

Unpolarized Nucleon TMDPDFs from LQCD

Based on arXiv:2211.02340 [hep-lat]

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① Introduction

② Framework

③ Numerical Results

④ Prospect

① Introduction

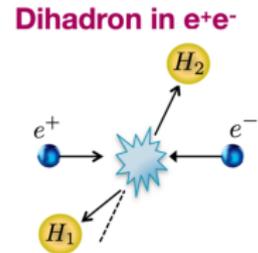
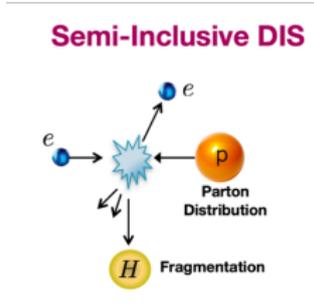
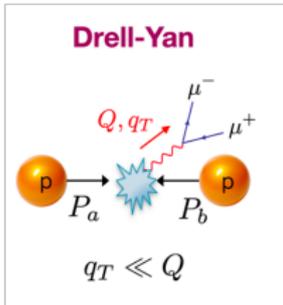
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Why Care about TMDPDFs

- Understanding the inner structure of hadrons.
- TMD processes are very important processes in high energy collisions, like SIDIS on EIC.
- TMDPDFs are important inputs for experiments on TMD processes.



Boussarie, et al. 2304.03302

Recent Progress

- Theoretical analysis
 - TMD factorization, evolution and resummation
Boussarie, et al. TMD handbook (2023) . . .
- Phenomenological extractions
 - Unpolarized
Bacchetta, et al. JHEP 06 (2017); Scimem, et al. JHEP 06 (2020); Bury, et al. JHEP 10 (2022); Bacchetta, et al. JHEP 10 (2022); Moos, et al. 2305.07473 . . .
 - Sivers
Bury, et al. PRL 126 (2021), JHEP 05 (2021); Fernando, et al. 2304.14328 . . .
 - Boer-Mulders
Zhang, et al. PRD 77 (2008), Lu, et al. PRD 81 (2010) . . .

Recent Progress

- Lattice calculation

- Lorentz-invariant approach: ratios of Mellin moments

Hagler, et al. EPL 88 (2009); Musch, et al. PRD 85 (2012); Yoon, et al. 1601.05717, PRD 96 (2017) . . .

- LaMET: formalism

Ji, et al. PLB 811 (2020), NPB 955 (2020), RMP 93 (2021); Ebert, et al. JHEP 04 (2022); Deng, et al. JHEP 09 (2022); Ji, et al. PRD 105 (2022), 2305.04416; Ó. del Río, et al. 2304.14440 . . .

- LaMET: soft functions

LPC, PRL 125 (2020); Shanahan, et al. PRD 104 (2021); Schlemmer, et al. JHEP 08 (2021); Li, et al. PRL 128 (2022); LPC, PRD 106 (2022), 2306.06488 . . .

- LaMET: renormalization and resummation

LPC, PRL 129 (2022); Su, et al. NPB 991 (2023) . . .

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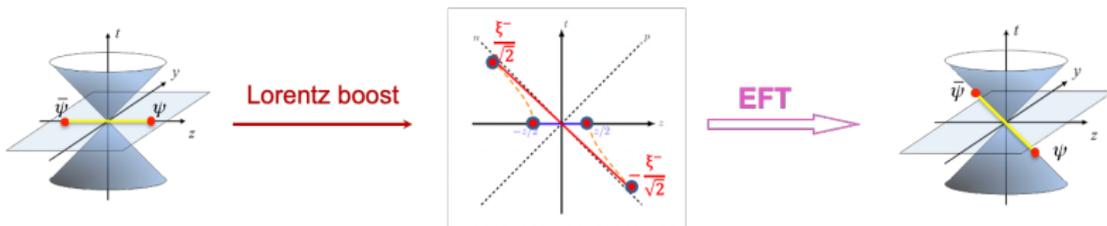
④ Prospect

LaMET

- Quasi distribution: equal time, P^z -dependent.
- Light-cone distribution: separate in time direction, universal.

$$\tilde{f}_\Gamma(x, b_\perp, \zeta_z, \mu) \sqrt{S_I(b_\perp, \mu)} = H_\Gamma\left(\frac{\zeta_z}{\mu^2}\right) e^{\frac{1}{2} \ln\left(\frac{\zeta_z}{\tau}\right) K(b_\perp, \mu)} \times f(x, b_\perp, \mu, \zeta) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{\zeta_z}, \frac{M^2}{(P^z)^2}, \frac{1}{b_{1\perp}^2 \zeta_z}\right)$$

Diagram labels: Intrinsic soft function, Collins-Soper kernel, Quasi distribution, Hard kernel, Light-cone distribution.

