

Harnessing intricacies of Jets for Breakthroughs in QCD at the Collider Frontier

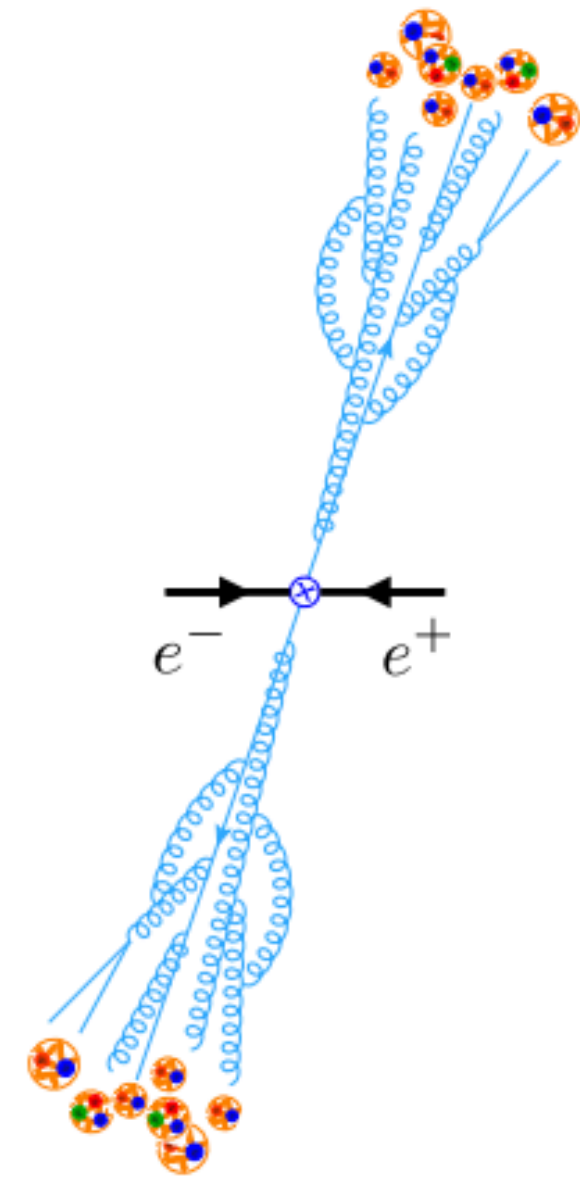
Kyle Lee
CTP, MIT

EICUG/ePIC meeting

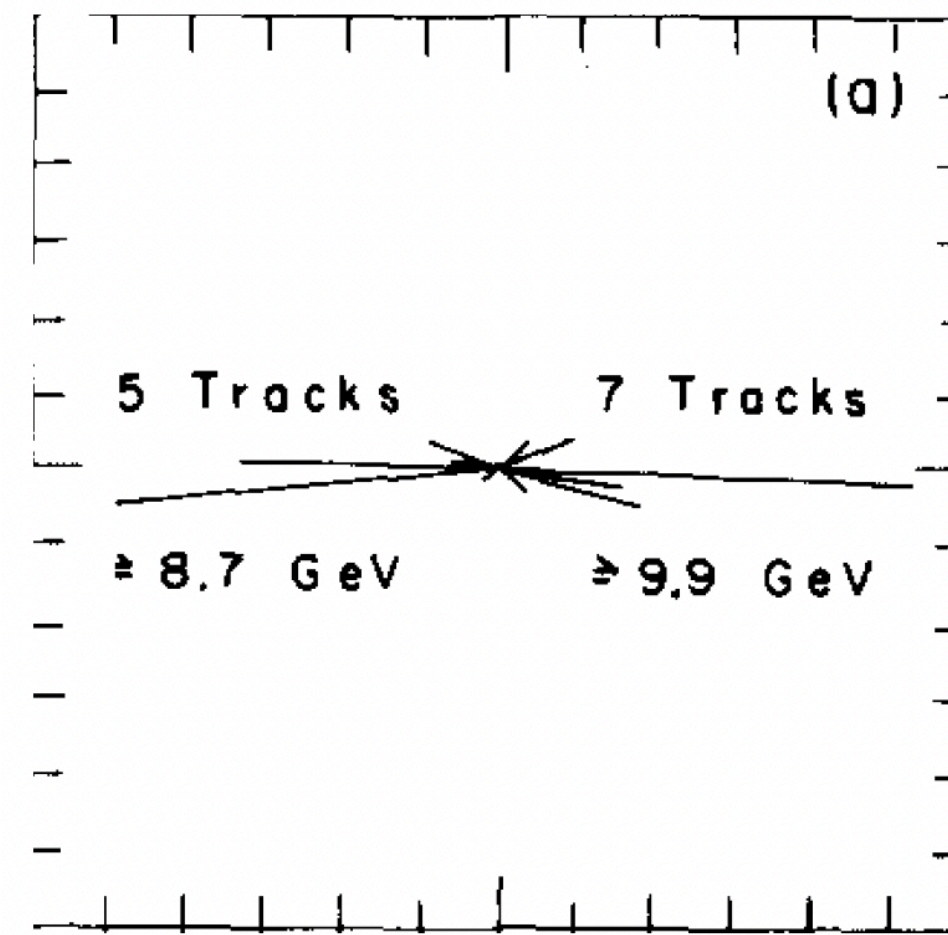
July 2024



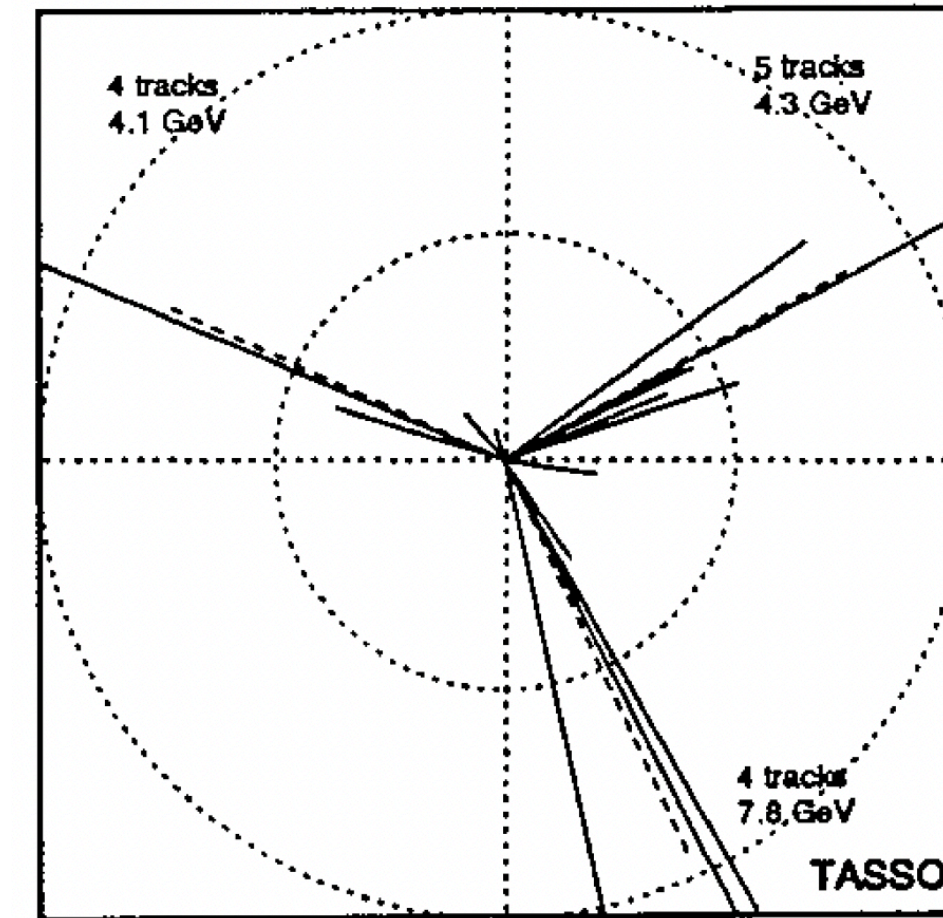
THE DAWN OF QCD: FROM PARTONS TO JETS



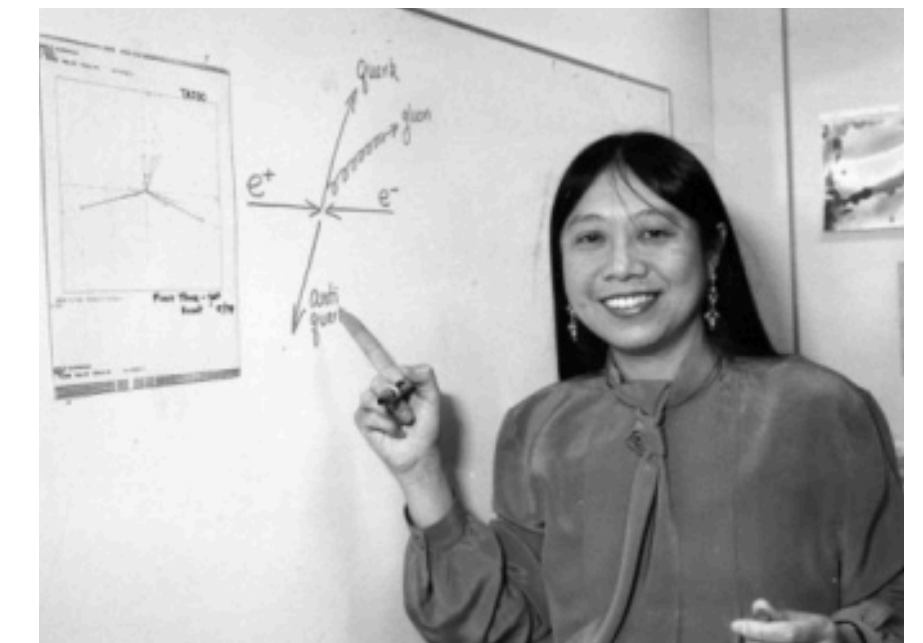
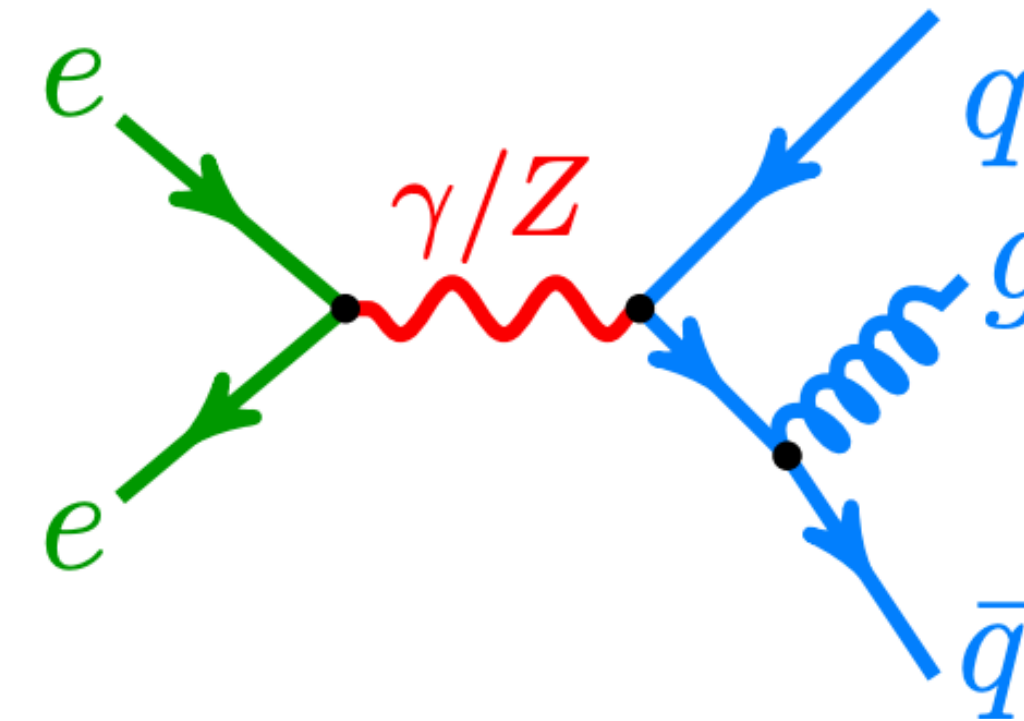
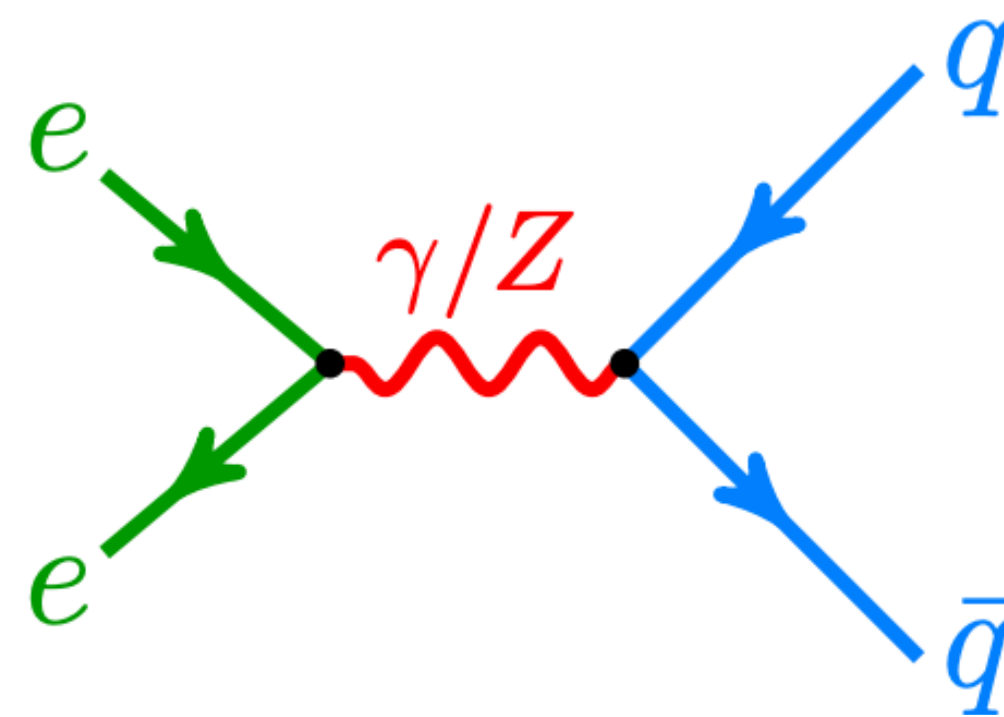
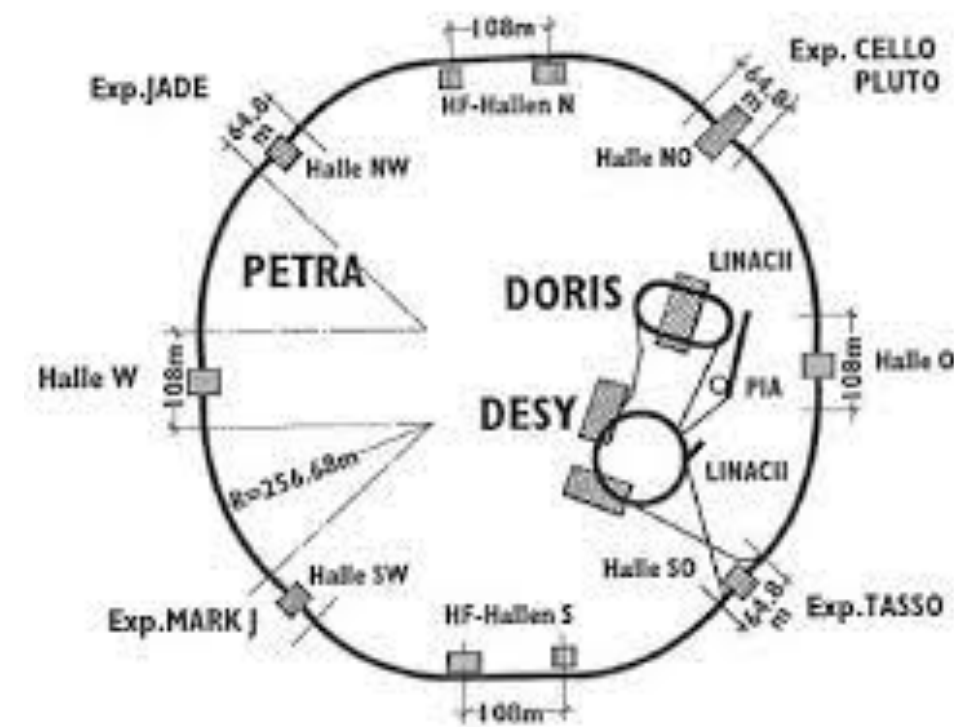
2-jet event



3-jet event



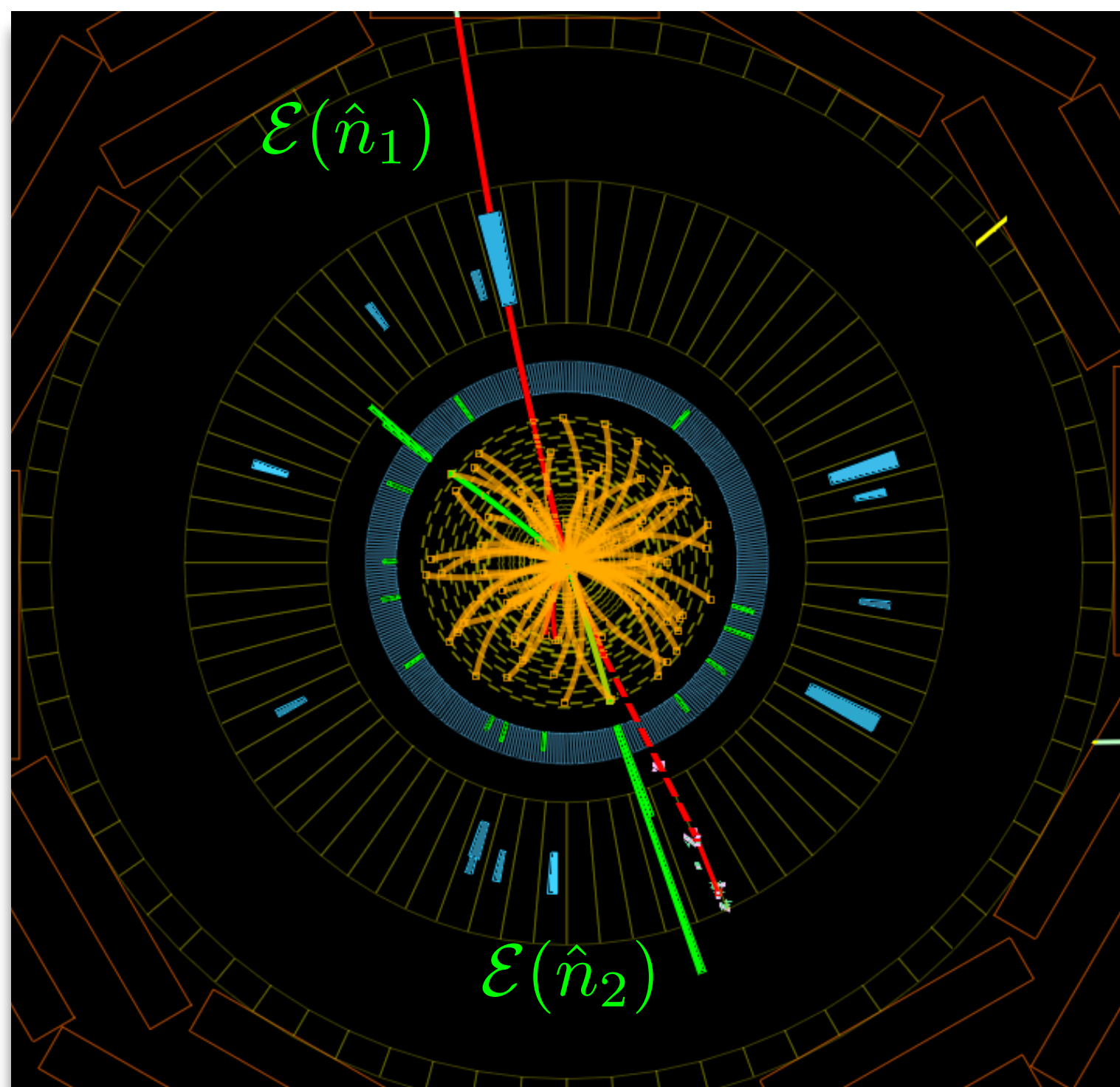
	$\sum P_i \text{CHARGE}$	TOTAL ENERGY
JET 1	4.3 GEV	7.4 GEV
JET 2	7.8	8.9
JET 3	4.1	11.1



Wu, Zoernig '79

Jets unveiled the **partonic nature** of QCD, playing an important role in the **confirmation of QCD** as the **theory of strong interactions!**

JETS AND ENERGY FLOW



Energy Flow Operators

$$\text{camera icon} = \mathcal{E}(\hat{n}) = \int_0^\infty dt \lim_{r \rightarrow \infty} r^2 n^i T_{0i}(t, r\hat{n})$$

$$\mathcal{E}(\hat{n})|X\rangle = \sum_a E_a \delta^{(2)}(\Omega_{\vec{p}_a} - \Omega_{\hat{n}}) |X\rangle$$

Basham, Brown, Ellis, Love, '78-79

Sveshnikov, Tkachov, '95

Korchemsky, Serman, '01



Serman '75

Serman, Weinberg '77

“Energy flow becomes the focus of computability”

