Using Neutron Stars to Understand Fundamental Physics

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arXiv: 2205.10283, 2302.02989, and 2304.07836

Where are we going? Part 1 of 4

1. Compact stars and the phase diagram of nature

2. New models for neutron star phenomenology

- 3. Studying the **lightest** observed compact star
- 4. Some possible directions of **future work**

Where are we going? Part 1 of 4

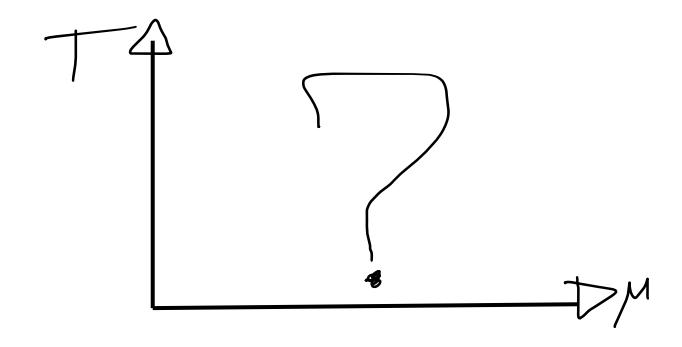
1. Compact stars and the **phase diagram** of nature

2. New models for neutron star phenomenology

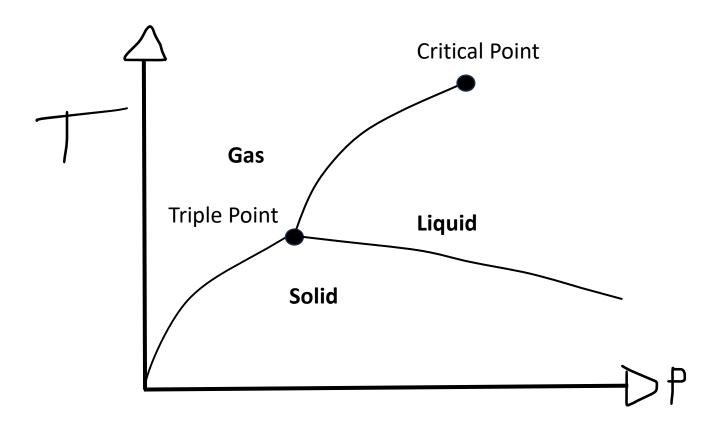
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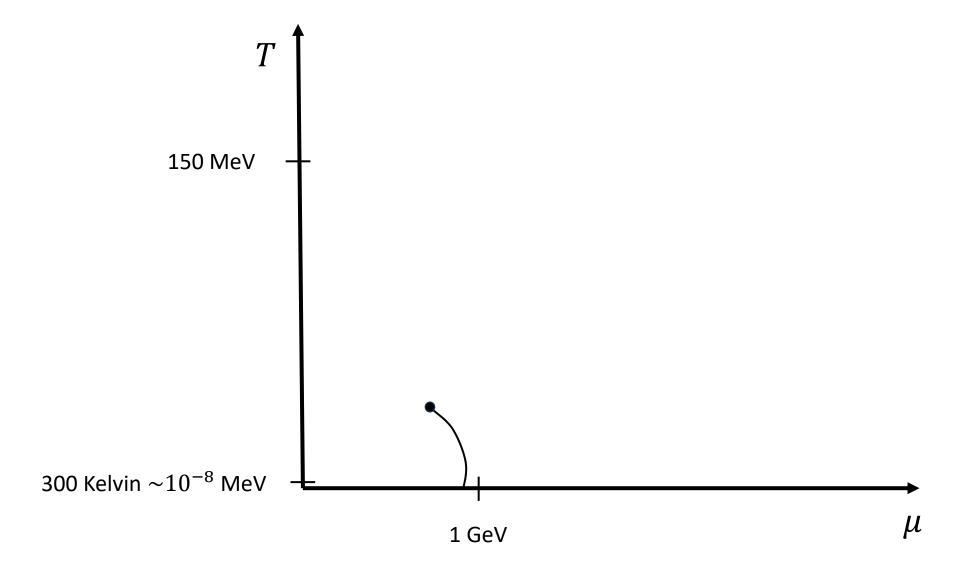
Fundamental Forces

