

# Investigation of Backward-Angle ( $u$ -channel) VCS and DVCS at the EIC

Zachary Sweger  
University of California, Davis



UNIVERSITY  
OF  
CALIFORNIA



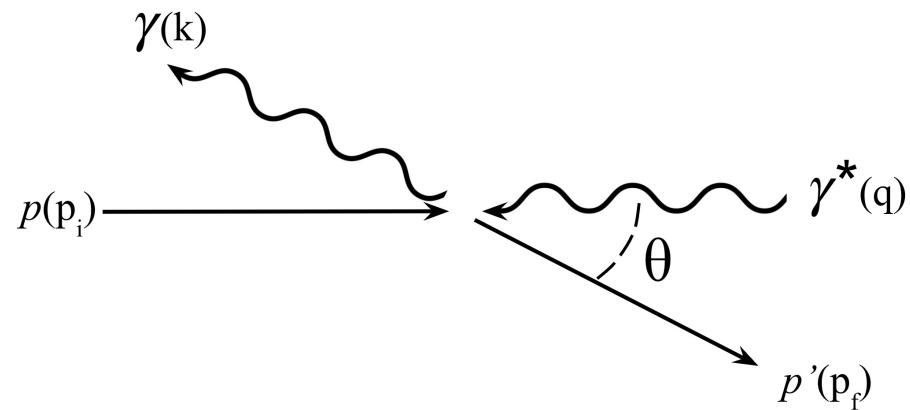
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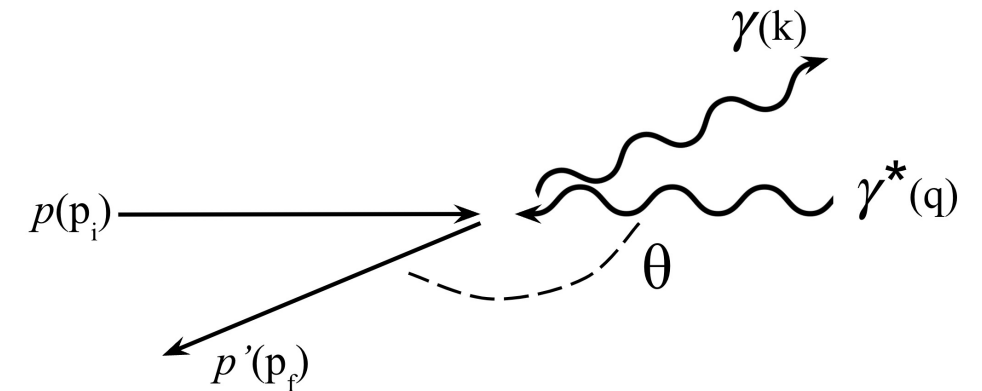
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## Forward Compton Scattering (COM Frame)



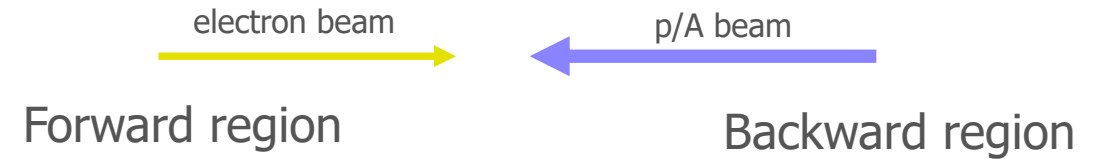
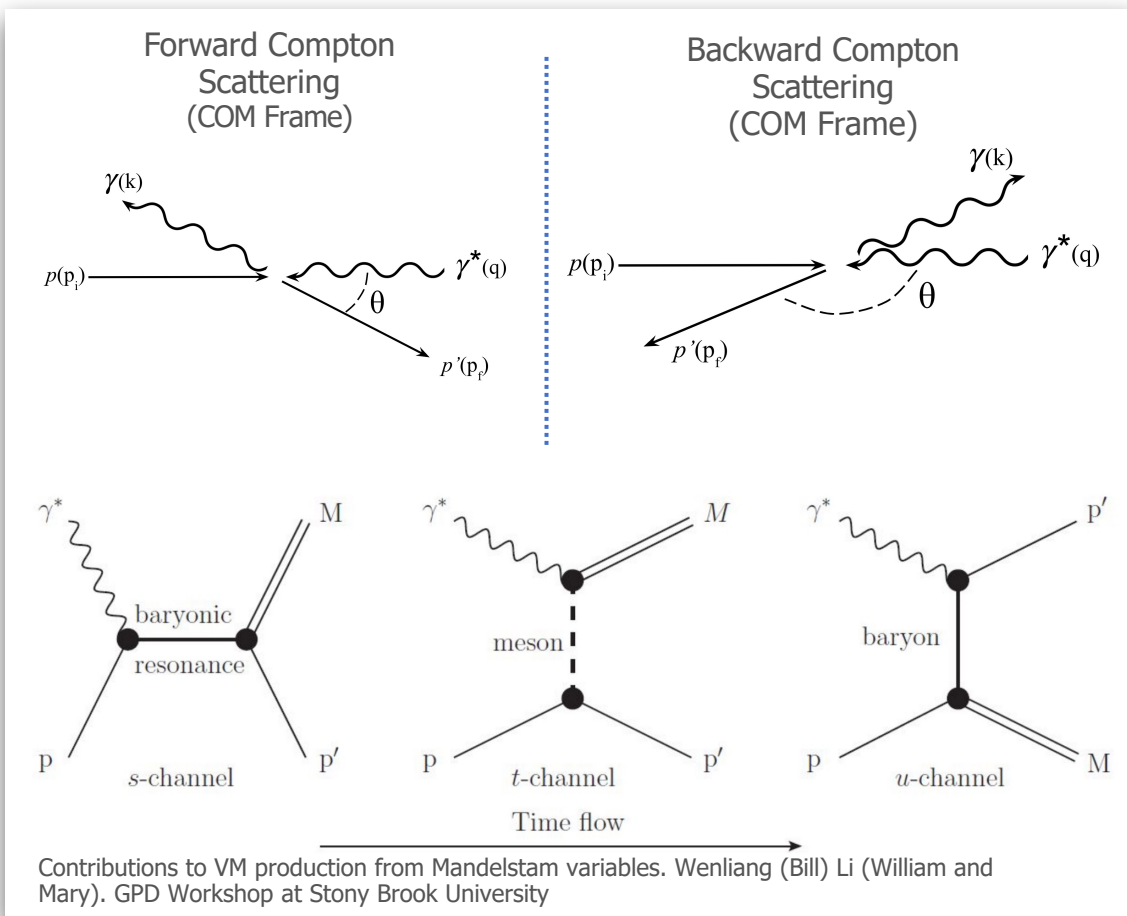
Glancing collision, small momentum transfer

