DIFFRACTIVE VECTOR MESON PRODUCTION IN THE COLOR GLASS CONDENSATE

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EIC Theory WG Meeting
03/23/2023

Good-Walker/Miettinen-Pumplin

M. L. Good and W. D. Walker, Phys. Rev. 120 (1960) 1857
H. I. Miettinen and J. Pumplin, Phys. Rev. D18 (1978) 1696

Discussing mainly diffractive scattering in p+p collisions, Miettinen and Pumplin ask two questions:

1. What are the states which diagonalize the diffractive part of the S-matrix, so that their interactions are described simply by absorption coefficients?

   Answer in their paper: States of the parton model (fixed number N, positions $\vec{b}_i$, fixed x)

2. What causes the large variations in the absorption coefficients at a given impact parameter, which are implied by the large cross section for diffractive production?

   Answer in their paper: Fluctuations in N, $\vec{b}_i$, x between the states. "Among the parton states which describe a high-energy hadron, there are some which are rich in wee partons, and are therefore likely to interact, while other states have few or no wee partons, and correspond to the transparent channels of diffraction."

Miettinen-Pumplin: Optical Model Formulation

H. I. Miettinen and J. Pumplin, Phys. Rev. D18 (1978) 1696

Target: Average optical potential

Beam particle: $|B\rangle = \sum_k C_k |\psi_k\rangle$ (linear combination of the eigenstates of diffraction $|\psi_k\rangle$)

With $\mathrm{Im}T = 1 - \mathrm{Re}S$ the imaginary part of the scattering amplitude operator, we have

$$\mathrm{Im}T|\psi_k\rangle = t_k|\psi_k\rangle$$

with t_k the probability for eigenstate $|\psi_k\rangle$ to interact with the target (absorption coefficients)

Normalize: $\langle B|B\rangle = \sum_k |C_k|^2 = 1$

Elastic scattering: $\langle B|\mathrm{Im}T|B\rangle = \sum_k |C_k|^2 t_k = \langle t\rangle$

Miettinen-Pumplin: Cross Sections

H. I. Miettinen and J. Pumplin, Phys. Rev. D18 (1978) 1696

Total cross section:

$$d\sigma_{\text{tot}}/d^2\vec{b} = 2\langle t\rangle$$

Elastic cross section:

$$d\sigma_{\text{el}}/d^2\vec{b} = \langle t\rangle^2$$

Incoherent diffractive cross section:

$$d\sigma_{\text{diff}}/d^2\vec{b} = \sum_k |\langle\psi_k|\text{Im}T|B\rangle|^2 - d\sigma_{\text{el}}/d^2\vec{b} = \sum_k |\langle\psi_k|\text{Im}T|\sum_i C_i|\psi_i\rangle|^2 - d\sigma_{\text{el}}/d^2\vec{b}$$

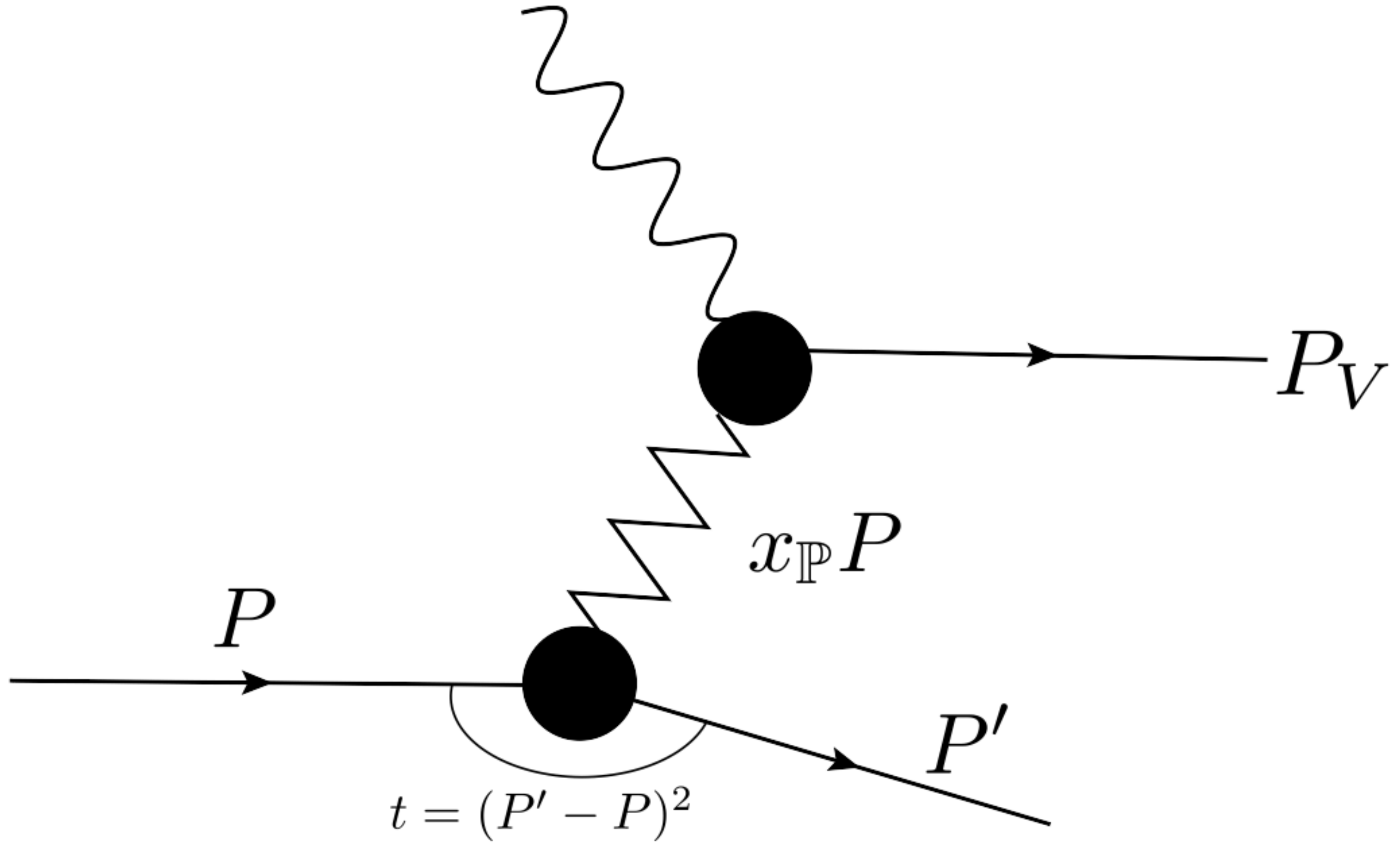
$$= \sum_{k,i} |\langle\psi_k|C_i t_i|\psi_i\rangle|^2 - d\sigma_{\text{el}}/d^2\vec{b} = \sum_{k,i} \delta_{ik}|C_i t_i|^2 - d\sigma_{\text{el}}/d^2\vec{b} = \sum_k |C_k|^2 t_k^2 - \langle t\rangle^2 = \langle t^2\rangle - \langle t\rangle^2$$

$$d\sigma_{\text{diff}}/d^2\vec{b} = \langle t^2\rangle - \langle t\rangle^2$$

Color Glass Condensate calculation

- We study diffractive production in e+p/A (not p+p)
- The projectile can be understood as a quark anti-quark dipole (splitting from the incoming virtual photon)
- The fluctuations are included in the target wave function: Fluctuating spatial distribution of the gluon fields (normalization fluctuations correspond to N fluctuations, spatial fluctuations to $\vec{b}_i$ fluctuations
  (see **Blaizot and Traini, 2209.15545 [hep-ph]** for the effect of fluctuations of the dipole size)

Diffractive vector meson production



- Coherent diffraction: $$\frac{d\sigma^{\gamma^* p \to Vp}}{dt} = \frac{1}{16\pi} \left| \left\langle A^{\gamma^* p \to Vp}\left(x_P, Q^2, \vec{\Delta}\right) \right\rangle \right|^2$$

sensitive to the average size of the target

- Incoherent diffraction: $$\frac{d\sigma^{\gamma^* p \to Vp^*}}{dt} = \frac{1}{16\pi} \left(\left\langle \left| A^{\gamma^* p \to Vp}\left(x_P, Q^2, \vec{\Delta}\right) \right|^2 \right\rangle - \left| \left\langle A^{\gamma^* p \to Vp}\left(x_P, Q^2, \vec{\Delta}\right) \right\rangle \right|^2 \right)$$

sensitive to fluctuations (including geometric ones)

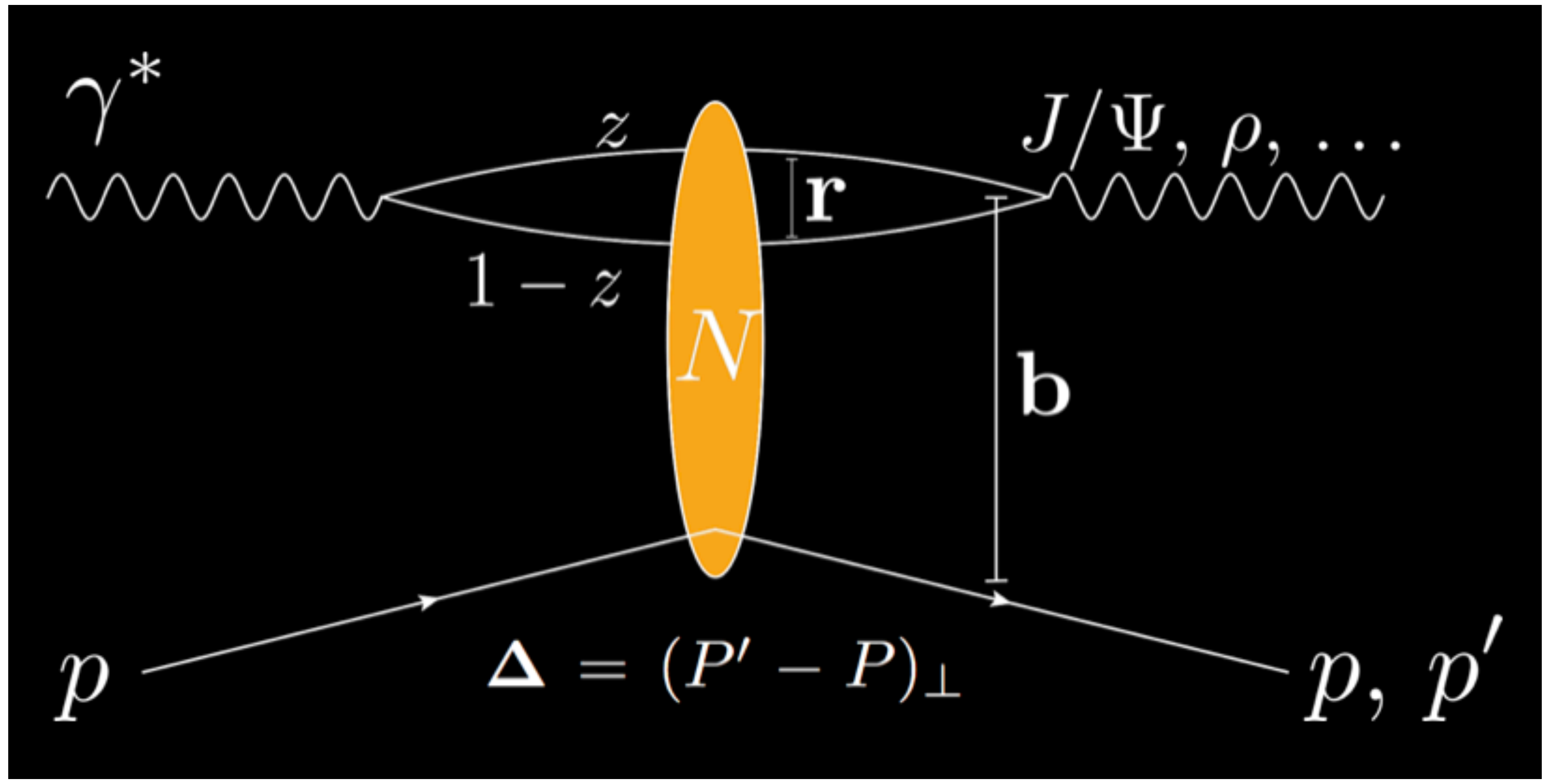
H. Kowalski, L. Motyka, G. Watt, Phys.Rev. D 74 (2006) 074016
A. Caldwell, H. Kowlaski, EDS 09, 190-192, e-Print: 0909.1254 [hep-ph]
M. L. Good and W. D. Walker, Phys. Rev. 120 (1960) 1857
H. I. Miettinen and J. Pumplin, Phys. Rev. D18 (1978) 1696
Y. V. Kovchegov and L. D. McLerran, Phys. Rev. D60 (1999) 054025
A. Kovner and U. A. Wiedemann, Phys. Rev. D64 (2001) 114002

Dipole picture: Scattering amplitude

H. Mäntysaari, B. Schenke, Phys. Rev. Lett. 117 (2016) 052301; Phys.Rev. D94 (2016) 034042

High energy factorization:

- $\gamma^* \to q\bar{q} : \psi^\gamma(r, Q^2, z)$
- $q\bar{q}$ dipole scatters with amplitude N
- $q\bar{q} \to V : \psi^V(r, Q^2, z)$



$$A \sim \int d^2b\, dz\, d^2r\, \psi^* \psi^V(\vec{r}, z, Q^2) e^{-i\vec{b}\cdot\vec{\Delta}} N(\vec{r}, x, \vec{b})$$

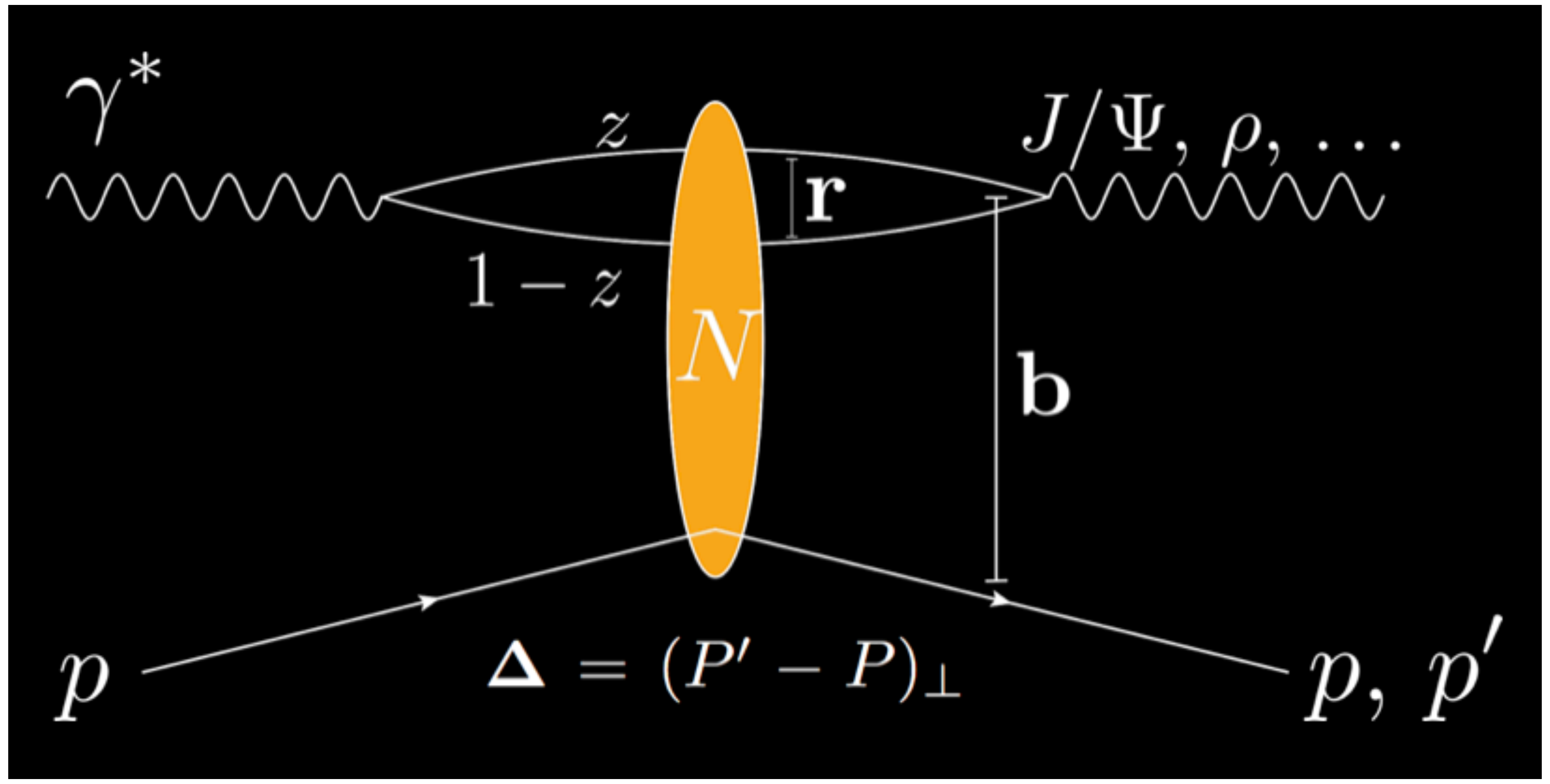
- Impact parameter ***b*** is the Fourier conjugate of transverse momentum transfer ***Δ*** → Access to spatial structure $(t = -\Delta^2)$

Color glass condensate formalism

H. Mäntysaari, B. Schenke, Phys.Rev.D 98 (2018) 3, 034013

Compute the Wilson lines using color charges whose correlator depends on $\vec{b}_\perp$

$$\langle \rho^a(\mathbf{b}_\perp)\rho^b(\mathbf{x}_\perp)\rangle = g^2\mu^2(x,\mathbf{b}_\perp)\delta^{ab}\delta^{(2)}(\mathbf{b}_\perp - \mathbf{x}_\perp)$$

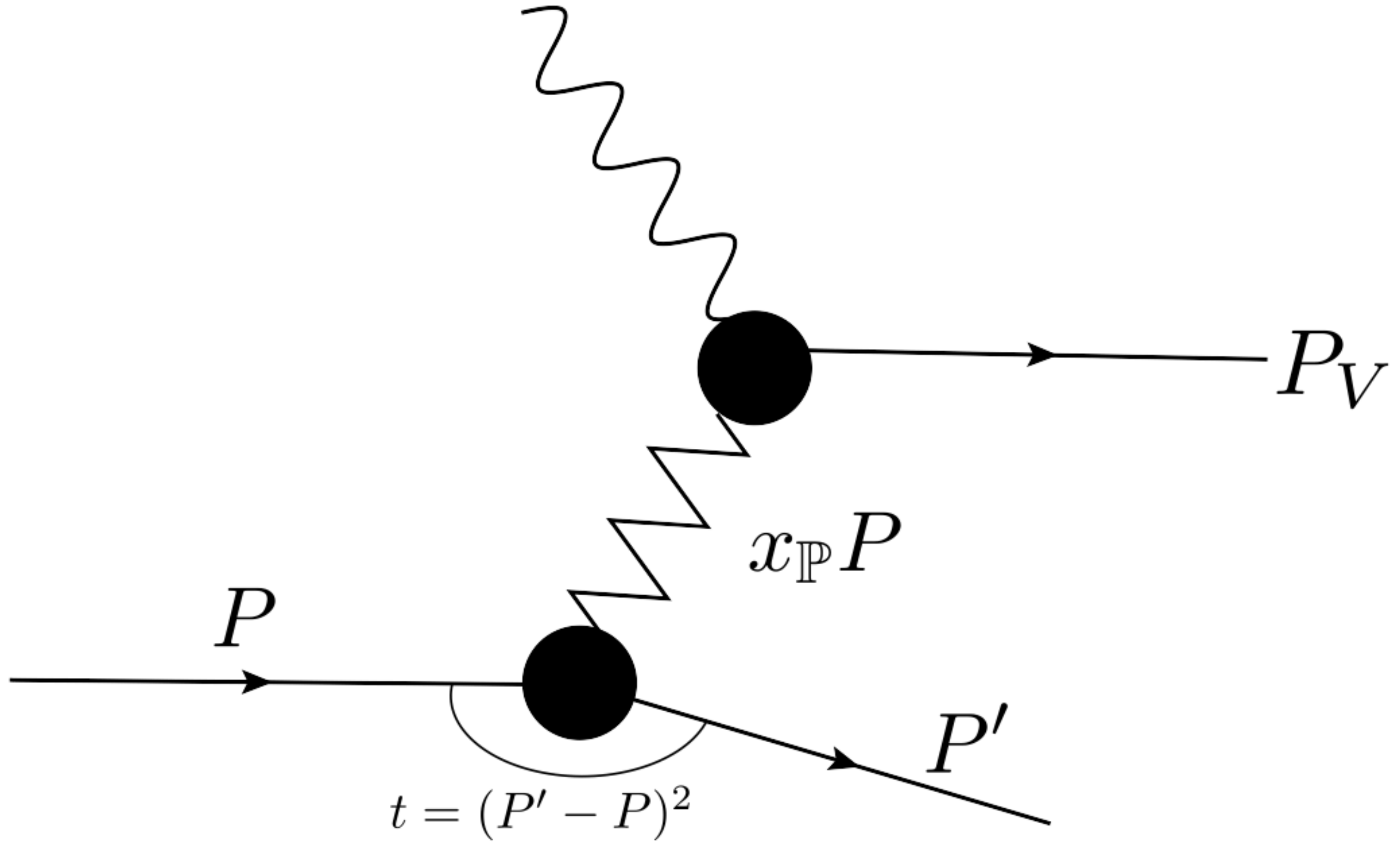


$$N(\vec{r}, x, \vec{b}) = N(\vec{x} - \vec{y}, x, (\vec{x} + \vec{y})/2) = 1 - \mathrm{Tr}(\mathrm{V}(\vec{\mathrm{x}})\mathrm{V}^\dagger(\vec{\mathrm{y}}))/\mathrm{N_c}$$

The trace appears at the level of the amplitude, because we project on a color singlet

$$A \sim \int d^2b\, dz\, d^2r\, \psi^*\psi^V(\vec{r}, z, Q^2) e^{-i\vec{b}\cdot\vec{\Delta}}[1 - Tr(V(\vec{x})V^\dagger(\vec{y}))/N_c]$$

Diffractive vector meson production



- Coherent diffraction: $$\frac{d\sigma^{\gamma^* p \to Vp}}{dt} = \frac{1}{16\pi} \left| \left\langle A^{\gamma^* p \to Vp} \left(x_P, Q^2, \vec{\Delta} \right) \right\rangle \right|^2$$

sensitive to the average size of the target

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sensitive to fluctuations (including geometric ones)

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Fluctuations in the target following the discussion in Blaizot and Traini, 2209.15545 [hep-ph]

Define

$$\hat{T}_p(\vec{b}) = \sum_i^{N_q} T_G(\vec{b}_i - \vec{b}) = \int d^2\vec{x}\, \hat{\rho}(\vec{x})\, T_G(\vec{x} - \vec{b})$$

T_G is the gluon distribution in a hot spot

$\hat{\rho}(\vec{x}) = \sum_i^{N_q} \delta(\vec{x} - \vec{b}_i)$ is the hot spot density operator in the transverse plane

The dipole cross section can be written as

$S = \exp\left[-\frac{1}{2}\sigma_{\mathrm{dip}}(x, \vec{r})\hat{T}_p(\vec{b})\right] \approx 1 - \frac{1}{2}\sigma_{\mathrm{dip}}(x, \vec{r})\hat{T}_p(\vec{b})$ in the weak field limit

The dipole cross section then is $\dfrac{d\sigma_{q\bar{q}}}{d^2\vec{b}} = 2[1 - S] = \sigma_{\mathrm{dip}}(x, \vec{r})\hat{T}_p(\vec{b})$

Fluctuations in the target

The dipole cross section then is $\frac{d\sigma_{q\bar{q}}}{d^2\vec{b}} = 2[1 - S] = \sigma_{\rm dip}(x, \vec{r})\hat{T}_p(\vec{b})$