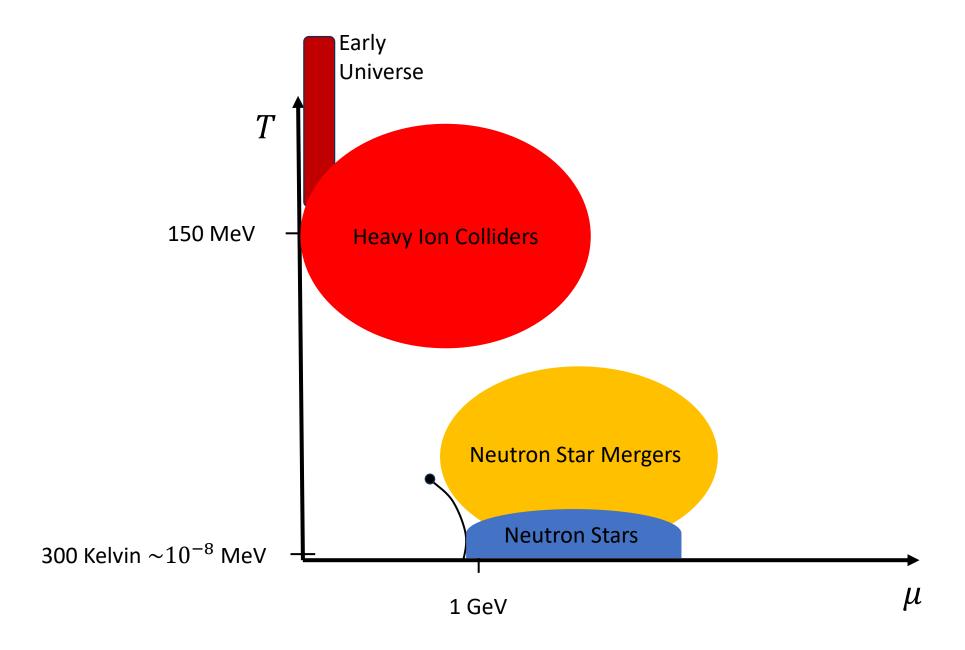
Gravity Electroweak Strong

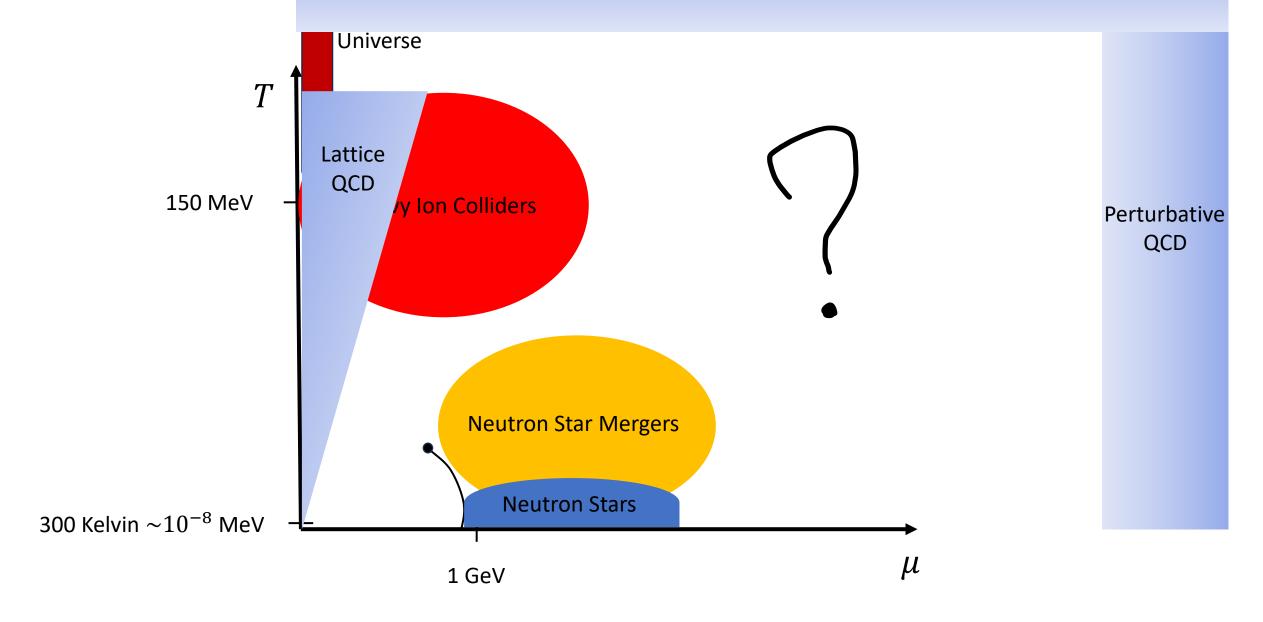


Earthly Matter: Water









Why else are neutron stars interesting?

All fundamental forces are relevant!

Strong

Stabilizes star against gravity

Gravity

Second most **dense** object in the universe ($^{\sim}10^{15}$ g/cm³)

