

# Precise charm- and bottom-quark mass determinations & multi-loop calculations to polarization functions

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- I. Introduction & Motivation
- II. Method & Calculation
- III. Analysis & Results
- IV. Summary & Conclusion

based on [Phys. Rev. D80 \(2009\)](#), [Nucl. Phys. B 778, 192 \(2007\)](#)

In collaboration with:

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# I. Introduction

## Motivation

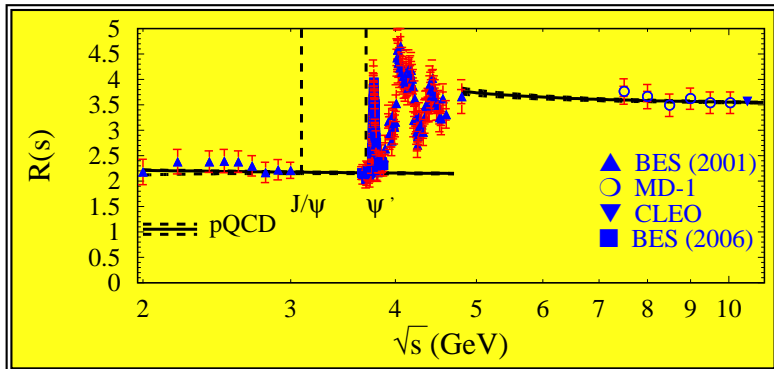
Precise determination of the **charm-** and **bottom-quark masses** important:

- **Quark masses** are fundamental parameters of the Standard Model  $\rightsquigarrow$  enter in many physical observables
- **Quark masses** play an important role in Higgs physics:  
e.g. Higgs decays:  
SM Higgs boson light  $\rightsquigarrow$  dominant decay into  $b\bar{b}$   
$$\Gamma(H \rightarrow b\bar{b}) = \frac{G_f M_h}{4\sqrt{2}\pi} m_b^2 (1 + \mathcal{O}(\alpha_s) + \dots), \quad \Gamma(H \rightarrow c\bar{c}) \sim m_c^2$$
- **Quark masses** relevant in flavor physics:  
e.g.  $B$  meson decays:  $\Gamma \propto m_b^5$ ,  $B \rightarrow X_u \ell \bar{\nu}$ ,  $B \rightarrow X_c \ell \bar{\nu}$   
Virtual **charm quarks**:  $K \rightarrow \pi \nu \bar{\nu}$ ,  $B \rightarrow X_s \gamma$
- Comparison with other methods, e.g. lattice methods  
 $\rightsquigarrow$  valuable, mutual cross-checks

## II. Method

Experiment:  $R$ -ratio

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{Hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$




## II. Method

### Theory

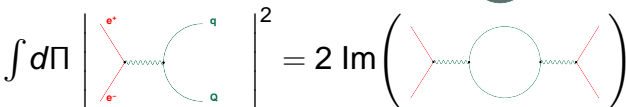
- Heavy quark correlator

$$\Pi^{\mu\nu}(q, j) = i \int dx e^{iqx} \langle 0 | T j^\mu(x) j^\nu(0) | 0 \rangle$$

Here:  $j^\mu(x)$  electromagnetic heavy quark vector current

$$\Pi^{\mu\nu}(q) = (-g^{\mu\nu} + q^\mu q^\nu / q^2) \Pi(q^2) \sim \text{diagram}$$
A diagram showing a heavy quark loop (represented by a grey circle) with a photon line (represented by a wavy line) entering and exiting the loop. The photon line is labeled with momentum q.

- $\int d\Pi \left| \text{diagram} \right|^2 = 2 \text{Im} \left( \text{diagram} \right)$

The equation shows the squared magnitude of a diagram (a photon line entering a quark loop) is equal to twice the imaginary part of the same diagram. The diagrams use red lines for quarks and green lines for photons.