This operator is diagonal in the basis of states $|\vec{b}_1, \ldots, \vec{b}_{N_q}\rangle$, where the $\vec{b}_i$ are the positions of the individual hot spots, frozen during the collision process:

These states can be considered the diffractive eigenstates

Coherent diffractive cross section:

$$\int d^2\vec{b}d^2\vec{b}'e^{-i\vec{\Delta}\cdot(\vec{b}-\vec{b}')} \left\langle \frac{d\sigma^{q\bar{q}}}{d^2\vec{b}} \right\rangle \left\langle \frac{d\sigma^{q\bar{q}}}{d^2\vec{b}'} \right\rangle = \langle \Sigma_{q\bar{q}}(\overrightarrow{\Delta}) \rangle^2$$

with $\Sigma_{q\bar{q}}(\overrightarrow{\Delta}) = \int d^2\vec{b}e^{-i\vec{\Delta}\cdot\vec{b}}\frac{d\sigma^{q\bar{q}}}{d^2\vec{b}}$ and $\langle \cdot \rangle$ is the average over the ground state wave function

Fluctuations in the target

Total diffractive cross section:

Allow all possible diffractive eigenstates $|\alpha\rangle$ as intermediate states (assume dilute limit here)

$$\int d^2\vec{b}d^2\vec{b}'e^{-i\vec{\Delta}\cdot(\vec{b}-\vec{b}')}\sigma^2_{\rm dip}\sum_\alpha \left|\langle\alpha|\hat{T}_p(\vec{b})|\psi_0\rangle\right|^2 = \langle\Sigma^2_{q\bar{q}}(\vec{\Delta})\rangle$$

in analogy to the optical model example

This also shows the relation to the density-density correlation function $\langle\hat{T}_p(\vec{b})\hat{T}_p(\vec{b}')\rangle$

and how we are sensitive to different distance scales via $\vec{b}-\vec{b}'$

See **Blaizot and Traini, 2209.15545 [hep-ph]** for a more detailed discussion

Model impact parameter dependence (proton, nucleon)

H. Mäntysaari, B. Schenke, Phys. Rev. Lett. 117 (2016) 052301; Phys.Rev. D94 (2016) 034042

1) Assume Gaussian proton shape:

$$T(\vec{b}) = T_{\mathrm{p}}(\vec{b}) = \frac{1}{2\pi B_{\mathrm{p}}} e^{-b^2/(2B_{\mathrm{p}})}$$

2) Assume Gaussian distributed and Gaussian shaped hot spots:

$$P(b_i) = \frac{1}{2\pi B_{\mathrm{qc}}} e^{-b_i^2/(2B_{\mathrm{qc}})}$$ (angles uniformly distributed)

$$T_{\mathrm{p}}(\vec{b}) = \frac{1}{N_{\mathrm{q}}} \sum_{i=1}^{N_{\mathrm{q}}} T_{\mathrm{G}}(\vec{b} - \vec{b}_i)$$ with N_{q} hot spots; $$T_{\mathrm{G}}(\vec{b}) = \frac{1}{2\pi B_{\mathrm{q}}} e^{-b^2/(2B_{\mathrm{q}})}$$

Diffractive J/ψ production in e+p at HERA

Nucleon parameters B_q, B_{qc}, can be constrained by e+p scattering data from HERA

Exclusive diffractive J/Ψ production in e+p:

Incoherent x-sec sensitive to fluctuations

H. Mäntysaari, B. Schenke, Phys. Rev. Lett. 117 (2016) 052301

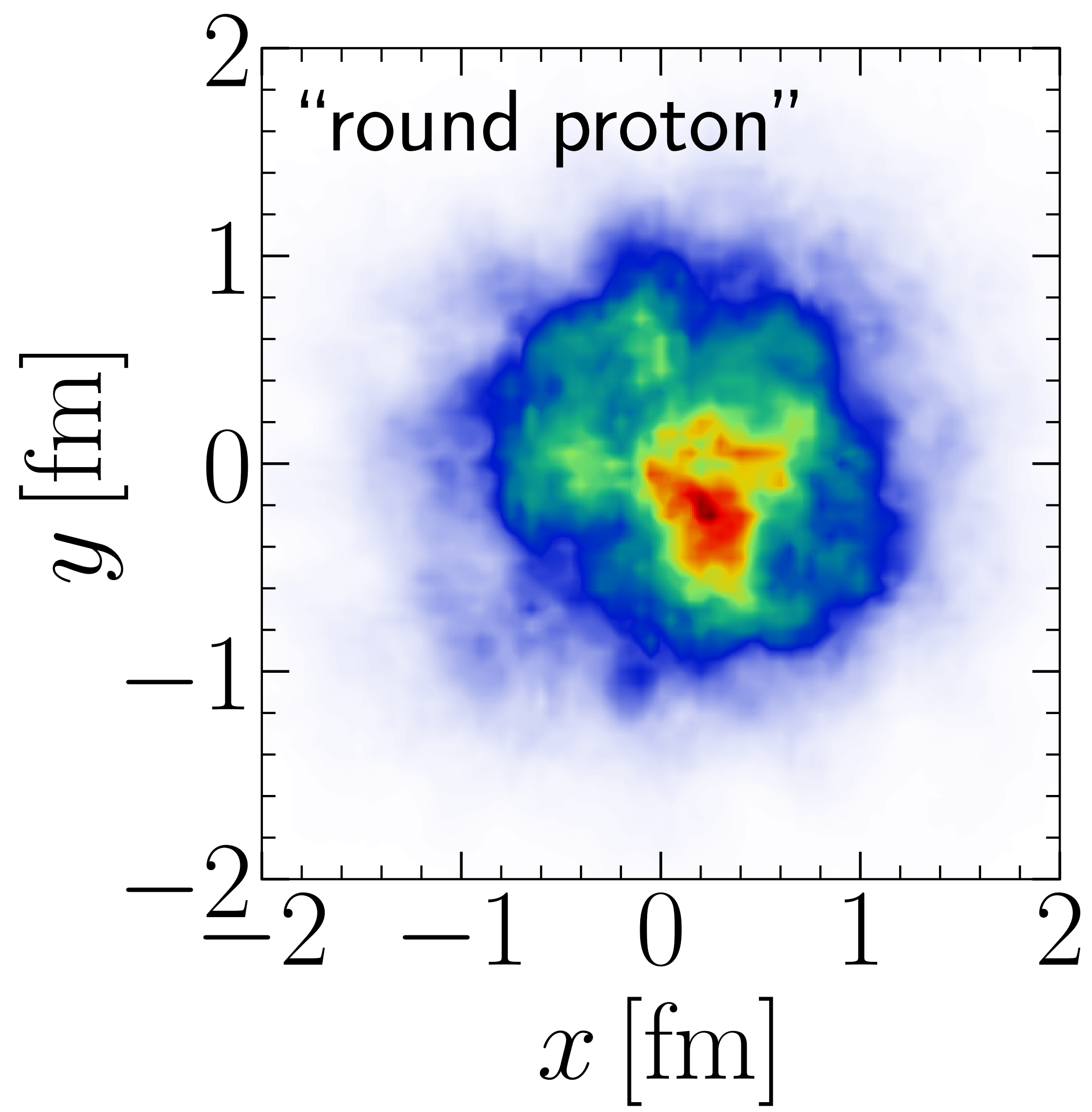
Phys.Rev. D94 (2016) 034042

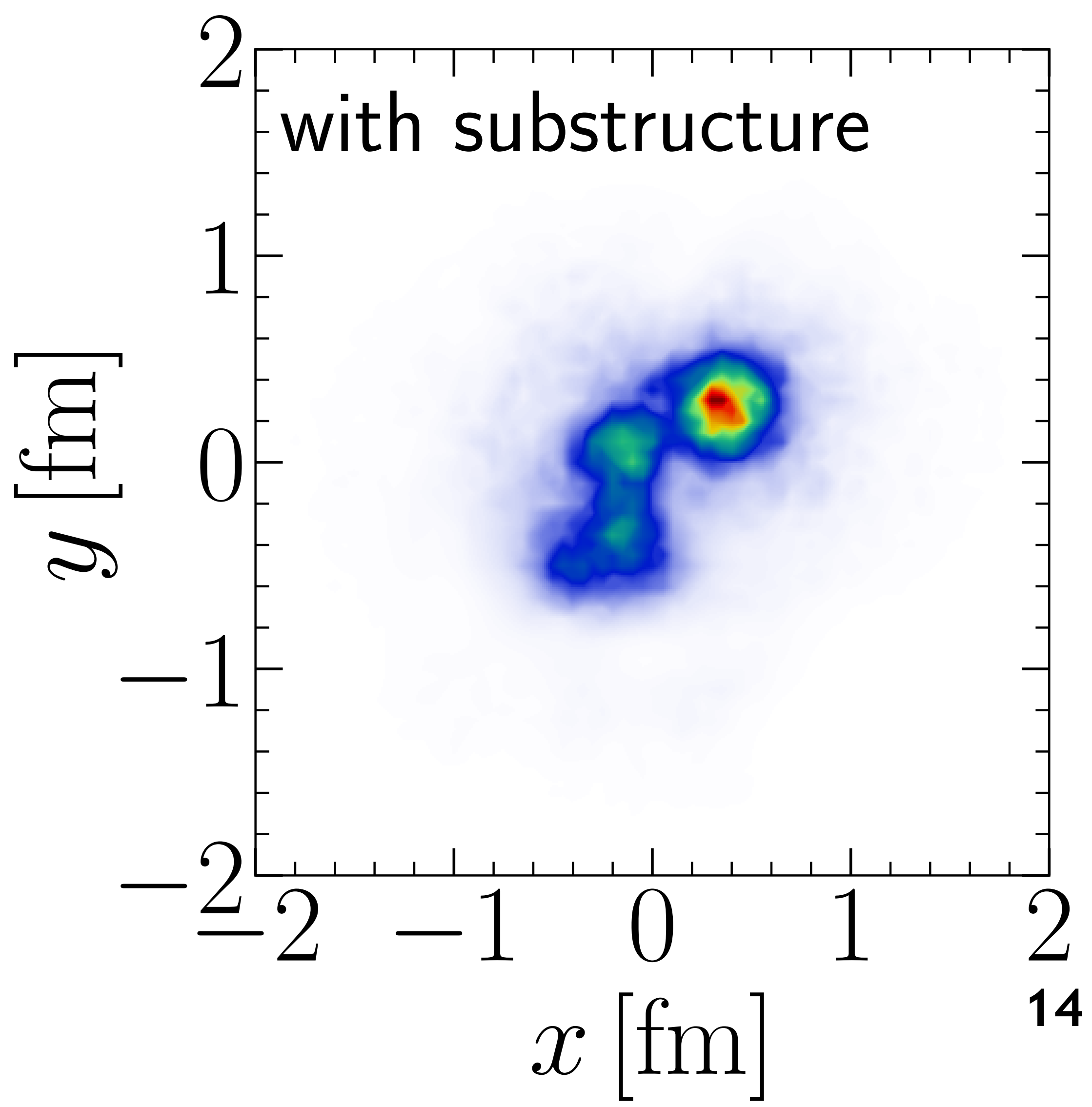
also see:

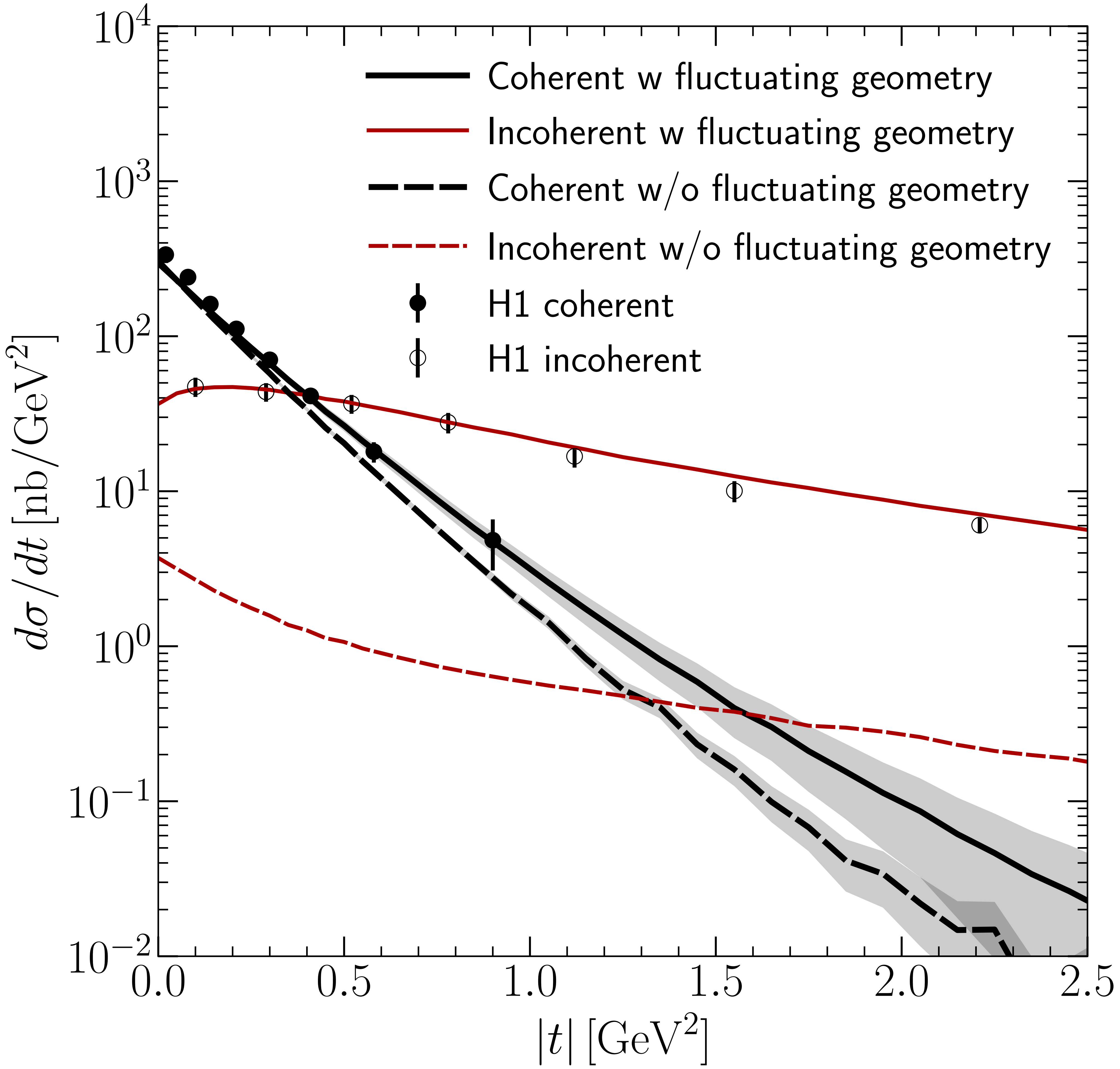
S. Schlichting, B. Schenke, Phys.Lett. B739 (2014) 313-319

H. Mäntysaari, Rep. Prog. Phys. 83 082201 (2020)

B. Schenke, Rep. Prog. Phys. 84 082301 (2021)

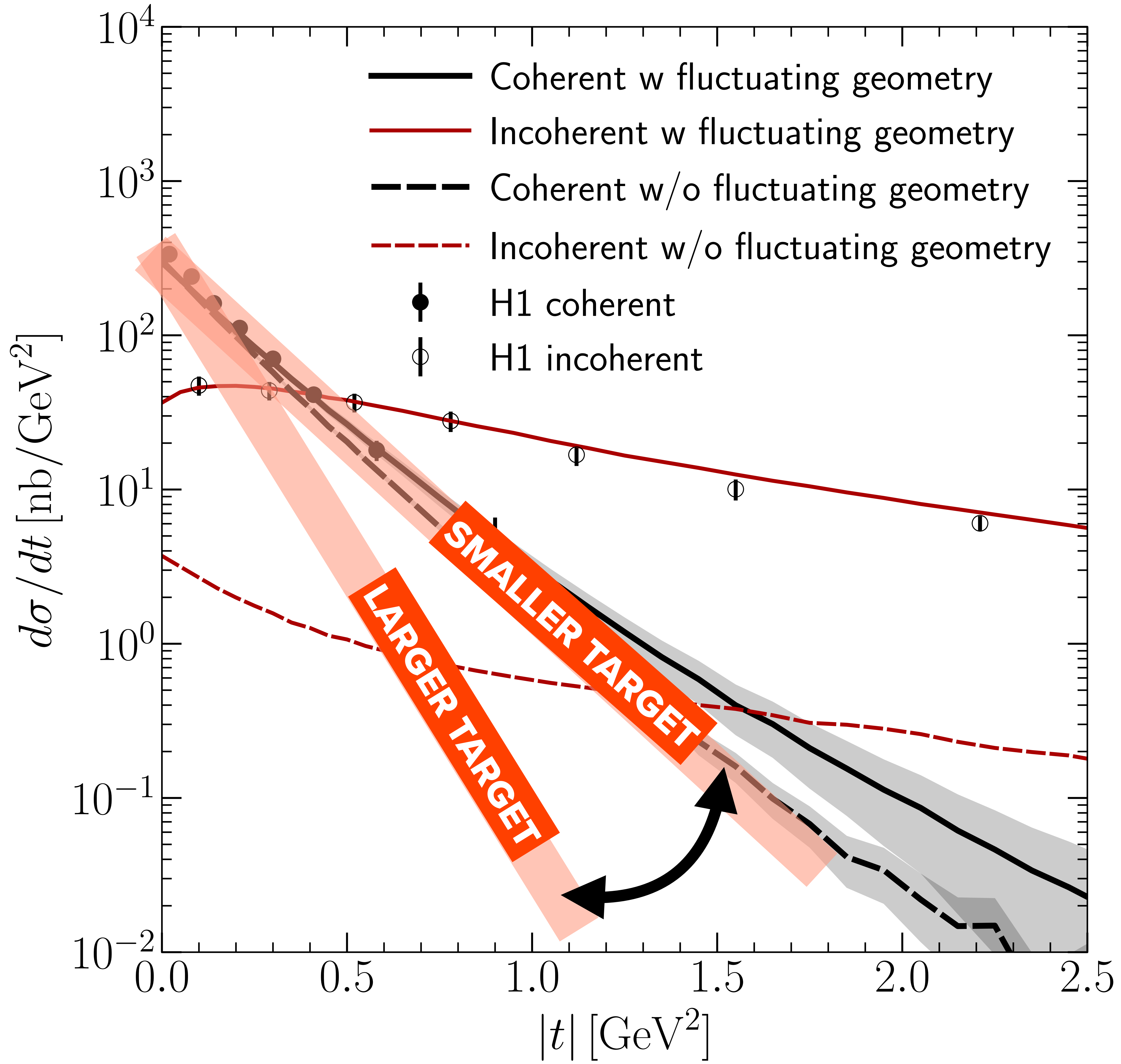






H1 Collaboration, Eur. Phys. J. C73 (2013) no. 6 2466

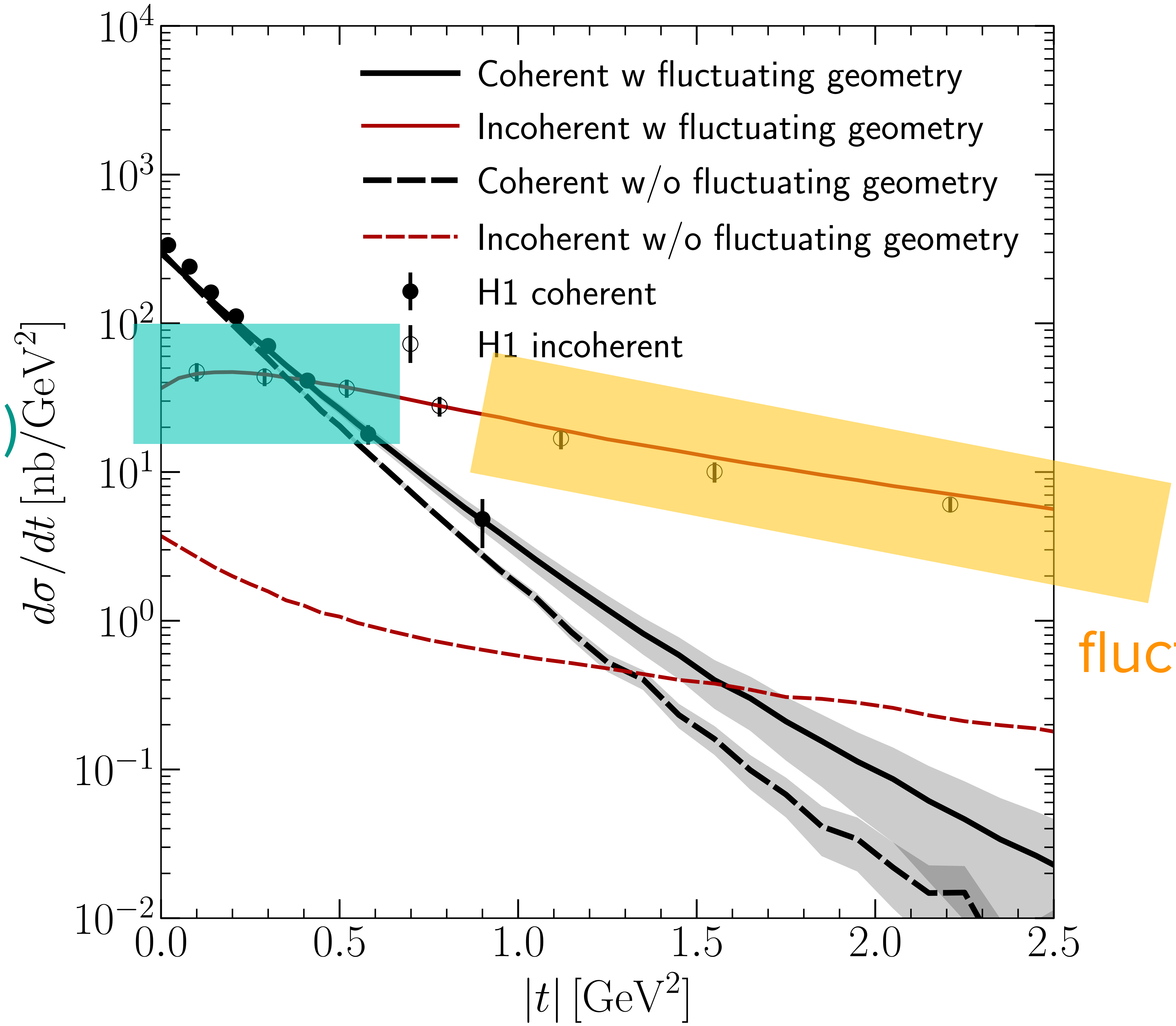
Information in the diffractive cross sections



