

Transversity PDFs and GPDs from lattice QCD

Martha Constantinou



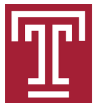
Temple University

CFNS Workshop

**TMDs: Towards a Synergy between
Lattice QCD & Global Analysis**

June 22, 2023

Topics



Transversity parton distribution functions from lattice QCD

PHYSICAL REVIEW D **99**, 114504 (2019)

Systematic uncertainties in parton distribution functions from lattice QCD simulations at the physical point

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou,⁴ Kyriakos Hadjiyiannakou,² Karl Jansen,⁵ Aurora Scapellato,^{1,6} and Fernanda Steffens⁷

★ **Twist-2 PDFs & GPDs:
“traditional calculations**

PHYSICAL REVIEW D **105**, 034501 (2022)

Transversity GPDs of the proton from lattice QCD

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou[Ⓧ],⁴ Kyriakos Hadjiyiannakou,^{1,2} Karl Jansen,⁵ Aurora Scapellato,⁴ and Fernanda Steffens⁶

Transversity parton distribution functions from lattice QCD

PHYSICAL REVIEW D **99**, 114504 (2019)

Systematic uncertainties in parton distribution functions from lattice QCD simulations at the physical point

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou,⁴ Kyriakos Hadjiyiannakou,² Karl Jansen,⁵ Aurora Scapellato,^{1,6} and Fernanda Steffens⁷

★ **Twist-2 PDFs & GPDs:
“traditional calculations**

PHYSICAL REVIEW D **105**, 034501 (2022)

Transversity GPDs of the proton from lattice QCD

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou[Ⓜ],⁴ Kyriakos Hadjiyiannakou,^{1,2} Karl Jansen,⁵ Aurora Scapellato,⁴ and Fernanda Steffens⁶

★ **Twist-3 distribution functions**

PHYSICAL REVIEW D **104**, 114510 (2021)

Parton distribution functions beyond leading twist from lattice QCD: The $h_L(x)$ case

Shohini Bhattacharya[Ⓜ],¹ Krzysztof Cichy[Ⓜ],² Martha Constantinou[Ⓜ],¹ Andreas Metz,¹ Aurora Scapellato[Ⓜ],¹ and Fernanda Steffens³

Transversity parton distribution functions from lattice QCD

PHYSICAL REVIEW D **99**, 114504 (2019)

Systematic uncertainties in parton distribution functions from lattice QCD simulations at the physical point

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou,⁴ Kyriakos Hadjiyiannakou,² Karl Jansen,⁵ Aurora Scapellato,^{1,6} and Fernanda Steffens⁷

★ **Twist-2 PDFs & GPDs: “traditional calculations**

PHYSICAL REVIEW D **105**, 034501 (2022)

Transversity GPDs of the proton from lattice QCD

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou[Ⓜ],⁴ Kyriakos Hadjiyiannakou,^{1,2} Karl Jansen,⁵ Aurora Scapellato,⁴ and Fernanda Steffens⁶

★ **Twist-3 distribution functions**

PHYSICAL REVIEW D **104**, 114510 (2021)

Parton distribution functions beyond leading twist from lattice QCD: The $h_L(x)$ case

Shohini Bhattacharya[Ⓜ],¹ Krzysztof Cichy[Ⓜ],² Martha Constantinou[Ⓜ],¹ Andreas Metz,¹ Aurora Scapellato[Ⓜ],¹ and Fernanda Steffens³

PHYSICAL REVIEW D **106**, 114512 (2022)

Generalized parton distributions from lattice QCD with asymmetric momentum transfer: Unpolarized quarks

+ Joshua Miller

Shohini Bhattacharya[Ⓜ],^{1,*} Krzysztof Cichy,² Martha Constantinou[Ⓜ],^{3,†} Jack Dodson,³ Xiang Gao,⁴ Andreas Metz,³ Swagato Mukherjee[Ⓜ],¹ Aurora Scapellato,³ Fernanda Steffens,³ and Yong Zhao⁴

★ **Twist-2 GPDs: new approach**

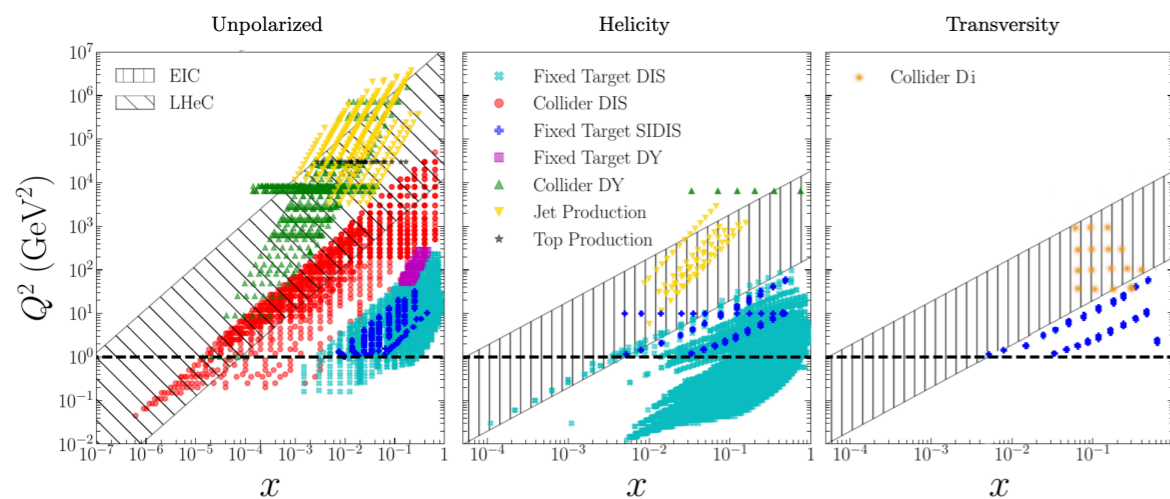


Distribution Functions

- ★ **Key universal non-perturbative tools for study of hadron structure**
- ★ **Global analyses: main source of information for distribution functions**
- ★ **Global fits improved: theoretical advances & new data**
- ★ **ambiguities due to limited kinematic regions, limited measurements, etc**

Distribution Functions

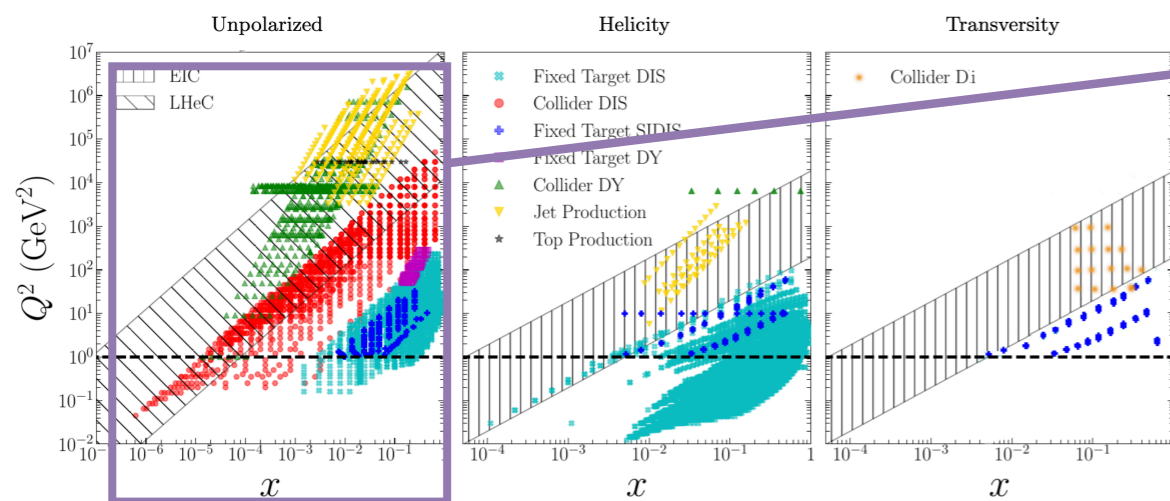
- ★ Key universal non-perturbative tools for study of hadron structure
- ★ Global analyses: main source of information for distribution functions
- ★ Global fits improved: theoretical advances & new data
- ★ ambiguities due to limited kinematic regions, limited measurements, etc



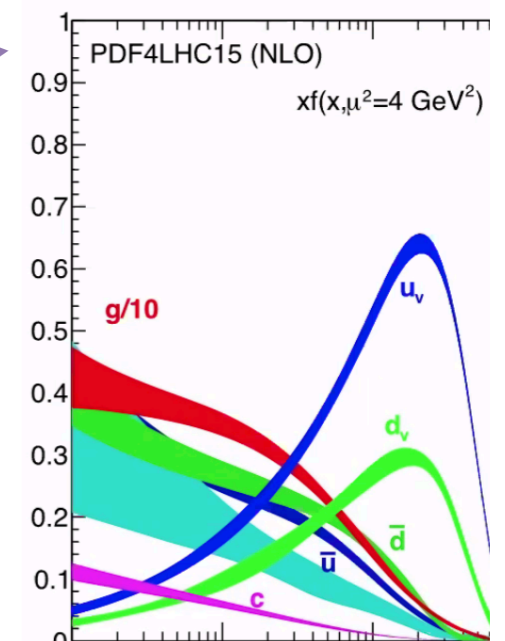
[Ethier & Nocera, *Ann. Rev. Nucl. Part. Sci.* 70 (2020) 1, arXiv:2001.07722]

Distribution Functions

- ★ Key universal non-perturbative tools for study of hadron structure
- ★ Global analyses: main source of information for distribution functions
- ★ Global fits improved: theoretical advances & new data
- ★ ambiguities due to limited kinematic regions, limited measurements, etc

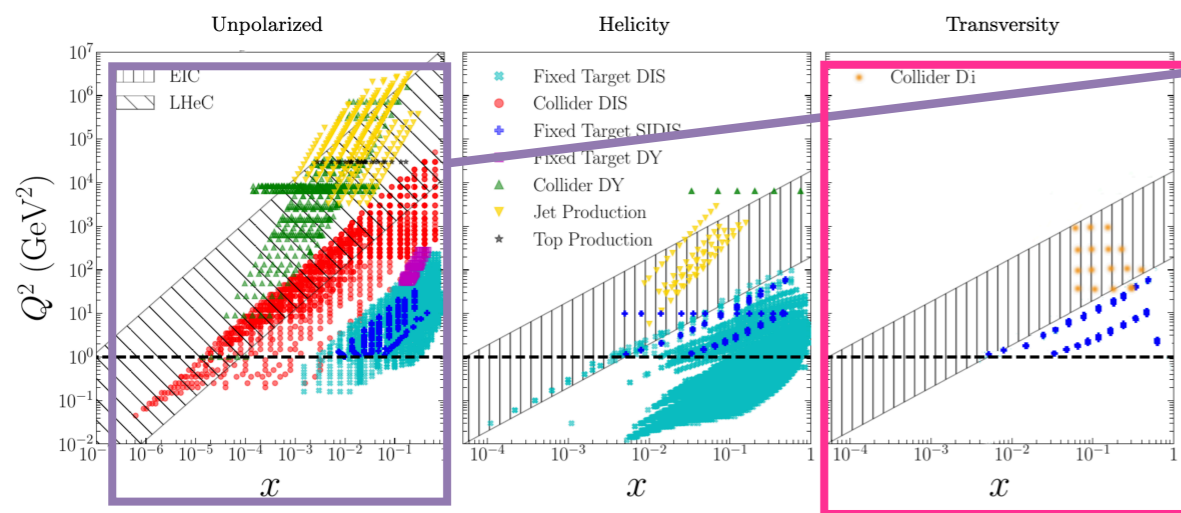


[Ethier & Nocera, *Ann. Rev. Nucl. Part. Sci.* 70 (2020) 1, arXiv:2001.07722]

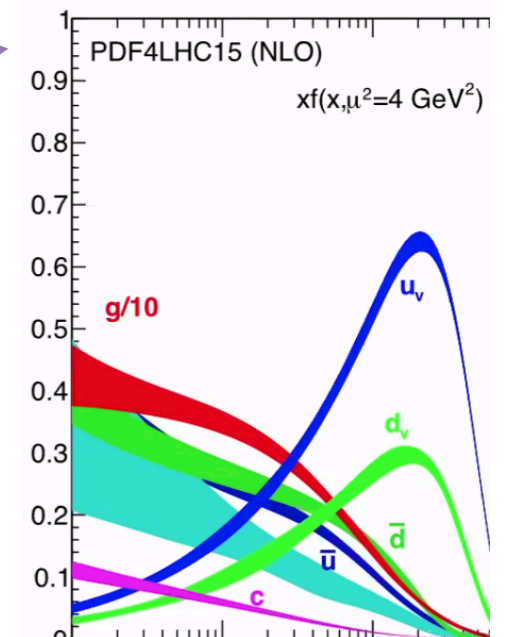
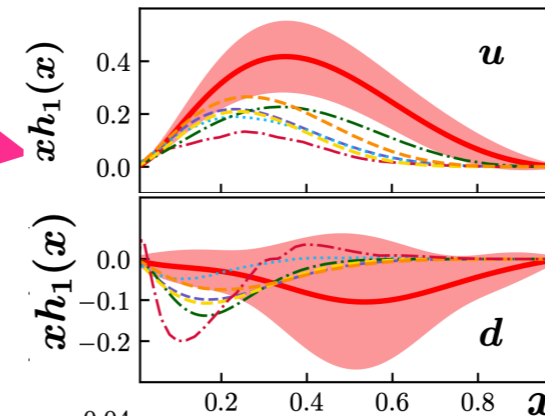


Distribution Functions

- ★ Key universal non-perturbative tools for study of hadron structure
- ★ Global analyses: main source of information for distribution functions
- ★ Global fits improved: theoretical advances & new data
- ★ ambiguities due to limited kinematic regions, limited measurements, etc

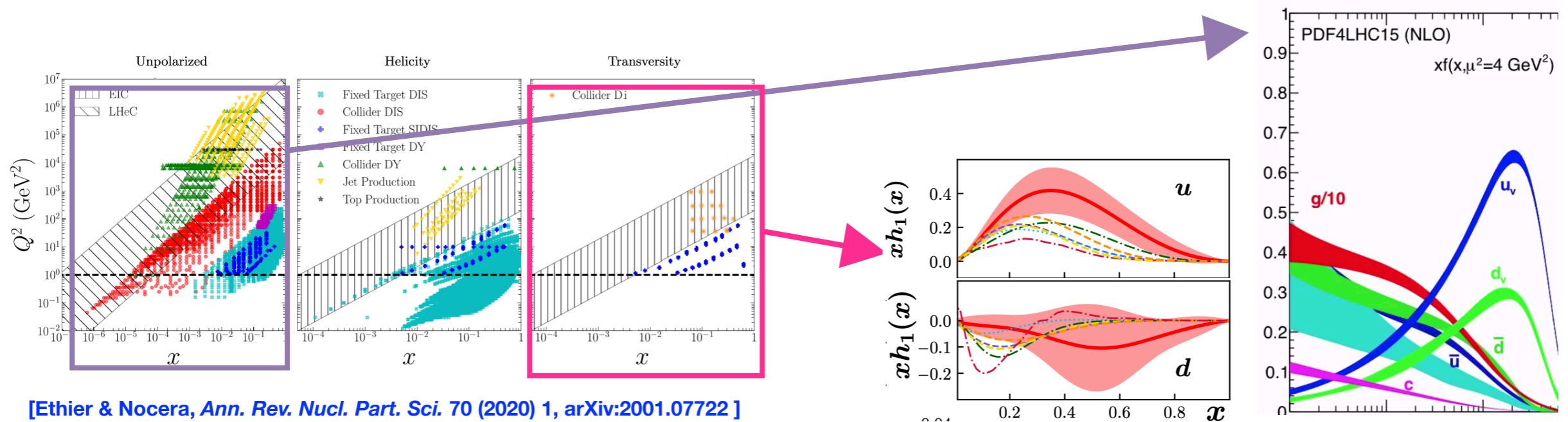


[Ethier & Nocera, *Ann. Rev. Nucl. Part. Sci.* 70 (2020) 1, arXiv:2001.07722]



Distribution Functions

- ★ Key universal non-perturbative tools for study of hadron structure
- ★ Global analyses: main source of information for distribution functions
- ★ Global fits improved: theoretical advances & new data
- ★ ambiguities due to limited kinematic regions, limited measurements, etc



Calculation from first principle (lattice QCD) can help in the estimation of PDFs

Twist-classification of PDFs, GPDs, TMDs

★ Twist: specifies the order in $1/Q$ at which the function enters factorization formula for a given observable

$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$

★ Twist-2: probabilistic densities - a wealth of information exists

★ Twist-3: poorly known, but very important:

- as sizeable as twist-2
- contain information about quark-gluon correlations inside hadrons
- appear in QCD factorization theorems for various observables (e.g. g_2)
- certain twist-3 PDFs are related to the TMDs
- physical interpretation (e.g. average force on partons inside hadron)

Twist-classification of PDFs, GPDs, TMDs

★ Twist: specifies the order in $1/Q$ at which the function enters factorization formula for a given observable $f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$

★ Twist-2: probabilistic densities - a wealth of information exists

★ Twist-3: poorly known, but very important:

- as sizeable as twist-2
- contain information about quark-gluon correlations inside hadrons
- appear in QCD factorization theorems for various observables (e.g. g_2)
- certain twist-3 PDFs are related to the TMDs
- physical interpretation (e.g. average force on partons inside hadron)

While twist-3 $f_i^{(1)}$ share some similarities with twist-2 $f_i^{(0)}$ in their extraction, there are several challenges both experimentally and theoretically

Distribution functions

Through non-local matrix elements
of fast-moving hadrons

Access of PDFs/GPDs on a Euclidean Lattice

[X. Ji, Phys. Rev. Lett. 110 (2013) 262002]

Matrix elements of nonlocal (equal-time) operators with **fast moving hadrons**

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$

$$\Delta = P_f - P_i$$

$$t = \Delta^2 = -Q^2$$

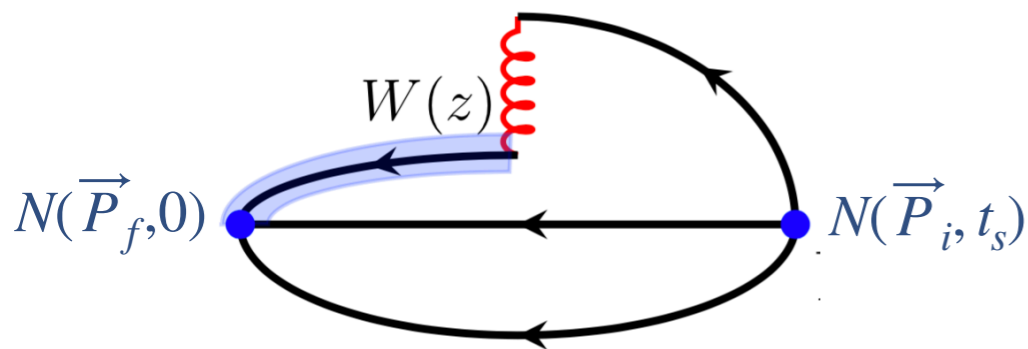
$$\xi = \frac{Q_3}{2P_3}$$

Access of PDFs/GPDs on a Euclidean Lattice

[X. Ji, Phys. Rev. Lett. 110 (2013) 262002]

Matrix elements of nonlocal (equal-time) operators with **fast moving hadrons**

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-i x P_3 z} \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z, 0) \Psi(0) | N(P_i) \rangle_{\mu}$$



Computationally intensive

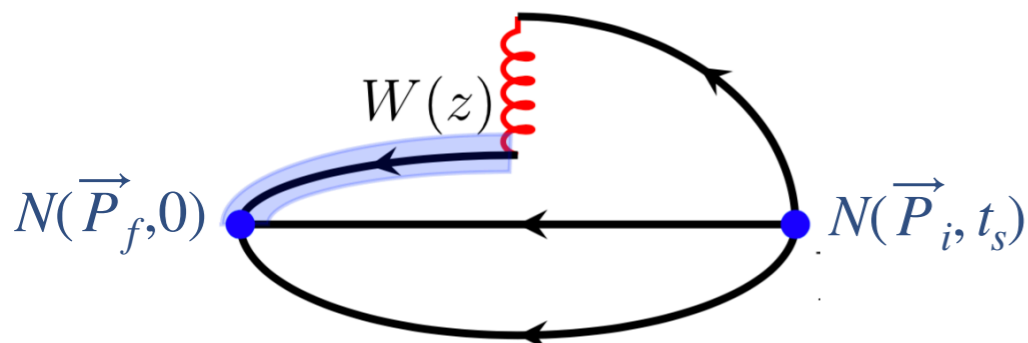
$$\Delta = P_f - P_i$$
$$t = \Delta^2 = -Q^2$$
$$\xi = \frac{Q_3}{2P_3}$$

Access of PDFs/GPDs on a Euclidean Lattice

[X. Ji, Phys. Rev. Lett. 110 (2013) 262002]

Matrix elements of nonlocal (equal-time) operators with **fast moving hadrons**

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-ixP_3z} \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z,0) \Psi(0) | N(P_i) \rangle_{\mu}$$

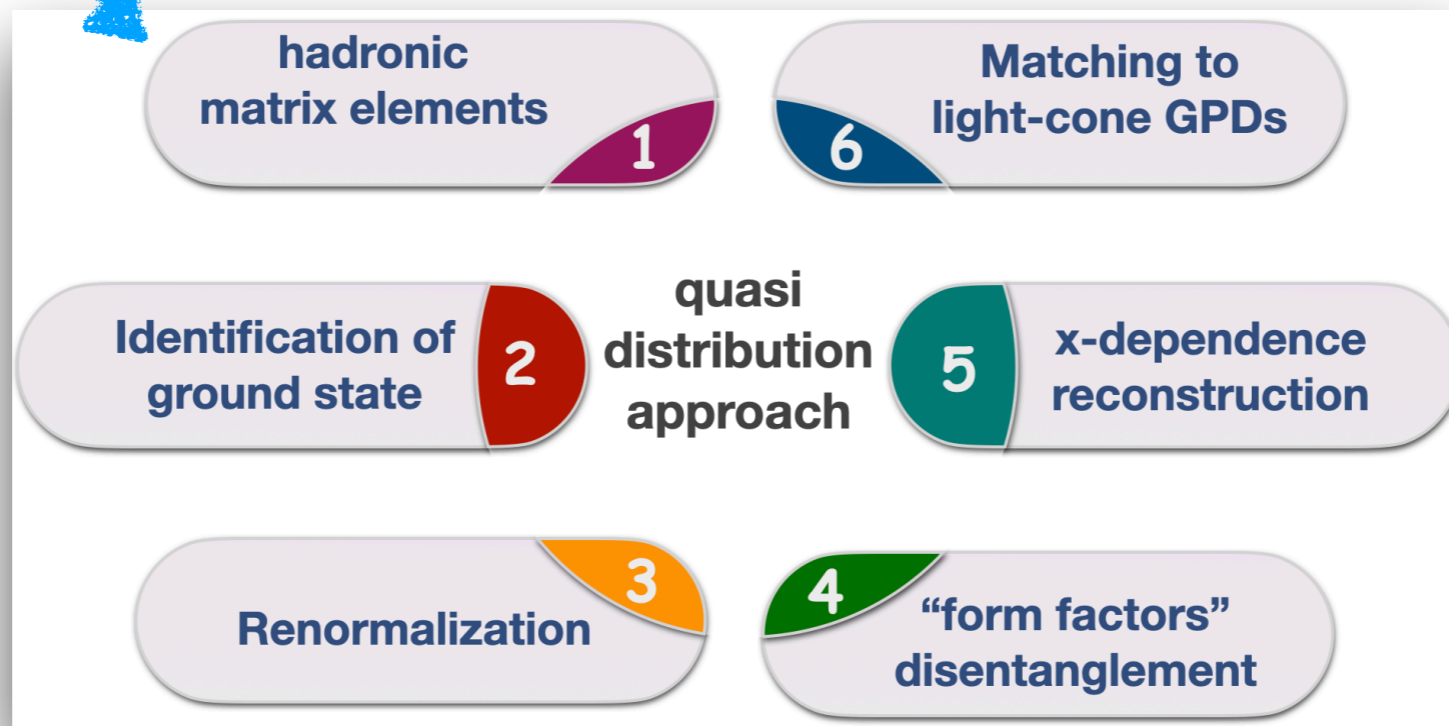


Computationally intensive

$$\Delta = P_f - P_i$$

$$t = \Delta^2 = -Q^2$$

$$\xi = \frac{Q_3}{2P_3}$$

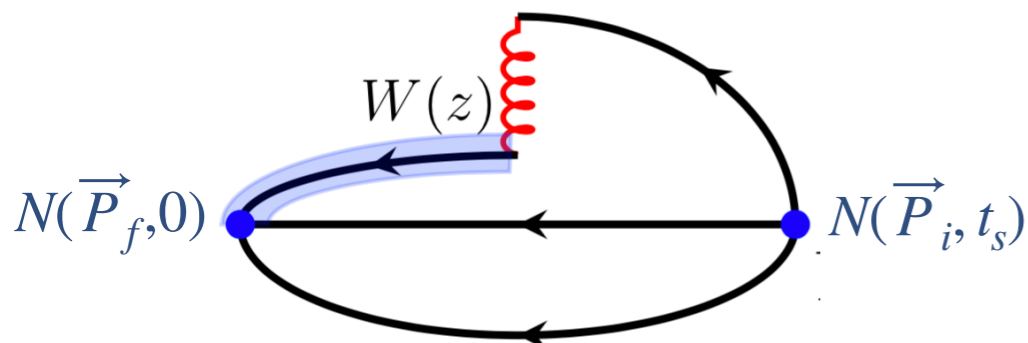


Access of PDFs/GPDs on a Euclidean Lattice

[X. Ji, Phys. Rev. Lett. 110 (2013) 262002]

