

Diffractive longitudinal structure function at the Electron Ion Collider

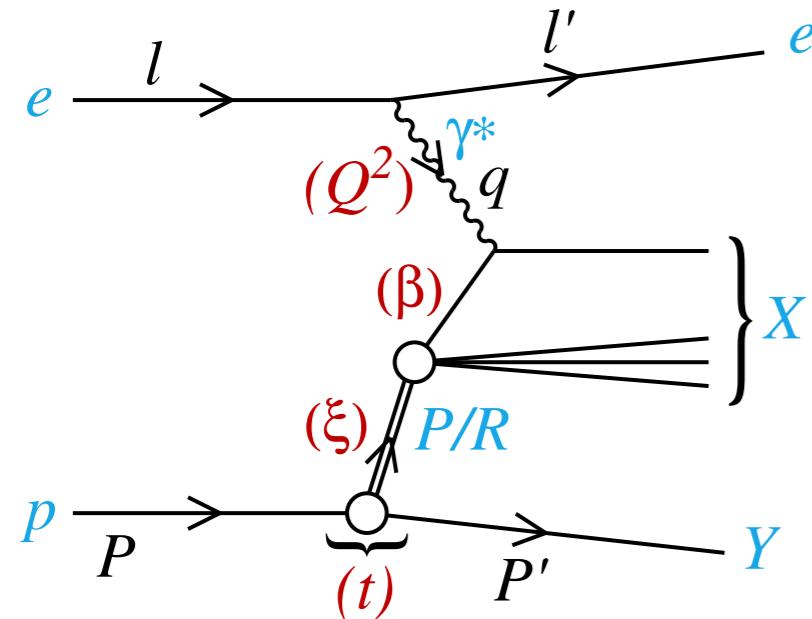
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Outline

- $F_L^{D(3)}$
 - Motivation: why is $F_L^{D(3)}$ interesting ?
 - H1 measurement
 - Proton tagging as a method for diffraction at EIC
 - Pseudodata simulation, energy beam scenarios
 - Extraction by linear fit. Kinematic range and precision
 - $F_L^{D(3)}$ and R ratio of longitudinal to transverse cross section

*Nestor Armesto, Paul Newman, Wojciech Słomiński, A.S., 2112.06839
also EIC Yellow Report*

Diffractive kinematics in DIS



Standard DIS variables:

electron-proton
cms energy squared:
 $s = (k + p)^2$

photon-proton
cms energy squared:
 $W^2 = (q + p)^2$

inelasticity
 $y = \frac{p \cdot q}{p \cdot k}$

Bjorken x
 $x = \frac{-q^2}{2p \cdot q}$
(minus) photon virtuality
 $Q^2 = -q^2$

Diffractive DIS variables:

$$x = \xi \beta$$

$$\xi \equiv x_{IP} = \frac{Q^2 + M_X^2 - t}{Q^2 + W^2}$$

$$\beta = \frac{Q^2}{Q^2 + M_X^2 - t}$$

$$t = (p - p')^2$$

momentum fraction of the
Pomeron w.r.t hadron

momentum fraction of parton
w.r.t Pomeron

4-momentum transfer squared

Diffractive cross section, structure functions

Diffractive cross section depends on 4 variables (ξ, β, Q^2, t):

$$\frac{d^4\sigma^D}{d\xi d\beta dQ^2 dt} = \frac{2\pi\alpha_{\text{em}}^2}{\beta Q^4} Y_+ \sigma_r^{\text{D}(4)}(\xi, \beta, Q^2, t)$$
$$Y_+ = 1 + (1 - y)^2$$

Reduced cross section depends on two **structure functions**:

$$\sigma_r^{\text{D}(4)}(\xi, \beta, Q^2, t) = F_2^{\text{D}(4)}(\xi, \beta, Q^2, t) - \frac{y^2}{Y_+} F_L^{\text{D}(4)}(\xi, \beta, Q^2, t)$$

Upon integration over t :

$$F_{2,L}^{\text{D}(3)}(\xi, \beta, Q^2) = \int_{-\infty}^0 dt F_{2,L}^{\text{D}(4)}(\xi, \beta, Q^2, t)$$

Dimensions:

$$[\sigma_r^{\text{D}(4)}] = \text{GeV}^{-2}$$

$$\sigma_r^{\text{D}(3)} \quad \text{Dimensionless}$$

When $y \ll 1$

$$\sigma_r^{\text{D}(4,3)} \simeq F_2^{\text{D}(4,3)}$$

Why $F_L^D(3)$ is interesting? $F_L^D(3)$ at HERA

Why F_L^D is interesting?

F_L^D vanishes in the parton model

Gets non-vanishing contributions in QCD

As in inclusive case, particularly sensitive to the diffractive **gluon density**

Expected large **higher twists**, provides test of the **non-linear, saturation** phenomena

Experimentally challenging...

Measurement requires several beam energies

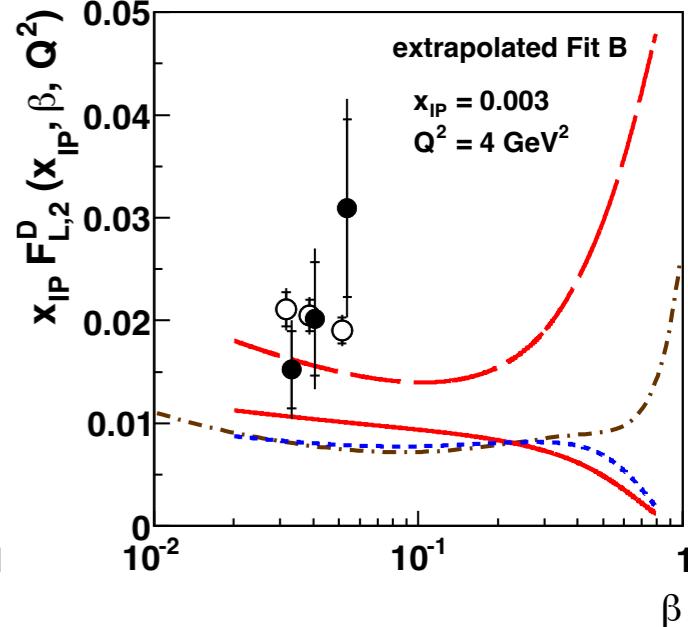
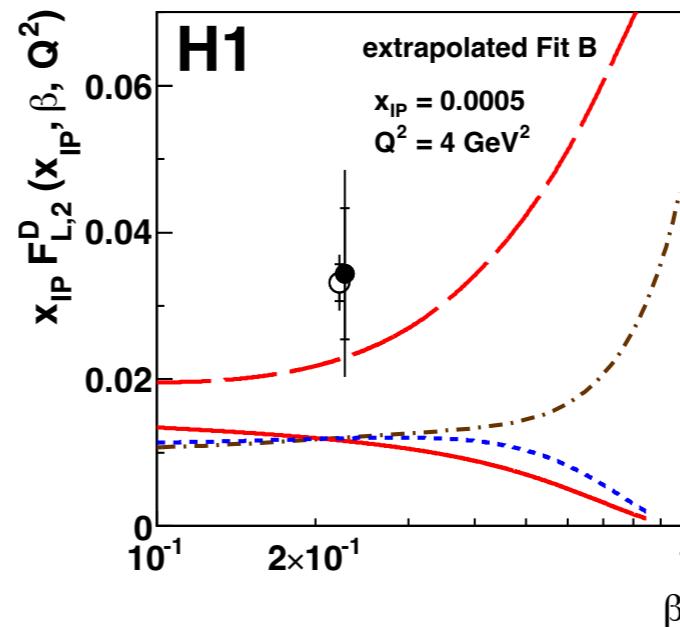
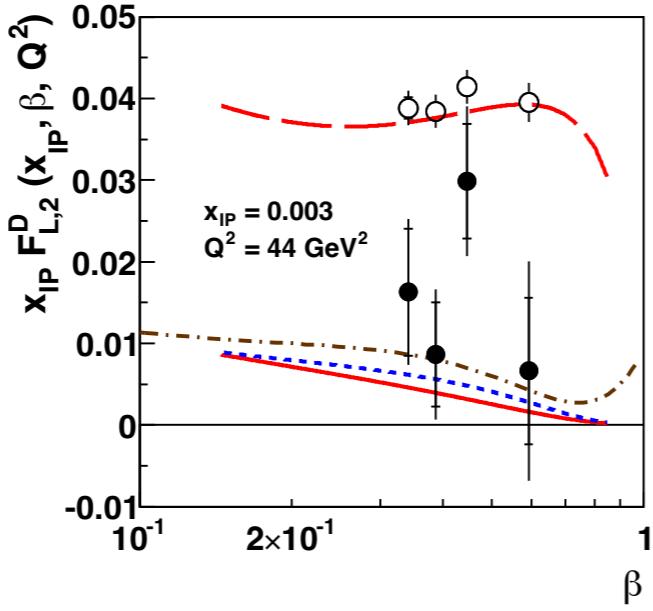
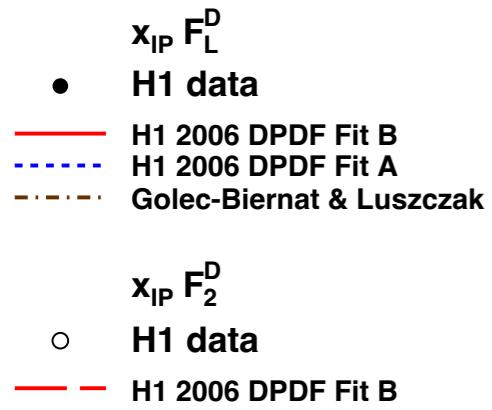
F_L^D strongest when $y \rightarrow 1$. Low electron energies

H1 measurement: 4 energies, $E_p = 920, 820, 575, 460$ GeV, electron beam $E_e = 27.6$ GeV

Large errors, limited by statistics at HERA

Careful evaluation of systematics. Best precision 4%, with uncorrelated sources as low as 2%

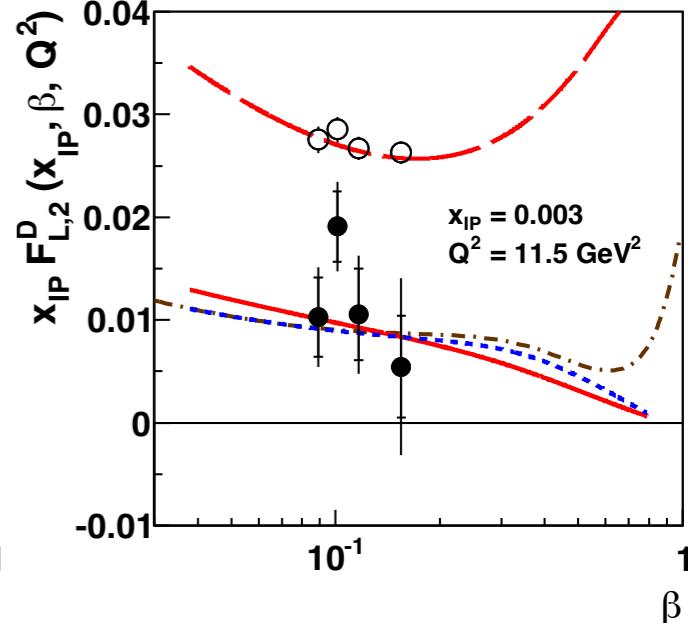
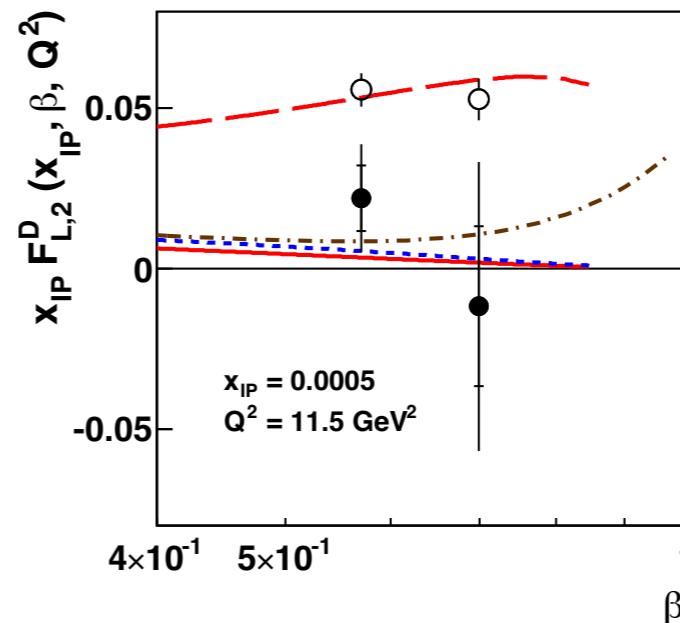
$F_L^D(3)$ at HERA



