

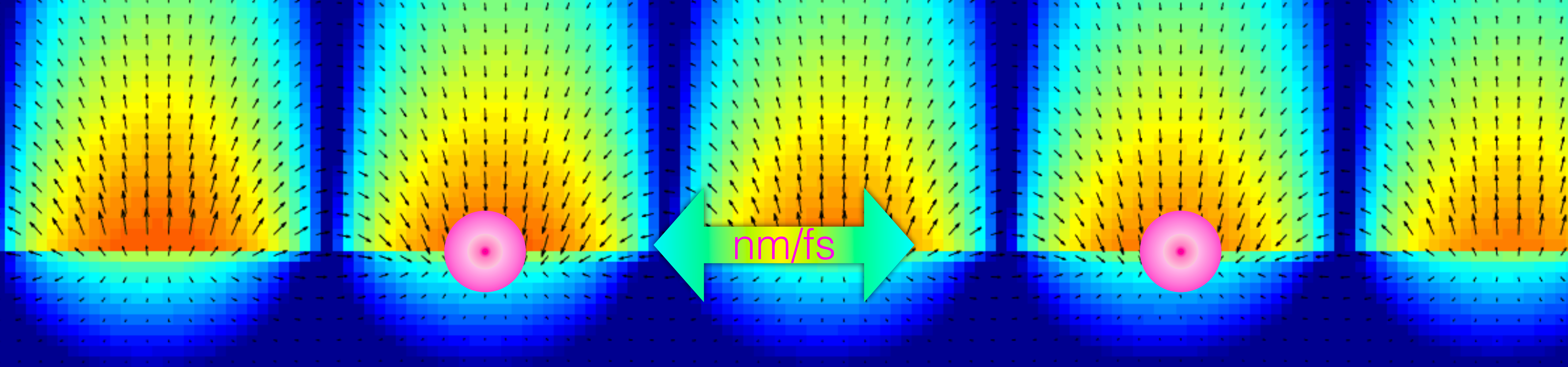
Plasmonic vortices as sources of angular momentum electron beams - and other good things



P3

Oct. 5, 2023

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Floquet engineering of nonlinear photoemission fields.

Vectorial electromagnetic fields beyond Maxwell?

Ultrafast photoelectron emission microscopy - how to see light.

Nanofemto electron imaging the spin-orbit interaction of light.

Vortices - topology, parity-time symmetry, and arrays.

Magnetolectric ($E \cdot B$) interaction.

Floquet engineering of metal surface states

$$\Psi_\alpha(x, t) = \exp\left(-\frac{i\epsilon_\alpha t}{\hbar}\right) \Phi_\alpha(x, t)$$

where $\Phi_\alpha(x, t) = \Phi_\alpha(x, t + T)$

$$\hbar\omega_L = \epsilon_\alpha$$

Floquet ladder

image potential states
do not preexist

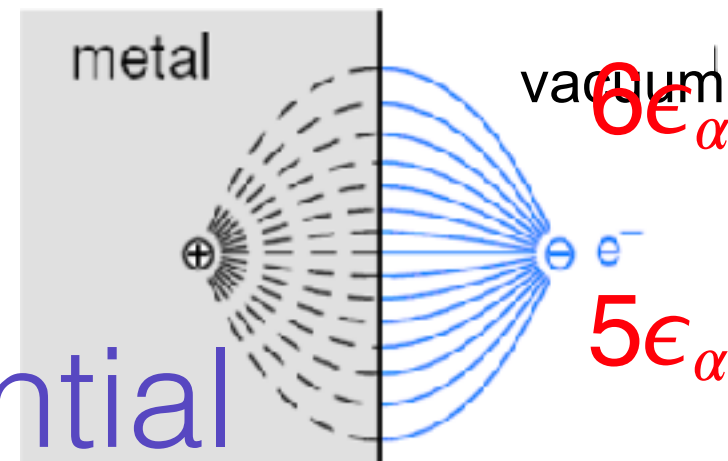
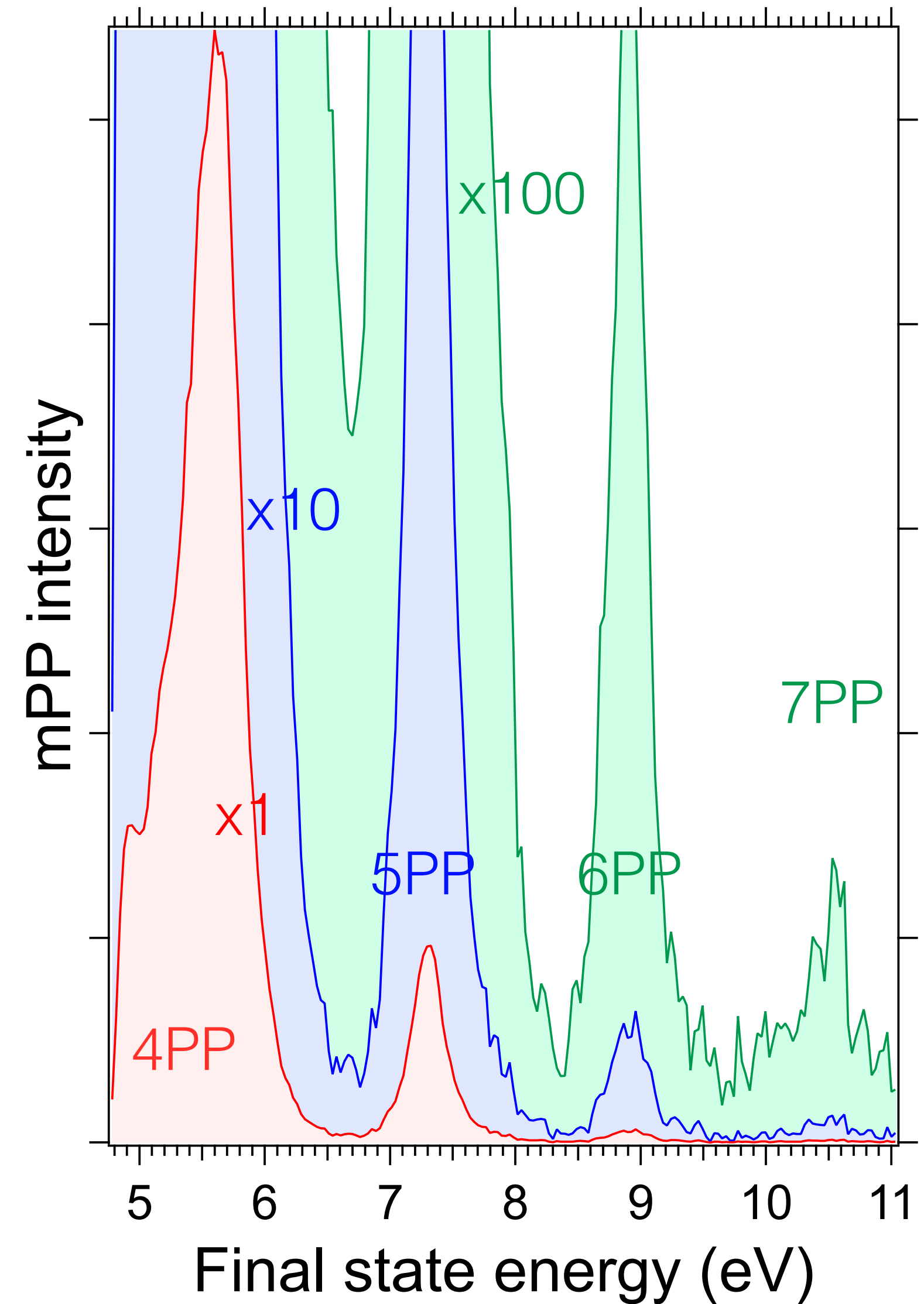
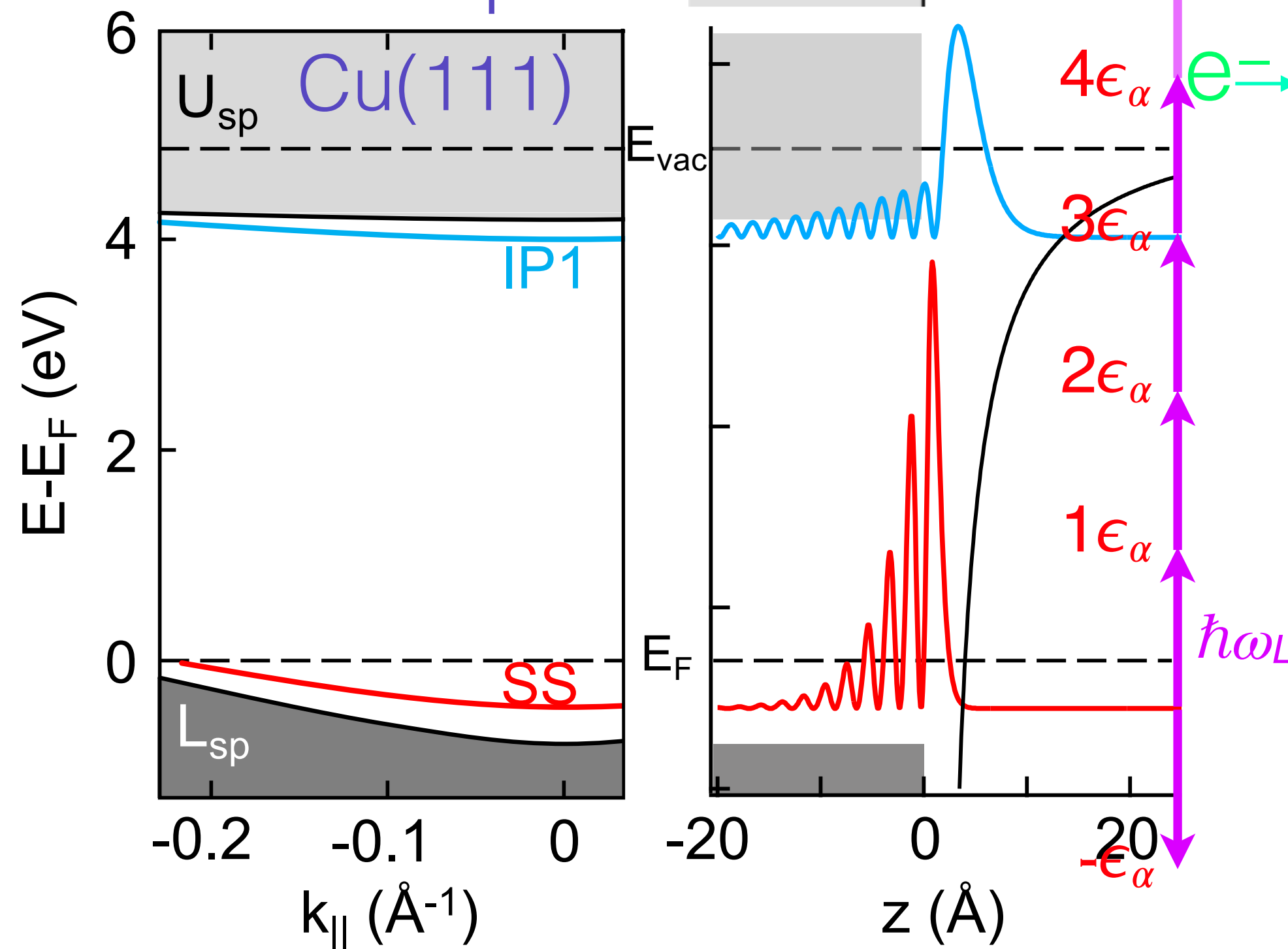
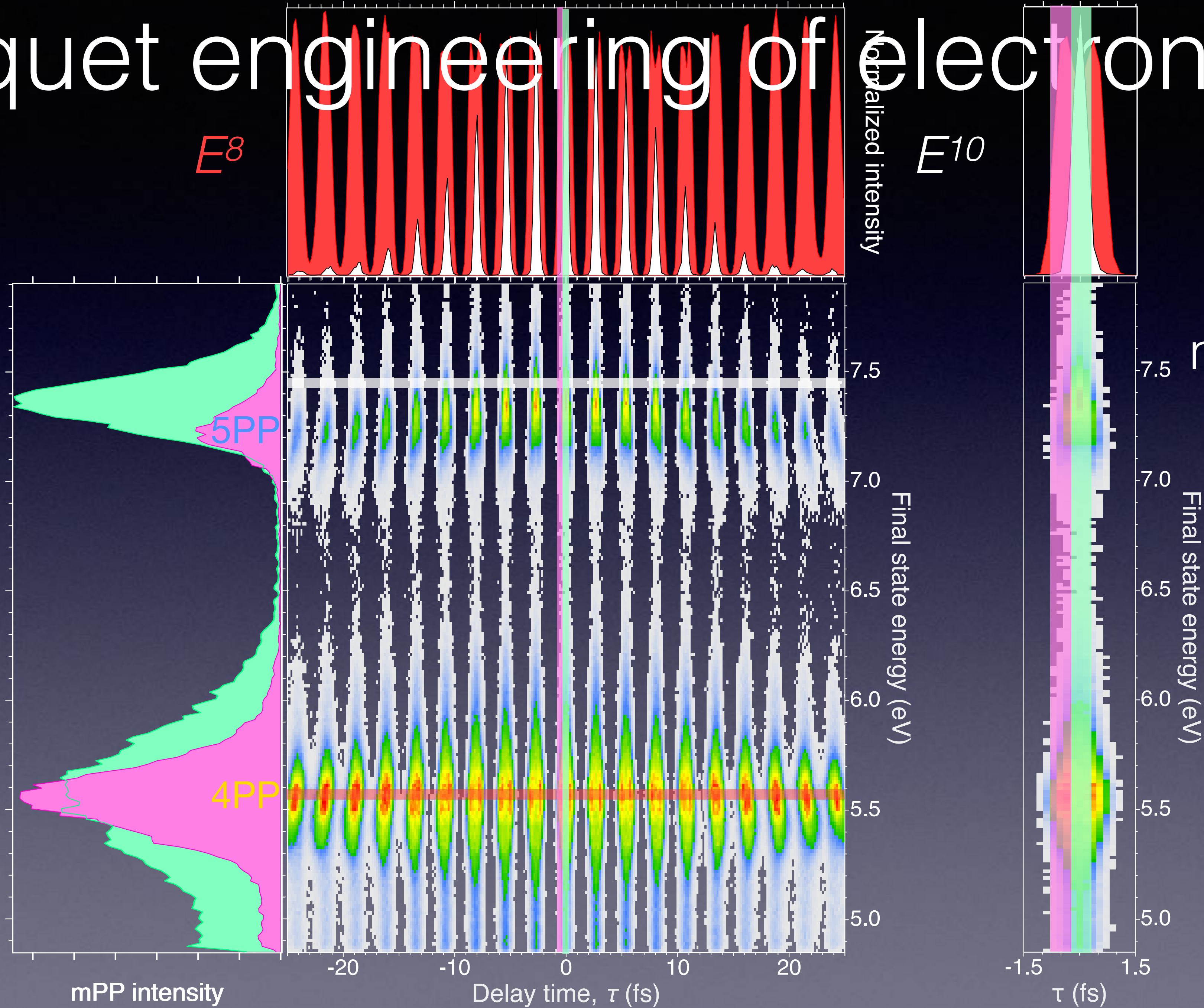


Image potential
lattice potential



Floquet engineering of electronic bands



How fast does matter respond?

