

Overview of WG Small-x evolution + NLO calculations

Zhongbo Kang
UCLA

 @ZhongboK

SURGE Collaboration Meeting and Workshop

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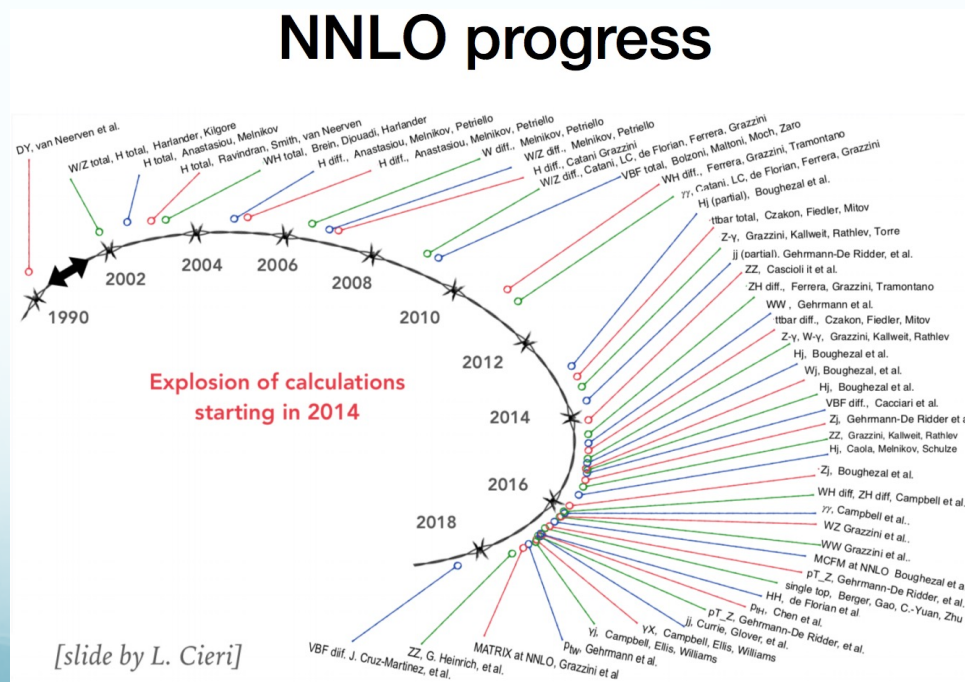


Working Group Members

- Faculty
 - Ian Balitsky (ODU & JLab)
 - Simon Caron-Huot (McGill)
 - Yuri Kovchegov (Ohio State)
 - Alex Kovner (Connecticut)
 - Jamal Jalilian-Marian (Baruch College)
 - Vladi Skokov (North Carolina)
 - Anna Stasto (Penn State)
 - Raju Venugopalan (BNL)
 - Yoshitaka Hatta (BNL)
 - Daniel Pitonyak (Lebanon Valley)
 - Matthew Sievert (New Mexico)
- Postdoc and students
 - Farid Salazar (LBL/UCLA)
 - ...
 - Please write to me zkang@ucla.edu to include your names

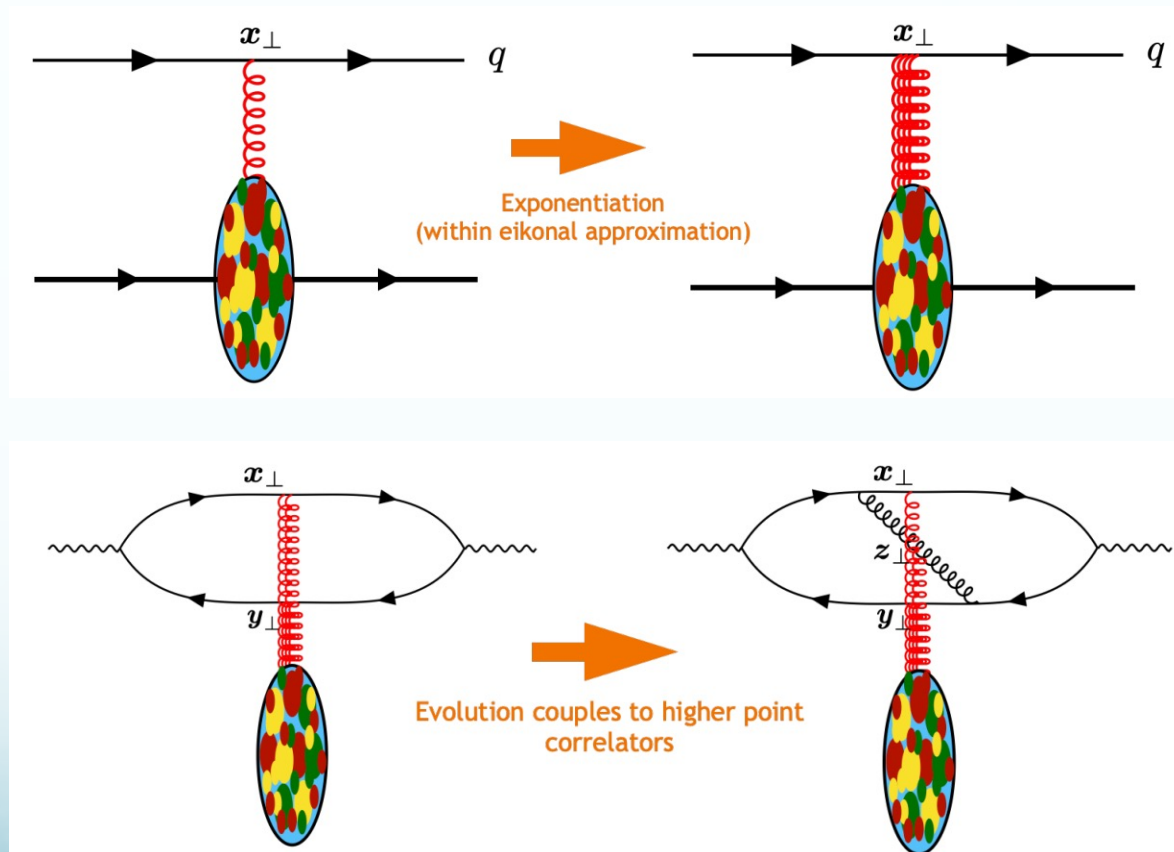
Why?