H1 Collaboration, Eur. Phys. J. C73 (2013) no. 6 2466

Information in the diffractive cross sections

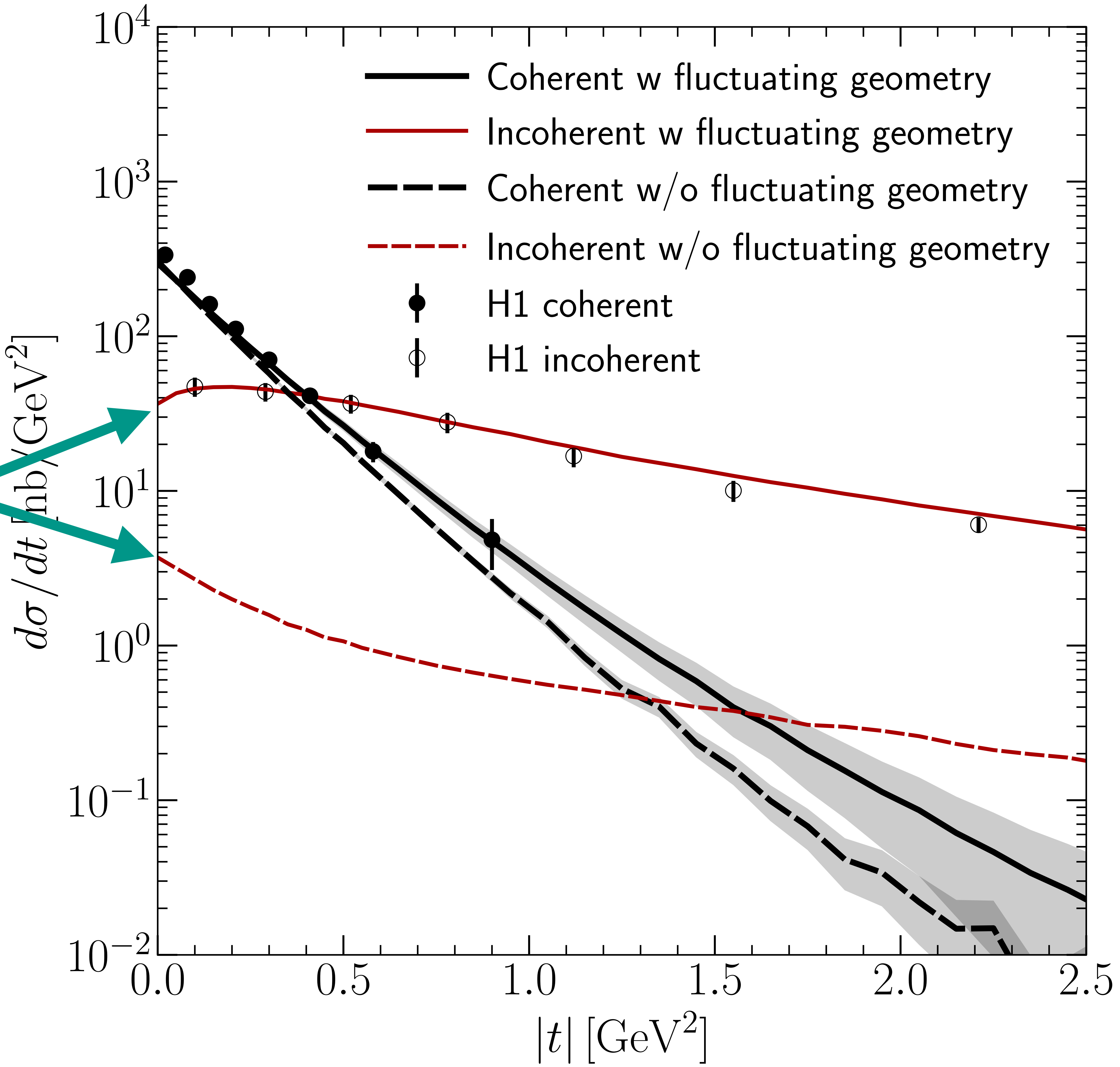
larger scale
fluctuations (>0.2 fm)



short scale
fluctuations (<0.2 fm)

H1 Collaboration, Eur. Phys. J. C73 (2013) no. 6 2466

Information in the diffractive cross sections

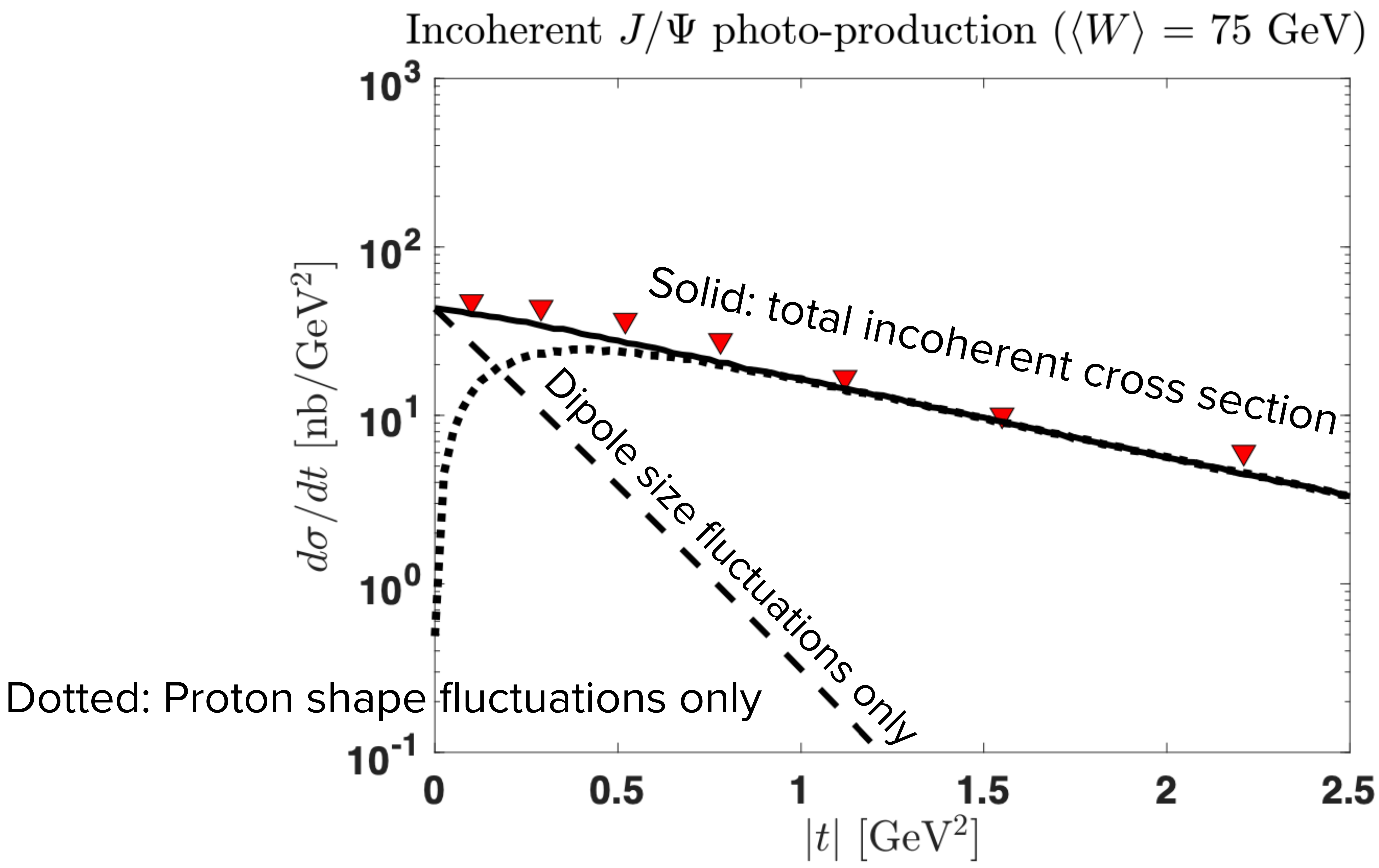


Does not go to zero

H1 Collaboration, Eur. Phys. J. C73 (2013) no. 6 2466

Dipole size fluctuations

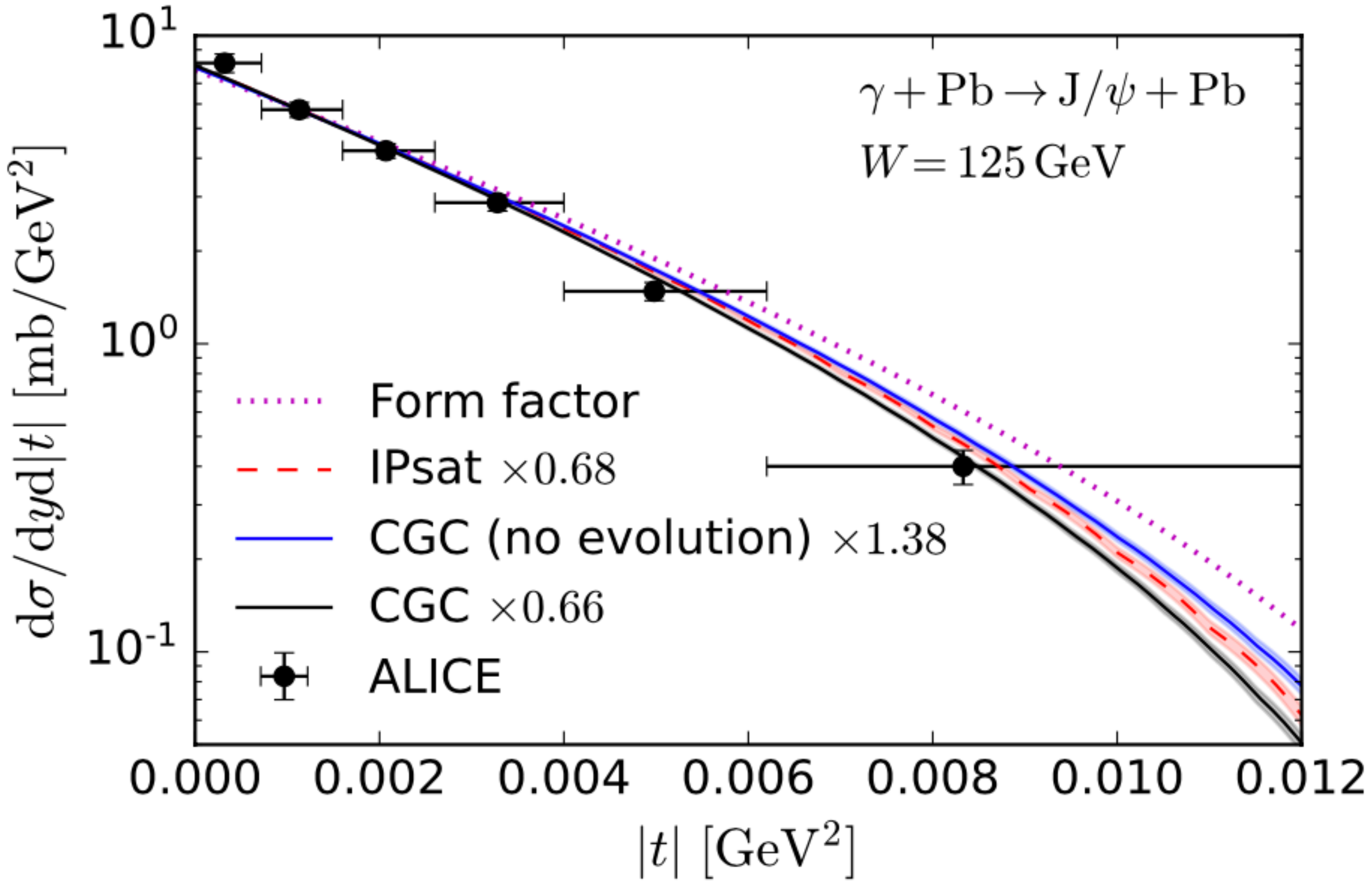
Blaizot and Traini, 2209.15545 [hep-ph]

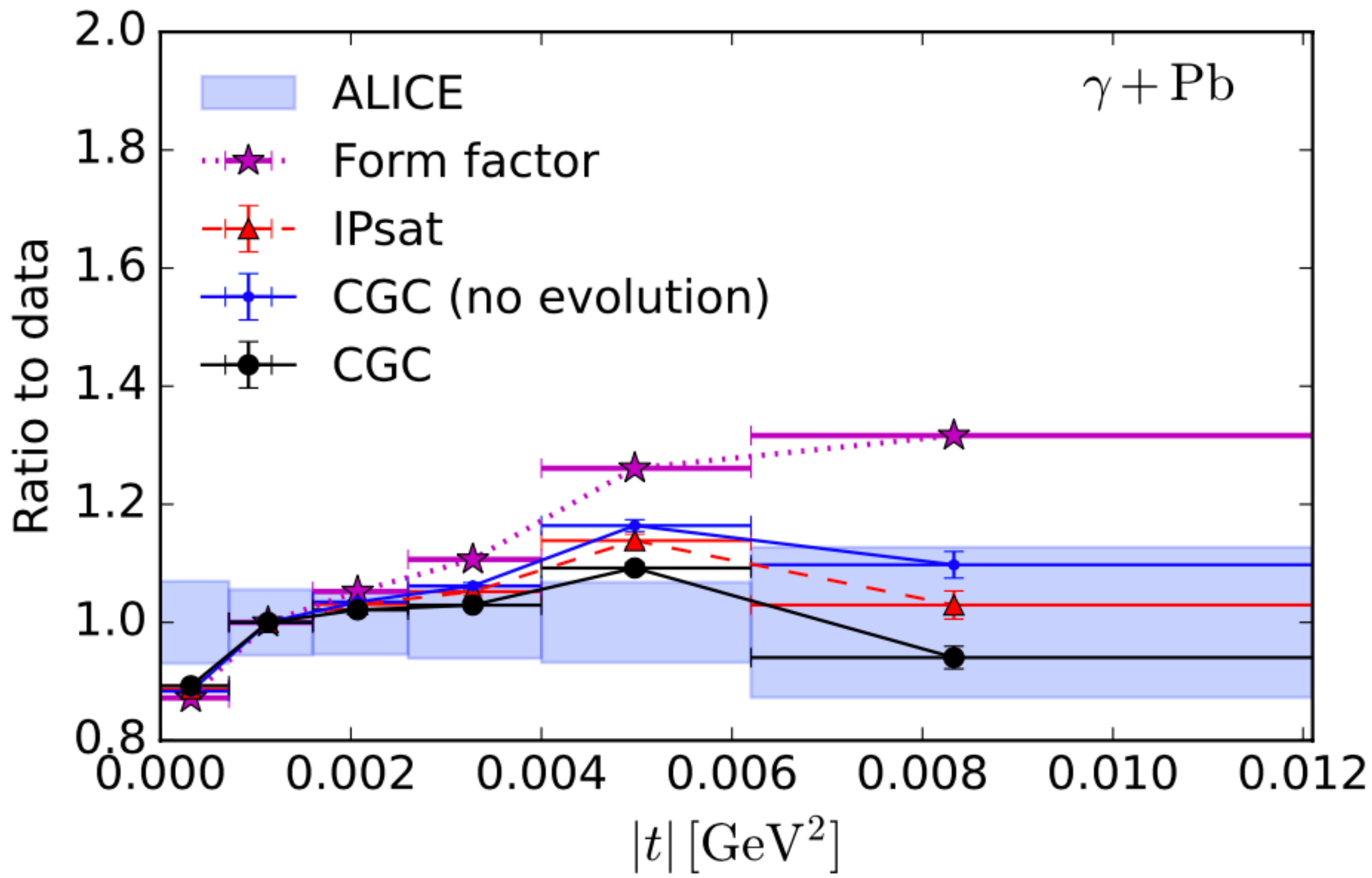


UPCs: γ+Pb measurement - Role of saturation effects

H. Mäntysaari, F. Salazar, B. Schenke, Phys.Rev.D 106 (2022) 7, 074019

Here, ALICE removed interference and photon k_T effects to get the γ+Pb cross section





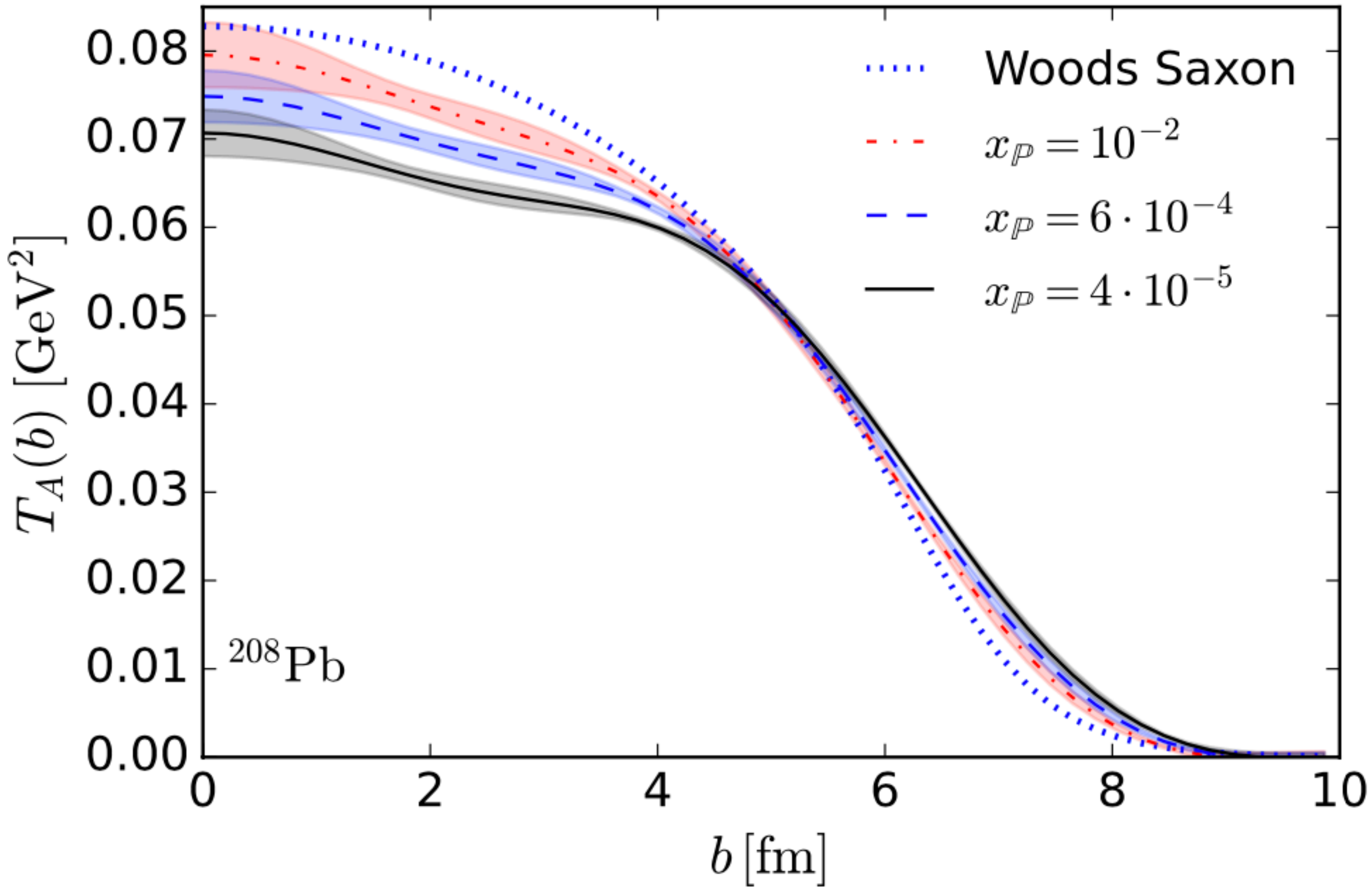
ALICE Collaboration, Phys.Lett.B 817 (2021) 136280

Saturation effects improve agreement with experimental data significantly

Saturation effects on nuclear geometry

H. Mäntysaari, F. Salazar, B. Schenke, Phys.Rev.D 106 (2022) 7, 074019

Fourier transform to coordinate space



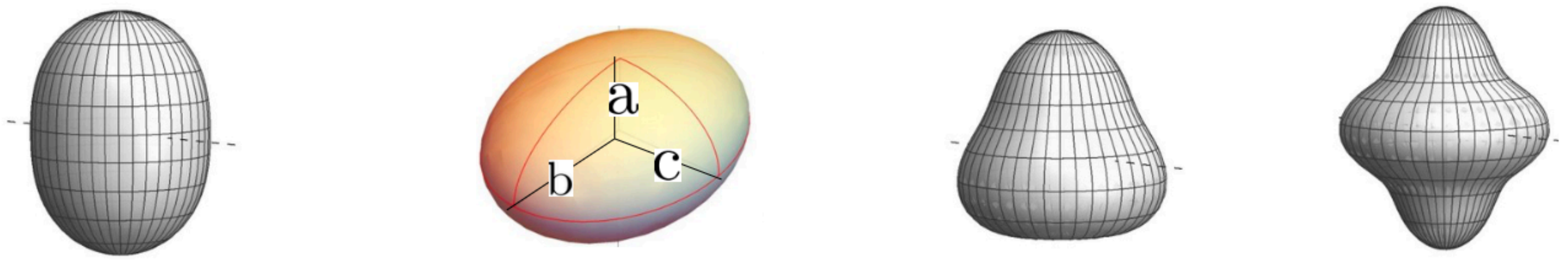
JIMWLK evolution leads to growth of the nucleus towards small x and depletion near the center (normalized so $\int d^2 b T_A(b) = 208$)

Effects of deformation on diffractive cross sections

H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, in progress

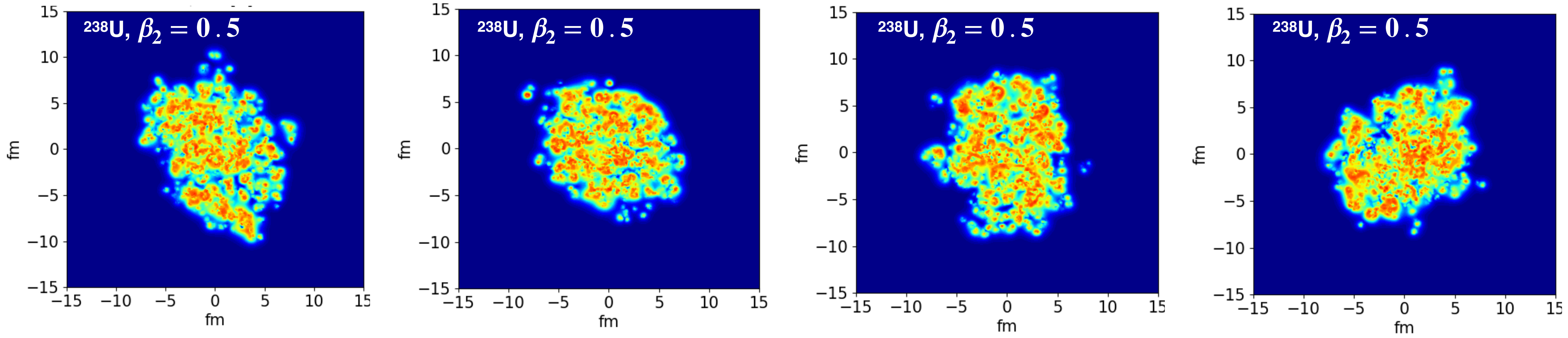
Implement deformation in the Woods-Saxon distribution:

$$\rho(r,\Theta,\Phi) \propto \frac{1}{1+\exp\left(\left[r-R(\Theta,\Phi)\right]/a\right)} \;, \quad R(\Theta,\Phi) = R_0\left[1+\underline{\beta_2}\left(\cos\underline{\gamma} Y_{20}(\Theta)+\sin\gamma Y_{22}(\Theta,\Phi)\right)+\underline{\beta_3} Y_{30}(\Theta)+\underline{\beta_4} Y_{40}(\Theta)\right]$$



