Nucleon structure study in domain wall fermion for physics beyond the standard model; chromo-electric dipole moment, proton decay operator, and (axial) vector form factor

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@ 2016 USQCD All-Hands Collaboration Meeting, April 29, 2016

Our goal

Understand nucleon structure and probe BSM physics

- (i) Nucleon vector and axial-vector form factors
- (ii) Proton decay matrix element

(iii) Electric dipole moments induced by quark chromo-electric dipole moments

We measure nucleon matrix elements using 2+1 flavor of chirally-symmetric domain wall fermions. For (i) form factor, and (ii) proton decay matrix elements,

we compute them at the physical point (m π =140MeV) on 48^3 x 96 lattice (Lm π =3.86).

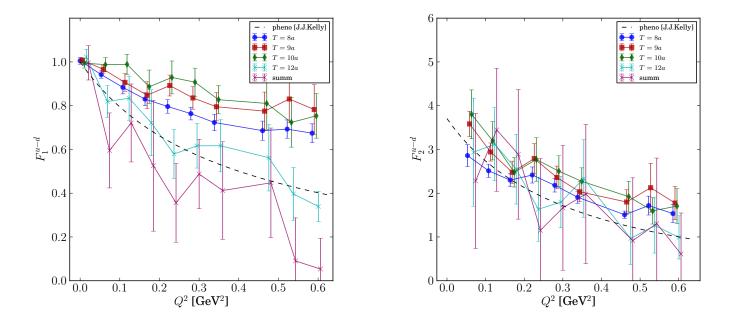
For (iii) we compute it on 32^3 x 64 lattice (m π ~ 170 MeV) and take a chiral extrapolation.

Gauge ensembles are generated by RBC/UKQCD collaborations.

Electromagnetic form factor of Nucleon

$$\langle N|V_{\mu}^{EM}|\bar{N}\rangle = F_1(q^2)\gamma_{\mu} + \frac{iF_2(q^2)}{2m_N}\frac{\sigma_{\mu\nu}}{2}q_{\nu}$$

2+1 Mobius DWF, Physical point, 48 x 96 (a=0.114 fm, Lmπ=3.86)

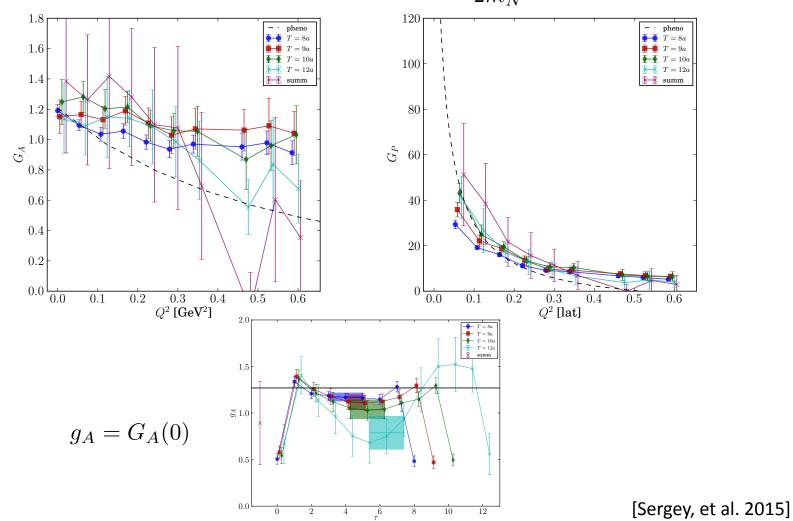


No disconnected diagram calculation, statistics is not good. A large excited state contamination is observed.

[[]Sergey, et al. 2015]

Axial form factor of Nucleon

 $\langle N|A_{\mu}^{EM}|\bar{N}\rangle = G_A(q^2)\gamma_{\mu}\gamma_5 + \frac{iG_P(q^2)}{2m_N}\gamma_5 q_{\nu}$



The axial charge radius is an important parameter for precision measurement of neutrino scattering off nucleons.