## Backward Compton Scattering (COM Frame)



Backscattering, large momentum transfer

# Backwards ( $u$ -channel) Compton Scattering



## Forward vs Backward DVCS

- Forward Production
  - $t$ -channel: low Mandelstam  $t$ , high  $u$
  - Momentum transfer to target is small
  - $\gamma$  is produced in backwards ( $e^-$ -going) direction
  - Proton in forward direction
    - Proton rapidity only slightly modified
- Backwards Production
  - $u$ -channel: low Mandelstam  $u$ , high  $t$
  - Momentum transfer to target is large
  - $\gamma$  produced in forwards ( $p$ -going) direction
  - Proton shifted in many units in rapidity

- DVCS can be parameterized in terms of

- $Q^2$

- $W = \sqrt{s} = \sqrt{(p + q)^2}$

- $|t| = |(p - p')^2|$

- $|u| = |(p - k)^2|$

- $\theta_{\text{CM}}$

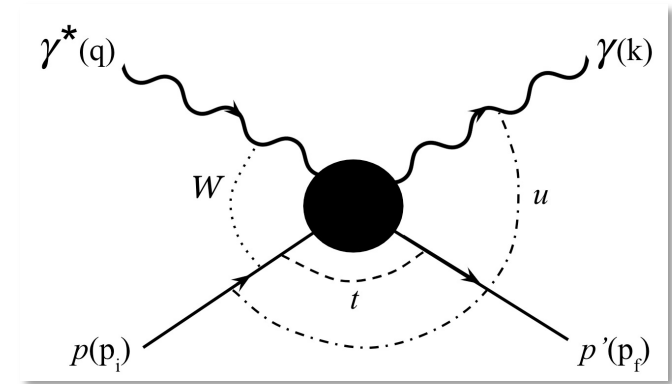
- $\phi$



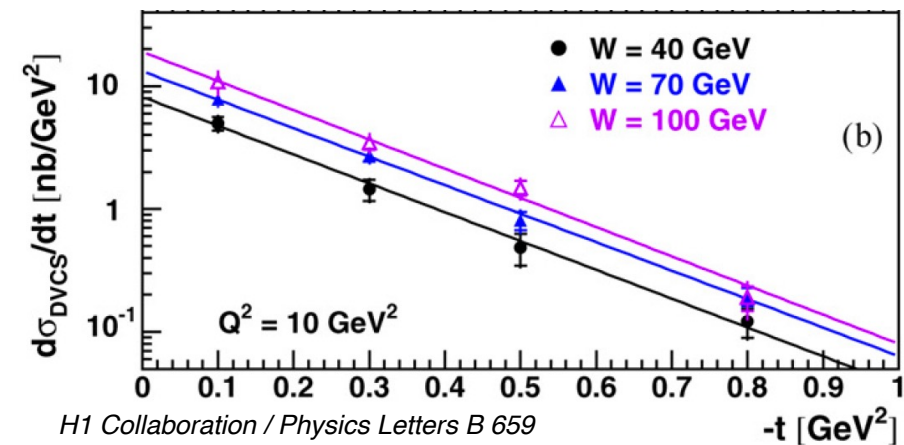
$t$ ,  $u$ , and  $\theta$  all parameterize the momentum transfer in the reaction. Only one is needed in the cross section

$\phi$  describes rotation of  $\gamma p$  plane relative to  $\gamma^* e^-$  plane. This is a polarization observable, but does not affect rapidity distributions that we're studying

$$\frac{d^4\sigma[ep \rightarrow e'p'\gamma]}{dQ^2 dW d\phi dt} = \Gamma(Q^2, W) \frac{d^2\sigma[\gamma^* p \rightarrow p' \gamma]}{d\phi dt}(Q^2, W, \phi, t)$$



# A $u$ -channel Peak?



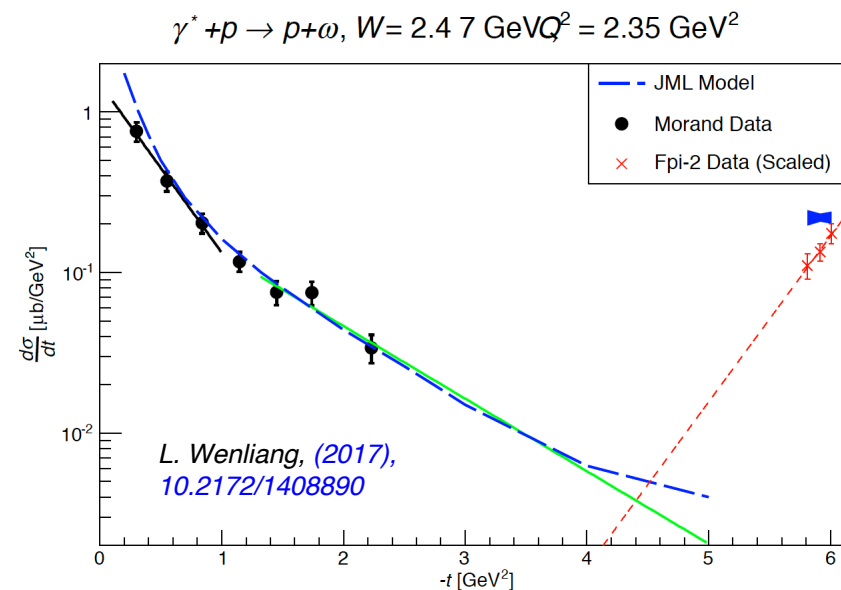
H1 Collaboration / Physics Letters B 659  
(2008) 796–806