Magnetolectric effect of pseudoscalar $E \cdot B$

- Topological field theory relates $E \cdot B \neq 0$ to magnetic monopoles and axion quasiparticles
- under this condition H induces polarization P , and D magnetization M
- $E \cdot B \neq 0$ is a material property when \mathcal{P} and \mathcal{T} are anti-symmetric, but \mathcal{PT} symmetry is preserved

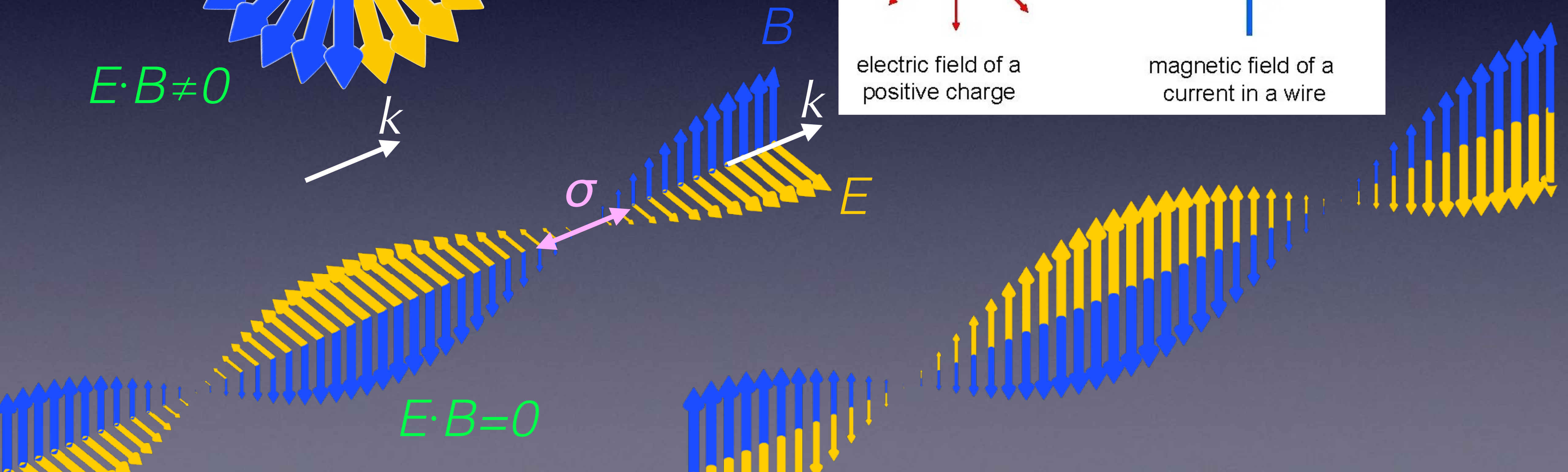
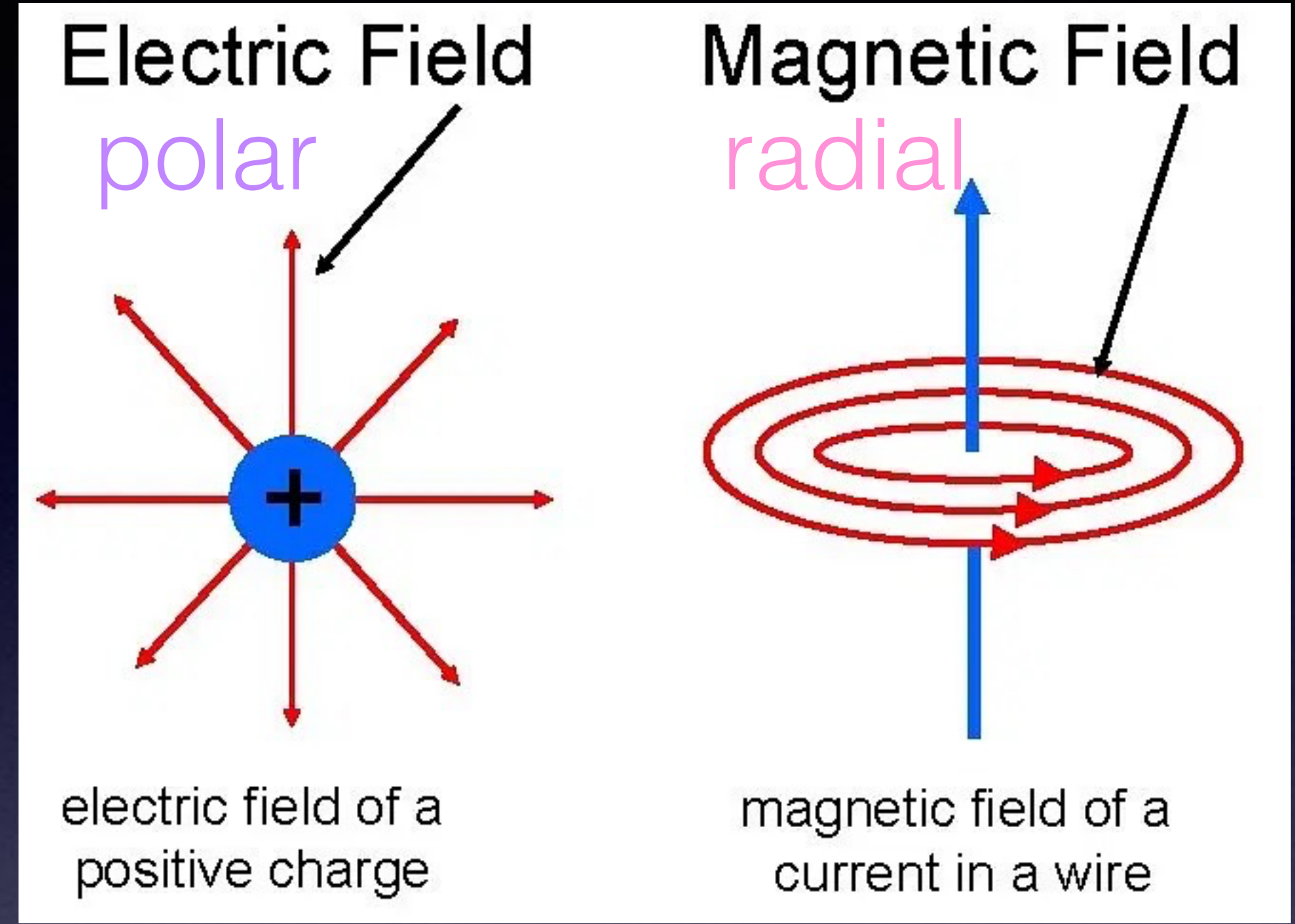
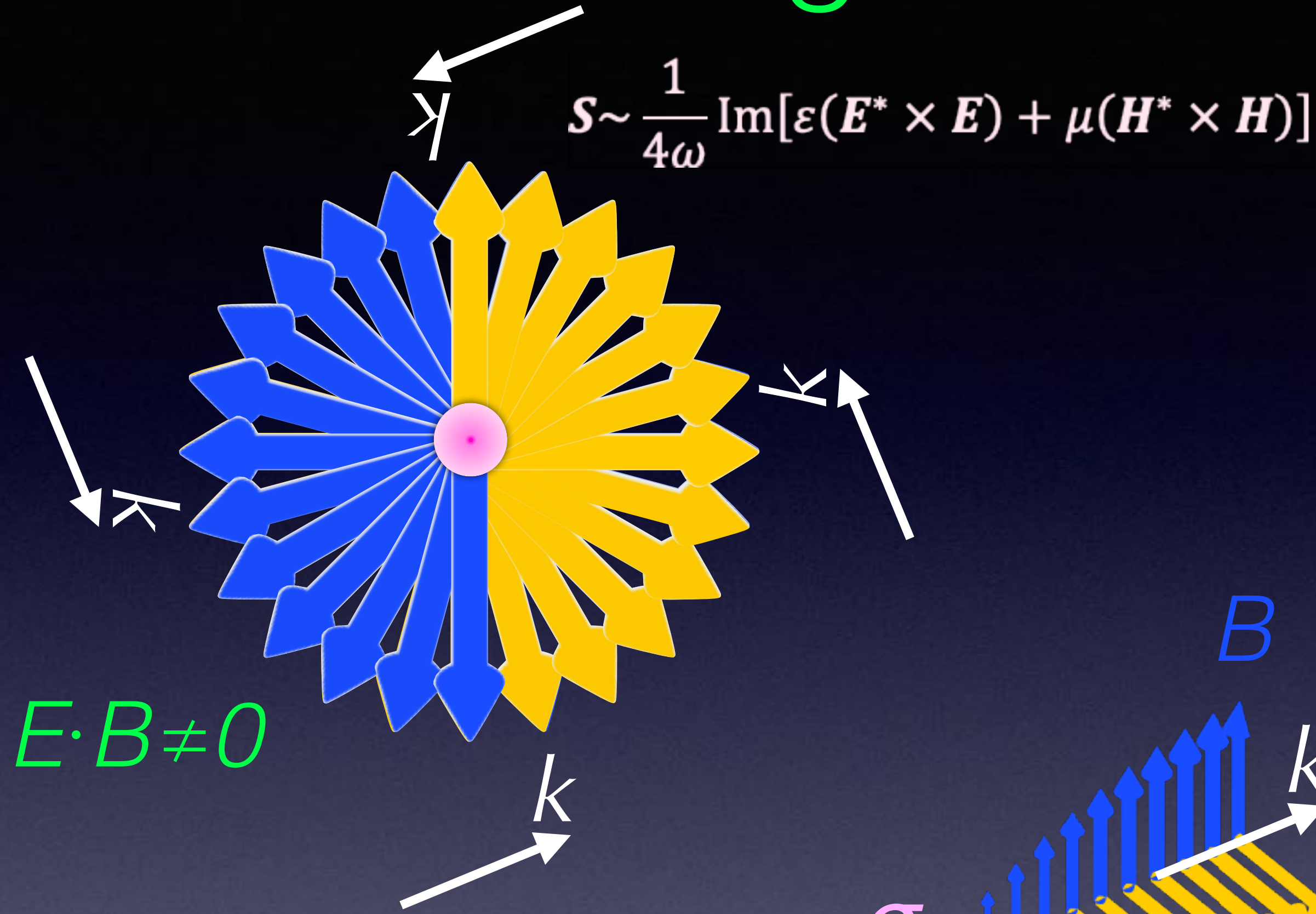
Maxwell's equations when $E \cdot B \neq 0$

Name	Equations
Gauss's law	$\nabla \cdot (\mathbf{E} - \kappa \theta c\mathbf{B}) = \frac{\rho_e}{\epsilon_0}$
Gauss's law for magnetism	$\nabla \cdot (c\mathbf{B} + \kappa \theta \mathbf{E}) = c\mu_0 \rho_m$
Faraday's law	$\nabla \times (-\mathbf{E} + \kappa \theta c\mathbf{B}) = \partial_{ct}(c\mathbf{B} + \kappa \theta \mathbf{E}) + \mu_0 \mathbf{J}_m$
Ampère–Maxwell law	$\nabla \times (c\mathbf{B} + \kappa \theta \mathbf{E}) = \partial_{ct}(\mathbf{E} - \kappa \theta c\mathbf{B}) + c\mu_0 \mathbf{J}_e$
Axion law	$(\square + m_A^2) \theta = -\kappa \mathbf{E} \cdot \mathbf{B}$

Magnetolectric effect

$$\begin{bmatrix} \mathbf{E}' \\ c\mathbf{B}' \end{bmatrix} = \begin{bmatrix} \mathbf{E} - \kappa \theta c\mathbf{B} \\ \kappa \theta \mathbf{E} + c\mathbf{B} \end{bmatrix} = \frac{1}{\cos \xi} \begin{bmatrix} \cos \xi & \sin \xi \\ -\sin \xi & \cos \xi \end{bmatrix} \begin{bmatrix} \mathbf{E} \\ c\mathbf{B} \end{bmatrix}$$

Coherent magnetolectric $E \cdot B$ interaction



Vectorial properties of optical vortex light

Spin angular momentum (SAM) rotating polarization

Orbital angular momentum (OAM) rotating wavefront

Maxwell's equations very

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad \text{versatile, but ..}$$

$$\nabla \cdot \mathbf{B} = 0$$

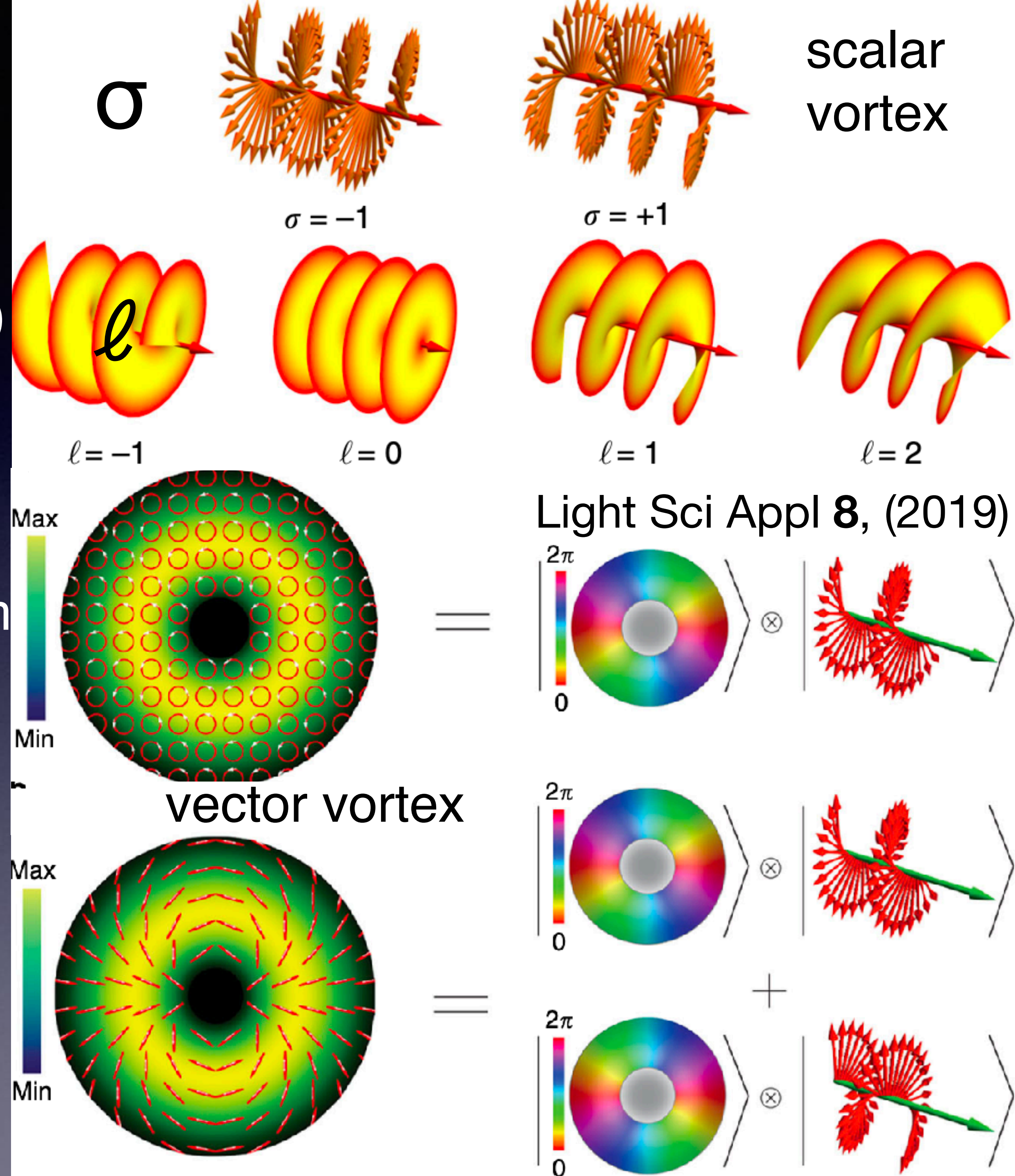
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \left(\mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right)$$

Conservation of angular momentum $\mathbf{J}=\mathbf{L}+\mathbf{S}$

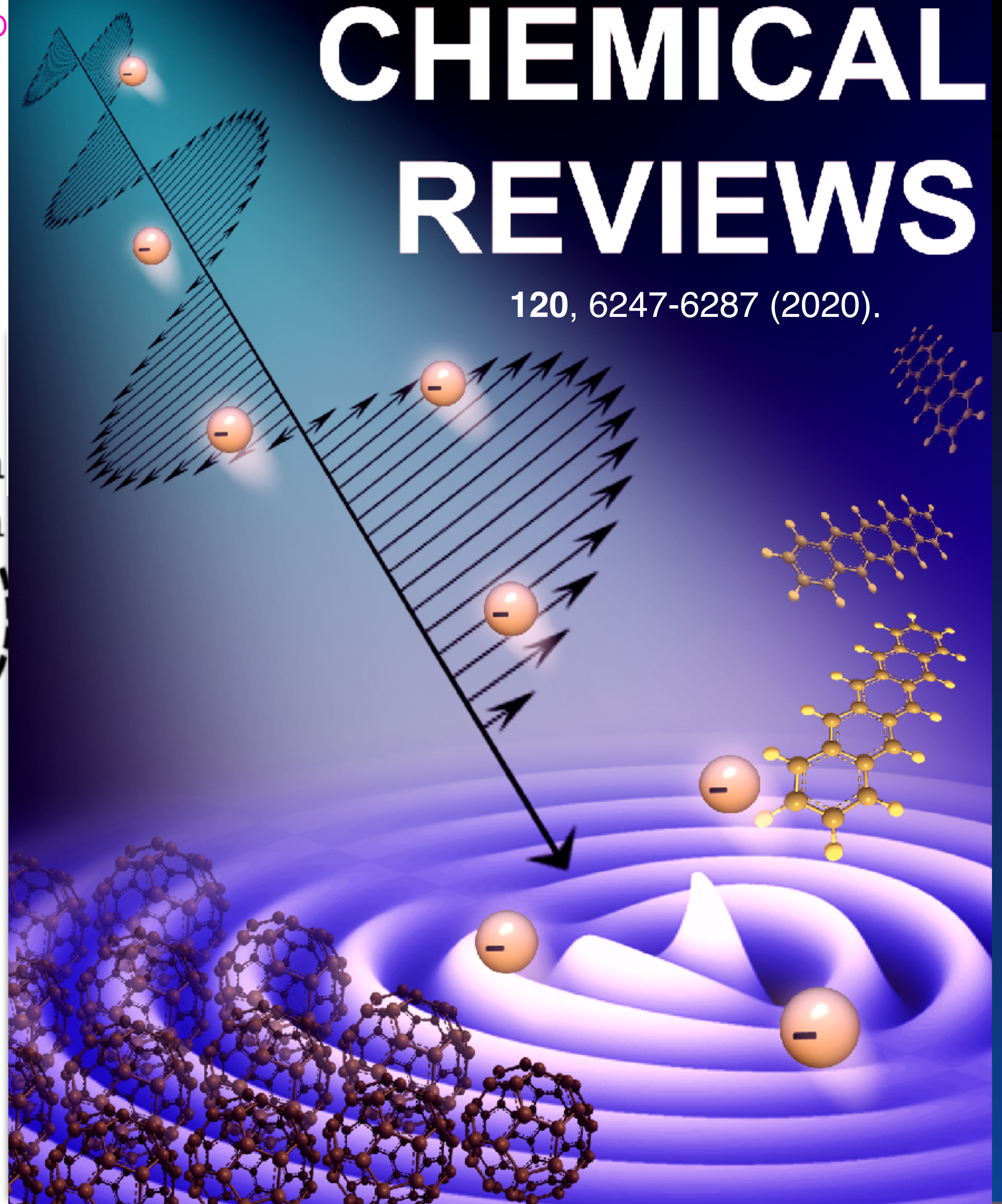
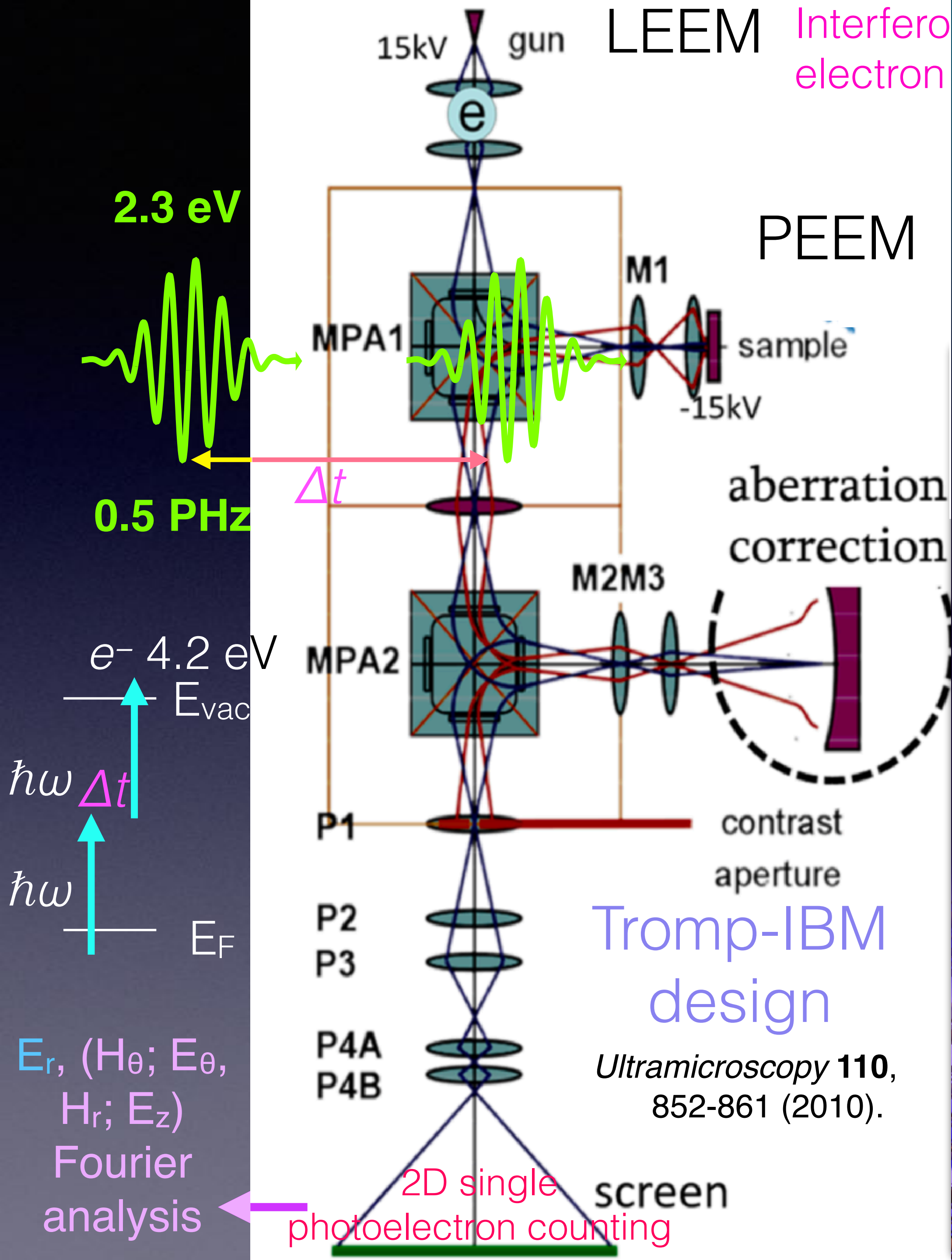
Spin-orbit interaction of light

