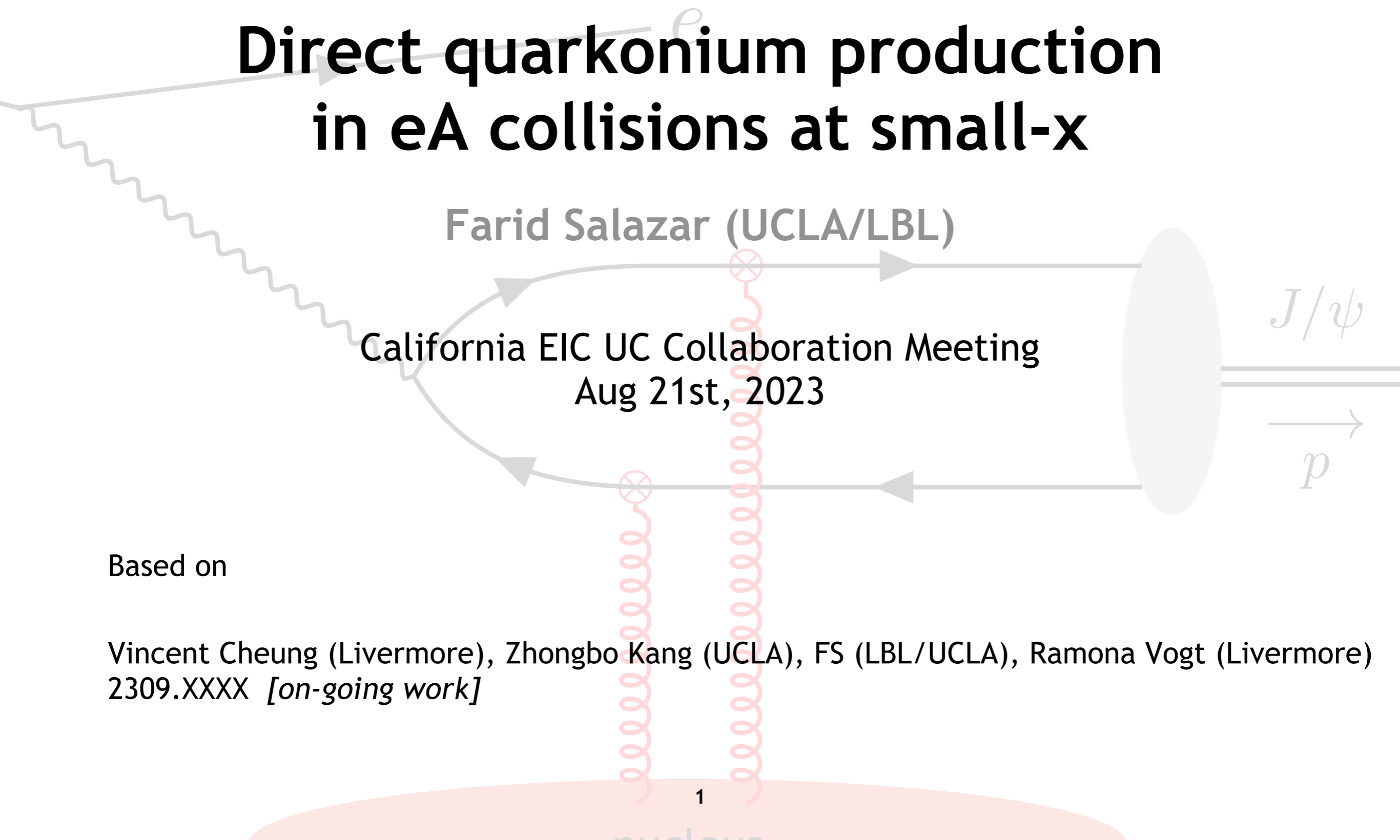


Direct quarkonium production in eA collisions at small-x

Farid Salazar (UCLA/LBL)

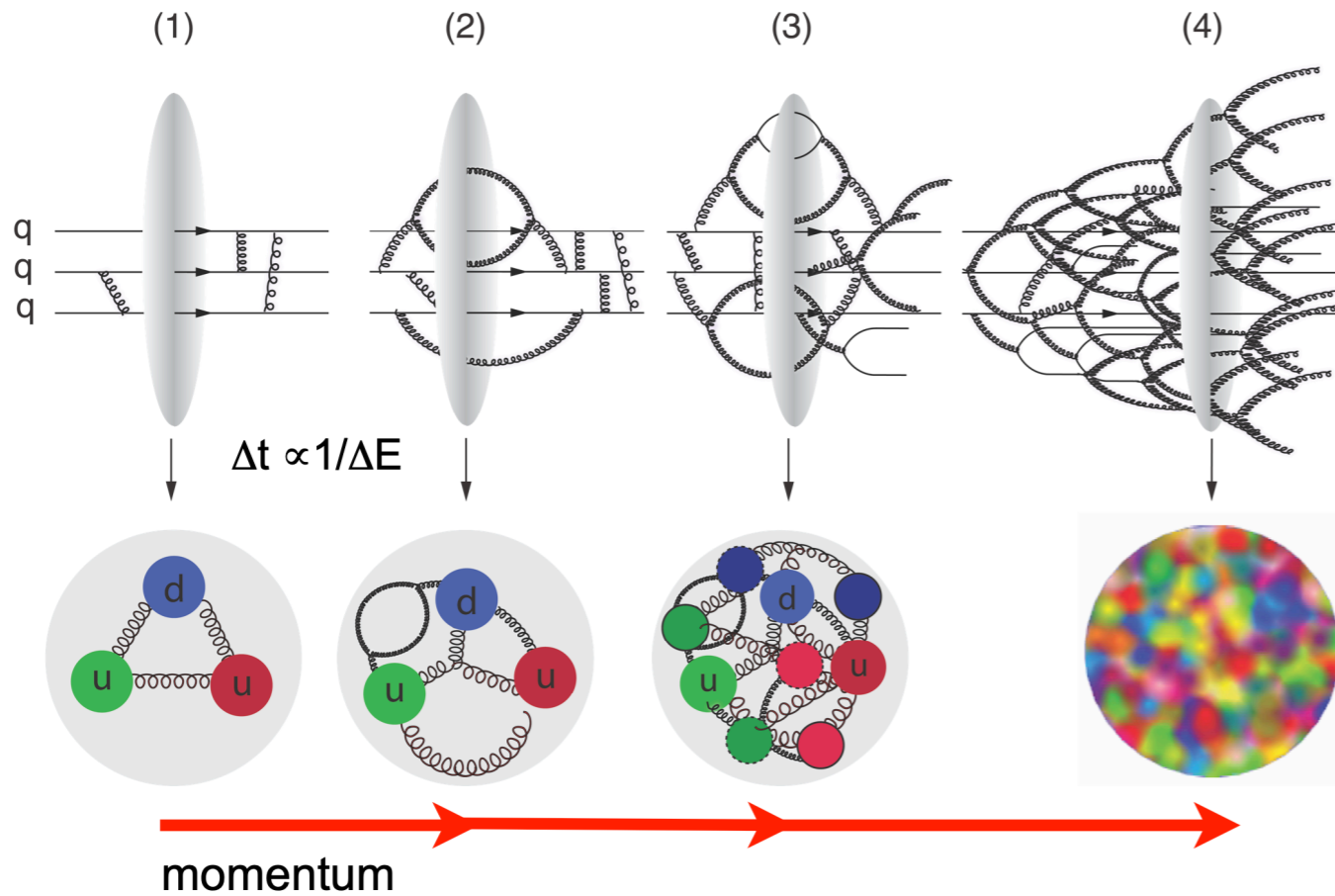
California EIC UC Collaboration Meeting
Aug 21st, 2023