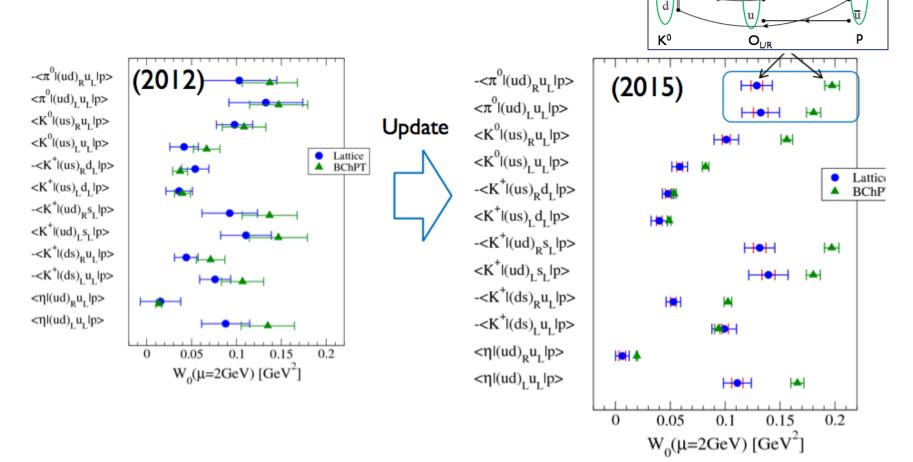
Statistical improvement for Nucleon form factor

Increase # of Eigenvectors, in all-mode averaging AMA, 500 [Blum, et al. 2013] -> 2,000 # of measurement per configs. 32 -> 32 x 4 (coherent) = 128 # of configs 20 -> 100 (We use coherent sequential trick [LHP, K. Orginos, et al.])

Goal : 4% statistical error on gA and 10% axial charge radius at the physical point. Having 3 different time separations between sink & source, we will carry out a careful study of the excited state contamination. Finite size correction should be sufficiently small compared to the target nucleon quantities ($m\pi L = 3.86$).

2. Proton decay matrix elements

[E. Shintani, et al.]



 $-\vec{p}$

 $t_{\rm s} = t_{\rm p} - t_{\rm PS}$

 $\vec{p} = 0$

 t_p

 $\frac{\overline{u}}{\overline{d}}$

 \vec{p}

 t_{PS}

Calculation done on $M\pi >= 300$ MeV, extrapolated linearly in quark mass and q^2 . Calculation @ M_{π} =140 MeV is proposed. The chiral bag model's prediction of orders of magnitude suppression at a small quark mass [A. Martin and G. C. Stavenga, 2012]

3. Nucleon EDM matrix elements

We will calculate these nucleon matrix element for CP-odd operators

$$\mathcal{L} = d^{\text{qEDM}} \bar{q} (\sigma \cdot F) \gamma_5 q + d^{\text{cEDM}} \bar{q} (\sigma \cdot G) \gamma_5 q + \cdots$$
$$\mathcal{L} = \frac{\bar{\theta}}{64\pi^2} G\tilde{G}, \qquad \bar{\theta} = \theta + \arg \det M$$

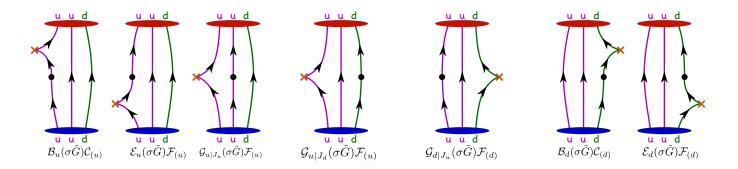
Our primary target : chromo EDM (cEDM) nucleon matrix element

The cEDM is a low-energy effective interaction which is induced by BSM physics The (c)EDM may be in the discovery reach of future EDM measurement experiments.