Degrees of freedom of light: Amplitude, phase, polarization, spin and orbital angular momentum with nonseparable phase and polarization wavefronts.



CHEMICAL REVIEWS

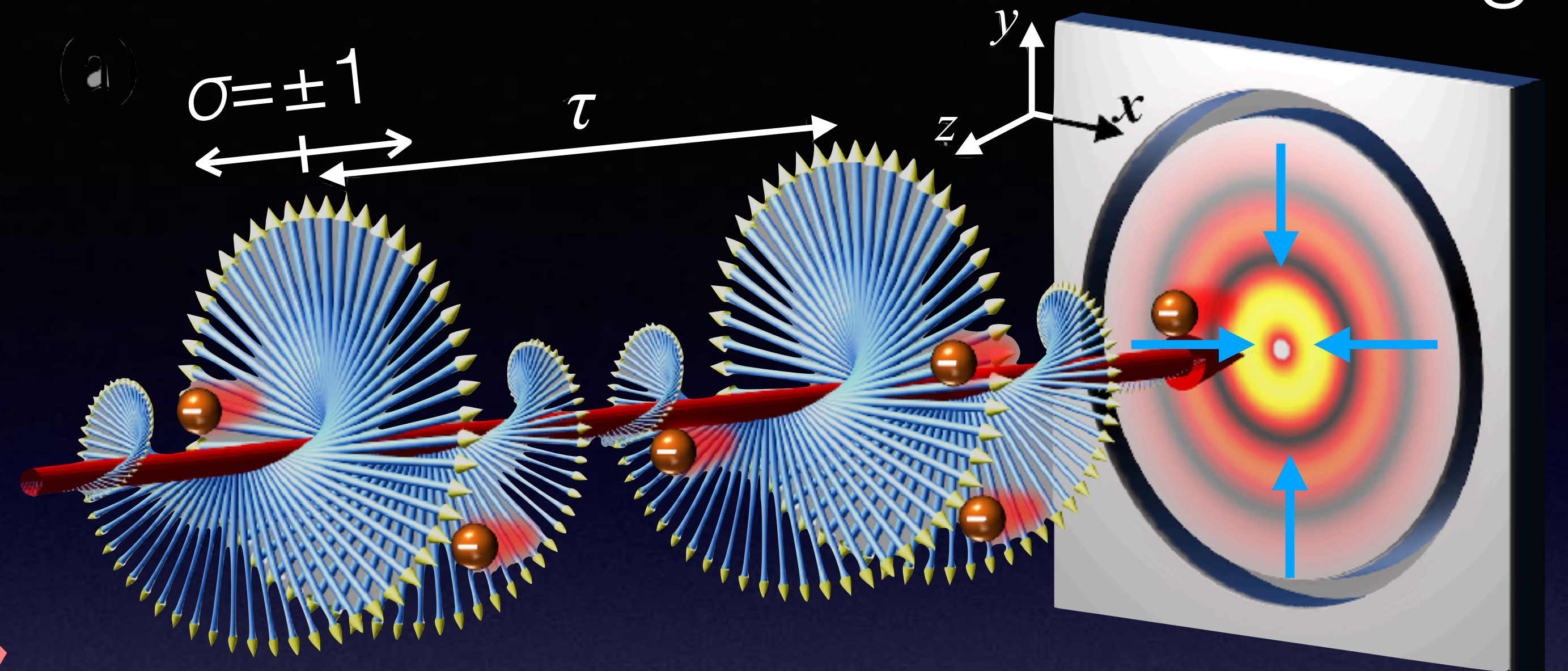
120, 6247-6287 (2020).



Plasmonic vortices-SO interaction of light

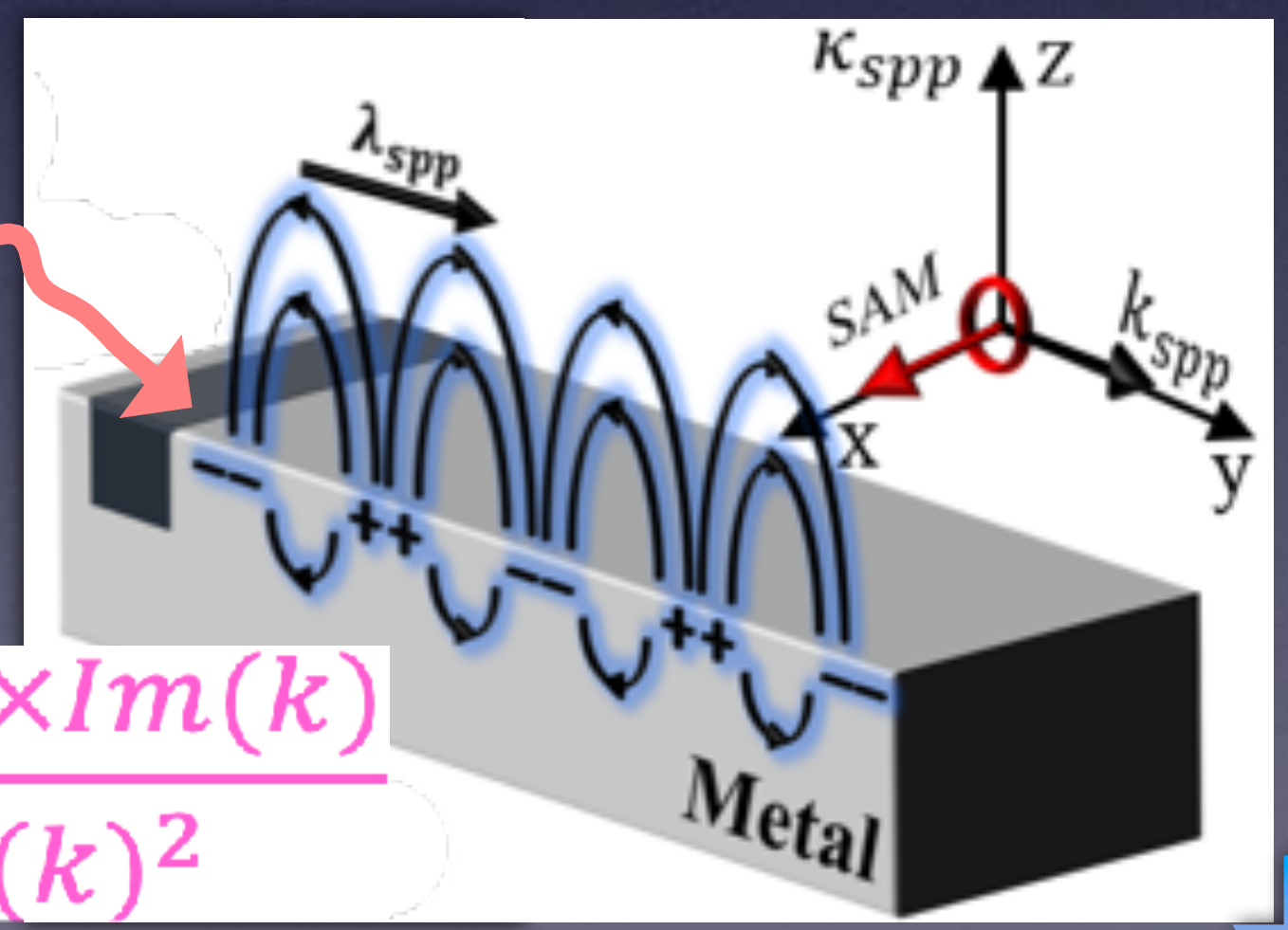


orbital angular momentum

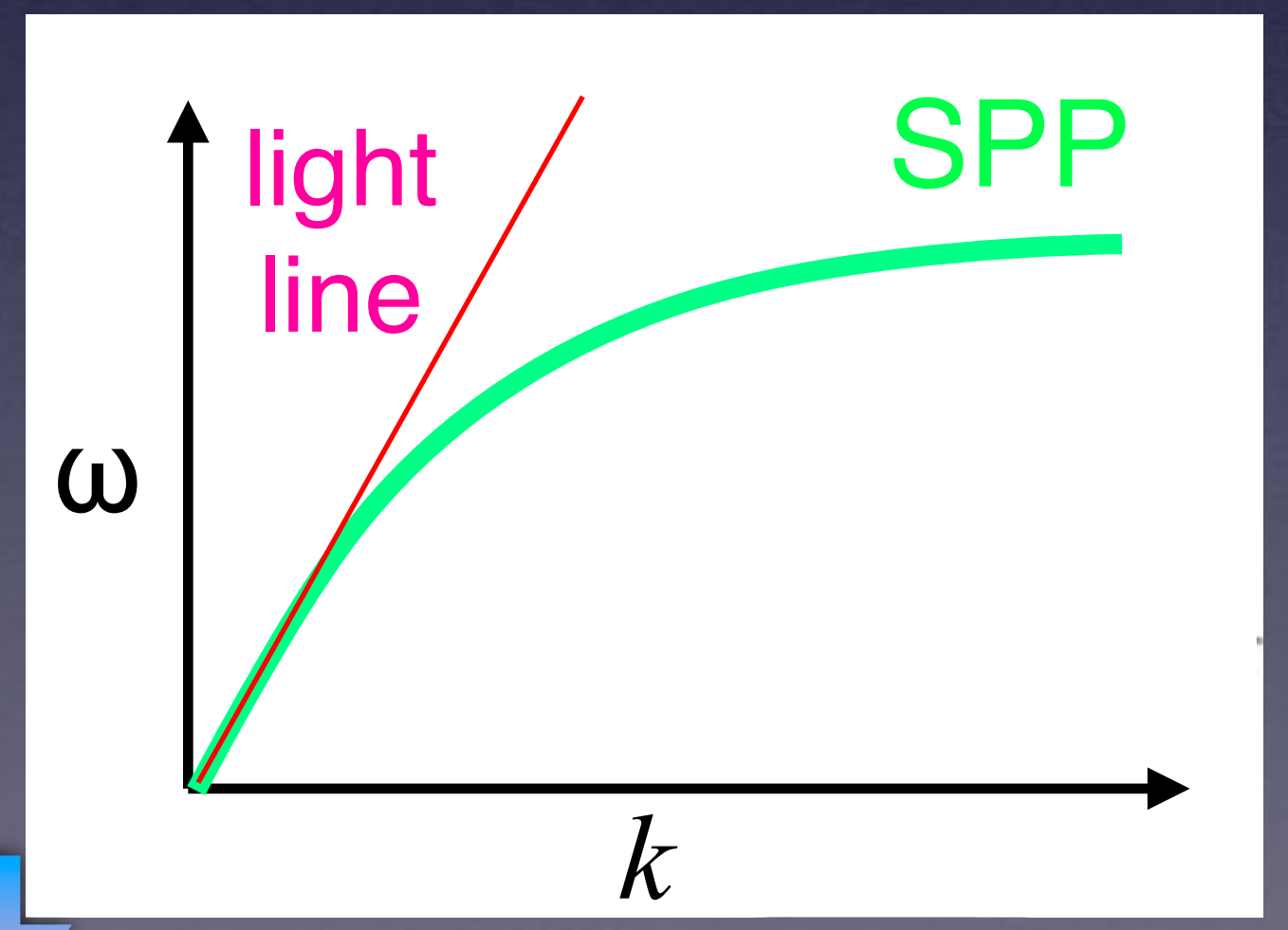


plasmonic spin Hall effect

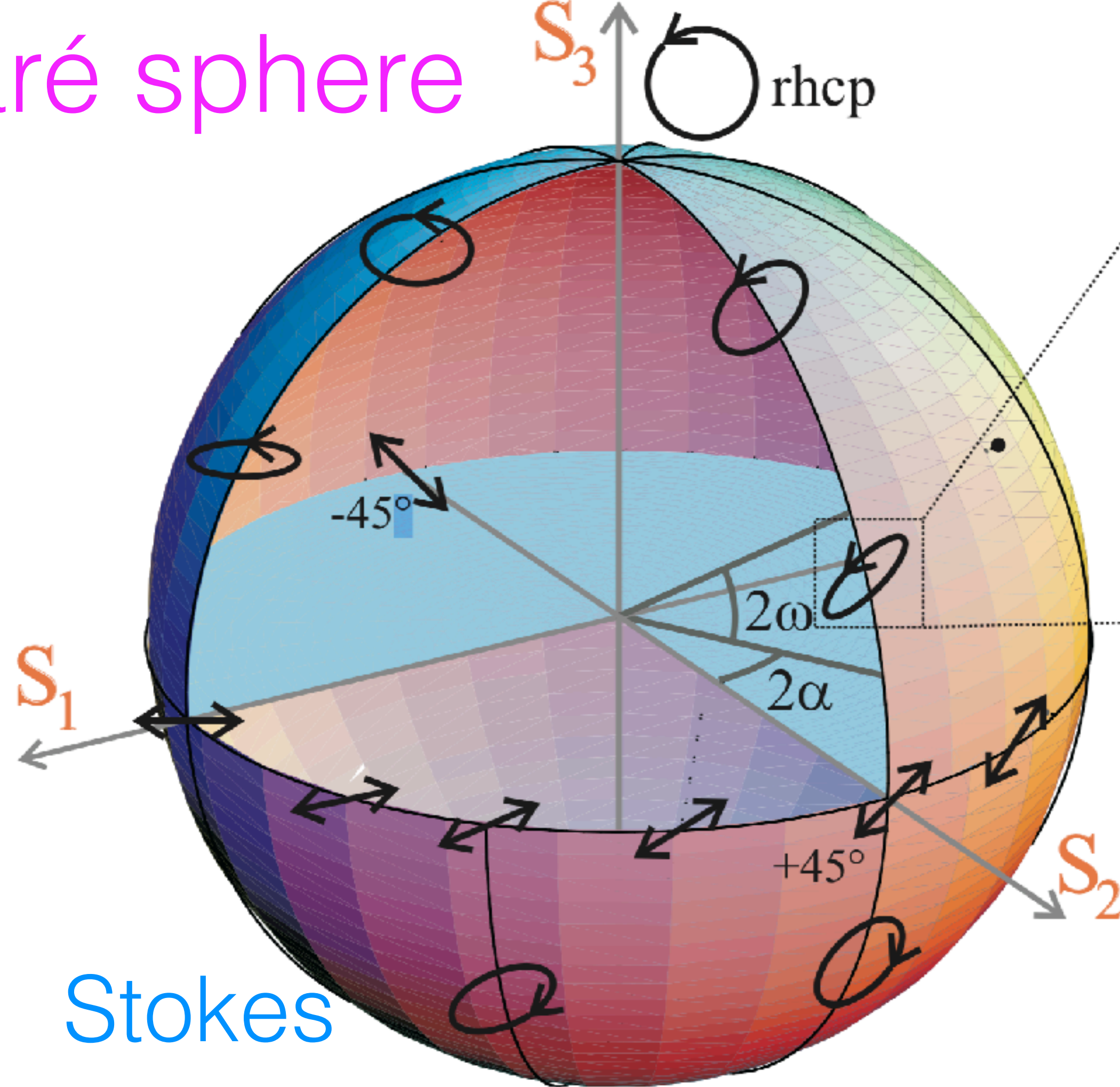
$$S \propto \frac{\text{Re}(k) \times \text{Im}(k)}{\text{Re}(k)^2}$$



