



Illinois Center for Advanced Studies of the Universe

U.S. DEPARTMENT OF ENERGY



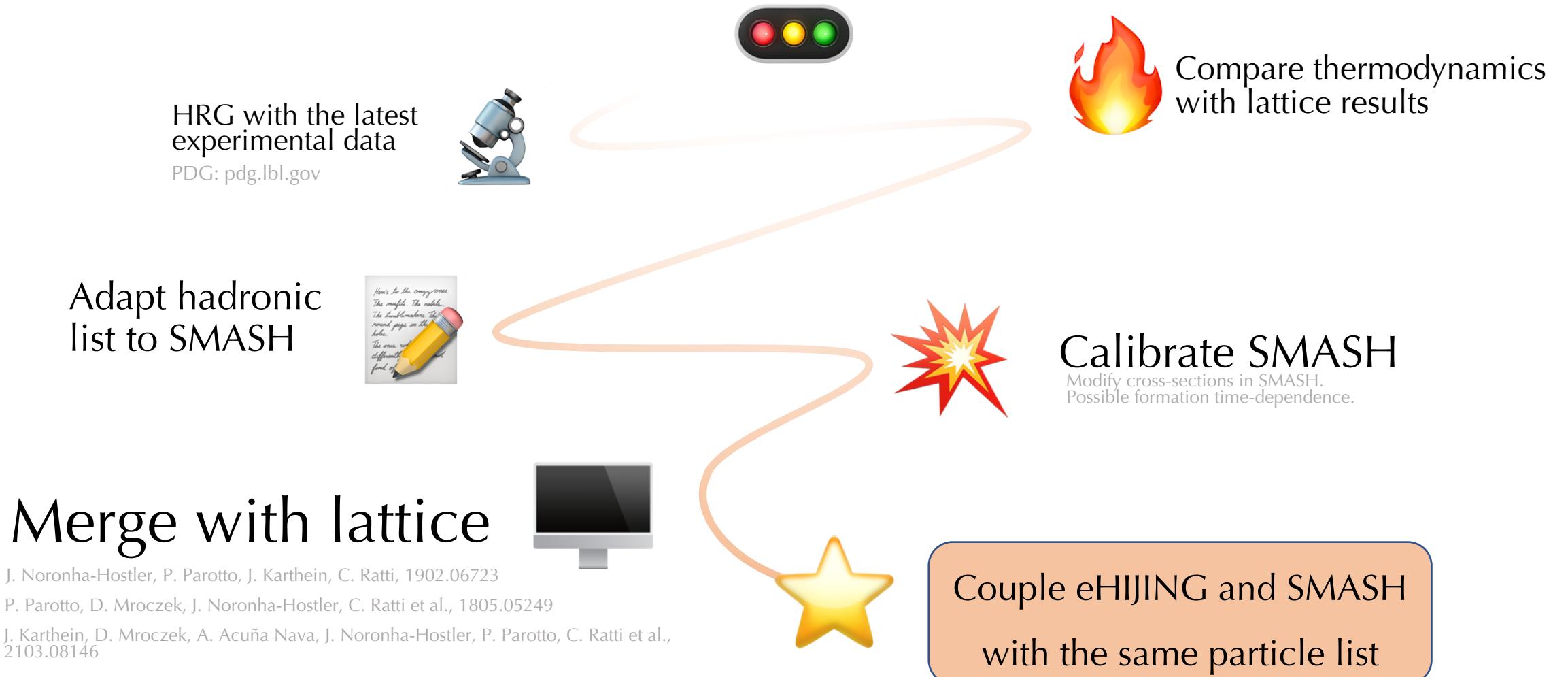
Using the latest resonances from PDG in SMASH and their impact on cross sections

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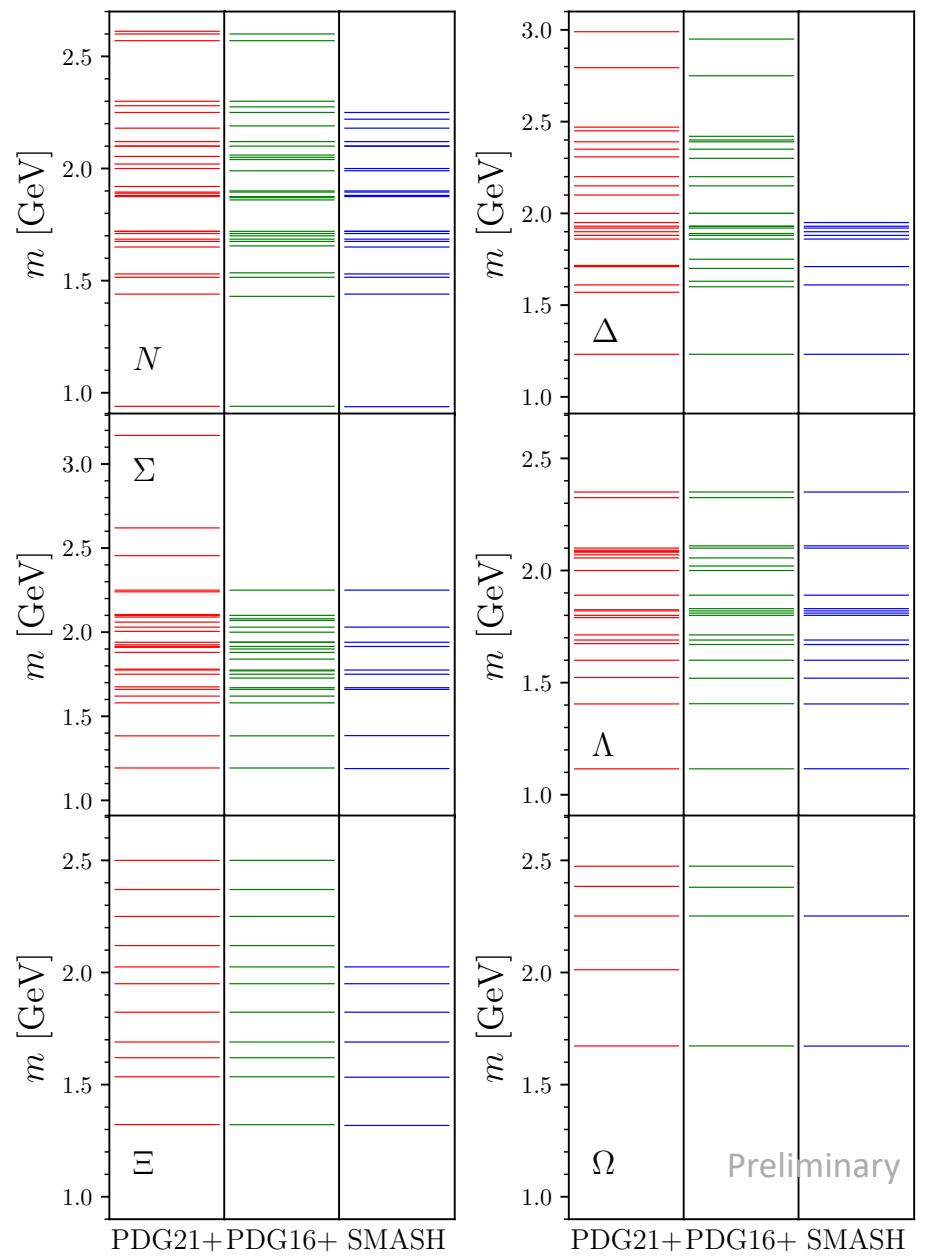
In collaboration with: R. Hirayama, J. Hammelmann,
J. Karthein, P. Parotto, J. Noronha-Hostler, H. Elfner,
C. Ratti, MUSES Collaboration

