

pfRICH event reconstruction considerations

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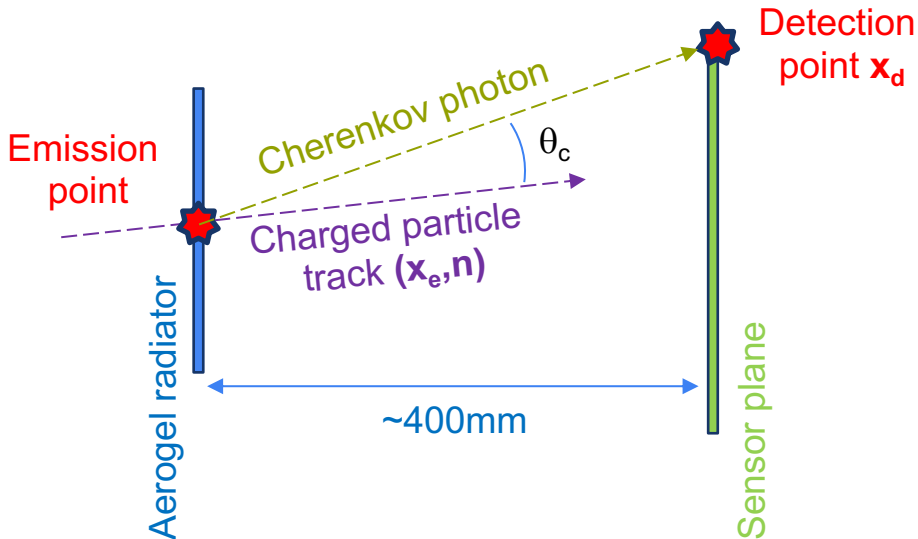
ePIC Collaboration meeting, Argonne, January 11, 2024

Disclaimer

- Talk will be focused on ring imaging rather than a “photon flash in the HRPPD window” timing
- Most part of the described algorithms was actually implemented for the March 2023 review as IRT 2.0
 - Will be indicated in case of the opposite
- Existing codes are algorithmic, combinatorial, χ^2 based
- This was all done in a standalone GEANT4 environment
- Porting to dd4hep is in the geometry description stage

Starting point

- Single track (and ideal tracking)
- Very thin radiator
 - No emission point uncertainty
- Single photon with a known λ
 - And therefore, a known $n(\lambda)$
- Ideal sensor plane
 - $\sigma_{xy} \sim 0$ (no detection point uncertainty)
- No mirrors & no refractive boundaries
 - Photon trajectory is a straight line
- No noise hits



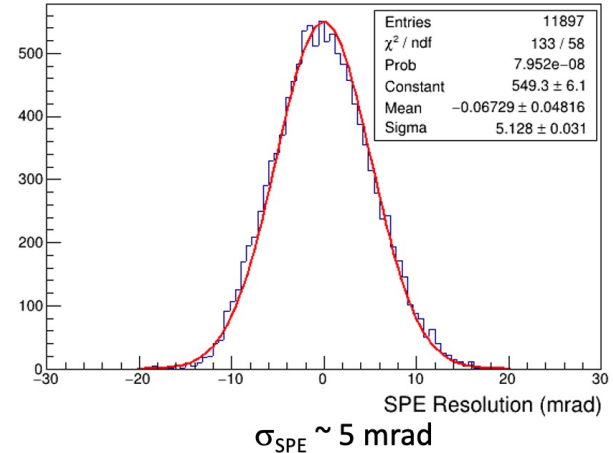
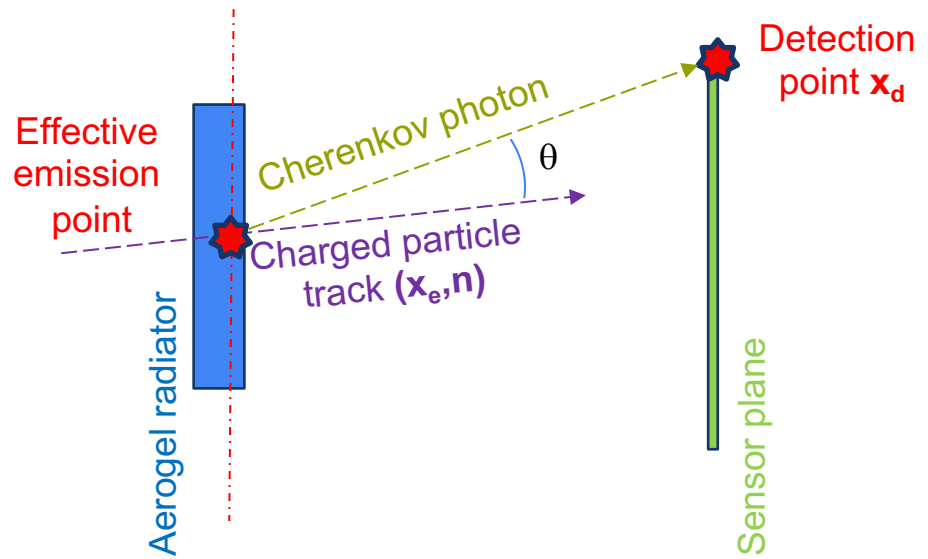
$$\cos \theta_c = \frac{\vec{n}(\vec{x}_d - \vec{x}_e)}{\|\vec{x}_d - \vec{x}_e\|} = \frac{1}{\beta n}$$

