

Chasing Hagedorn's Dream

Mapping the QCD Phase Diagram at RHIC

Berndt Mueller

Brookhaven National Laboratory
& Duke University

USQCD Meeting

BNL

29 April 2016

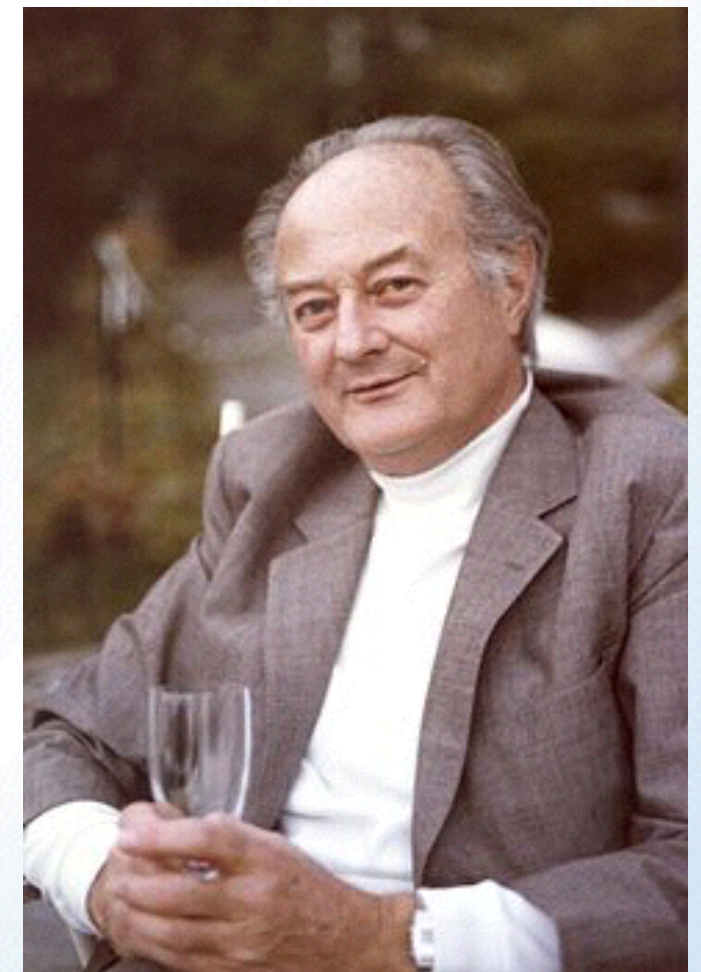
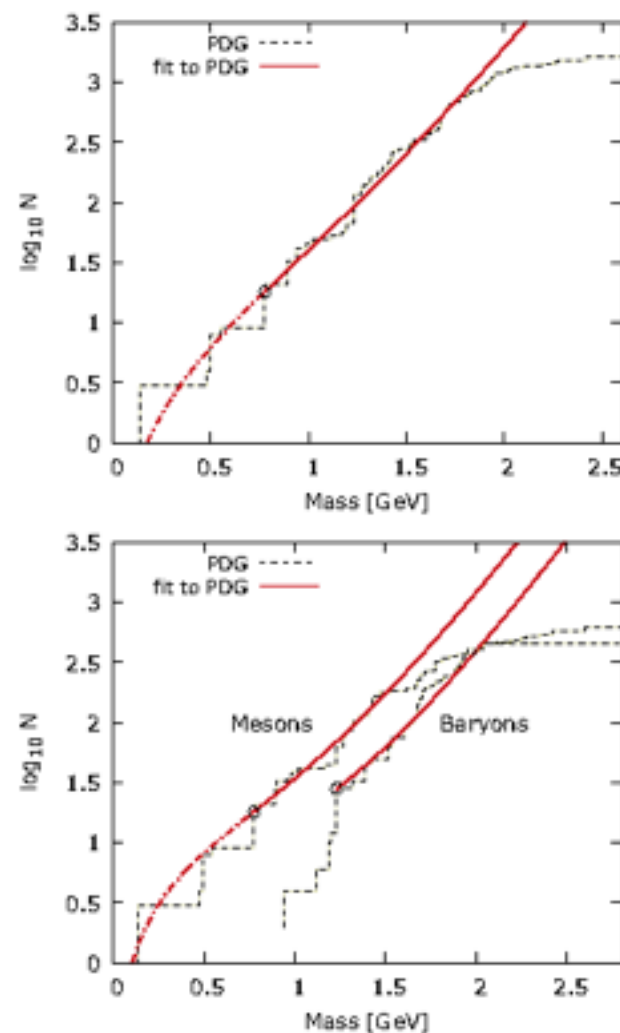


1965

was a momentous year



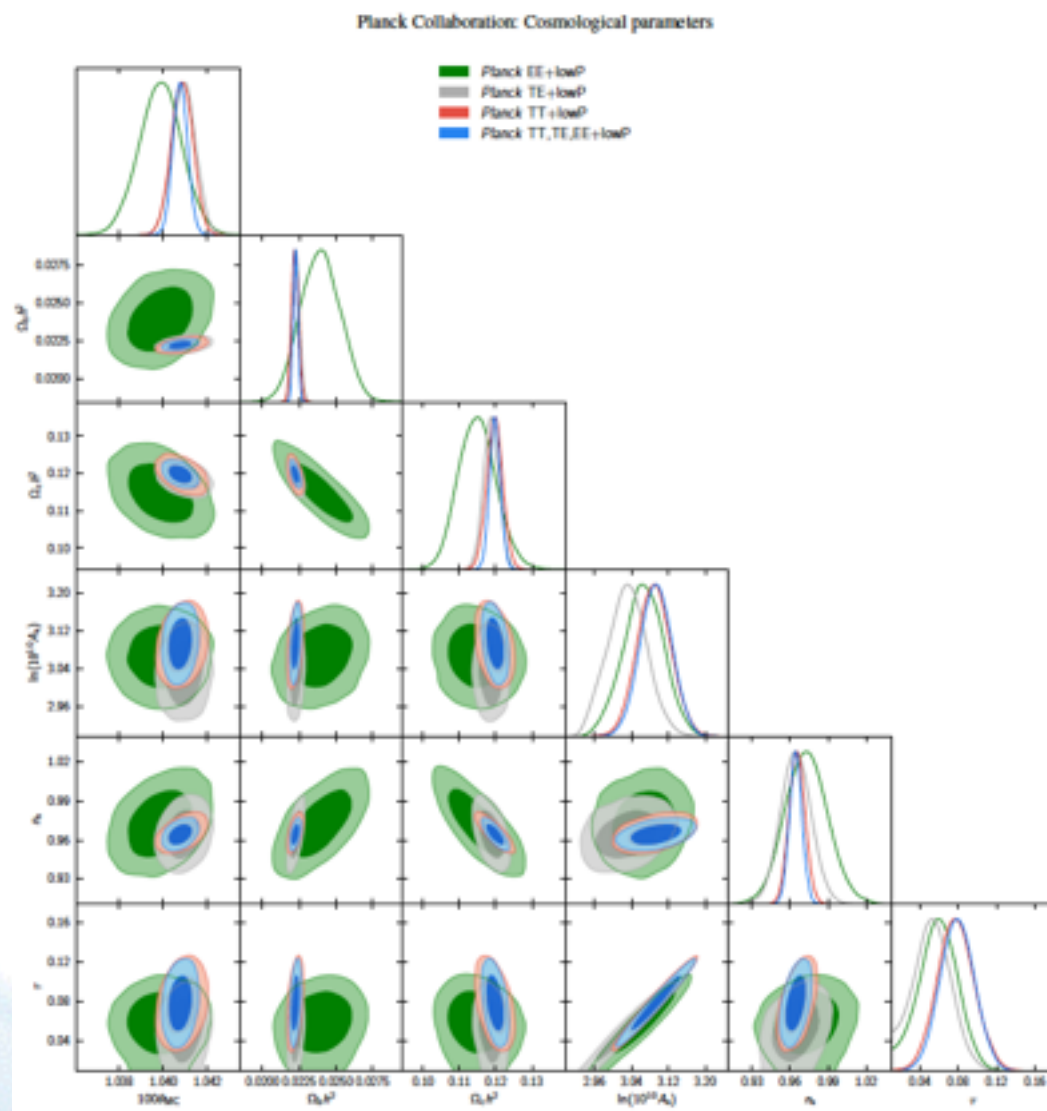
cosmic microwave background



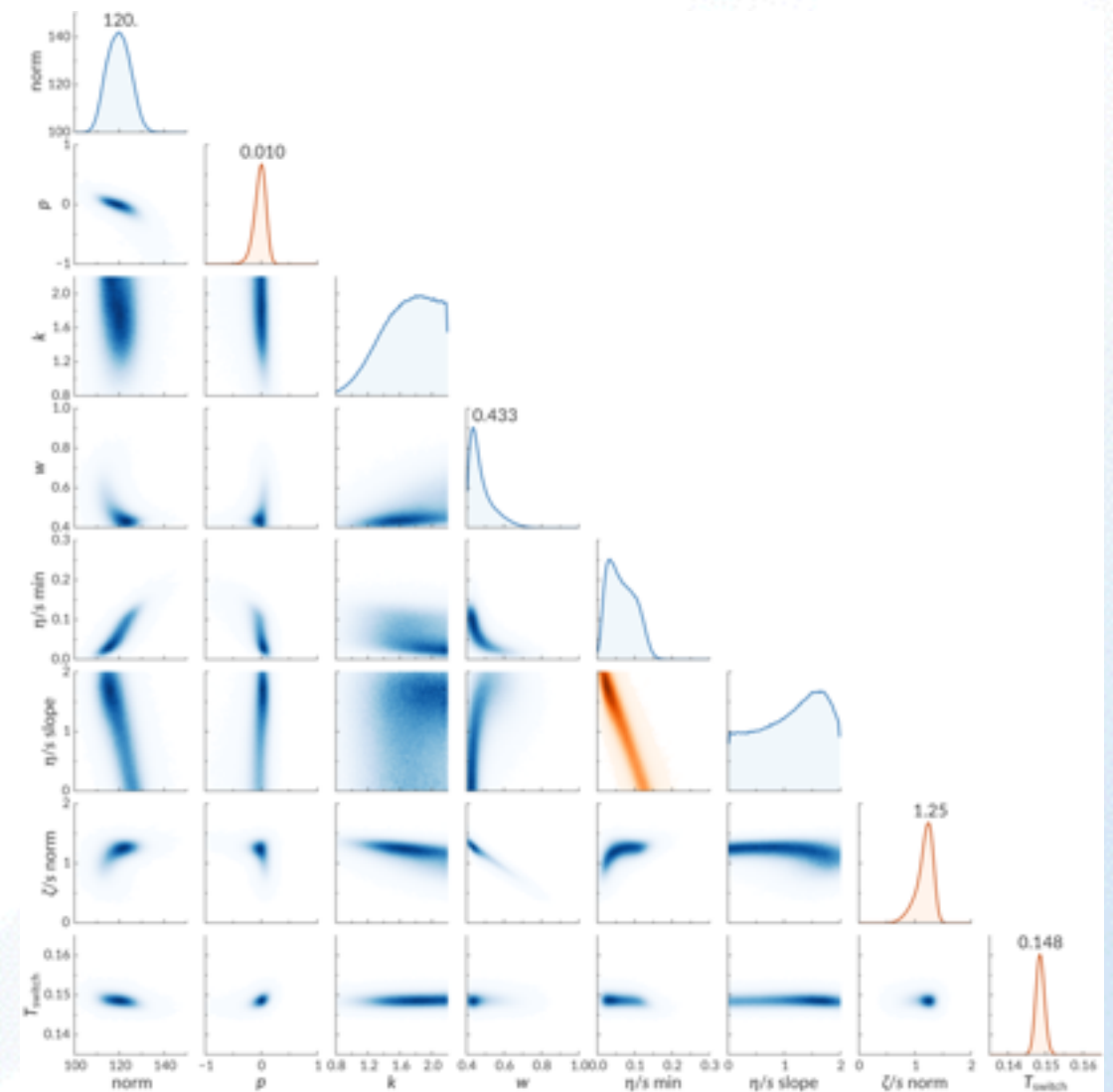
Hagedorn's temperature "limit"

51 years later

almost unimaginable progress

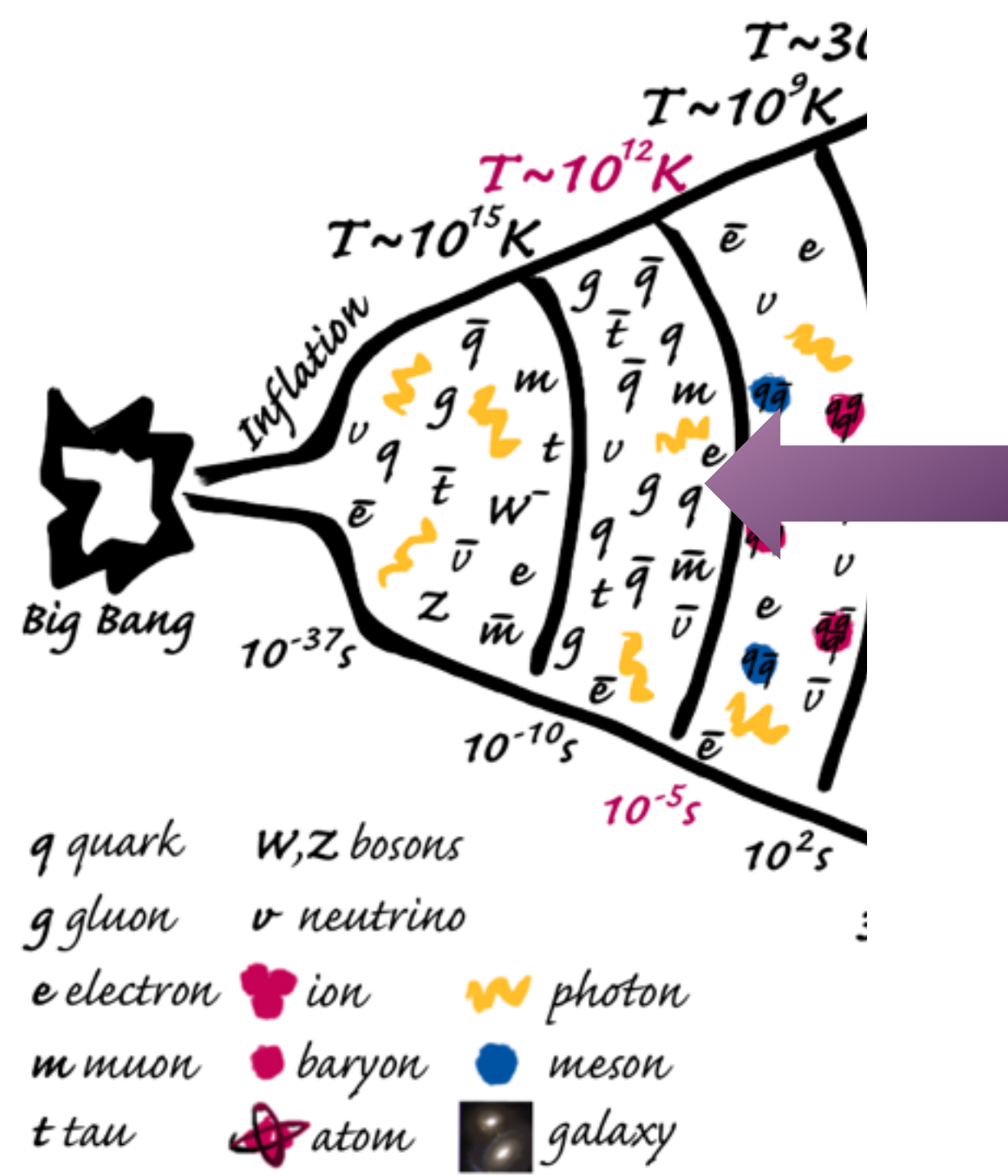


Planck: cosmological parameters



MADAI: Multiparameter analysis of heavy ion data

Evolution of the universe



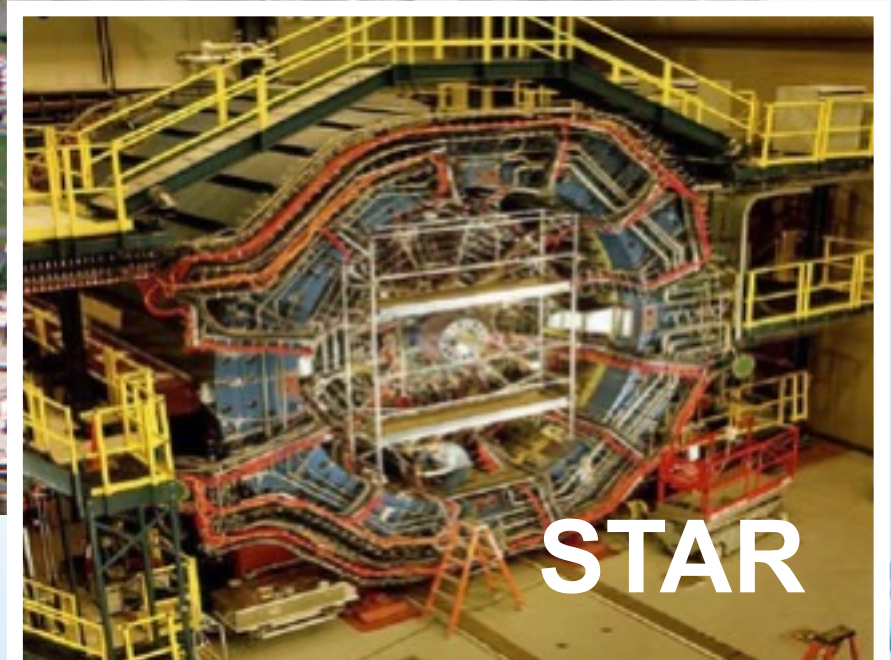
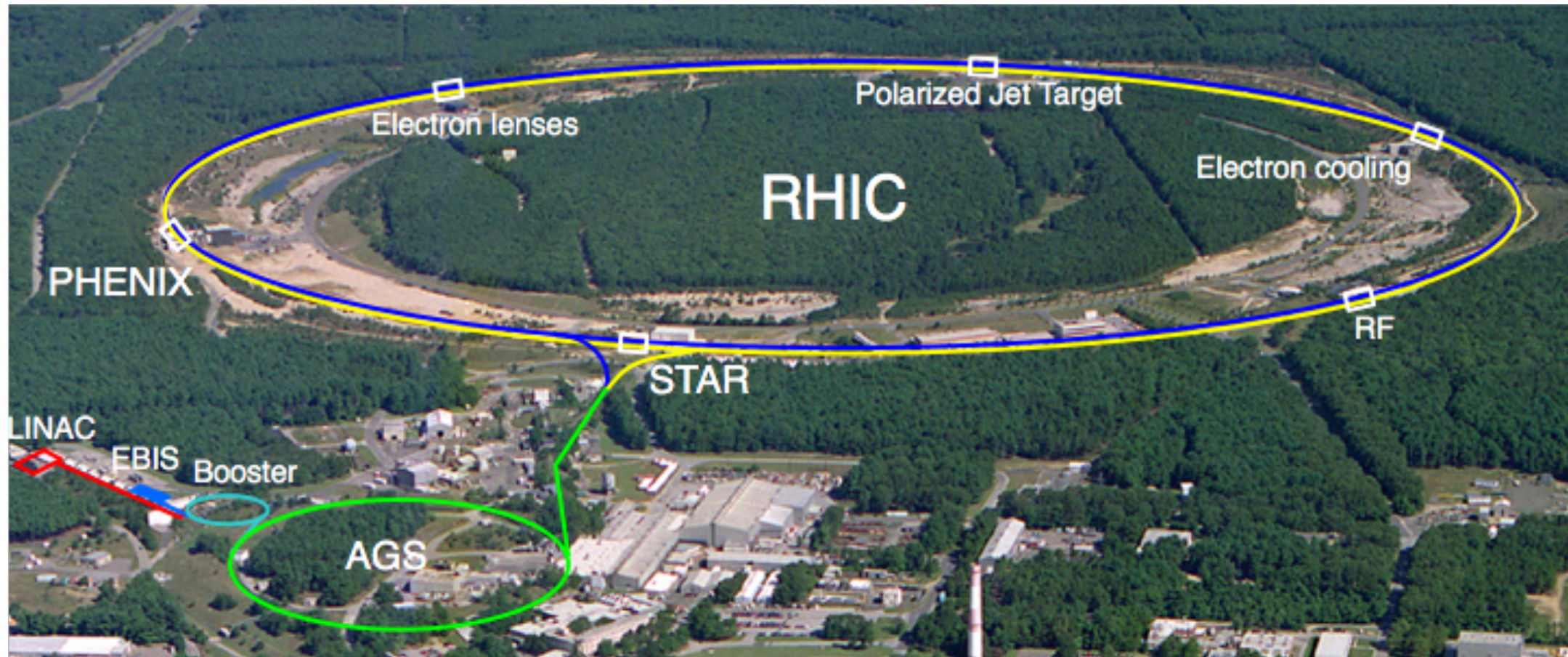
after the Big Bang

