

# *Impact of finite magnetic field and volume on the susceptibilities of conserved charges*

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under the supervision of

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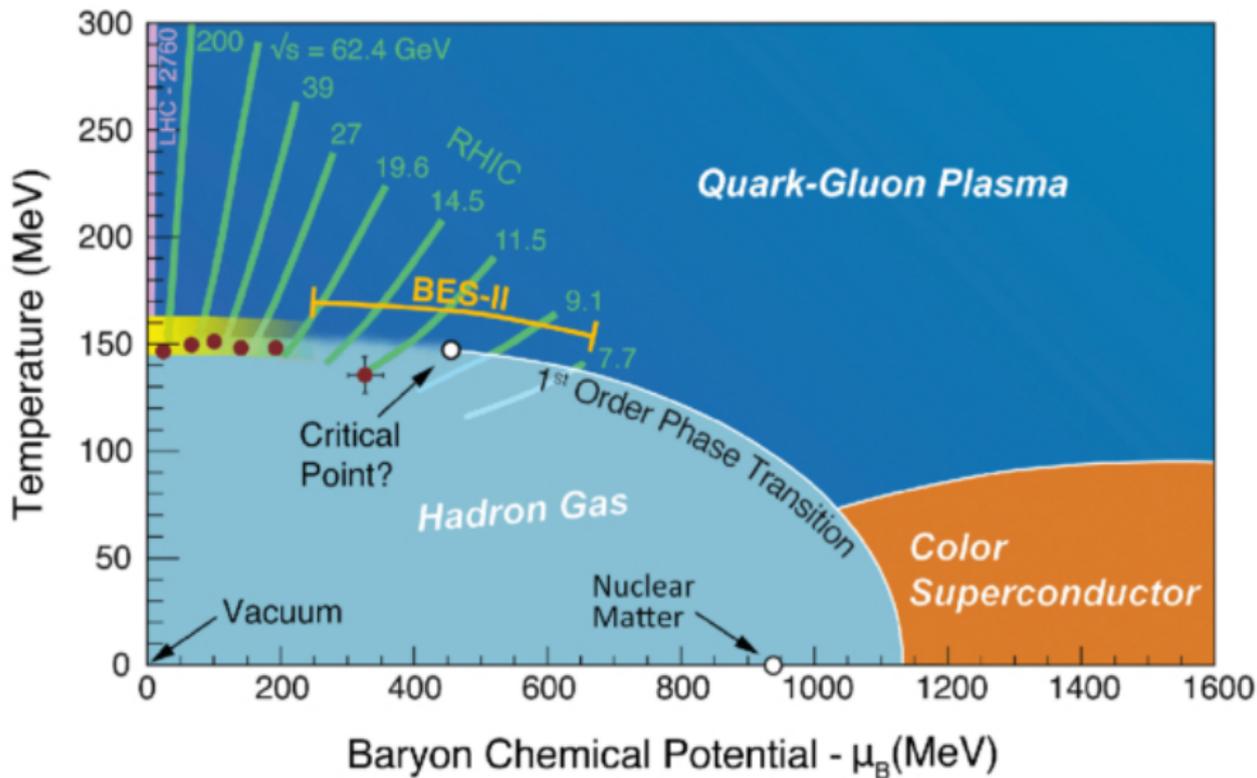
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# Outline

- 1 *Introduction*
- 2 *Methodology*
- 3 *Magnetic field and volume effects*
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# QCD Phase Diagram



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# Chiral $SU(3)$ Quark Mean Field Model

Lagrangian density

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{q0} + \mathcal{L}_{qm} + \mathcal{L}_{\Sigma\Sigma} + \mathcal{L}_{VV} + \mathcal{L}_{SB} + \mathcal{L}_{\Delta m} + \mathcal{L}_h. \quad (1)$$

- $\mathcal{L}_{q0}$  is the free part of massless quarks.
- $\mathcal{L}_{qm}$  quark meson interaction term.
- $\mathcal{L}_{\Sigma\Sigma}$  scalar meson self-interaction term ( $\sigma, \zeta, \chi$  and  $\delta$  fields).
- $\mathcal{L}_{VV}$  vector meson self-interaction term ( $\omega, \rho$  and  $\phi$  fields).
- $\mathcal{L}_{SB}, \mathcal{L}_{\Delta m}$  and  $\mathcal{L}_h$  are explicit symmetry breaking terms.