spin-momentum locking



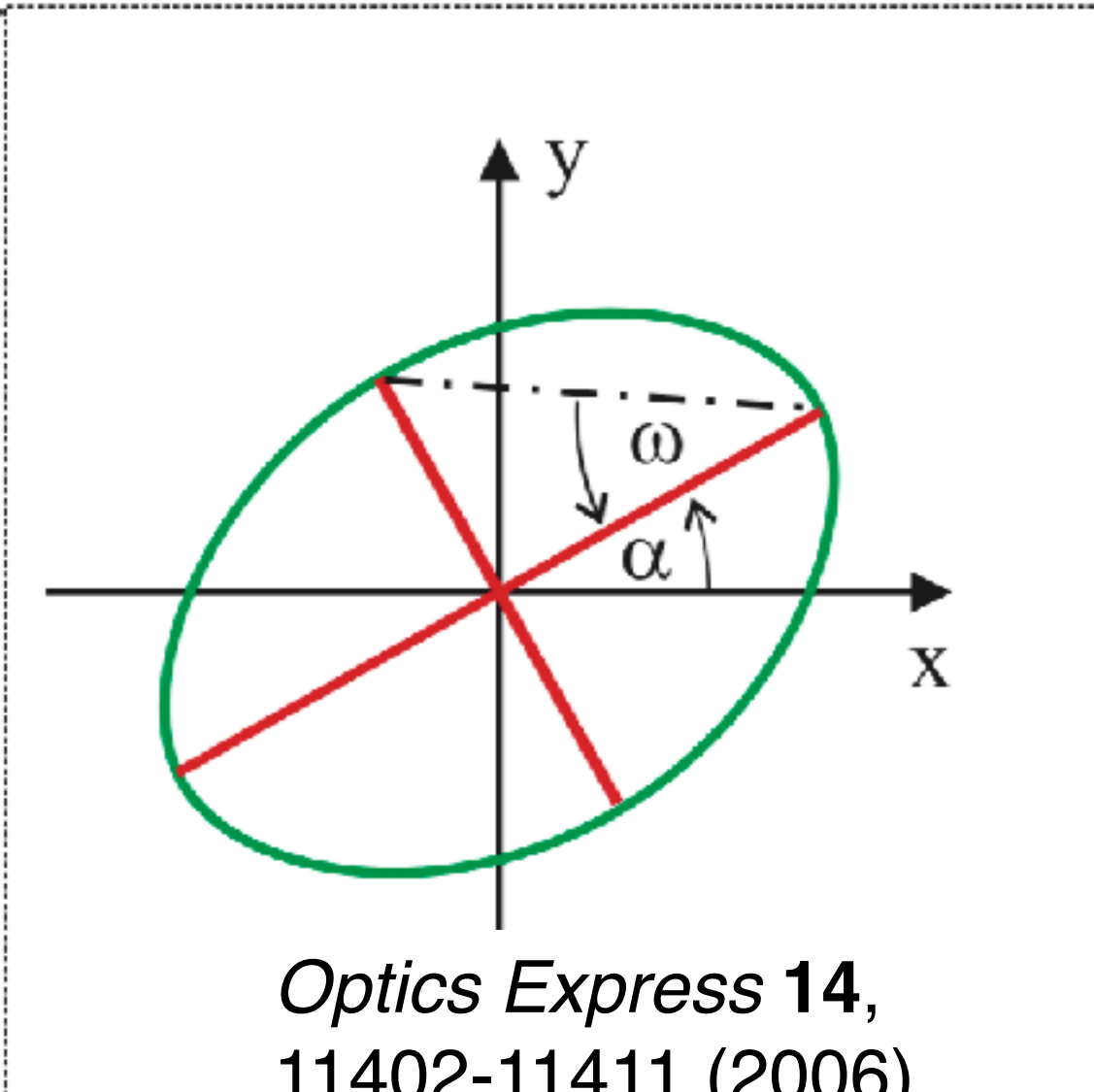
Poincaré sphere



Stokes parameters

\leftrightarrow	$S_0 = E_x ^2 + E_y ^2,$	$= I,$
\leftrightarrow	$S_1 = E_x ^2 - E_y ^2,$	$= I_{0^\circ} - I_{90^\circ}$
\nearrow	$S_2 = E_x^* E_y + E_x E_y^*$	$= I_{45^\circ} - I_{135^\circ}$
\circlearrowright	$S_3 = -i(E_x^* E_y - E_x E_y^*)$	$= I_{\text{rhcp}} - I_{\text{lhcp}}$

\circlearrowleft lhcp Pole - C-point - S-normal



Optics Express 14, 11402-11411 (2006).

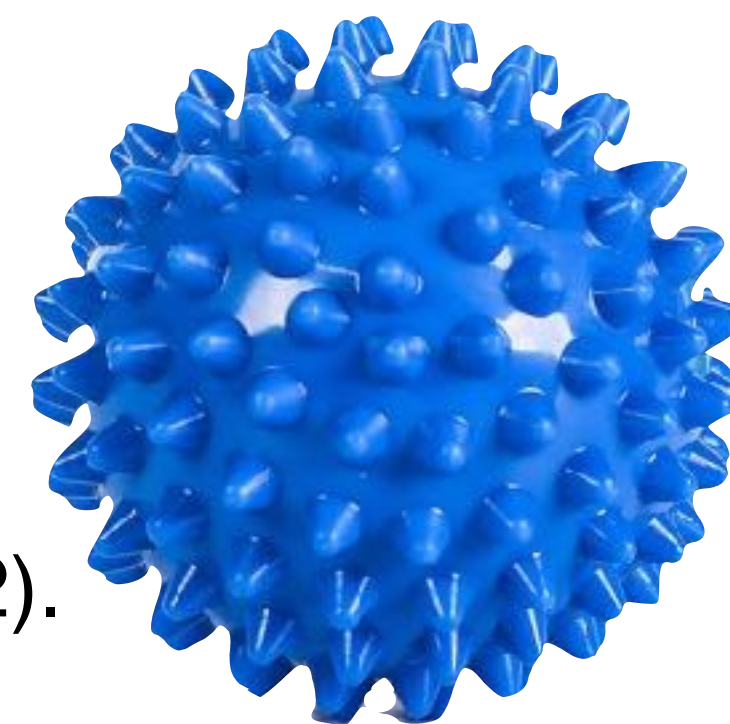
$$Q = \frac{1}{4\pi} \iint d^2r \bar{S} \cdot \left(\frac{\partial \bar{S}}{\partial x} \times \frac{\partial \bar{S}}{\partial y} \right)$$

Topological spin textures

Equator - L-line- S-in-plane

skyrmion
 $Q=1$

Appl. Phys. Rev. 9, 011420 (2022).

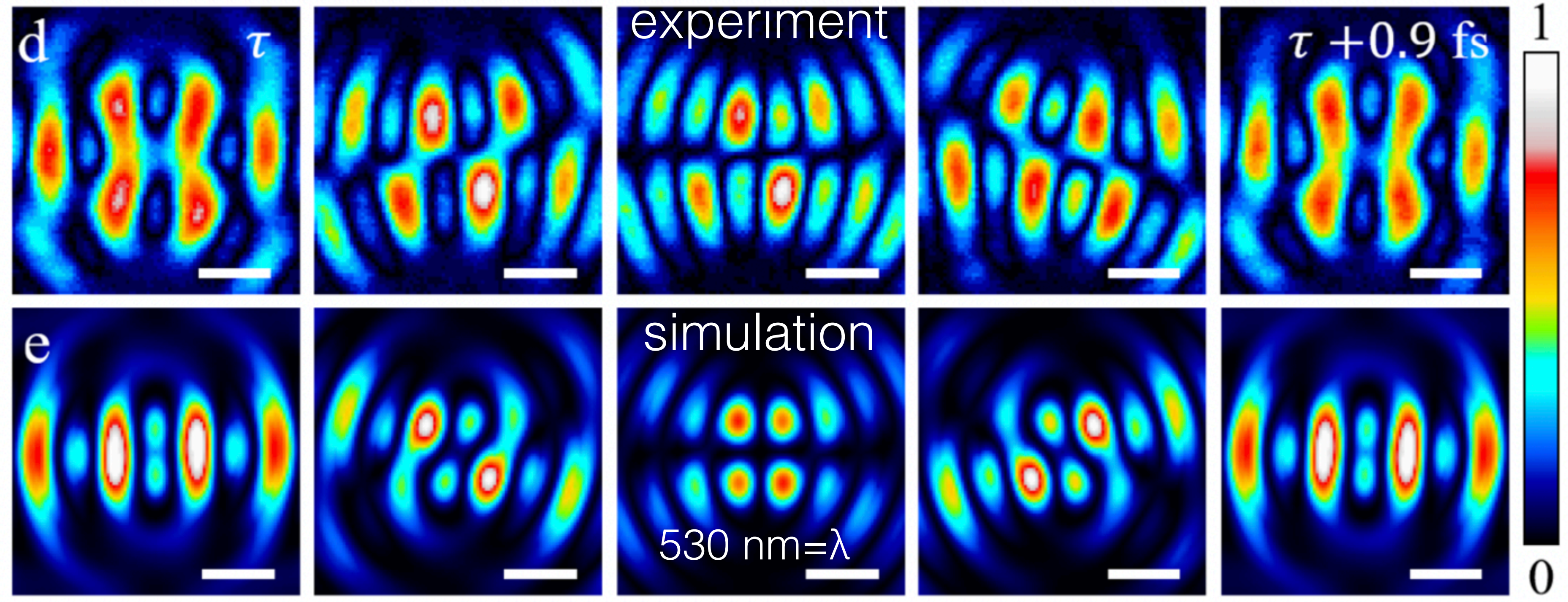
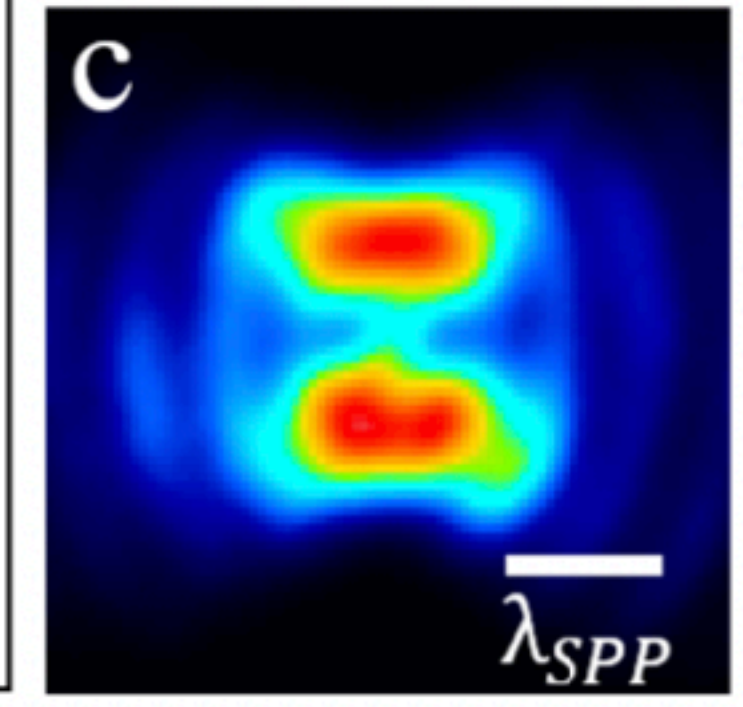
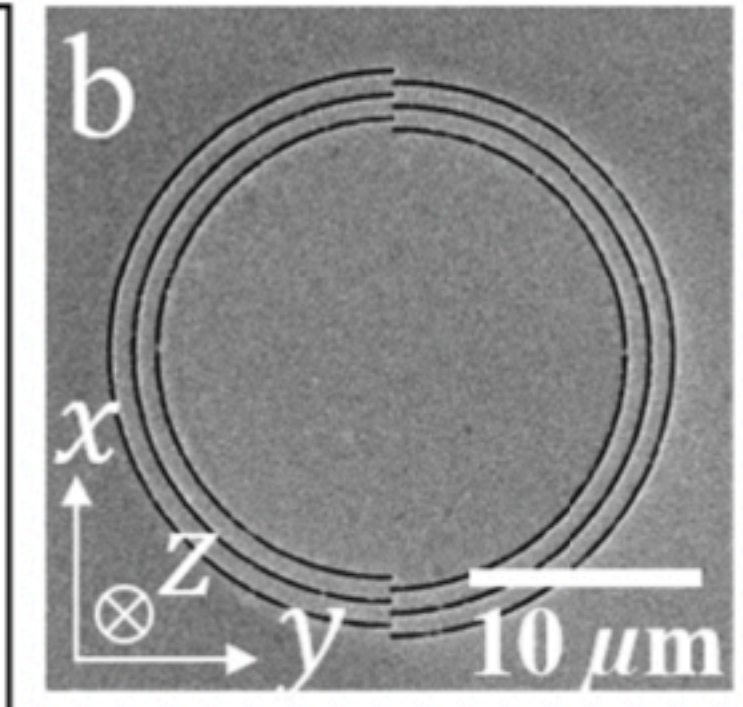
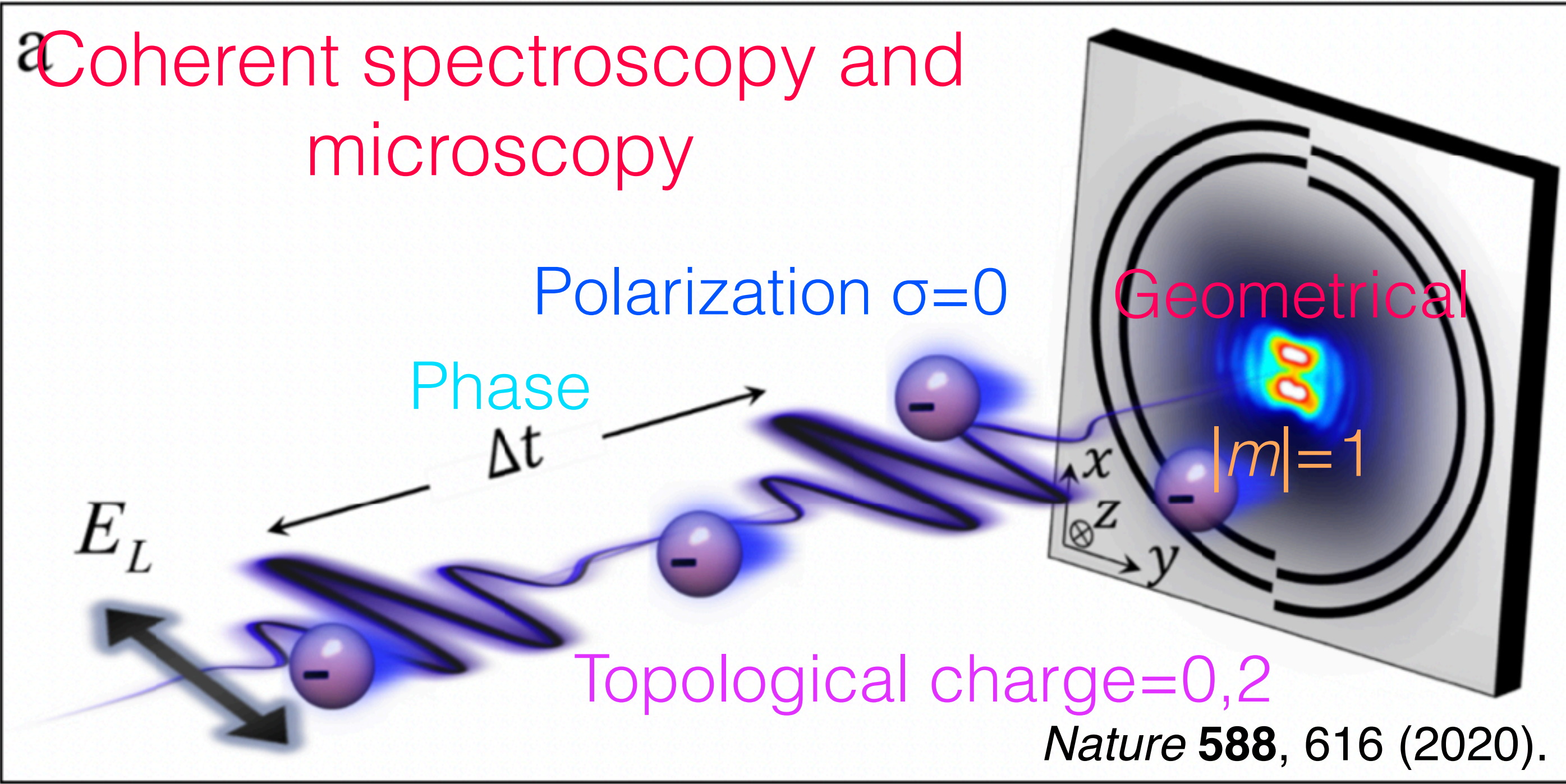