Based on

Vincent Cheung (Livermore), Zhongbo Kang (UCLA), FS (LBL/UCLA), Ramona Vogt (Livermore)
2309.XXXX [on-going work]

Quarkonium as a tool to search for gluon saturation



Gluon occupancy is high at small- x (high-energy) and it saturates

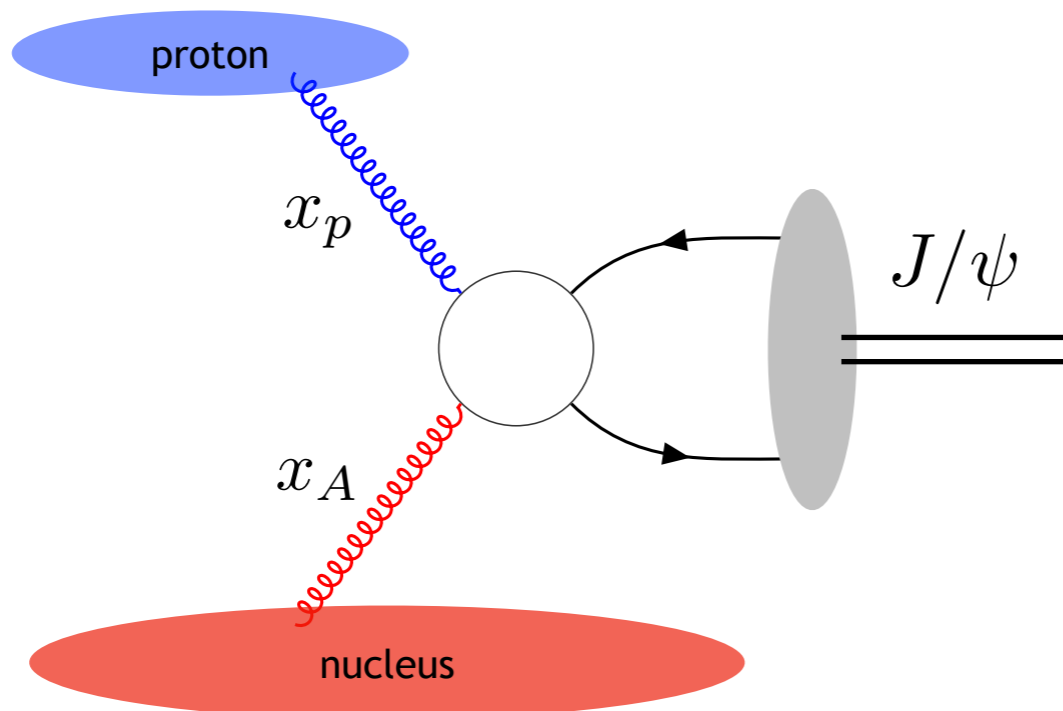


Gluons with transverse momentum $k_{\perp} \lesssim Q_s$ (saturation scale) are suppressed



Imprint on particle production in high-energy collisions at low k_{\perp}

$$Q_s^2 \propto \frac{A^{1/3}}{x^{\lambda}} \sim M_{J/\psi}^2$$



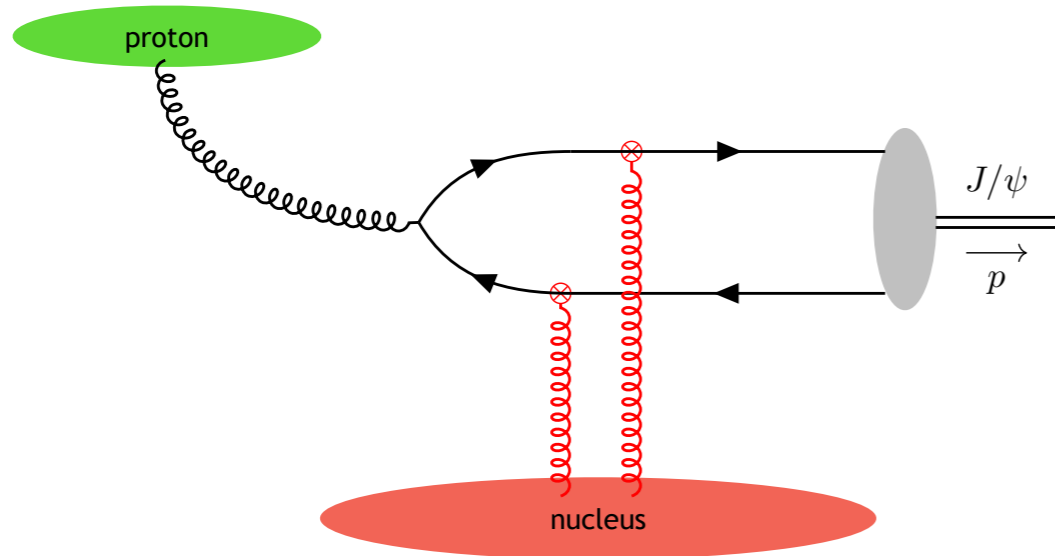
$$x_p = \sqrt{\frac{M_{J/\psi}^2 + P_{\perp}^2}{s}} e^Y \quad x_A = \sqrt{\frac{M_{J/\psi}^2 + P_{\perp}^2}{s}} e^{-Y}$$

Forward production $Y \gg 1$

$$x_A \ll 1$$

Quarkonium production in proton-nucleus collisions

CGC meets NRQCD



Non-relativistic QCD

Non-perturbative LDME

$$\frac{d\sigma^{J/\psi}}{d\mathbf{p}_\perp^2 d\eta} = \sum_{\kappa} \langle \mathcal{O}_{\kappa}^{J/\psi} \rangle \frac{d\hat{\sigma}^{\kappa}}{d\mathbf{p}_\perp^2 d\eta}$$

Decompose contribution into specific quantum state of the heavy quark pair

$$\kappa = 2S+1 L_J^{[c]}$$

S (spin), L (angular momentum), J (total angular momentum), c (color state)

Contributing to J/ψ production: ${}^3S_1^{[1]}, {}^1S_0^{[8]}, {}^3S_1^{[8]}, {}^3P_J^{[8]}$

Short-distance coefficients

$$\frac{d\hat{\sigma}^{\kappa}}{d\mathbf{p}_\perp^2 d\eta} = g(x_p, \mathbf{k}_\perp) \otimes \tilde{\Gamma}^{\kappa}(\mathbf{p}_\perp; \mathbf{l}_\perp, \mathbf{l}'_\perp, \mathbf{k}_\perp) \otimes \tilde{\mathcal{G}}^{\kappa}(x_A, \mathbf{p}_\perp; \mathbf{l}_\perp, \mathbf{l}'_\perp)$$

Proton UGD/TMD Perturbative factor Nuclear-dependent (CGC)

Kang, Ma, Venugopalan, Zhang (JHEP 2013)