$$m = \frac{p}{c\beta\gamma}$$

➔ Substitute momentum & velocity in the formula and check PDG for this mass

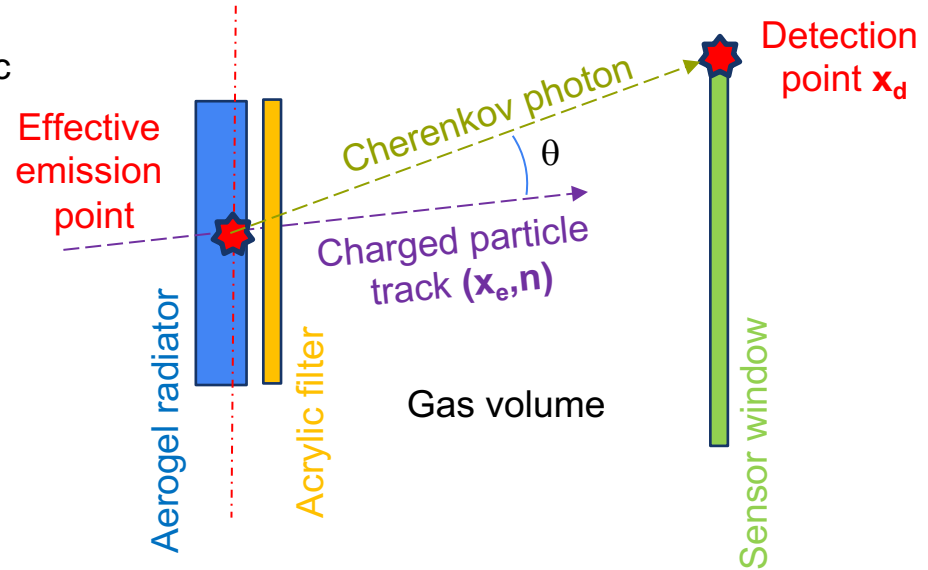
Unrecoverable errors

- Emission point uncertainty
 - Aerogel radiator has a finite thickness
 - We have no means to measure the actual emission point of each photon
 - Wavelength-dependent absorption in the radiator
- Detection point uncertainty
 - Sensor plane has a finite resolution
- Chromatic effects
 - We do not know wavelengths on per-photon basis
 - Have to deal with $\langle \lambda \rangle$ and RMS
 - There is a $n(\lambda)$ dependency ...
 - ... which translates into a spread of expected emission angles for a fixed particle momentum



Refraction on optical media boundaries

- Photon path is not really a straight line
 - A sequence of *refractions* on aerogel, gas, acrylic filter and HRPPD window boundaries
- Taken care of by the IRT algorithm
 - No analytic solution exists in a generic case
 - Iterations (a 2D Newton-Gauss method)
 - Same technique used for optical paths involving *reflections* on mirror surfaces
 - At the end of the day what matters is the θ_c angle estimate (and to some extent ϕ_c) *at the anticipated emission vertex*



Building a χ^2 statistics in this simple case

- Prefer to work in a *measurement space*
 - Where “measurement” is a single photon emission angle θ_c with respect to the track

$$\chi_H^2 = \frac{[\theta_H(p, n) - \theta_c]^2}{\sigma_\theta^2} \quad \text{for a given PID hypothesis } H \text{ (e.g. a pion)}$$

- This quantity should be distributed as a tabulated χ^2 with *one* degree of freedom *for a correct PID hypothesis ...*
 - In other words: a cumulative quantity (CCDF) should be a flat distribution between 0 and 1
- ... and systematically biased towards higher (less probable) values *for wrong hypotheses*
 - In other words: a CCDF plot would tend to produce a spike close to 0