- In comparison with conventional collinear pQCD computations
 - NNLO and beyond
- Higher-order computations in the small-x regime is much more challenging
 - NLO becomes state of the art only in the last couple of years
- This also means a lot of new opportunities



Complexity

- In the small- x regime, one has to simultaneously include
 - Coherent multiple scattering to all order (*not existing in the usual collinear factorization nor in the more advanced TMD factorization*)
 - Radiative corrections order by order



Courtesy of Farid Salazar

Challenges

- NLO computations in the small- x regime is extremely complex
- Numerical implementation is also quite complicated
- One would like to have interpolation between small x and moderate x to describe the data in the full kinematic region

Main goals of WG

- Perform NLO calculations at small x
- Include corrections that allow computations at moderate x
- Develop NLO+NLL numerical packages for phenomenology

Milestones of WG

- Year 1: Skokov, Stasto, Venugopalan
 - Assess and formulate numerical implementation of dipole and quadrupole correlators NLLx BK and JIMWLK
- Year 2-3: Kang, Venugopalan, Jallian-Marian, Kovner
 - Perform NLO computations for dihadron production in e+A and assess its status in p+A and formulate strategy for small-x evolution and resummation at NLLx and beyond
 - Calculate corrections to small x observables such as single inclusive particle production in DIS to extend saturation studies to high Q^2 and mid rapidity
- Year 3: Stasto, Venugopalan, Skokov, Kovchegov, Pitonyak, Sievert
 - Numerical implementation of NLLx BK

Milestones of WG

- Year 4: [Jalilian-Marian, Kovner](#)
 - Based on the Born Oppenheimer approximation and other perturbative techniques, develop a consistent unified evolution NLO framework applicable to small as well as moderate values of x .
 - Study saturation in single inclusive particle production in DIS on nuclei from very forward (dominated by small x gluons) to mid rapidity at EIC (dominated by moderate x gluons)
- Year 4-5: [Kang, Venugopalan](#)
 - Study and implement dijet and dihadron production at NLO in p+A collisions
- Year 4-5: [Hatta](#)
 - Calculate single and double spin asymmetries in dijet production at small- x to NLO including the quark and gluon OAM contributions and make predictions for the EIC

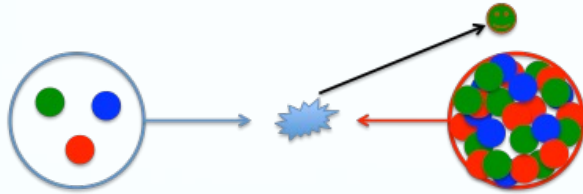
Current status

- Arguably, NLO computations in the color-glass condensate / gluon saturation regime is one of the most active research directions in the field and a lot of progress has been made

- “Single inclusive hadron production in p+A collisions at the NLO” by Liu, Kang, Liu [3] and by Shi, Wang, Wei, Xiao [4]
- “Single inclusive jet production in p+A collisions at the NLO” by Liu, Xie, Kang, Liu [5] and by Wang, Chen, Gao, Shi, Wei, Xiao [6]
- “Dijet photoproduction in γ +A collisions at the NLO” by Taels, Altinoluk, Beuf, Marquet [7]
- “Dijet impact factor in e+A collisions at the NLO” by Caucal, Salazar, Venugopalan [8–10]
- “Single inclusive hadron production in e+A collisions at the NLO” by Bergabo, Jalilian-Marian [11]
- “Dihadron production in e+A collisions at the NLO” by Bergabo, Jalilian-Marian [12, 13] and by Iancu, Mulian [14]
- “Diffractive dihadron production in γ^* +A collisions at the NLO” by Fucilla, Grabovsky, Li, Szymanowski, Wallon [15]
- ...