DWF : Chiral symmetry is very important No mixing with mass-proportional chromo-magnetic dipole (clover) term c.f. Wilson fermions We will calculate 4-point function with cEDM operator and electromagnetic current (9 different types of diagrams)

$$C_{4pt}(T,t,0) = \langle N(T) | J^{EM}_{\mu}(t) \left(\int d^4x \bar{q} G_{\nu\rho} \sigma^{\nu\rho} \gamma_5 q \right) | N(0) \rangle$$

$$\langle u_N(p')|J^{\mu}|u_N(p)\rangle_{CP} = \gamma^{\mu}F_1(q^2) + \sigma^{\mu\nu}q_{\nu}\frac{F_2(q^2)}{2m_N} + i\gamma_5\frac{\sigma^{\mu\nu}q_{\nu}}{2m_N}F_3(q^2)$$



We use sequential-source for propagators for cEDM insertions (X) and EM current (.).

For cEDM three- and four- point functions, we have cross-checked against background method with quark propagator

$$D^{-1} \to (D + \epsilon \mathcal{O})^{-1} \sim D^{-1} - \epsilon D^{-1} \mathcal{O} D^{-1}$$

Preliminary study : CP mixing angle (α) induced by cEDM operator

$$\sum_{s} u_N(p,s)\bar{u}_N(p,s) = E_N\gamma_0 - i\vec{p}\vec{\gamma} + m_N e^{i\alpha\gamma_5}$$
$$\operatorname{Tr}\frac{1+\gamma_t}{2}\gamma_5 \left\langle N(T)\sum_{x} \mathcal{O}_{\mathrm{cEDM}}(x)\bar{N}(0) \right\rangle \propto \alpha$$

 $T = 8a \approx 1.13 \text{ fm and varied the time coordinate of the cEDM insertion (see Fig. 2). Sur$ prisingly, gauge field smearing does not have any effect on the gauge noise; we hope that**CEDM operator insertionatime**(**t**toedml) dependence of**Nucleons 2ps**for the stochastic**Left: HEX smearing, Right:**Will so approximately 12% with only 512 sloppy samples on 16 gauge configutions. An analogous study is underway for the cEDM operator with the application of**t_source=0, t_sink=8**gradient (Wilson) flow for a range of flow times, from which we will select the three valuesthe flow time for the renormalization study. Preliminary study : CP mixing angle (α) induced by cEDM operator

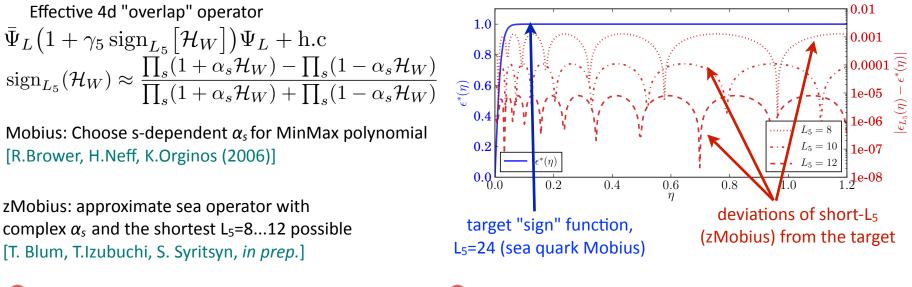
proton_pos5,U

Wilson flow time dependence of total result and its error with the sum of the cEDM insertion (see Fig. 2). Sur-Ongoing calculation: currentillsQCD1 allocation for a range of the operators smeared with a wide range of 2+1 DWF, 32x64, 170 M gradient flow times will not suffer from increased noise. It is reassuring that the stochastic Study of the renormalization schemes using Wilson who will select the three values the flow time for the renormalization study.