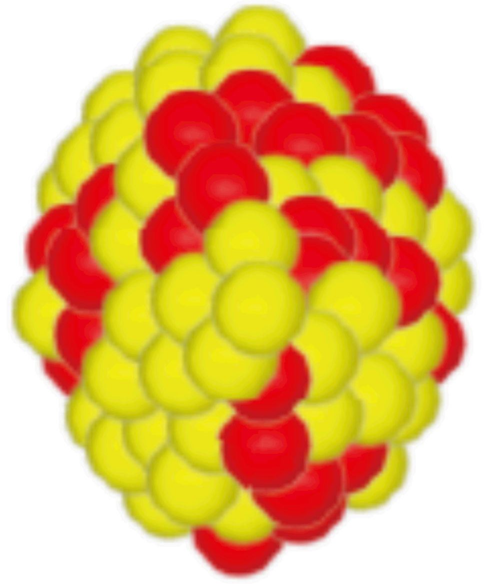
from G. Giacalone

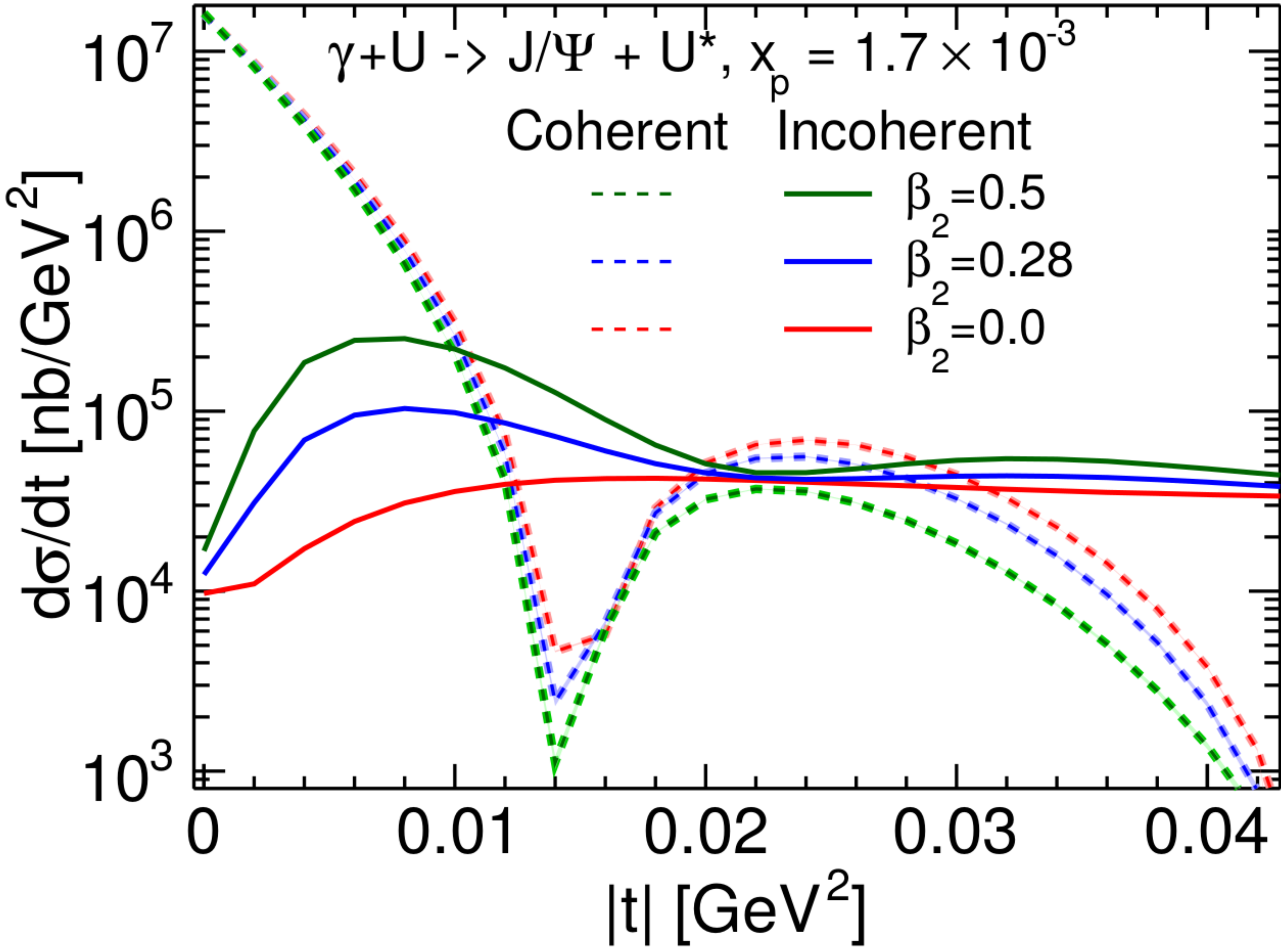
Deformed nuclei exhibit larger fluctuation in the transverse projection:



Effects of deformation on diffractive cross sections: Uranium

H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, in progress

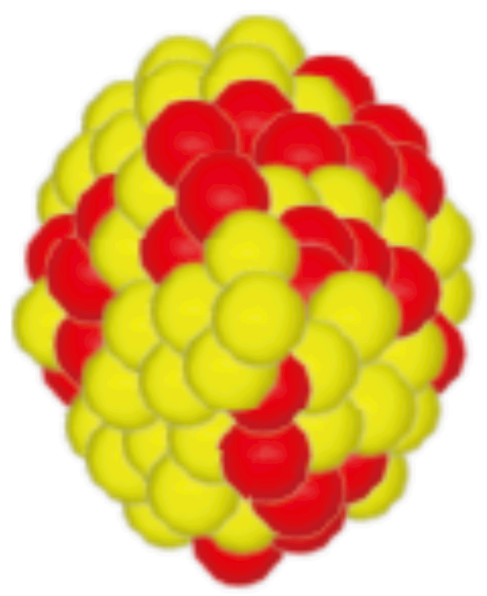




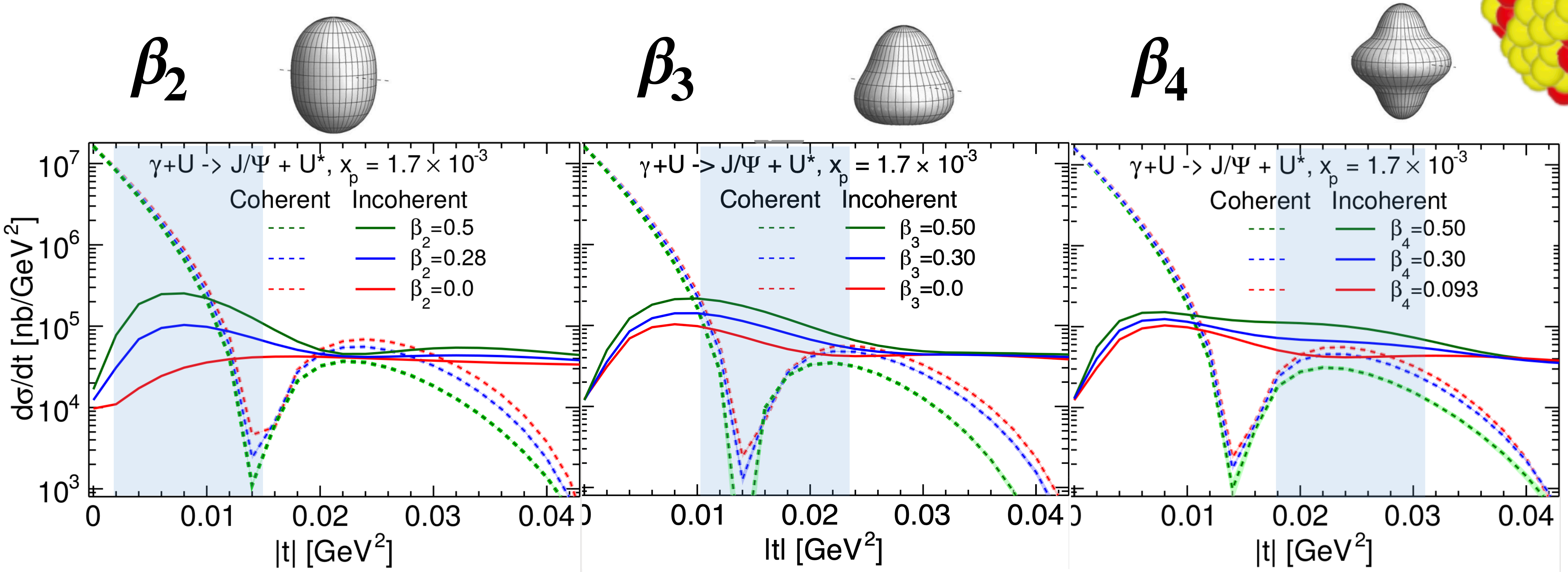
Deformation of the nucleus affects incoherent cross section at small $|t|$ (large length scales)

This observable provides direct information on the small x structure

Effects of deformation on diffractive cross sections: Uranium



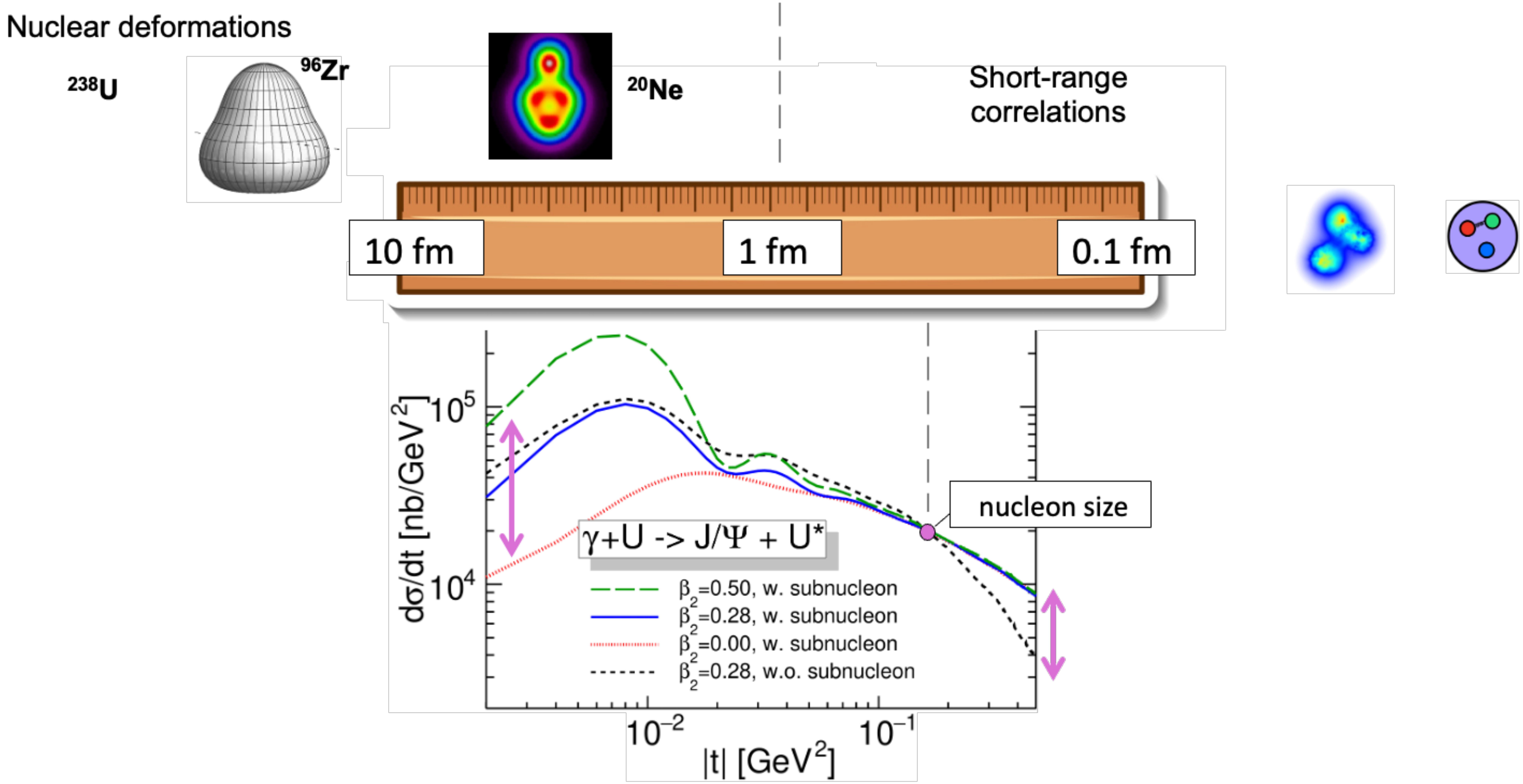
H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, arXiv:2303.04866



- β_2 , β_3 and β_4 modify fluctuations at different length scales: Change incoherent cross section in different |t| regions
- Different values of deformation do not affect the location of the first minimum of the coherent cross sections (average size remains the same)

Multi-scale sensitivity

H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, arXiv:2303.04866



Chiral effective field theory (low-energy QCD) ⟷ **?** ⟷ **CGC effective field theory (high-energy QCD)**

slide from G. Giacalone

Some points for discussion

- Coherent production in event with breakup:
  - We assume that we have clean coherent diffraction: over the course of the interaction, the nucleus remains in its ground state.
    So, as Spencer asked, why do we see a coherent scattering signal in events where the target clearly broke up?
  - What are the time scales? The excitation could happen long after the scattering, not affecting the fact that it was coherent (can it happen way before and the scattering happen with the excited nucleus? probably not, as there is no time for that to happen).
- Small $|t|$:
  - Miettinen and Pumplin say: "We clarified the reason for the *catastrophic failure* of the additive quark models (relativistic as well as nonrelativistic) in predicting the |t| dependence of diffractive production." Does the model fail? Does all data show the lack of a dip towards $|t| \to 0$?
  - Spencer says: "As |t| decreases, the energy transfer to the nucleus decreases, and, as $|t| \to 0$ there is insufficient energy transferred to excite the nucleus, so incoherent interactions become impossible."

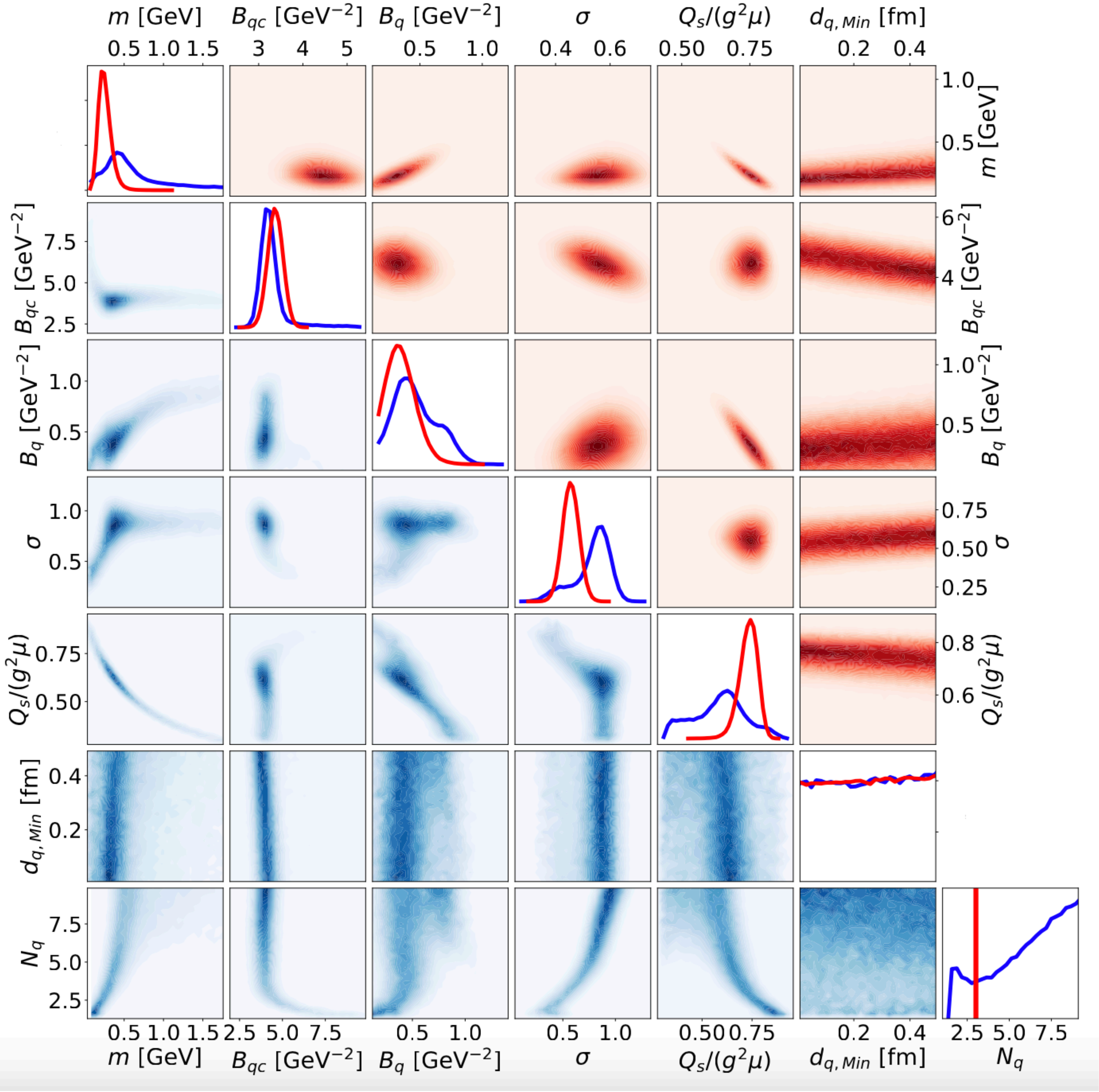
SUMMARY

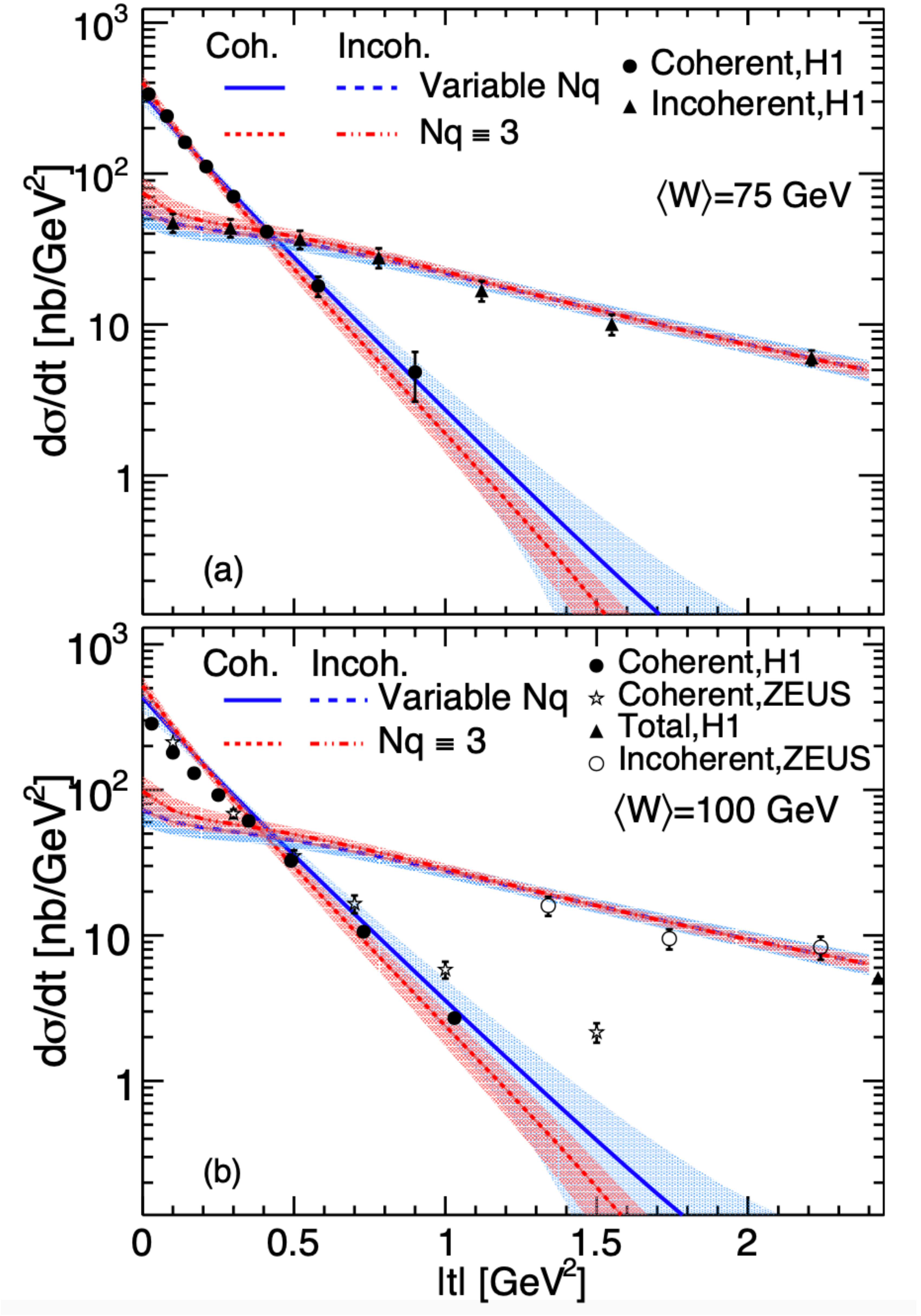
- Scattering amplitude for diffractive vector meson production:
  - Color singlet final state (color trace on the level of the amplitude)
  - Coherent: Target average on the level of the amplitude
  - Total diffractive: Target average on the level of the cross section
- $|t|$-differential incoherent cross section is sensitive to fluctuations at different length scales: Strong effect of deformation, nucleon, and sub-nucleon fluctuations
- Low $|t|$ incoherent cross section does not go to zero - number (or normalization) fluctuations (also dipole size fluctuations)

BACKUP

Extracting parameters using Bayesian inference

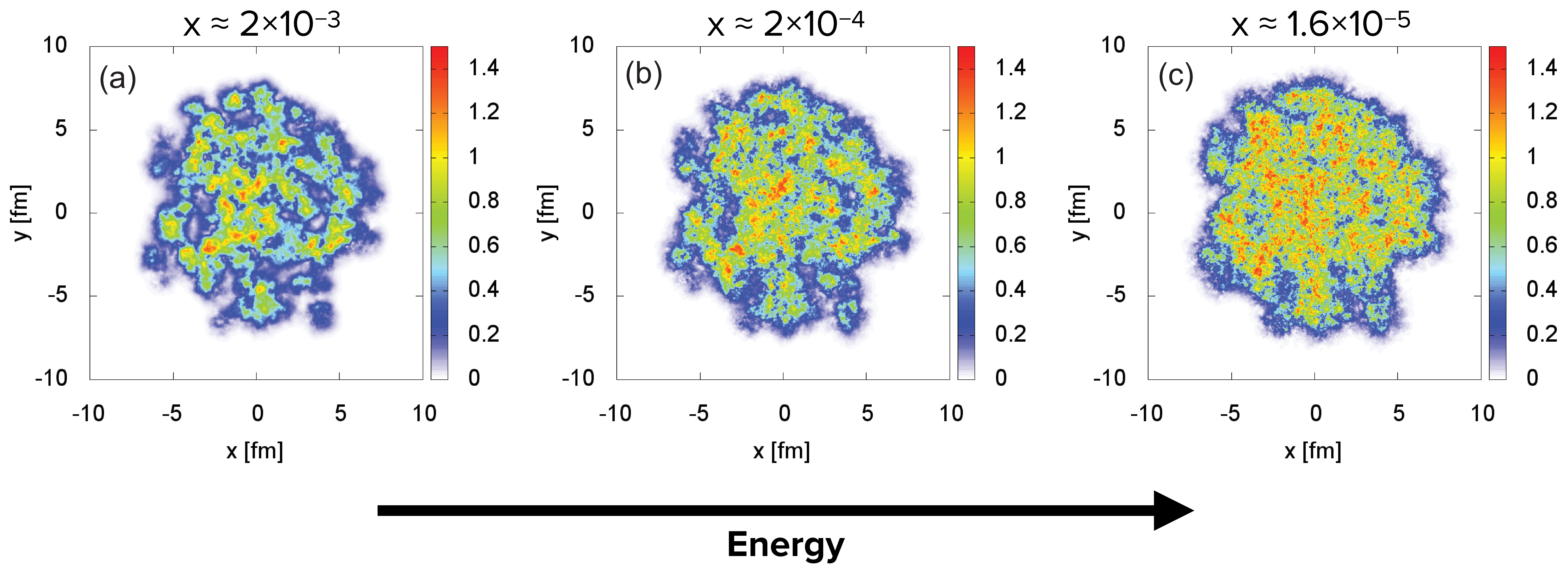
H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, Phys.Lett.B 833 (2022) 137348





Questions:

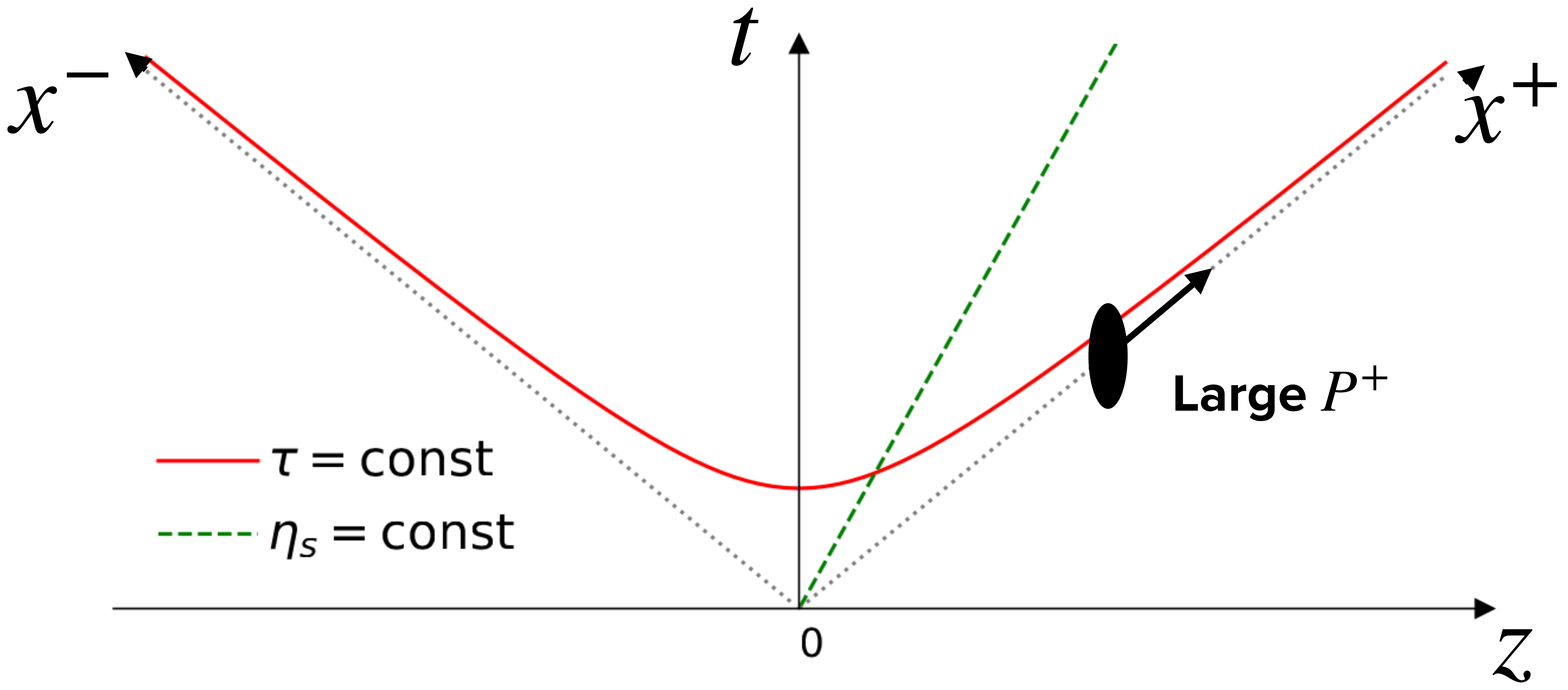
How does energy evolution affect the nuclear structure?