Matrix elements of nonlocal (equal-time) operators with **fast moving hadrons**

$$\tilde{q}_{\Gamma}^{\text{GPD}}(x, t, \xi, P_3, \mu) = \int \frac{dz}{4\pi} e^{-ixP_3z} \langle N(P_f) | \bar{\Psi}(z) \Gamma \mathcal{W}(z,0) \Psi(0) | N(P_i) \rangle_{\mu}$$

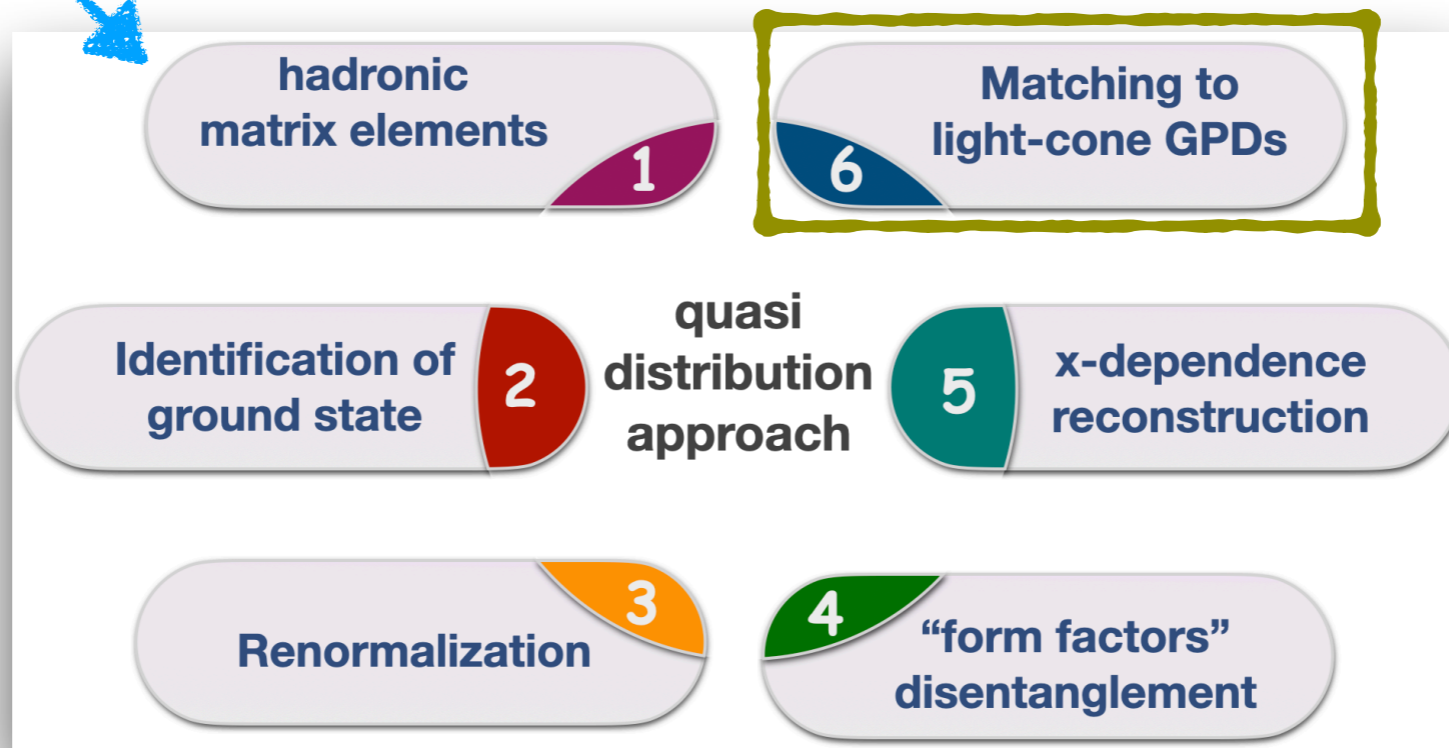


Computationally intensive

$$\Delta = P_f - P_i$$

$$t = \Delta^2 = -Q^2$$

$$\xi = \frac{Q_3}{2P_3}$$



Transversity PDFs

PHYSICAL REVIEW D **98**, 091503(R) (2018)

Rapid Communications

Transversity parton distribution functions from lattice QCD

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou,⁴
Karl Jansen,⁵ Aurora Scapellato,^{1,6} and Fernanda Steffens⁷

PHYSICAL REVIEW D **99**, 114504 (2019)

Systematic uncertainties in parton distribution functions from lattice QCD simulations at the physical point

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou,⁴ Kyriakos Hadjiyiannakou,² Karl Jansen,⁵
Aurora Scapellato,^{1,6} and Fernanda Steffens⁷

Parameters of calculations



★ Nf=2 twisted mass fermions with a clover term @ physical point

[Extended Twisted Mass Collaboration, Phys. Rev. D 95, 094515 (2017)]

$\beta = 2.10, c_{\text{SW}} = 1.57751, a = 0.0938(3)(2) \text{ fm}, r_0/a = 5.32(5)$

$48^3 \times 96, L \approx 4.5 \text{ fm}$

$a\mu = 0.0009$

$m_\pi = 0.1304(4) \text{ GeV}$

$m_\pi L = 2.98$

$m_N = 0.932(4) \text{ GeV}$

Parameters of calculations



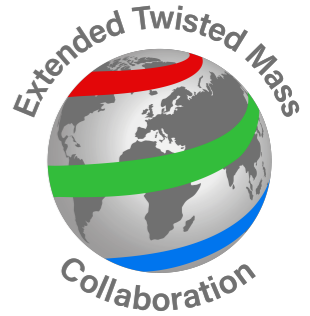
★ Nf=2 twisted mass fermions with a clover term @ physical point

[Extended Twisted Mass Collaboration, Phys. Rev. D 95, 094515 (2017)]

| | |
|---|--|
| $\beta=2.10, c_{\text{SW}}=1.57751, a=0.0938(3)(2) \text{ fm}, r_0/a=5.32(5)$ | |
| $48^3 \times 96, L \approx 4.5 \text{ fm}$ | $a\mu = 0.0009$ $m_\pi = 0.1304(4) \text{ GeV}$ $m_\pi L = 2.98$ $m_N = 0.932(4) \text{ GeV}$ |

| $P_3 = \frac{6\pi}{L}$ | | | | $P_3 = \frac{8\pi}{L}$ | | | | | $P_3 = \frac{10\pi}{L} \sim 1.38 \text{ GeV}$ | | | | |
|------------------------|-------------------|-----------------|-------------------|------------------------|-------------------|-----------------|-----------------|-------------------|---|-------------------|-----------------|-----------------|-------------------|
| Ins. | N_{conf} | N_{HP} | N_{meas} | Ins. | N_{conf} | N_{HP} | N_{LP} | N_{meas} | Ins. | N_{conf} | N_{HP} | N_{LP} | N_{meas} |
| γ^3 | 100 | 16 | 9600 | γ^3 | 425 | 1 | 16 | 38250 | γ^3 | 811 | 1 | 16 | 72990 |
| γ^0 | 50 | 16 | 4800 | γ^0 | 425 | 1 | 16 | 38250 | γ^0 | 811 | 1 | 16 | 72990 |
| $\gamma^5\gamma^3$ | 65 | 16 | 6240 | $\gamma^5\gamma^3$ | 425 | 1 | 16 | 38250 | $\gamma^5\gamma^3$ | 811 | 1 | 16 | 72990 |
| σ^{3j} | 50 | 16 | 9600 | σ^{3j} | 425 | 1 | 16 | 38250 | σ^{3j} | 811 | 1 | 16 | 72990 |

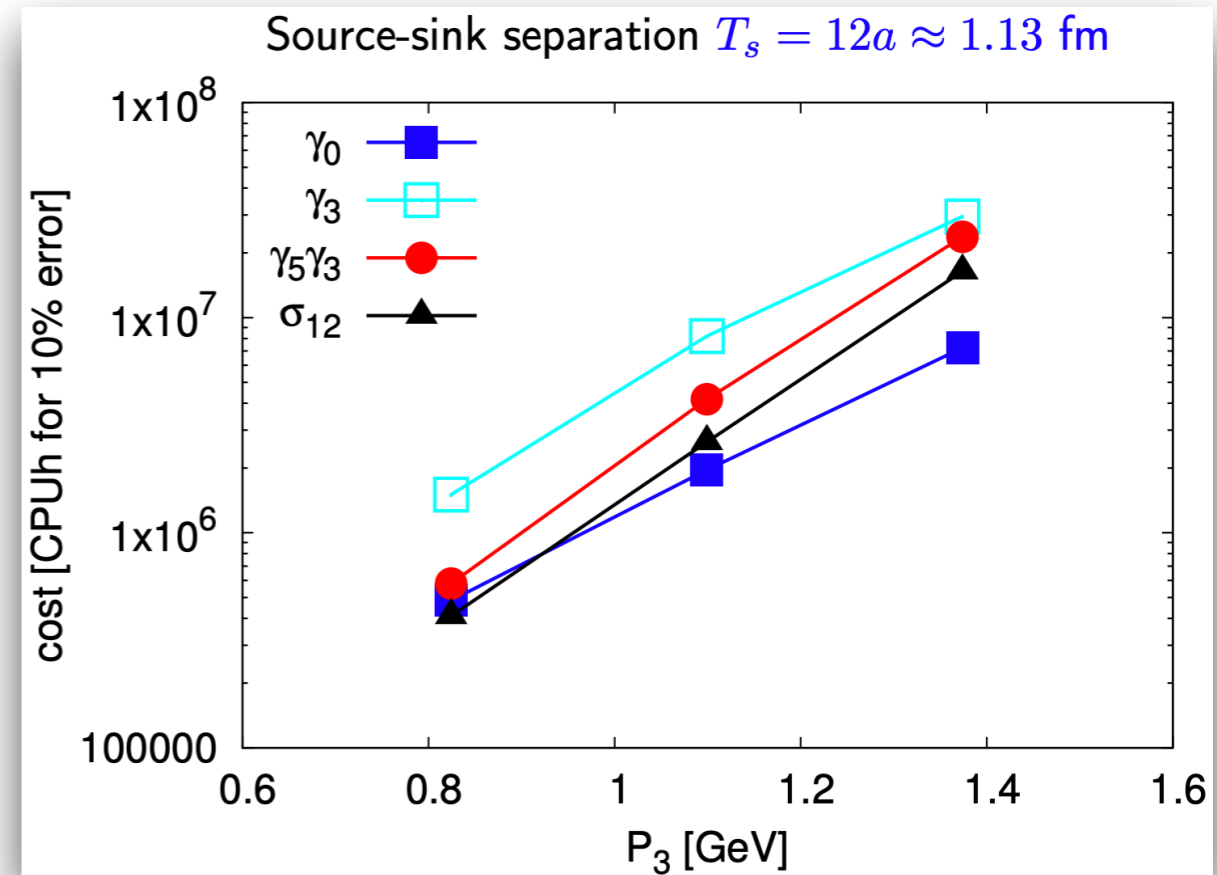
Parameters of calculations



★ Nf=2 twisted mass fermions with a clover term @ physical point

[Extended Twisted Mass Collaboration, Phys. Rev. D 95, 094515 (2017)]

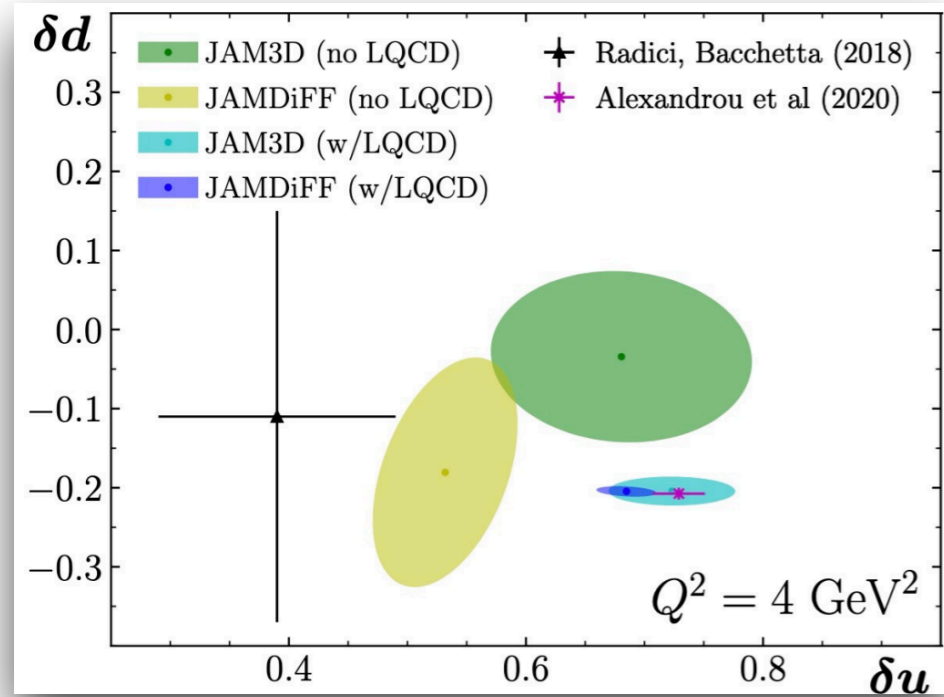
| | |
|---|--|
| $\beta=2.10, c_{\text{SW}}=1.57751, a=0.0938(3)(2) \text{ fm}, r_0/a=5.32(5)$ | |
| $48^3 \times 96, L \approx 4.5 \text{ fm}$ | $a\mu = 0.0009$ $m_\pi = 0.1304(4) \text{ GeV}$ $m_\pi L = 2.98$ $m_N = 0.932(4) \text{ GeV}$ |



| $P_3 = \frac{6\pi}{L}$ | | | | $P_3 = \frac{8\pi}{L}$ | | | | | $P_3 = \frac{10\pi}{L} \sim 1.38 \text{ GeV}$ | | | | |
|------------------------|-------------------|-----------------|-------------------|------------------------|-------------------|-----------------|-----------------|-------------------|---|-------------------|-----------------|-----------------|-------------------|
| Ins. | N_{conf} | N_{HP} | N_{meas} | Ins. | N_{conf} | N_{HP} | N_{LP} | N_{meas} | Ins. | N_{conf} | N_{HP} | N_{LP} | N_{meas} |
| γ^3 | 100 | 16 | 9600 | γ^3 | 425 | 1 | 16 | 38250 | γ^3 | 811 | 1 | 16 | 72990 |
| γ^0 | 50 | 16 | 4800 | γ^0 | 425 | 1 | 16 | 38250 | γ^0 | 811 | 1 | 16 | 72990 |
| $\gamma^5\gamma^3$ | 65 | 16 | 6240 | $\gamma^5\gamma^3$ | 425 | 1 | 16 | 38250 | $\gamma^5\gamma^3$ | 811 | 1 | 16 | 72990 |
| σ^{3j} | 50 | 16 | 9600 | σ^{3j} | 425 | 1 | 16 | 38250 | σ^{3j} | 811 | 1 | 16 | 72990 |



Transversity PDF

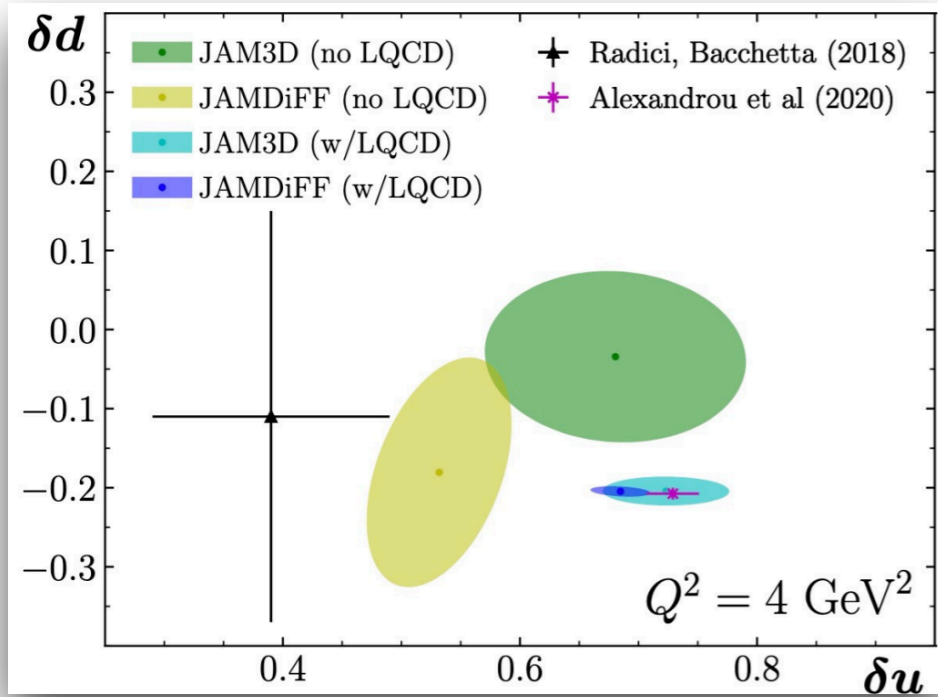


[Daniel Pitonyak, GHP 2023]

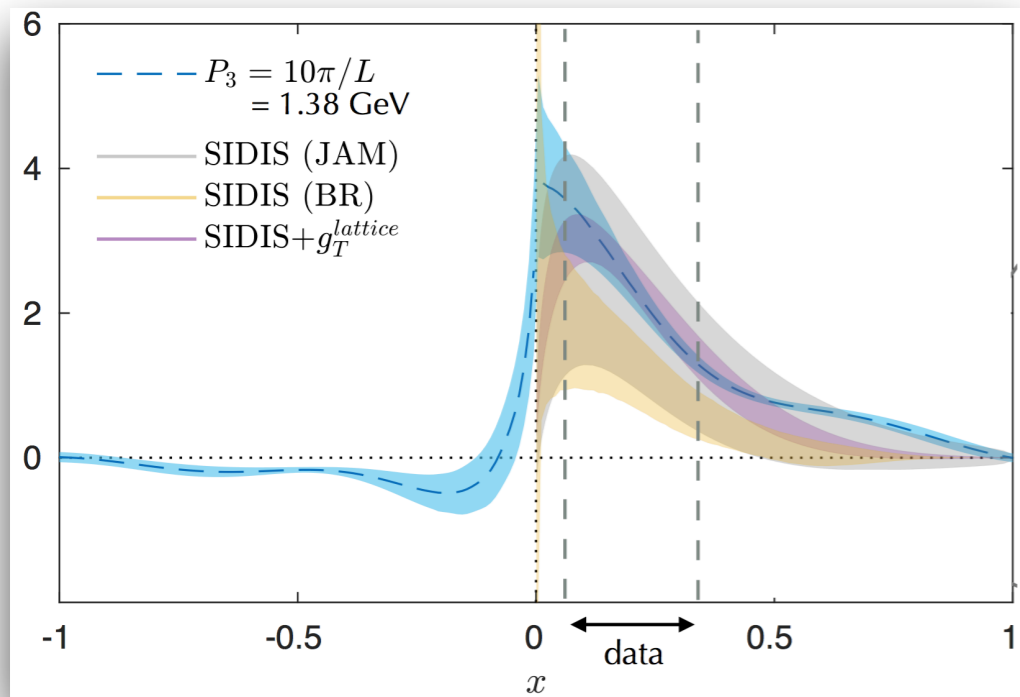
★ Opportunity for lattice QCD to complement global analysis

Transversity PDF

★ Opportunity for lattice QCD to complement global analysis



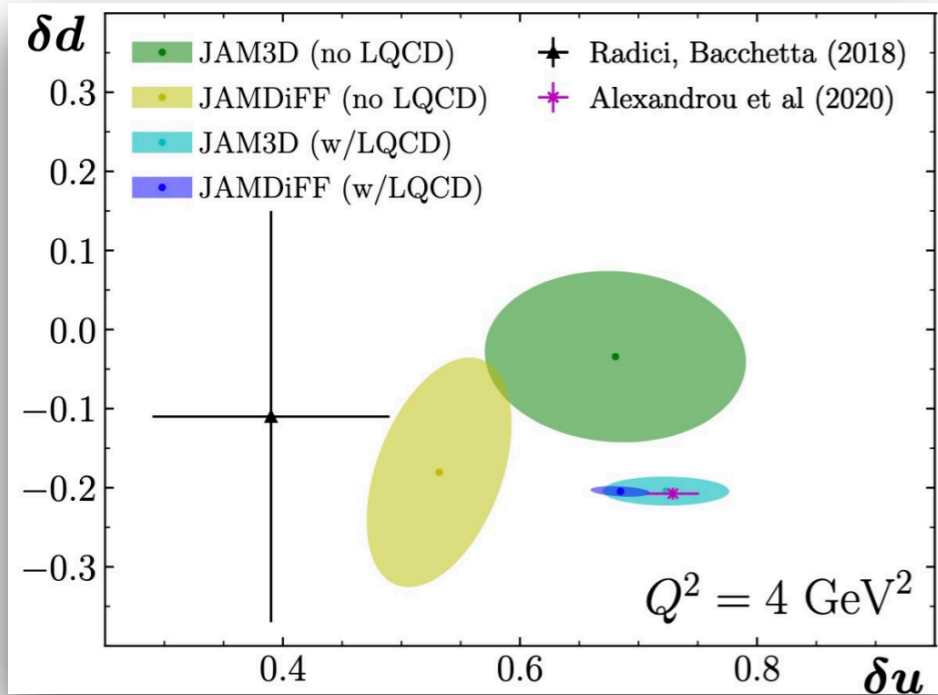
[Daniel Pitonyak, GHP 2023]



[C. Alexandrou et al. (ETMC), PRD 98 (2018) 091503 (R)]

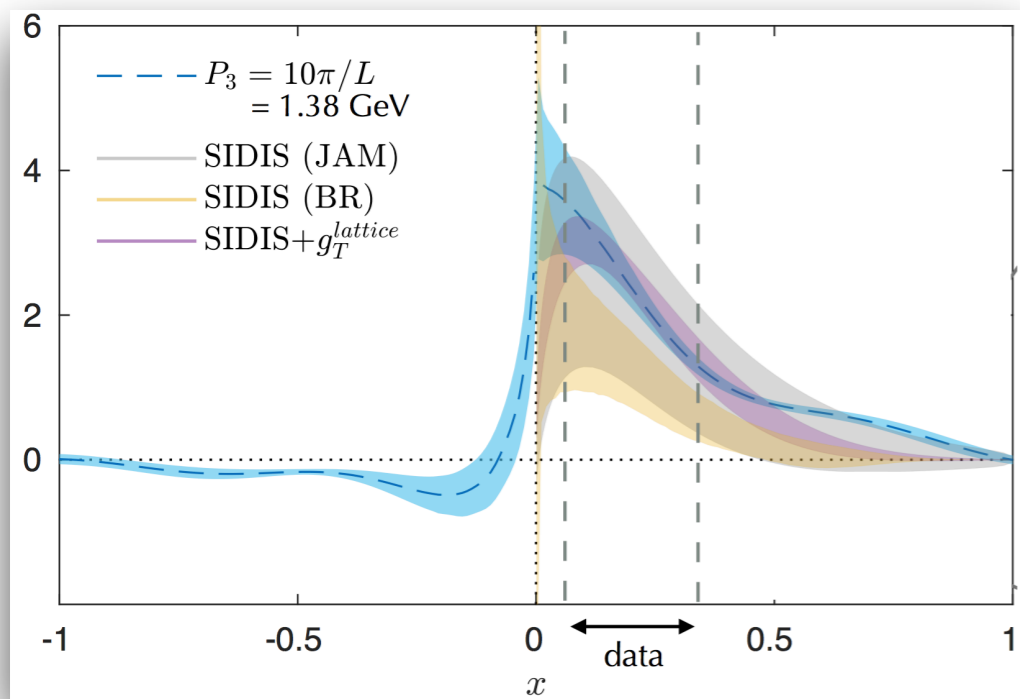
[C. Alexandrou et al. (ETMC), PRD 99 (2019) 114504]

Transversity PDF



[Daniel Pitonyak, GHP 2023]