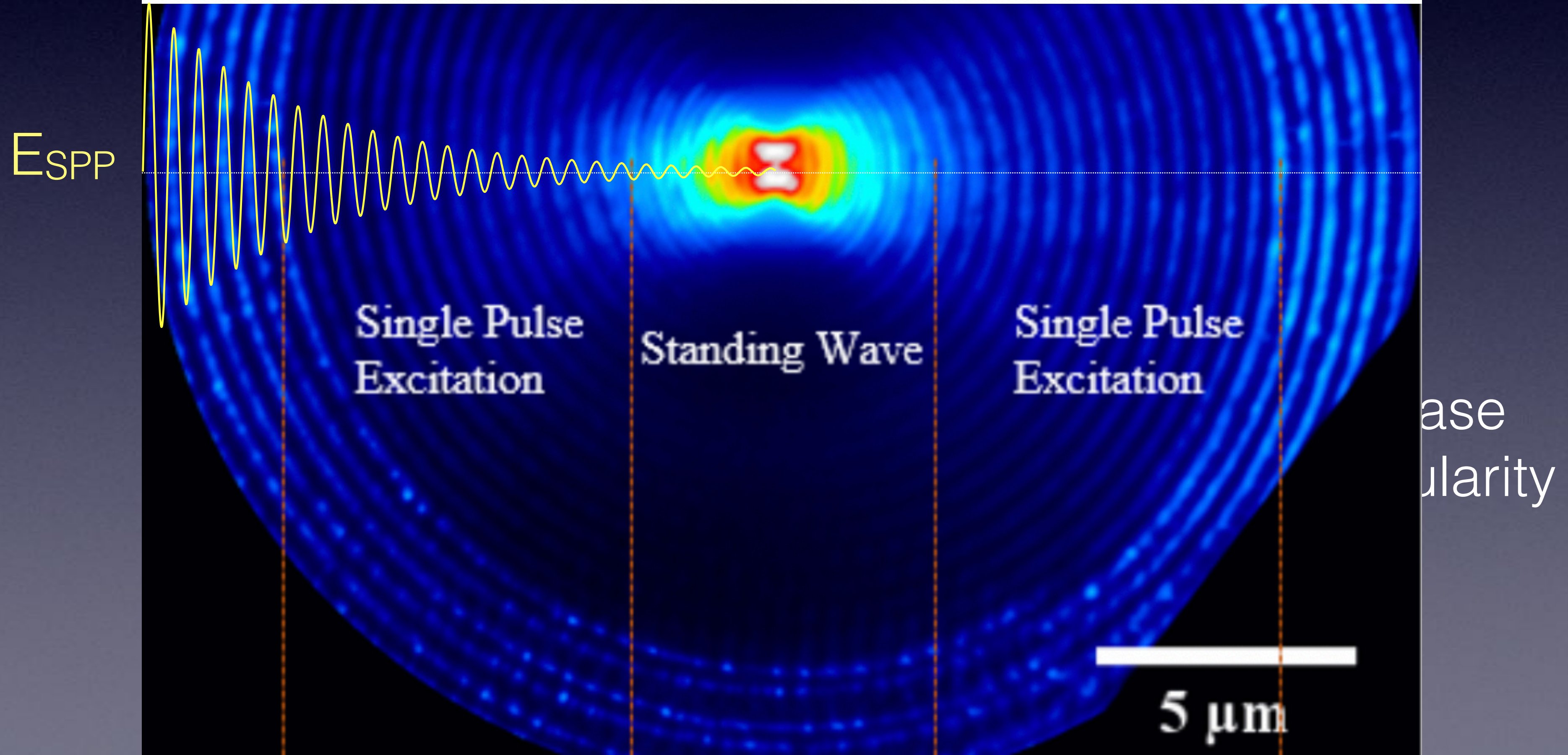
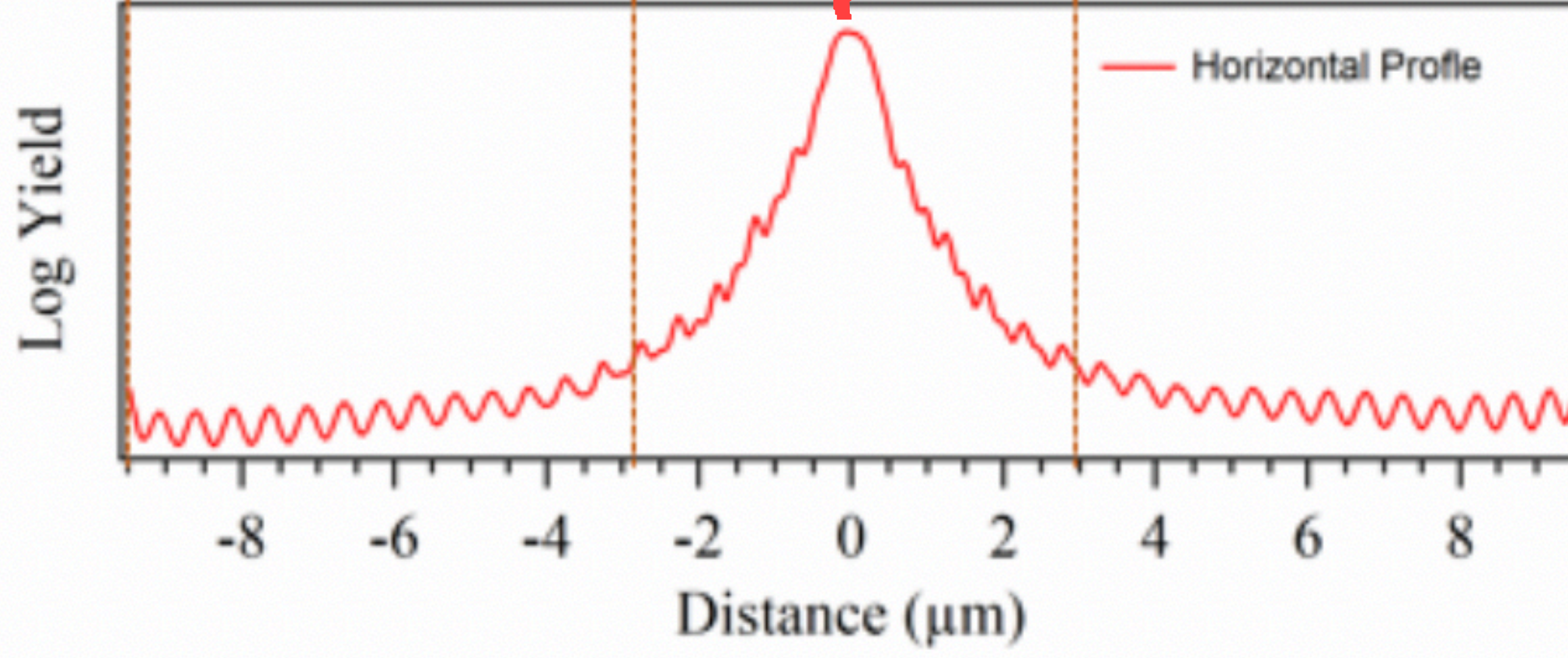
meron

$\frac{1}{2}$ skyrmion
 $Q=\frac{1}{2}$

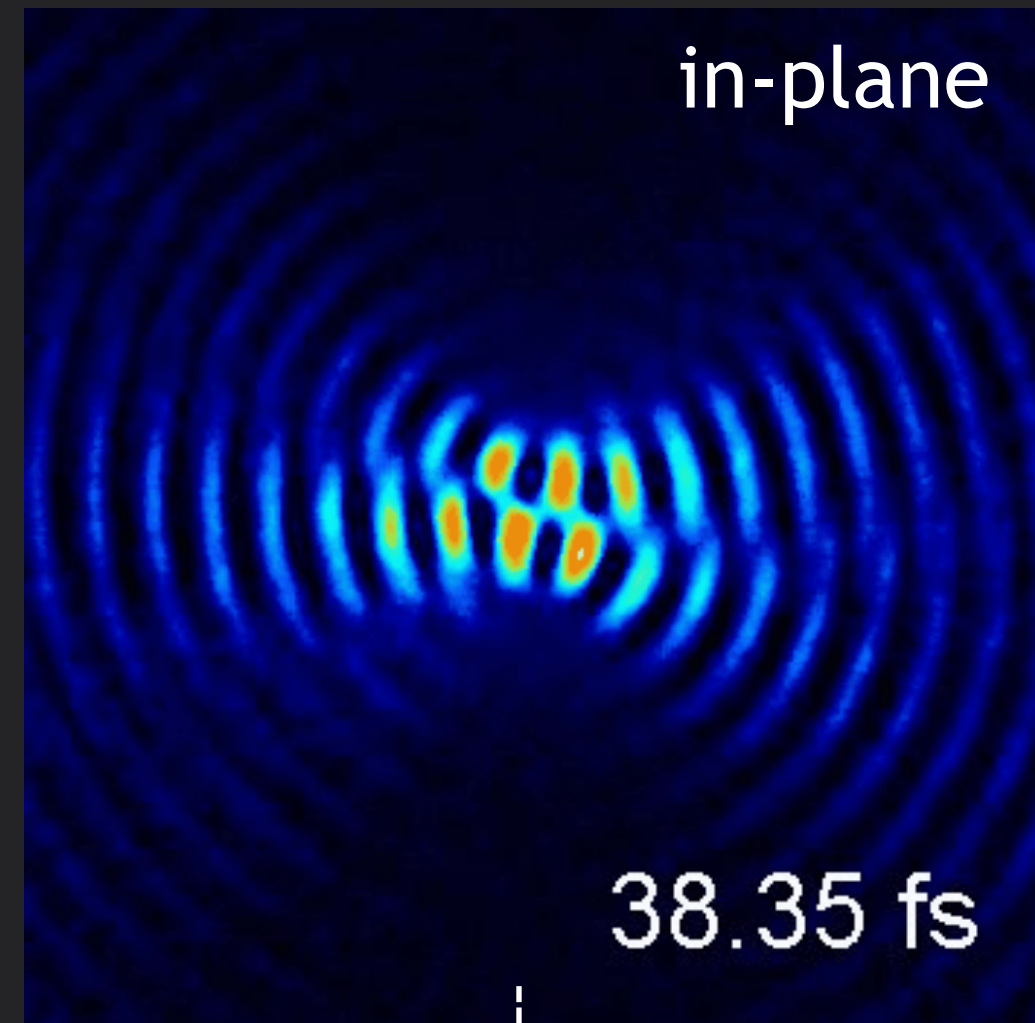
Nature 588, 616 (2020).



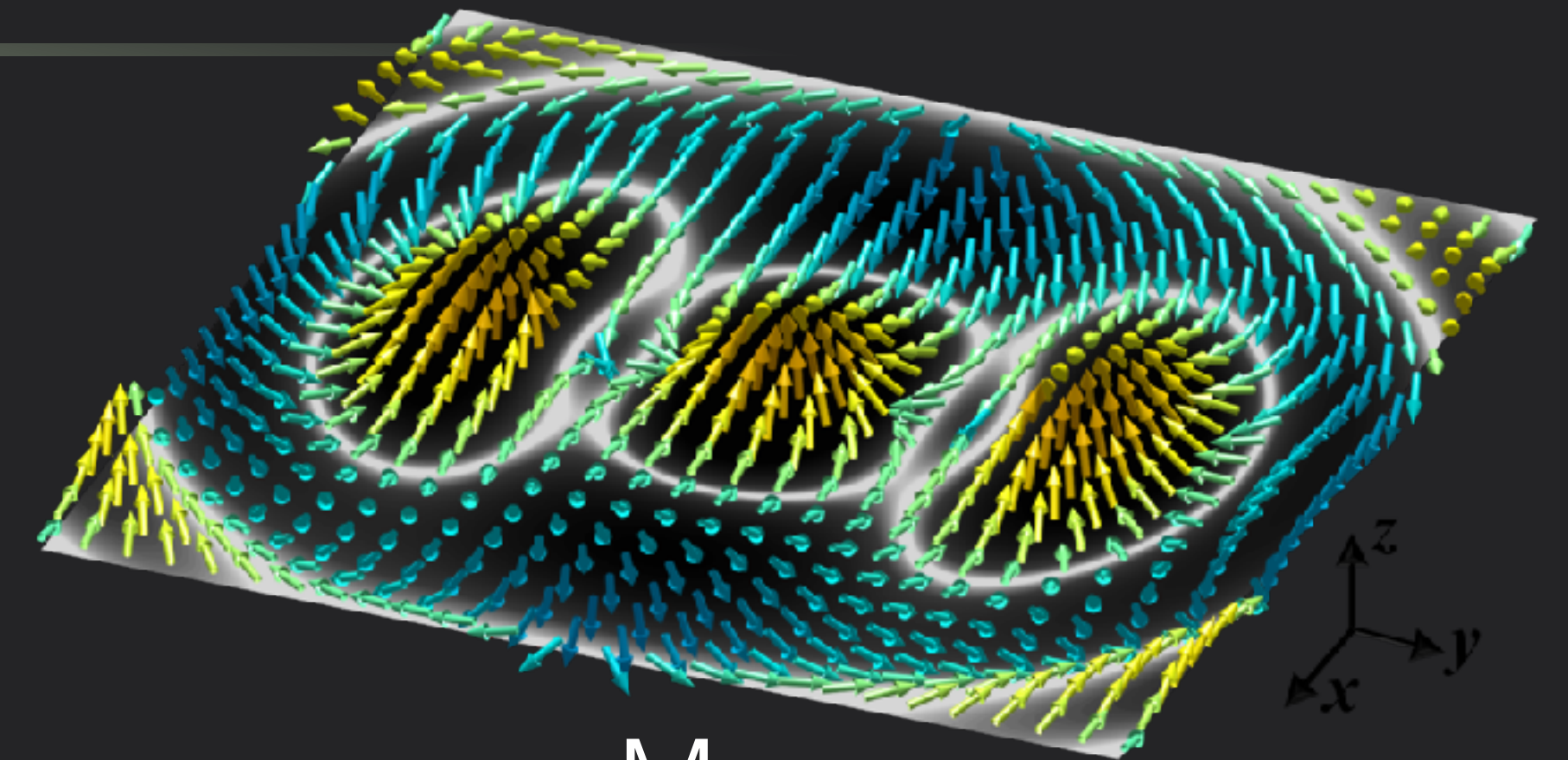
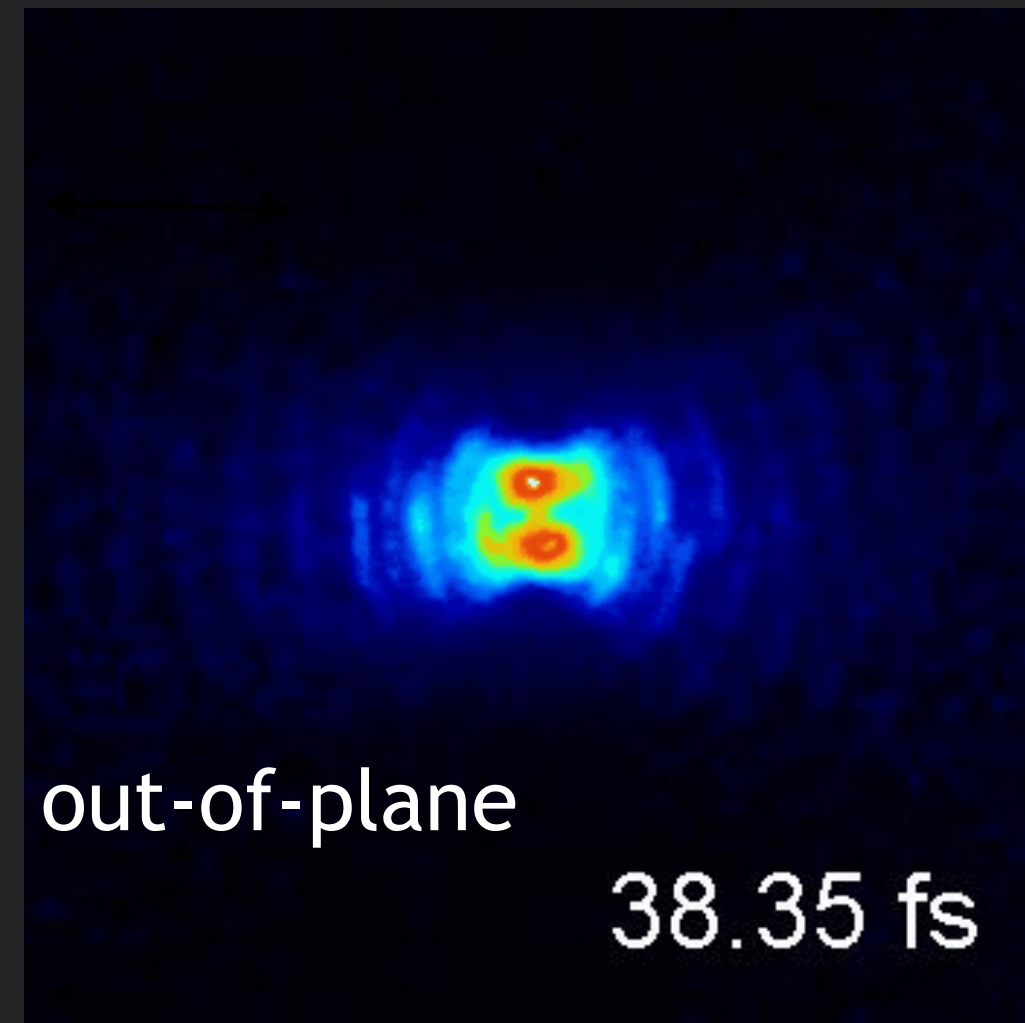




