# Non-Gaussian cumulants of conserved charge fluctuations

Sayantan Sharma

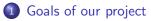


Project Members: Frithjof Karsch, Swagato Mukherjee, Hiroshi Ohno, Peter Petreczky and Patrick Steinbrecher

USQCD All Hands Meeting 2016.

Sayantan Sharma USQCD All Hands Meeting 2016, BNL

# Outline

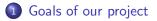






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# Outline



2 Recent developments

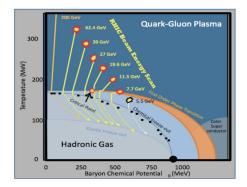


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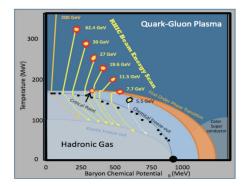
### QCD Phase Diagram: Status



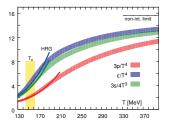
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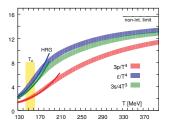


- QCD phase diagram largely unexplored inspite of intense efforts in last decades. → major focus of the Beam Energy Scan phase II experiments planned at the Relativistic Heavy Ion Collider at BNL.
- Themes: Existence of critical end-point, hydrodynamic modeling of the QCD medium formed in the experiments to understand experimental results.



- Lattice studies have given an Equation of state in QCD at  $\mu_B = 0$  in the continuum limit. [HotQCD and Budapest-Wuppertal Collaboration]
- $T \sim 140$  MeV, QCD can be described as Hadron Resonance Gas model but near chiral crossover  $T_c \sim 154$  MeV, HRG picture breaks down.

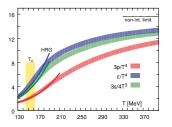
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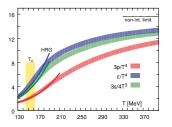
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- What happens to the HRG picture at finite density like the QCD medium formed in the experiments?
- Can we bracket the position of the critical end-point in QCD phase diagram?
- Can we understand the critical behaviour due to the light quarks in the crossover region?

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•  $\chi_6^B$  can already constrain QCD pressure in the regime approximated by Hadron Resonance gas model.

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# Challenges

- The Baryon no. susceptibilities can be expressed in terms of Quark no. susceptibilities (QNS).
- QNS  $\chi_{ij}$ 's can be written as derivatives of the Dirac operator. Example : $\chi_2^u = \frac{T}{V} \langle Tr(D_u^{-1}D_u^{''} - (D_u^{-1}D_u^{'})^2) + (Tr(D_u^{-1}D_u^{'}))^2 \rangle$ .  $\chi_{11}^{us} = \frac{T}{V} \langle Tr(D_u^{-1}D_u^{'}D_s^{-1}D_s^{'}) \rangle$ .

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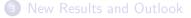
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- Extending to higher orders?
  - Matrix inversions increasing with the order
  - Delicate cancellation between a large number of terms for higher order QNS.

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# Outline







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#### A new method to introduce $\mu$

• The staggered fermion matrix used at finite  $\mu$  [Hasenfratz, Karsch ,83]

$$D(\mu)_{xy} = \sum_{i=1}^{3} \eta_{i}(x) \left[ U_{i}^{\dagger}(y) \delta_{x,y+\hat{i}} - U_{i}(x) \delta_{x,y-\hat{i}} \right] + \eta_{4}(x) \left[ e^{\mu a} U_{4}^{\dagger}(y) \delta_{x,y+\hat{4}} - e^{-\mu a} U_{4}(x) \delta_{x,y-\hat{4}} \right]$$

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• One can also add  $\mu$  coupled to the conserved number density as in the continuum.

$$D(0)_{xy} - rac{\mu a}{2} \eta_4(x) \Big[ U_4^{\dagger}(y) \delta_{x,y+\hat{4}} + U_4(x) \delta_{x,y-\hat{4}} \Big] \; .$$

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# • Linear method: $D' = \sum_{x,y} N(x, y)$ , and D'' = D''' = D''' = 0

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- Linear method:  $\chi_n$  have additional zero-T artifacts.  $\rightarrow$  explicit counter terms needed for  $\chi_{2,4}$ , discussed in detail [Gavai & Sharma, 15]
- In Exp method: counter terms already at the Lagrangian level. We use this method for  $\chi_n^B$ , n = 2, 4.

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• For any order *n*, the artifacts  $\sim \mathcal{O}(a^{n-4})$ .

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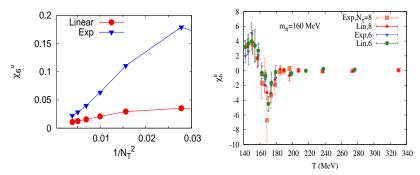
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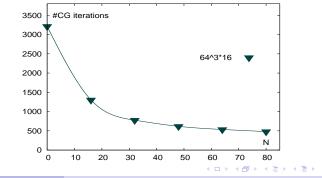
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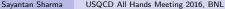
# Speeding up the inversions with deflation

• Calculating explicitly the lowest eigenvalues improves performance of the fermion inverter

 $D^{-1}|R\rangle = \sum_{i=1}^{N} 1/\lambda_i |\psi_i\rangle \langle \psi_i | R \rangle + CG$  Inversion.