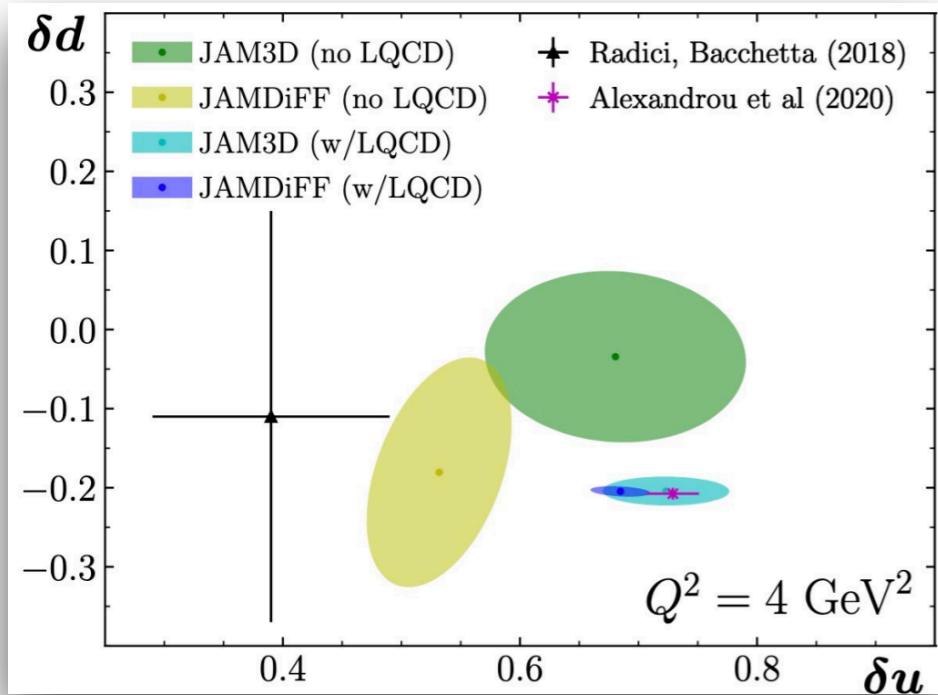
- ★ Opportunity for lattice QCD to complement global analysis
- ★ Field moves towards addressing systematic uncertainties
- ★ Disconnected contributions for flavor decomposition



[C. Alexandrou et al. (ETMC), PRD 98 (2018) 091503 (R)]

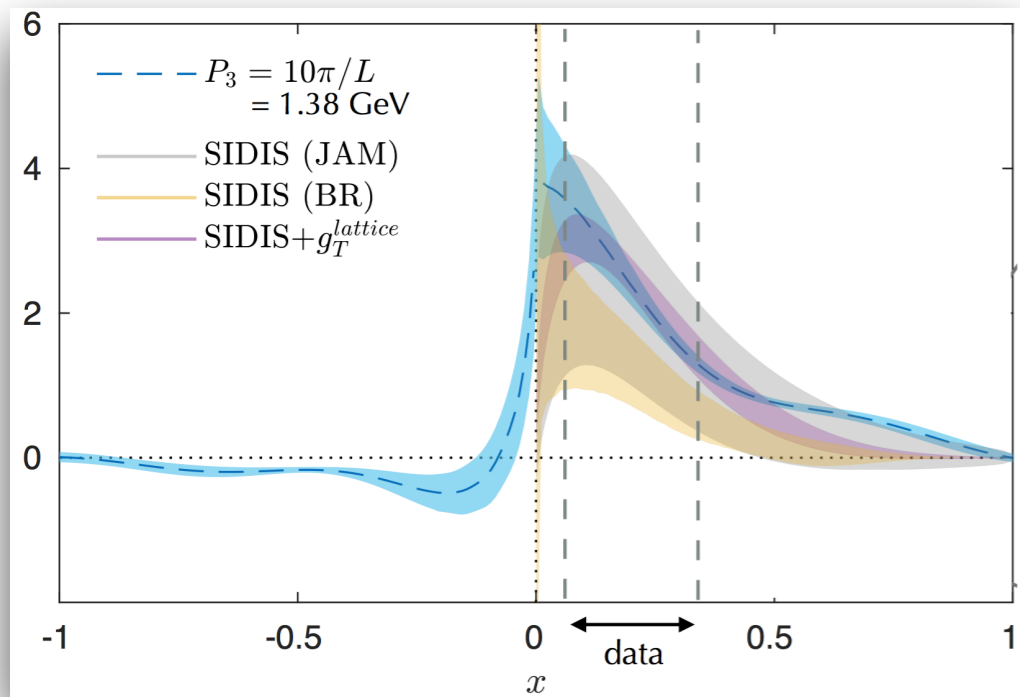
[C. Alexandrou et al. (ETMC), PRD 99 (2019) 114504]

Transversity PDF



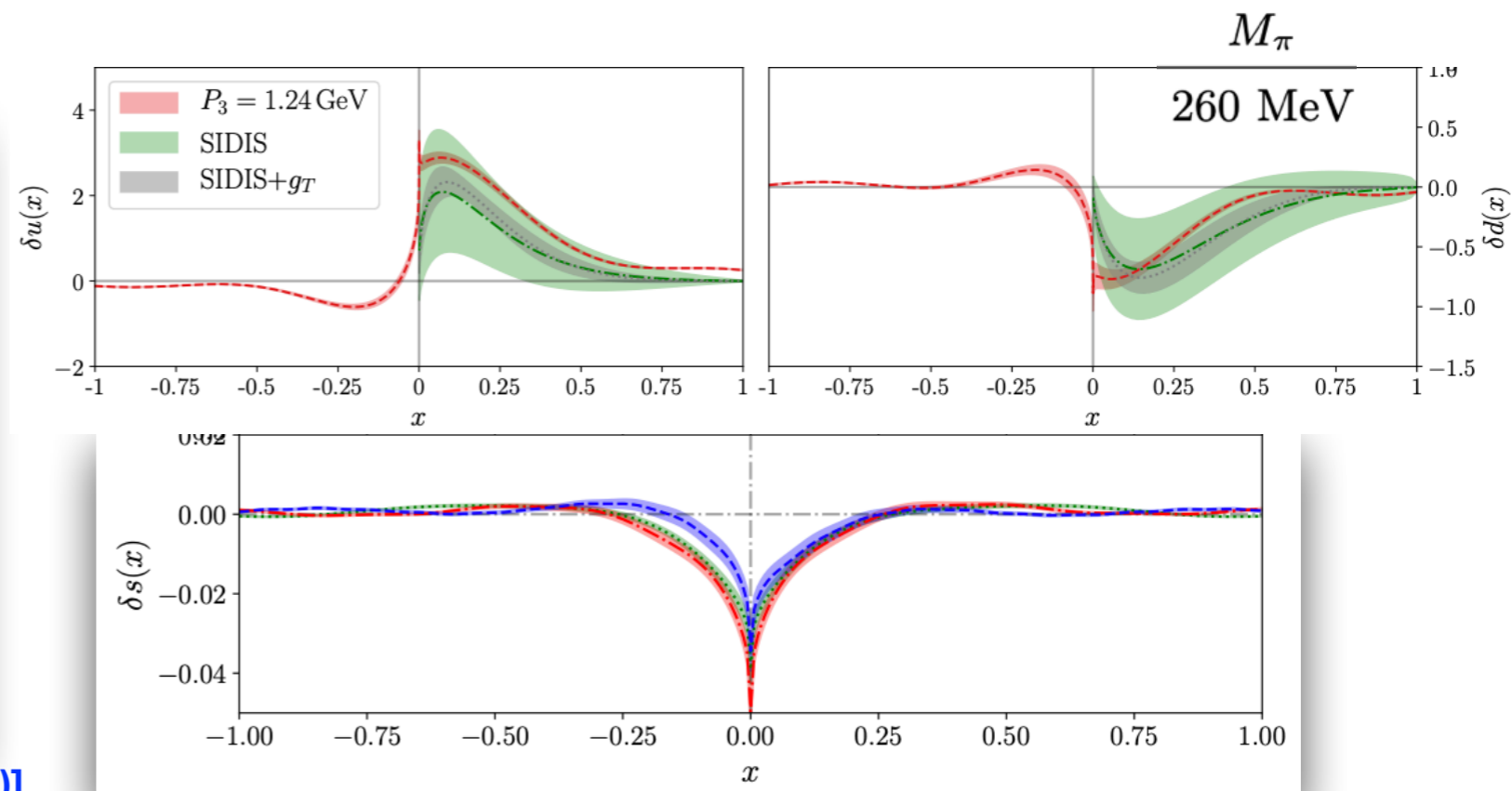
[Daniel Pitonyak, GHP 2023]

- ★ Opportunity for lattice QCD to complement global analysis
- ★ Field moves towards addressing systematic uncertainties
- ★ Disconnected contributions for flavor decomposition



[C. Alexandrou et al. (ETMC), PRD 98 (2018) 091503 (R)]

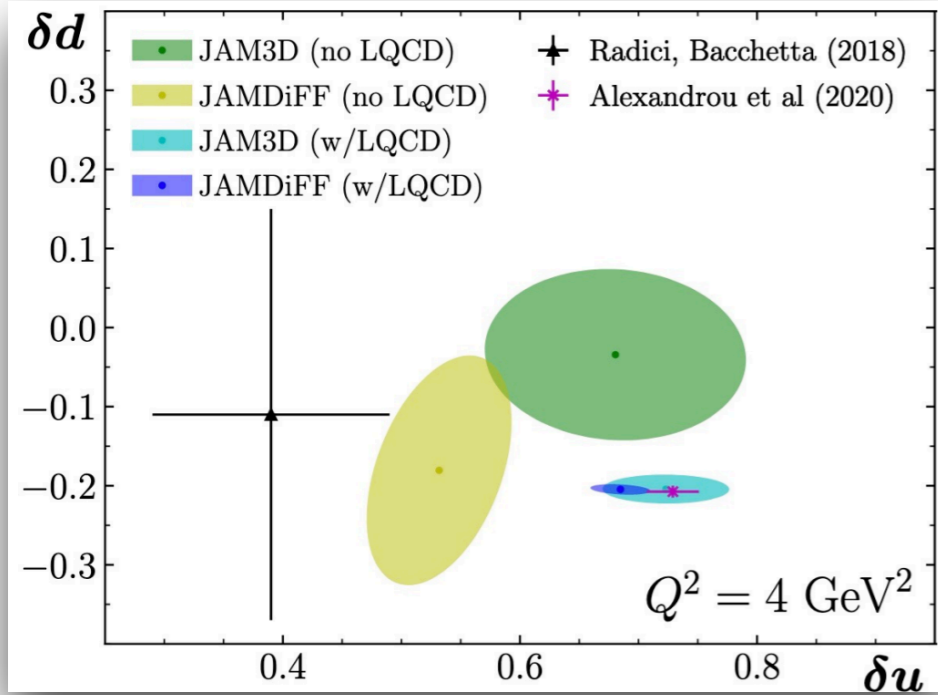
[C. Alexandrou et al. (ETMC), PRD 99 (2019) 114504]



[C. Alexandrou et al., (ETMC), PRD 104 (2021) 5, 054503]

Transversity PDF

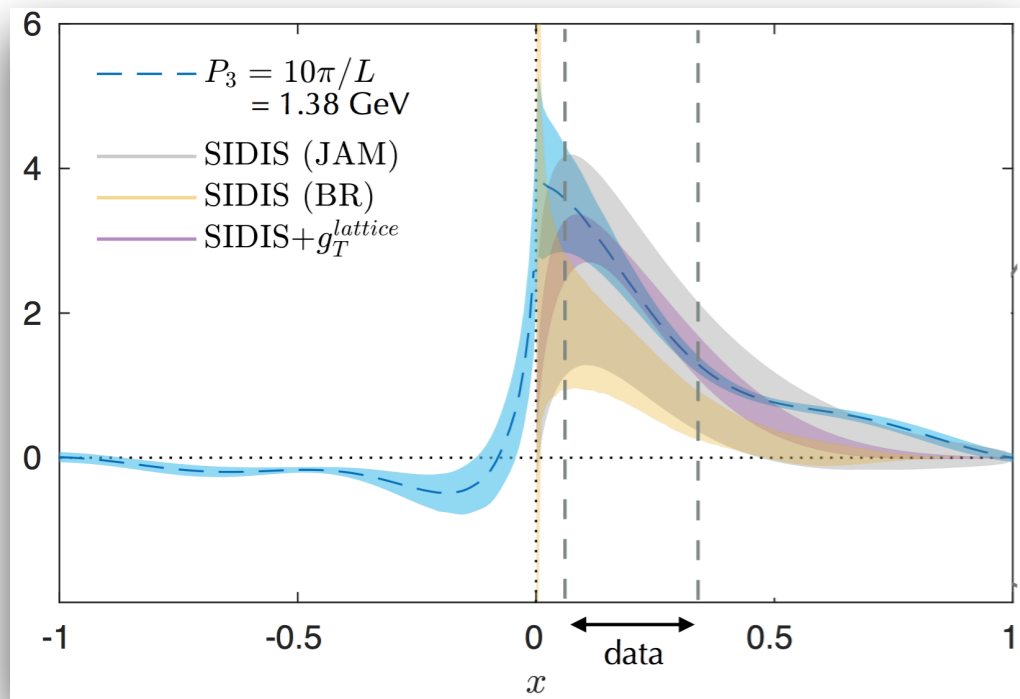
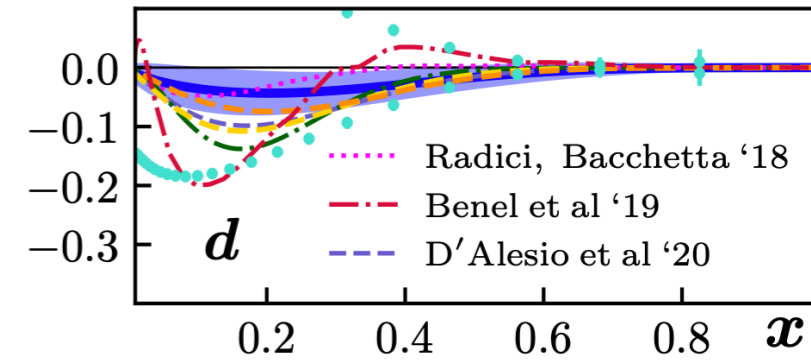
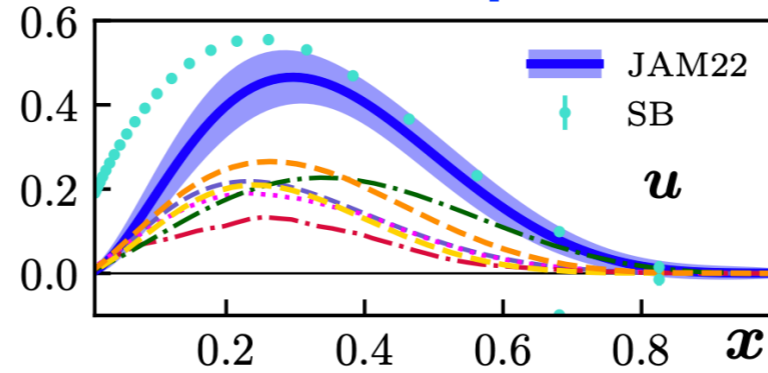
★ Opportunity for lattice QCD to complement global analysis



[Daniel Pitonyak, GHP 2023]

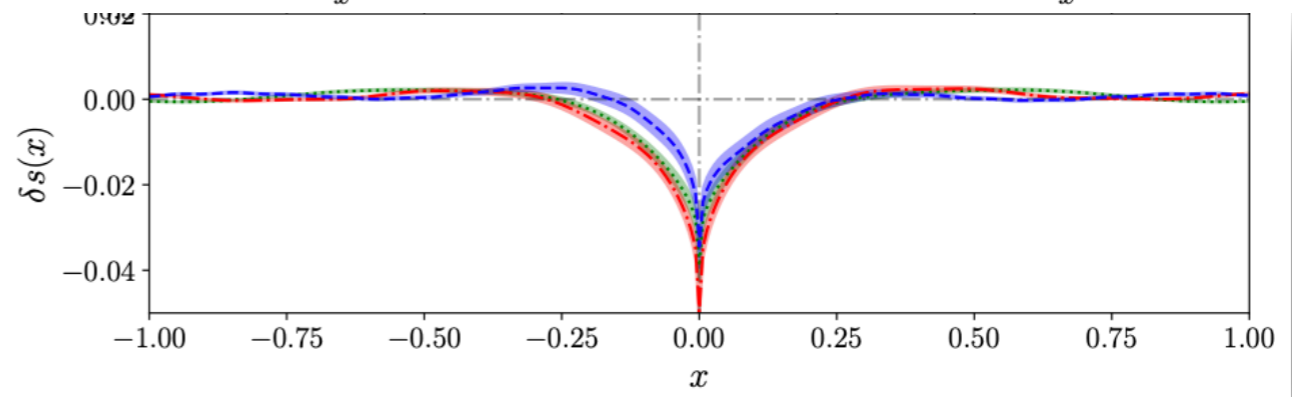
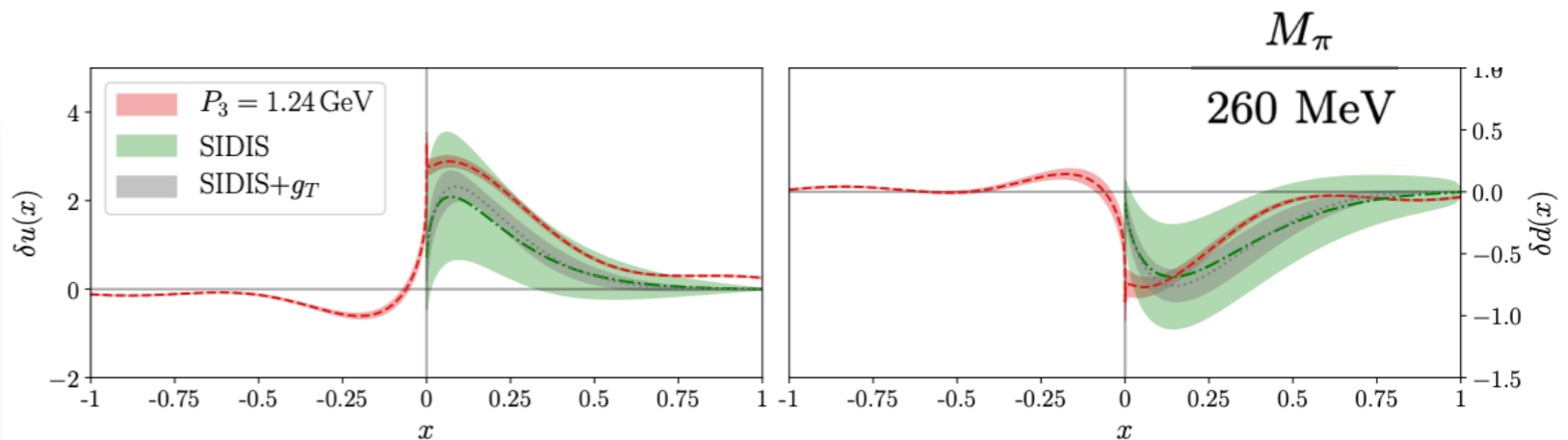
$xh_1(x)$

[JAM Collaboration, PRD 106 (2022) 3, 034014]



[C. Alexandrou et al. (ETMC), PRD 98 (2018) 091503 (R)]

[C. Alexandrou et al. (ETMC), PRD 99 (2019) 114504]



[C. Alexandrou et al., (ETMC), PRD 104 (2021) 5, 054503]



Transversity GPDs

PHYSICAL REVIEW D **105**, 034501 (2022)

Transversity GPDs of the proton from lattice QCD

Constantia Alexandrou,^{1,2} Krzysztof Cichy,³ Martha Constantinou⁴,⁴ Kyriakos Hadjiyiannakou,^{1,2}
Karl Jansen,⁵ Aurora Scapellato,⁴ and Fernanda Steffens⁶

Parameters of calculations



★ $N_f=2+1+1$ twisted mass fermions with a clover term;

[Extended Twisted Mass Collaboration, Phys. Rev. D 104, 074515 (2021), arXiv:2104.13408]

| Name | β | N_f | $L^3 \times T$ | a [fm] | M_π | $m_\pi L$ |
|----------|---------|--------------|------------------|----------|---------|-----------|
| cA211.32 | 1.726 | u, d, s, c | $32^3 \times 64$ | 0.093 | 260 MeV | 4 |

| P_3 [GeV] | Δ [$\frac{2\pi}{L}$] | $-t$ [GeV ²] | ξ | N_{confs} | N_{meas} |
|-------------|-------------------------------|--------------------------|-------|--------------------|-------------------|
| 0.83 | (0,2,0) | 0.69 | 0 | 519 | 4152 |
| 1.25 | (0,2,0) | 0.69 | 0 | 1315 | 42080 |
| 1.67 | (0,2,0) | 0.69 | 0 | 1753 | 112192 |
| 1.25 | (0,2,2) | 1.02 | 1/3 | 417 | 40032 |
| 1.25 | (0,2,-2) | 1.02 | -1/3 | 417 | 40032 |

Collaboration

Parameters of calculations



★ Nf=2+1+1 twisted mass fermions with a clover term;

[Extended Twisted Mass Collaboration, Phys. Rev. D 104, 074515 (2021), arXiv:2104.13408]

| Name | β | N_f | $L^3 \times T$ | a [fm] | M_π | $m_\pi L$ |
|----------|---------|--------------|------------------|----------|---------|-----------|
| cA211.32 | 1.726 | u, d, s, c | $32^3 \times 64$ | 0.093 | 260 MeV | 4 |

| P_3 [GeV] | Δ [$\frac{2\pi}{L}$] | $-t$ [GeV ²] | ξ | N_{confs} | N_{meas} |
|-------------|-------------------------------|--------------------------|-------|--------------------|-------------------|
| 0.83 | (0,2,0) | 0.69 | 0 | 519 | 4152 |
| 1.25 | (0,2,0) | 0.69 | 0 | 1315 | 42080 |
| 1.67 | (0,2,0) | 0.69 | 0 | 1753 | 112192 |
| 1.25 | (0,2,2) | 1.02 | 1/3 | 417 | 40032 |
| 1.25 | (0,2,-2) | 1.02 | -1/3 | 417 | 40032 |

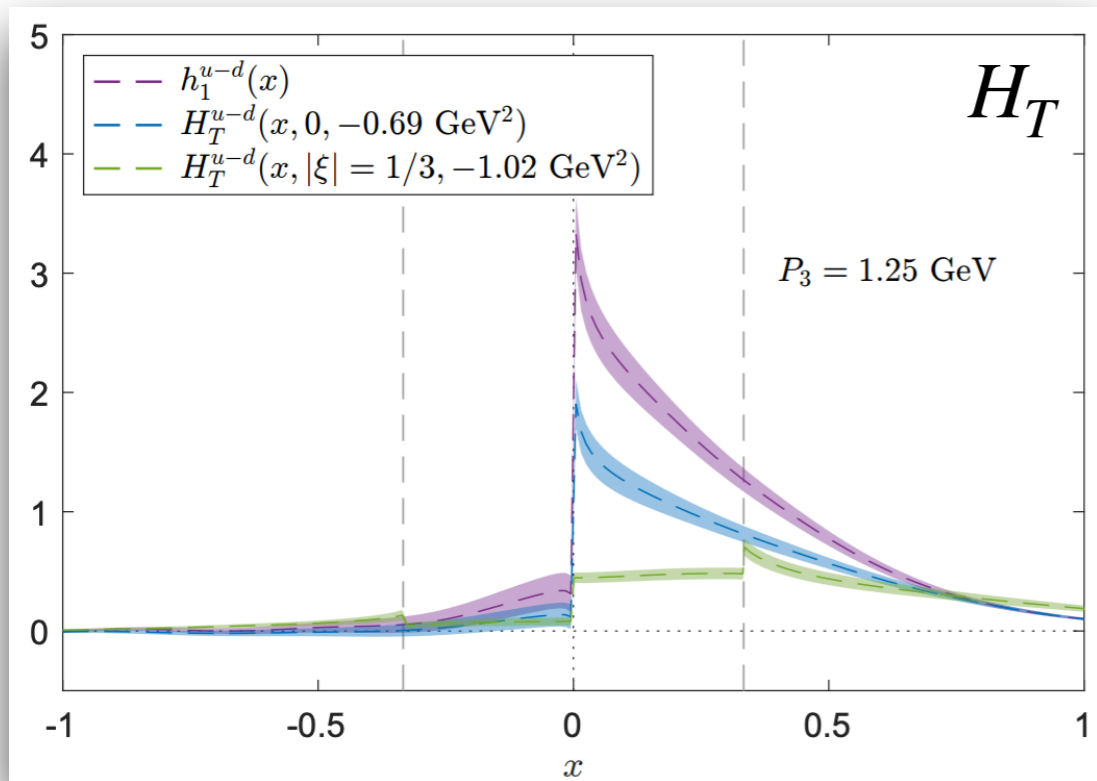
★ Parametrization

[Meissner et al., JHEP 08 (2009) 056]

$$\begin{aligned}
 h_T^j(\Gamma_\nu, z, P_f, P_i) = & \langle\langle \sigma^{3j} \rangle\rangle F_{HT}(z, \xi, t, P_3) + \frac{i}{2m} \langle\langle \gamma^3 \Delta_j - \gamma^j \Delta_3 \rangle\rangle F_{ET}(z, \xi, t, P_3) \\
 & + \frac{P_3 \Delta_j - P_j \Delta_3}{m^2} \langle\langle \hat{1} \rangle\rangle F_{\tilde{HT}}(z, \xi, t, P_3) + \frac{1}{m} \langle\langle \gamma^3 P_j - \gamma^j P_3 \rangle\rangle F_{\tilde{ET}}(z, \xi, t, P_3)
 \end{aligned}$$

