



Covariant Extension of Generalized Parton Distributions using Artificial Neural Networks

Pietro Dall'Olio

Collaborators: J. M. Morgado, J. Rodríguez Quintero, C. Mezrag, P. Sznajder

Probing the Frontiers of Nuclear Physics with AI at the EIC

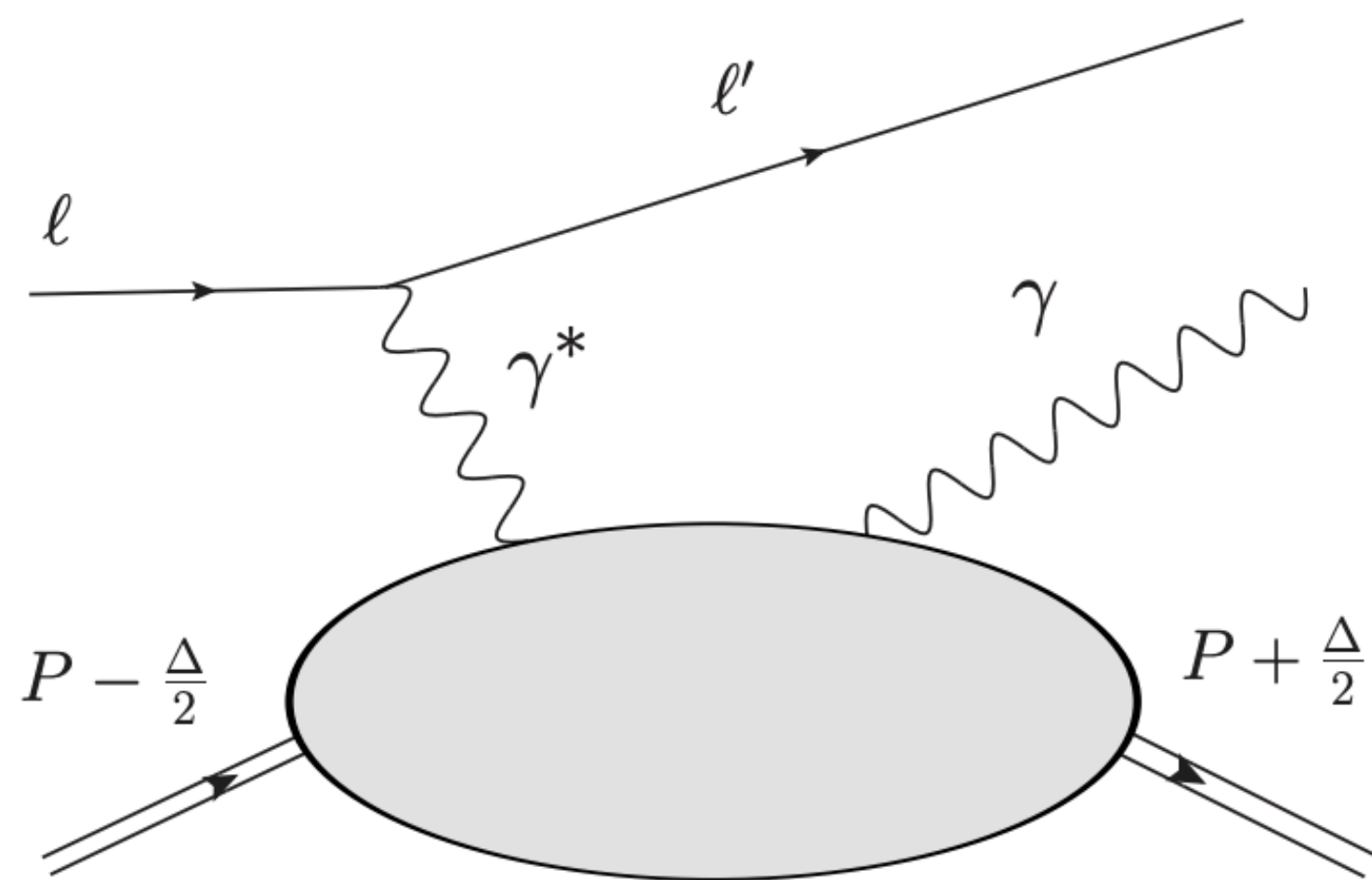
26 Sept 2023

Outline

- Generalized Parton Distributions
- Covariant Extension via Radon Transform (RT) inversion
- RT inversion using Artificial Neural Networks
- Results on analytical models
- Conclusions and outlook

Generalized Parton Distribution (GPD)

Exclusive processes (DVCS)



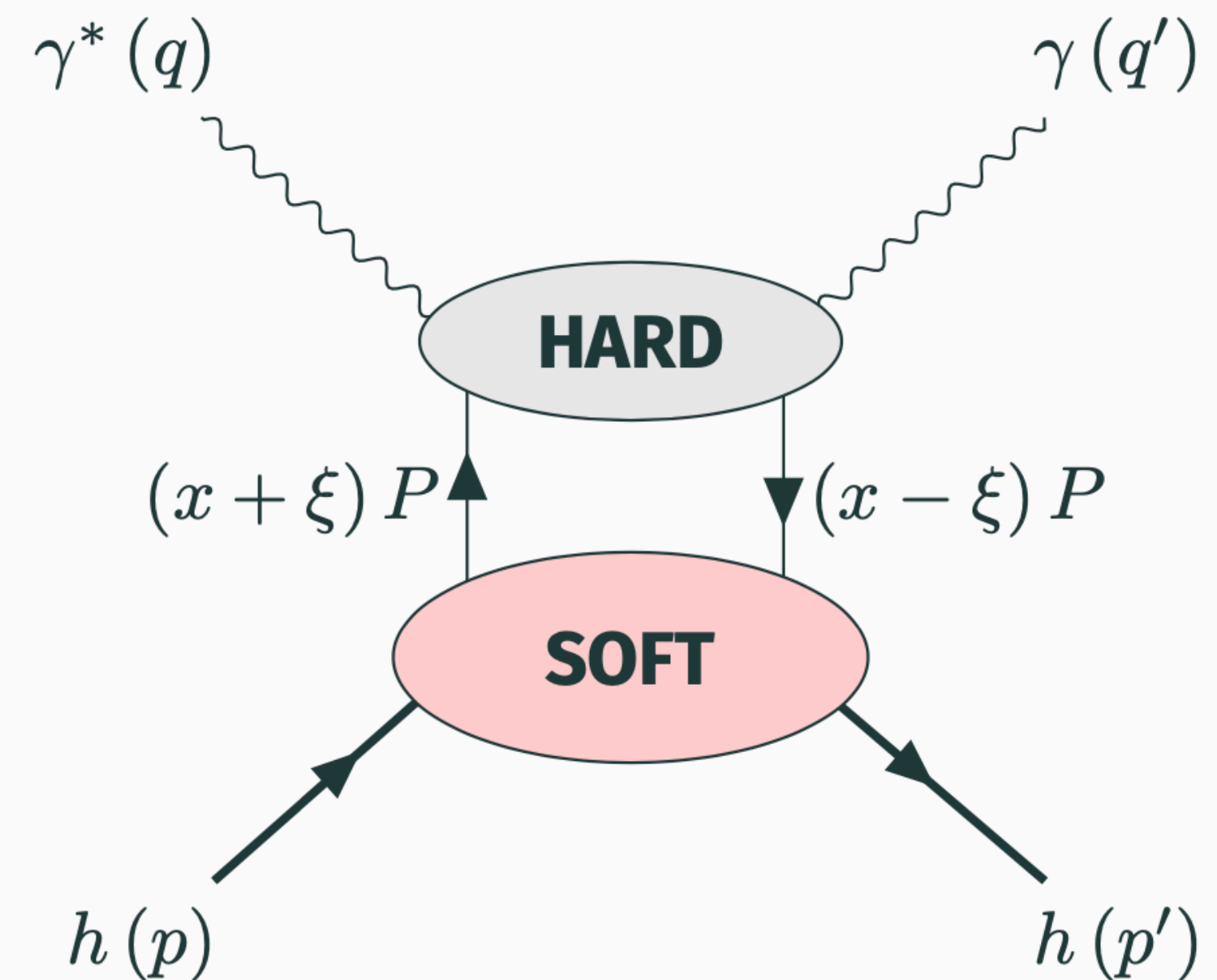
$$Q^2 \rightarrow \infty$$



$$P = \frac{p + p'}{2}, \quad \Delta = p' - p$$

$$\xi = \frac{(p' - p)^+}{2P^+} = -\frac{\Delta^+}{2P^+}$$

$$t = \Delta^2$$



Factorization of the amplitude

$$\mathcal{M}(\xi, t; Q^2) = \sum_{p=q,g} \int_{-1}^1 \frac{dx}{\xi} K^p \left(\frac{x}{\xi}, \frac{Q^2}{\mu_F^2}, \alpha_s(\mu_F) \right) \mathcal{H}^p(x, \xi, t; \mu_F) + O(1/Q^2)$$

hard / perturbative

soft / non perturbative **GPD**

Ex: Chiral even twist-2 operators

Nucleon quark GPD

$$\begin{aligned} \mathcal{H}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\mathbf{H}^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + \mathbf{E}^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right] \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{H}}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \gamma^5 \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{\mathbf{H}}^q(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{\mathbf{E}}^q(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right] \end{aligned}$$

Nucleon gluon GPD

$$\begin{aligned} \mathcal{H}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) G_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\mathbf{H}^g(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + \mathbf{E}^g(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right] \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{H}}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) \tilde{G}_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{\mathbf{H}}^g(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{\mathbf{E}}^g(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right] \end{aligned}$$

Ex: Chiral even twist-2 operators

Nucleon quark GPD

$$\begin{aligned}\mathcal{H}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[H^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right]\end{aligned}$$

$$\begin{aligned}\tilde{\mathcal{H}}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \gamma^5 \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{H}^q(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{E}^q(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right]\end{aligned}$$

Nucleon gluon GPD

$$\begin{aligned}\mathcal{H}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) G_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[H^g(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^g(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right]\end{aligned}$$

$$\begin{aligned}\tilde{\mathcal{H}}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) \tilde{G}_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{H}^g(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{E}^g(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right]\end{aligned}$$

quark GPD of spin zero hadron

$$H_\pi^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle \pi(P + \Delta/2) | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | \pi(P - \Delta/2) \rangle \Big|_{z^+=0, z^\perp=0}$$

Ex: Chiral even twist-2 operators

Nucleon quark GPD

$$\begin{aligned} \mathcal{H}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[H^q(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^q(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right] \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{H}}^q(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}^q(-z/2) \gamma^+ \gamma^5 \psi^q(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{H}^q(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{E}^q(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right] \end{aligned}$$

Nucleon gluon GPD

$$\begin{aligned} \mathcal{H}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) G_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[H^g(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E^g(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\mu} \Delta_\mu}{2M} u(p) \right] \end{aligned}$$

$$\begin{aligned} \tilde{\mathcal{H}}^g(x, \xi, t) &= \frac{1}{2} \int dz^- e^{ixP^+z^-} \langle P + \Delta/2 | G^{\mu+}(-z/2) \tilde{G}_\mu^+(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \frac{1}{2P^+} \left[\tilde{H}^g(x, \xi, t) \bar{u}(p') \gamma^+ \gamma^5 u(p) + \tilde{E}^g(x, \xi, t) \bar{u}(p') \frac{\gamma^5 \Delta^+}{2M} u(p) \right] \end{aligned}$$

quark GPD of spin zero hadron

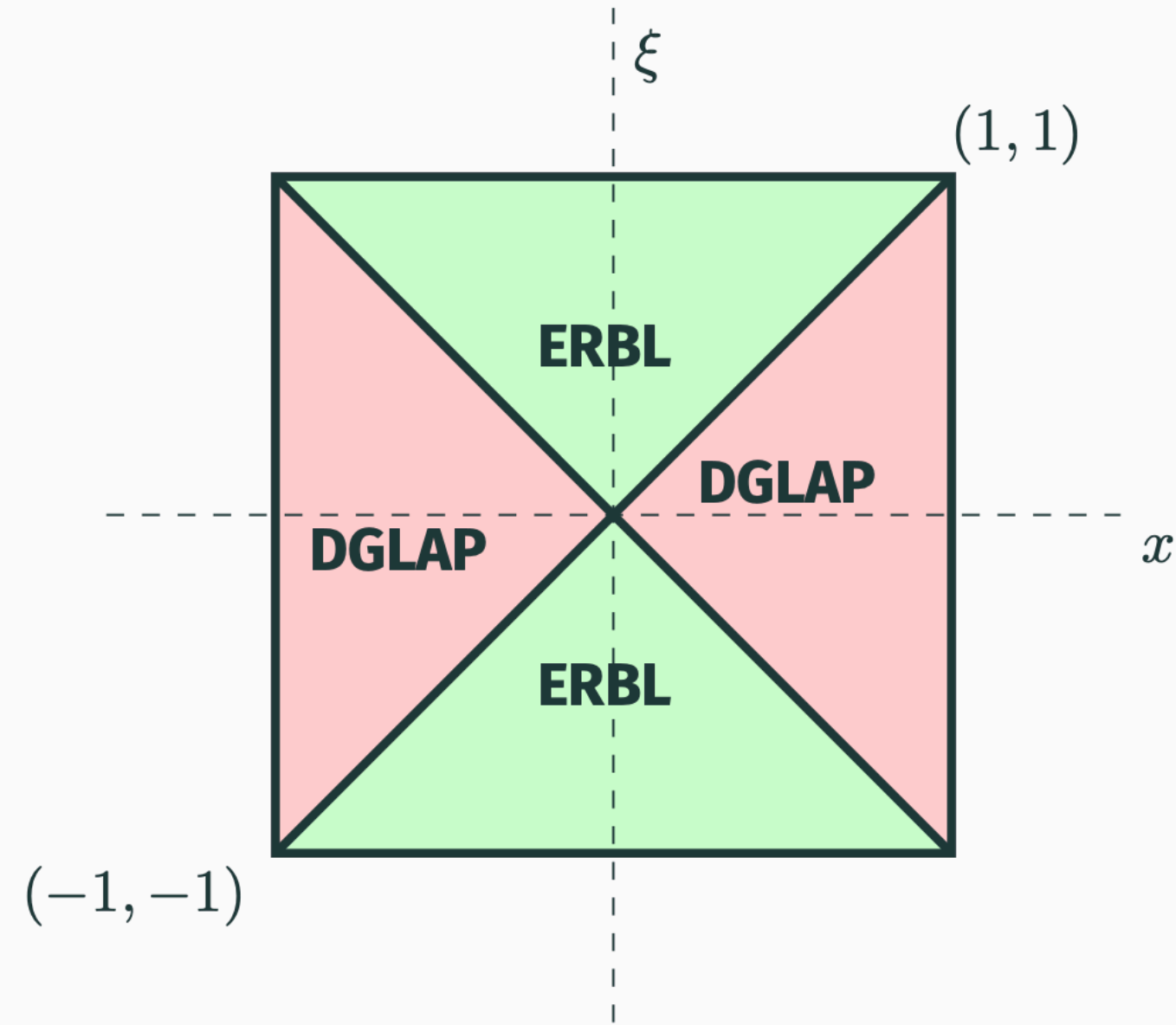
$$H_\pi^q(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle \pi(P + \Delta/2) | \bar{\psi}^q(-z/2) \gamma^+ \psi^q(z/2) | \pi(P - \Delta/2) \rangle \Big|_{z^+=0, z^\perp=0}$$



No Wilson line in light-cone gauge $A^+ = 0$

GPD properties

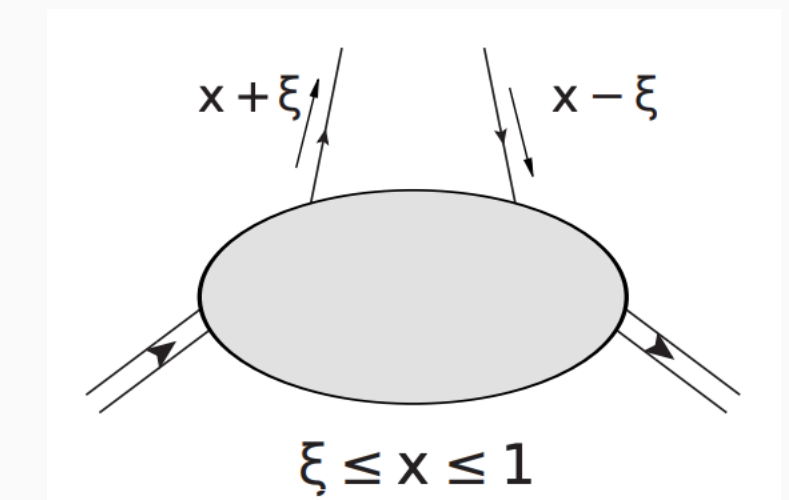
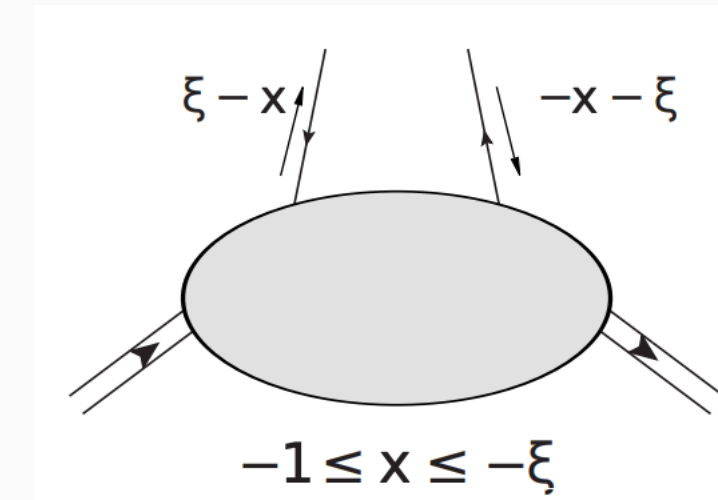
- **Support** $(x, \xi) \in [-1, 1] \otimes [-1, 1]$



DGLAP

$$|x| > |\xi|$$

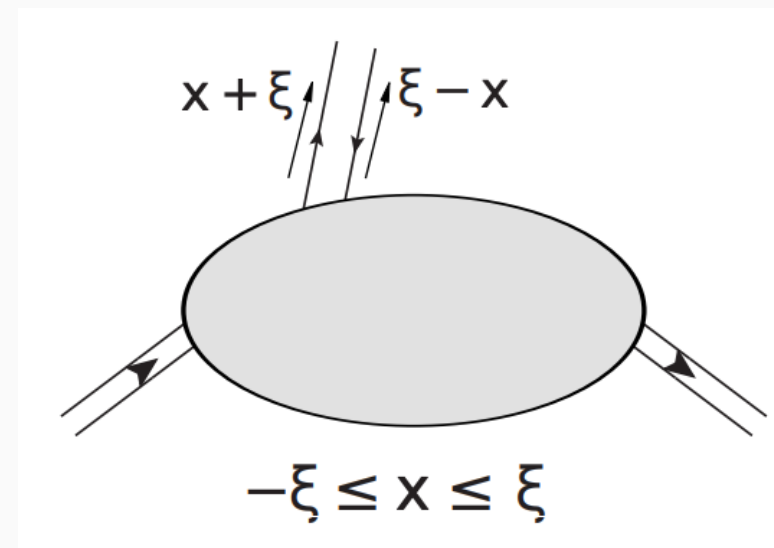
Emission/absorption quark ($x > 0$) or antiquark ($x < 0$)



ERBL

$$|x| < |\xi|$$

Emission of quark/antiquark pair



- ξ -parity

$$H^q(x, -\xi, t) = H^q(x, \xi, t) \quad \text{Time inversion symmetry}$$

- Polynomiality

$$\mathcal{A}_m(\xi, t) = \int_{-1}^1 dx x^m H^q(x, \xi, t) = \sum_{\substack{k=0 \\ k \text{ even}}}^{m+1} C_{k,m}(t) \xi^k \quad \text{Lorentz symmetry}$$

- Positivity

$$\left| H^q(x, \xi, t=0) \right| \leq \sqrt{q \left(\frac{x+\xi}{1+\xi} \right) q \left(\frac{x-\xi}{1-\xi} \right)}, \quad |x| > |\xi| \quad \text{Hilbert space norm}$$

- Form factors

$$\int_{-1}^1 dx H^q(x, \xi, t) = F^q(t), \quad \int_{-1}^1 dx x H^q(x, \xi, t) = A^q(t) + \xi^2 C^q(t)$$

- **Forward limit**

$$H^q(x,0,0) = q(x)\theta(x) - \bar{q}(-x)\theta(-x)$$

- **Hadron 3D tomography**

$$\rho^q(x, b_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-ib_\perp \cdot \Delta_\perp} H^q(x,0, -\Delta_\perp^2)$$

Probability density of finding a parton with longitudinal momentum fraction x and position b_\perp in the transverse plane.

- **Forward limit**

$$H^q(x,0,0) = q(x)\theta(x) - \bar{q}(-x)\theta(-x)$$

- **Hadron 3D tomography**

$$\rho^q(x, b_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-ib_\perp \cdot \Delta_\perp} H^q(x,0, -\Delta_\perp^2)$$

Probability density of finding a parton with longitudinal momentum fraction x and position b_\perp in the transverse plane.

It is hard to build GPDs from first principles that satisfy both **positivity** and **polynomiality**

- **Forward limit**

$$H^q(x,0,0) = q(x)\theta(x) - \bar{q}(-x)\theta(-x)$$

- **Hadron 3D tomography**

$$\rho^q(x, b_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} e^{-ib_\perp \cdot \Delta_\perp} H^q(x,0, -\Delta_\perp^2)$$

Probability density of finding a parton with longitudinal momentum fraction x and position b_\perp in the transverse plane.

It is hard to build GPDs from first principles that satisfy both **positivity** and **polynomiality**

Truncation in **Overlap** representation of **LCWF** is consistent only in **DGLAP** region

$$H^q(x, \xi, t) \Big|_{|x| > |\xi|} = \sum_{N,\beta} \sqrt{1 - \xi^2}^{2-N} \int [d\bar{x}]_N [d^2\bar{\mathbf{k}}_\perp]_N \delta(x - \bar{x}_i) \psi_{N,\beta}^*(x_i^{out}, \mathbf{k}_{i\perp}^{out}) \psi_{N,\beta}(x_i^{in}, \mathbf{k}_{i\perp}^{in})$$

Positivity is built in, but for **polynomiality** GPD must be extended to **ERBL** region

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from **DGLAP** to **ERBL** preserving polynomiality?