Gravitational waves from merger

Electromagnetism

Strongest **magnetic fields** in the universe

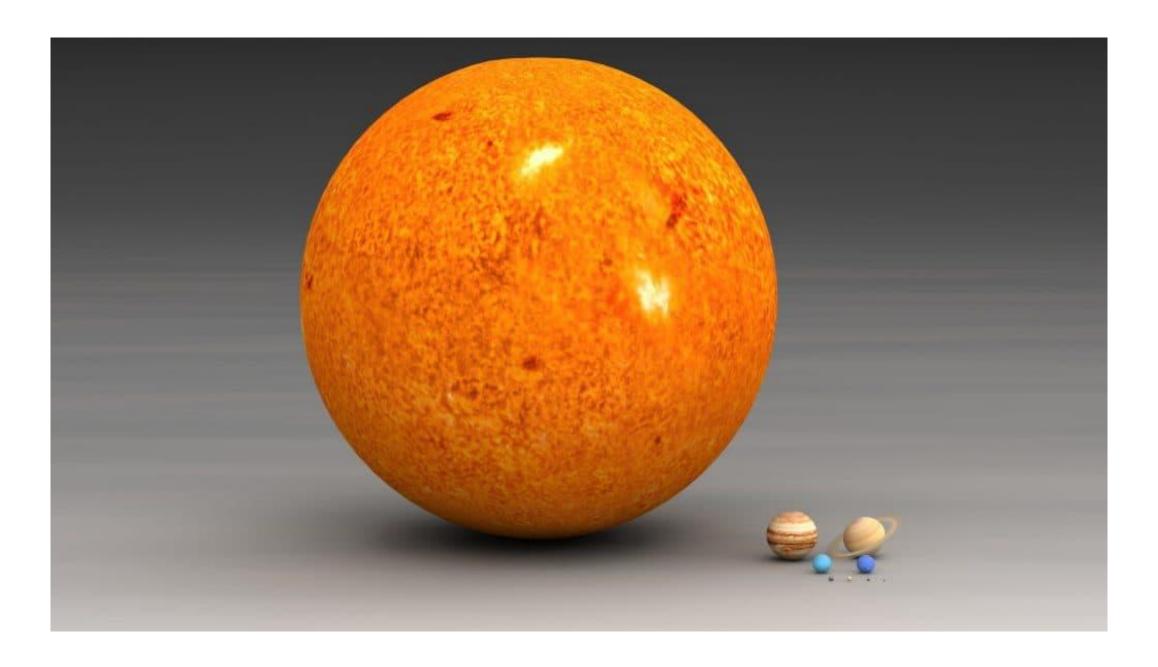
Charge **neutrality**

Weak

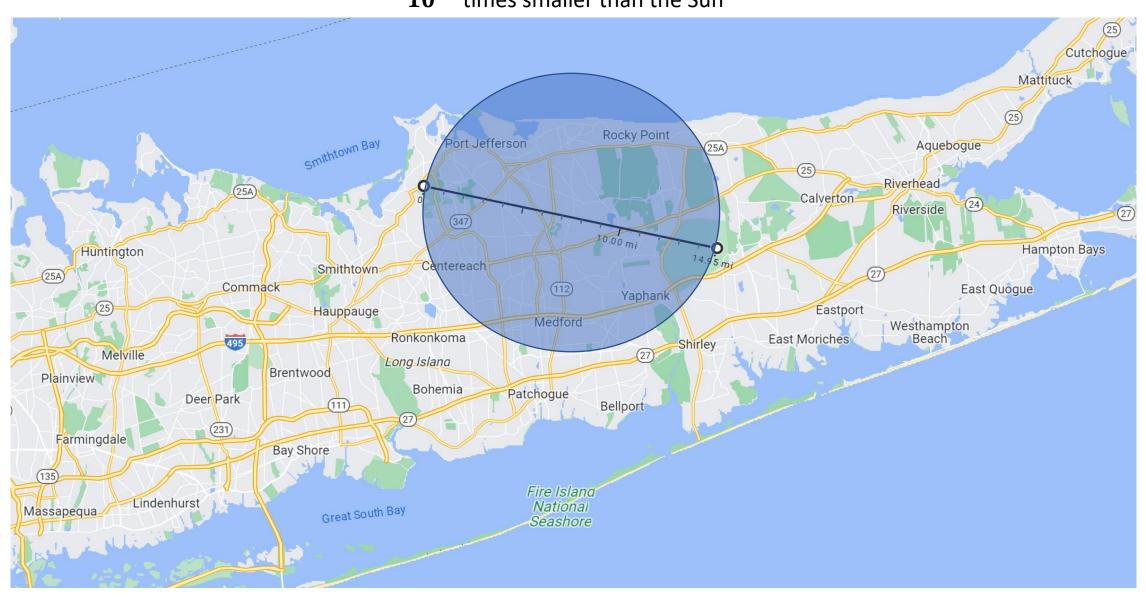
Isospin equilibration

→ Bulk viscosity

Nucleosynthesis



10¹⁴ times smaller than the Sun



Conclusions: Part 1 of 4

- Neutron stars are...
 - **extremely** compact
 - used as a "laboratory" to explore some region of the **phase diagram** of matter
 - cool and dense (μ/T large)

QCD cannot be solved exactly → phenomenological models needed

Where are we going? Part 2 of 4

1. Compact stars and the phase diagram of nature

2. New models for neutron star **phenomenology**

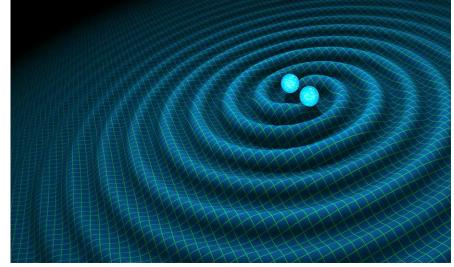
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Why develop new phenomenological models?

Neutron star mergers probe regions of neutron-rich matter

 Most equations of state are calibrated to isospin-symmetric matter (50% neutrons) and extrapolated to ~90% neutron matter



R. HURT/CALTECH-JPL

Let's calibrate models to neutron-rich matter!

Summary: Part 2 of 4

- Want better microscopic models of nuclear matter that are constrained by:
 - Properties of neutron matter (how?)
 - Saturation properties of **isospin-symmetric** nuclear matter
 - Observations of neutron star structure

Relativistic Mean-Field Theory (RMFT)

Advantages	Limitations
Relativistic theory (e.g., $v_{sound} < c$) \odot	Not a controlled approximation 🕾
Tractable calculations ©	Reasonable to about 6 times nuclear saturation density (~1fm ⁻³) ⊕ ⊖
Finite temperature and out of equilibrium physics included ©	No phase transition to deconfined quarks 😊
Microscopic information available ©	Coupling constants need to be fit to something ⊕⊛

$$\mathcal{L} = \mathcal{L}_{ ext{N}} + \mathcal{L}_{ ext{M}} + \mathcal{L}_{l}$$

$$\mathcal{L}_{N} = \bar{\psi} \left[i \gamma^{\mu} \partial_{\mu} - m_{N} + g_{\sigma} \langle \sigma \rangle - g_{\omega} \langle \omega_{0} \rangle - \frac{g_{\rho}}{2} \tau_{3} \langle \rho_{03} \rangle \right] \psi$$

$$\mathcal{L}_{M} = -\frac{1}{2}m_{\sigma}^{2}\langle\sigma\rangle^{2} - b\frac{M}{3}(g_{\sigma}\langle\sigma\rangle)^{3} - \frac{c}{4}(g_{\sigma}\langle\sigma\rangle)^{4} + \frac{1}{2}m_{\omega}^{2}\langle\omega_{0}\rangle^{2} + \frac{1}{2}m_{\rho}^{2}\langle\rho_{03}\rangle^{2} + \Lambda(g_{\rho}\langle\omega_{0}\rangle\langle\rho_{03}\rangle)^{2}$$

$$\mathcal{L}_l = \bar{\psi}_e \left(i \gamma^\mu \partial_\mu - m_e \right) \psi_e$$

$$\mathcal{L} = \mathcal{L}_{ ext{N}} + \mathcal{L}_{ ext{M}} + \mathcal{L}_{l}$$

$$\mathcal{L}_{N} = \bar{\psi} \left[i \gamma^{\mu} \partial_{\mu} - m_{N} + g_{\sigma} \langle \sigma \rangle - g_{\omega} \langle \omega_{0} \rangle - \frac{g_{\rho}}{2} \tau_{3} \langle \rho_{03} \rangle \right] \psi$$

$$\mathcal{L}_{M} = -\frac{1}{2}m_{\sigma}^{2}\langle\sigma\rangle^{2} - \frac{b}{3}\frac{M}{3}(g_{\sigma}\langle\sigma\rangle)^{3} - \frac{c}{4}(g_{\sigma}\langle\sigma\rangle)^{4} + \frac{1}{2}m_{\omega}^{2}\langle\omega_{0}\rangle^{2} + \frac{1}{2}m_{\rho}^{2}\langle\rho_{03}\rangle^{2} + \Lambda(g_{\rho}\langle\omega_{0}\rangle\langle\rho_{03}\rangle)^{2}$$

$$\mathcal{L}_l = \bar{\psi}_e \left(i \gamma^\mu \partial_\mu - m_e \right) \psi_e$$