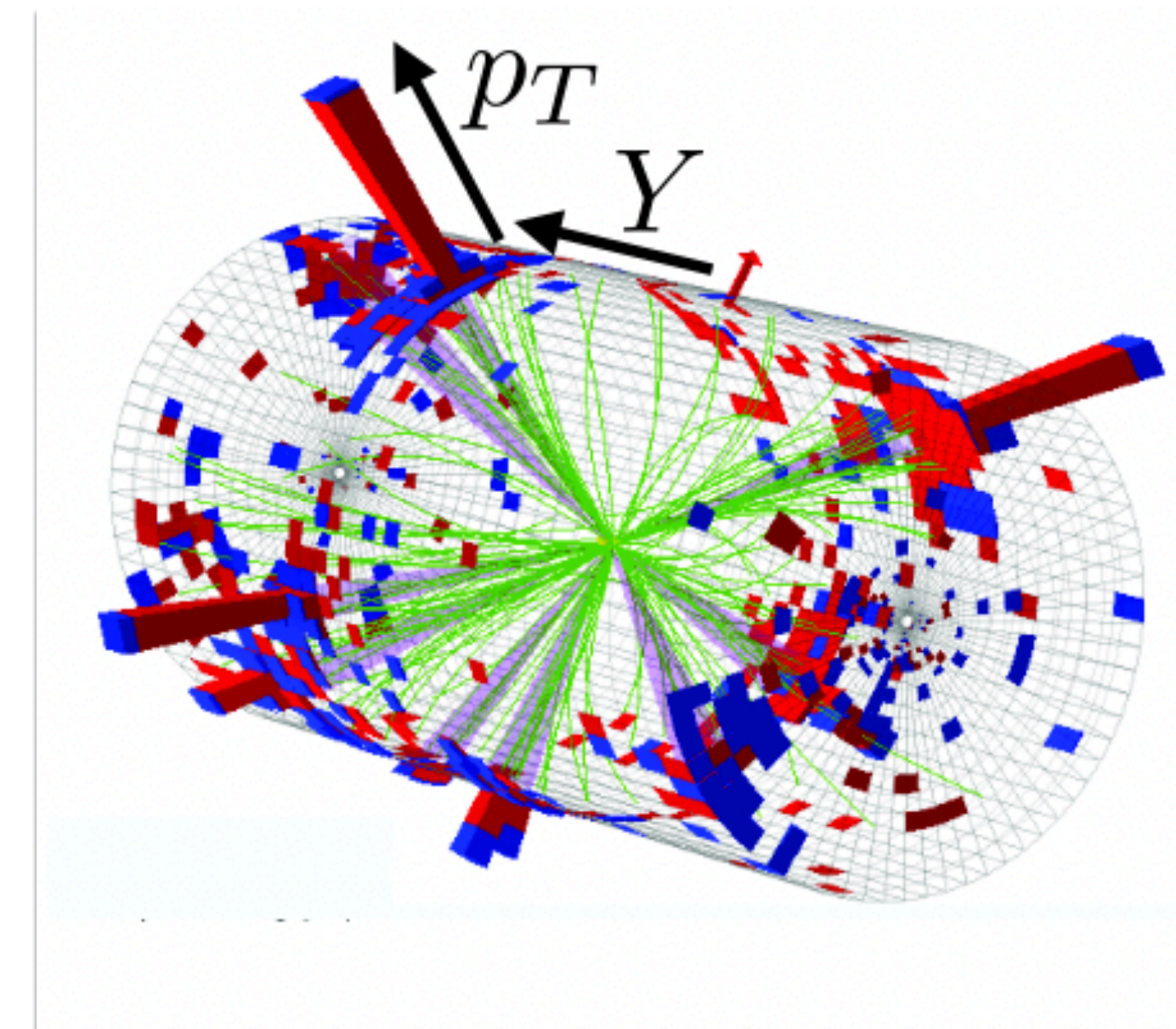
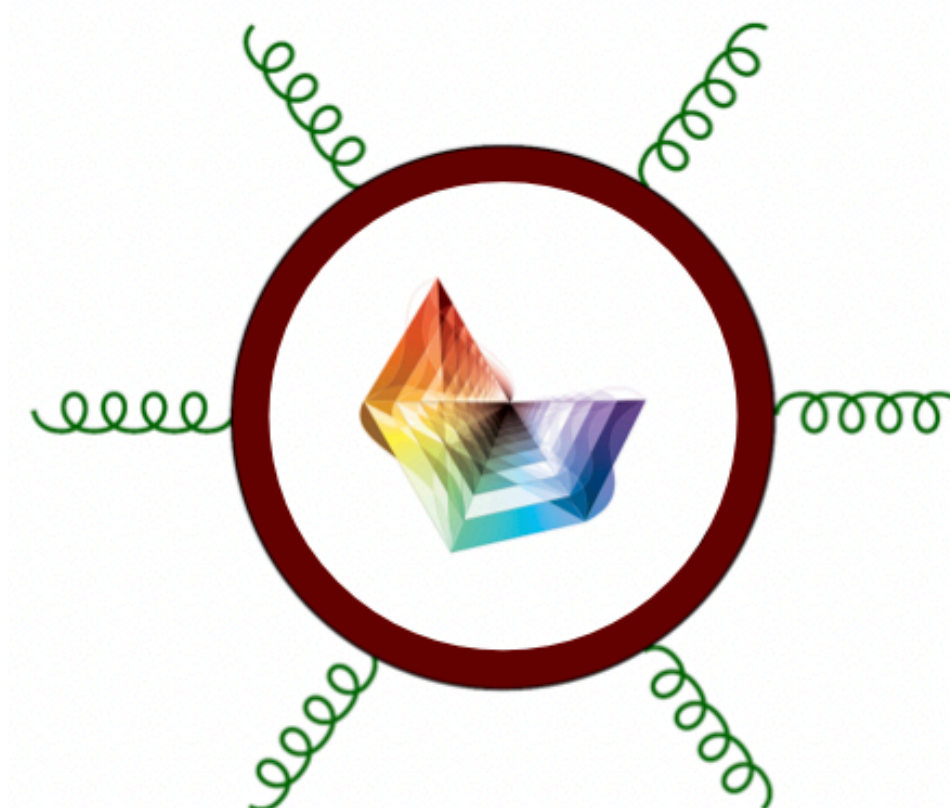
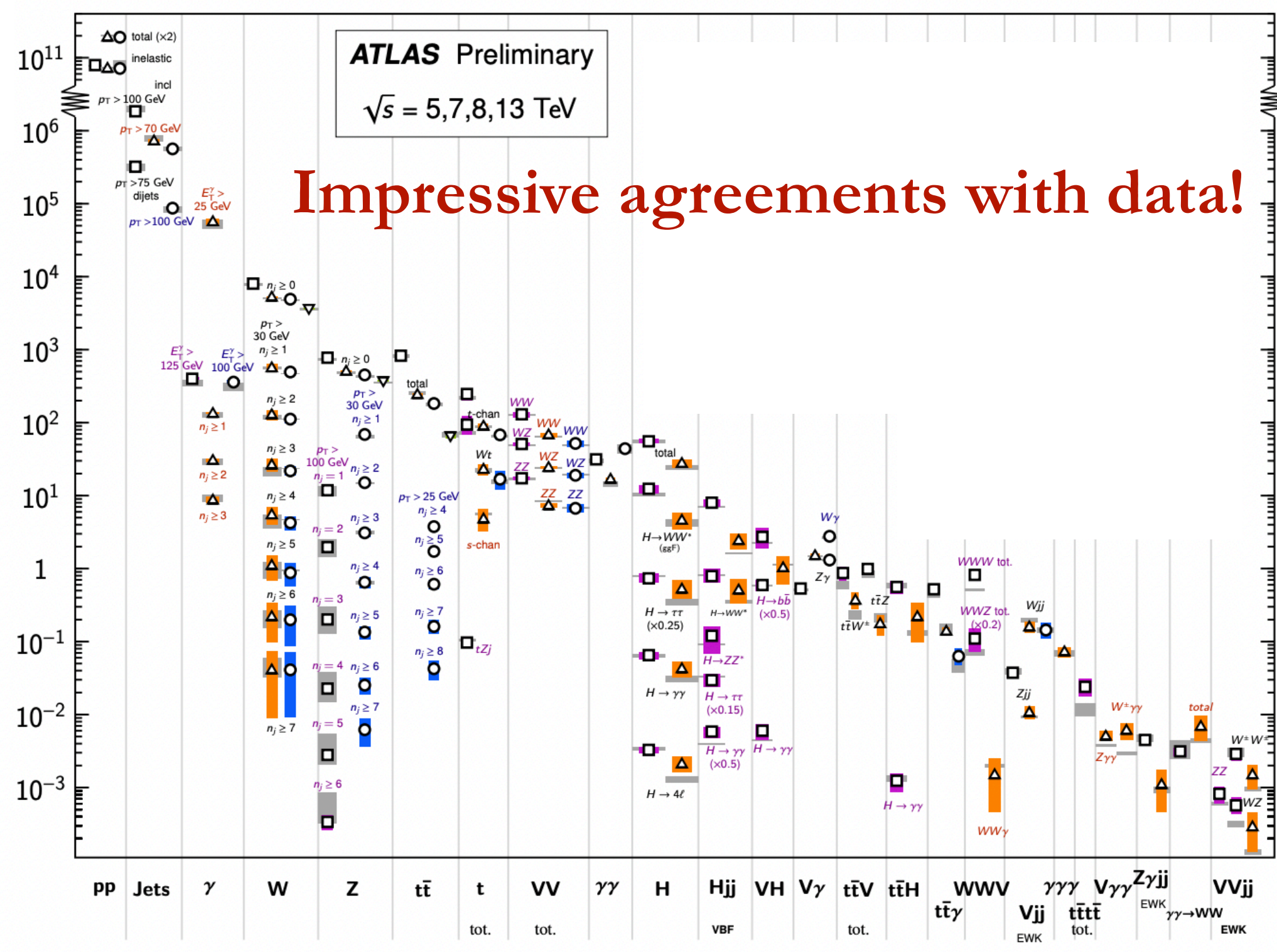
Serman-Weinberg jets played a crucial role in formulating the first IRC definition to study energy flow, or jets



JETS AT COLLIDERS

- The effort to achieve precise predictions of jet cross sections has driven important theoretical developments in Quantum Field Theory

cross-section of SM processes



9-jet at the LHC

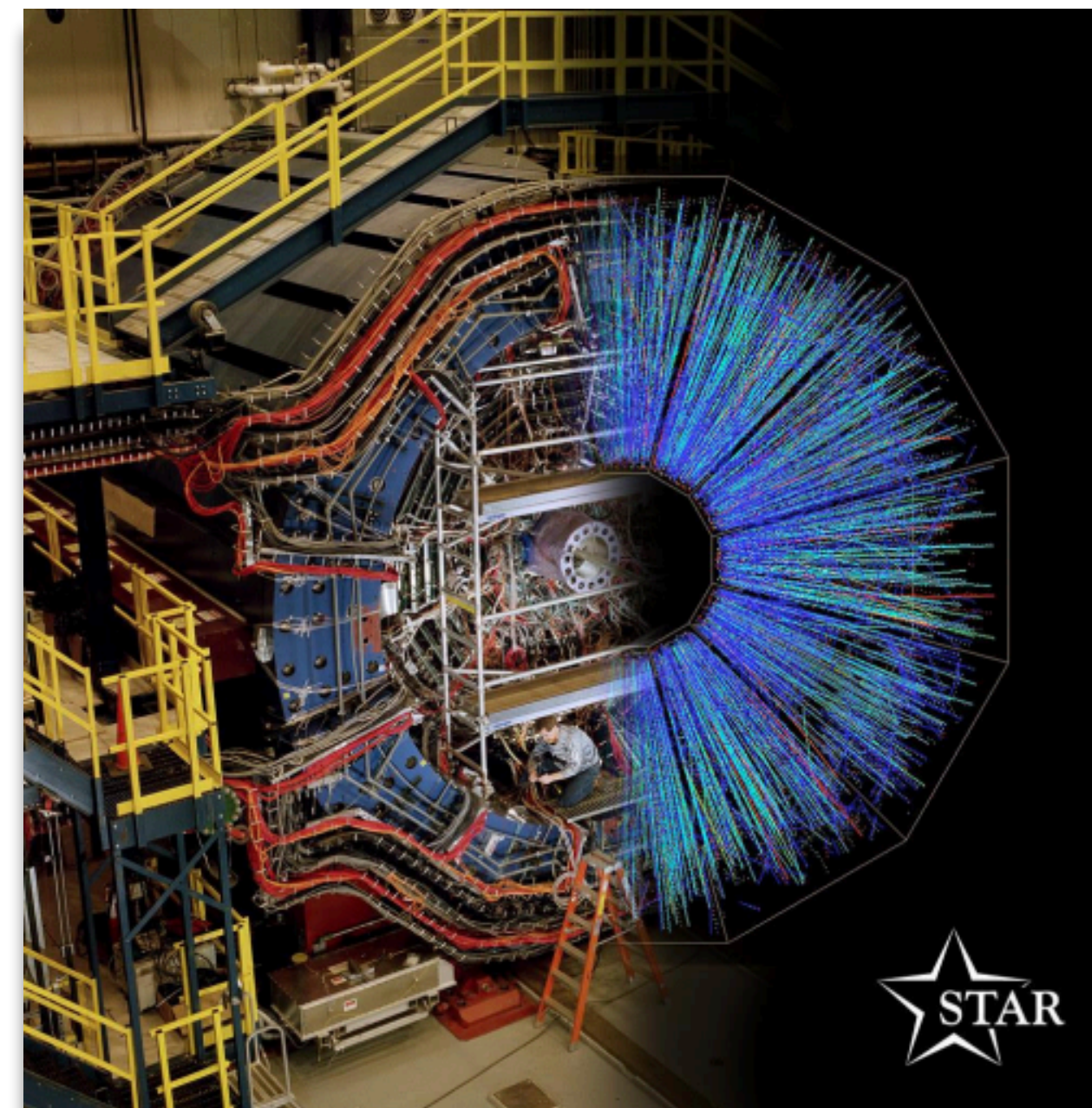
- Field of jet physics (energy flow) have always been intricately connected to the success of the collider physics program!

EXCITING COLLIDER PHYSICS ERA



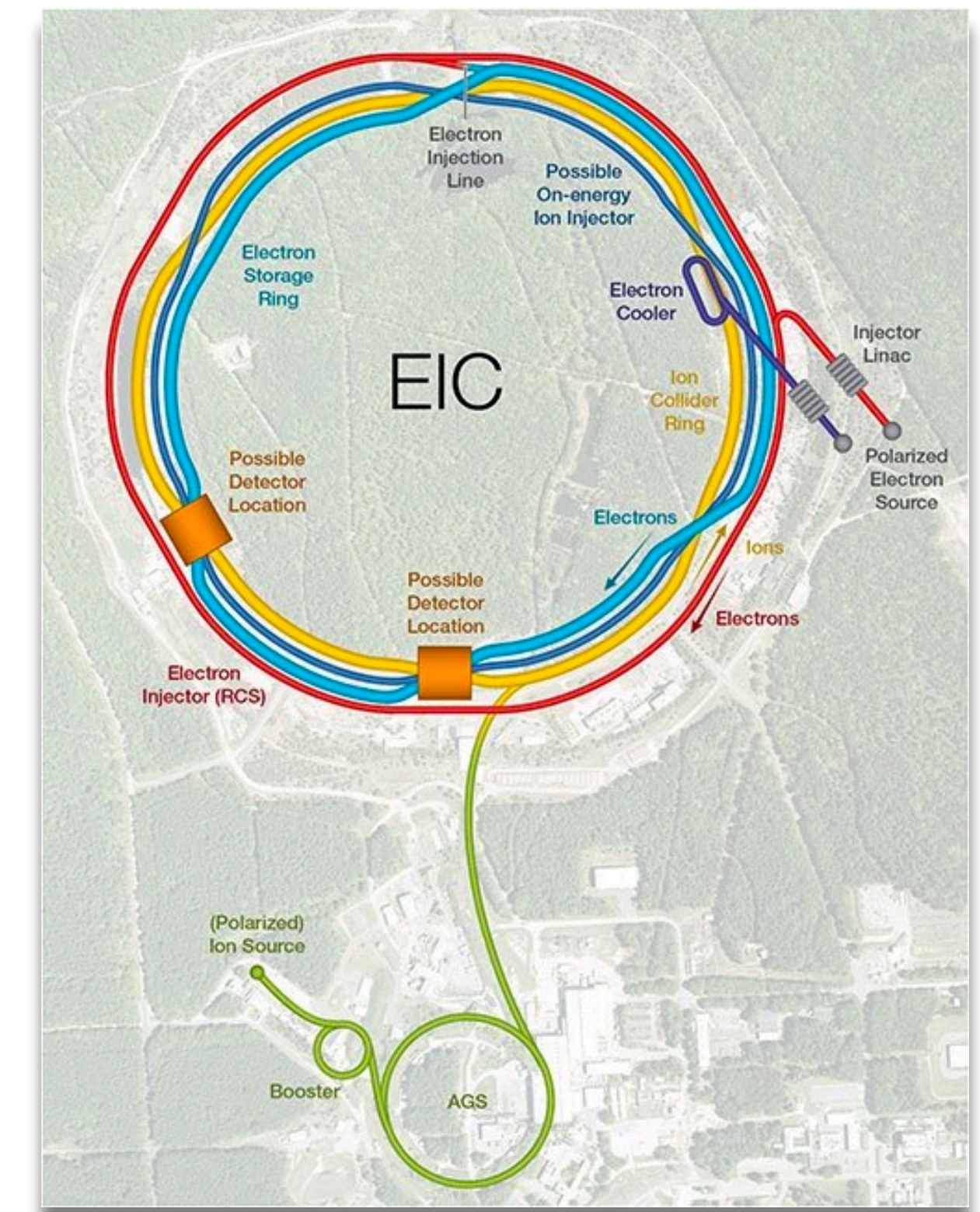
LHC, 2008 - Present

Run 3 running! ✨



RHIC, 2000 - Present

sPHENIX: 2024- ✨



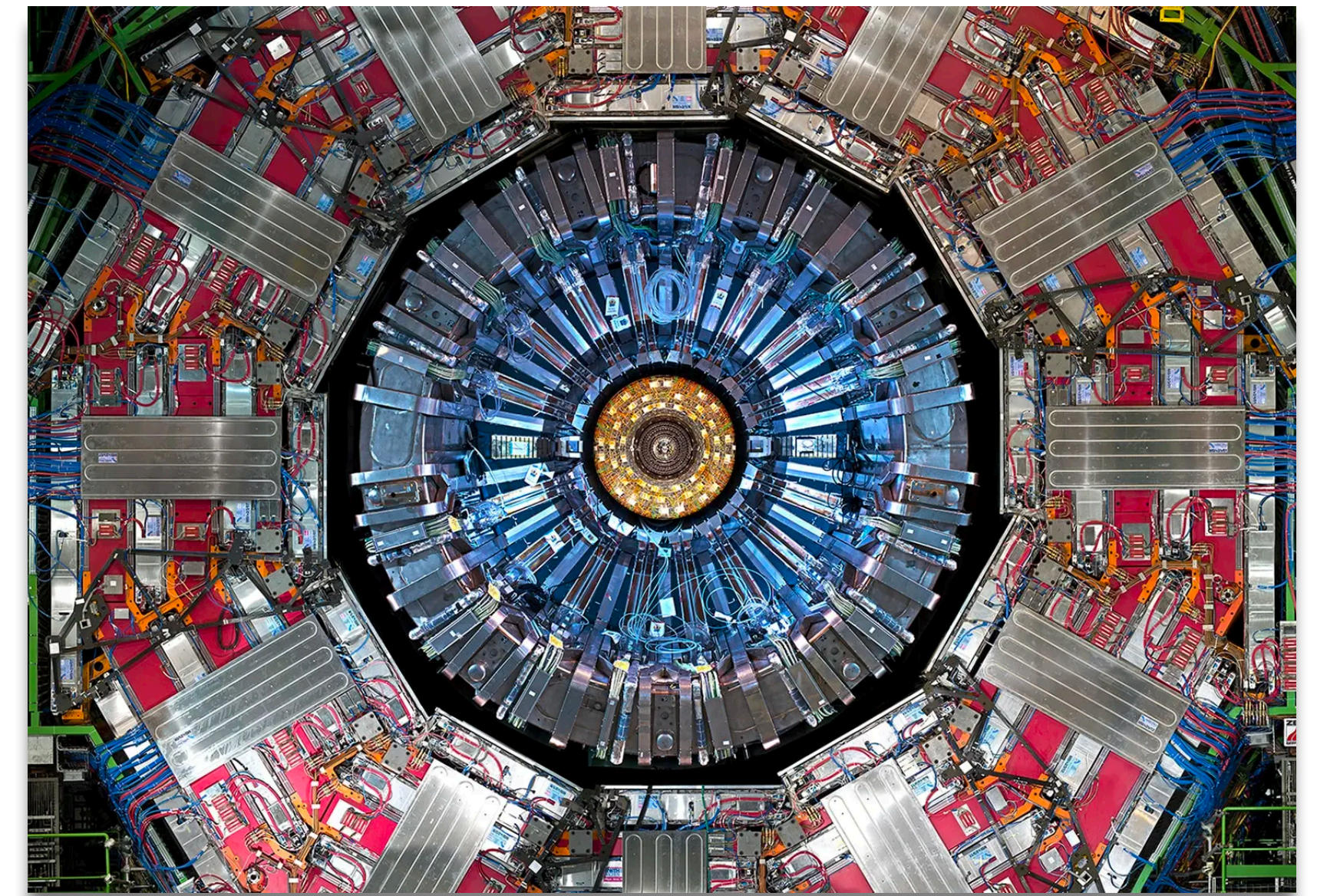
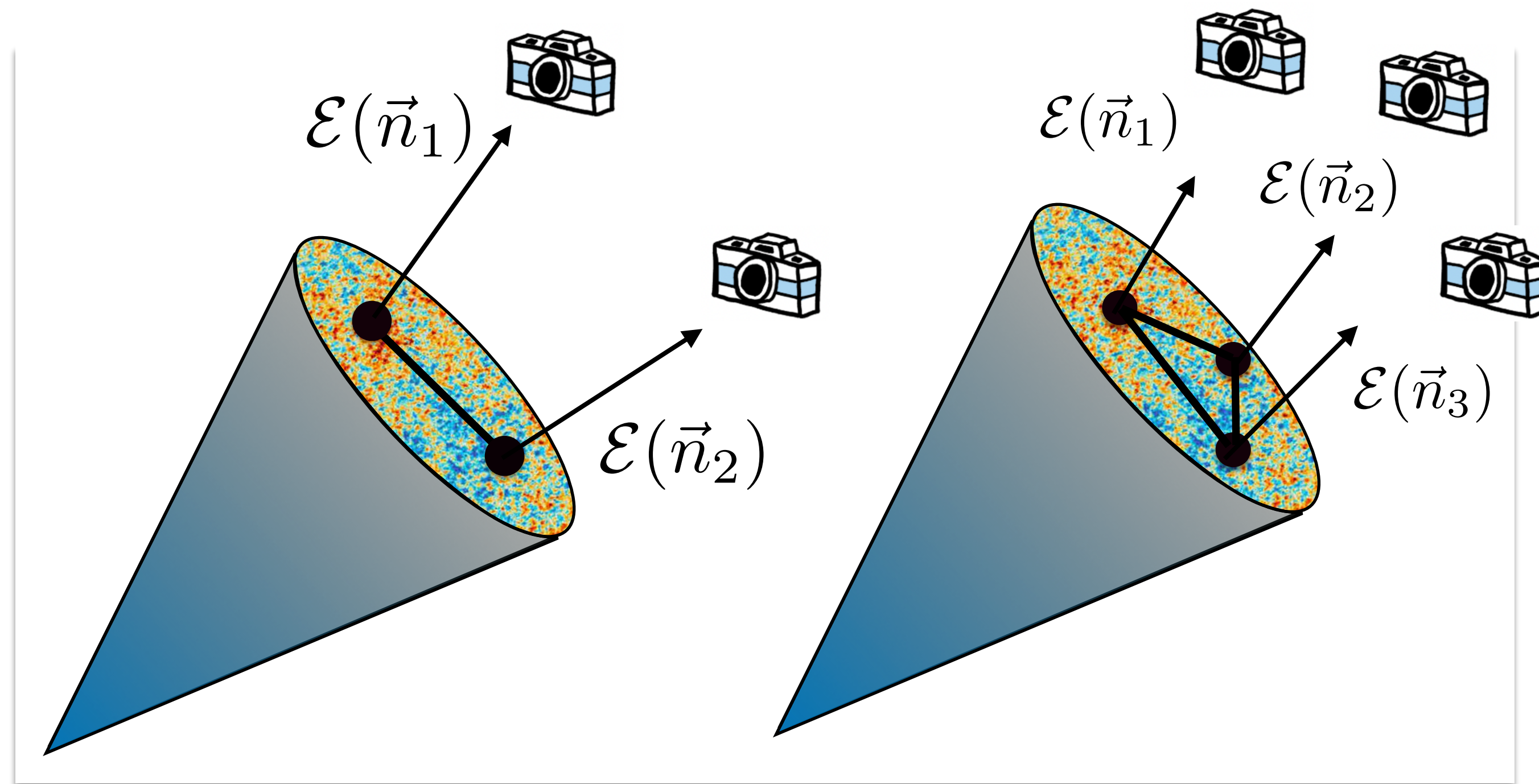
EIC, 2030s- ✨

- **Jets at colliders give us the means to probe field theory in data!**

How can we harness jets to continue making breakthroughs in collider frontier?

