Probing energy, intensity and cosmic frontiers through Z′

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Advancing searches for New Physics
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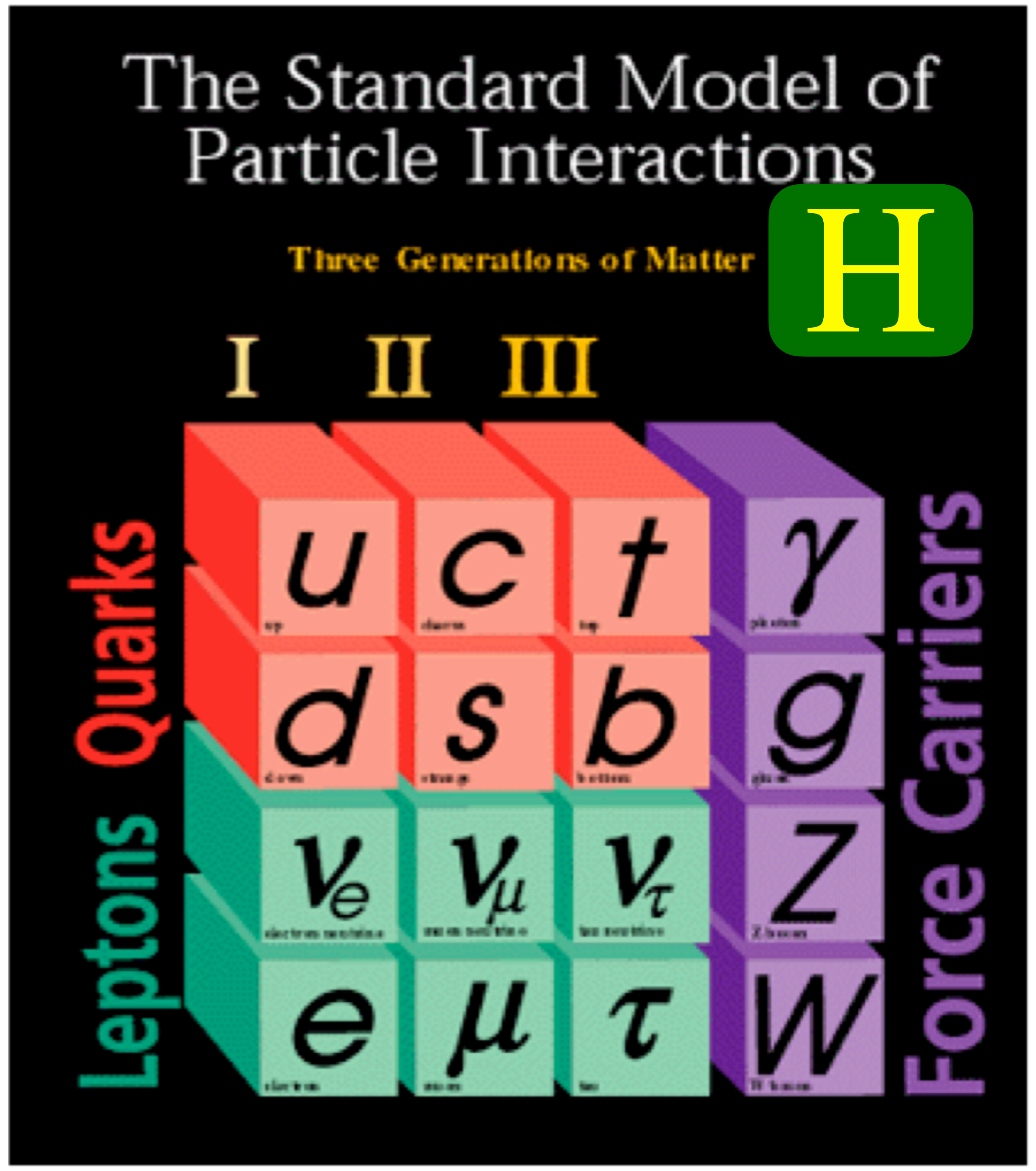
Introduction



Over the decades experiments have found each and every missing pieces

Verified the facts that they belong to this family

Finally at the Large Hadron collider Higgs has been observed

→ Its properties must be verified

Strongly established with interesting shortcomings

Few of the very interesting anomalies :

Tiny neutrino mass and flavor mixings

Relic abundance of dark matter . . .

SM can not explain them

Birth of a new idea : generation of neutrino mass

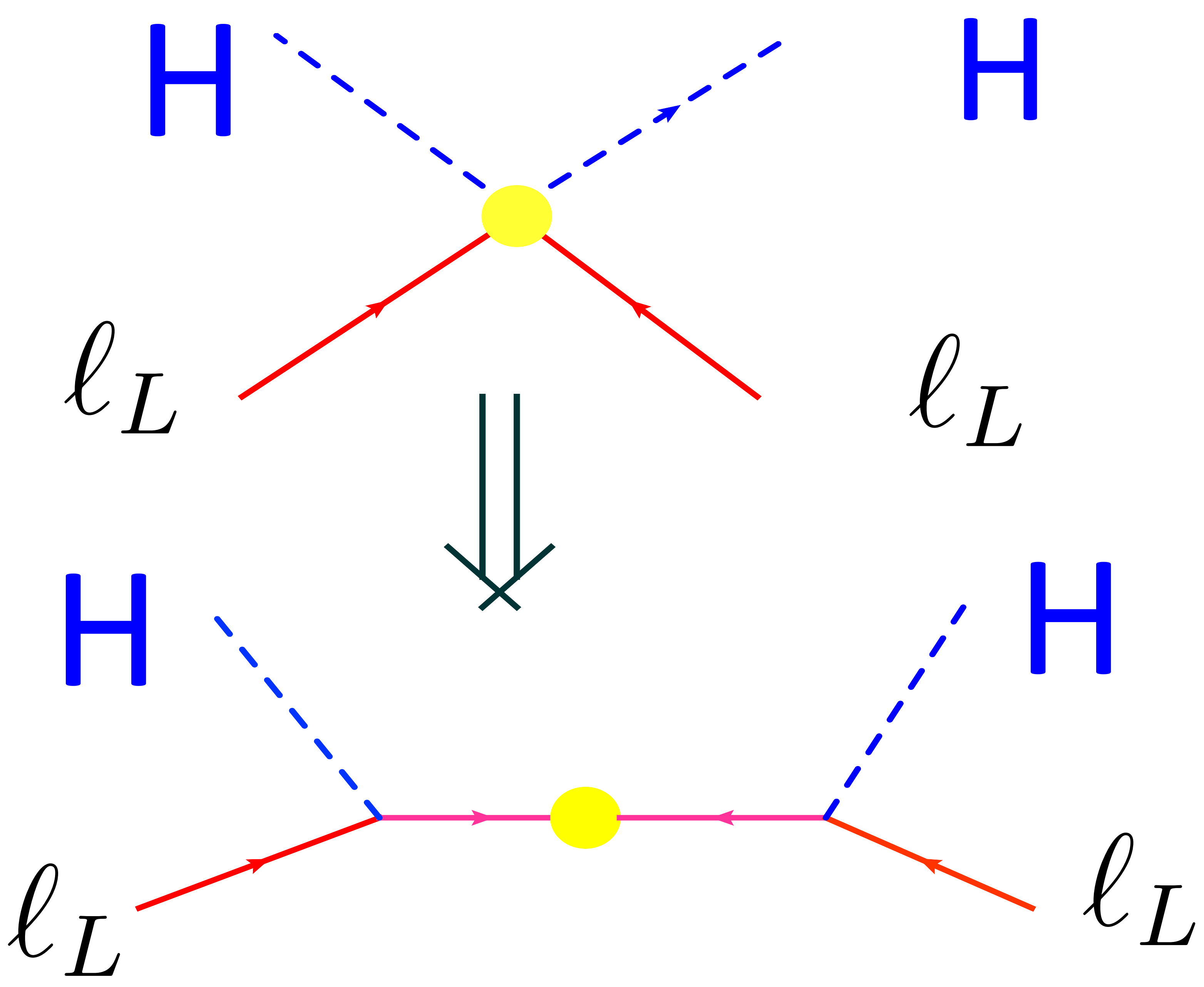
Weinberg Operator in SM (d=5), PRL 43, 1566(1979)