Roadmap to coupling SMASH and eHIJING

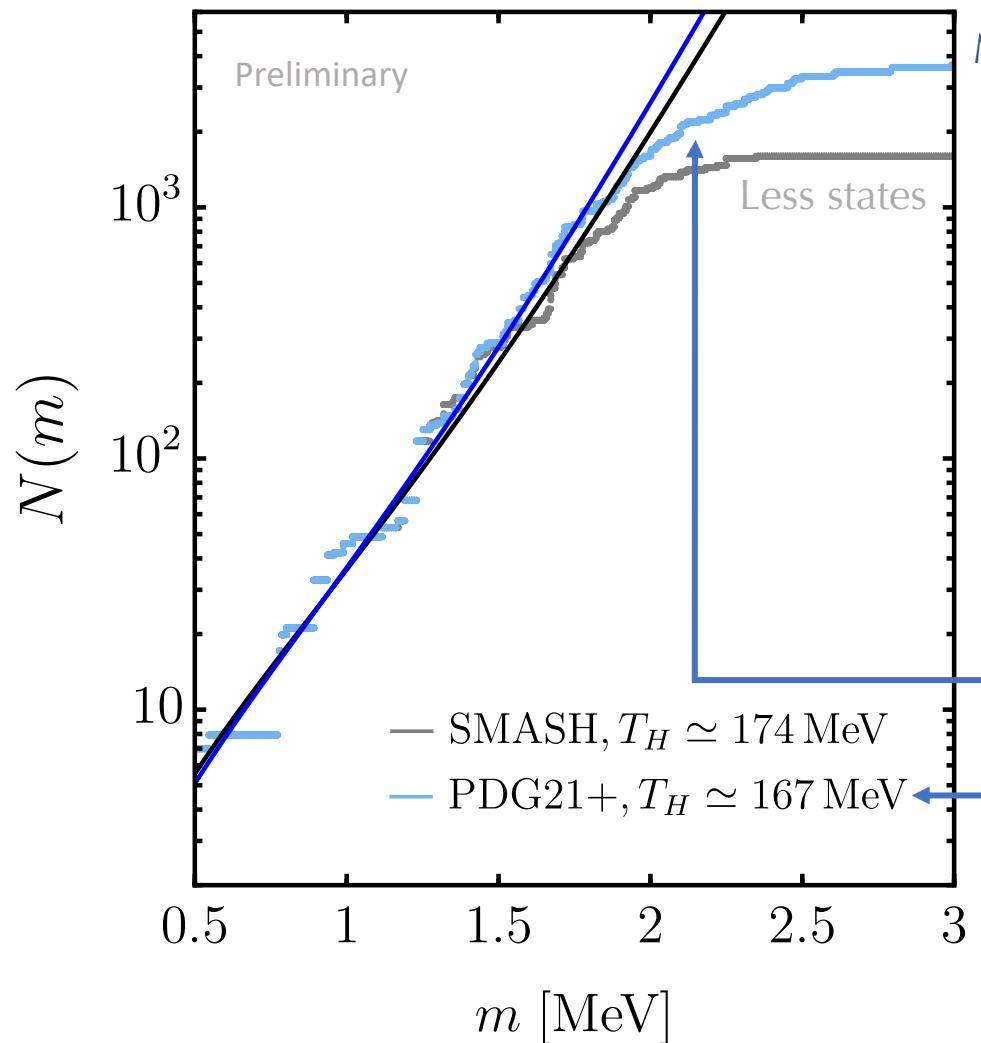


Latest PDG data

- Centralized hadronic database based on the PDG
- Tracks several particle properties, e.g., mass, width, isospin, etc.
- Has all hadrons and their reported branching ratios
- 760 particles
- Updated branching ratios vs. PDG16+
- Contains *-**** particles



A lower limiting temperature



$$N(m) = \int_0^m \rho(m') dm'; \quad \rho(m) \sim m^{-a} e^{m/T_H}$$

Hagedorn observed an exponentially rising spectrum which led to the conclusion there was a limiting hadron temperature

Adding more hadronic states supports the exponentially rising spectrum

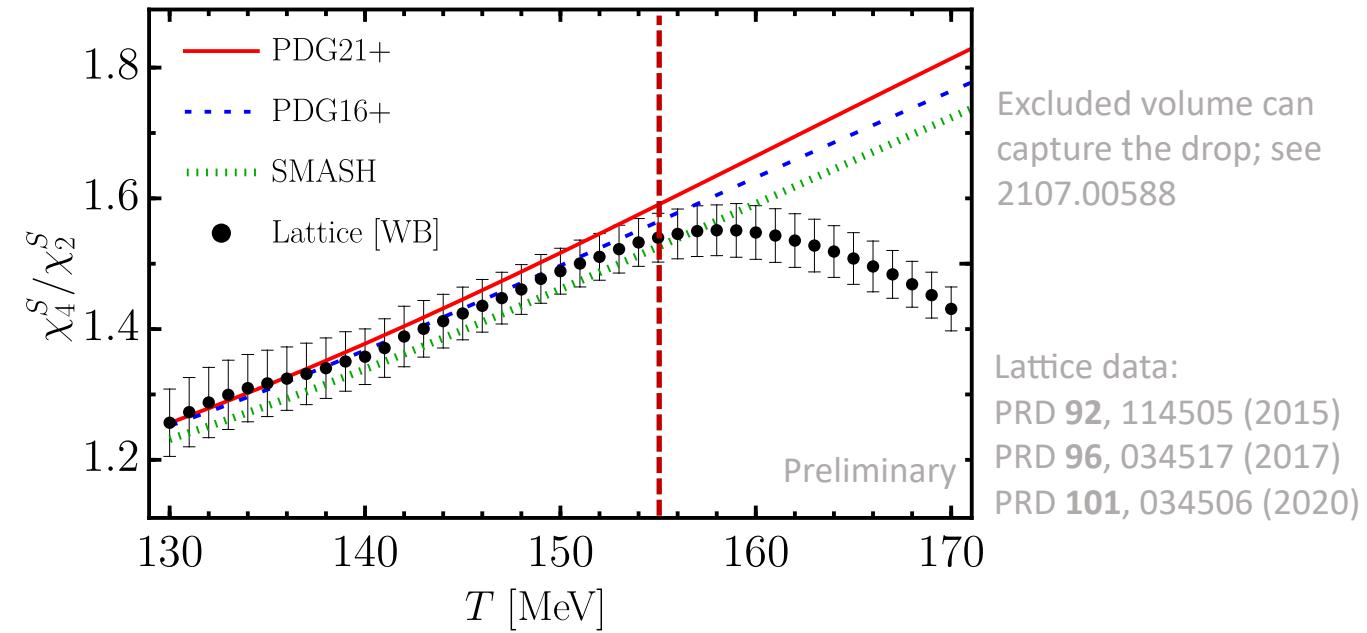
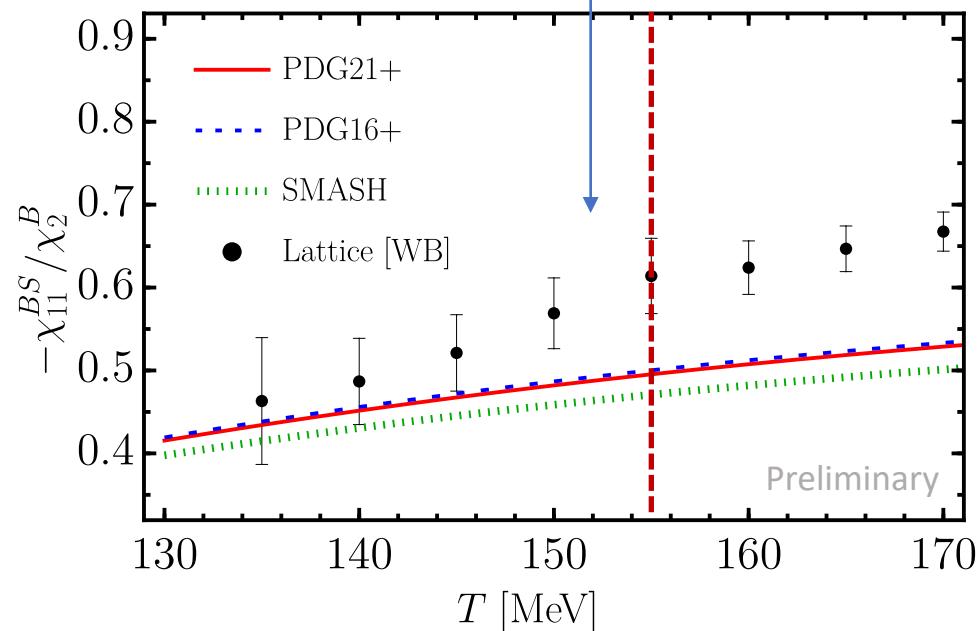
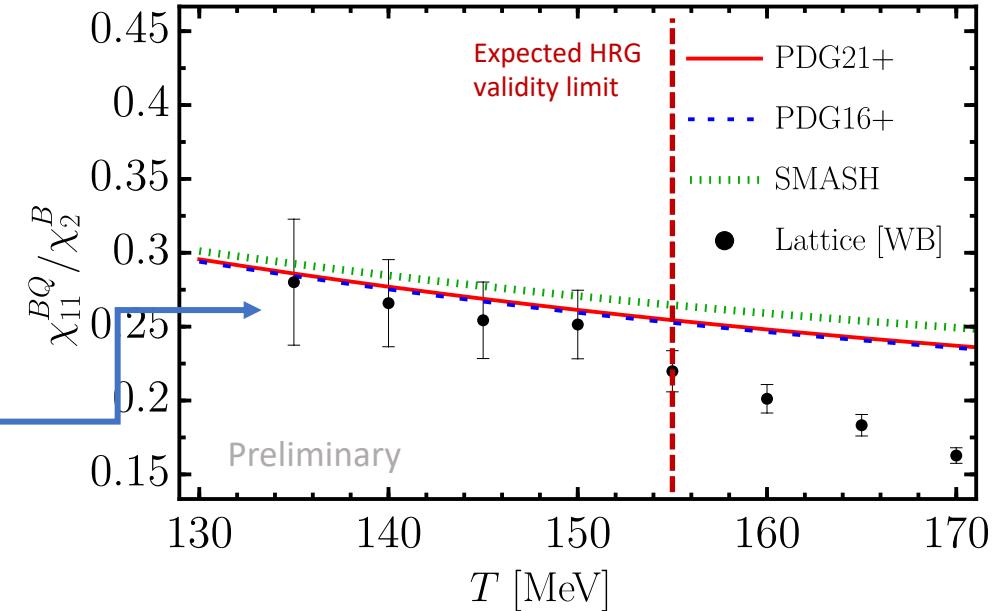
The extracted temperature becomes closer to the pseudo-critical temperature from LQCD