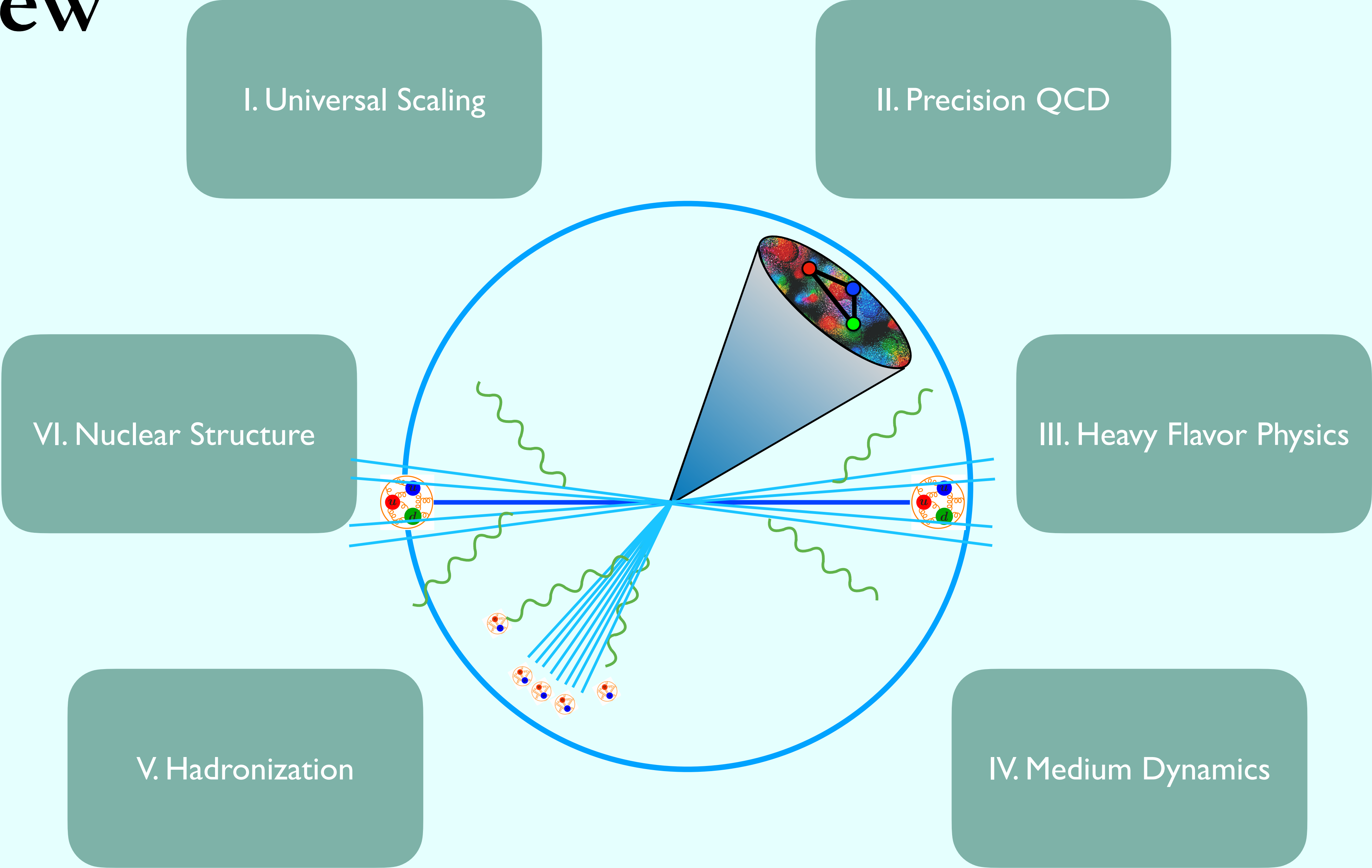
JET SUBSTRUCTURE: STUDYING ENERGY FLOW WITHIN JETS

- **Modern detectors with spectacular angular resolution** gives us an **unprecedented opportunity to peer into the energy flow within jets**

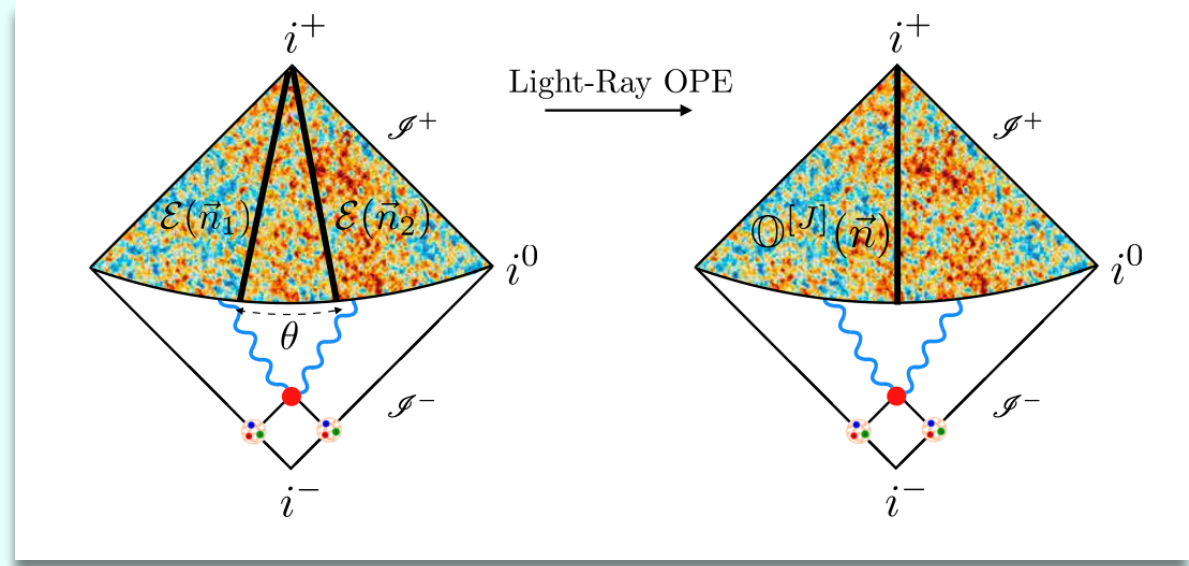


- **Relative to inclusive jet cross-section, or one-point energy correlation, jet substructure gives us opportunity to study multi-point correlations of energy within jets**

Overview



Overview



I. Universal Scaling
QFT perspective of jet substructure

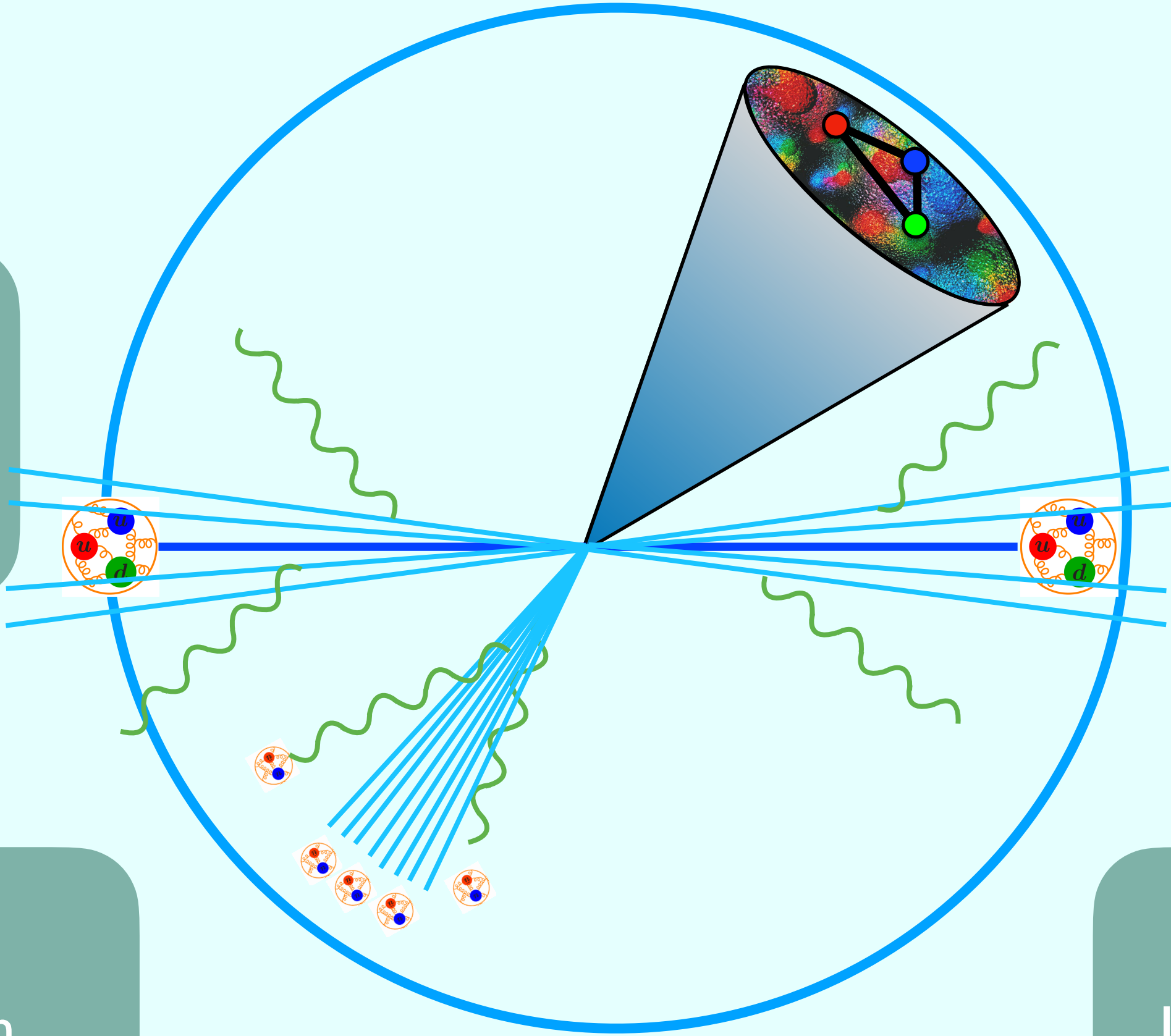
II. Precision QCD

VI. Nuclear Structure

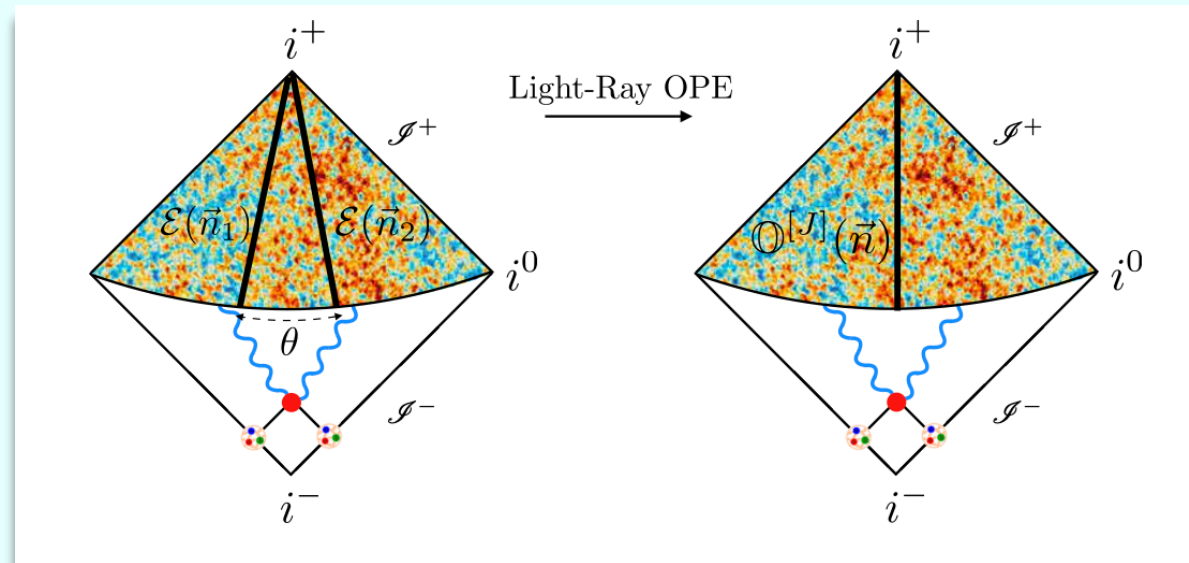
III. Heavy Flavor Physics

V. Hadronization

IV. Medium Dynamics

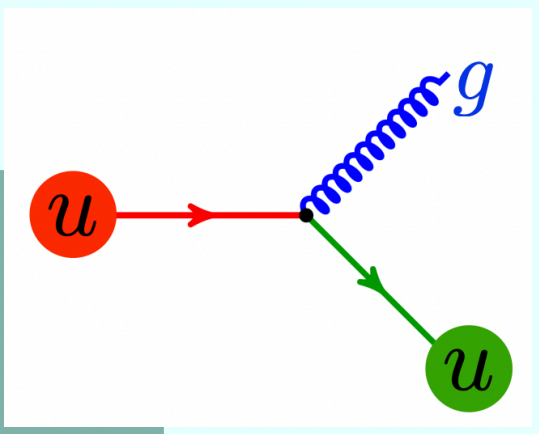


Overview



I. Universal Scaling
QFT perspective of jet substructure

II. Precision QCD
Precise determination of α_s

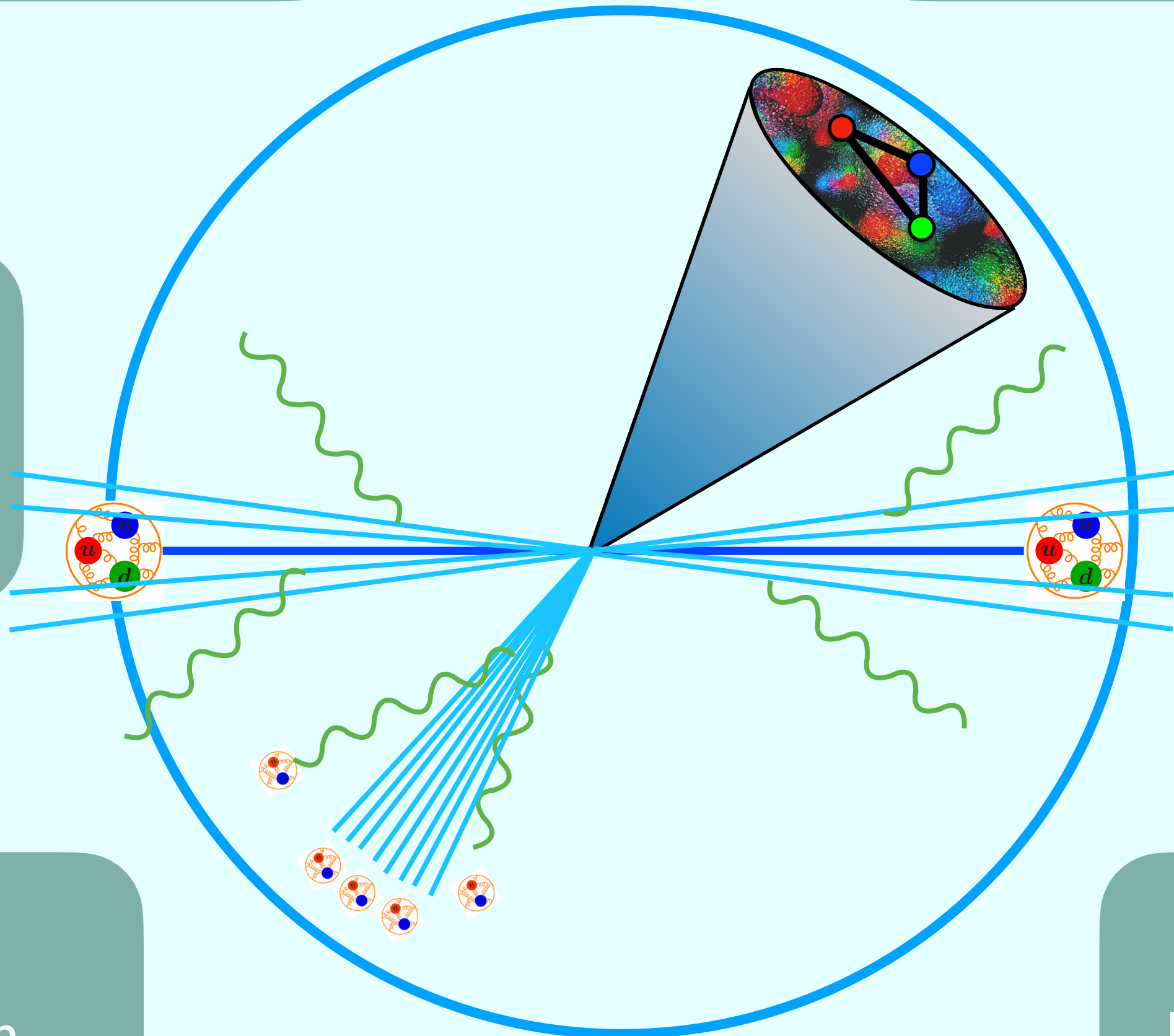


VI. Nuclear Structure

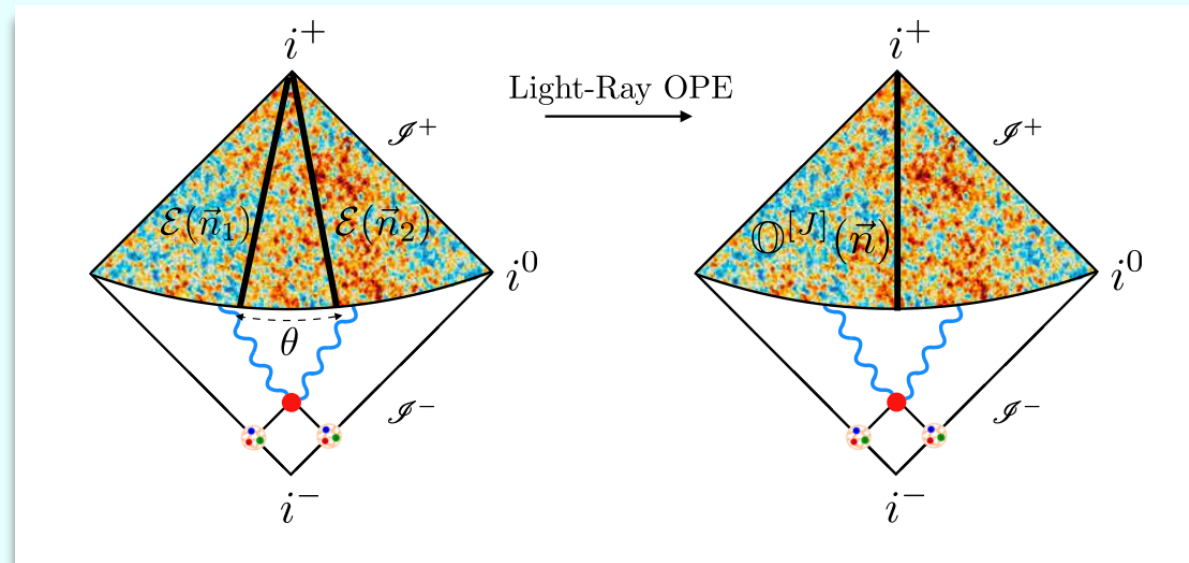
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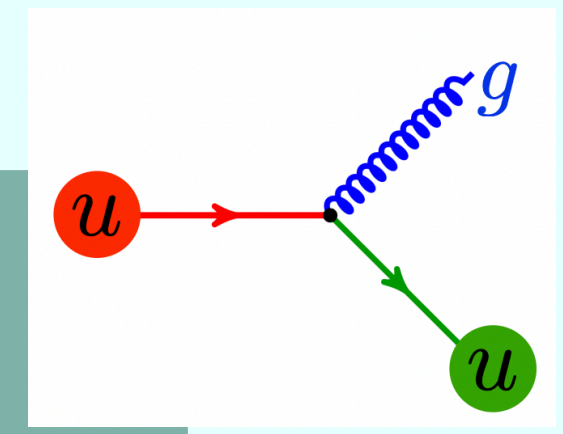