Locating Meron (half-skyrmion) with PEEM



Fourier Filtering



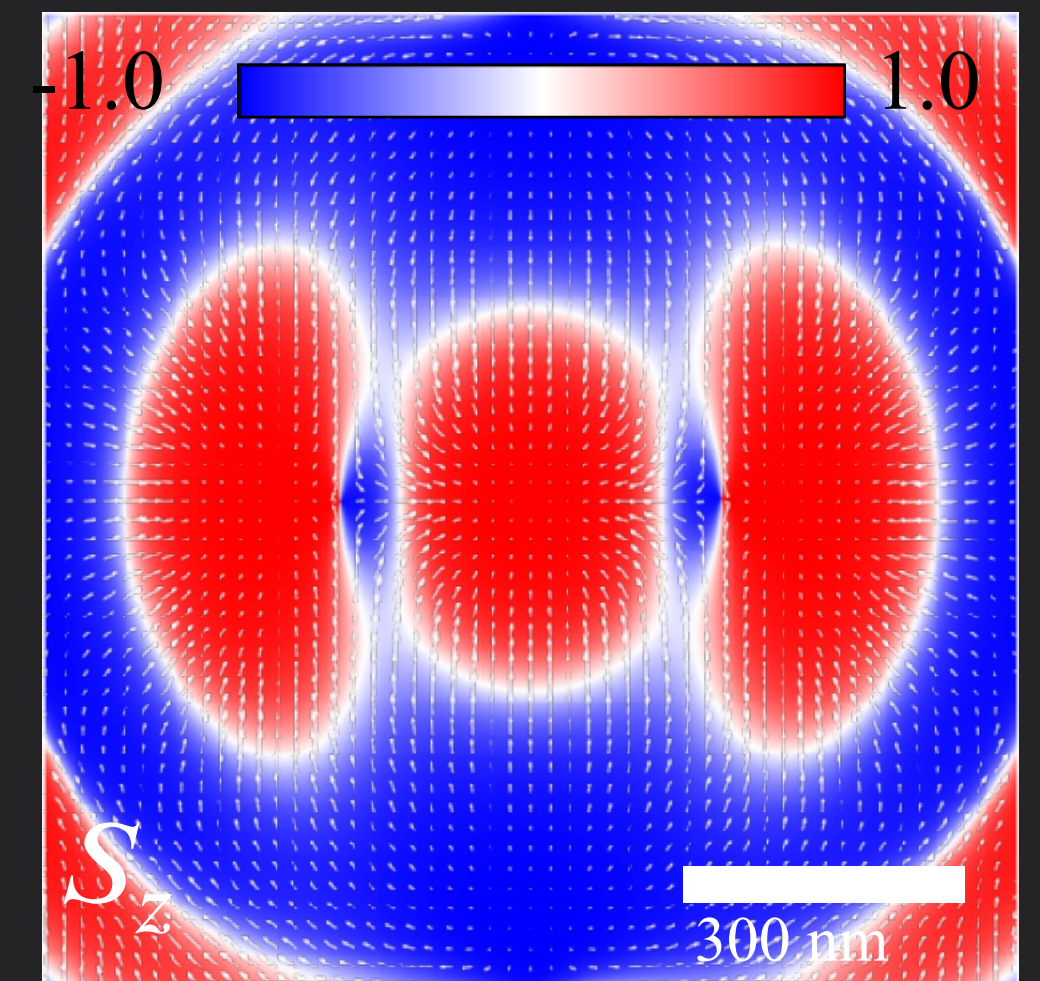
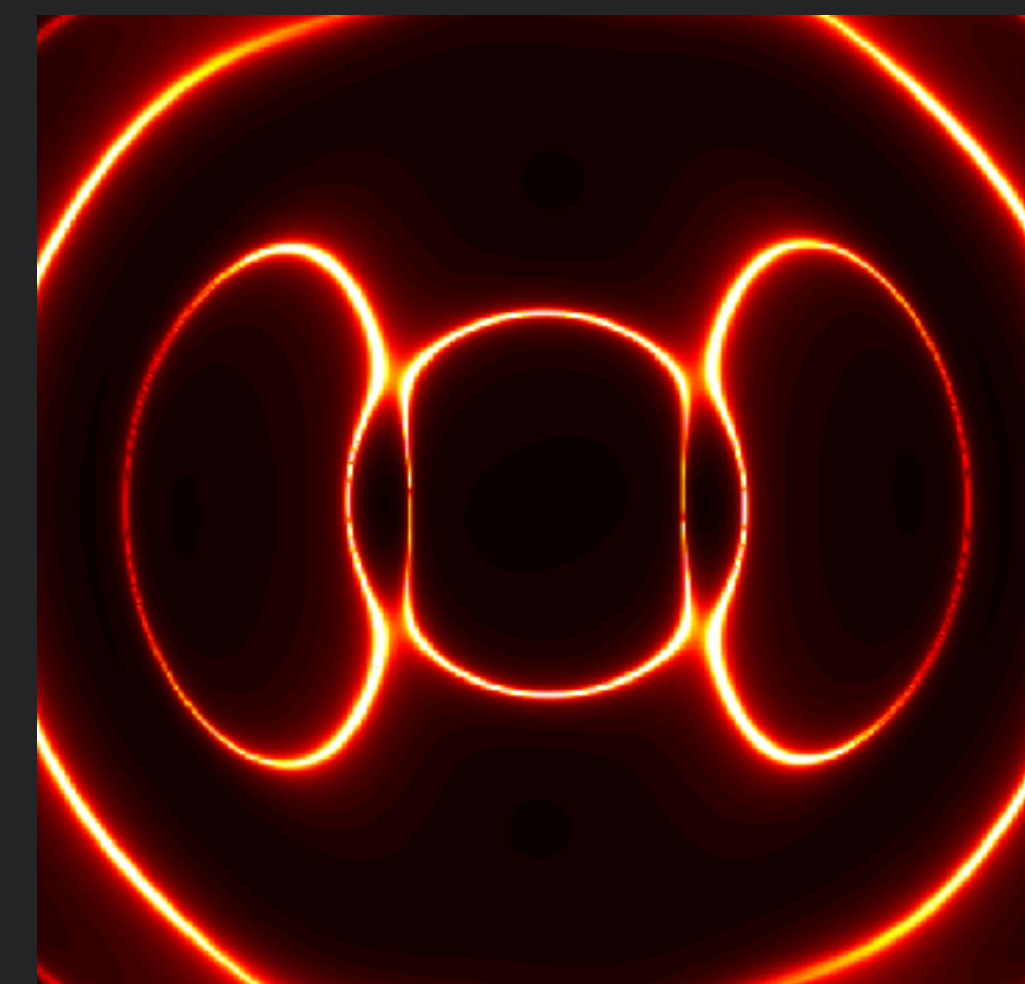
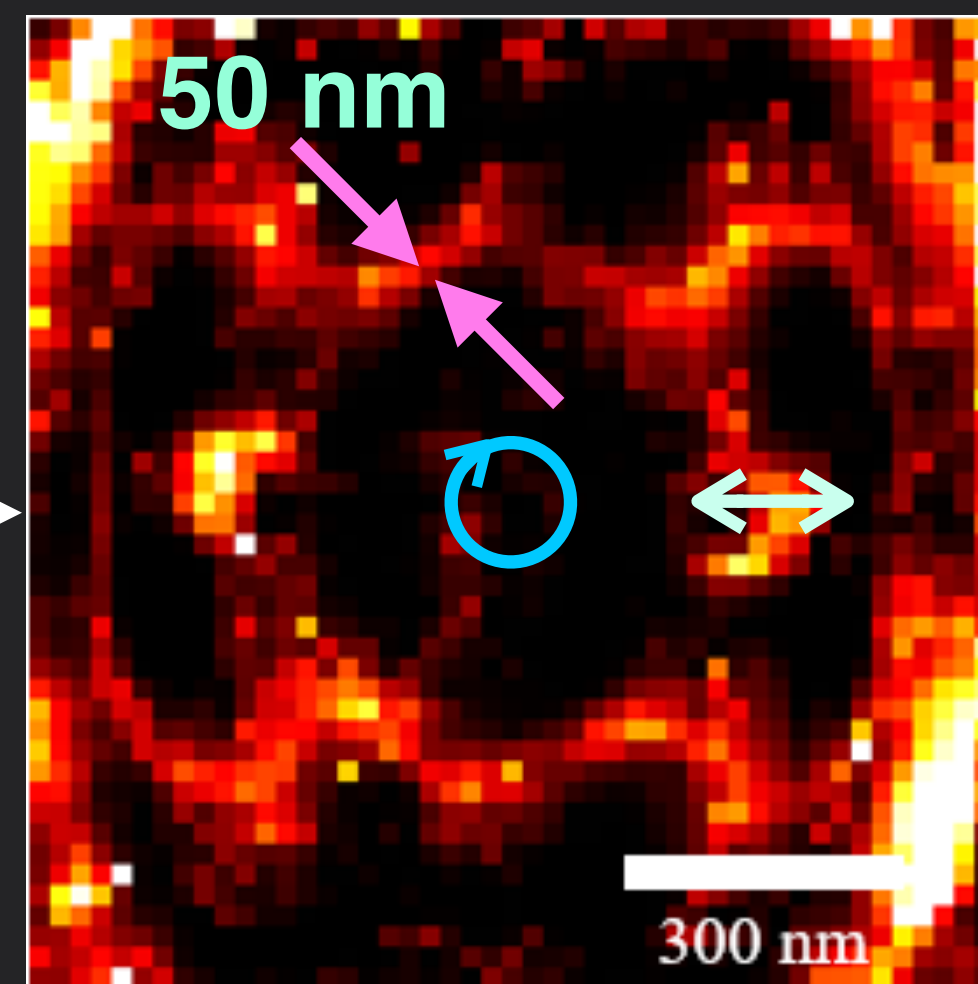
Nature **588**, 616 (2020)

Nature Reviews Physics **4**, 562 (2022)

Horn Schunck Method



Optical flow



Artif. Intell. **17**, 185 (1981).

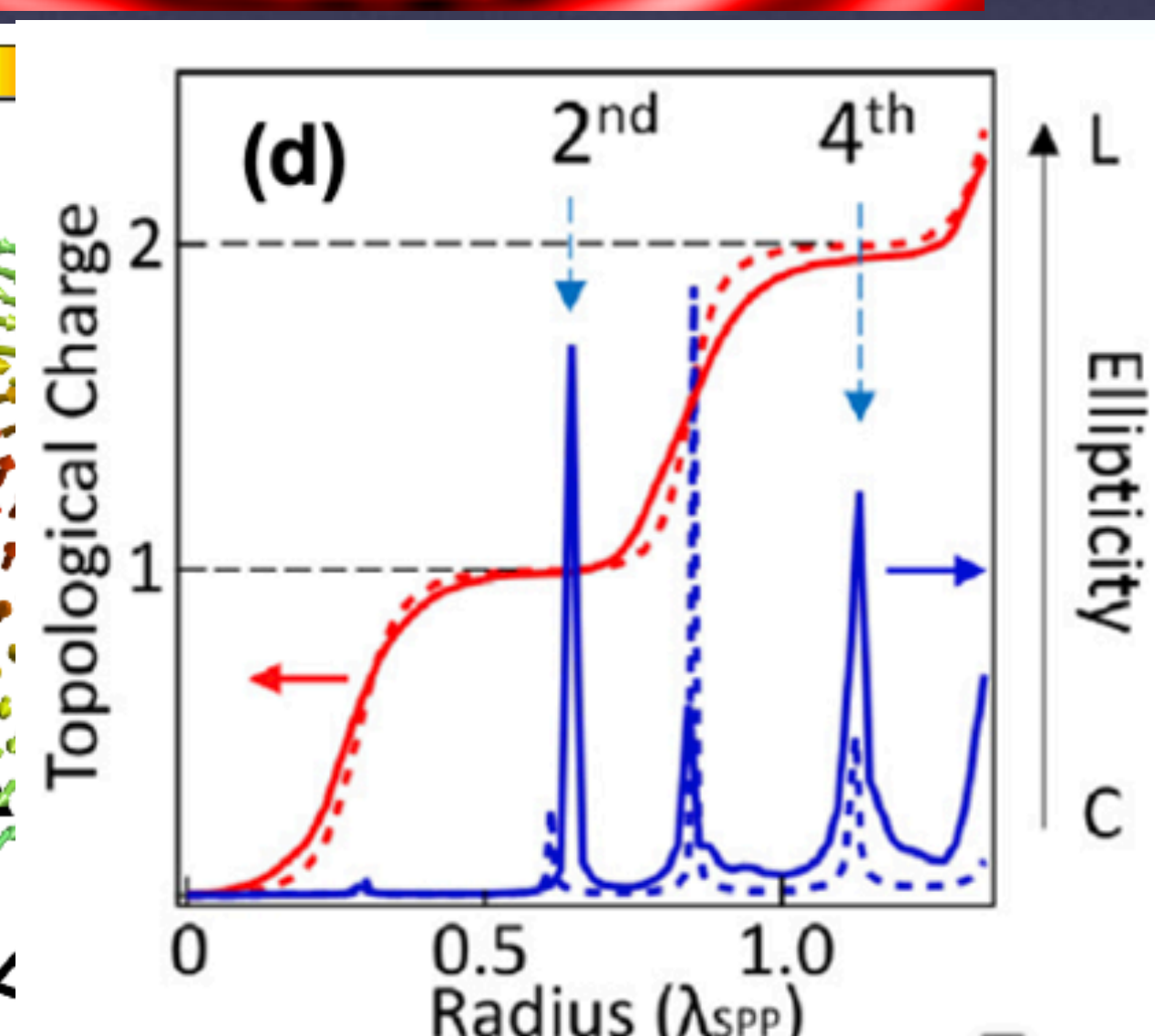
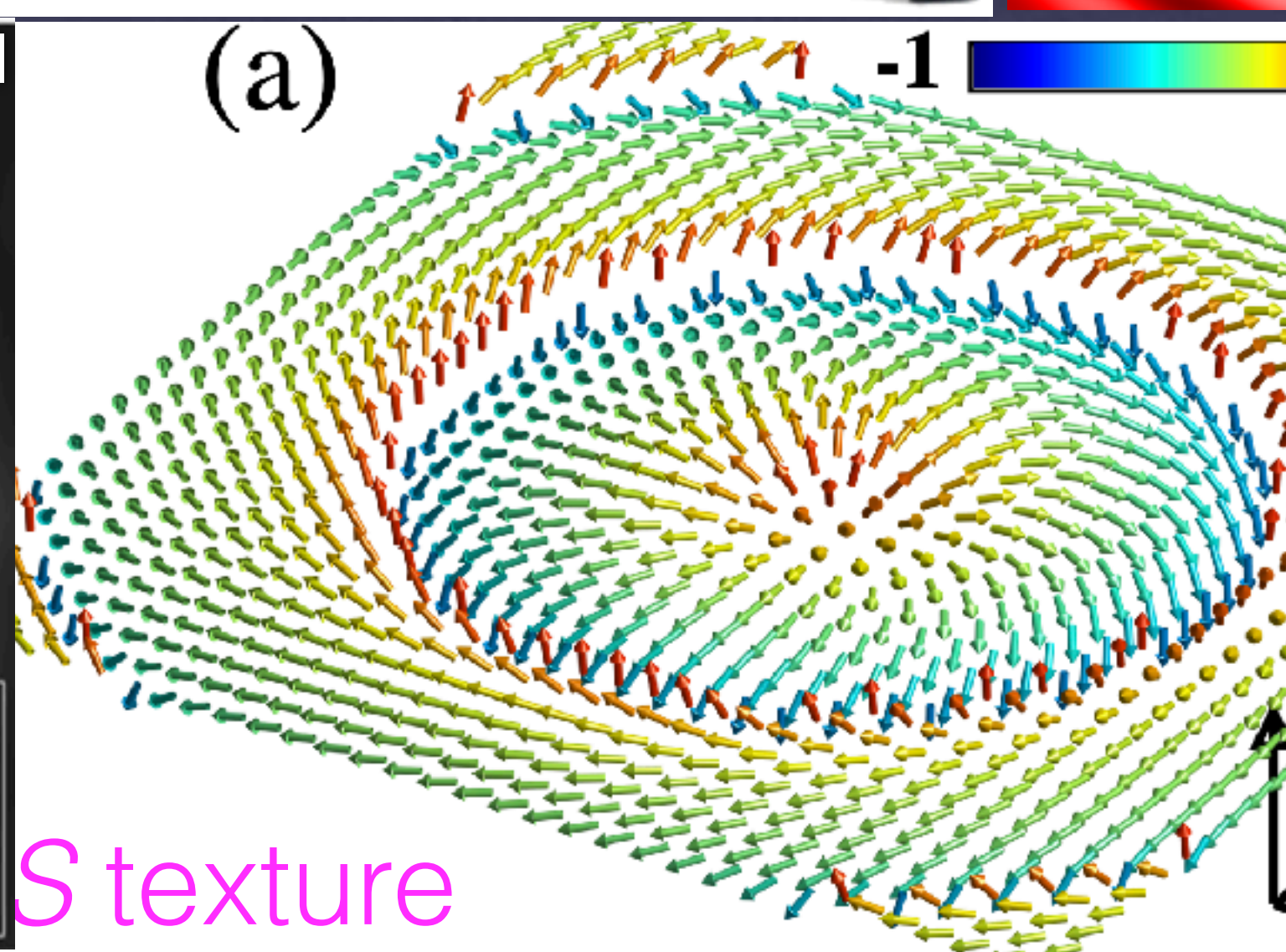
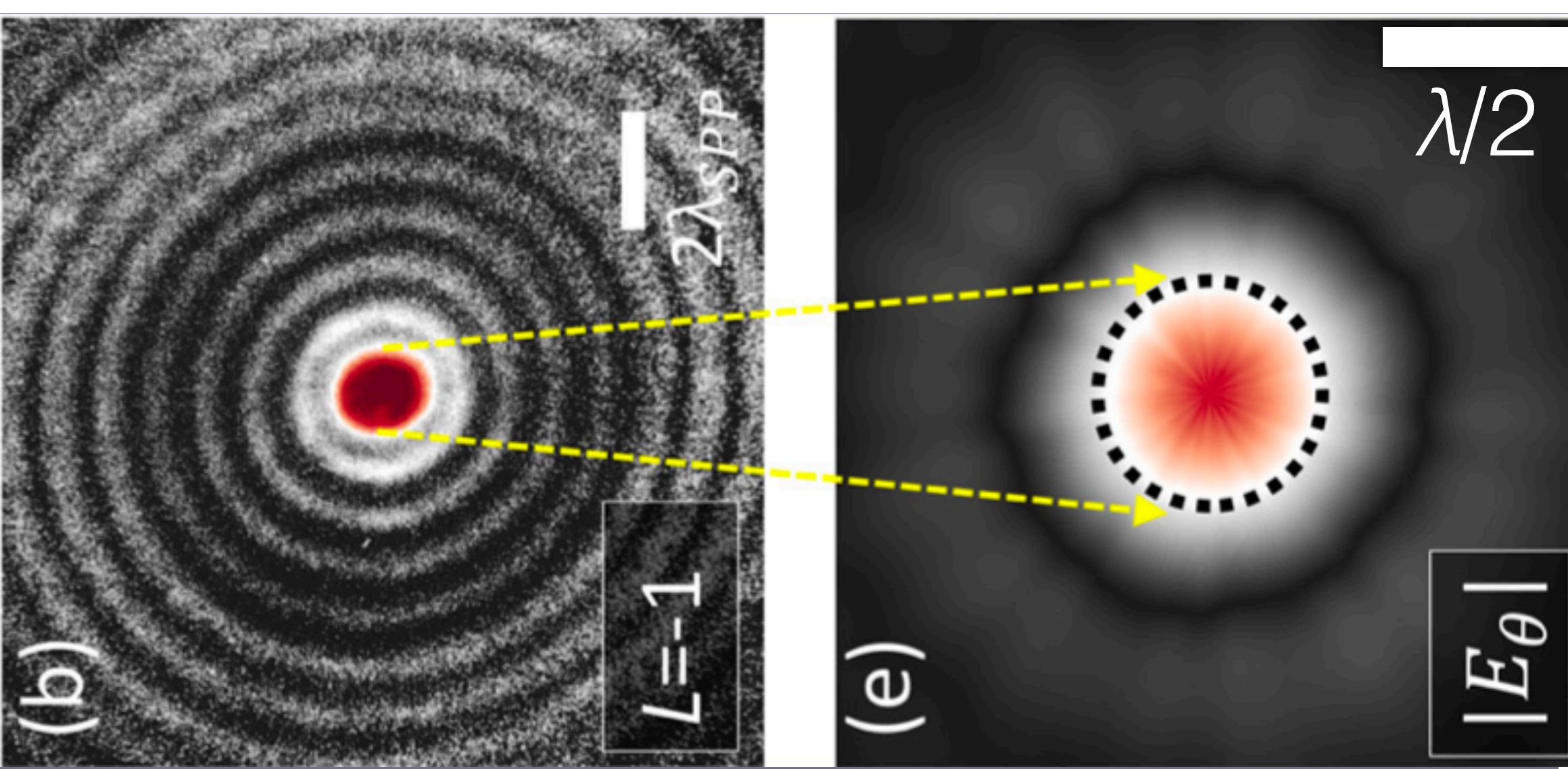
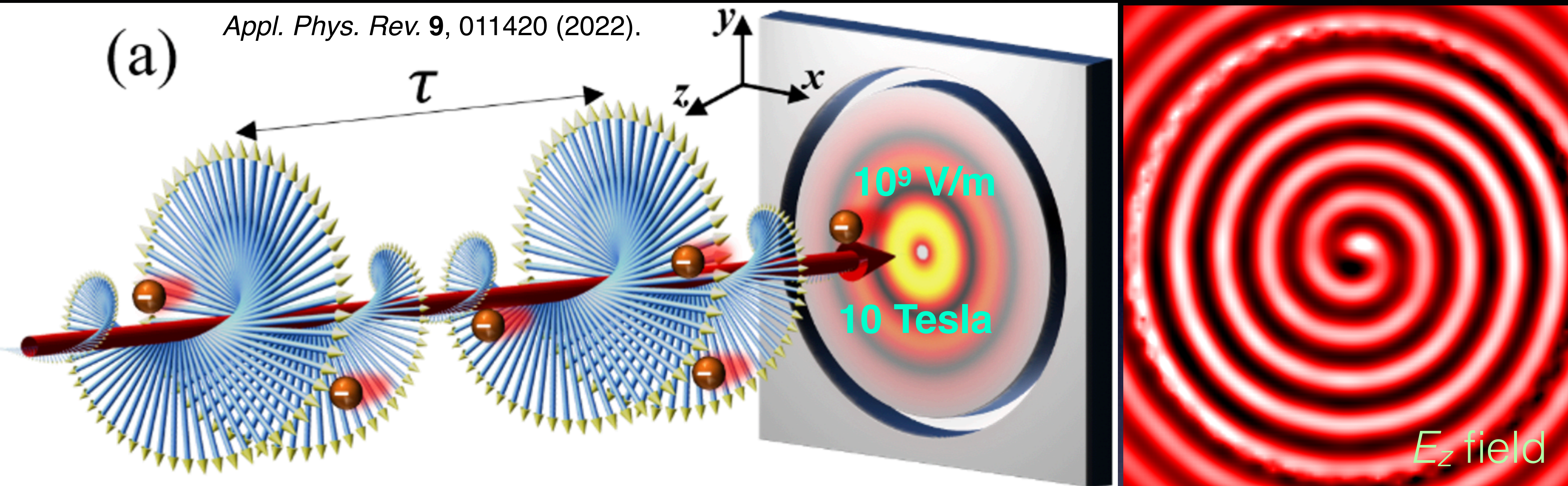
$$F_{HS} \sim \iint (I_x v_x + I_y v_y + I_t) dx dy$$