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from **DGLAP** to **ERBL** preserving polynomiality?

$$\langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma_\mu \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^2=0} = \int d\beta d\alpha e^{-i\beta(P \cdot z) + i\alpha \frac{\Delta \cdot z}{2}} \left[2P_\mu f(\beta, \alpha, t) - \Delta_\mu g(\beta, \alpha, t) \right]$$

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from **DGLAP** to **ERBL** preserving polynomiality?

$$\langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma_\mu \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^2=0} = \int d\beta d\alpha e^{-i\beta(P \cdot z) + i\alpha \frac{\Delta \cdot z}{2}} \left[2P_\mu f(\beta, \alpha, t) - \Delta_\mu g(\beta, \alpha, t) \right]$$

$$\begin{aligned} \Rightarrow H(x, \xi, t) &= \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma^+ \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0} \\ &= \int_\Omega d\beta d\alpha (f(\beta, \alpha, t) + \xi g(\beta, \alpha, t)) \delta(x - \beta - \xi\alpha) \end{aligned}$$

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from **DGLAP** to **ERBL** preserving polynomiality?

$$\langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma_{\mu} \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^2=0} = \int d\beta d\alpha e^{-i\beta(P \cdot z) + i\alpha \frac{\Delta \cdot z}{2}} \left[2P_{\mu} f(\beta, \alpha, t) - \Delta_{\mu} g(\beta, \alpha, t) \right]$$

$$\begin{aligned} \Rightarrow H(x, \xi, t) &= \frac{1}{2} \int \frac{dz^{-}}{2\pi} e^{ixP^{+}z^{-}} \langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma^{+} \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^{+}=z_{\perp}=0} \\ &= \int_{\Omega} d\beta d\alpha (f(\beta, \alpha, t) + \xi g(\beta, \alpha, t)) \delta(x - \beta - \xi\alpha) \end{aligned}$$

Radon transform

- Integral over line parametrized by x, ξ
- It guarantees polynomiality

Covariant extension via Radon transform inversion

Is it possible to extend the GPD from **DGLAP** to **ERBL** preserving polynomiality?

$$\langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma_\mu \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^2=0} = \int d\beta d\alpha e^{-i\beta(P \cdot z) + i\alpha \frac{\Delta \cdot z}{2}} \left[2P_\mu f(\beta, \alpha, t) - \Delta_\mu g(\beta, \alpha, t) \right]$$

$$\Rightarrow H(x, \xi, t) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle P + \Delta/2 | \bar{\psi}(-z/2) \gamma^+ \psi(z/2) | P - \Delta/2 \rangle \Big|_{z^+=z_\perp=0}$$

$$= \int_{\Omega} d\beta d\alpha (f(\beta, \alpha, t) + \xi g(\beta, \alpha, t)) \delta(x - \beta - \xi\alpha)$$

Double Distributions (DD)

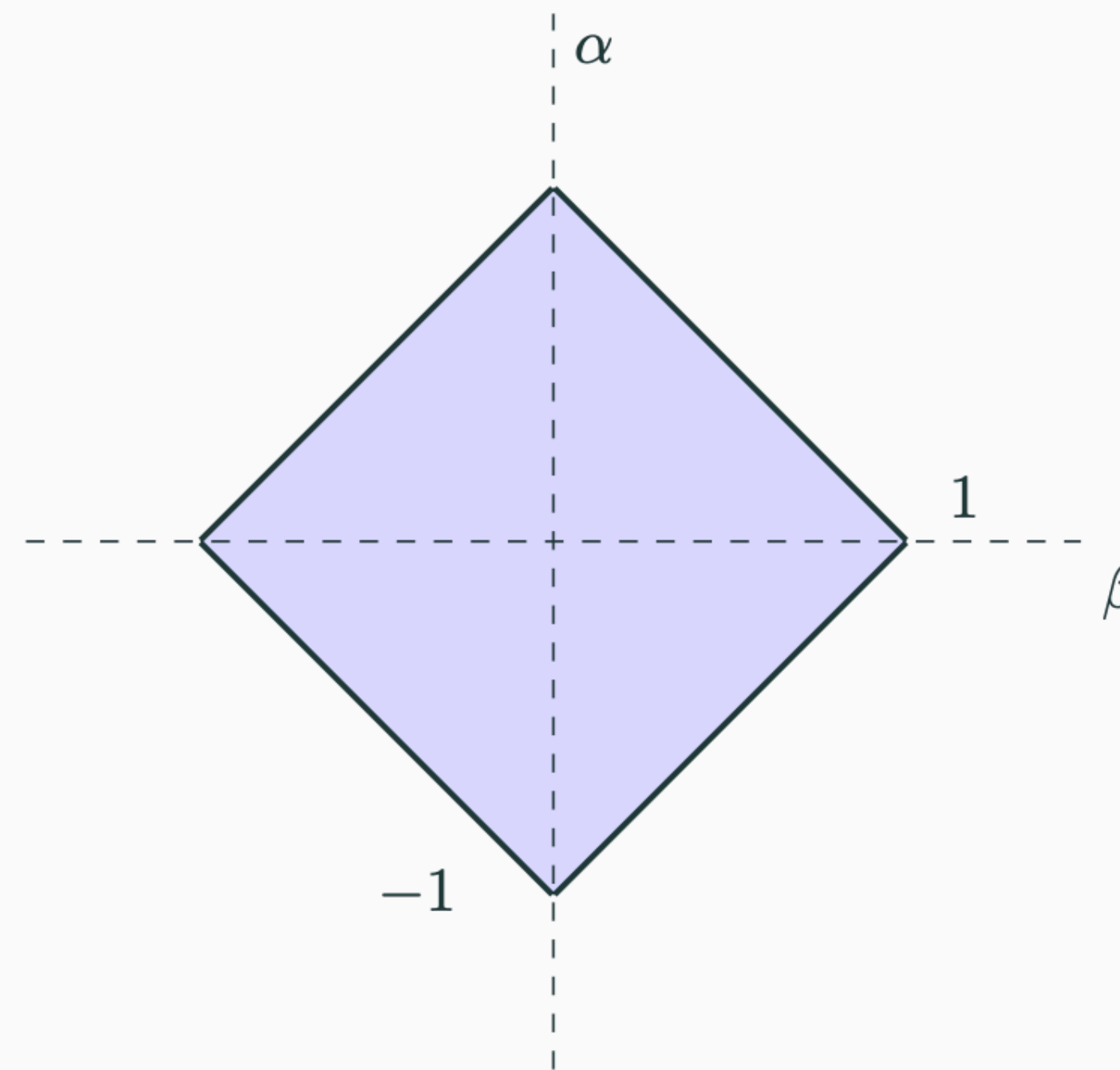
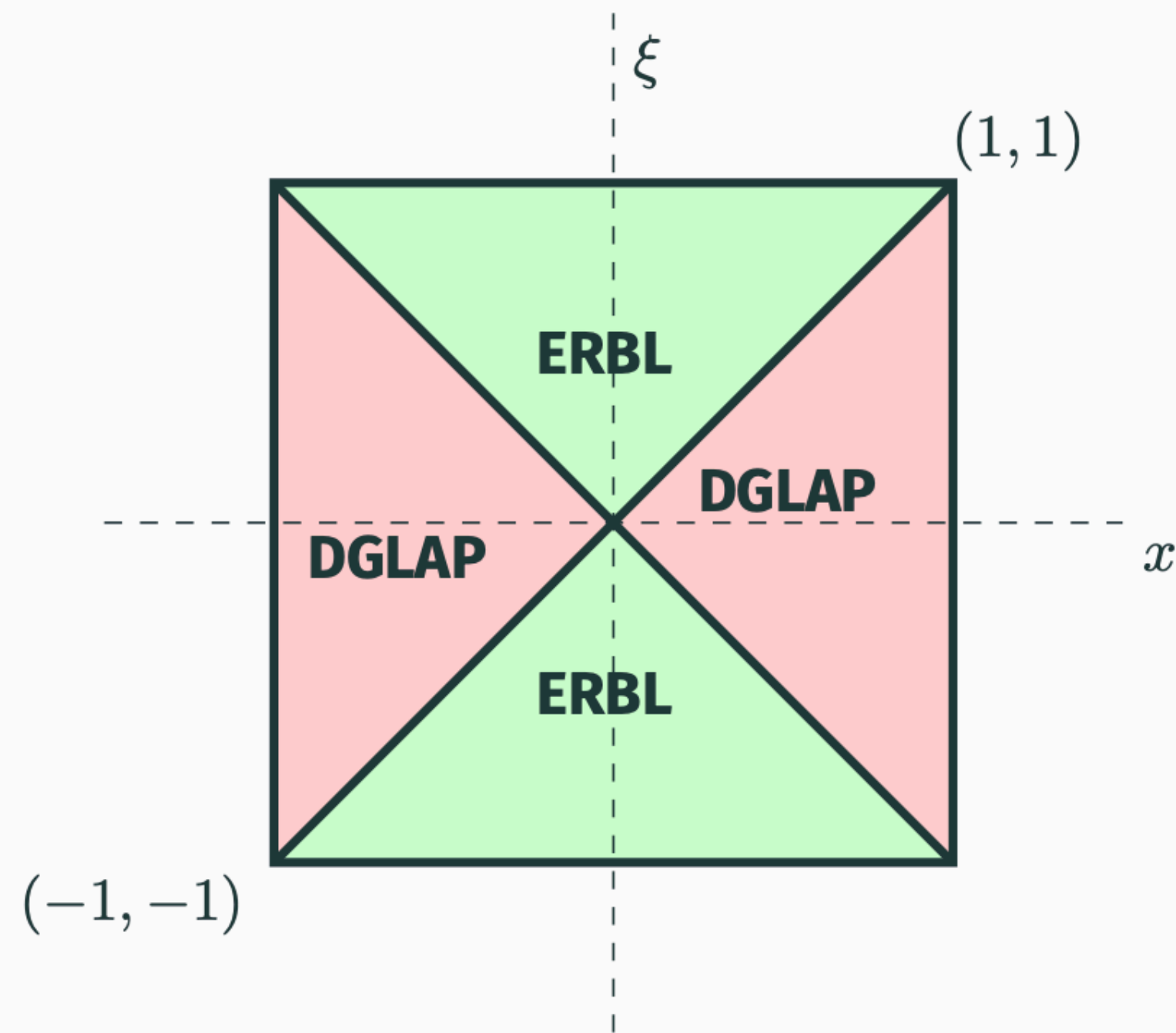
Not uniquely fixed

Radon transform

- Integral over line parametrized by x, ξ
- It guarantees polynomiality

$$H^q(x, \xi) = \mathcal{R}[h] = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha) + \theta(|\xi| - |x|) \mathbf{D}\text{-terms}$$

Support: $\Omega = \{(\beta, \alpha) \mid |\beta| + |\alpha| \leq 1\} = \Omega^+ \theta(\beta) + \Omega^- \theta(-\beta)$

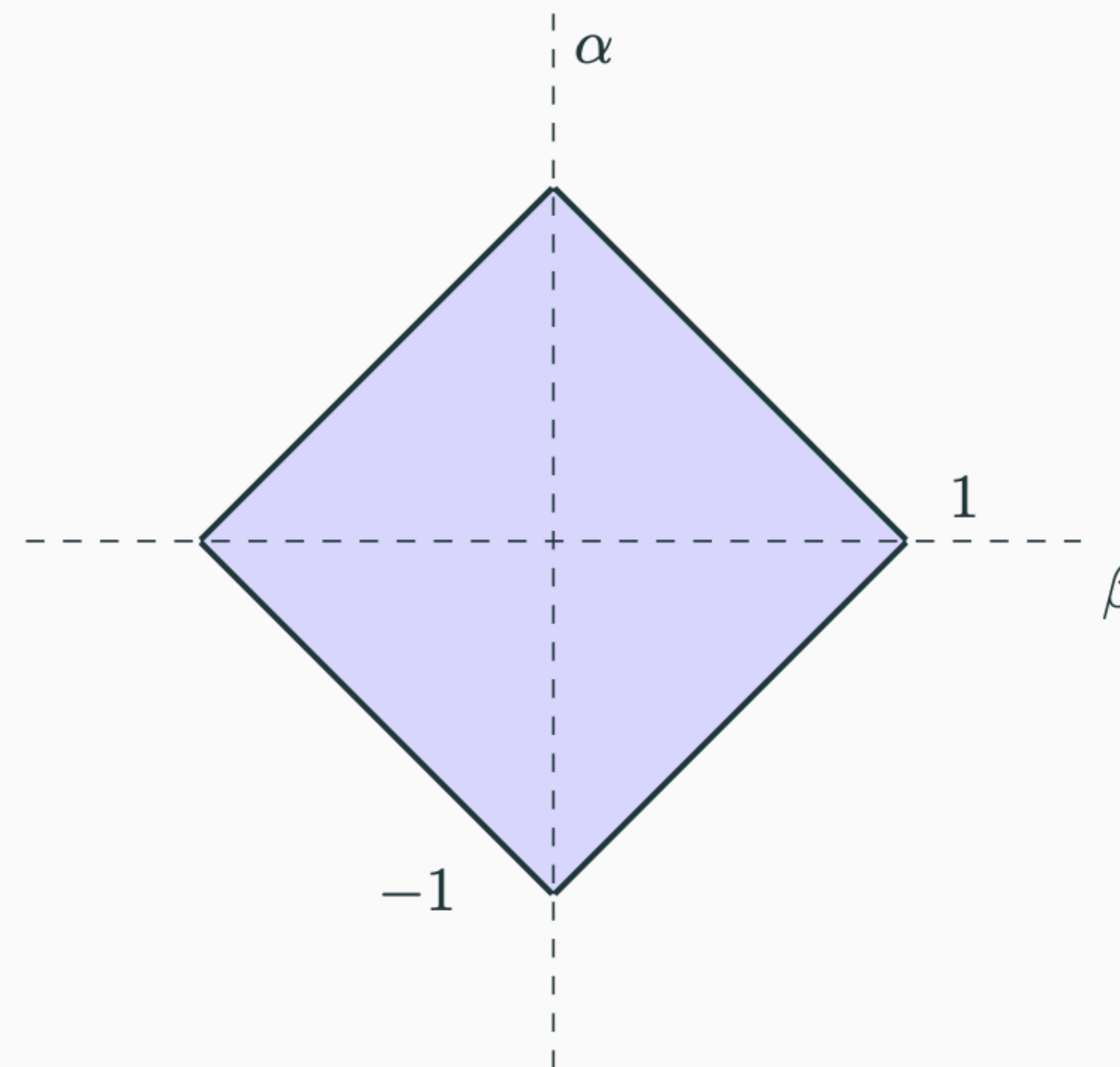
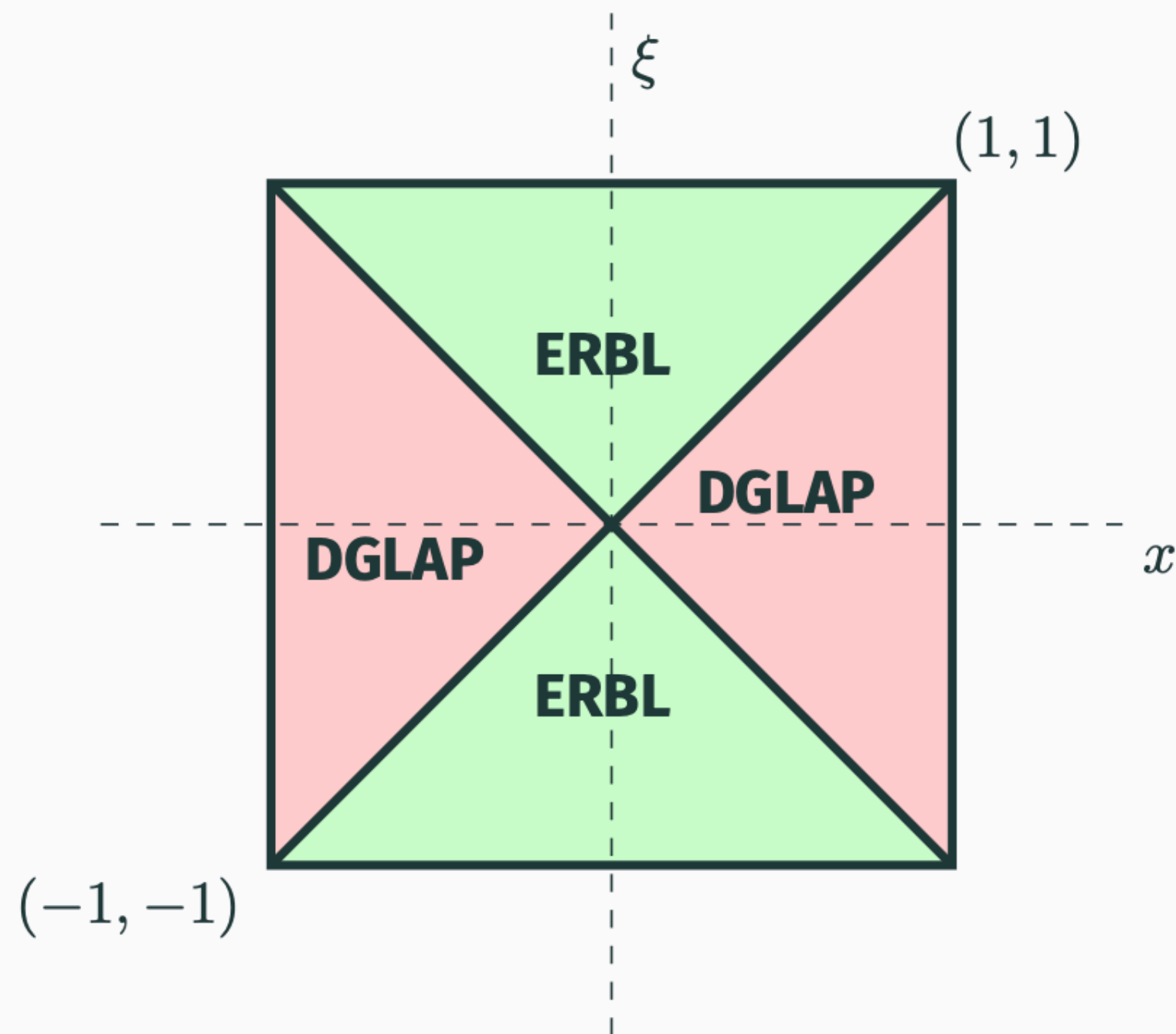


line $\alpha = -\frac{\beta}{\xi} + \frac{x}{\xi}$

$$H^q(x, \xi) = \mathcal{R}[h] = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha) + \theta(|\xi| - |x|) \text{ D-terms}$$

- Cannot be fixed by DGLAP data.
- Its DD has support only in $\beta = 0$
- Contribute only to $C_{m,m+1}(t)$

Support: $\Omega = \{(\beta, \alpha) \mid |\beta| + |\alpha| \leq 1\} = \Omega^+ \theta(\beta) + \Omega^- \theta(-\beta)$

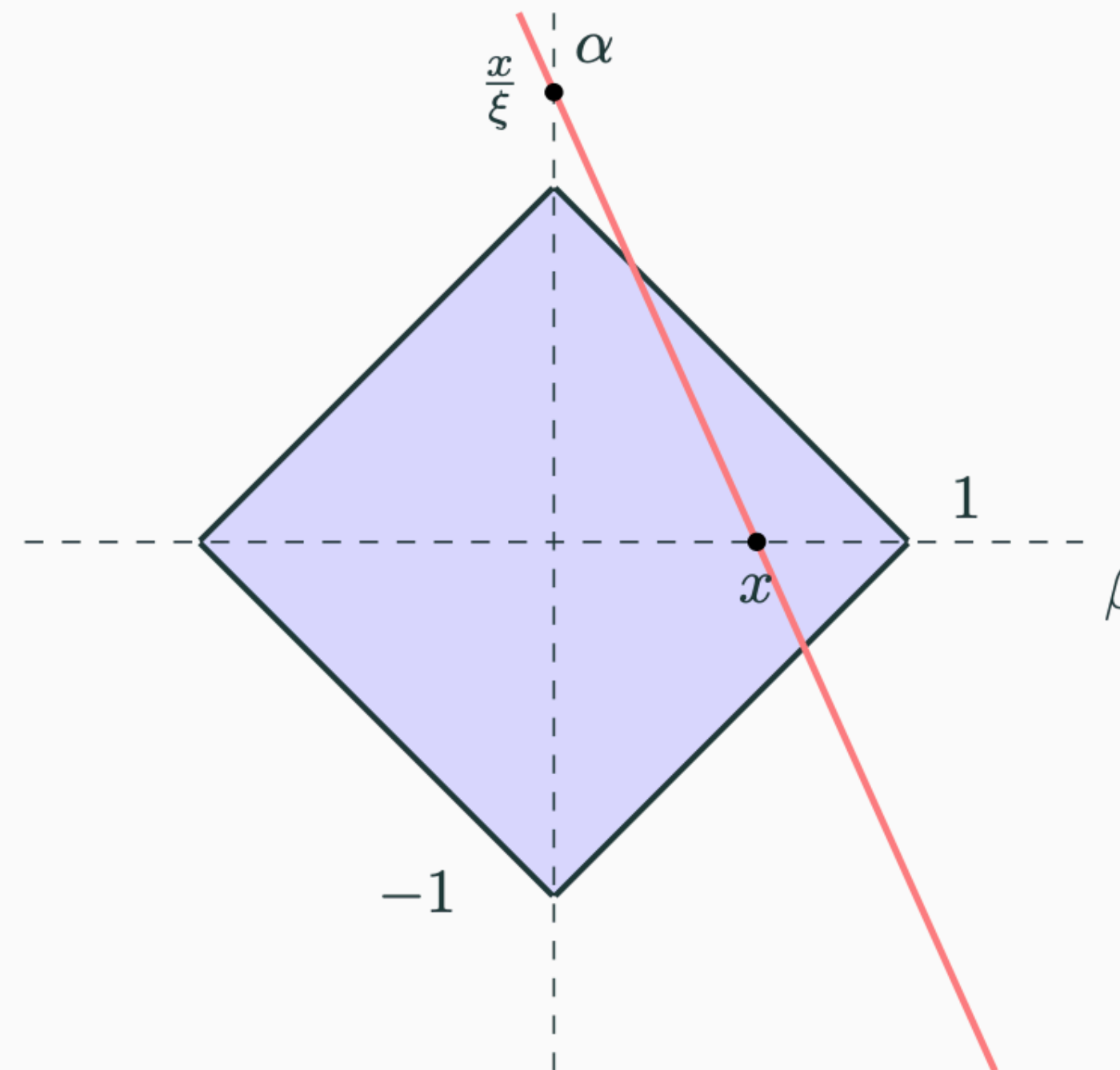
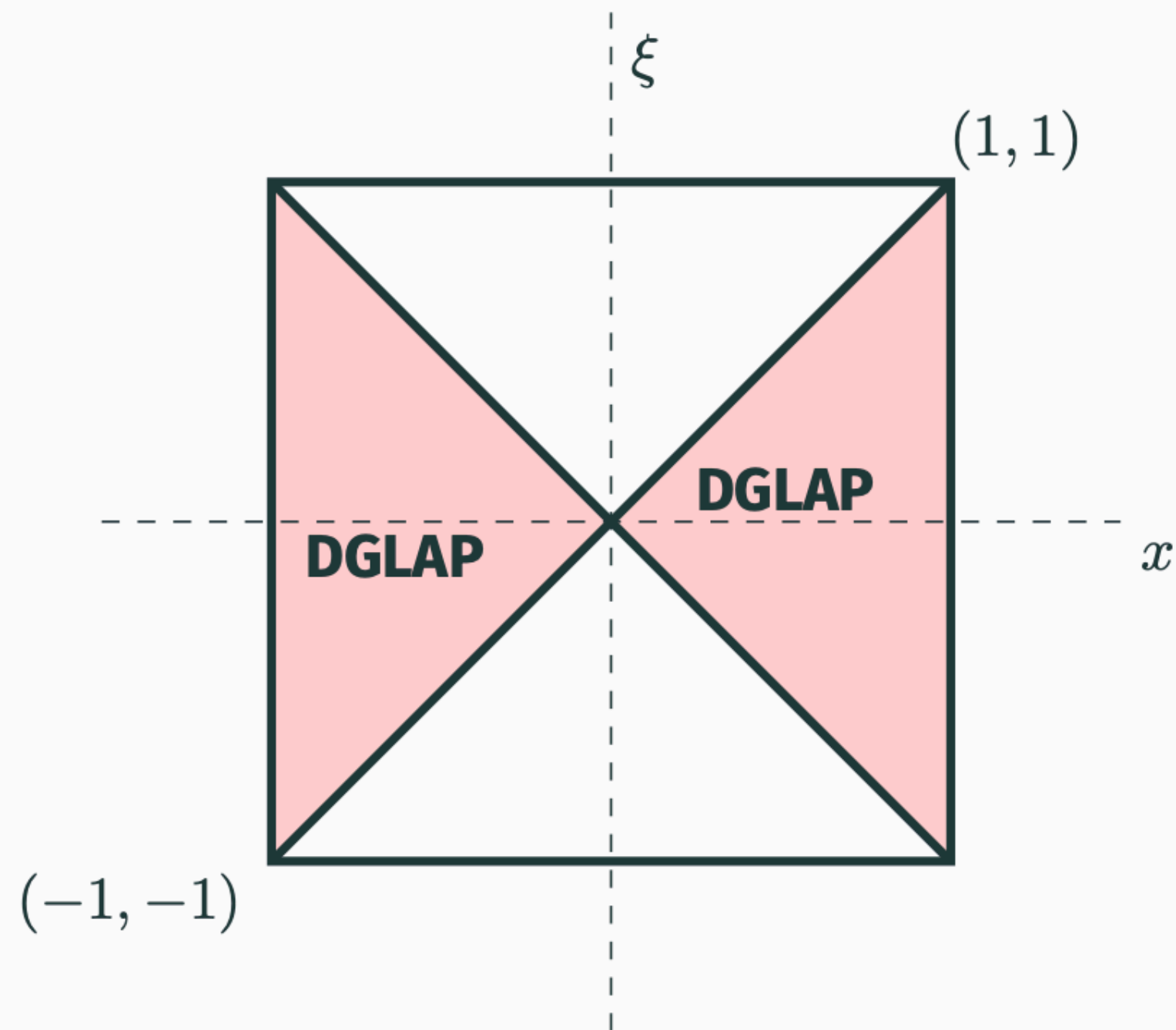


line $\alpha = -\frac{\beta}{\xi} + \frac{x}{\xi}$

$$H^q(x, \xi) = \mathcal{R}[h] = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha) + \theta(|\xi| - |x|) \text{ D-terms}$$

- Cannot be fixed by DGLAP data.
- Its DD has support only in $\beta = 0$
- Contribute only to $C_{m,m+1}(t)$

Support: $\Omega = \{(\beta, \alpha) \mid |\beta| + |\alpha| \leq 1\} = \Omega^+ \theta(\beta) + \Omega^- \theta(-\beta)$