Measurements of σ_r^D consistent with predictions from the models

Extracted F_L^D has a tendency to be higher than the predictions, though compatible with model predictions within errors

Overall: $0 < F_L^D < F_2^D$



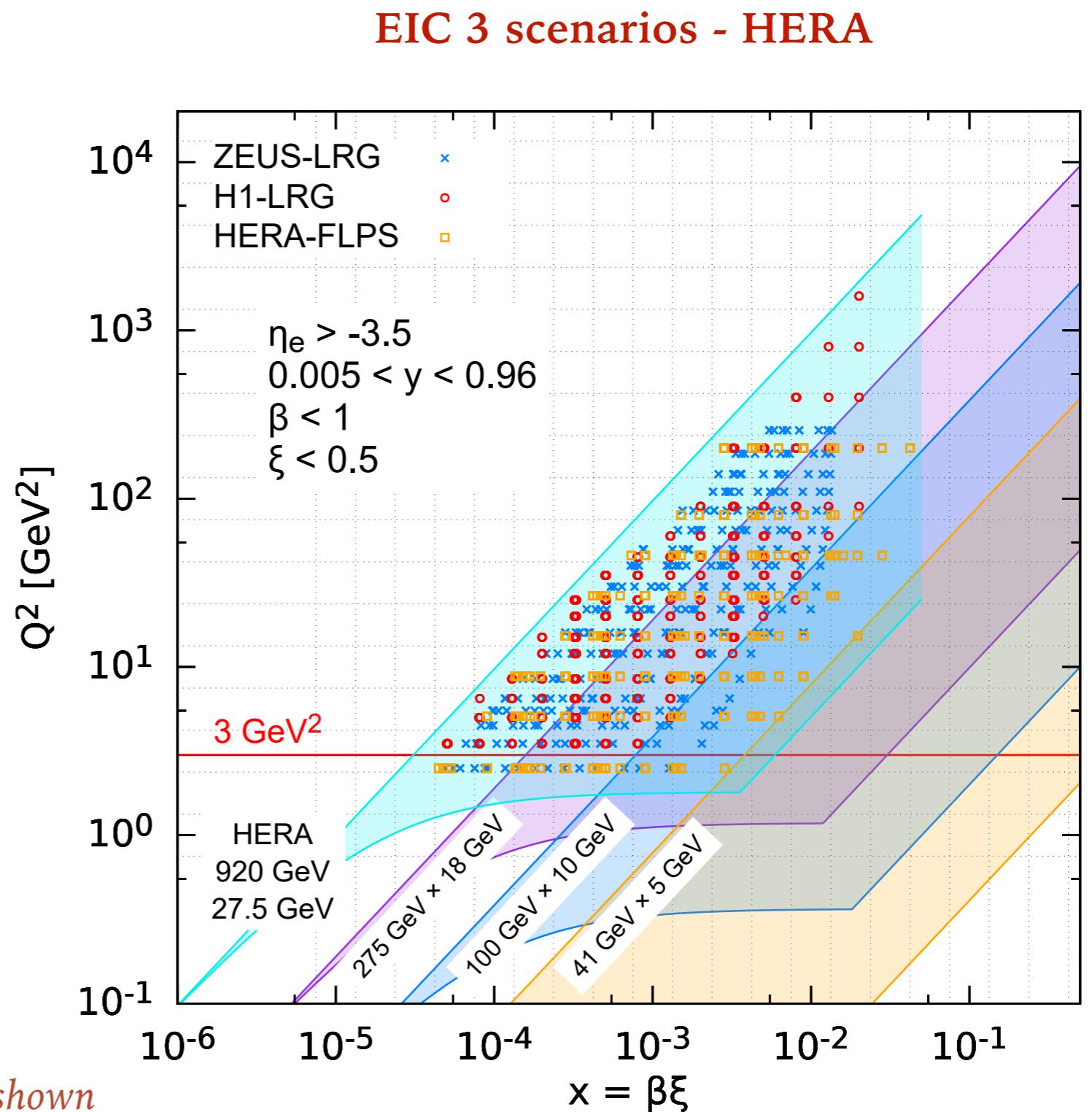
Phase space (x, Q^2) EIC-HERA

EIC can operate at various energy combinations

Can cover wide range of x

Large instantaneous luminosity

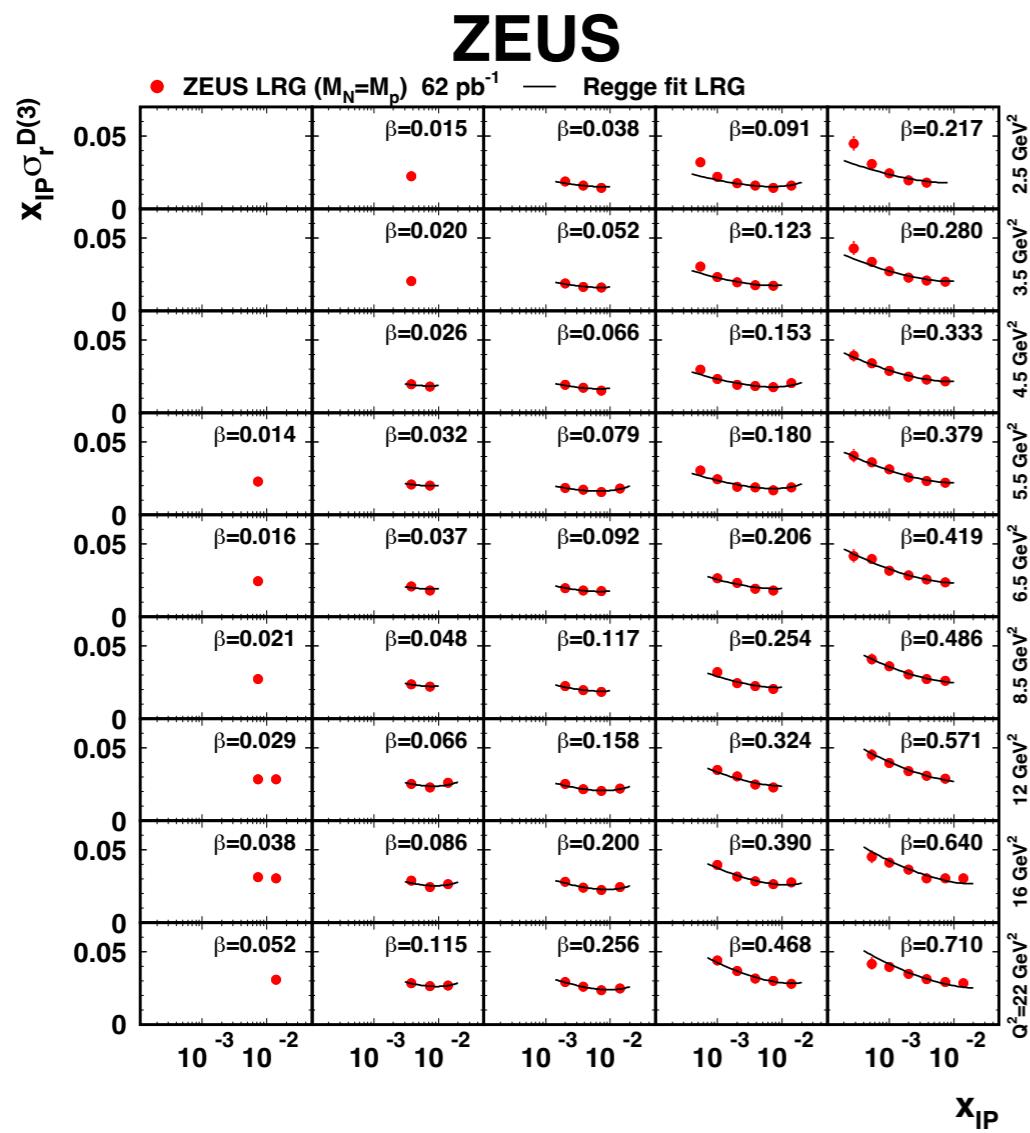
Statistics should not be a limiting factor



Measurement methods: LRG vs LP

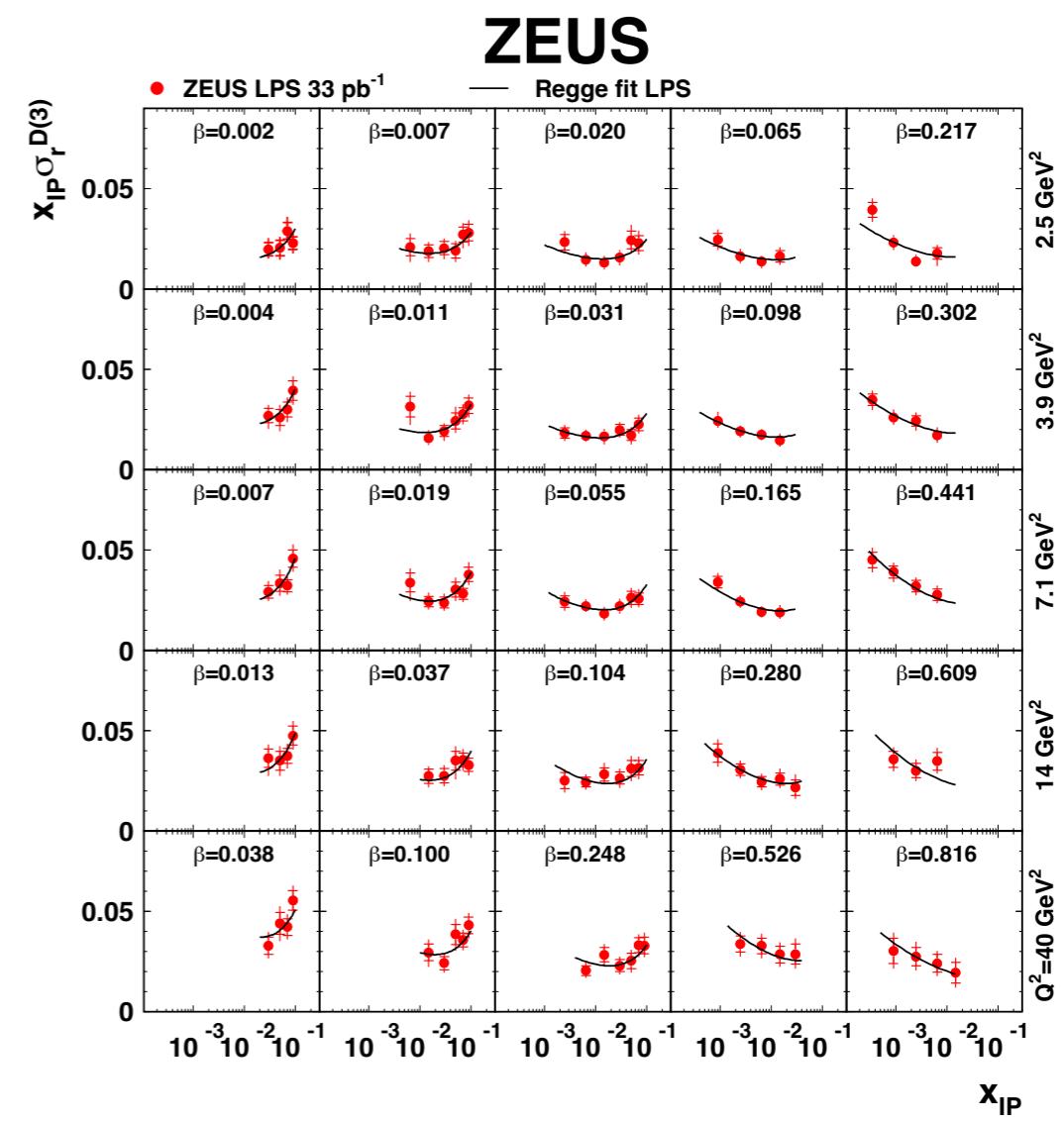
Large Rapidity Gap method:

request a large rapidity gap (ex. ZEUS 2009
 $\xi < 0.02$)

