

Based on V. Guzey, M. Rinaldi, S. Scopetta, M. Strikman, M. Viviani, PRL 129 (2022) 242503

Outline:

- Nuclear shadowing
- Coherent J/ ψ production in eA scattering on He-4 and He-3
- Summary and Outlook

Nuclear shadowing

• Nuclear shadowing is general phenomenon of high-energy scattering \rightarrow nuclear cross section < sum of nucleon cross sections.

• Present in soft hadron-nucleus scattering, Bauer, Spital, Yennie, Pipkin, Rev. Mod. Phys. 50 (1978) 261 and hard processes with nuclei, H. Paukkunen, talk this morning.

• In the target rest frame, understood as multiple interactions of projectile with target nucleons \rightarrow destructive interference of amplitudes with N=1,2,..A nucleons \rightarrow Gribov-Glauber theory of shadowing, Glauber, PRD 50 (1955) 242; Gribov, Sov. Phys. JETP 29 (1969) 483; Frankfurt, Strikman, Phys. Rept. 160 (1988) 235; Piller, Weise, Phys. Rept. 330 (2000) 1; Armesto, J. Phys. G 32 (2006) R367.

• For hard processes, QCD factorization theorem for N=2 term to relate shadowing to diffraction on nucleon at the level of parton distributions, Frankfurt, Strikman, EPJA 5 (1999) 293

• Interactions with $3 \le N \le A$ modeled using soft physics \rightarrow predictions for nuclear parton distributions (nPDFs) at small x (leading twist model of nuclear shadowing), Frankfurt, Guzey, Strikman, Phys. Rept. 512 (2012) 255.

• Alternative, complementary point of view: shadowing is mixture of leading and higher twist effects in dipole picture with saturation, Kowalski, Lappi, Venugopalan, PRL 100 (2008) 022303, or a purely HT effect, Qiu, Vitev, PRL 93 (2004) 262301.

LT shadowing in coherent J/ ψ production

• Amplitude of coherent J/ ψ production on nucleus as a series of interactions with N=1, 2, 3, ... nucleons:



• The lower part is the gluon density in the leading twist model of shadowing:

$$xf_{j/A}(x, Q_0^2) = Axf_{j/N}(x, Q_0^2) - 8\pi A(A-1) \Re e \frac{(1-i\eta)^2}{1+\eta^2} B_{\text{diff}} \int_x^{0.1} dx_{\mathbb{P}} \beta f_j^{D(3)}(\beta, Q_0^2, x_{\mathbb{P}}) \longrightarrow \frac{\text{Proton diffractive PDFs}}{\text{PDFs}}$$

$$\times \int d^2 b \int_{-\infty}^{\infty} dz_1 \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)\int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z')} \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)\int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z')} \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)\int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z')} \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)\int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z')} \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)\int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z')} \int_{z_1}^{\infty} dz_2 \rho_A(\vec{b}, z_1) \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)} \int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z') \rho_A(\vec{b}, z_2) e^{i(z_1-z_2)x_{\mathbb{P}}m_N} e^{-\frac{A}{2}(1-i\eta)\sigma_{\text{soft}}^j(x, Q_0^2)} \int_{z_1}^{z_2} dz' \rho_A(\vec{b}, z') \rho_$$

Coherent J/ ψ photoproduction in UPCs@LHC

• Predictions of LT model nicely agree with LHC data on coherent J/ψ photoproduction on heavy nuclei in Pb-Pb ultraperipheral collisions (UPCs).

• Measured cross section converted to nuclear suppression factor

 $S_{\rm Pb}(x) = \sqrt{\frac{\sigma_{\gamma A \to J/\psi A}(W_{\gamma p})}{\sigma_{\gamma A \to J/\psi A}^{\rm IA}(W_{\gamma p})}} = \kappa_{A/N} \frac{xg_A(x,\mu^2)}{Axg_N(x,\mu^2)}$

- → direct evidence of large gluon shadowing, Guzey, Kryshen, Strikman, Zhalov, PLB 726 (2013) 290, Guzey, Zhalov, JHEP 1310 (2013) 207
- LT nuclear shadowing depends on impact parameter b \rightarrow broadening of gluon distribution in b-space \rightarrow shifts t-dependence of coherent cross section,

Guzey, Strikman, Zhalov, PRC 95 (2017) 2, 025204 \rightarrow confirmed by ALICE, Acharya et al., PLB 817 (2021) 1, 136280



Coherent J/ ψ production on light nuclei

- Systematics of A dependence of nuclear shadowing and nPDFs
- Theoretically clean since modeling of interactions with N≥3 nucleons is numerically small
- Possibility to single out N = 2,3 terms by studying t-dependence. In particular, nuclear shadowing shifts t-dependence of coherent cross Section, Levin, Strikman, Sov. J. Nucl. Phys. 23 (1976) 216.
- State-of-the-art calculations of nuclear structure.