L-line map

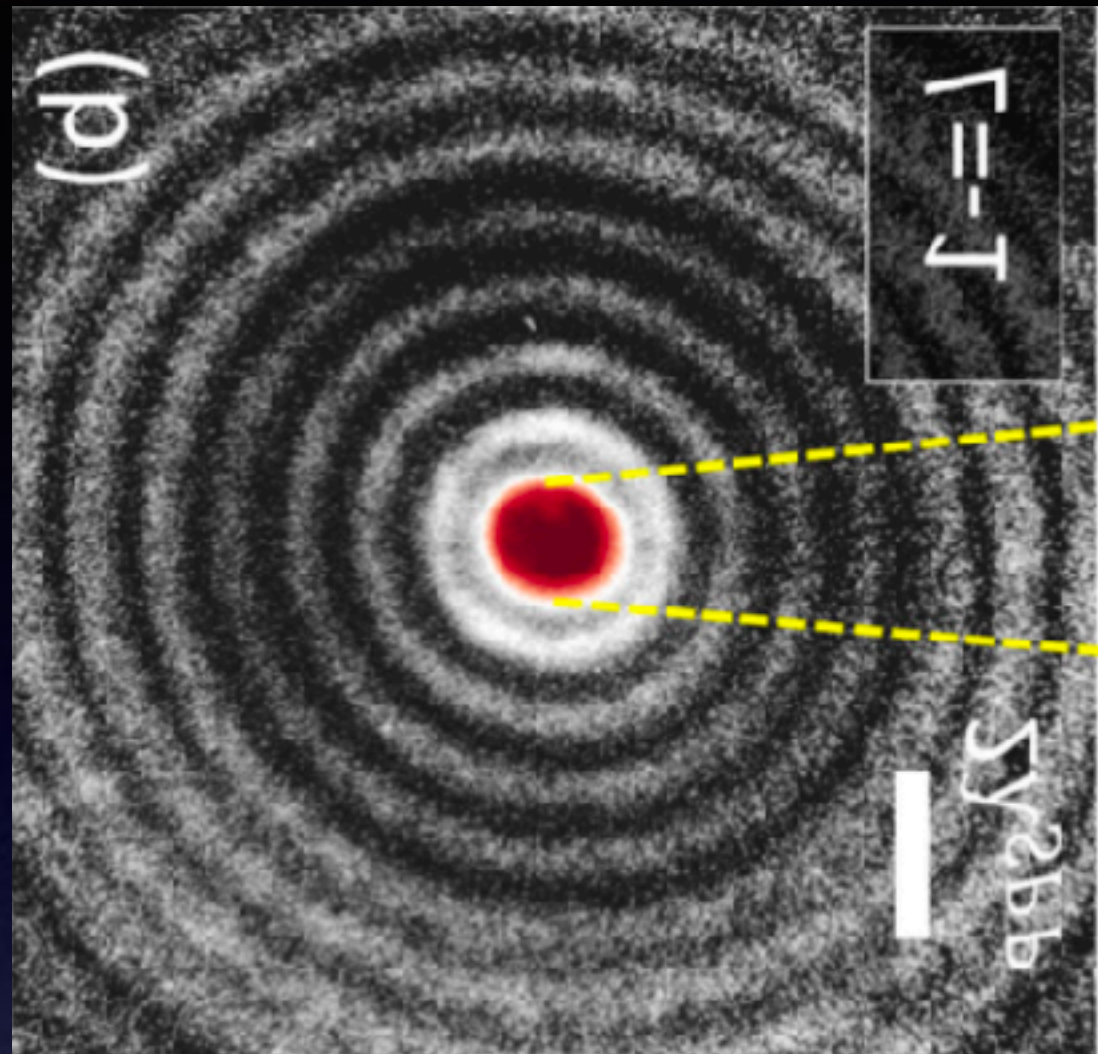
S_z
13

Spin-angular momentum ($\sigma=\pm 1$)+geometric charge ($m=0$)

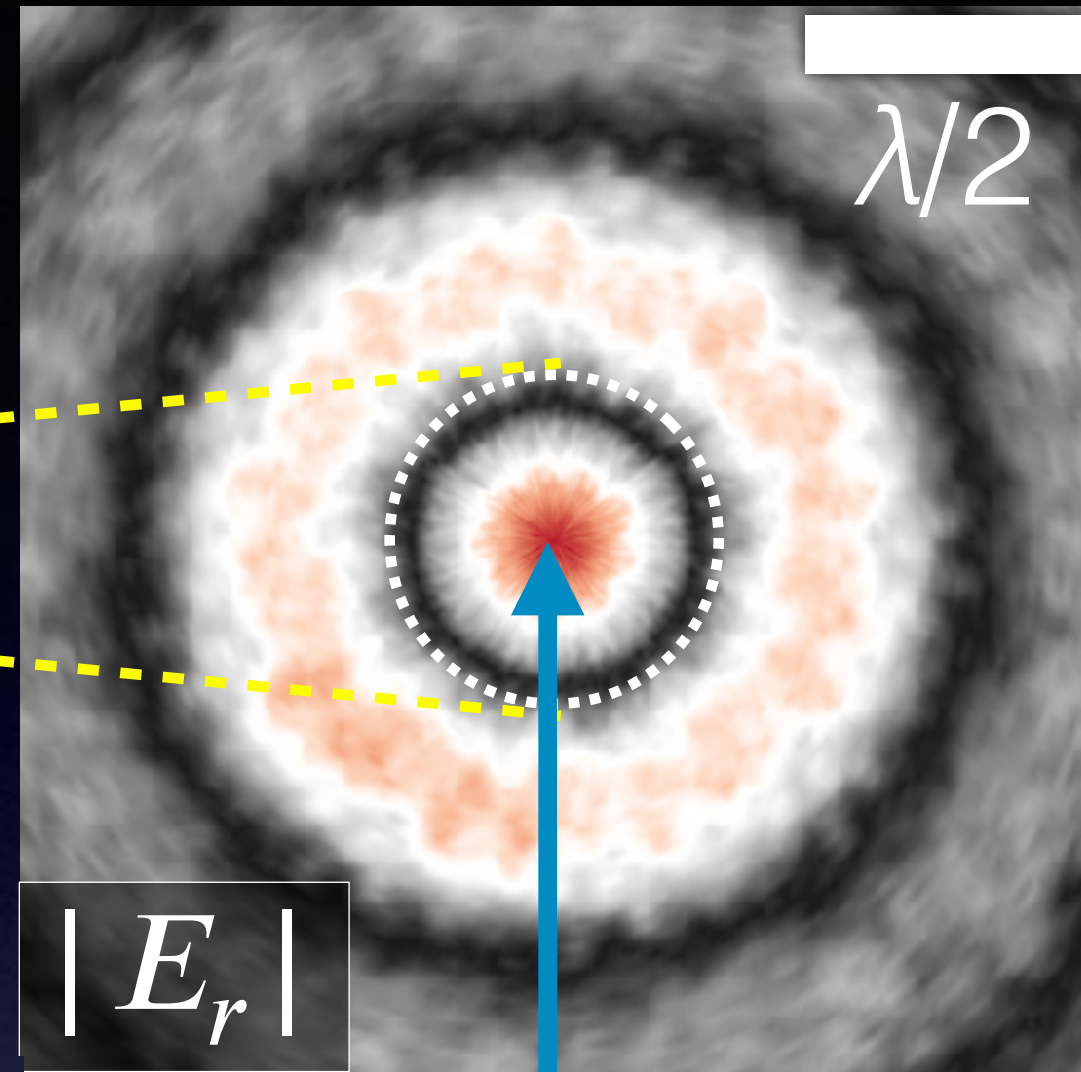
Appl. Phys. Rev. **9**, 011420 (2022).