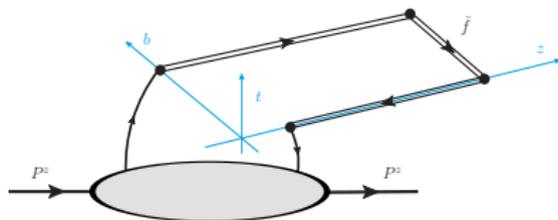


Matrix Element on the Lattice

$$\tilde{f}_\Gamma(x, b_\perp, P^z, \mu) \equiv \lim_{\substack{a \rightarrow 0 \\ L \rightarrow \infty}} \int \frac{dz}{2\pi} e^{-iz(xP^z)} \times \frac{\tilde{h}_\Gamma^0(z, b_\perp, P^z, a, L)}{\sqrt{Z_E(2L+z, b_\perp, a)} Z_O(1/a, \mu, \Gamma)}$$

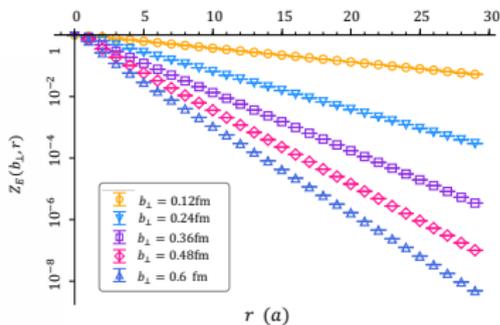
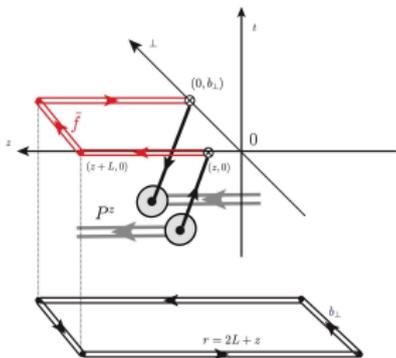
$$\tilde{h}_\Gamma^0(z, b_\perp, P^z, a, L) = \langle P^z | \tilde{O}_{\Gamma, \square}^0(z, b_\perp, P^z; L) | P^z \rangle$$

$$\tilde{O}_{\Gamma, \square}^0(z, b_\perp, L) \equiv \bar{\psi}(b_\perp \hat{n}_\perp) \Gamma U_{\square, L}(b_\perp \hat{n}_\perp, z \hat{n}_z) \psi(z \hat{n}_z)$$



Renormalization: Divergence from Wilson links

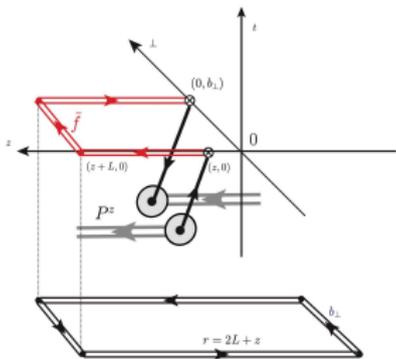
$$\tilde{h}_\Gamma(z, b_\perp, P^z, a, \mu) = \lim_{L \rightarrow \infty} \frac{\tilde{h}_\Gamma^0(z, b_\perp, P^z, a, L)}{\sqrt{Z_E(2L+z, b_\perp, a)Z_O(1/a, \mu, \Gamma)}}$$



Zhang, et al. PRL 129 (2022)

Renormalization: Divergence from quark-gauge link vertices

$$\tilde{h}_\Gamma(z, b_\perp, P^z, a, \mu) = \lim_{L \rightarrow \infty} \frac{\tilde{h}_\Gamma^0(z, b_\perp, P^z, a, L)}{\sqrt{Z_E(2L+z, b_\perp, a)} Z_O(1/a, \mu, \Gamma)}$$



- Extract Z_O by matching the renormalized results with perturbation theory at short distance.

$$Z_O(1/a, \mu, \Gamma) =$$

$$\lim_{L \rightarrow \infty} \frac{\tilde{h}_\Gamma^0(z, b_\perp, 0, a, L)}{\sqrt{Z_E(2L+z, b_\perp, a)} \tilde{h}_\Gamma^{\overline{\text{MS}}}(z, b_\perp, \mu)}$$

Zhang, et al. PRL 129 (2022)

RG Resummation: Hard kernel

The hard kernel involves some log terms, the higher orders cannot be ignored when $\zeta_z/\mu^2 = (2xP^z/\mu)^2$ is not close to 1.

$$\ln H_\Gamma^{(1)}\left(\frac{\zeta_z}{\mu^2}\right) = \frac{\alpha_s C_F}{2\pi} \left(-2 + \frac{\pi^2}{12} + \ln \frac{\zeta_z}{\mu^2} - \frac{1}{2} \ln^2 \frac{\zeta_z}{\mu^2} \right)$$

$$\ln H_\Gamma^{(2)}\left(\frac{\zeta_z}{\mu^2}\right) = \alpha_s^2 \left[c_2 - \frac{1}{2} (\gamma_C^{(2)} - \beta_0 c_1) \ln \frac{\zeta_z}{\mu^2} - \frac{1}{4} \left(\Gamma_{\text{cusp}}^{(2)} - \frac{\beta_0 C_F}{2\pi} \right) \ln^2 \frac{\zeta_z}{\mu^2} - \frac{\beta_0 C_F}{24\pi} \ln^3 \frac{\zeta_z}{\mu^2} \right]$$