Phenomenology example: single hadron

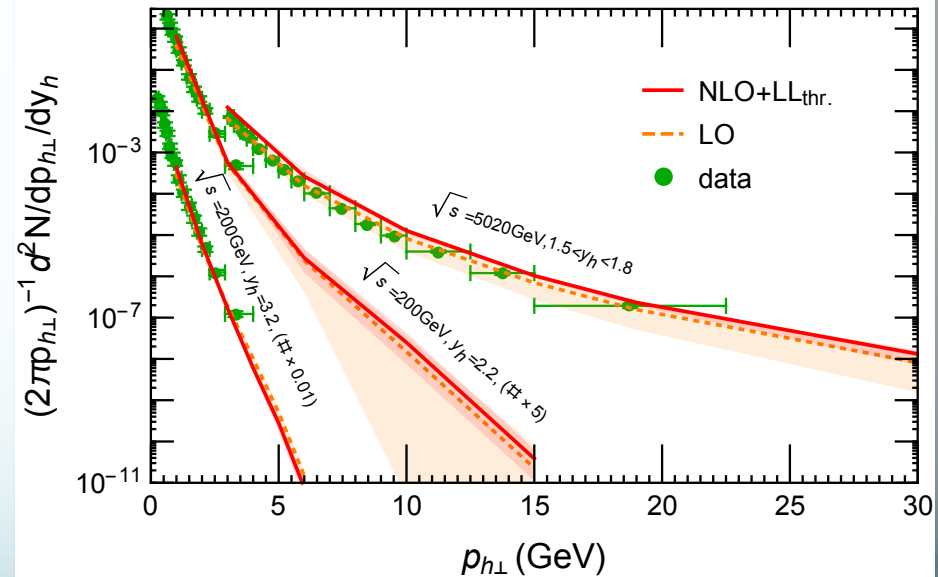
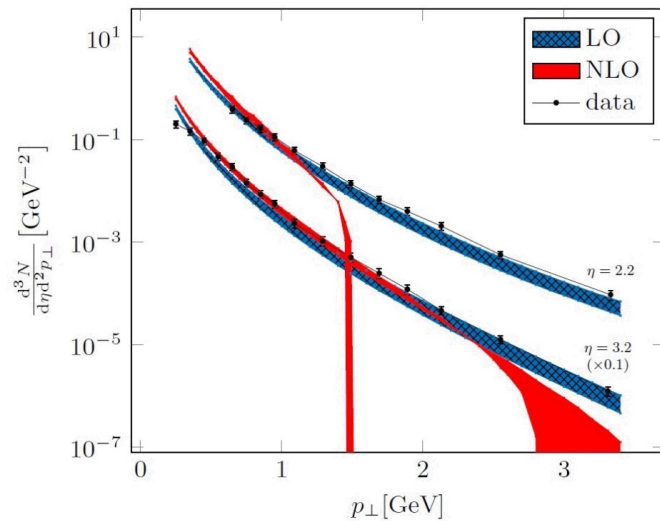
- Forward single hadron production in p+A collisions
 - Lesson: threshold resummation $x_1 \sim 1$ and small-x resummation $x_2 \ll 1$