Are observables in high energy e+A scattering sensitive to nuclear deformation?

B. Schenke, S. Schlichting, Phys.Rev.C 94 (2016) 4, 044907

The scenario: Hadron moving at high momentum



Probe hadron (or nucleus) moving with large P^+ at scale $x_0 P^+$ with $x_0 \ll 1$

Separate partonic content based on longitudinal momentum $k^+ = xP^+$

Large $x > x_0$: Static and localized color sources ρ

Dynamic color fields

The moving color sources generate a current, independent of light cone time z^+:

$$J^{\mu,a}(z) = \delta^{\mu +}\rho^a(z^-, z_T)$$

a is the color index of the gluon

This current generates delocalized dynamical fields $A^{\mu,a}(z)$ described by the Yang-Mills equations

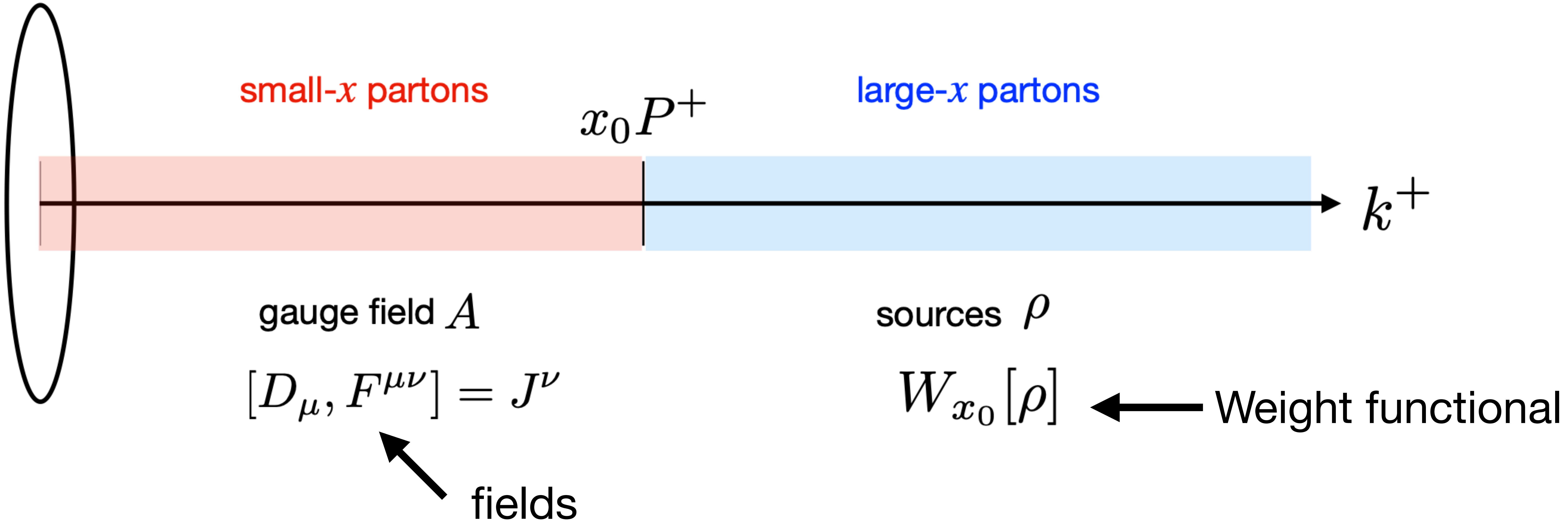
$$[D_\mu, F^{\mu\nu}] = J^\nu$$

with $D_\mu = \partial_\mu + igA_\mu$ **and** $F_{\mu\nu} = \frac{1}{ig}[D_\mu, D_\nu] = \partial_\mu A_\nu - \partial_\nu A_\mu + ig[A_\mu, A_\nu]$

These fields A are the small $x < x_0$ degrees of freedom

They can be treated classically, because their occupation number is large $\langle AA \rangle \sim 1/\alpha_s$

Color Glass Condensate (CGC): Sources and fields



When $x \lesssim x_0$ the path integral $\langle \mathcal{O} \rangle_\rho$ is dominated by classical solution and we are done

For smaller x we need to do quantum evolution

Wilson lines

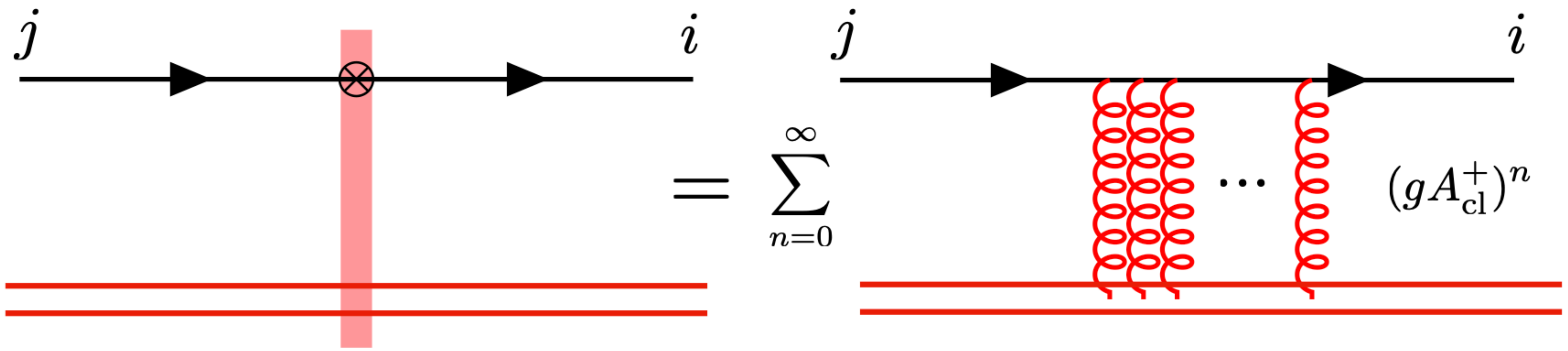
Interaction of high energy color-charged probe with large k^- momentum (and small $k^+ = \frac{k_T^2}{2k^-}$)

with the classical field of a nucleus can be described in the **eikonal approximation:**

The scattering rotates the color, but keeps k^-, transverse position $\vec{x}_T$, and any other quantum numbers the same.

The color rotation is encoded in a light-like Wilson line, which for a quark probe reads

$$V_{ij}(\vec{x}_T) = \mathscr{P}\left(ig \int_{-\infty}^{\infty} A^{+,c}(z^-, \vec{x}_T)\, t^c_{ij}\, dz^- \right)$$



MULTIPLE INTERACTIONS NEED TO BE RESUMMED, BECAUSE $A^+ \sim 1/g$

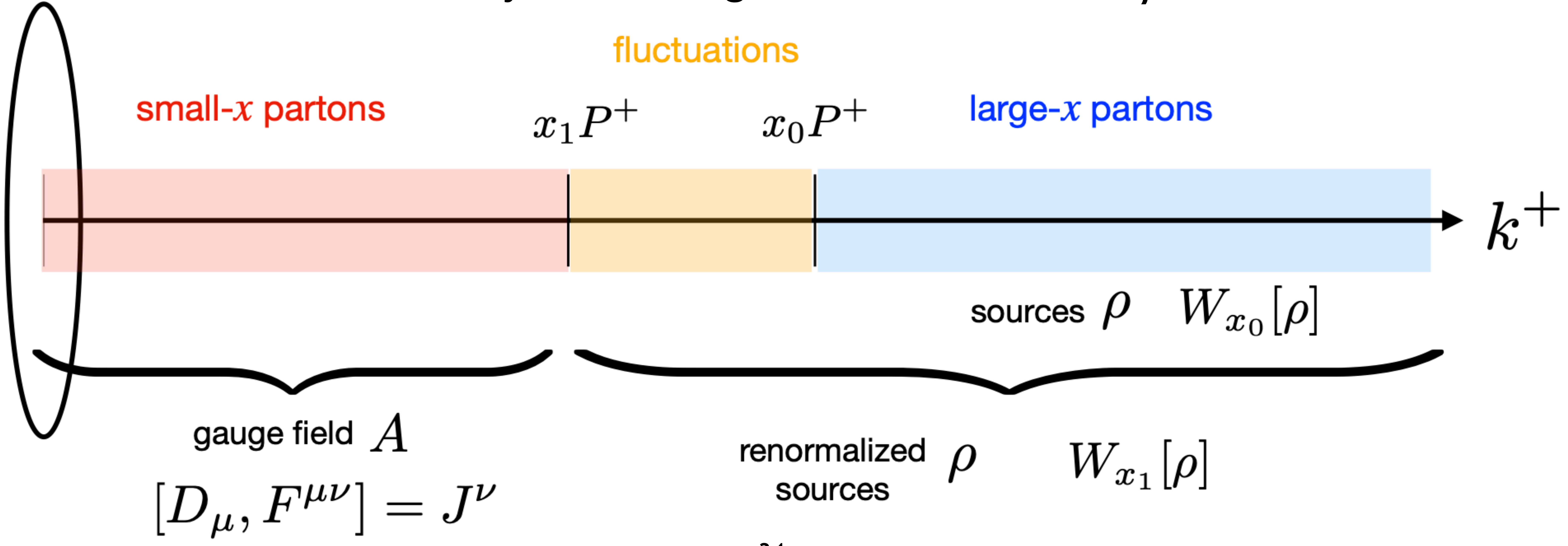
JIMWLK evolution

Jalilian-Marian, J.; Kovner, A.; McLerran, L.D.; Weigert, H., Phys. Rev. D 1997, 55, 5414–5428, [hep-ph/9606337]
Jalilian-Marian, J.; Kovner, A.; Weigert, H., Phys. Rev. D 1998, 59, 014015, [hep-ph/9709432]
Kovner, A.; Milhano, J.G.; Weigert, H., Phys. Rev. D 2000, 62, 114005, [hep-ph/0004014]
Iancu, E.; Leonidov, A.; McLerran, L.D., Nucl. Phys. A 2001, 692, 583–645,[hep-ph/0011241]
Iancu, E.; Leonidov, A.; McLerran, L.D., Phys. Lett. B 2001, 510, 133–144, [hep-ph/0102009]
Ferreiro, E.; Iancu, E.; Leonidov, A.; McLerran, L., Nucl. Phys. A 2002, 703, 489–538, [hep-ph/0109115]

LO Small-x evolution resums logarithmically enhanced terms $\sim \alpha_s \ln(x_0/x)$

$$\frac{dW_x[\rho]}{d\ln(1/x)} = -\mathcal{H}_{\text{JIMWLK}} W_x[\rho]$$

Physically, one absorbs the quantum fluctuations in the interval $[x_0 - dx, x_0]$ into stochastic fluctuations of the color sources by redefining the color sources ρ



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Evolution is done using the Langevin formulation of the JIMWLK equations on the level of Wilson lines

K. Rummukainen and H. Weigert Nucl. Phys. A739 (2004) 183; T. Lappi, H. Mäntysaari, Eur. Phys. J. C73 (2013) 2307

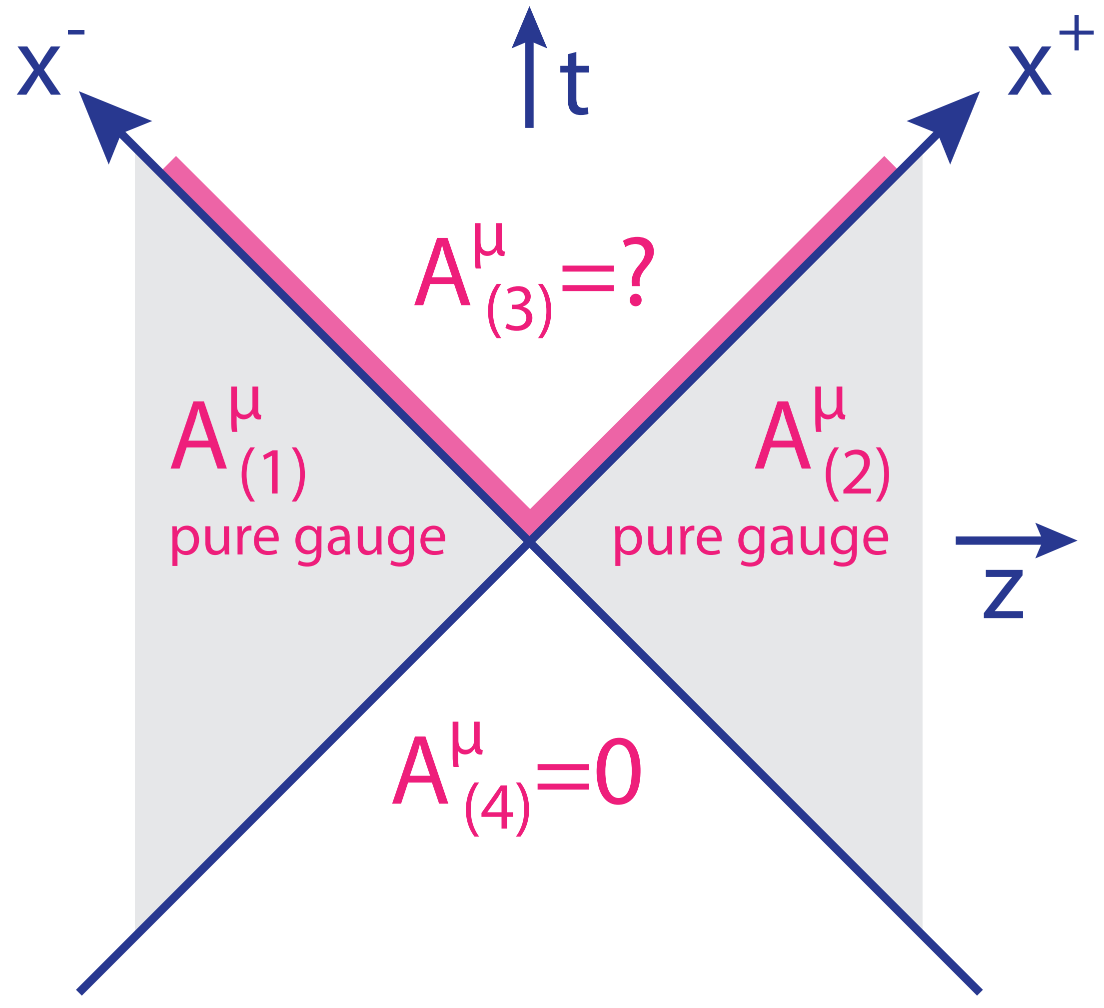
Long distance tales are tamed by imposing a regulator in the JIMWLK kernel, *m*

S. Schlichting, B. Schenke, Phys.Lett. B739 (2014) 313-319

Connection between the initial state of heavy ion collisions and the EIC

- These Wilson lines are the building blocks of the CGC
- In heavy ion collisions, one can compute the initial state by determining Wilson lines after the collision from the Wilson lines of the colliding nuclei
- At the EIC (and HERA, and in UPCs), cross sections will be calculated as convolutions of Wilson line correlators with perturbatively calculable and process-dependent impact factors
- This allows the computation of rather direct constraints for the initial state of heavy ion collisions from electron-nucleus (γ-nucleus) or electron-proton collisions

Heavy ion collision



Compute gluon fields after the collision using light cone gauge:
$A^+ = 0$ for a right moving nucleus, $A^- = 0$ for a left moving nucleus

gauge transformation: $A_\mu(x) \to V(x)\left(A_\mu(x) - \frac{i}{g}\partial_\mu\right)V^\dagger(x)$

using our Wilson lines $V^\dagger(x^-, \mathbf{x}_\perp) = \mathscr{P}\exp\left(-ig\int_{-\infty}^{x^-} dz^- A^+(z^-, \mathbf{x}_\perp)\right)$ (for the right moving nucleus)

Then, the gauge fields read (choosing $A^\mu = 0$ for the quadrant for $x^- < 0$ and $x^+ < 0$)

$$A^i(x) = \theta(x^+)\theta(x^-)\alpha^i(\tau, \mathbf{x}_\perp) + \theta(x^-)\theta(-x^+)\alpha^i_P(\mathbf{x}_\perp) + \theta(x^+)\theta(-x^-)\alpha^i_T(\mathbf{x}_\perp)$$

$$A^\eta(x) = \theta(x^+)\theta(x^-)\alpha^\eta(\tau, \mathbf{x}_\perp)$$

with $\alpha^i_P(\mathbf{x}_\perp) = \frac{1}{ig}V_P(\mathbf{x}_\perp)\partial^i V^\dagger_P(\mathbf{x}_\perp)$ and $\alpha^i_T(\mathbf{x}_\perp) = \frac{1}{ig}V_T(\mathbf{x}_\perp)\partial^i V^\dagger_T(\mathbf{x}_\perp)$

$A^\tau = 0$, because we chose Fock-Schwinger gauge $x^+A^- + x^-A^+ = 0$

Heavy ion collision

Plugging this ansatz

$$A^i(x) = \theta(x^+)\theta(x^-)\alpha^i(\tau, \mathbf{x}_\perp) + \theta(x^-)\theta(-x^+)\alpha_P^i(\mathbf{x}_\perp) + \theta(x^+)\theta(-x^-)\alpha_T^i(\mathbf{x}_\perp)$$

$$A^\eta(x) = \theta(x^+)\theta(x^-)\alpha^\eta(\tau, \mathbf{x}_\perp)$$

into YM equations leads to singular terms on the boundary from derivatives of θ-functions

Requiring that the singularities vanish leads to the solutions

$$\alpha^i = \alpha_P^i + \alpha_T^i \qquad \alpha^\eta = -\frac{ig}{2}\left[\alpha_{Pj}, \alpha_T^j\right] \qquad \begin{aligned}\partial_\tau \alpha^i &= 0 \\ \partial_\tau \alpha^\eta &= 0\end{aligned}$$

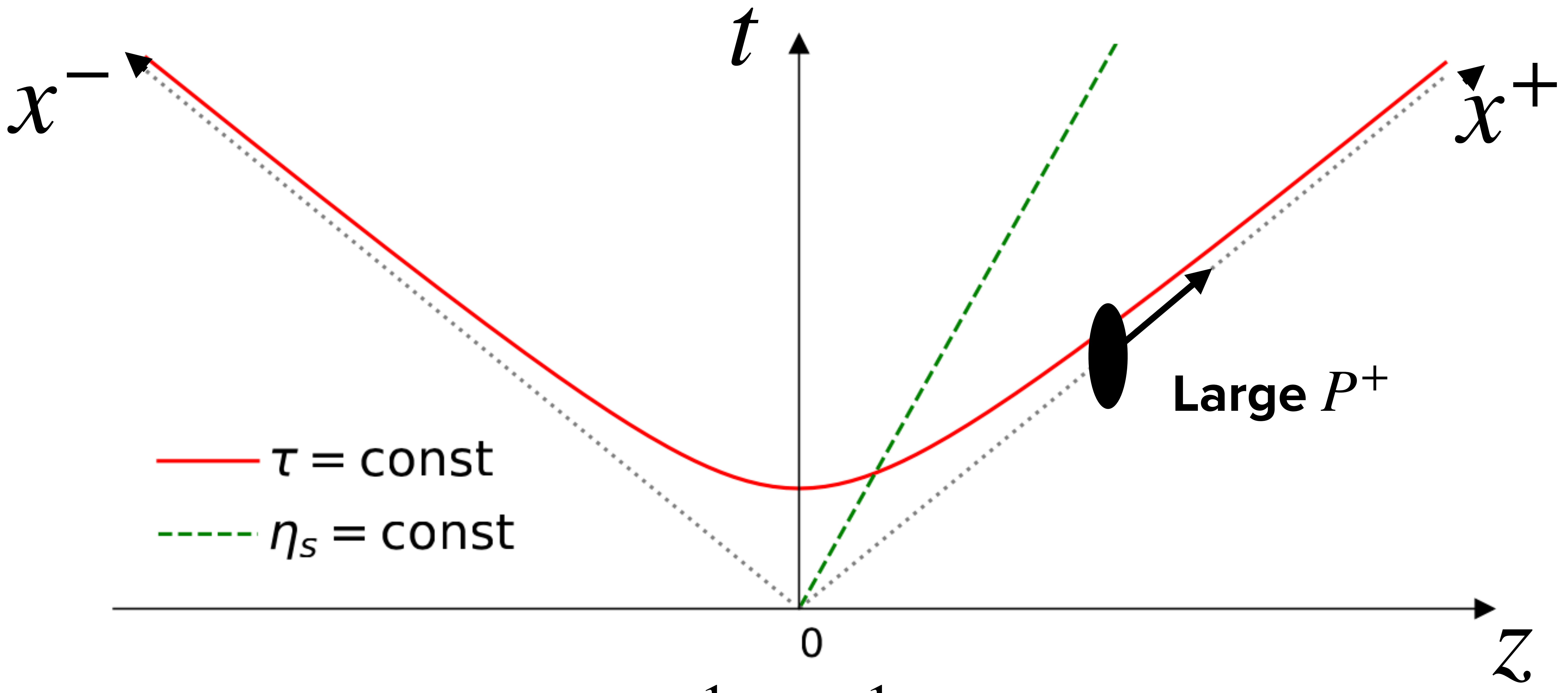
These are the gauge fields in the forward light cone.

We can compute $T^{\mu\nu}$ from it, providing an initial condition for hydrodynamics.

Geometry, fluctuations, ...

