{beam} \\ \text{polarization} \\ 11 \end{array}$$

Summary

- New luminosity spectrometer design implemented in DD4HEP.
- With new design (sweeper magnet + thin converter), pileup is greatly suppressed.
- Preliminary estimates for energy resolutions from trackers and CALs have been given.
 - <u>Synergy with taggers</u>: special low-luminosity runs are planned to cross calibrate spectrometer CAL and taggers. Full detection of $e+p \rightarrow e+p+\gamma$

Next steps

- Check the fringe-field strength of Lumi dipole magnets to confirm < 10 Gauss @ beams.
- Optimize placements/sizes of trackers and CAL.
- Further investigate feasibility & benefit of an extended He² gas chamber.
- Implement a clusterizer to more appropriately extract hit centers & deal with residual pileup.

Backup

Tracker Energy reconstruction

Slope of track obtained from the least-squares linear regression formula

$$\tan \theta = \frac{N \sum z_i y_i - \sum z_i \sum y_i}{N \sum z_i^2 - (\sum z_i)^2} \qquad E_{e\pm} = \frac{p_T}{\sin \theta} = \frac{0.3 \int B_x \, dz}{\sin \theta} \qquad E_{\gamma} = E_{e-} + E_{e+}$$



Generated $E_{\chi} = 10 \text{ GeV}$

More tracking layers leads to more multiple scattering (of course). •

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CAL Energy reconstruction

Crystal Ball fit used: Gaussian + power law tail

 $E_{gen} = 10 \text{ GeV}$



No observable difference in CAL energy resolution with these configurations.

Tracker relative Energy reconstruction



Bethe-Heitler

$$\frac{d\sigma}{dE_{\gamma}} = 4\alpha Z^2 r_e^2 \frac{E'_e}{E_{\gamma} E_e} \left(\frac{E_e}{E'_e} + \frac{E'_e}{E_e} - \frac{2}{3}\right) \left(\ln\frac{4E_p E_e E'_e}{m_p m_e E_{\gamma}} - \frac{1}{2}\right)$$

Bethe-Heitler