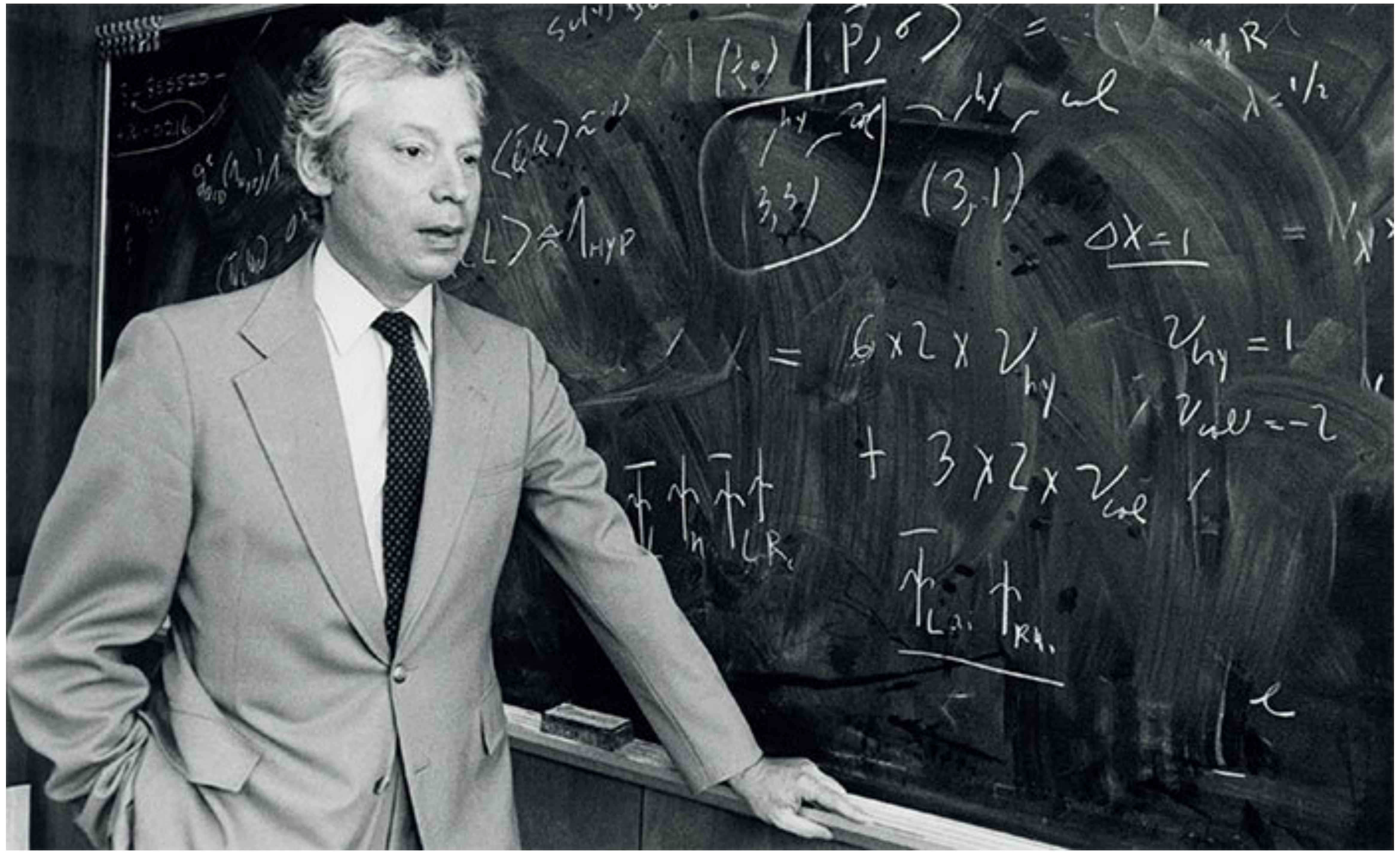
$$\frac{\overline{\ell_L} H \overline{\ell_L^c}^T H}{M}$$

within the Standard Model

Steven Weinberg : 1933 – 2021

The dimension 5 operator can be realized in the following ways





Majorana mass term is generated by the breaking of the lepton numbers by 2 units.

Scenarios

Neutrino mass
Tree level
SM + Particles
Gauge extension
Singlet, triplet fermions
Triplet scalars
Left – right
general U(1)
Seesaw, inverse seesaw
I, II, III
Quantum level
I, II, III . . .
Higher dimensional operators

Different frontiers

Energy frontier : Scientists build partcile acclerators to explore high energy scale to explore new phenomena after the subatomic collisions .

Intensity frontier : Highly intense beams from accelerators are used to to investigate the ultra rare processes of nature .

Cosmic frontier : Astrophysicists use the cosmos as the laboratory to investigate the fundamental laws of physics from a complementary point of view of particle accelerator .

Particle content

Dobrescu, Fox; Cox, Han, Yanagida; AD, Okada, Raut;
Chiang, Cottin, AD, Mandal; AD, Takahashi, Oda, Okada AD, Dev, Okada

	SU(3)$_c$	SU(2)$_L$	U(1)$_Y$			$U(1)_X$
q_L^i	**3**	**2**	$+1/6$	x_q	$=$	$\frac{1}{6}x_H + \frac{1}{3}x_\Phi$
u_R^i	**3**	**1**	$+2/3$	x_u	$=$	$\frac{2}{3}x_H + \frac{1}{3}x_\Phi$
d_R^i	**3**	**1**	$-1/3$	x_d	$=$	$-\frac{1}{3}x_H + \frac{1}{3}x_\Phi$
ℓ_L^i	**1**	**2**	$-1/2$	x_ℓ	$=$	$-\frac{1}{2}x_H - x_\Phi$
e_R^i	**1**	**1**	-1	x_e	$=$	$-x_H - x_\Phi$
H	**1**	**2**	$+1/2$	x'_H	$=$	$\frac{1}{2}x_H$
N_R^i	**1**	**1**	0	x_ν	$=$	$-x_\Phi$
Φ	**1**	**1**	0	x'_Φ	$=$	$2x_\Phi$

$m_{Z'} = 2\, g_X v_\Phi$

x_H, x_Φ will appear the coupling with Z'

B − L case

$x_H = 0,\ x_\Phi = 1$

3 generations of SM singlet right handed neutrinos (anomaly free)

Charges before the anomaly cancellations

Charges after Imposing the anomaly cancellations

$U(1)_X$ breaking

$$\mathcal{L}_Y \supset -\sum_{i,j=1}^{3} Y_D^{ij}\overline{\ell_L^i} H N_R^j - \frac{1}{2}\sum_{i=k}^{3} Y_N^k \Phi \overline{N_R^{k\,c}} N_R^k + \text{h.c.},$$

$$m_D^{ij} = \frac{Y_D^{ij}}{\sqrt{2}} v_h$$

$$m_{N^i} = \frac{Y_N^i}{\sqrt{2}} v_\Phi$$

$$m_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \quad m_\nu \simeq -M_D M_N^{-1} M_D^T$$

Seesaw mechnism

Z' interactions

Interaction between the quarks and Z' $$\mathcal{L}^q = -g'(\overline{q}\gamma_\mu q^q_{x_L} P_L q + \overline{q}\gamma_\mu q^q_{x_R} P_R q) Z'_\mu$$

Interaction between the leptons and Z' $$\mathcal{L}^\ell = -g'(\overline{\ell}\gamma_\mu q^\ell_{x_L} P_L \ell + \overline{e}\gamma_\mu q^\ell_{x_R} P_R e) Z'_\mu$$

$q^f_{x_L} \neq q^f_{x_R}$ affects the phenomenology

Partial decay width

Charged fermions $$\Gamma(Z' \to 2f) = N_c \frac{M_{Z'}}{24\pi} \left(g^f_L\left[g', x_H, x_\Phi\right]^2 + g^f_R\left[g', x_H, x_\Phi\right]^2 \right)$$

light neutrinos $$\Gamma(Z' \to 2\nu) = \frac{M_{Z'}}{24\pi}\, g^\nu_L\left[g', x_H, x_\Phi\right]^2$$

heavy neutrinos $$\Gamma(Z' \to 2N) = \frac{M_{Z'}}{24\pi}\, g^N_R\left[g', x_\Phi\right]^2 \left(1 - 4\frac{m_N^2}{M_{Z'}^2}\right)^{\frac{3}{2}}$$

Properties of the model and phenomenology

New particles

Z' boson   Heavy Majorana Neutrino(RHN)

$U(1)_X$ Higgs boson

Phenomenology

Z' boson production and decay

Z' boson mediated processes   Heavy neutrino production

$U(1)_X$Higgs phenoemenology : Vacuum Stability

BSM scalar production, decay into RHN pair

Dark Matter   Collider phenomenology

Dev, Pilaftsis; Iso, Okada, Orikasa
Orikasa, Okada, Yamada; Dev, Mohapatra, Zhang

Leptogenesis and many more

Z' boson production and heavy neutrino phenomenology

$\nu-\mathcal{N}$, $\nu-$ e, e^-e^+ scattering; proton, electron beam dump and dark photon search