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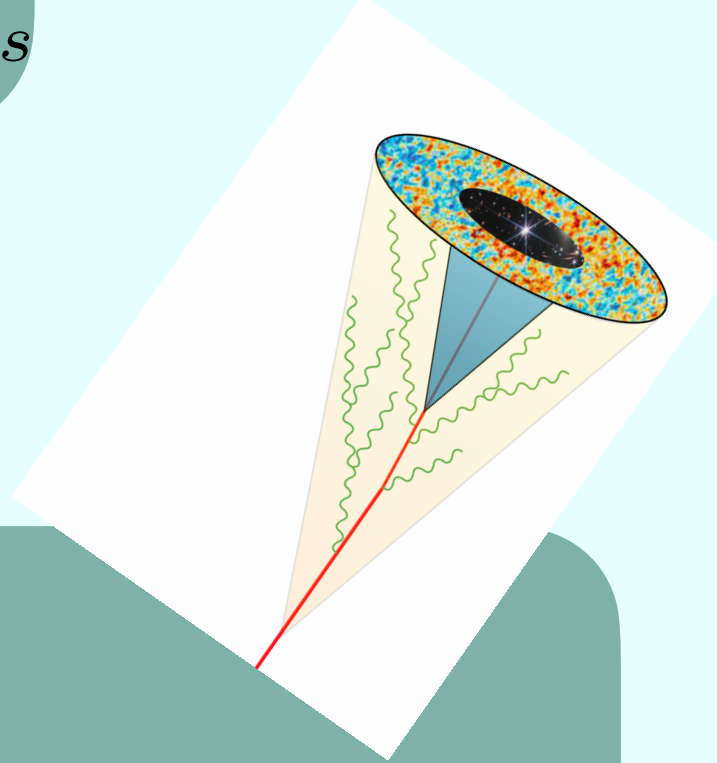
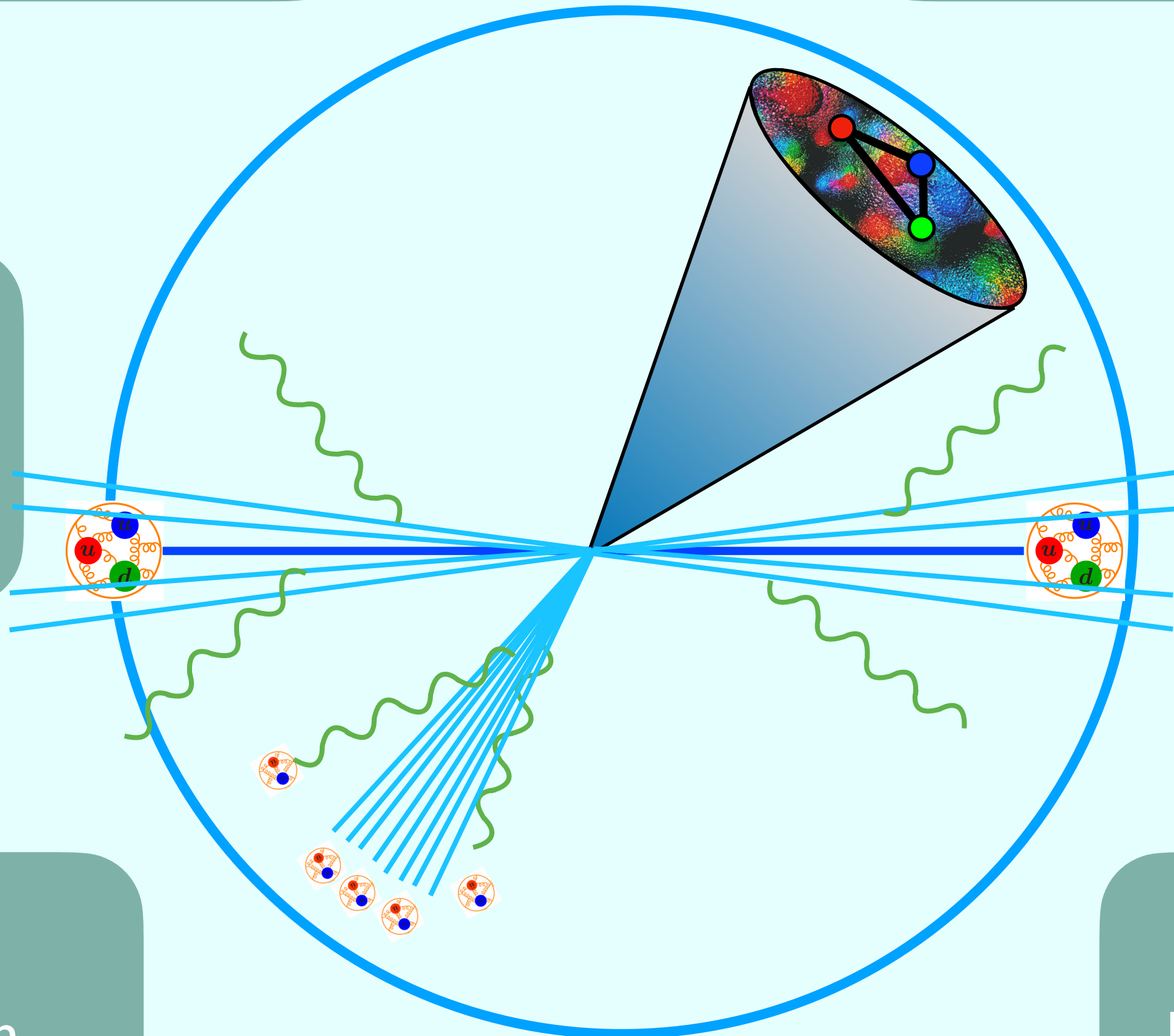


VI. Nuclear Structure

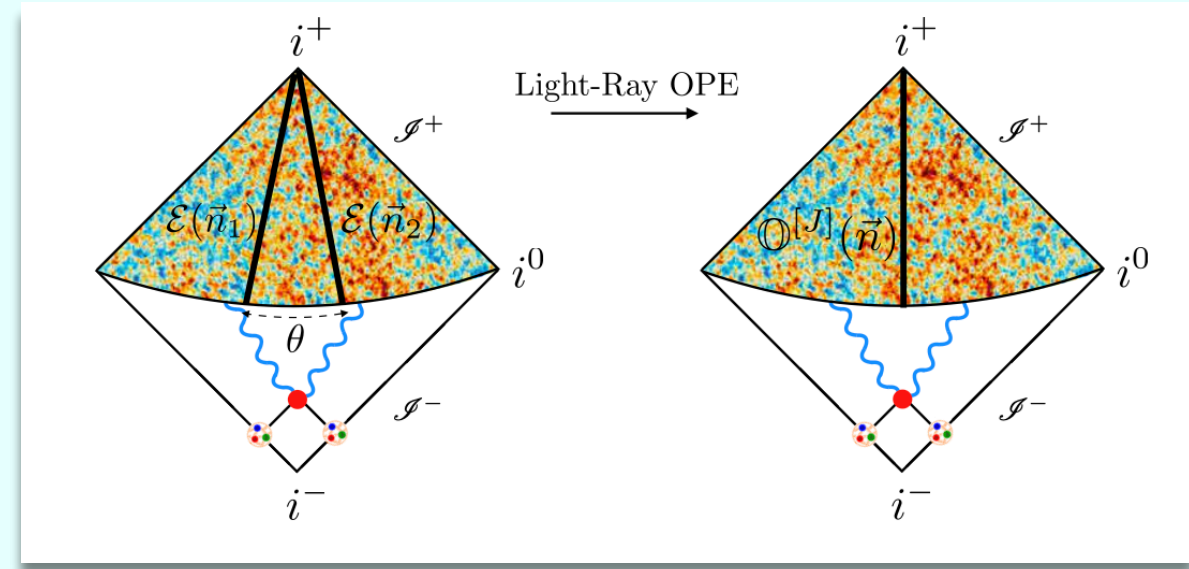
III. Heavy Flavor Physics
Revealing dead-cone

V. Hadronization

IV. Medium Dynamics

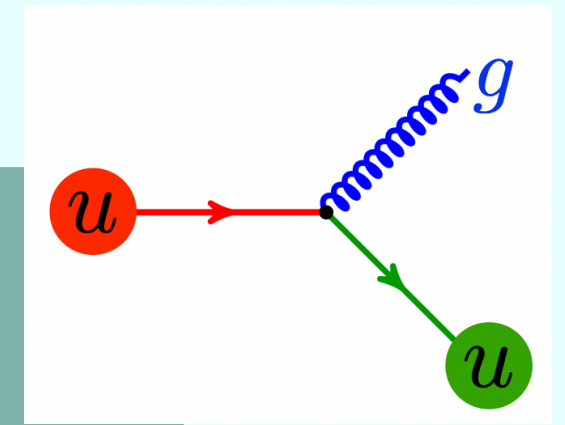


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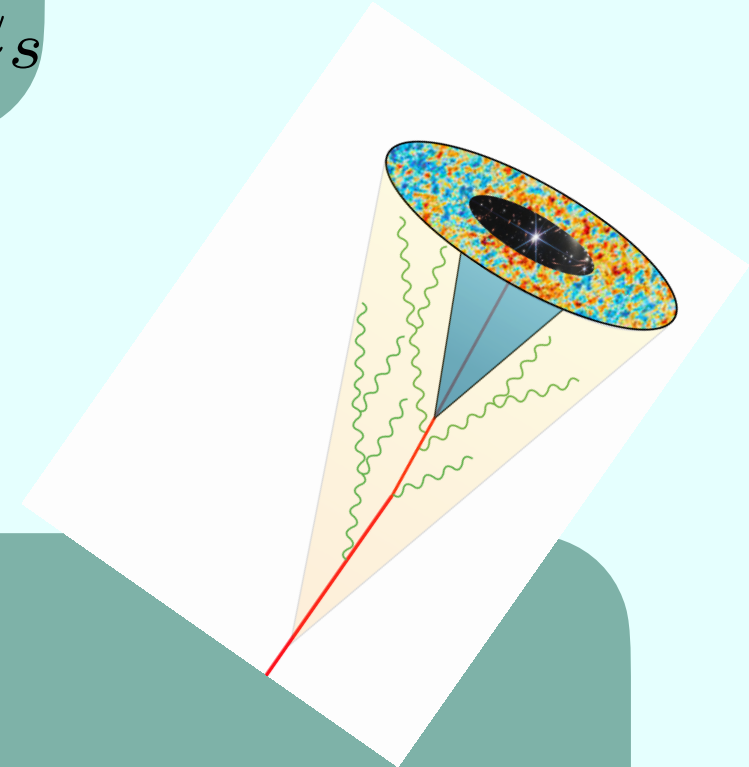
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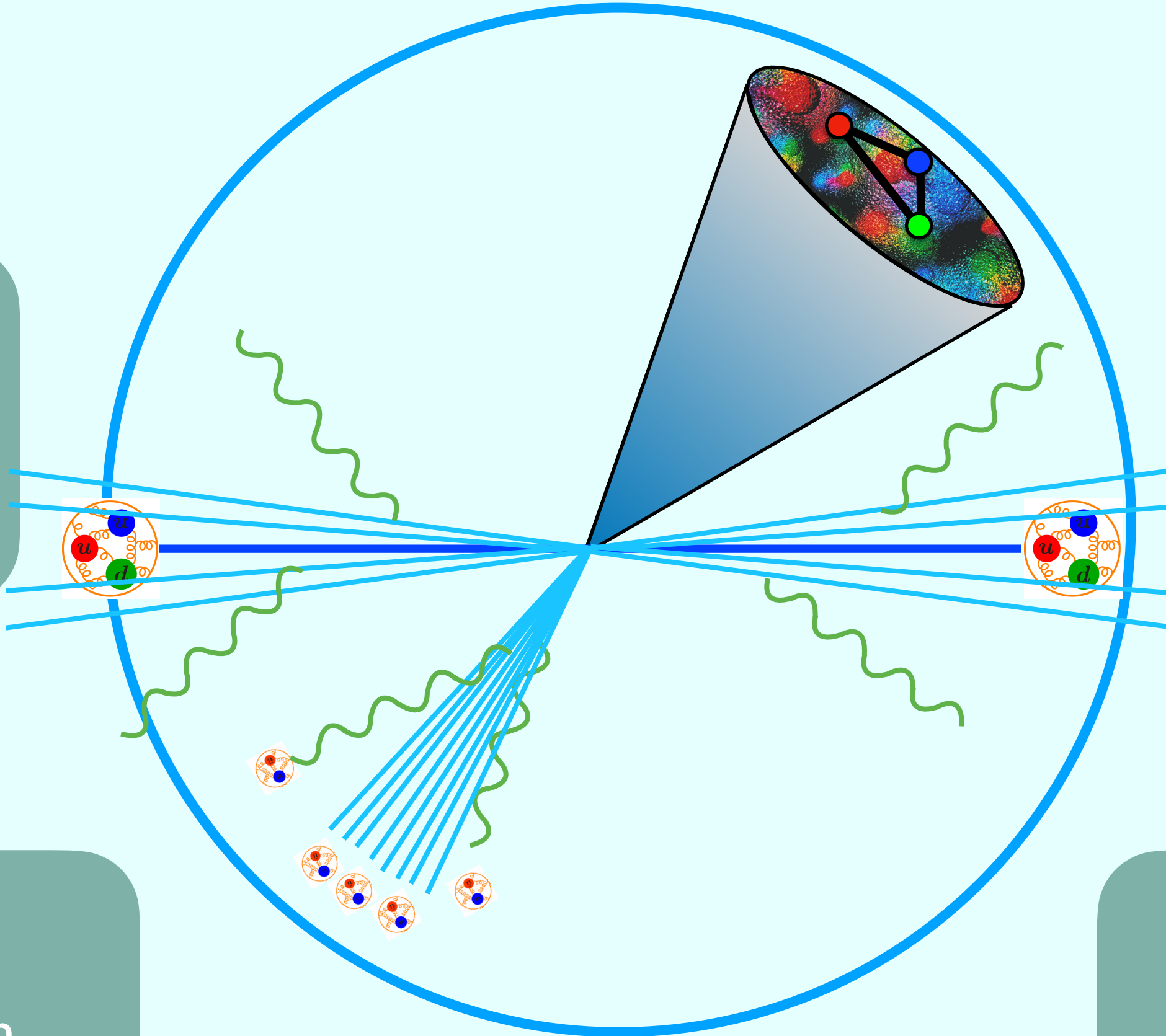
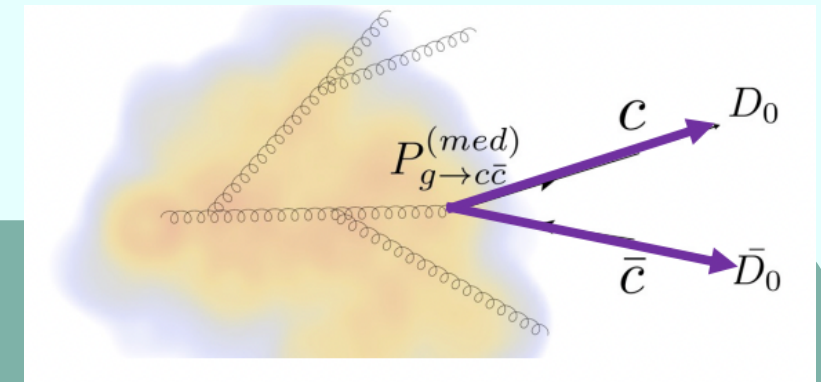
VI. Nuclear Structure

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Revealing dead-cone

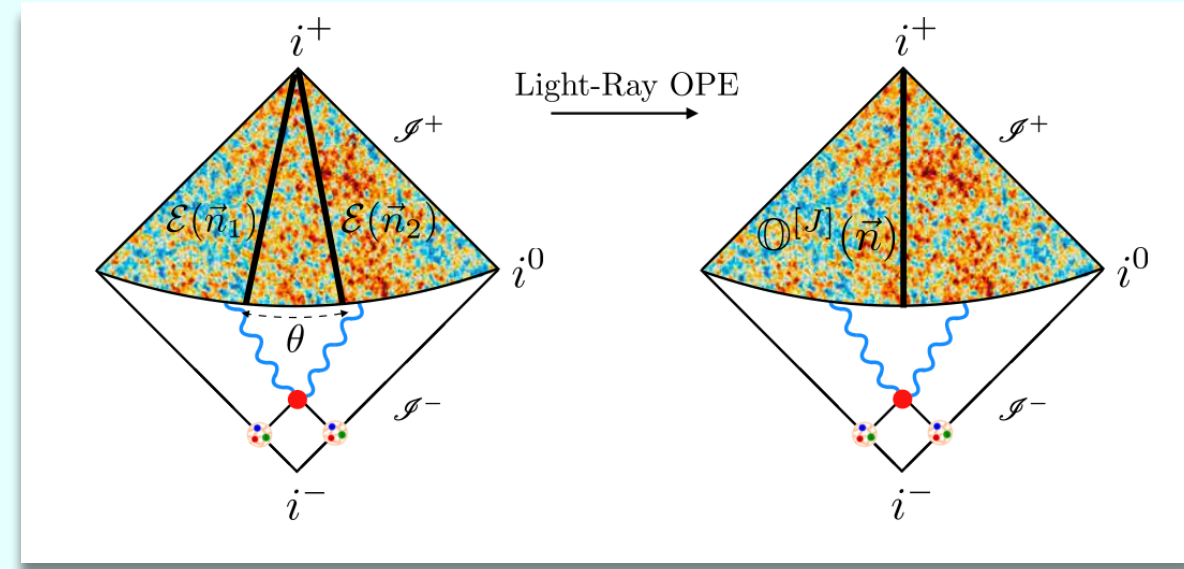


V. Hadronization

IV. Medium Dynamics
Revealing medium scale and modifications

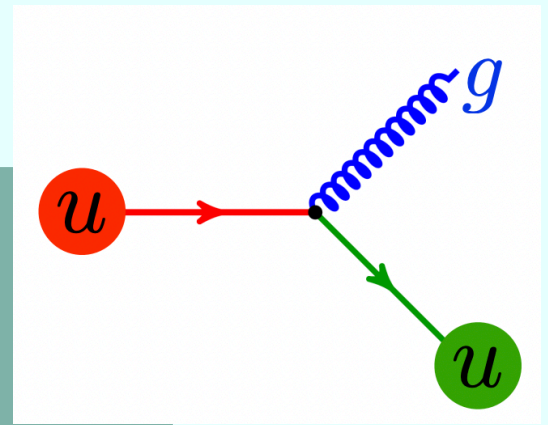


Overview



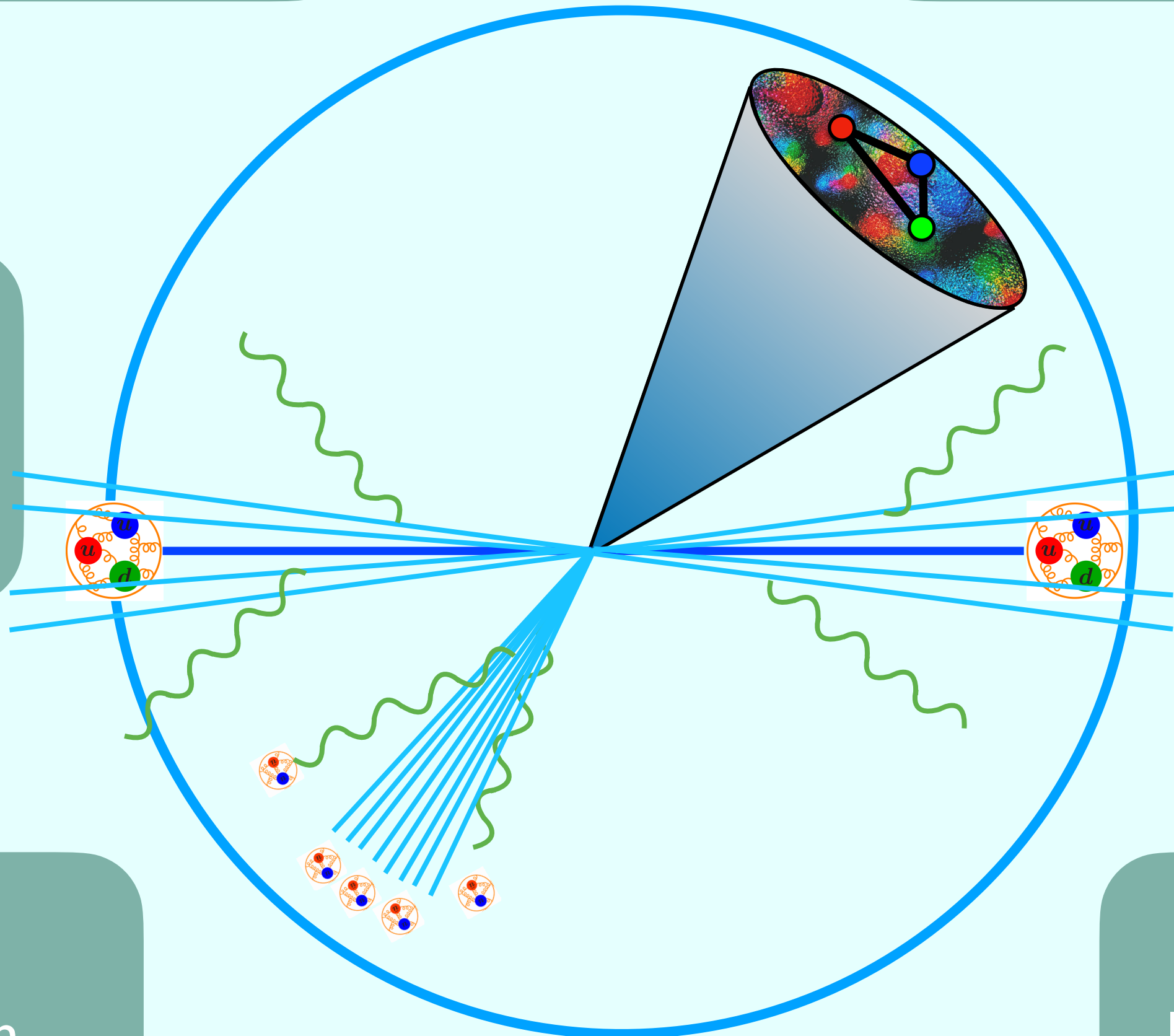
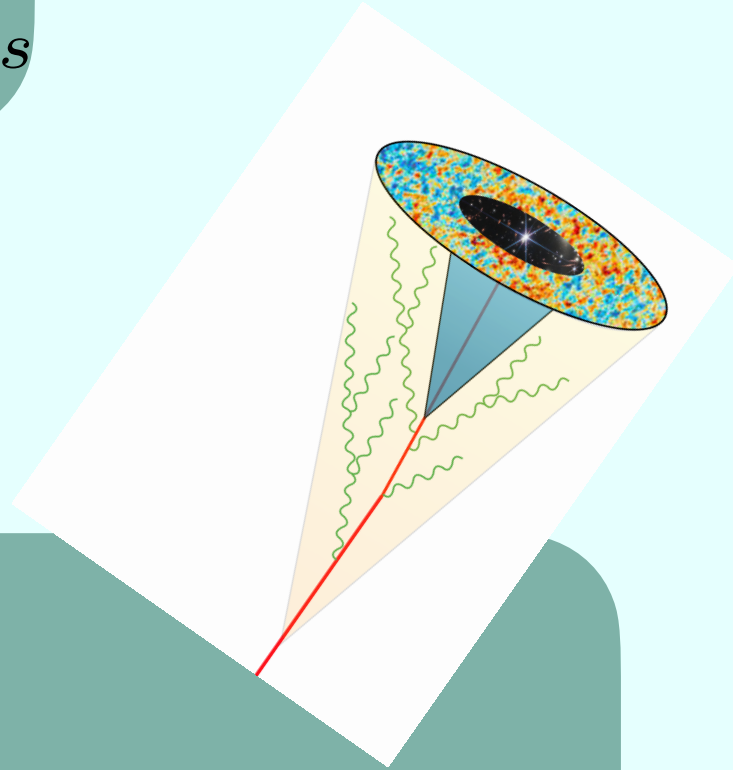
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 QFT perspective of jet substructure

II. Precision QCD
 Precise determination of α_s



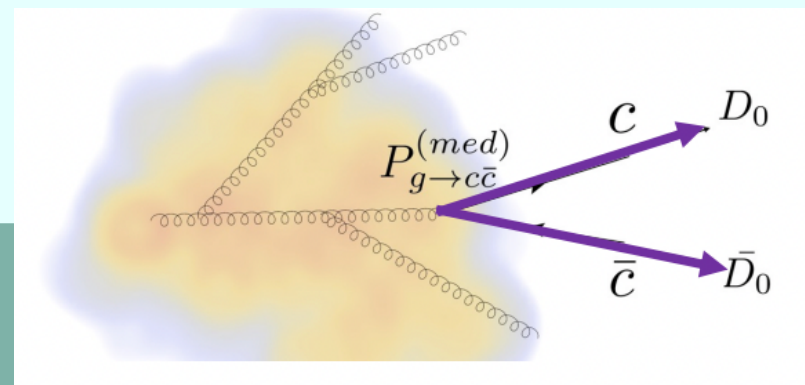
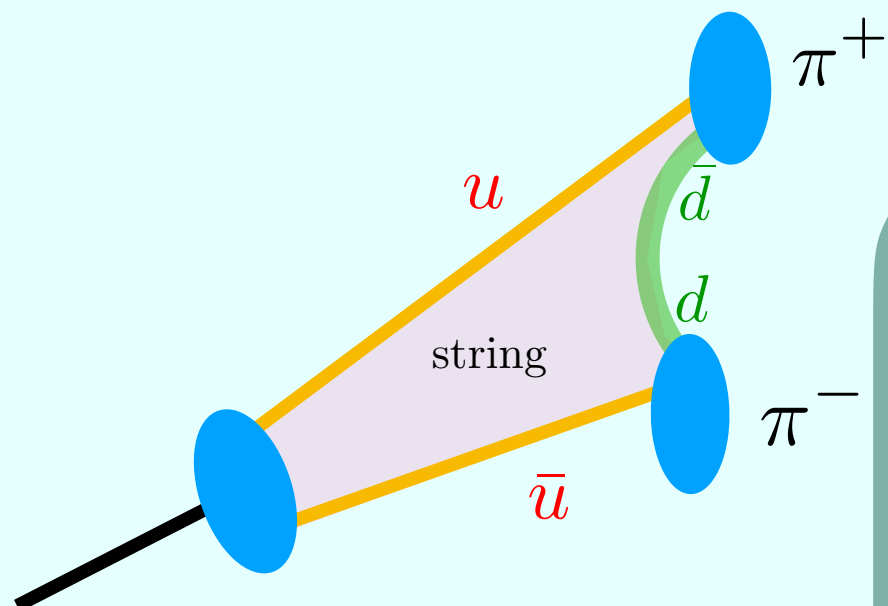
VI. Nuclear Structure