# Chiral $SU(3)$ Quark Mean Field Model

Thermodynamical potential density

$$\Omega = \sum_{i=u,d,s} \frac{-2k_B T \gamma_i}{(2\pi)^3} \int_0^\infty d^3 k [\ln(1 + e^{-(E_i^*(k) - \nu_i)/k_B T}) + \ln(1 + e^{-(E_i^*(k) + \nu_i)/k_B T})] - \mathcal{L}_M, \quad (2)$$

- $\mathcal{L}_M = \mathcal{L}_{\Sigma\Sigma} + \mathcal{L}_{VV} + \mathcal{L}_{SB}$ .
- $E_i^*(k) = \sqrt{m_i^{*2} + k^2}$  is the effective single particle energy of quarks.
- $m_i^* = -g_\sigma^i \sigma - g_\zeta^i \zeta - g_\delta^i \delta + m_{i0}$  is effective constituent quark mass.
- $\nu_i^* = \mu_i - g_\omega^i \omega - g_\phi^i \phi - g_\rho^i \rho$  is effective chemical potential.
- $\gamma_i$  is spin degeneracy factor for quarks ( $\gamma_i=3$ ) and electrons ( $\gamma_i=1$ ).

# *Polyakov Chiral SU(3) quark mean field model*

Polyakov loop

$$\Phi(\tilde{x}) = (\text{Tr}_c L)/N_C, \quad (3)$$

and its conjugate

$$\bar{\Phi}(\tilde{x}) = (\text{Tr}_c L^\dagger)/N_C. \quad (4)$$

Total lagrangian density

$$\mathcal{L}_{\text{PCQMF}} = \mathcal{L}_{\text{eff}} -$$

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Total lagrangian density

$$\mathcal{L}_{\text{PCQMF}} = \mathcal{L}_{\text{eff}} - \mathcal{U}(\Phi(\tilde{x}), \bar{\Phi}(\tilde{x}), \mathbf{T}), \quad (5)$$

Modified thermodynamical potential density

$$\begin{aligned} \Omega_{\text{PCQMF}} = & -2k_B T \sum_{u,d,s} \int_0^\infty \frac{d^3 k}{(2\pi)^3} [\ln(1 + e^{-3(E_i^*(\mathbf{k}) - \nu_i)/k_B T} \\ & + 3\Phi e^{-(E_i^*(\mathbf{k}) - \nu_i)/k_B T} + 3\bar{\Phi} e^{-2(E_i^*(\mathbf{k}) - \nu_i)/k_B T} + \ln(1 + e^{-3(E_i^*(\mathbf{k}) + \nu_i)/k_B T} \\ & + 3\bar{\Phi} e^{-(E_i^*(\mathbf{k}) + \nu_i)/k_B T} + 3\Phi e^{-2(E_i^*(\mathbf{k}) + \nu_i)/k_B T})] + \mathcal{U}(\Phi, \bar{\Phi}, \mathbf{T}), \quad (6) \end{aligned}$$

# Polyakov Chiral $SU(3)$ quark mean field model

here,  $\mathcal{U}(\Phi(\tilde{x}), \bar{\Phi}(\tilde{x}), T)$  is temperature dependent Polyakov loop effective potential,

$$\frac{\mathcal{U}(\Phi, \bar{\Phi}, T)}{T^4} = -\frac{a(T)}{2}\bar{\Phi}\Phi + b(T)\ln[1 - 6\bar{\Phi}\Phi + 4(\bar{\Phi}^3 + \Phi^3) - 3(\bar{\Phi}\Phi)^2], \quad (7)$$

with T-dependent parameters:

$$a(T) = a_0 + a_1\left(\frac{T_0}{T}\right) + a_2\left(\frac{T_0}{T}\right)^2, \quad b(T) = b_3\left(\frac{T_0}{T}\right)^3. \quad (8)$$

$a_0$	$a_1$	$a_2$	$b_3$
3.51	-2.47	15.2	-1.75

*Table 1:* Parameters in Polyakov effective potential

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# Magnetic field and volume effects

- The total thermodynamical potential is altered and the term giving the contribution of quarks and antiquarks interaction is written as

$$\Omega_{q\bar{q}} = - \sum_{i=u,d,s} \frac{|q_i| BT}{2\pi} \sum_{k=0}^{\infty} \alpha_k \int_{-\infty}^{\infty} \frac{dp_z}{2\pi} (\ln g_i^+ + \ln g_i^-). \quad (9)$$