Nuclear Physics Data to Constrain Our Model

Fit relativistic mean-field theory to symmetric and neutron matter



- Isospin-Symmetric Matter
 - Binding energy
 - Pressure
 - Incompressibility

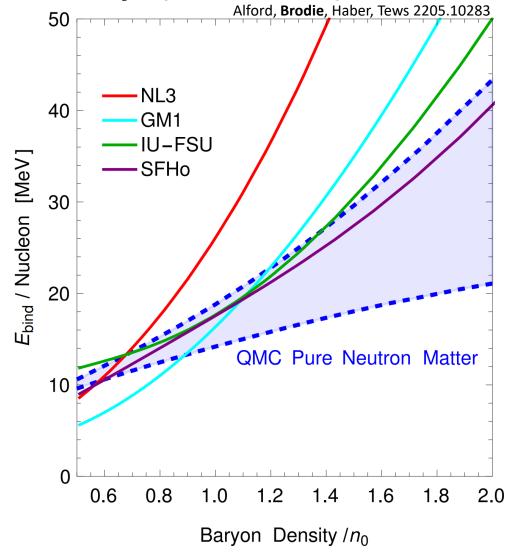


- Neutron Matter
 - → Chiral Effective Field Theory

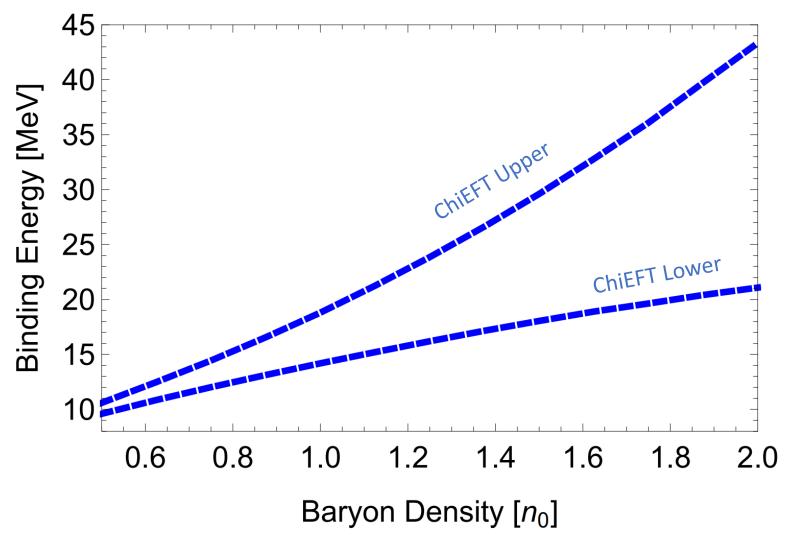
Chiral Effective Field Theory (ChiEFT)

- Based on the <u>symmetries</u> of QCD with nucleon and pion degrees of freedom
- <u>Controlled</u> approximation to QCD valid at low densities
- Theory fitted to data from scattering experiments

Commonly used relativistic mean-field theories are **inconsistent** with ChiEFT for neutron matter

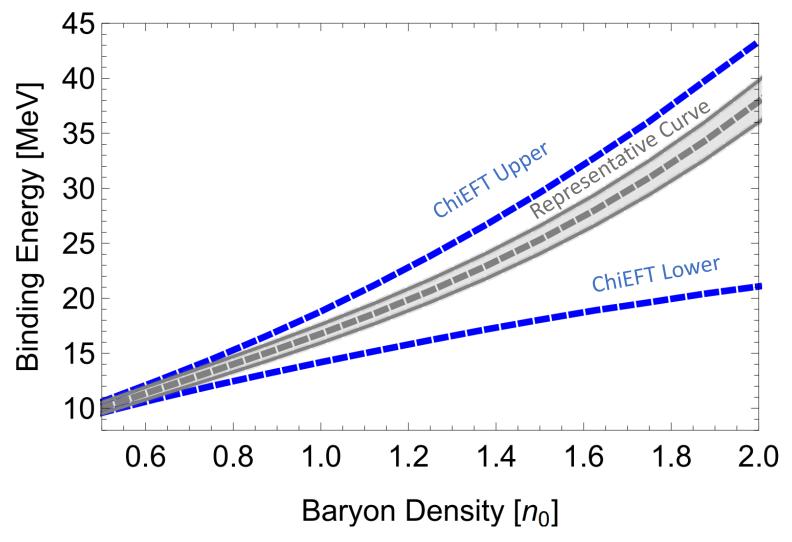


Uncertainty in Chiral Effective Field Theory



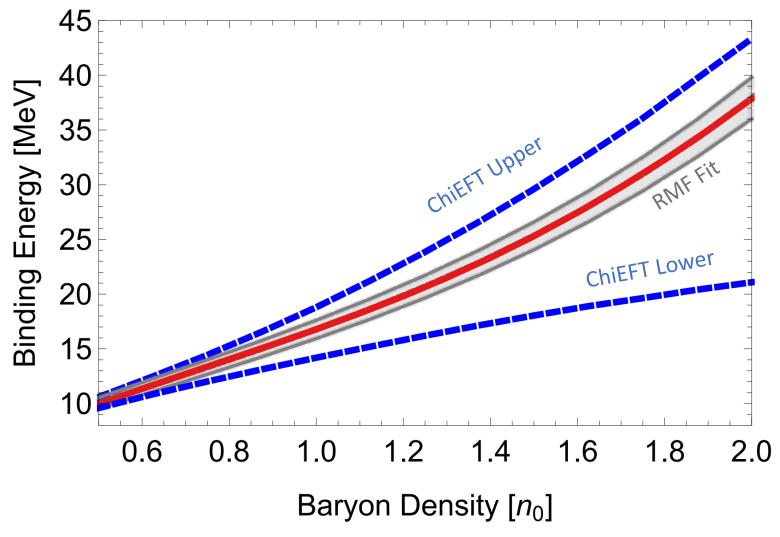
ChiEFT data from Tews et al. (arXiv:1801.01923)