All the rest is built on a simple basic principle: we construct *static* χ^2 estimates for various PID hypotheses (say $\pi/K/p$) for a given set of hit-to-track associations, and the smallest χ^2 wins

Multiple photons per track

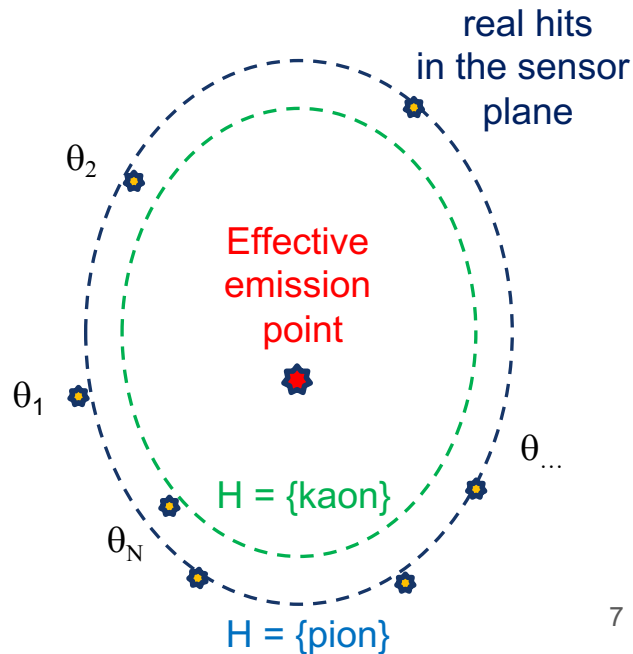
- Assuming each photon is a single independent measurement, the χ^2 formulation is extended in a trivial way:

$$\chi_H^2 = \sum_{k=1}^{nhits} \frac{[\theta_H(p, n) - \theta_c^k]^2}{\sigma_\theta^2}$$

for a given PID hypothesis H (e.g. a pion)

- If there were no Rayleigh scattered photons, this quantity should be distributed as a tabulated χ^2 with $nhits$ degrees of freedom *for a correct* PID hypothesis and be systematically biased towards higher (less probable) values *for wrong hypotheses*

In practical terms, this gives one a $1/\sqrt{N}$ factor in a *track-level* Cherenkov angle resolution, of an order of ~ 1.5 mrad for pFRICH



Accounting for a Poisson term

- Omitting certain technical issues, an expected number of “true” Cherenkov photons $\langle N^{\text{expected}} \rangle$ for a given combination of momentum and a mass hypothesis is known ($N_{pe} = N_0 L \sin^2 \theta_c$), and can be added to the overall χ^2 estimate as a so-called χ^2_λ term (Baker & Cousins notation), which accounts for a mismatch between the expected and the detected number of photons:

$$\chi_H^2 = \sum_{k=1}^{nhits} \frac{[\theta_H(p, n) - \theta_c^k]^2}{\sigma_\theta^2} + \chi_{\lambda(H)}^2$$

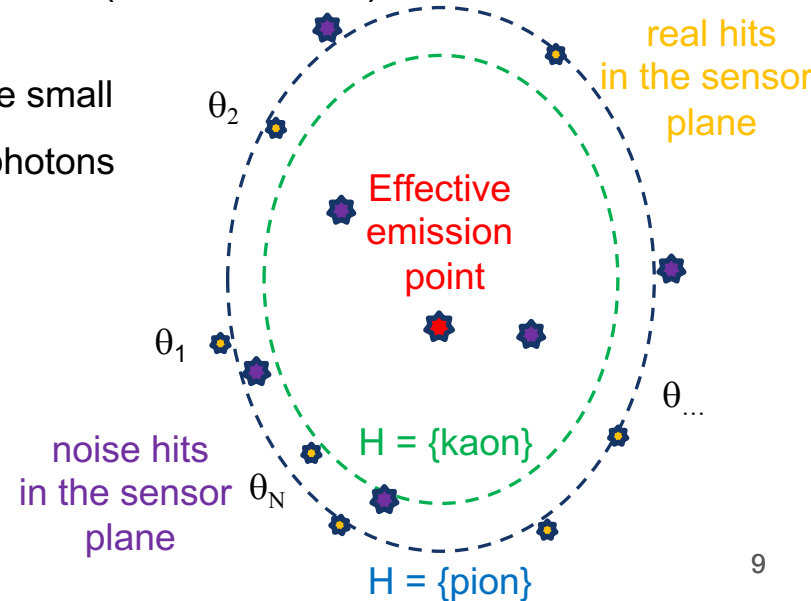
where

$$\chi_{\lambda(H)}^2 = 2N^{nhits} \ln(N^{nhits} / \langle N^{\text{expected}}(H) \rangle)$$