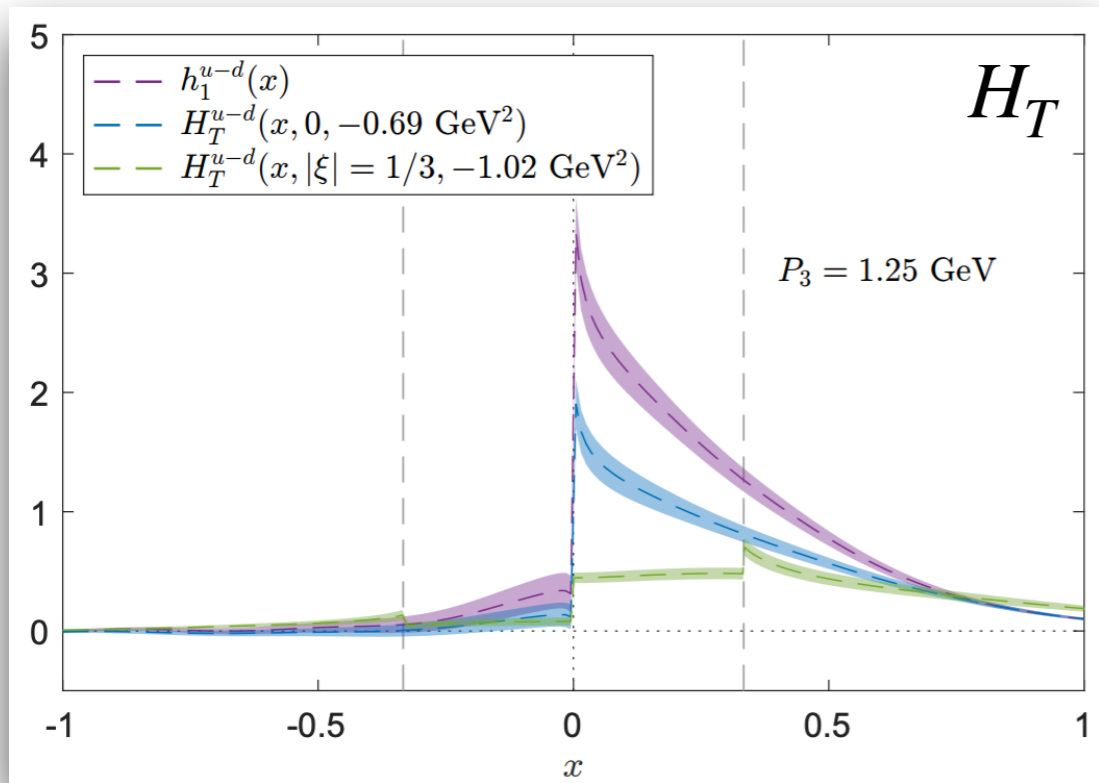
x-dependent transversity GPDs

x-dependent transversity GPDs



- ★ ERBL: decrease of H_T as $-t$ increases
- ★ DGLAP region shows less sensitivity in $-t$

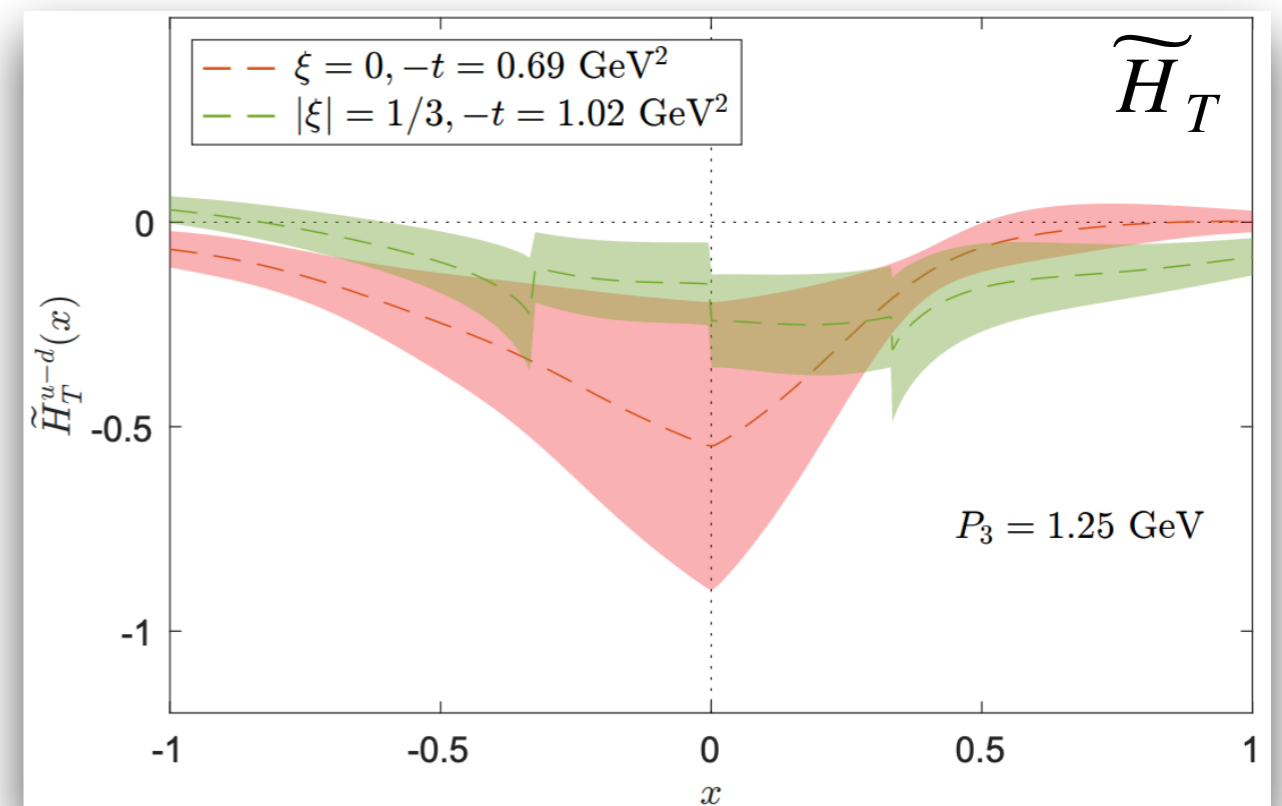
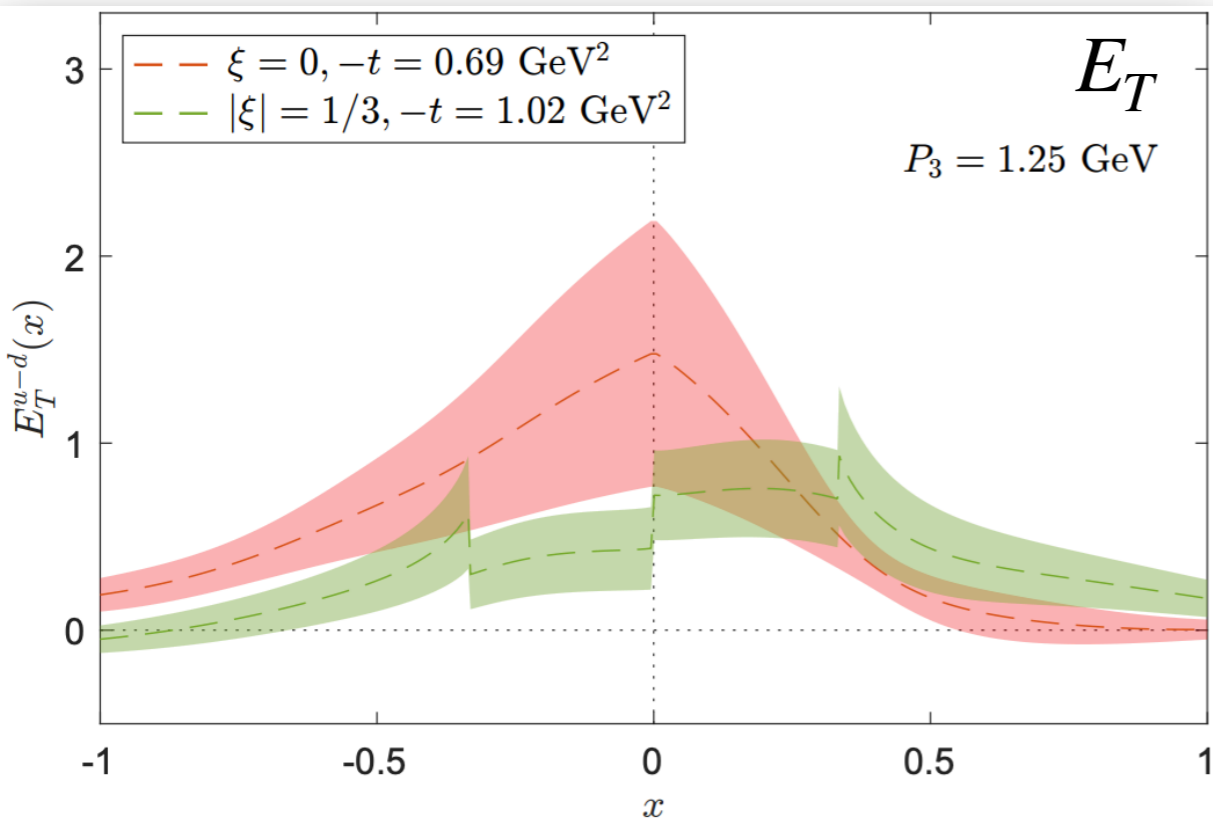
x-dependent transversity GPDs



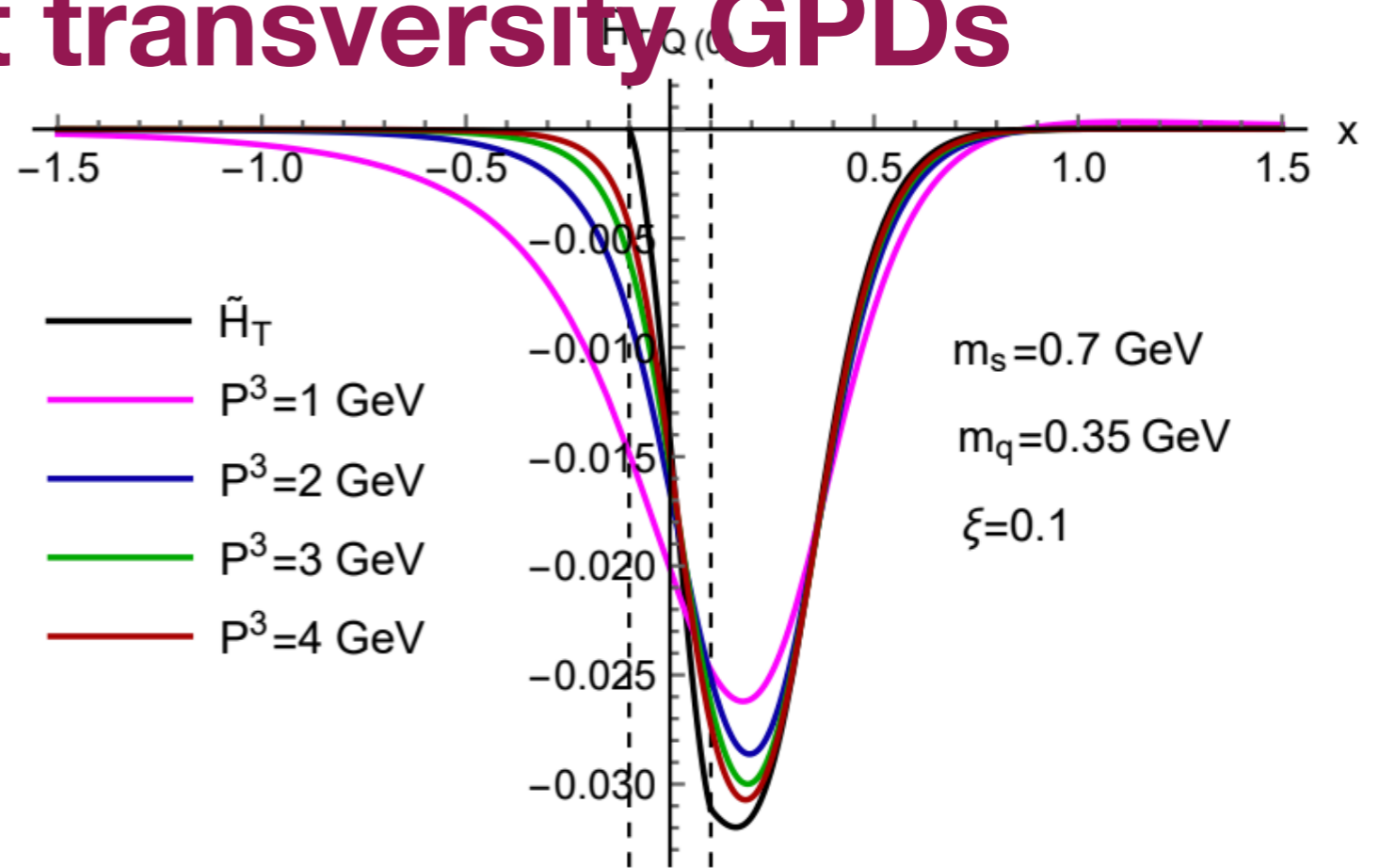
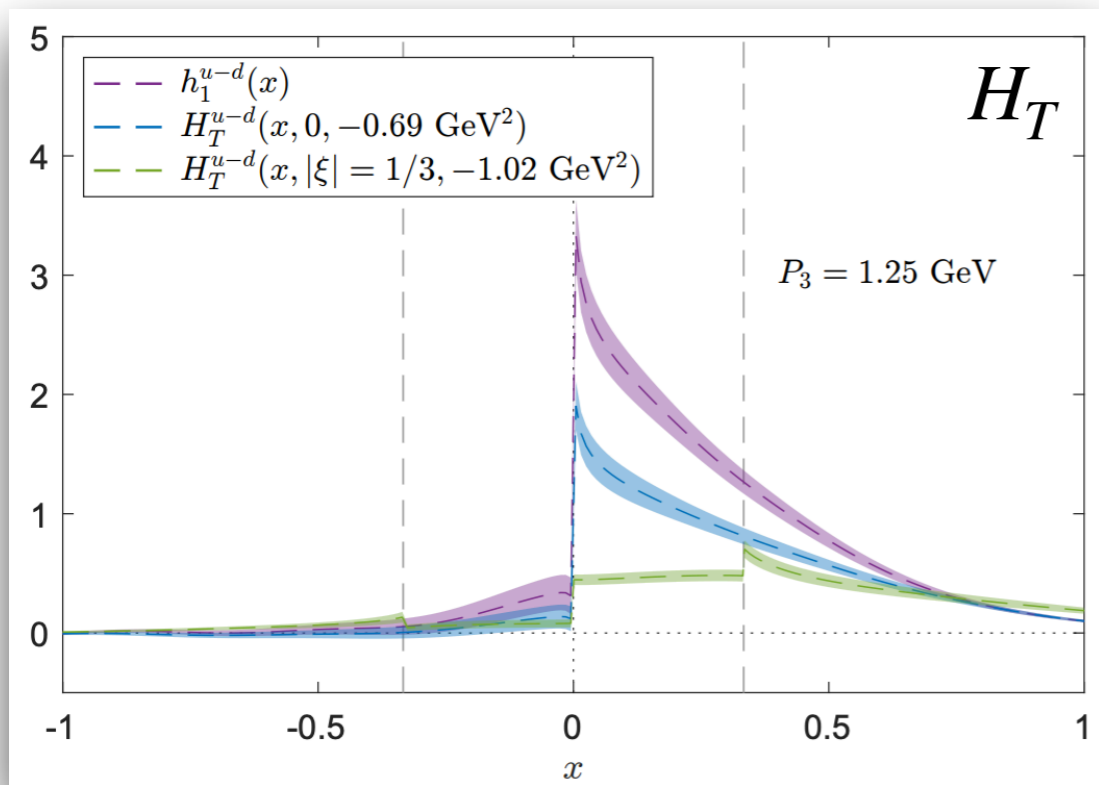
- ★ ERBL: decrease of H_T as $-t$ increases
- ★ DGLAP region shows less sensitivity in $-t$

- ★ Qualitatively E_T , \widetilde{H}_T show approximate symmetry of quark & antiquark regions
- ★ $\widetilde{H}_T < 0$ as in scalar diquark model

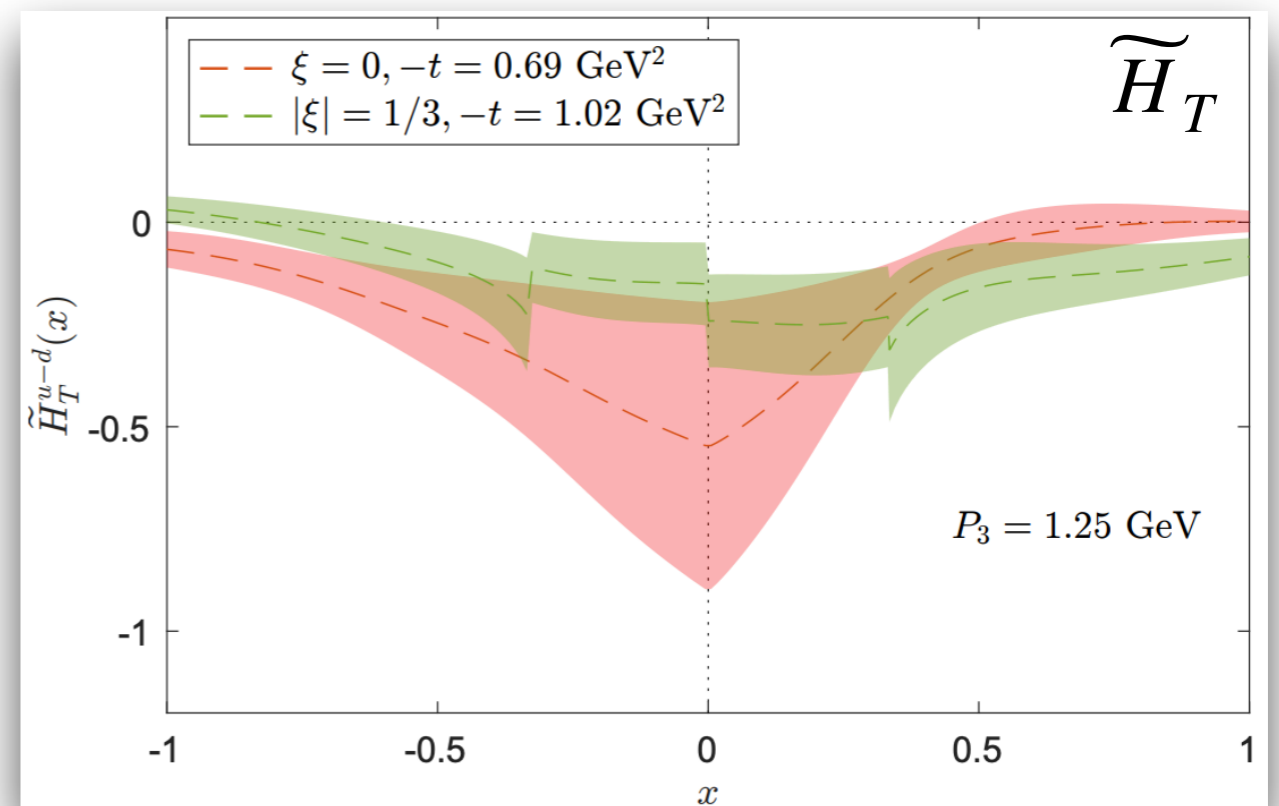
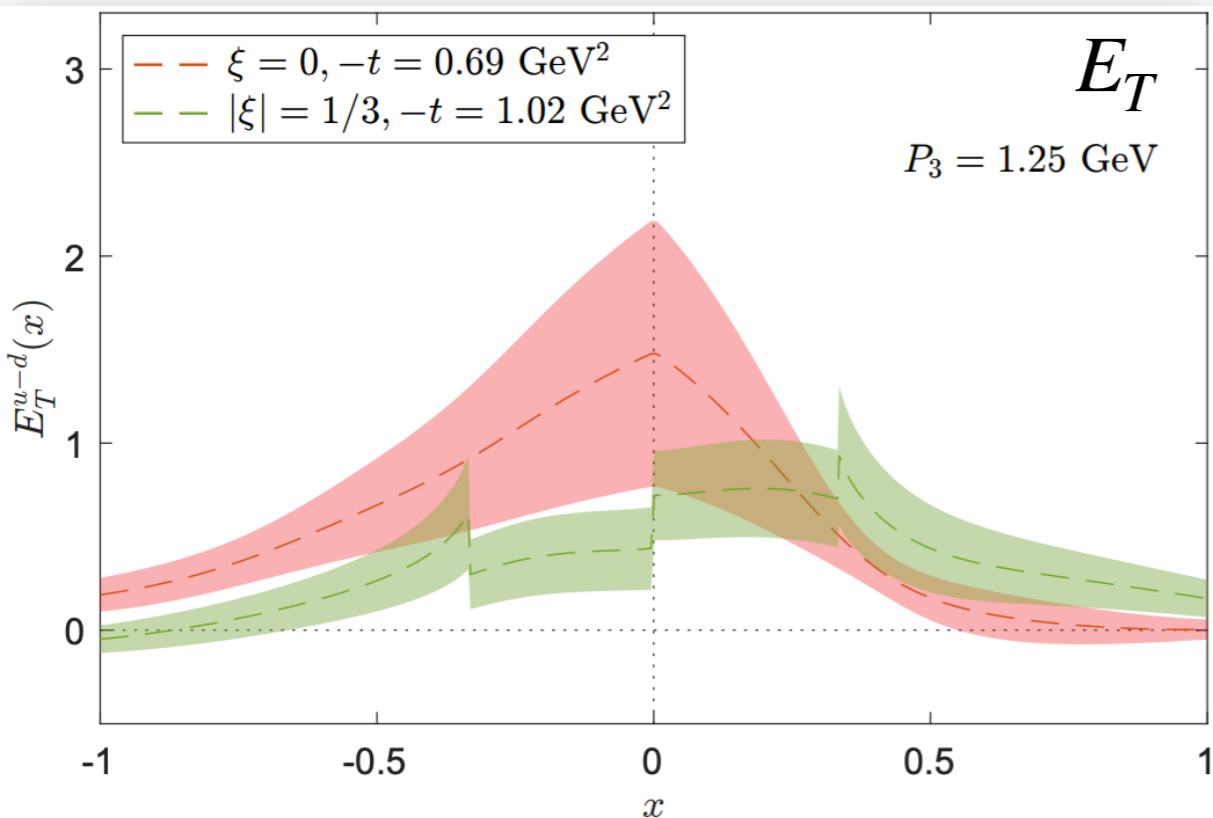
[S. Bhattacharya et al, Phys. Rev. D 102, 054021 (2020)]



x-dependent transversity GPDs



[S. Bhattacharya et al, Phys. Rev. D 102, 054021 (2020)]



x-dependent transversity GPDs

- ★ $[2\widetilde{H}_T + E_T](0)$ lowest Mellin moment:

transverse spin-flavor dipole moment in unpolarized target k_T

[M. Burkardt, Phys. Rev. D 72, 094020 (2005)]

- ★ $[2\widetilde{H}_T + E_T](0)$ second Mellin moment:

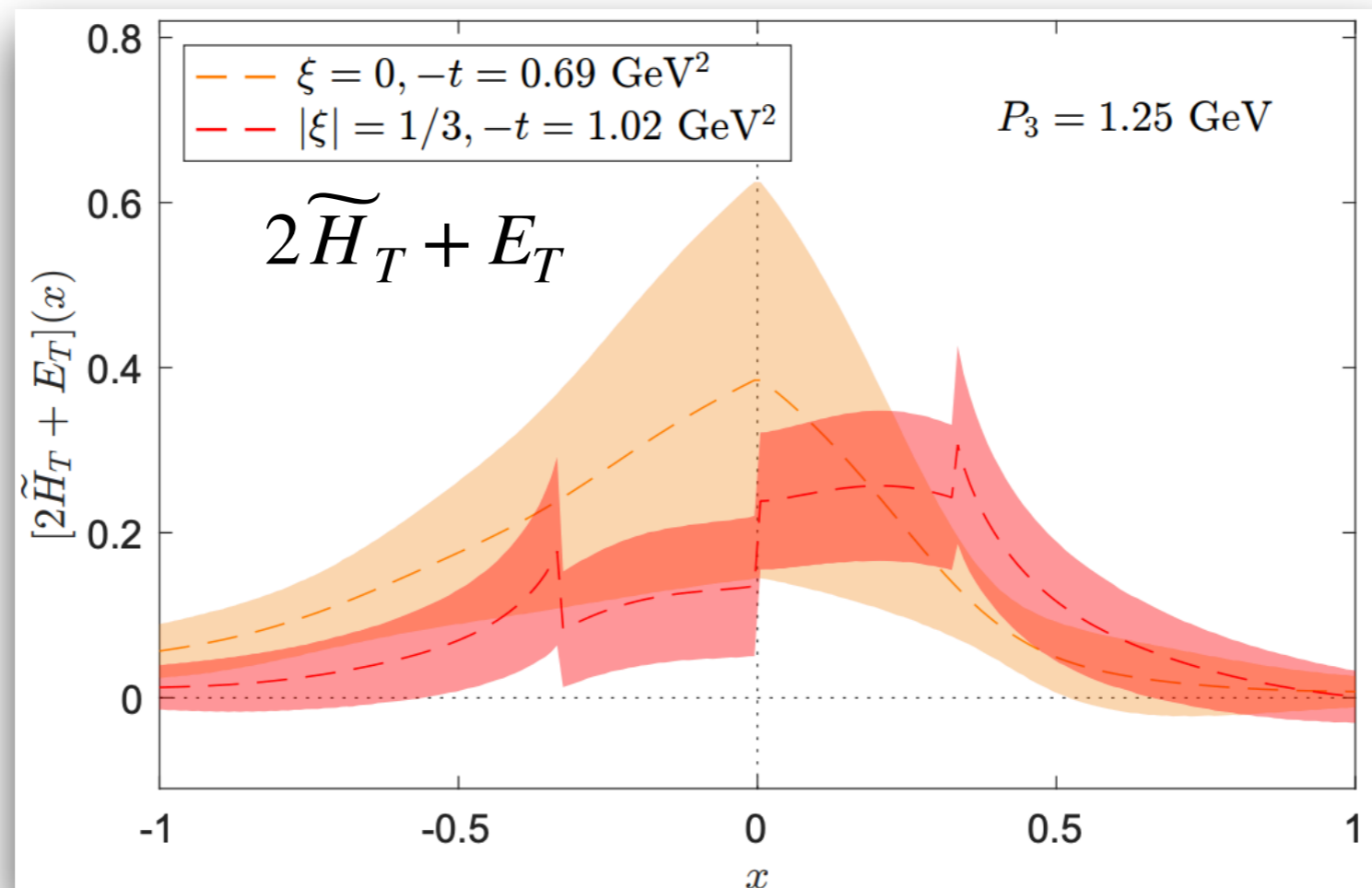
related to transverse-spin quark angular momentum in unpolarized proton

x-dependent transversity GPDs

- ★ $[2\widetilde{H}_T + E_T](0)$ lowest Mellin moment:
transverse spin-flavor dipole moment in unpolarized target k_T