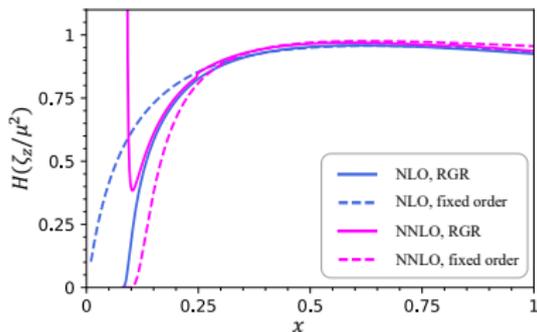
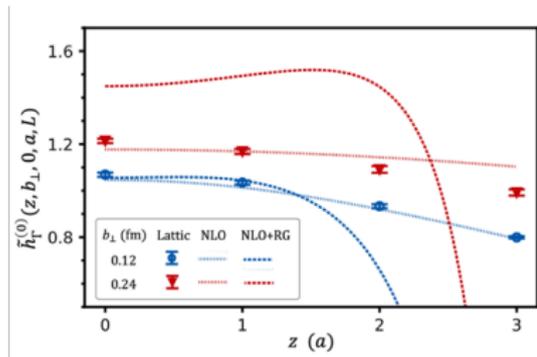
Thus, evolve the fixed order result from $\mu_0 = 2xP^z$ to $\mu = 2$ GeV.

$$H\left(\zeta_z/\mu^2\right) = H\left(\zeta_z/\mu_0^2\right) \exp \left[\int_{\mu_0}^{\mu} \frac{d\mu}{\mu} \left(\Gamma_{\text{cusp}}^{(1)} \ln \frac{\zeta_z}{\mu^2} \alpha_s(\mu) + \gamma_C^{(1)} \alpha_s(\mu) + \Gamma_{\text{cusp}}^{(2)} \ln \frac{\zeta_z}{\mu^2} \alpha_s^2(\mu) + \dots \right) \right]$$

Ji, et al. PLB 811 (2020), Su, et al. NPB 991 (2023)

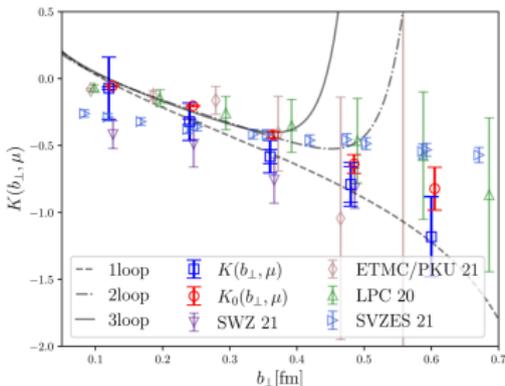
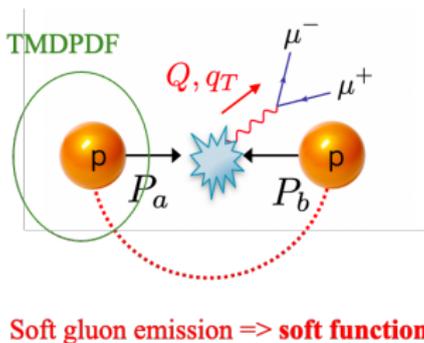
RG Resummation

All perturbation results should be modified with RG resummation in the consistent way.



Soft Functions: CS kernel

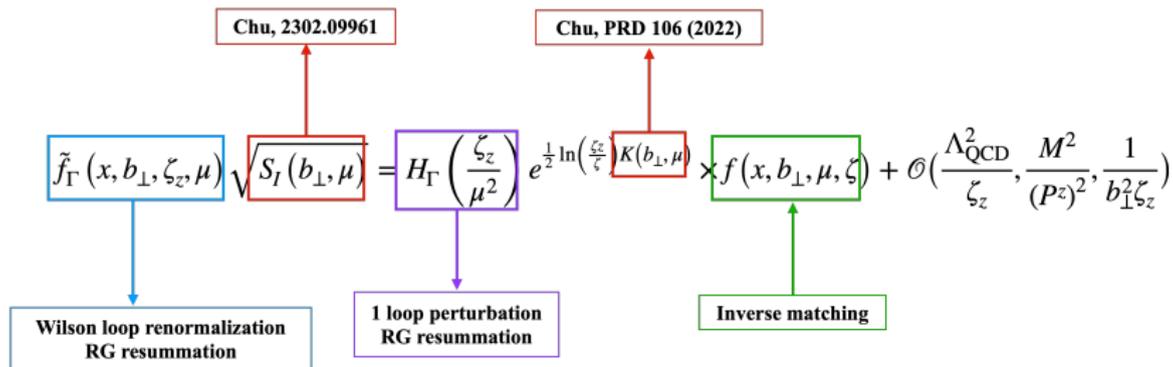
Rapidity dependent: Collins-Soper kernel.