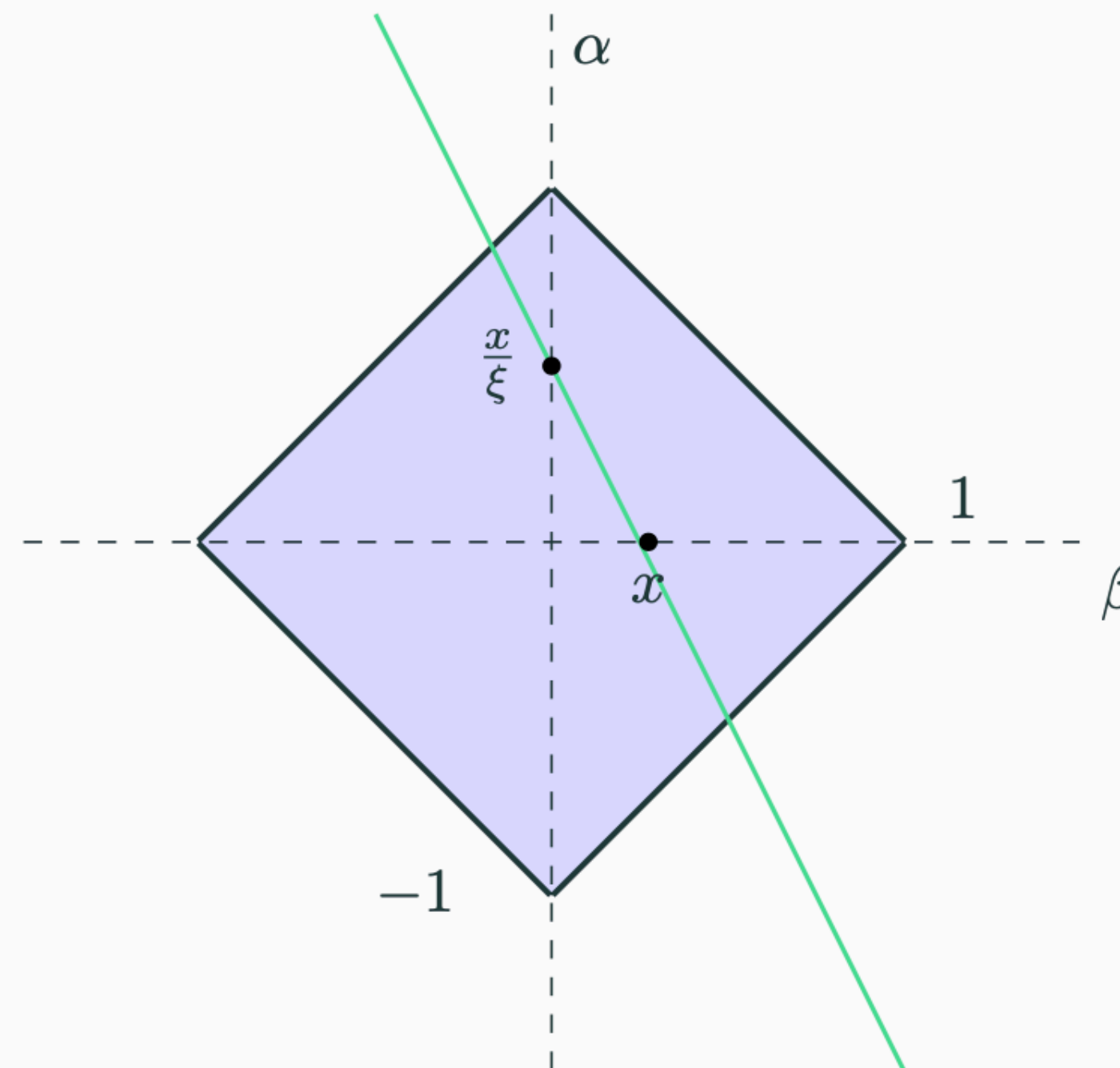
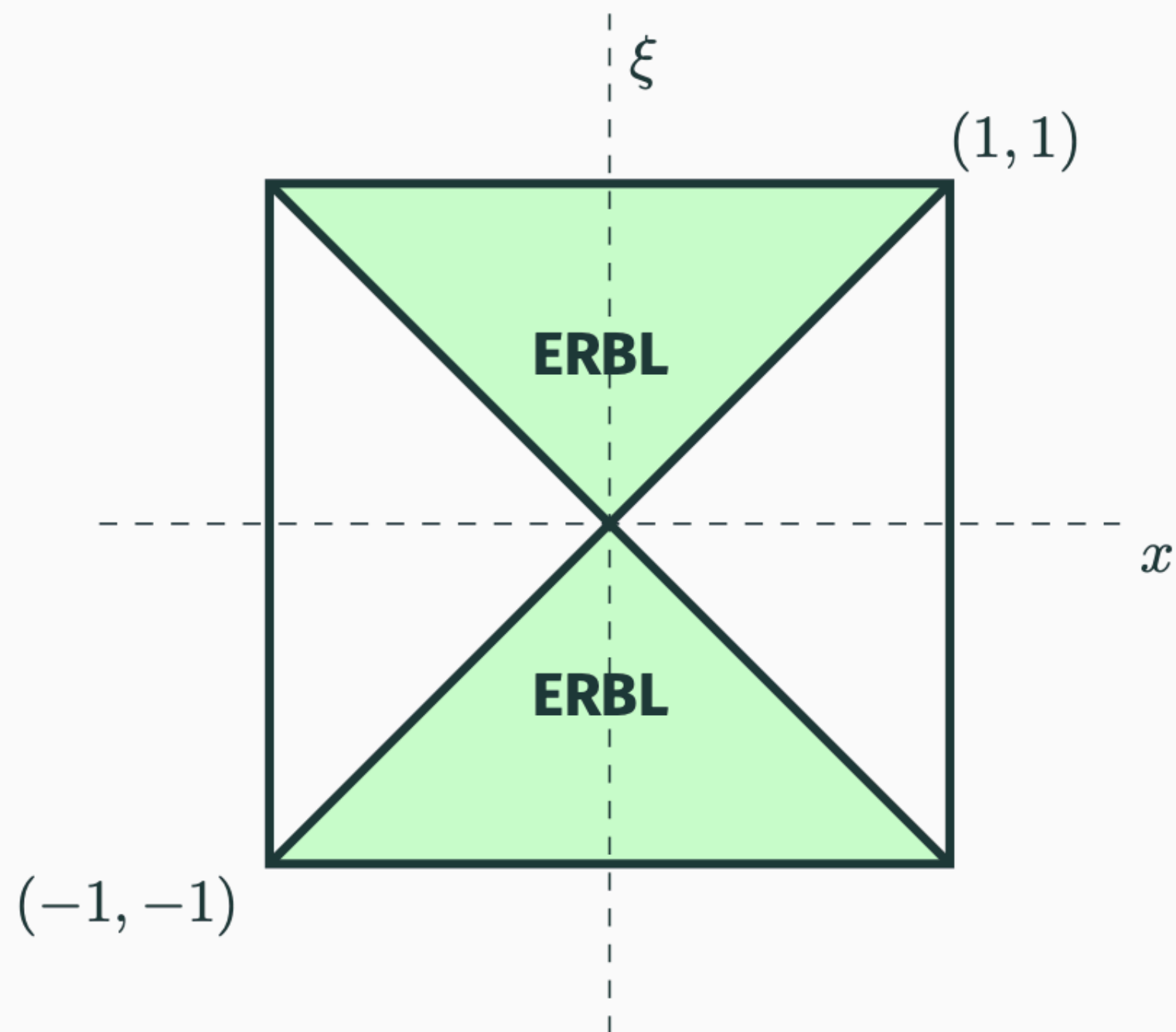


line $\alpha = -\frac{\beta}{\xi} + \frac{x}{\xi}$

$$H^q(x, \xi) = \mathcal{R}[h] = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha) + \theta(|\xi| - |x|) \text{ D-terms}$$

- Cannot be fixed by DGLAP data.
- Its DD has support only in $\beta = 0$
- Contribute only to $C_{m,m+1}(t)$

Support: $\Omega = \{(\beta, \alpha) \mid |\beta| + |\alpha| \leq 1\} = \Omega^+ \theta(\beta) + \Omega^- \theta(-\beta)$



line $\alpha = -\frac{\beta}{\xi} + \frac{x}{\xi}$

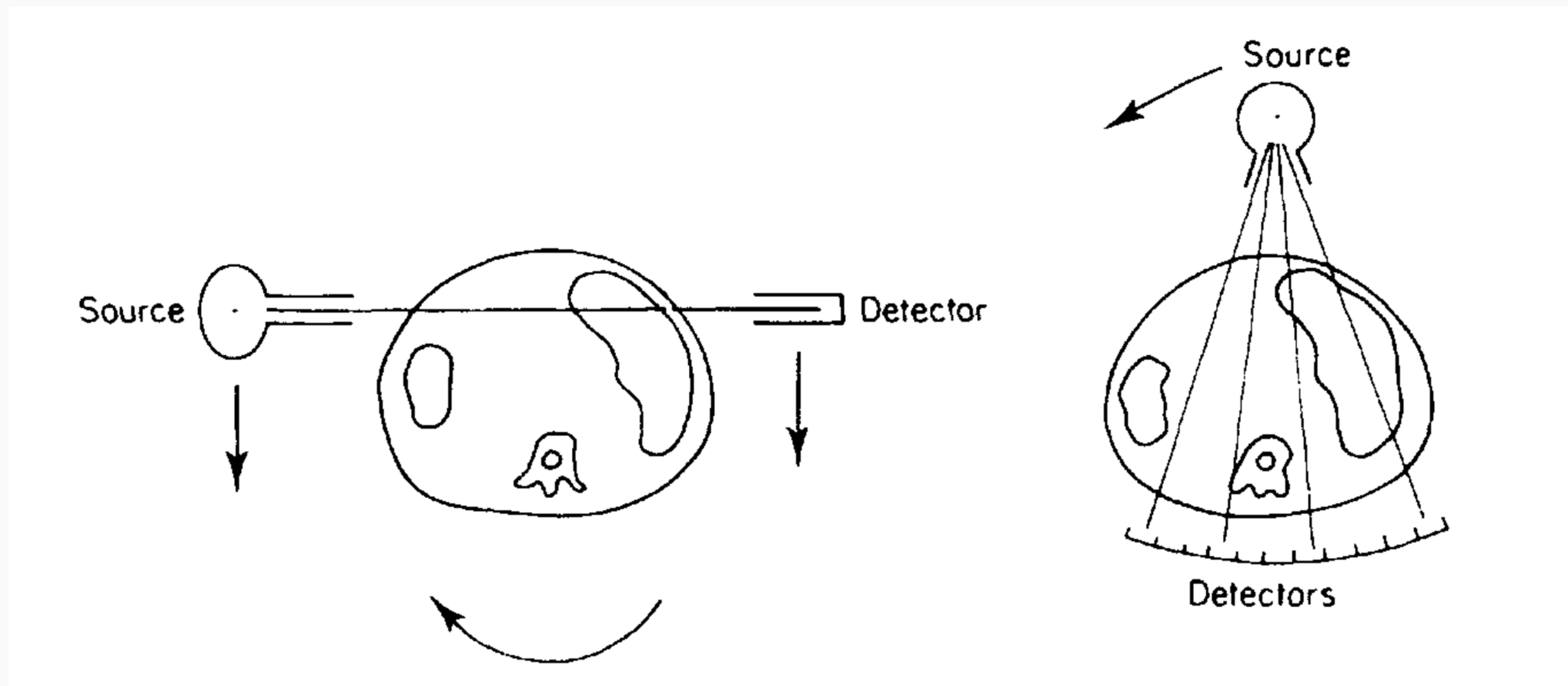
Covariant Extension \longleftrightarrow Radon transform inversion

Is it possible to invert the Radon transform knowing the GPD only in **DGLAP**?

Covariant Extension \longleftrightarrow Radon transform inversion

Is it possible to invert the Radon transform knowing the GPD only in **DGLAP**?

Answer from **Computerized Tomography**



Parallel scanning

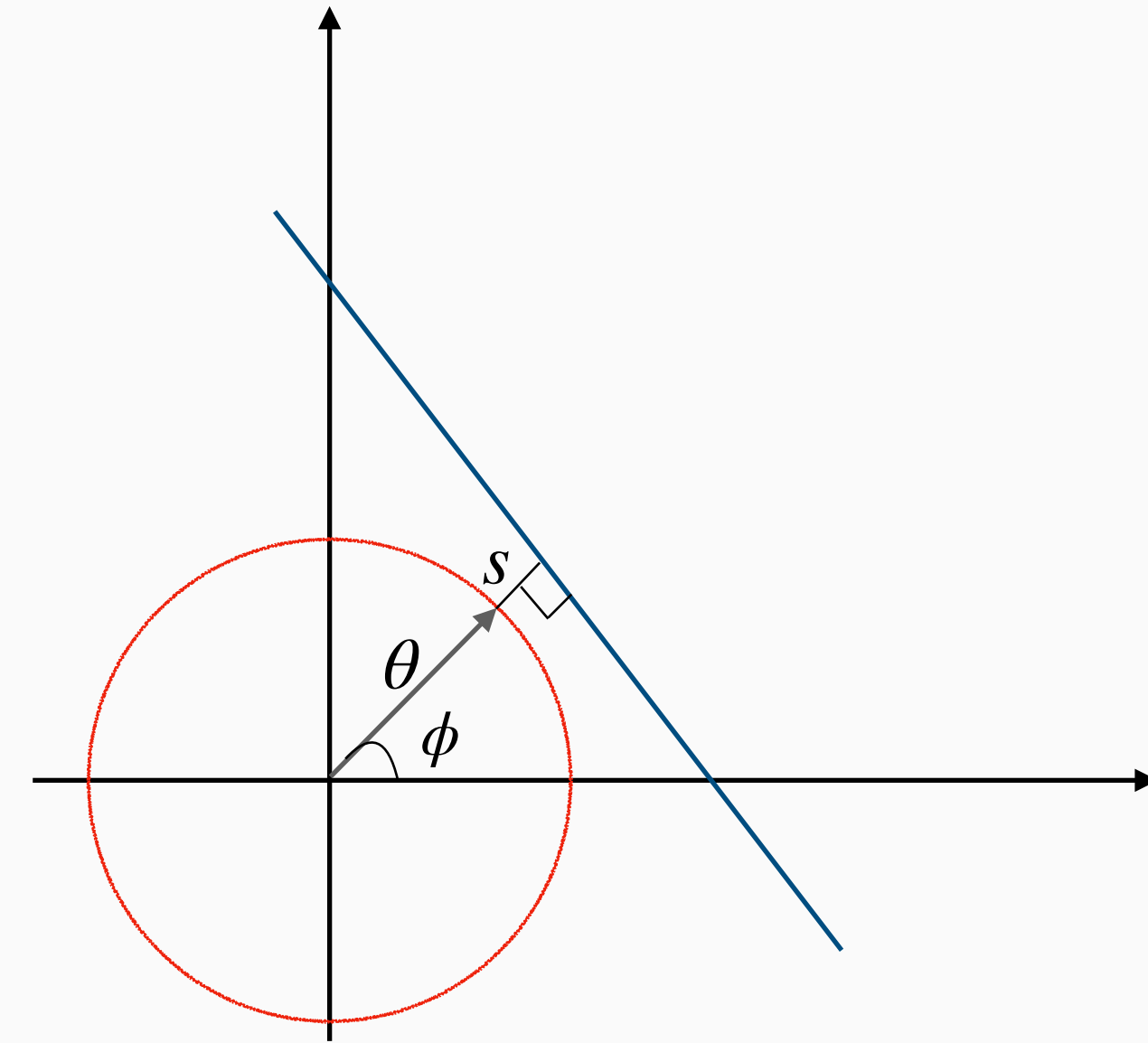
Fan-beam scanning

$$\frac{\Delta I}{I} = f(x) dx$$

$$\frac{I_{out}}{I_{in}} [L] = \exp \left\{ - \int_L f(x) dx \right\}$$

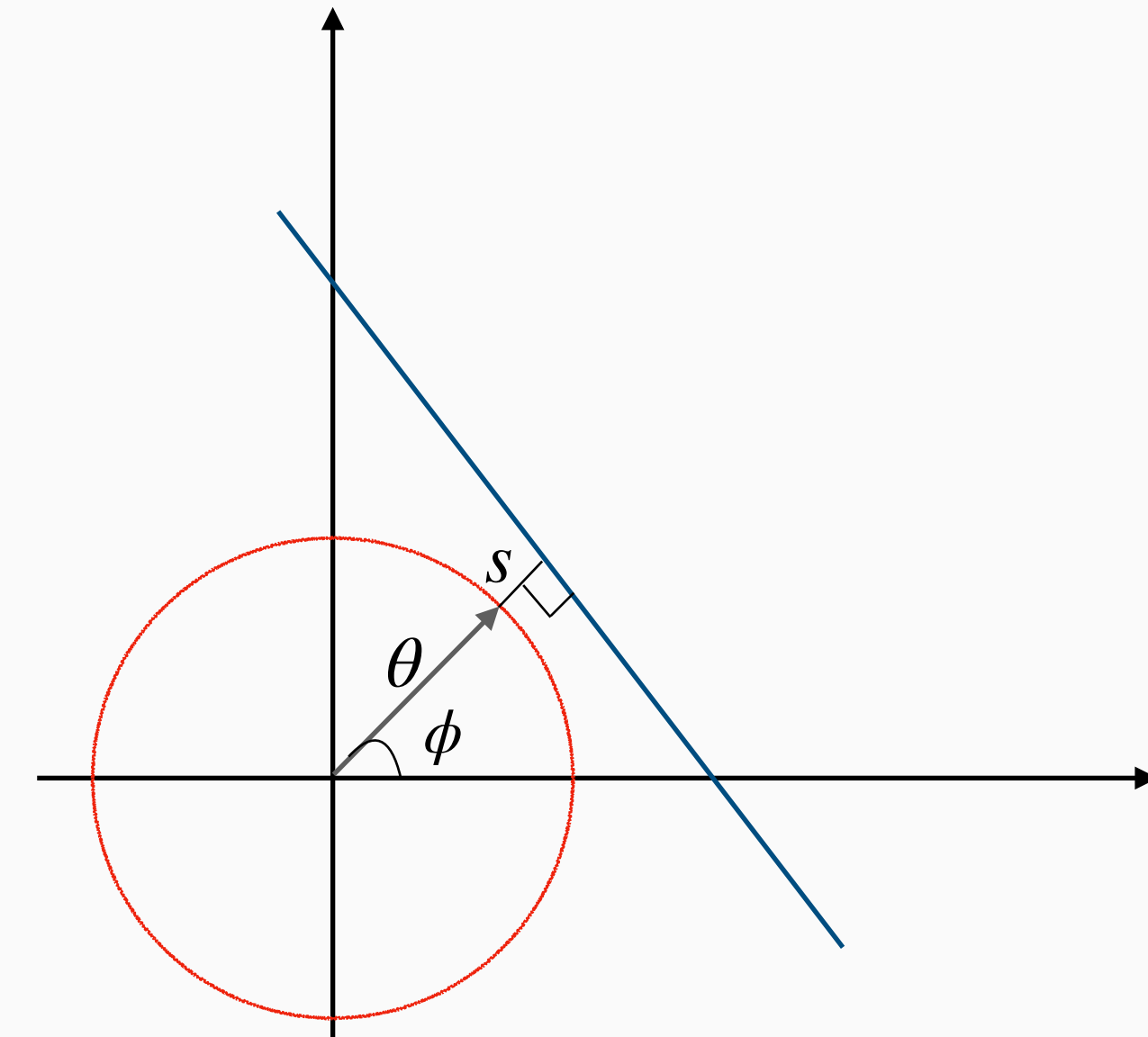
$$\mathcal{R}f(\theta, s) \equiv \int_{z \cdot \theta = s} dz f(z) \quad f \in \mathcal{S}(\mathbb{R}^n), \quad \theta \in S^{n-1}, \quad s \in \mathbb{R}$$

GPD: $z \rightarrow (\beta, \alpha), \quad \theta \rightarrow (\cos(\phi), \sin(\phi)), \quad x = \frac{s}{\cos(\phi)}, \quad \xi = \tan(\phi)$



$$\mathcal{R}f(\theta, s) \equiv \int_{z \cdot \theta = s} dz f(z) \quad f \in \mathcal{S}(\mathbb{R}^n), \quad \theta \in S^{n-1}, \quad s \in \mathbb{R}$$

GPD: $z \rightarrow (\beta, \alpha), \quad \theta \rightarrow (\cos(\phi), \sin(\phi)), \quad x = \frac{s}{\cos(\phi)}, \quad \xi = \tan(\phi)$



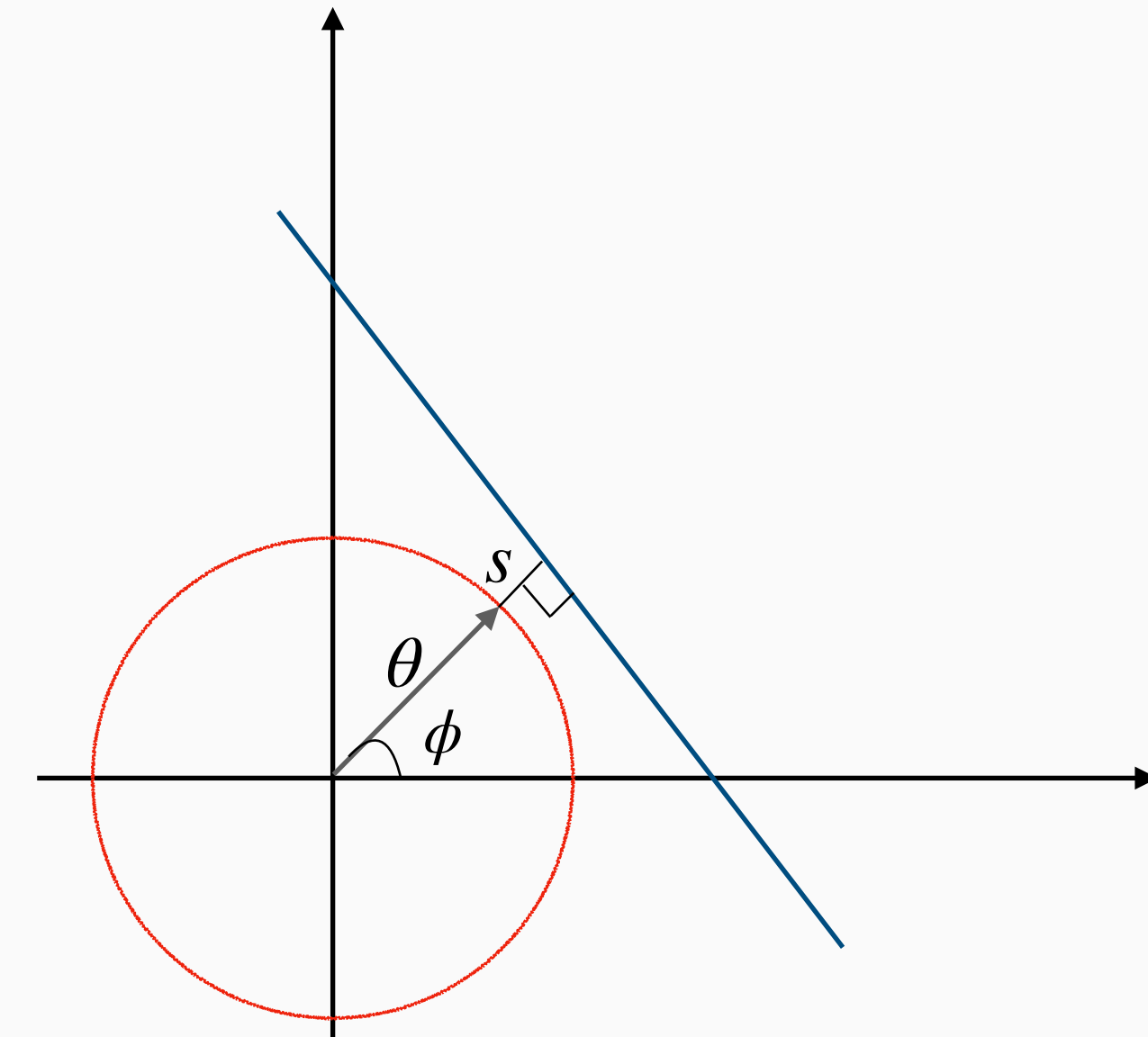
Uniqueness theorems:

[F. Natterer, *The Mathematics of Computerized Tomography*]

Given $f \in \mathcal{S}(\mathbb{R}^n)$, if $\mathcal{R}f(\theta, s) = 0 \quad \forall$ hyperplane $z \cdot \theta = s$ that does not intersect a compact convex set $K \subset \mathbb{R}^n \Rightarrow f(z) = 0$ outside K

$$\mathcal{R}f(\theta, s) \equiv \int_{z \cdot \theta = s} dz f(z) \quad f \in \mathcal{S}(\mathbb{R}^n), \quad \theta \in S^{n-1}, \quad s \in \mathbb{R}$$

GPD: $z \rightarrow (\beta, \alpha), \quad \theta \rightarrow (\cos(\phi), \sin(\phi)), \quad x = \frac{s}{\cos(\phi)}, \quad \xi = \tan(\phi)$



Uniqueness theorems:

[F. Natterer, *The Mathematics of Computerized Tomography*]

Given $f \in \mathcal{S}(\mathbb{R}^n)$, if $\mathcal{R}f(\theta, s) = 0 \quad \forall$ hyperplane $z \cdot \theta = s$ that does not intersect a compact convex set $K \subset \mathbb{R}^n \Rightarrow f(z) = 0$ outside K

→ **DD** is uniquely fixed by **GPD** in **DGLAP** (except for $\beta = 0$)

Given f in \mathbb{R}^2 compactly-supported and locally summable, $(\theta_0, s_0) \in S^1 \times \mathbb{R}$, U_0 open neighborhood of θ_0 :

$$\mathcal{R}f(\theta, s) = 0 \quad \forall s > s_0, \theta \in U_0 \quad \Rightarrow \quad f(z) = 0 \quad \text{if } z \cdot \theta_0 > s_0$$

Given f in \mathbb{R}^2 compactly-supported and locally summable, $(\theta_0, s_0) \in S^1 \times \mathbb{R}$, U_0 open neighborhood of θ_0 :

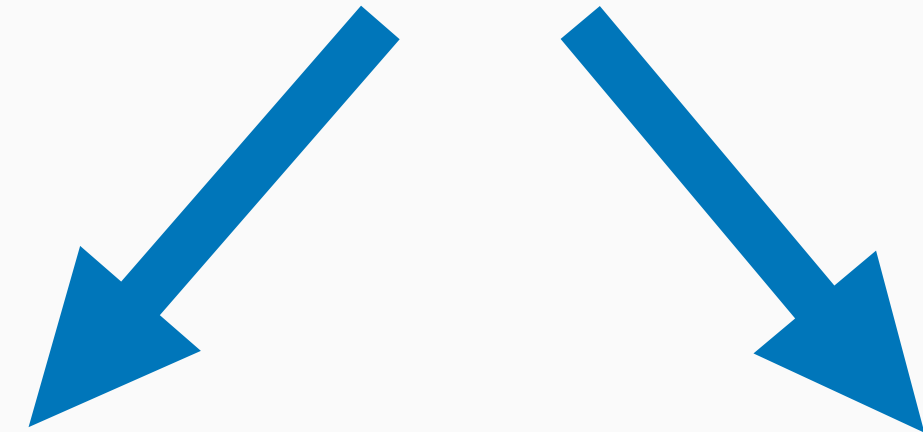
$$\mathcal{R}f(\theta, s) = 0 \quad \forall s > s_0, \theta \in U_0 \quad \Rightarrow \quad f(z) = 0 \quad \text{if } z \cdot \theta_0 > s_0$$