Accounting for noise hits

From this point on we assume that a hit-to-track association is *not* known

- We only deal with hits which are within a $\pm 3\sigma$ band (think of ± 15 mrad) around a nominal Cherenkov θ angle *for at least one of the PID hypotheses*
- Omitting some technical complications, for a case of e.g. a π/K separation this defines two “circular” bands on the sensor plane, where such hits can be located
- A sheer area of these two bands, and a noise rate per cm^2 (as simulated!) defines an expected number of noise hits $\langle N^{\text{bg}} \rangle$
 - HRPPDs: $\sim \text{kHz}/\text{cm}^2$ & $< 50\text{ps}$ timing \rightarrow expect $\langle N^{\text{bg}} \rangle$ to be small
 - Plus, another small contribution from Rayleigh scattered photons
- Each PID hypothesis can either account each of the hits as a “true” or as a “background” one, pushing highest χ^2 hits into a separate (background) χ^2_{λ} term one by one
 - In practical terms, a “wrong” hypothesis will incur an additional penalty because of existence of too many fake noise hits (and possibly too few true ones)



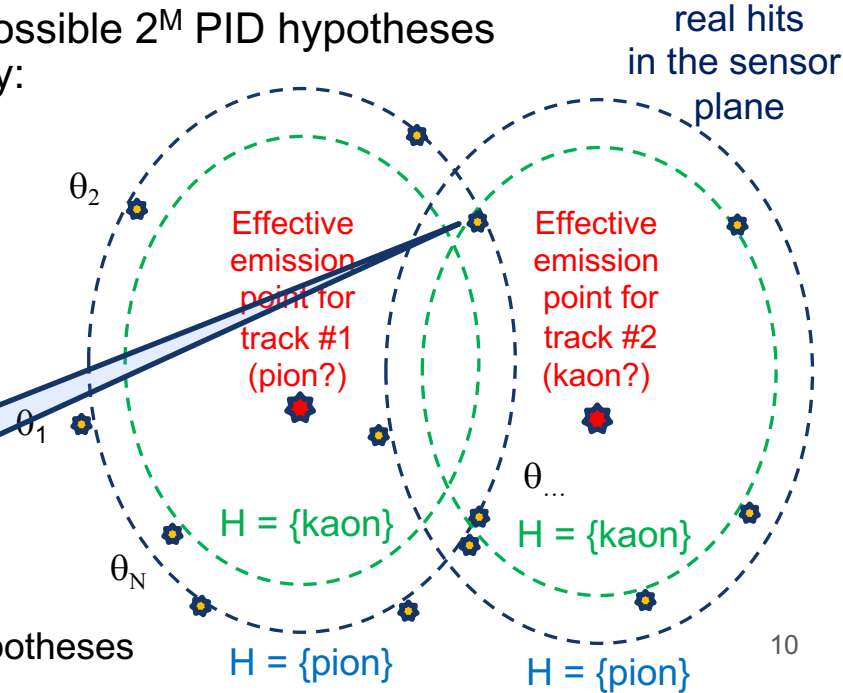
Overlapping rings

- The algorithm works on the *event* level; let's assume there are M tracks
- We only deal with hits which are within $\pm 3\sigma$ band (think of ± 15 mrad) around a nominal Cherenkov θ angle *for at least one of the PID hypotheses for at least one track*
- Assuming e.g. a π/K separation case, each of the possible 2^M PID hypotheses *event-level combinations* $\{H\}$ is evaluated separately:

$$\chi_H^2 = \sum_{i=1}^{mtracks} \sum_{k=1}^{nhits(i)} \frac{[\theta_H(p_i, n) - \theta_c^{ik}]^2}{\sigma_\theta^2}$$

- Tracks compete for hits
 - Various options for a conflict resolution
 - Assignment can be different for different sets of PID hypotheses