2307.09737

Fermionic pair production form the Z' : $\mathcal{A}_{\rm FB,\ LR,\ FB-LR}$

Bhabha scattering

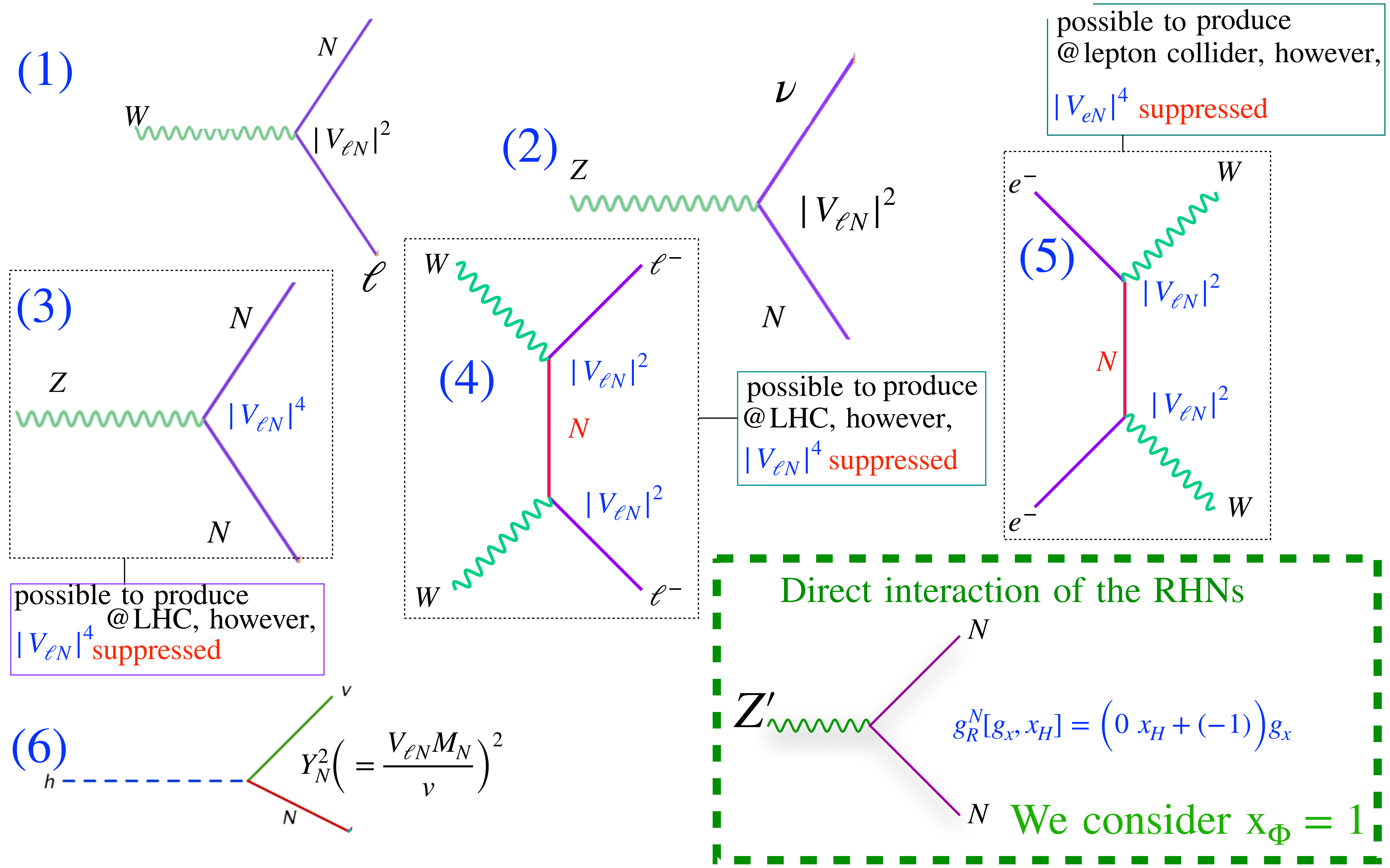
2104.10902

Production modes of the RHNs at the colliders : pp, e^-e^+, e^-p

Flavor eigenstate can be expressed in terms of the mass eigenstate

$$\nu_\ell \simeq U_{\ell m}\nu_m + V_{\ell n}N_n$$

PMNS matrix ($U_{\ell m}$) $M_D M_N^{-1}$ ($V_{\ell n}$)

