Processes with light ions

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in collaboration w C. Weiss JLab LDRD on spectator tagging (non-exhaustive)

Supported by

lons: physics objectives









Neutron structure

- flavor decomposition of quark PDFs/GPDs/TMDs
- flavor structure of the nucleon sea
- singlet vs non-singlet QCD evolution, leading/higher-twist effects

Nuclear interactions in QCD

- medium modification of quark/gluon structure
- QCD origin of short-range nuclear force
- non-nucleonic components
- nuclear gluons

Imaging nuclear bound states

- imaging of quark-gluon degrees of freedom in nuclei through GPDs
- clustering in nuclei
 → talk Hyde

Coherence and saturation

- interaction of high-energy probe with coherent quark-gluon fields
 - \rightarrow talks Kovchegov, Schenke, Venugopalan

Nuclear structure: A tale with two sides

- Low-energy nuclear structure enters as input in a whole range of calculations for EIC
 - But high-energy nuclear scattering
 - \rightarrow light-front quantization

Novel studies of nuclear structure at EIC

- Short-range correlations
- Non-nucleonic components
- Breakup, Final-state interactions

EIC: unique opportunities

- Detection of nuclear breakup (spectator tagging) / coherent scattering
 - Enabled by the EIC far-forward detectors / secondary focus
 - Additional **control** over initial configuration
 - ↔ averaging over all configs in inclusive scattering
 - → on-shell extrapolation [Sargsian, Strikman '05; WC, Weiss 19+]
 - \rightarrow differential study of medium modifications
 - \rightarrow needs modelling of **final-state interactions**
- Light ions
 - \rightarrow polarization; different spins d, ³H(e), ⁴He
 - \rightarrow precision predictions with ab initio wave functions
- ► Lots of potential for novel measurements ↔ more challenging in fixed target
- But needs much more theory input!



High-energy scattering with nuclei [Frankfurt,Strikman 80s+]

► Interplay of two scales: high-energy scattering and low-energy nuclear structure. Virtual photon probes nucleus at fixed lightcone time $x^+ = x^0 + x^3$



- Scales can be separated using methods of light-front quantization and QCD factorization
- ► Tools for high-energy scattering known from *ep*
- Nuclear input: light-front momentum densities, spectral functions, overlaps with specific final states in breakup/tagging reactions
 - framework known for deuteron, can be extended to few-body A > 2 but challenges
 - still low-energy nuclear physics, just formulated differently
- Structure functions (IA) $F_A \propto F_N(\tilde{x}, Q^2) S(\alpha_p, p_{pT})$

Light-front nuclear structure

- LF gives rise to
 - Relativistic spin effects
 - Dynamical variables in LF boosted rest frame of on-shell A-nucleon state
- ▶ We currently match to the NR (ab initio) wf for light nuclei
 - Ab initio / EFT based methods could be extended to light-front quantization
 - Mean-field formulation for medium/heavy nuclei exist [Blunden, Burkardt, Miller 99]
- ► For A > 2 Poincaré invariance and cluster separability imposes non-trivial constraints [Sokolov 70s, Lev et al 90s]
 - Formalism is known [Sokolov packing operators], but technical
 - Trivial for deuteron
 - Applications to ³He [Rinaldi, Pace, Salme, Scopetta et al]

Opportunity to involve the low-energy community in EIC physics

Reaction Frameworks

Tagged spectator reactions

- focused on DIS on deuteron
- more complicated reactions (SIDIS, incoherent exclusive) possible
- A>2, but FSI a lot more complicated
- ▶ (untagged) (SI)DIS on ³He [Rinaldi,Pace,Salme,Scopetta et al]
- Coherent/incoherent DVCS on ⁴He
 [Guzey et al; Liuti, Taneja; Perugia/Orsay: Fucini, Scopetta et al.;...]
- ► Polarized deuteron yields additional tensor polarized structure functions/observables → spin-orbit, FSI

Polarized FSI

On-shell extrapolation of F_{2N}

$$F_{2d} = [2(2\pi)^3] S_d(\alpha_p, p_{pT})[\text{unpol}] F_{2n}(\tilde{x}, Q^2)$$



Detailed simulations for EIC [Jentsch,Tu,Weiss, PRC 21]

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Polarized structure function g_{1n} : longitudinal asymmetry

Pole extrapolation of double spin asymmetry



- Nominator $d\sigma_{||} \equiv \frac{1}{4} \left[d\sigma(+\frac{1}{2},+1) - d\sigma(-\frac{1}{2},+1) - d\sigma(+\frac{1}{2},-1) + d\sigma(-\frac{1}{2},-1) \right]$
- Denominators: 2-state

$$d\sigma_2 \equiv rac{1}{4} \sum_{\Lambda_e} \left[\mathrm{d}\sigma(\Lambda_e, +1) + \mathrm{d}\sigma(\Lambda_e, -1) \right]$$

• $\lambda_d = 0$ state **not needed**!