$$x_1 \sim \frac{p_\perp}{\sqrt{s}} e^{+y} \sim \mathcal{O}(1)$$

$$x_2 \sim \frac{p_\perp}{\sqrt{s}} e^{-y} \ll 1$$

A. Stasto, B.-W. Xiao, D. Zaslavsky,
Phys. Rev. Lett. **112** (2014) 012302 BRAHMS $\eta = 2.2, 3.2$



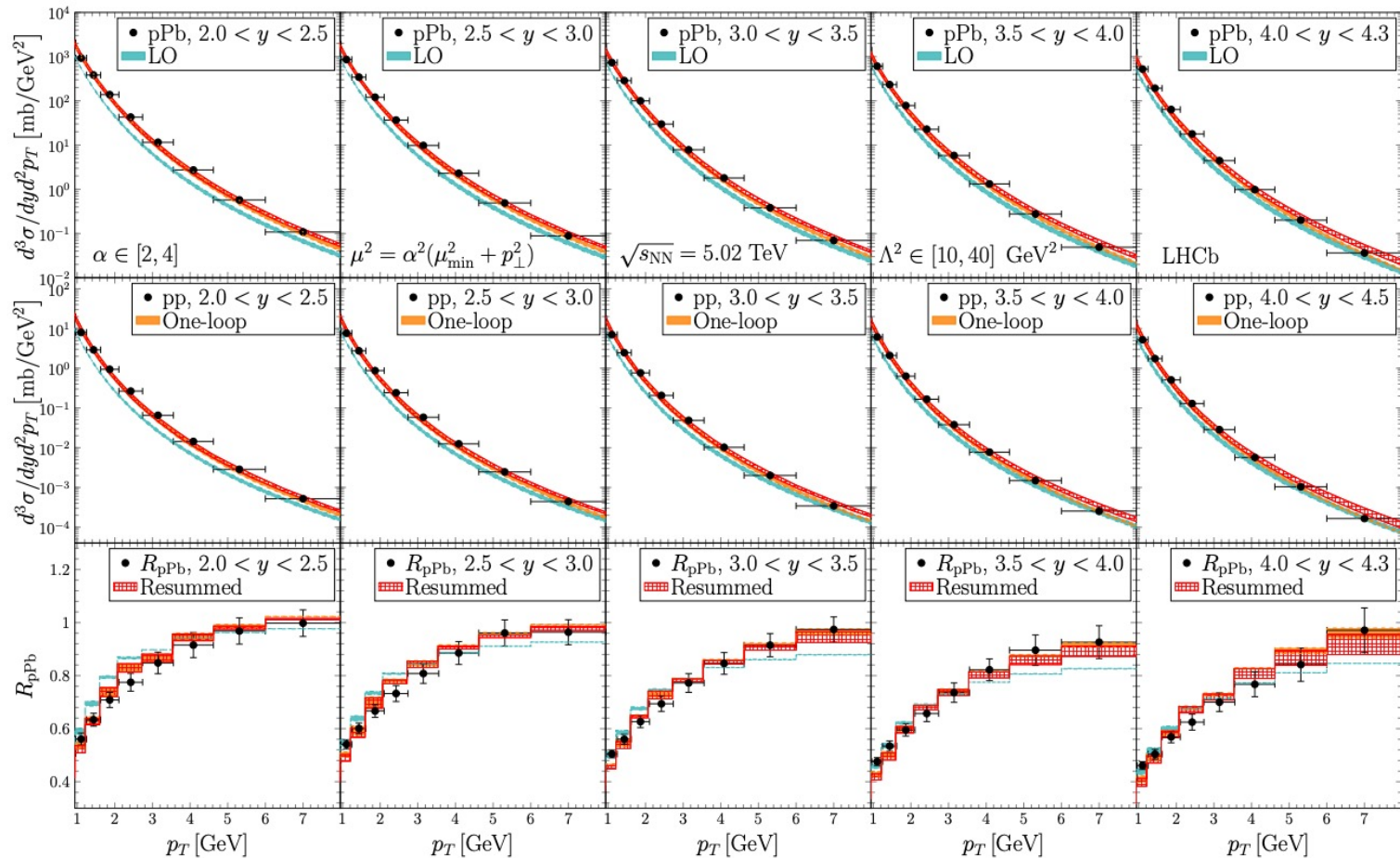
Kang, Liu, Liu, 2004.11990, PRD 20

Full phenomenology

PHYSICAL REVIEW LETTERS **128**, 202302 (2022)

Pursuing the Precision Study for Color Glass Condensate in Forward Hadron Productions

Yu Shi^{1,2,*} Lei Wang^{1,2,†} Shu-Yi Wei^{1,3,‡} and Bo-Wen Xiao^{1,4,§}



Phenomenology example: single jet

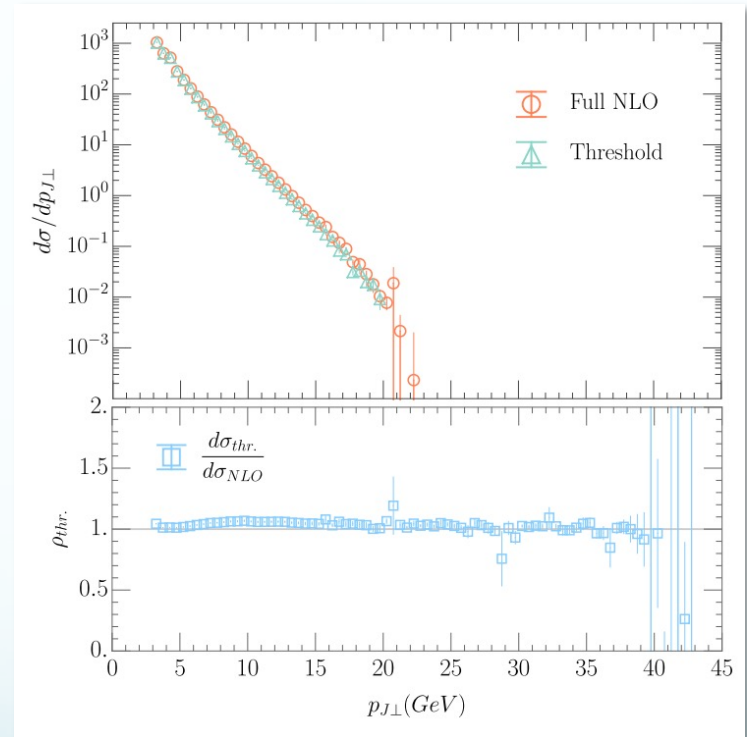
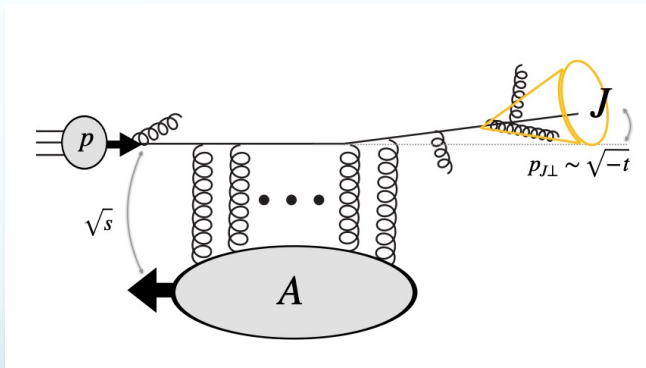
- Forward single jet production in p+A collisions
 - Complete NLO with full jet algorithm implementation
 - Analytical computations for jets with small radius R

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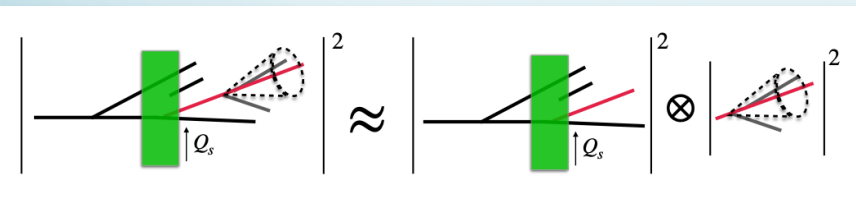
RECEIVED: April 20, 2022
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Single inclusive jet production in pA collisions at NLO in the small- x regime