~ 70000 x hotter than
the center of the sun

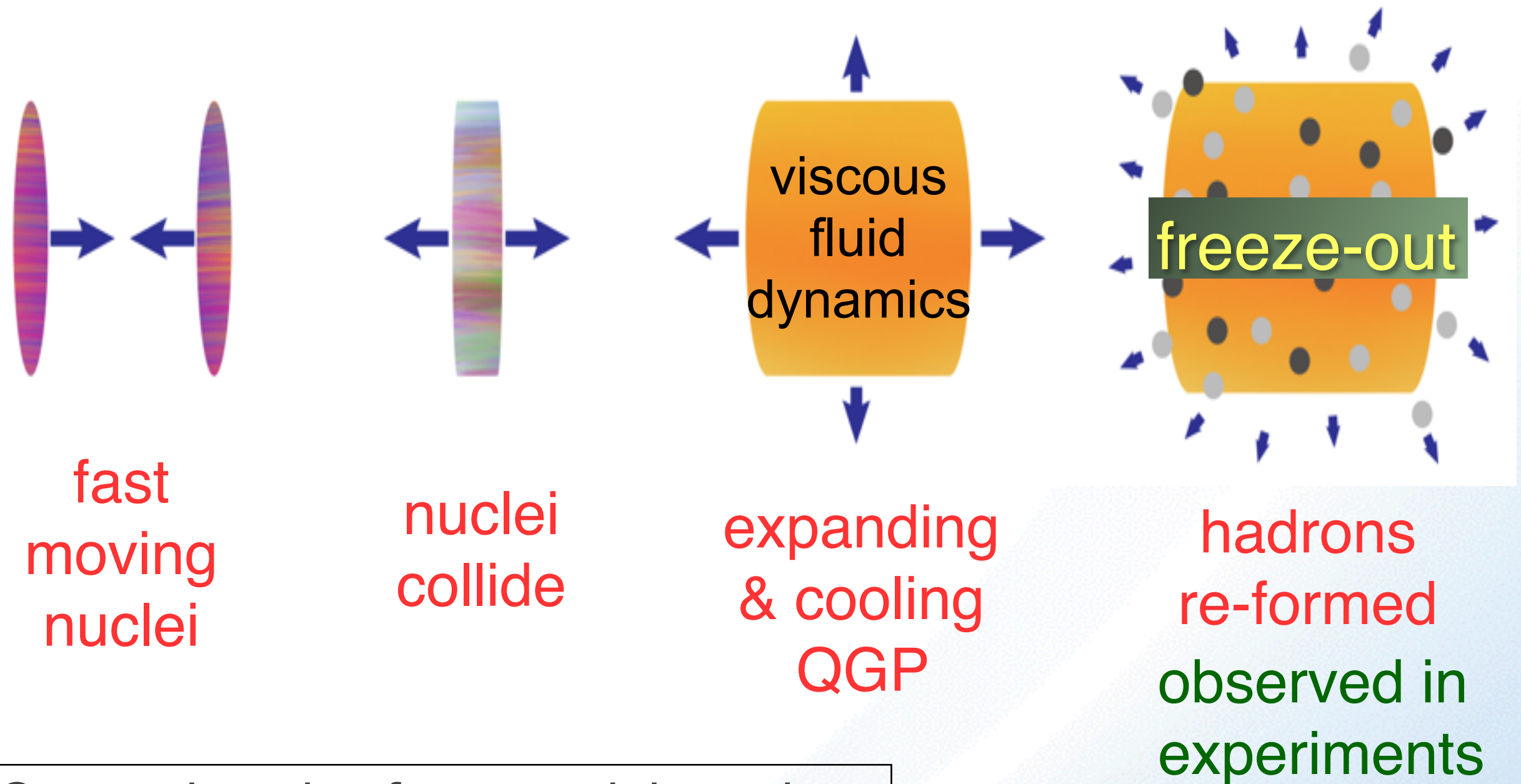


based on 'Hist

The Relativistic Heavy Ion Collider



Recreating QGP on Earth in **Little bangs**



Comprehensive framework based on transport theory provides description of reaction from start to finish

Equation of State of QCD Matter

Equation of State

EOS of flowing matter has conservative and dissipative contributions:

$$\begin{aligned} T_{\mu\nu} &= T_{\mu\nu}^{(\text{cons})} + T_{\mu\nu}^{(\text{diss})} \\ &= \epsilon u_\mu u_\nu + p(u_\mu u_\nu - g_{\mu\nu}) \\ &\quad + \eta \left(\partial_\mu u_\nu + \partial_\nu u_\mu - \frac{2}{3} g_{\mu\nu} \partial_\alpha u^\alpha \right) + \zeta \partial_\alpha u^\alpha (g_{\mu\nu} - u_\mu u_\nu) \end{aligned}$$

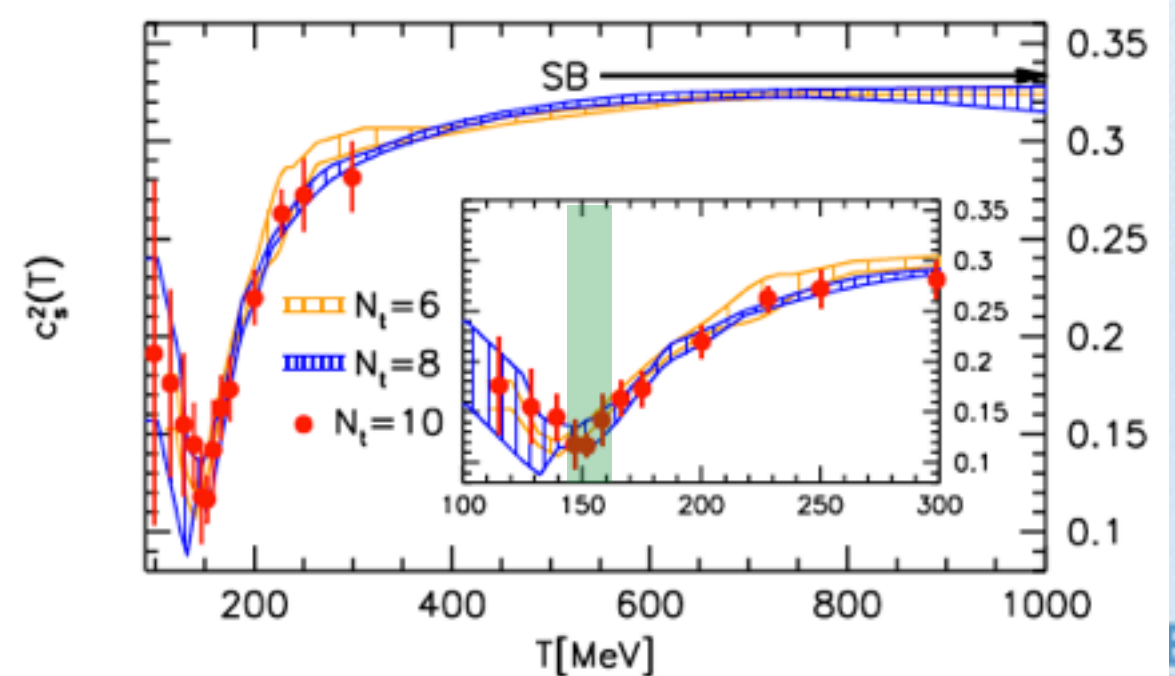
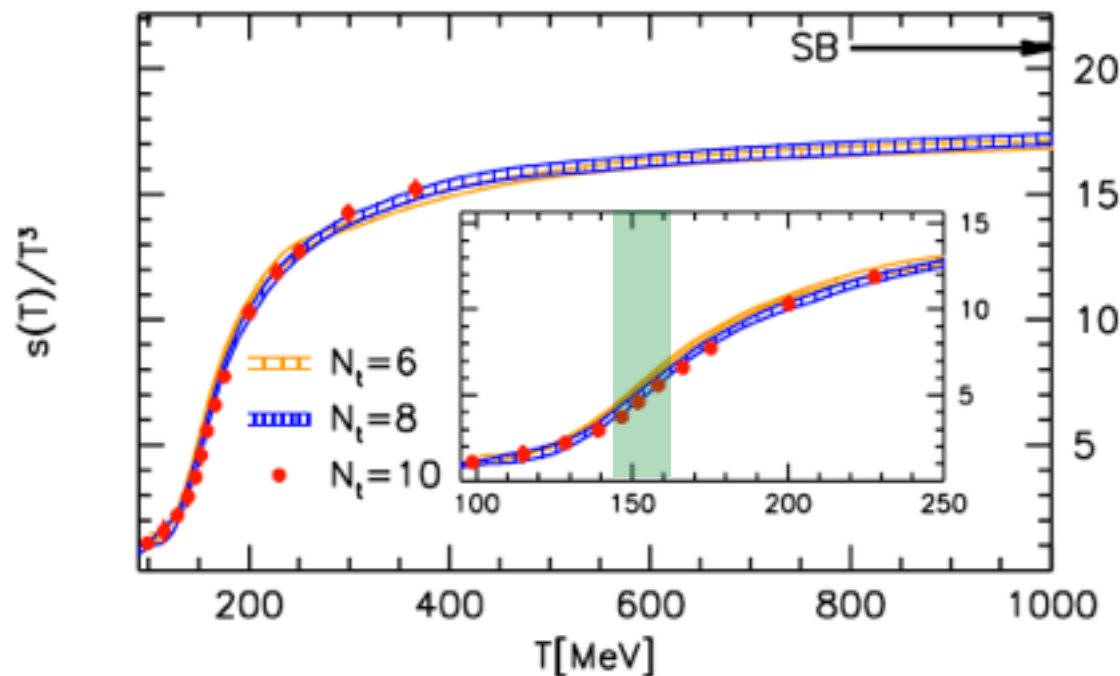
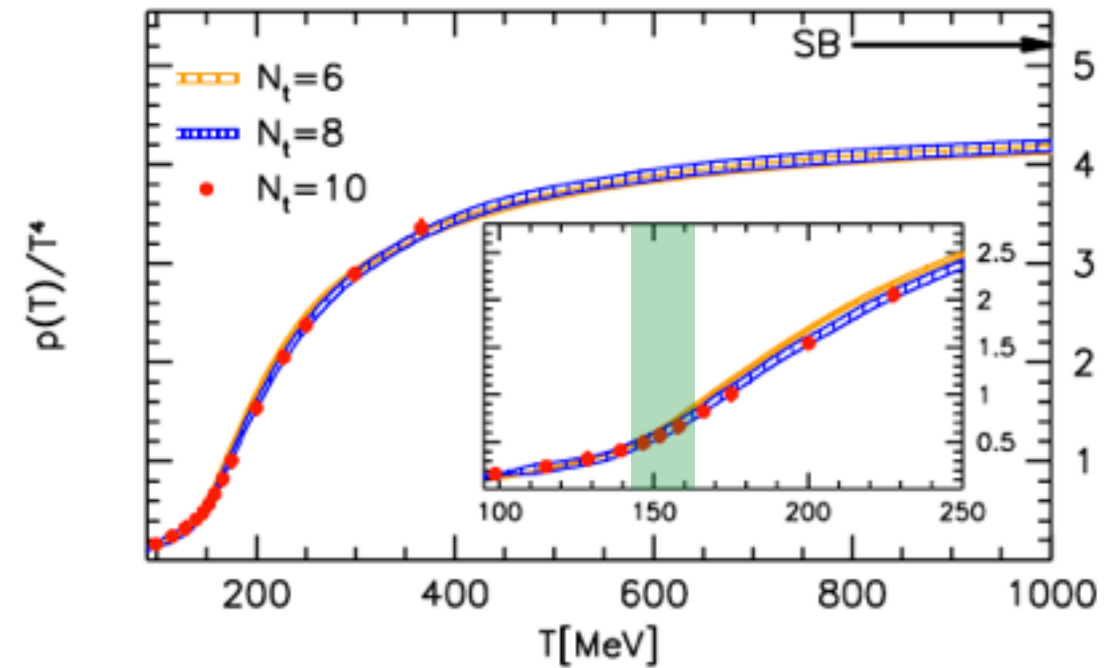
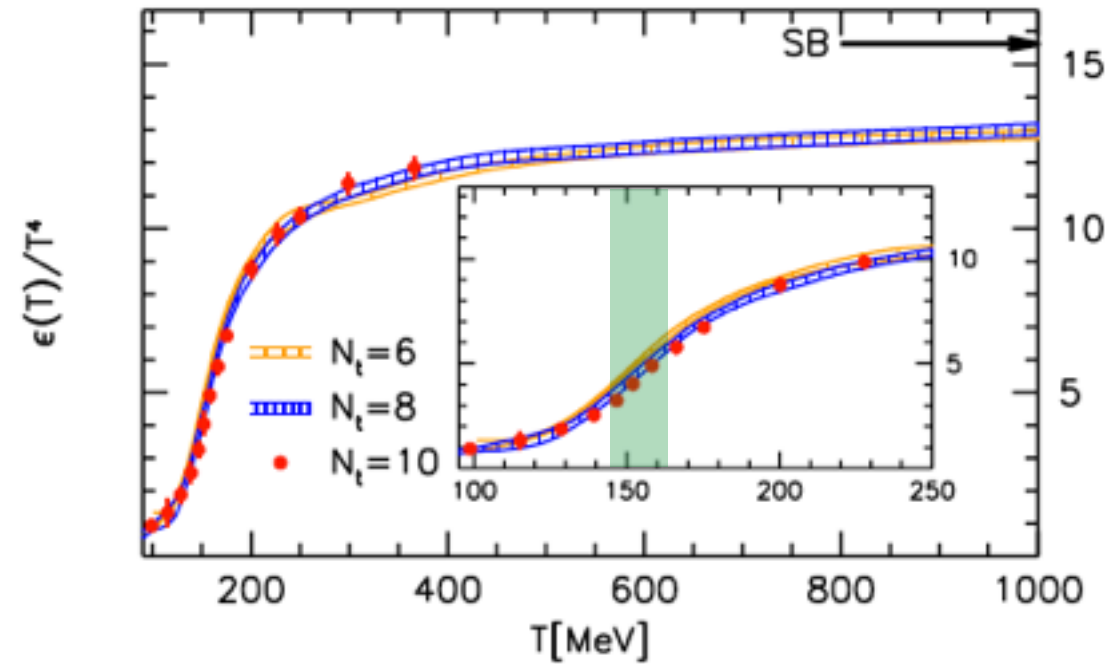
When $\zeta(\partial_\alpha u^\alpha) > p$, the matter becomes unstable and cavitates.

In general, $T_{\mu\nu}$ is a dynamical quantity that relaxes to its equilibrium value on a time scale τ_π that itself is related to the viscosity.

While the shear viscosity η has a lower quantum bound, the bulk viscosity ζ vanishes for conformally invariant matter.

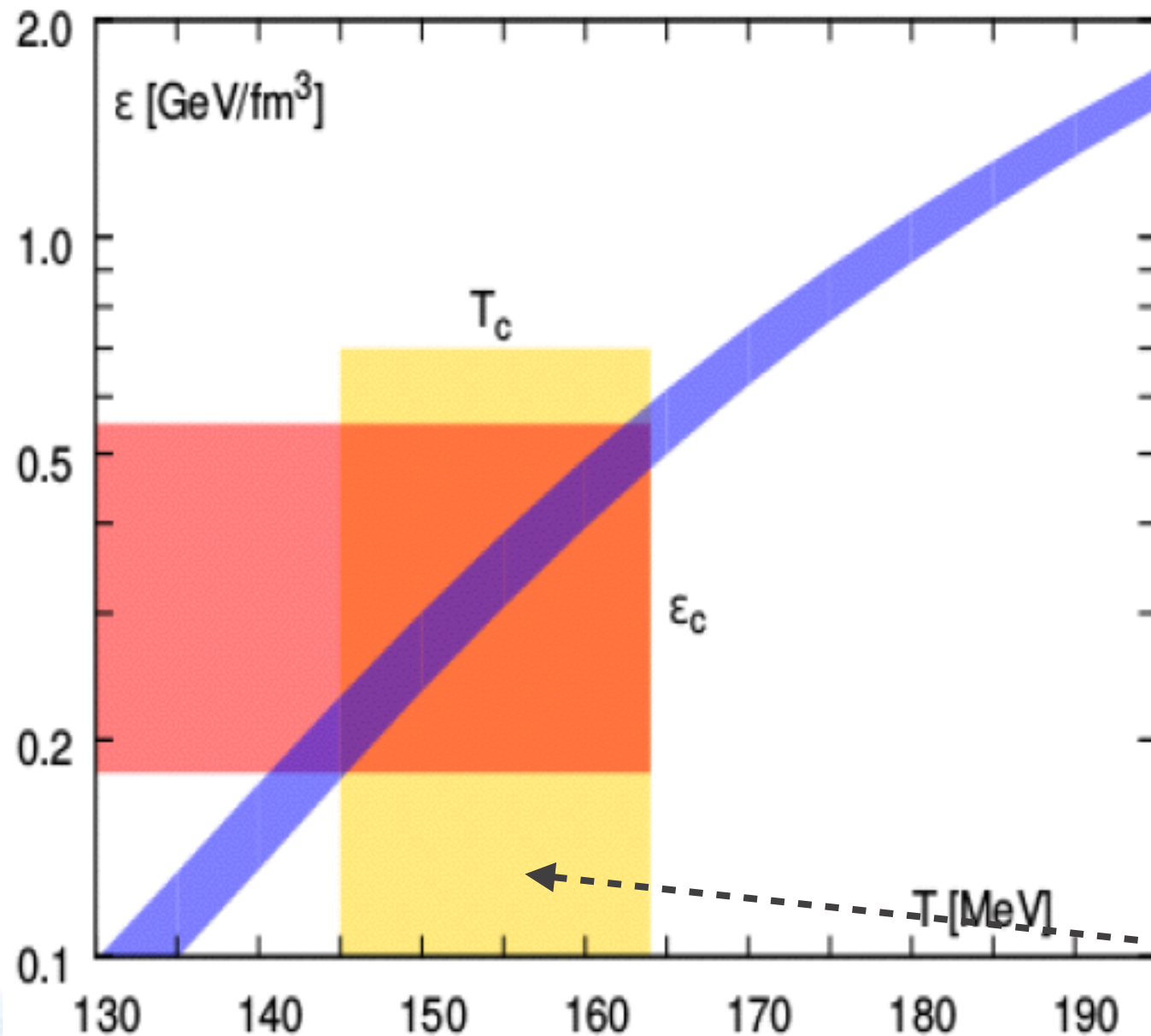
QCD EOS at $\mu_B = 0$

Results (true quark masses, continuum extrapolated) have converged; full agreement found between groups (HotQCD, Wuppertal-Budapest) using different quark actions.