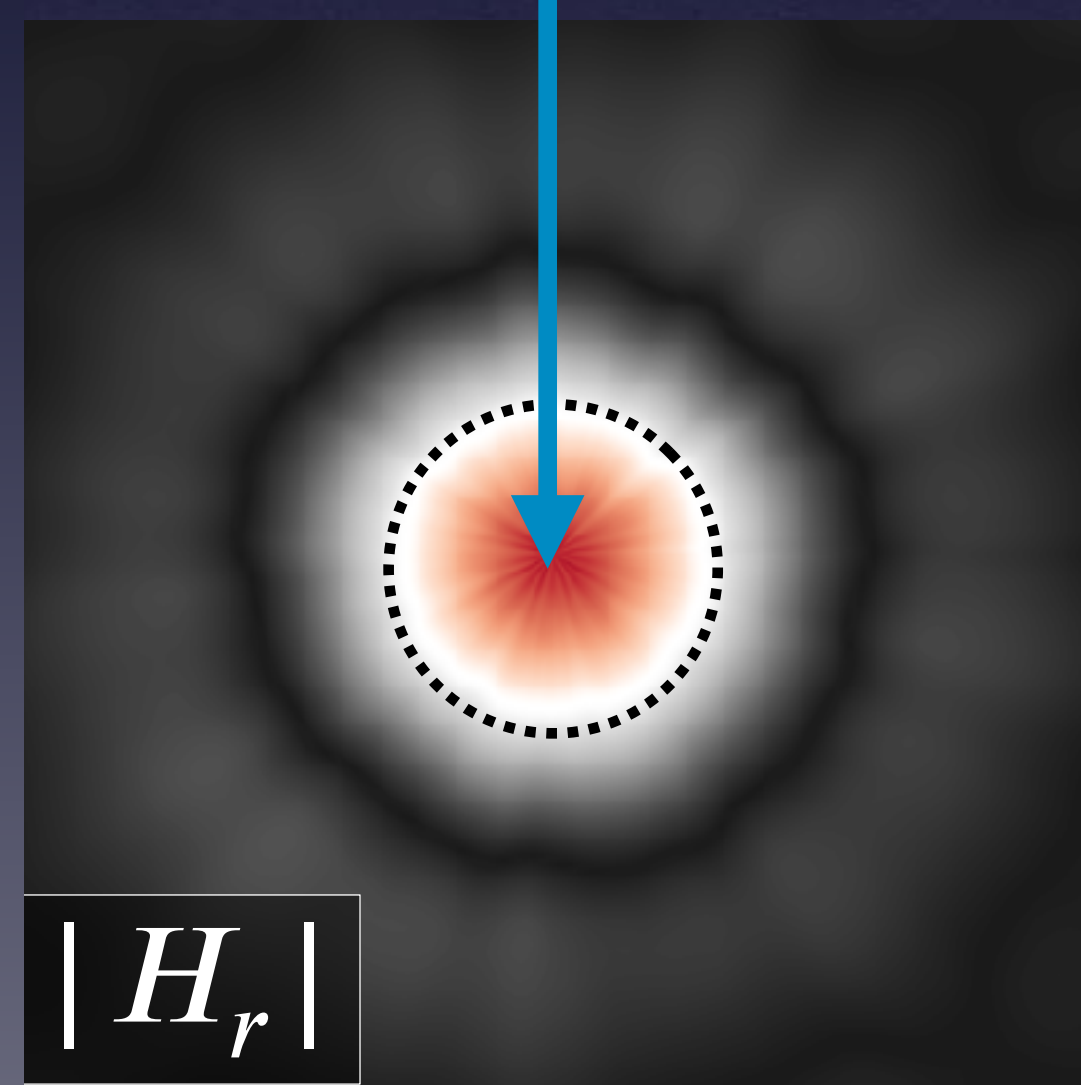
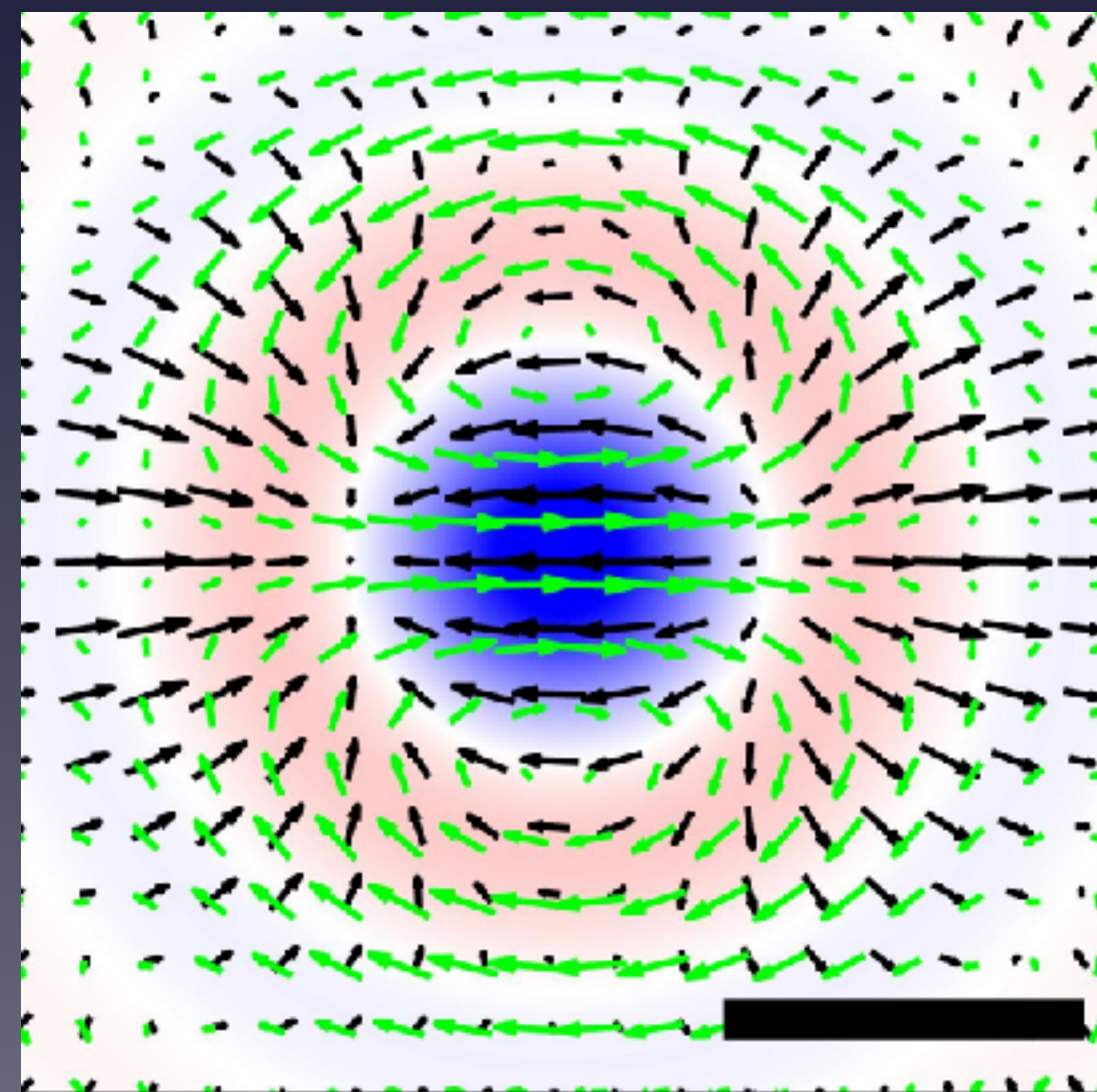
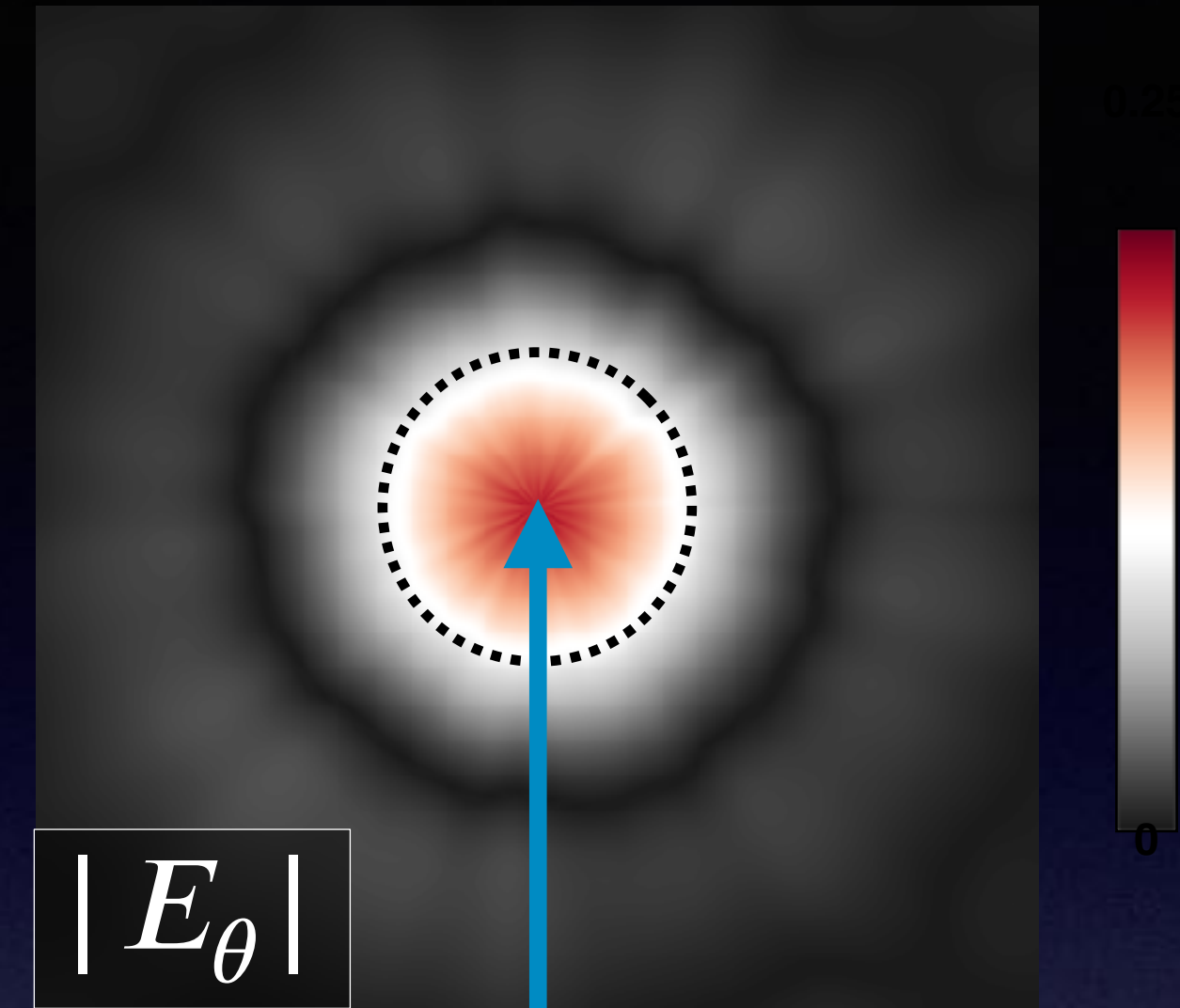
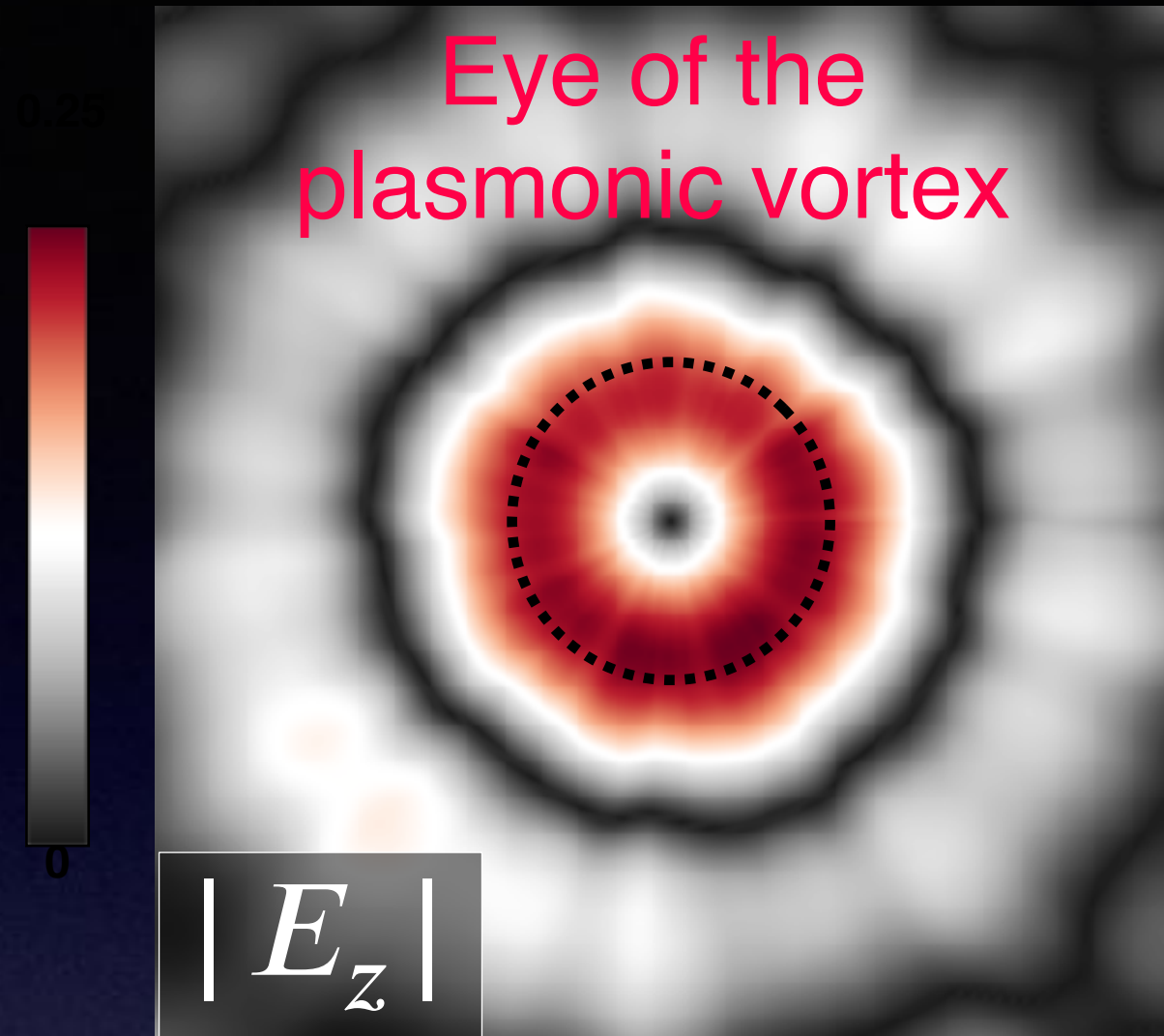
PEEM Image



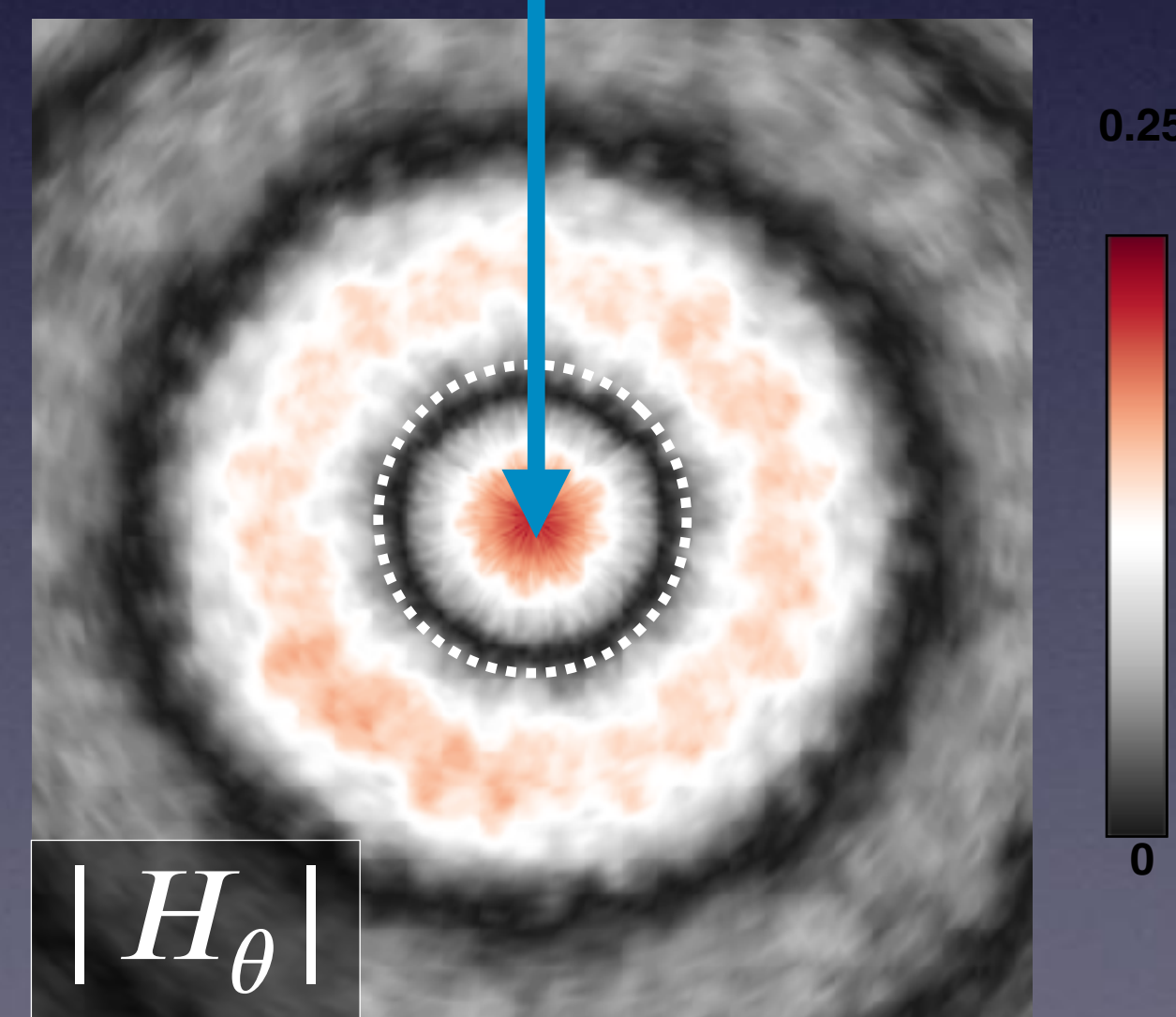
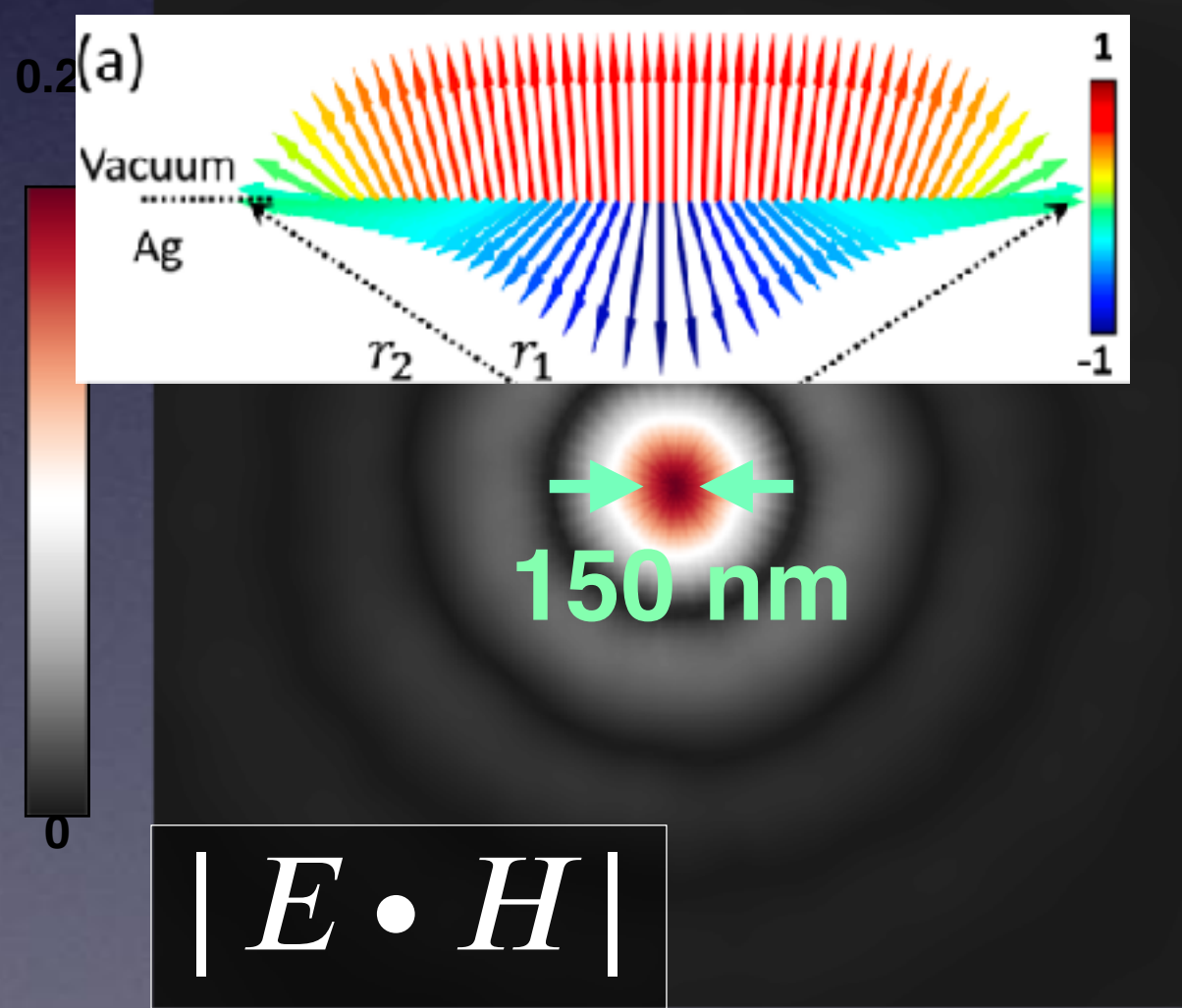
FT analysis



Maxwell



magnetic monopole



$\longleftrightarrow 2\lambda \text{ SPP} \longrightarrow$

Magnetolectric field

Conclusions

- We have demonstrated imaging of vectorial plasmonic fields on $<\lambda/2$ length scale with attosecond precision
- This enables study of topological spin texturing by spin-orbit interaction of light. The textures are stable on second to attosecond pulse duration time.
- we can study spin-charge entanglement, \mathcal{PT} symmetry breaking, non-reciprocal and non-Hermitian physics, in trivial and bianisotropic condensed matter
- The structured fields can induce Poincaré dressing of topologically trivial matter (*Nature Reviews Physics* **4**, 562 (2022))
- Where did the angular momentum go?
- Can we use structured light to probe axion physics in topological matter or even cosmos?