Hao-yu Liu,^{a,b} Kexin Xie,^c Zhong-Bo Kang,^{d,e,f} and Xiaohui Liu^{a,b,g}



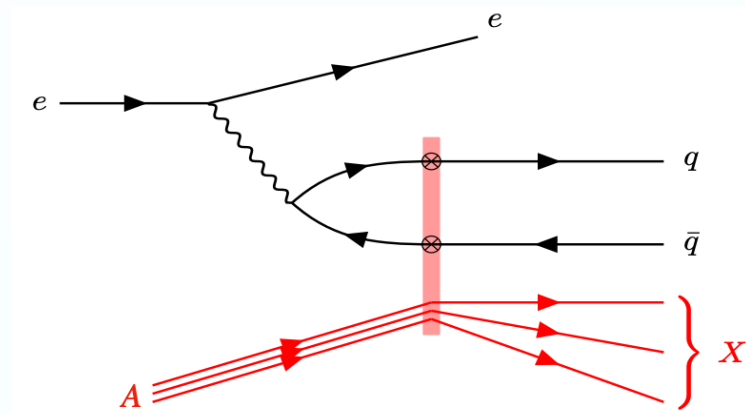
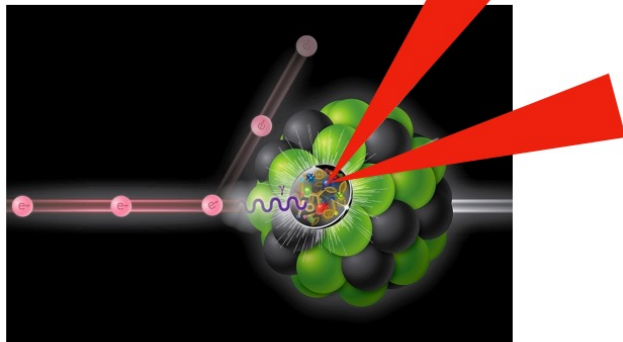
Consistent with Wang, Chen, Gao, Shi, Wei, Xiao, PRD 23



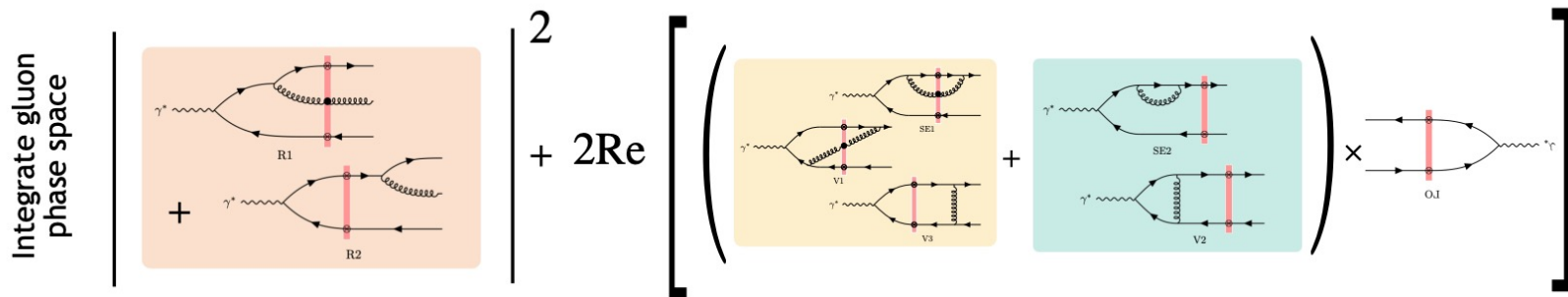
Phenomenology example: dijet in e+A

- Dijet production in e+A collisions at NLO

F. Gelis, J. Jalilian-Marian (2003)



P. Caucal, FS, and R. Venugopalan. *JHEP* 11 (2021) 222

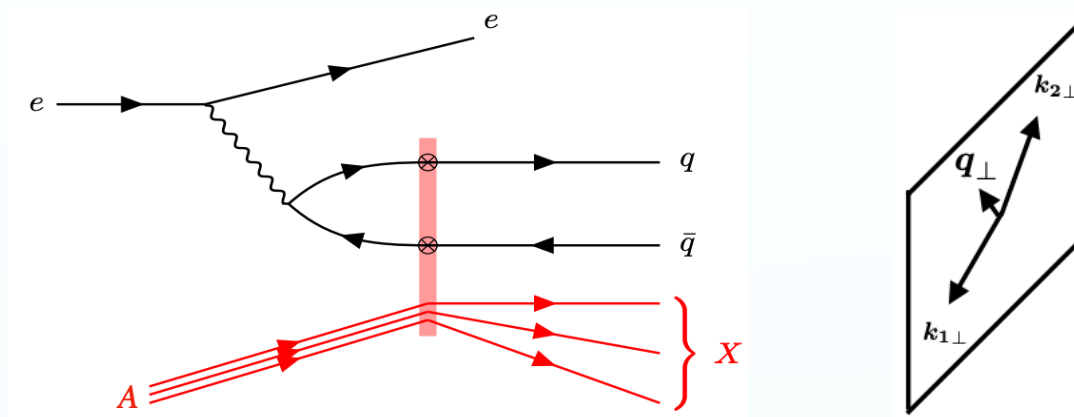


Caucal, Salazar, Schenke, Stebel,
Venugopalan, 21, 22, 23

Courtesy of Farid Salazar

One very interesting question

- What happens when the dijet are produced back-to-back in the transverse plane?
 - Forwards jets but back-to-back in transverse plane