Uncertainty in Chiral Effective Field Theory



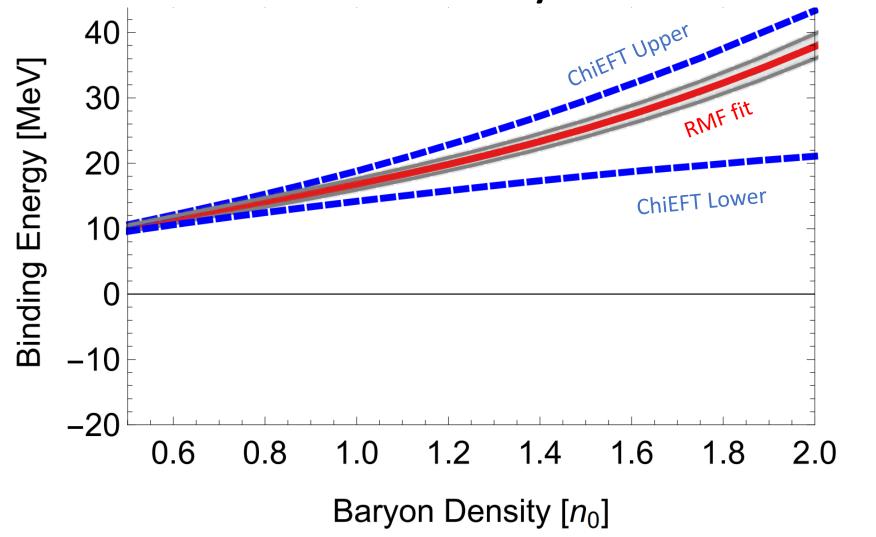
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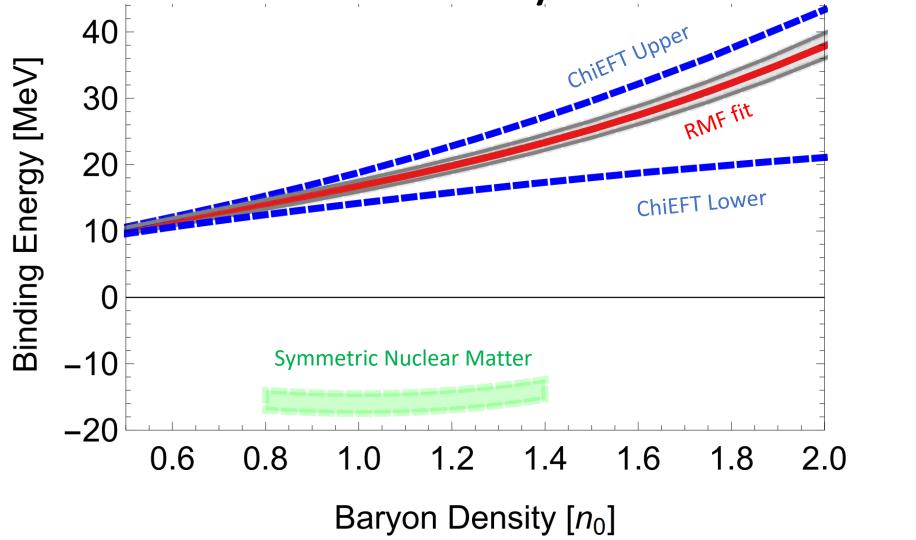


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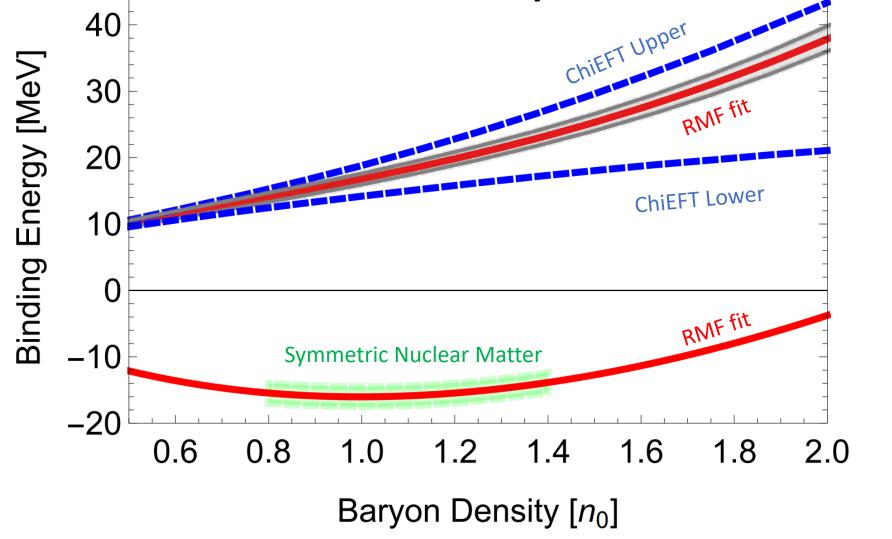
Simultaneous Nuclear Physics Constraints



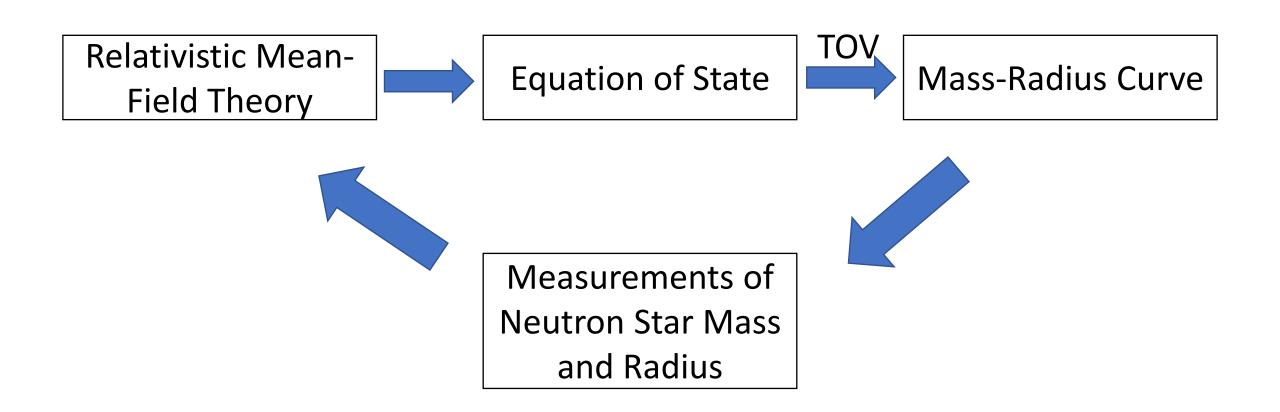
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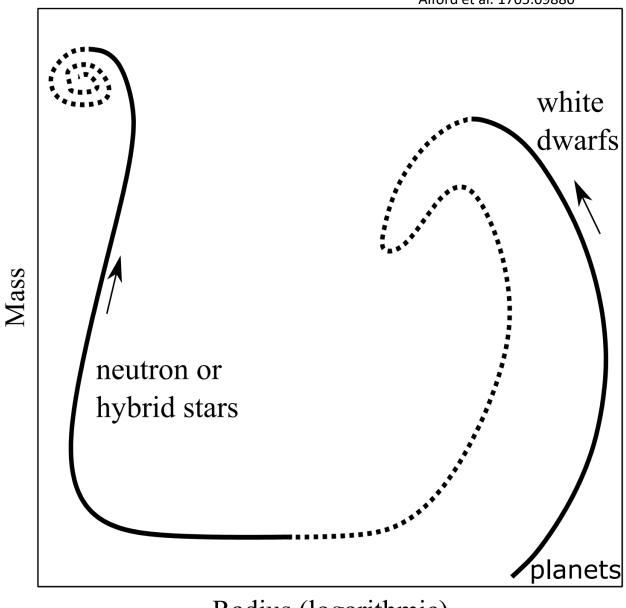