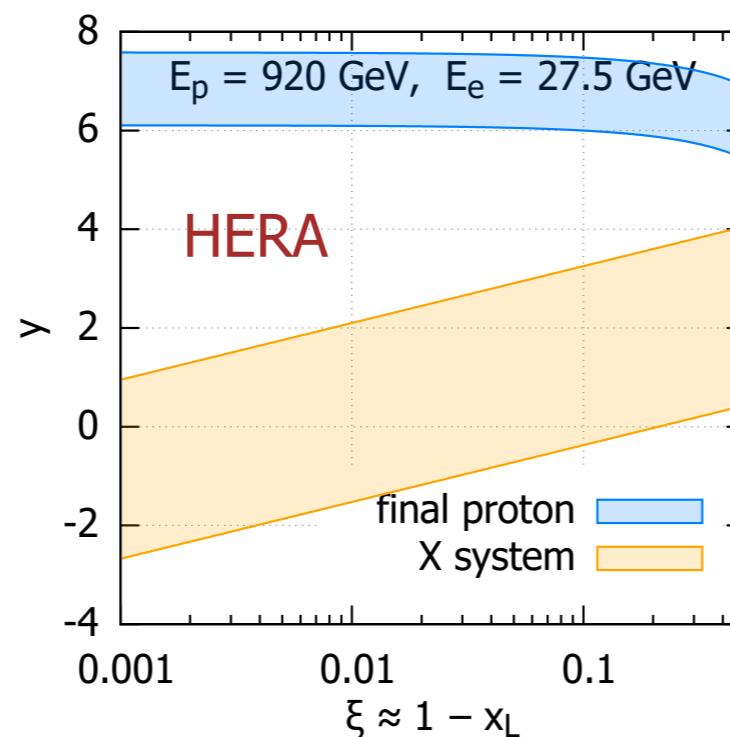
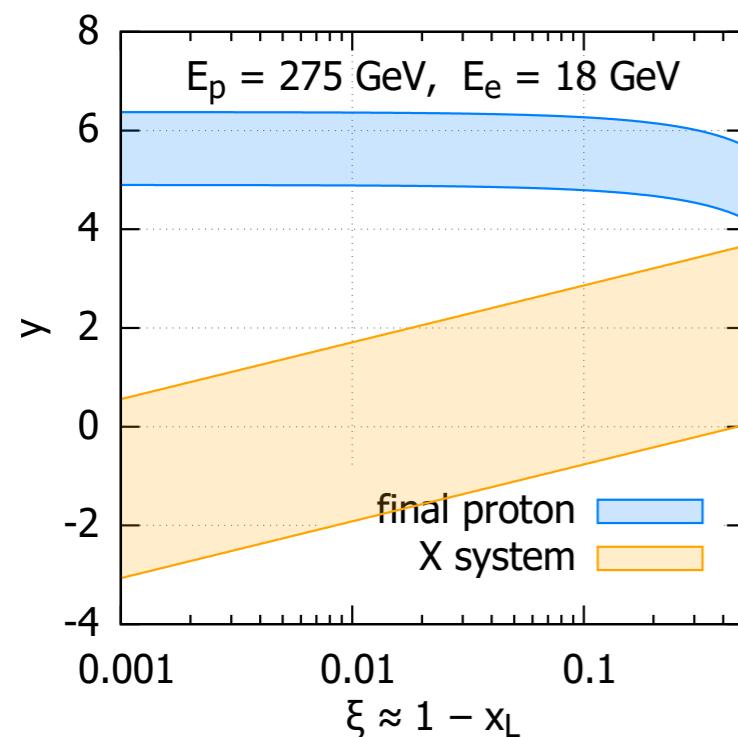
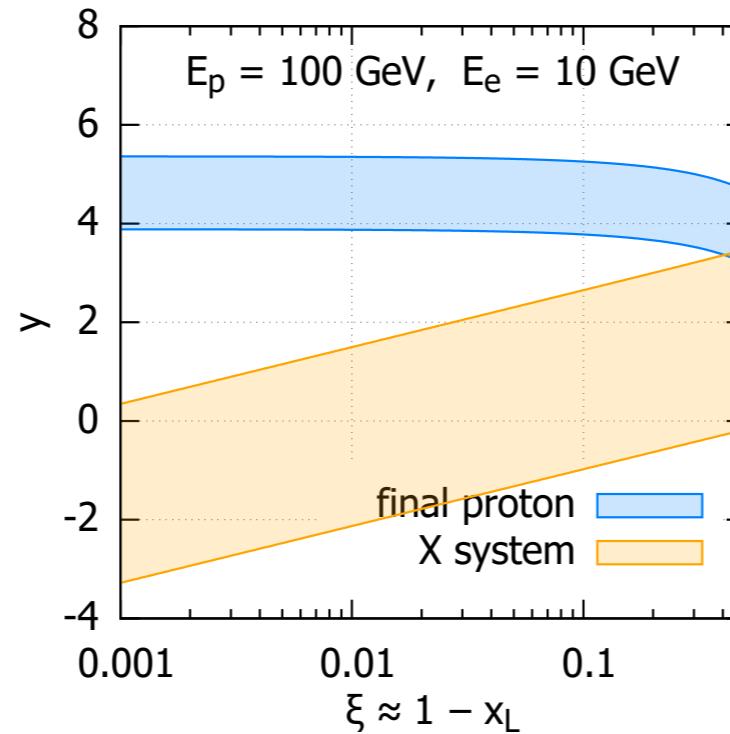
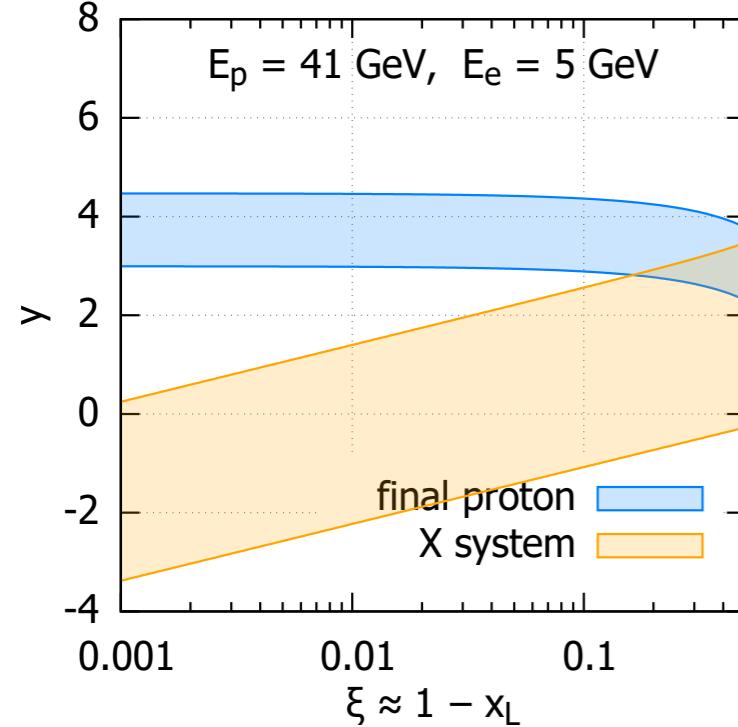


Proton Tagging (Leading Proton) method:

detection of a leading proton (ex. Leading Proton Spectrometer in ZEUS, Forward Proton Spectrometer in H1, can go to higher $\xi < 0.1$)



Rapidity range at EIC in diffraction



Rapidity range of proton and undecayed diffractive system X

$$p_T^{proton} < 4 \text{ GeV}$$

$$0.1 < \beta < 0.9$$

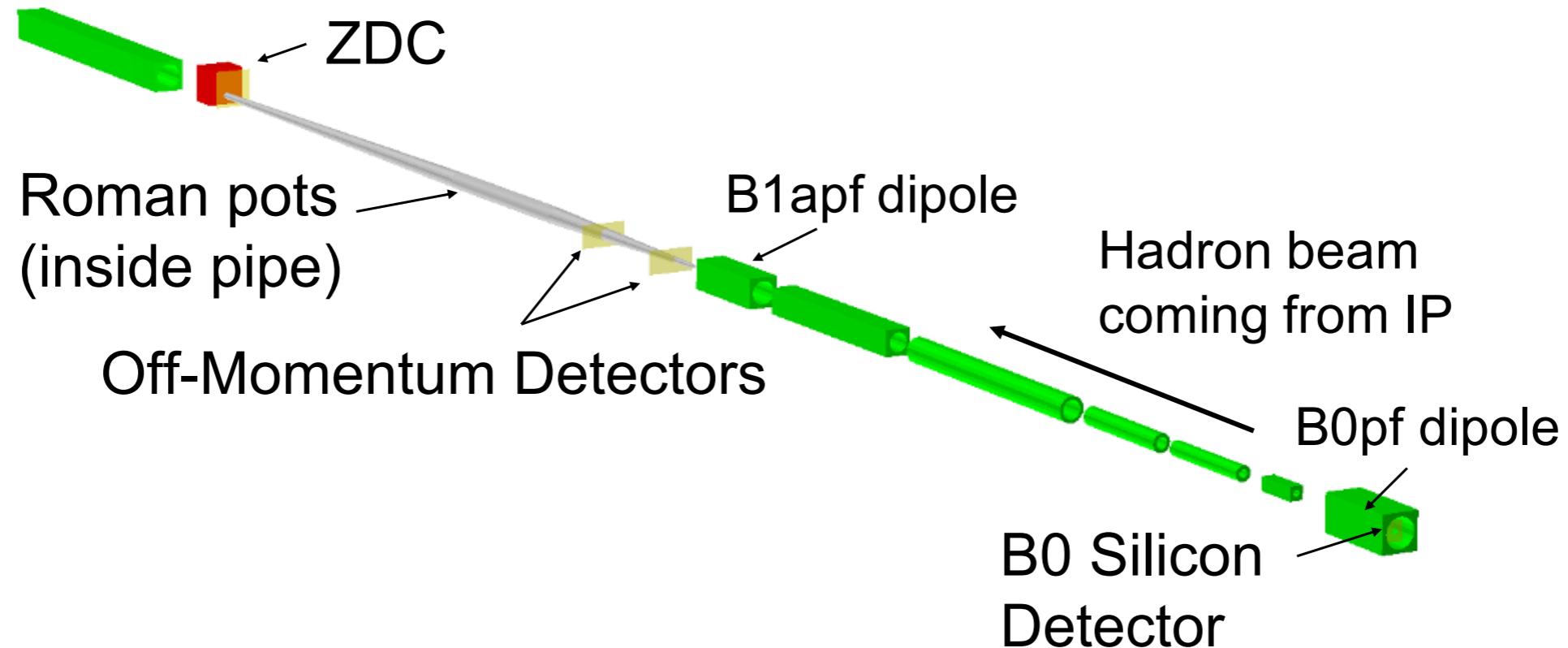
$$0.005 < y < 0.96$$

HERA: LRG method reliable for gaps > 3 units of rapidity

EIC: fairly large gaps ($\Delta\eta \geq 4$) exist for smallest ξ and largest s

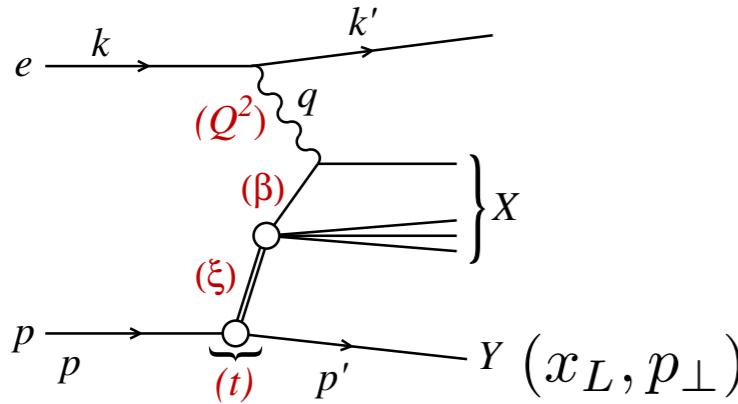
However, through most region LRG method may be challenging at EIC