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{Hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 12\pi \text{Im} [\Pi(q^2 = s + i\epsilon)]$$

- With the help of dispersion-relations:

$$\Pi(q^2) = \Pi(q^2 = 0) + \frac{q^2}{12\pi^2} \int ds \frac{R(s)}{s(s - q^2)}$$

- **Exp. moments** are related to derivatives of  $\Pi(q^2)$  at  $q^2 = 0$

## II. Method

Relation: Theory  $\iff$  Experiment

- **Exp. moments** are related to derivatives of  $\Pi(q^2)$  at  $q^2 = 0$ :

$$\frac{12\pi}{n!} \left( \frac{d}{dq^2} \right)^n \Pi(q^2) \Big|_{q^2=0} = \mathcal{M}_n^{\text{exp}} = \int \frac{ds}{s^{n+1}} R^{\text{exp}}(s)$$

- In terms of **expansion coefficients**:

$$\Pi(q^2) = \frac{3Q_f^2}{16\pi^2} \sum_n \bar{C}_n^v \left( \frac{q^2}{4m^2} \right)^n, \quad \begin{array}{l} Q_f: \text{charge of quark} \\ m = m(\mu) : \overline{\text{MS}} \text{ mass} \end{array}$$

$\bar{C}_n^v$  can be calculated perturbatively

- **First and higher derivatives** of  $\Pi(q^2)$  allow direct determination of the  $\overline{\text{MS}}$  charm- and bottom-quark mass:

$$\bar{m}(\mu) = \frac{1}{2} \left( Q_f^2 \frac{9}{4} \bar{C}_n^v \mathcal{M}_n^{\text{exp}} \right)^{1/(2n)}$$

← Theory

← Experiment

c-quarks: Novikov et al. '78; b-quarks: Reinders et al. '85

$\bar{C}_n^v$  depend on the quark mass through  $\log(m(\mu)^2/\mu^2)$

## II. Calculation

Pert. calculation of expansion coefficients

### ■ Sample diagrams

$$\Pi^{\mu\nu}(q, j) = \text{tree} + \text{QCD-corrections} \rightarrow \text{1-loop} + \text{2-loop} + \text{3-loop} + \text{4-loop}$$

### ■ Expansion diagrammatically:

$$\text{tree} \rightarrow \text{1-loop} + q^2 \left( \text{2-loop} + \text{3-loop} + \dots \right) \dots$$

↪ One-scale multi-loop integrals in pQCD

- 3-loop (order  $\alpha_s^2$ ) coefficients  $\bar{C}_n$  up to  $n=8$  Chetyrkin, Kühn, Steinhauser 96  
up to higher moments  $n \sim 30$  Czakon et al. 06; Maierhöfer, Maier, Marquard 07  
for correlators  $VV, AA, PP, SS$

# II. Calculation

Techniques, IBP, MI

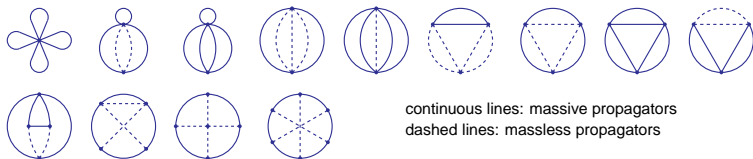
Computation consists of two steps:

- First step:

Reduction to a small set of master integrals,  
Integration by parts techniques

- Second step:

Computation of master integrals  
Here: 13 master integrals



continuous lines: massive propagators  
dashed lines: massless propagators

Solution with high precision numeric Y. Schröder, A. Vuorinen  
with difference equation method S. Laporta

Subsequently with independent method:  $\epsilon$ -finite basis

K.G. Chetyrkin, M. Faisst, C.S., M. Tentyukov

other contributions: D.J. Broadhurst; S. Laporta; B.A. Kniehl, A.V. Kotikov; Y. Schröder, M. Steinhauser

Analytical results in sufficient deep order

# II. Calculation

## Results at 4-loops

### R-ratio method:

#### – Vector case:

- first moments  $\overline{C}_0, \overline{C}_1$

K. G. Chetyrkin, J. H. Kühn, C.S.'06; R. Boughezal, M. Czakon, T. Schutzmeier'06

- second moment  $\overline{C}_2$  A. Maier, P. Maierhöfer, P. Marquard'08

- third moment  $\overline{C}_3$  A. Maier, P. Maierhöfer, P. Marquard, A.V. Smirnov '09 ← new

- fourth moment  $\overline{C}_{4,\dots,10}$  Y. Kiyo, A. Maier, P. Maierhöfer, P. Marquard '09 ← new

Lattice method: Replace  $\mathcal{M}_n^{exp}$  by lattice sim. HPQCD+ K. Chetyrkin, J. Kühn, M. Steinhauser, C.S.