III. Heavy Flavor Physics
 Revealing dead-cone

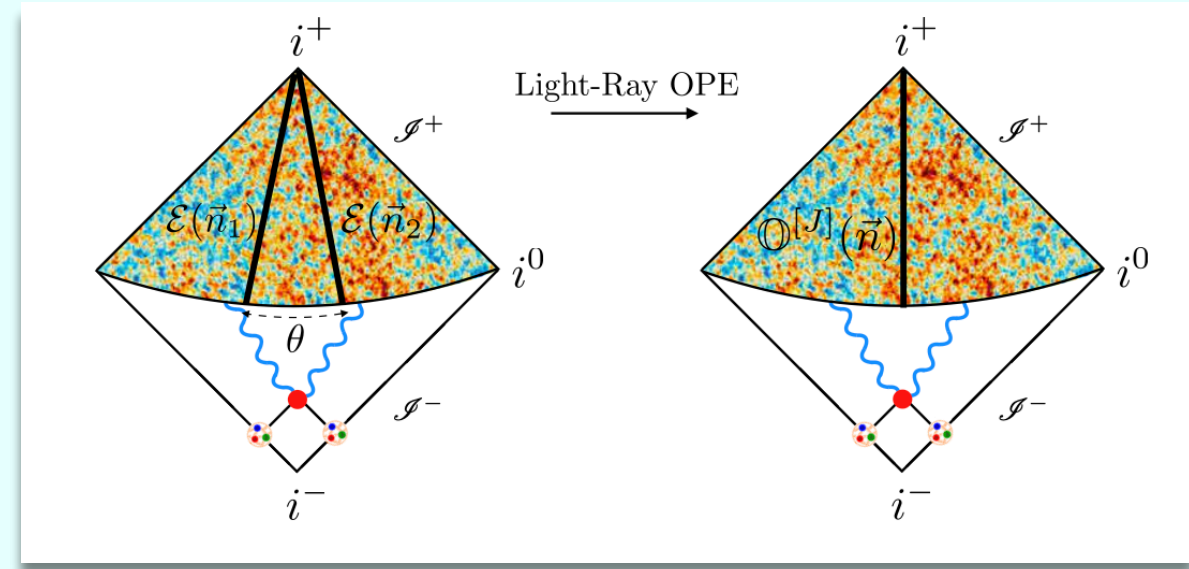


V. Hadronization
 Discrimination between hadronization mechanisms

IV. Medium Dynamics
 Revealing medium scale and modifications

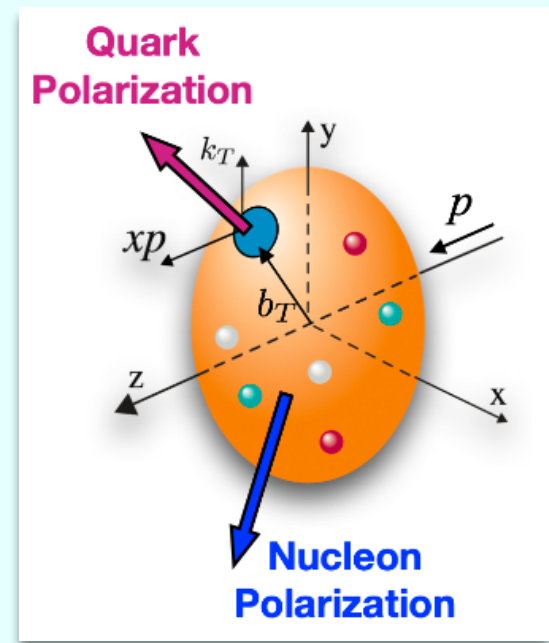
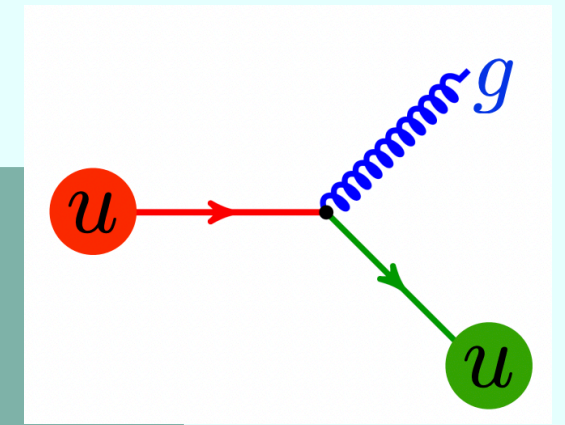


Overview



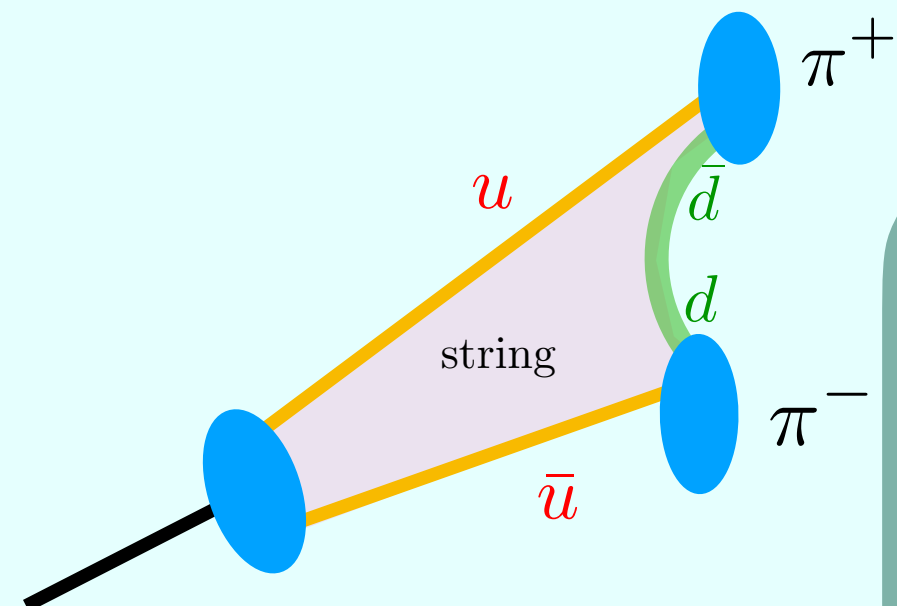
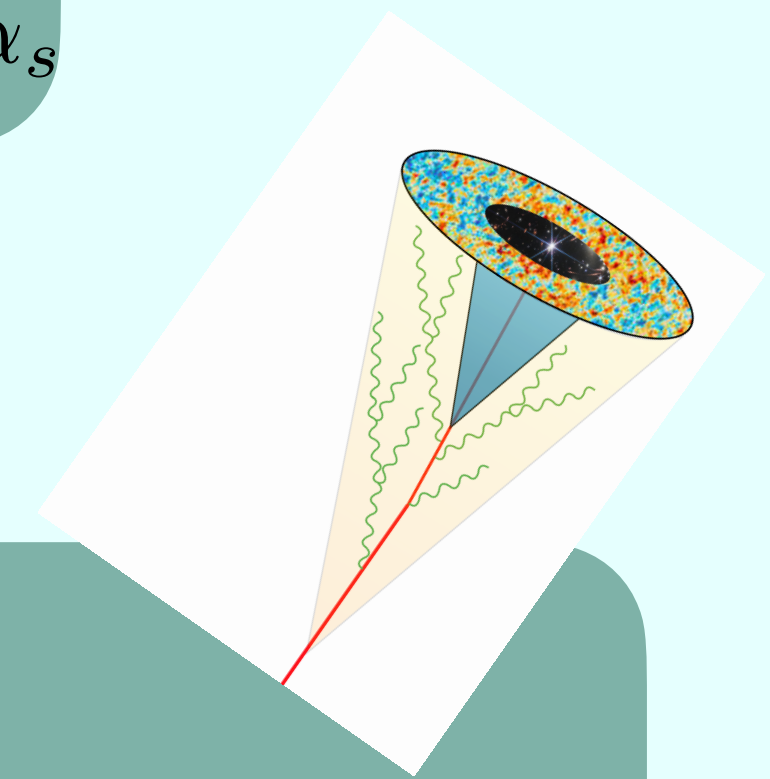
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QFT perspective of jet substructure

II. Precision QCD
Precise determination of α_s



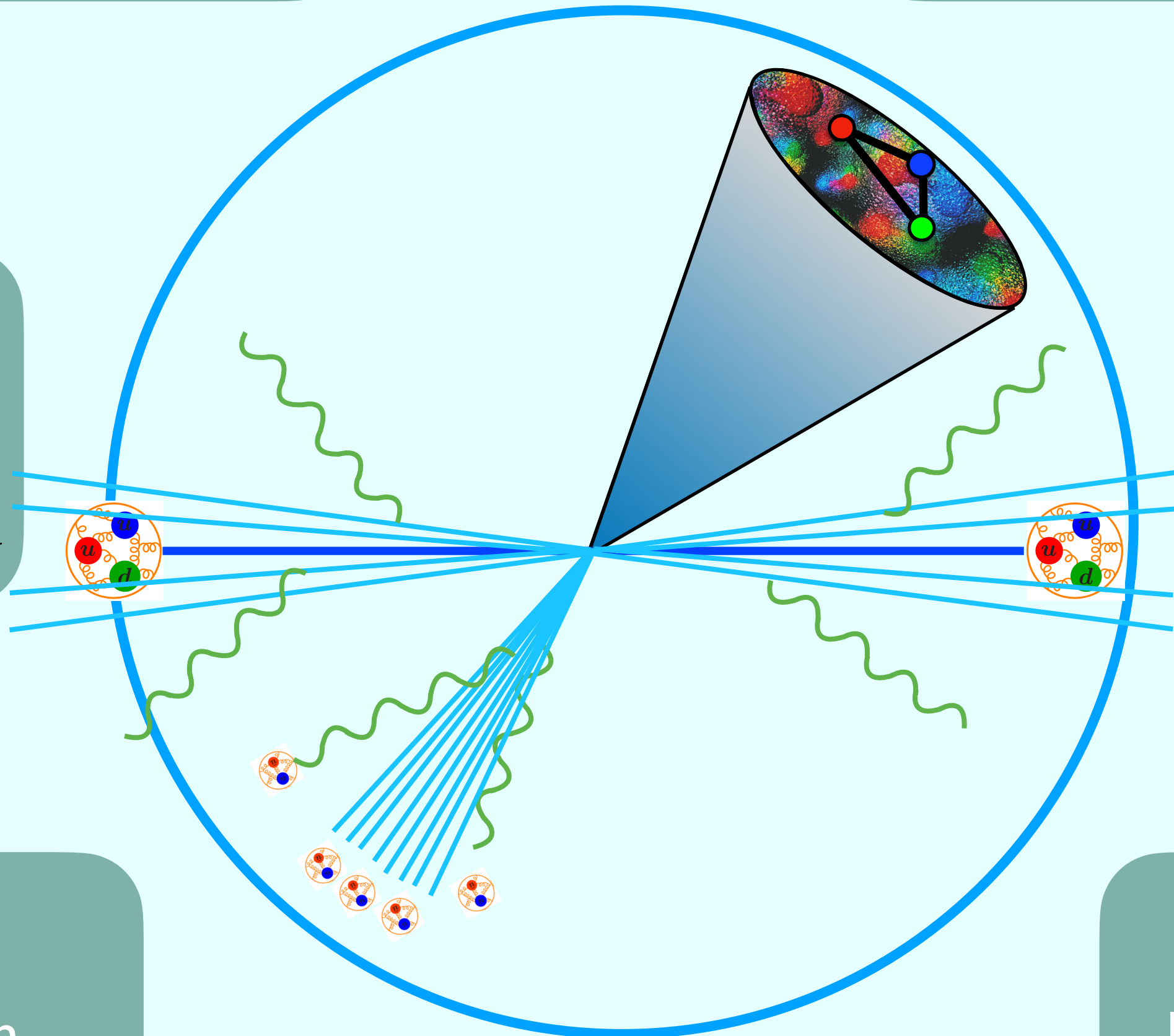
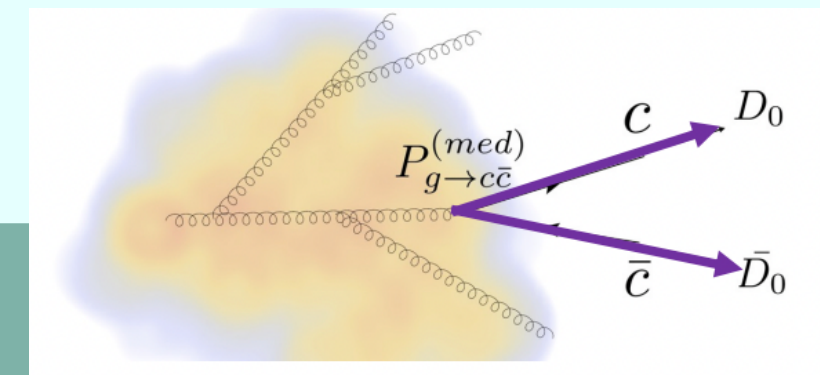
VI. Nuclear Structure
Nuclear tomography study

III. Heavy Flavor Physics
Revealing dead-cone



V. Hadronization
Discrimination between hadronization mechanisms

IV. Medium Dynamics
Revealing medium scale and modifications



Overview

I. Universal Scaling

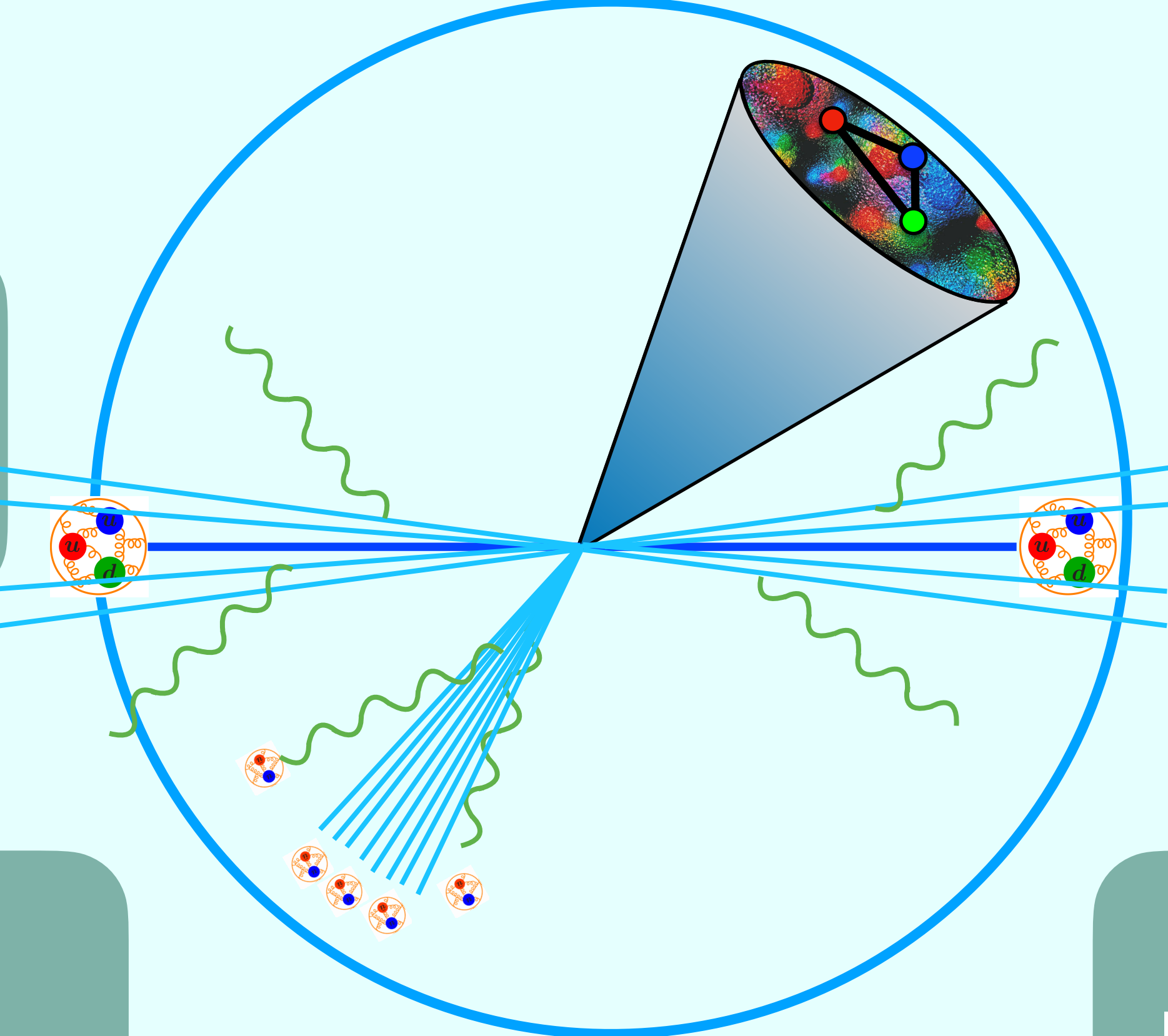
II. Precision QCD

VI. Nuclear Structure

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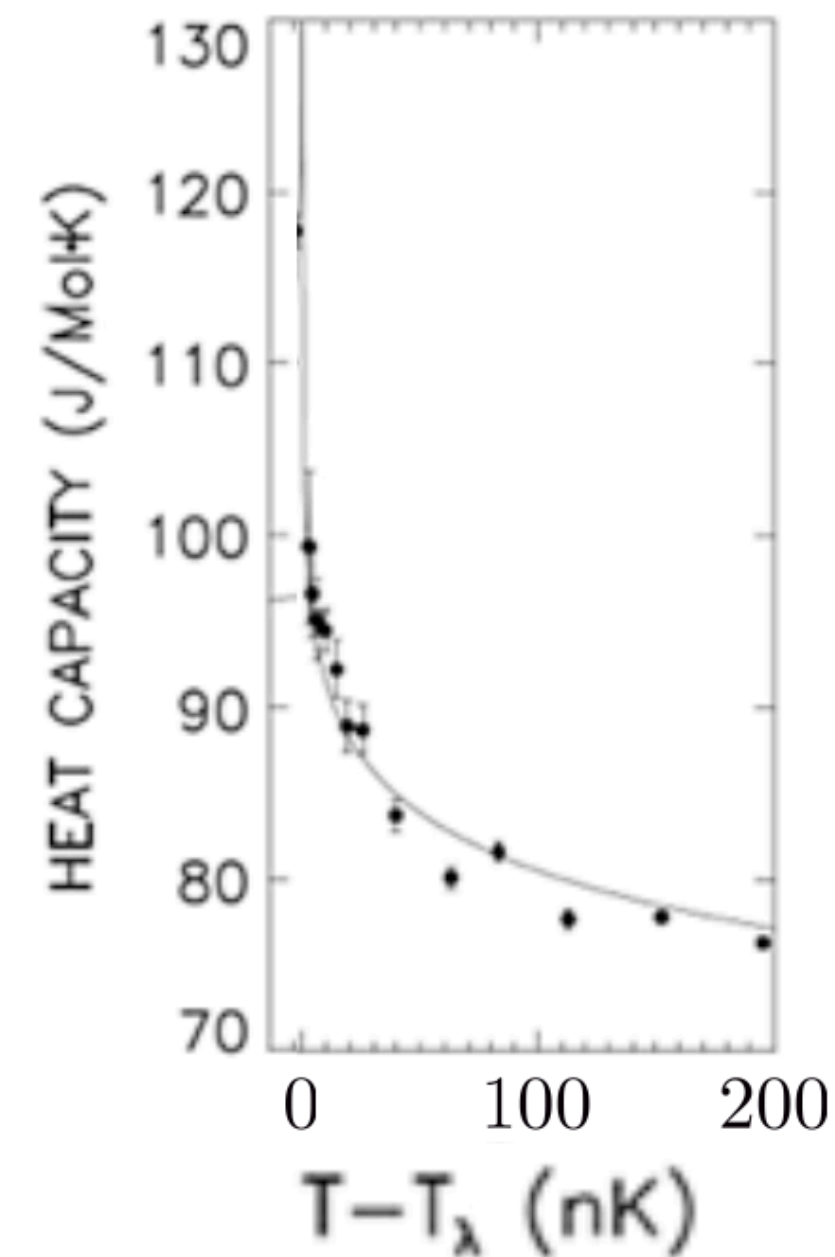
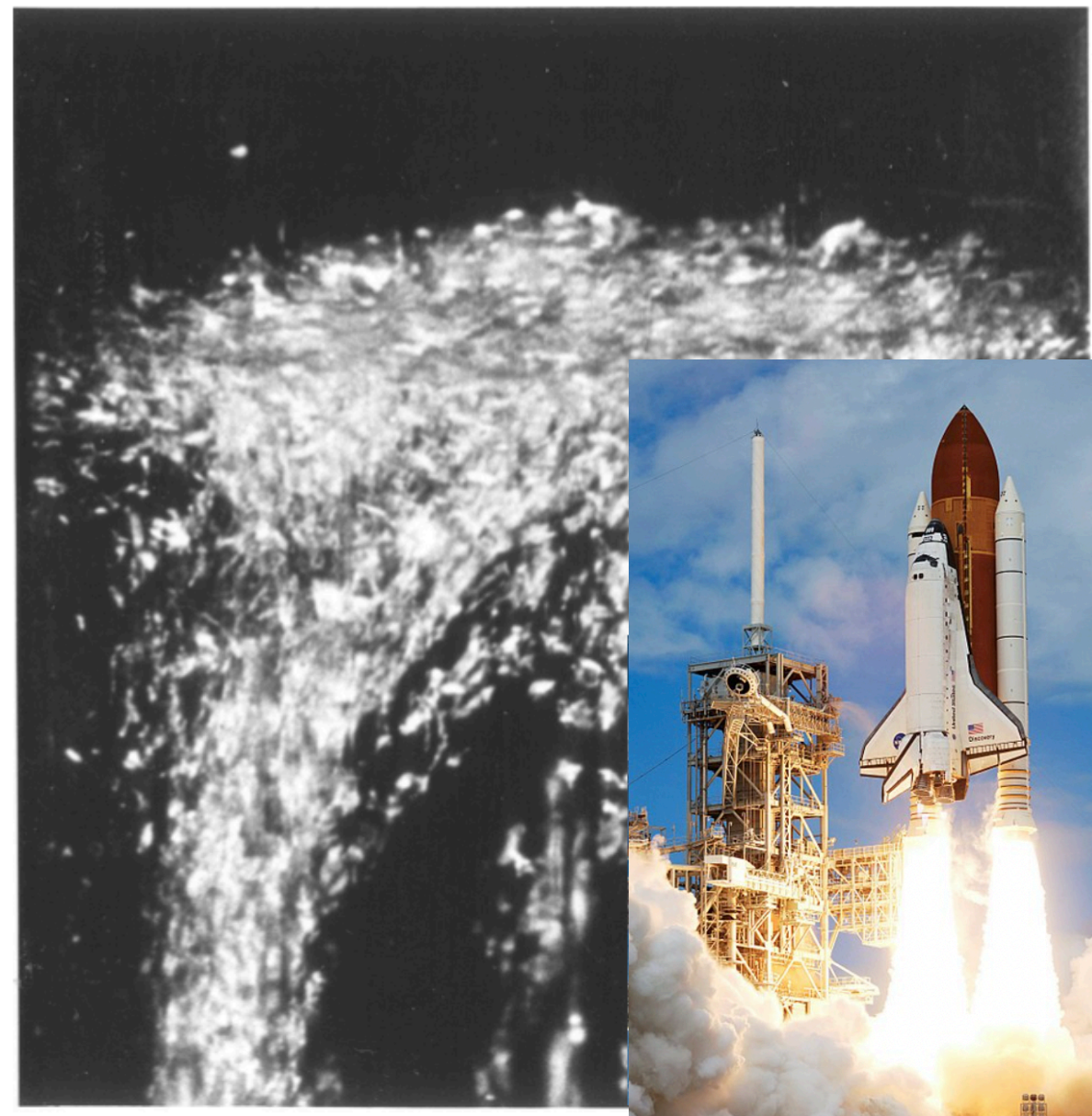