→ **DD** is uniquely fixed (except for $\beta = 0$) by knowing **GPD** in:

$$x \in [-1, 1], \quad \xi \in [0, \lambda x], \quad 0 < \lambda \leq 1 \quad (\text{H. Moutarde})$$

In experiments $\xi \lesssim 0.2$

Numerical Inversion of Radon Transform



Artificial Neural Networks

H. Dutrieux et al. *Eur.Phys.J.C.* 82 (2022)

Finite Elements Methods

N. Chouika et al. *Eur.Phys.J.C.* 77 (2017)

J.M. Morgado et al. *Phys.Rev.D* 105 (2022)

Numerical Inversion of Radon Transform

Artificial Neural Networks

H. Dutrieux et al. *Eur.Phys.J.C.* 82 (2022)

Finite Elements Methods

N. Chouika et al. *Eur.Phys.J.C.* 77 (2017)

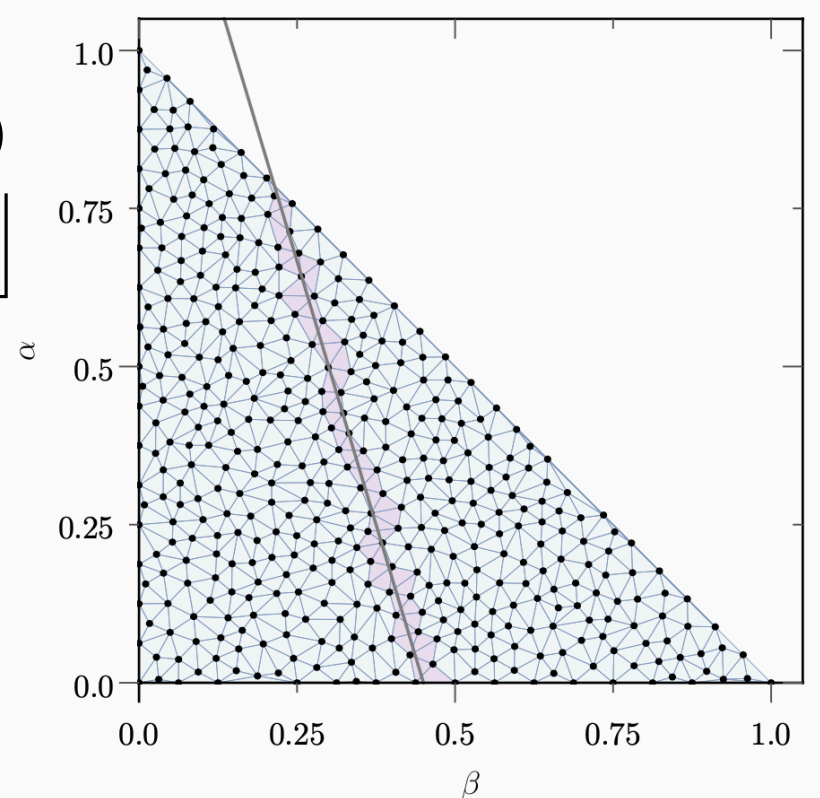
J.M. Morgado et al. *Phys.Rev.D* 105 (2022)

Lagrange Polynomial

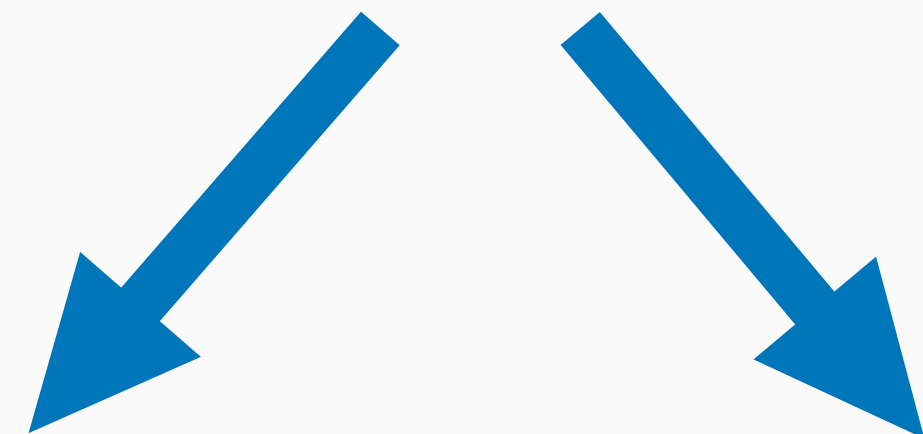
$$H_i \equiv H(x_i, \xi_i) = \sum_j h_j \int_{\Omega^+} d\beta d\alpha \delta(x_i - \beta - \alpha \xi_i) v_j(\beta, \alpha)$$

$\downarrow R_{ij}$

$h = (R^T R)^{-1} R^T H$

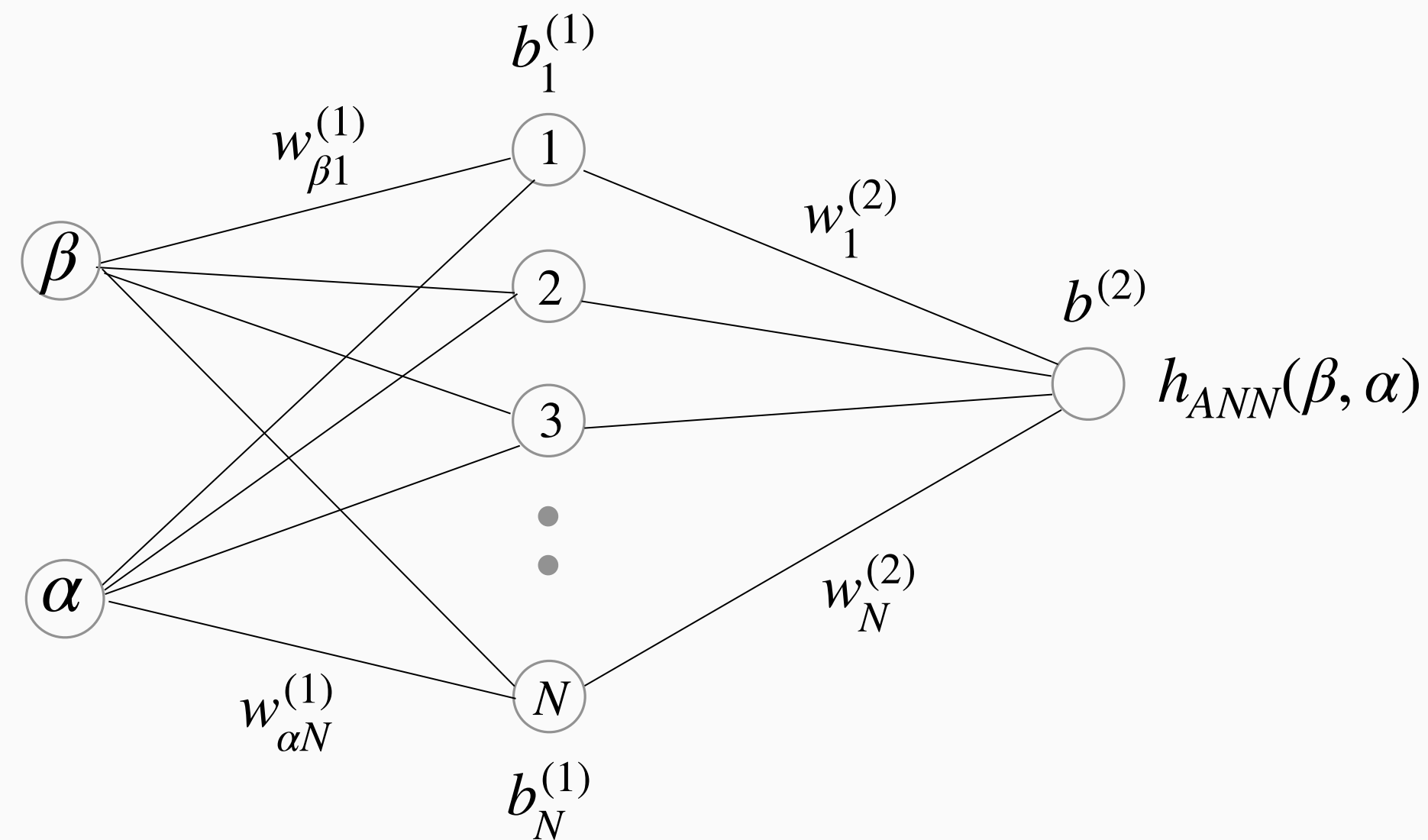


Numerical Inversion of Radon Transform



Artificial Neural Networks

H. Dutrieux et al. *Eur.Phys.J.C.* 82 (2022)



Unbiased parametrization of DD

Finite Elements Methods

N. Chouika et al. *Eur.Phys.J.C.* 77 (2017)

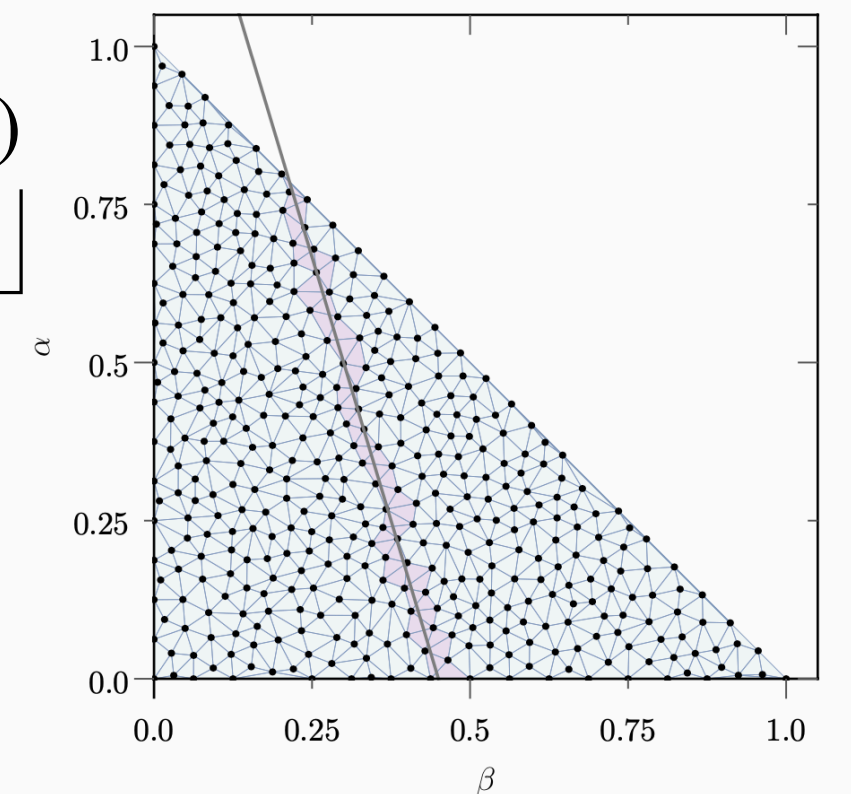
J.M. Morgado et al. *Phys.Rev.D* 105 (2022)

$$H_i \equiv H(x_i, \xi_i) = \sum_j h_j \int_{\Omega^+} d\beta d\alpha \delta(x_i - \beta - \alpha \xi_i) v_j(\beta, \alpha)$$

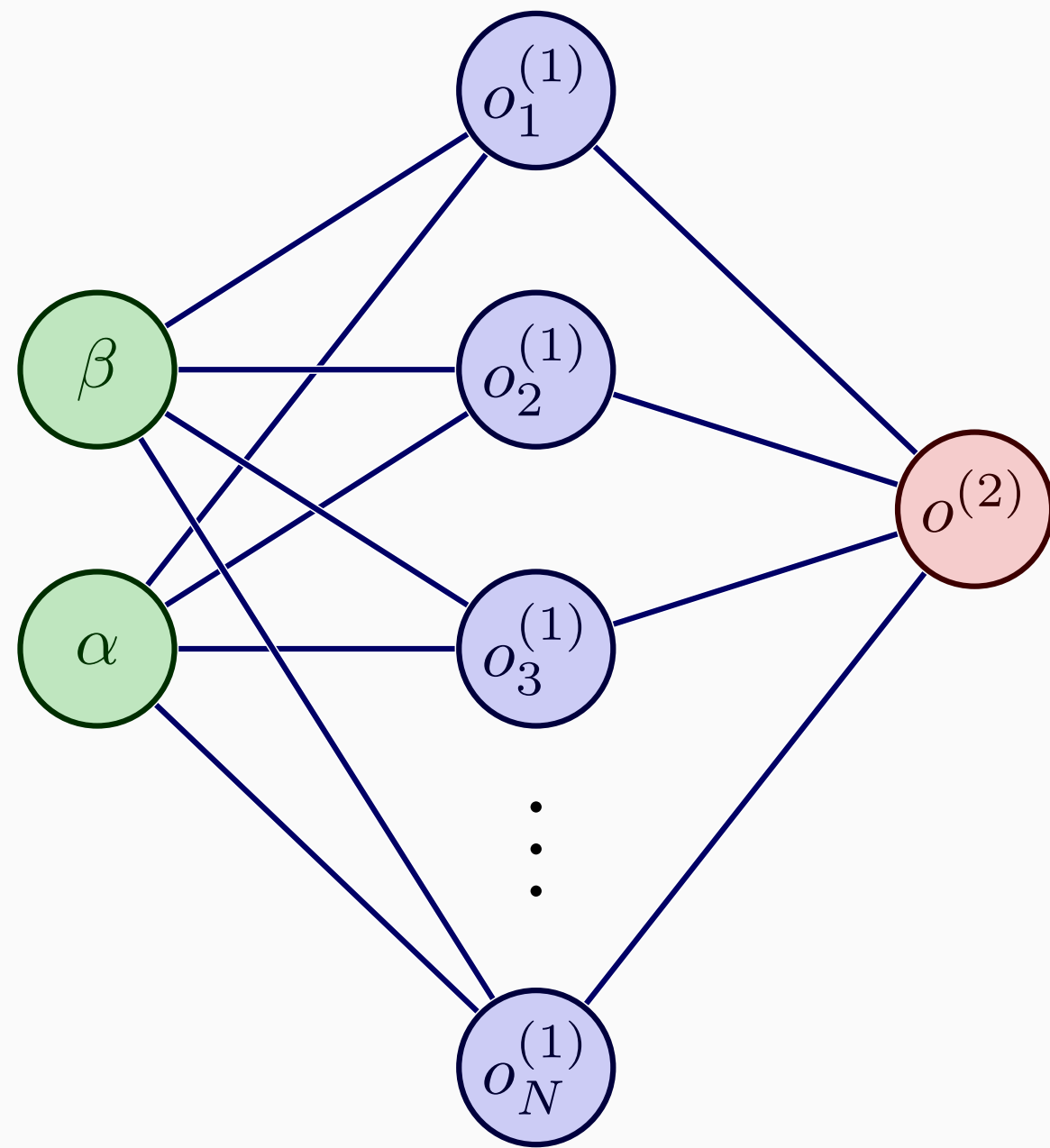
Lagrange Polynomial

R_{ij}

$$h = (R^T R)^{-1} R^T H$$



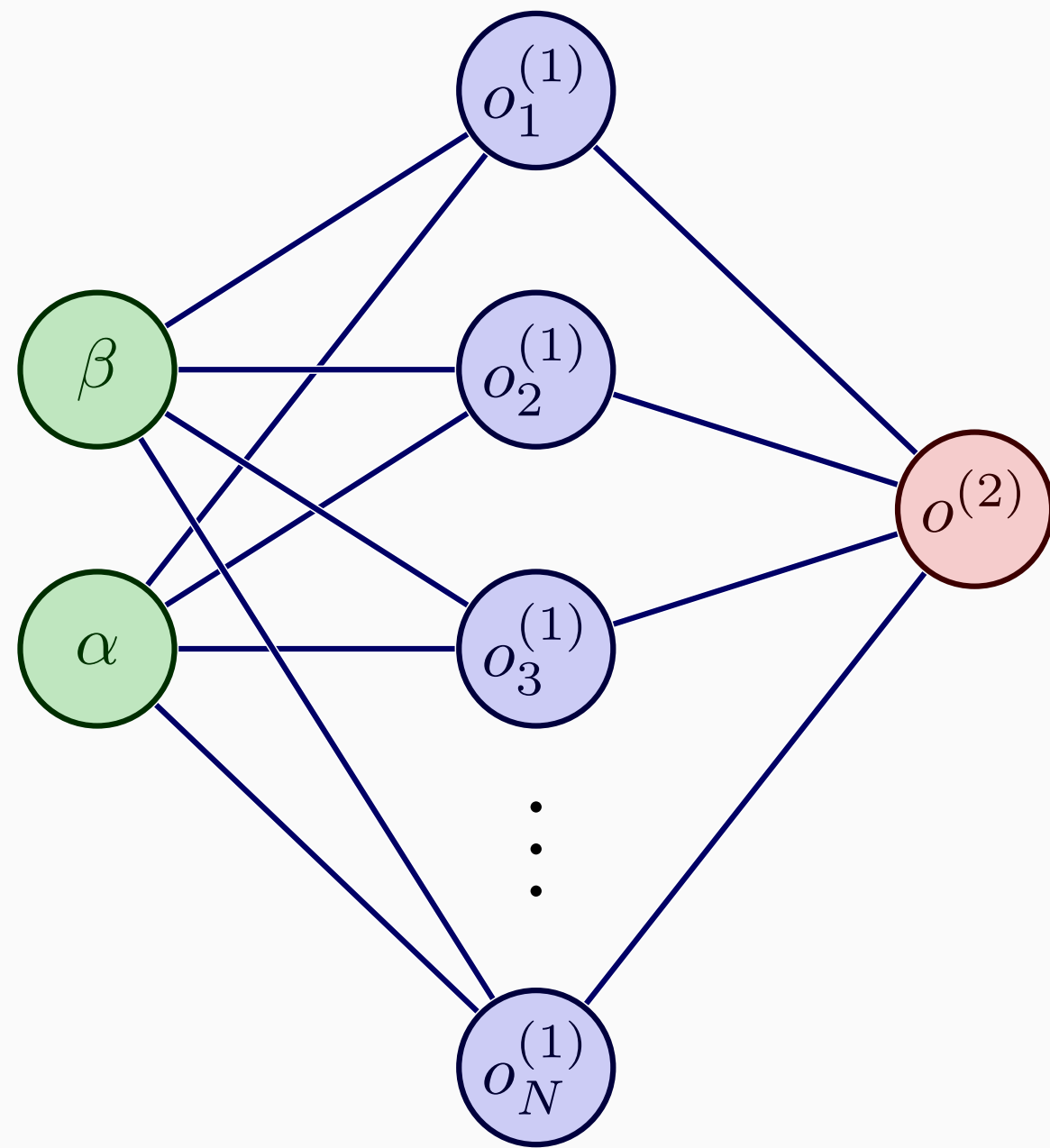
Parametrizing $h(\beta, \alpha)$ using **Artificial Neural Networks (ANN)**



$$\begin{aligned} = h_{ANN}(\beta, \alpha) &= \sum_{i=1}^N w_i^{(2)} o_i^{(1)} + b^{(2)} \\ &= \sum_{i=1}^N w_i^{(2)} \left[\sigma \left(w_{\beta i}^{(1)} \beta + w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) + \sigma \left(w_{\beta i}^{(1)} \beta - w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) \right] + b^{(2)} \end{aligned}$$

- Adam optimizer (gradient descent)
- Dropout regularization

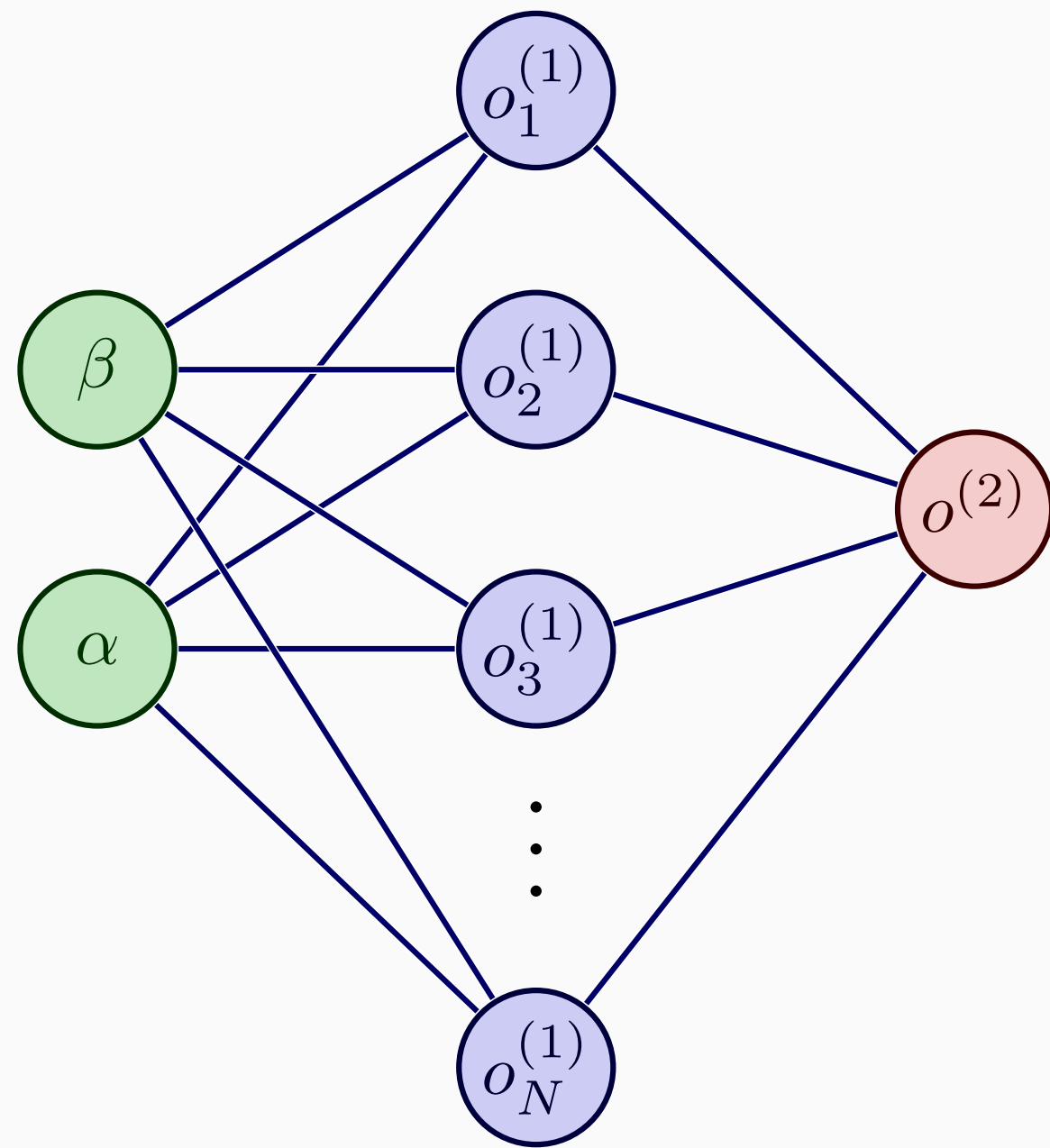
Parametrizing $h(\beta, \alpha)$ using **Artificial Neural Networks (ANN)**



$$\begin{aligned}
 = h_{ANN}(\beta, \alpha) &= \sum_{i=1}^N w_i^{(2)} o_i^{(1)} + b^{(2)} & h(\beta, -\alpha) = h(\beta, \alpha) \leftrightarrow H(x, -\xi) = H(x, \xi) \\
 &= \sum_{i=1}^N w_i^{(2)} \left[\sigma \left(w_{\beta i}^{(1)} \beta + w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) + \sigma \left(w_{\beta i}^{(1)} \beta - w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right) \right] + b^{(2)}
 \end{aligned}$$

- Adam optimizer (gradient descent)
- Dropout regularization

Parametrizing $h(\beta, \alpha)$ using **Artificial Neural Networks (ANN)**



$$\begin{aligned}
 = h_{ANN}(\beta, \alpha) &= \sum_{i=1}^N w_i^{(2)} o_i^{(1)} + b^{(2)} & h(\beta, -\alpha) = h(\beta, \alpha) \leftrightarrow H(x, -\xi) = H(x, \xi) \\
 &= \sum_{i=1}^N w_i^{(2)} \left[\underset{\downarrow}{\sigma \left(w_{\beta i}^{(1)} \beta + w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right)} + \underset{\uparrow}{\sigma \left(w_{\beta i}^{(1)} \beta - w_{\alpha i}^{(1)} \alpha + b_i^{(1)} \right)} \right] + b^{(2)} \\
 &\quad \sigma(x) = \frac{1}{1 + e^{-x}}
 \end{aligned}$$

- Adam optimizer (gradient descent)
- Dropout regularization

Algorithm

- Initialize the ANN parameters (randomly).
- Given a sampling set of GPD values $H_i(x_i, \xi_i)$ in the DGLAP region, numerically evaluate the RT along each line (x_i, ξ_i) using $h_{ANN}(\beta, \alpha)$ as DD. $\mathcal{R}h_{ANN}(x_i, \xi_i) = \hat{H}(x_i, \xi_i)$.

