

### (Di-)jets & Di-hadrons from HI studies to ep and eA [some liberties taken on definitions]

1<sup>st</sup> International Workshop on a 2nd Detector for the Electron-Ion Collider

Helen Caines, Wright Lab, Yale



### **Di-jets and di-hadrons in Yellow Report**

Too many topics to cover in 20 mins includes:

Access to the gluon Wigner distribution, Probing the linearly polarized Weizsaecker-Williams gluon TMDs Probing the gluon Sivers function Exploring the (un)polarized hadronic structure of the photon fraction (x)

Studies of hadronization and cold nuclear matter properties

- Constraining (un)polarized quark and gluon PDFs at moderate to high momentum
  - In this talk I will focus on topics on which I have expertise/ knowledge as a HI physicist

plus studies with recent results from RHIC and LHC





### **Kinematics and Properties of Jets at EIC**



### **Kinematics and Properties of Jets at EIC**



### **Kinematics and Properties of Di-jets at EIC**





## **Underlying Event at RHIC**



At mid rapidity UE for jet  $p_T > 15$  GeV/c Average charged particle density = 0.4-0.6 Mean  $p_T = 0.5 - 0.7 \text{ GeV/c}$ Largely independent of jet  $p_T$ Significant dependence on low p<sub>T</sub> cut-off







# **Underlying Event at EIC**



- Potential solution to use off-axis technique

### UE for jet $p_T > 5 \text{ GeV/c}$

Average charged particle density = 0.1 - 0.2

Mean  $p_T = 1-1.2 \text{GeV/c}$ (p<sub>T</sub> of UE increased because using Breit frame)

For jet with R= 0.4: Charged particle contamination from UE ~80 MeV

B. Page et al. PRD 101, 072003 (2020)





### **Detector Requirements for (Di-)jets at EIC**

Hermetic detector

Tracking:

- efficient with excellent momentum and angular resolution
- large range of rapidity [jets only back-to-back in azimuth]
- single hadrons  $p_T > 100 \text{ MeV/c}$
- $3\sigma \pi/K/p$  B: up to 7 GeV/c, C: up to 10 GeV/c, F: up to 50 GeV/c, [species dependent FF]

### Electromagnetic calorimetry:

- Resolution: B:  $\sigma(E)/E \approx 2\%/\sqrt{E} \oplus 1 3\%$  [drives jet performance capabilities], C:  $\sigma(E)/E \approx 10 12\%/\sqrt{E} \oplus 1$
- 3% [sufficient, take advantage of excellent tracking], F:  $\sigma(E)/E \approx 4.5/\sqrt{E}$

### Hadron calorimetry:

- Neutral hadron isolation [important for jet energy scale and resolution, especially F and B] - Resolution: B:  $\sigma(E)/E \approx 50\%/\sqrt{E} \oplus 6\%$ , C:  $\sigma(E)/E \approx 85\%/\sqrt{E} \oplus 7\%$ , F:  $\sigma(E)/E \approx 35\%/\sqrt{E}$
- Minimum energy: 500 MeV/c

Helen Caines - Yale

### (similar for di-ha

Yellow Report

### **Gluon Saturation**



 $Q_s \propto A^{1/3}$  : More suppression for heavier nuclei Forward for low-x

Helen Caines - Yale

Rapid increase of gluon density as lower x - saturation when splitting rate matches rate of recombination

- saturation scale at Q<sub>s</sub>

D.Kharzeev et al. NPA 748, 627 (2005)





### **Gluon Saturation**



Helen Caines - Yale

### **Di-pi Correlations at RHIC**

Suppression in pA: - Dependence on A - Dependence on EA (related to but not b) - Only at low pt - No broadening



STAR: PRL 129, 092501 (2022) PHENIX: PRL 107, 172301 (2011)

### **Di-pi Correlations at RHIC**



Suppression in dAu:

- Only at low p<sub>T</sub>
- Only in central events

Hints of reaching saturation at RHIC 2024: unique opportunity with STAR forward upgrades

Helen Caines - Yale

### Suppression in pA: - Dependence on A - Dependence on EA (related to but not b) - Only at low p<sub>T</sub> - No broadening



### **Opportunities at the EIC**



- Similar Moderate-Q<sup>2</sup>-low x - Similar collision energy - Complimentary probes (e vs p)

- High Q<sup>2</sup> - low x - Complimentary probes (e vs p)

- Study wide range of ion beams from deuterons to heavy nuclei (Au, Pb, U)

Definitive measurements at EIC?