[M. Burkardt, Phys. Rev. D 72, 094020 (2005)]

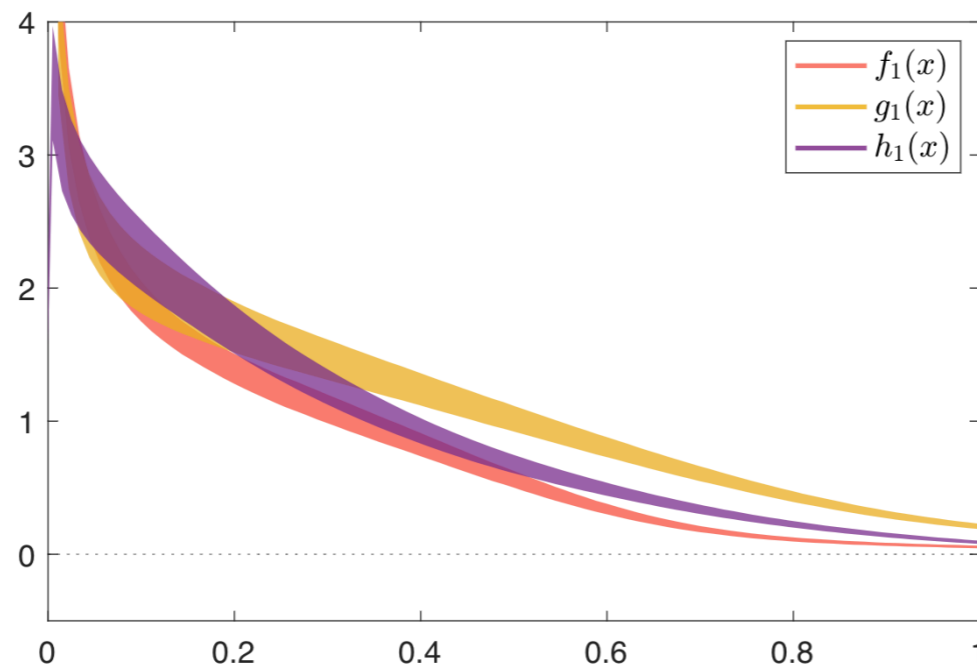
- ★ $[2\widetilde{H}_T + E_T](0)$ second Mellin moment:
related to transverse-spin quark angular momentum in unpolarized proton



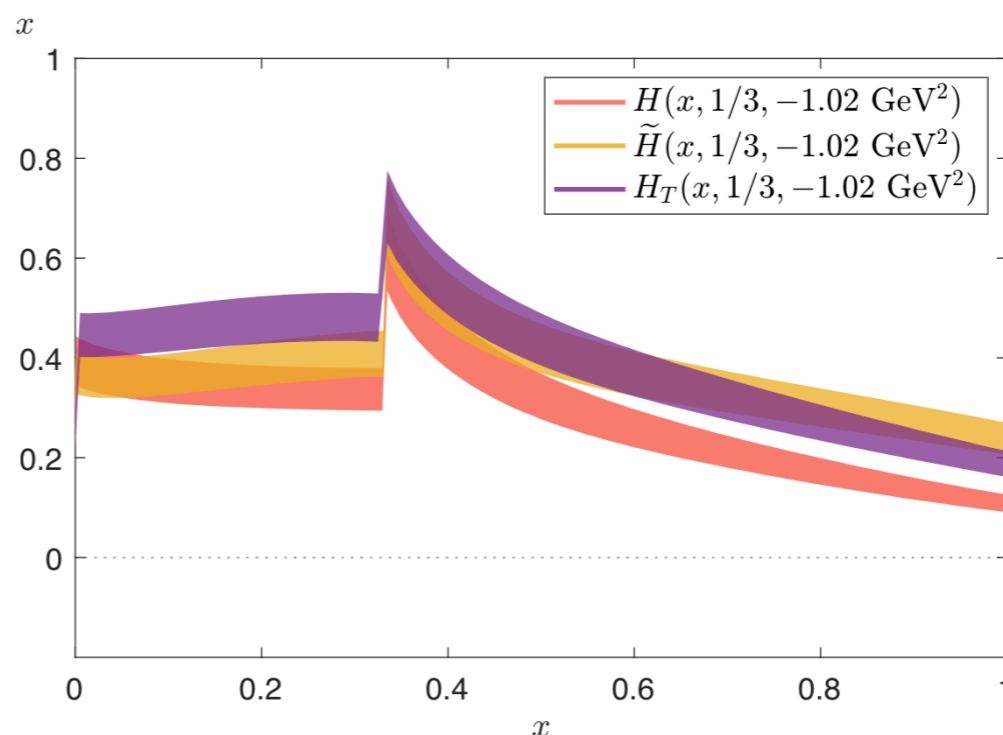
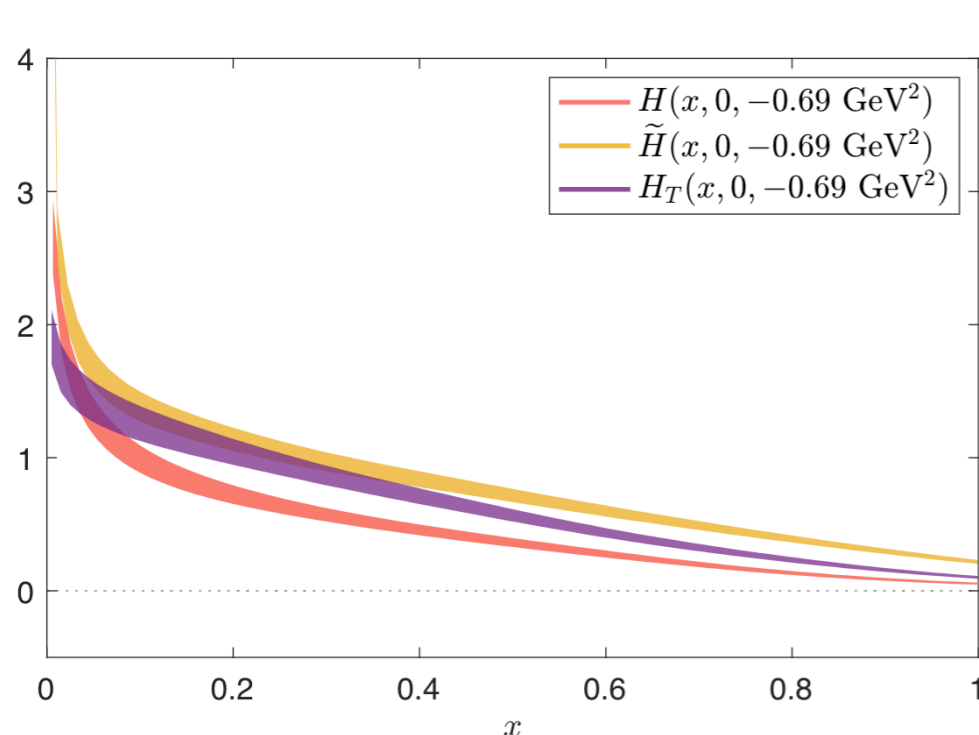
- ★ distribution for $|\xi| = 1/3$ tends to be systematically lower than the one at $\xi = 0$

x-dependent transversity GPDs

- ★ Qualitative understanding of GPDs and their relations
- ★ Qualitative understanding of ERBL and DGLAP regions



★ Relations can be identified for the t -dependence of GPDs



What can we currently check using lattice results?



What can we currently check using lattice results?

Address system. effects
through sum rules

$$\int_{-1}^1 dx H_T(x, \xi, t) = \int_{-\infty}^{\infty} dx H_{Tq}(x, \xi, t, P_3) = A_{T10}(t), \quad \int_{-1}^1 dx x H_T(x, \xi, t) = A_{T20}(t),$$

$$\int_{-1}^1 dx E_T(x, \xi, t) = \int_{-\infty}^{\infty} dx E_{Tq}(x, \xi, t, P_3) = B_{T10}(t), \quad \int_{-1}^1 dx x E_T(x, \xi, t) = B_{T20}(t),$$

[S. Bhattacharya et al., PRD 102, 054021 (2020)]

Sum rules not imposed
in calculation

$$\int_{-1}^1 dx \tilde{H}_T(x, \xi, t) = \int_{-\infty}^{\infty} dx \tilde{H}_{Tq}(x, \xi, t, P_3) = \tilde{A}_{T10}(t), \quad \int_{-1}^1 dx x \tilde{H}_T(x, \xi, t) = \tilde{A}_{T20}(t),$$

$$\int_{-1}^1 dx \tilde{E}_T(x, \xi, t) = \int_{-\infty}^{\infty} dx \tilde{E}_{Tq}(x, \xi, t, P_3) = 0. \quad \int_{-1}^1 dx x \tilde{E}_T(x, \xi, t) = 2\xi \tilde{B}_{T21}(t).$$

What can we currently check using lattice results?

Address system. effects through sum rules

$$\int_{-1}^1 dx H_T(x, \xi, t) = \int_{-\infty}^{\infty} dx H_{Tq}(x, \xi, t, P_3) = A_{T10}(t), \quad \int_{-1}^1 dx x H_T(x, \xi, t) = A_{T20}(t),$$

$$\int_{-1}^1 dx E_T(x, \xi, t) = \int_{-\infty}^{\infty} dx E_{Tq}(x, \xi, t, P_3) = B_{T10}(t), \quad \int_{-1}^1 dx x E_T(x, \xi, t) = B_{T20}(t),$$

[S. Bhattacharya et al., PRD 102, 054021 (2020)]

Sum rules not imposed in calculation

$$\int_{-1}^1 dx \tilde{H}_T(x, \xi, t) = \int_{-\infty}^{\infty} dx \tilde{H}_{Tq}(x, \xi, t, P_3) = \tilde{A}_{T10}(t), \quad \int_{-1}^1 dx x \tilde{H}_T(x, \xi, t) = \tilde{A}_{T20}(t),$$

$$\int_{-1}^1 dx \tilde{E}_T(x, \xi, t) = \int_{-\infty}^{\infty} dx \tilde{E}_{Tq}(x, \xi, t, P_3) = 0. \quad \int_{-1}^1 dx x \tilde{E}_T(x, \xi, t) = 2\xi \tilde{B}_{T21}(t).$$

★ Results

$$\int_{-2}^2 dx H_{Tq}(x, 0, -0.69 \text{ GeV}^2, P_3) = \{0.65(4), 0.64(6), 0.81(10)\}, \quad \int_{-2}^2 dx H_{Tq}(x, \frac{1}{3}, -1.02 \text{ GeV}^2, 1.25 \text{ GeV}) = 0.49(5),$$

$$\int_{-1}^1 dx H_T(x, 0, -0.69 \text{ GeV}^2) = \{0.69(4), 0.67(6), 0.84(10)\}, \quad \int_{-1}^1 dx H_T(x, \frac{1}{3}, -1.02 \text{ GeV}^2) = 0.45(4),$$

$$\int_{-1}^1 dx x H_T(x, 0, -0.69 \text{ GeV}^2) = \{0.20(2), 0.21(2), 0.24(3)\}, \quad \int_{-1}^1 dx x H_T(x, \frac{1}{3}, -1.02 \text{ GeV}^2) = 0.15(2).$$

$$A_{T10}(-0.69 \text{ GeV}^2) = \{0.65(4), 0.65(6), 0.82(10)\}, \quad A_{T10}(-1.02 \text{ GeV}^2) = 0.49(5)$$

- lowest moments the same between quasi-GPDs and GPDs
- Values of moments decrease as t increases
- Higher moments suppressed compared to the lowest

New parametrization of GPDs

PHYSICAL REVIEW D **106**, 114512 (2022)

Generalized parton distributions from lattice QCD with asymmetric momentum transfer: Unpolarized quarks

Shohini Bhattacharya^{1,*}, Krzysztof Cichy², Martha Constantinou^{3,†}, Jack Dodson³, Xiang Gao⁴, Andreas Metz³,
Swagato Mukherjee¹, Aurora Scapellato³, Fernanda Steffens⁵ and Yong Zhao⁴

Parameters of calculations



★ $N_f=2+1+1$ twisted mass fermions with a clover term;

[Extended Twisted Mass Collaboration, Phys. Rev. D 104, 074515 (2021), arXiv:2104.13408]

| Name | β | N_f | $L^3 \times T$ | a [fm] | M_π | $m_\pi L$ |
|----------|---------|--------------|------------------|----------|---------|-----------|
| cA211.32 | 1.726 | u, d, s, c | $32^3 \times 64$ | 0.093 | 260 MeV | 4 |

| frame | P_3 [GeV] | Δ [$\frac{2\pi}{L}$] | $-t$ [GeV ²] | ξ | N_{ME} | N_{confs} | N_{src} | N_{tot} |
|-------|-------------|--|--------------------------|-------|----------|-------------|-----------|-----------|
| N/A | ± 1.25 | (0,0,0) | 0 | 0 | 2 | 731 | 16 | 23392 |
| symm | ± 0.83 | ($\pm 2, 0, 0$), ($0, \pm 2, 0$) | 0.69 | 0 | 8 | 67 | 8 | 4288 |
| symm | ± 1.25 | ($\pm 2, 0, 0$), ($0, \pm 2, 0$) | 0.69 | 0 | 8 | 249 | 8 | 15936 |
| symm | ± 1.67 | ($\pm 2, 0, 0$), ($0, \pm 2, 0$) | 0.69 | 0 | 8 | 294 | 32 | 75264 |
| symm | ± 1.25 | ($\pm 2, \pm 2, 0$) | 1.39 | 0 | 16 | 224 | 8 | 28672 |
| symm | ± 1.25 | ($\pm 4, 0, 0$), ($0, \pm 4, 0$) | 2.76 | 0 | 8 | 329 | 32 | 84224 |
| asymm | ± 1.25 | ($\pm 1, 0, 0$), ($0, \pm 1, 0$) | 0.17 | 0 | 8 | 429 | 8 | 27456 |
| asymm | ± 1.25 | ($\pm 1, \pm 1, 0$) | 0.33 | 0 | 16 | 194 | 8 | 12416 |
| asymm | ± 1.25 | ($\pm 2, 0, 0$), ($0, \pm 2, 0$) | 0.64 | 0 | 8 | 429 | 8 | 27456 |
| asymm | ± 1.25 | ($\pm 1, \pm 2, 0$), ($\pm 2, \pm 1, 0$) | 0.80 | 0 | 16 | 194 | 8 | 12416 |
| asymm | ± 1.25 | ($\pm 2, \pm 2, 0$) | 1.16 | 0 | 16 | 194 | 8 | 24832 |
| asymm | ± 1.25 | ($\pm 3, 0, 0$), ($0, \pm 3, 0$) | 1.37 | 0 | 8 | 429 | 8 | 27456 |
| asymm | ± 1.25 | ($\pm 1, \pm 3, 0$), ($\pm 3, \pm 1, 0$) | 1.50 | 0 | 16 | 194 | 8 | 12416 |
| asymm | ± 1.25 | ($\pm 4, 0, 0$), ($0, \pm 4, 0$) | 2.26 | 0 | 8 | 429 | 8 | 27456 |

Symmetric frame
very expensive
computationally

Collaboration

Theoretical setup

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Goals

- (A) A_i are related to the standard H, E GPDs $\mathcal{H}_0^s(A_i^s; z) = A_1 + \frac{z(\Delta_1^2 + \Delta_2^2)}{2P_3} A_6$
- (B) Extraction of standard GPDs using A_i obtained from any frame
- (C) quasi-GPDs may be redefined (Lorentz covariant) inspired by light-cone:

Theoretical setup

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Goals

(A) A_i are related to the standard H, E GPDs $\mathcal{H}_0^s(A_i^s; z) = A_1 + \frac{z(\Delta_1^2 + \Delta_2^2)}{2P_3} A_6$

(B) Extraction of standard GPDs using A_i obtained from any frame

(C) quasi-GPDs may be redefined (Lorentz covariant) inspired by light-cone:

$$H(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = A_1 + \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3$$

$$E(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = -A_1 - \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3 + 2A_5 + 2P_{avg,sla} \cdot z A_6 + 2\Delta_{sla} \cdot z A_8$$

Theoretical setup

$$F_{\lambda,\lambda'}^\mu = \bar{u}(p', \lambda') \left[\frac{P^\mu}{M} A_1 + z^\mu M A_2 + \frac{\Delta^\mu}{M} A_3 + i\sigma^{\mu z} M A_4 + \frac{i\sigma^{\mu\Delta}}{M} A_5 + \frac{P^\mu i\sigma^{z\Delta}}{M} A_6 + \frac{z^\mu i\sigma^{z\Delta}}{M} A_7 + \frac{\Delta^\mu i\sigma^{z\Delta}}{M} A_8 \right] u(p, \lambda)$$

Goals

(A) A_i are related to the standard H, E GPDs $\mathcal{H}_0^s(A_i^s; z) = A_1 + \frac{z(\Delta_1^2 + \Delta_2^2)}{2P_3} A_6$

(B) Extraction of standard GPDs using A_i obtained from any frame

(C) quasi-GPDs may be redefined (Lorentz covariant) inspired by light-cone:

$$H(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = A_1 + \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3$$

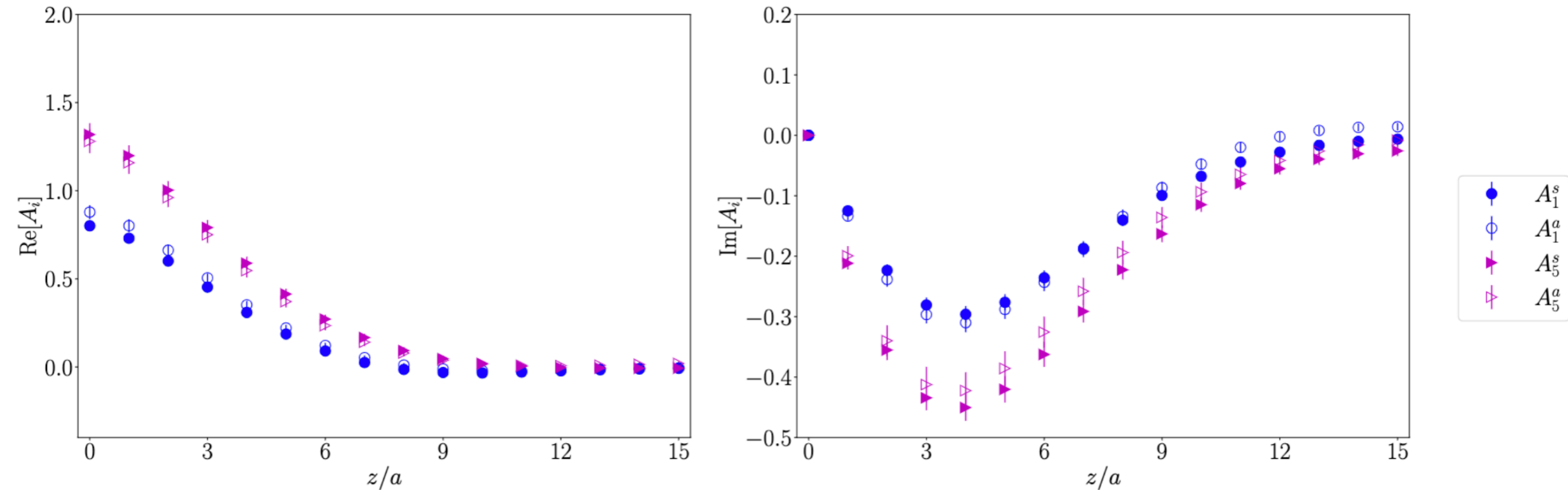
$$E(z \cdot P, z \cdot \Delta, t = \Delta^2, z^2) = -A_1 - \frac{\Delta_{sla} \cdot z}{P_{avg,sla} \cdot z} A_3 + 2A_5 + 2P_{avg,sla} \cdot z A_6 + 2\Delta_{sla} \cdot z A_8$$

(A) Proof-of-concept calculation ($\xi = 0$):

- symmetric frame: $\vec{p}_f^s = \vec{P} + \frac{\vec{Q}}{2}, \quad \vec{p}_i^s = \vec{P} - \frac{\vec{Q}}{2} \quad -t^s = \vec{Q}^2 = 0.69 \text{ GeV}^2$

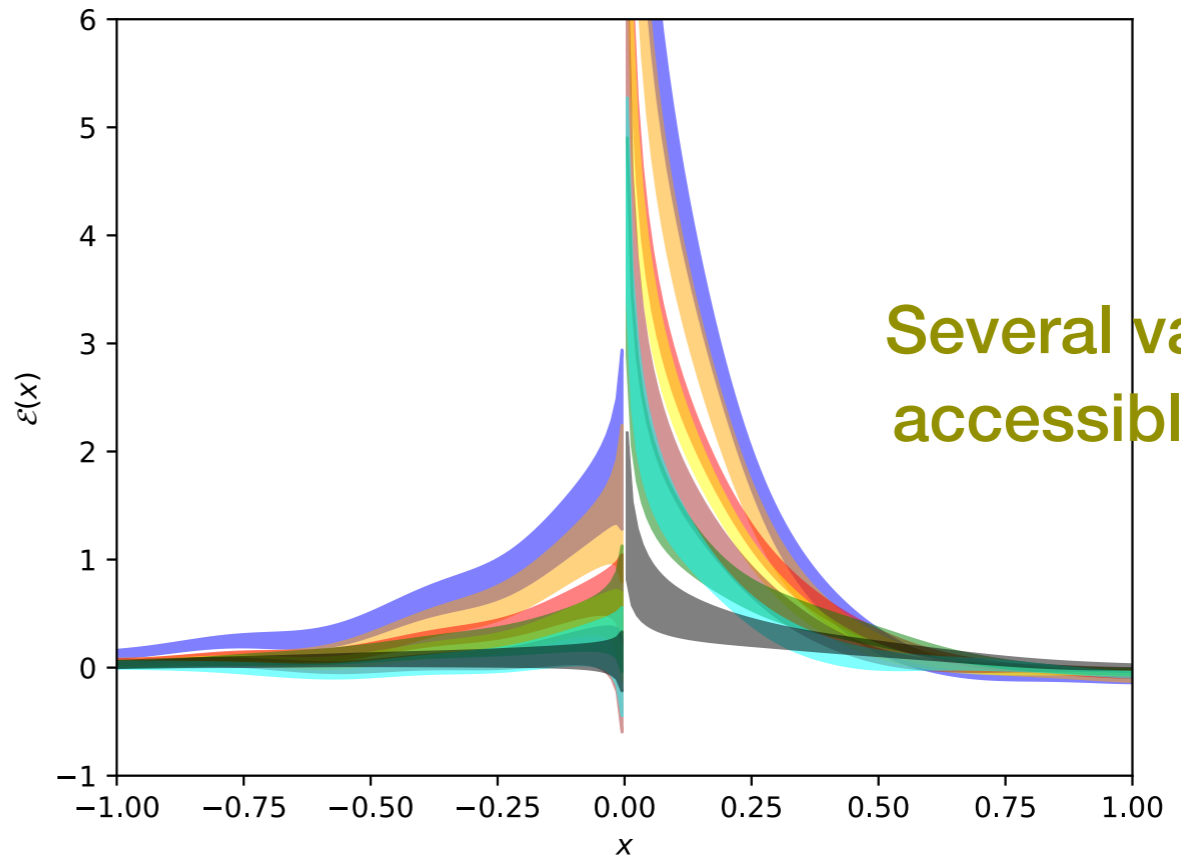
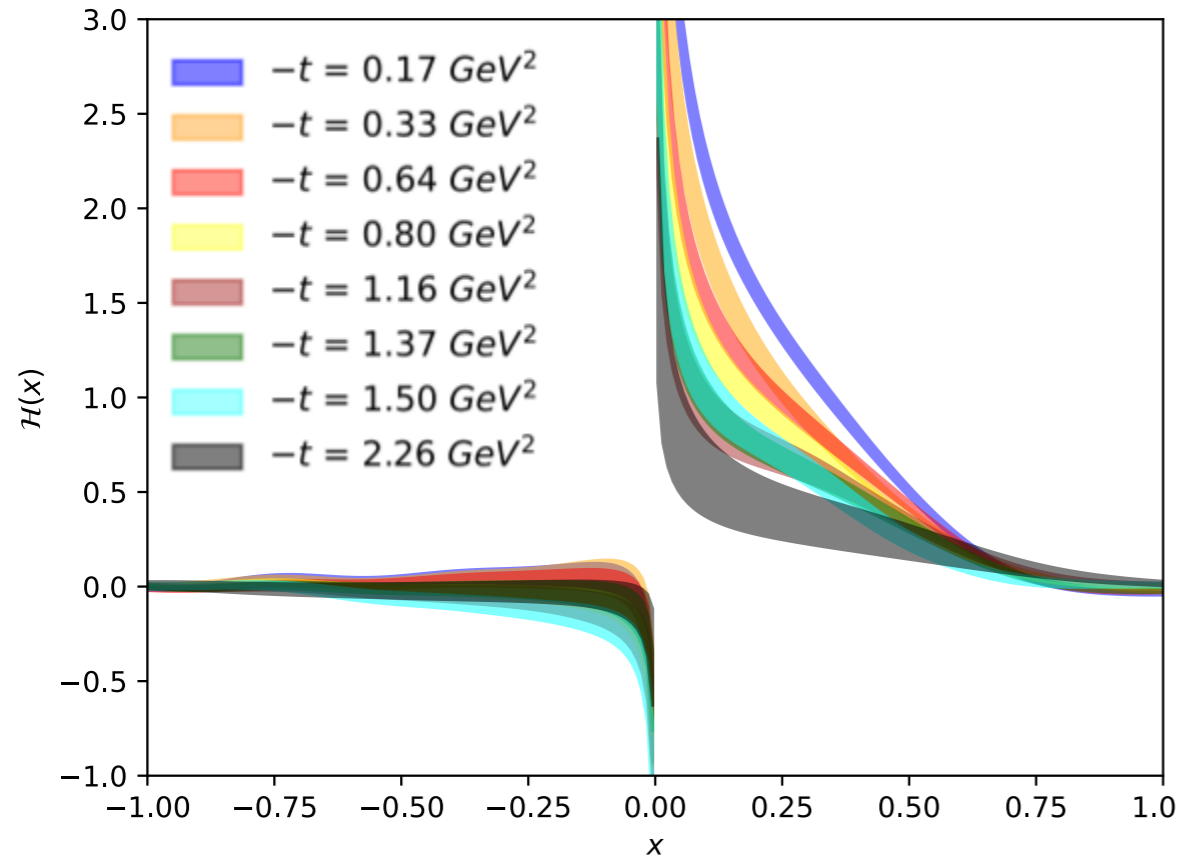
- asymmetric frame: $\vec{p}_f^a = \vec{P}, \quad \vec{p}_i^a = \vec{P} - \vec{Q} \quad t^a = -\vec{Q}^2 + (E_f - E_i)^2 = 0.65 \text{ GeV}^2$