- Total effective energy of the quarks is modified as

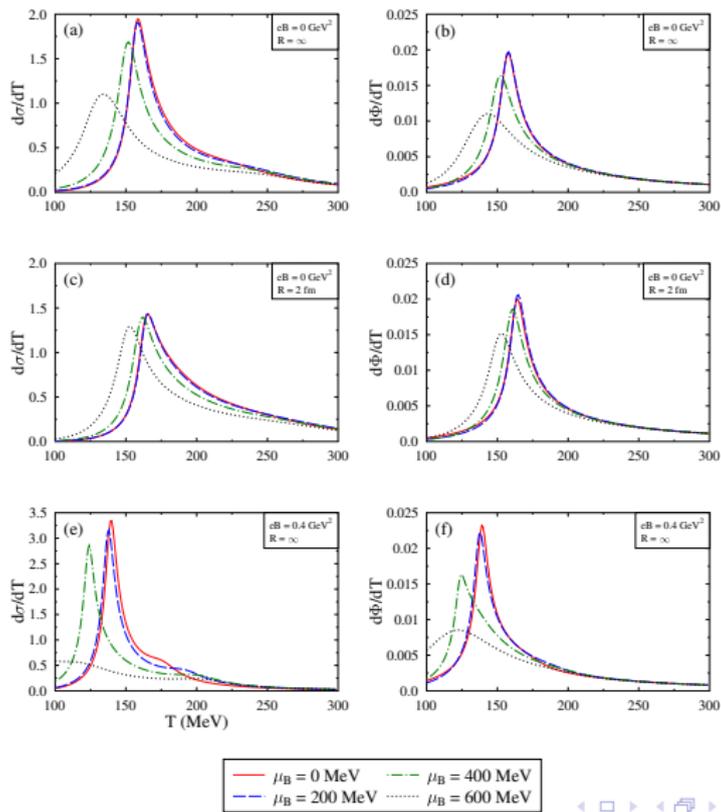
$$E_i^* = \sqrt{p_z^2 + m_i^{*2} + |q_i| (2n + 1 - \Upsilon) B}, \quad (10)$$

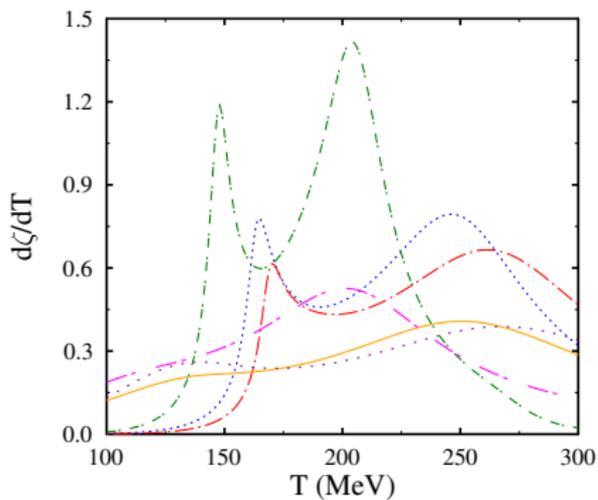
- The impact of finite size effect is assimilated in the model by using the approximation method defined in by introducing a lower momentum cut-off,  $p_{min}$  [MeV] =  $\pi/R$  [MeV] =  $\Lambda$ , where  $R$  is the length of a cubic volume.

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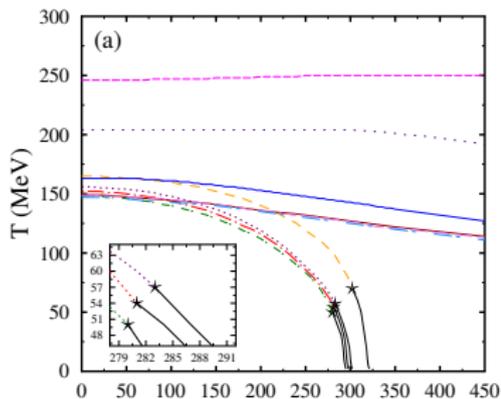
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# Results