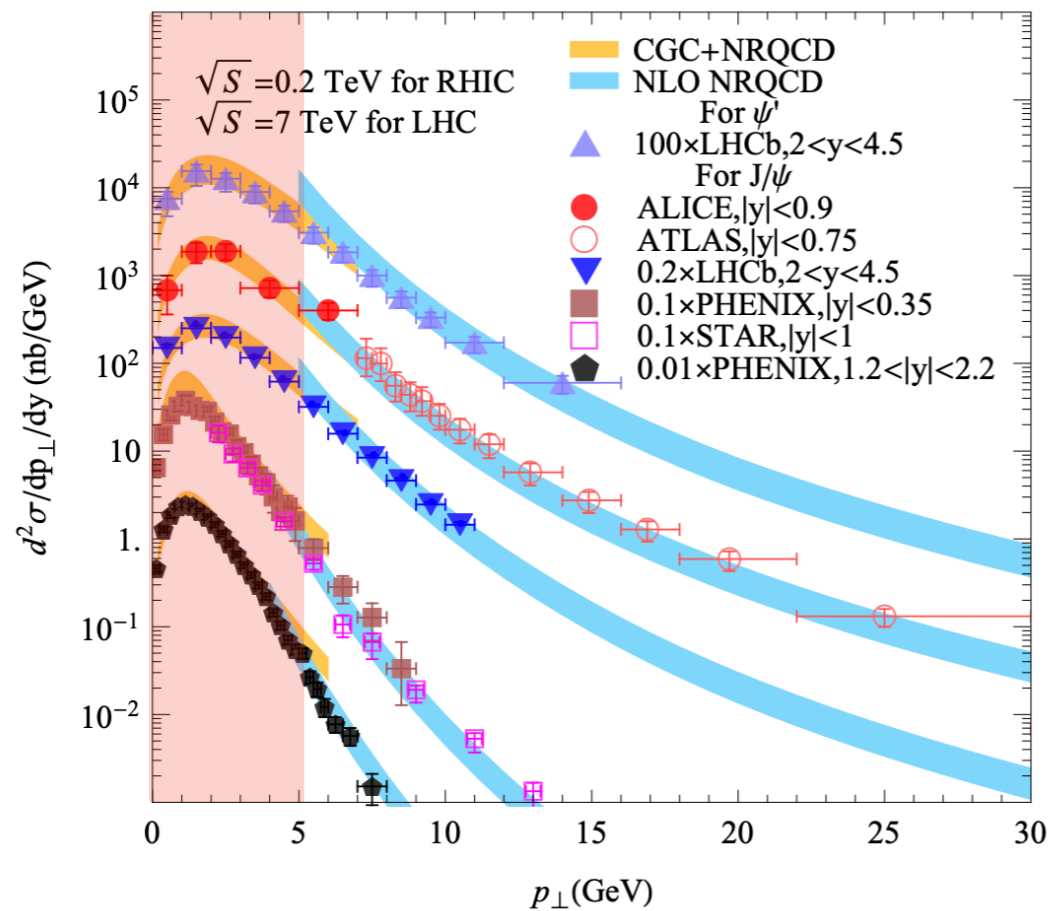
Confronting to RHIC and LHC data

Ma, Venugopalan (PRL 2014)

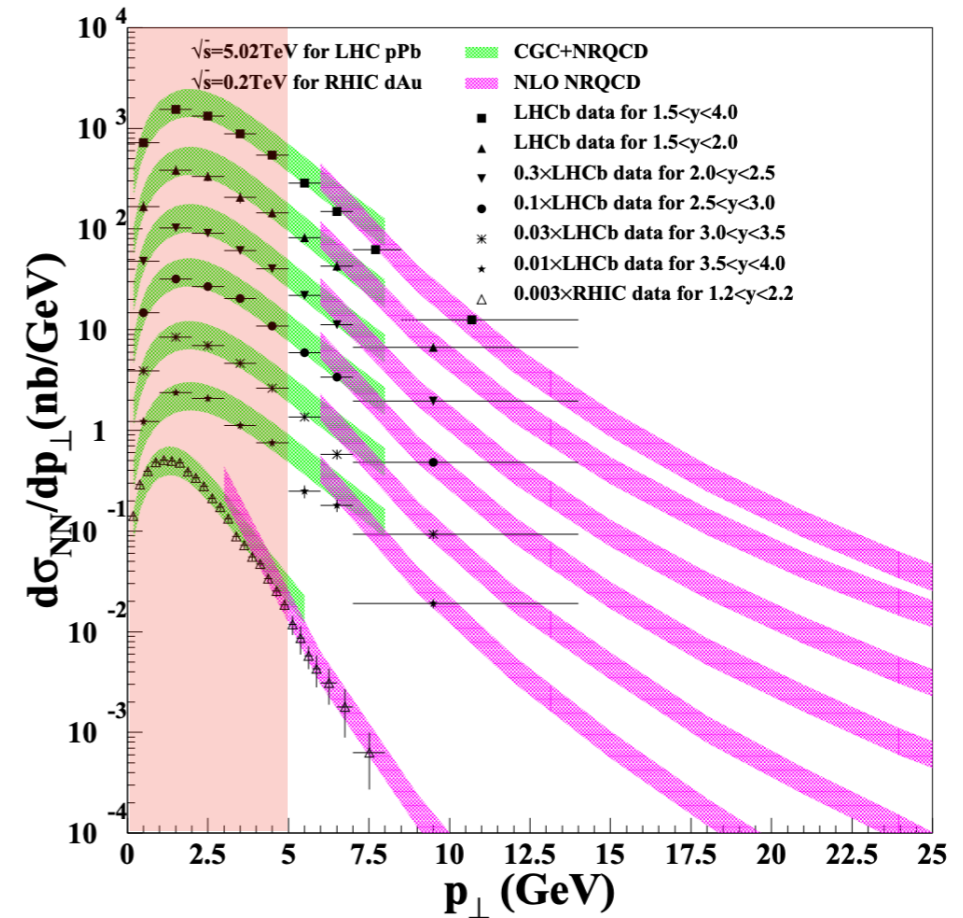
Ma, Venugopalan, Zhang (PRD 2015)

Transverse momentum distribution

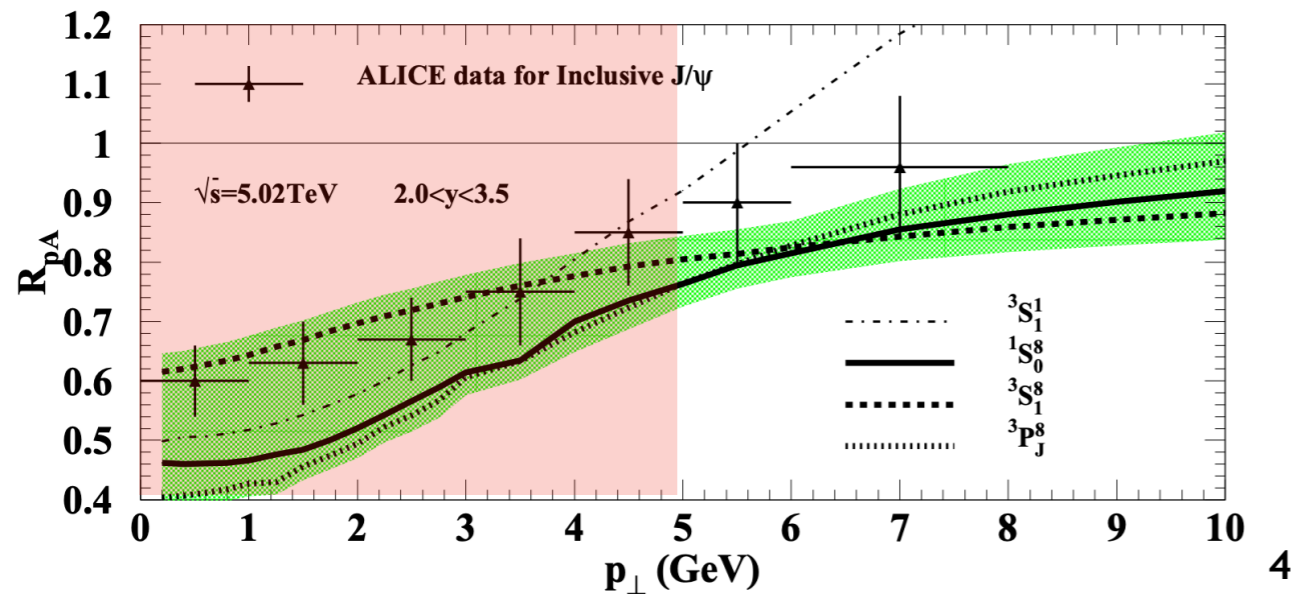
Proton-proton



Proton-nucleus



Nuclear modification Ratio



CGC provides good description of experimental data at low p_T ($p_{\perp} \lesssim Q_s$)