Simultaneous Nuclear Physics Constraints

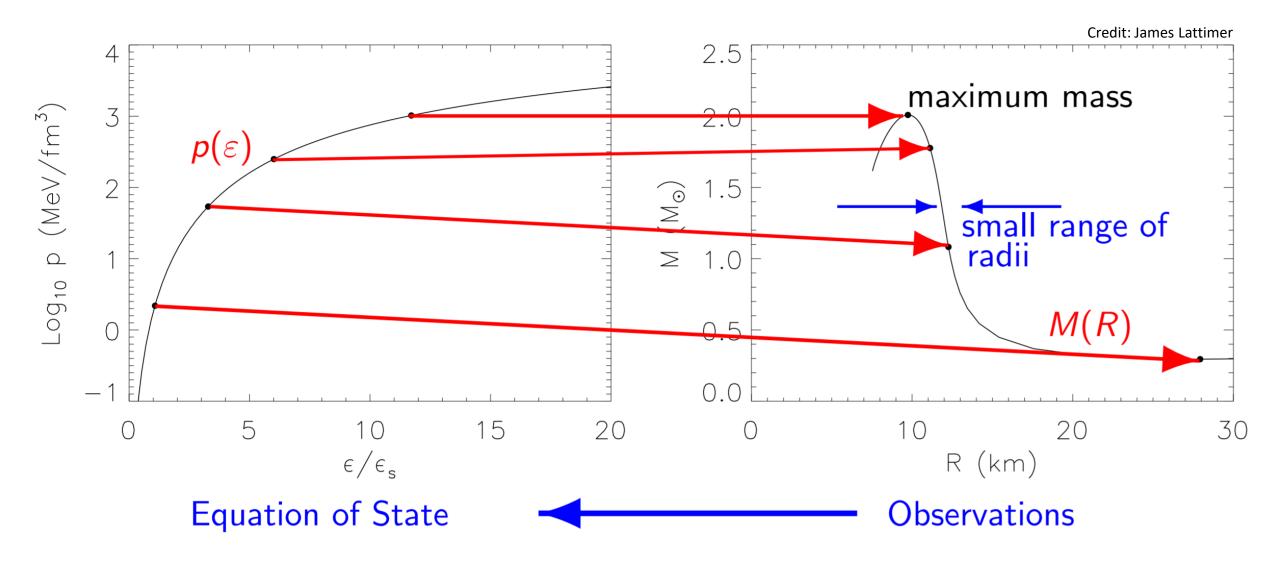


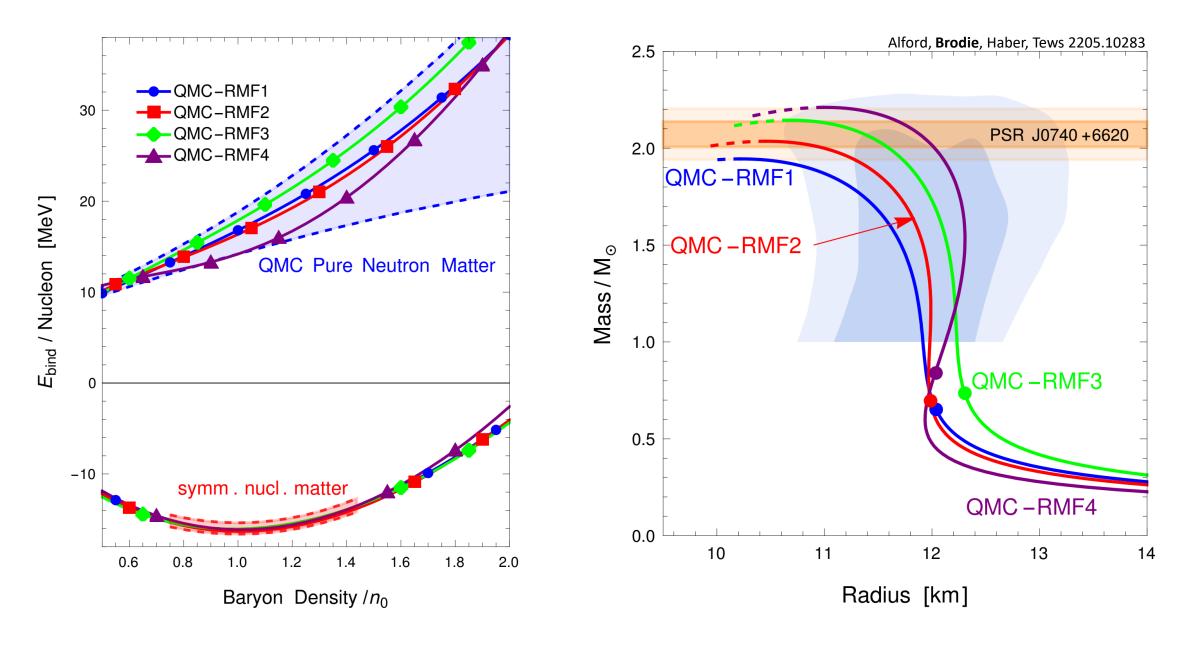
How can neutron stars constrain our model?