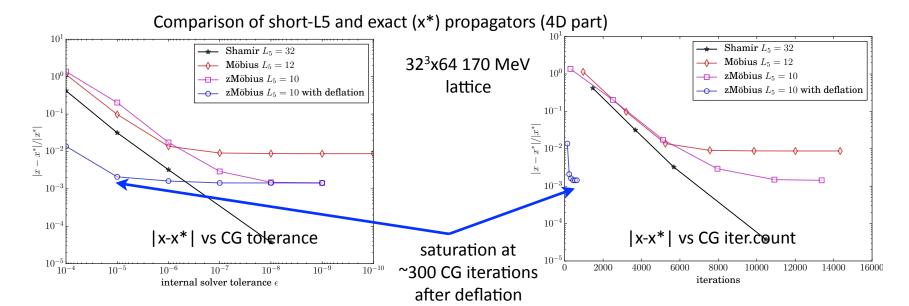
zMobius acceleration of chiral fermions

cheaper AMA approximation

smaller eigenvectors for 5D deflation



acceleration of exact solve through iterative refinement ("zMADCG")



Cost of One Solve zMobius vs Overlap

- » On Physical point ensemble, 48 cube DWF,
 - » ~ 20 k CG iteration with Ls=24 DWF
 - » Becomes ~220 CG iteration with Ls=10 zMobius, sloppy solve
 - » This is roughly equivalent to <u>5.6 k Wilson-mult, or 2.8 k Wilson CG iteration</u>
 - One overlap ~ 600 Wilson polynomial to construct. Cost of sloppy zMobius solve is equivalent of 5.6 / (2*0.6) = 4.7 outer CG of overlap.
 If low-precision outer CG of overlap is 50, zMobius is more than 10 times faster.
- » We will consider low mode substitutions of ChiQCD to further reduce cost, thanks to SPC

Request: Summary table

Table 1: CPU costs for cEDM operator based on the current calculation.

$32^3 \times 64$	Count	M Jpsi core*hours
Propagators, sloppy	888	0.100
Propagators, exact	88	0.0528
TOTAL per cfg		0.153
TOTAL	$\times 150$	15.3

Table 2: CPU costs for nucleon axial vector form factors based on current CPS tests.

Count	M Jpsi core*hours
141	0.151
13	0.0697
	0.221
$\times 100$	22.1
	141 13

Table 3: CPU costs for proton decay matrix elements based on current CPS tests.

$48^3 \times 96$	Count	M Jpsi core*hours
Propagators, sloppy	224	0.0762
Propagators, exact	7	0.0119
TOTAL per cfg		0.0881
TOTAL	$\times 100$	8.81

Total computing: 46.21 M Jpsi core hours on Fermilab clusters

512 TB tape storage, equivalent to 1.84 M Jpsi core hours at Fermilab 64 TB disk storage,

equivalent to 2.56 M Jpsi core hours at Fermilab

(50.61 M Jpsi core hours total computing and storage)

16 % of proton decay calculation will be removed due to overlapping

(See SPC questions and our answers)

Total request change: 50.61 -> 47.4 [M Jpsi core hours]

SPC questions

1) In your proposal you specify the statistical uncertainties you expect to obtain for gA and the Dirac radius. What are your expectations for the systematic errors?

Our central value for pilot calculation for gA on the physical point 48 cube ensemble turns out to be lower the experimental values. The plateau for gA indicates significant excited state contamination, while the value from the summation method is consistent with experiment but with a larger statistical error (50%).

About the charge radii, the statistics are not sufficient to analyze the form factor in particular with larger time separation. In both cases, additional statistics will be needed to reduce their uncertainties, which also enables us to carry out a careful analysis of excited state contamination and other systematics. Volume is one of the largest existing lattice (48^3 x 96), which corresponding to $m\pi L = 3.86$, which should be sufficient to suppress finite volume effects to below the target statistical error for the nucleon quantities. As for the discretization error, we have 64^3 already generated years ago, which we could use to remove the systematic error from discretization, which would be small for DWF, in following calculations in next years.

For proton decay matrix element, the calculation on the physical quark mass has a strong motivation due to the chiral bag model's prediction of orders of magnitude suppression at a small quark mass [5]. Thus, calculations at the physical quark mass would remove the claimed huge uncertainty due to chiral fit and extrapolations.

2) In your proposal you request resources to calculate EDMs, Nucleon form factors, and proton decay matrix elements. How much overlap is there in the computation of the three different quantities? How would the computational cost for this project change, if you were to perform the calculations sequentially?