Far forward detectors at EIC



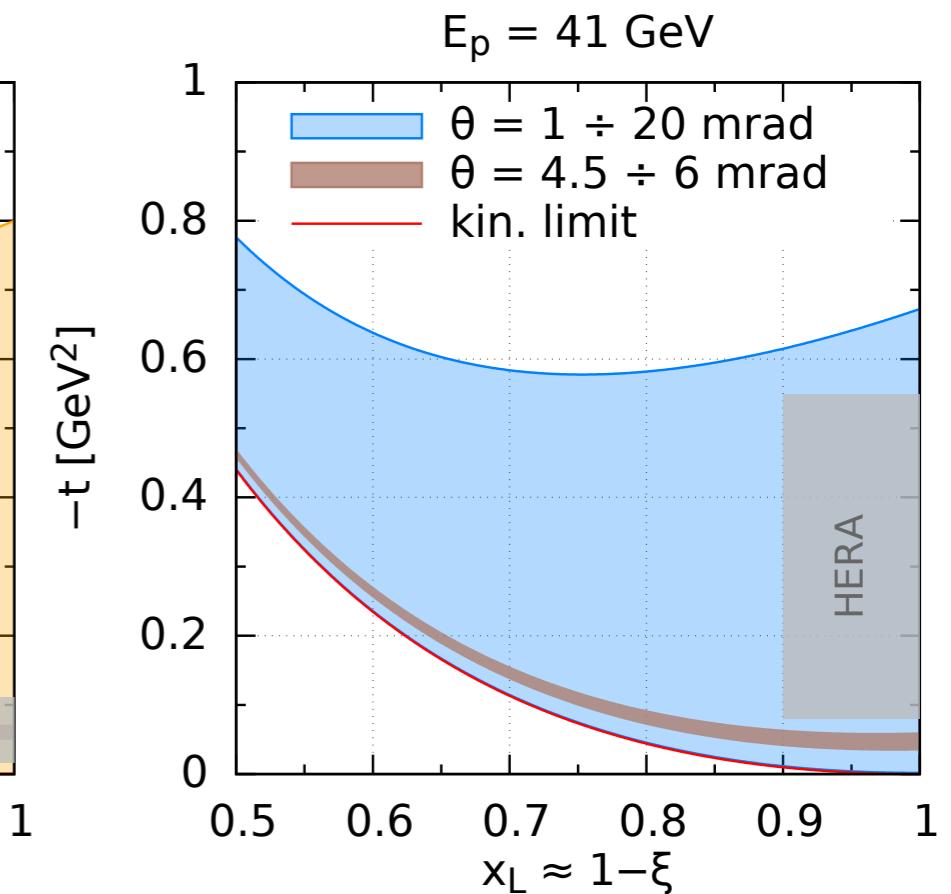
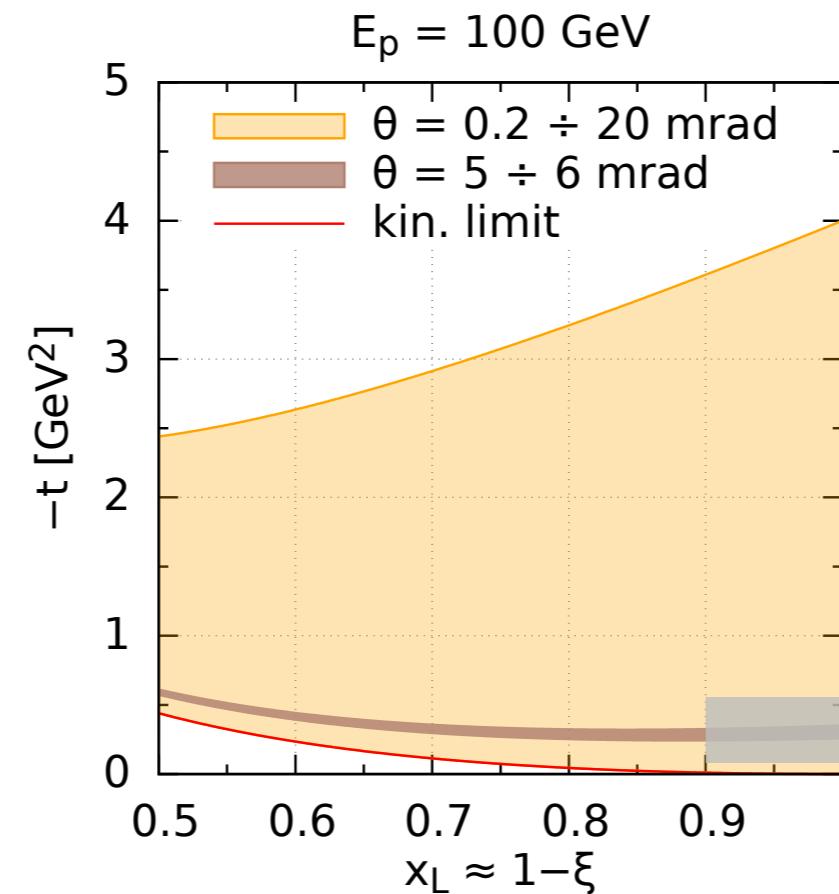
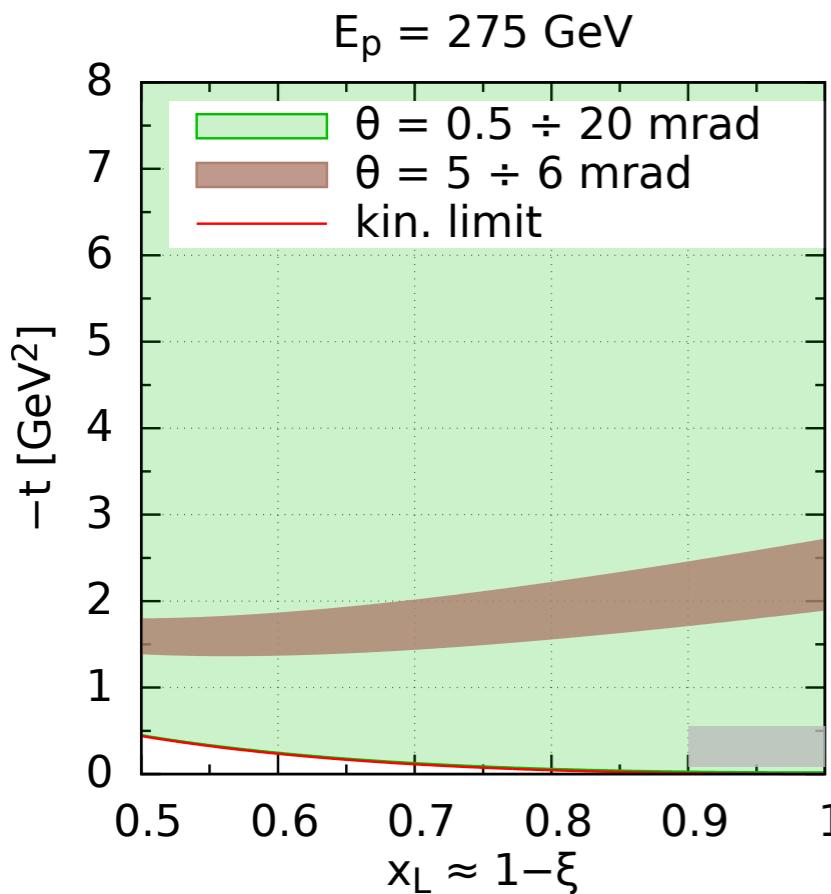
Detector	Angle	Position [m]
ZDC	$\theta < 5.5$ mrad	37.5
Roman Pots	$0.5 < \theta < 5.0$ mrad	26.0, 28.0
Off-momentum detectors	$\theta < 5.0$ mrad	22.5, 25.5
B0	$6.0 < \theta < 20.0$ mrad	$5.4 < z < 6.4$

Final proton tagging



Small angle acceptance i.e. Roman pots

(x_L, p_\perp, θ) measured in LAB, collinear (e, p) frame



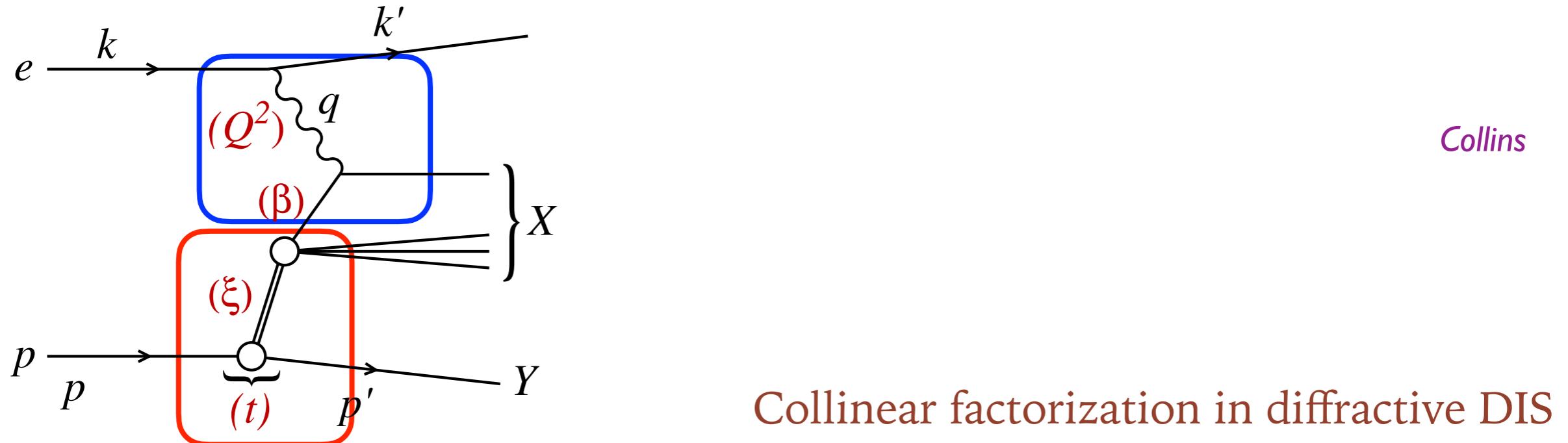
Much better than at HERA

Best way to select diffractive events through proton tagging

$$t = -\frac{p_\perp^2}{x_L} - \frac{(1 - x_L)^2}{x_L} m_p^2$$

Pseudodata generation: collinear factorization for diffraction

Use the collinear factorization for the description of HERA and pseudodata simulation



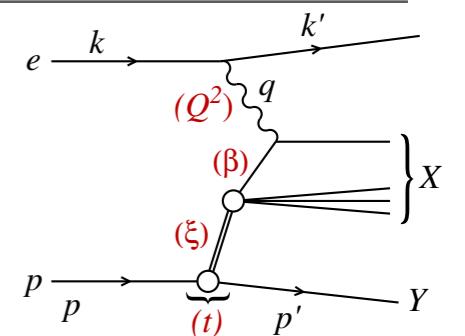
$$F_{2/L}^{D(4)}(\beta, \xi, Q^2, t) = \sum_i \int_{\beta}^1 \frac{dz}{z} C_{2/L,i} \left(\frac{\beta}{z}, Q^2 \right) f_i^D(z, \xi, Q^2, t)$$

- Diffractive cross section can be factorized into the convolution of the perturbatively calculable **partonic cross sections** and **diffractive parton distributions** (DPDFs).
- The DPDFs are **non-perturbative** objects, but evolved perturbatively with DGLAP

Pseudodata generation: model for diffractive structure functions

Regge factorization with **Pomeron** terms works for small $\xi < 0.01$

At higher ξ additional exchanges ‘**Reggeons**’ need to be included



$$f_i^{D(4)}(z, \xi, Q^2, t) = f_{IP}^p(\xi, t) f_i^{IP}(z, Q^2) + f_{IR}^p(\xi, t) f_i^{IR}(z, Q^2)$$

Pomeron

Reggeon

Regge type flux:

$$f_{IP,IR}^p(\xi, t) = A_{IP,IR} \frac{e^{B_{IP,IR} t}}{\xi^{2\alpha_{IP,IR}(t)-1}}$$

For t-integrated case

$$f_i^{D(3)}(z, \xi, Q^2) = \phi_{IP}^p(\xi) f_i^{IP}(z, Q^2) + \phi_{IR}^p(\xi) f_i^{IR}(z, Q^2)$$

Trajectory:

$$\alpha_{IP,IR}(t) = \alpha_{IP,IR}(0) + \alpha'_{IP,IR} t.$$

Integrated flux:

$$\phi_{IP,IR}^p(\xi) = \int dt f_{IP,IR}^p(\xi, t)$$

Pomeron PDFs obtained via NLO DGLAP evolution starting at initial scale $\mu_0^2 = 1.8 \text{ GeV}^2$

$$z f_i(z, \mu_0^2) = A_i z^{B_i} (1-z)^{C_i} \quad i=q,g$$

Pseudodata generation: energy choice

$$\sigma_{\text{red}}^{\text{D}(3)} = F_2^{\text{D}(3)}(\beta, \xi, Q^2) - Y_L F_L^{\text{D}(3)}(\beta, \xi, Q^2) \quad \text{Integrated over t-momentum transfer}$$

$$Y_L = \frac{y^2}{Y_+} = \frac{y^2}{1 + (1 - y)^2}$$

Can disentangle $F_2^{\text{D}(3)}$ from $F_L^{\text{D}(3)}$ by varying energy and performing the linear fit.

$$y = \frac{Q^2}{xs} = \frac{Q^2}{\beta\xi s} \quad \text{Need to vary the energy } \sqrt{s} \text{ to change } y \text{ for fixed } (\beta, \xi, Q^2)$$

EIC energies for **electron** and **proton**:

$$E_e = 5, 10, 18 \text{ GeV}$$

$$E_p = 41, 100, 120, 165, 180, 275 \text{ GeV}$$

		$E_p [\text{GeV}]$					
		41	100	120	165	180	275
$E_e [\text{GeV}]$	5	29	45	49	57	60	74
	10	40	63	69	81	85	105
	18	54	85	93	109	114	141