Radius (logarithmic)





Conclusions: Part 2 of 4

Developed **QMC-RMF1,2,3,4**: models of nuclear matter that are constrained by <u>isospin-symmetric</u> and <u>neutron</u> matter and by observations of <u>neutron star</u> structure

Ready for use in neutron star merger simulations: compose.obspm.fr/eos/297

- Tabulated over a range of <u>temperatures</u>, <u>densities</u>, and <u>proton fractions</u>
- Provides an <u>equation of state</u> (p,ε,s,μ,...)
- Provides <u>particle dispersion relations</u> and <u>effective masses</u>

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nature astronomy

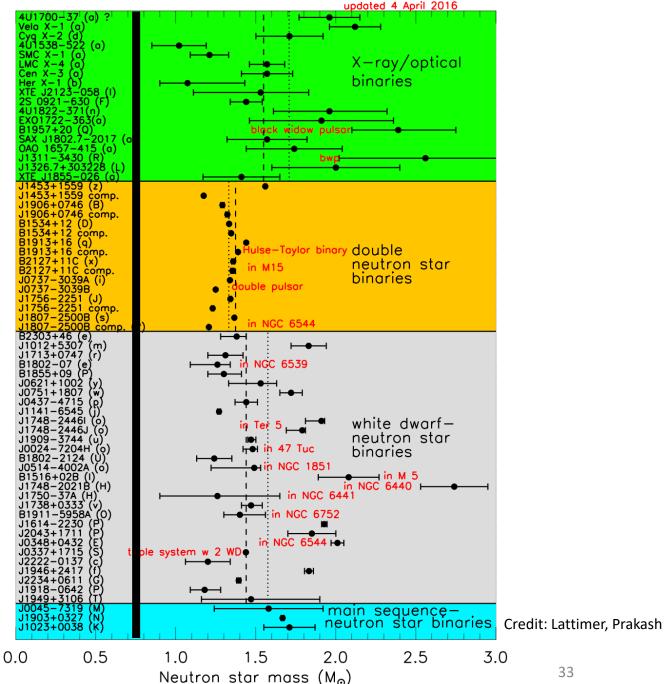
Article

Doroshenko et al., 2022

A strangely light neutron star within a supernova remnant

$$M = 0.77^{+0.20}_{-0.17} \,\mathrm{M}_{\odot}$$

$$R = 10.4^{+0.86}_{-0.78} \text{ km}$$

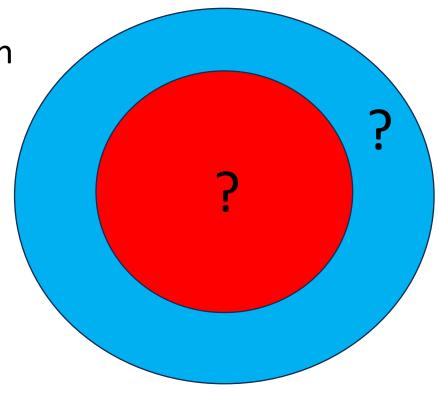


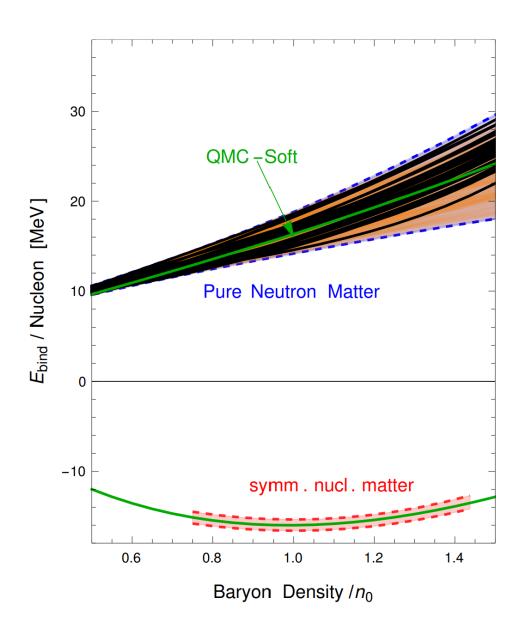
Studying the Lightest Compact Star

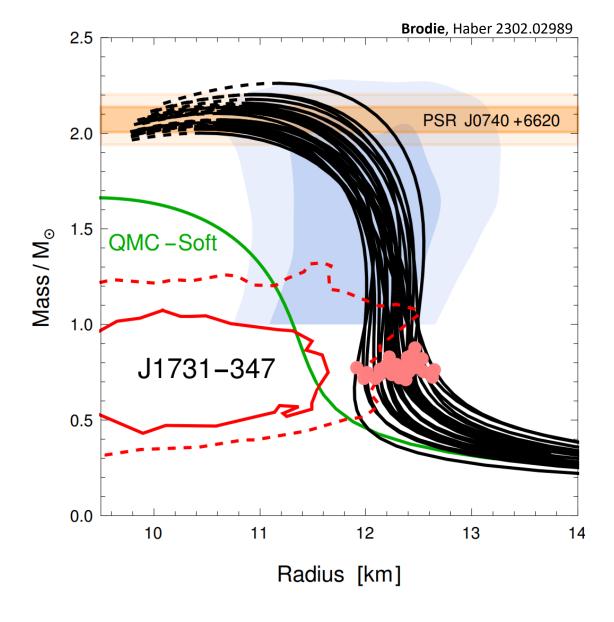
$$M = 0.77^{+0.20}_{-0.17} \,\mathrm{M}_{\odot} \qquad R = 10.4^{+0.86}_{-0.78} \,\mathrm{km}$$

Caution: Parameter estimation & formation mechanism

- What is the **composition** of this star?
 - Quark star
 - Neutron star
 - Hybrid star
 - Etc...