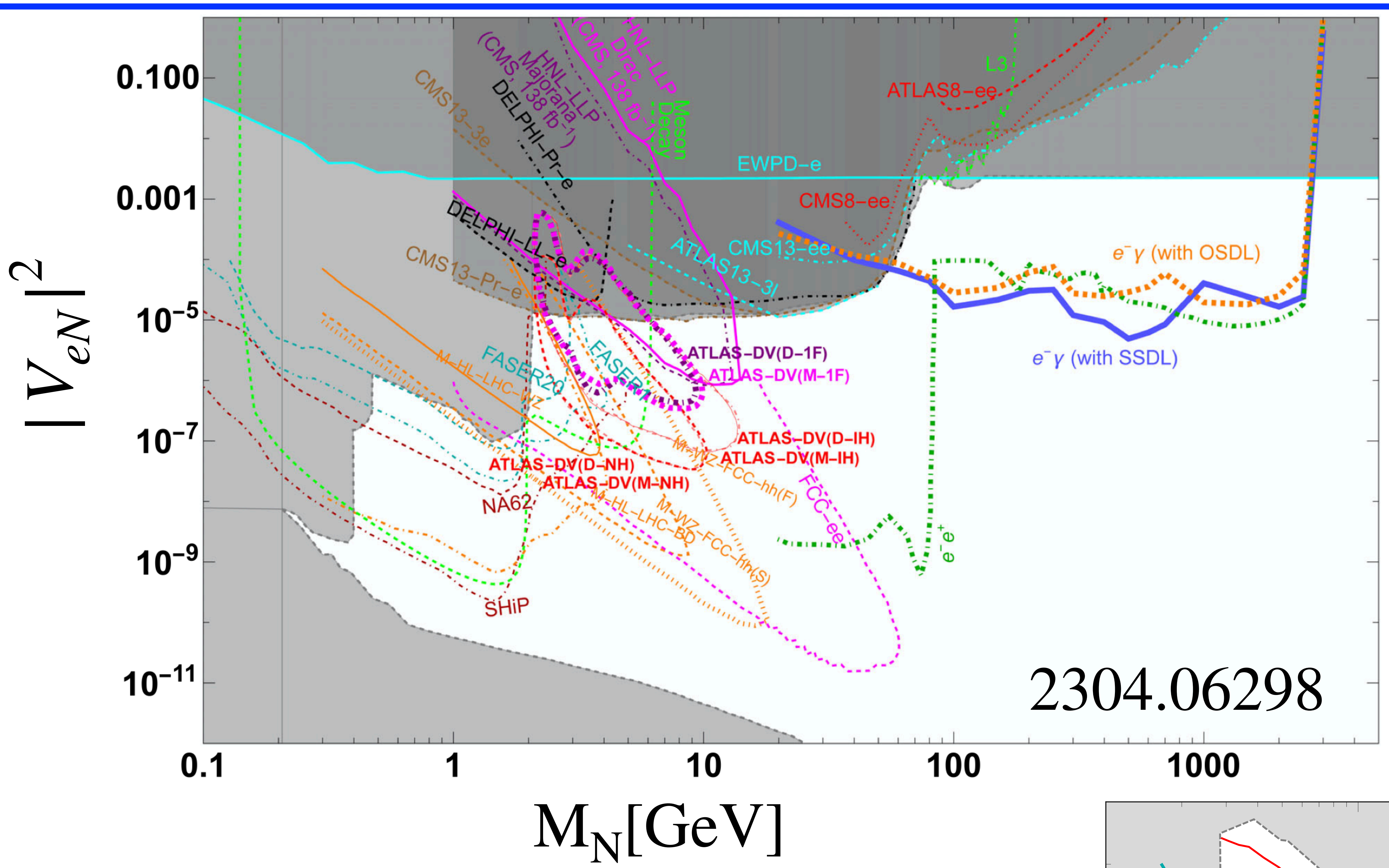


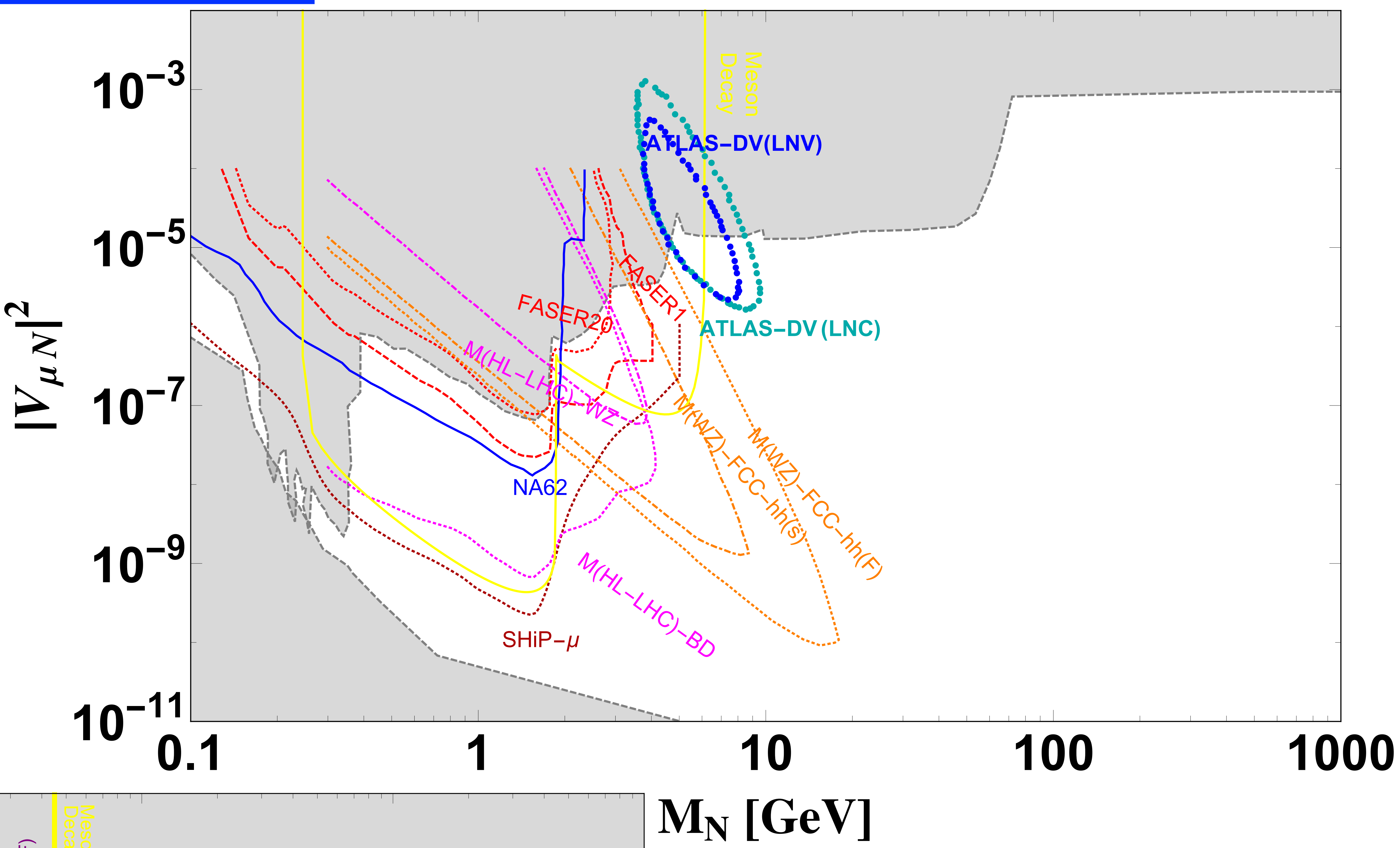
Production processes

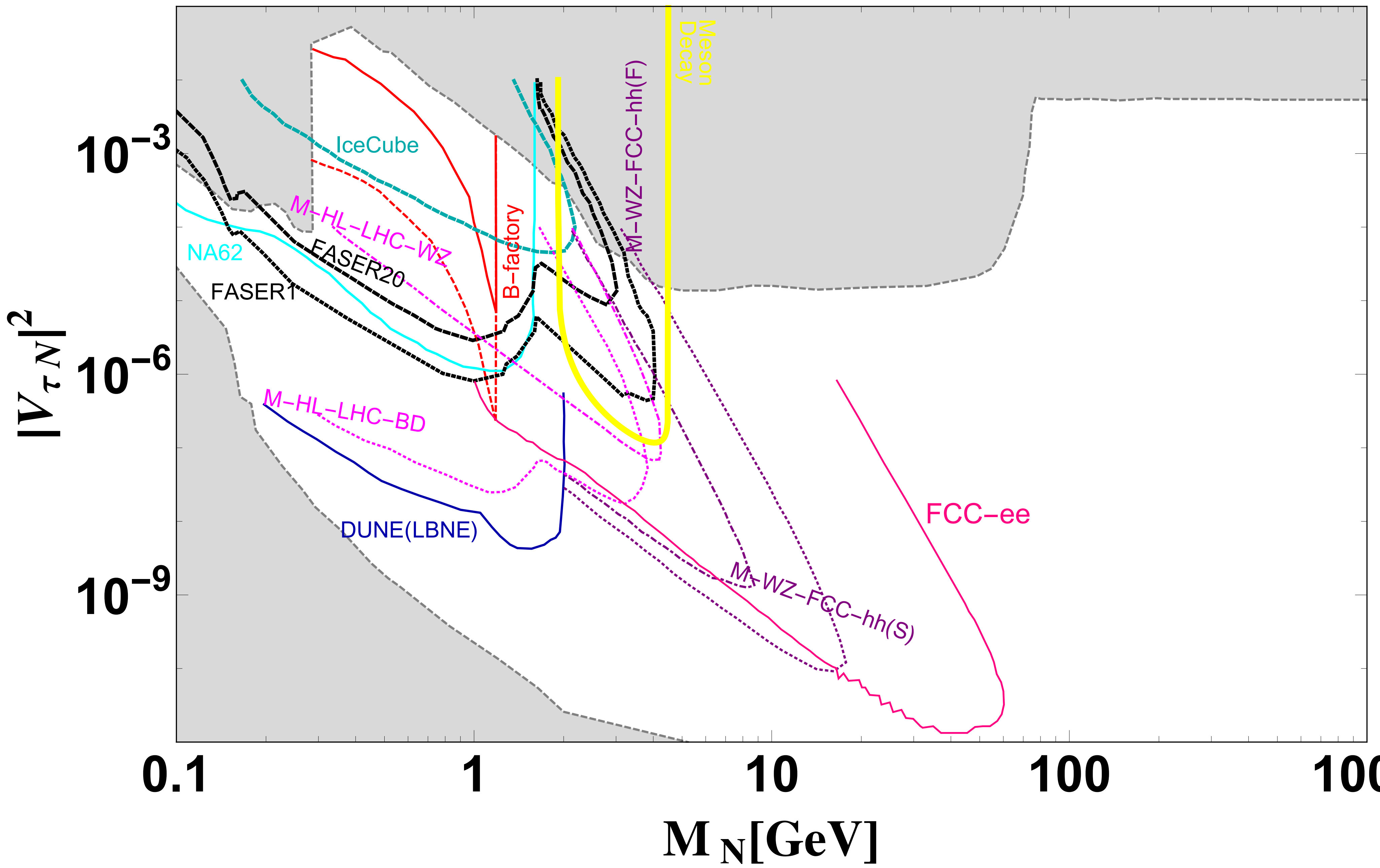
Decay

Existing and prospective bounds on the mixings

1502.06541 1805.00070 1908.09562

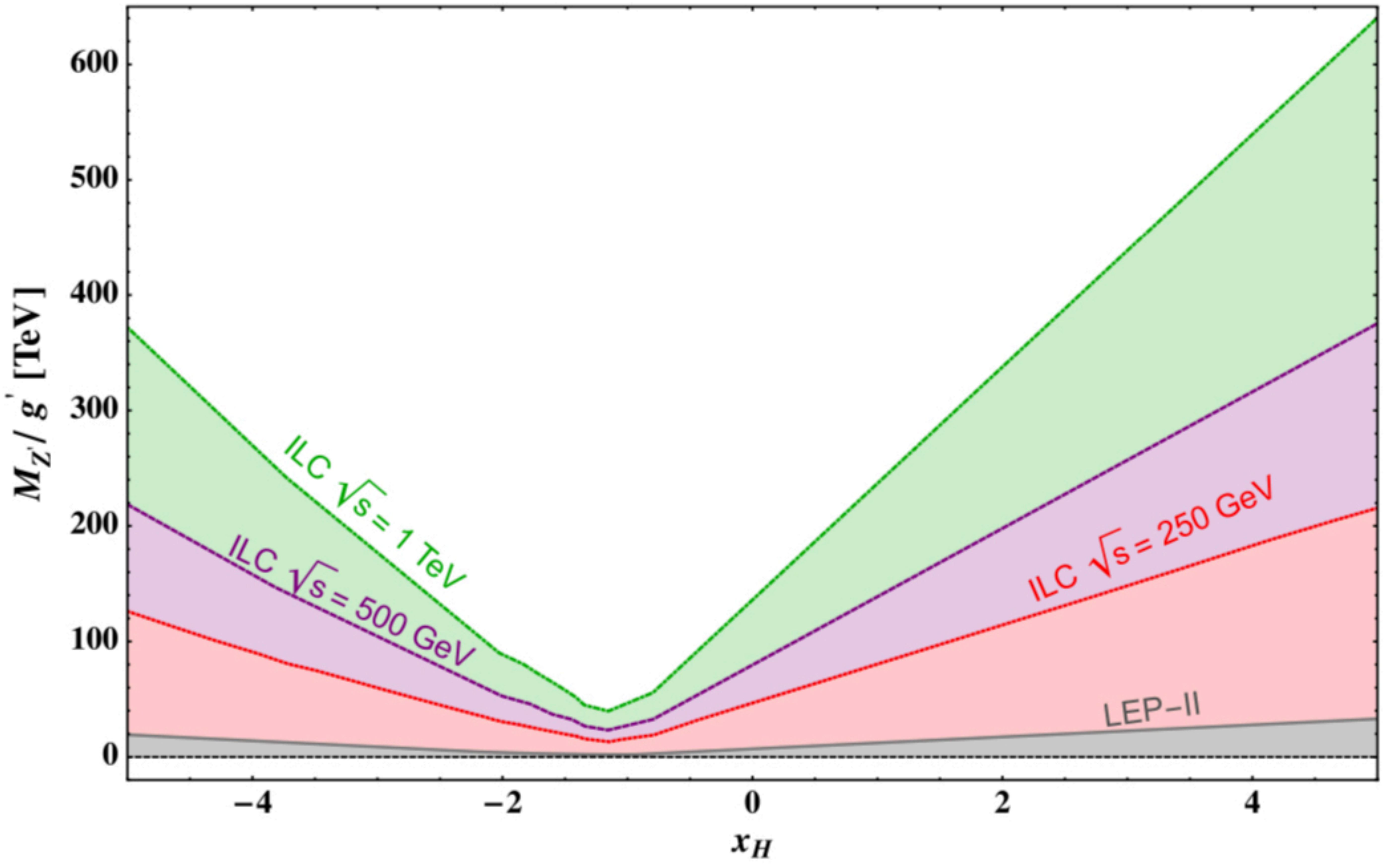






Limits on the model parameters

Considering the limit $M_{Z'} >> \sqrt{s}$ and appling effective theory we find the limits on $\frac{M_{Z'}}{g'}$ using LEP – II (1302.3415) and (prospective) ILC (1908.11299) :



$$\frac{\pm 4\pi}{(1+\delta_{ef})(\Lambda_{AB}^{f\pm})^2}(\bar{e}\gamma_\mu P_A e)(\bar{f}\gamma_\mu P_B f)$$

Z' exchange matrix element for our process

$$\frac{(g')^2}{{M_{Z'}}^2 - s}[\bar{e}\gamma_\mu({x_\ell}' P_L + {x_e}' P_R)e][\bar{f}\gamma_\mu(x_{f_L} P_L + x_{f_R} P_R)f]$$