- All the information on geometry and nucleon and sub-nucleon fluctuations is contained in the distribution of color charges $\rho^a_{P/T}(x^\mp, \mathbf{x}_\perp)$
- Typically, use the MV model, which gives
  $\langle \rho^a(\mathbf{b}_\perp)\rho^b(\mathbf{x}_\perp)\rangle = g^2\mu^2(x, \mathbf{b}_\perp)\delta^{ab}\delta^{(2)}(\mathbf{b}_\perp - \mathbf{x}_\perp)$
- The color charge distribution $g^2\mu(x, \mathbf{b}_\perp)$ depends on the longitudinal momentum fraction x and the transverse position $\mathbf{b}_\perp$. The latter needs to be modeled, the former can be modeled or obtained from e.g. JIMWLK evolution
- We factorize $\mu(x, \mathbf{b}_\perp) \sim T(\mathbf{b}_\perp)\mu(x)$ and constrain the impact parameter $\mathbf{b}_\perp$ dependence using input from a process sensitive to geometry, such as diffractive VM production
- The cross section for that process can be expressed with the Wilson lines of the target
  The same quantities we have used to initialize the heavy ion collision

Color sources

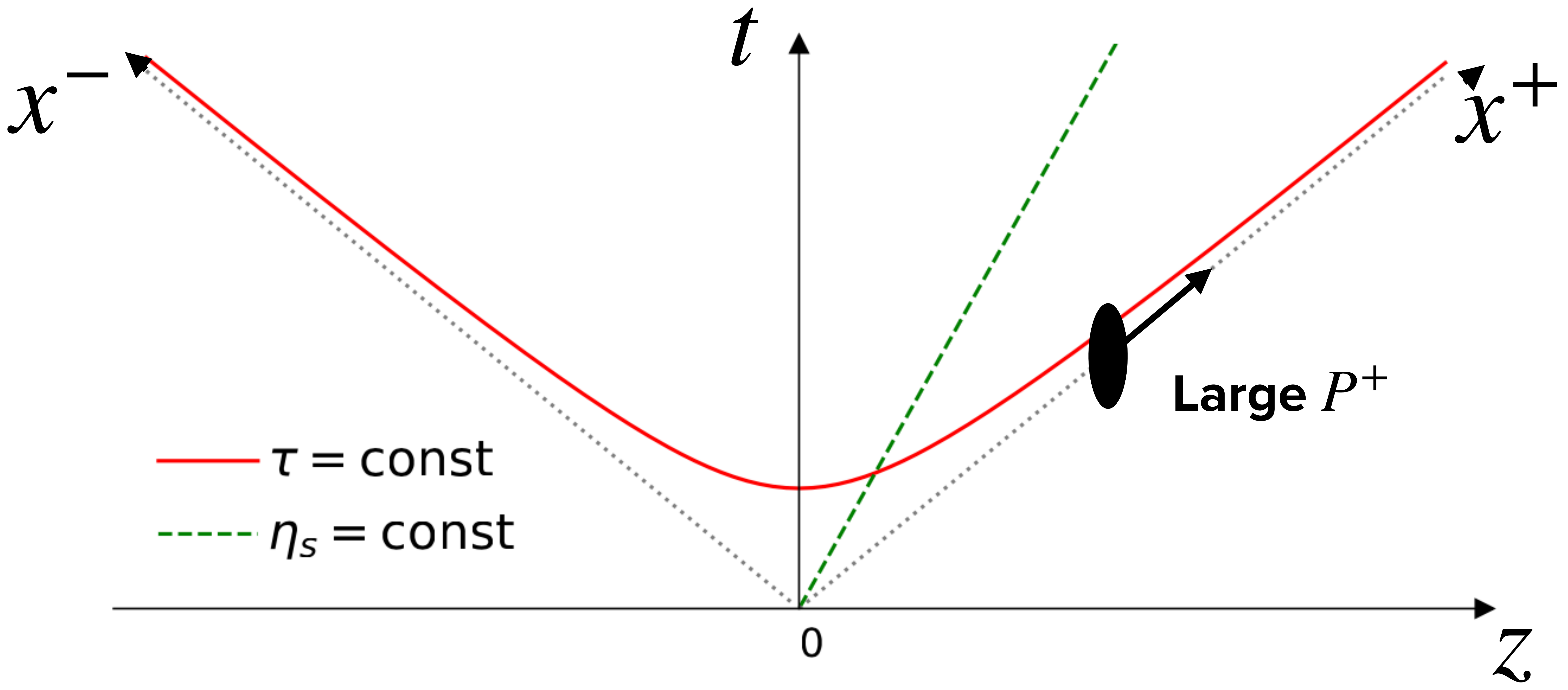


How localized are these sources? $\Delta z^- \sim \frac{1}{k^+} = \frac{1}{xP^+}$

What is the resolution scale of the probe? $\frac{1}{x_0 P^+} > \frac{1}{xP^+}$ **for** $x > x_0$

→Color sources look fully localized to the probe in z^-

Color sources



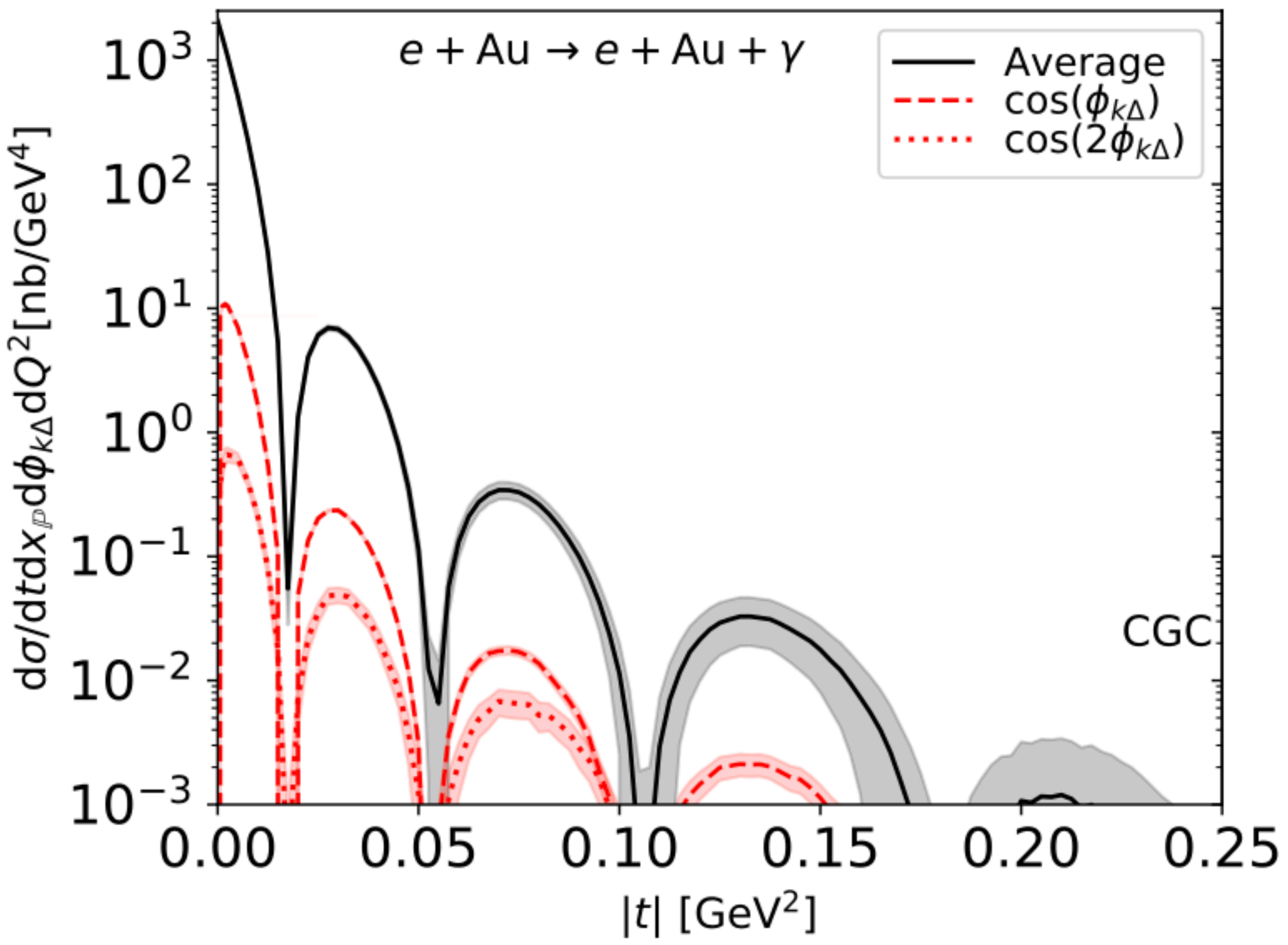
How fast do they evolve? $\Delta z^+ \sim \frac{1}{k^-} = \frac{2k^+}{k_T^2} = \frac{2xP^+}{k_T^2}$ (because $a_\mu b^\mu = a^+b^- + a^-b^+ - \vec{a}_T \cdot \vec{b}_T$)

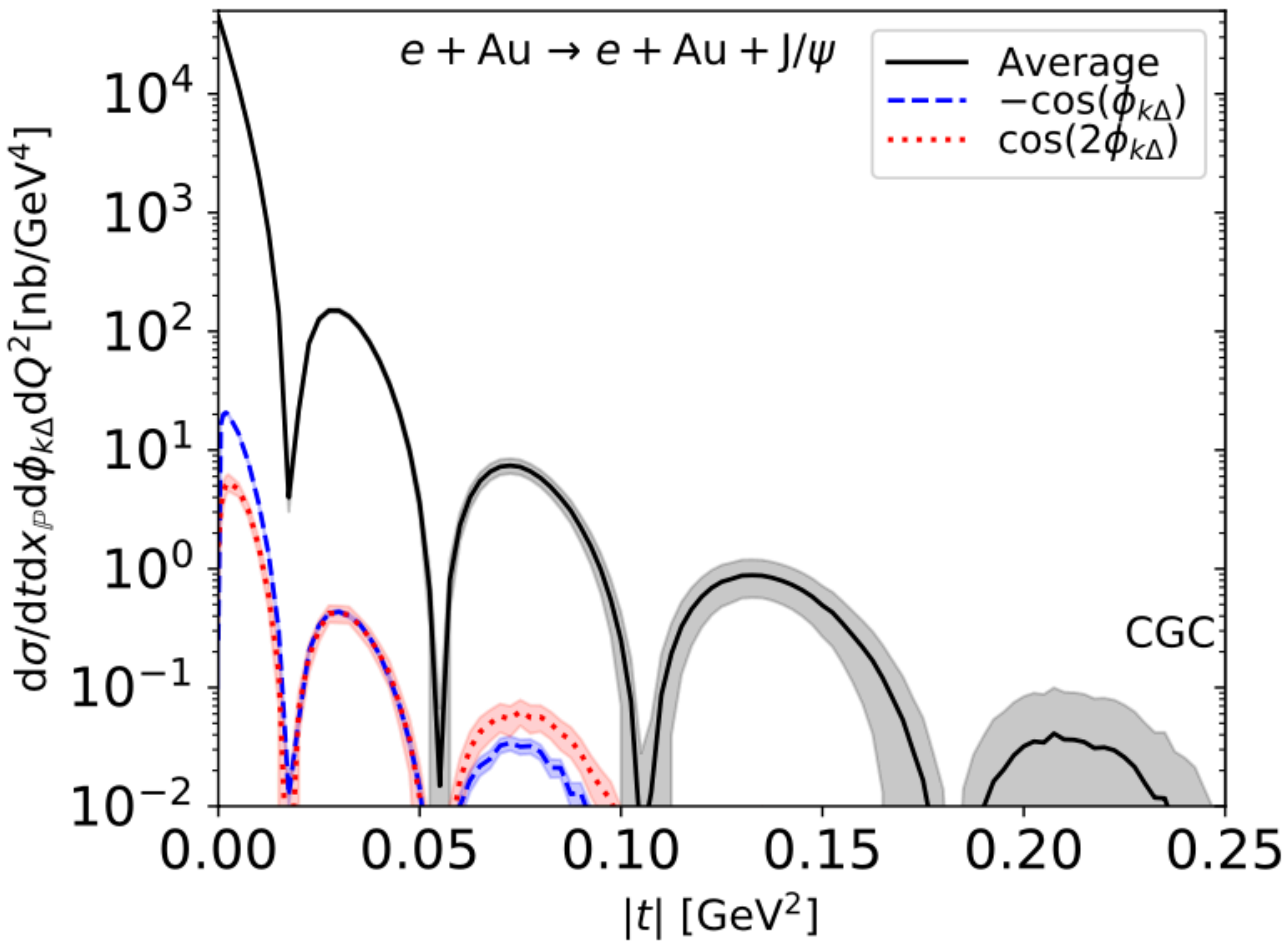
What is the time scale of the probe? $\tau \approx \frac{2x_0P^+}{k_T^2} < \frac{2xP^+}{k_T^2}$

→Color sources look static to the probe in light cone time z^+

Predictions for e-Au at the future EIC

DVCS and exclusive J/ψ : Spectra and azimuthal modulations



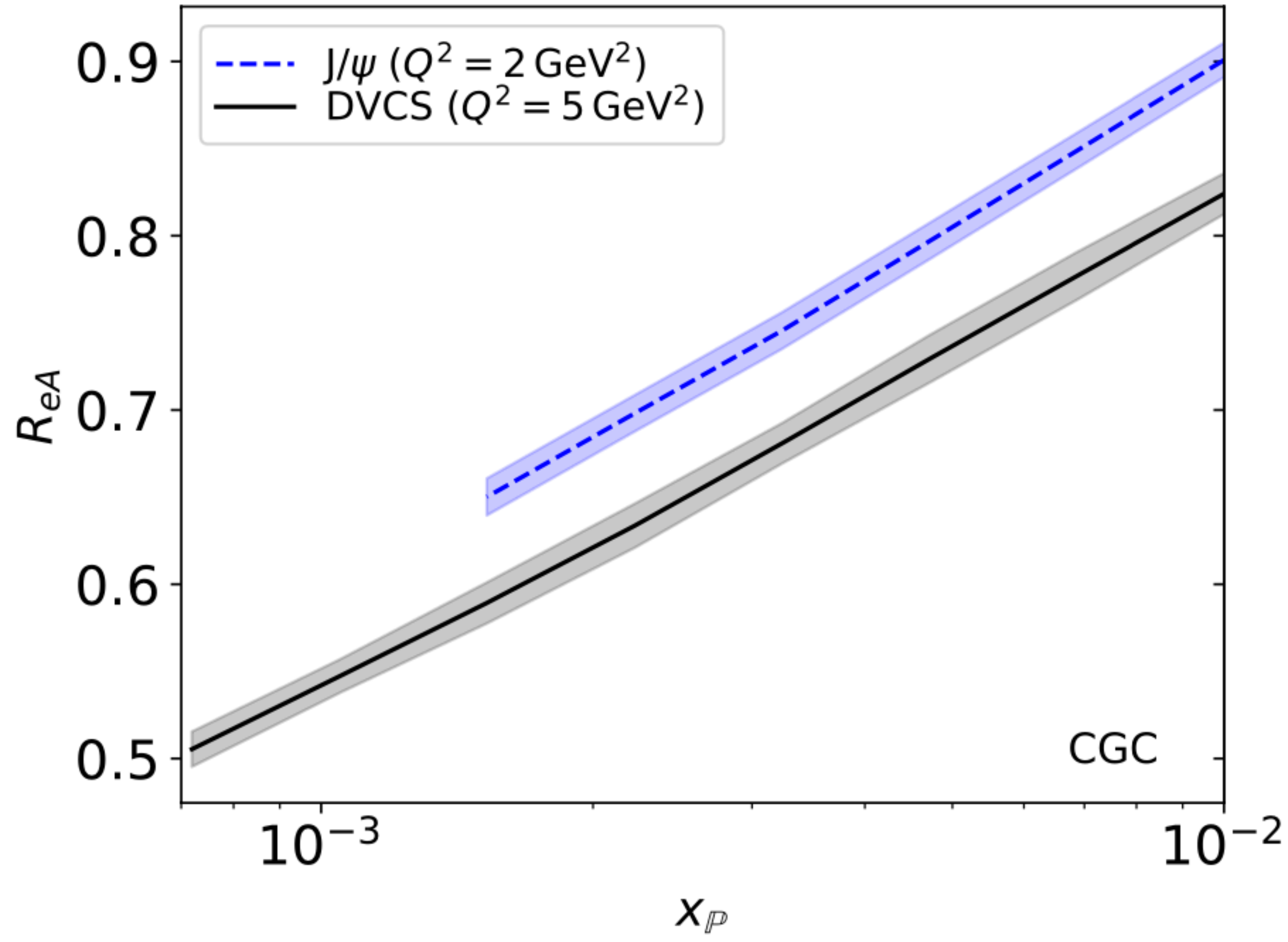


Characteristic dips in spectra due to Woods-Saxon nuclear profile

Azimuthal modulations v_n a few percent for DVCS, and less than 1% for J/ψ

Predictions for e-Au at the future EIC

Nuclear suppressions factor for DVCS and exclusive J/ψ



$$R_{eA} = \left. \frac{\mathrm{d}\sigma^{e+A\to e+A+V}/\mathrm{d}t\mathrm{d}Q^2\mathrm{d}x_{\mathbb{P}}}{A^2\mathrm{d}\sigma^{e+p\to e+p+V}/\mathrm{d}t\mathrm{d}Q^2\mathrm{d}x_{\mathbb{P}}} \right|_{t=0}$$

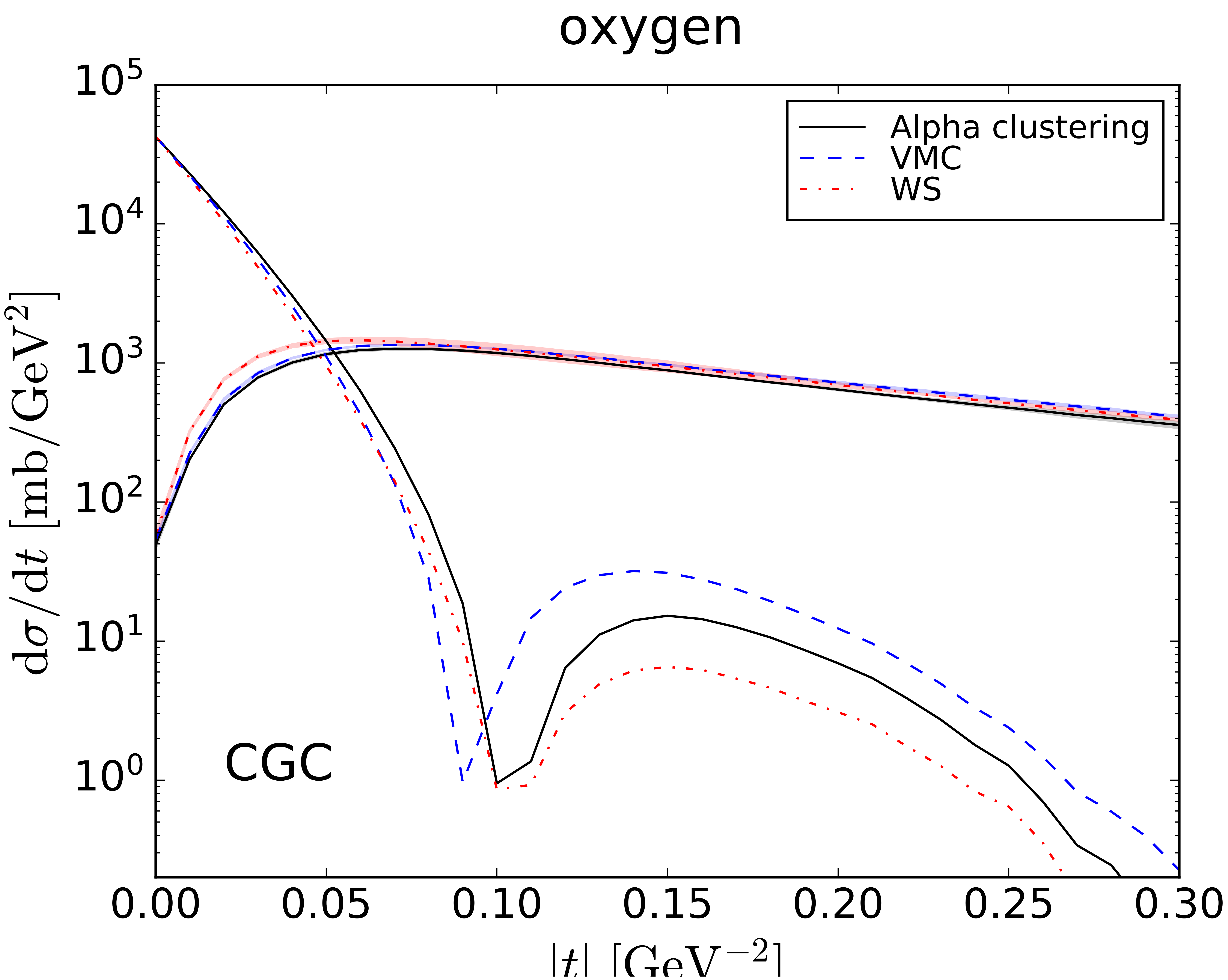
Expect $R_{eA} = 1$ in the dilute limit.