## Typical Description of DVCS cross section

- Differential cross section at fixed  $Q^2$  and  $W$  is typically modeled using an exponential of the form  $e^{-b|t|}$
- The Fourier transform of this differential cross section encodes information about the proton GPDs in impact-parameter space
- So why care about cross section at very high  $|t|$ ?

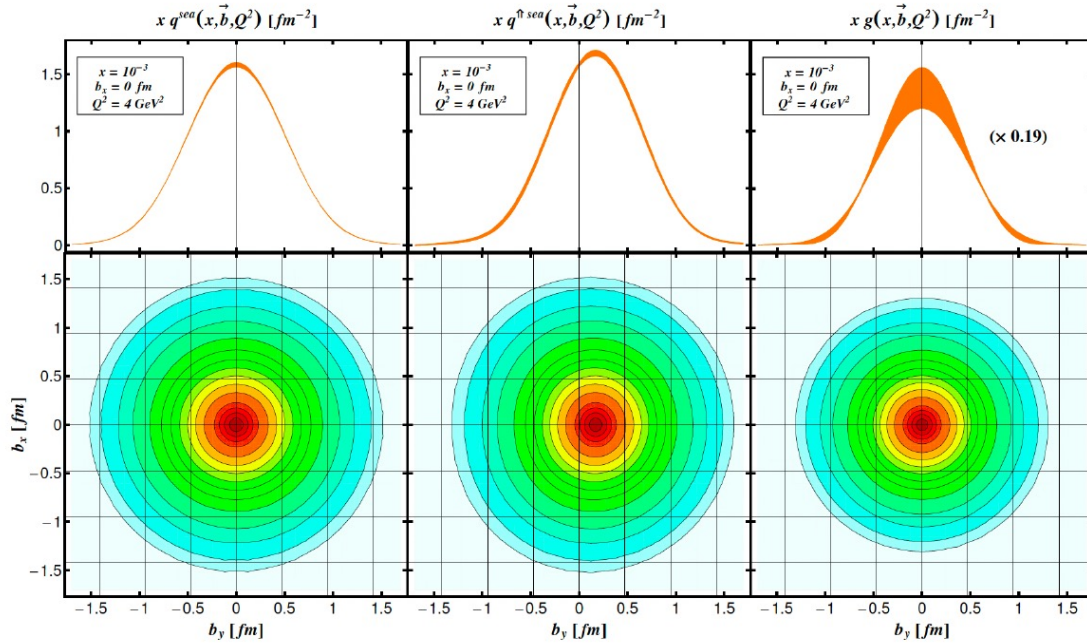
## Non-trivial Behavior at High $t$

- We should start from the assumption that we should not expect photon production (DVCS) cross sections to be wildly different from vector-meson production cross sections (vector-meson dominance)
- Cross sections for vector (and non-vector) mesons see similar exponential drop-off in  $|t|$ , BUT also an exponential rise at the highest  $|t|$  values
- This is from  $u$ -channel contributions which may also be expected in DVCS



# Meaning of $t$ -channel Cross Section

Yellow Report, R. Abdul Khalek et al.,  
[arXiv:2103.05419](https://arxiv.org/abs/2103.05419).

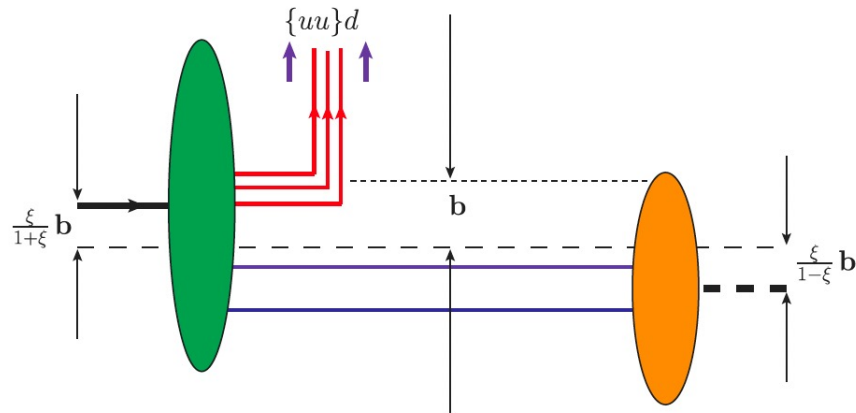


Forward DVCS cross section  $\rightarrow$  proton GPDs

- Differential cross section as a function of  $t$  encodes information about proton GPDs
- GPDs can be translated into an impact-parameter description of the proton via a Fourier transform in  $t$
- Thus the forward DVCS cross section is meaningfully related to the parton structure of the proton

**Figure 7.46:** Impact parameter distributions at  $x = 0.001$  and  $Q^2 = 4 \text{ GeV}^2$  for unpolarized sea quarks in an unpolarized proton (left), a transversely polarized proton (middle), and for unpolarized gluons in an unpolarized proton (right), obtained from a combined fit to the HERA collider data and EIC pseudodata [23]. Top row: IPDs at fixed  $b_x = 0$  as a function of  $b = b_y$ . Bottom row: density plots of IPDs in the  $(b_x, b_y)$ -plane.





$$\text{ERBL : } x_3 = w_3 - \xi \geq 0; \quad x_1 + x_2 = \xi - w_3 \geq 0;$$

B. Pire, K. Semenov-Tian-Shansky, and L. Szymanowski,  
*Phys. Rept.* 940, 1 (2021), [arXiv:2103.01079](https://arxiv.org/abs/2103.01079)  
[hep-ph].