Matching the above equations we obtain

$$M_{Z'}^2 - s \geq \frac{g'^2}{4\pi}|x_{e_A} x_{f_B}|(\Lambda_{AB}^{f\pm})^2$$

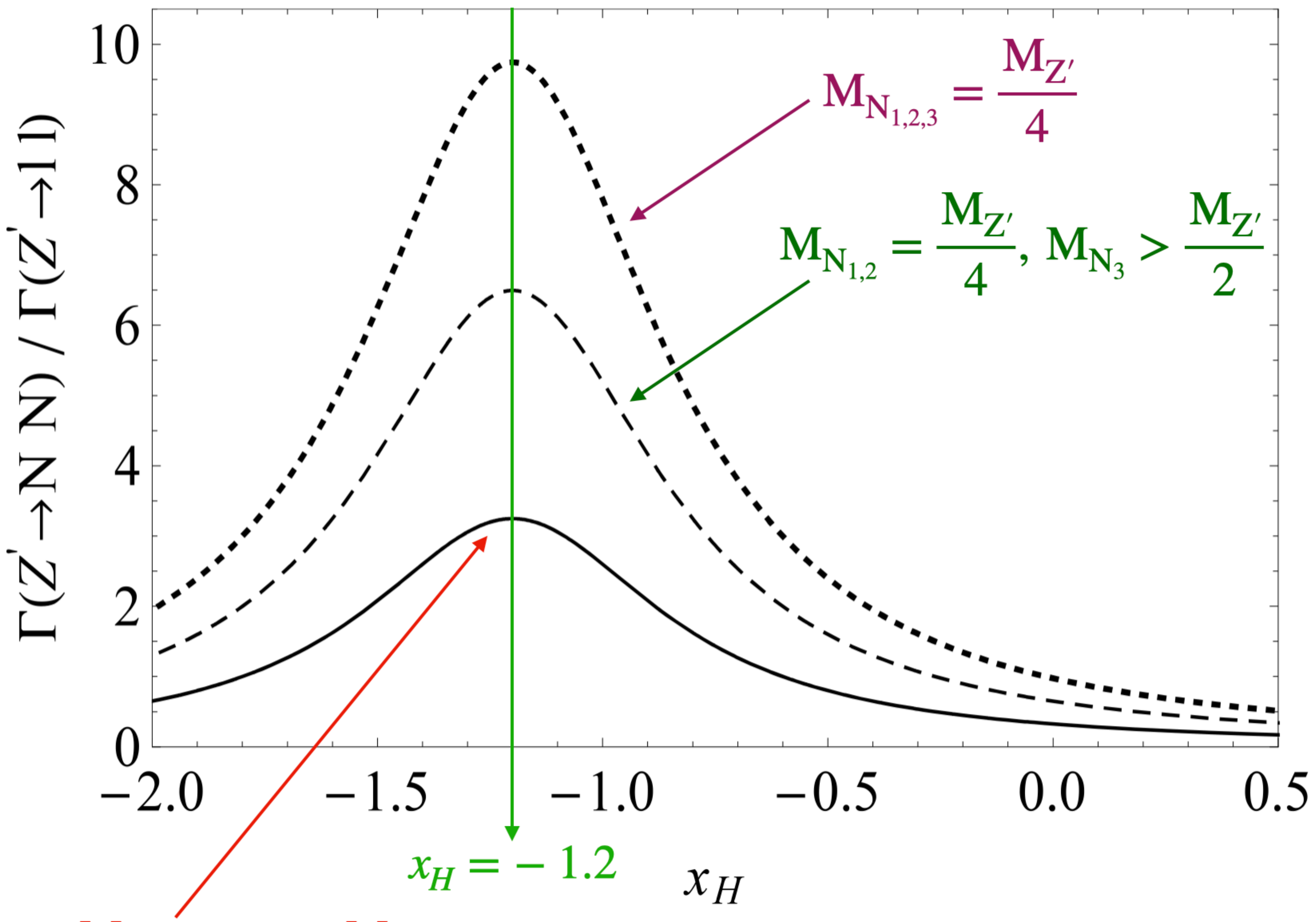
Indicates a large VEV scale can be probed from LEP – II to ILC1000 via ILC250 and ILC500

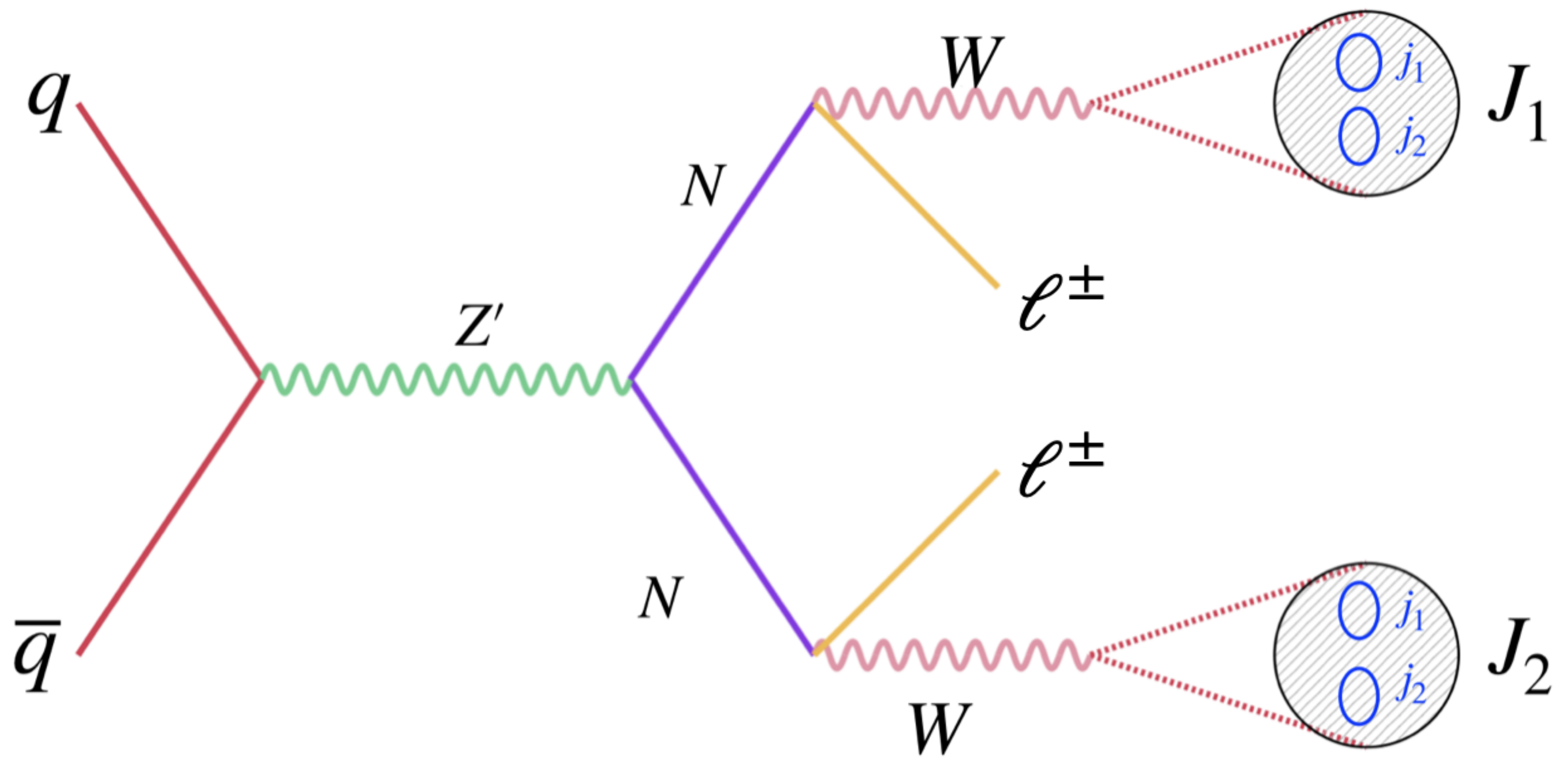
Shows limits on $M_{Z'}$ vs g' for LEP – II, ILC250, ILC500 and ILC1000

Limits on $M_{Z'}$ and g' can also be obtained from dilepton and dijet searches at the LHC

$$g' = \sqrt{g_{\rm Model}^2\left(\frac{\sigma_{\rm ATLAS}^{\rm Obs.}}{\sigma_{\rm Model}}\right)}$$

Pair production of RHN from Z'