Chu(LPC), PRD 106 (2022)

Inverse Matching

Now we are prepared to get the physical TMDPDF.



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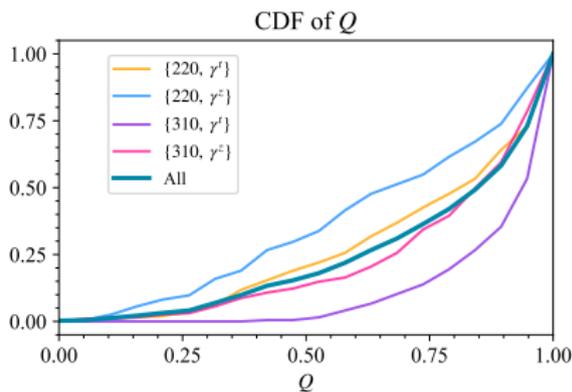
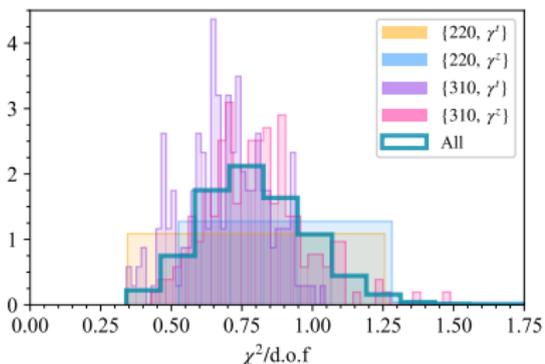
③ Numerical Results

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Lattice Setup

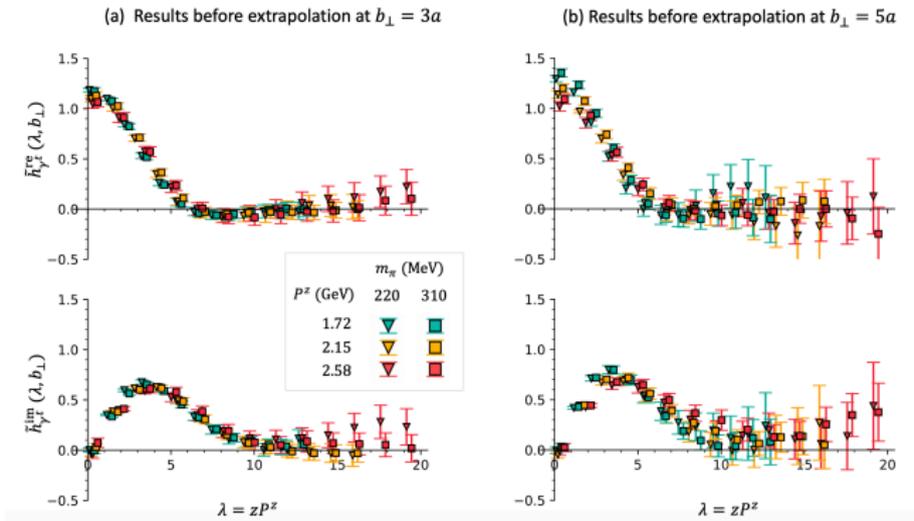
- 2+1+1 flavors of HISQ action by MILC ($a=0.12$ fm);
- Lattice volume: $n_s^3 \times n_t = 48^3 \times 64$;
- Valance pion mass: 310 MeV, 220 MeV;
- Gamma structure: γ^t and γ^z ;
- Hadron momentum: 1.72 GeV, 2.15 GeV, 2.58 GeV;
- Separation in the momentum direction:
 $z \in [0, 12]a = [0, 1.44]$ fm;
- Separation in the transverse direction:
 $b \in [1, 5]a = [0.12, 0.6]$ fm.

Ground State Fit: Fit quality



- $\chi^2/\text{d.o.f.}$ mainly distribute in the reasonable range.
- Almost no bad fits with p-value smaller than 0.05.

Renormalized Quasi Distribution in the Coordinate Space



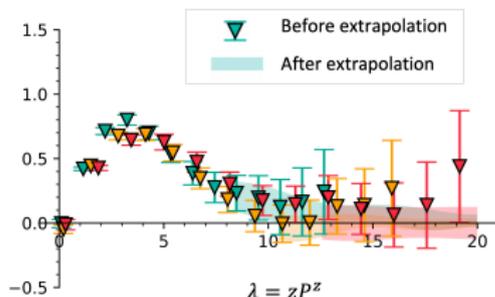
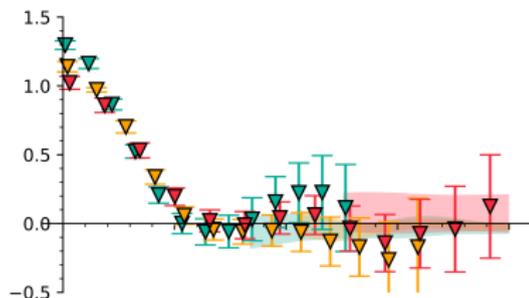
- Truncation at the finite distance.
- Tails are close to zero.