HRG susceptibilities vs LQCD

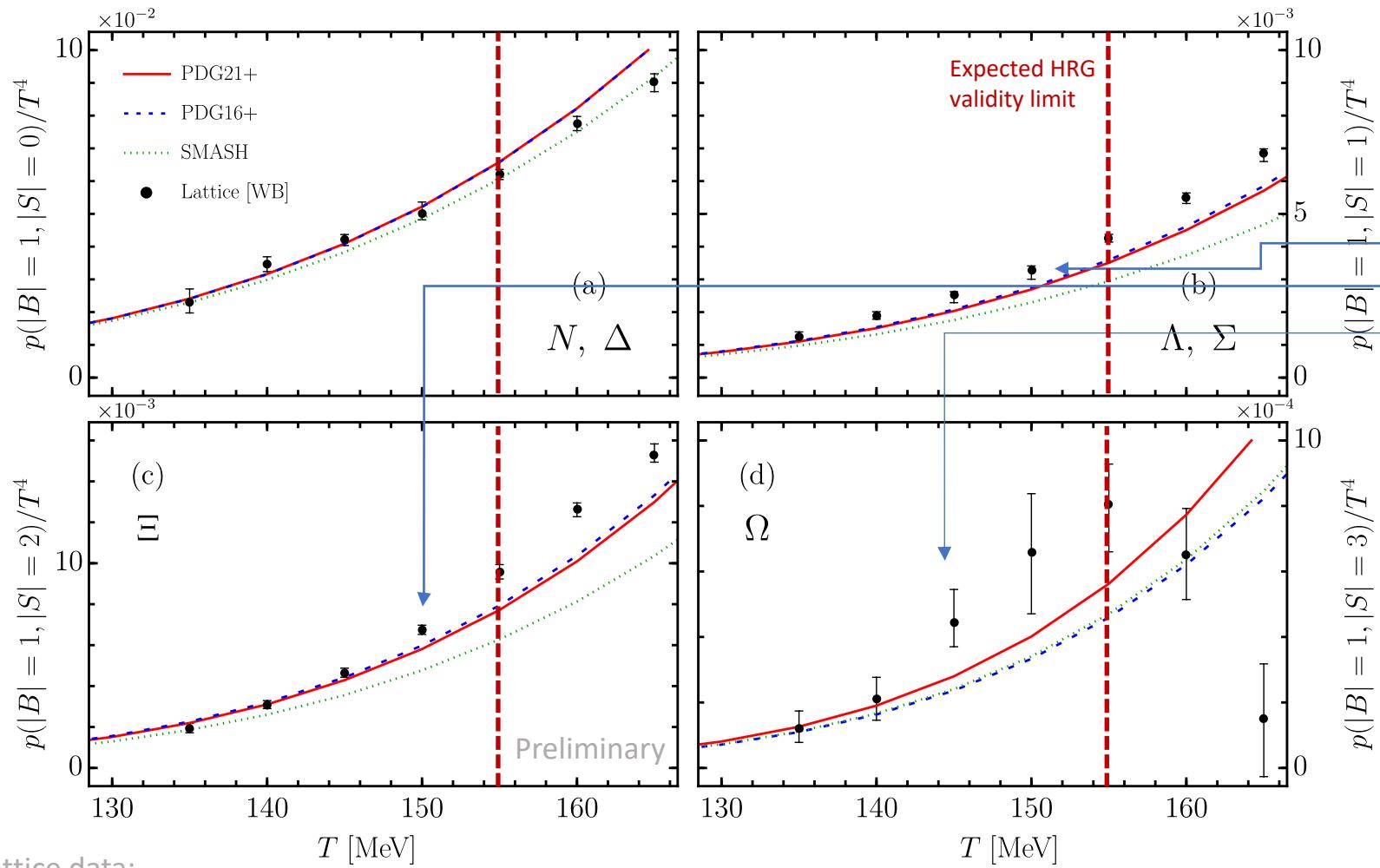
Disagreement with lattice data hints at missing strange resonances ($\Lambda, \Sigma, \Xi, \Omega$)

A flavor-dependent excluded volume could improve this comparison; see 2107.00588

The new list is consistent with both the previous PDG2016+; more strange particles could improve this too



HRG partial pressures vs LQCD



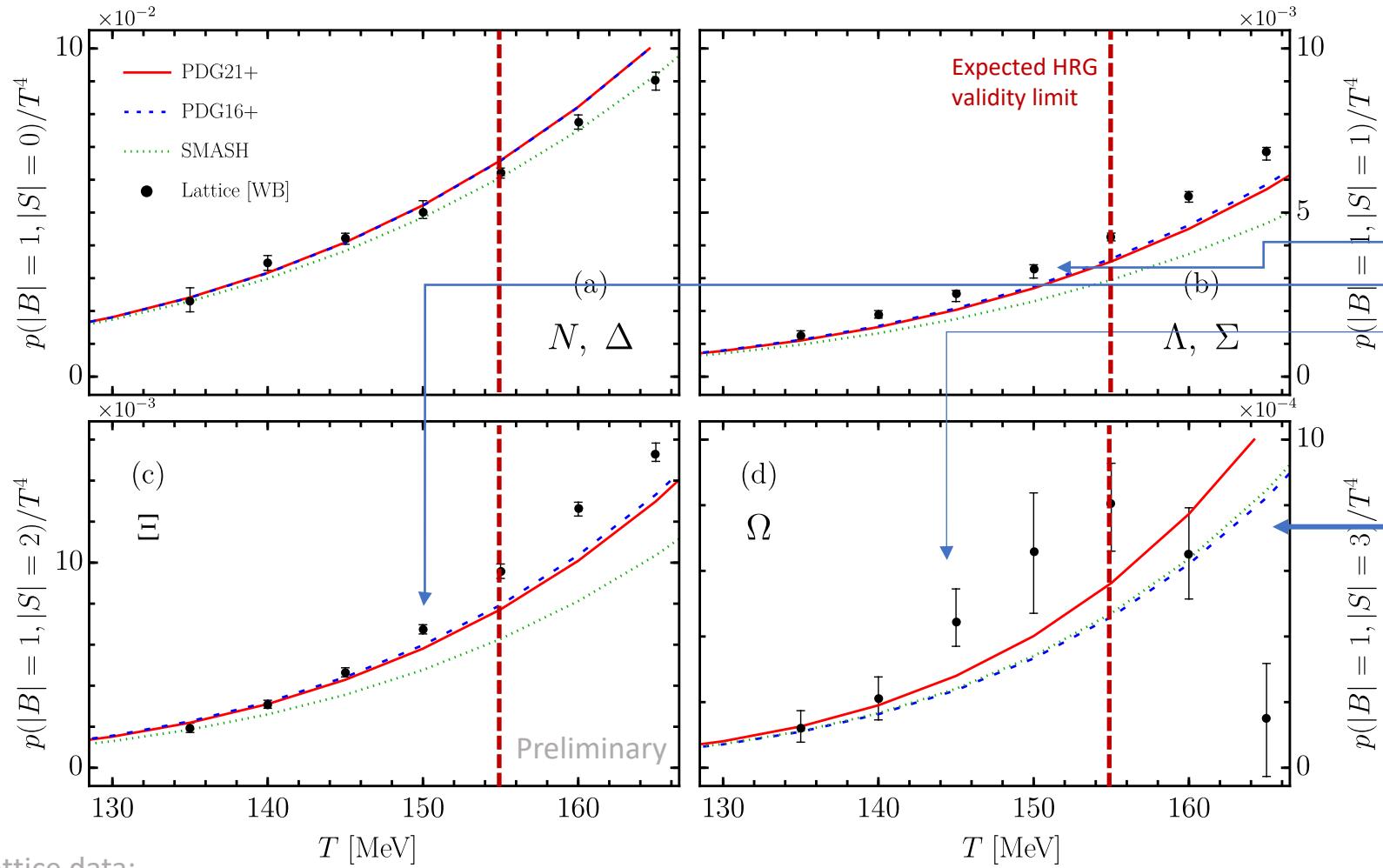
The new PDG2021+ list is in agreement with the previous results from PDG2016+.