- Update the ANN parameters using Adam optimization algorithm in order to minimize

$$MSE = \frac{1}{N_{sample}} \sum_{i=1}^{N_{sample}} \left(H_i - \hat{H}_i \right)^2$$

- Iterate until convergence.

Testing with analytical models

Nakanishi based model for pion

N.Chouika et al. *Phys.Lett.B*:780(2018)

$$H(x, \xi, t = 0) = \begin{cases} 30 \frac{(1-x)^2(x^2 - \xi^2)}{(1 - \xi^2)^2}, & |x| > \xi \\ 15 \frac{(1-x)(\xi^2 - x^2)(x + 2x\xi + \xi^2)}{2\xi^3(1 + \xi)^2}, & |x| < \xi \end{cases}$$

$$H(x, \xi, t = 0) = (1 - x) \int_{\Omega^+} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha)$$
$$h(\beta, \alpha) = \frac{15}{2} (1 - 3(\alpha^2 - \beta^2) - 2\beta)$$

Testing with analytical models

Nakanishi based model for pion

N.Chouika et al. *Phys.Lett.B*:780(2018)

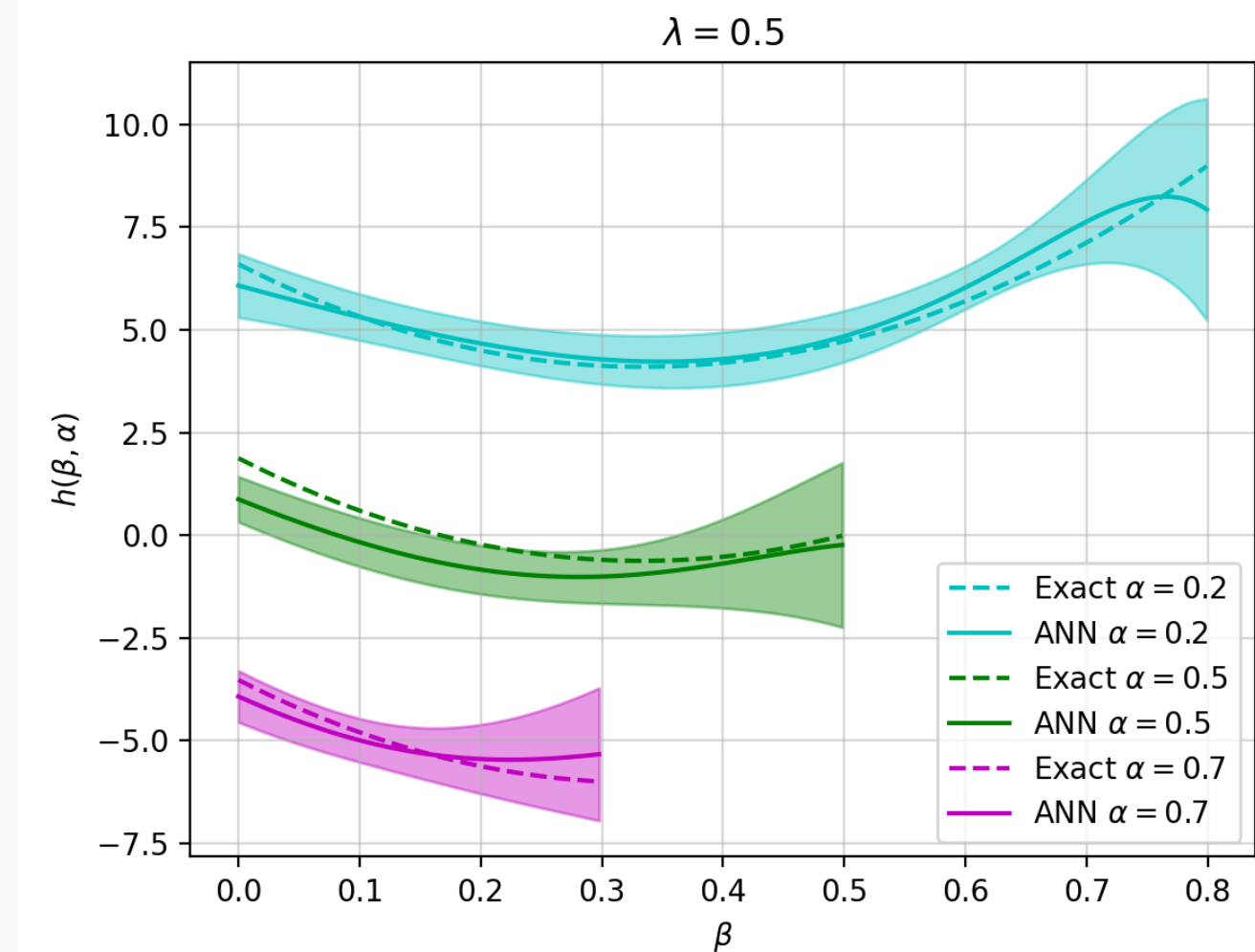
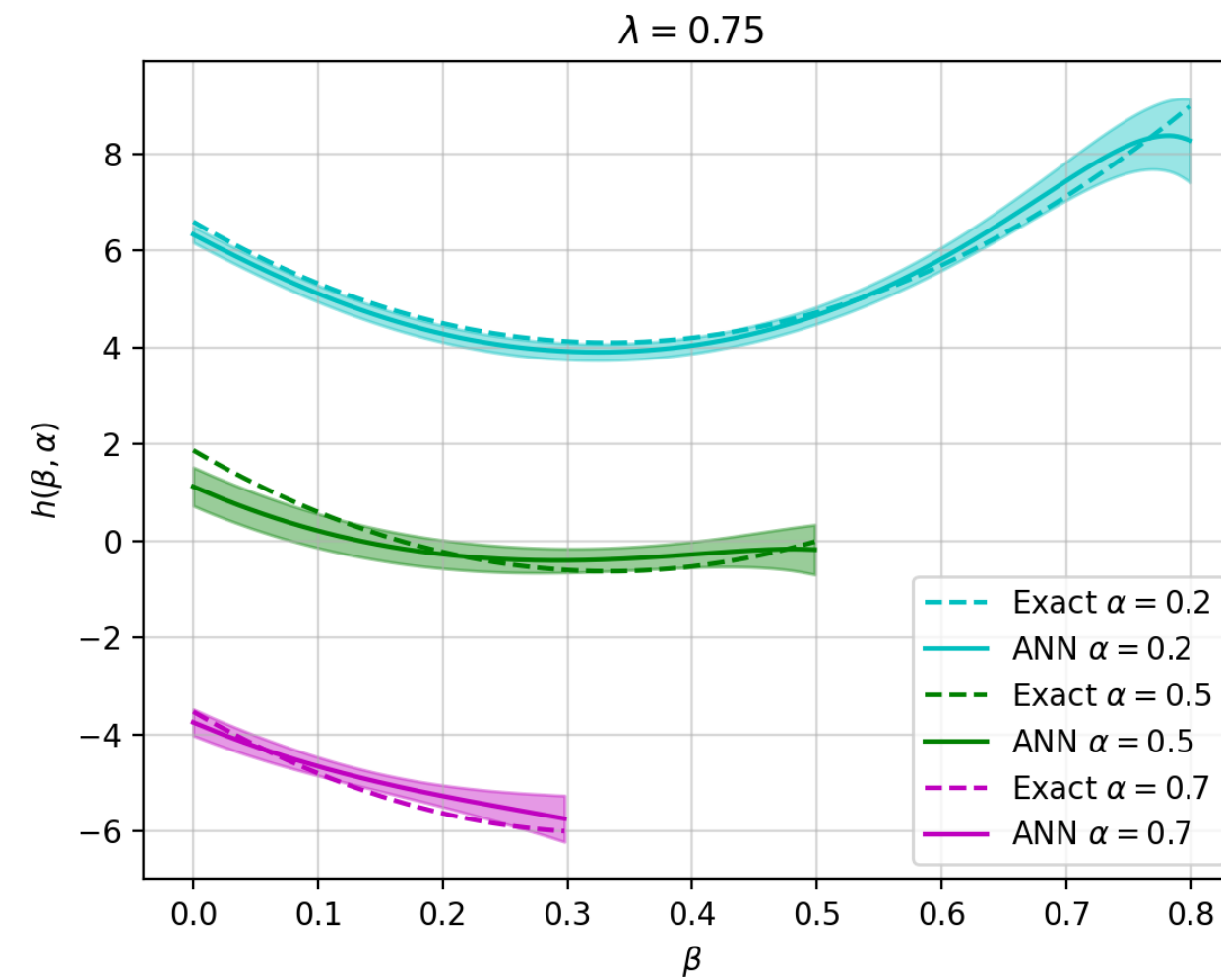
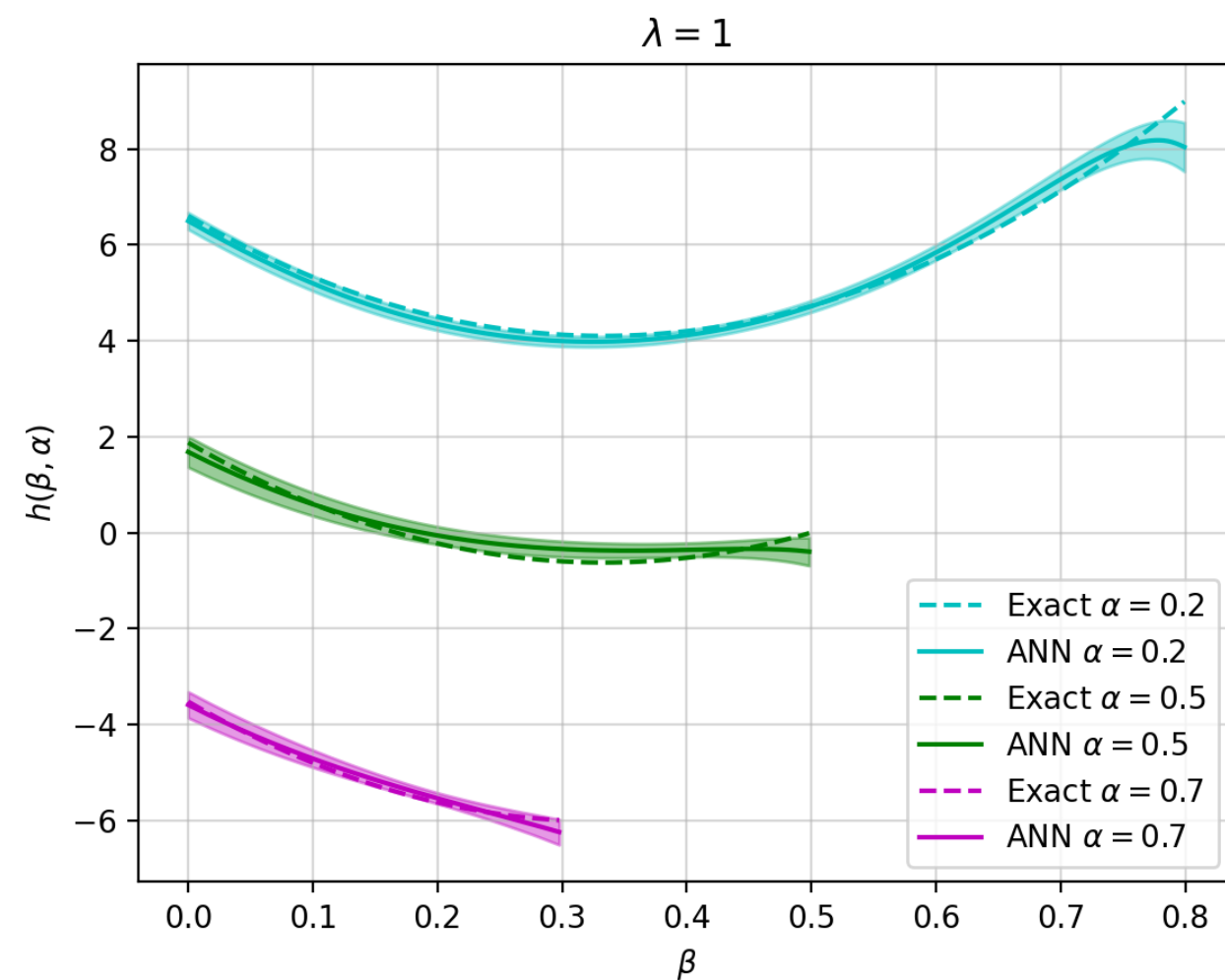
$$H(x, \xi, t = 0) = \begin{cases} 30 \frac{(1-x)^2(x^2 - \xi^2)}{(1 - \xi^2)^2}, & |x| > \xi \\ 15 \frac{(1-x)(\xi^2 - x^2)(x + 2x\xi + \xi^2)}{2\xi^3(1 + \xi)^2}, & |x| < \xi \end{cases}$$

$$H(x, \xi, t = 0) = (1 - x) \int_{\Omega^+} d\beta d\alpha \delta(x - \beta - \alpha\xi) h(\beta, \alpha)$$

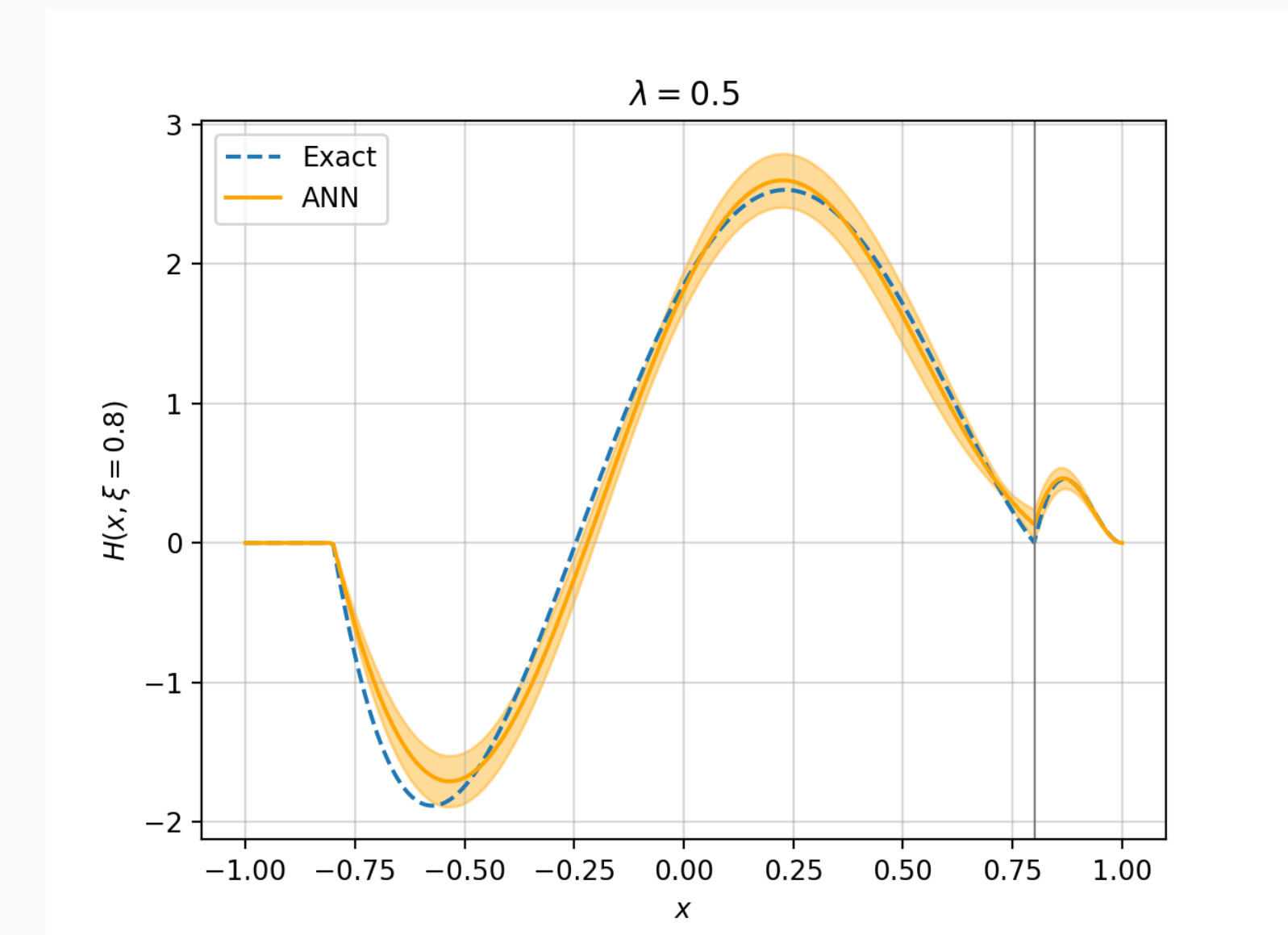
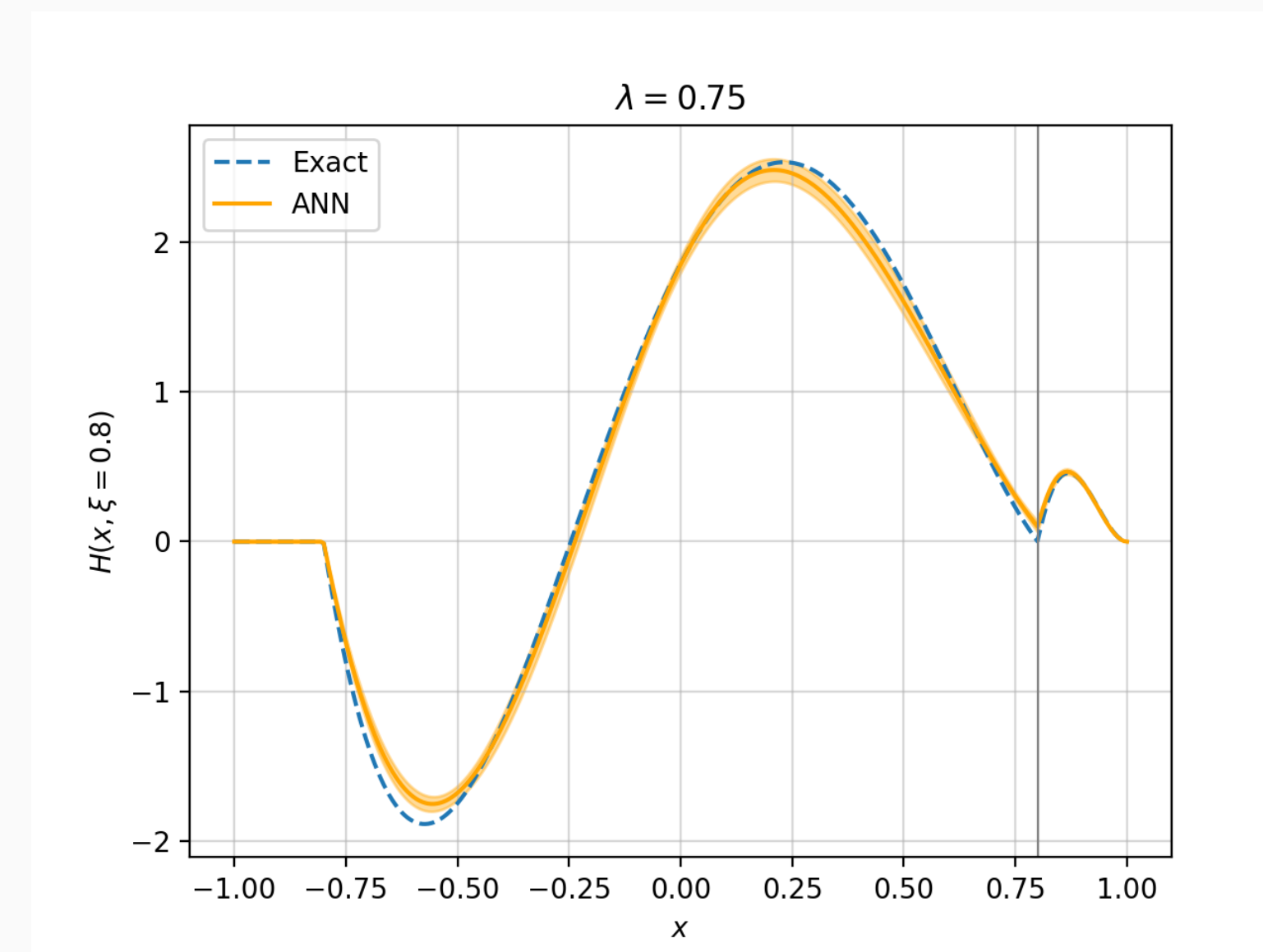
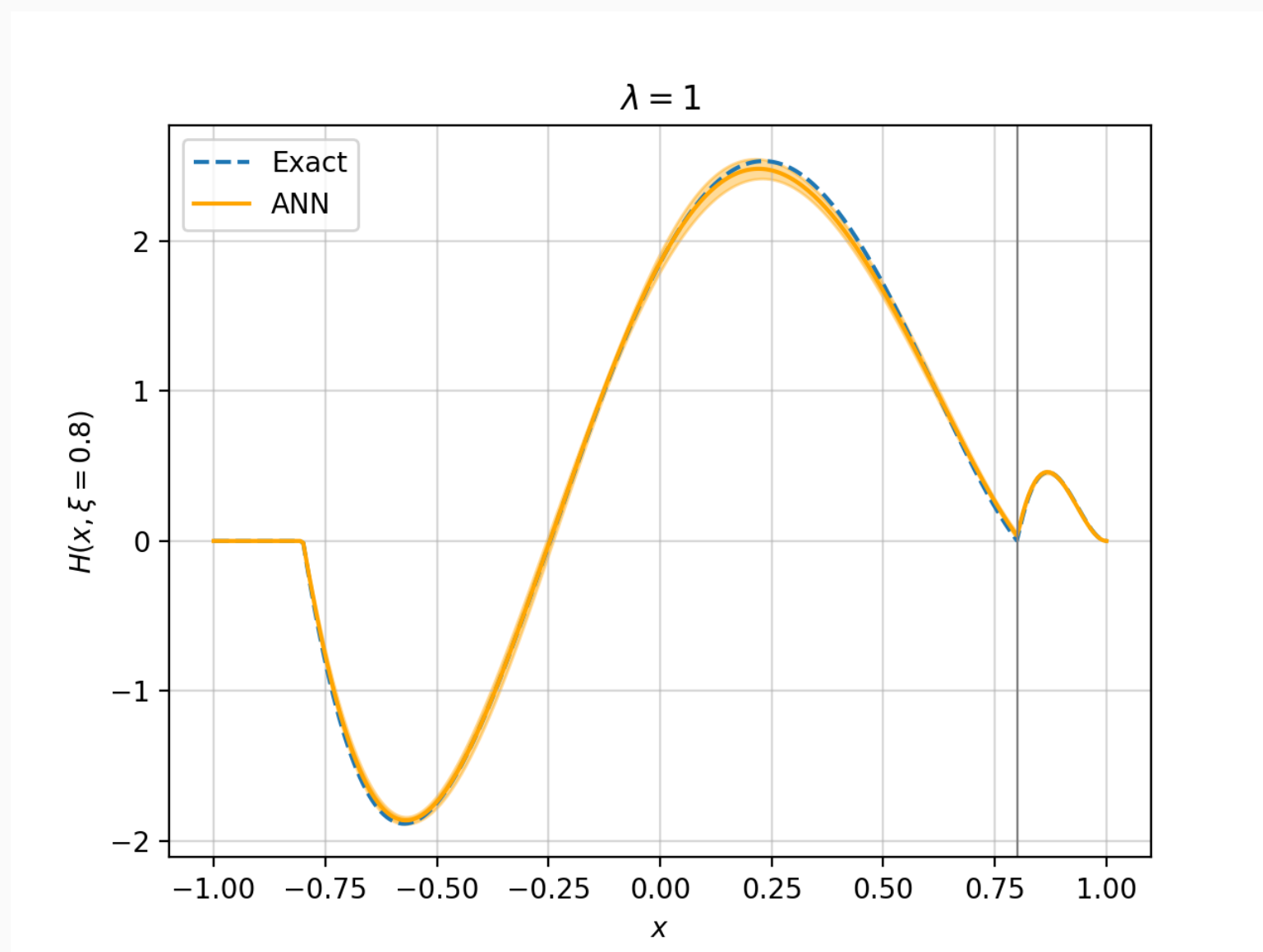
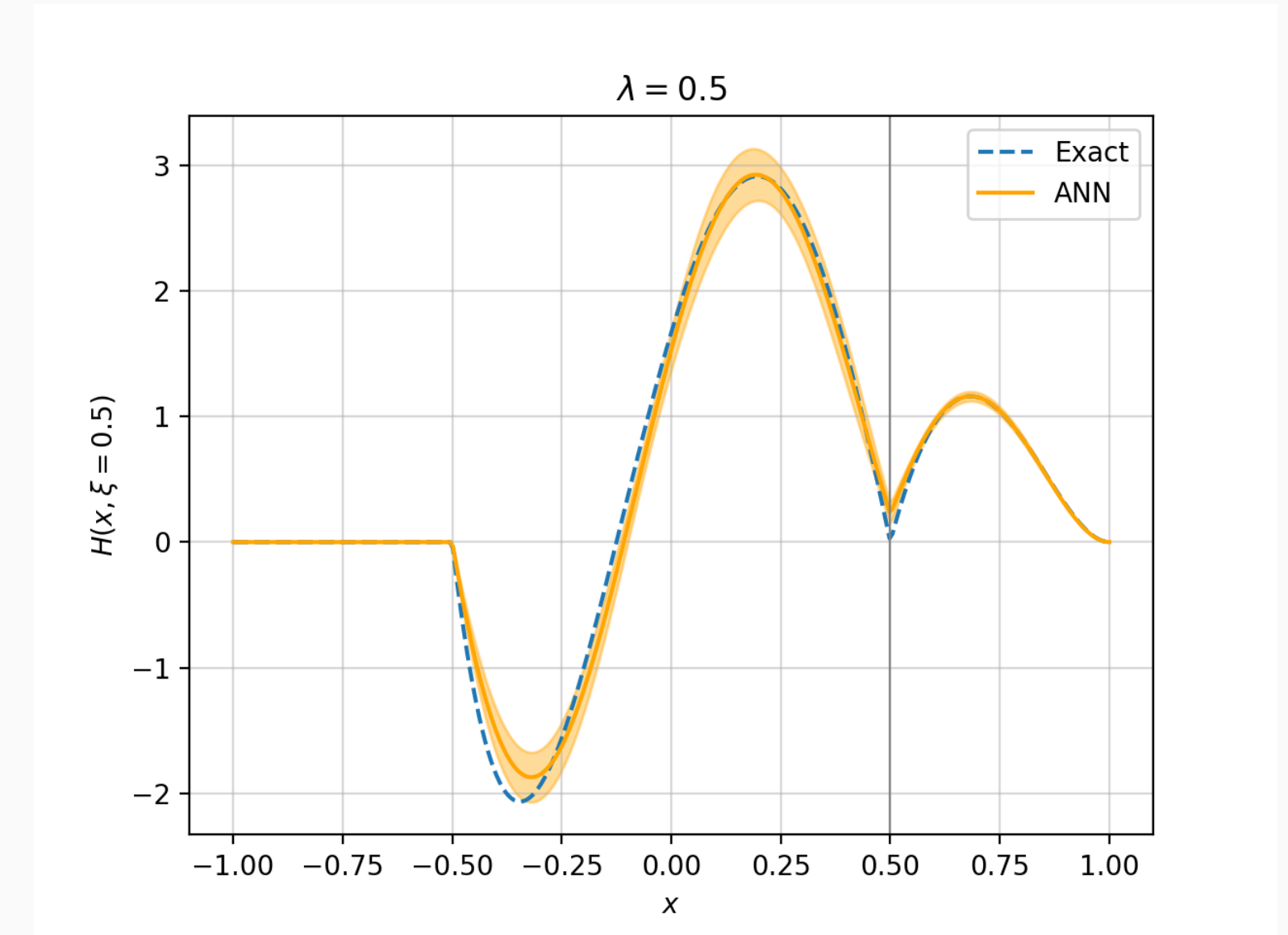
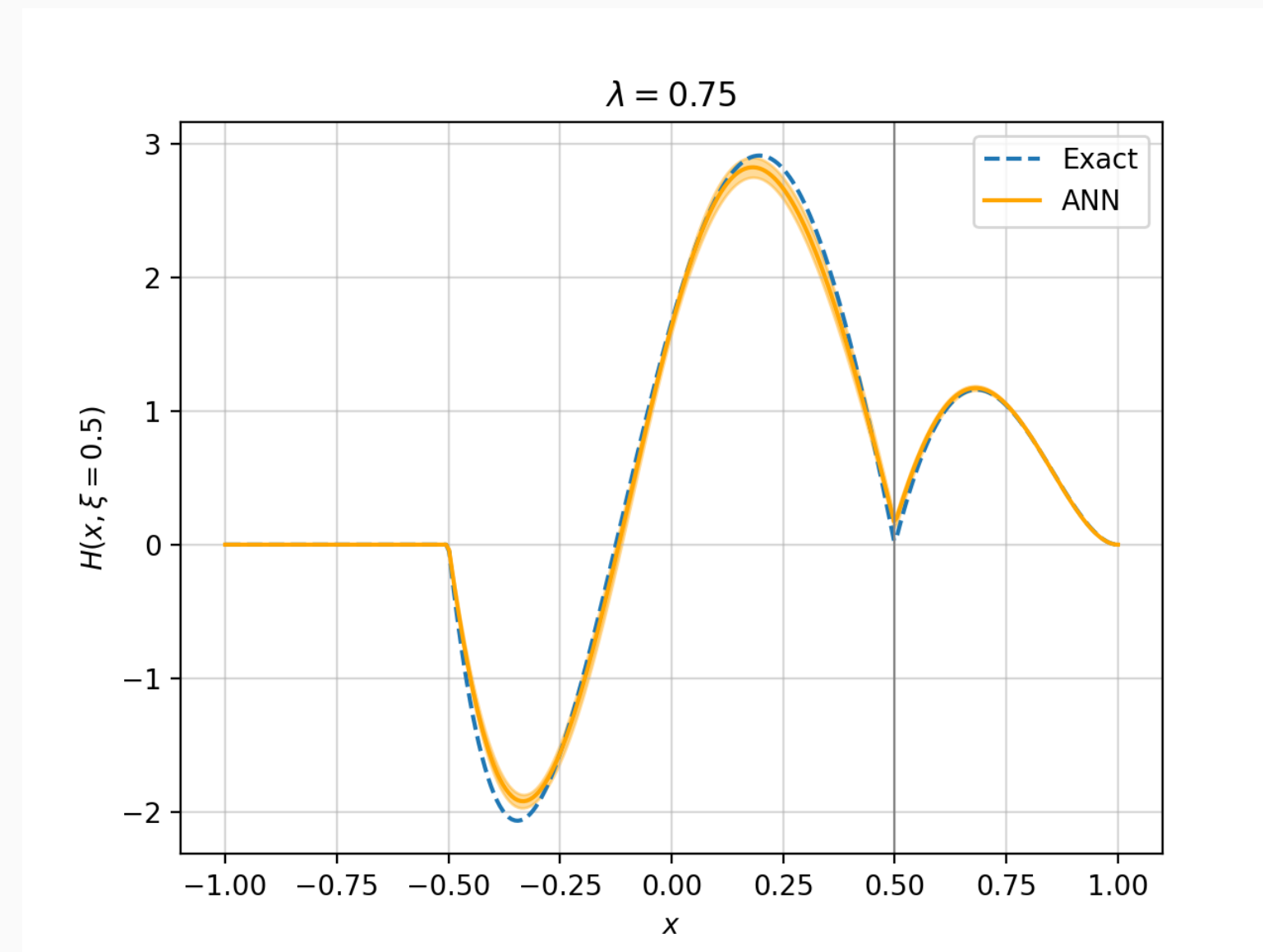
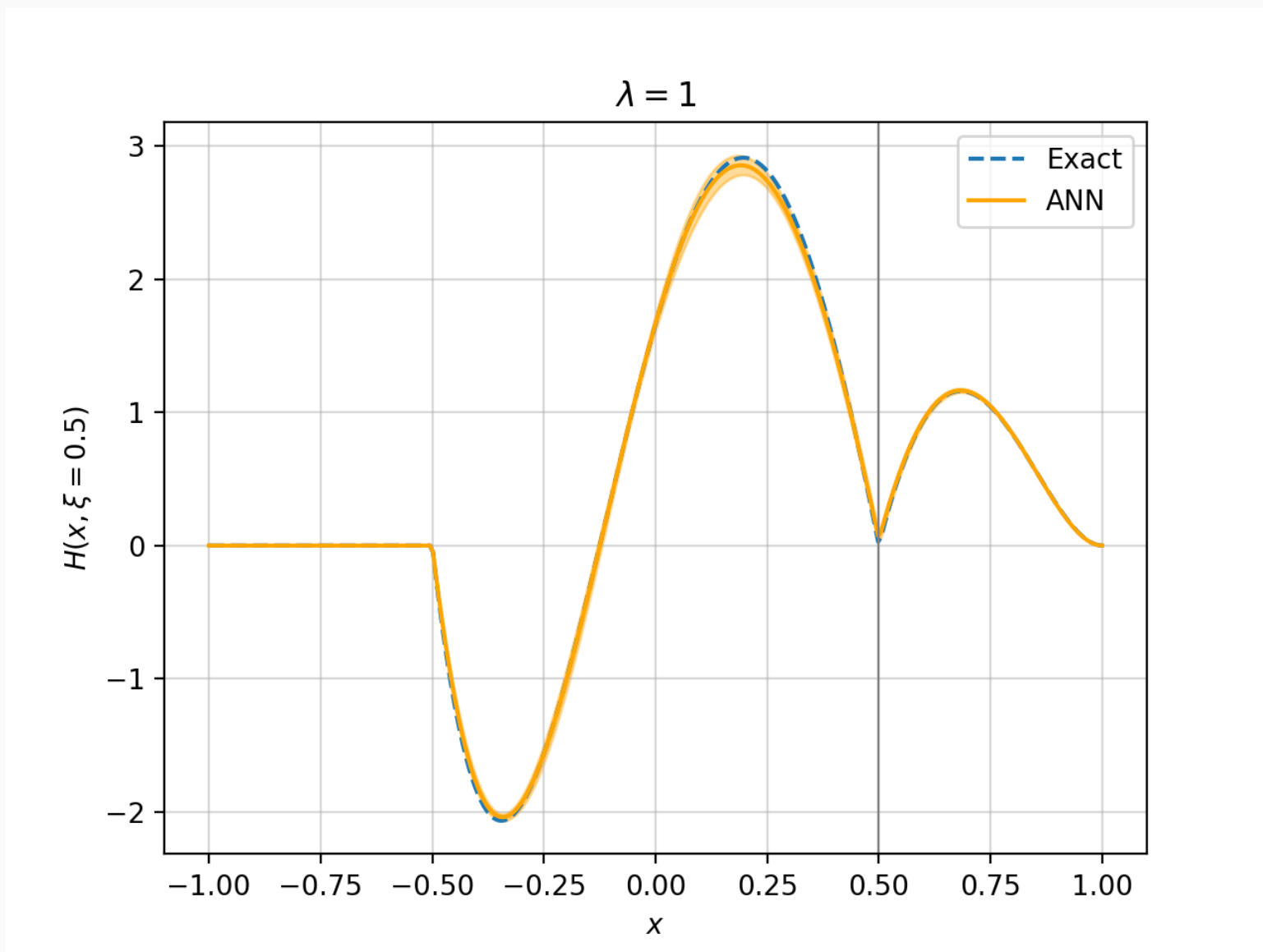
$$h(\beta, \alpha) = \frac{15}{2} (1 - 3(\alpha^2 - \beta^2) - 2\beta)$$