Extrapolation

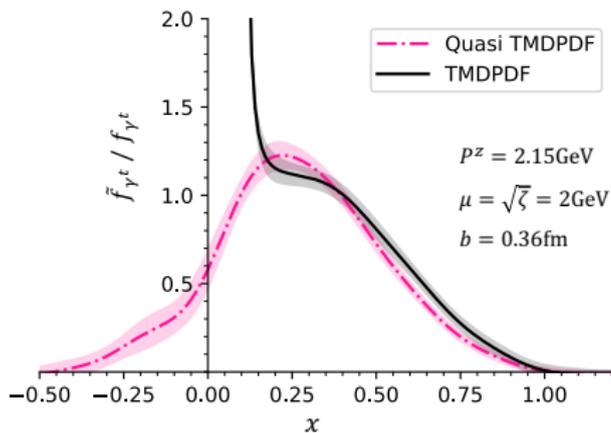
The extrapolation form in the coordinate space originated from the power law behavior in the endpoint regions.

Ji, et al. NPB 964 (2021)

$$\tilde{h}_{\Gamma,\text{extra}}(\lambda) = \left[\frac{m_1}{(-i\lambda)^{n_1}} + e^{i\lambda} \frac{m_2}{(i\lambda)^{n_2}} \right] e^{-\lambda/\lambda_0}$$



Inverse Matching

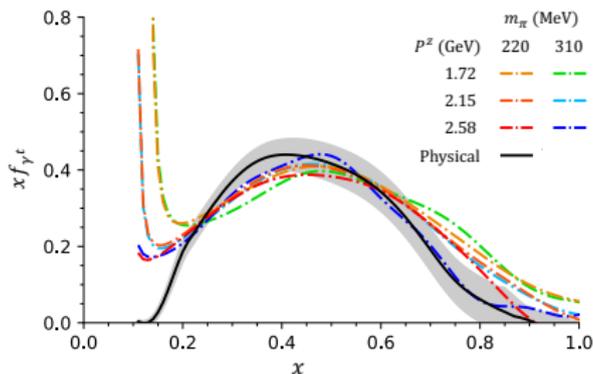


- Quasi and light-cone distributions are consistent in the reliable region.
- Near $x = 0$, the RG resummed hard kernel blows up.

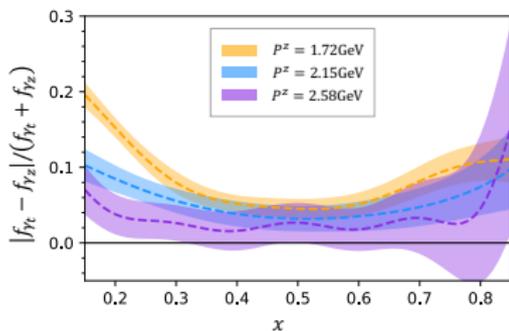
Physical Limit

We did a combined extrapolation to achieve the physical limit.

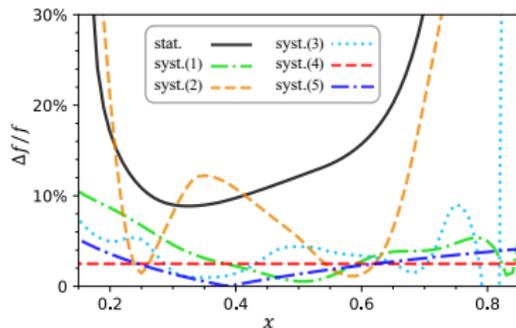
$$f_{\Gamma}(x, b_{\perp}, \mu, \zeta; m_{\pi}, P^z) = f_{\Gamma}(x, b_{\perp}, \mu, \zeta) \Big|_{\substack{m_{\pi} \rightarrow m_{\pi, \text{phy}} \\ P^z \rightarrow \infty}} \times \left[1 + d_0 (m_{\pi}^2 - m_{\pi, \text{phy}}^2) + \frac{d_1}{(P^z)^2} \right]$$