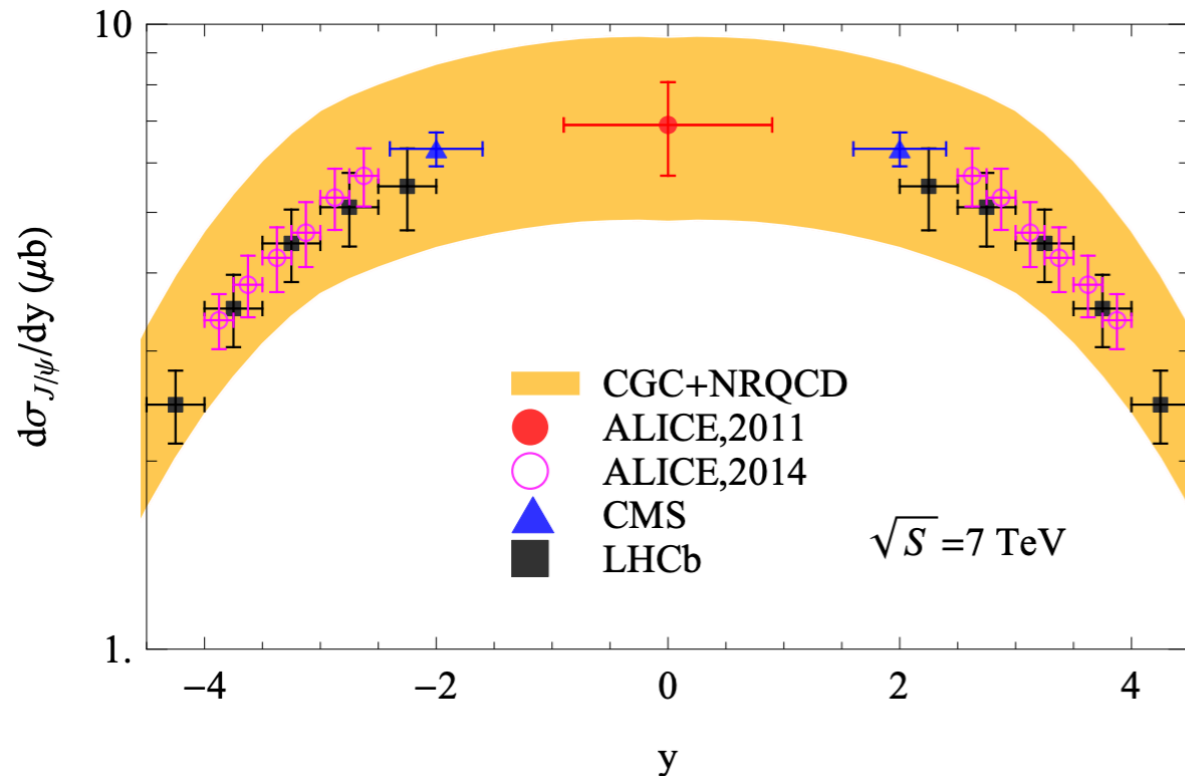
Confronting to RHIC and LHC data

Ma, Venugopalan (PRL 2014)

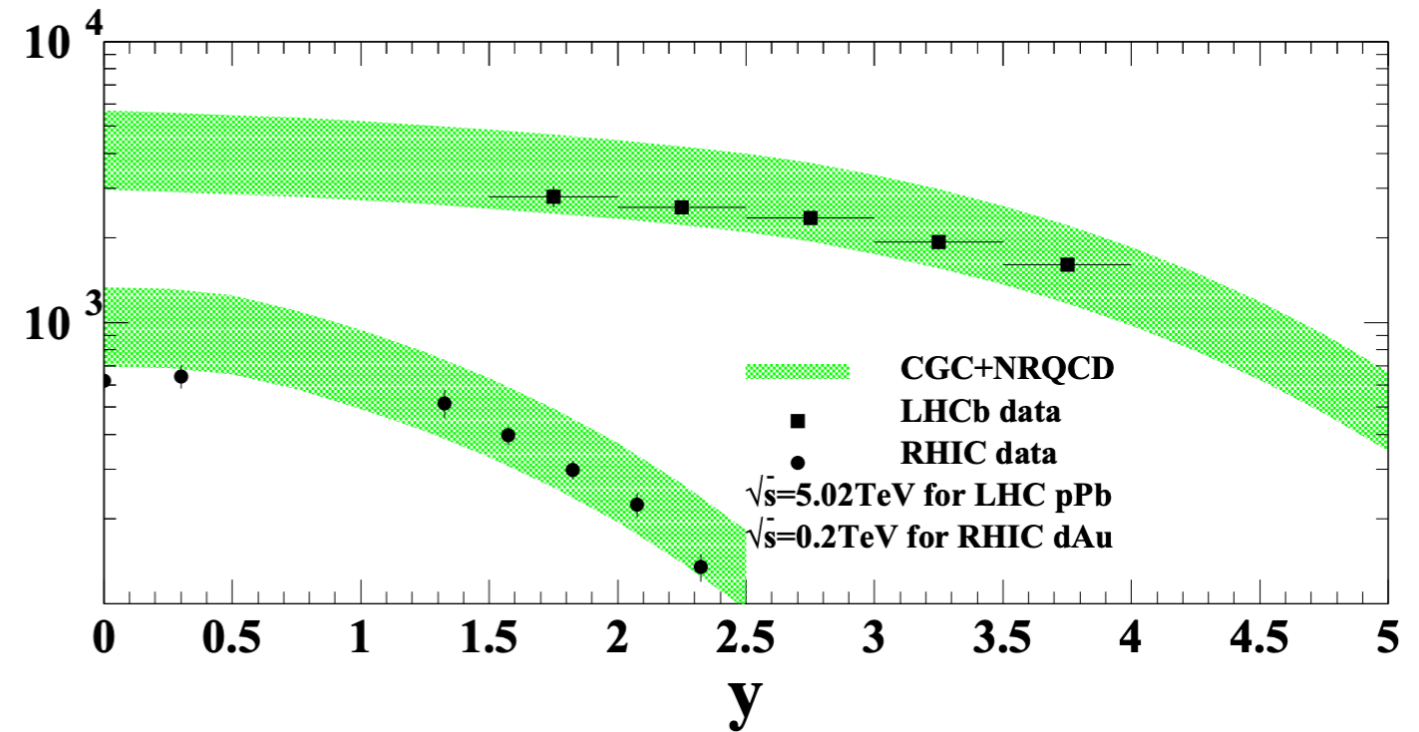
Ma, Venugopalan, Zhang (PRD 2015)