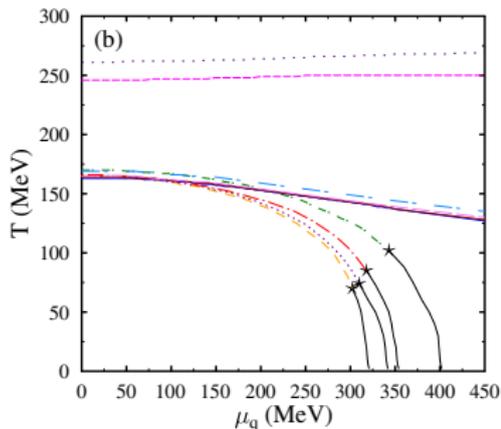




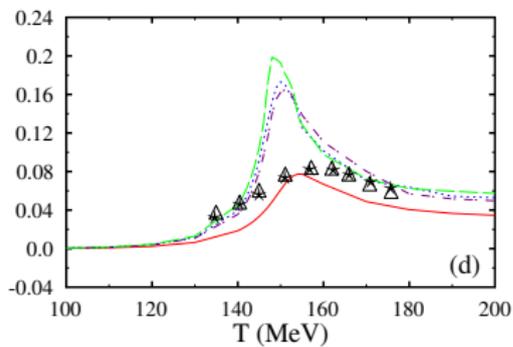
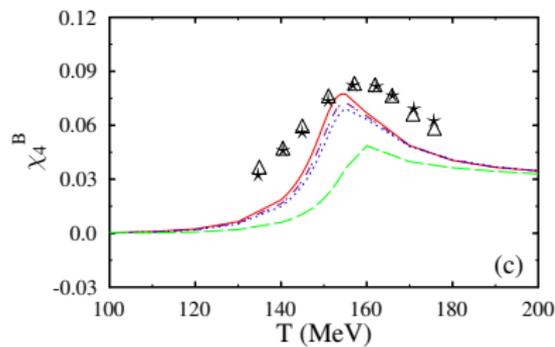
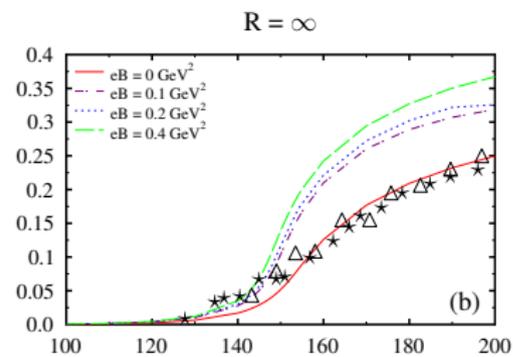
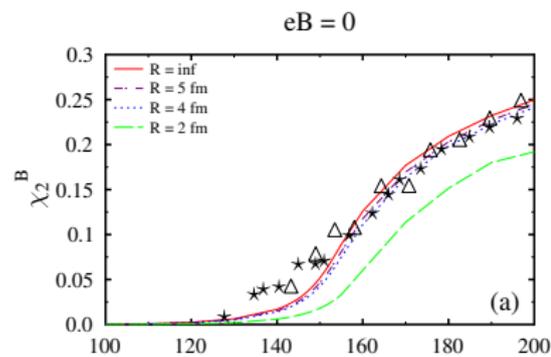
- ⋯  $\mu_q=0, eB = 0 \text{ GeV}^2$  and  $R = \infty$
- -  $\mu_q=0, eB = 0 \text{ GeV}^2$  and  $R = 2 \text{ fm}$
- -  $\mu_q=0, eB = 0.4 \text{ GeV}^2$  and  $R = \infty$
- $\mu_q=350 \text{ MeV}, eB = 0 \text{ GeV}^2$  and  $R = \infty$
- ⋯  $\mu_q=350 \text{ MeV}, eB = 0 \text{ GeV}^2$  and  $R = 2 \text{ fm}$
- -  $\mu_q=350 \text{ MeV}, eB = 0.4 \text{ GeV}^2$  and  $R = \infty$