Error Estimation



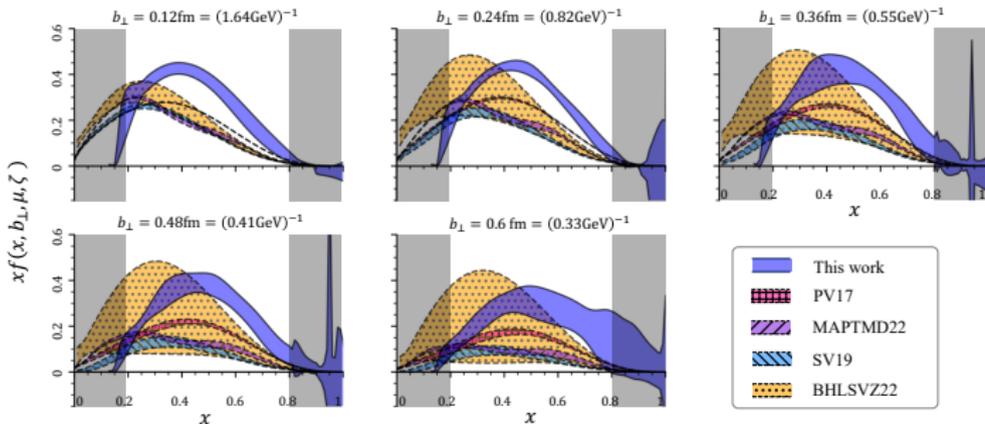
- Asymptotic behavior when P^z increases.
- The ratio shows the higher power corrections.
- There are large higher power corrections near the endpoint regions.



Source of systematical uncertainties:

- Power corrections;
- Physical extrapolation;
- Long distance extrapolation;
- Intrinsic soft function;
- Collins-Soper kernel.

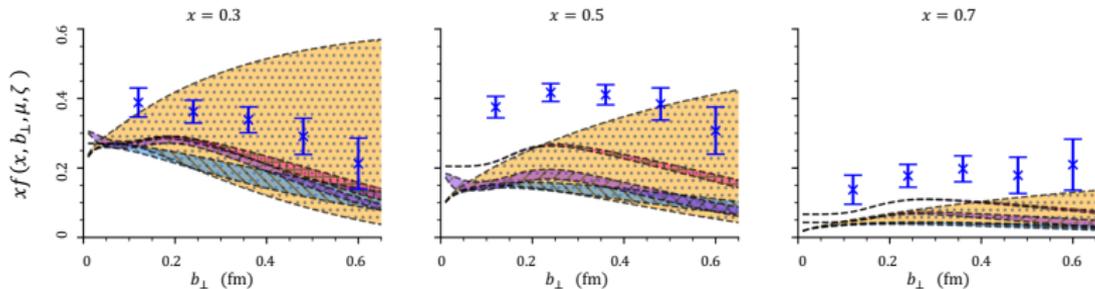
Comparison with Pheno: x -dependence



Bacchetta, et al. JHEP 06 (2017); Scimemi, et al. JHEP 06 (2020); Bury, et al. JHEP 10 (2022); Bacchetta, et al. JHEP 10 (2022)

- Preliminary lattice results consist with pheno results qualitatively.
- Shaded area cannot be predicted with LaMET.

Comparison with Pheno: b -dependence



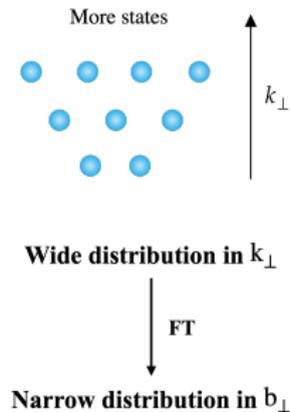
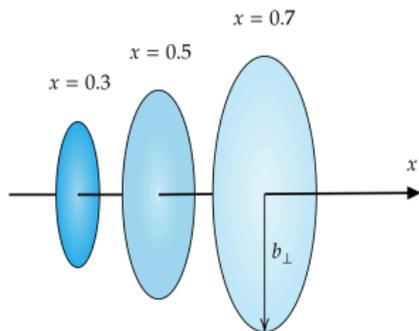
Bacchetta, et al. JHEP 06 (2017); Scimemi, et al. JHEP 06 (2020); Bury, et al. JHEP 10 (2022); Bacchetta, et al. JHEP 10 (2022)

- Both lattice and pheno results show a decreasing behavior in the b -dependence.
- Larger x corresponds to longer correlation length.

Discussion

Saturation: there are a lot of partons with small x fractions, they need to occupy more states with large transverse momentum.

Our results: large x — long correlation length



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Improvements

Direction	Prospect
Discretization effects	Continuum limit, Lattice pert. theory
Transverse behavior	Small b problem, Confinement
Power correction	Larger γ factor
More	Higher loop calculation