(Pseudo-) Critical temperature

Transition between hadron gas and quark-gluon plasma is a **cross-over** at $\mu_B = 0$ and for small μ_B . Precise value of T_c depends on the quantity used to define it.



Pseudo-critical temperature from chiral susceptibility peak:

$$T_c = 154 \pm 9 \text{ MeV}$$

critical energy density:

$$\epsilon_c = 0.18 - 0.50 \text{ GeV/fm}^3$$

$$\epsilon_c = (1.2 - 3.3) \rho_{\text{nuclear}}$$

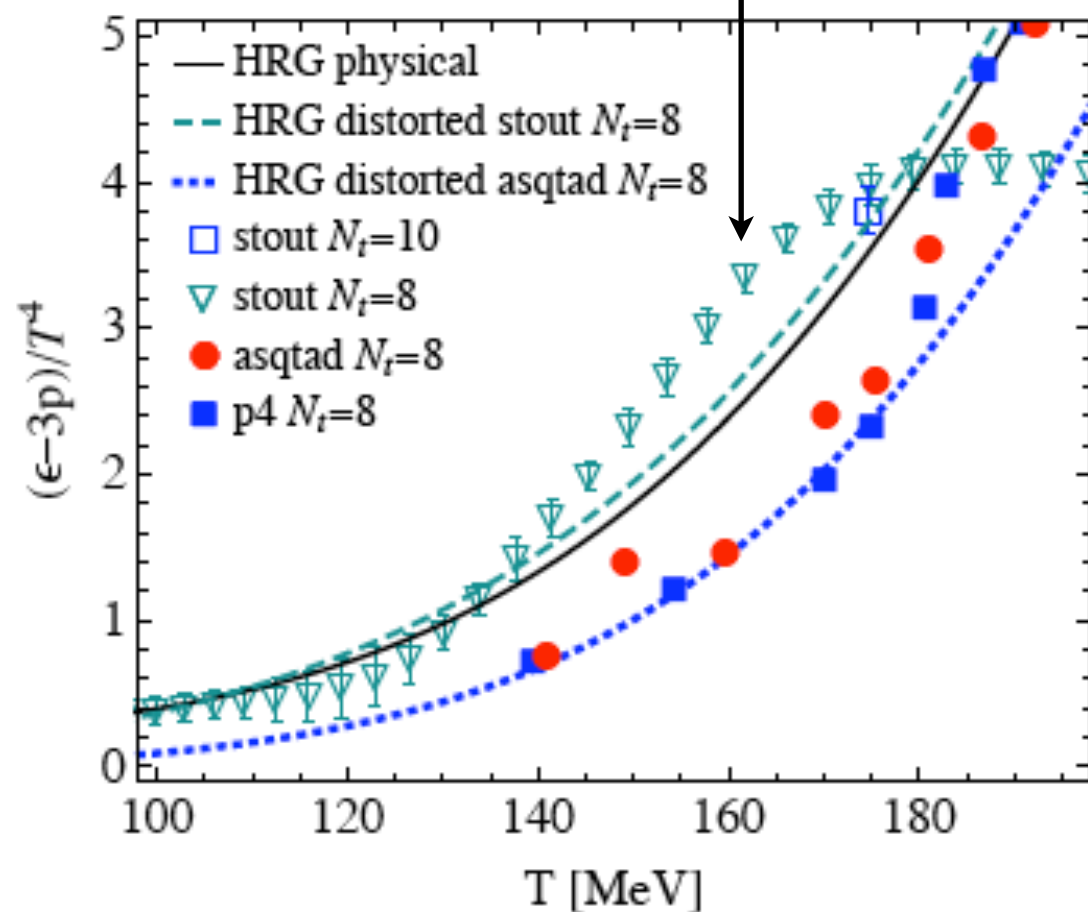
Uncertainty in T_c , not width of cross-over region!

HotQCD: arXiv:1407.6387

Hadron mass spectrum

Below T_c , the quantity $(\epsilon-3p)/T^4$ measures the level density of massive hadronic excitations of the QCD vacuum.

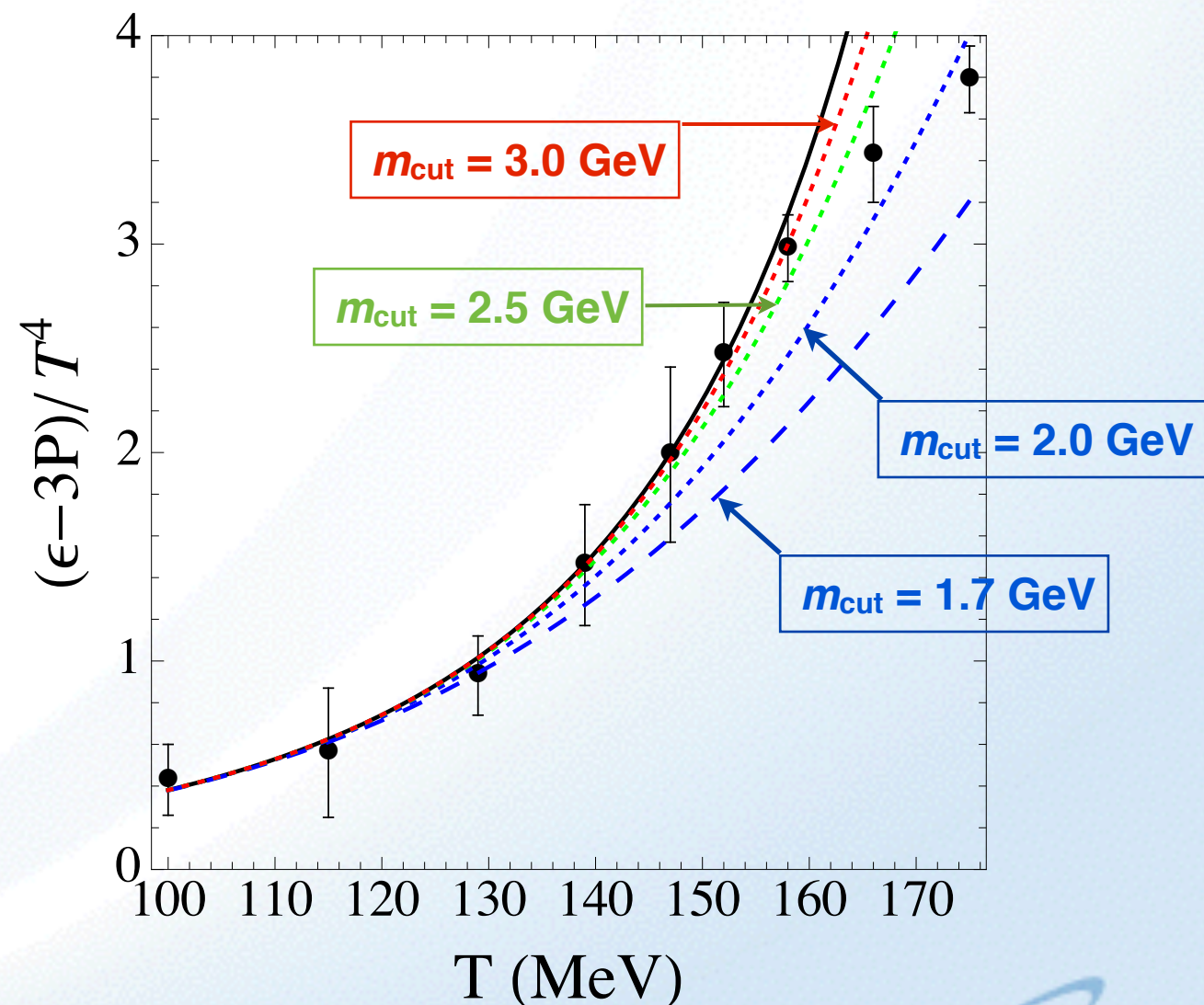
Lines: Hadron resonance gas using only PDG resonances
Data points: Lattice QCD
LQCD lies **above** HRG for $T > 140$ MeV
Indicates additional hadron resonances



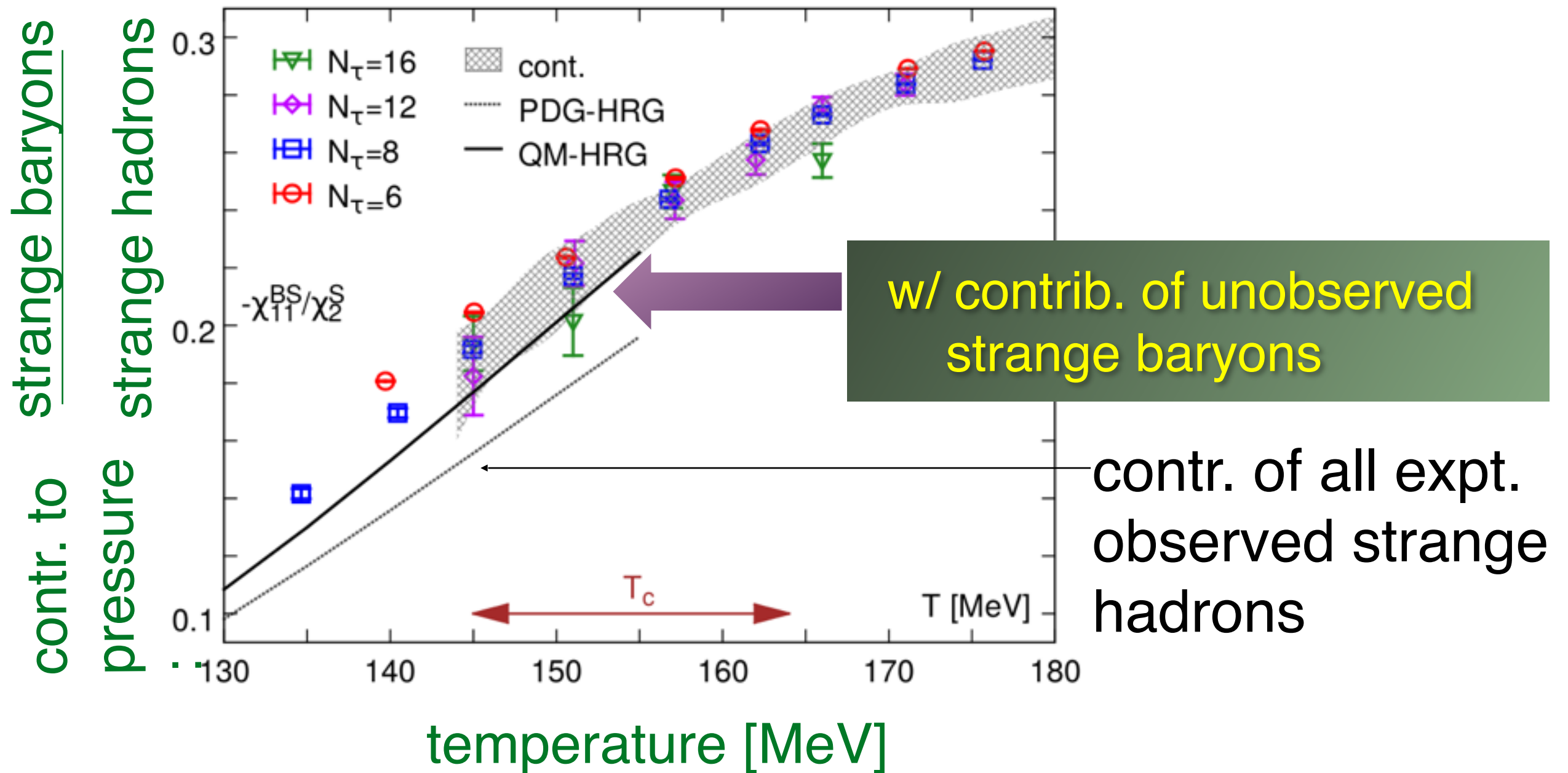
Hagedorn spectrum ($T_H \approx 180$ MeV):

$$\rho_H(m) = \frac{A e^{m/T_H}}{(m^2 + m_0^2)^{5/4}}$$

In good agreement with lattice results
Hadrons up to 3 GeV mass contribute



Lattice evidence for strange baryons

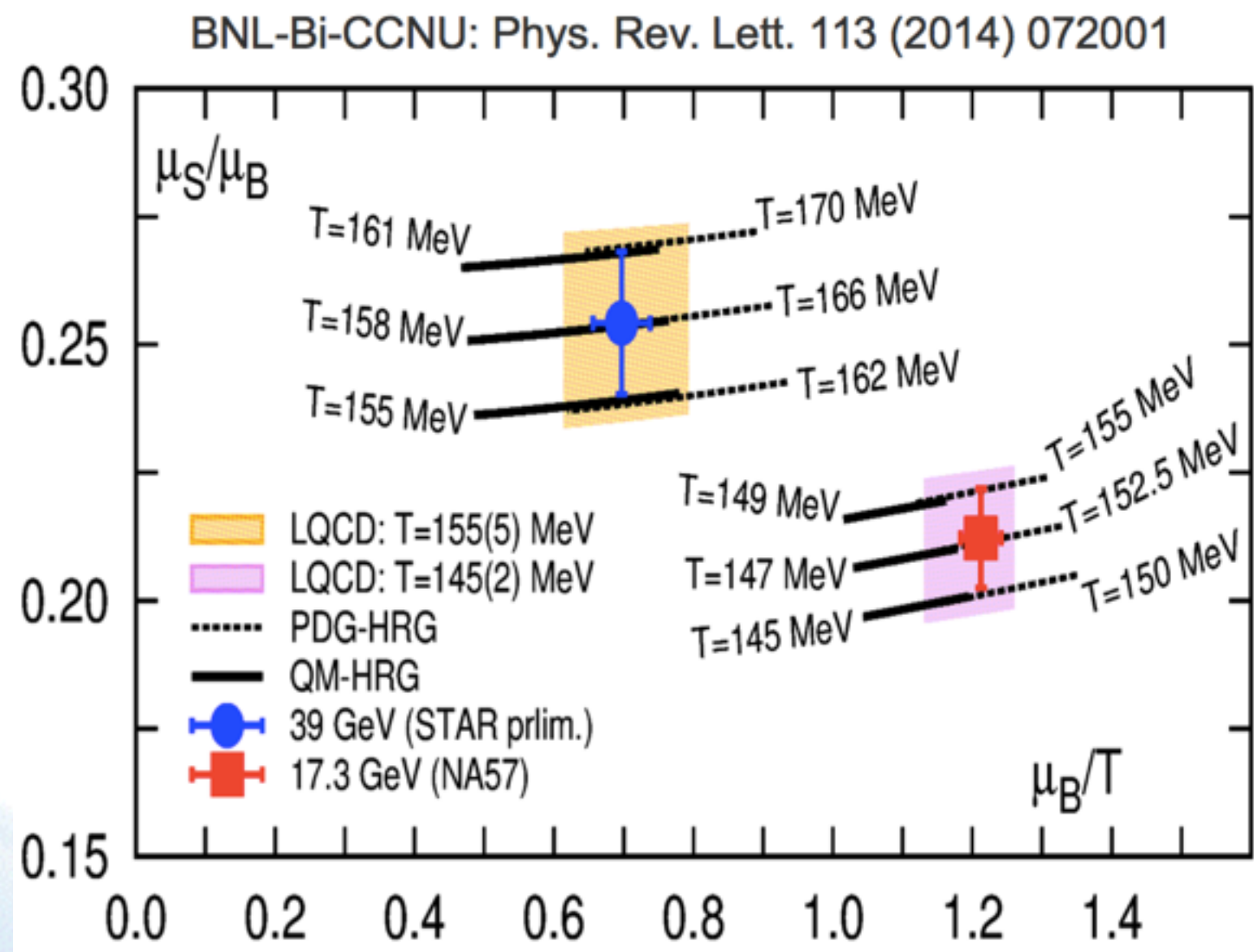


indirect signatures in heavy-ion experiments

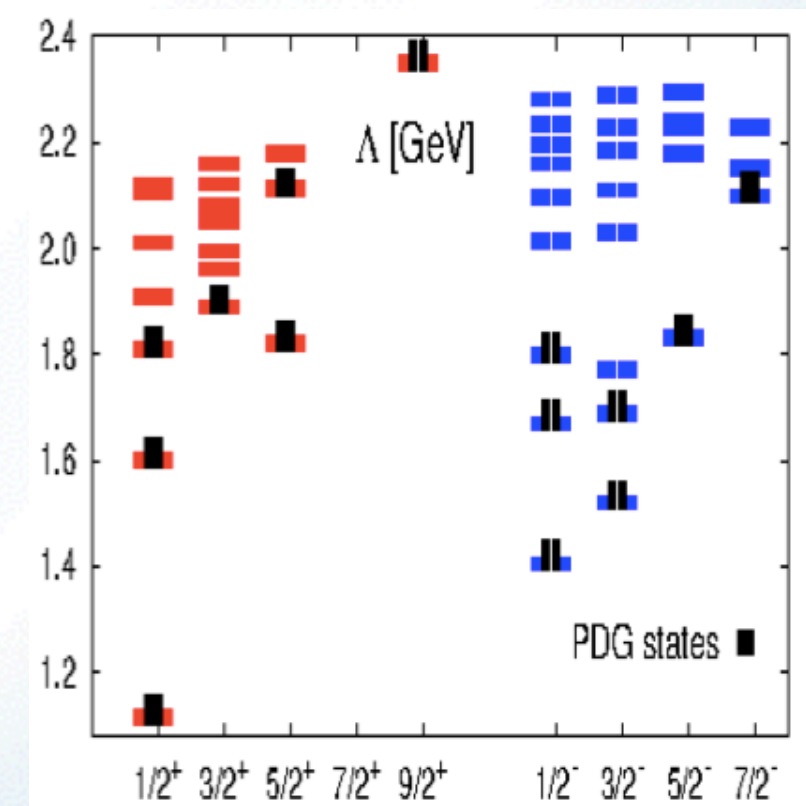
BNL-Bielefeld-CCNU: Phys. Rev. Lett. 113, 072001 (2014)

Probing the baryon spectrum

Consistency of μ_s/μ_B and μ_B/T with chemical composition of emitted hadrons and Lattice QCD requires additional strange baryon resonances beyond those in the PDG tables.