Rapidity distribution and nuclear modification

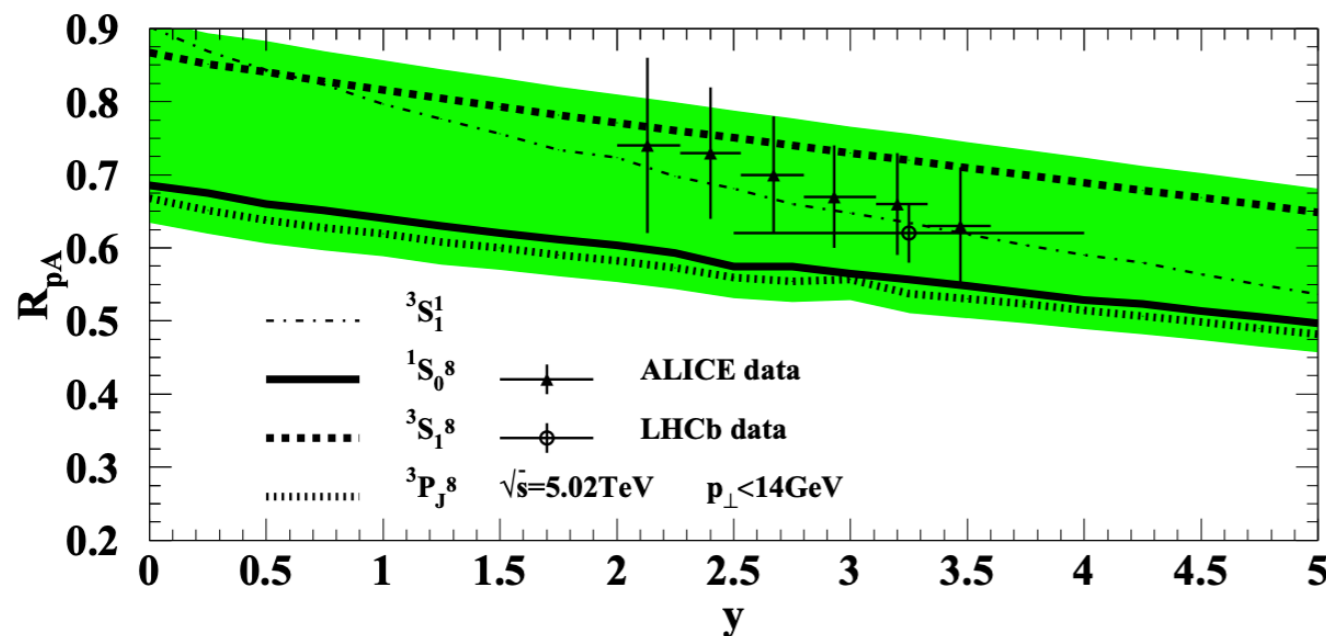
Proton-proton



Proton-nucleus



Nuclear modification Ratio

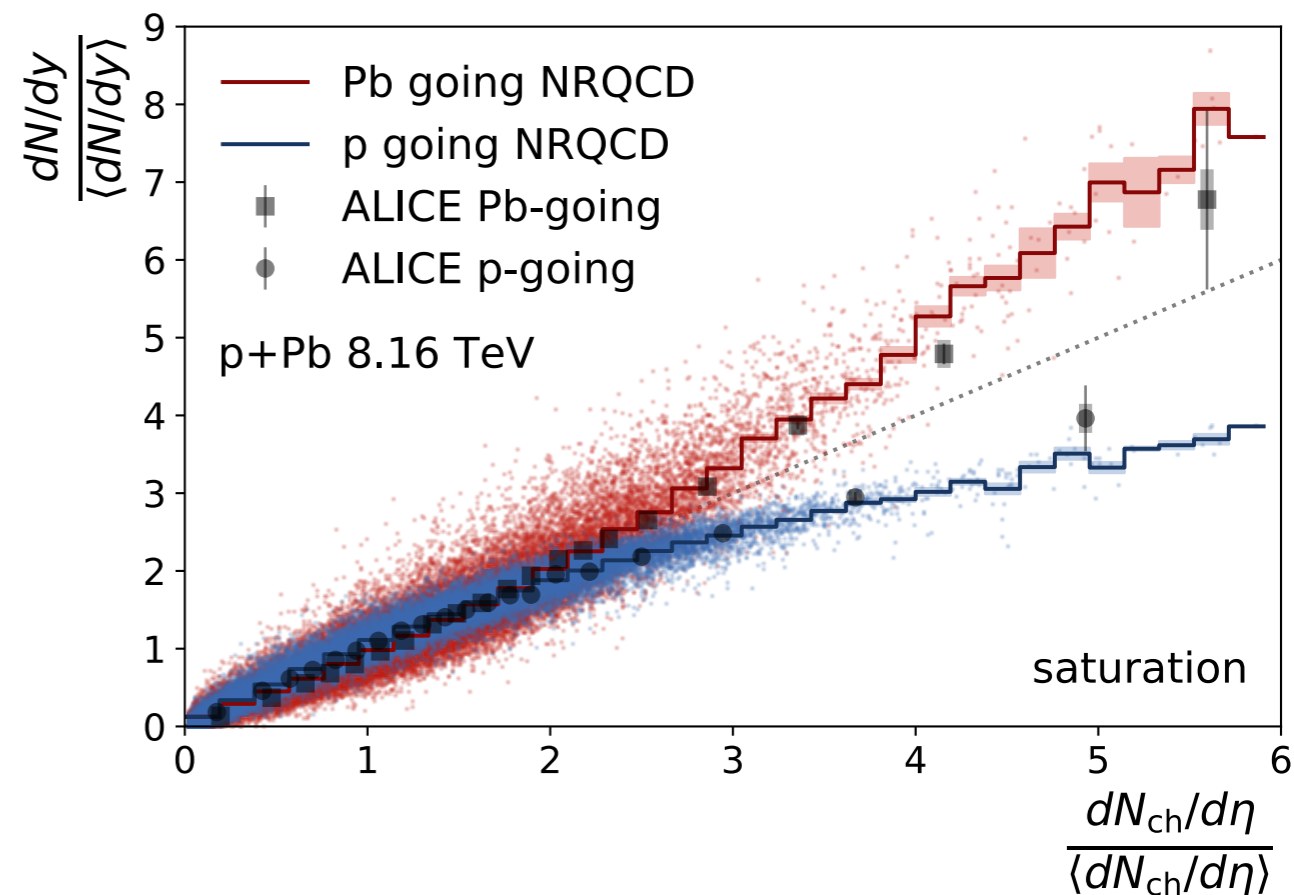
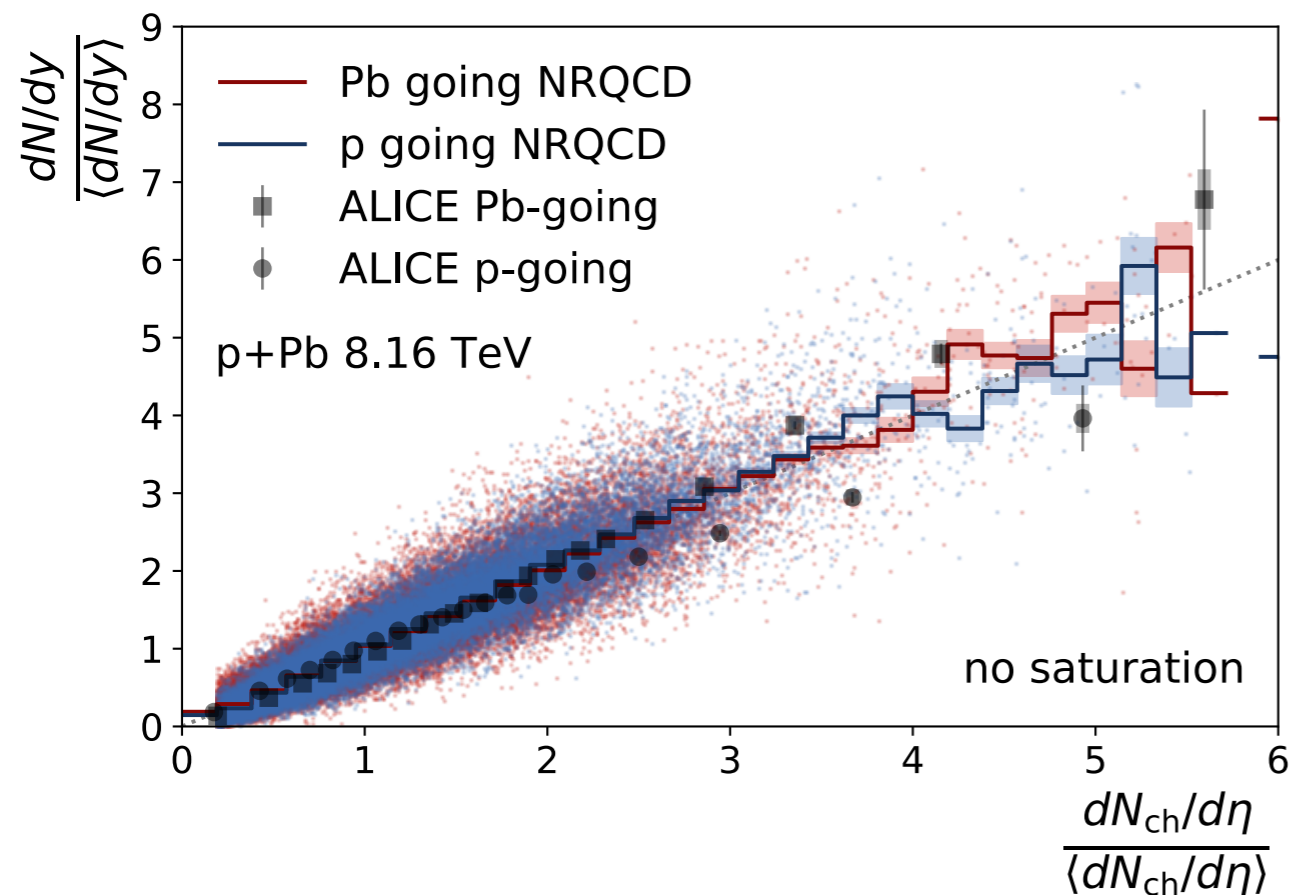