Disagreement with lattice data hints at missing strange resonances ($\Lambda, \Sigma, \Xi, \Omega$)

see KLF Collaboration proposal at JLAB,
2207.10779

$$\begin{aligned} \frac{p}{T^4} = & \phi_0 + \phi_{01} \cosh(\mu_S/T) \\ & + \phi_{10} \cosh(\mu_B/T) \\ & + \phi_{11} \cosh(\mu_B/T - \mu_S/T) \\ & + \phi_{12} \cosh(\mu_B/T - 2\mu_S/T) \\ & + \phi_{13} \cosh(\mu_B/T - 3\mu_S/T) \end{aligned}$$

HRG partial pressures vs LQCD



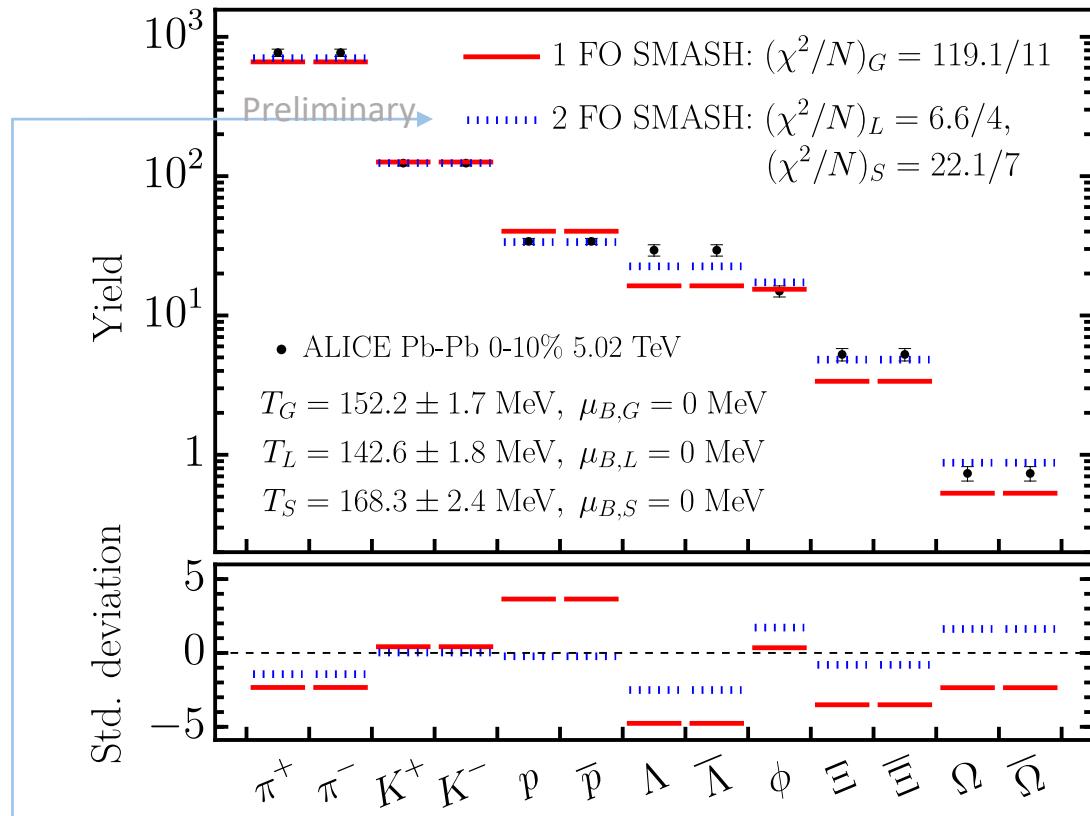
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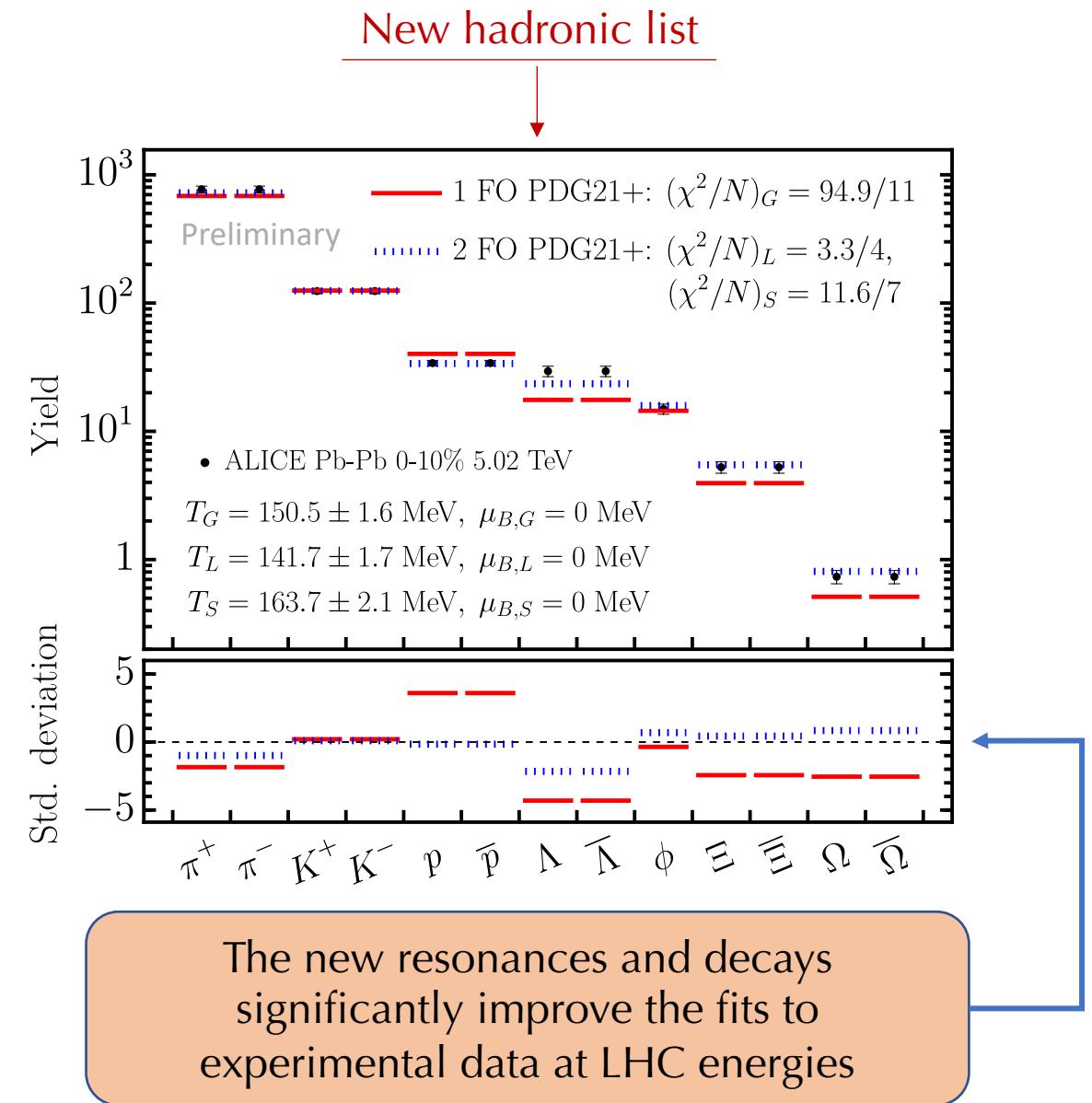
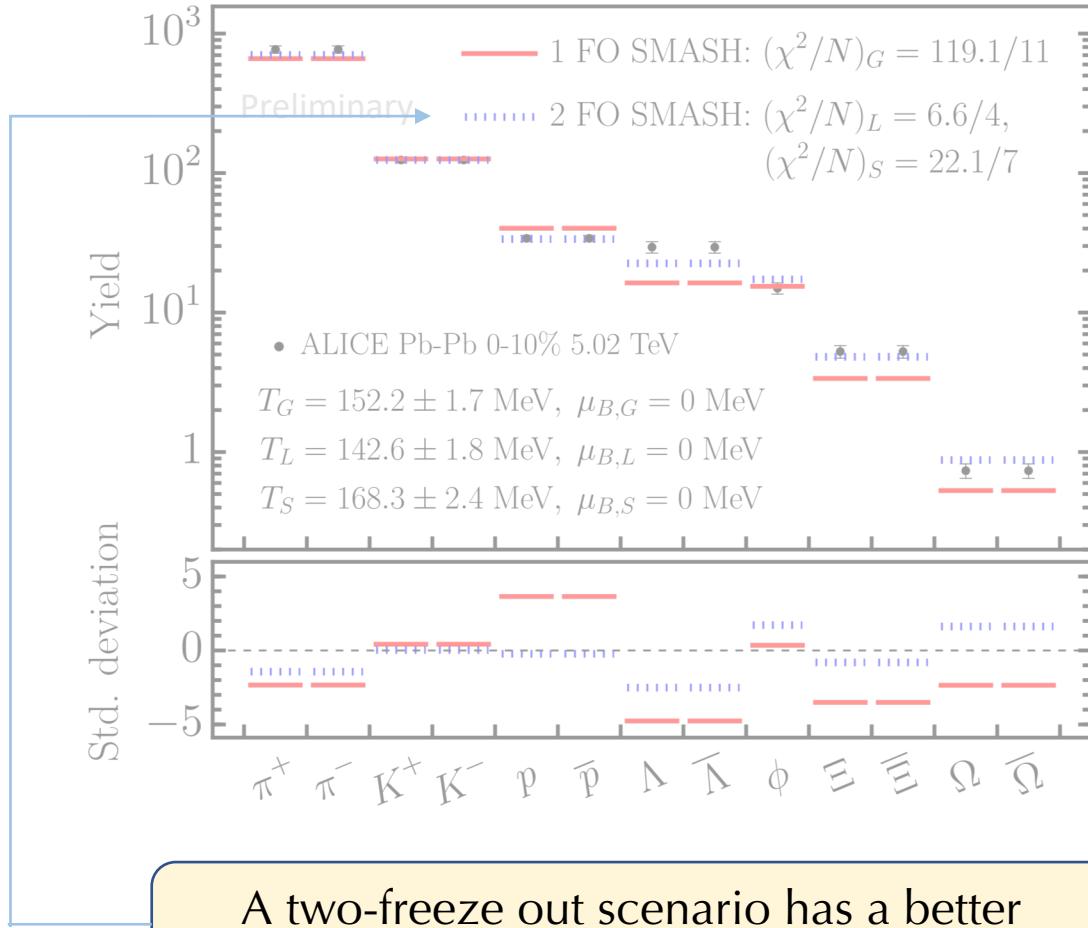
New high-confidence Ω baryon contributes pronouncedly to the triple-strange partial pressure