▶ Impulse approximation yields in the Bjorken limit $\left[\alpha_{p} = \frac{2p_{p}^{-}}{p_{+}^{+}}\right]$

$$\mathcal{A}_{||,2} = \frac{d\sigma_{||}}{d\sigma_{2}} [\phi_{h} \operatorname{avg}] \approx \mathcal{D}_{i}(\alpha_{p}, |p_{pT}|) \mathcal{A}_{||n} = \mathcal{D}_{2}(\alpha_{p}, |p_{pT}|) \frac{D_{||}g_{1n}(\tilde{x}, Q^{2})}{2(1 + \epsilon R_{n})F_{1n}(\tilde{x}, Q^{2})}$$

Nuclear structure factor \mathcal{D}_2

- Quantifies neutron depolarization due to nuclear structure
- Depends on spectator kinematics α_p , p_{pT}
- $D_2 = \Delta S_d$ [pure +1]/ S_d [pure +1] has probabilistic interpretation O_d / O_d



WC, C. Weiss, PLB ('19); PRC ('20)

- Bounds: $-1 \leq \mathcal{D}_2 \leq 1$
- Due to lack of OAM $\mathcal{D}_2\equiv 1$ for $p_T=0$
- Clear contribution from D-wave at finite recoil momenta
- ▶ D₂ close to unity at small recoil momenta
 - \rightarrow relevant for on-shell extrapolation

Tensor polarized observable A_{zz} [Frankfurt, Strikman '83]

► Analogue of $A_{LL} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-}$ for vector polarization (~ S_L) is the tensor asymmetry

$$A_{zz} = \frac{\sigma^+ + \sigma^- - 2\sigma^0}{\sigma^+ + \sigma^- + \sigma^0} [\phi_h \operatorname{avg}]$$

 \rightarrow **no** electron polarization required.

- ▶ Requires all three polarization states; $A_{zz} \in [-2, 1]$ and $\sim T_{LL}$
- ► \vec{ed} : A_{zz} measured in different ranges of Bjorken $x = 2 \frac{Q^2}{2(p_d q)}$, $x \in [0, 2]$: ■ elastic x = 2 [" T_{20} "]: NIKHEF ('98), ...
 - quasi-elastic x > 1: NIKHEF ('99), BLAST @ MIT-Bates ('17), future JLab12
 - DIS inclusive x < 1 [" b_1 "]: HERMES @ DESY ('05), future JLab12 → HERMES measured very small values

Maximize A_{zz} with spectator tagging

► Tensor polarization is sensitive to unpolarized quark distributions, partonic factor cancels out → ratio of LF densities remains

$$A_{zz}(\alpha_p, \boldsymbol{p}_T) = -\frac{\frac{f_0(k)f_2(k)}{\sqrt{2}} + \frac{f_2^{2}(k)}{4}}{f_0^{2}(k) + f_2^{2}(k)} (3\cos 2\theta_k + 1) \qquad \alpha_p = \frac{2p_p^+}{p_D^+} = \left(1 + \frac{k^3}{\sqrt{m^2 + k^2}}\right); \quad \boldsymbol{p}_{pT} = \boldsymbol{k}_T$$



• Maximal A_{zz} at $f_2(k) = \sqrt{2} f_0(k)$, not the *S* wave node!

• $A_{zz} = 1$ at $\alpha_p = 1$ ($\theta_k = \pi/2$) \rightarrow pure $m = \pm 1$) $A_{zz} = -2$ at $p_{pT} = 0$ ($\theta_k = 0, \pi$) \rightarrow pure m = 0

\rightarrow Constraints on deuteron *D*-wave

FSI in DIS: physical pictures

FSI can be min/maximized depending on spectator kinematics



subasymptotic regime: low Q^2 , high x



DIS regime, intermediate \boldsymbol{x}



 rescattering of resonance-like structure with spectator nucleon in eikonal approximation [Deeps,BONuS].

WC,M. Sargsian arXiv:1704.06117

- FSI between slow hadrons from the DIS products and spectator nucleon, fast hadrons hadronize after leaving the nucleus.
- Shadowing in DIS $x \ll 10^{-1}$
- For decreasing x the FSI become more and more low-energy NN scattering

CFNS 2nd detector workshop

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Leading twist nuclear shadowing

⁴He coherent J/Ψ production



- ► Can be studied on heavy and light nuclei, large effect
- Gluonic imaging of nuclei $g_A(x, Q^2)$
- Likely gluons belong to more than 1 nucleon
- Shadowing one nucleon at a time in diffractive J/ψ production on light nuclei
- k-body FF enter in amplitude \rightarrow nuclear wf

- Light ions address many aspects of EIC physics program
- Nuclear physics plays an important role
- ▶ Nuclear breakup and coherent scattering enable novel measurements
- Requires lots of theory input
- Several channels, aspects require work
- Complementarity with second detector