Backward DVCS cross section  $\rightarrow$  partonic correlations and baryon number?

- Recent (2021) work by Pire, Shansky and Szymanowski works to formulate a similarly meaningful interpretation of the backward cross section
- In this work they argue that backward reactions provide access to the location in impact parameter space of di-quark and three-quark (shown at right) clusters
- In backward reactions the baryon number follows these clusters to form a “new” baryon

“**baryon-to-meson (and baryon-to-photon) TDAs** share common features both with baryon DAs and with GPDs and encode a conceptually close physical picture. They **characterize partonic correlations inside a baryon and give access to the momentum distribution of the baryonic number inside a baryon**. Similarly to GPDs, TDAs – after the Fourier transform in the transverse plane – represent valuable information on the transverse location of hadron constituents.”

## Modeling $u$ -channel DVCS

- We presuppose a peak at backward angles ( $u=u_0$ ) as is seen in meson production
- EIC will provide an opportunity to measure this peak if it exists, a task that is challenging in fixed-target experiments due to the softness of the photons produced
- **The strategy: exploit similarities to  $t$ -channel**

$$\frac{d\sigma}{dt}(t) \sim \exp(-B|t - t_0|) \longrightarrow \frac{d\sigma}{du}(u) \sim \exp(-D|u - u_0|)$$

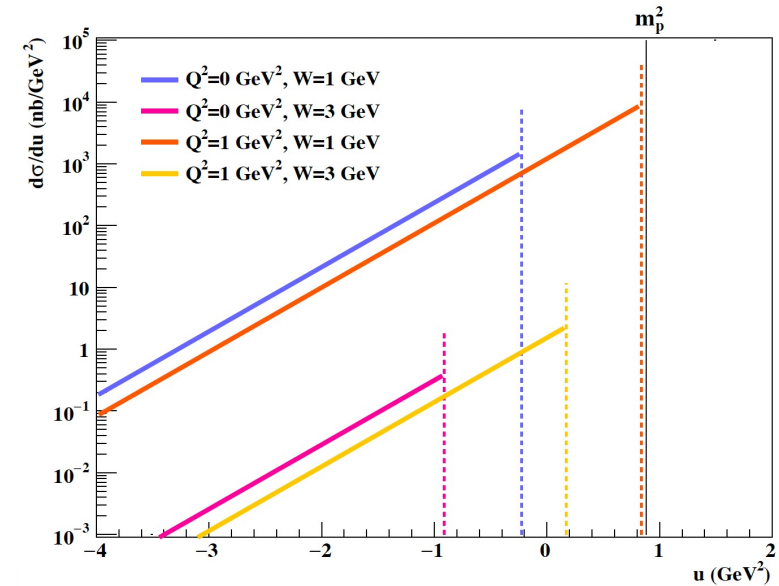
- $B$  and  $D$  are related to the size of production region which differs in  $t$  and  $u$  channels due to role of meson vs baryon exchange trajectories
- $D$  has not been measured for backward DVCS, so for our models we test values measured for backward  $\omega$

production

L. Wenliang, (2017), [10.2172/1408890](#).

D. Cebra, Z. Sweger, X. Dong, Y. Ji, and S. R. Klein, *Phys. Rev. C* **106**, 015204 (2022).

Exclusive/Diffractive/Tagging Meeting





## Modeling $W$ -Dependence

- Backward physics is dominated by Regge-exchange trajectories for which the cross sections typically scale with  $W^{-\alpha}$  where  $\alpha > 0$
- In our backward  $\omega/\rho$  paper, we used a data-driven  $(W^2 - m_p^2)^{-2.4}$  dependence

*D. Cebra, Z. Sweger, X. Dong, Y. Ji, and S. R. Klein,  
[Phys. Rev. C 106, 015204 \(2022\)](#).*

- Several sources suggest rough  $(W^2 - m_p^2)^{-2}$  scaling which is what we start from.

*G. Laveissière et al., [Physical Review C 79 \(2009\), 10.1103/physrevc.79.015201](#).*

*S. J. Brodsky, F. J. Llanes-Estrada, and A. P. Szczepaniak,  
[Phys. Rev. D 79, 033012 \(2009\)](#).*

*W. B. Li et al. (Jefferson Lab FTT Collaboration),  
[Phys. Rev. Lett. 123, 182501 \(2019\)](#).*

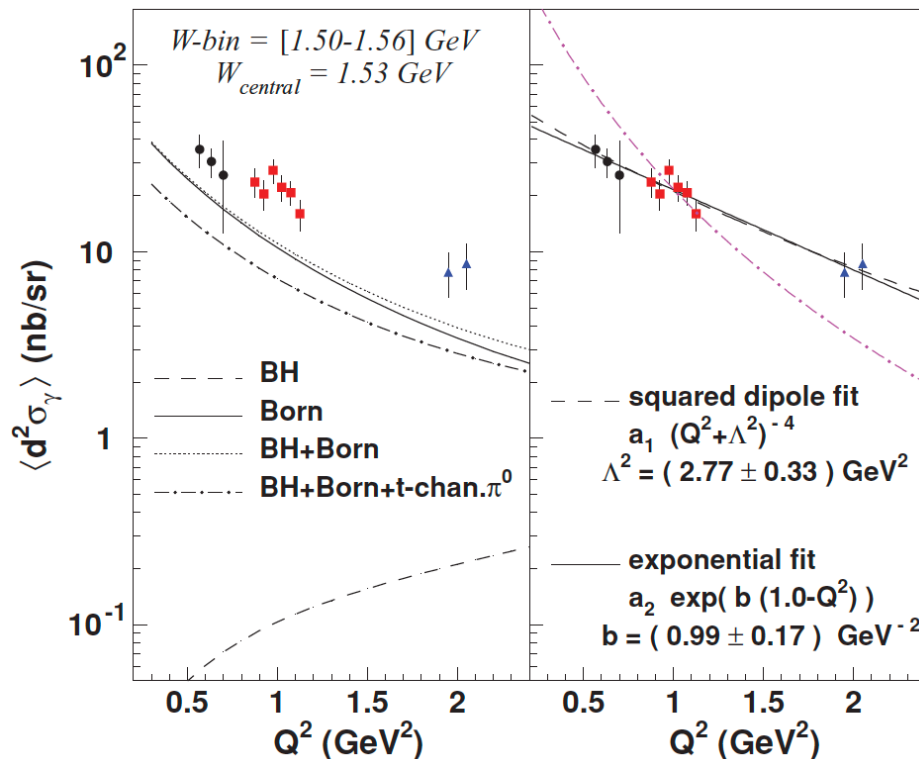
$$\frac{d\sigma}{du}(W, u) \sim \frac{1}{(W^2 - m_p^2)^2} \exp(-D|u - u_0|)$$