Thermal model yields



A two-freeze out scenario has a better agreement with experimental data

Thermal model yields



Modeling the list with intermediate states

1 → 2 decays needed for
SMASH

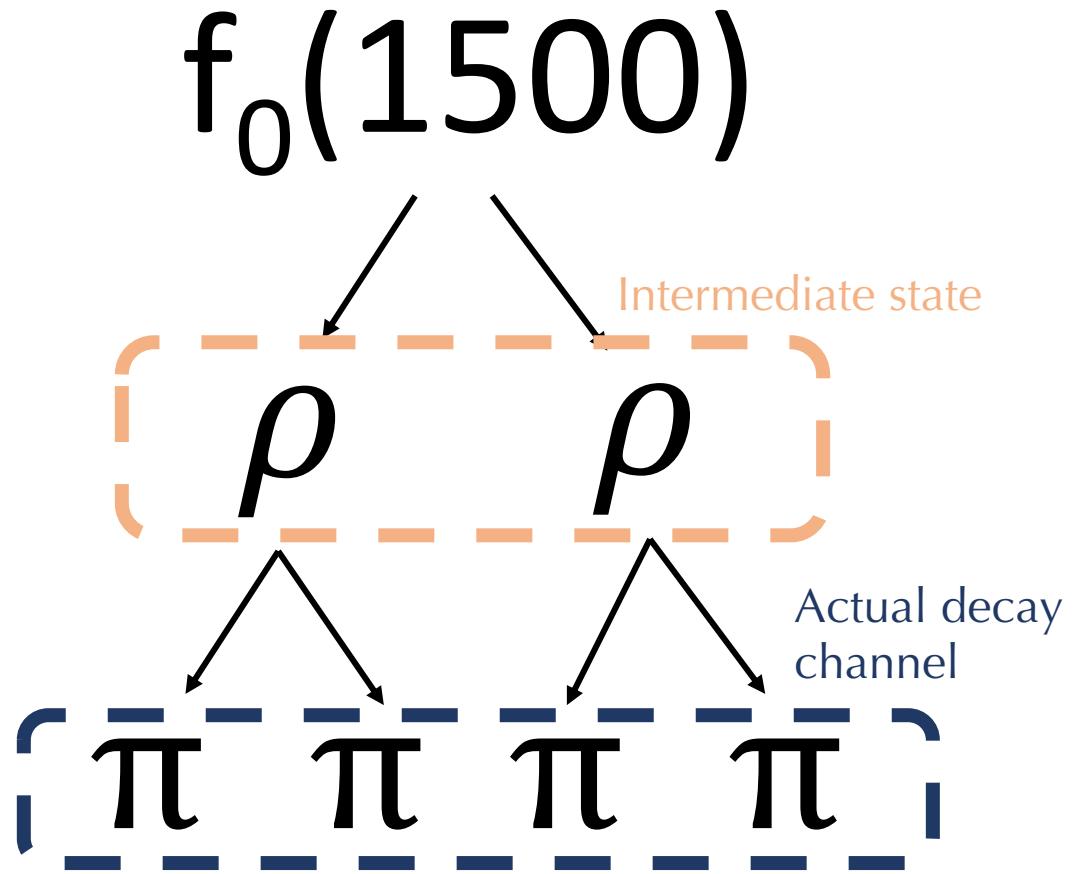
$$d_{\text{trans}} < \sqrt{\frac{\sigma_{\text{tot}}}{\pi}}$$

J. Weil et al., PRC **94** (2016) 054905
D. Oliinychenko et al., SMASH-transport (2021),
<https://doi.org/10.5281/zenodo.5796168>