SCALING BEHAVIOR IN QFT

- Why is the study of jet substructure of interest in QFT?
- QFTs display **universal scaling** behaviors when operators approach one another



Wilson '70



Euclidean Operator Product Expansion

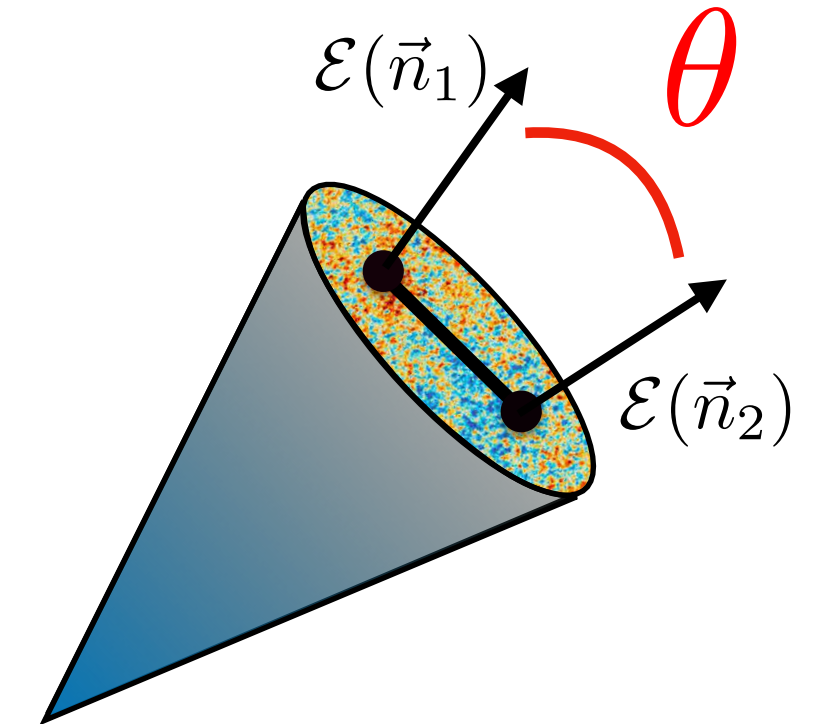
$$\mathcal{O}(x)\mathcal{O}(0) = \sum x^{\gamma_i} c_i \mathcal{O}_i$$

- **Critical phenomena** give us access to **universal scaling behavior** as **Euclidean operators** are brought together

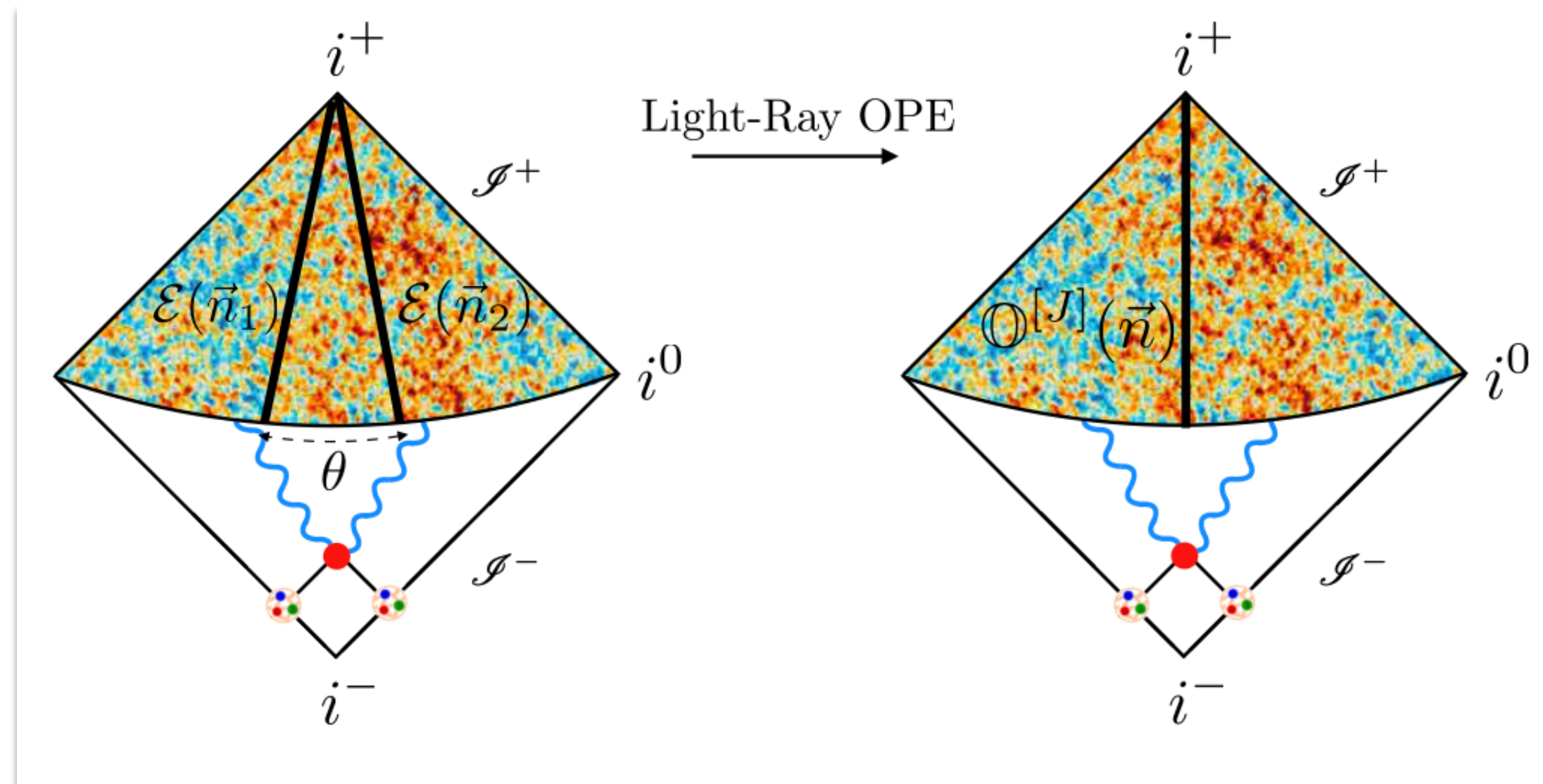
UNIVERSAL LORENTZIAN SCALING WITHIN JETS

- **Jet substructure** describes the limit where **energy flow operators** are brought together, thus probing the **OPE limit of Lorentzian operators**

⇒ **Profound field theory predictions within jets!**



Hofman, Maldacena '08



Light-ray Operator Product Expansion

$$\mathcal{E}(\hat{n}_1) \mathcal{E}(\hat{n}_2) \sim \sum \theta^{\gamma(3)-2} \mathcal{O}_i(\hat{n}_1)$$

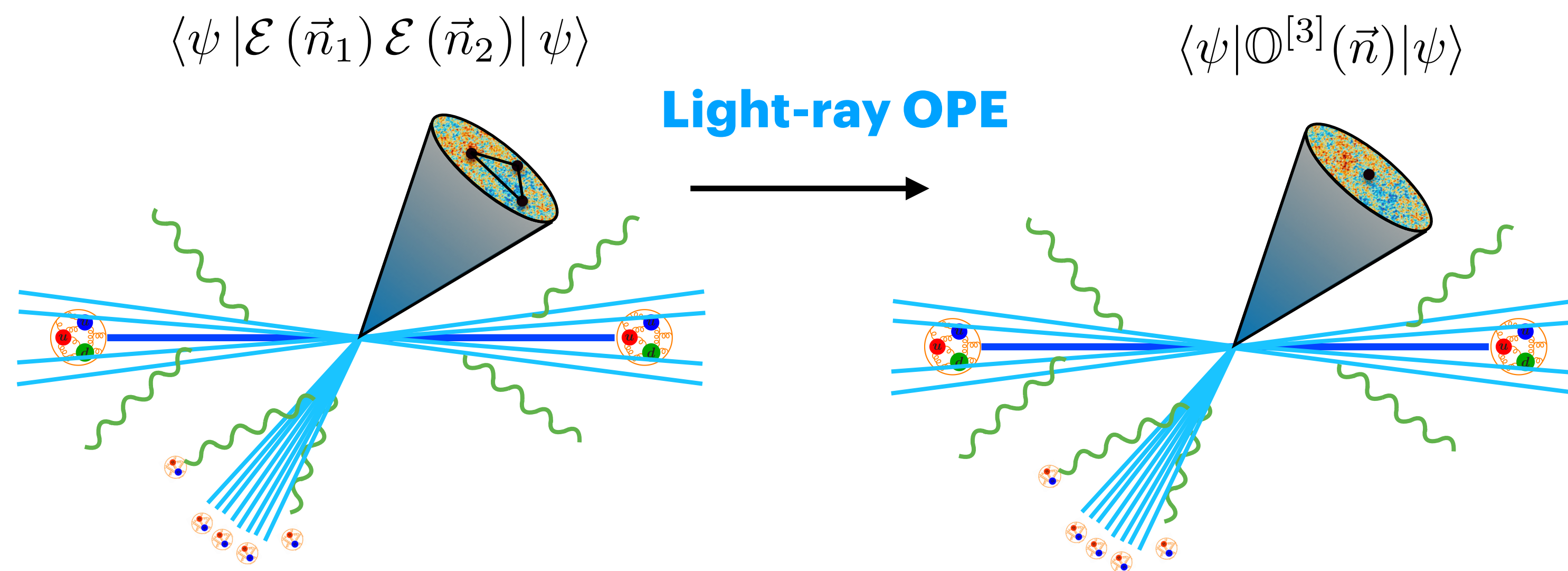
$$\mathcal{E}(\hat{n}) = \int_0^\infty dt \lim_{r \rightarrow \infty} r^2 n^i T_{0i}(t, r\hat{n})$$

$$\mathcal{E}(\hat{n})|X\rangle = \sum_a E_a \delta^{(2)}(\Omega_{\vec{p}_a} - \Omega_{\hat{n}})|X\rangle$$

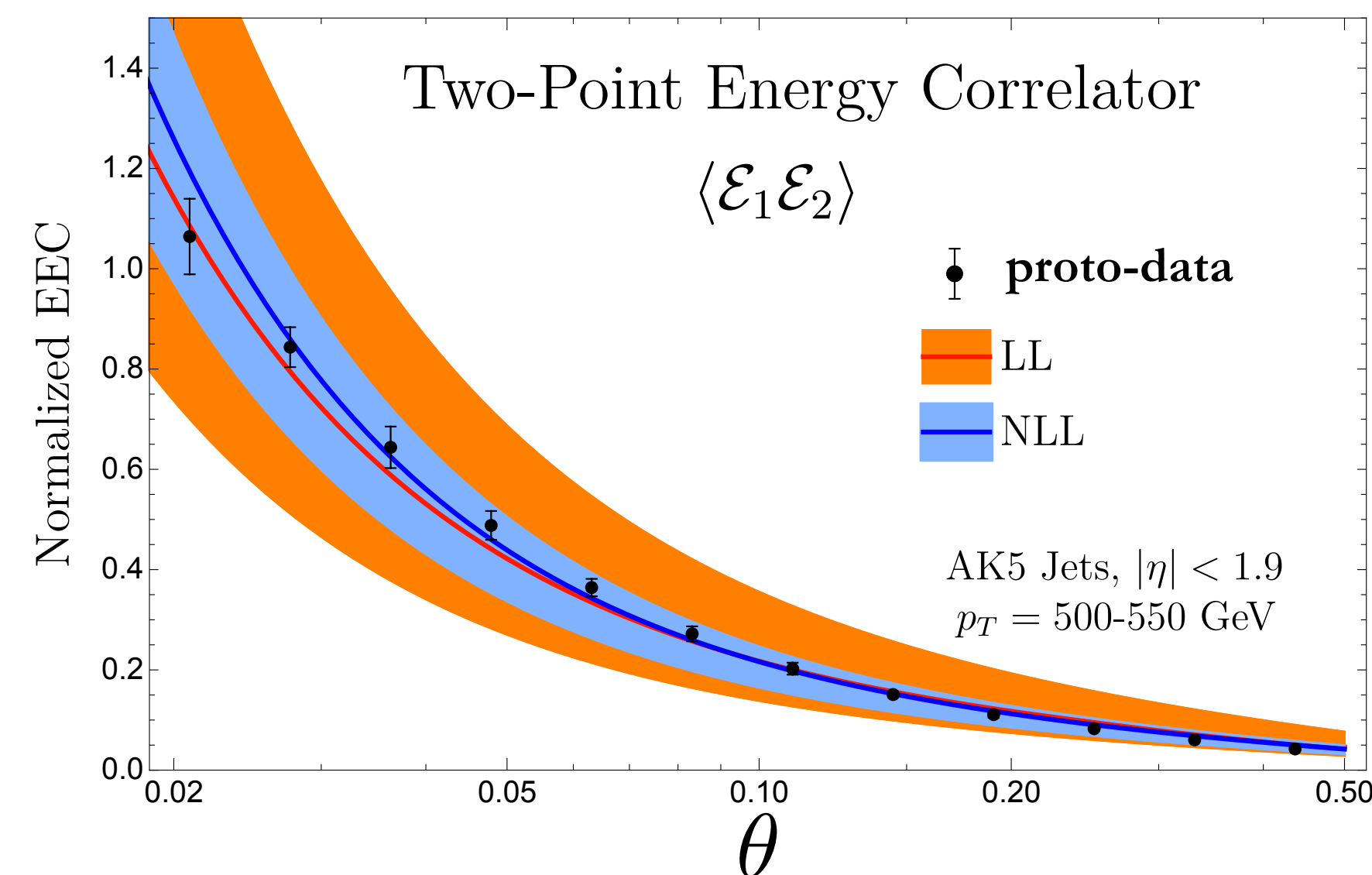
CAN THIS UNIVERSAL SCALING OF THE FIELD THEORY BE OBSERVED IN JETS???

UNIVERSAL SCALING BEHAVIOR IN JETS!

- In QCD, we developed the proper **framework** to observe the **universal scaling behavior within jets!**



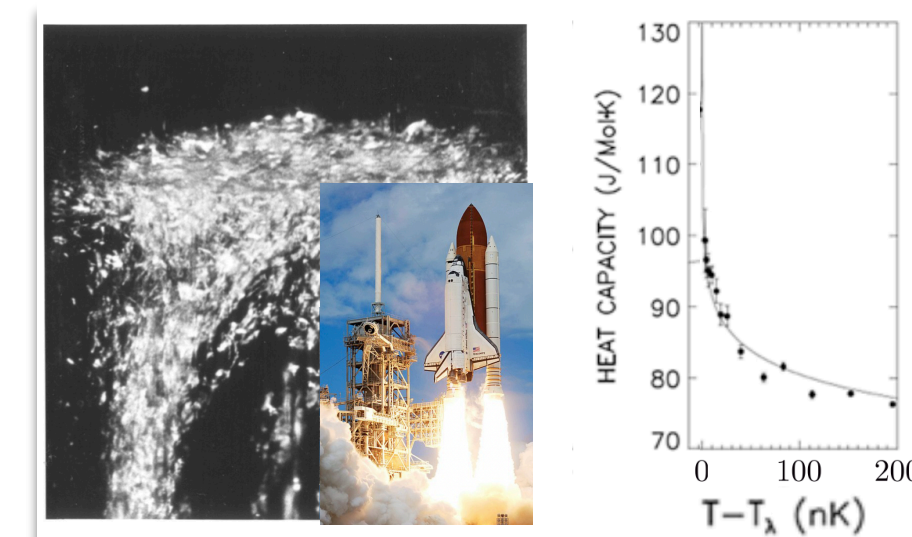
Komiske, Moul, Thaler, Zhu '22
 KL, Meçaj, Moul '22



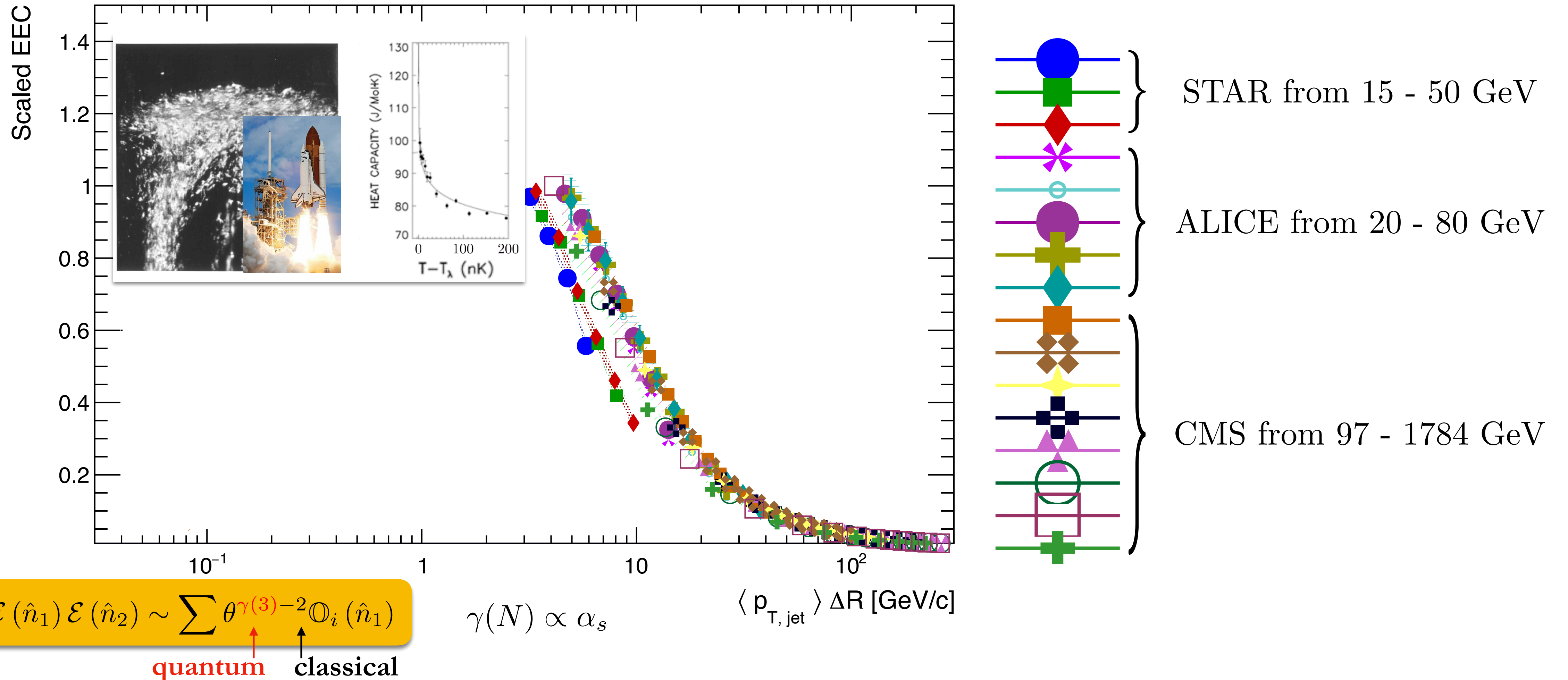
QCD factorization:

$$\frac{d\sigma^{pp \rightarrow \text{jet}(\mathcal{E}\mathcal{E})X}}{dp_T d\eta d\theta} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes \mathcal{G}_c^{\text{EEC}}(\theta)$$

Λ_{QCD} p_T $p_T R$
 $p_T \theta$



SCALING FROM 15 GEV TO 2 TEV IN DATA!



- Universal scaling of QCD operators revealed in data from ALICE, CMS, and STAR, from 15 GeV to 1784 GeV!