16 % of proton decay calculation will be overlapping with nucleon form factor calculation. We like to emphasis that this overlap is besides the shared cost that already computed eigenvectors at Argonne g-2 calculation, which we will use to leverage both the proton decay and the form factor calculations. 3) Since you are planning to calculate gA for which there already exists an accurate measurement, have you considered performing a "blind analysis" to prevent any inadvertent bias? To blind your analysis, you could add an overall off-set factor to the correlation functions that would be kept unknown to the people doing the analysis until the systematic error analysis is finalized.

We would certainly consider the blind analysis. In fact, we have already been exercising the blind analysis in the V_us determination of tau-inclusive decay to remove possible human prejudices (Blum's proposal).

4) With the new resources at JLab being as yet unspecified, we would like to know if you are in a position to use them efficiently if they are a) cpu, b) GPU, c) KNL. If you are not, that is fine, but it will help in our allocation decisions to know this information from every proposal.

Our main measurement program is based on Qlua code, and we are ready to use CPU resources, at J-lab. For b) GPU, c) KNL, we prefer KNL due to the existing efficient code based on Grid of Peter Boyle. For deflation, we would prefer to have 96GB/node for KNL.

5) The SPC would like you to explore possibilities for coordinating or collaborating with the chiQCD collaboration which also has a hadron structure program on the same or similar lattices.

We have a concerns about the potential systematic error due to the partially quenched effect in chQCD's overlap, whose Dirac kernel may be significantly different from that of sea quark. We note that the 4D kernel inside the approximated sign function of sea quark action (Mobius parameter is set, b-c = 1) is different from that in overlap (Neuberger) (b-c=0), so the unitarity violation may not be removed simply by adjusting quark mass at the finite lattice spacing.

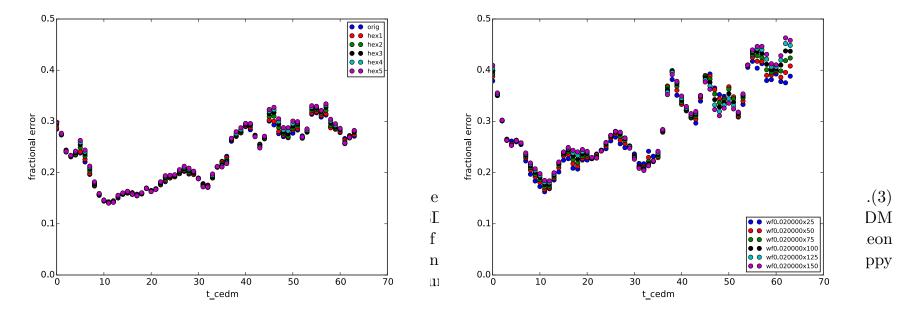
Thanks to SPC's suggestion we started productive conversations with Keh-Fei Liu. At this point, both of parties strongly feel having two independent strategies and calculations would be better and healthy especially considering about relatively premature states of nucleon matrix element calculations compared to meson calculations. It would be good to try further various explorations of methods as well as different quark discretizations (MDWF or Overlap). We do, however, think it would be very beneficial to learn from each other, especially about methods, to maximize the outcome of the precious USQCD resources.

Thank you

Backup

Preliminary study : CP mixing angle (α) induced by cEDM operator

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cEDM operator insertion insertion time is will also hold for the gradient flow and the operators smeared with a wide range of the stochastic t_source=0, s_sink=8 $T = 8a \approx 1.13$ fm and varied the time coordinate of the cEDM insertion (see Fig. 2). Surprisingly, gauge field smearing does not have any effect on the gauge noise; we hope that this will also hold for the gradient flow and the operators smeared with a wide range of the stochastic precision for α_u is approximately 12% with only 512 sloppy samples on 16 gauge configurations. An analogous study is underway for the cEDM operator with the application of gradient (Wilson) flow for a range of flow times, from which we will select the three values the flow time for the renormalization study.