Rapidity distributions are integrated over p_{\perp} , low p_{\perp} dominates the bulk of the cross-section

Confronting to RHIC and LHC data

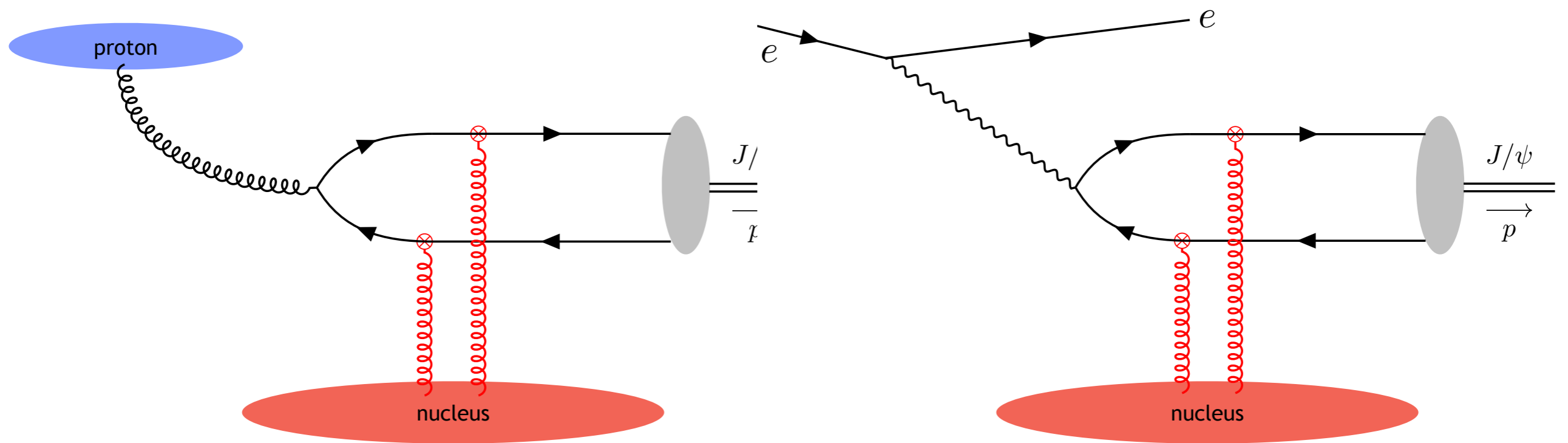
J/ψ multiplicity vs charged hadron multiplicity

FS, Schenke, Soto-Ontoso [PLB 2022]



Sub-nuclear fluctuations in hotspots size and saturation scale provide a natural framework to generate different multiplicity classes that describe well LHC data

What about electron-nucleus deep inelastic scattering?



Replace the proton projectile by an electron

Reconstruct kinematics of the “projectile” photon
Electromagnetic probe -> cleaner theoretical calculation
Possible to measure at the future EIC

Surprisingly, calculations [in CGC] hadn't been done yet, not even at LO...

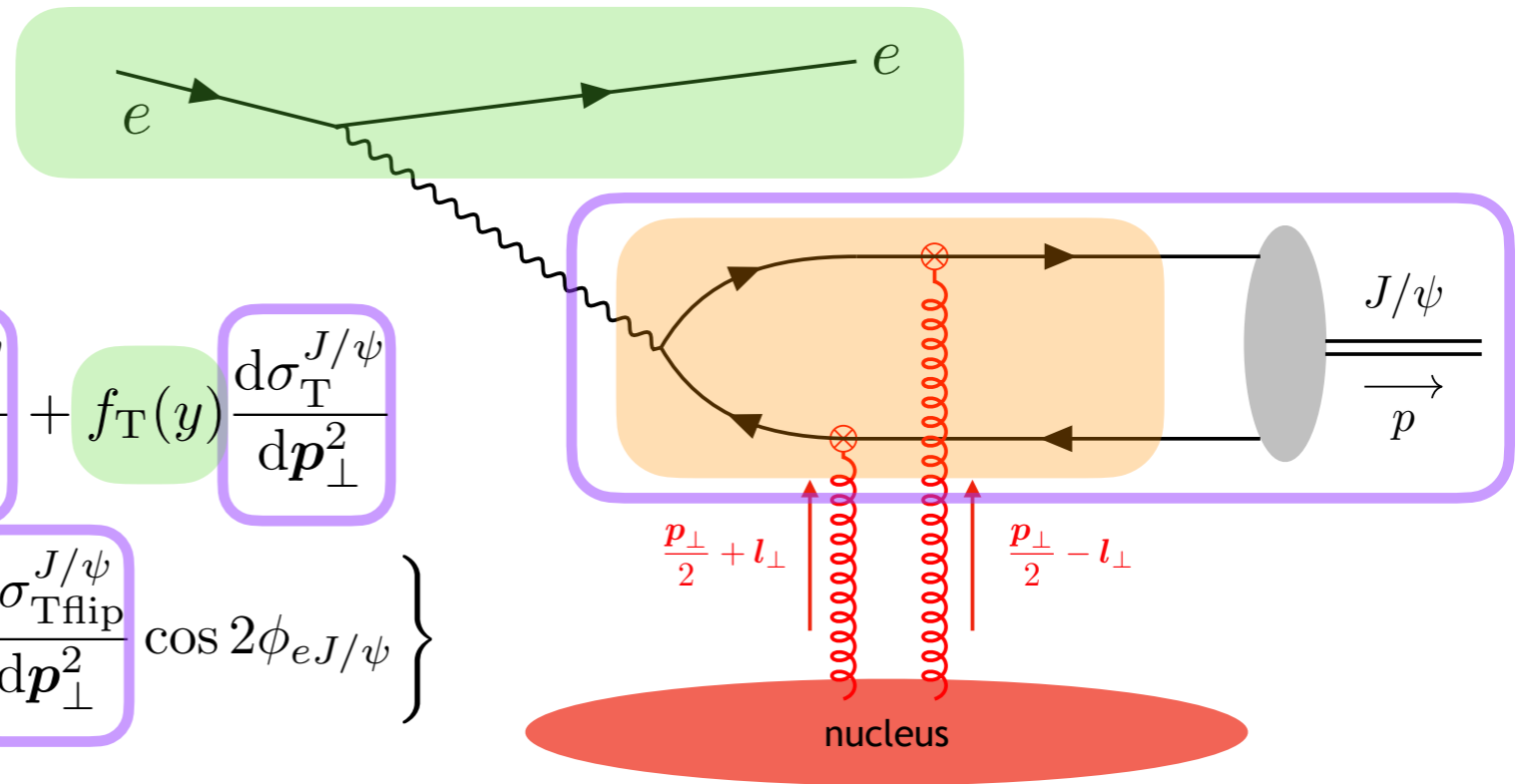
Motivated our work in Cheung, Kang, FS, Vogt [2309.XXXX]