$$N_{neurons} = 10^2, \quad N_{sample} = 10^4$$

$$\xi \in [0, \lambda x]$$



$$\xi \in [0, \lambda x]$$



Goloskokov–Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha),$$

$$h_{GK}(\beta, \alpha) = \frac{\Gamma(n+2)}{2^{n+1}\Gamma^2(n+1)} \frac{[(1-\beta)^2 - \alpha^2]^n}{(1-\beta)^{2n+1}}$$

Goloskokov–Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha),$$

$$h_{GK}(\beta, \alpha) = \frac{\Gamma(n+2)}{2^{n+1}\Gamma^2(n+1)} \frac{[(1-\beta)^2 - \alpha^2]^n}{(1-\beta)^{2n+1}}$$

Valence distribution $q_{val}^u(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^2 c_j \beta^{j/2}, \quad n=1, \delta=0.48$

Goloskokov–Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha),$$

$$h_{GK}(\beta, \alpha) = \frac{\Gamma(n+2)}{2^{n+1}\Gamma^2(n+1)} \frac{[(1-\beta)^2 - \alpha^2]^n}{(1-\beta)^{2n+1}}$$

Valence distribution $q_{val}^u(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^2 c_j \beta^{j/2}, \quad n=1, \quad \delta=0.48$

Sea distribution $q_{sea}(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^3 c_j \beta^{j/2}, \quad n=2, \quad \delta=1.1$

Goloskokov–Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha), \quad h_{GK}(\beta, \alpha) = \frac{\Gamma(n+2)}{2^{n+1}\Gamma^2(n+1)} \frac{[(1-\beta)^2 - \alpha^2]^n}{(1-\beta)^{2n+1}}$$

Valence distribution $q_{val}^u(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^2 c_j \beta^{j/2}, \quad n=1, \delta=0.48$

Sea distribution $q_{sea}(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^3 c_j \beta^{j/2}, \quad n=2, \delta=1.1$

$$H_{sea}(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) \text{sign}(\beta) q_{sea}(|\beta|) h_{GK}(|\beta|, \alpha) \longrightarrow H_{sea}(-x, \xi) = -H_{sea}(x, \xi)$$

Goloskokov–Kroll model

S.V. Goloskokov, P. Kroll, Eur.Phys.J.C. 50 (2007)

$$H(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) q(\beta) h_{GK}(\beta, \alpha), \quad h_{GK}(\beta, \alpha) = \frac{\Gamma(n+2)}{2^{n+1}\Gamma^2(n+1)} \frac{[(1-\beta)^2 - \alpha^2]^n}{(1-\beta)^{2n+1}}$$

Valence distribution $q_{val}^u(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^2 c_j \beta^{j/2}, \quad n=1, \delta=0.48$

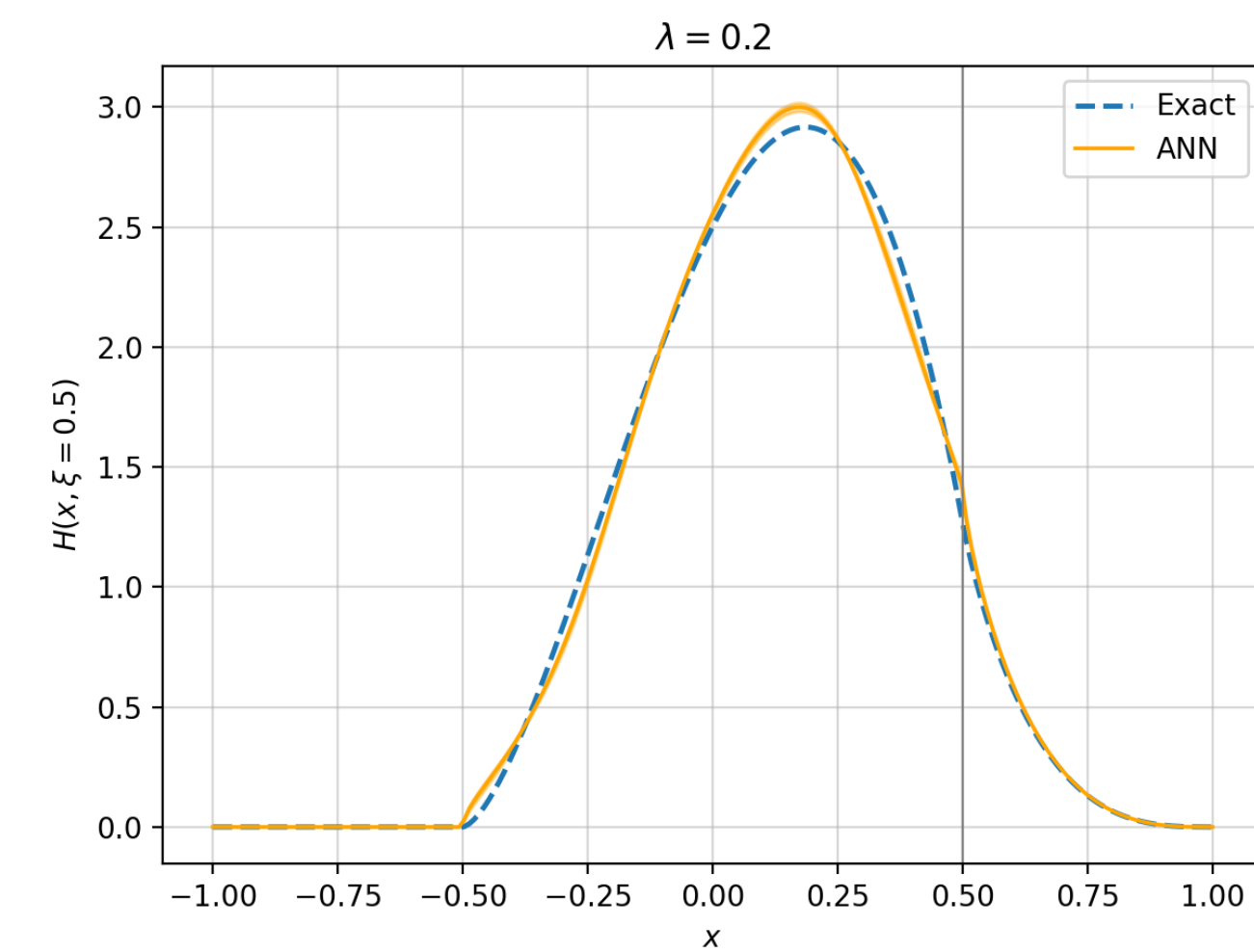
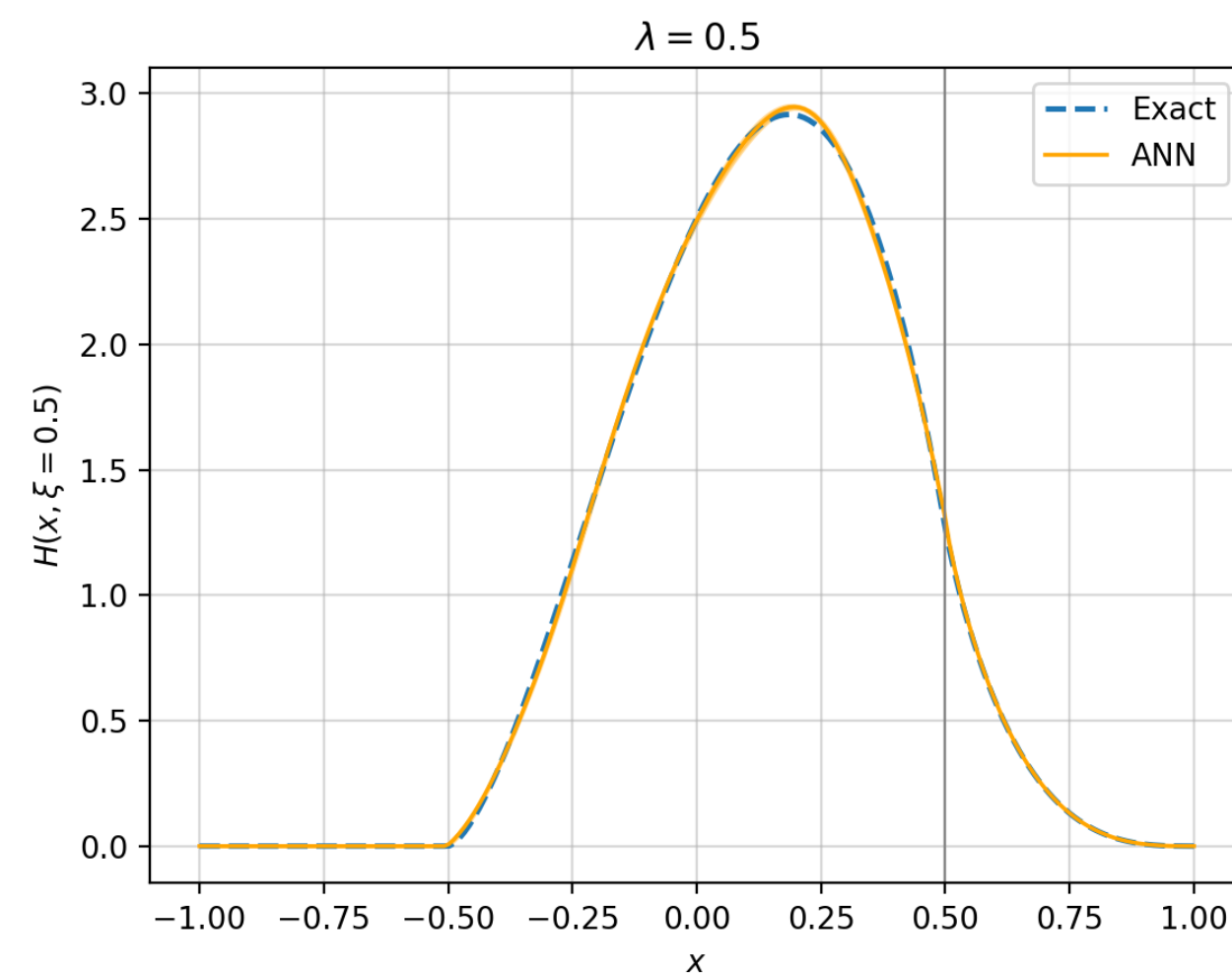
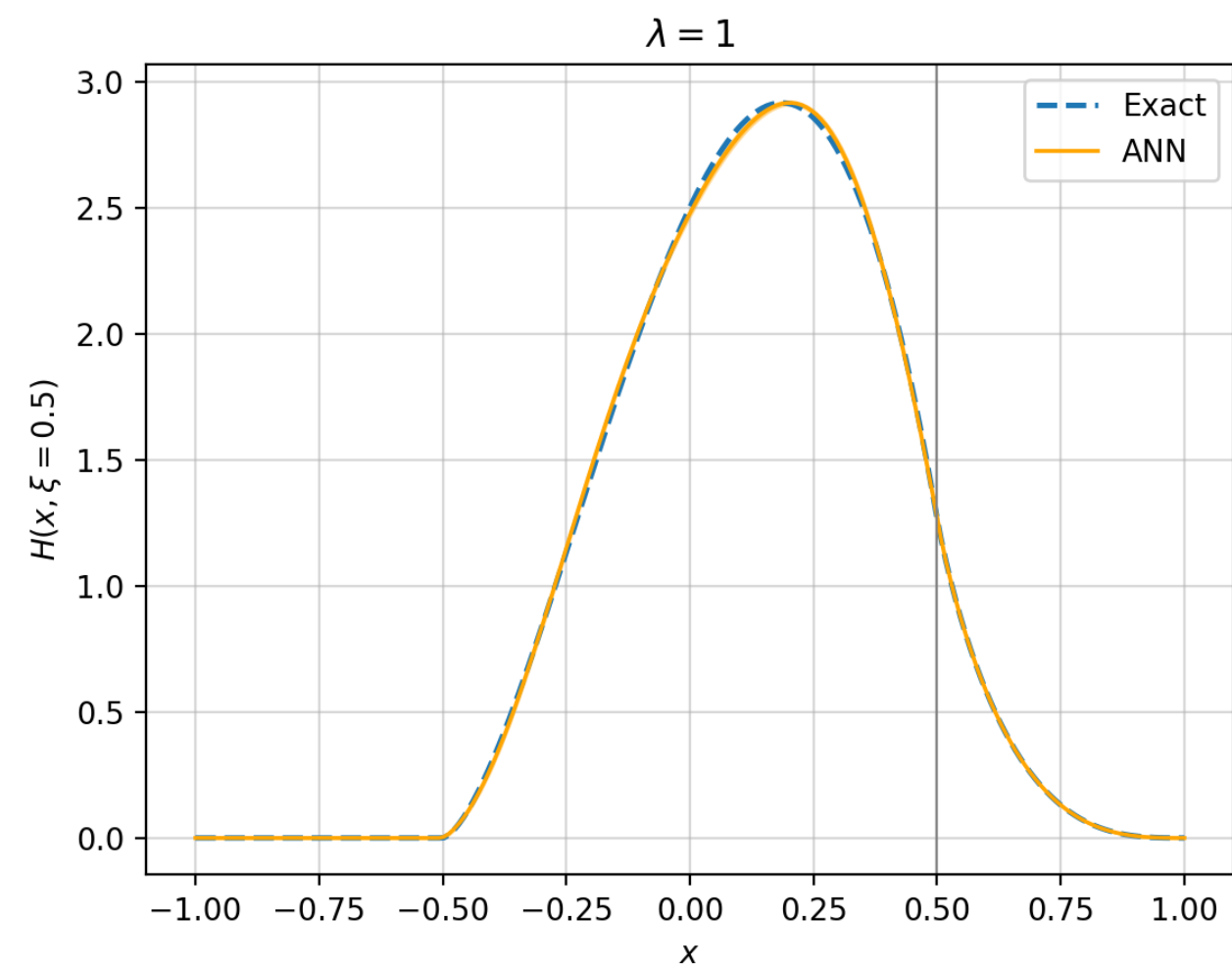
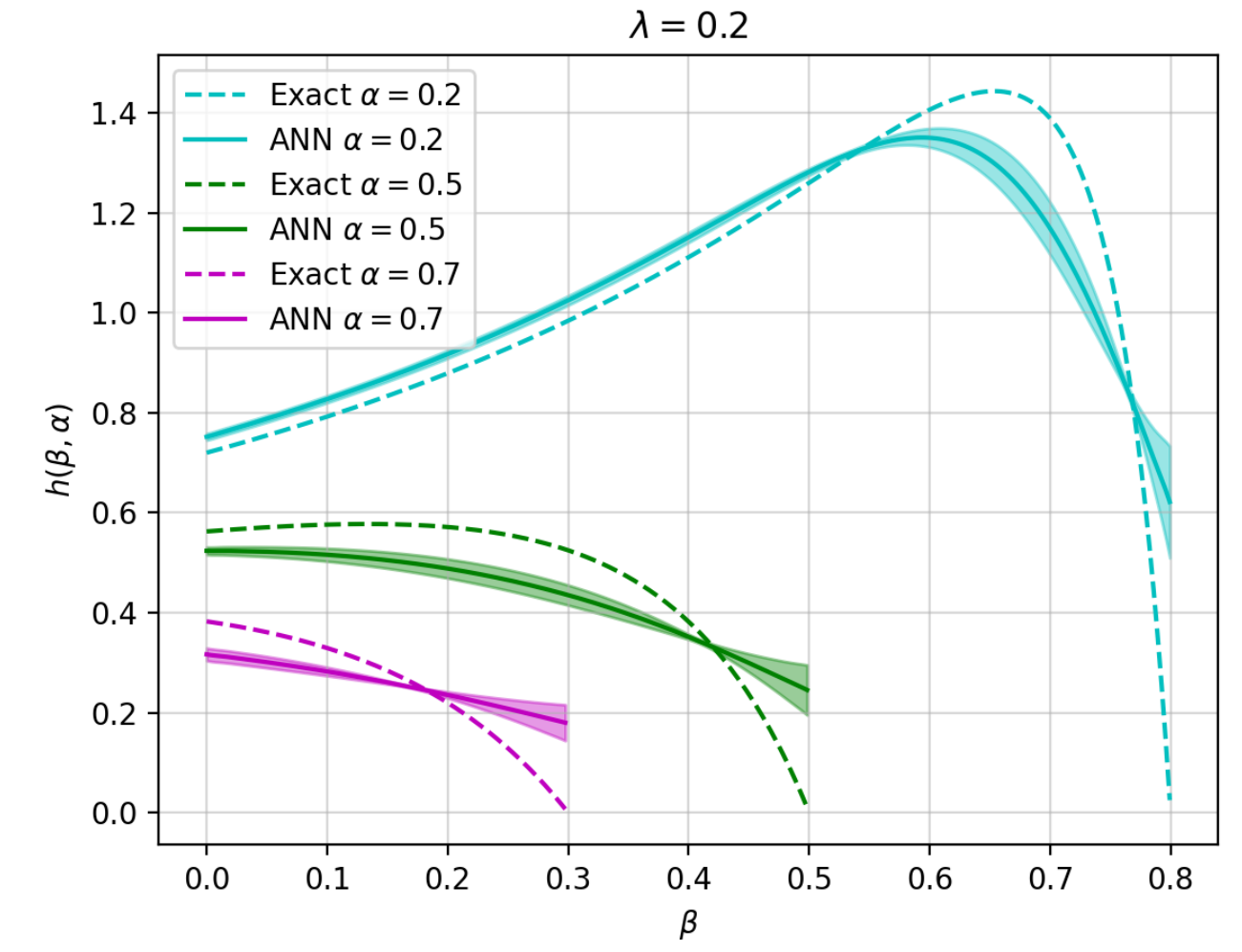
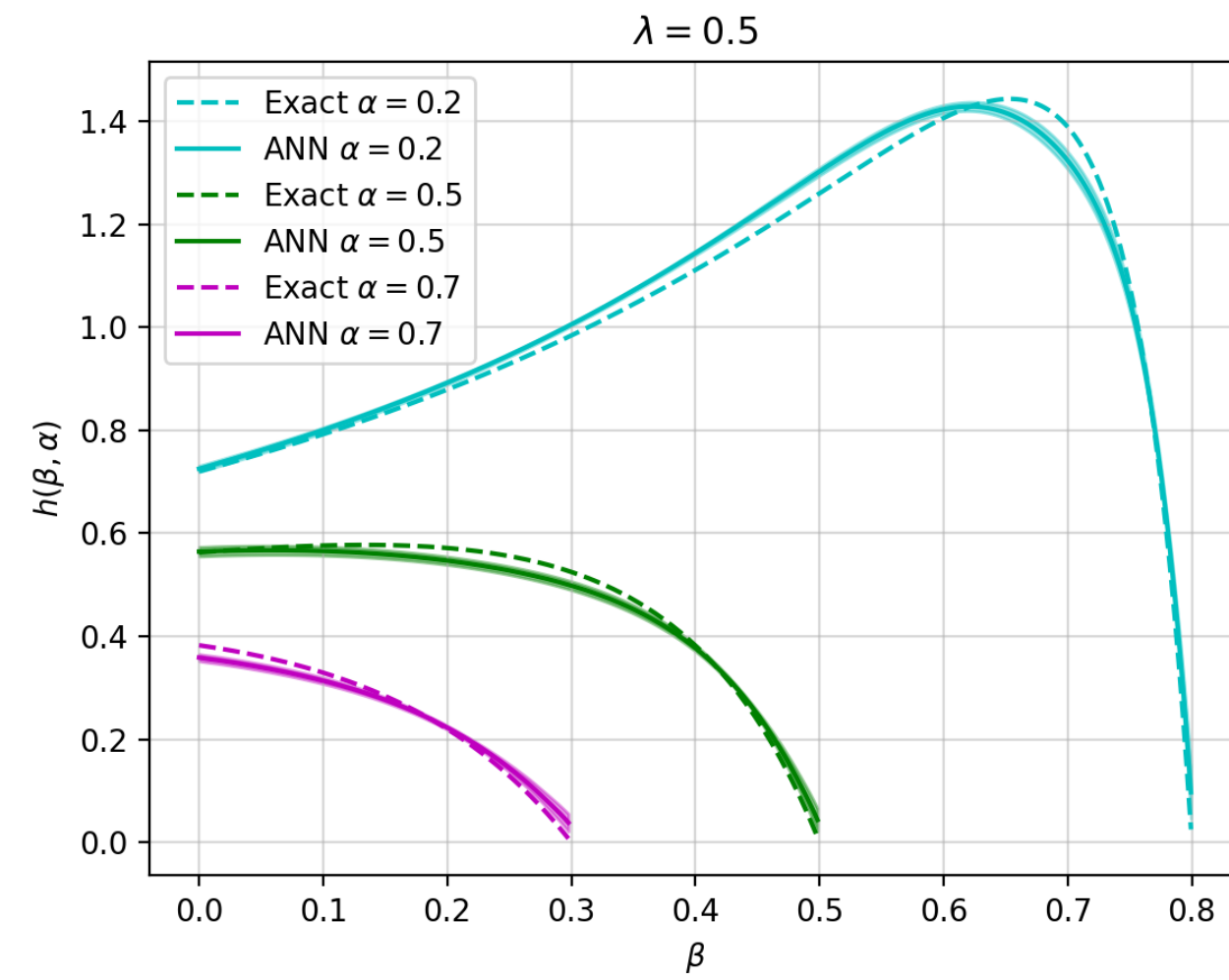
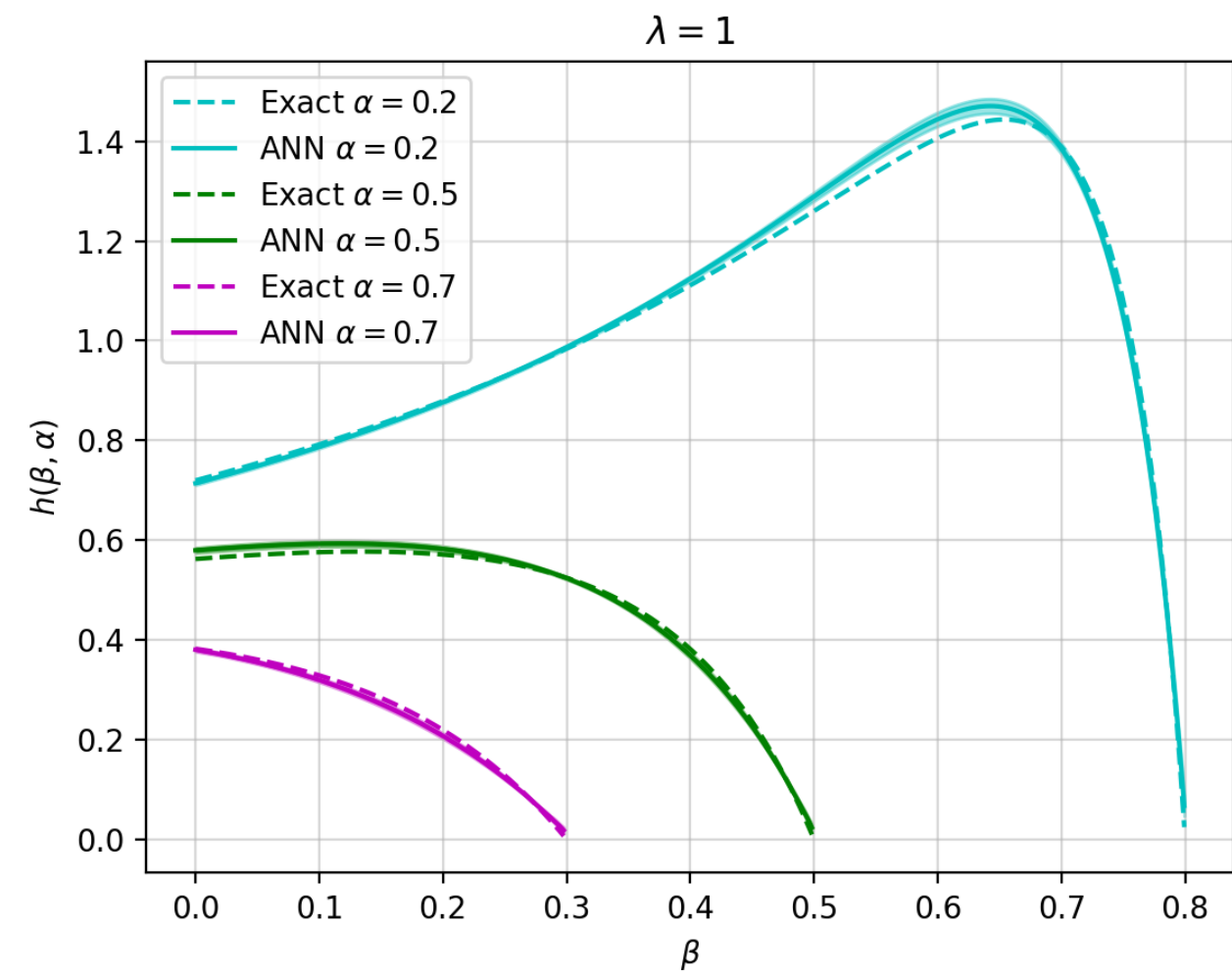
Sea distribution $q_{sea}(\beta) = \beta^{-\delta}(1-\beta)^{2n+1} \sum_{j=0}^3 c_j \beta^{j/2}, \quad n=2, \delta=1.1$