Model 3 and 4-body decays
with intermediate states

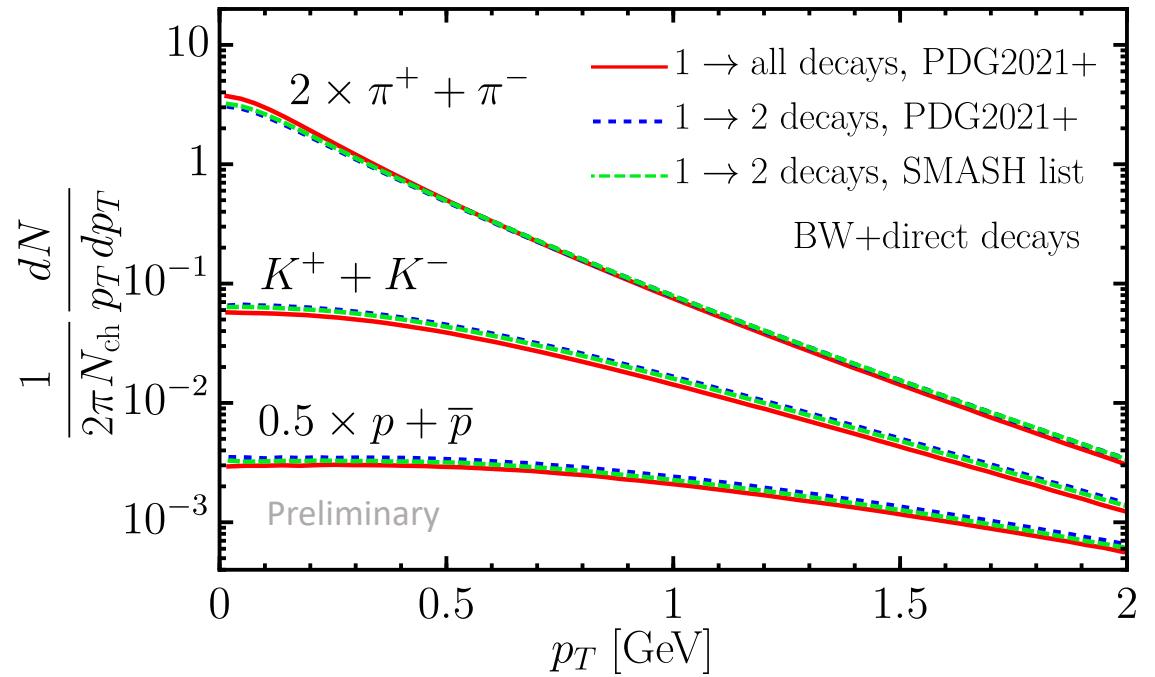


SMASH input:
1. Particle list
2. Decay modes



Identified particles spectra

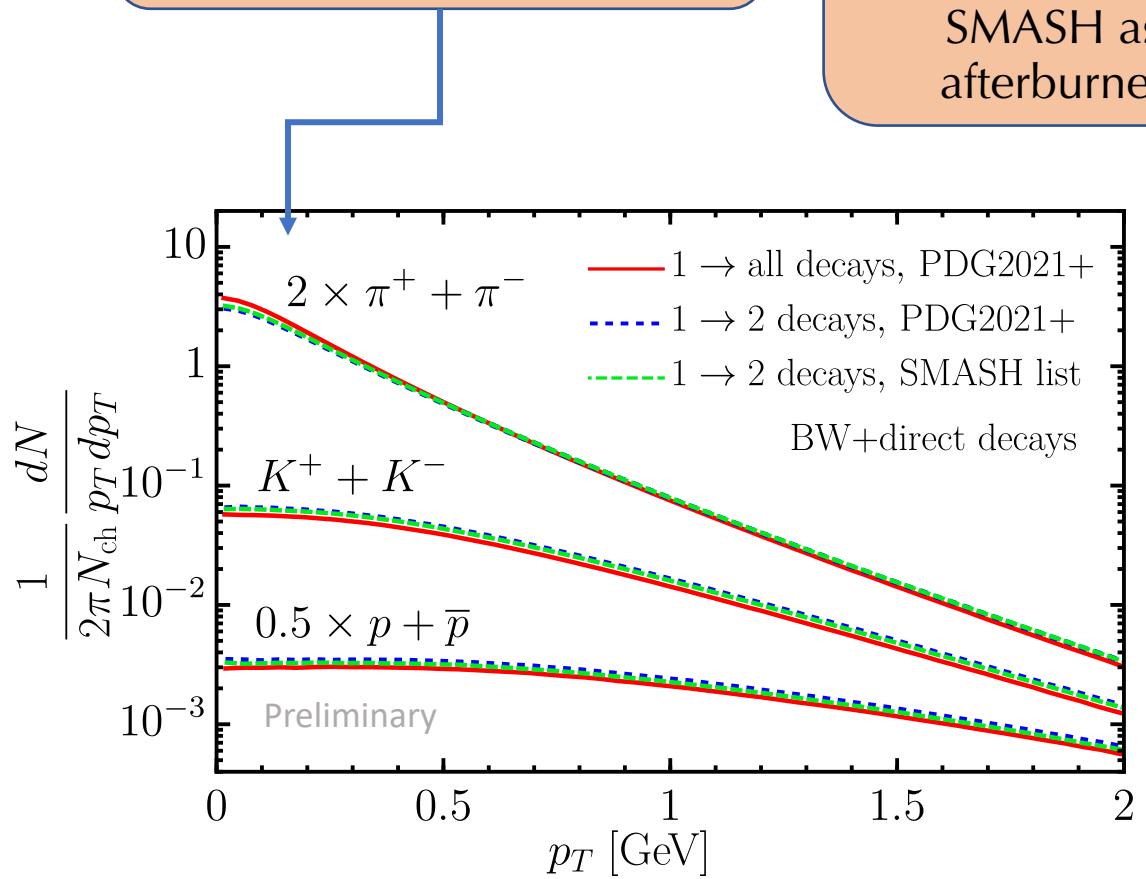
The addition of more resonances and modification of decay channels has an effect on the particle spectra



Identified particles spectra

The adaptation of the list to
only have $1 \rightarrow 2$ decays
causes the slope to change

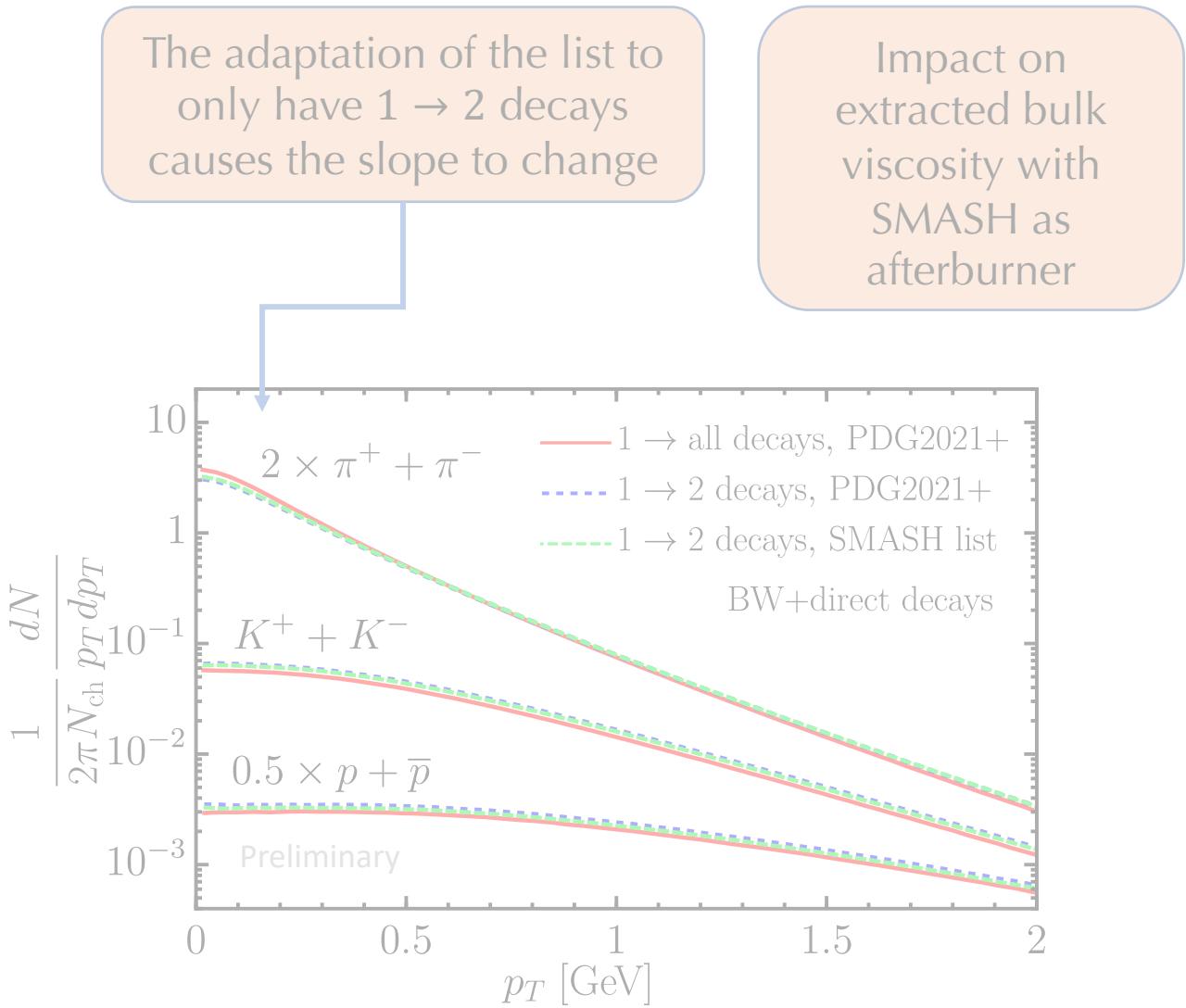
Impact on
extracted bulk
viscosity with
SMASH as
afterburner



Identified particles spectra

BW+direct decays $\pi^+ + \pi^-$	
SMASH list	0.548 ± 0.001
PDG2021+ (1 → 2 decays)	0.551 ± 0.001
PDG2021+ (1 → all decays)	0.523 ± 0.001
BW+SMASH $\pi^+ + \pi^-$	
SMASH list	0.5463 ± 0.0001
PDG2021+ (1 → 2 decays)	0.5467 ± 0.0001
Experiment	0.56965 ± 0.02505
Preliminary	
BW+direct decays $K^+ + K^-$	
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Data: PRC 101, 044907 (2020)