Quarkonium production in electron-nucleus collisions

CGC meets NRQCD

Direct J/ψ production in eA collision

$$\frac{d\sigma^{J/\psi}}{dx_{Bj} dy d\mathbf{p}_\perp^2 d\phi_{eJ/\psi}} = \frac{\alpha_{em}}{2\pi^2 y x_{Bj}} \left\{ f_L(y) \frac{d\sigma_L^{J/\psi}}{d\mathbf{p}_\perp^2} + f_T(y) \frac{d\sigma_T^{J/\psi}}{d\mathbf{p}_\perp^2} \right. \\ \left. + f_{TL}(y) \frac{d\sigma_{TL}^{J/\psi}}{d\mathbf{p}_\perp^2} \cos \phi_{eJ/\psi} + f_{Tflip}(y) \frac{d\sigma_{Tflip}^{J/\psi}}{d\mathbf{p}_\perp^2} \cos 2\phi_{eJ/\psi} \right\}$$



Direct J/ψ production in $\gamma^* A$ collision

$$\frac{d\sigma_\lambda^{J/\psi}}{d\mathbf{p}_\perp^2} = \sum_\kappa \langle \mathcal{O}_\kappa^{J/\psi} \rangle \frac{d\hat{\sigma}_\lambda^\kappa}{d\mathbf{p}_\perp^2} \quad \left(\frac{d\hat{\sigma}_\lambda^\kappa}{d\mathbf{p}_\perp^2} = \int \frac{d^2\mathbf{l}_\perp}{2\pi} \int \frac{d^2\mathbf{l}'_\perp}{2\pi} \tilde{\Gamma}_\lambda^\kappa(\mathbf{p}_\perp, Q; \mathbf{l}_\perp, \mathbf{l}'_\perp) \tilde{\mathcal{G}}_Y^\kappa(\mathbf{p}_\perp; \mathbf{l}_\perp, \mathbf{l}'_\perp) \right)$$

LDME for J/ψ production

Short-distance coefficients

Spin and polarization-dependent perturbative factor (20 functions)

Nuclear-dependent CGC distribution (Octet and Singlet)

Convolution manifestly breaks kT factorization

Quarkonium production in electron-nucleus collisions

kT factorization and the Weizsäcker Williams gluon distribution

Our result for Short-distance coefficients in CGC + NRQCD

$$\frac{d\hat{\sigma}_\lambda^\kappa}{d\mathbf{p}_\perp^2} = \int \frac{d^2\mathbf{l}_\perp}{2\pi} \int \frac{d^2\mathbf{l}'_\perp}{2\pi} \tilde{\Gamma}_\lambda^\kappa(\mathbf{p}_\perp, Q; \mathbf{l}_\perp, \mathbf{l}'_\perp) \tilde{\mathcal{G}}_Y^\kappa(\mathbf{p}_\perp; \mathbf{l}_\perp, \mathbf{l}'_\perp)$$

Improved TMD limit

$$Q_s^2 \ll Q^2 + M_{J/\psi}^2$$

Expansion following
Boussarie, Mehtar-Tani
(PRD 2021)

$$\frac{d\hat{\sigma}_\lambda^\kappa}{d\mathbf{p}_\perp^2} = \mathcal{H}_{\lambda, \alpha\alpha'}^\kappa(Q, \mathbf{p}_\perp) xG^{\alpha\alpha'}(x, \mathbf{p}_\perp)$$

Satisfies (generalized)
kT factorization

TMD limit

$$Q_s^2 \ll Q^2 + M_{J/\psi}^2 \quad \text{and} \quad \mathbf{p}_\perp^2 \ll Q^2 + M_{J/\psi}^2$$

Reproduces results by
Bacchetta, Boer, Pisano and Tael
(EPJC 2018) within TMD factorization

$$\frac{d\hat{\sigma}_\lambda^\kappa}{d\mathbf{p}_\perp^2} = H_{\lambda, \alpha\alpha'}^\kappa(Q) xG^{\alpha\alpha'}(x, \mathbf{p}_\perp)$$

$$xG^{\alpha\alpha'}(x, \mathbf{p}_\perp)$$

Weizsäcker-Williams
UGD or TMD at small-x

See e.g. Dominguez,
Marquet, Xiao, Yuan
(PRD 2013)

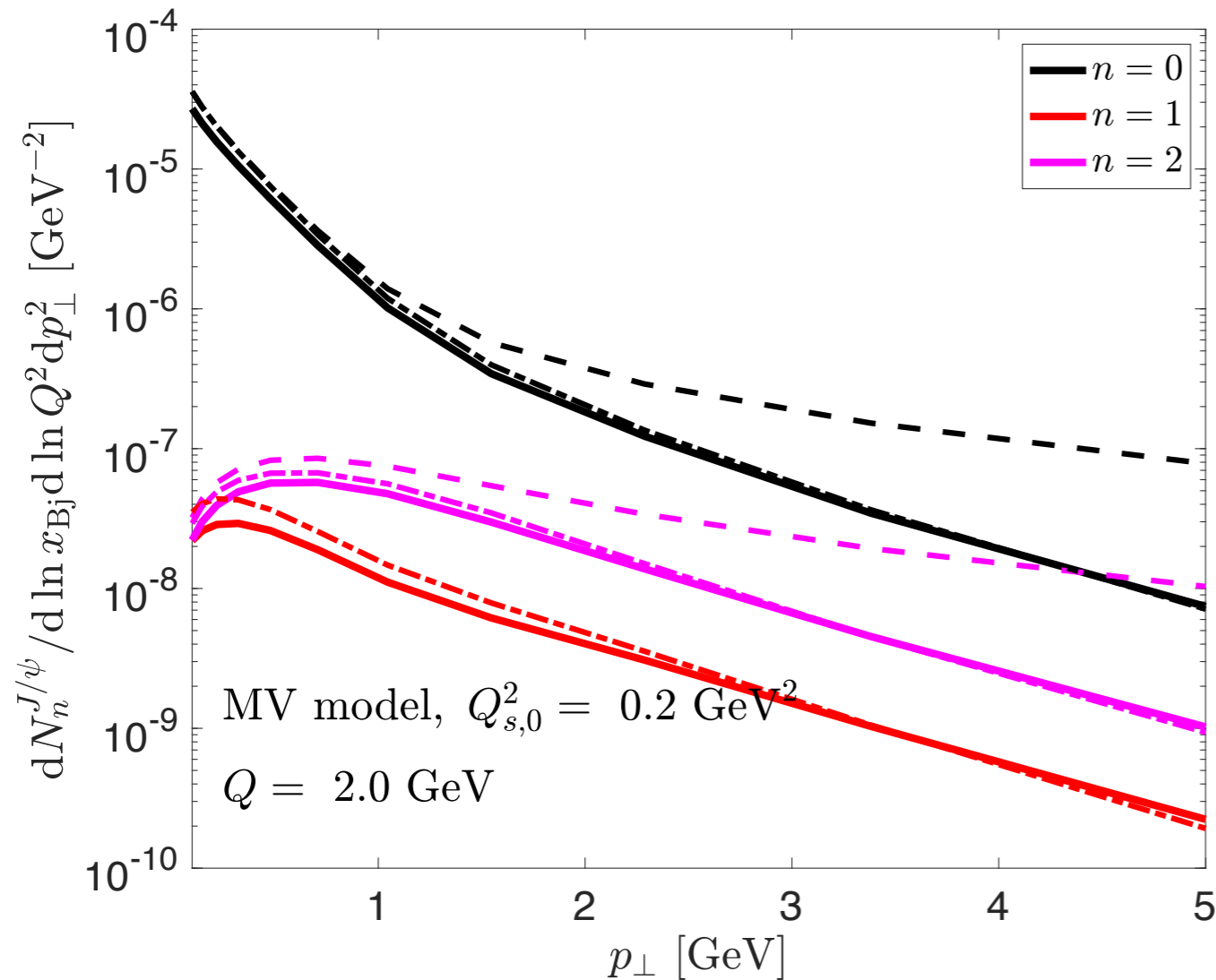
kT factorization breaking grows with saturation scale Q_s^2
(i.e. larger nuclei or larger energies)