- In this region, Sudakov logarithms become important
- Connections with TMD factorization
- Important progress and insights have been made

Caucal, Salazar, Schenke, Stebel,
Venugopalan, 23

Observation and question

- One first perform NLO computations
 - Extremely challenging, so it requires truly remarkable efforts
- One then takes the back-to-back limit to achieve both Sudakov + small- x logarithm resummation
 - Also not easy



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108 pages

Dijet impact factor in DIS at next-to-leading order in the Color Glass Condensate

Paul Caucal,^a Farid Salazar^{a,b,c} and Raju Venugopalan^a



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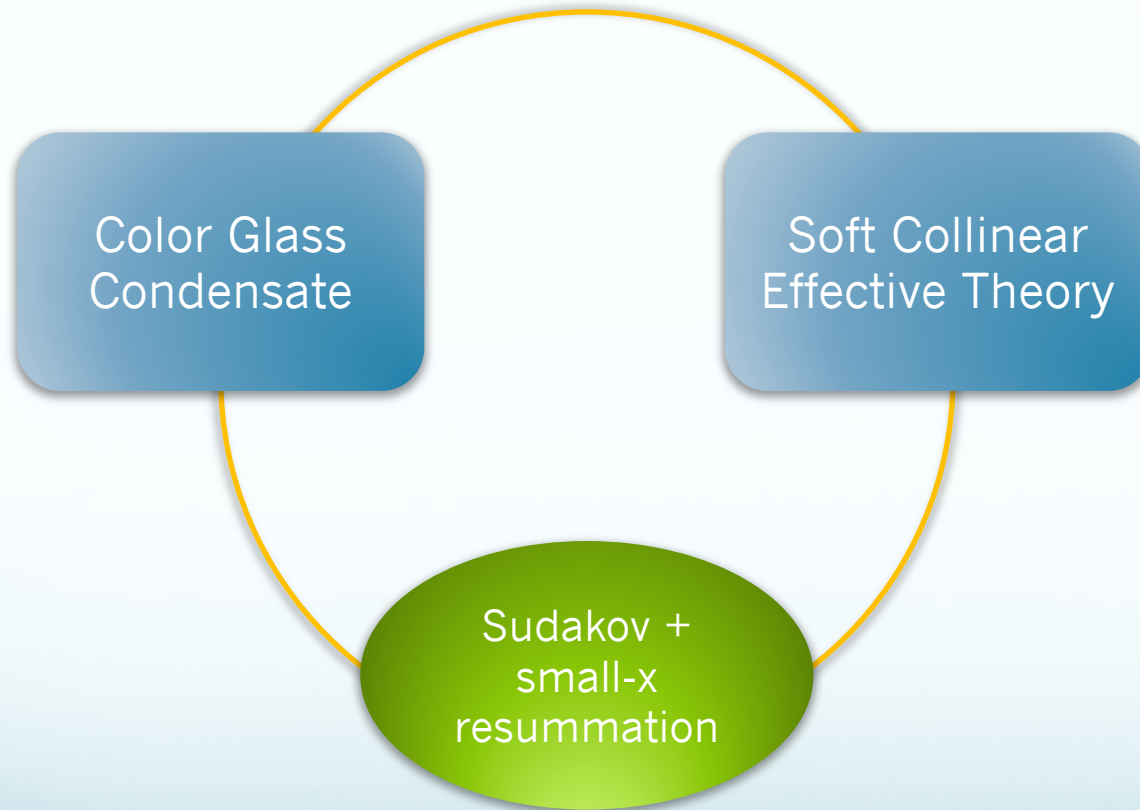
**Back-to-back inclusive dijets in DIS at small x :
Sudakov suppression and gluon saturation at NLO**

Paul Caucal,^a Farid Salazar,^{b,c,d,e} Björn Schenke^a and Raju Venugopalan^a

- Is it possible to deal with the back-to-back limit from the very beginning?

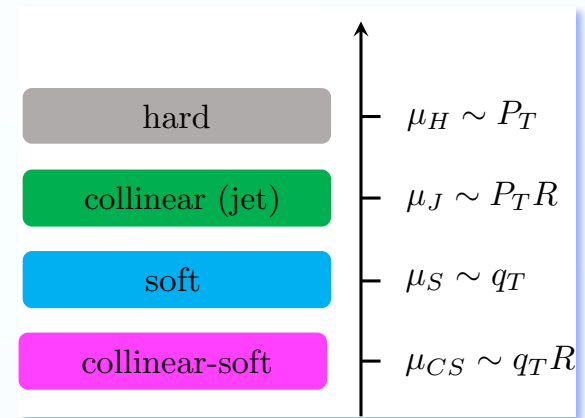
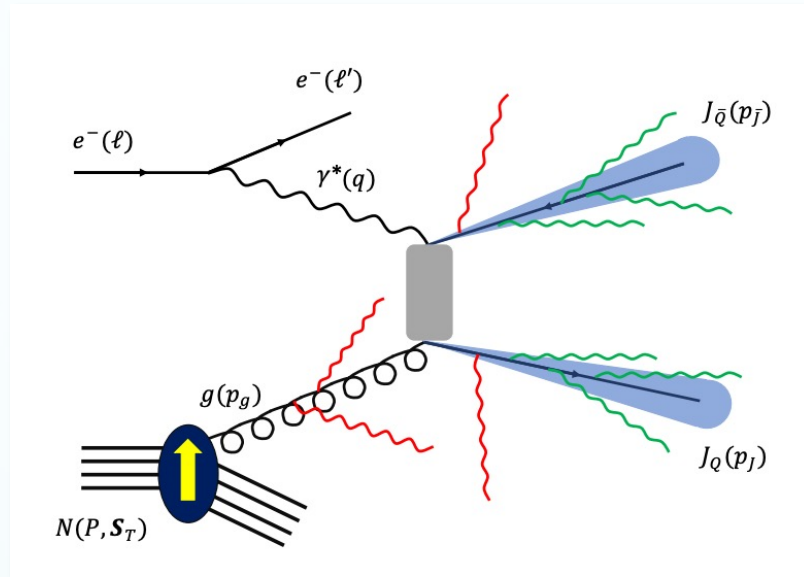
Maybe: CGC + SCET