$$H_{sea}(x, \xi) = \int_{\Omega} d\beta d\alpha \delta(x - \beta - \alpha\xi) \text{sign}(\beta) q_{sea}(|\beta|) h_{GK}(|\beta|, \alpha) \longrightarrow H_{sea}(-x, \xi) = -H_{sea}(x, \xi)$$

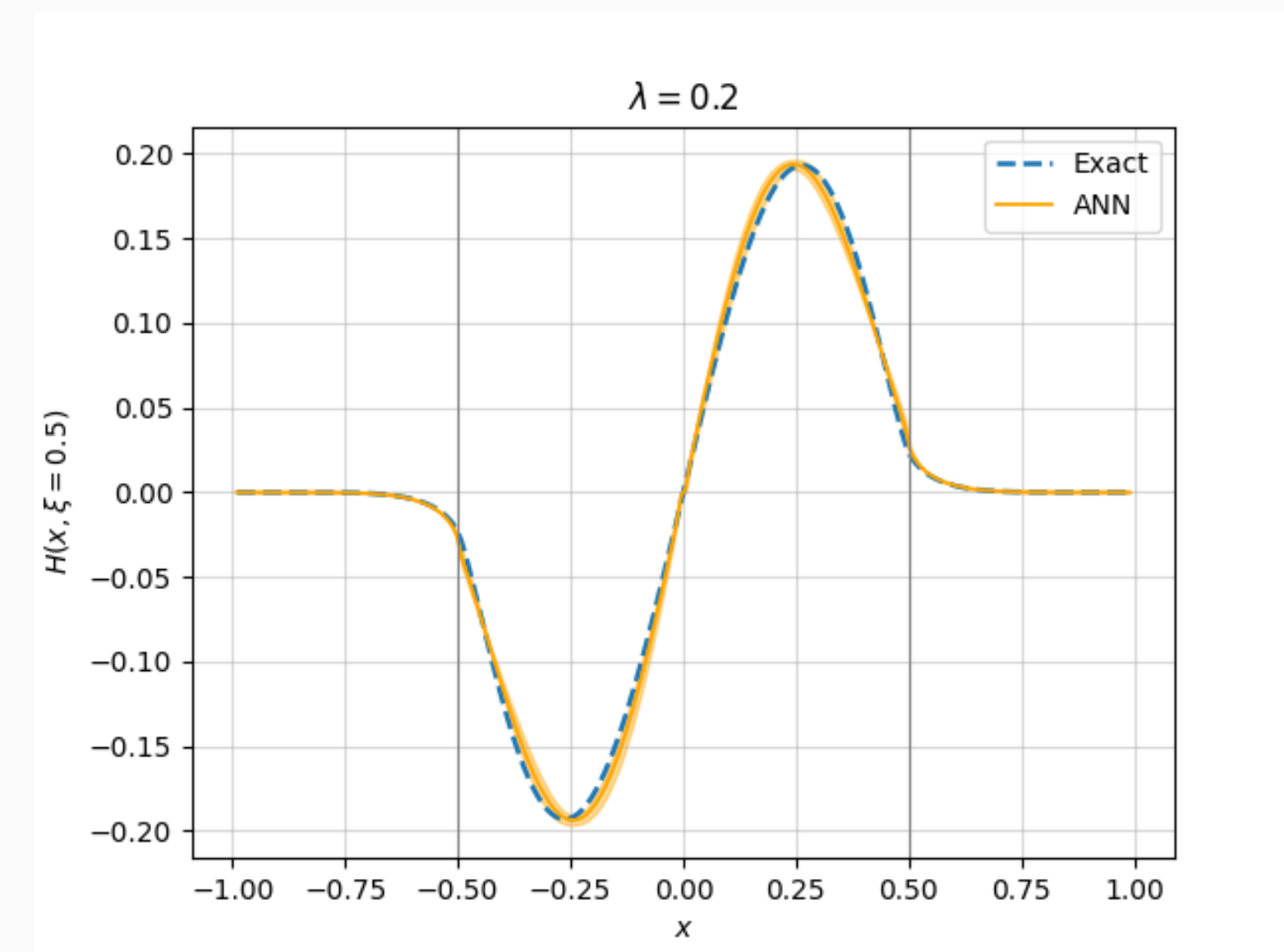
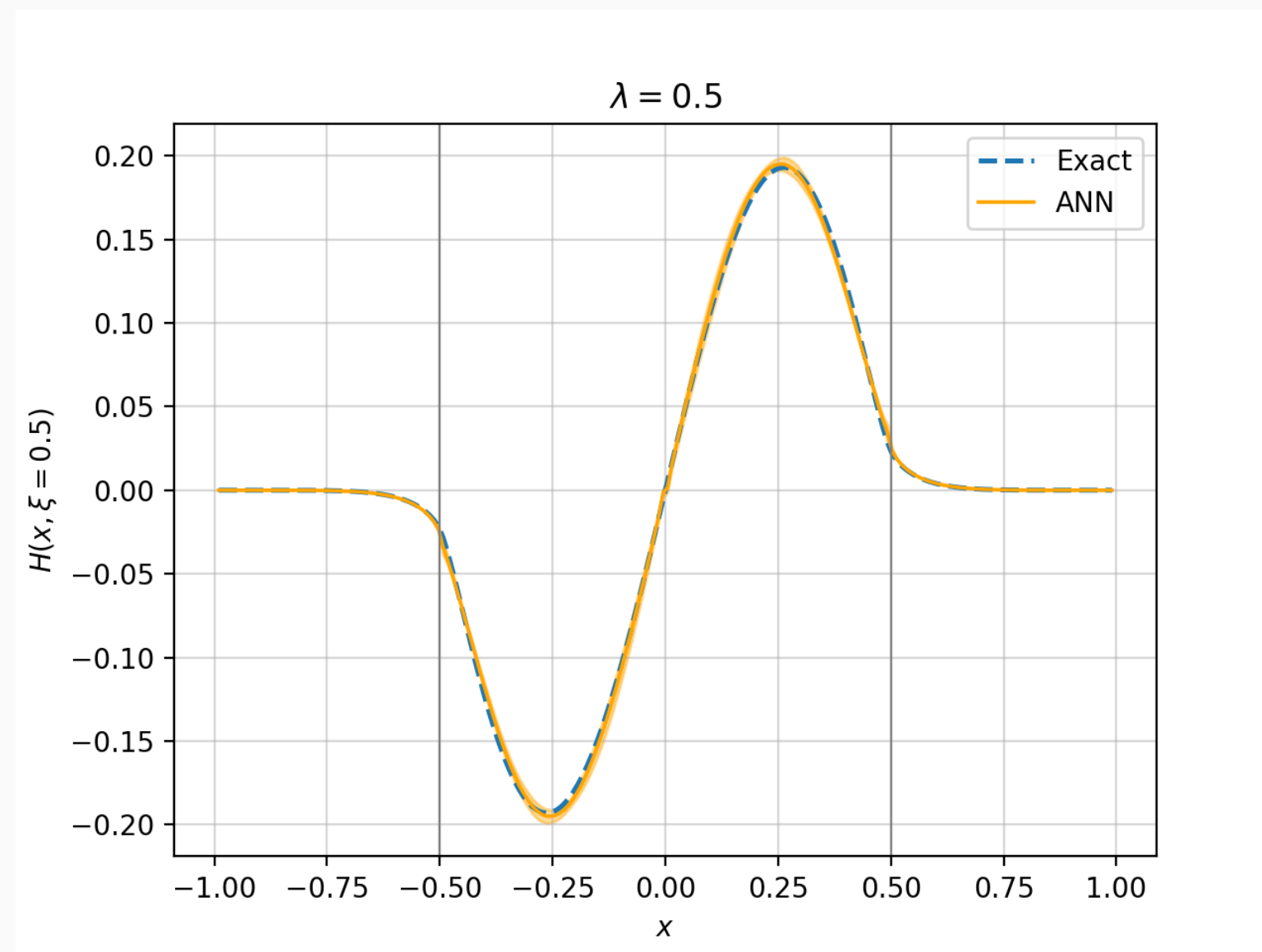
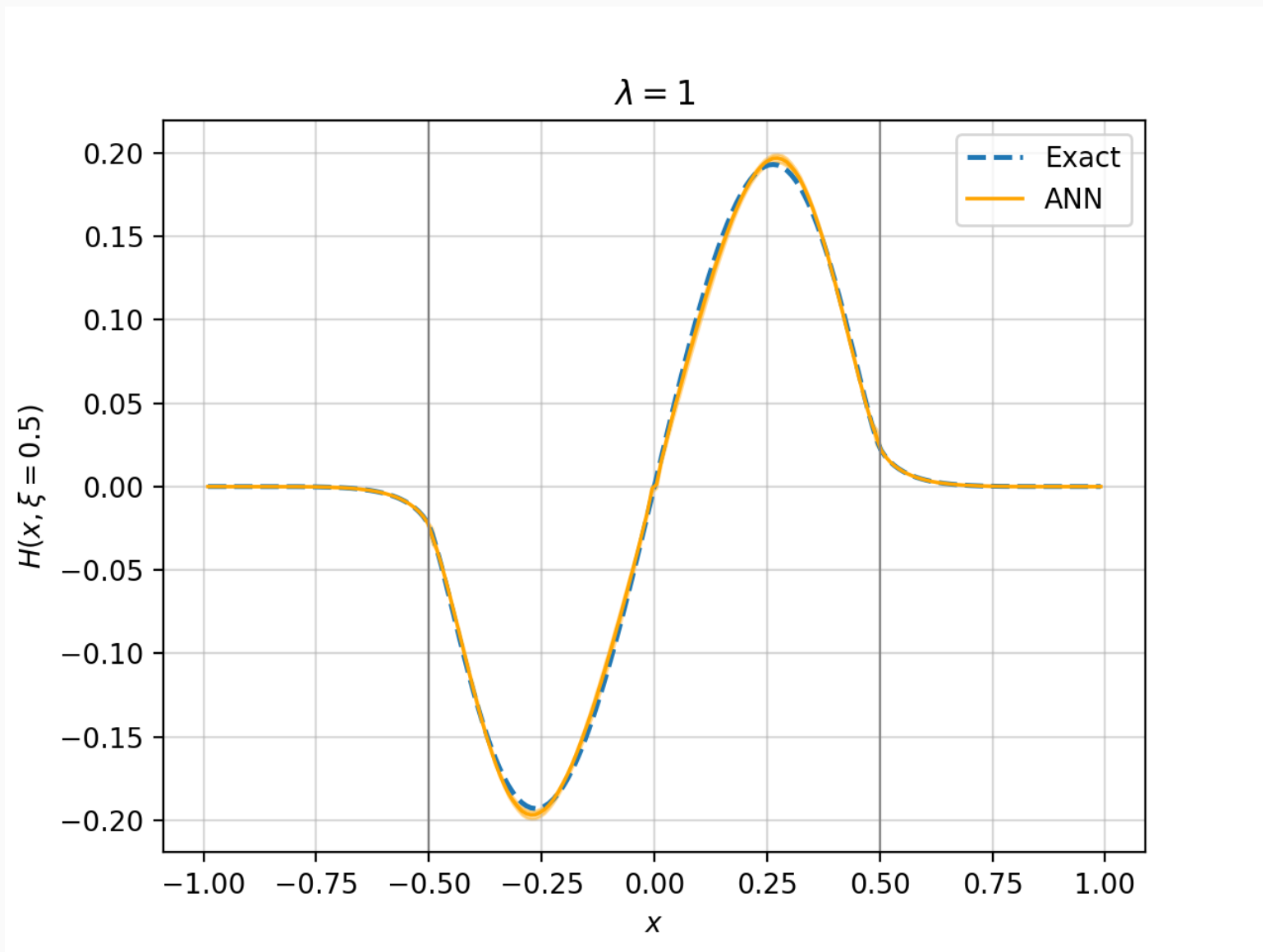
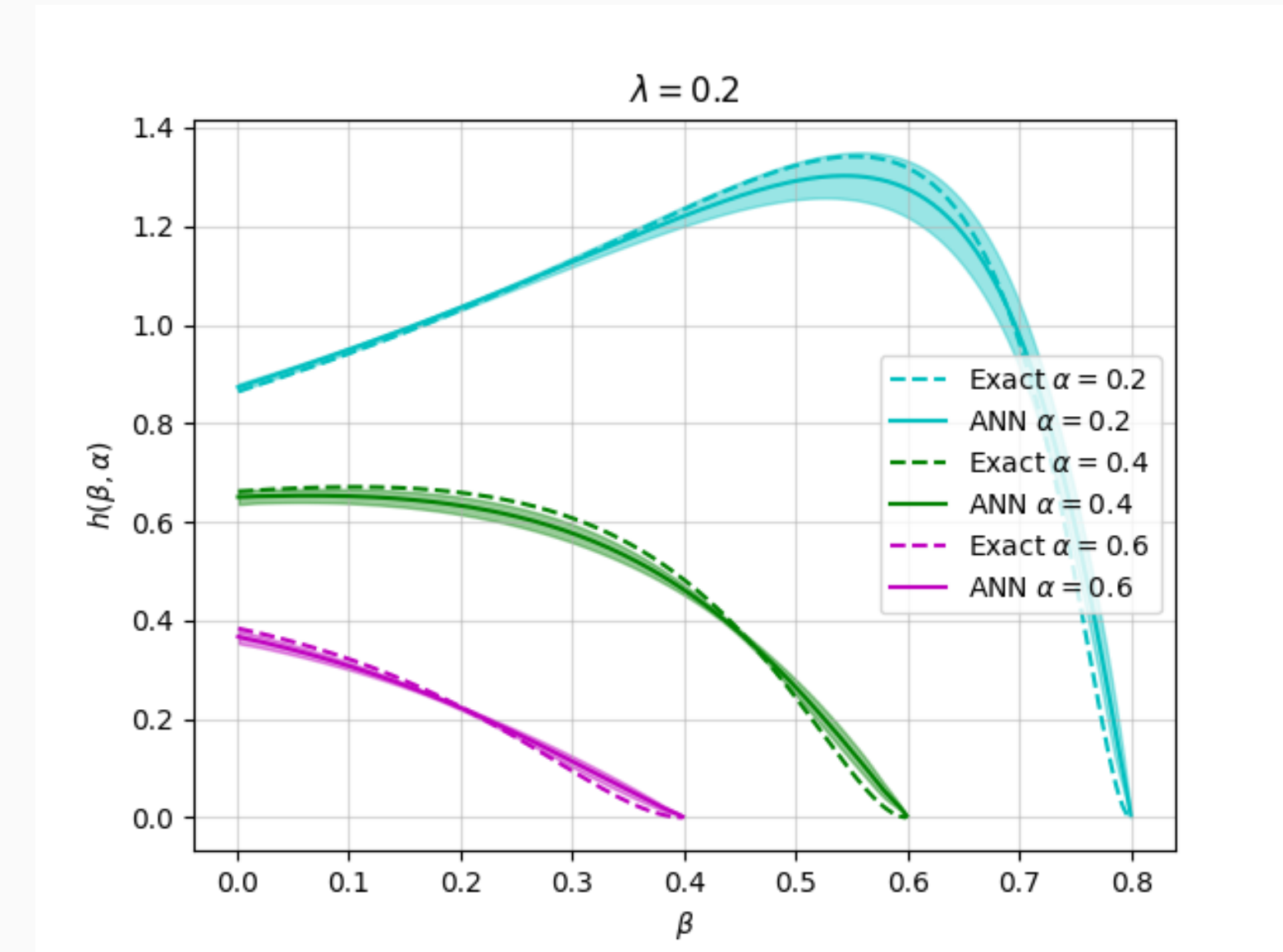
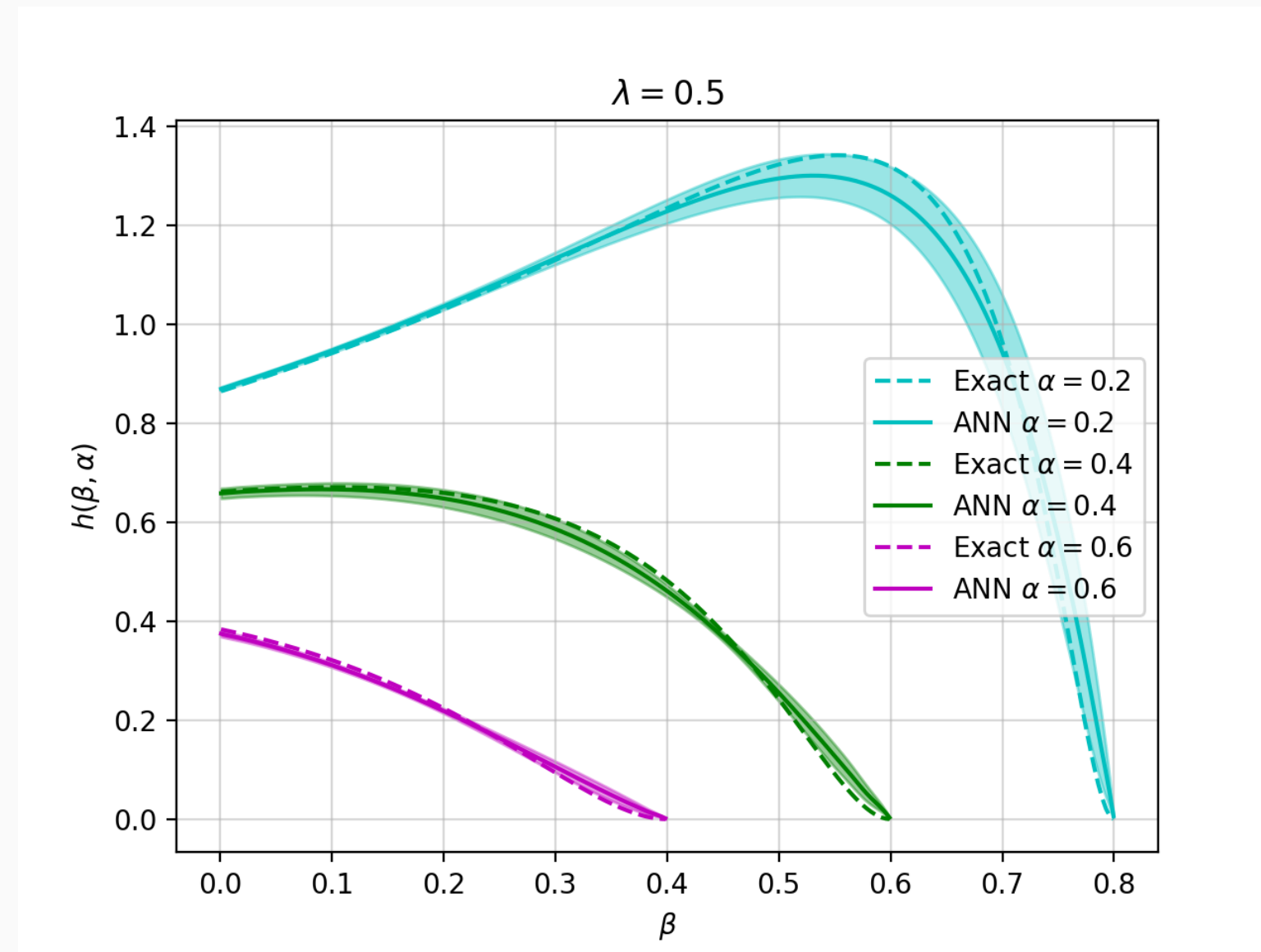
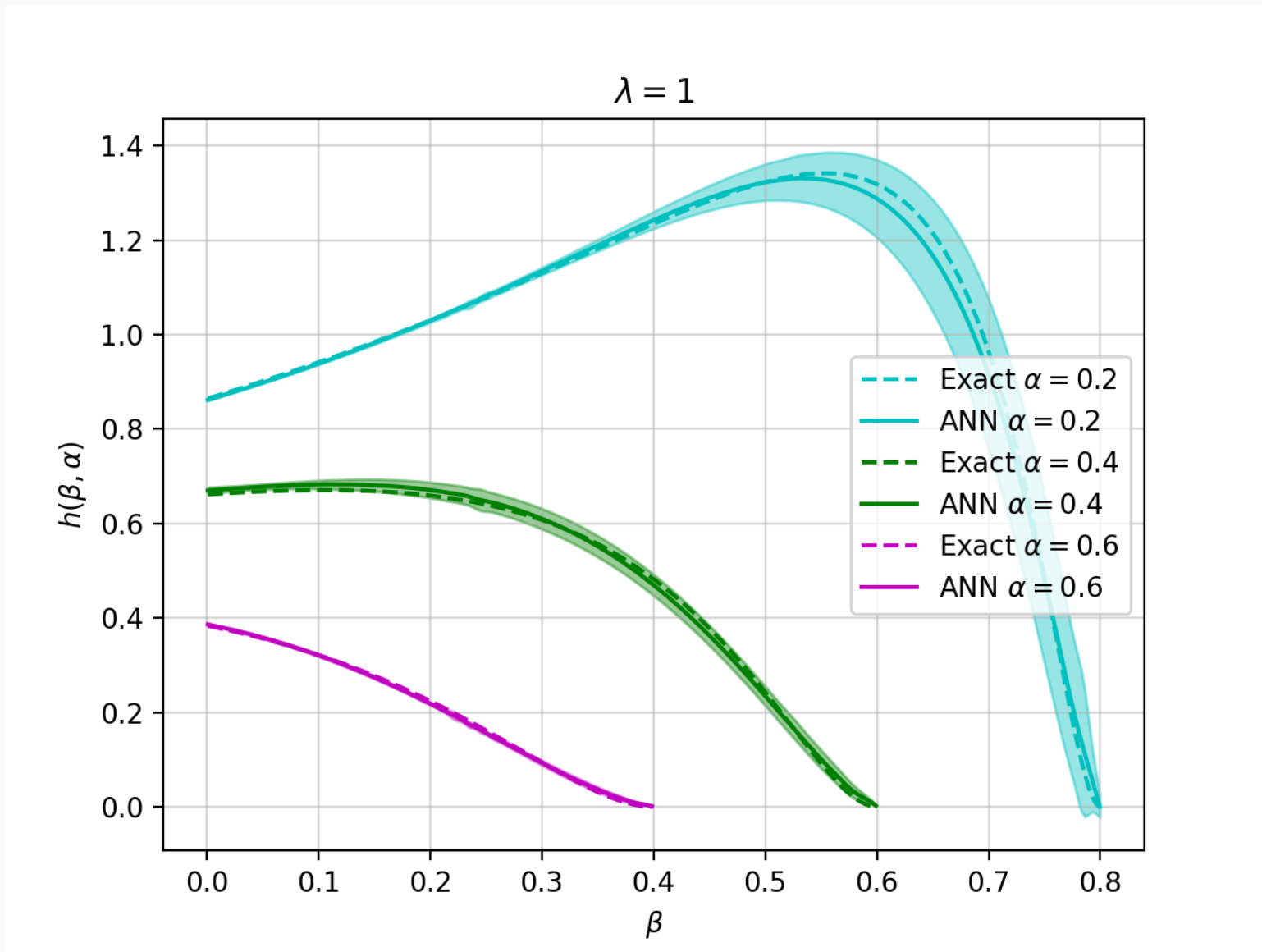
$$N_{neurons} = 25, \quad N_{sample} = 5 \times 10^3$$