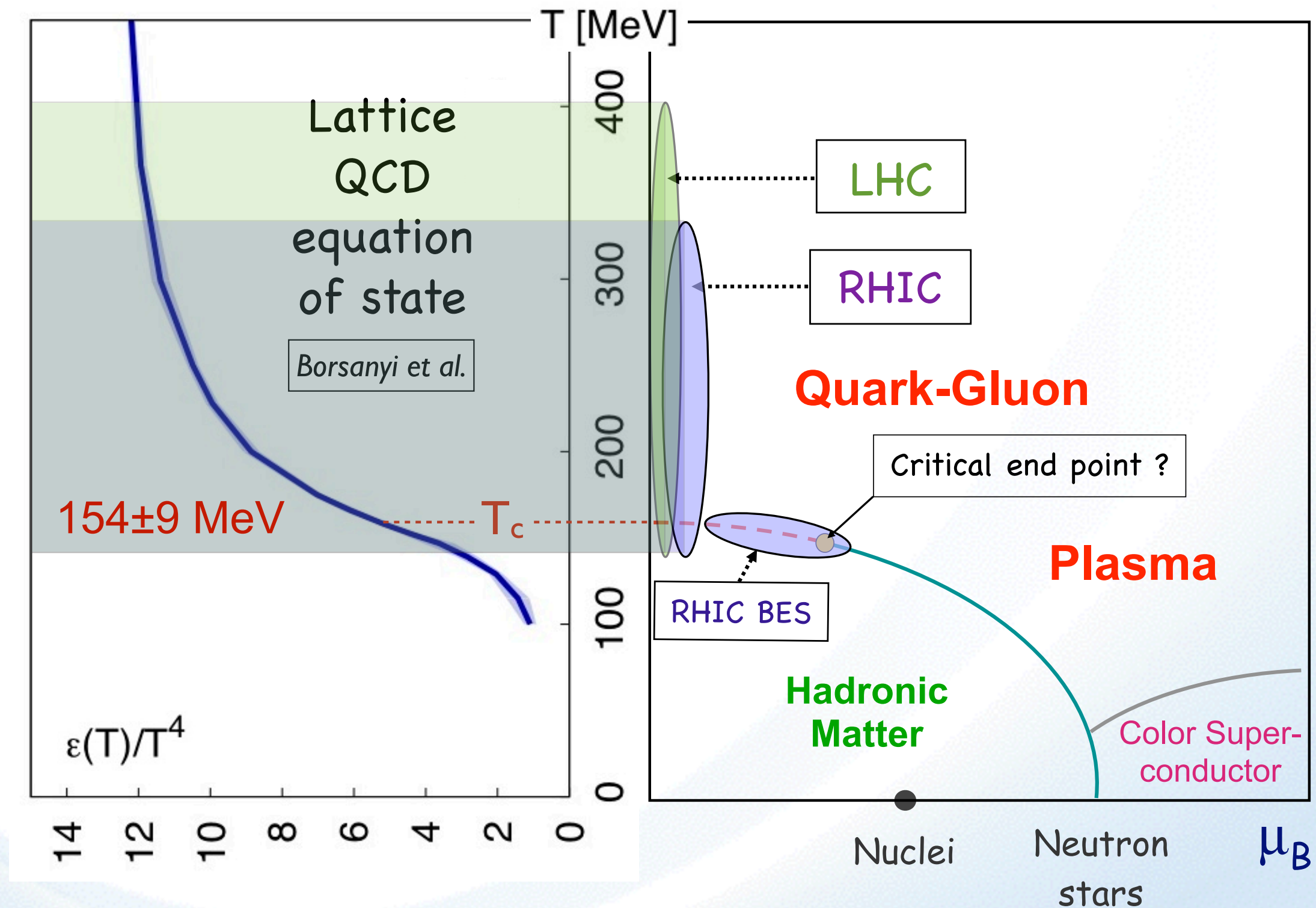


Quark model states of strange baryons



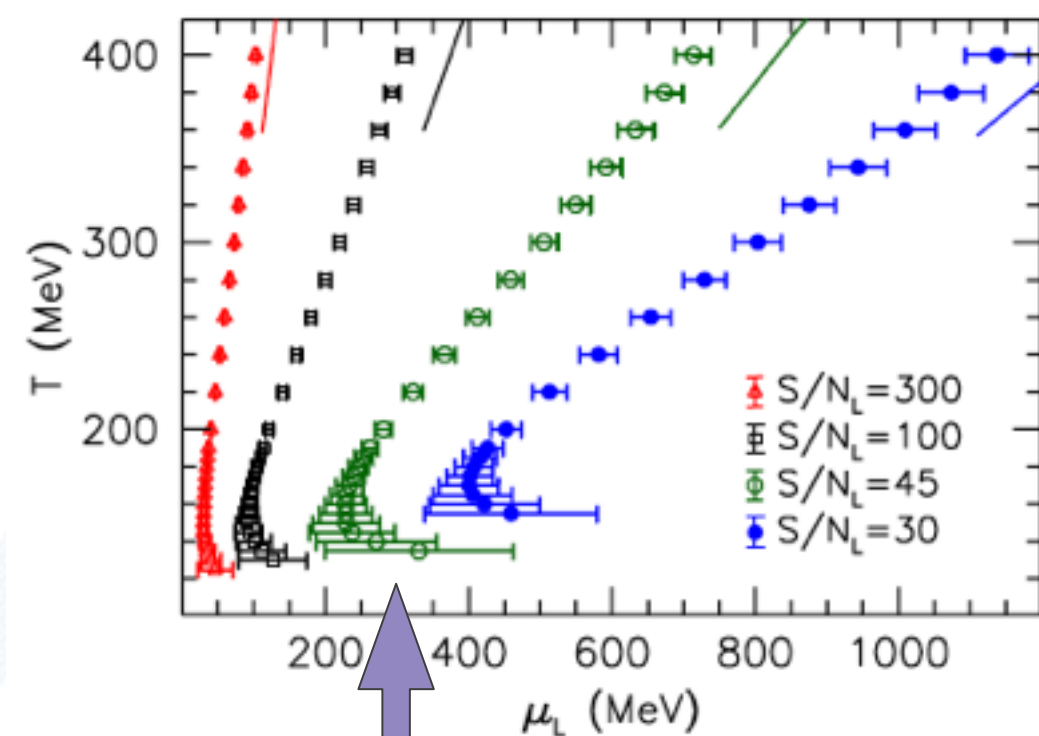
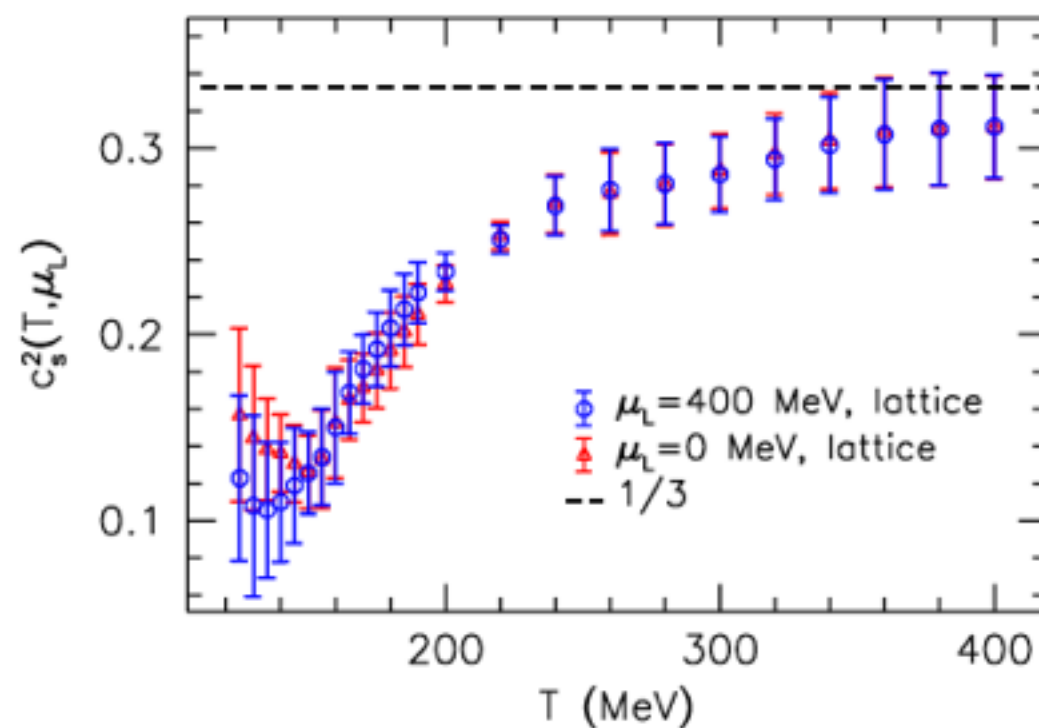
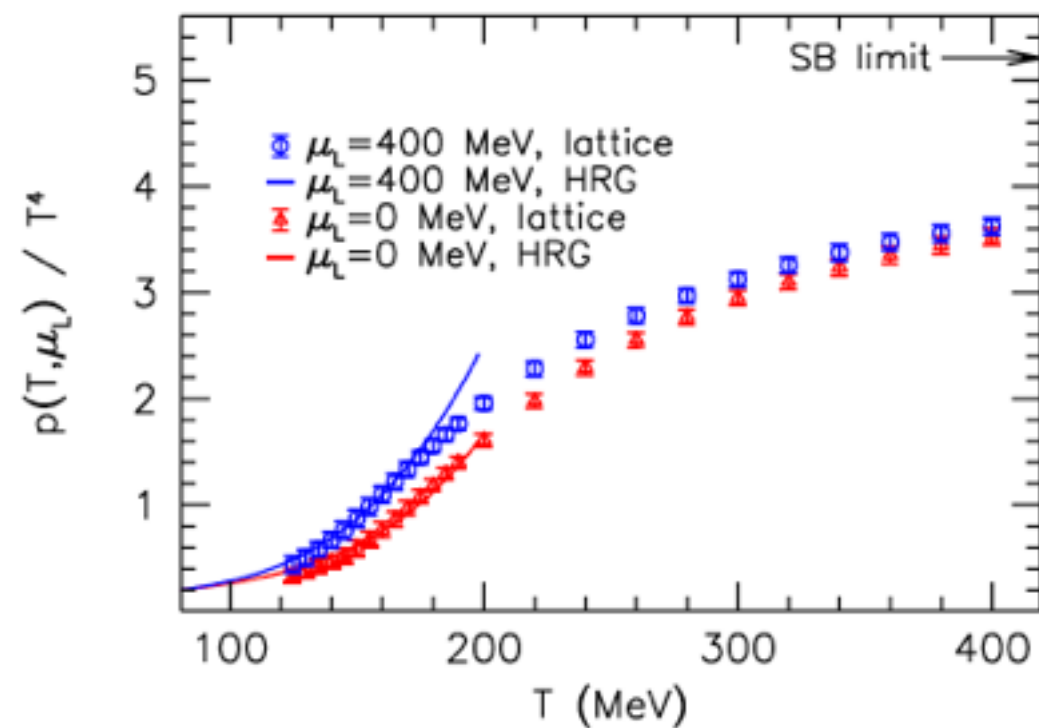
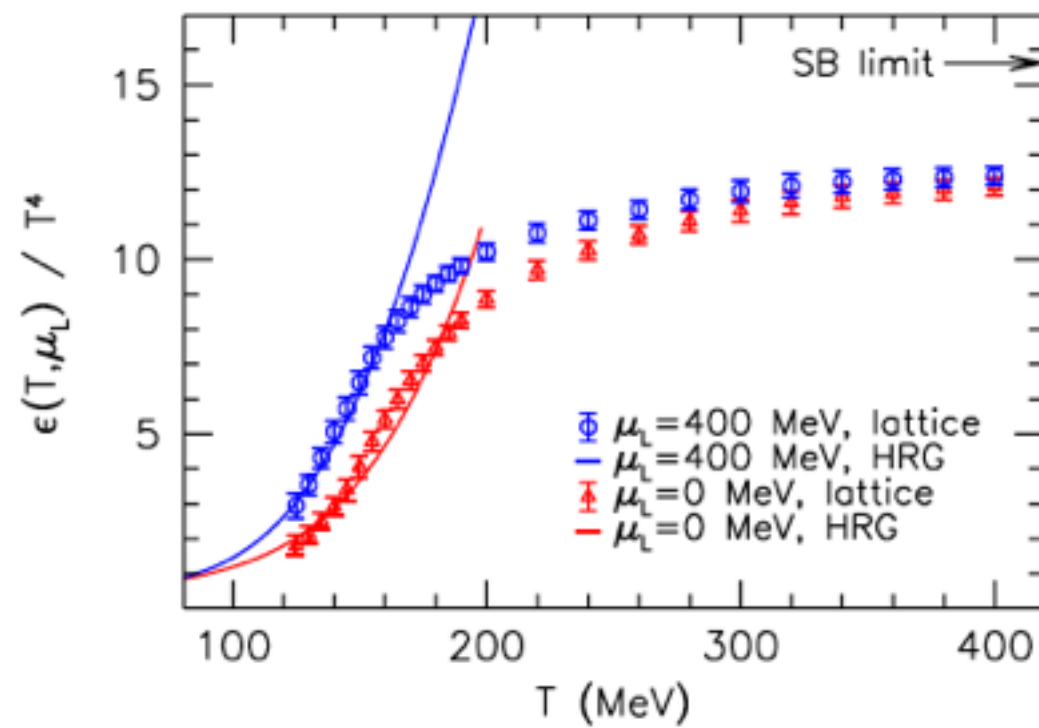
Probing the QCD Phase Boundary

QCD Phase Diagram



QCD EOS at $\mu_B \neq 0$

Borszanyi et al., arXiv:1204:6710



Approximate trajectories in QCD phase diagram

Thermodynamic fluctuations

Susceptibilities measure thermodynamic fluctuations.
Interesting because they exhibit singularities at a critical point.
Fluctuations of **conserved quantities** (charge Q , baryon number B ,...) cannot be changed by local final-state processes.

Expt.: mean: M_Q
variance: σ_Q^2
skewness: S_Q
kurtosis: κ_Q

$$\sqrt{s} \Leftrightarrow (T, \mu_B)$$

Lattice gauge theory:

$$\chi_n^X(T, \mu_X) = \frac{\partial^n (p(T, \mu_X)/T^4)}{\partial (\mu_X/T)^n}$$

Ratios are independent of the (unknown) freeze-out volume:

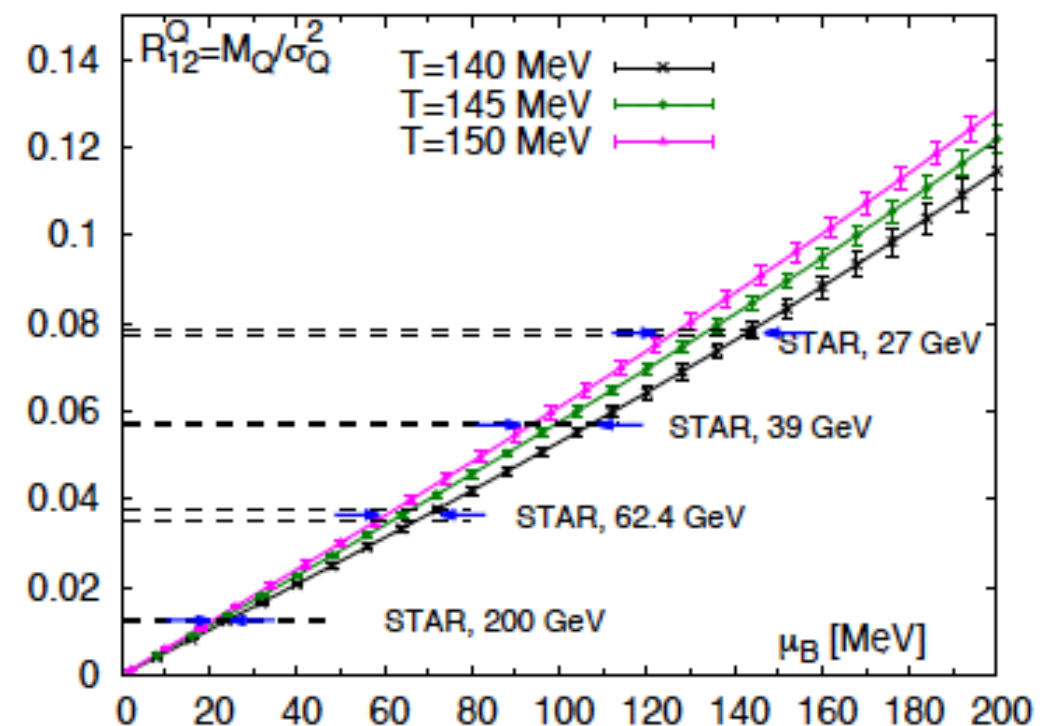
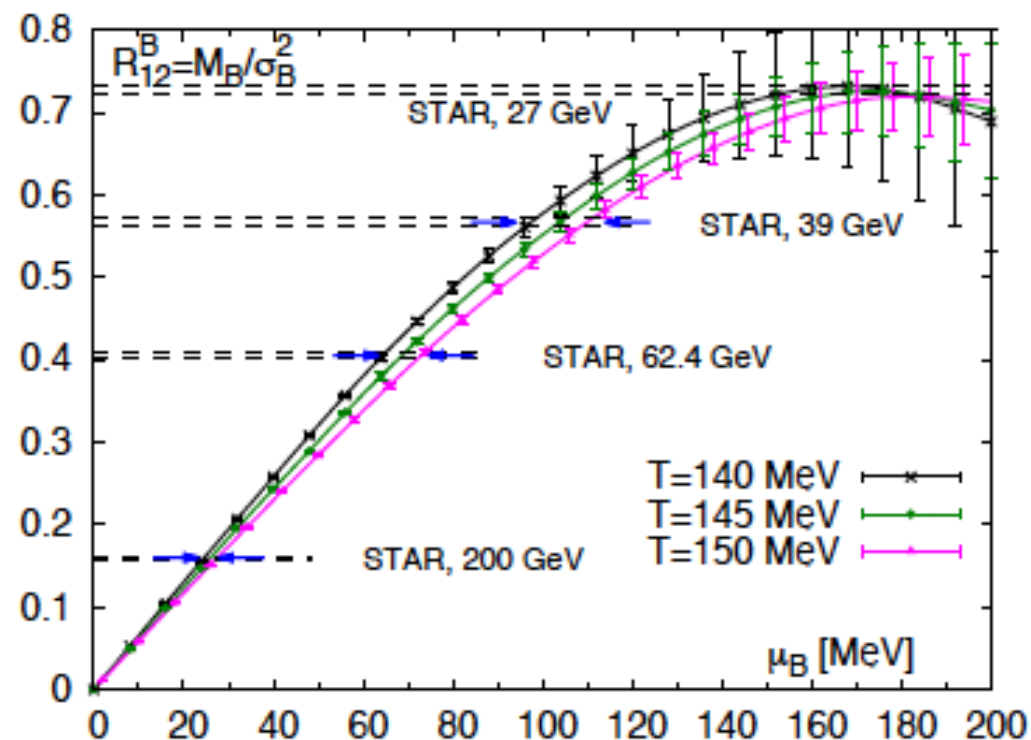
$$\frac{M_Q(\sqrt{s})}{\sigma_Q^2(\sqrt{s})} = \frac{\chi_1^Q(T, \mu_B)}{\chi_2^Q(T, \mu_B)} \quad \frac{S_Q(\sqrt{s}) \sigma_Q^3(\sqrt{s})}{M_Q(\sqrt{s})} = \frac{\chi_3^Q(T, \mu_B)}{\chi_1^Q(T, \mu_B)}$$

Chemical freeze-out

... from fluctuations of conserved quantum numbers (Q , B):

Borsanyi et al. Wuppertal-Budapest Coll. Phys.Rev.Lett. 111, 062005 (2013); Phys.Rev.Lett. 113, 052301 (2014)