S-17 all 17 combinations

S-9 9 - bold red

S-5 5 - green (EIC preferred)

Pseudodata generation

Binning and cuts

Uniform logarithmic binning, 4 bins per order of magnitude in each β, Q^2, ξ

Bins in (ξ, β, Q^2) , common to at least four beam setups

$Q^2 > 3 \text{ GeV}^2$ both H1 and ZEUS fits indicate deterioration of fits for low Q^2

$0.96 > y > 0.005$ expected coverage of the experiment

Simulations

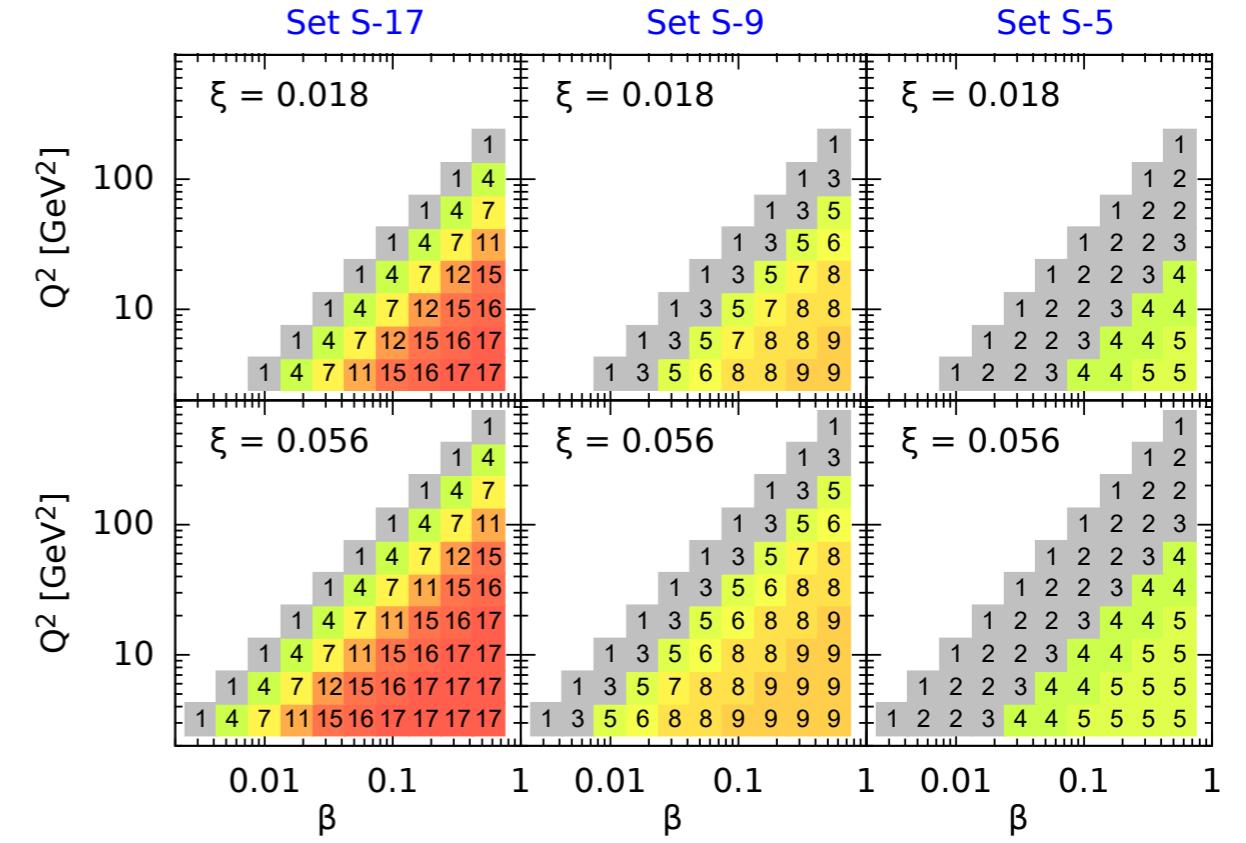
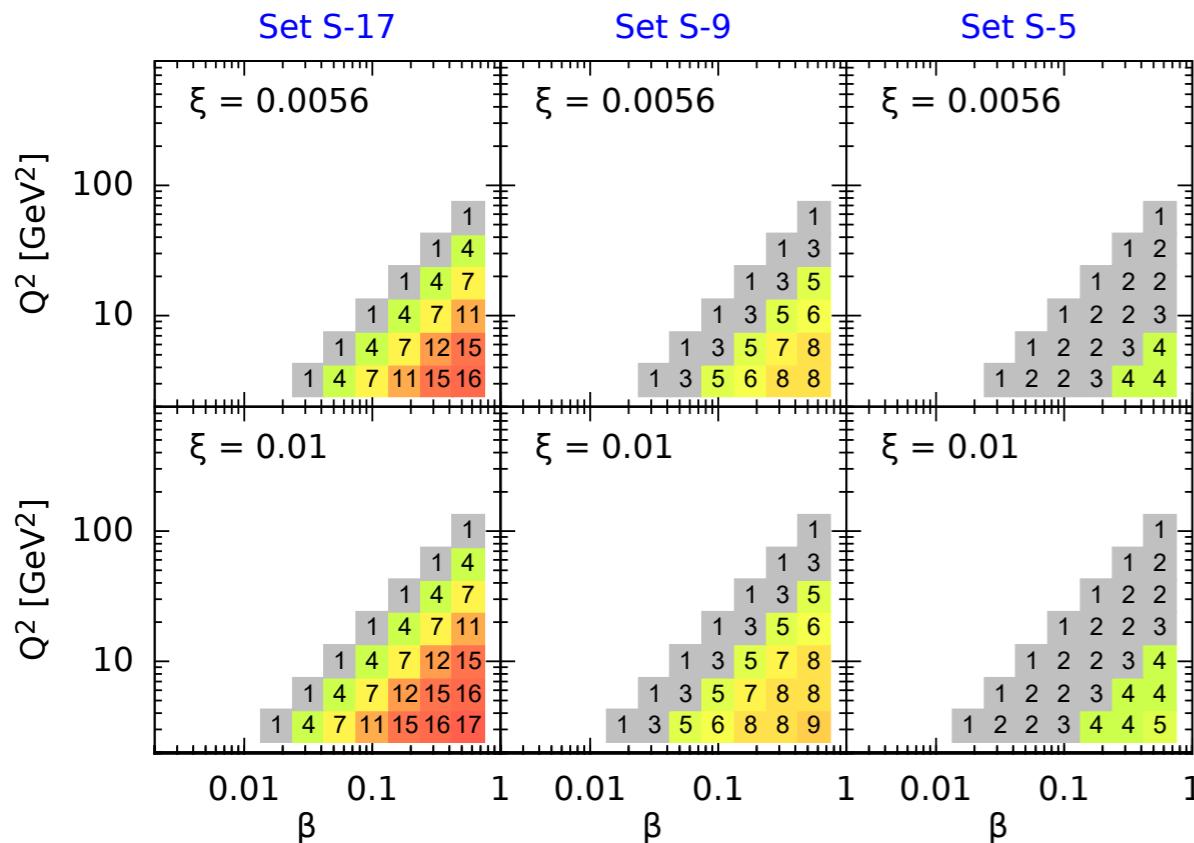
Cross section generation from ZEUS-SJ diffractive PDFs evolved with DGLAP

Assumed $\delta_{\text{sys}} = 1\text{-}2\%$, extrapolated from HERA 2% uncorrelated systematics;
normalization/correlated systematics negligible effect on extraction of F_L^D

δ_{stat} from 10 fb^{-1} integrated luminosity

Several random samples are generated

Kinematic range and number of points



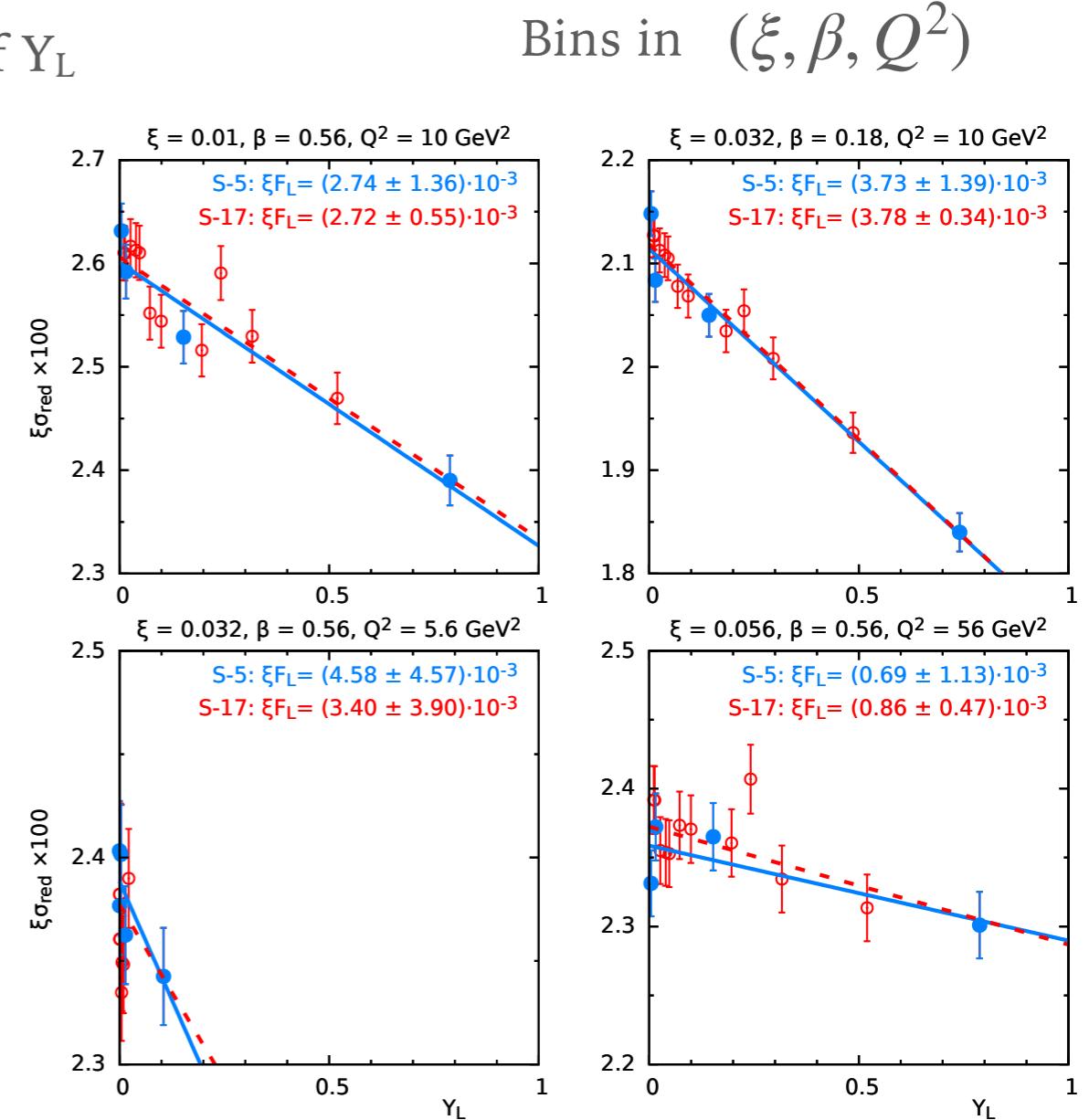
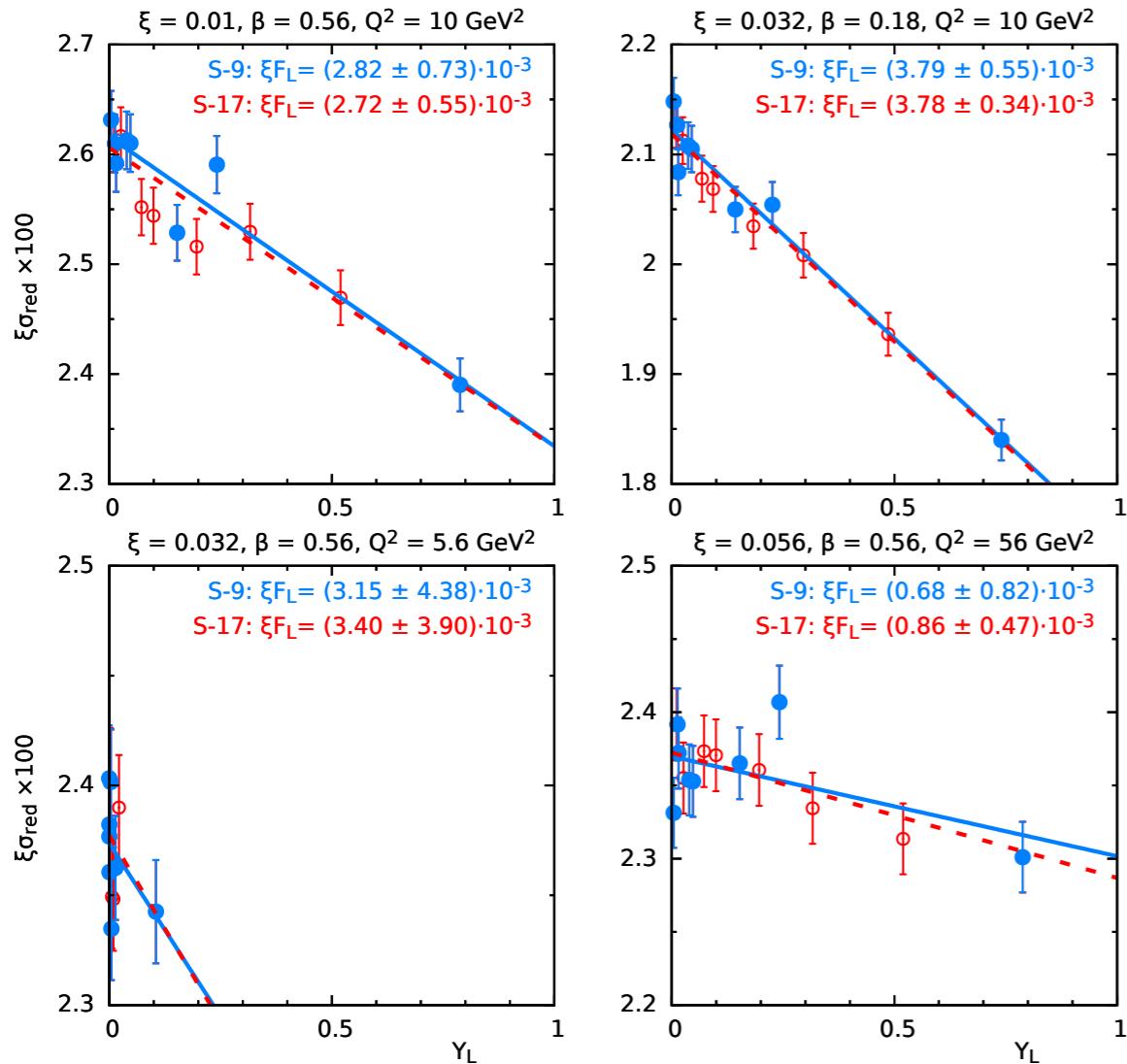
Count of different beam energy combinations for S-17, S-9, S-5