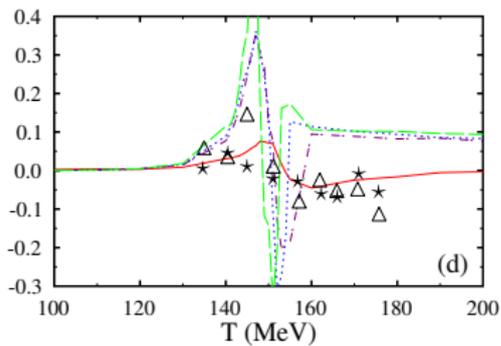
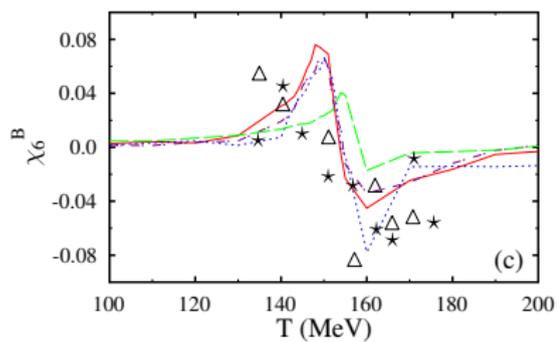
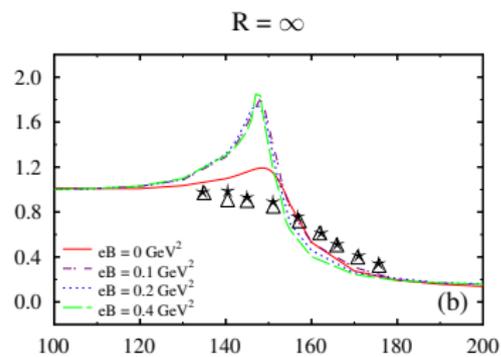
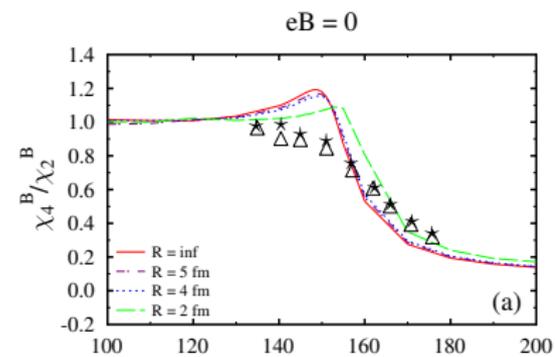
- crossover  $eB = 0 \text{ GeV}^2$
- ⋯ crossover  $eB = 0.1 \text{ GeV}^2$
- - - crossover  $eB = 0.2 \text{ GeV}^2$
- · - crossover  $eB = 0.4 \text{ GeV}^2$
- deconfinement  $eB = 0 \text{ GeV}^2$
- deconfinement  $eB = 0.1 \text{ GeV}^2$
- deconfinement  $eB = 0.2 \text{ GeV}^2$
- deconfinement  $eB = 0.4 \text{ GeV}^2$
- ★ CEP
- 1st order
- s quark transition  $eB=0$
- ⋯ s quark transition  $eB = 0.4 \text{ GeV}^2$