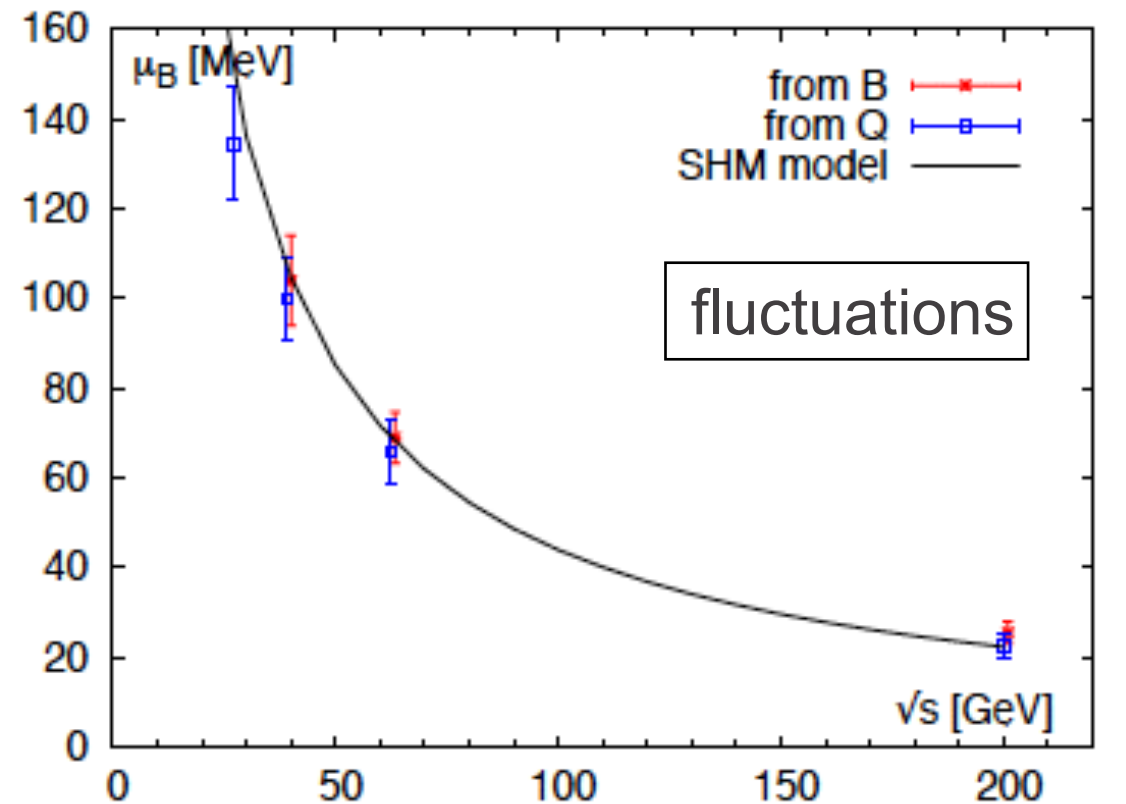
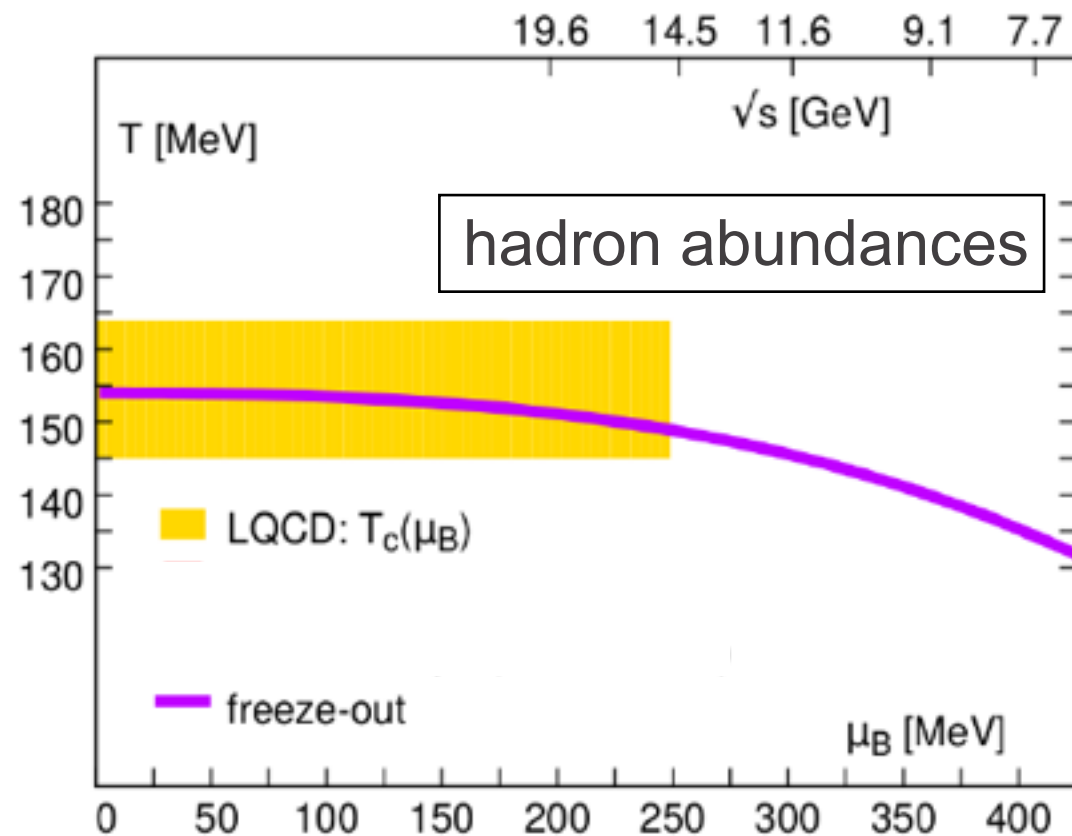
use M/σ^2 both in the baryon and in the charge sector



Compare lattice results with the STAR data for the fluctuation ratios in the temperature range 140–150 MeV permits to read off μ_B .

Both methods are consistent with each other and with the measured baryon/antibaryon ratios, if additional strange baryon states beyond those in the PDG tables (e.g. in the quark model) are accounted for.

Chemical freeze-out



Consistency of freeze-out parameters from mean hadron abundances and from fluctuations (Q , B) opens the door to search for a critical point in the QCD phase diagram by looking for enhanced critical fluctuations as function of beam energy.

The “perfect” fluid

Viscous hydrodynamics

Hydrodynamics = effective theory of energy and momentum conservation

$$\text{energy-momentum tensor} = \text{ideal fluid} + \text{dissipation}$$

$$\partial_\mu T^{\mu\nu} = 0 \quad \text{with} \quad T^{\mu\nu} = (\varepsilon + P)u^\mu u^\nu - P g^{\mu\nu} + \Pi^{\mu\nu}$$

$$\tau_\Pi \left[\frac{d\Pi^{\mu\nu}}{d\tau} + \left(u^\mu \Pi^{\nu\lambda} + u^\nu \Pi^{\mu\lambda} \right) \frac{du^\lambda}{d\tau} \right] = \eta \left(\partial^\mu u^\nu + \partial^\nu u^\mu - \text{trace} \right) - \Pi^{\mu\nu}$$

Input: Equation of state $P(\varepsilon)$, shear viscosity, initial conditions $\varepsilon(x,0)$, $u^\mu(x,0)$

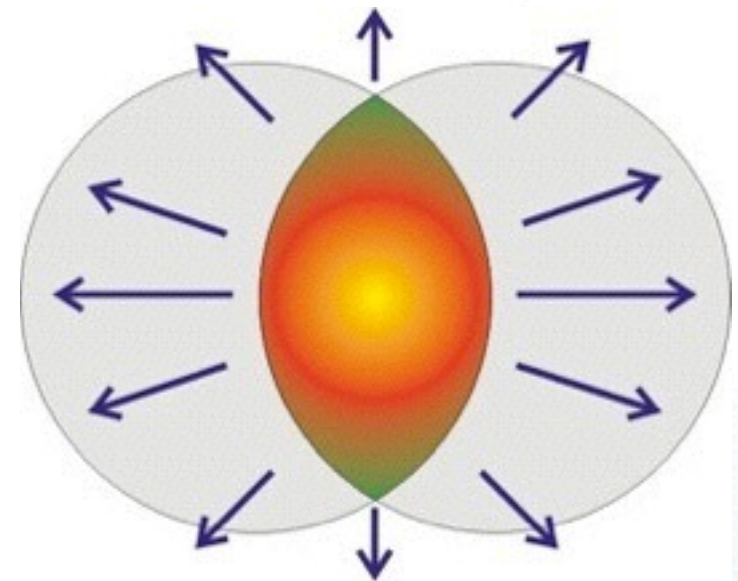
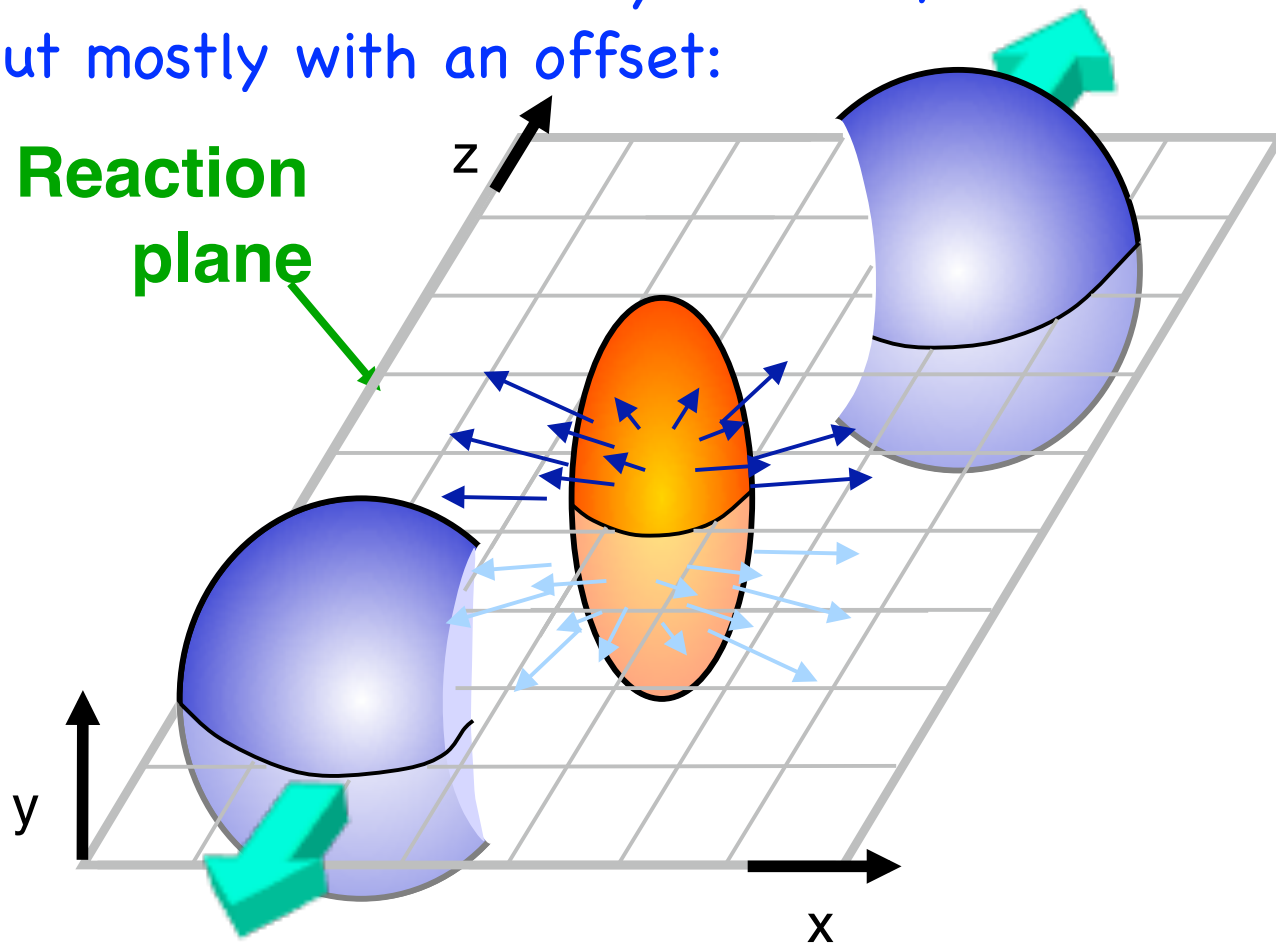
Shear viscosity η is normalized by density: **kinematic viscosity** η/ρ .

Relativistically, the appropriate normalization factor is the **entropy density** $s = (\varepsilon+P)/T$, because the particle density is not conserved: η/s .

Elliptic flow

- two nuclei collide rarely head-on, but mostly with an offset:

Reaction plane



only matter in the overlap area gets compressed and heated

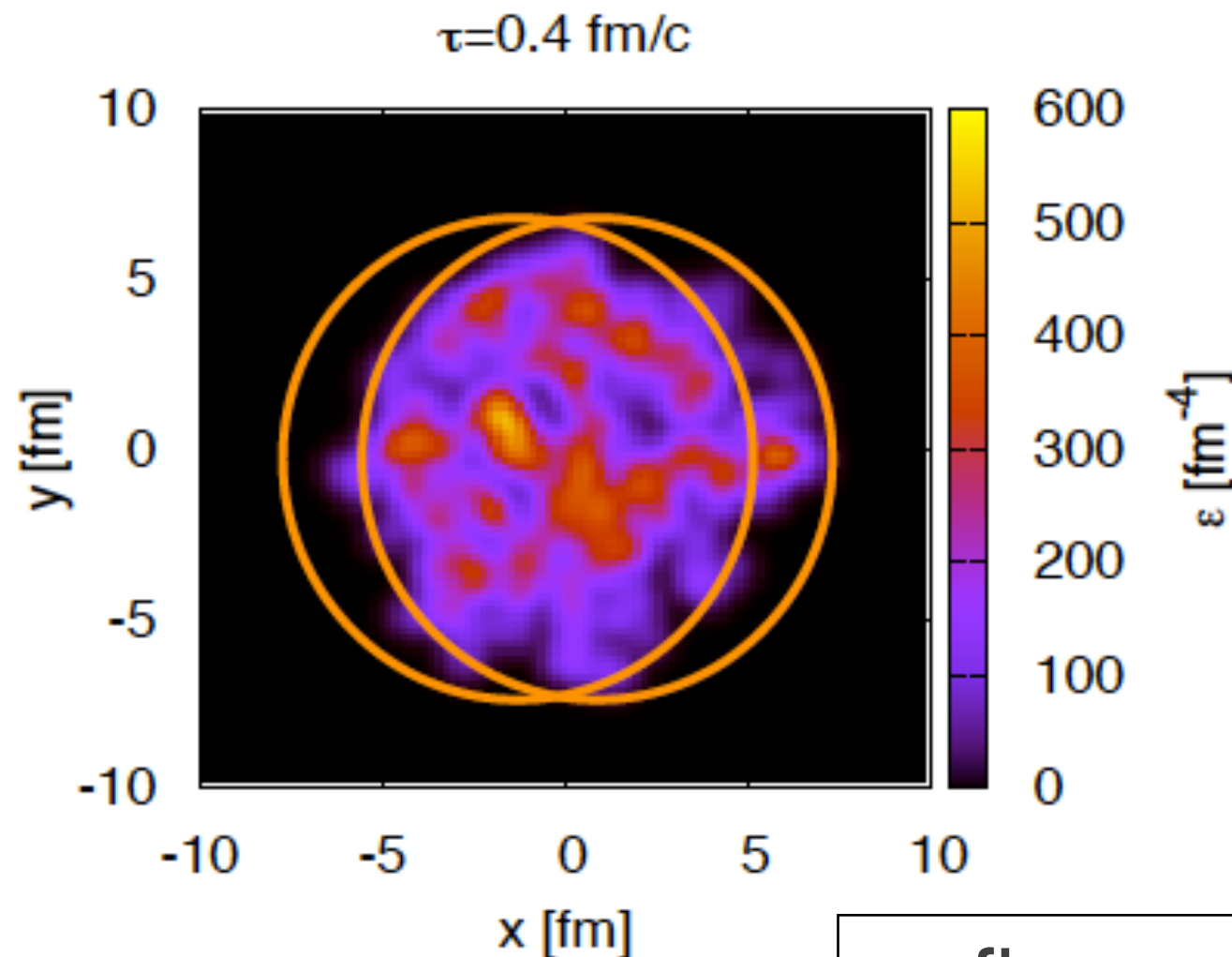
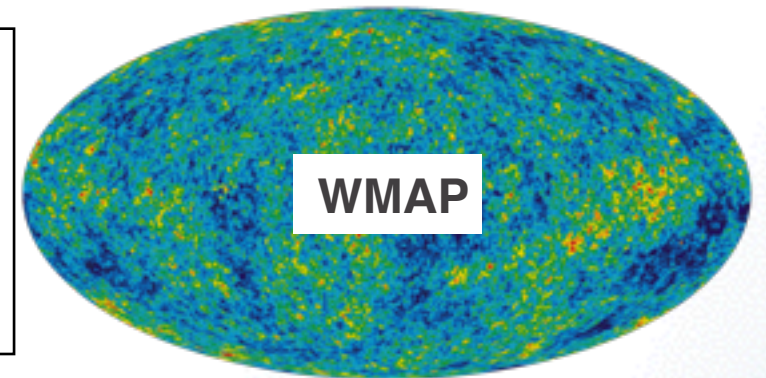
$$2\pi \frac{dN}{d\phi} = N_0 \left(1 + 2 \sum_n v_n(p_T, \eta) \cos n(\phi - \psi_n(p_T, \eta)) \right)$$

anisotropic flow coefficients

event plane angle

Event-by-event fluctuations

Initial state generated in A+A collision is grainy
event plane \neq reaction plane
 \Rightarrow eccentricities $\varepsilon_1, \varepsilon_2, \varepsilon_3, \varepsilon_4$, etc. $\neq 0$

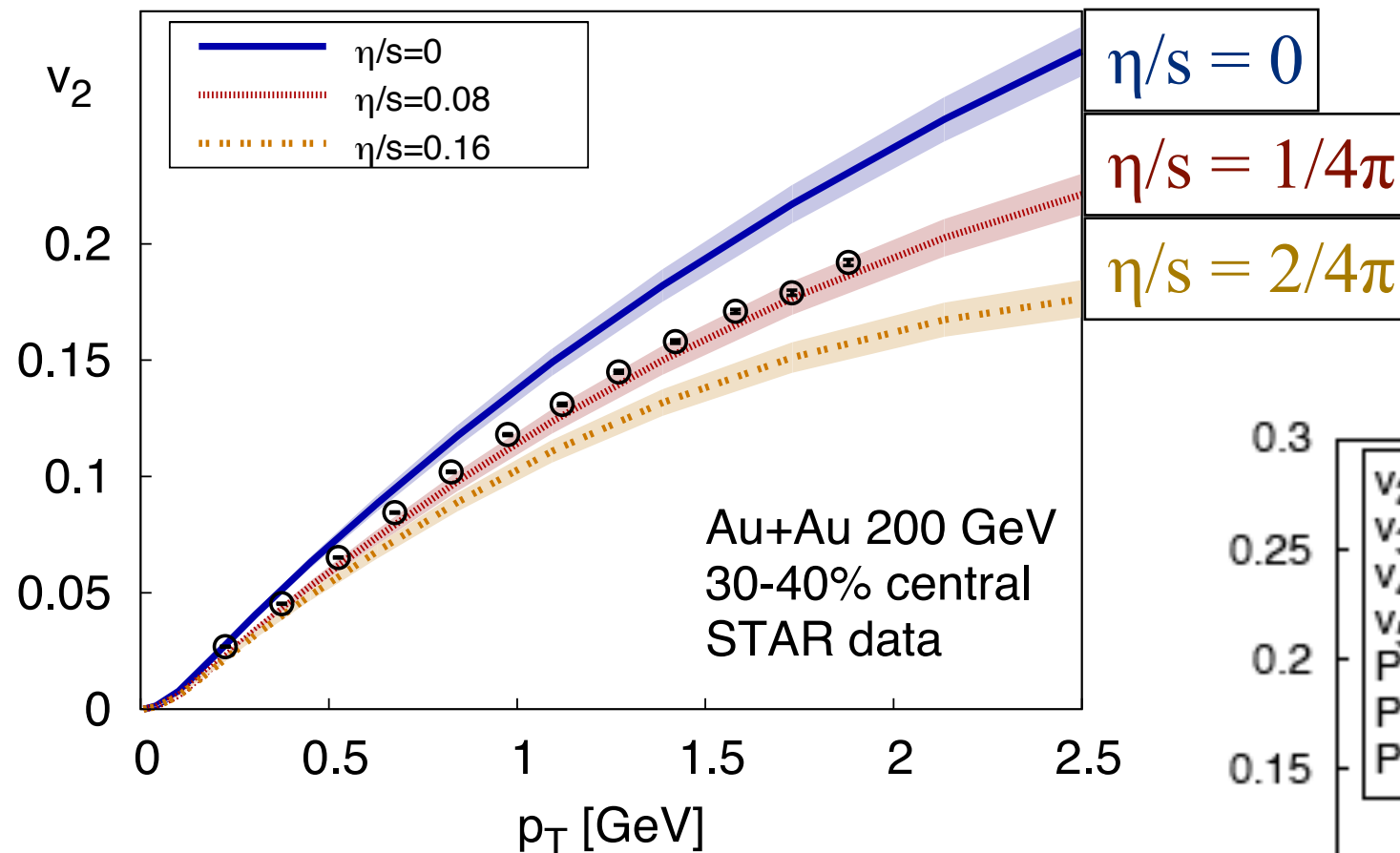


Idea: Energy density fluctuations in transverse plane from initial state quantum fluctuations. These thermalize to different temperatures locally and then propagate hydrodynamically to generate angular flow velocity fluctuations in the final state.