Mäntysaari, Venugopalan. 1712.02508

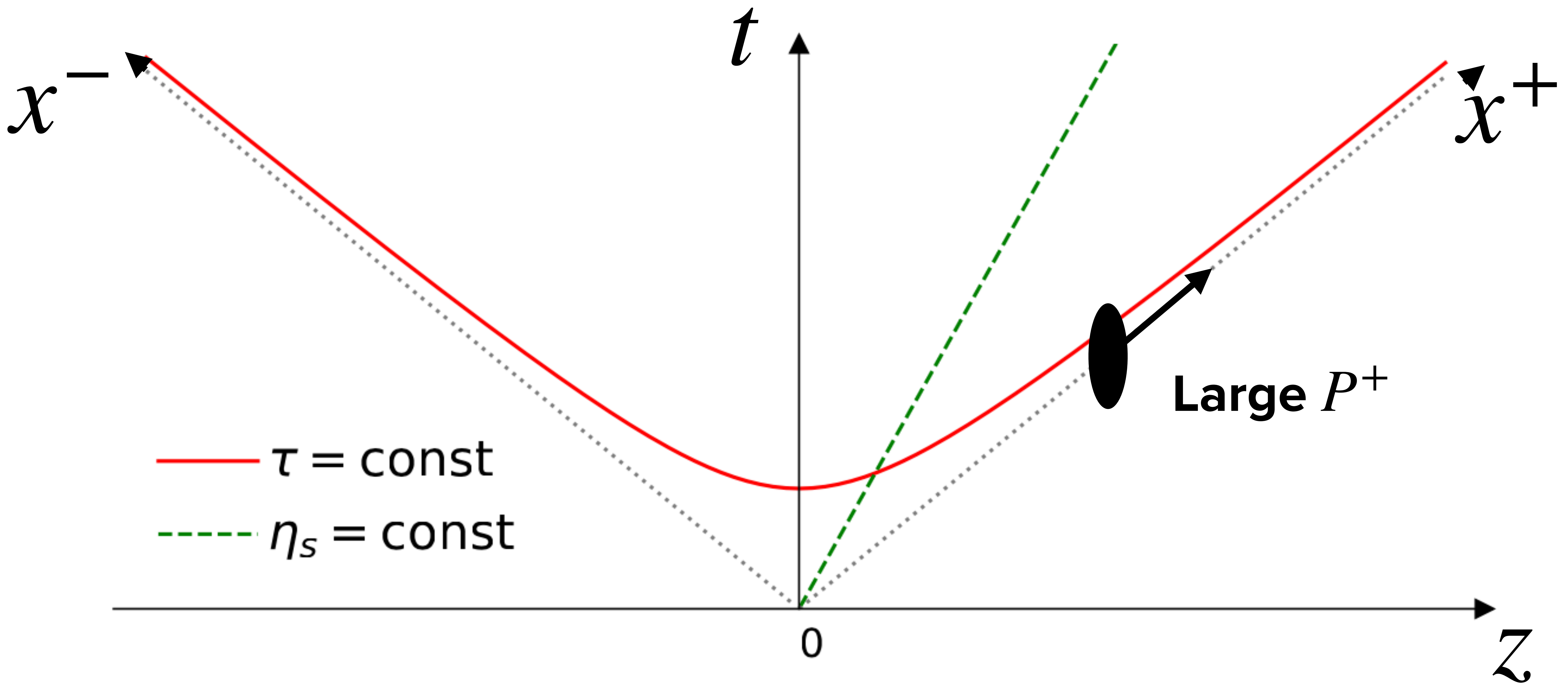
Significant suppression that evolves with energy/ $x_{\mathbb{P}}$

Larger suppression for DVCS due to larger dipole contributions.

e+O: Oxygen wave function dependence



Light cone



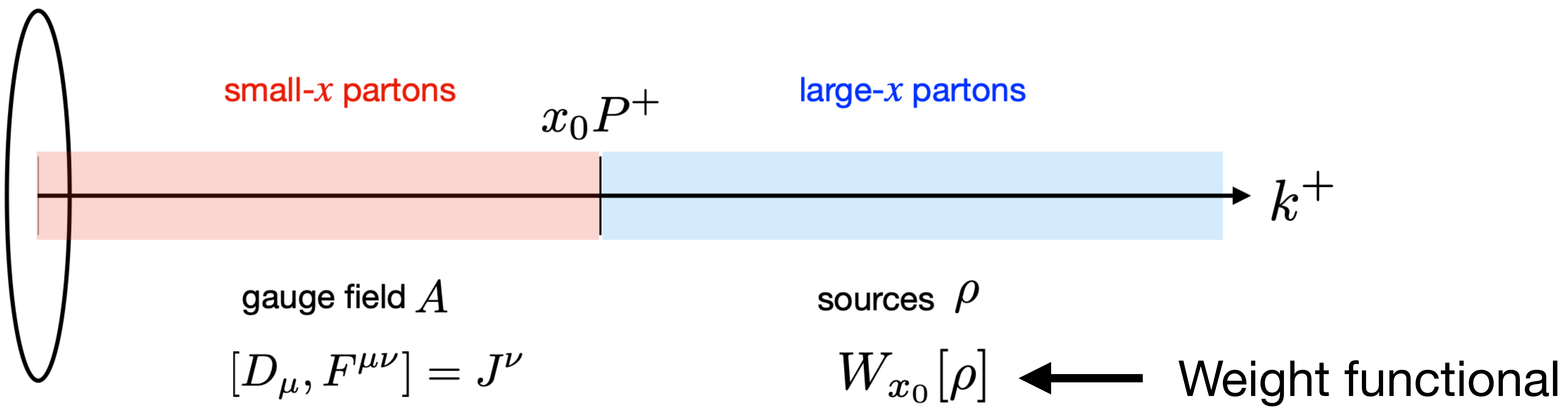
Light cone coordinates $v^{\pm} = (v^0 \pm v^3)/\sqrt{2}$

In the future light cone define $x^+ = \dfrac{\tau}{\sqrt{2}}e^{+\eta}$, and $x^- = \dfrac{\tau}{\sqrt{2}}e^{-\eta}$

or inverted $\tau = \sqrt{2x^+x^-}$, and $\eta = \dfrac{1}{2}\ln\left(\dfrac{x^+}{x^-}\right)$

Weight functional

https://arxiv.org/pdf/hep-ph/0406169.pdf



What is the weight functional?

Need to model. E.g. the McLerran-Venugopalan model:

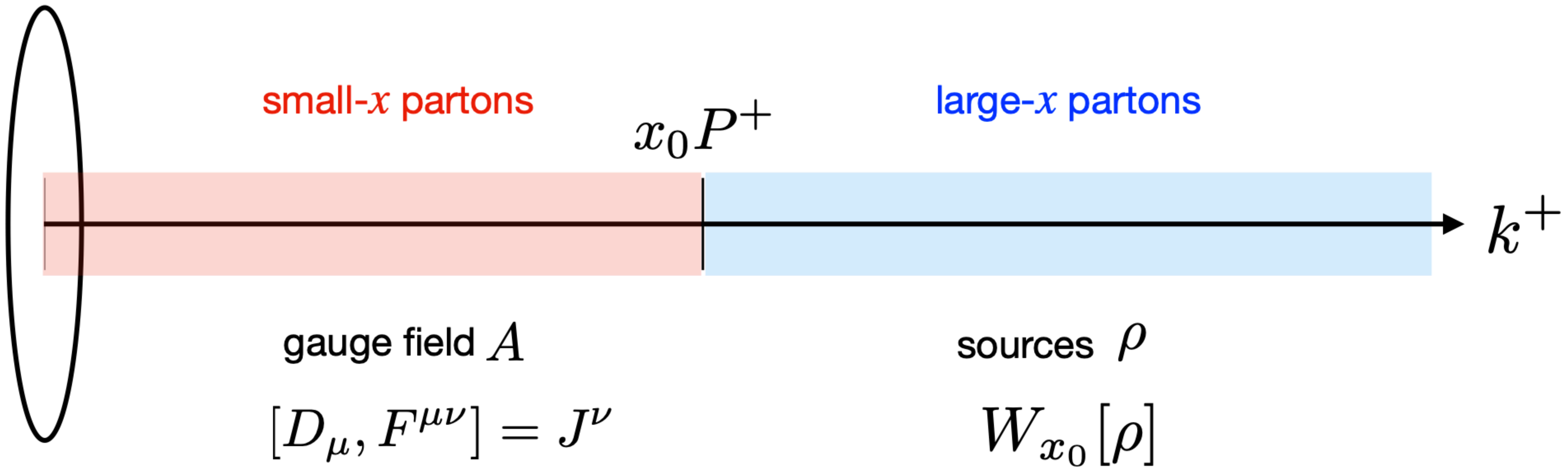
Assume a large nucleus, invoke central limit theorem. All correlations of ρ^a are Gaussian

$$W_{x_0}[\rho] = \mathcal{N} \exp\left(-\frac{1}{2} \int dx^- d^2x_T \frac{\rho^a(x^-, x_T)\rho^a(x^-, x_T)}{\lambda_{x_0}(x^-)} \right)$$

where $\lambda_{x_0}(x^-)$ is related to the transverse color charge density distribution of the nucleus

Weight functional

https://arxiv.org/pdf/hep-ph/0406169.pdf



…where $\lambda_{x_0}(x^-)$ is related to the transverse color charge density distribution of the nucleus

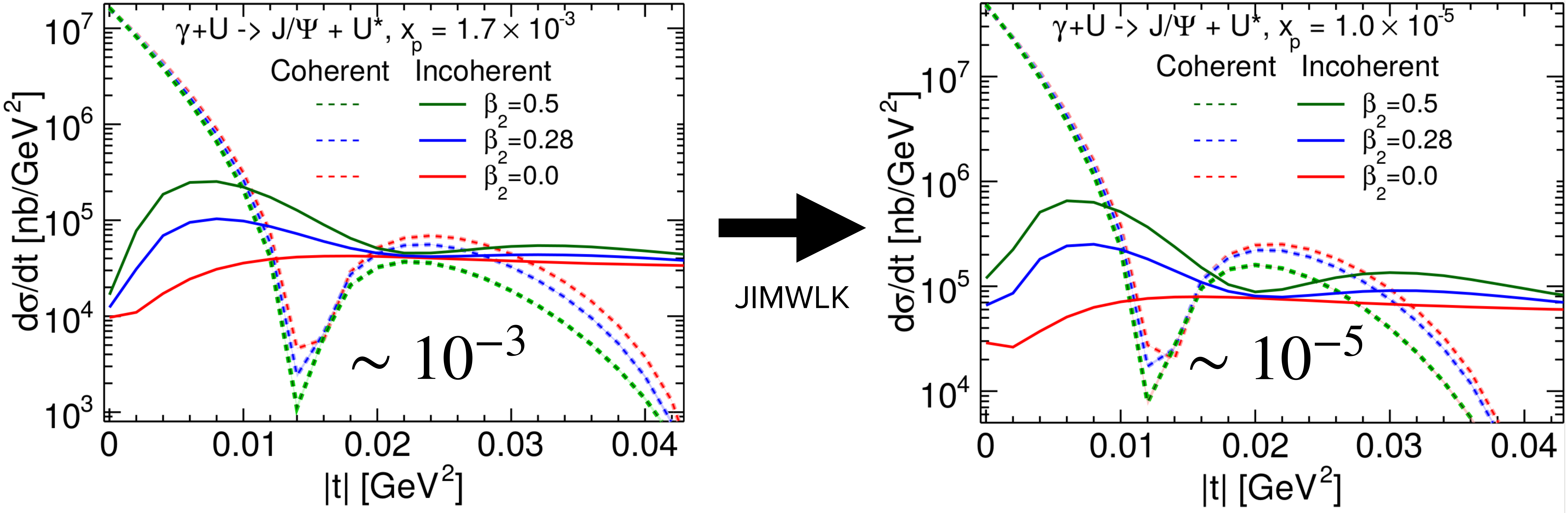
$$\mu^2 = \int dx^- \lambda_{x_0}(x^-) = \frac{(g^2 C_F)(AN_c)}{\pi R_A^2} \frac{1}{N_c^2 - 1} = \frac{g^2 A}{2\pi R_A^2} \sim A^{1/3}$$

normalized per color degree of freedom

That color charge density is related to Q_s, the saturation scale.

Towards smaller *x*: Do deformation effects survive?

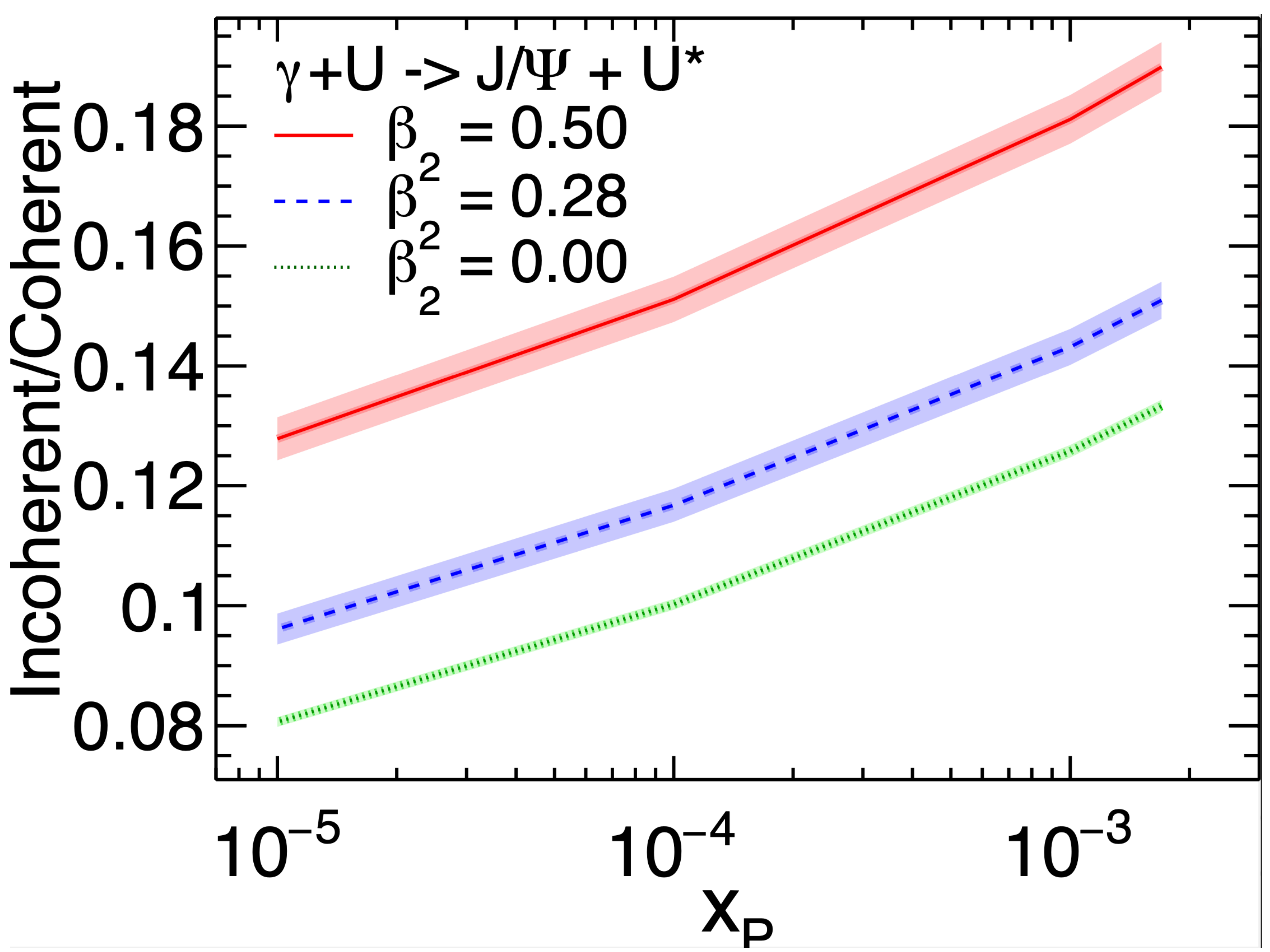
H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, in progress

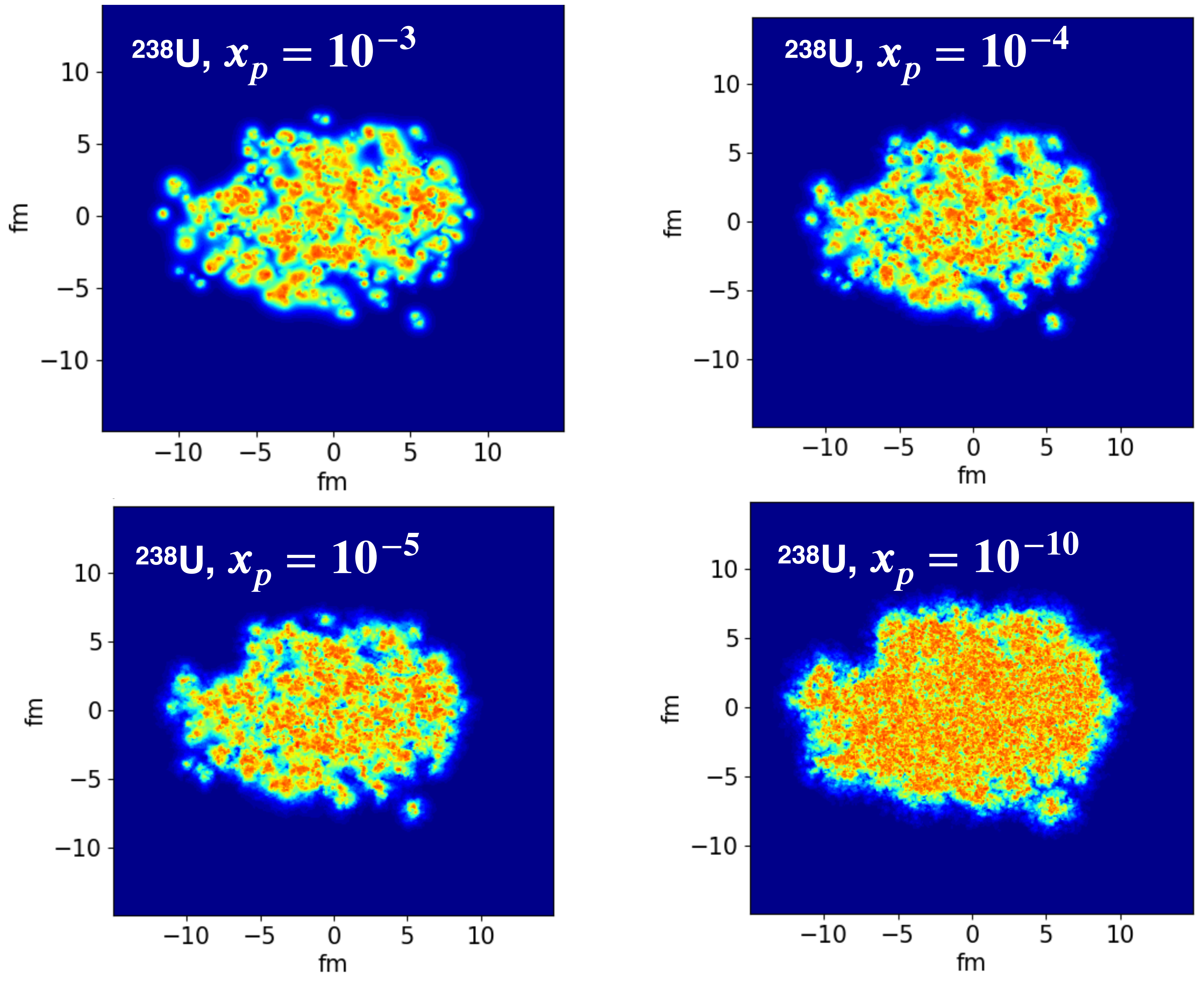


Some changes in the cross section, but deformation effects survive

Towards smaller *x*: Incoherent / coherent ratio

H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, in progress

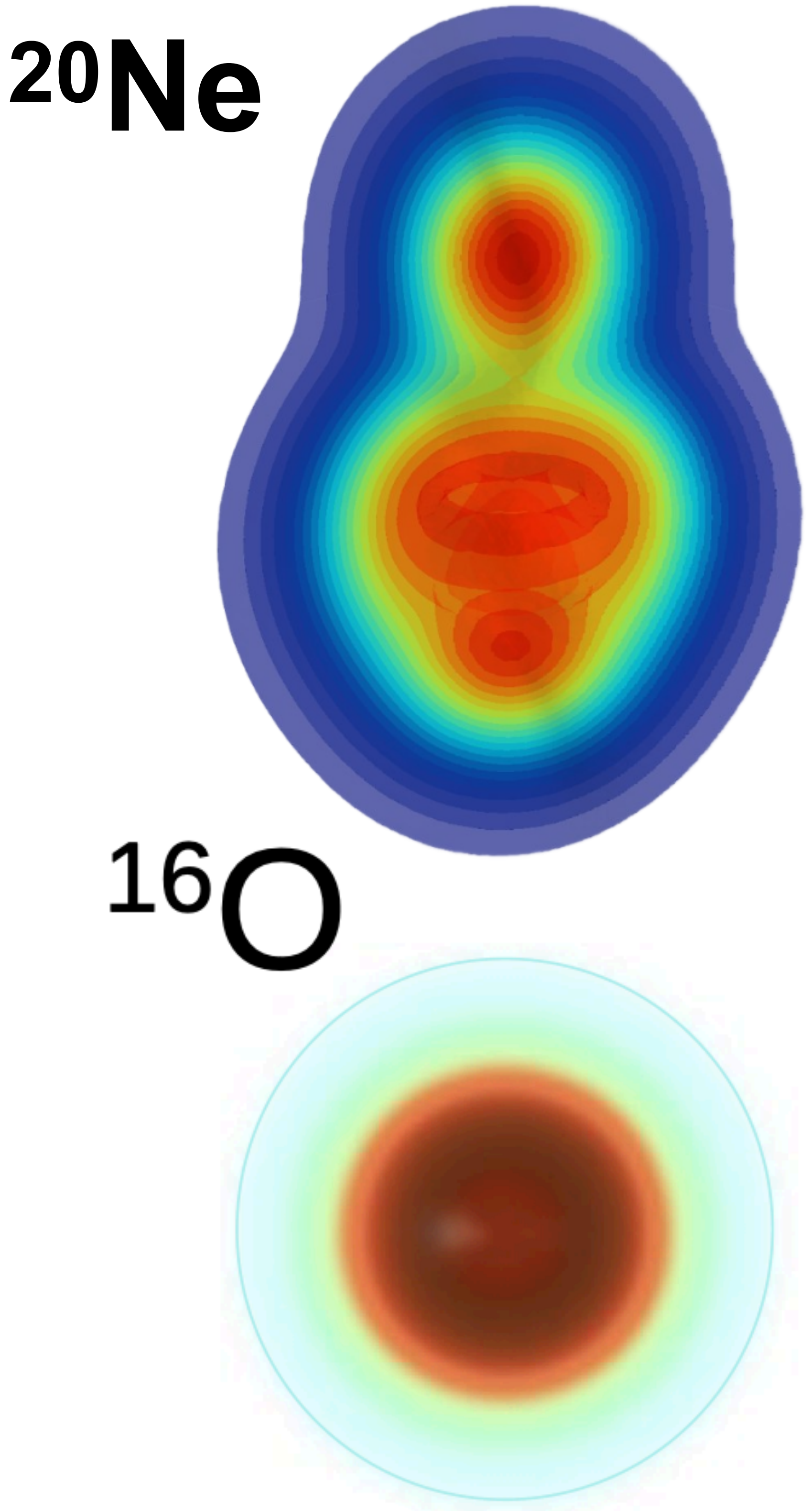


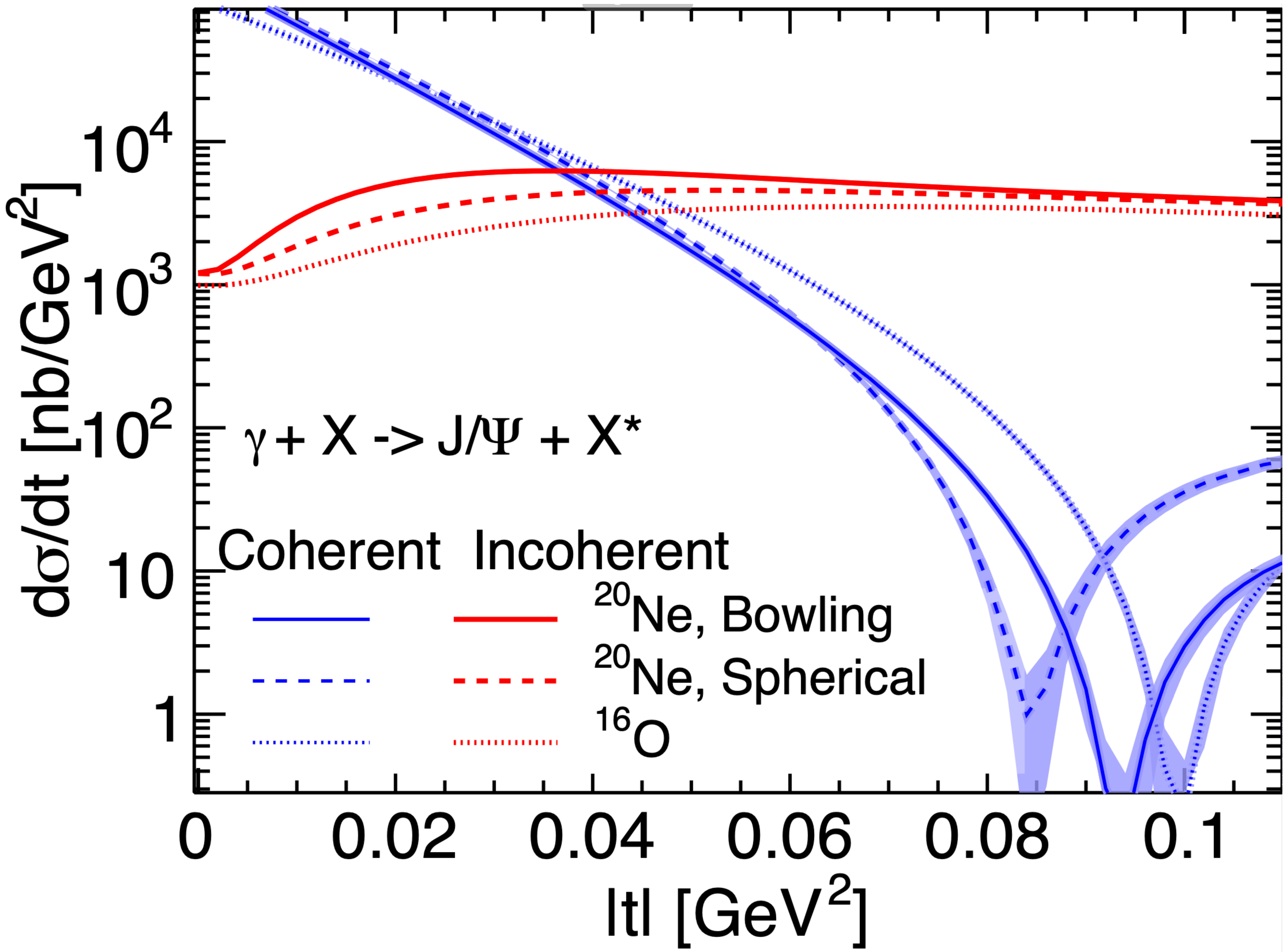


- Both cross sections grow for decreasing *x*
- Because fluctuations are reduced, incoherent/coherent ratio decreases
- Effects of deformation not noticeably reduced