# Our Backward DVCS Model: $Q^2$ -Dependence

## Backward VCS in Resonance Region

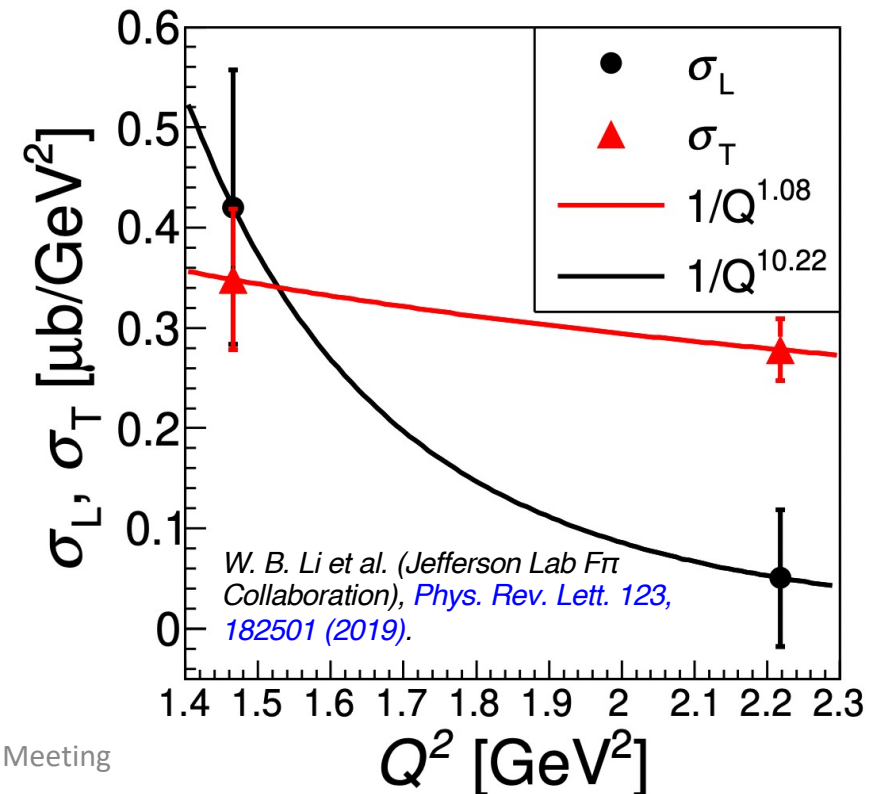
- There is some limited data available for this
- For backward VCS in the resonance region, JLab measured  $(Q^2 + 2.77 \text{ GeV}^2)^{-4}$  dependence

G. Laveissière et al., *Physical Review C* 79 (2009),  
10.1103/physrevc.79.015201.



## Backward $\omega$ Production Above Resonance

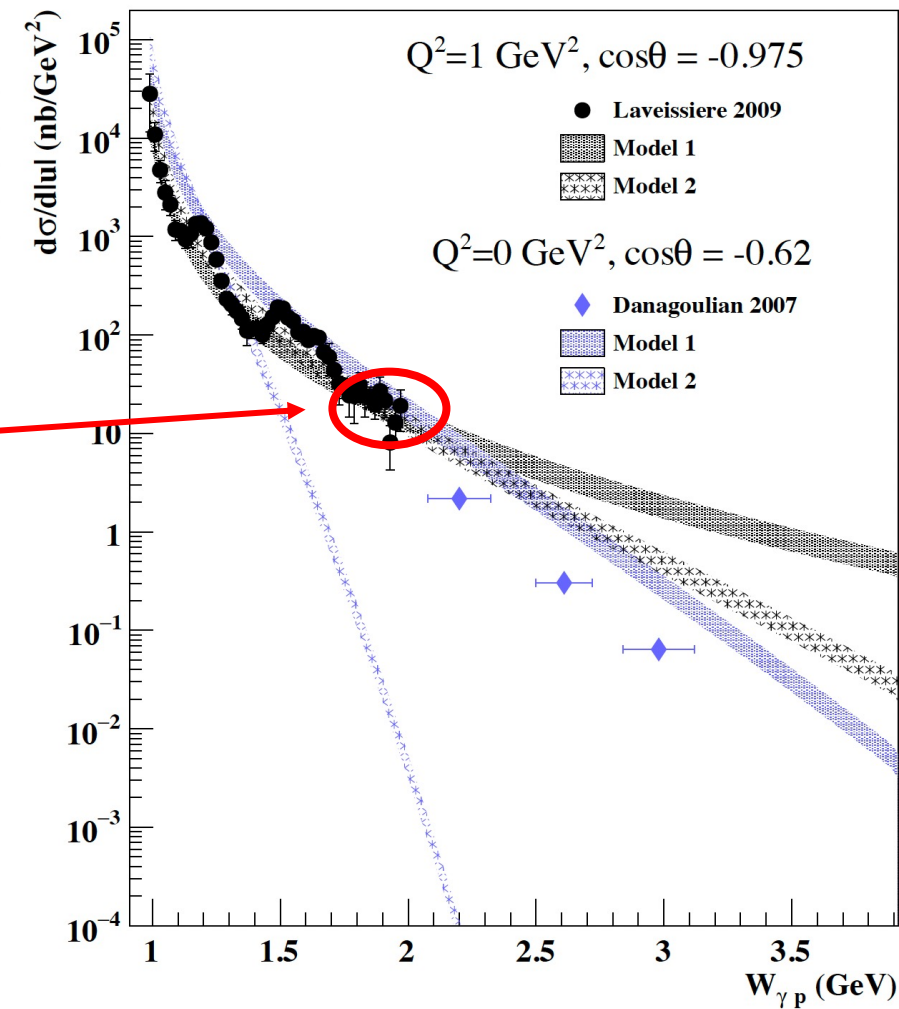
- Polarization-dependent cross section
- $Q^2$ -dependence is much softer for transversely polarized photons.
- Needs to be explored further in our simulations