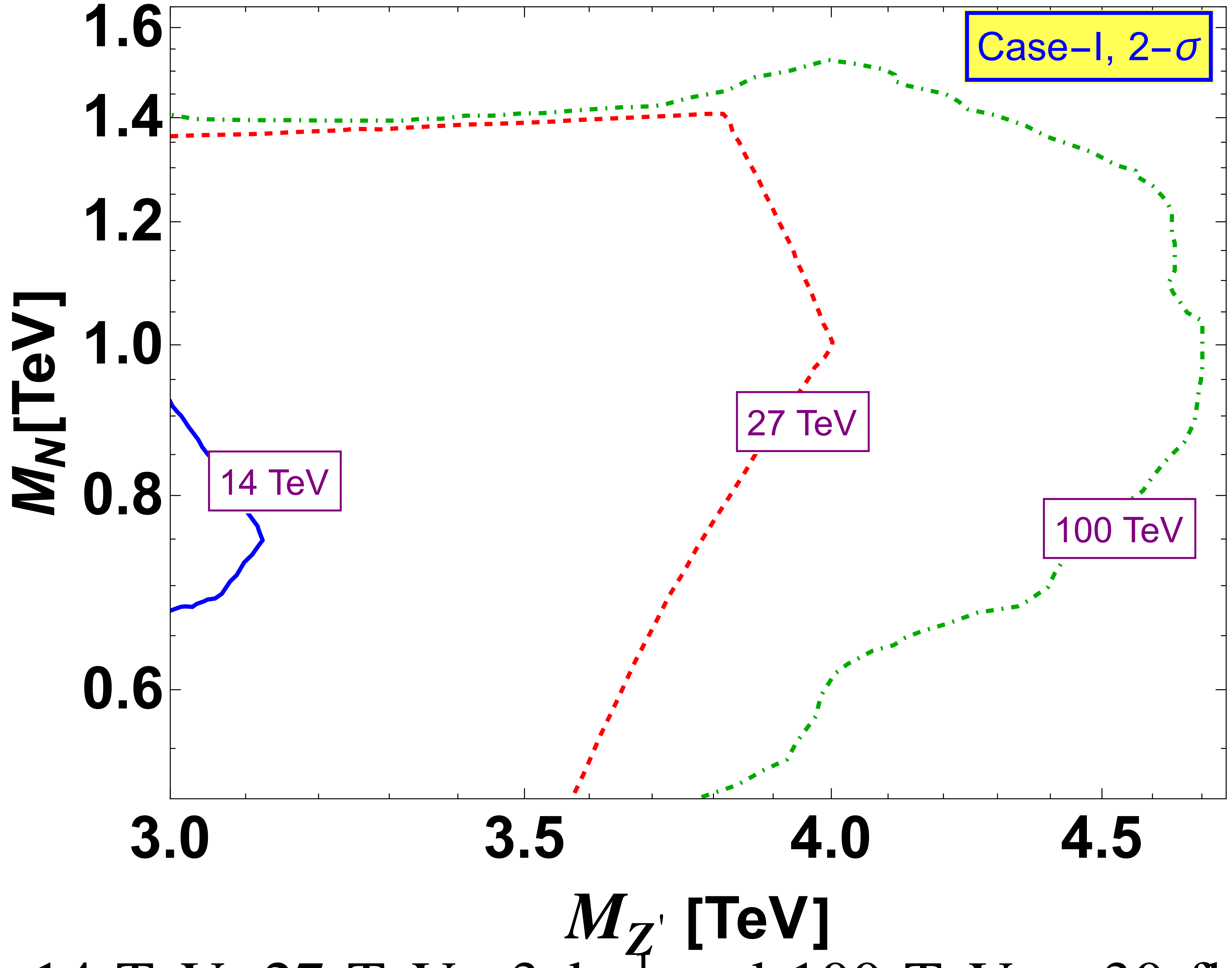


Cuts : $p_T^J > 100$ GeV, $p_T^{\ell_1,\ell_2} > 100$ GeV,

jet – mass window 15 GeV, b – veto

Backgorunds :

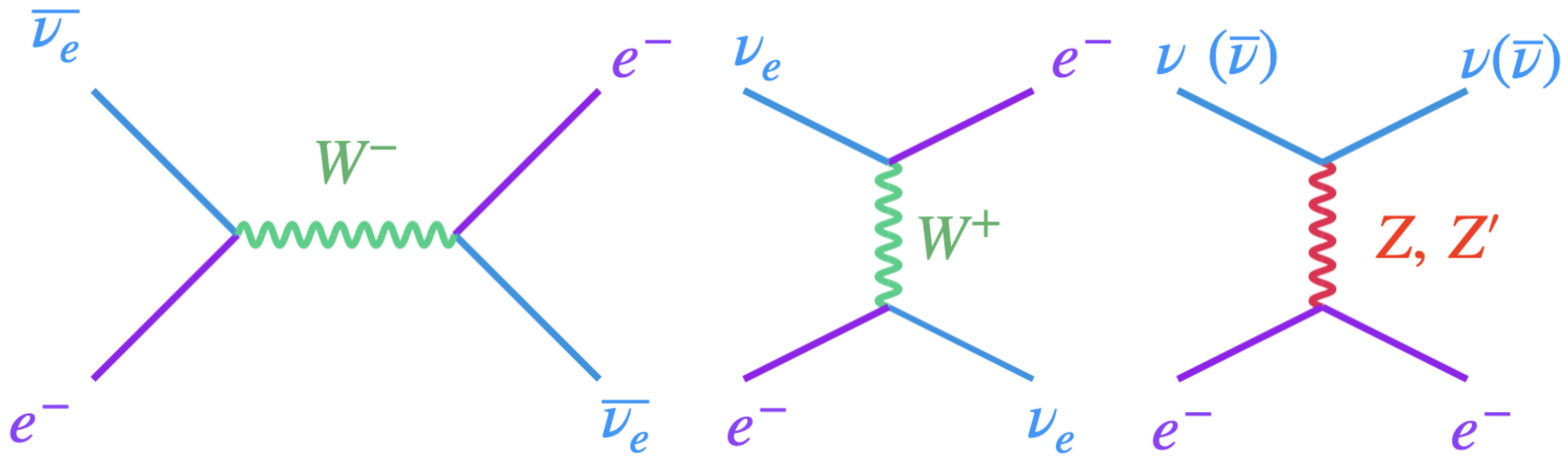
$W^\pm W^\pm jj,\ W^\pm Zjj,\ W^\pm W^\mp Zj,\ t\bar{t}W^\pm,\ t\bar{t}z$



14 TeV, 27 TeV : 3ab^{-1} and 100 TeV at 30 fb^{-1}

Neutrino electron scattering

2111.08767



(1) (2)

The interactions between charged leptons and Z'

	ℓ_{L_i}	e_{R_i}	N_{R_α}	H	Φ
$SU(2)_L$	**2**	**1**	**1**	**2**	**1**
$U(1)_Y$	$-1/2$	-1	0	$1/2$	0
$U(1)_X$	$-\frac{1}{2}x_H - x_\Phi$	$-x_H - x_\Phi$	$-x_\Phi$	$\frac{1}{2}x_H$	$2x_\Phi$

$$-\mathcal{L}^{\ell}_{\text{int}} = g_X(\overline{\ell_L}Q^{\ell}_X\gamma^\mu Z'_\mu \ell_L + \overline{\ell_R}Q^{e_R}_X\gamma^\mu Z'_\mu \ell_R)$$

SM : The interactions of the leptons with Z and W bosons

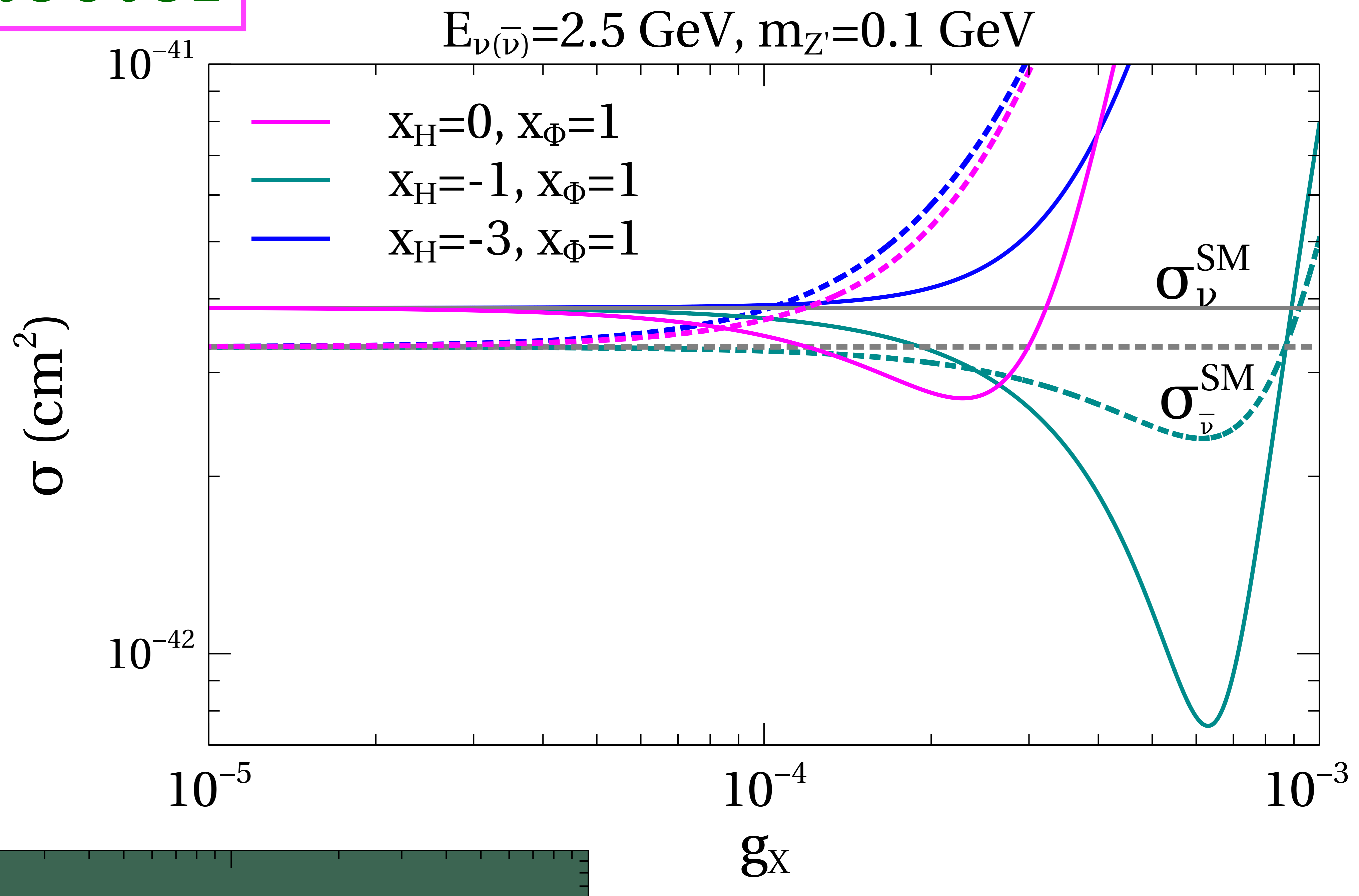
Scattering Process	a_1	a_2
$\nu_e e \to \nu_e e$	$\sin^2\theta_w + 1/2$	$\sin^2\theta_w$
$\bar{\nu}_e e \to \bar{\nu}_e e$	$\sin^2\theta_w$	$\sin^2\theta_w + 1/2$
$\nu_\beta e \to \nu_\beta e$	$\sin^2\theta_w - 1/2$	$\sin^2\theta_w$
$\bar{\nu}_\beta e \to \bar{\nu}_\beta e$	$\sin^2\theta_w$	$\sin^2\theta_w - 1/2$