Comparison of A_i in two frames



- ★ A_1, A_5 dominant contributions
- ★ Full agreement in two frames for both Re and Im parts of A_1, A_5
- ★ A_3, A_4, A_8 zero at $\xi = 0$
- ★ A_2, A_6, A_7 suppressed (at least for this kinematic setup and $\xi = 0$)

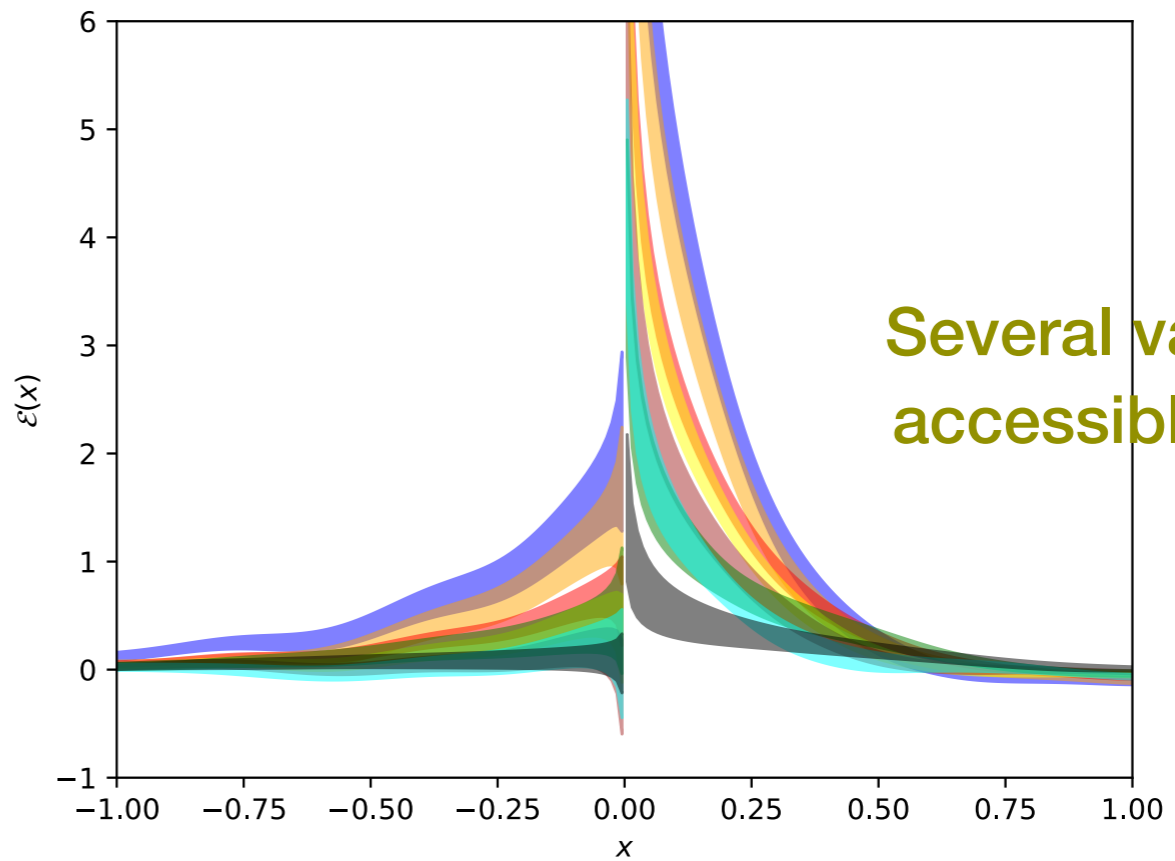
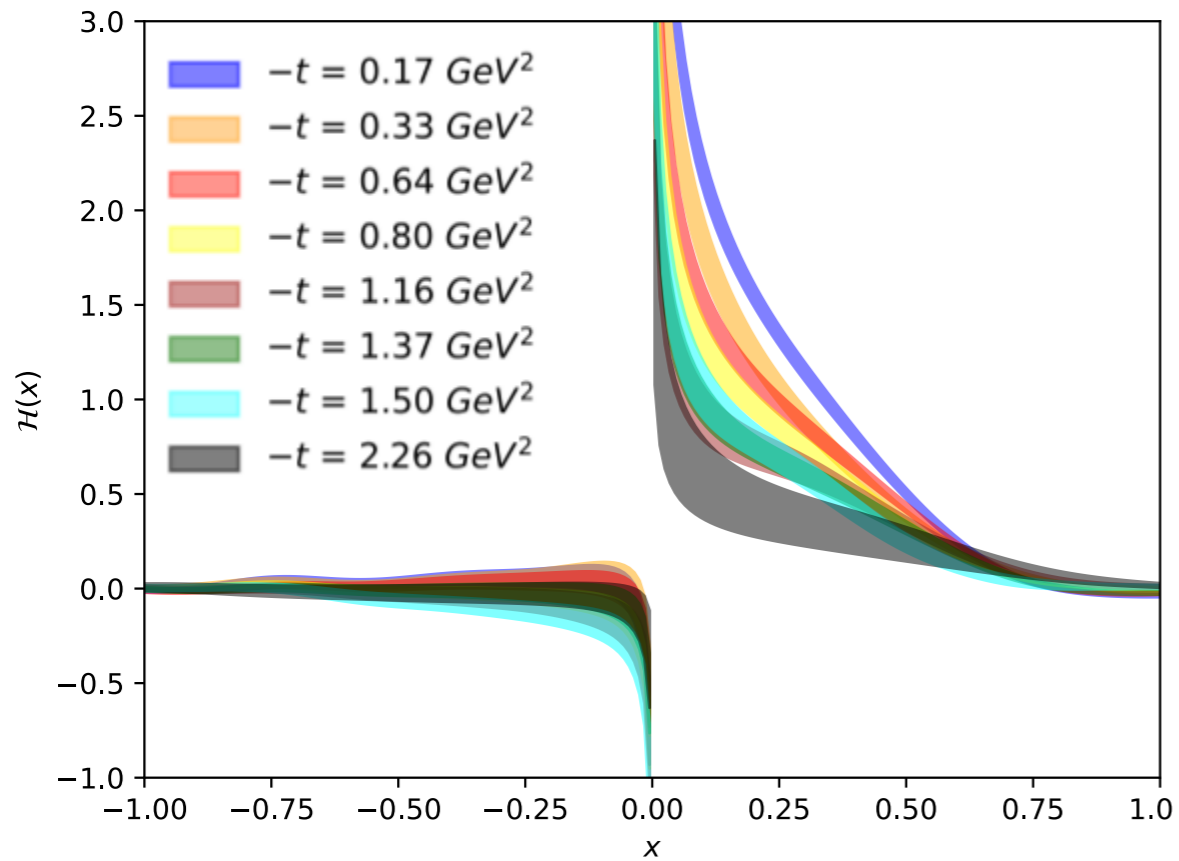
H, E light-cone GPDs



Several values of $-t$ accessible at once

★ Anti-quark region susceptible to systematic uncertainties

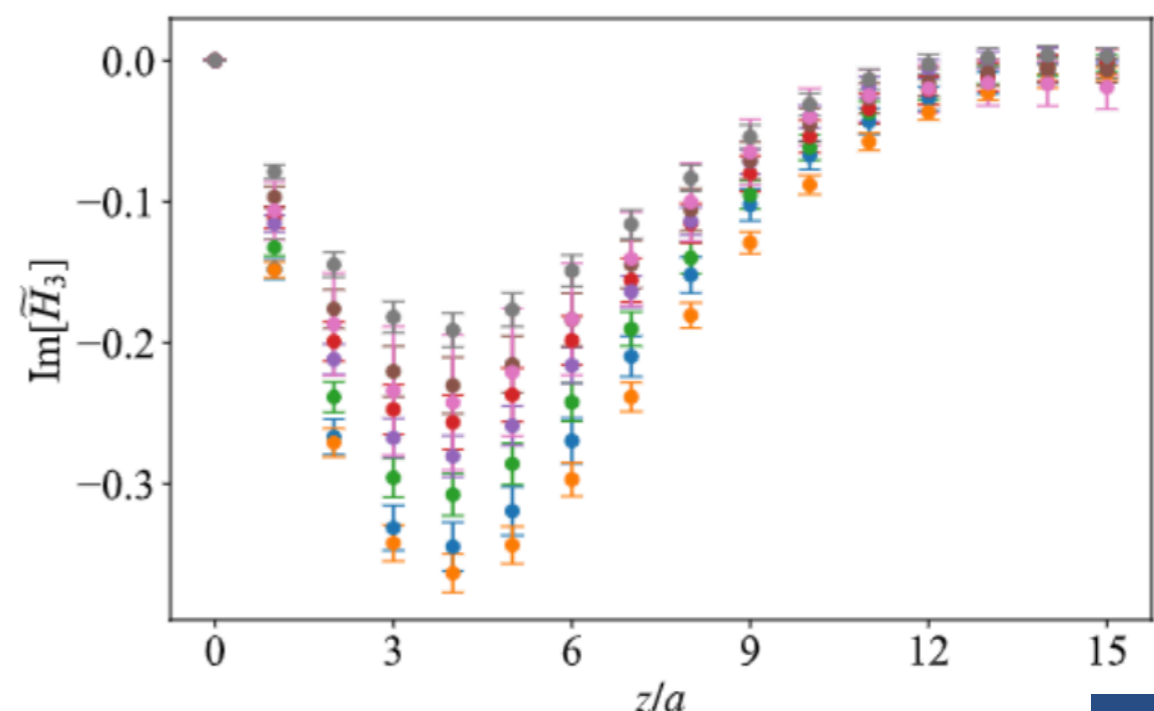
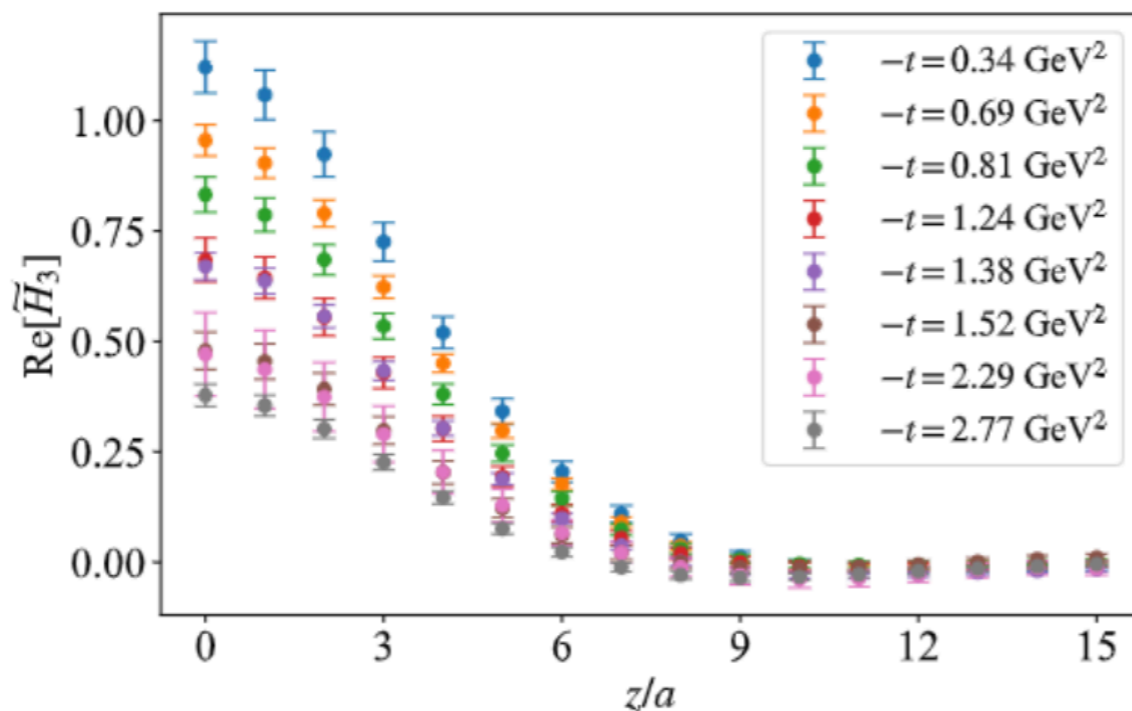
H, E light-cone GPDs



Several values of $-t$ accessible at once

★ Anti-quark region susceptible to systematic uncertainties

Similar analysis: helicity GPDs



Exploration of twist-3 PDFs & GPDs

PHYSICAL REVIEW D **102**, 111501(R) (2020)

Rapid Communications

Editors' Suggestion

Insights on proton structure from lattice QCD: The twist-3 parton distribution function $g_T(x)$

Shohini Bhattacharya¹, Krzysztof Cichy², Martha Constantinou¹, Andreas Metz¹,
Aurora Scapellato² and Fernanda Steffens³

PHYSICAL REVIEW D **104**, 114510 (2021)

Parton distribution functions beyond leading twist from lattice QCD: The $h_L(x)$ case

Shohini Bhattacharya¹, Krzysztof Cichy², Martha Constantinou¹, Andreas Metz¹,
Aurora Scapellato¹ and Fernanda Steffens³

arXiv:2306.05533v1 [hep-lat] 8 Jun 2023

Chiral-even axial twist-3 GPDs of the proton from lattice QCD

Shohini Bhattacharya^{1,2}, Krzysztof Cichy³, Martha Constantinou¹,
Jack Dodson¹, Andreas Metz¹, Aurora Scapellato¹, Fernanda Steffens⁴

h_L twist-3 PDF

- ★ h_L decouples from usual DIS (chiral odd) - not trivial to extract
- ★ most elusive of the three twist-3 PDFs ($e(x)$, $g_T(x)$, $h_L(x)$)
- ★ h_L may be accessed via:
 - double-polarized Drell-Yan process
[R. Jaffe, PRL 67 (1991) 552-555; Y. Koike et al., PLB 668 (2008) 286, arXiv:0805.2289]
 - di-hadron single spin asymmetries (CLAS)
[Gliske et al., PRD 90 (2014) 11, 114027, arXiv:1408.5721; A. Vossen, CIPANP 2018, arXiv: 1810.02435]
 - single-inclusive particle production in proton-proton collisions
[Y. Koike et al., PLB 759 (2016) 75, arXiv:1603.07908]

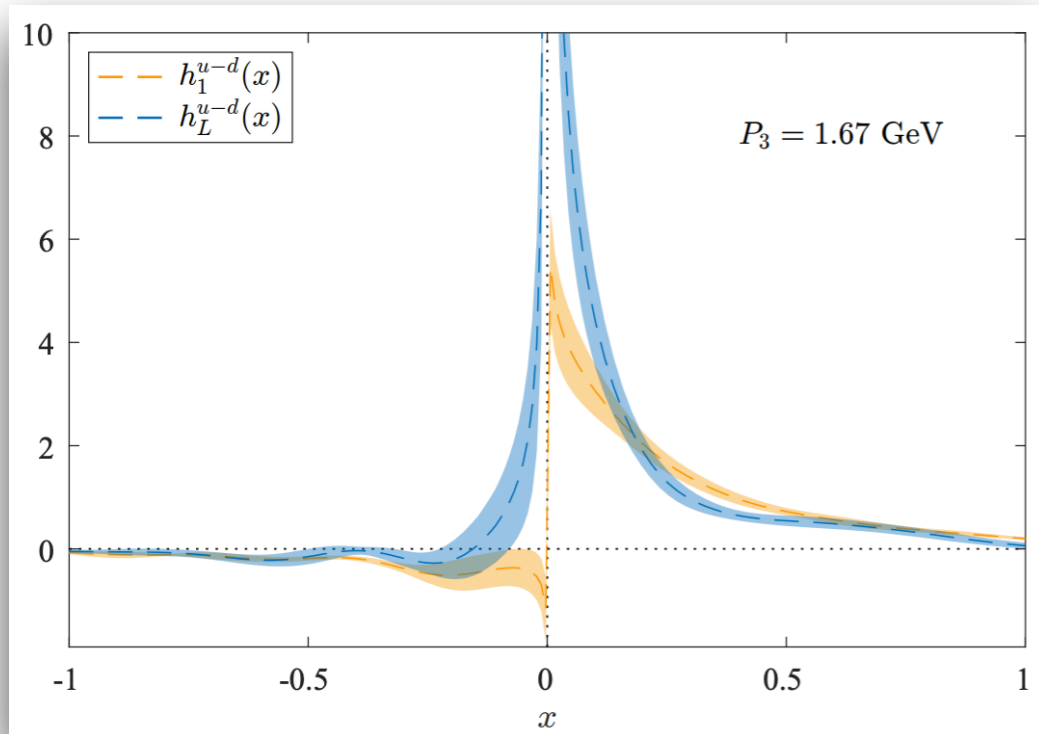
h_L twist-3 PDF

- ★ h_L decouples from usual DIS (chiral odd) - not trivial to extract
- ★ most elusive of the three twist-3 PDFs ($e(x)$, $g_T(x)$, $h_L(x)$)
- ★ h_L may be accessed via:
 - double-polarized Drell-Yan process
[R. Jaffe, PRL 67 (1991) 552-555; Y. Koike et al., PLB 668 (2008) 286, arXiv:0805.2289]
 - di-hadron single spin asymmetries (CLAS)
[Gliske et al., PRD 90 (2014) 11, 114027, arXiv:1408.5721; A. Vossen, CIPANP 2018, arXiv: 1810.02435]
 - single-inclusive particle production in proton-proton collisions
[Y. Koike et al., PLB 759 (2016) 75, arXiv:1603.07908]

- ★ Parametrization [Meissner et al., JHEP 08 (2009) 056]

$$F^{[\sigma^{+-}\gamma_5]} = \bar{u}(p') \left(\gamma^+ \gamma_5 \tilde{H}'_2 + \frac{P^+ \gamma_5}{M} \tilde{E}'_2 \right) u(p)$$

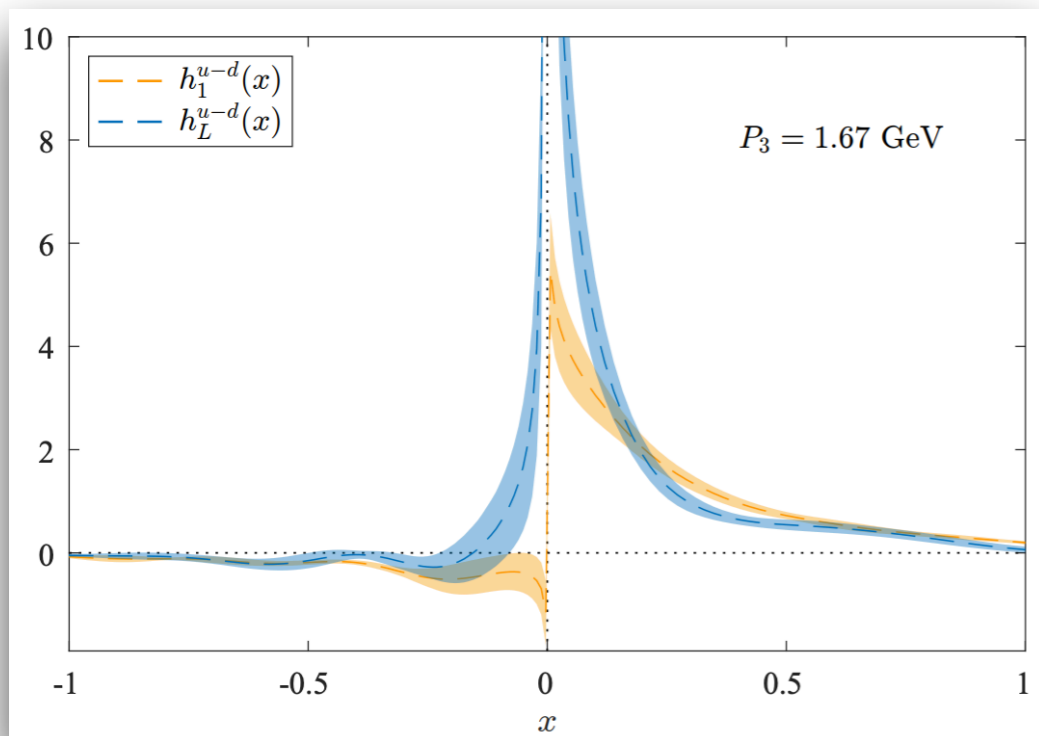
First lattice calculation of x-dependent GPDs



★ Twist-3 h_L as sizable as twist-2 h_1

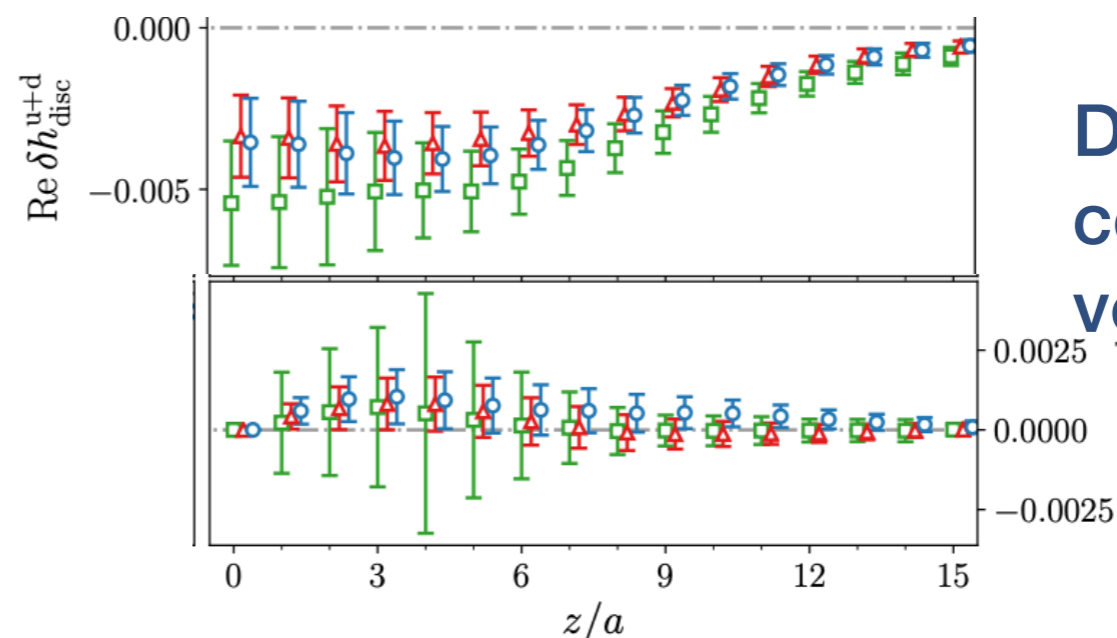
$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$