\Rightarrow flows $v_1, v_2, v_3, v_4, \dots$

Elliptic flow “measures” η_{QGP}

Schenke, Jeon, Gale, PRL 106 (2011) 042301

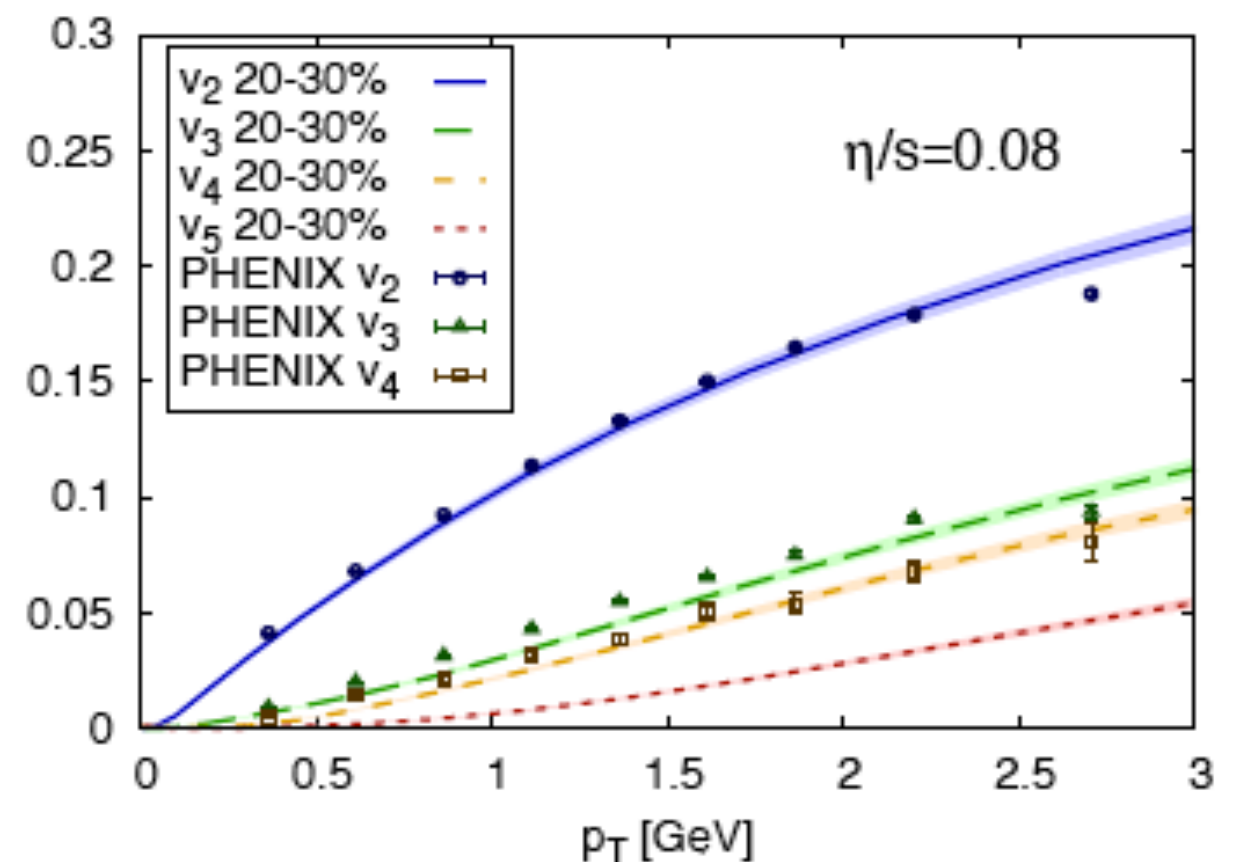


Schenke, Jeon, Gale, PRC 85 (2012) 024901

Universal strong coupling limit of non-abelian gauge theories with a gravity dual:

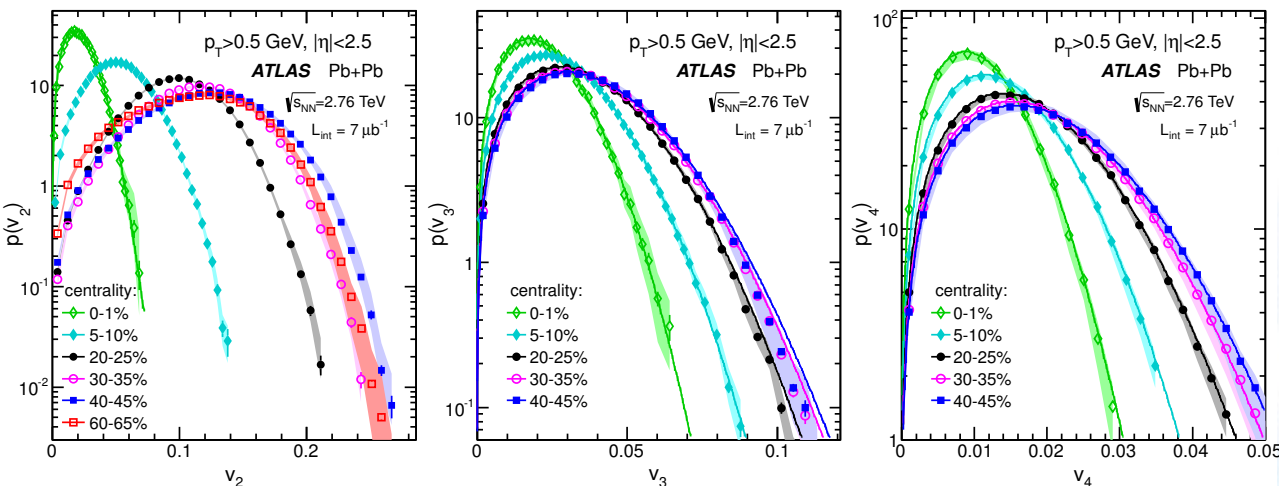
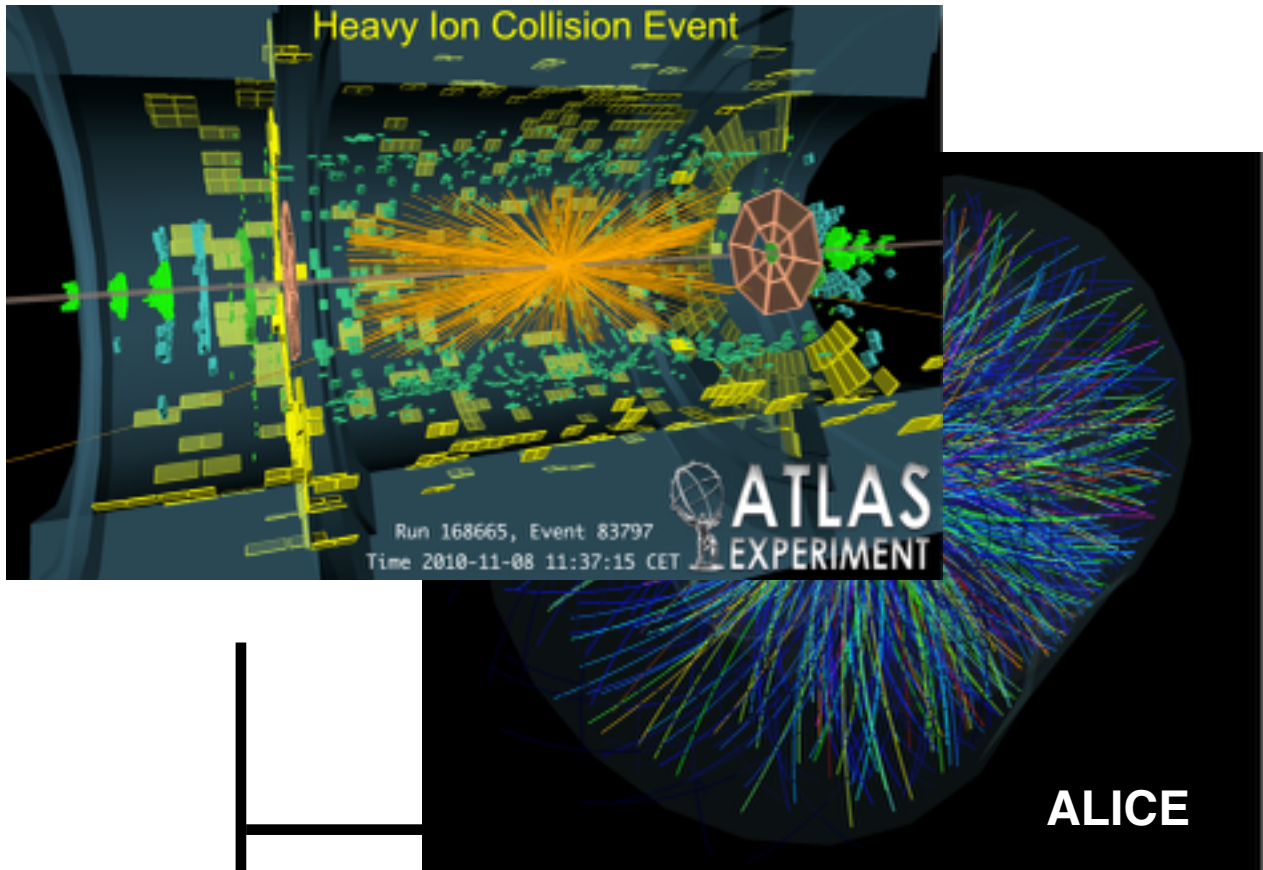
$$\eta/s \rightarrow 1/4\pi$$

aka: the “perfect” liquid

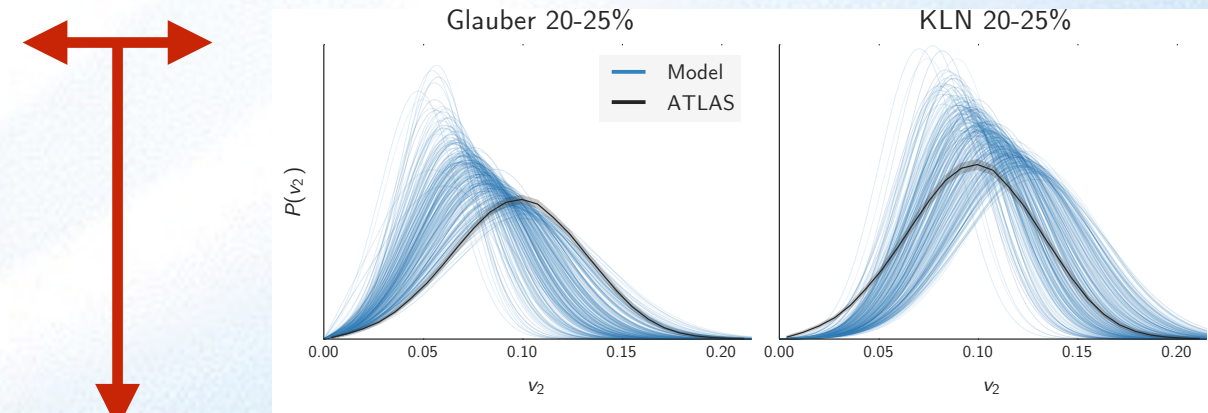
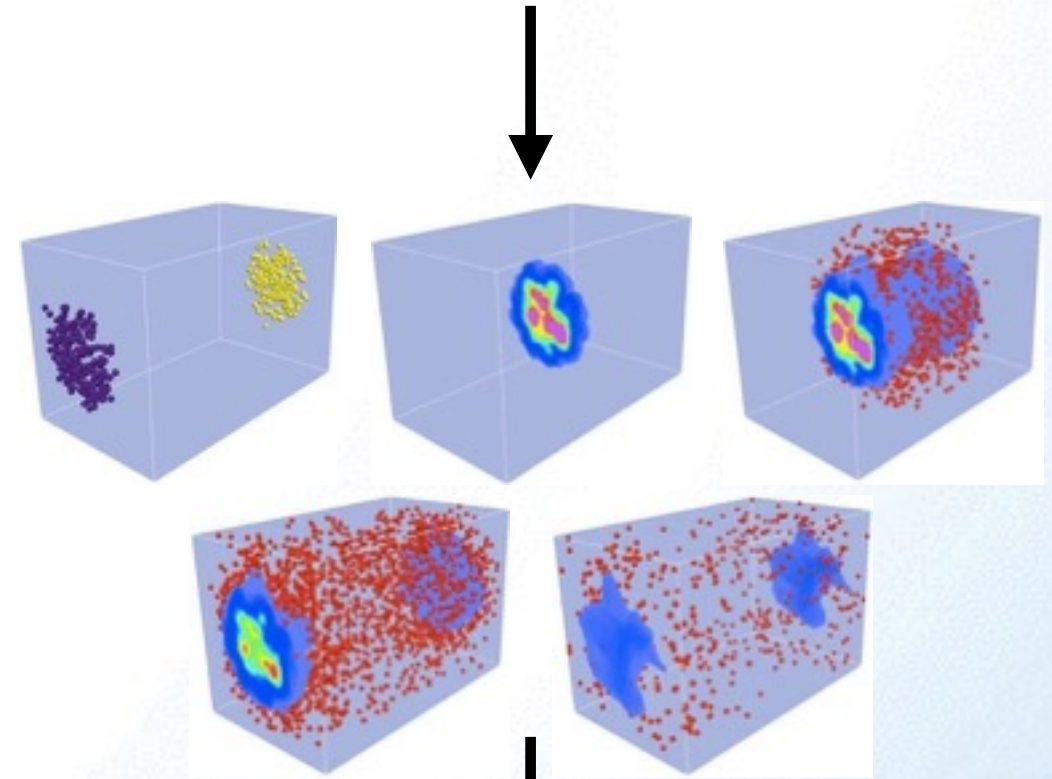


Discovery by model-data comparison

Data:

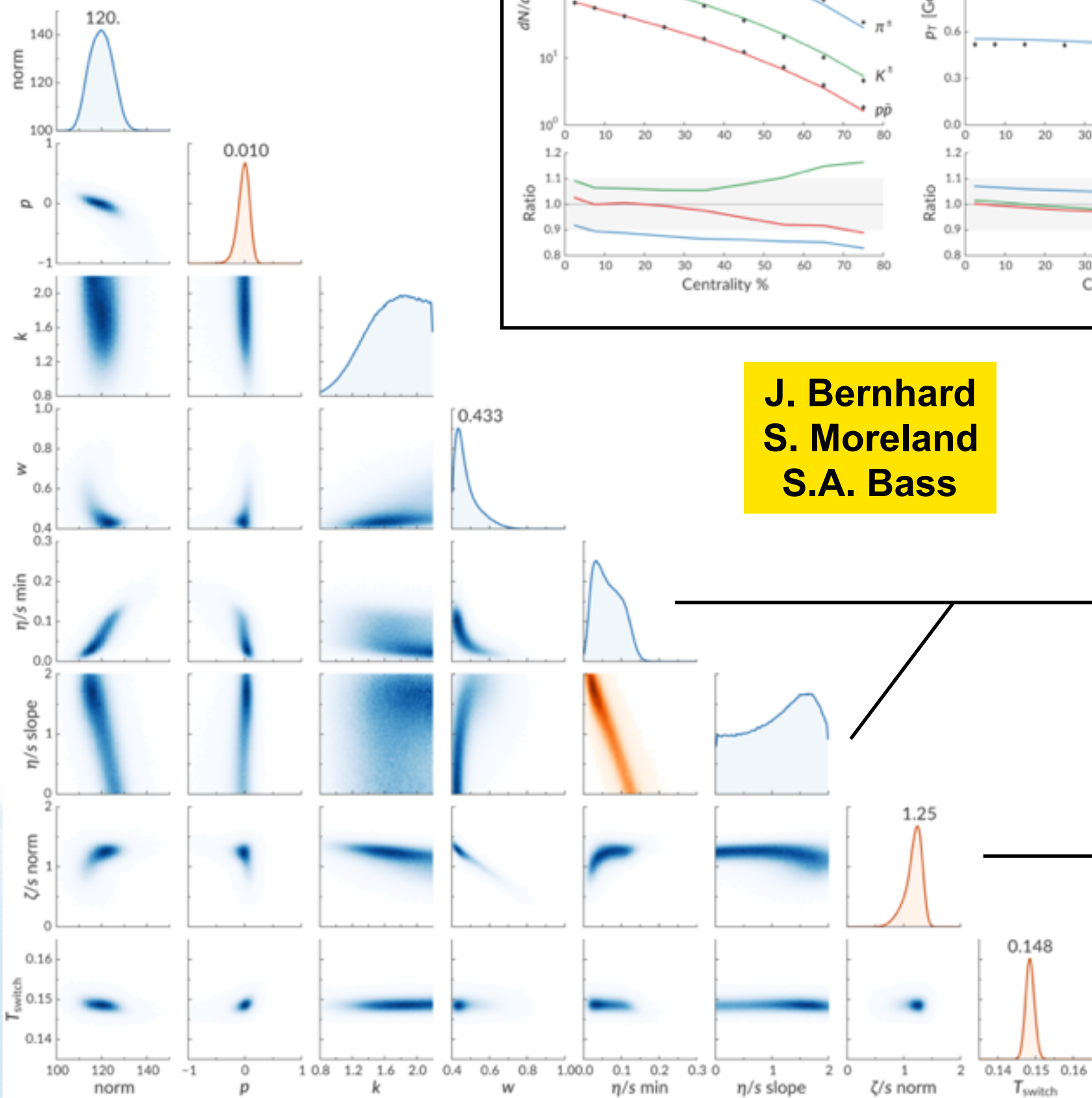


Model:
initial conditions, τ_0 , η/s , ζ/s ,

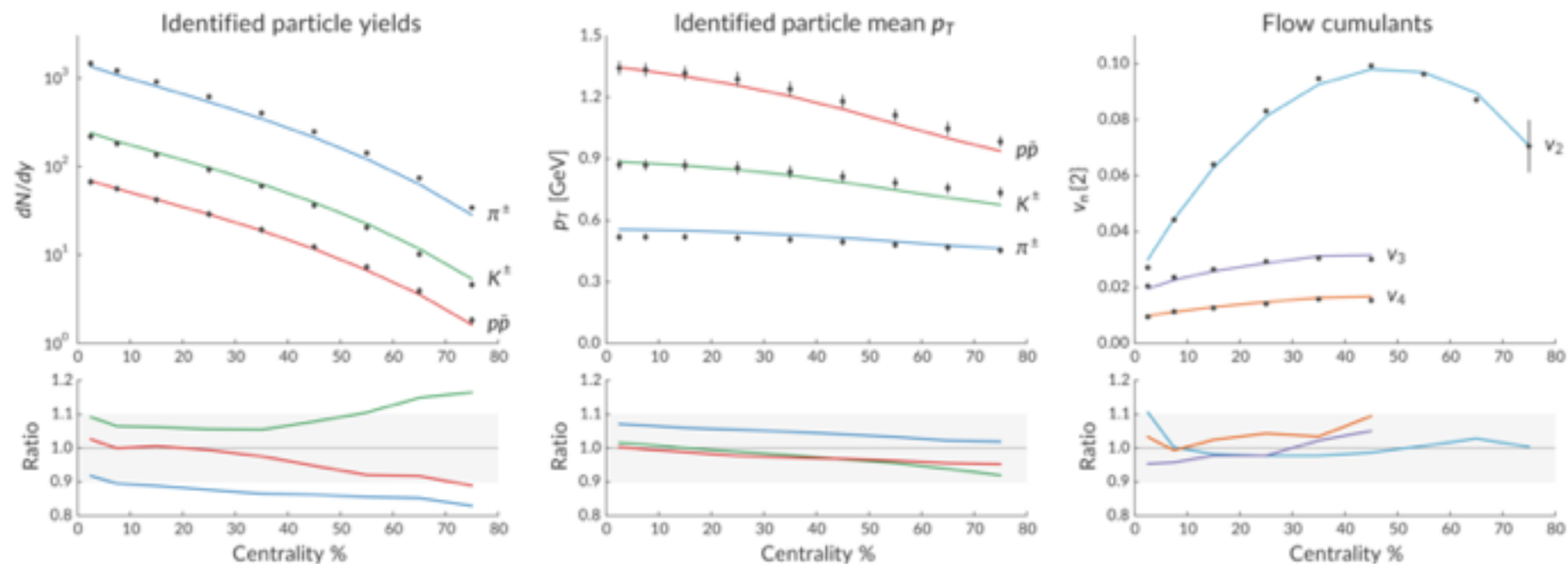


extracted QGP properties: η/s , ...

Calibrated posterior distributions:

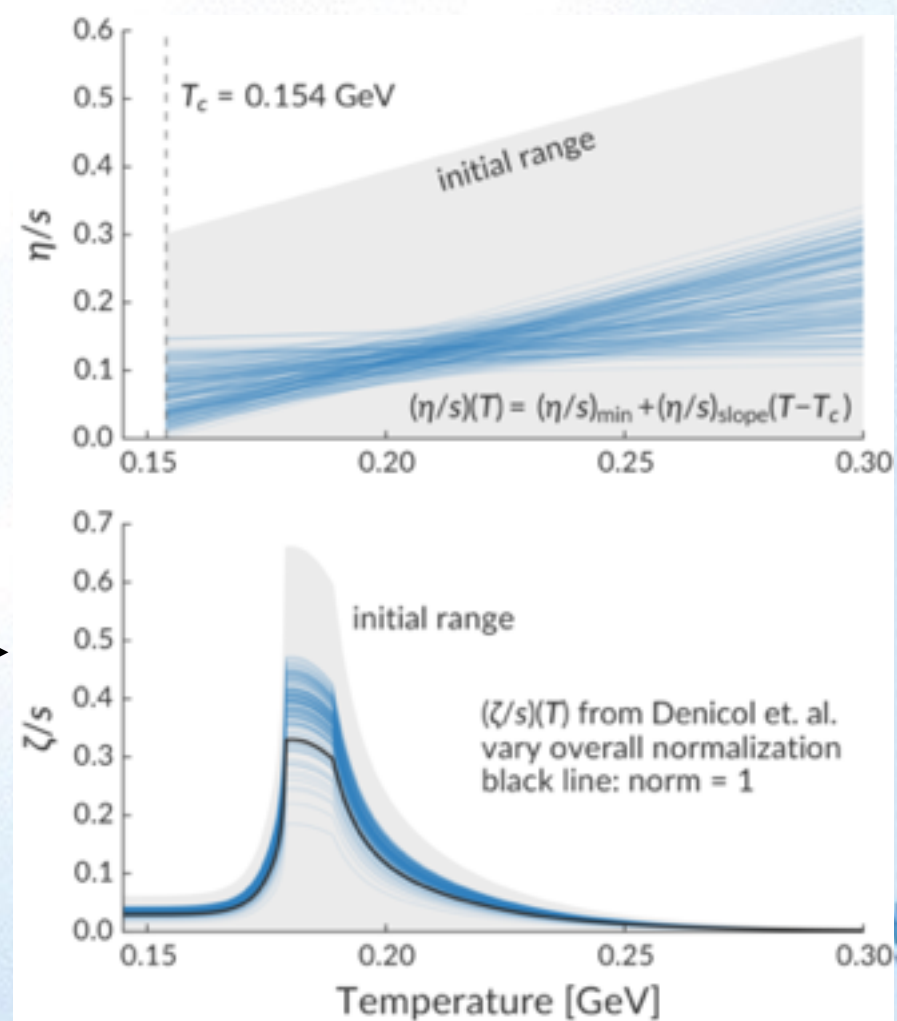


Explicit model calculations (no emulator):



J. Bernhard
S. Moreland
S.A. Bass

Temperature-dependent viscosities from the calibrated posterior:



Probing the Quark-Gluon Plasma

Hot QCD matter properties

Which **properties of hot QCD matter** can we hope to determine and how ?

Easy
for
LQCD

$$T_{\mu\nu} \iff \varepsilon, p, s$$

Equation of state: spectra, coll. flow, fluctuations

$$\eta = \frac{1}{T} \int d^4x \langle T_{xy}(x) T_{xy}(0) \rangle$$

Shear viscosity: anisotropic collective flow

Very
Hard
for
LQCD

$$\left. \begin{aligned} \hat{q} &= \frac{4\pi^2 \alpha_s C_R}{N_c^2 - 1} \int dy^- \langle U^\dagger F^{a+i}(y^-) U F_i^{a+}(0) \rangle \\ \hat{e} &= \frac{4\pi^2 \alpha_s C_R}{N_c^2 - 1} \int dy^- \langle i U^\dagger \partial^- A^{a+}(y^-) U A^{a+}(0) \rangle \\ \kappa &= \frac{4\pi \alpha_s}{3N_c} \int d\tau \langle U^\dagger F^{a0i}(\tau) t^a U F^{b0i}(0) t^b \rangle \end{aligned} \right\}$$

Momentum/energy diffusion:
parton energy loss, jet fragmentation

Hard
for
LQCD

$$\Pi_{\text{em}}^{\mu\nu}(k) = \int d^4x e^{ikx} \langle j^\mu(x) j^\nu(0) \rangle$$

QGP Radiance: Lepton pairs, photons

Easy
for
LQCD

$$m_D = -\lim_{|x| \rightarrow \infty} \frac{1}{|x|} \ln \langle U^\dagger E^a(x) U E^a(0) \rangle$$

Color screening: Quarkonium states

Future of RHIC

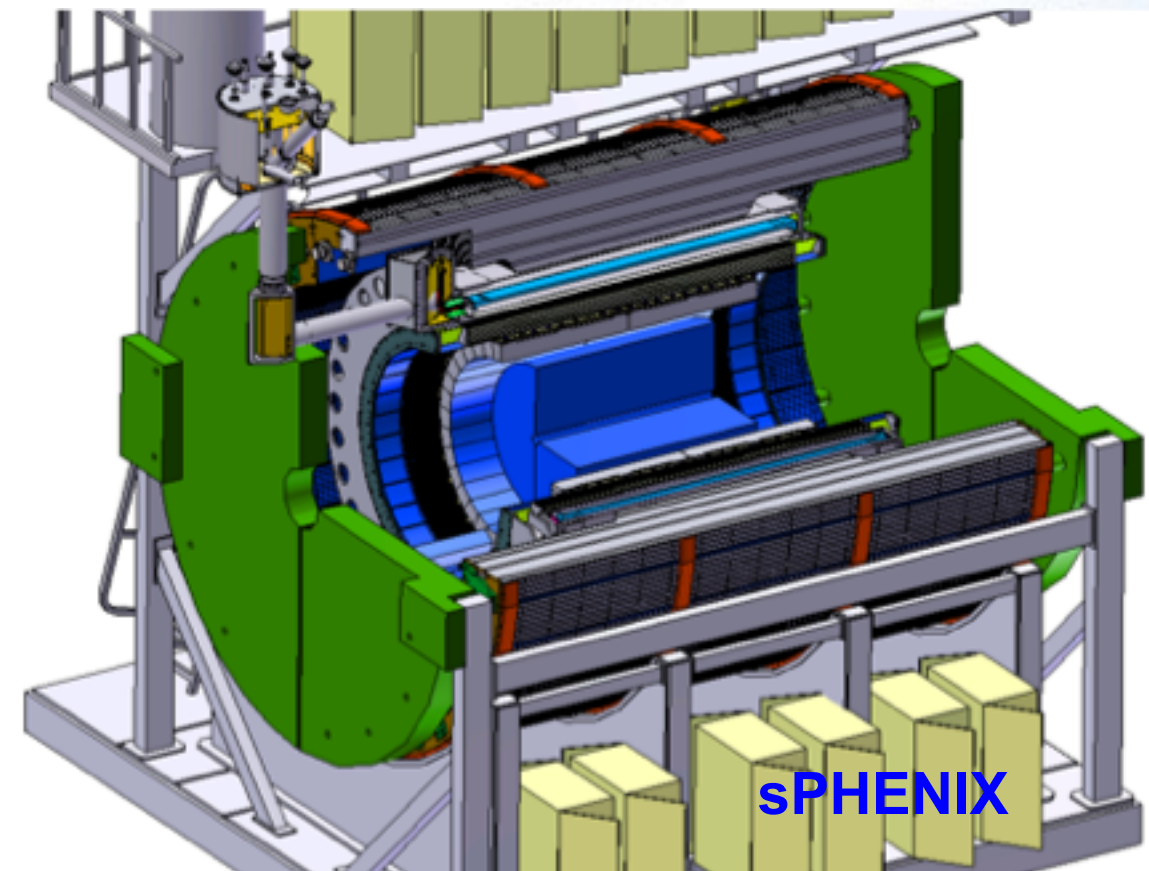
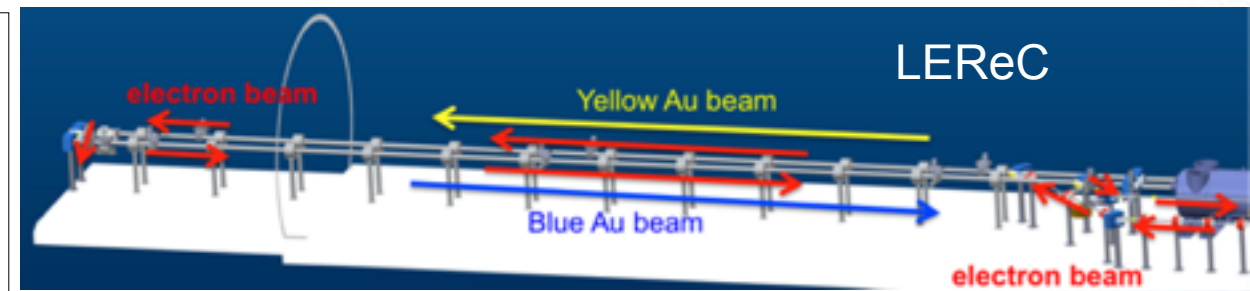
Completing the RHIC science mission

Status: RHIC-II configuration is complete

- Vertex detectors in STAR (HFT) and PHENIX
- Luminosity reaches 25x design luminosity

Plan: Complete RHIC mission in 3 campaigns:

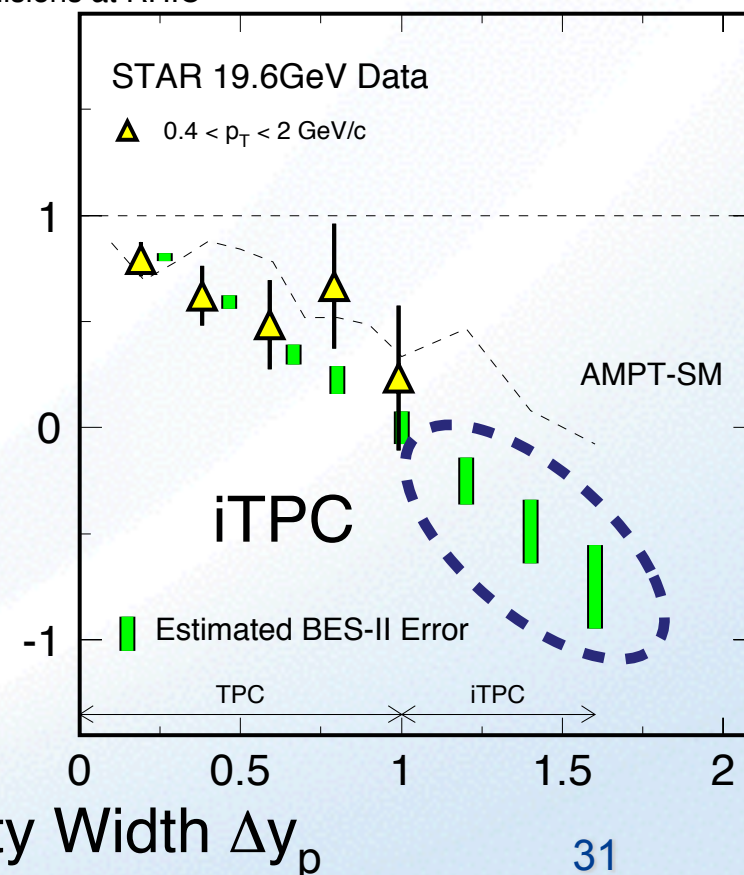
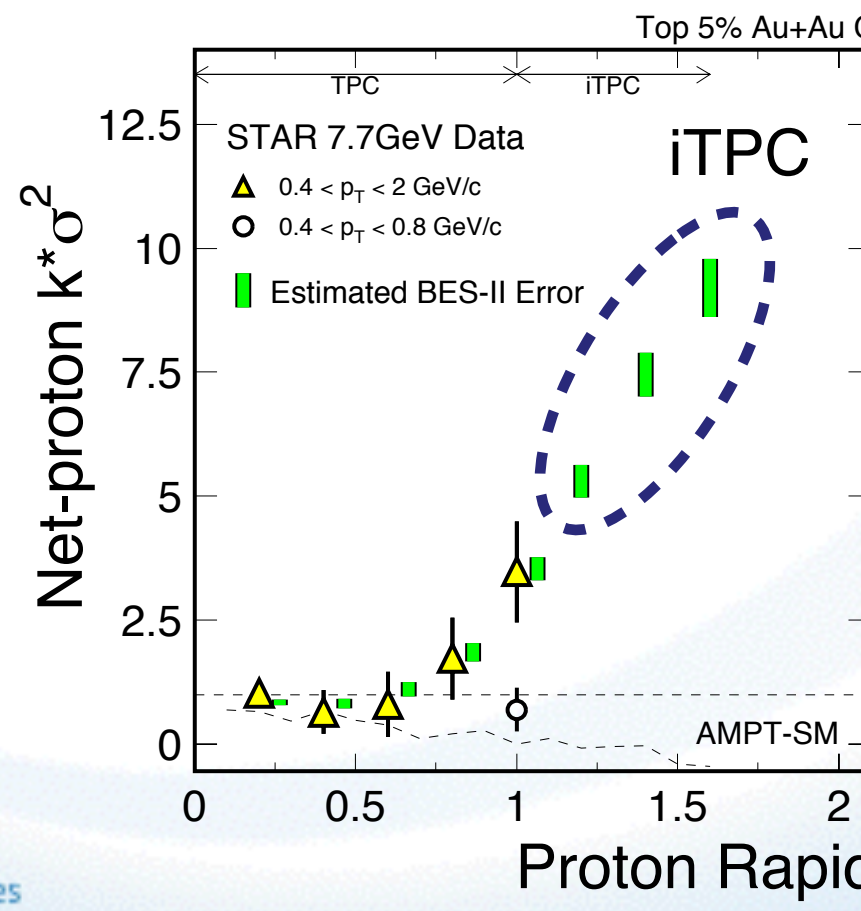
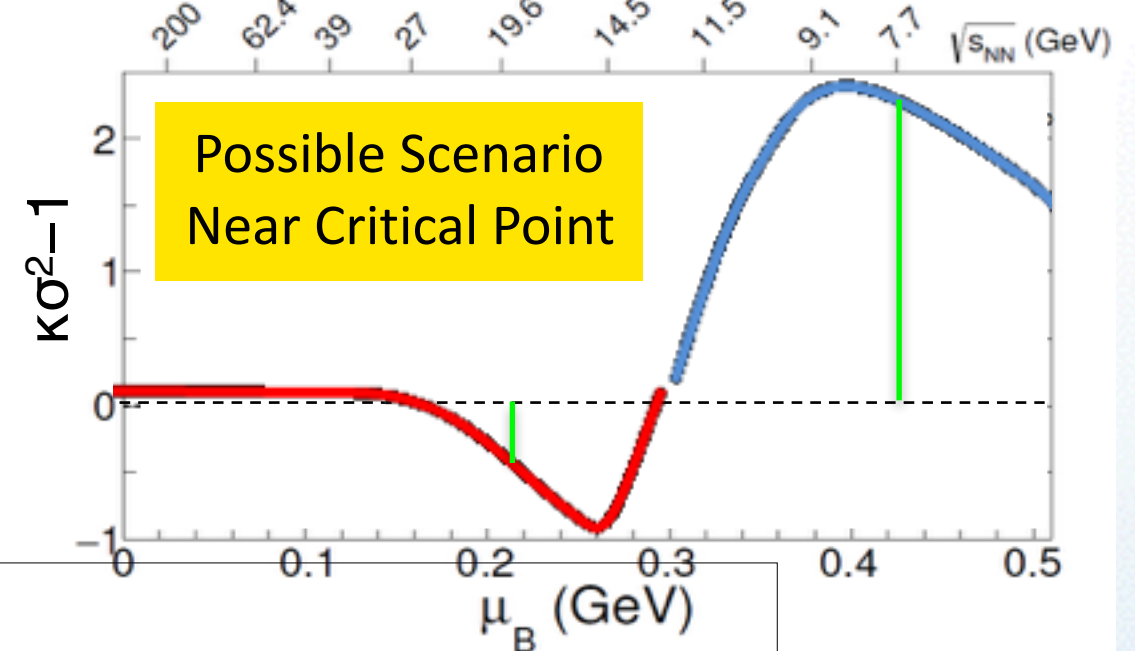
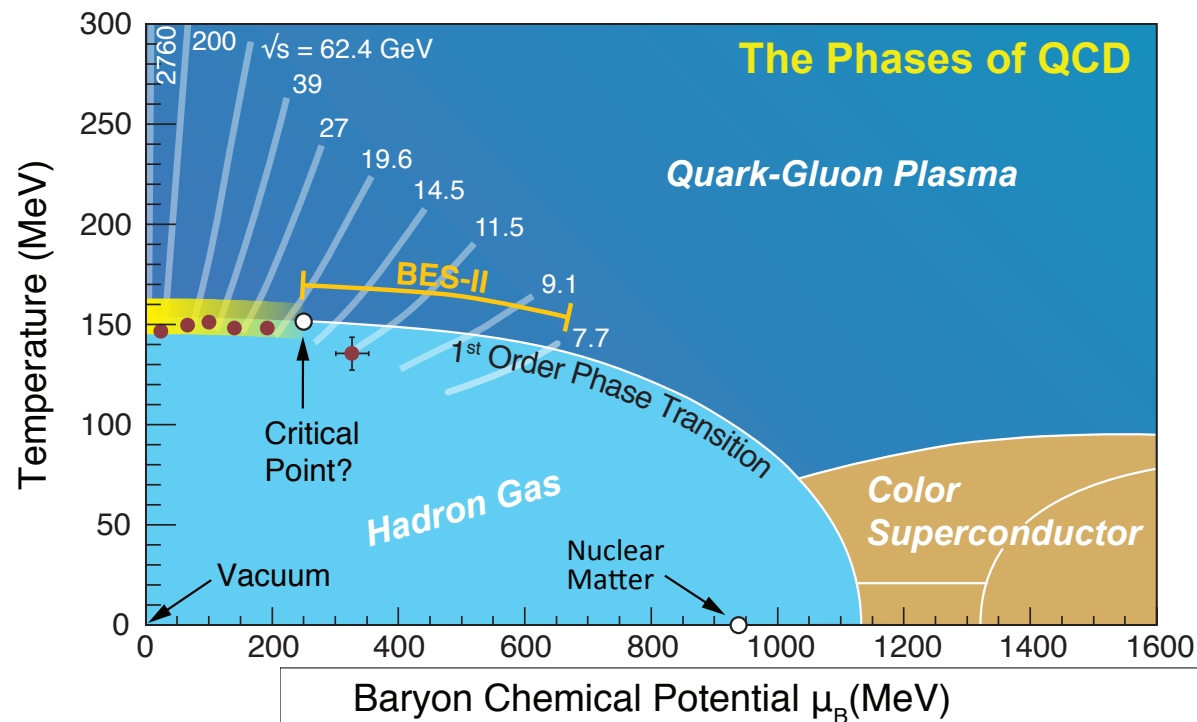
- 2014–17: Heavy flavor probes of the QGP using the micro-vertex detectors;
Transverse spin physics
- 2018: Install low energy e-cooling
- 2019/20: High precision scan of the QCD phase diagram & search for critical point
- 2021: Install sPHENIX
- 2022-23: Probe perfect liquid QGP with precision measurements of jet quenching and Upsilon suppression
- Transition to eRHIC?