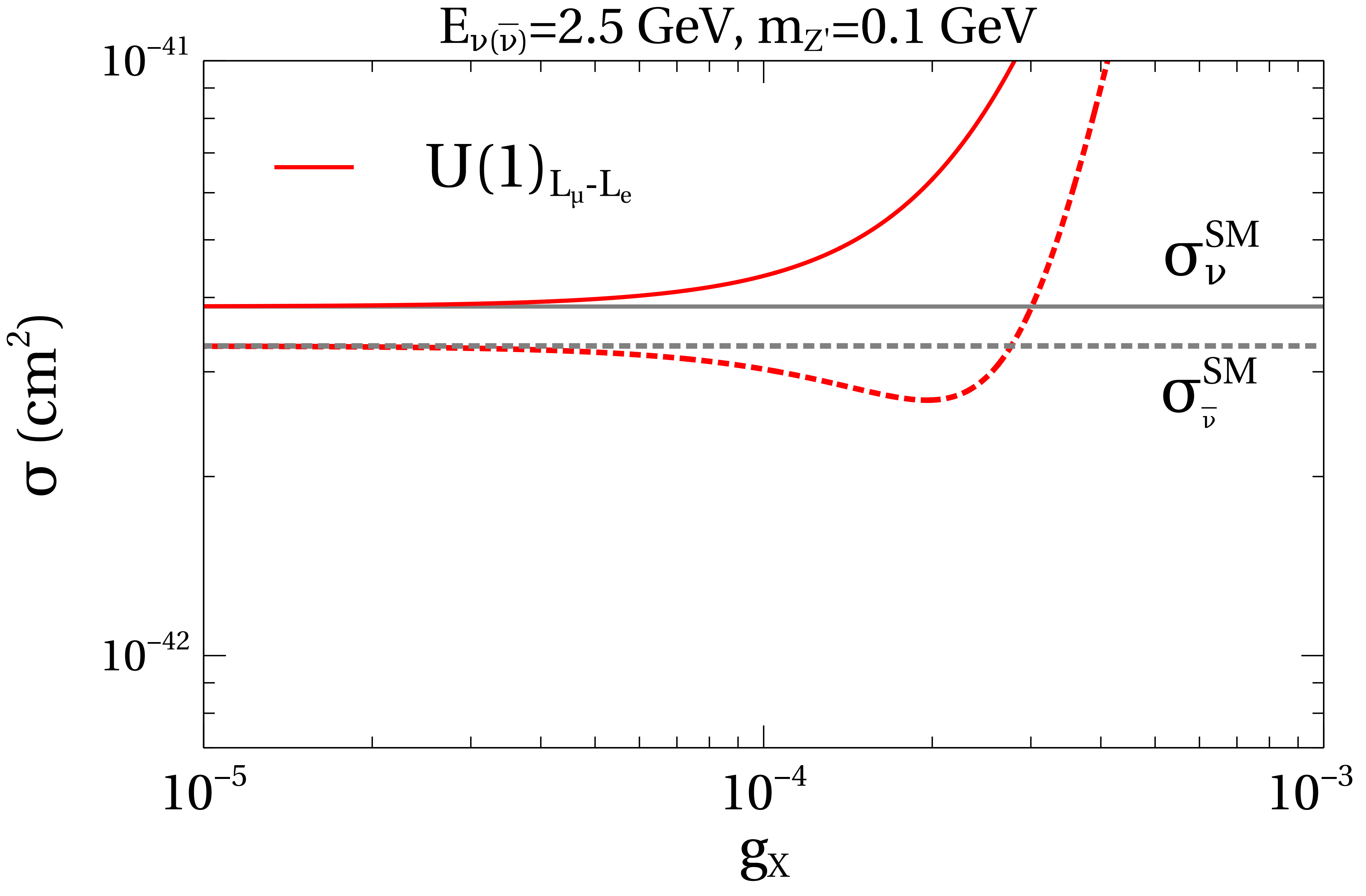
$$\frac{d\sigma(\nu e)}{dT} = \left.\frac{d\sigma(\nu e)}{dT}\right|_{\text{SM}} + \left.\frac{d\sigma(\nu e)}{dT}\right|_{Z'} + \left.\frac{d\sigma(\nu e)}{dT}\right|_{\text{int}}$$

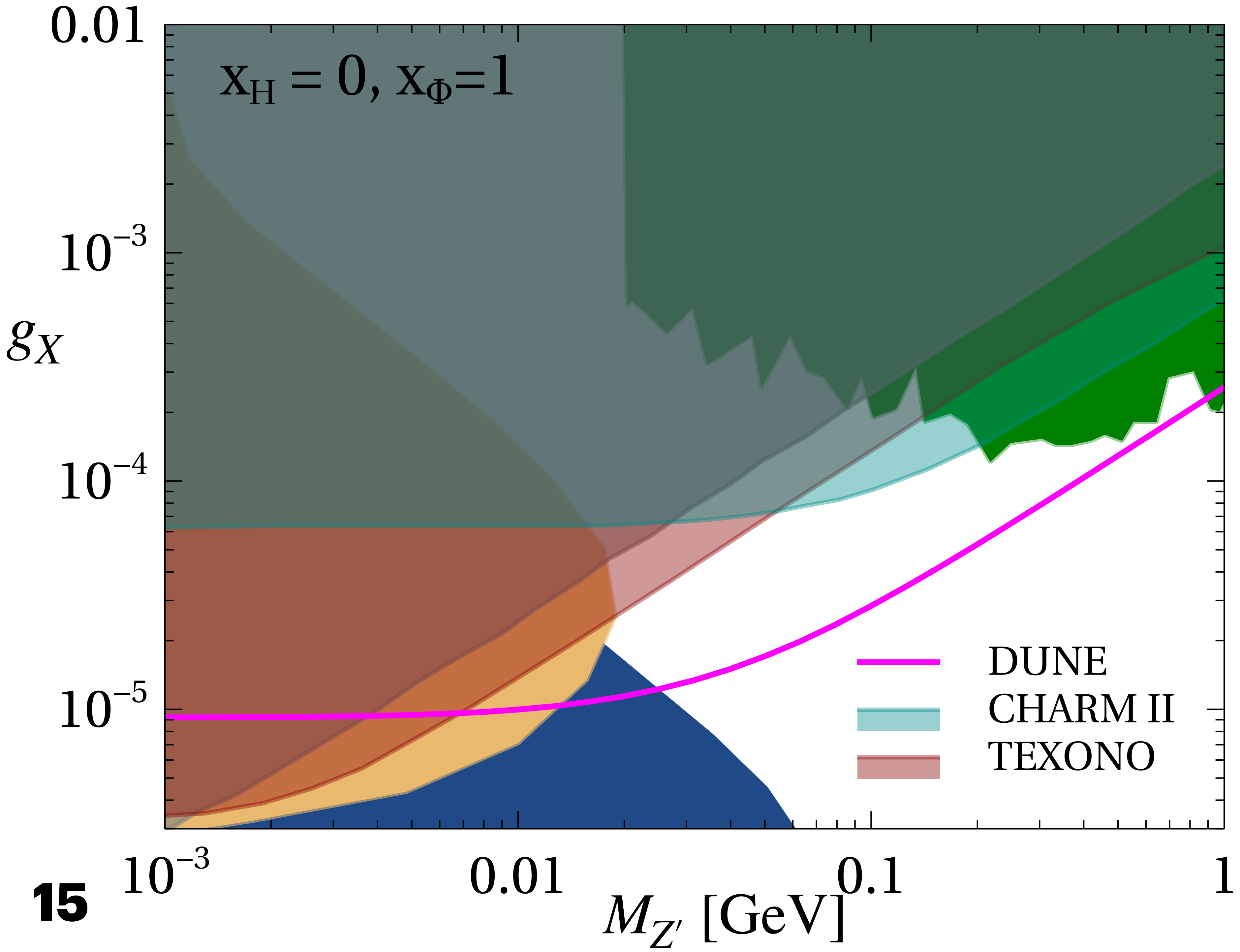
$$c_1 = -1/2 + 2\sin^2\theta_W,\ c_2 = -1/2$$