## Energy Loss to p(d)-Au Medium at RHIC

![](_page_16_Figure_1.jpeg)

I<sub>pAu</sub> for recoil jet Normalization per high p<sub>T</sub> trigger object

Suppression on near and away-side similar; inconsistent with Eloss

![](_page_16_Picture_6.jpeg)

![](_page_16_Picture_7.jpeg)

# Energy Loss to p(d)-Au Medium at RHIC

![](_page_17_Figure_1.jpeg)

I<sub>pAu</sub> for recoil jet Normalization per high p<sub>T</sub> trigger object

> Suppression on near and away-side similar; inconsistent with Eloss

 $R_{dAu}$  for  $\pi^0$ N<sub>bin</sub> determined by forcing R<sub>dAu</sub> γ to unity (colorless object no interaction with medium) Qualitatively consistent with predictions of Eloss in small systems

A. Huss et al. PRC103, 054903 (2021). W. Ke, I. Vitev arXiv:2204.00634

Helen Caines - Yale

1<sup>st</sup> International Workshop on a 2nd Detector for the EIC - May 17-20

![](_page_17_Figure_8.jpeg)

## Energy Loss to p(d)-Au Medium at RHIC

![](_page_18_Figure_1.jpeg)

What will we see at EIC?

 $R_{dAu}$  for  $\pi^0$ N<sub>bin</sub> determined by forcing R<sub>dAu</sub> γ to unity (colorless object no interaction with medium)

### Qualitatively consistent with predictions of Eloss in small systems

A. Huss et al. PRC103, 054903 (2021). W. Ke, I. Vitev arXiv:2204.00634

Helen Caines - Yale

1<sup>st</sup> International Workshop on a 2nd Detector for the EIC - May 17-20

I<sub>pAu</sub> for recoil jet Normalization per high p<sub>T</sub> trigger object

inconsistent with Eloss

![](_page_18_Figure_10.jpeg)

### Jet Azimuthal Broadening at the EIC

![](_page_19_Figure_1.jpeg)

### Lepton-jet correlations in DIS in eA

 $\hat{q}$ L=0GeV<sup>2</sup>

 $\hat{q}$ L=0.2GeV<sup>2</sup>

 $\hat{q}$ L=0.8GeV<sup>2</sup>

Multiple interactions of hard scattered parton as it exits the target nucleus generates p<sub>T</sub>broadening/deflection

- similar concept used to explain Cronin effect in pA collisions

\_Δφ 0.5

### Sensitivity to Eloss in Cold nuclear matter

X. Liu et al. PRL 122, 192003 (2019)

![](_page_19_Picture_13.jpeg)

![](_page_19_Picture_14.jpeg)

![](_page_19_Picture_15.jpeg)

# **Energy-Energy Correlators**

The ultimate di-hadron correlation

![](_page_20_Picture_2.jpeg)

![](_page_20_Figure_6.jpeg)

![](_page_21_Figure_0.jpeg)

### Good theory-date agreements

 $R_L^*p_T \sim constant * \Lambda_{QCD}$ 

Universal turning point across large range in jet pt

Helen Caines - Yale

1<sup>st</sup> International Workshop on a 2nd Detector for the EIC - May 17-20

measurements of the shapes and scalings of  $\langle \Psi | \mathcal{E}(\vec{n}_1) \mathcal{E}(\vec{n}_2) \cdots \mathcal{E}(\vec{n}_k) | \Psi \rangle$ , providing new ins modynamics (QCD). In this *Letter*, we use rigorous factorization theorem for the light-ra <del>nomentum jets at the LHC. Us</del>ing the light-ra multi-point correlators can be computed from

### New preliminary

resolve the quantum scaling dimensions of Q Anti- $k_{T}$  ch-particle jets, factorization theorem for the light-ray density developments in the study of energy correlat 40 < A substructure stud

> Introduction.—The Large Hadron Collider (LHC) pro vides an opportunity **COD** (**N** Quantum field theory is general, and quantum chromodynamics (QCD) in partic ular, at unprecedent AR Rergy scales, and with moder resolution detectors [1, 2]. Due to the phenomenon of asymptotic freedom [3-6], this gain in energy is particular larly advantageous, as it enables QCD to be studied i the perturbative regime, where first principles calcula tions are currently possible.

One of the major achievements, which re-invigorate the study of QCD at the LHC, was the introduction of experimentally robust infrared safe jet algorithms, mos notably the anti- $k_T$  algorithm [7–10], that allow for the identification of high transverse momentum,  $p_T$ , jets i hadronic collisions. The inclusive production of such jet an bertuel is of or bus factorization theorems [11 17], whose perturbative components have been computeallowing for

 $H_{I}$ 

 $10^{-2}$ 

ALI-PKEL-338340

stions abou QCD, namely understanding the Lorentzian dynamics of and gluons, and the nature of their real-time con nt in STARions, at a misence de 2023 he distribu jetsAbyOEthR in Churzst Torresof Pere 2023 within jets, known as jet substructure [25, 26]. Jet sub structure has been extraordinarily successful as a 5 ne

 $[\rho \mathbb{O}^{[J]}]$ . We co sults for the ne aring with CM

## **Energy-Energy Correlators at the EIC**

![](_page_22_Figure_1.jpeg)

 $\theta^2_L p_T L \sim 1$  $\theta L \propto 1/\sqrt{p_T} L \propto 1/(\sqrt{pT} A^{1/6})$ 

K.Devereaux et al. arXiv:2303.08143v1

![](_page_22_Picture_7.jpeg)