- SCET/TMD factorization is suitable for Sudakov resummation
- CGC effective theory is suitable for small-x resummation



TMD factorization for dijet

- TMD factorization derived using SCET: back-to-back dijet in e+p



$$\frac{d\sigma}{dq_T dP_T} \propto \mathbf{H}_{\gamma^* g \rightarrow q\bar{q}}(P_T) f_{g/A}(x_a, k_{aT})$$

$$\otimes \mathbf{S}^{\text{global}}(q_T) \otimes \mathbf{S}_c^{\text{CS}}(q_T R) \mathbf{S}_d^{\text{CS}}(q_T R) \mathbf{J}_c(P_T R) \mathbf{J}_d(P_T R)$$

See Kang, Reiten, Shao, Terry, 2012.01756,
Castillo, Echevarria, Makris, Scimemi, 2008.07531


Conjecture at the small x region

- Sudakov logs can be obtained via RG equations for gluon TMD, hard function, soft function, collinear soft function, and jet function


$$\begin{aligned} \frac{d\sigma^{UU}}{dQ^2 dy d^2\mathbf{q}_T dy_J d^2\mathbf{p}_T} &= H(Q, y, p_T, y_J, \mu_h) \int_0^\infty \frac{bdb}{2\pi} J_0(b q_T) f_{g/N}(x, \mu_{b*}) \\ &\times \exp \left[- \int_{\mu_{b*}}^{\mu_h} \frac{d\mu}{\mu} \Gamma^h(\alpha_s) - 2 \int_{\mu_{b*}}^{\mu_j} \frac{d\mu}{\mu} \Gamma^{j\mathcal{Q}}(\alpha_s) - \int_{\mu_{b*}}^{\mu_{cs}} \frac{d\mu}{\mu} (\bar{\Gamma}^{cs\mathcal{Q}}(\alpha_s) + \bar{\Gamma}^{cs\bar{\mathcal{Q}}}(\alpha_s)) \right] \\ &\times \exp[-S_{\text{NP}}(b, Q_0, n \cdot p_g)] , \end{aligned} \quad (2.75)$$

See for example, Kang, Reiten, Shao, Terry, 2012.01756

- Then you could expand gluon TMD in terms of WW gluon distribution in the small-x limit



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Nuclear Physics B 921 (2017) 104–126



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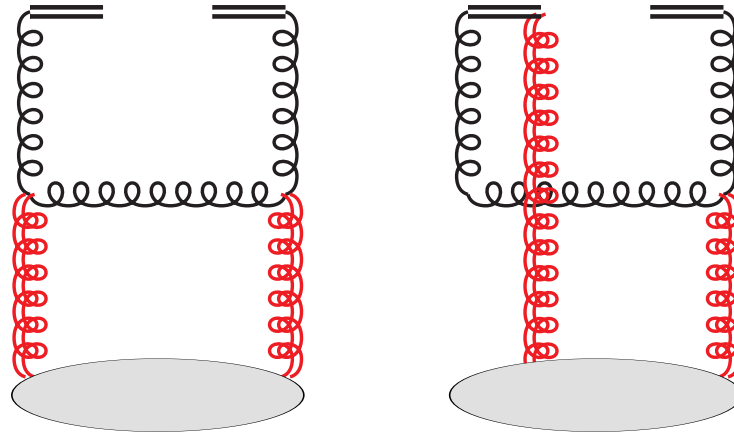
www.elsevier.com/locate/nuclphysb

Transverse momentum dependent parton distributions at small- x

Bo-Wen Xiao ^a, Feng Yuan ^{b,*}, Jian Zhou ^c

Gluon TMDs at small x

- Work by Xiao, Yuan, Zhou 2017 + Zhou 2019



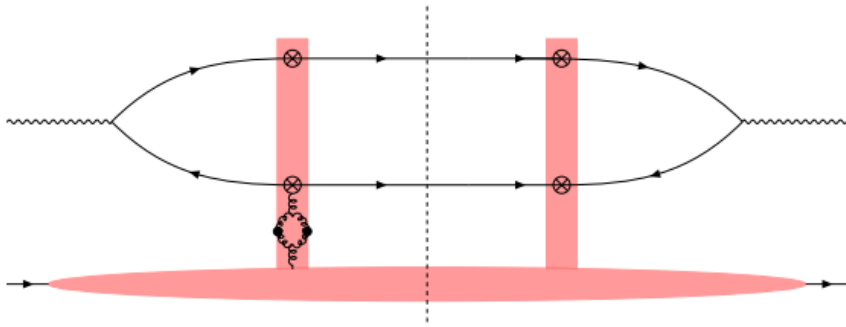
$$\begin{aligned}
 xG(x, k_{\perp}, \zeta_c^2 = \mu_F^2 = Q^2) &= -\frac{2}{\alpha_s} \int \frac{d^2x_{\perp} d^2y_{\perp}}{(2\pi)^2} e^{ik_{\perp} \cdot b_{\perp}} \\
 &\times \mathcal{H}(\alpha_s(Q)) \text{Exp} \left\{ -\int_{\mu_b}^Q \frac{d\mu}{\mu} \left(A \ln \frac{Q^2}{\mu^2} + B \right) \right\} \\
 &\times \mathcal{F}(x_{\perp}, y_{\perp}),
 \end{aligned}$$