Only points with more than 4 combinations are taken for F_L extraction

Set-17: 364, set-9: 285, set-5: 160 values of F_L

$F_L^D(3)$ extraction

$\sigma_r = F_2(\xi, \beta, Q^2) - Y_L F_L(\xi, \beta, Q^2)$ as a function of Y_L



Uncorrelated systematics 1%

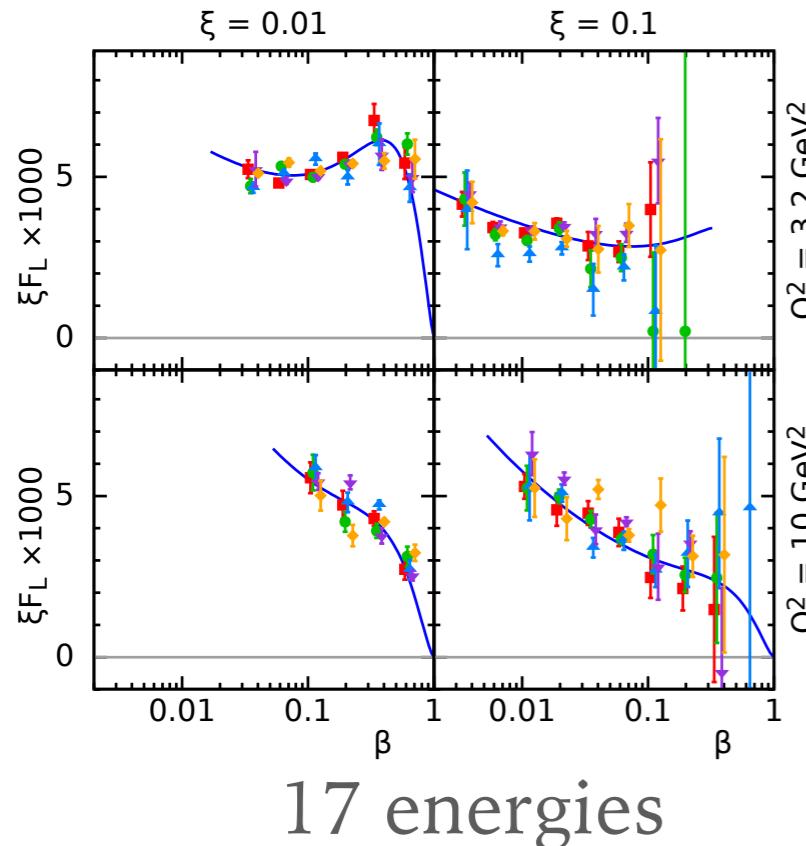
Differences between S-17 and S-9, S-5 small

Increase in error bar on the extraction when smaller number of energy points

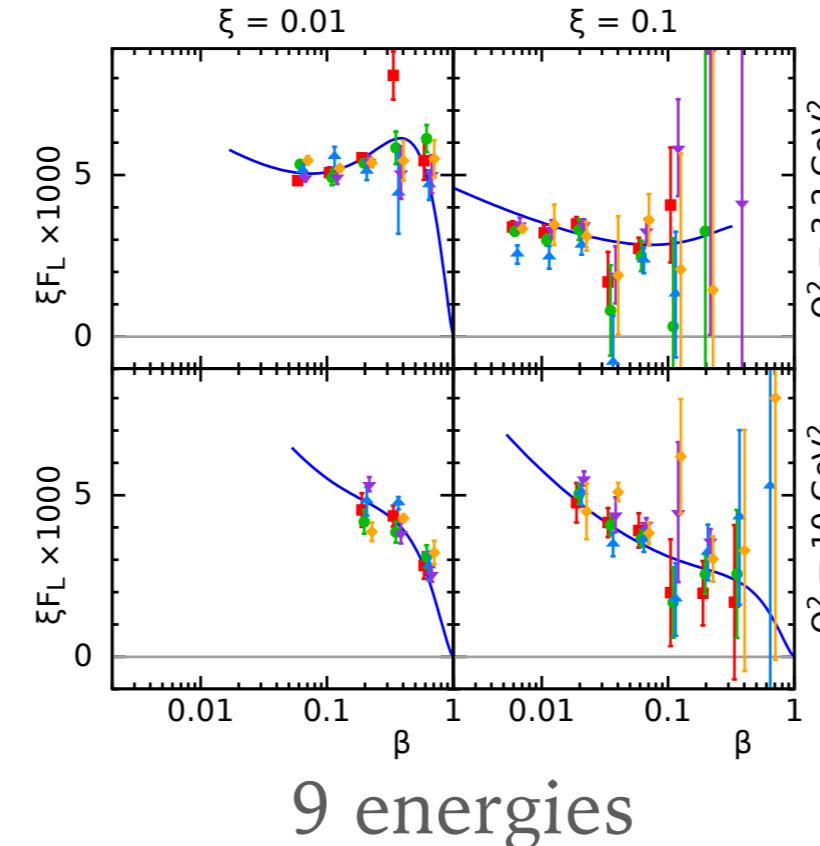
Largest errors for bins with shortest range of Y_L

Simulated measurement of $F_L^D(3)$ vs β in bins of (ξ, Q^2)

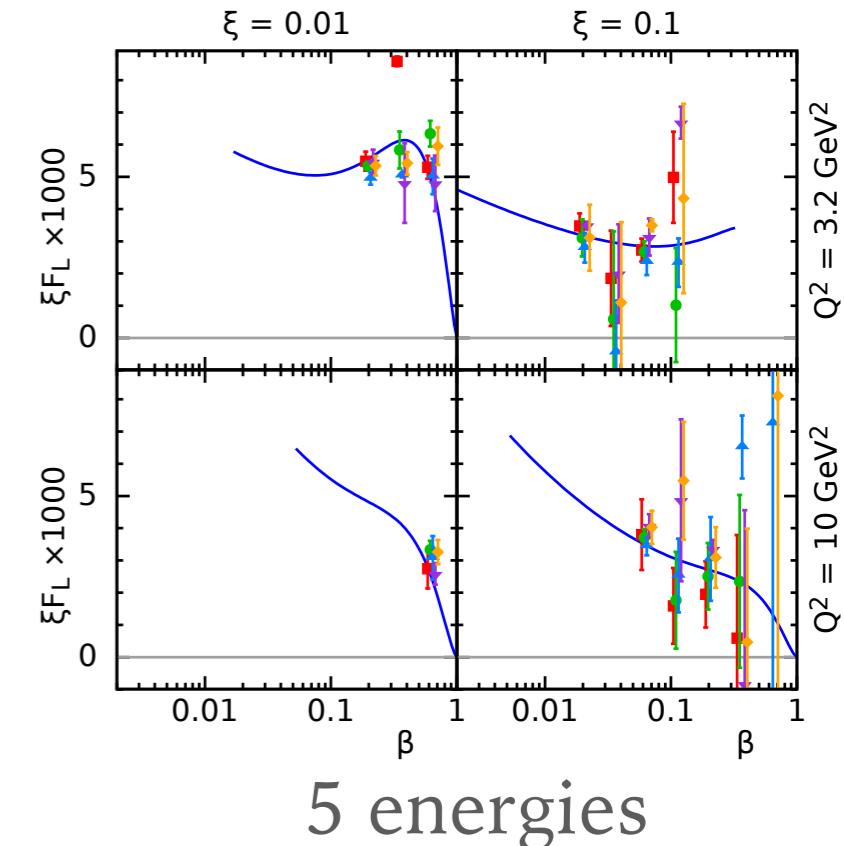
Uncorr. systematic error 1%, 5 MC samples to illustrate fluctuations



17 energies



9 energies



5 energies

Small differences between S-17 and S-9, small reduction to range and increase in uncertainties.

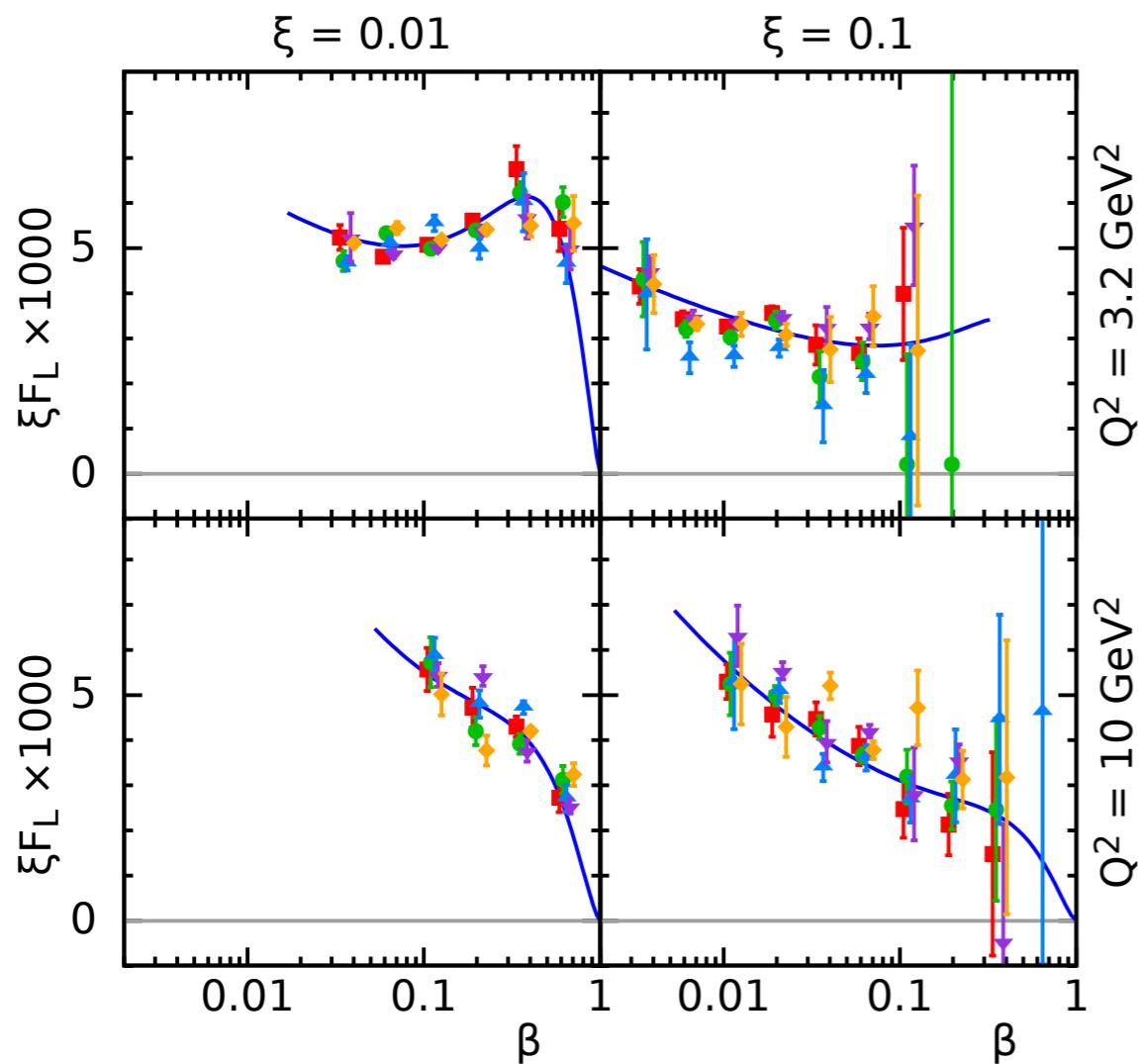
More pronounced reduction in range and higher uncertainties in S-5.