First lattice calculation of x-dependent GPDs



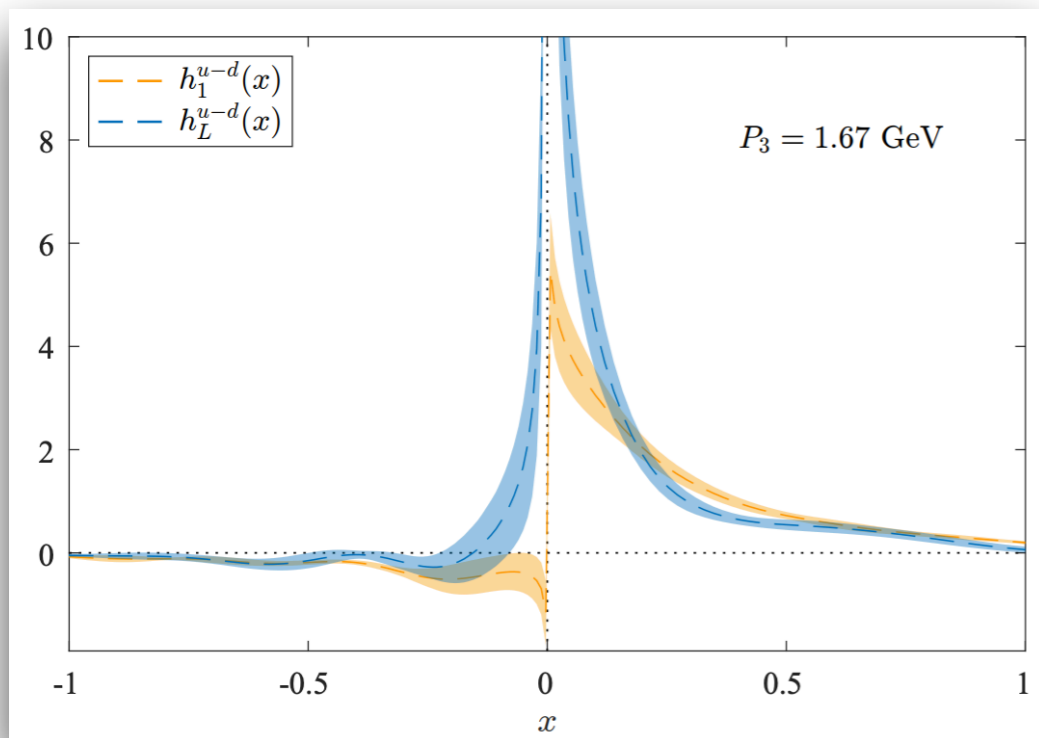
★ Twist-3 h_L as sizable as twist-2 h_1

$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$



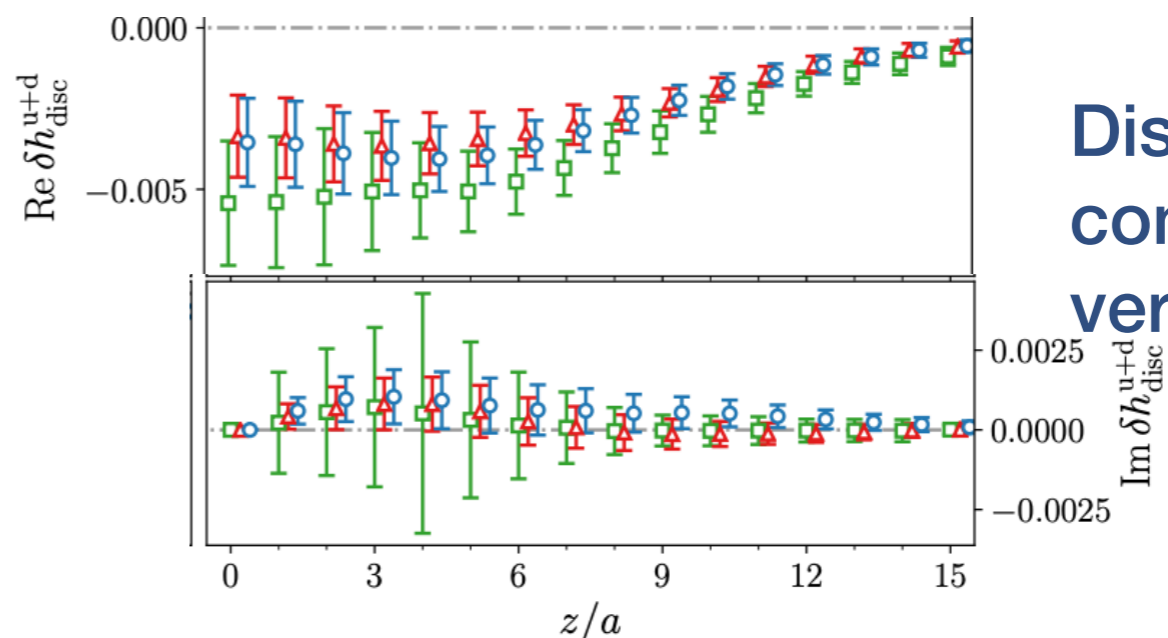
Disconnected contributions very small

First lattice calculation of x-dependent GPDs

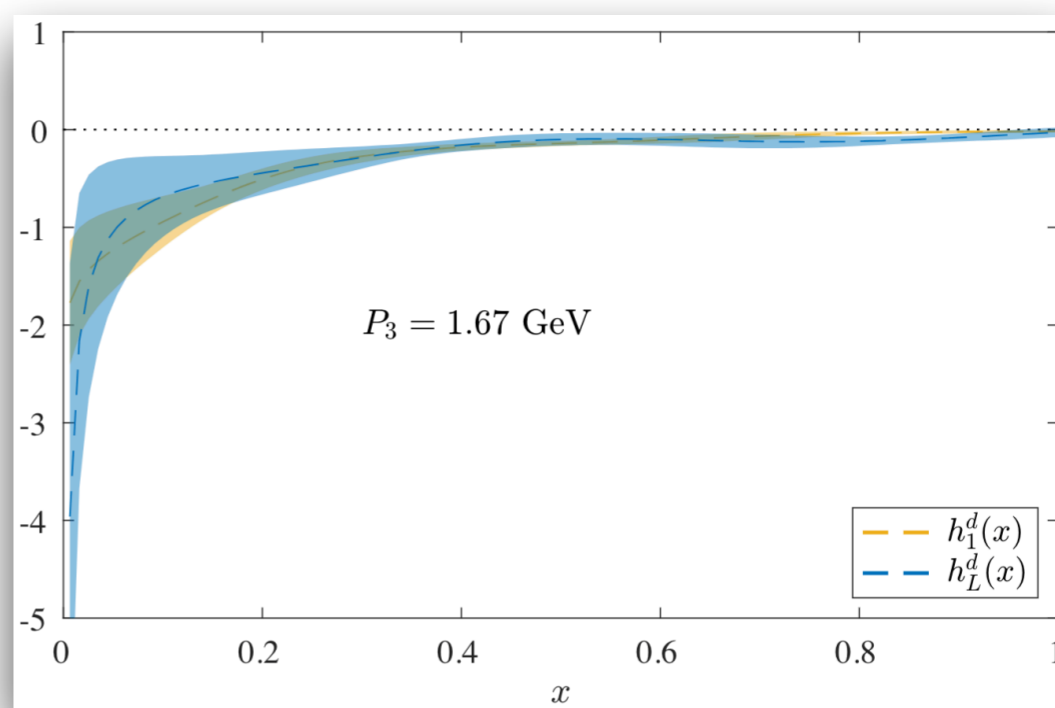
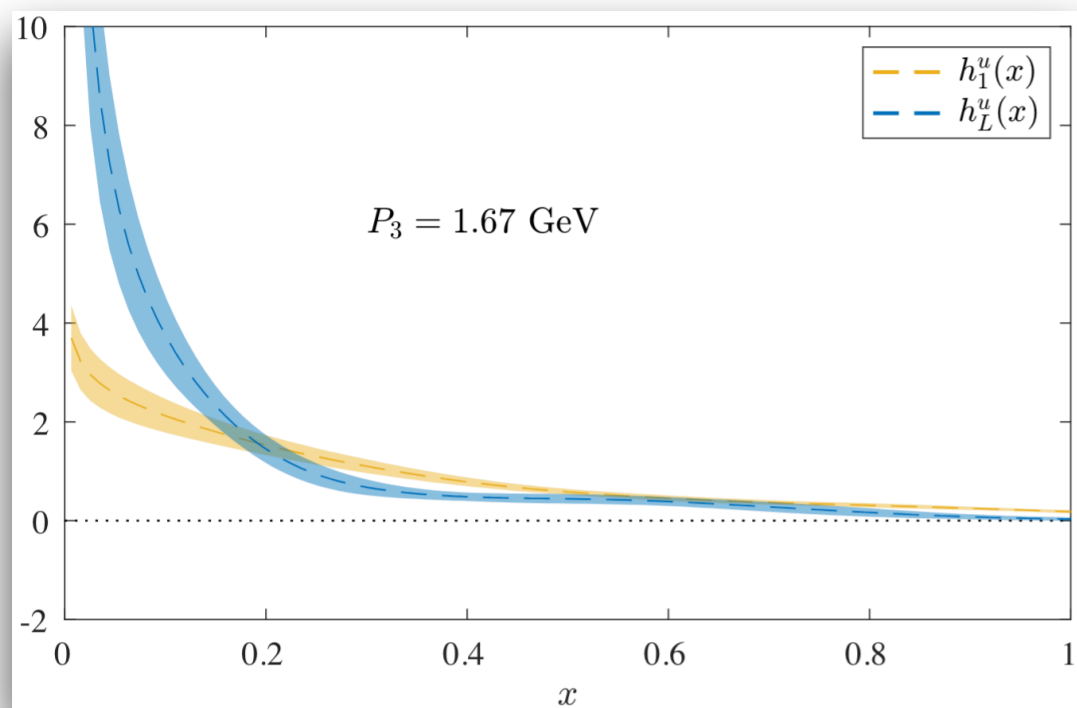


★ Twist-3 h_L as sizable as twist-2 h_1

$$f_i = f_i^{(0)} + \frac{f_i^{(1)}}{Q} + \frac{f_i^{(2)}}{Q^2} \dots$$



★ Role of up and down quarks:



★ Up-quarks dominant in both h_L, h_1

x-dependence: $h_{1,L}^u \sim 2h_{1,L}^d$

Twist-3 $h_L(x)$ PDF

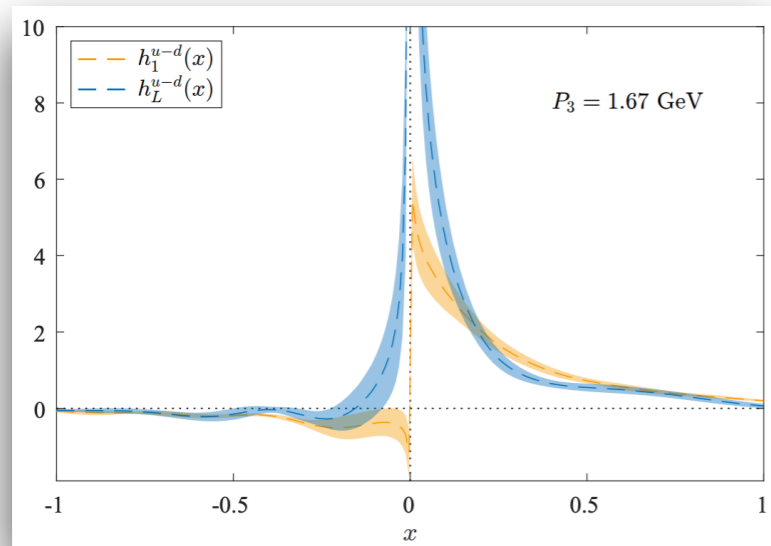
Burkhardt-Cottingham (quasi- h_L)

$$\int_{-1}^1 dx \tilde{h}_L(x, P_3) = \int_{-1}^1 dx \tilde{h}_1(x, P_3) = g_T$$

Twist-3 $h_L(x)$ PDF

Burkhardt-Cottingham (quasi- h_L)

$$\int_{-1}^1 dx \tilde{h}_L(x, P_3) = \int_{-1}^1 dx \tilde{h}_1(x, P_3) = g_T$$

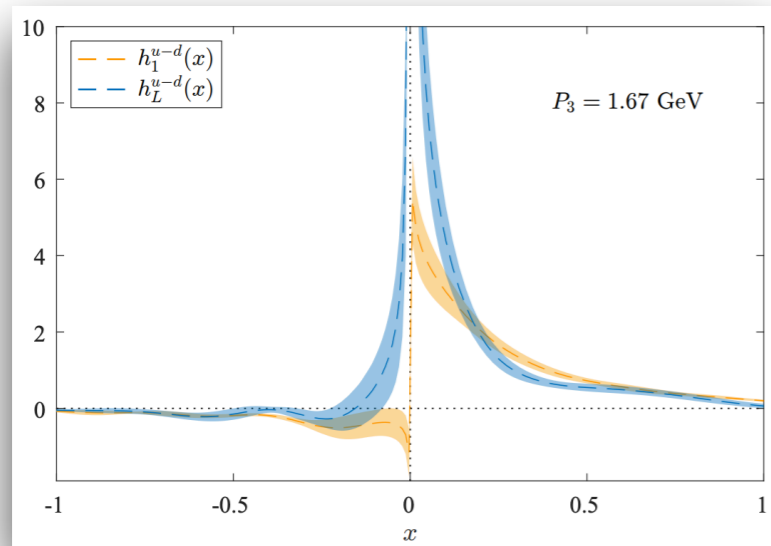


$$\int dx \tilde{h}_L(x, 1.67 \text{ GeV}) = 1.03(16), \quad \int dx \tilde{h}_1(x, 1.67 \text{ GeV}) = 0.94(10)$$

Twist-3 $h_L(x)$ PDF

Burkhardt-Cottingham (quasi- h_L)

$$\int_{-1}^1 dx \tilde{h}_L(x, P_3) = \int_{-1}^1 dx \tilde{h}_1(x, P_3) = g_T$$



Wandzura-Wilczek approximation

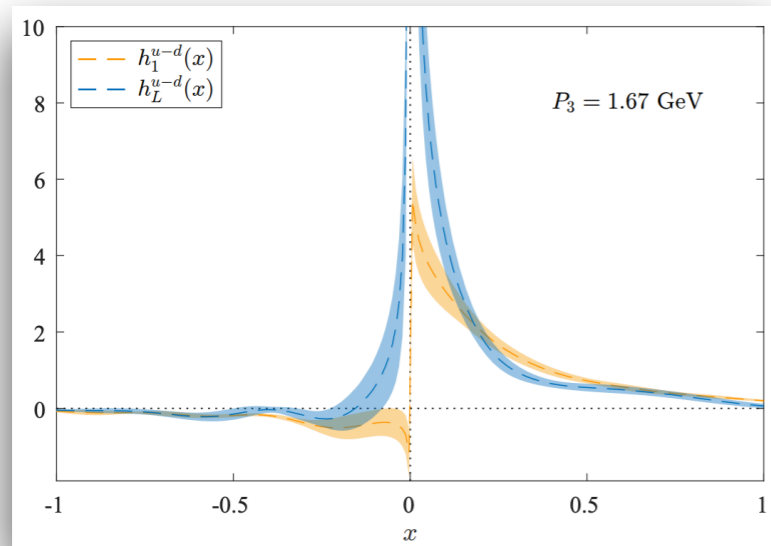
$$h_L(x) = h_L^{\text{WW}}(x) + h_L^{\text{twist-3}}(x) = 2x \int_x^1 dy \frac{h_1(y)}{y^2} + h_L^{\text{twist-3}}(x)$$

$$\int dx \tilde{h}_L(x, 1.67 \text{ GeV}) = 1.03(16), \quad \int dx \tilde{h}_1(x, 1.67 \text{ GeV}) = 0.94(10)$$

Twist-3 $h_L(x)$ PDF

Burkhardt-Cottingham (quasi- h_L)

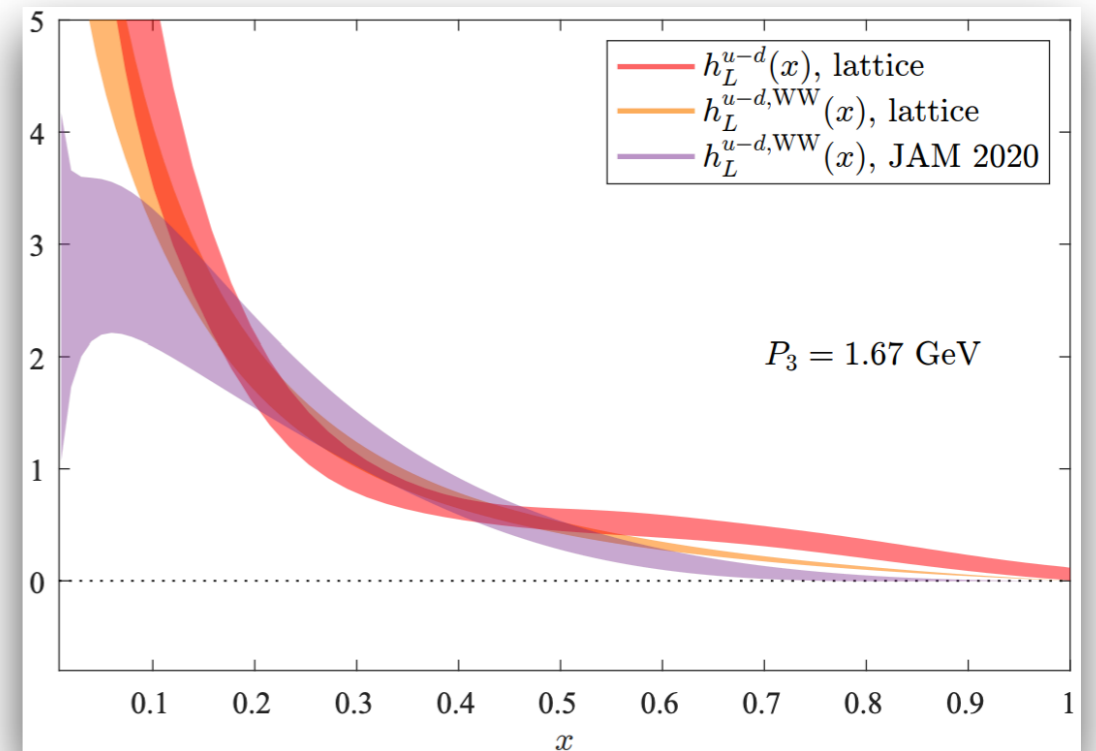
$$\int_{-1}^1 dx \tilde{h}_L(x, P_3) = \int_{-1}^1 dx \tilde{h}_1(x, P_3) = g_T$$



$$\int dx \tilde{h}_L(x, 1.67 \text{ GeV}) = 1.03(16), \quad \int dx \tilde{h}_1(x, 1.67 \text{ GeV}) = 0.94(10)$$

Wandzura-Wilczek approximation

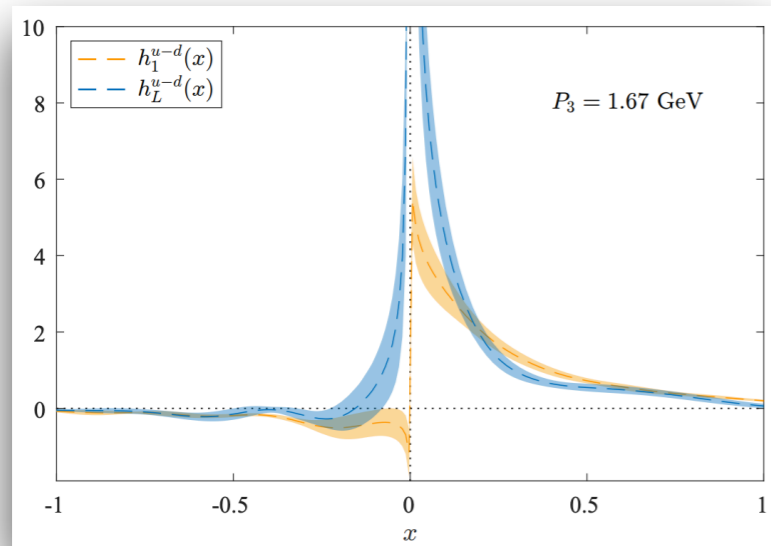
$$h_L(x) = h_L^{\text{WW}}(x) + h_L^{\text{twist-3}}(x) = 2x \int_x^1 dy \frac{h_1(y)}{y^2} + h_L^{\text{twist-3}}(x)$$



Twist-3 $h_L(x)$ PDF

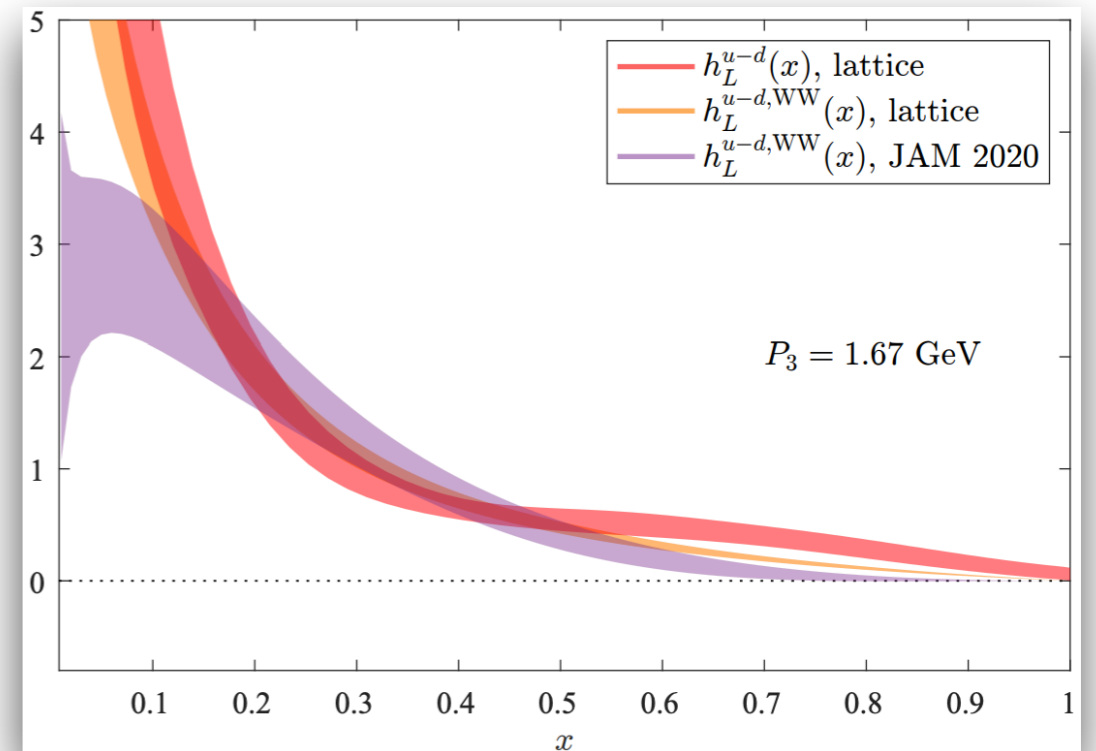
Burkhardt-Cottingham (quasi- h_L)

$$\int_{-1}^1 dx \tilde{h}_L(x, P_3) = \int_{-1}^1 dx \tilde{h}_1(x, P_3) = g_T$$