RHIC remains a unique discovery facility

Critical fluctuations in BES-II

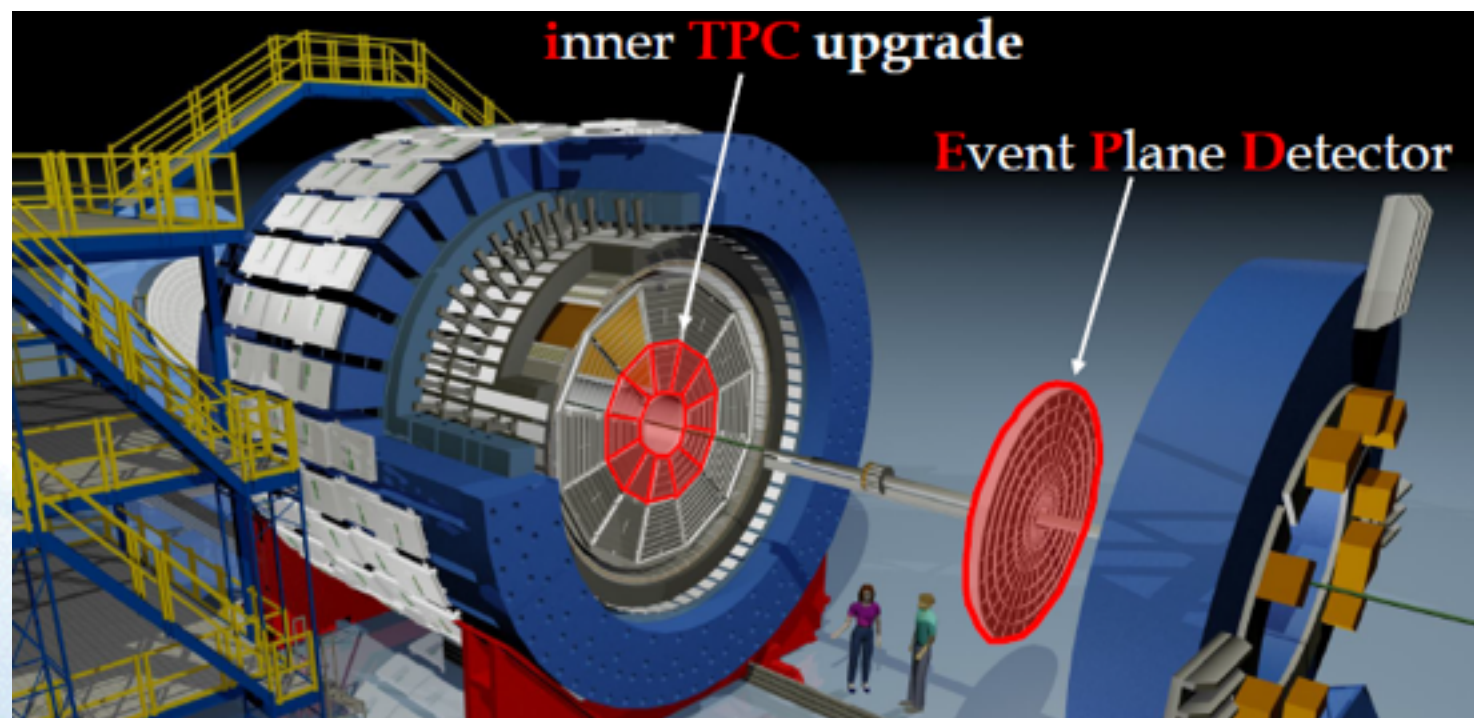
Model independent structure of
net baryon number kurtosis



STAR Upgrades for BES II

iTPC upgrade (2018)

Replace inner TPC Sectors
Extend rapidity coverage
Better particle ID
Low p_T coverage

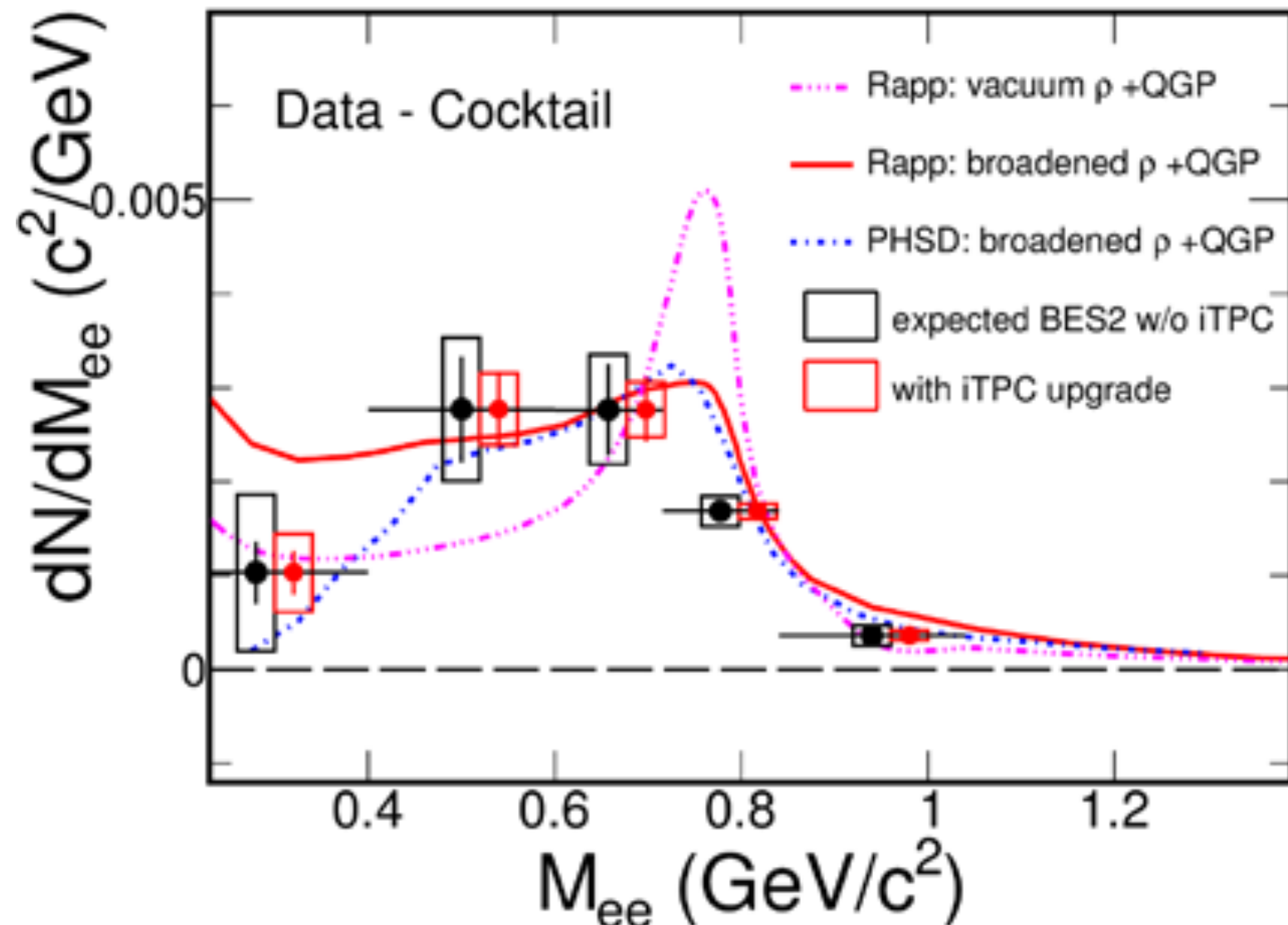


Event Plane Detector (2018)

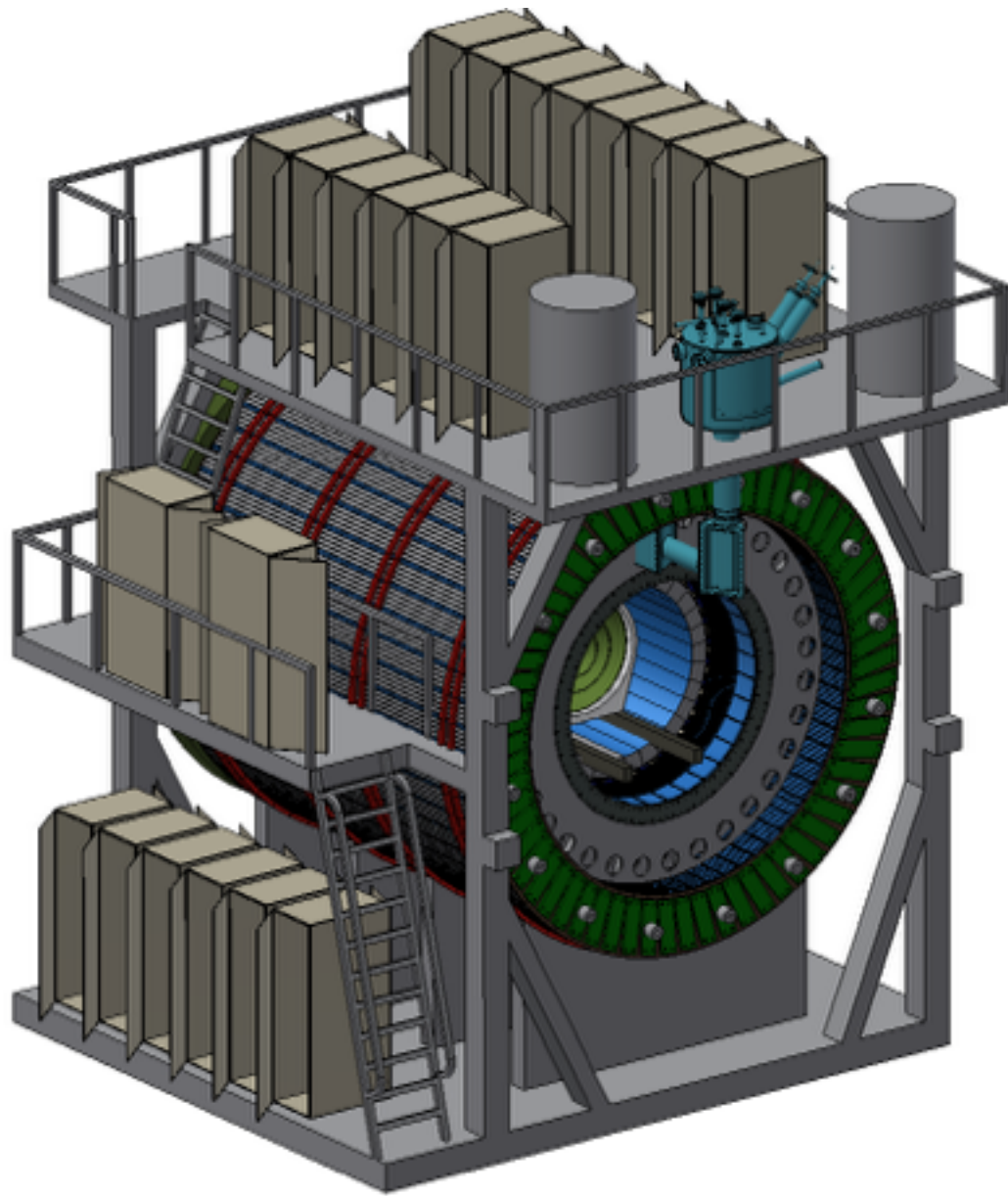
Improved Event Plane Resolution
Centrality definition
Improved trigger
Background rejection

Photon spectral function

The iTPC will also enable a more precise measurement of the photon spectral function at masses below 1 GeV

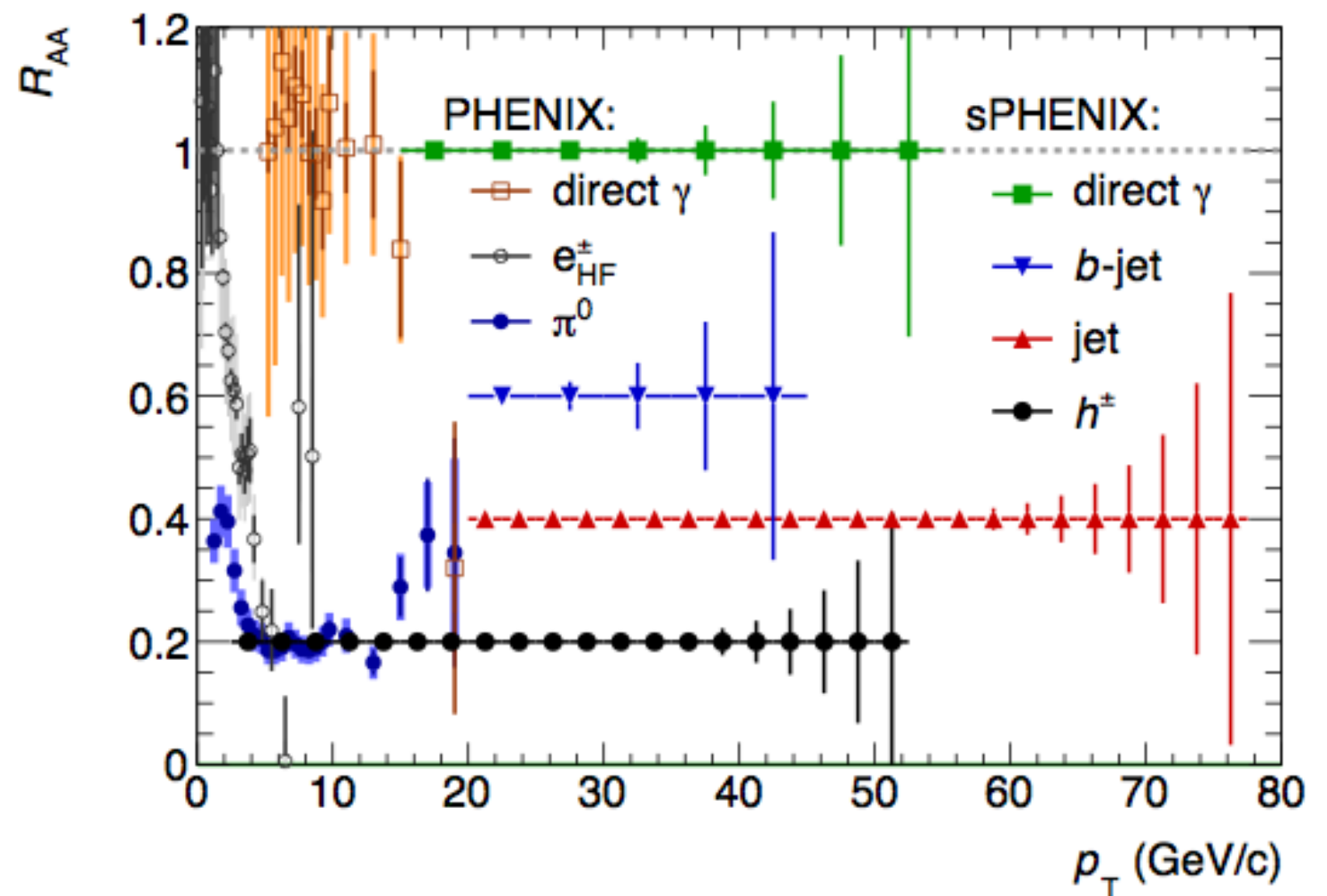


sPHENIX



Built around the BaBar solenoid –
undergoing testing at BNL

A large-acceptance detector to make precision measurements of hard probes (jets and Upsilon states) of the QGP at strongest coupling (near T_c). Replacing PHENIX and upgradable to EIC detector.



Completing the RHIC science mission

- **A unique forefront science program with tremendous discovery potential that is ONLY possible with RHIC:**
- **Quantify the transport properties of the QGP *near* T_c using heavy quarks as probes (together with LHC)**
- **Measure gluon and sea quark contributions to proton spin and explore coupled momentum-spin dynamics of QCD**
- **High statistics map of the QCD phase diagram, including possible discovery of a critical point**
- **Probe internal structure of the *most liquid* QGP using fully reconstructed jets and resolved Upsilon states as probes (together with LHC)**
- **Refine the physics program of an EIC with studies of *polarized* pp and pA collisions in forward kinematics**
- **RHIC enabled R&D to retire major risks of eRHIC design**