X. Ji, Y. Liu and Y. Su, 2305.04416

Ó. del Río and A. Vladimirov, 2304.14440

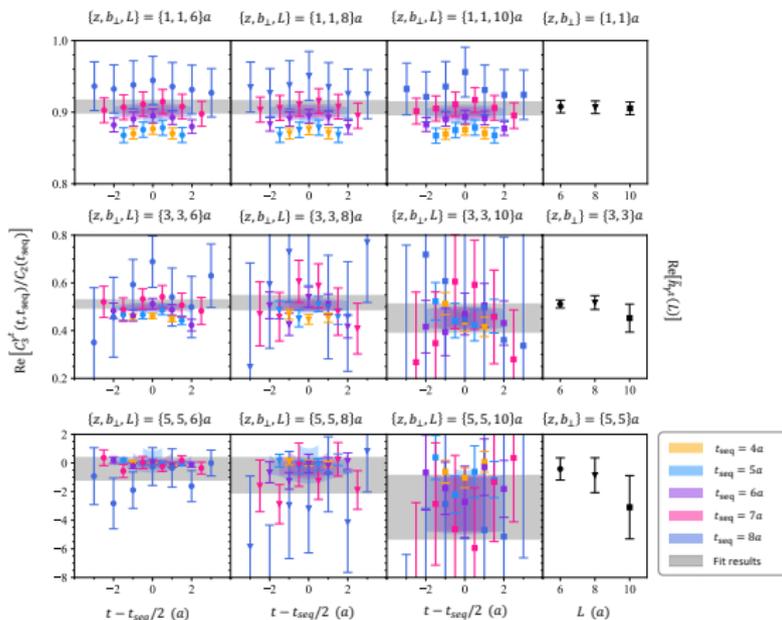
Other TMDPDFs

Leading Quark TMDPDFs  Nucleon Spin  Quark Spin

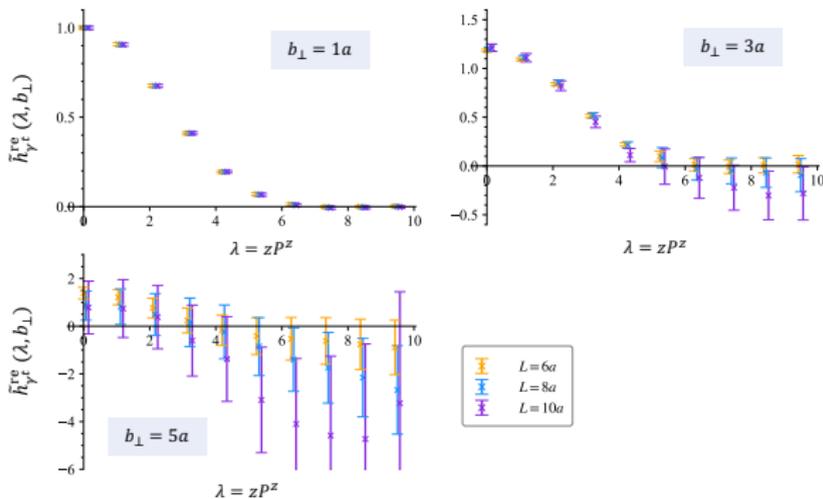
		Quark Polarization		
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Nucleon Polarization	U	$f_1 = \text{○} \begin{matrix} \bullet \\ \uparrow \end{matrix}$ Unpolarized		$h_1^\perp = \text{○} \begin{matrix} \uparrow \\ \bullet \end{matrix} - \text{○} \begin{matrix} \downarrow \\ \bullet \end{matrix}$ Boer-Mulders
	L		$g_{1L} = \text{○} \begin{matrix} \bullet \\ \rightarrow \end{matrix} - \text{○} \begin{matrix} \bullet \\ \leftarrow \end{matrix}$ Helicity	$h_{1L}^\perp = \text{○} \begin{matrix} \uparrow \\ \rightarrow \end{matrix} - \text{○} \begin{matrix} \uparrow \\ \leftarrow \end{matrix}$ Worm-gear
	T	$f_{1T}^\perp = \text{○} \begin{matrix} \uparrow \\ \bullet \end{matrix} - \text{○} \begin{matrix} \downarrow \\ \bullet \end{matrix}$ Sivers	$g_{1T}^\perp = \text{○} \begin{matrix} \uparrow \\ \bullet \end{matrix} - \text{○} \begin{matrix} \uparrow \\ \leftarrow \end{matrix}$ Worm-gear	$h_{1T}^\perp = \text{○} \begin{matrix} \uparrow \\ \bullet \end{matrix} - \text{○} \begin{matrix} \uparrow \\ \rightarrow \end{matrix}$ Pretzelosity

Thank You

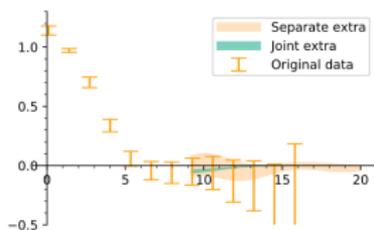
Backup: Ground State Fit: Consistent results with different L



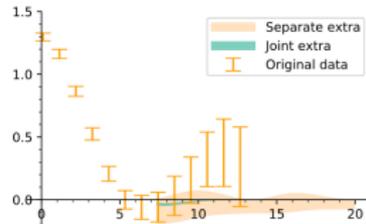
Backup: L -dependence



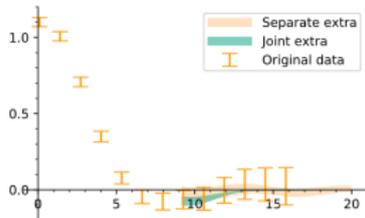
Backup: Extrapolation in the coordinate space