$$\mathcal{F}(x_{\perp}, y_{\perp}) = \langle \text{Tr}[\partial_{\perp}^i U(x_{\perp}) U^{\dagger}(y_{\perp}) \partial_{\perp}^i U(y_{\perp}) U^{\dagger}(x_{\perp})] \rangle$$

Recent insight

- Quantum correction to Shockwave formalism

Caucal, Salazar, Schenke, Stebel, Venugopalan, 2304.03304



$$\hat{G}_Y^{(1)}(\mathbf{r}_{bb'}) = -\alpha_s \beta_0 \left[\frac{1}{\epsilon} + \text{finite} \right] \hat{G}_Y^{(0)}(\mathbf{r}_{bb'})$$

- Gluon TMD in small-x

$$f_{3/A}^{\text{TMD}}(x_A, \mathbf{r}_\perp, M, z) = \exp \left[- \int_{\mu_{r_1}}^{\mu} \frac{d\mu'}{\mu'} \frac{\alpha_s}{\pi} \left(C_A \ln \frac{k_2}{\mu^2} - \frac{\beta_0}{z} \right) \right] * F_Y^{\text{NW}}(\mathbf{r}_\perp, \mathbf{R}_\perp)$$

$$G_Y^{\text{bare}}(r) = z \overset{\text{renormalized}}{G}_Y \quad z = + \frac{\alpha_s}{4\pi} \beta_0 \frac{1}{\epsilon}$$

$$\Rightarrow \mu \frac{d}{d\mu} G_Y = \gamma_G G_Y$$

$$\text{Where } \gamma_G = -z^{-1} \mu \frac{d}{d\mu} z$$

$$= -\mu \frac{d}{d\mu} \left[\frac{\alpha_s}{4\pi} \beta_0 \frac{1}{\epsilon} \right]$$

$$\Downarrow \mu \frac{d}{d\mu} \alpha_s = (-2\epsilon) \alpha_s + \dots$$

$$= -(-2\epsilon) \frac{\alpha_s}{4\pi} \beta_0 \frac{1}{\epsilon} = 2 \frac{\alpha_s}{4\pi} \beta_0$$

$$= -\frac{\alpha_s}{\pi} \left[-\frac{\beta_0}{2} \right]$$

This exactly gives us

$$G_Y(r, \mu) = \exp \left[- \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \frac{\alpha_s}{\pi} \left(-\frac{\beta_0}{z} \right) \right] G_Y(r, \mu_b)$$

$$\mu_b = \frac{c_0}{r}$$

$$c_0 = z e^{-\gamma_G \epsilon}$$

Checking with Higgs production

- gg Higgs at small transverse momentum
 - Start with a TMD factorization
 - Replace gluon TMD in the nuclear target by the small-x expansion
 - One obtains the correct formula now has both Sudakov + small-x logs

For $PA \rightarrow \text{Higgs} + X$ production, we'll then have

$$d\sigma \propto f_{g/p}^{\text{TMD}}(x_p, \underline{r}_\perp, \mu, \xi_1) f_{g/A}^{\text{TMD}}(x_A, \underline{R}_\perp, \mu, \xi_2) H(Q)$$

↓
replace this TMD in the target
by its expansion in terms of WW
gluon distribution

at NLL

$$f_{g/p}^{\text{TMD}}(x_p, \underline{r}_\perp, \mu, \xi_1) = f_{g/p}(x_p, \mu, \underline{r}_\perp) \times \exp\left[-\int_{\mu_{r_L}}^{\mu} \frac{d\mu'}{\mu'} \frac{\alpha_s}{\pi} \left(C_A \ln \frac{\xi_1}{\mu^2} - \frac{\beta_0}{2}\right)\right]$$

for gluon TMD on the nuclear target

$$f_{g/A}^{\text{TMD}}(x_A, \underline{R}_\perp, \mu, \xi_2) = \exp\left[-\int_{\mu_{R_L}}^{\mu} \frac{d\mu'}{\mu'} \frac{\alpha_s}{\pi} \left(C_A \ln \frac{\xi_2}{\mu^2} - \frac{\beta_0}{2}\right)\right] \times F_{\gamma}^{\text{NW}}(r_L, R_L)$$

Mueller, Xiao, Yuan, 2013

Summary

- WG on small-x evolution + NLO calculations should be extremely fun
- We expect to have a lot of exciting developments in the coming years
- We will try our best to work as a team to accomplish milestones we have promised to DOE
- We welcome interested postdocs and students to join us
- Check talks by our WG members:
 - Wednesday: Jamal, Ian, Tomasz
 - Thursday: Simon, Vladimir, Anna, Farid

Thank you!