- We combine these contributions to yield the form:

$$\frac{d\sigma}{du}(Q^2, W, u) \approx \frac{A \exp(-D|u - u_0|)}{(W^2 - m_p^2)^2 (Q^2 + \Lambda^2)^4 / \text{GeV}^8}$$

- In order to anchor the amplitude, we can fit this model to 11 VCS ( $Q^2=1 \text{ GeV}^2$ ) data points from JLab from  $1.77 < W < 1.96 \text{ GeV}$  (above strong resonances)
- Where
  - $\Lambda^2 = 2.77 \text{ GeV}^2$
  - Model 1:  $D = 2.4 \text{ GeV}^{-2}$ ,  $A = 32 \mu\text{b/GeV}^2$
  - Model 2:  $D = 21.8 \text{ GeV}^{-2}$ ,  $A = 65 \mu\text{b/GeV}$



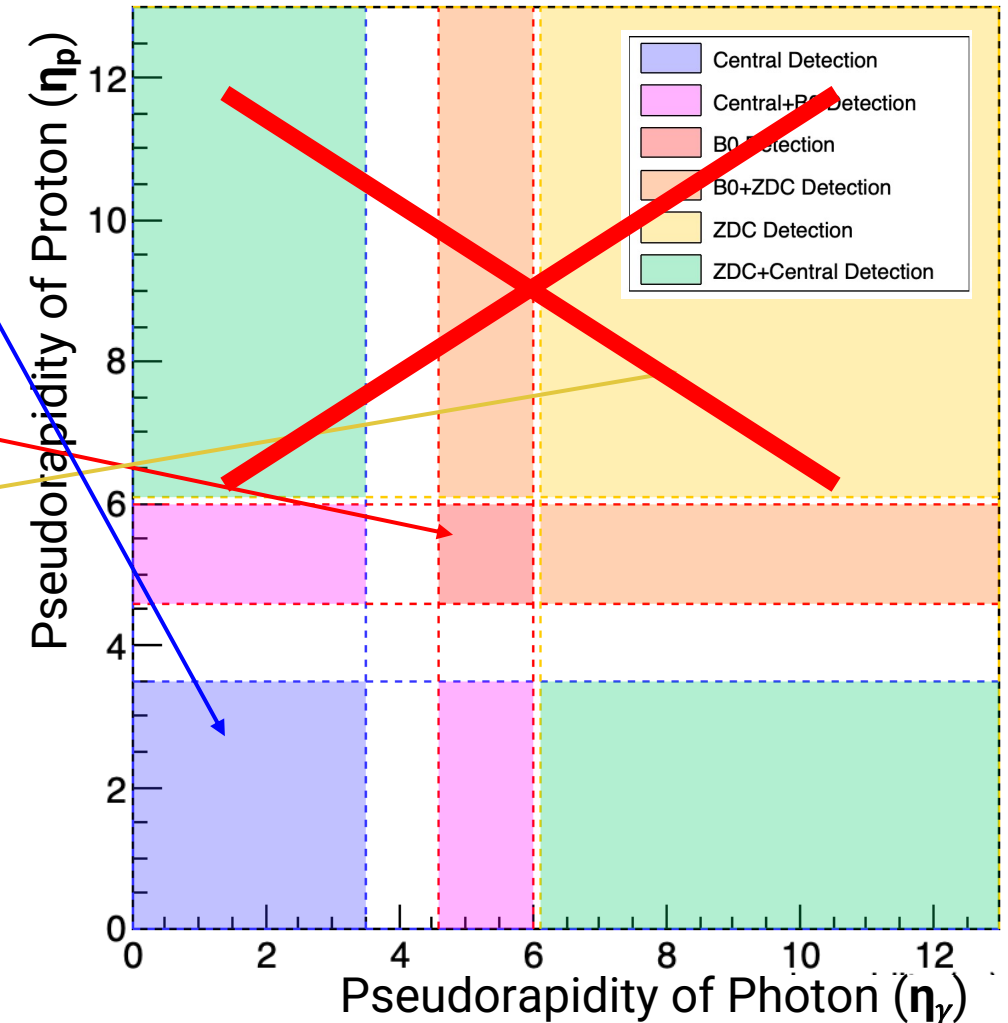
G. Laveissiere et al., *Physical Review C* 79 (2009),  
[10.1103/physrevc.79.015201](https://doi.org/10.1103/physrevc.79.015201).