Comparing Neon and oxygen

H. Mäntysaari, B. Schenke, C. Shen, W. Zhao, in progress



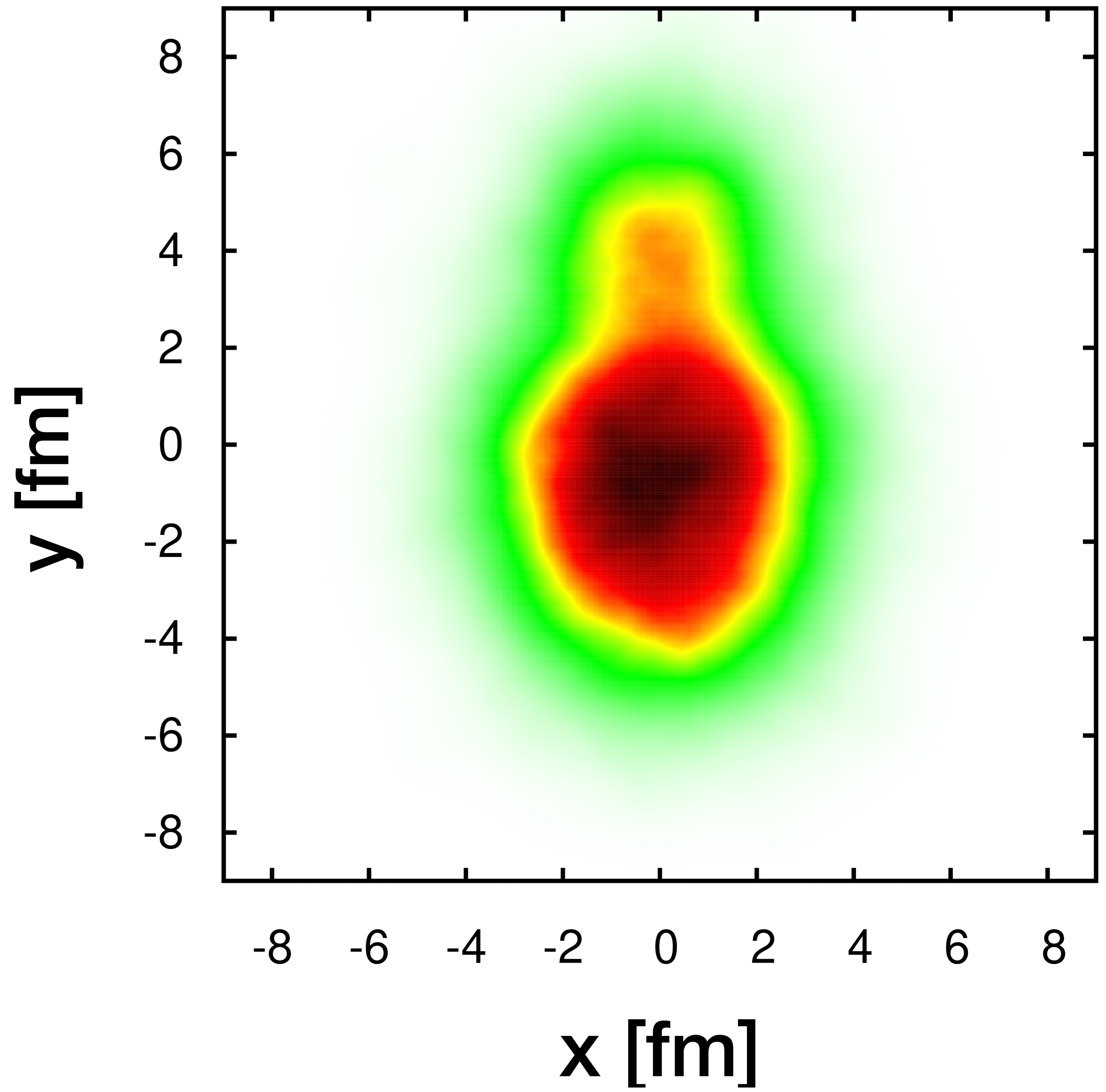


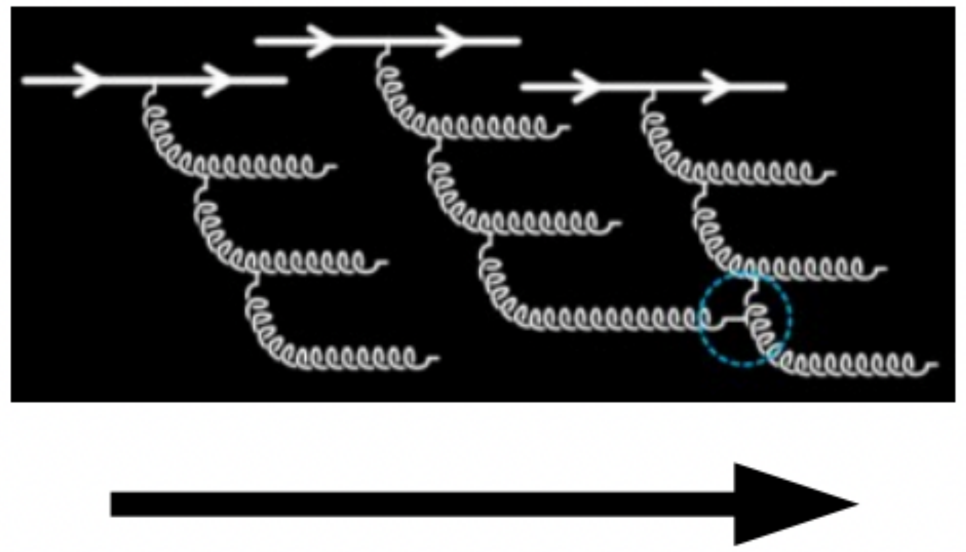
- Incoherent cross section at small |t| captures the deformation of ^{20}Ne
- Significant difference between ^{20}Ne and ^{16}O diffractive cross sections

Neon - JIMWLK evolution

G. Giacalone, B. Schenke, S. Schlichting, P. Singh, in progress

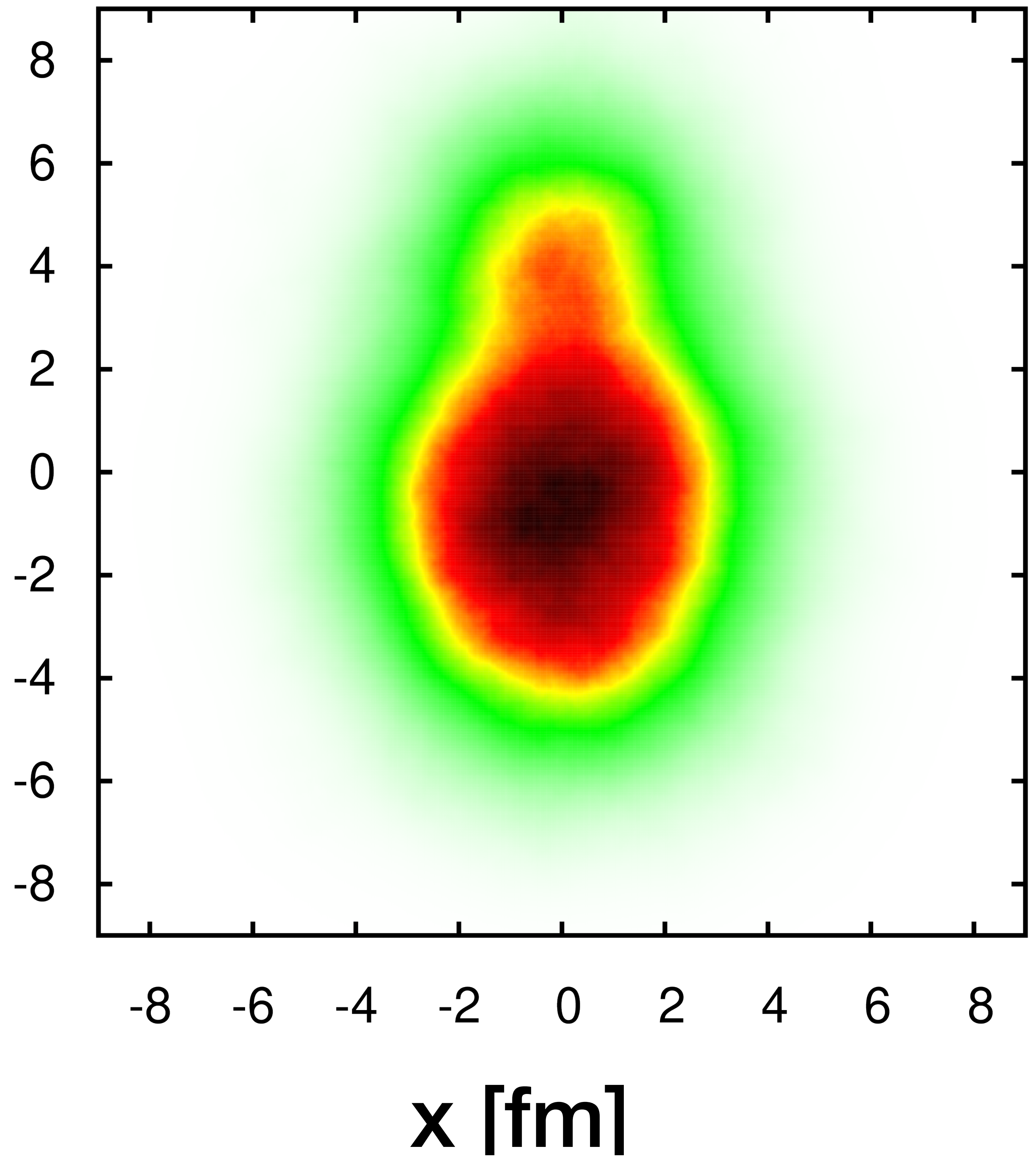
$Y = 0$





two orders of magnitude in x

$Y = 4.6$

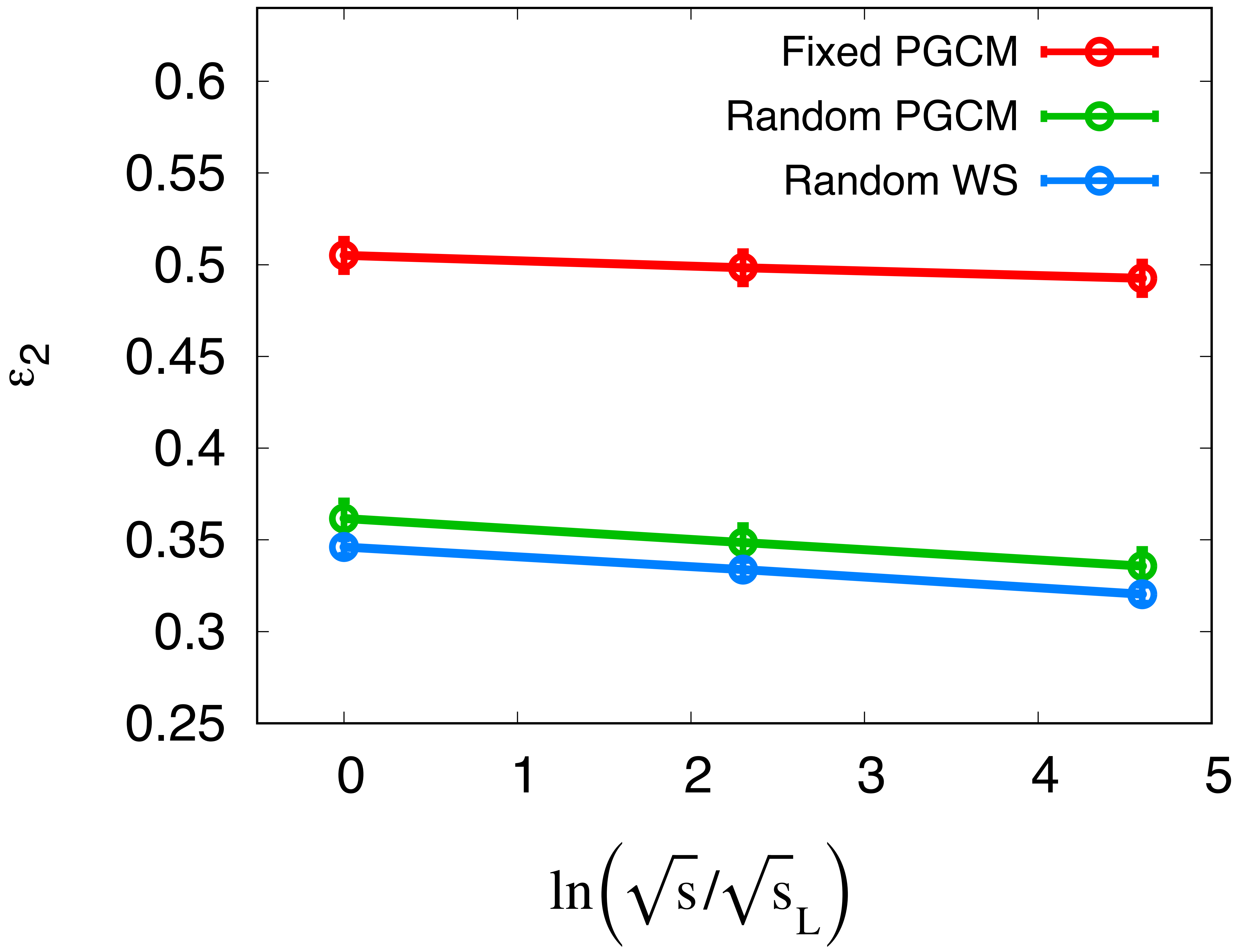


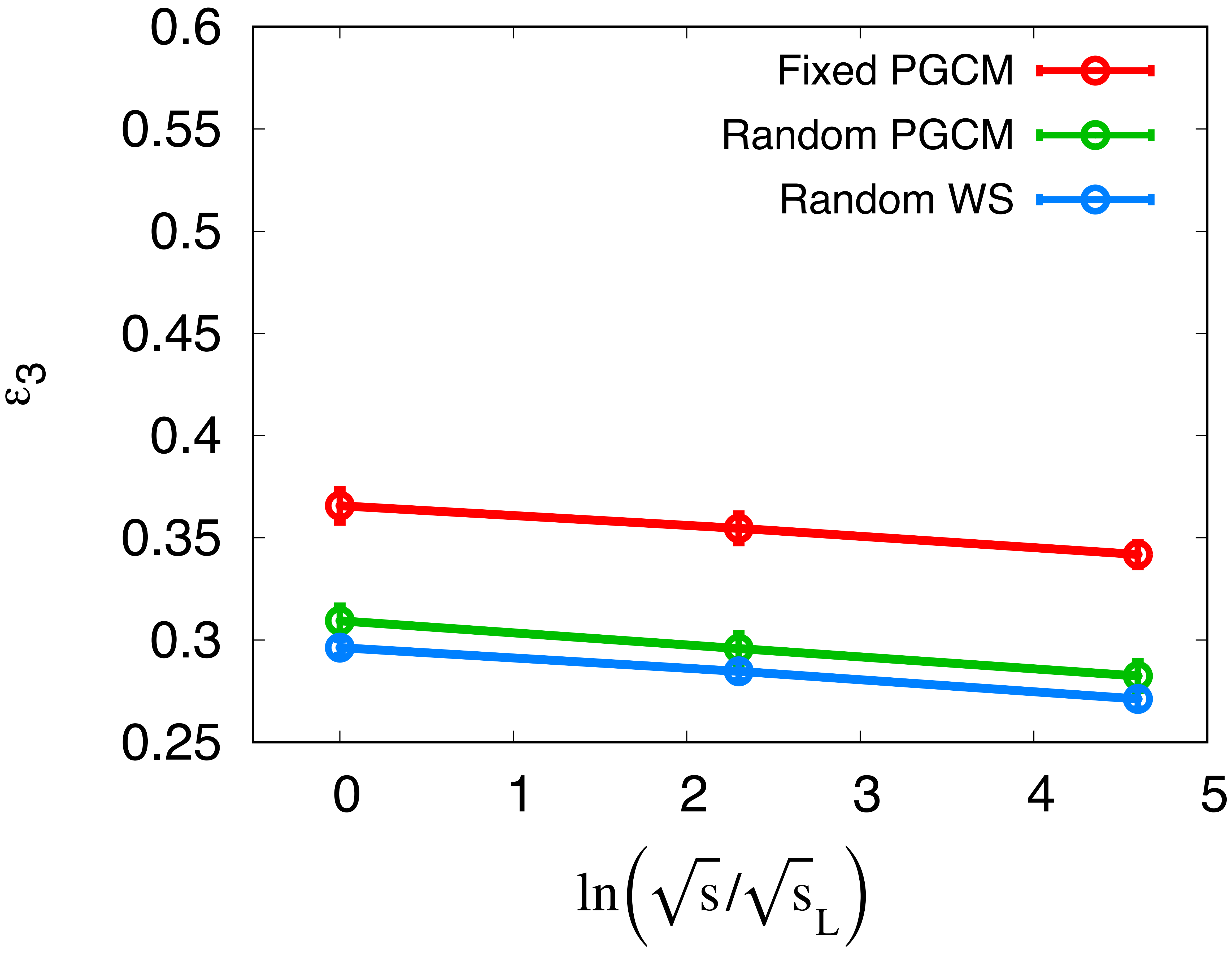
- Small-x evolution does not melt the bowling pin shape

Neon+Neon collisions - JIMWLK evolution

G. Giacalone, B. Schenke, S. Schlichting, P. Singh, in progress

- After the collision at different energies (x), measure the spatial eccentricities



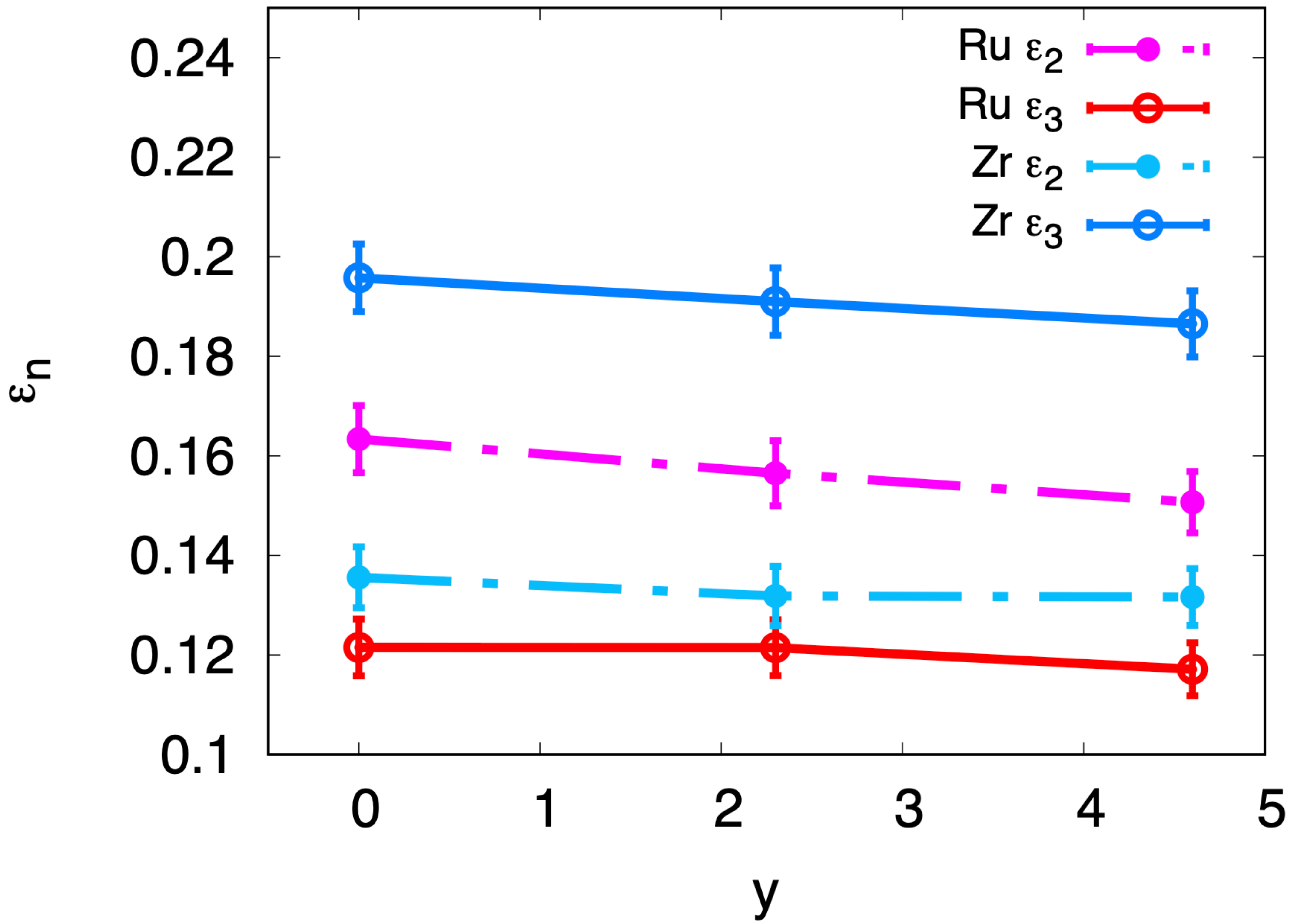


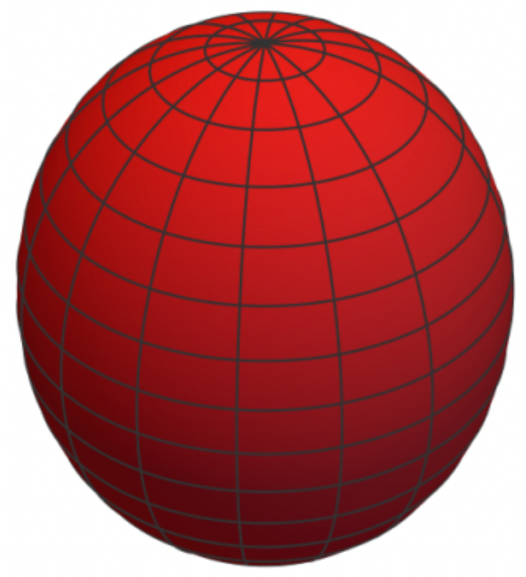
PGCM: Projected Generator Coordinate Method

- Expected reduction - smoother distributions, but no large change

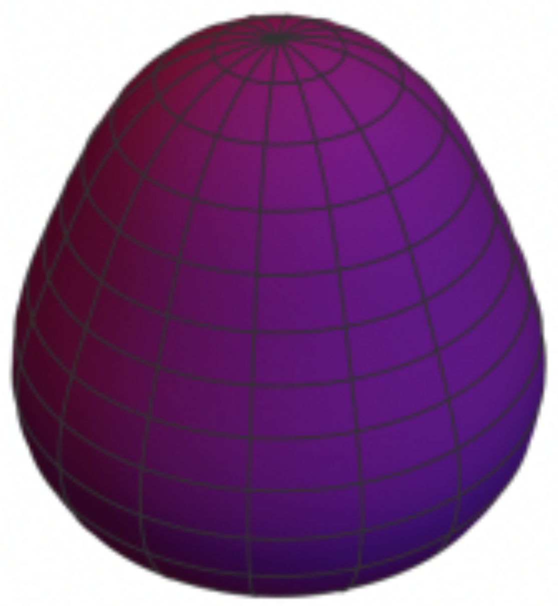
Isobar shapes - JIMWLK evolution

G. Giacalone, B. Schenke, S. Schlichting, P. Singh, in progress





$^{96}_{44}$Ru



$^{96}_{40}$Zr