$$h_{GK}(\beta, \alpha) \simeq \frac{h_{ANN}(\beta, \alpha)}{\int_{-1+|\beta|}^{1-|\beta|} d\alpha h_{ANN}(\beta, \alpha)}$$

Valence DD/GPD



Sea DD/GPD

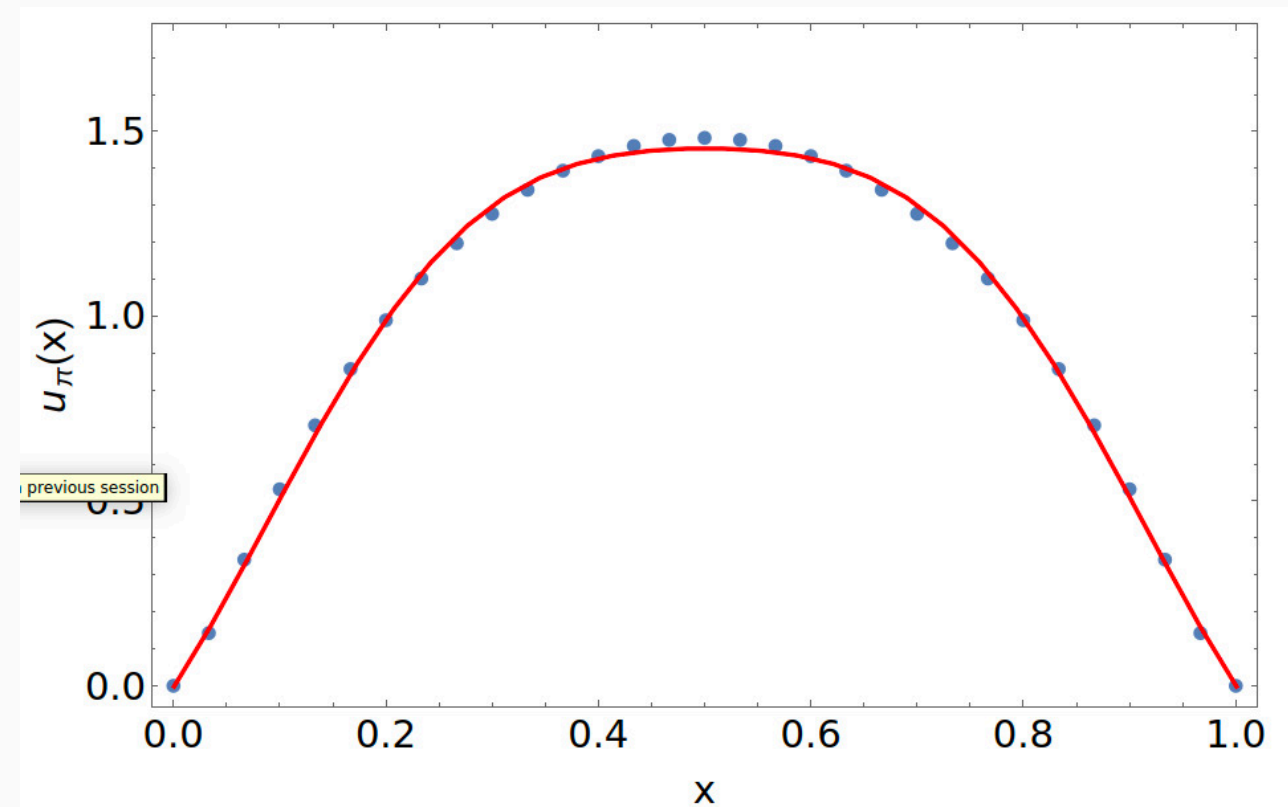


Conclusions and Outlook

- GPD can be extended from DGLAP (and proper subsets) to ERBL inverting its RT
- ANNs are a good tool for inverting RT (Unbiased parametrization)
- Testing the method on experimental results (EIC)

- Applying **ANNs** to **PDF** reconstruction from a sequence of **Mellin** moments
(with K. Raya and J. Quintero)

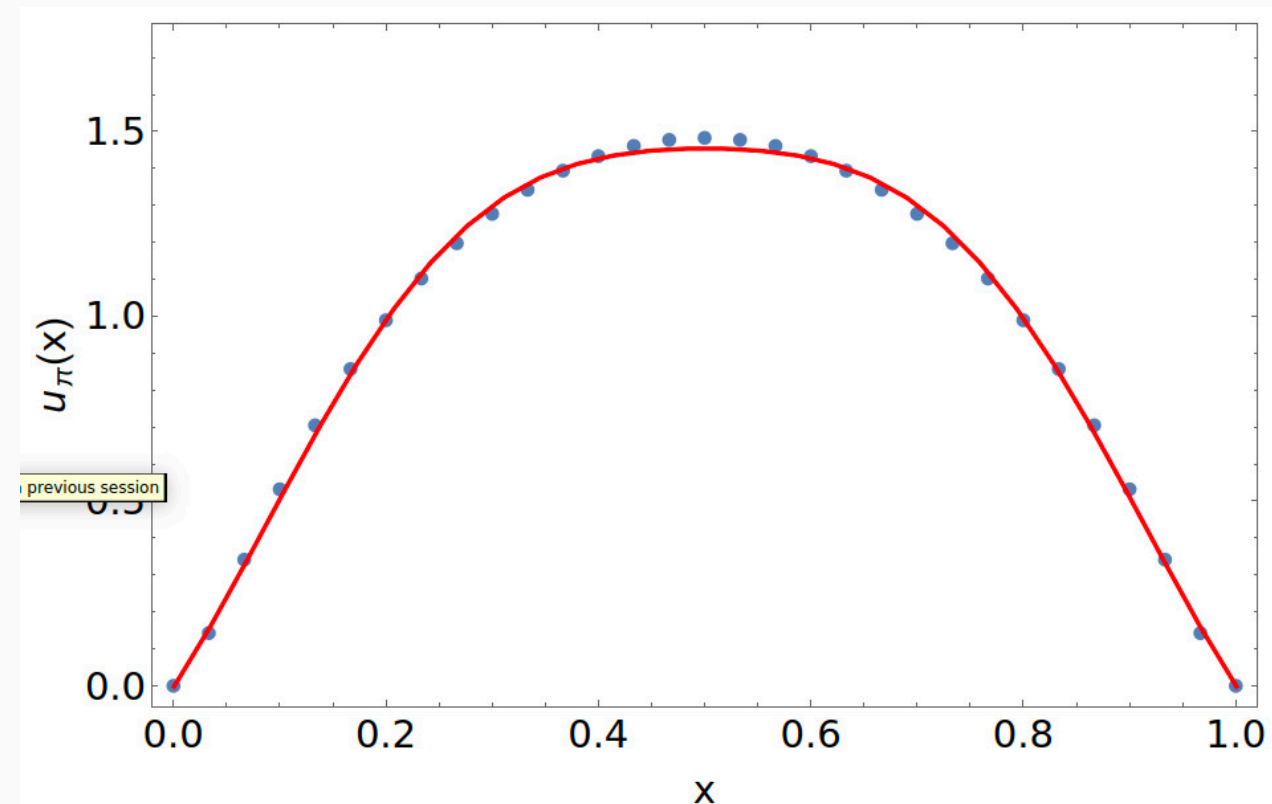
- Applying ANNs to PDF reconstruction from a sequence of Mellin moments
(with K. Raya and J. Quintero)



It works when PDF is symmetric, i.e. for the valence pion PDF (isospin symmetry)

- Applying ANNs to PDF reconstruction from a sequence of Mellin moments

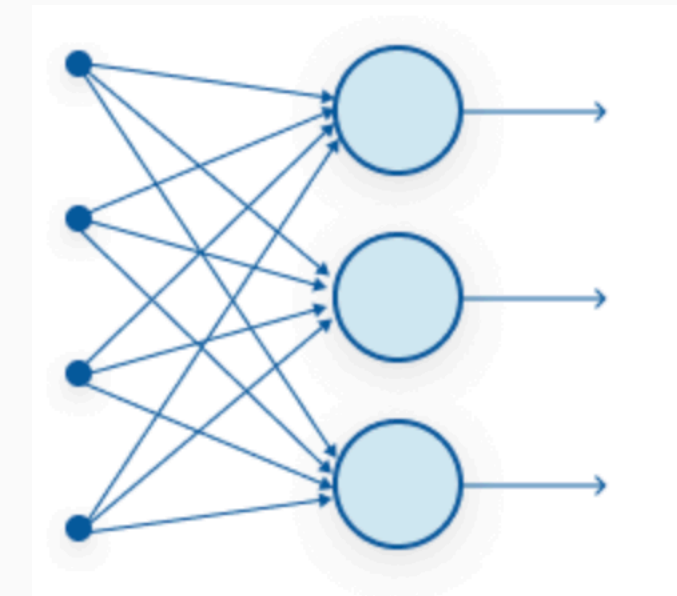
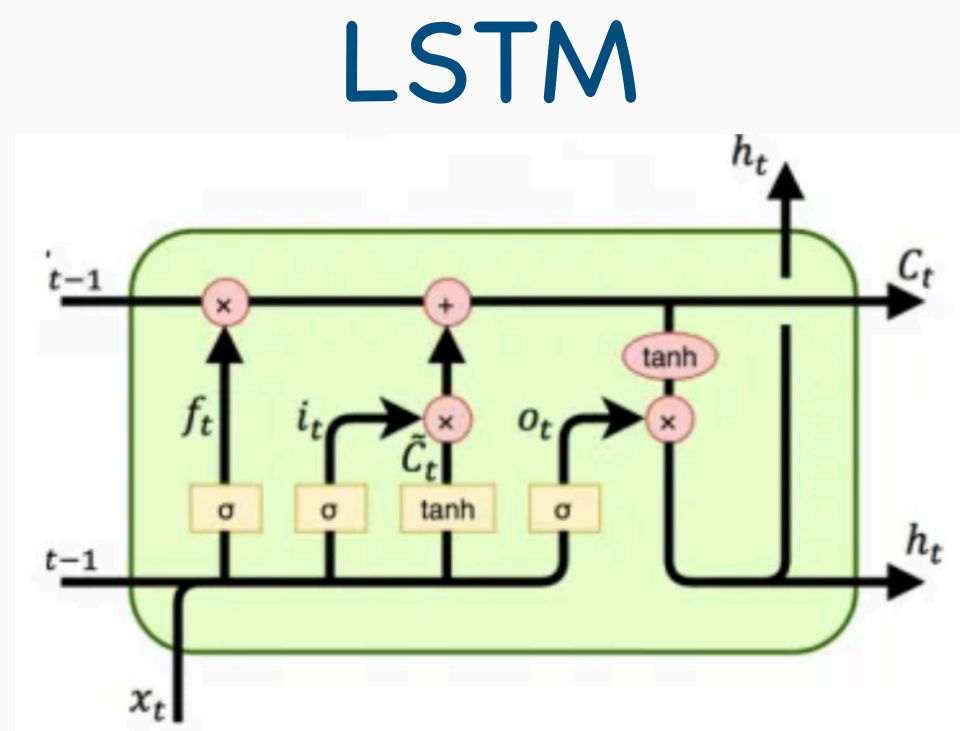
(with K. Raya and J. Quintero)



It works when PDF is symmetric, i.e. for the valence pion PDF (isospin symmetry)

Alternative

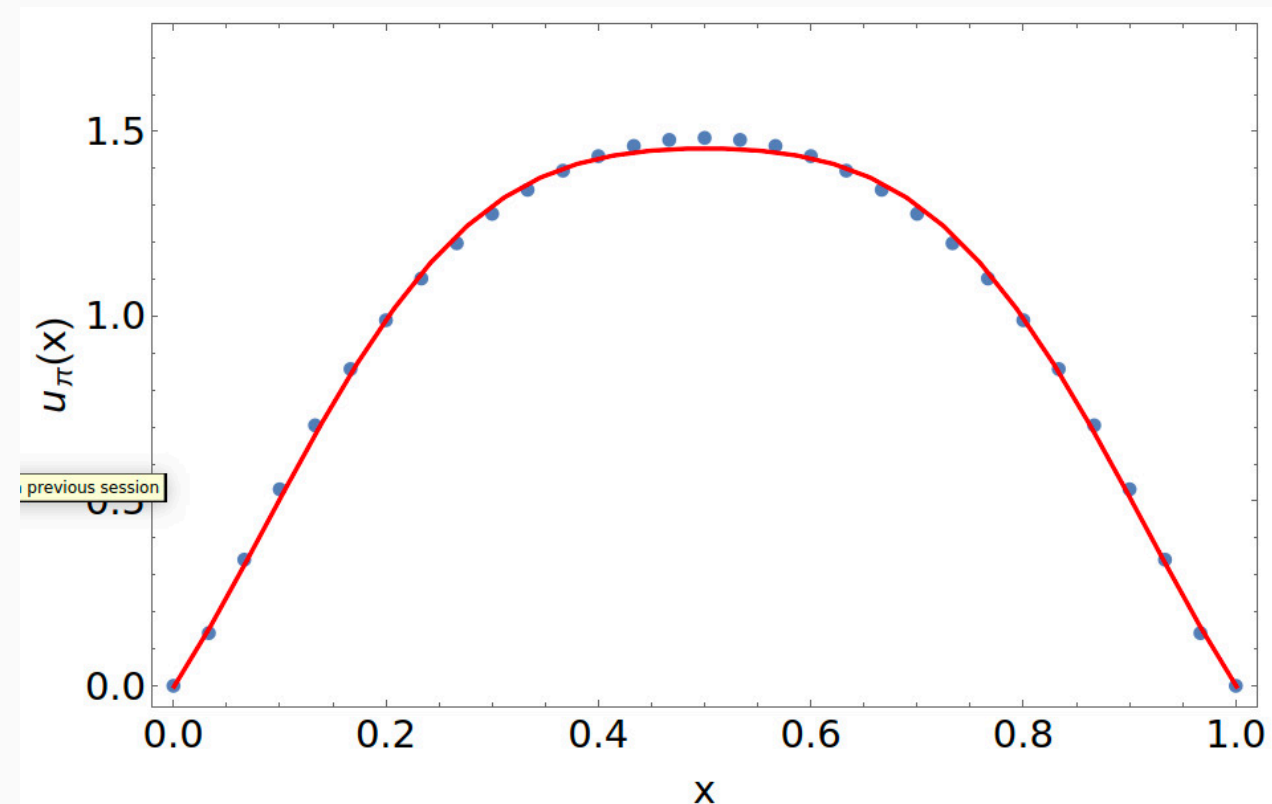
$[A_1, A_2, \dots, A_n]$



$[q(x_1), q(x_2), \dots, q(x_m)]$

- Applying ANNs to PDF reconstruction from a sequence of Mellin moments

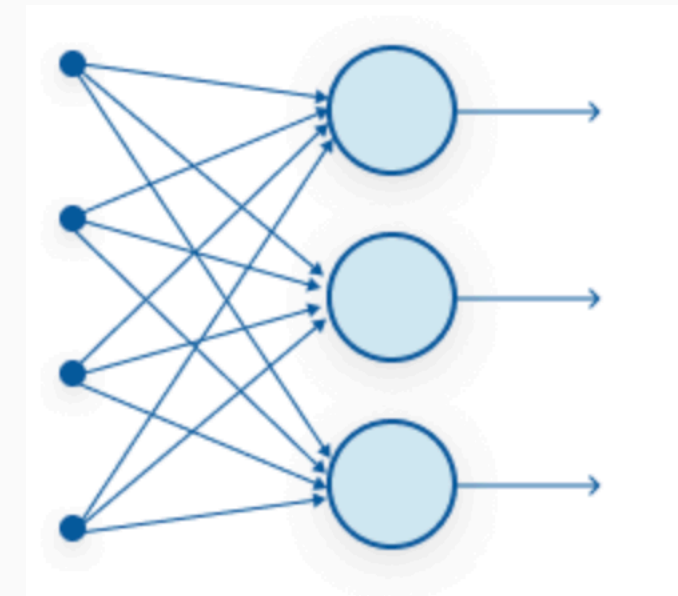
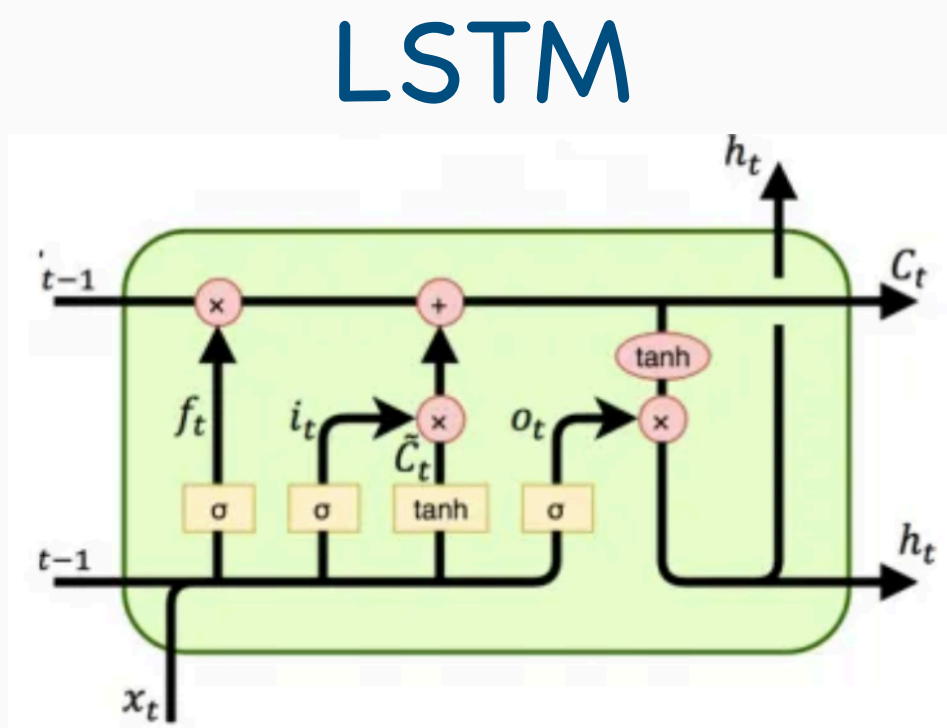
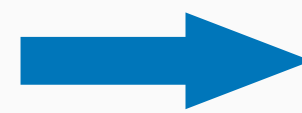
(with K. Raya and J. Quintero)



It works when PDF is symmetric, i.e. for the valence pion PDF (isospin symmetry)

Alternative

$[A_1, A_2, \dots, A_n]$



$[q(x_1), q(x_2), \dots, q(x_m)]$

Memory loss problem



Trying with Transformer