Feb 27, 2023

• No excited nuclear states \rightarrow easy to select coherent events.



Coherent J/ ψ production on light nuclei

• Amplitude as a series of interactions with k=1, 2, 3,... nucleons:



$$\frac{d\sigma_{\gamma^*A \to J/\psi A}}{dt}(t) = \frac{d\sigma_{\gamma^*N \to J/\psi N}}{dt}(t=0) \begin{vmatrix} A\Phi_1(t)e^{(B_0(x)/2)t} + \sum_{k=2}^{A} F_k(t) \end{vmatrix}^2$$
1-body form factor
$$\begin{array}{c} \text{slope of } J/\psi \\ \text{cross section,} \\ B_0=4.5 \text{ GeV}^{-2} \end{array}$$
2-body form factor

• Factors for k=2, 3,... nucleons:

$$F_k(t) = \left(-\frac{1}{8\pi^2}\right)^{k-1} \binom{A}{k} \frac{\langle \sigma^k \rangle}{\langle \sigma \rangle} \frac{(1-i\eta)^k}{1-i\eta_0} \int \prod_{l=1}^k d^2 \vec{q_l} f(q_l) \Phi_k(\vec{q},\vec{q_l}) \,\delta^2\left(\sum_{l=1}^k \vec{q_l} - \vec{q}\right)$$

Input of LT model of shadowing:

$$\frac{\langle \sigma^2 \rangle}{\langle \sigma \rangle} = \sigma_2(x), \qquad \frac{\langle \sigma^3 \rangle}{\langle \sigma \rangle} = \sigma_3(x)\sigma_2(x)$$

From diffraction@HERA

From soft models

k-body form factor as overlap of nuclear wave functions

t-dependence of soft amplitude with slope B=6 GeV⁻²

Coherent J/ ψ production in eA on He-4

• Normalized differential $\gamma^* + ^4\text{He} \rightarrow J/\psi + ^4\text{He}$ cross section at x_B=10⁻³ as function of t with $\sigma_2=25\pm15\%$ mb, $\sigma_3=30-50$ mb, and $\Phi_{2,3}$ using realistic nuclear wave function with N⁴LO chiral potential, Marcucci, Dohet-Eraly, Girlanda, Gnech, Kievsky, Viviani, Front. Phys. 8 (2020) 69; Entem, Machleidt, Nosyk, PRC 96 (2017) 2, 024004



• 2-body contribution (shadowing) shifts minimum $|t|=0.45 \text{ GeV}^2 \rightarrow |t|=0.27 \text{ GeV}^2$

• The minimum is filled due to real part of scattering amplitudes $\eta_0 \neq \eta \neq 0$.

Coherent J/ ψ production in eA on He-4

• Comparison of calculations at $x_B = 10^{-3}$ and $x_B = 0.05$:



• At $x_B=0.05$, nuclear shadowing negligibly small \rightarrow 1-body (IA) gives full result.

Coherent J/ ψ production in eA on He-3

• Repeated calculations for He-3 using realistic nuclear wf with AV18 potential, Marcucci, Dohet-Eraly, Girlanda, Gnech, Kievsky, Viviani, Front. Phys. 8 (2020) 69; Wiringa, Stoks, Schiavilla, PRC 51 (1995) 38 and 3-body forces, Pudliner, Pandharipande, Carlson, Wiringa, PRL 74 (1995) 4396

• Normalized differential γ^* + ³He $\rightarrow J/\psi$ + ³He cross section at x_B=10⁻³ vs. t



• Pattern similar to He-4.

Summary and Outlook:

• Measurements of coherent J/ψ production on light nuclei at finite t complement UPCs of heavy ions at the LHC.

- Light nuclei enable one to probe nuclear shadowing one nucleon at a time with good theoretical control of nuclear structure.
- Dominant effect by model-independent 2-nucleon contribution \rightarrow test of nuclear shadowing models.
- LT model of nuclear shadowing predicts shift of the t-dependence of coherent J/ψ production to smaller $|t| \rightarrow$ broadening of the gluon distribution in impact parameter space.
- Formalism can be readily applied to estimate coherent DVCS on light nuclei.