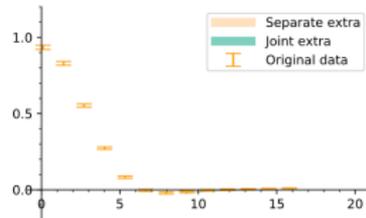
(a) $m_\pi = 220\text{MeV}$, $P^2 = 2.15\text{GeV}$, $b_\perp = 5a$



(b) $m_\pi = 220\text{MeV}$, $P^2 = 1.72\text{GeV}$, $b_\perp = 5a$



(c) $m_\pi = 220\text{MeV}$, $P^2 = 2.15\text{GeV}$, $b_\perp = 3a$



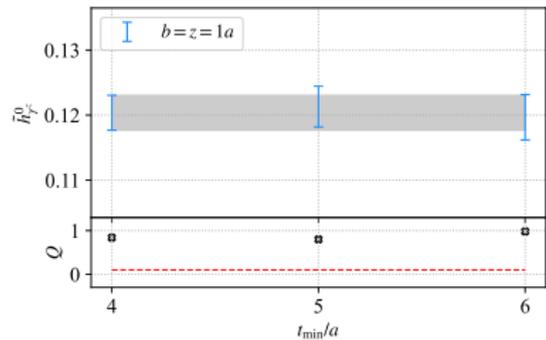
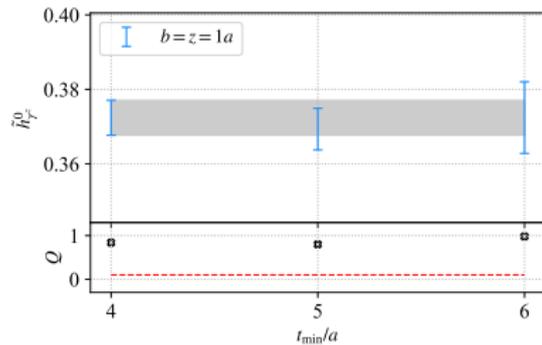
(d) $m_\pi = 220\text{MeV}$, $P^2 = 2.15\text{GeV}$, $b_\perp = 1a$

Backup: Extrapolation Parameters

Fit results:

	m_1	m_2	n_1	n_2	λ_0	$\chi^2/\text{d.o.f}$
$b = 1a$	-8.8(3.7)	0.25(50)	0.909(39)	1.13(74)	2.63(38)	1.0
$b = 2a$	-6.3(2.9)	-3.9(5.5)	0.943(61)	2.37(68)	4.1(1.1)	1.1
$b = 3a$	-66(60)	-78(76)	0.89(12)	1.71(31)	2.42(85)	1.4
$b = 4a$	-8.0(4.4)	-3.3(2.9)	0.801(78)	1.55(38)	4.3(1.6)	0.75
$b = 5a$	-8.5(10)	-3.8(5.3)	0.84(16)	1.22(44)	4.4(2.8)	0.57
Joint fit	-	-	0.887(28)	1.65(12)	2.53(28)	1.2

Backup: Ground state fit stability



Backup: Lattice setup

a12m130: We use the valence tadpole improved clover fermion on the hypercubic (HYP) smeared $2 + 1 + 1$ flavors MILC configurations with highly improved staggered quark (HISQ) sea and 1-loop Symanzik improved gauge action.

Backup: Dispersion Relation

$$E^2 = m^2 + c_1 \cdot P^2 + c_2 \cdot P^4$$

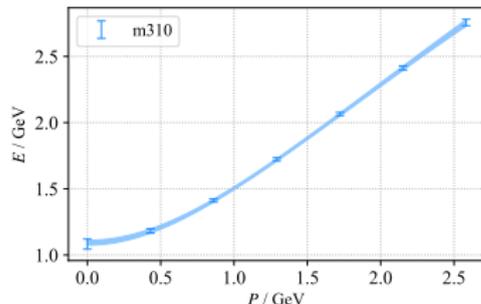
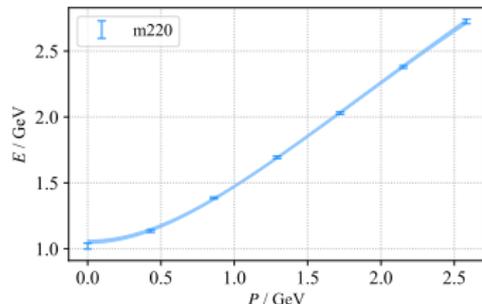


Figure 2: Dispersion relations.

m220: $Q = 0.19$, $c_1 = 1.081(20)$, $c_2 = -0.0548(96)$;

m310: $Q = 1$, $c_1 = 1.087(28)$, $c_2 = -0.053(13)$.