Quarkonium production in electron-nucleus collisions

pT dependence: CGC vs kT factorization (TMD and ITMD)

$$dN_n = \frac{1}{S_\perp} \int \frac{d\phi_{eJ/\psi}}{2\pi} d\sigma^{J/\psi} \cos(n\phi_{eJ/\psi})$$

$$e + p \rightarrow e + J/\psi + X$$



Solid line= full CGC
Dashed-solid= Improved TMD
Dashed = TMD

kT factorization (ITMD)

$\ln(Q_s^2/p_\perp^2)$ for $p_\perp^2 \lesssim Q_s^2$ TMD saturated

$1/p_\perp^2$ for $Q_s^2 \ll p_\perp^2 \ll Q^2 + M_{J/\psi}^2$ TMD

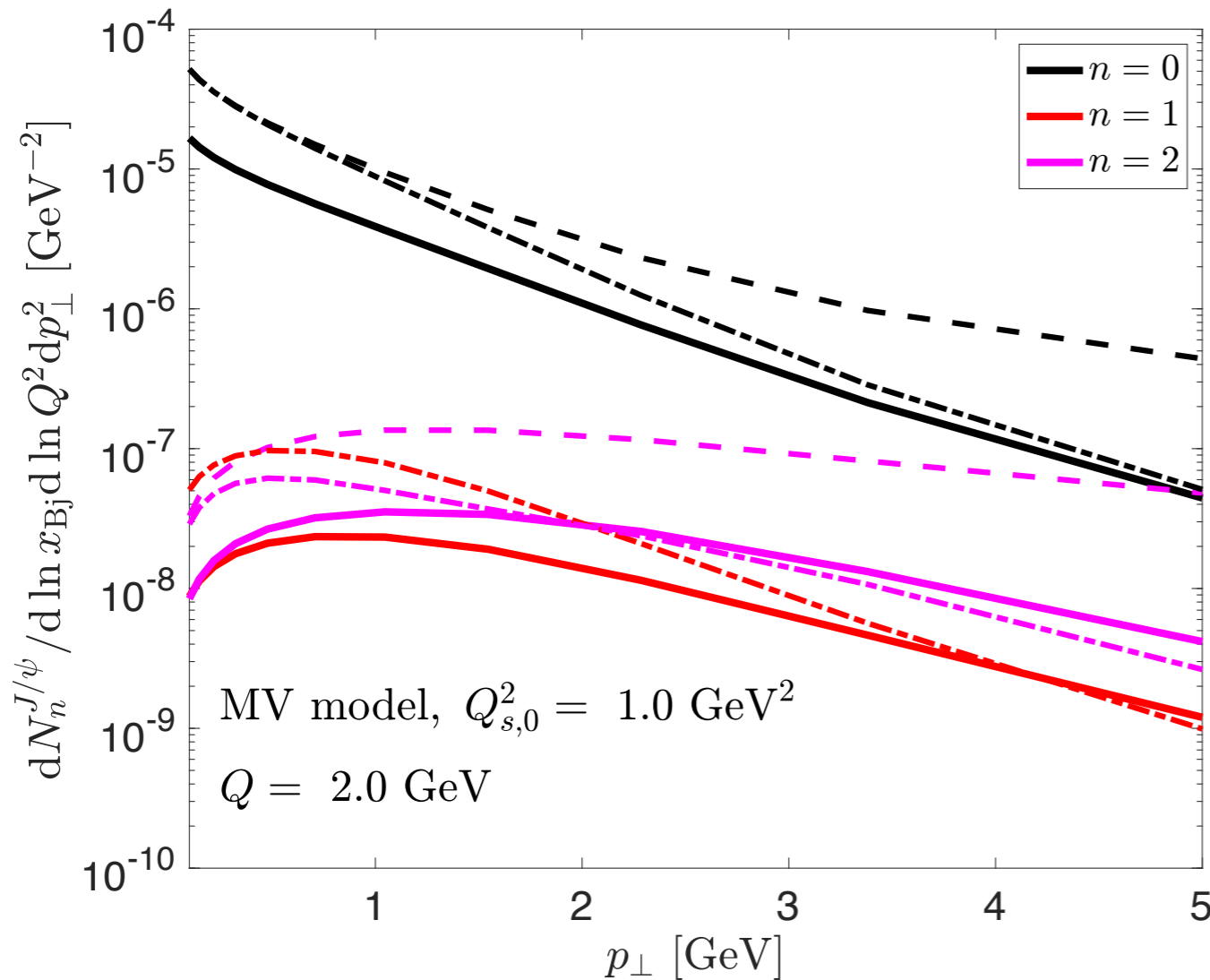
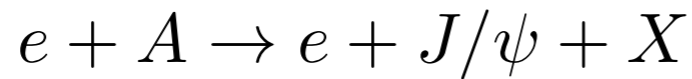
$1/p_\perp^4$ for $p_\perp \gg Q^2 + M_{J/\psi}^2$ collinear

Small saturation scale ->
little factorization kT breaking

Quarkonium production in electron-nucleus collisions

pT dependence: CGC vs kT factorization (TMD and ITMD)

$$dN_n = \frac{1}{S_\perp} \int \frac{d\phi_{eJ/\psi}}{2\pi} d\sigma^{J/\psi} \cos(n\phi_{eJ/\psi})$$



Solid line= full CGC
 Dashed-solid= Improved TMD
 Dashed = TMD

kT factorization (ITMD)

$\ln(Q_s^2/p_\perp^2)$ for $p_\perp^2 \lesssim Q_s^2$ TMD saturated

$1/p_\perp^2$ for $Q_s^2 \ll p_\perp^2 \ll Q^2 + M_{J/\psi}^2$ TMD

$1/p_\perp^4$ for $p_\perp \gg Q^2 + M_{J/\psi}^2$ collinear

Full CGC

Further suppression at low p_\perp
 when $Q_s^2 \sim Q^2 + M_{J/\psi}^2$ due to kT
 factorization breaking

Summary and Outlook

- Past

CGC + NRQCD provides good descriptions of rapidity and p_{\perp} distribution, nuclear modification ratio, and multiplicity-dependence in high-energy pp and pA collisions at RHIC and LHC

- Present

We computed direct quarkonium production in ep and eA, and carried out a numerical study focusing on deviations from kT factorization

- Future

Can we provide a good description of HERA data? UPCs?

Predictions for the EIC

Study polarized J/ψ production

Extend our results to NLO [ZK, FS and Emilie Li on-going work]