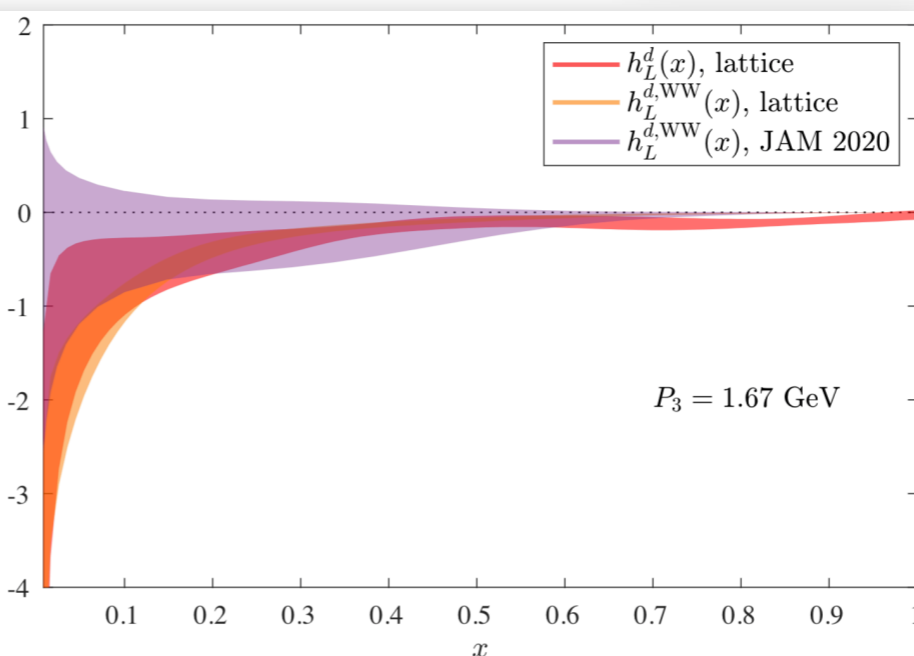
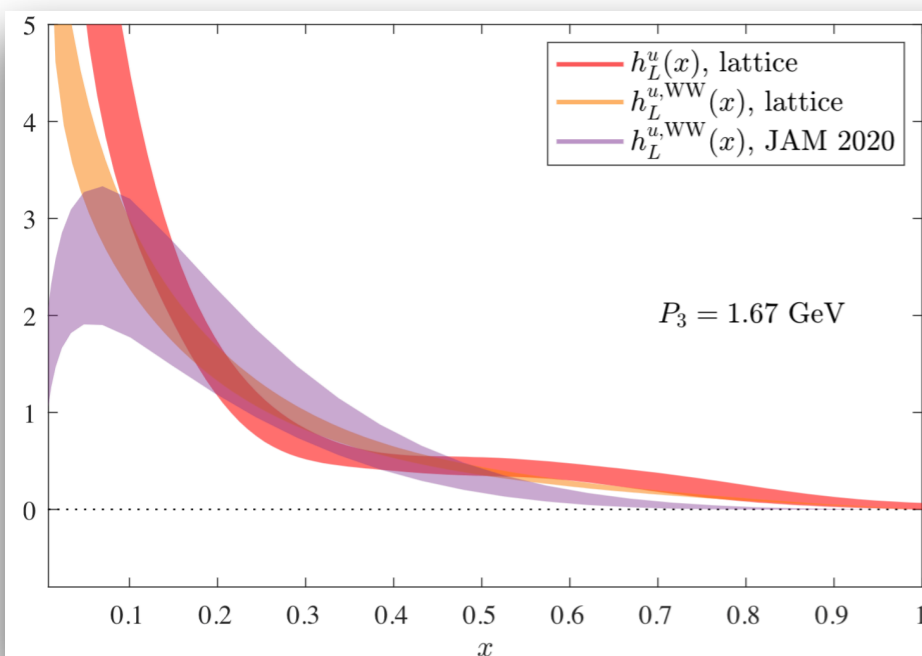


Wandzura-Wilczek approximation

$$h_L(x) = h_L^{\text{WW}}(x) + h_L^{\text{twist-3}}(x) = 2x \int_x^1 dy \frac{h_1(y)}{y^2} + h_L^{\text{twist-3}}(x)$$



$$\int dx \tilde{h}_L(x, 1.67 \text{ GeV}) = 1.03(16), \quad \int dx \tilde{h}_1(x, 1.67 \text{ GeV}) = 0.94(10)$$



- h_L^u dominant - diff. between h_L & h_L^{WW}
- $h_L^d < 0$ and decays faster than h_L^u

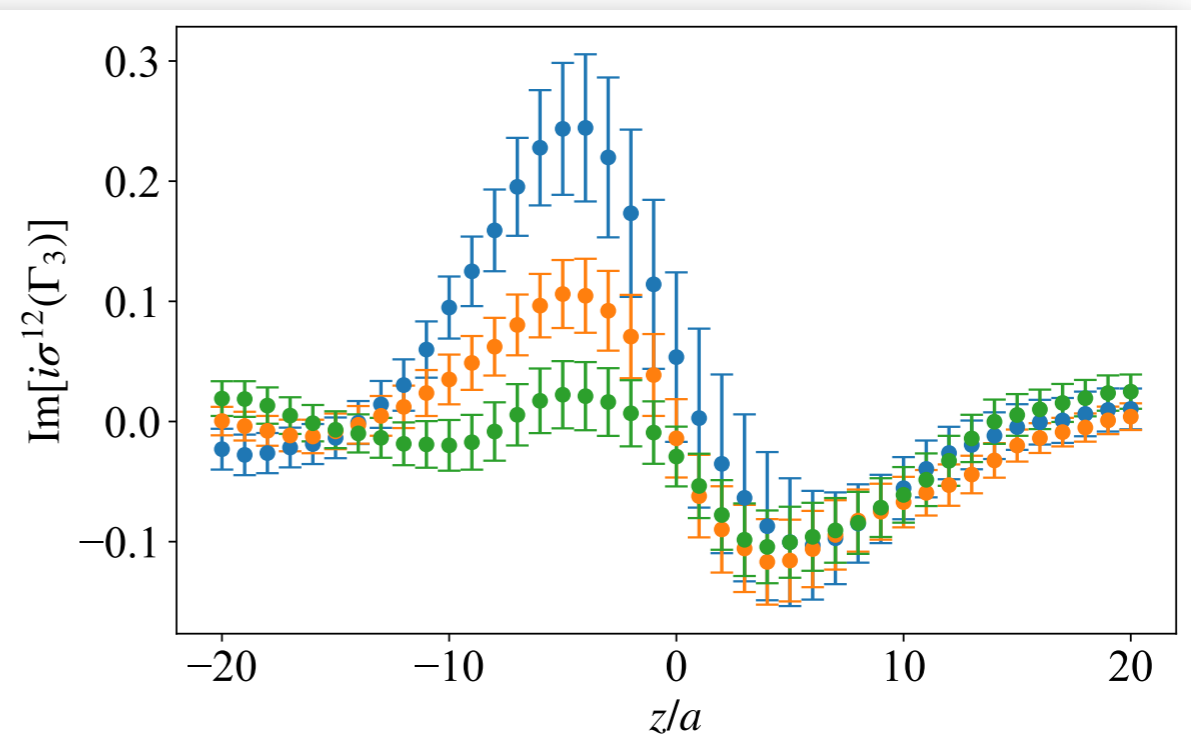
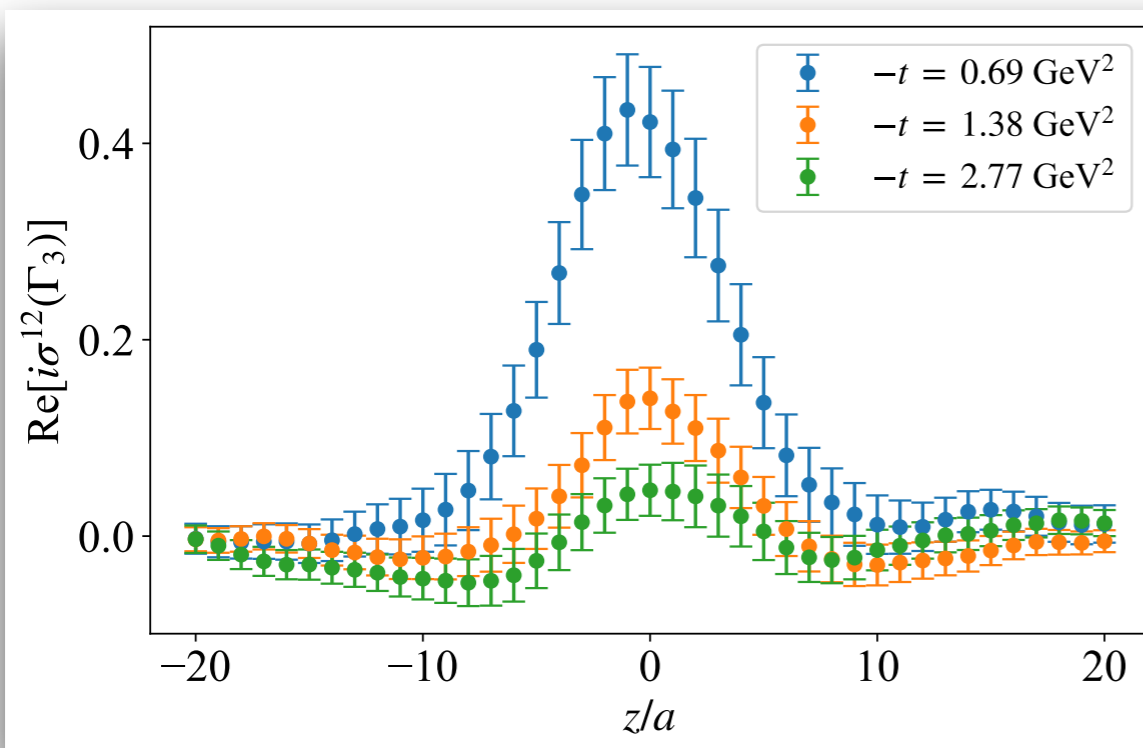
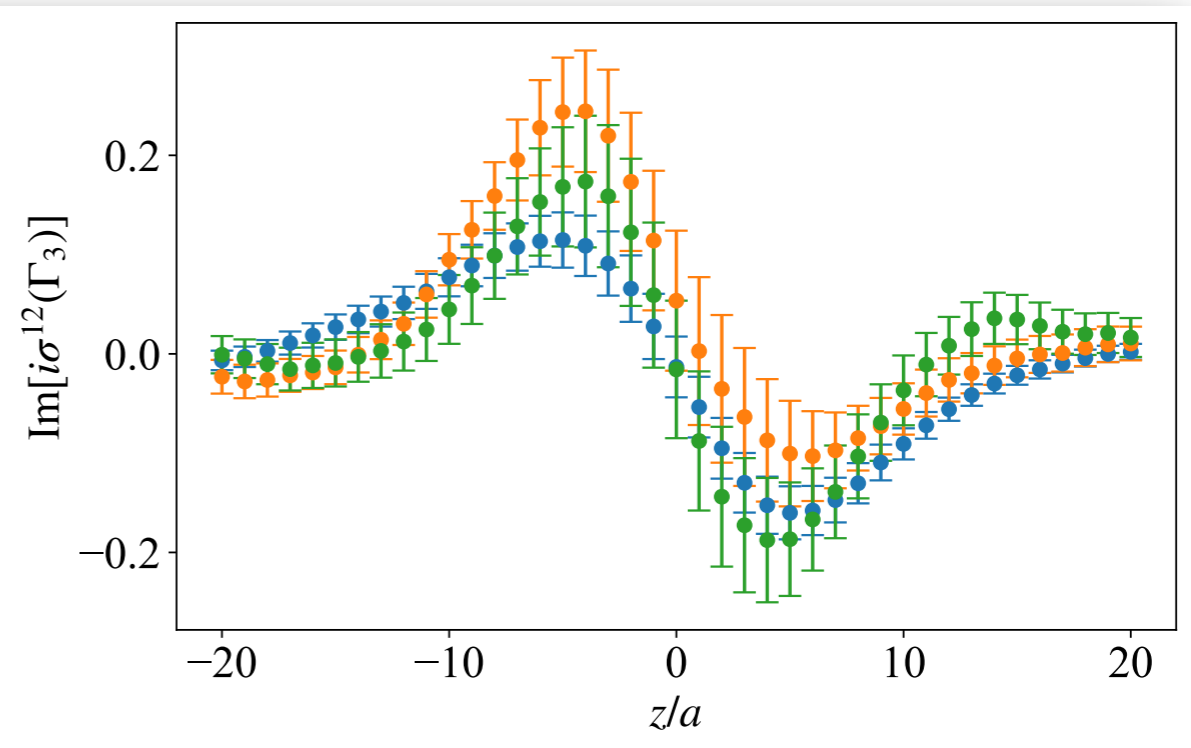
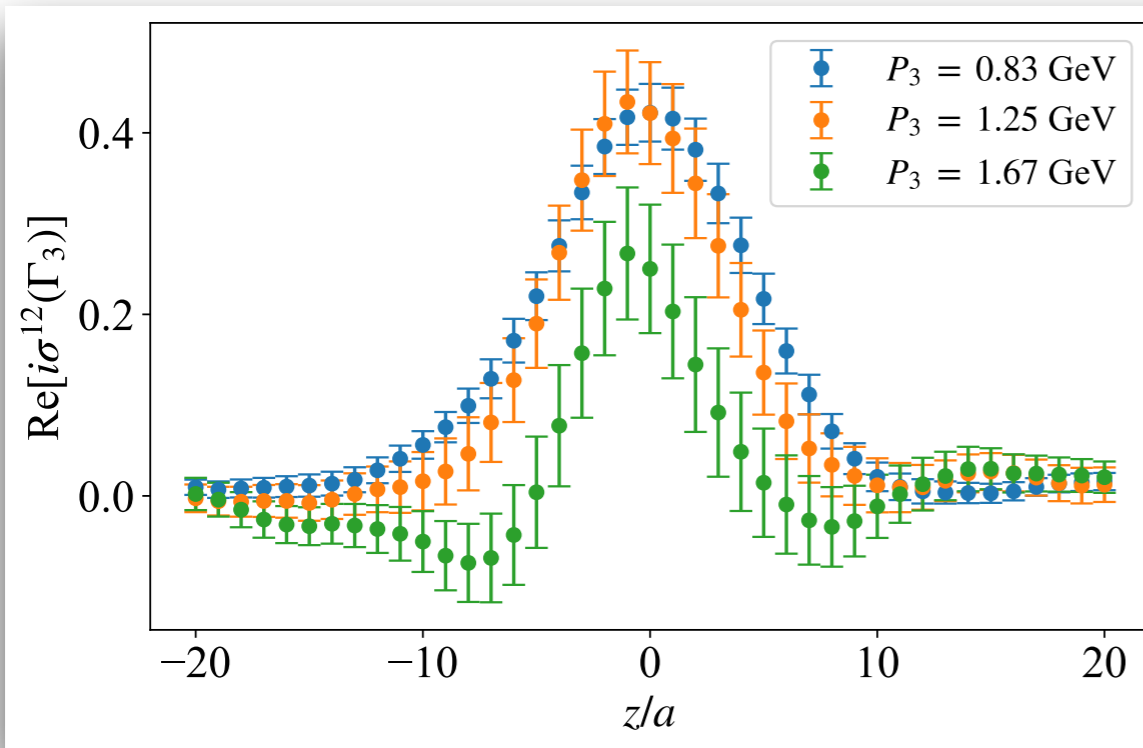
Extension to twist-3 tensor GPDs

arXiv:2306.05533v1 [hep-lat] 8 Jun 2023

Chiral-even axial twist-3 GPDs of the proton from lattice QCD

Shohini Bhattacharya^{1,2}, Krzysztof Cichy³, Martha Constantinou¹,
Jack Dodson¹, Andreas Metz¹, Aurora Scapellato¹, Fernanda Steffens⁴

Extension to twist-3 tensor GPDs



How to lattice QCD data fit into the overall effort for hadron tomography

How to lattice QCD data fit into the overall effort for hadron tomography

- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence

How to lattice QCD data fit into the overall effort for hadron tomography

- ★ Lattice data may be incorporated in global analysis of experimental data and may influence parametrization of t and ξ dependence



QUARK-GLUON TOMOGRAPHY COLLABORATION



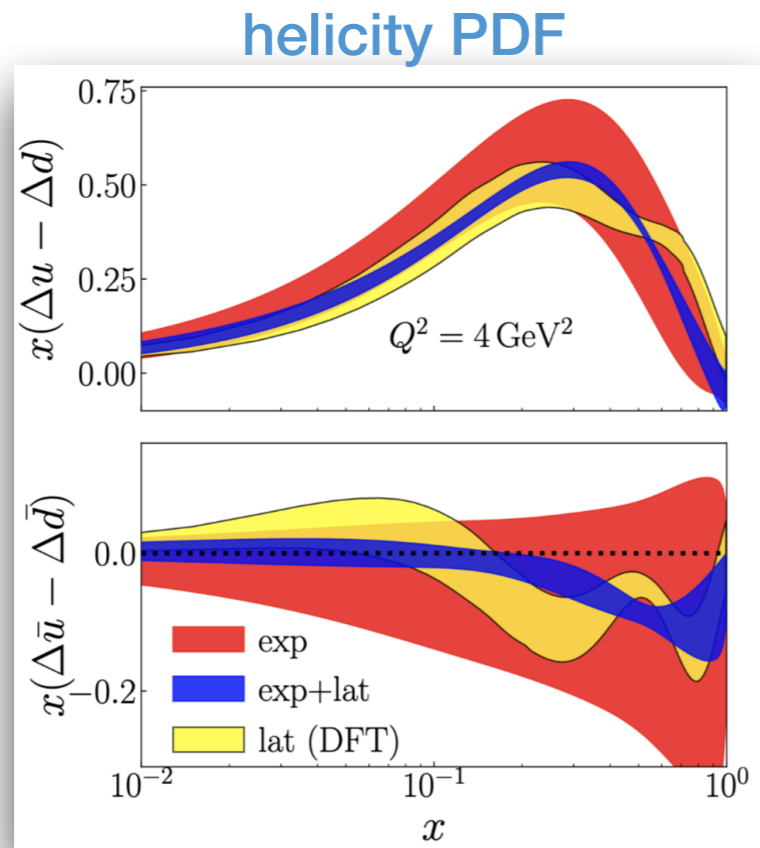
U.S. DEPARTMENT OF
ENERGY

Office of
Science

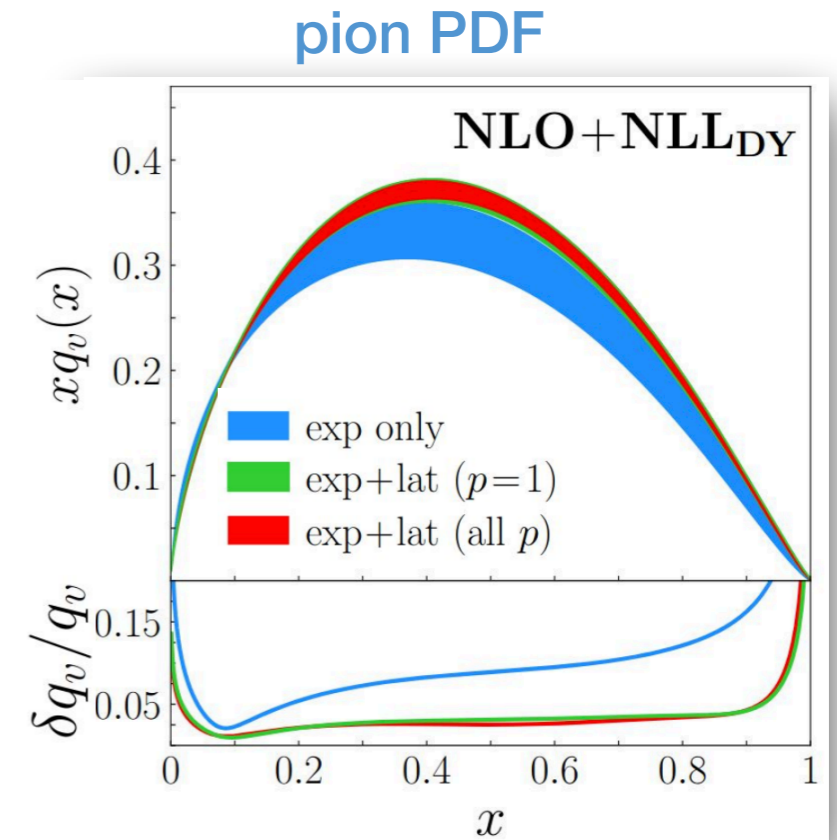
Award Number:
DE-SC0023646

1. **Theoretical studies** of high-momentum transfer processes using perturbative QCD methods and study of GPDs properties
2. **Lattice QCD** calculations of GPDs and related structures
3. **Global analysis** of GPDs based on experimental data using modern data analysis techniques for inference and uncertainty quantification

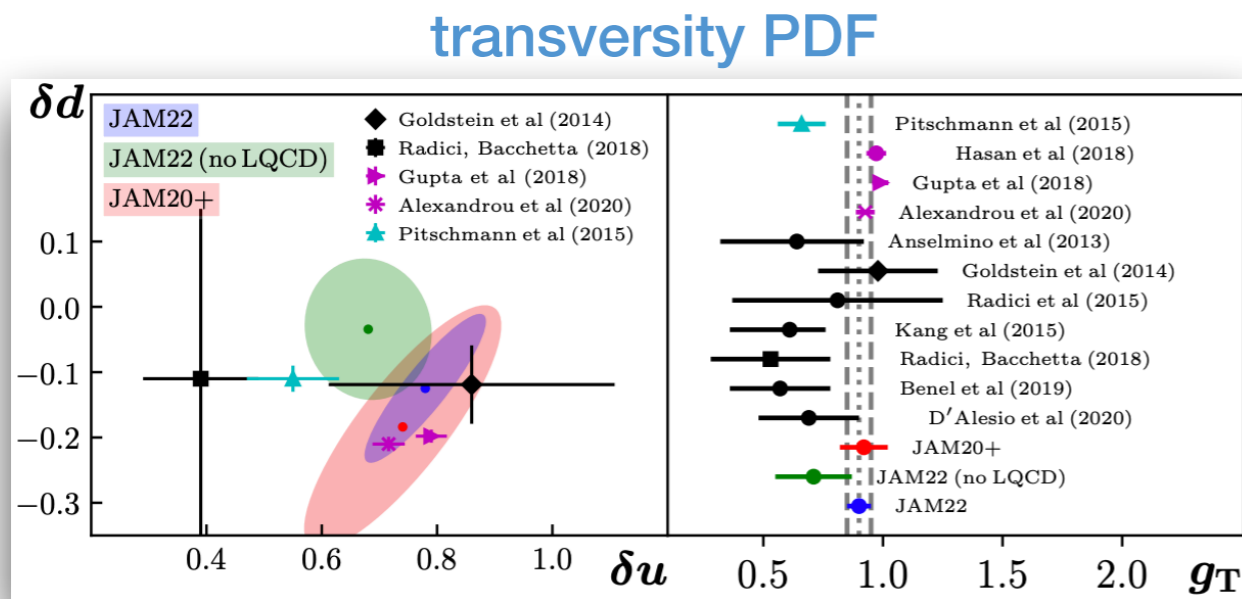
Synergies: constraints & predictive power of lattice QCD



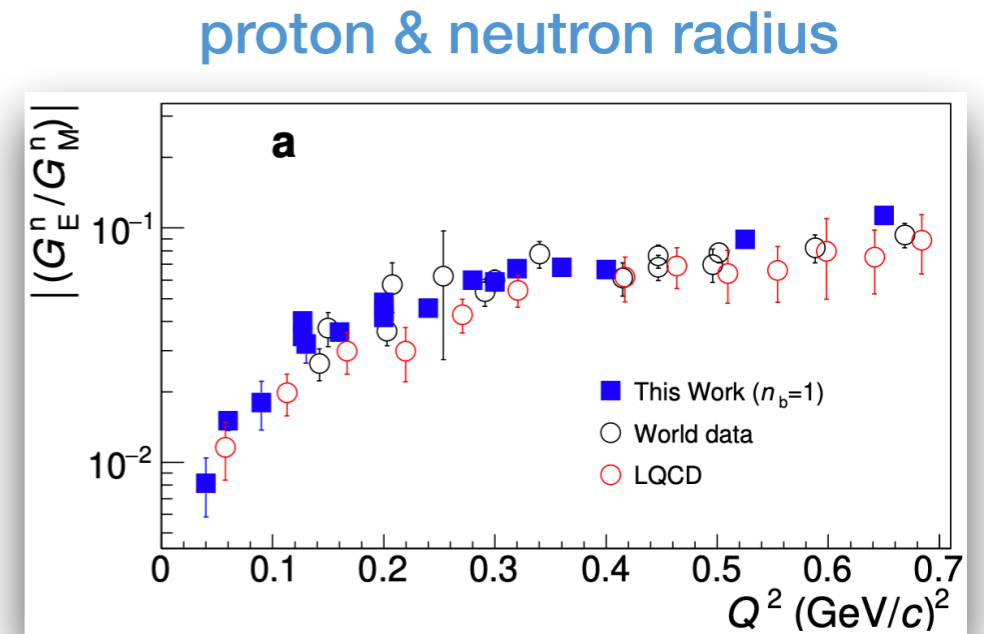
[JAM & ETMC, PRD 103 (2021) 016003]



[JAM/HadStruc, PRD105 (2022) 114051]



[JAM, PRD 106 (2022) 3, 034014]



[Atac et al., Nature Comm. 12, 1759 (2021)]

Summary

- ★ Lattice QCD data on GPDs will play an important role in the pre-EIC era and can complement experimental efforts of JLab@12GeV
- ★ New proposal for Lorentz invariant decomposition has great advantages:
 - significant reduction of computational cost
 - access to a broad range of t and ξ
- ★ Future calculations have the potential to transform the field of GPDs
- ★ Calculation of twist-3 distribution functions is promising
- ★ Synergy with phenomenology is an exciting prospect!

Summary

- ★ Lattice QCD data on GPDs will play an important role in the pre-EIC era and can complement experimental efforts of JLab@12GeV
- ★ New proposal for Lorentz invariant decomposition has great advantages:
 - significant reduction of computational cost
 - access to a broad range of t and ξ
- ★ Future calculations have the potential to transform the field of GPDs
- ★ Calculation of twist-3 distribution functions is promising
- ★ Synergy with phenomenology is an exciting prospect!

Thank you



DOE Early Career Award (NP)
Grant No. DE-SC0020405



QUARK-GLUON
TOMOGRAPHY
COLLABORATION



Award Number:
DE-SC0023646