An extraction of F_L^D possible with EIC-favored set of energy combinations

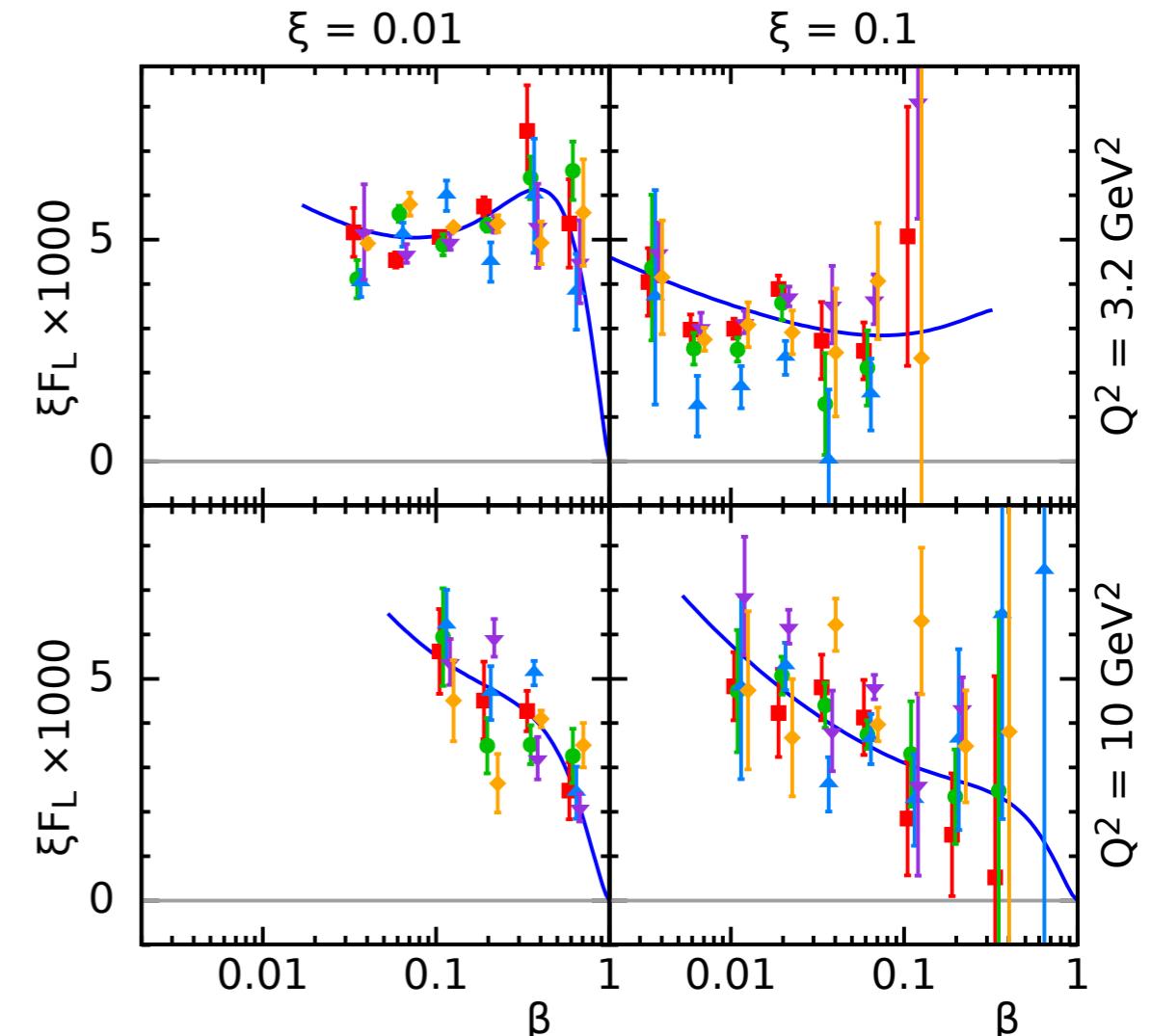
Simulated measurement of $F_L^D(3)$ vs β in bins of (ξ, Q^2)

S-17

$\delta_{\text{sys}} = 1 \%$



$\delta_{\text{sys}} = 2 \%$



Change from 1% to 2% results in roughly twice large error bars

Statistical errors negligible

$F_L^{D(3)}$ fit accuracy

Estimate the accuracy of extraction for $F_L^{D(3)}$

Generate several MC samples of pseudodata
and perform fits

Use direct arithmetic averaging

average

$$\bar{v} = \frac{S_1}{N}$$

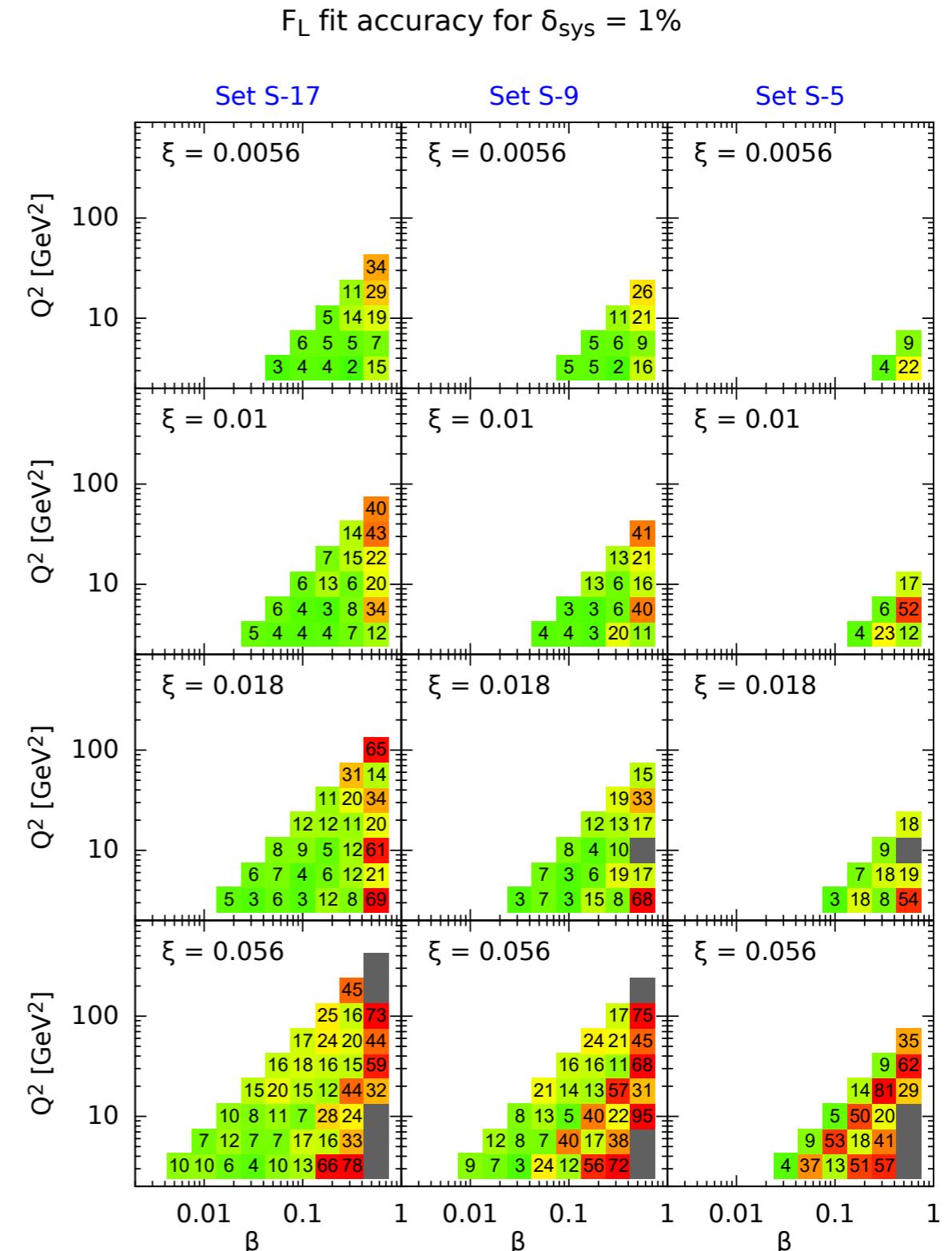
$$S_n = \sum_{i=1}^N v_i^n$$

variance

$$(\Delta v)^2 = \frac{S_2 - S_1^2/N}{N-1}$$

Where v_i is the value of F_L^D

in Monte Carlo sample i



$R^D = F_L^D / F_T^D$ ratio of longitudinal to transverse

Ratio of cross section for longitudinally polarized photons to cross sections for transverse polarized photons

$$R^{D(3)} = F_L^{D(3)} / F_T^{D(3)}$$

$$F_T^{D(3)} = F_2^{D(3)} - F_L^{D(3)}$$

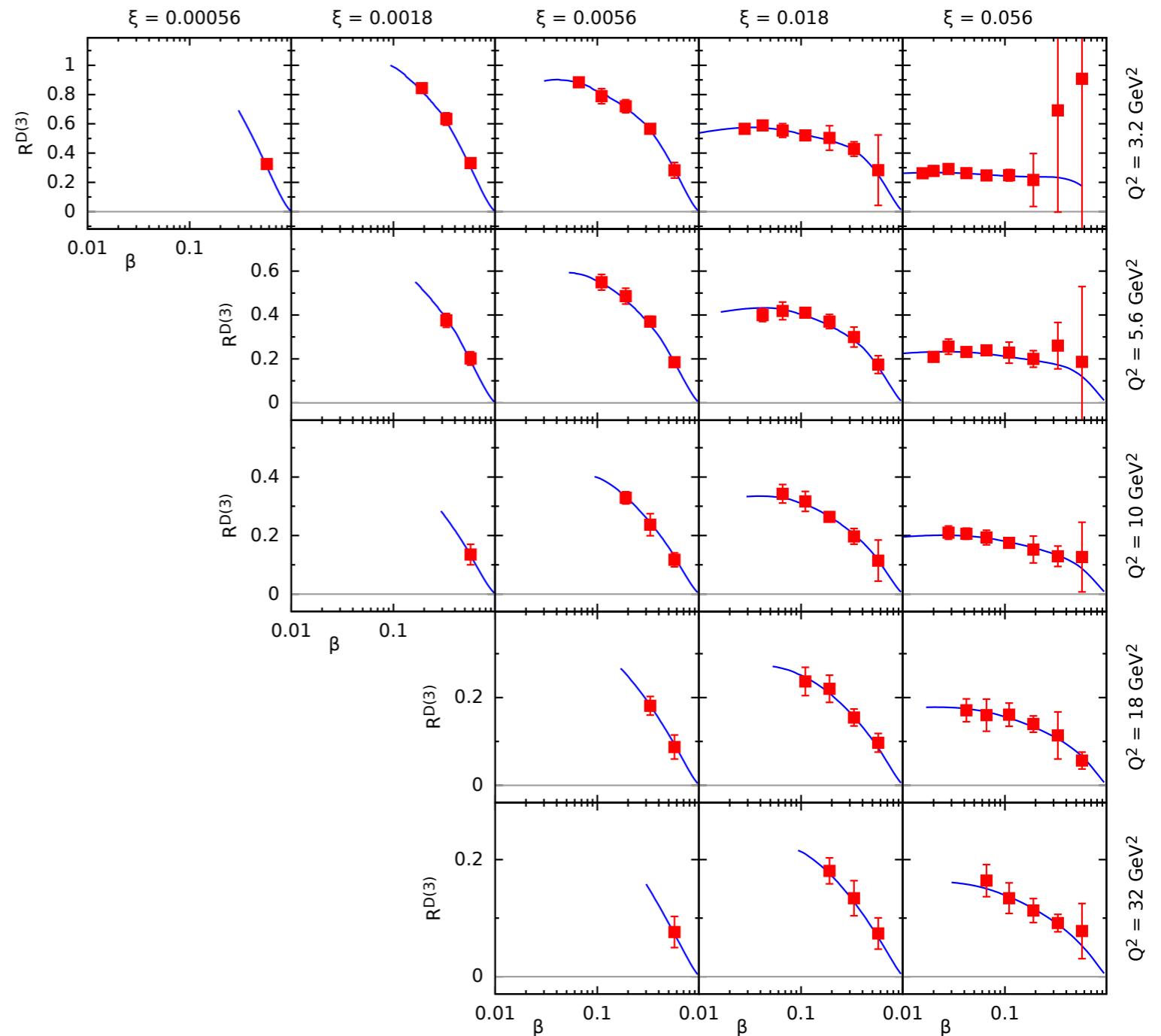
$$\sigma_{\text{red}}^{D(3)} = [1 + (1 - Y_L)R^{D(3)}]F_T^{D(3)}$$