A. Danagouliau et al. (Jefferson Lab Hall A Collaboration),  
*Phys. Rev. Lett.* 98, 152001 (2007)

There are three detector regions of interest for backwards production

- **Central Region (endcap & barrel):  $|\eta| < 3.5$** 
  - ✓ Charged-particle tracking ( $\rho$ )
  - ✓ Electromagnetic calorimetry ( $\omega$ )
- **B0 Magnets:  $4.6 < \eta < 6.0$** 
  - ✓ Charged-particle tracking ( $\rho$ )
  - ? Electromagnetic calorimetry ( $\omega$ )
- **ZDC:  $\eta > 6.215$ – $5.991$** 
  - ✓ Electromagnetic calorimetry ( $\omega$ )

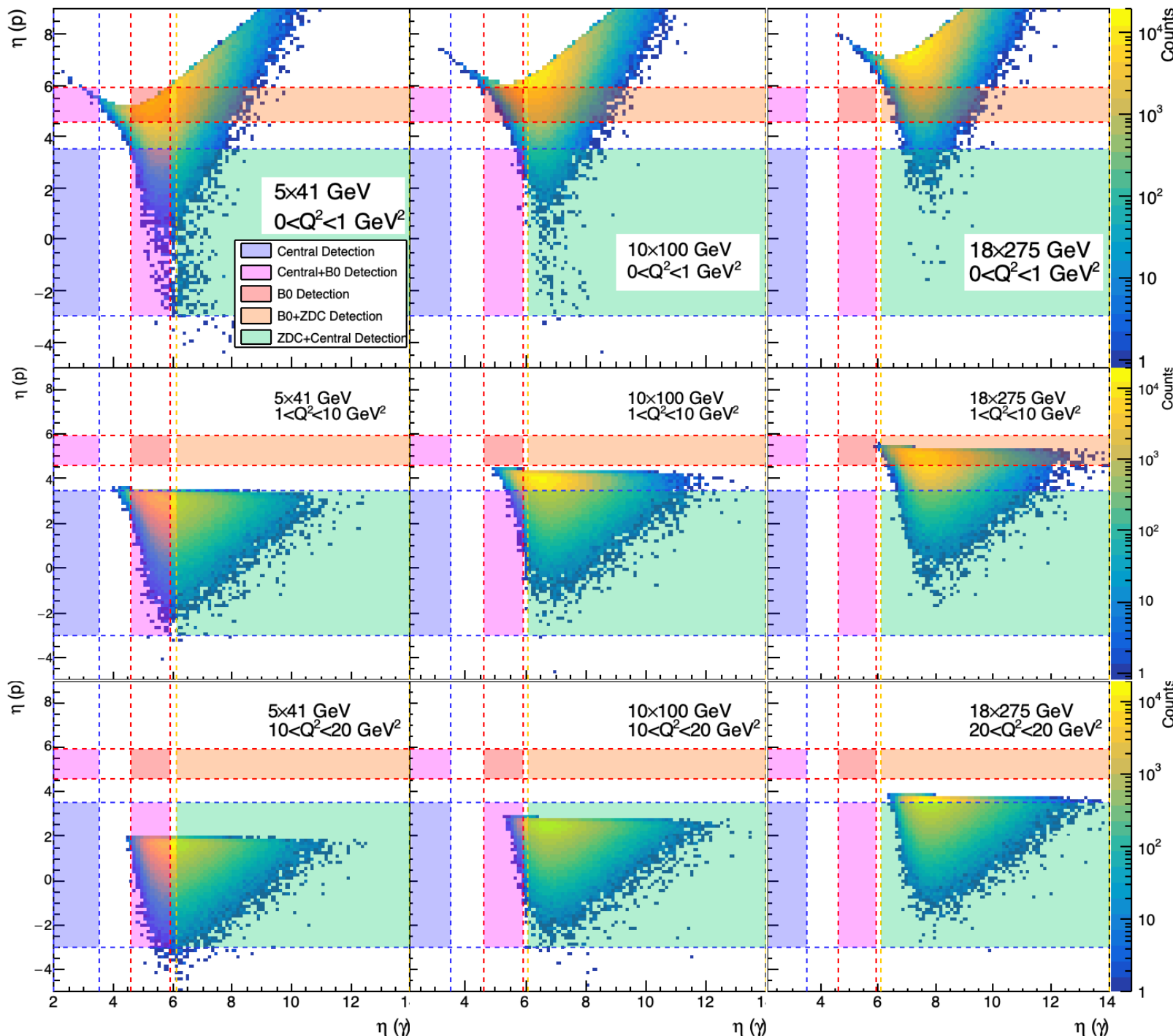
Proton may be detected in central or B0 detectors but not the ZDC. Photons may be detected in central, B0 or ZDC