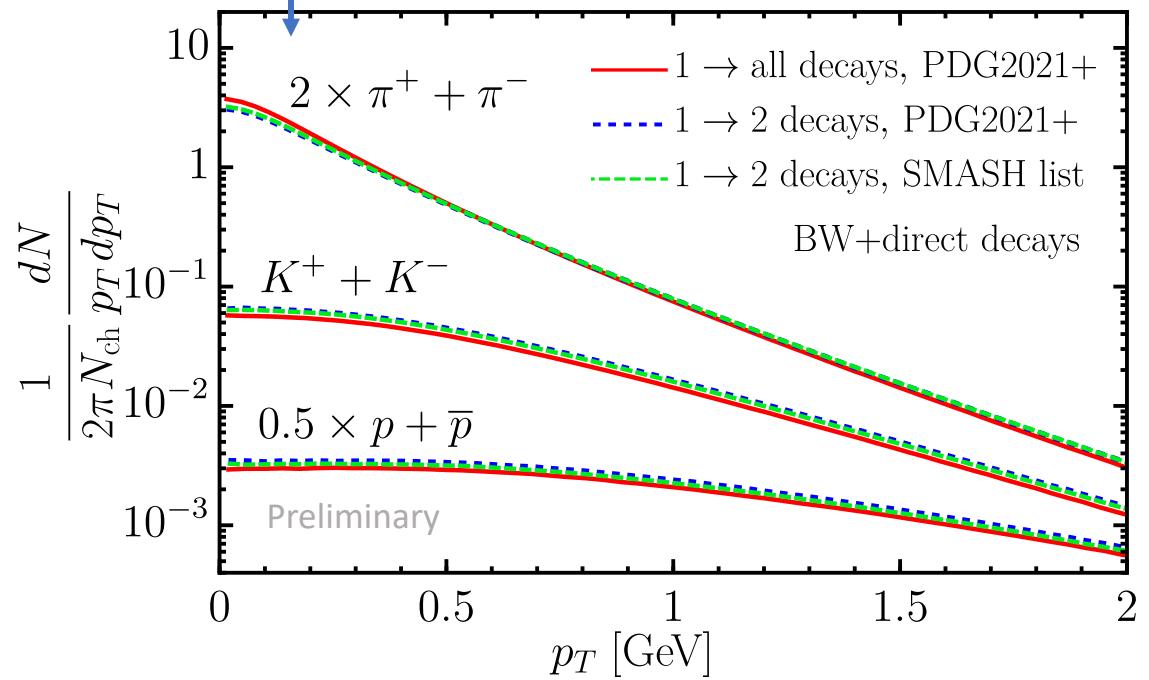
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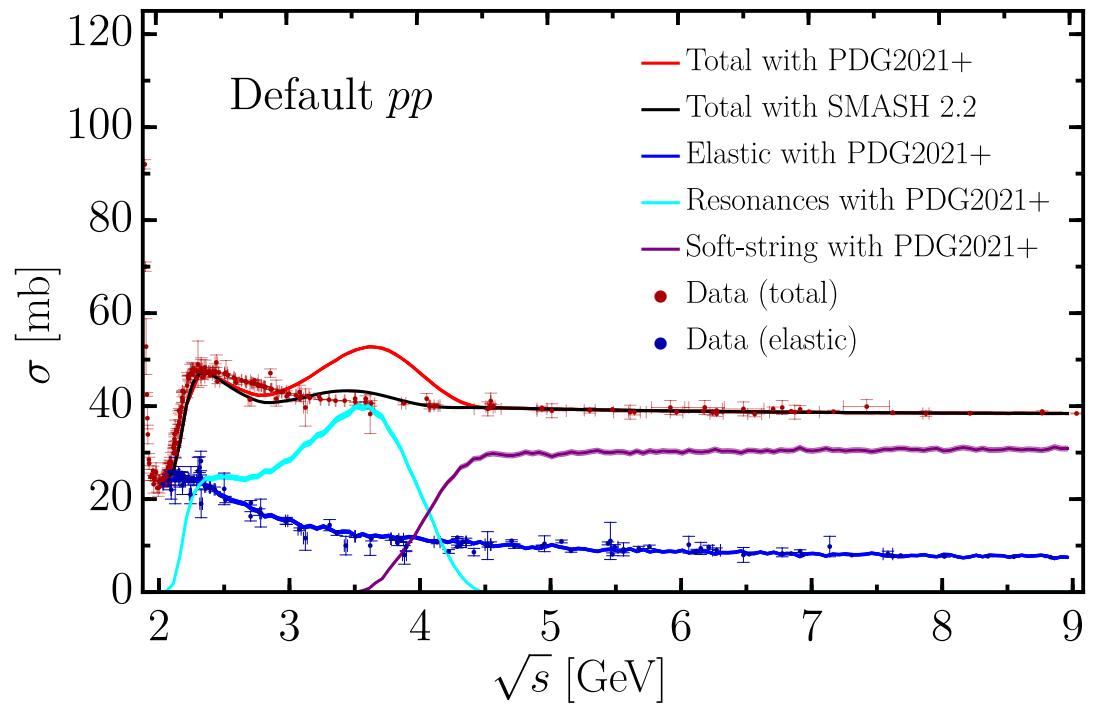
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Although changing the list does not affect the $\langle p_T \rangle$ too much, it impacts spectra slightly

The mean-transverse momentum spectra is mainly affected by using 1 → 2 or 1 → all decays

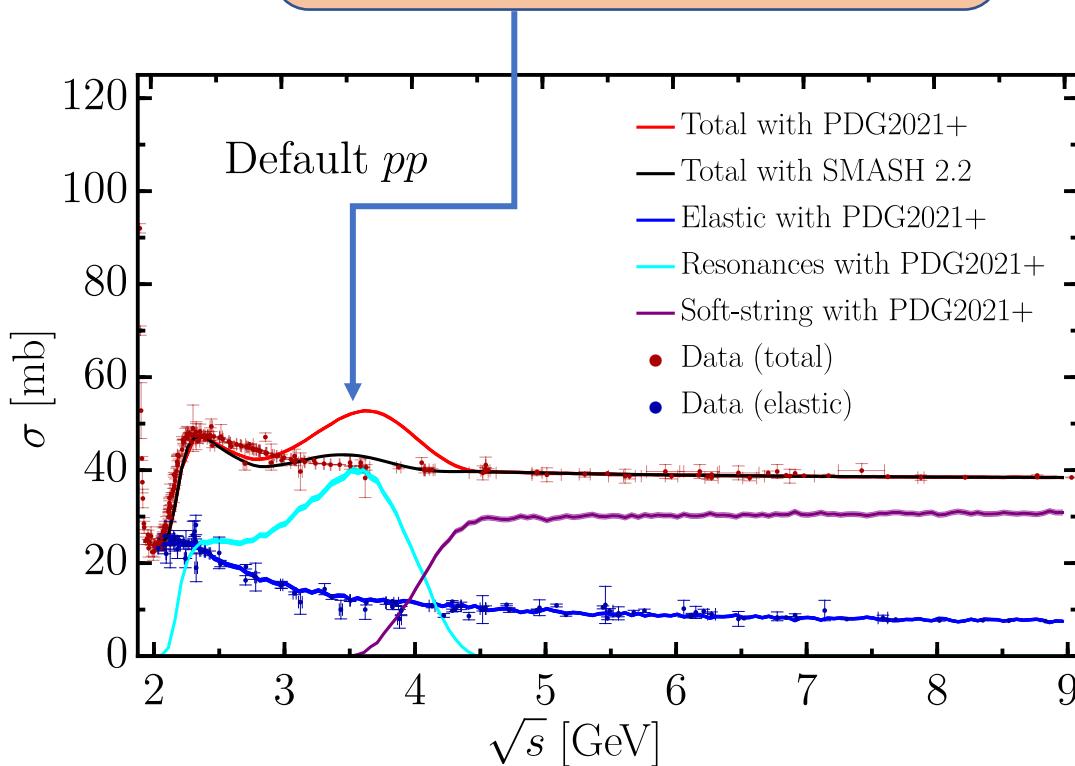


Cross section rescaling



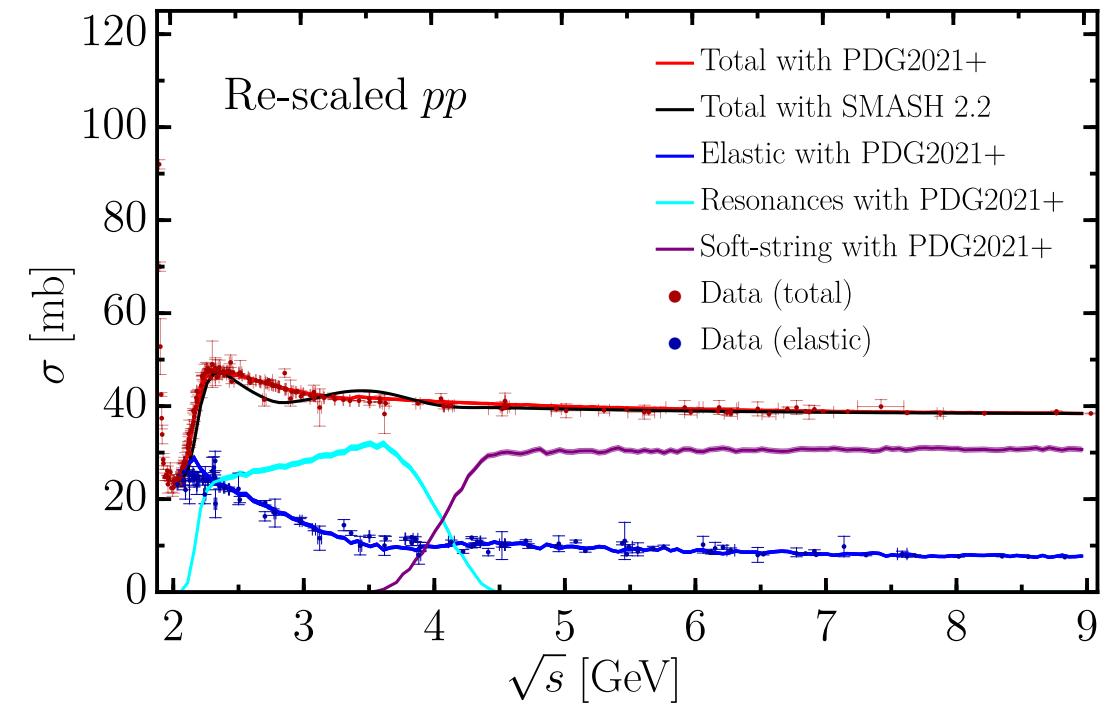
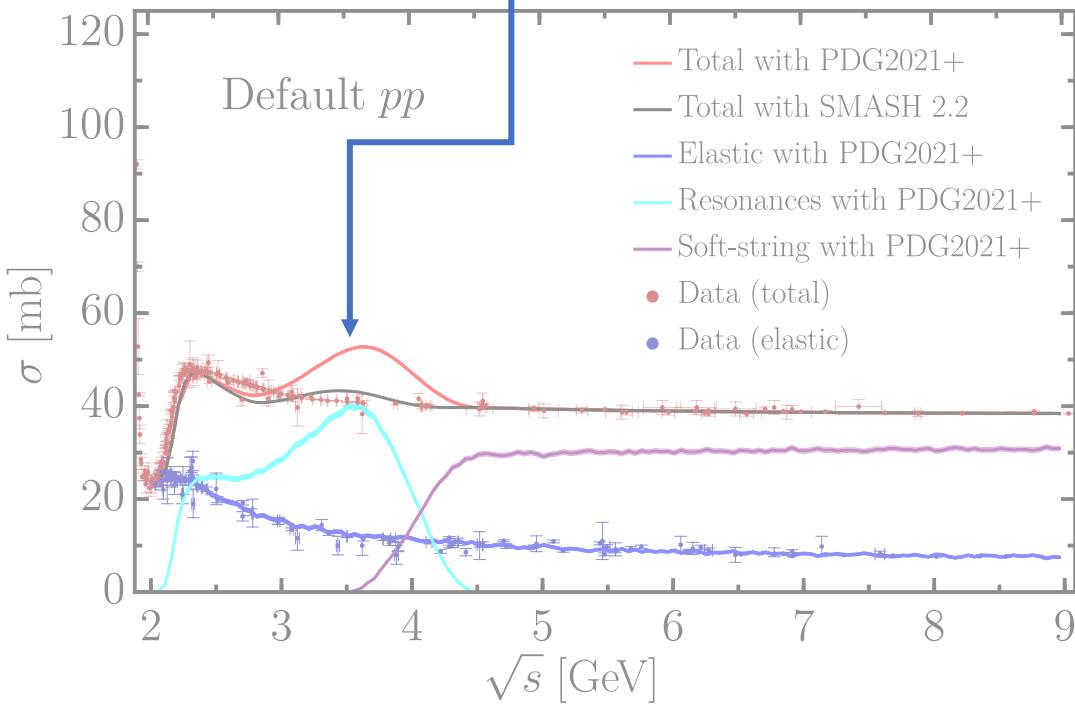
Cross section rescaling