T : KE of out going electron

DUNE near detector







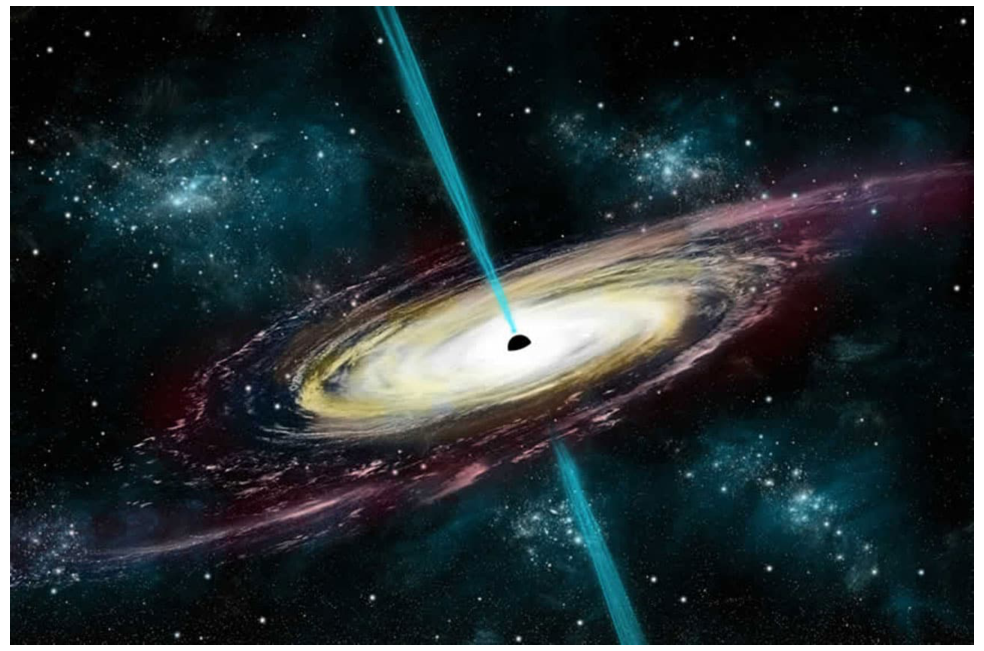
Total cross section of $\overset{(-)}{\nu}_\mu - e$ scattering for SM and U(1).

ν_μ : Solid, $\overline{\nu_\mu}$: Dashed

2111.08767 (DUNE), 2206.12676 (Beam dump)

Interesting propspects of neutrinos

with ShivaSankar et. al.
2308.14483



Neutrinos can be orginiated from the astrophysical sources

Galactic/ Extragalactic sources

Active galactic nuclei : A compact region at the center of the galaxy with high luminosity

Most luminous source of EM radiation. It has beeen observed in radiowaves and X – rays.

Gamma ray burst : Sudden intense flashes of the gamma rays

Highly energetic   Nature is unknown $(\nu\bar{\nu} \to e^- e^+)$

Are the neutrinos emitted by these *distant* sources?   If emitted then how to study?

To study distant objects we need a telescope.

Basically we built it at the south pole of the earth and put it 1 km under the ground(ice) : IceCube detector.

Neutrios came the distant sources scattered with the electrons in the ice producing Cherenkov light.

In 2013, they detected few PeV neutrinos. In 2018, they observed energetic neutrinos.

In course of time they have measured high energy cosmic neutrinos between 100 TeV to 10 PeV from extra galactic sources and cosmic ray muons.

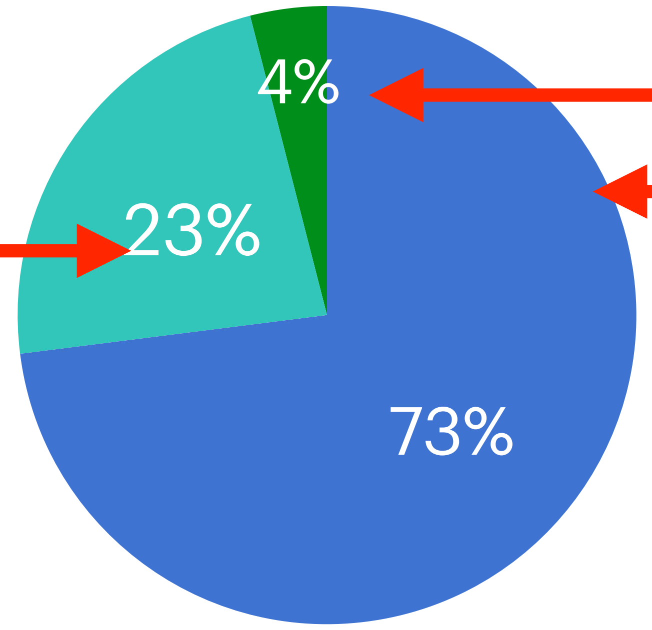
Neutrinos are charge neutral : no bending under the electromagnetic field.

So the neutrinos do not scatter and can be directly → pointed back to the source.

Dark matter and neutrinos

ν – DM scattering

(AGN/ blazar, IceCube)



Weekly interacting

Such models can successfully fit a potential Dark matter candidate ($N_{1/3}$) adding discrete symmetry Z_2

Depending upon the model parameters

Relic

Direct search

Z′ portal, Higgs portal

Mono – jet, mono – photon

LHC ILC

Feebly interacting

DM abundance is obtained by the annihilation and decay of the SM particles

DM yield increases with the increase in coupling

UV Freeze – in : DM produced by the SM particles mediated by portal

IR Freeze – in : DM interacts with visible sector via a renormalizable operator production dominates at low T

Z′ mediated processes

RHN DM

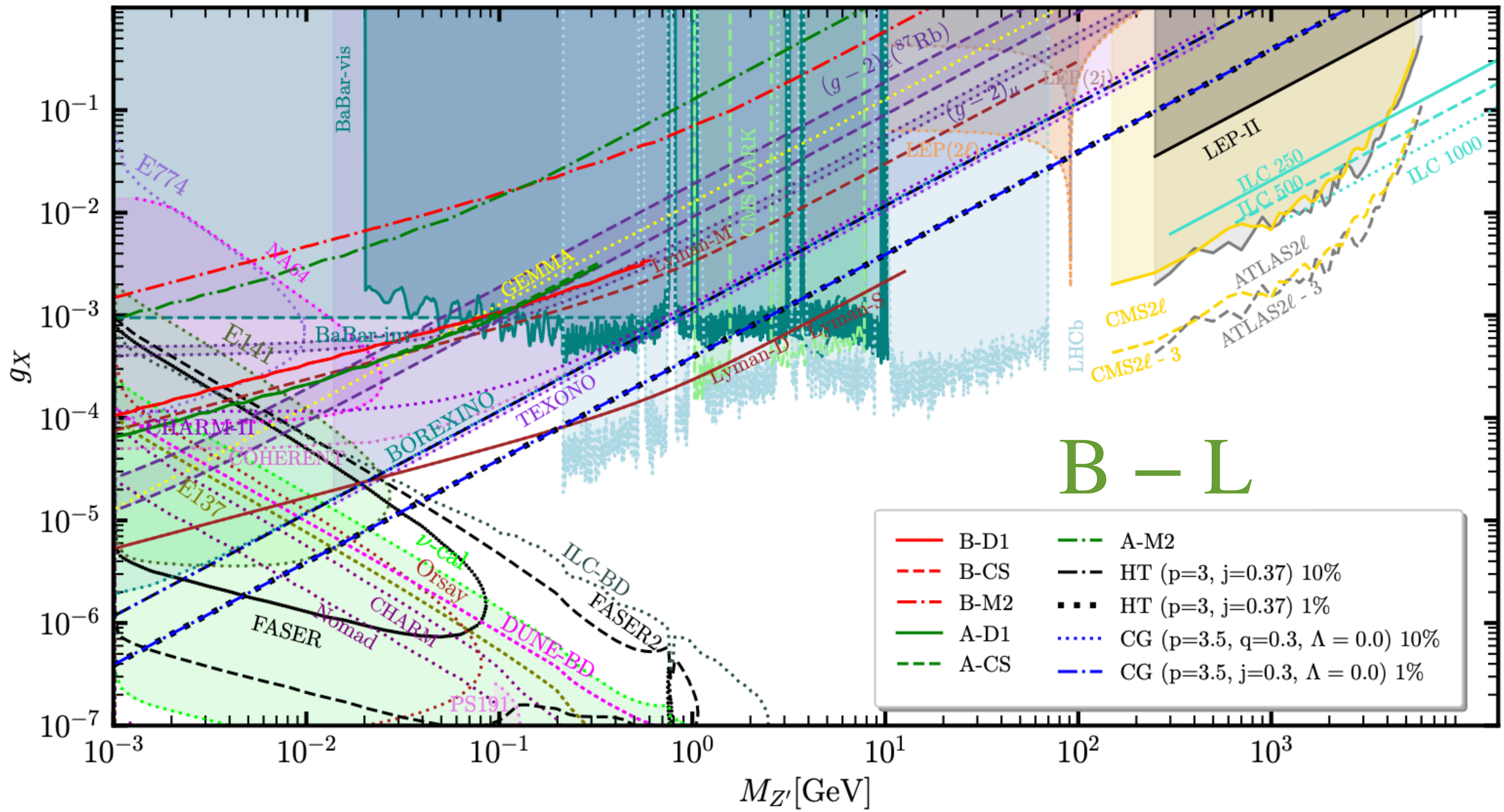
SM fermions annihilate into RHN DM through Z′

RHN DM in MeV scale Z′ in TeV scale

Pair produced from Z′

RHN DM decays into positrons to explain

511 keV INTEGRAL anomaly



For different charges and flavor models : 2308.14483

Conclusions

We are looking for a scenario where which can explain a variety of beyond the SM sceanrios .

The proposal for the generation of the tiny neutrino mass, from the seesaw mechanism, under investigation at the energy frontier .

Mnay aspects can be addressed in these scenarios which could connect three interesting frontiers of physics

The motivations of these works is to find a new particle, Z' a new force carrier as a part of the of the new physics search including various BSM aspects .