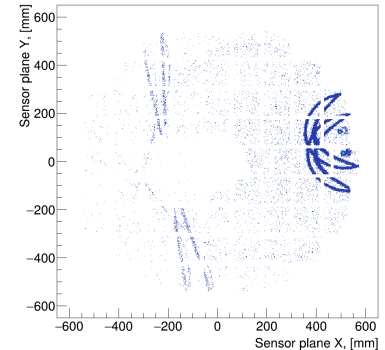
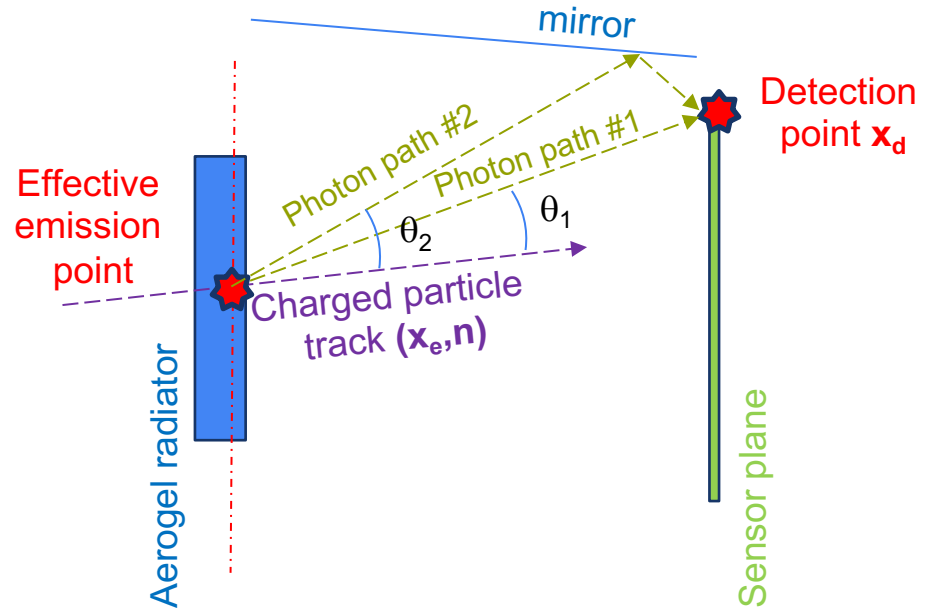
A possible conflict for a $\{\pi, K\}$ event candidate



Mirrors

- pfRICH has inner and outer conical mirrors, and (optionally) four small flat funneling mirrors per HRPPD sensor
 - This defines up to 20 possible optical paths for any pair of emission and detection points
- During a hit-to-track association process, the algorithm loops through all optical paths, calculates a Cherenkov emission angle, and picks up a path which gives a best match *for a PID hypothesis presently considered for this track*

Except for an obvious computational overhead, this does not change anything in the rest of the χ^2 evaluation process



Tracking information

- Small bending in the magnetic field is accounted in a trivial way (a track parameterization at the location of effective emission point is taken)
- Finite tracking system resolution can technically be accounted via replacing a static χ^2 evaluation (with no free parameters) by a full MINUIT-like pass where a track state vector at the effective emission point is added as a separate term to the χ^2 ansatz, with its inverse covariance matrix C as a metric:

$$\chi_H^2 = \sum_{k=1}^{nhits} \frac{[\theta_H(p, n) - \theta_c^k]^2}{\sigma_\theta^2} + \langle X_{track} | C^{-1} | X_{track}^T \rangle$$

where

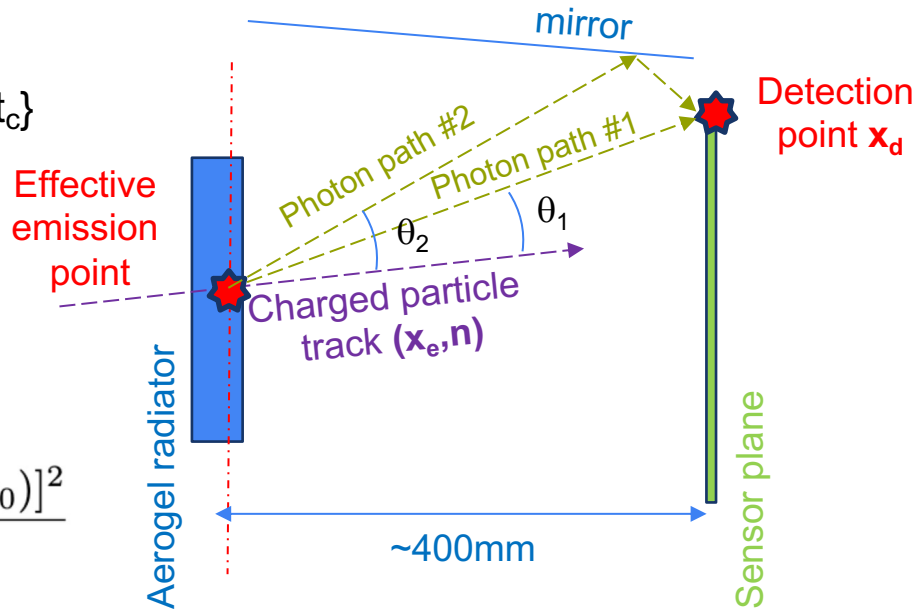
$$X_{track} = \{x, y, sx, sy, 1/p\}_{@aerogel}$$