# Backward DVCS Acceptances

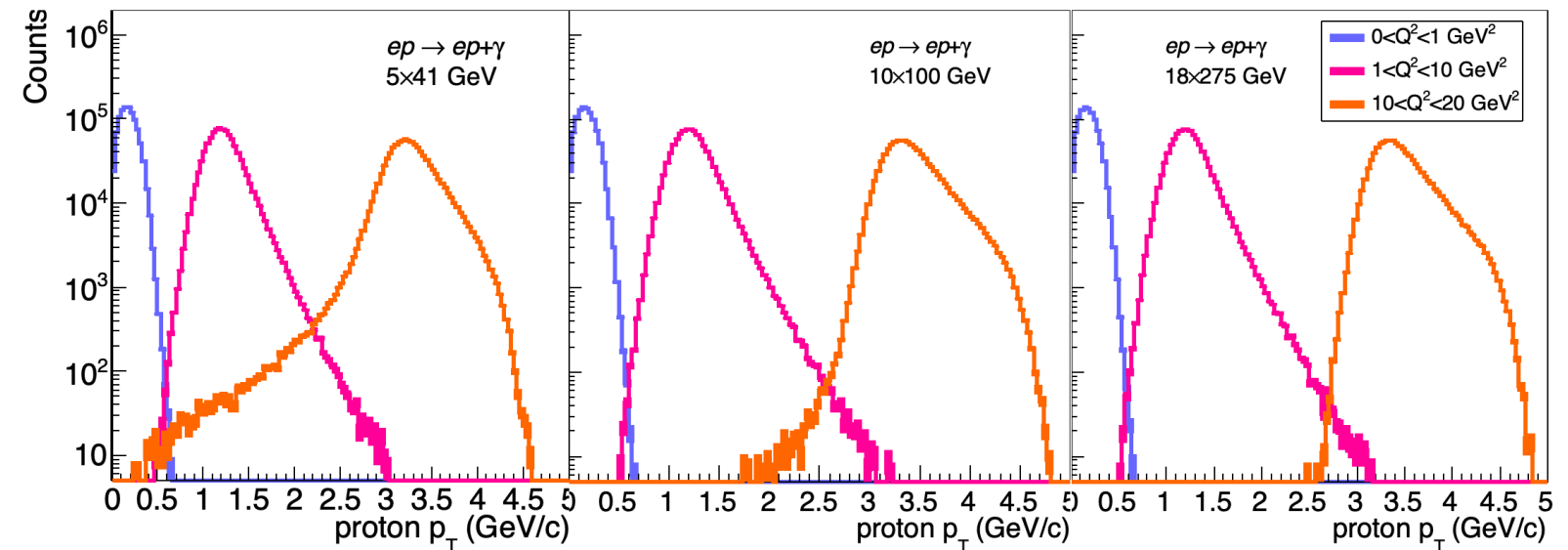
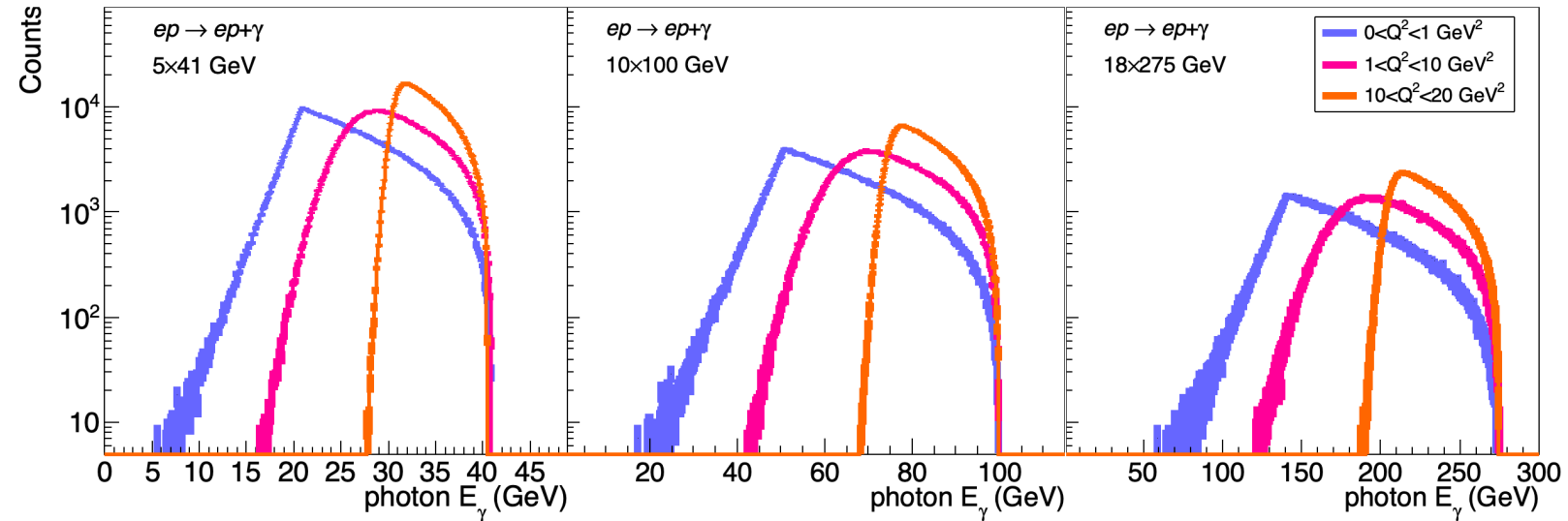
- These simulations used Model 2 for  $W > 1.3$  GeV
- At low collision energies, the photon will be seen in the B0 and ZDC
- At high energies, the ZDC is critical
- At very low  $Q^2$ , the proton will be seen mostly in the B0
- At high  $Q^2$ , the proton lands almost exclusively in the central detector region
- Geometric efficiency estimates:

Proton beam energy	$p + \gamma$ eff. $Q^2 < 1$	$p + \gamma$ eff. $1 < Q^2 < 10$	$p + \gamma$ eff. $10 < Q^2 < 20$
41 GeV w/o B0	12%	35%	44%
41 GeV w/ B0	62%	88%	86%
100 GeV w/o B0	23%	30%	97%
100 GeV w/ B0	29%	30%	97%
275 GeV w/o B0	46%	68%	68%
275 GeV w/ B0	46%	68%	68%



# Kinematics of Final-State Particles

- Final-state photons in the B0 and ZDC will be between 10 and 275 GeV
- Protons from low- $Q^2$  events will have low  $p_T$
- Moderate  $p_T$  for high- $Q^2$  events will aid detection but the potentially rapid drop-off of the cross section with  $Q^2$  may prevent this





## Conclusions and Outlook

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- Backward production is interesting, in its own right, for encoding unique information about proton GPDs, an active and evolving topic of research
- Measurements at the EIC can improve our understanding of these  $u$ -channel mechanisms as well as their contribution to the  $t$ -channel background near-threshold
- Our backward DVCS model is still under development
- Early simulations demonstrate importance of B0 and ZDC calorimetry especially for high-energy photons.
- Next step are to
  - Finalize and compare models
  - Generate event samples for simulations
  - Write up results

Thank you for your attention!

[zwsweager@ucdavis.edu](mailto:zwsweager@ucdavis.edu)