THE SPECTRUM OF A JET

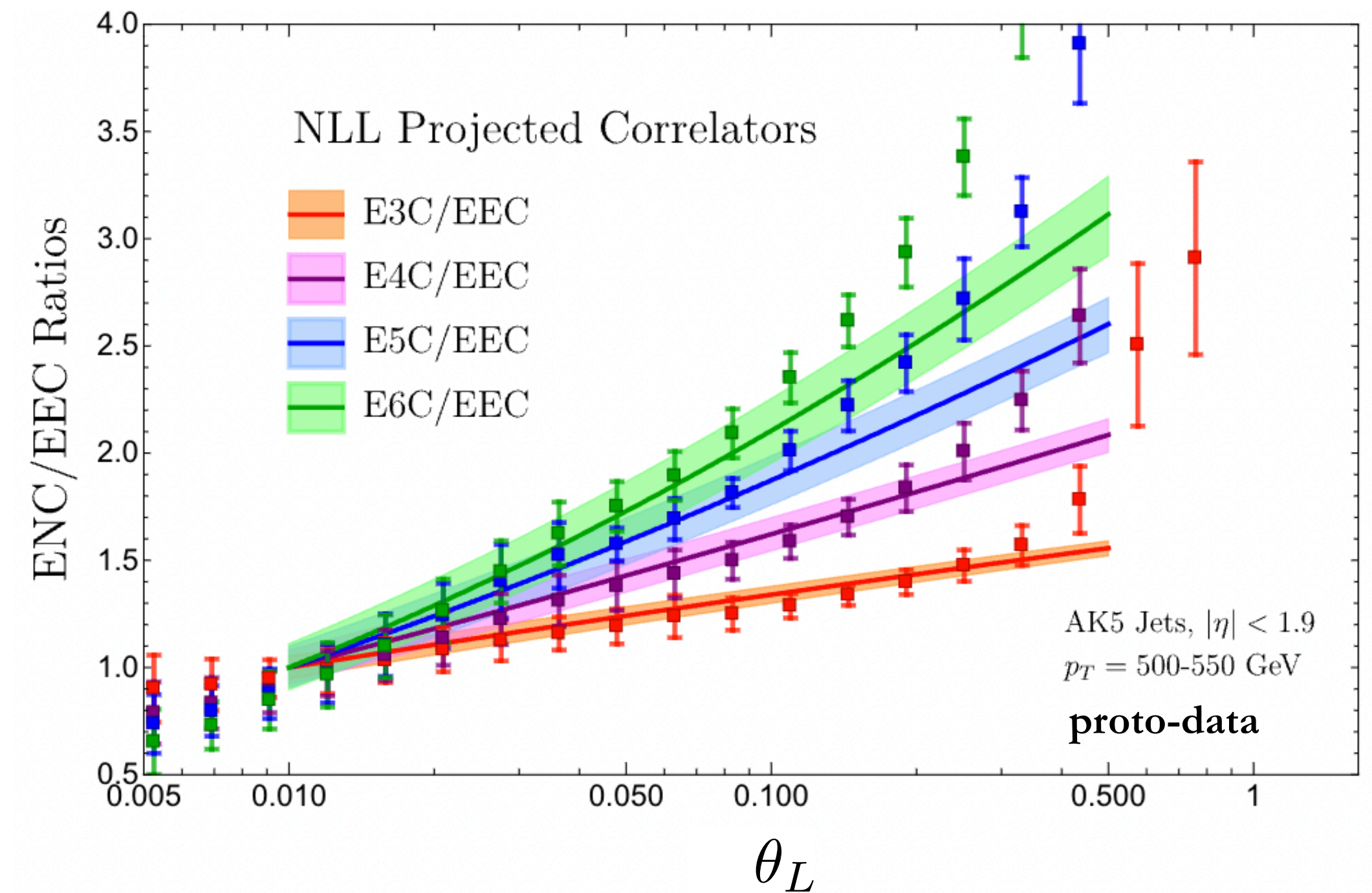
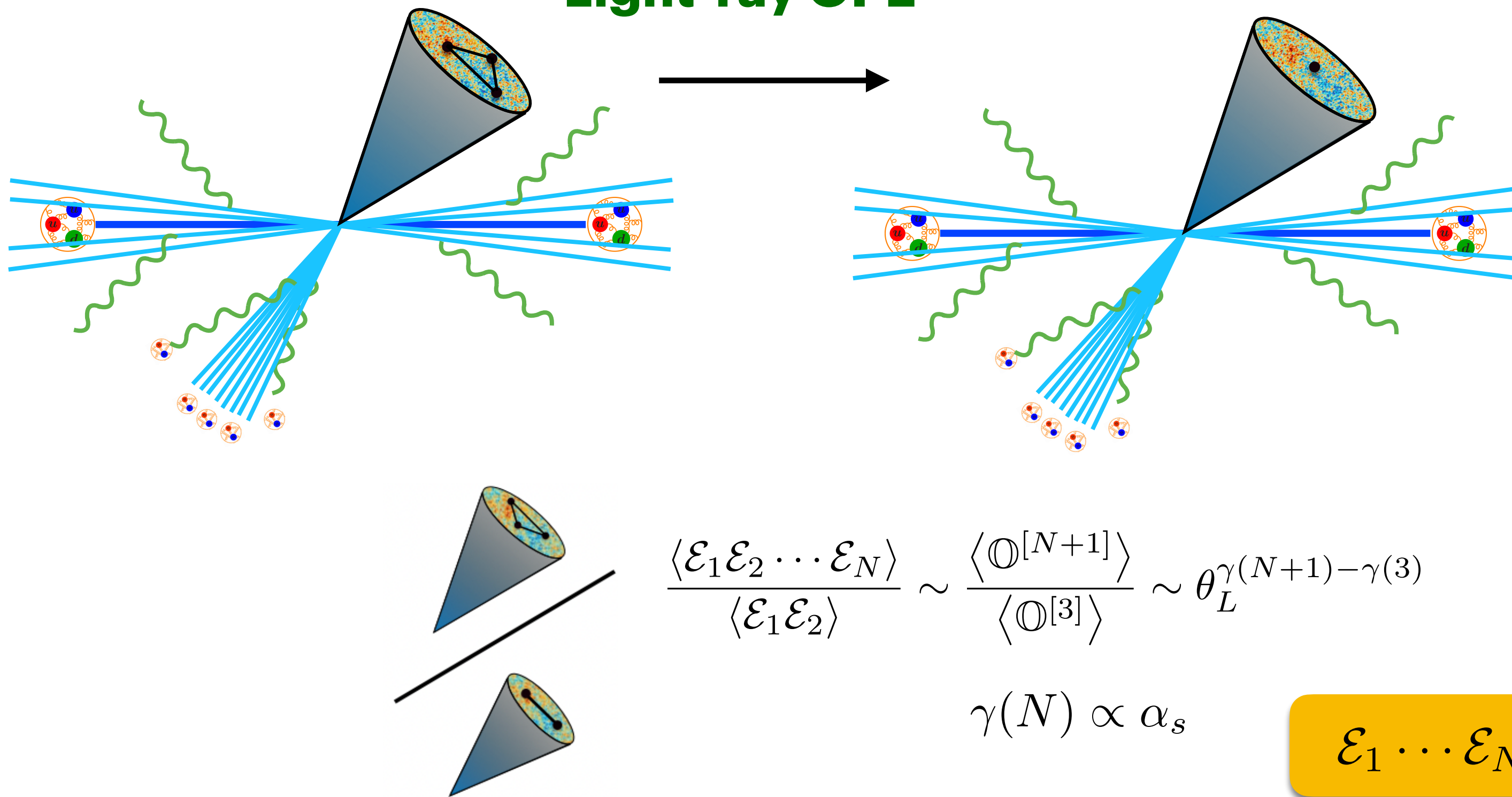
- The **light-ray OPE** can be iteratively applied to **N-point correlators**, predicting their **anomalous scaling behavior with N**

Chen, Moutl, Zhang, Zhu `20
 KL, Meçaj, Moutl `22
 Chen, Gao, Li, Xu, Zhang, Zhu `23

$$\langle \psi | \mathcal{E}(\vec{n}_1) \cdots \mathcal{E}(\vec{n}_N) | \psi \rangle$$

Light-ray OPE

$$\langle \psi | \mathbb{O}^{[J]}(\vec{n}) | \psi \rangle$$

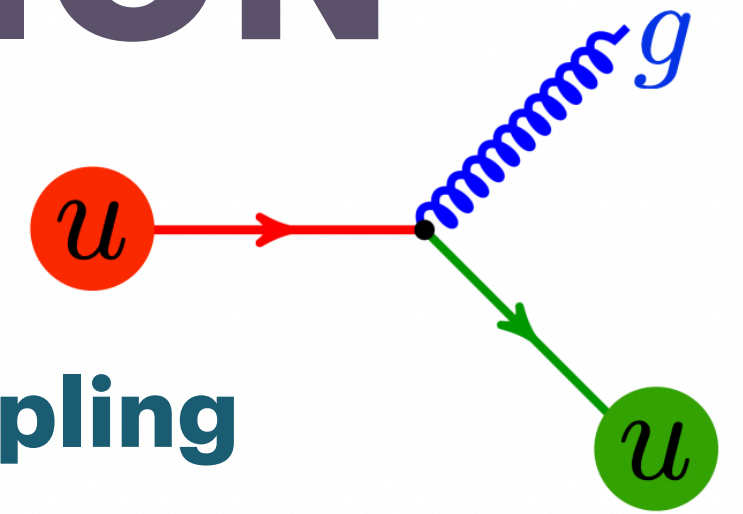


$$\mathcal{E}_1 \cdots \mathcal{E}_N \sim \theta^{\gamma(N+1) - 2} \mathbb{O}_i^{[N+1]}$$

↑ quantum ↑ classical

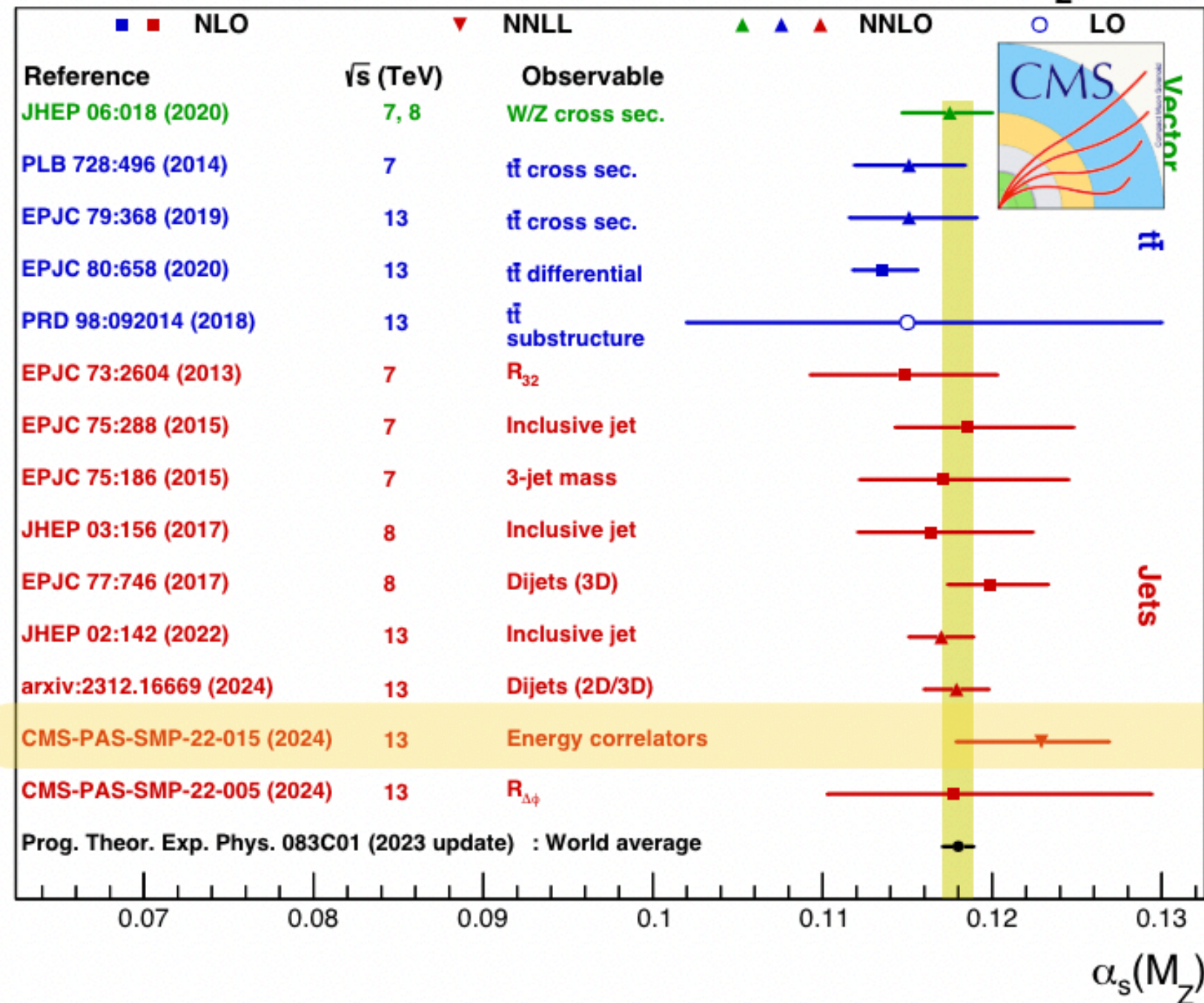
STRONG COUPLING DETERMINATION

- **How strong is the Strong Force?** In comparison, EM coupling: $\alpha_e = 0.0072973525693(11)$



Quarks are never free, and thus it is **very hard to measure their coupling**

Summary of $\alpha_s(M_Z)$



CMS collaboration carried out most precise determination of the strong coupling constant for jet substructure

$$\alpha_s(m_Z) = 0.1229^{+0.0040}_{-0.0050}$$

CMS Collaboration '23

⇒ 4% uncertainty

Energy Correlators in Jet

This yielded the worlds most precise α_S measurement from jet substructure: $\alpha_S = 0.1229^{+0.0040}_{-0.0050}$.

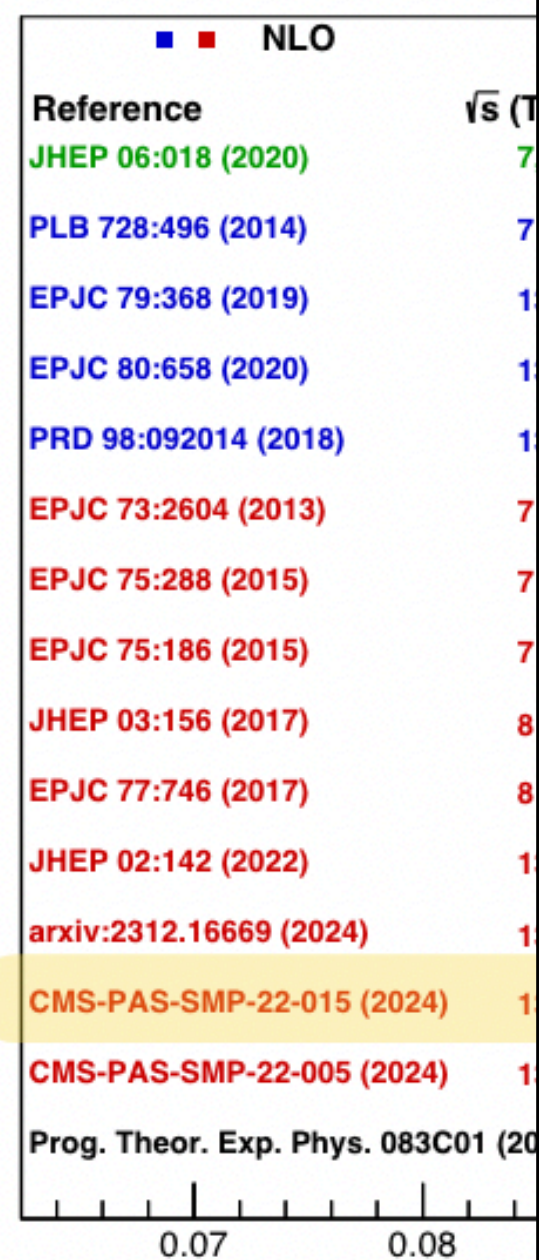
ROAD TO IMPROVED PRECISION

Road to precision

1. Measurements on Tracks

2. Power corrections

3. Improved perturbative accuracy



tion

3

This yielded the worlds most precise α_S measurement from jet substructure: $\alpha_S = 0.1229^{+0.0040}_{-0.0050}$.

ROAD TO IMPROVED PRECISION

Road to precision

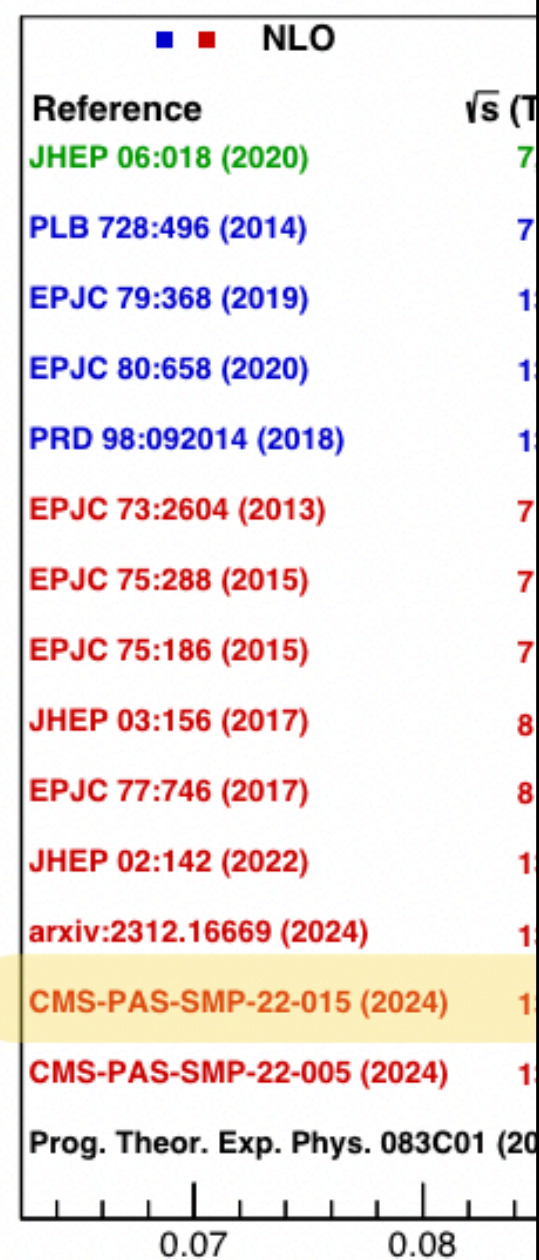
1. Measurements on Tracks

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← Backup slides

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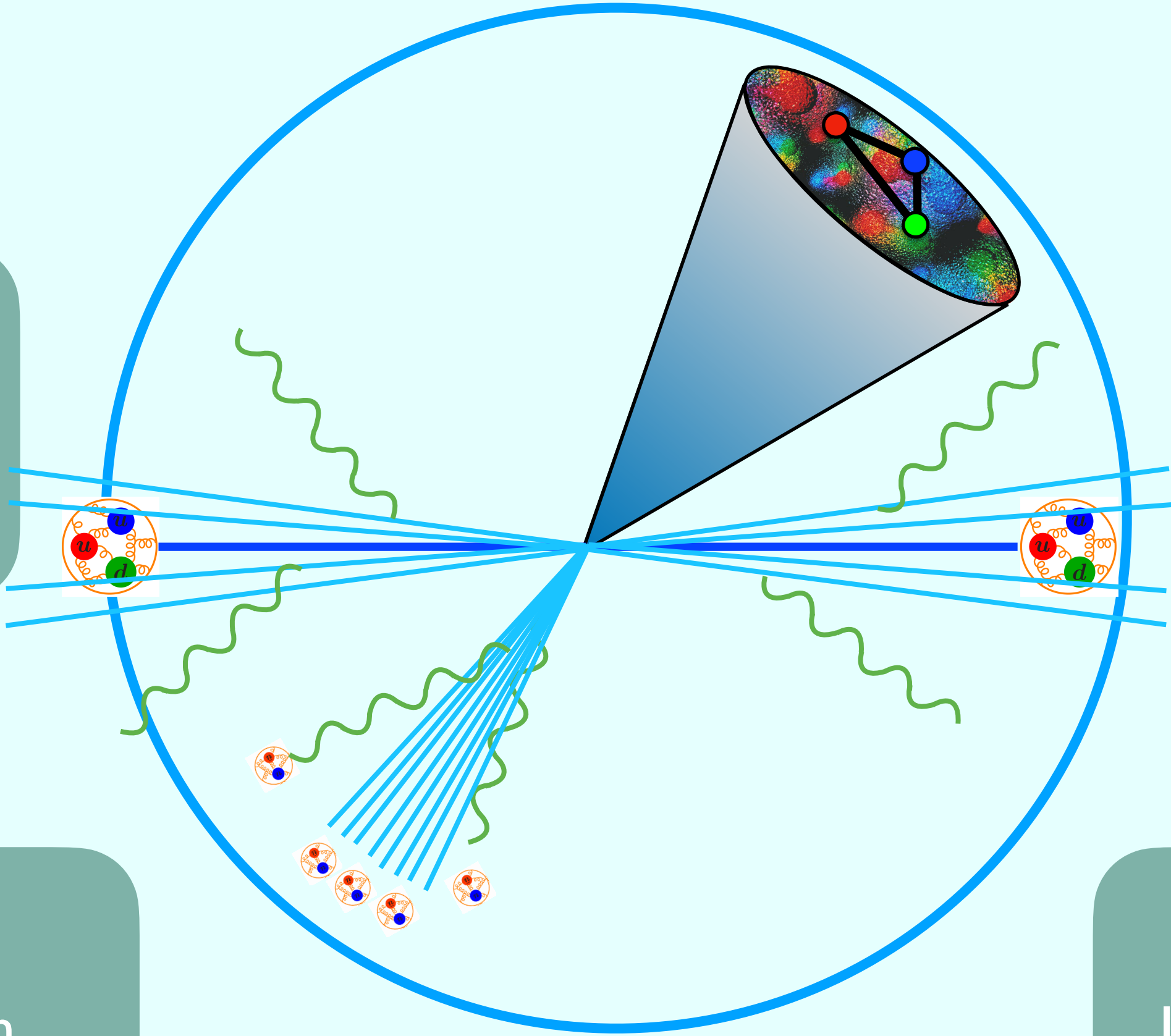
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III. Heavy Flavor Physics