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Cross sections and string fragmentation

SMASH implements string excitation and fragmentation for hard processes

A string forms hadrons by producing quark-antiquark pairs
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In SMASH a cross section scaling factor controls *when* hadrons are allowed to interact, i.e., normally at the point of closest approach

$$f_\sigma(t) = (1 - f_0) \left(\frac{t - t_{\text{prod}}}{t_{\text{form}} - t_{\text{prod}}} \right)^\alpha + f_0$$

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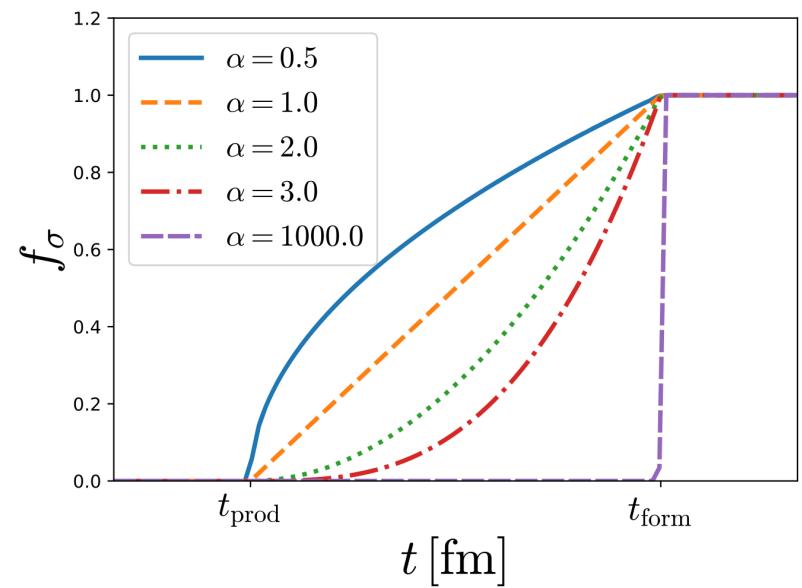
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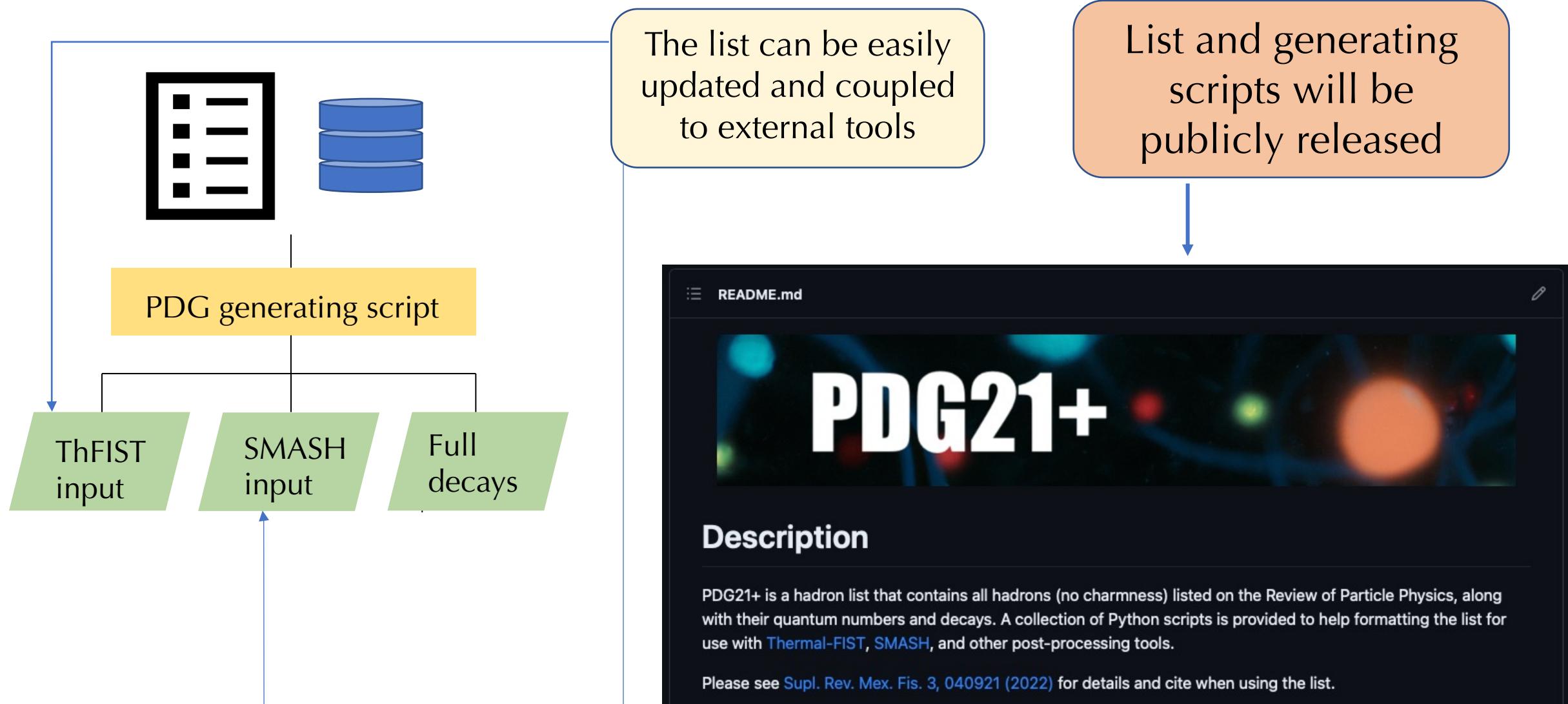
The scaling factor can be a function of time, to mimic a continuous formation process

SMASH on scaling factor:
J. Phys. G **47**, 065101 (2020)

Previous study using GiBUU:
Nucl. Phys. A **801**, 68 (2008)



Open-source code



Conclusions

- Lattice hints at additional strange hadronic states
- A new list, PDG21+, was built with the latest experimental data available
- The new resonances come with a need to retune cross sections
- Final state hadronic interactions can be studied as functions of formation times
- If SMASH is used as a hadron scattering phase evolver, one wants a consistent treatment of the particle list → updated SMASH particle list

Outlook

- We have to make sure to use the same resonances in SMASH and eHIJING
- We can test the formation time dependence of final state interactions by modifying the scaling factor
- Previous studies have found that indeed time dependent cross-sections seproduce data more accurately
- Extensions of these studies can be tested on HERMES and EMC data at low energies
- Further extrapolation to cover EIC can be done by coupling to PYTHIA