Though unlikely, this may actually improve the tracking estimate itself

Timing information

- Measurements become 2D vectors: $\{\theta_c\} \rightarrow \{\theta_c, t_c\}$
 - Where t_c is a time measured at the sensor plane
- Timing is used in both hit-to-track association for a given mass hypothesis, and in the χ^2 ansatz:

$$\chi_H^2 = \sum_{k=1}^{nhits} \frac{[\theta_H(p, n) - \theta_c^k]^2}{\sigma_\theta^2} + \sum_{k=1}^{nhits} \frac{[t_H(p) - (t_c^k - t_0)]^2}{\sigma_t^2}$$

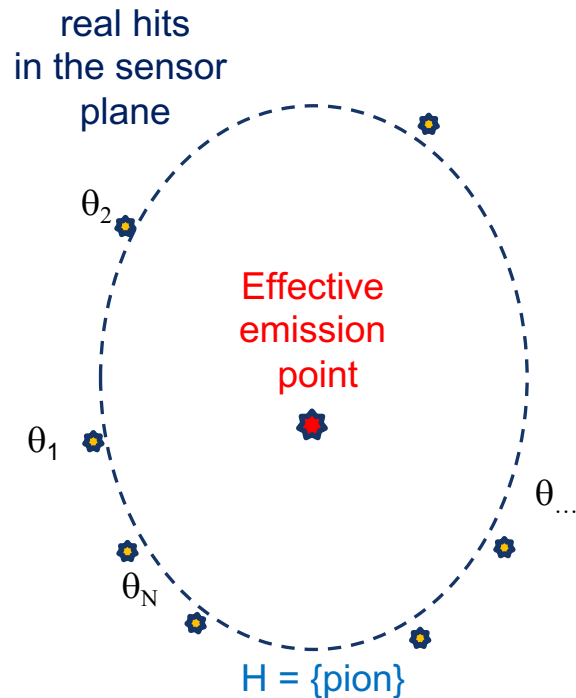


Path#1 and path#2 will not only have a different θ_c , but a substantially different length (and therefore, a very different - compared to a ~50ps resolution - flight time between the emission and detection points)

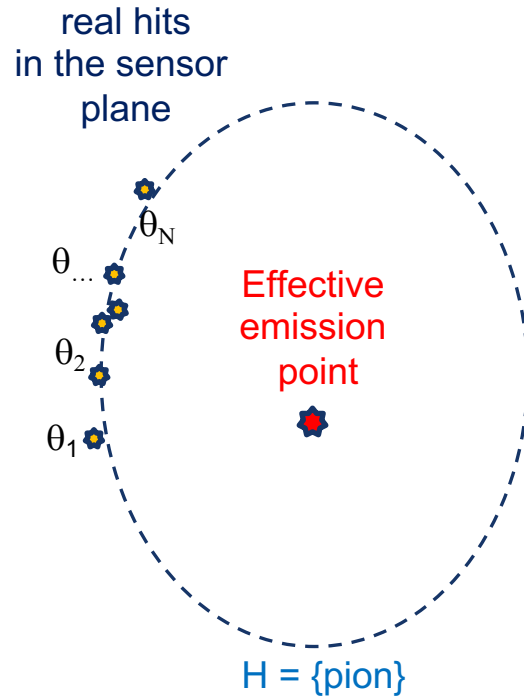
- Presently only a *static* χ^2 evaluation
 - Assume t_0 is known within few dozens of ps, resulting in an “effective” hit timing resolution $\sigma_t \sim 50ps$
- Apparently one can add a t_0 estimate in a (linearized) MINUIT-like fashion

Other considerations

➤ Make use of a signature χ^2 ?



➤ Make use of binning in ϕ_c ?



Both configurations are “unlikely”, but the presently available algorithm does not capture this