#### – Pseudoscalar case:

- first moments  $\overline{C}_0, \overline{C}_1, \overline{C}_2$  I. Allison, E. Dalgic, C.T.H. Davies, E. Follana, R.R. Horgan, K. Hornbostel, G.P. Lepage, C. McNeile, J. Shigemitsu, H. Trotter, R.M. Woloshyn, K.G. Chetyrkin, J.H. Kühn, M. Steinhauser, C.S. 08

- third moment  $\overline{C}_3$  A. Maier, P. Maierhöfer, P. Marquard'08

- fourth moment  $\overline{C}_4$  A. Maier, P. Maierhöfer, P. Marquard, A.V. Smirnov '09

- fifth moment  $\overline{C}_{5,\dots,10}$  Y. Kiyo, A. Maier, P. Maierhöfer, P. Marquard '09

#### – Axial-vector and scalar case:

- first moments  $\overline{C}_0, \overline{C}_1$  C. S.'08

- third moment  $\overline{C}_3$  A. Maier, P. Maierhöfer, P. Marquard, A.V. Smirnov '09

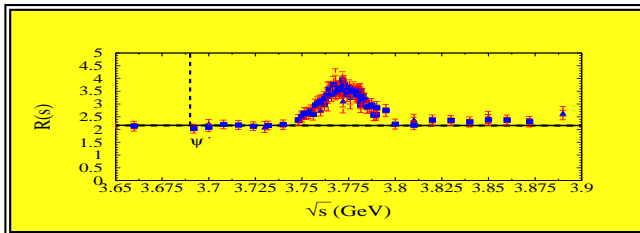
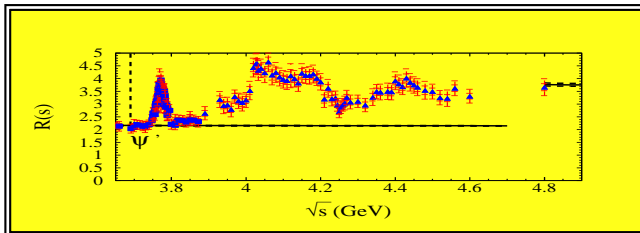
- fourth moment  $\overline{C}_{4,\dots,10}$  Y. Kiyo, A. Maier, P. Maierhöfer, P. Marquard '09



# III. Analysis

R-ratio

Determine:  $\mathcal{M}_n^{\text{exp}} = \int \frac{ds}{s^{n+1}} R^{\text{exp}}(s)$



# III. Analysis

Extraction of the exp. moments from  $R(s)$  (charm quark case)

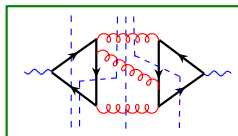
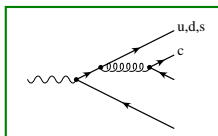
Determine:  $\mathcal{M}_n^{\text{exp}} = \int \frac{ds}{s^{n+1}} R^{\text{exp}}(s) = \mathcal{M}_n^{\text{res}} + \mathcal{M}_n^{\text{thr}} + \mathcal{M}_n^{\text{cont}}$

For charm quarks:

$\mathcal{M}_n^{\text{res}}$ : Contains:  $J/\psi, \psi(2S)$  treated in narrow width approximation

$$R^{\text{res}}(s) = \frac{9\pi M_R \Gamma_{ee}}{\alpha^2} \left( \frac{\alpha}{\alpha(s)} \right)^2 \delta(s - M_R^2)$$

$\mathcal{M}_n^{\text{thr}}$ : BES data ( $\sqrt{s} \geq 3.73$  GeV) subtract background from  $R_{uds}$ ,



$\bar{R}$  from data below 3.73 GeV,  $\sqrt{s}$ -dependence from theory

$\mathcal{M}_n^{\text{cont}}$ : pQCD above  $\sqrt{s} \geq 4.8$  GeV ,

spare data,

$R(s)$  with full quark mass dependence rhad: R. Harlander, M. Steinhauser '02

# III. Analysis & Results

Determination of the charm quark mass from  $R(s)$

## ■ Charm quark mass:

$$\mu = (3 \pm 1) \text{ GeV} \quad \alpha_s(M_Z) = 0.1189 \pm 0.002$$

$n$	$m_c(3 \text{ GeV})$	exp	$\alpha_s$	$\mu$	np	total
1	0.986	0.009	0.009	0.002	0.001	0.013
2	0.976	0.006	0.014	0.005	0.000	0.016
3	0.978	0.005	0.015	0.007	0.002	0.017
4	1.004	0.003	0.009	0.031	0.007	0.033