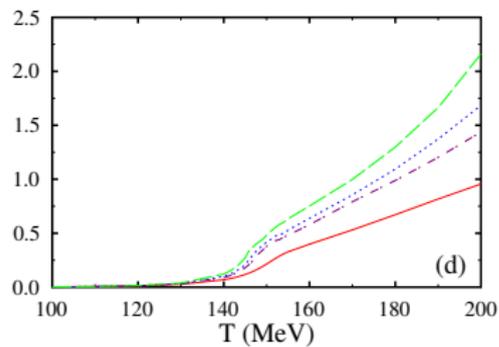
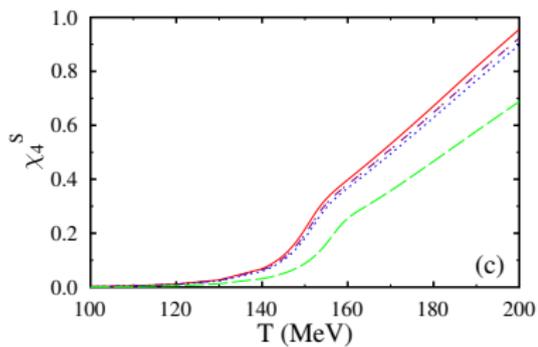
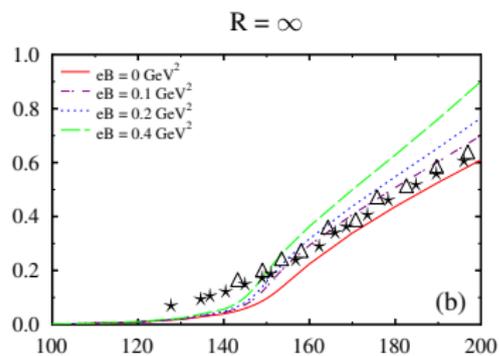
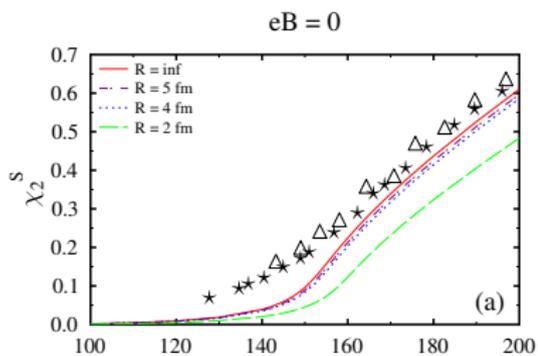
- crossover  $R = \infty$
- ⋯ crossover  $R = 5 \text{ fm}$
- - - crossover  $R = 4 \text{ fm}$
- · - crossover  $R = 2 \text{ fm}$
- deconfinement  $R = \infty$
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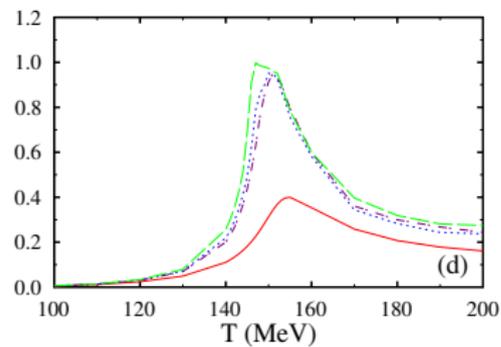
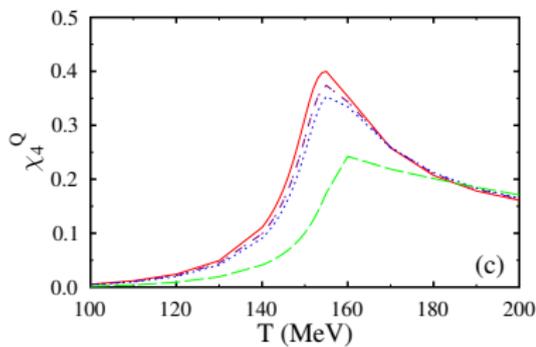
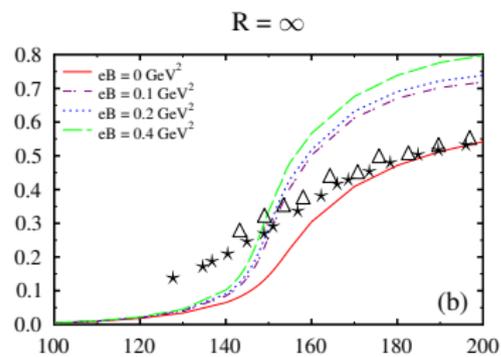
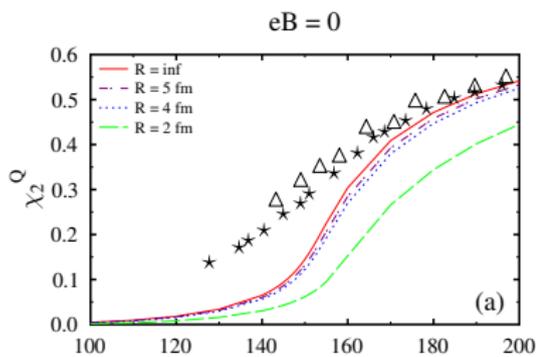
★ lattice data for ( $N_\tau = 8$ )  
 △ lattice data for ( $N_\tau = 12$ )



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# Summary

- We have analyzed the impact of finite volume and external magnetic field on the thermodynamic properties using Polyakov loop extended chiral SU(3) quark mean field model in the asymmetric quark matter.
- The impact of external magnetic field and finite system size on the phase diagram of QCD have been investigated by inspecting the variation of scalar and vector fields.
- Susceptibilities of conserved charges are found to be enhanced in the regime of critical-point.
- These fluctuations can be deduced from event-by-event inspection of the experimental data and hence play significant role in determination of CEP.

*Thank you!*