Different form of reduced cross section

Alternative fit has different sensitivities to the uncertainties

Systematics 1%

Averaged over 10 MC samples:
reduced fluctuations



Summary and outlook

- Investigated potential of EIC for the longitudinal structure function in diffraction : $F_L^{D(3)}$
- Important quantity, sensitive to diffractive gluon density (saturation, higher twists...). Only one extraction at HERA by H1, large errors. Challenging measurement.
- Three scenarios: 17, 9, 5 energy combinations. Pseudodata from DGLAP, assumed 1-2% systematics, 10 fb^{-1} integrated luminosity. Extraction via linear fit to reduced cross section
- Scenarios S-17 and S-9 do not differ much, S-5 reduced kinematic range
- Precision in a given bin of (Q^2, ξ, β) correlates strongly with range in inelasticity y , dominated by systematics.
- **Overall: good prospects for the measurement of $F_L^{D(3)}$ at EIC even with 5 energy combinations**
- **Outlook:**
 - Potential for 4-dim diffractive structure function measurement in large t -range
 - Pomeron/Reggeon component separation

Backup

Simulations of $\sigma^D(4)$

What can we learn about the t-dependence of the diffractive structure function?

Diffractive cross section depends on 4 variables (ξ, β, Q^2, t) :

$$\frac{d^4 \sigma^D}{d\xi d\beta dQ^2 dt} = \frac{2\pi\alpha_{\text{em}}^2}{\beta Q^4} Y_+ \sigma_r^{D(4)}(\xi, \beta, Q^2, t)$$
$$Y_+ = 1 + (1 - y)^2$$

$$d\sigma^{ep \rightarrow eXY}(\beta, \xi, Q^2, t) = \sum_i \int_\beta^1 dz \, d\hat{\sigma}^{ei} \left(\frac{\beta}{z}, Q^2 \right) f_i^D(z, \xi, Q^2, t)$$

Ansatz for DPDFs:

$$f_i^{D(4)}(z, \xi, Q^2, t) = f_{IP}^p(\xi, t) f_i^{IP}(z, Q^2) + f_{IR}^p(\xi, t) f_i^{IR}(z, Q^2)$$

Pomeron *Reggeon*

At HERA Reggeon part could not be extracted precisely.

Is it possible to disentangle Pomeron/Reggeon at EIC ?

Simulations of $\sigma^D(4)$

High luminosity and excellent possibility of proton tagging

Prospect of high quality data for $\sigma_r^{D(4)}(\xi, \beta, Q^2, t)$

From the ZEUS-SJ fit

$$\xi\varphi_P(\xi, t) \propto \xi^{-0.22} e^{-7|t|}$$

$$\xi\varphi_R(\xi, t) \propto \xi^{0.6+1.8|t|} e^{-2|t|}$$

Very different slopes in t for
Reggeon and Pomeron

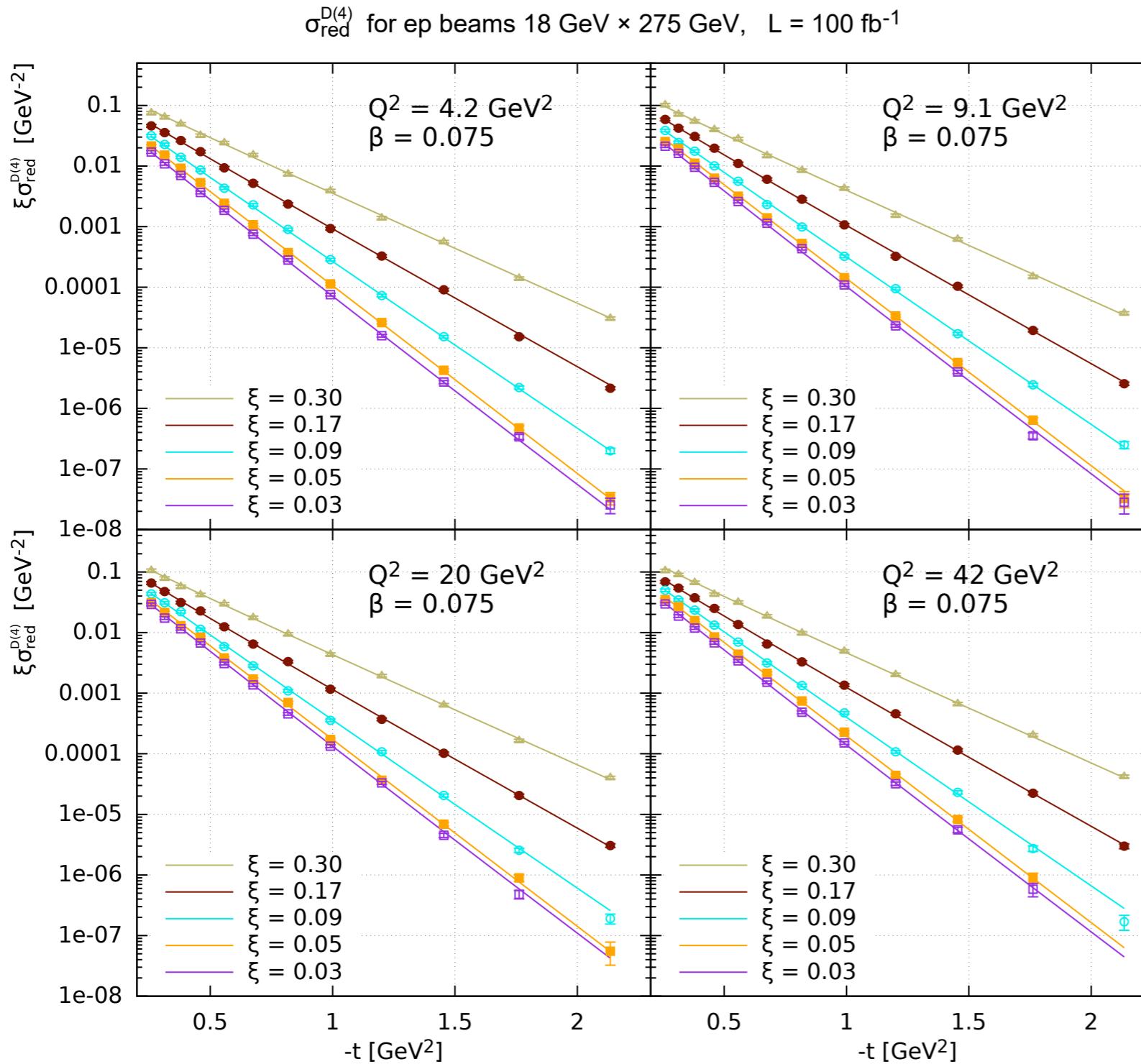
Extrapolation of reduced cross section using ZEUS SJ fit

Random smearing with errors:

Systematic: $\delta_{\text{sys}} = 5\%$

Statistics: δ_{stat} from integrated luminosity 10 fb^{-1}

$\sigma^D(4)$ vs t



$$E_e = 18 \text{ GeV}$$

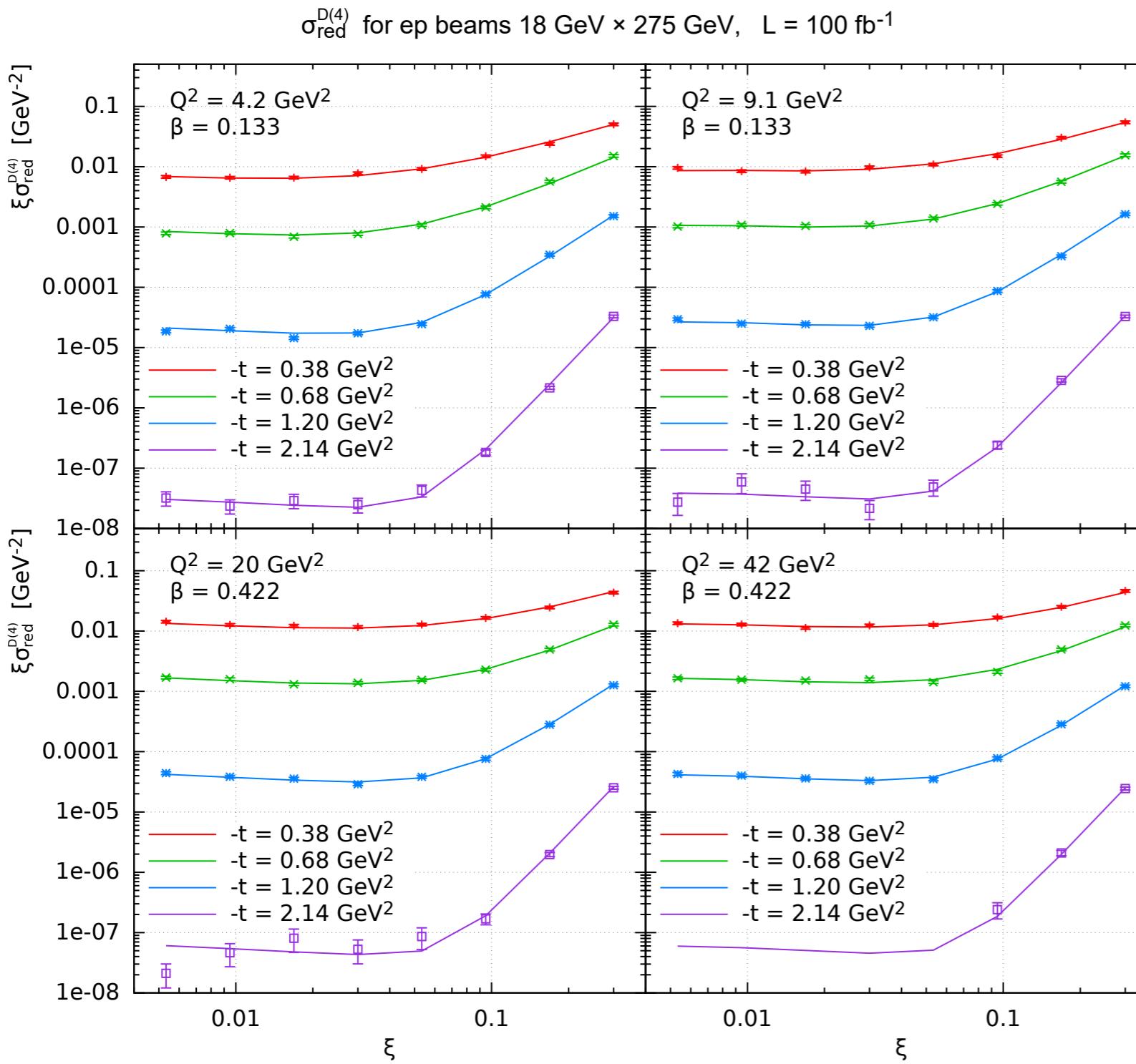
$$E_p = 275 \text{ GeV}$$

Lines-extrapolation

Points-simulation

Prospects for the very good measurement of the t -slope as a function of ξ

$\sigma^D(4)$ vs ξ



$E_e = 18 \text{ GeV}$

$E_p = 275 \text{ GeV}$

Lines-extrapolation

Points-simulation

Prospects for the very good measurement of the t-slope as a function of ξ

Double-slope structure