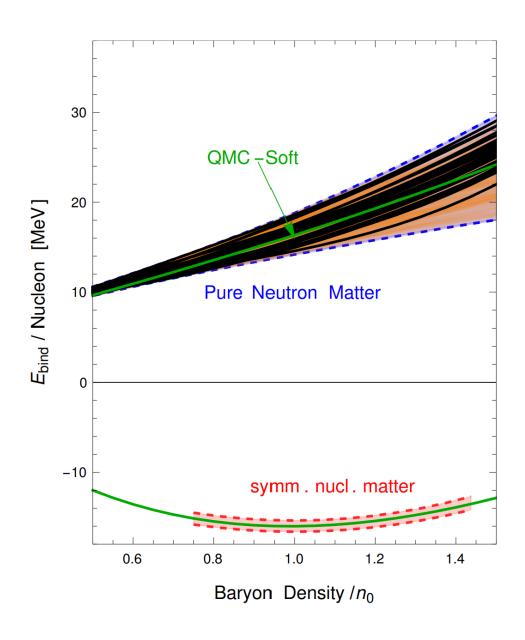


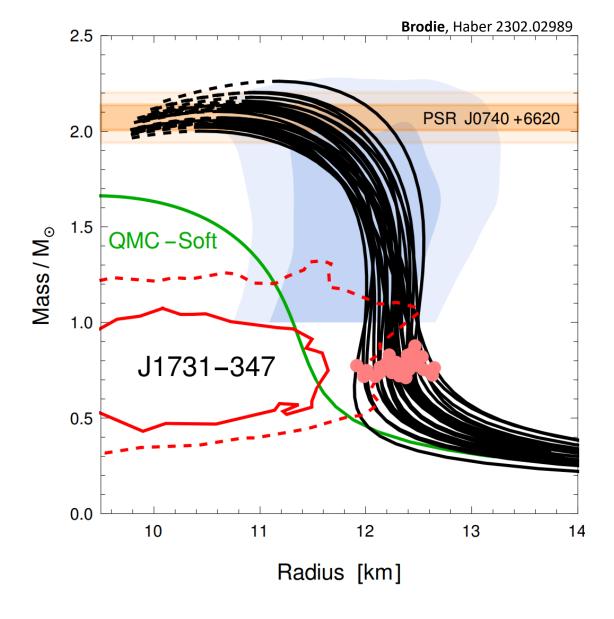




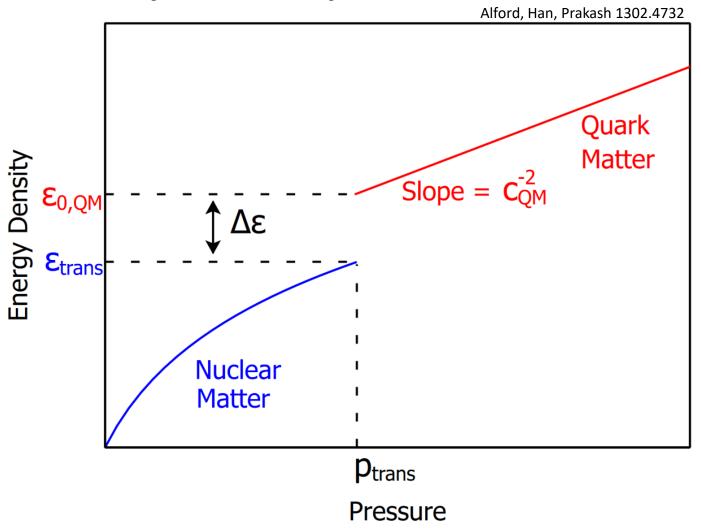
Nuclear models only reach 2σ consistency with the HESS J1731-347 measurement

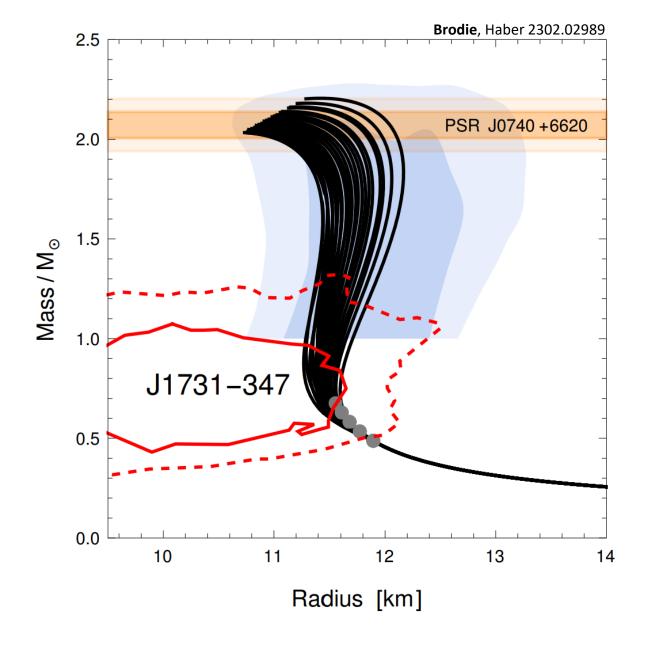
- Possibly because:
 - Measurement uncertainties underestimated
 - Additional nuclear physics not included in the family of RMFs used
 - Nucleons are not the only relevant degree of freedom





Hybrid Equation of State





$$c_{QM}^{2} = 0.48$$

$$n_{tr} \in \{2n_{0}, 2.4n_{0}\}$$

$$\frac{\Delta \epsilon}{\epsilon} \in \{0.004, 0.151\}$$

Conclusions: Part 3 of 4

If HESS J1731-347 measurement is credible

 Most compact stars could be hybrid stars with a nuclear matter "mantle" and a quark matter "core"

Nuclear models do not have to be as "stiff"

• Future observations could help constrain low-energy nuclear theory

Parts 1-3 Conclusions

- Neutron stars can be used as a "laboratory" to explore part of the phase diagram of matter
- Developed new models for neutron stars constrained by nuclear physics and astrophysics
- Most compact stars could be hybrid stars with a nuclear matter "mantle" and a quark matter "core"

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Future work

- Weak interaction transport properties
 - Urca processes with arbitrary neutrino distribution
 - Strangeness equilibration in quark matter at merger temperatures

- Interested in the rest of the phase diagram of matter
 - Is there is a critical endpoint?
 - Is chiral symmetry restored at deconfinement?
 - Other phases of matter?