## ■ Remarkable consistency between $n = 1, 2, 3, 4$

## ■ Result: $n=1$ : $m_c(3 \text{ GeV})=0.986(13) \text{ GeV}$

$$m_c(m_c)=1.279(13) \text{ GeV}$$

Theo. uncertainty by truncation error comparable

# III. Analysis

Extraction of the exp. moments from  $R(s)$  (bottom quark case)

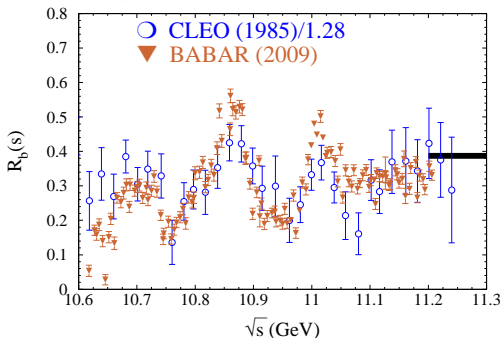
$\mathcal{M}_n^{th}$ : analog to charm case, only  $n_f = 5$

$\mathcal{M}_n^{np}$ : negligible

$\mathcal{M}_n^{res}$ :  $\Upsilon(1S), \Upsilon(2S), \Upsilon(3S), \Upsilon(4S)$  (PDG)

$\mathcal{M}_n^{thr.}$ : BABAR, CLEO data up to 11.24 GeV

Improvements based on recent BABAR results ← new



Systematic  
experimental  
error  $\sim 3.5\%$

$\mathcal{M}_n^{cont}$ : pQCD above 11.24 GeV

# III. Analysis & Results

Determination of the bottom quark mass from  $R(s)$

## ■ Bottom quark masses:

$$\mu = (10 \pm 5) \text{ GeV}; \quad \alpha_s(M_Z) = 0.1189 \pm 0.002$$

$n$	$m_b(10 \text{ GeV})$	exp	$\alpha_s$	$\mu$	total
1	3.597	0.014	0.007	0.002	0.016
2	3.610	0.010	0.012	0.003	0.016
3	3.619	0.008	0.014	0.006	0.018
4	3.631	0.006	0.015	0.020	0.026

## ■ Consistency and stability between $n = 1, 2, 3, 4$

## ■ Result: $n=2$ : $m_b(10 \text{ GeV}) = 3.610(16) \text{ GeV}$

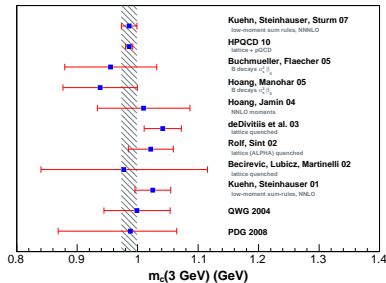
$$m_b(m_b) = 4.163(16) \text{ GeV}$$

Theo. uncertainty by truncation error comparable

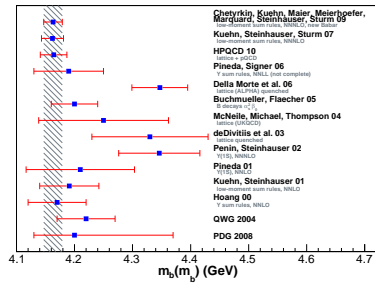
## ■ Well consistent with KSS 2007

# III. Comparison

## charm-quarks



## bottom-quarks



## IV. Summary & Conclusion

- Precise determination of the charm- and bottom-quark mass can be obtained from the experimentally measured  $R$ -ratio in combination with heavy quark current correlators computed in continuum perturbation theory
- Calculation of expansion coefficients of polarization functions up to NNNLO
- Analysis of the  $R$ -ratio and extraction of charm- and bottom-quark masses
- Final results  
quark masses :
  - Charm-mass:  $m_c(3 \text{ GeV}) = 0.986(13) \text{ GeV}$   $e^+e^- + \text{pQCD}$
  - Bottom-mass:  $m_b(10 \text{ GeV}) = 3.610(16) \text{ GeV}$   $e^+e^- + \text{pQCD}$