## **Energy-Energy Correlators at the EIC**

![](_page_23_Figure_1.jpeg)

 $\theta^2_L p_T L \sim 1$  $\theta L \propto 1/\sqrt{p_T} L \propto 1/(\sqrt{pT} A^{1/6})$ 

# radiation by changing power E<sup>n</sup>

K.Devereaux et al. arXiv:2303.08143v1

![](_page_23_Picture_7.jpeg)

## **Energy-Energy Correlators at the EIC**

![](_page_24_Figure_1.jpeg)

![](_page_24_Picture_3.jpeg)

EECs versatile tools with many other applications

K.Devereaux et al. arXiv:2303.08143v1

![](_page_24_Picture_10.jpeg)

### **Charge Correlation Ratio**

Access nonperturbative hadronization process with this di-hadron correlation - Count leading and subleading particle in jet with SS and OS

$$\boldsymbol{r}_{c} \equiv \frac{N_{CC} - N_{C\overline{C}}}{N_{CC} + N_{C\overline{C}}}$$

Random/No correlation:  $r_c = 0$ 

![](_page_25_Figure_4.jpeg)

### Can also look at species correlations

Helen Caines - Yale

![](_page_25_Figure_9.jpeg)

Y.T.Chien PRD 105, L051502 (2022)

![](_page_25_Picture_11.jpeg)

### **Charge Correlation Ratio**

First preliminary results from H1

$$\boldsymbol{r_c} \equiv \frac{N_{CC} - N_{C\overline{C}}}{N_{CC} + N_{C\overline{C}}}$$

### Take one step further and look at rc for first and second split of jet using SoftDrop

Models reproduce trends of the data

### Need EIC to look at PID correlations

![](_page_26_Figure_9.jpeg)

## **Exploiting Multifold**

Correction tool just beginning to be used

Machine learning driven

- Unbinned
- Simultaneously unfolds multiple observables  $\rightarrow$  Correlation information is retained

1/N<sub>jet</sub> dN<sub>jet</sub>/dM Use example collinear drop: 0.2 Probes soft component 0.0 STAR initial measured  $\Delta M$ 0 Theory prefer M<sup>2</sup>-M<sup>2</sup><sub>q</sub>)/p<sup>2</sup><sub>T</sub> Because using multifold could calculate without redoing analysis

Could be powerful tool at EIC

0.8

0.6

0.4

<sup>2</sup>/GeV]

![](_page_27_Figure_9.jpeg)

![](_page_27_Picture_11.jpeg)

Andreassen et al. PRL 124, 182001 (2020) Chien and Stewart JHEP 2020, 64 (2020). STAR: Y. Song DIS 2023

![](_page_27_Picture_14.jpeg)

### Summary

Many studies build on those occurring at RHIC, LHC and HERA

Exploit tools being developed now and over next few years

Wealth of di-jet and di-hadron correlations studies to be done at EIC only scratched the surface in this talk

Success of the studies rely on meeting or exceeding detector requirements - As ePIC design becomes set in stone will probably need D2 to do this

1<sup>st</sup> International Workshop on a 2nd Detector for the EIC - May 17-20

![](_page_28_Picture_9.jpeg)

### **Kinematics and propert**

![](_page_29_Figure_1.jpeg)

FIG. 11. [color online] The average number of particles in a jet as a function of the transverse momentum of the jet for all stable particles and only charged particles for minimum particle  $p_T s$  of 250 and 500 MeV/c. Also shown are the RMS variations for all particles with  $p_{\rm T} > 250 \text{ MeV/c}$  and charged particles with  $p_{\rm T} > 500 \text{ MeV/c.}$ 

![](_page_29_Figure_5.jpeg)

B. Page et al. PRD 101, 072003 (2020)

![](_page_29_Figure_7.jpeg)

![](_page_29_Picture_8.jpeg)

![](_page_29_Picture_9.jpeg)

### **No broadening in pA at STAR** Results: no broadening of acoplanarity with **EA**<sub>BBC</sub>

![](_page_30_Figure_1.jpeg)

Helen Caines - Yale

![](_page_30_Picture_4.jpeg)

### **The Breit Frame**

The Breit or "brick wall" frame is particularly useful for jet analyses It is oriented such that for the DIS process  $\gamma * q \rightarrow q'$ , the virtual photon and interacting quark collide head-on along the z-axis It is then boosted such that t the virtual photon four-momentum (0,0,0,-Q)This boost means Incoming quark  $p_z = Q/2$ Scattered quark  $p_z = -Q/2$ Proton remnant  $p_z = (1-x)Q/(2x)$ . Gives a natural separation between jets from struck quark and those associated with proton remnant. Working in the Breit frame has the effect of suppressing contributions from the L.O. subprocess, as the scattered quark has zero transverse momentum by construction

![](_page_31_Figure_6.jpeg)

![](_page_31_Picture_7.jpeg)

![](_page_32_Figure_0.jpeg)

![](_page_32_Picture_4.jpeg)

![](_page_32_Picture_6.jpeg)

![](_page_32_Picture_7.jpeg)