- We have developed highly optimized codes based on Ritz and Lanczos algorithms for CPU's and GPU respectively.
- Current volumes  $N_s = 4N_{\tau}$ , already approaches a plateau for N = 80 for  $T \sim 145$  MeV. Typical N = 192 256.





#### Performance of our codes

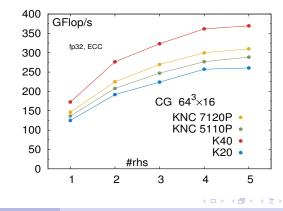
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- Currently highly optimized codes for Intel Knights Corner extended to Knights Landing. [Mainly led by Patrick Steinbrecher, graduate student since 15]



# Outline

Goals of our project

2 Recent developments

3 New Results and Outlook

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#### New data analyzed 2015-16

	$32^3  imes 8$		$24^3  imes 6$	
<b>T</b> [MeV]	$\beta$	# analyzed	$\beta$	# analyzed
135	6.245	104420	5.980	68000
140	6.285	104480	6.015	120790
145	6.315	107480	6.045	120770
150	6.354	108030	6.080	30080
155	6.390	108580	6.120	23546
160	6.423	119290	6.150	31164
165	6.445	122340	6.170	20000
170	6.474	141780	6.200	138470
175	6.500	142960	6.225	125280

• Results with physical quark mass

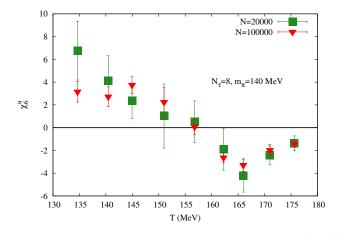
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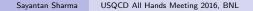
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- Results with physical quark mass
- Deflation + Multiple Right hand side technique+ special care of noisy operators → a speedup of 30 allowed for analysis of extensive set of configurations.

#### Main Outcome : Sixth order cumulants

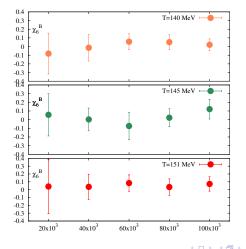
•  $\chi_6^u$  signal improved with the statistics *N*. We observe the negative dip just above  $T_c \rightarrow$  signal of O(4) criticality?





## Main Outcome : Sixth order cumulants

- Improvement visible already in χ<sup>B</sup><sub>6</sub> at the lowest temperatures by increasing number of configurations,
- We aim to increase statistics needed to reduce errors by a factor two.

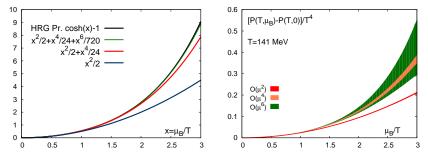


# Constraining EoS

 In a regime where Hadron Resonance gas is anticipated to be a good description of QCD, including χ<sup>B</sup><sub>6</sub> term already reproduces P(μ<sub>B</sub>) within 5% accuracy.

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- In a regime where Hadron Resonance gas is anticipated to be a good description of QCD, including χ<sup>B</sup><sub>6</sub> term already reproduces P(μ<sub>B</sub>) within 5% accuracy.
- Improve errors on our current data to observe this
  - $\rightarrow$  increase statistics twofold this year.



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#### Breakdown of HRG

• Breakdown of HRG+ onset of criticality can be already constrained with  $\chi_6^B$ .

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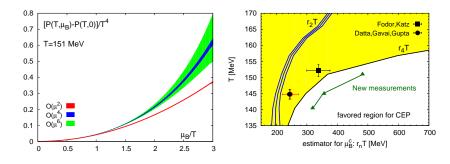
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- Near critical point all terms in the Taylor expansion nearly equal  $\rightarrow$  need to improve the errors!
- Our data gives a preliminary bound on the location of critical point from radius of convergence estimates,  $r_{2n} \equiv \sqrt{2n(2n-1)\left|\frac{\chi_{B_{12}}^B}{\chi_{B_{12}}^B}\right|}$ .

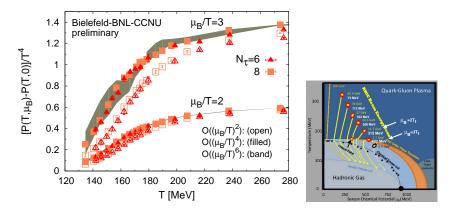


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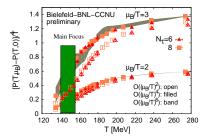
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### EoS away from criticality

• The pressure for T > 160 MeV which is an important input for the hydrodynamic modeling of the plasma already constrained by  $\chi_B^6$  even for highest  $\mu_B/T$ .

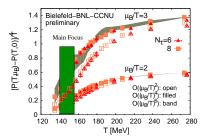


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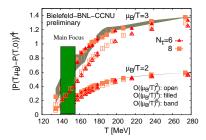
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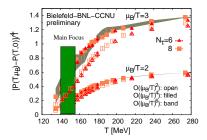
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- Improving statistics on  $\chi_6^B$  in the hadron phase will already improve the EoS.
- Analysis of  $\chi_8^B$  is also crucial to estimate the errors on the EoS measured with the sixth order cumulants.
- Higher order cumulants will also help in bracketing the possible QCD critical end-point which is one of the focus of BES-II experiments.