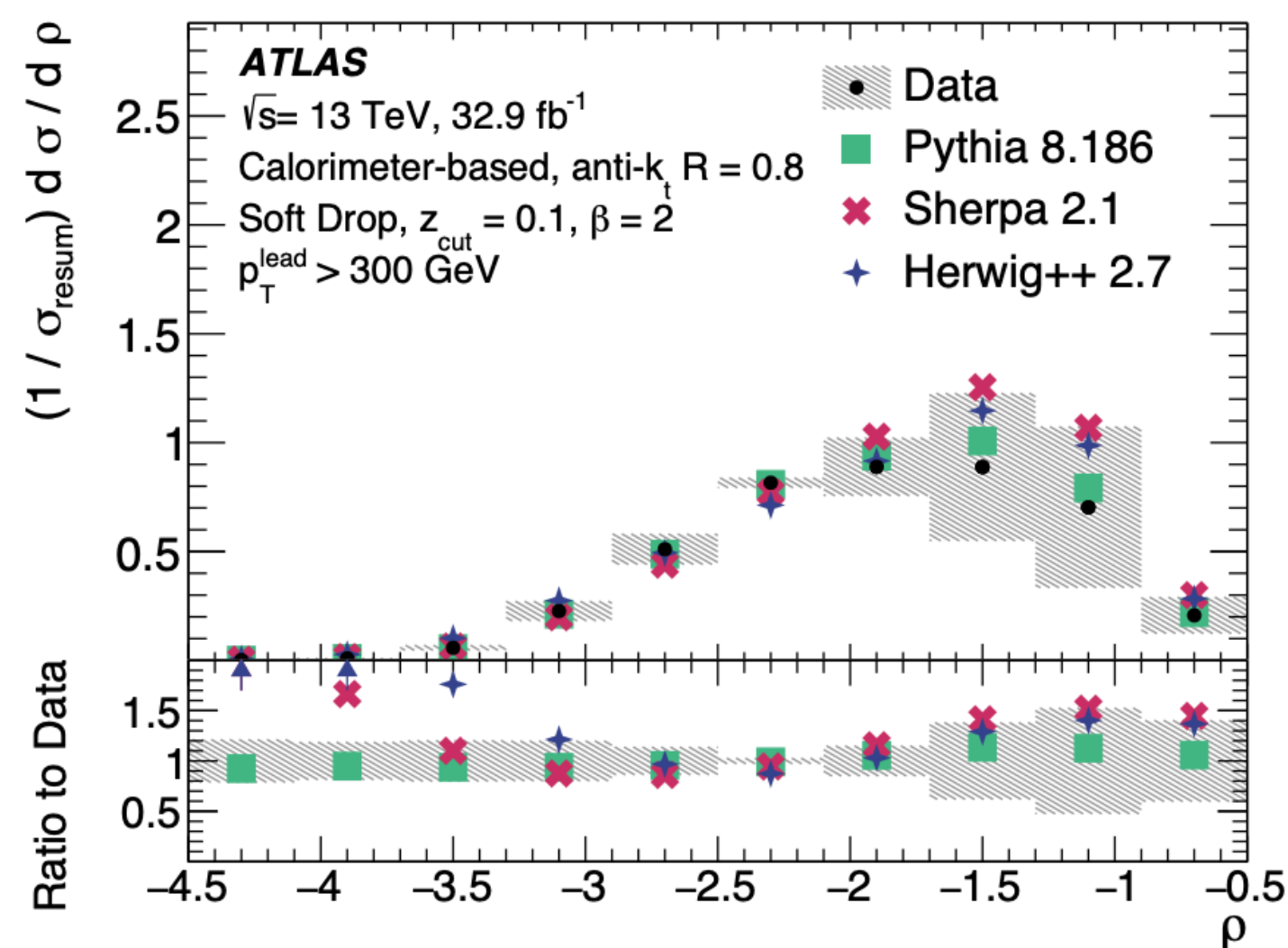
V. Hadronization

IV. Medium Dynamics

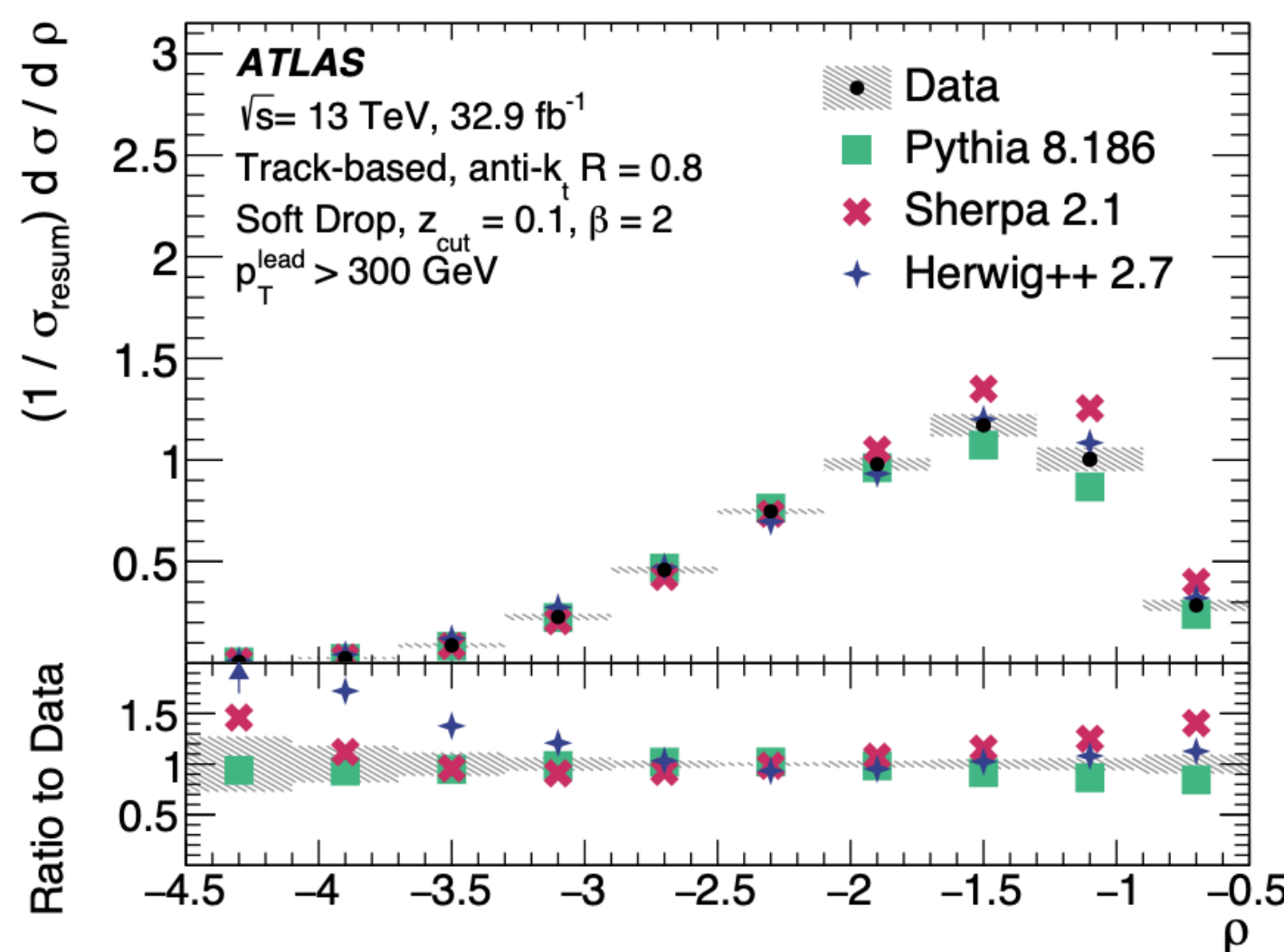


MEASURING TRACKS

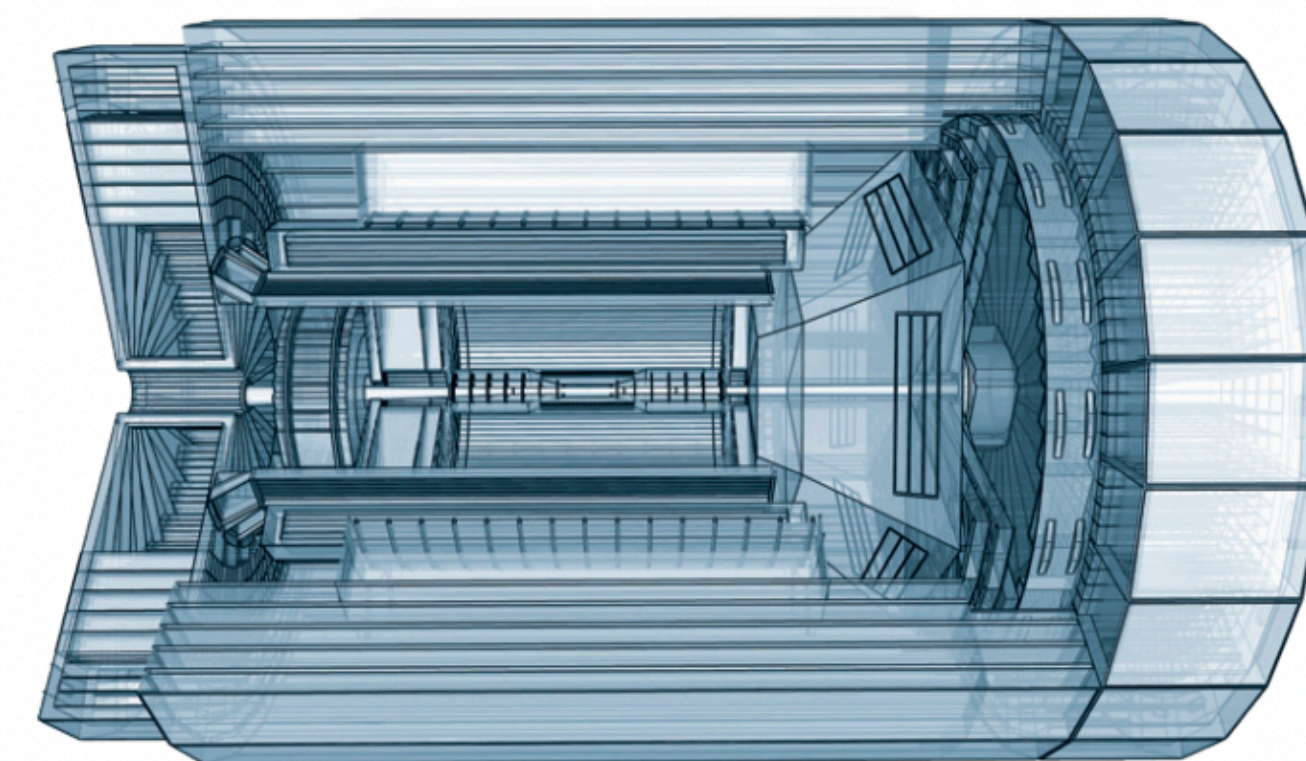
- Measuring tracks provides much more precise experimental results



All particles



Tracks



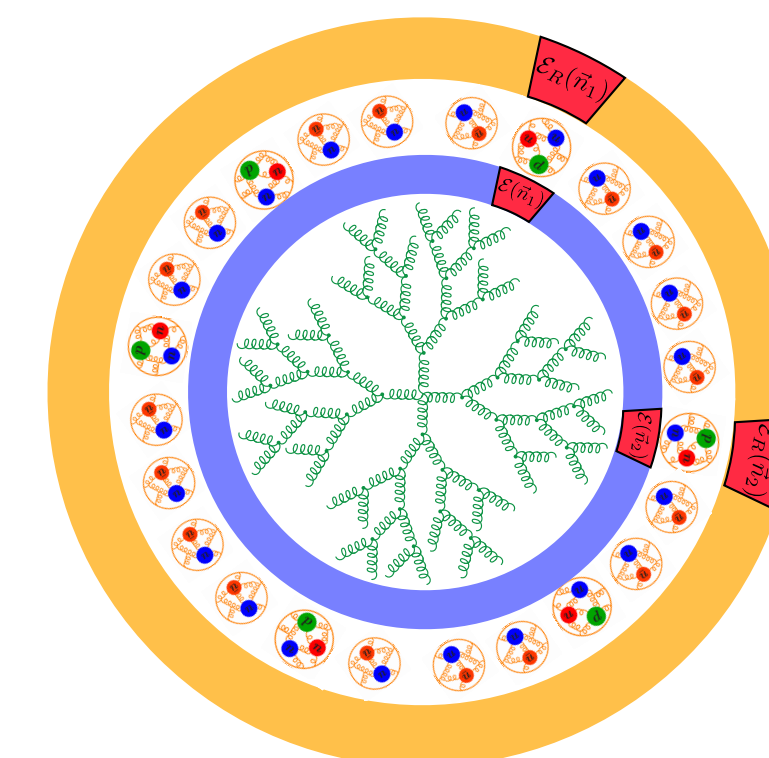
EIC detector

EIC will have state-of-the-art tracking systems!

- Depend on quantum numbers of final state hadrons other than energy

⇒ not computable purely from perturbation theory

We need QCD factorization

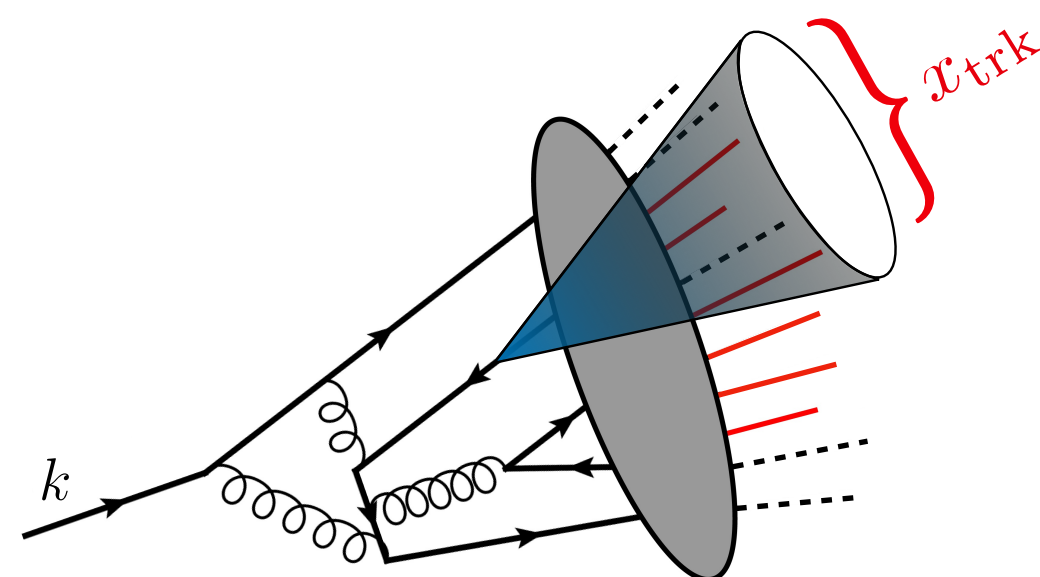


TRACK INSIDE JETS

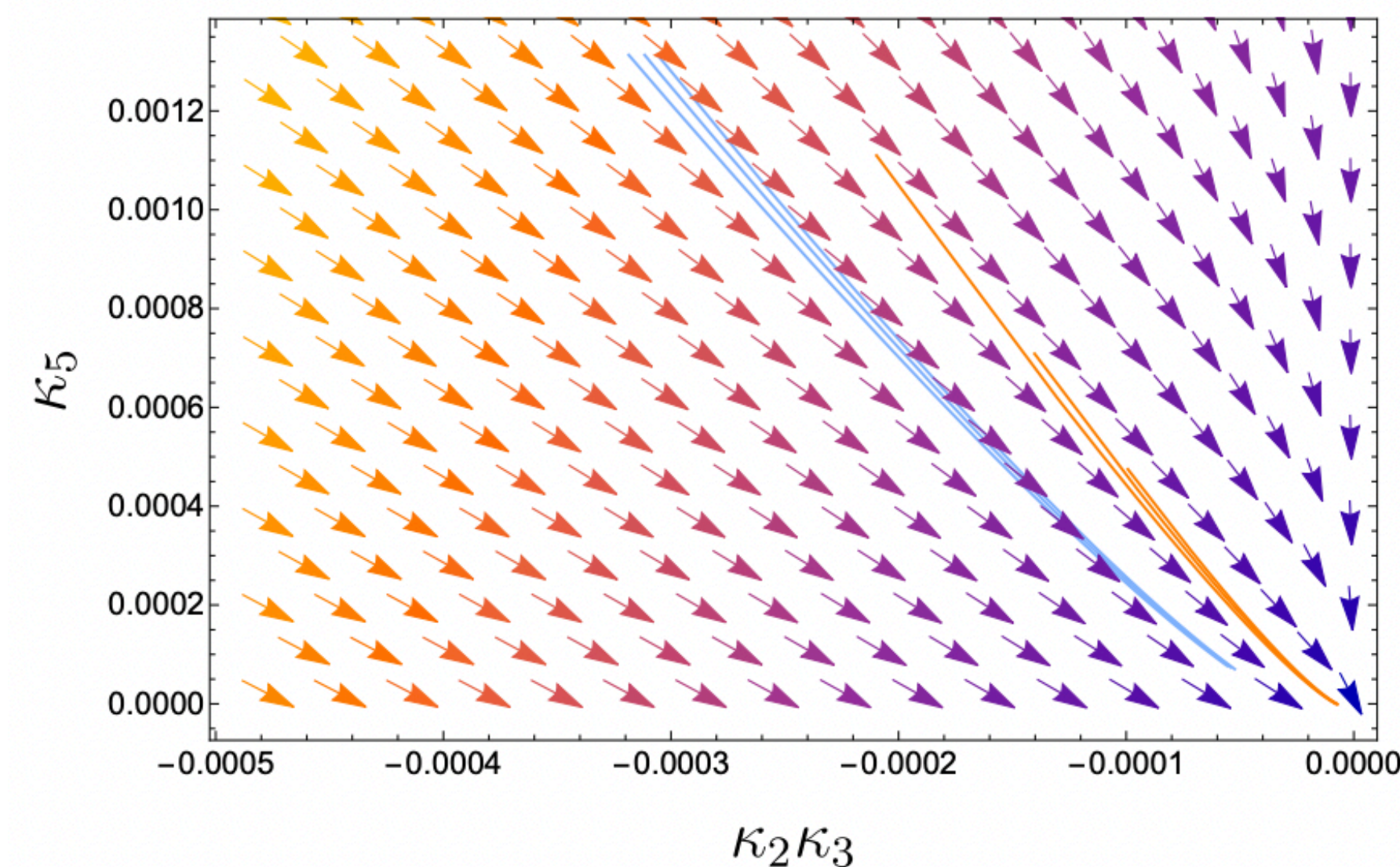
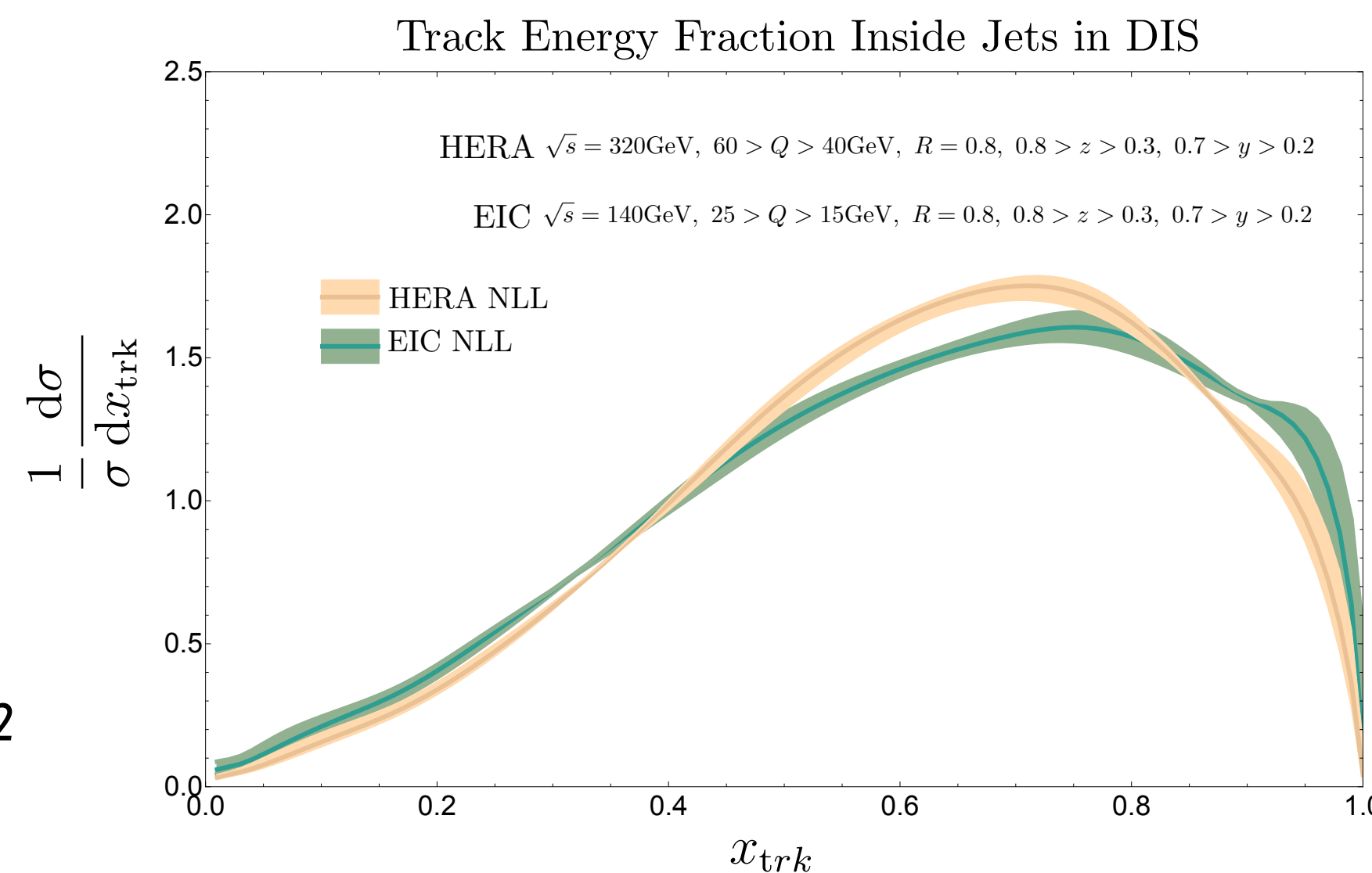
QCD factorization:

Requires separation of parts that are **perturbative** from **universal non-perturbative functions**

- **Non-perturbative Track functions** describe the total energy fraction of charged hadrons from a fragmenting quark or a gluon state

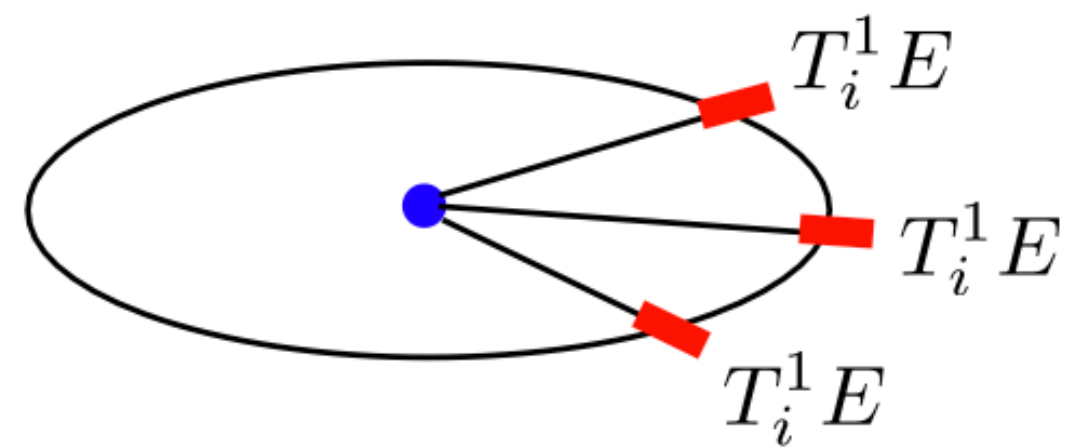


Chang, Procura, Thaler, Waalewijn '13
 Jaarsma, Li, Mout, Waalewijn, Zhu et al '21, 22
 KL, Mout, Ringer, Waalewijn '23
 KL, Mout '23



ENERGY CORRELATORS ON TRACK

- Track function formalism provides the essential **matching** between **partonic and hadronic detectors**



$$\langle \mathcal{E}_R(\vec{n}_1) \mathcal{E}_R(\vec{n}_2) \cdots \mathcal{E}_R(\vec{n}_k) \rangle = \sum_{i_1, i_2, \dots, i_k} T_{i_1}(1) \cdots T_{i_k}(1) \langle \mathcal{E}_{i_1}(\vec{n}_1) \mathcal{E}_{i_2}(\vec{n}_2) \cdots \mathcal{E}_{i_k}(\vec{n}_k) \rangle + \text{contact terms}$$

- Only depends on the **“moments”** of track functions \implies Only involves **NP numbers**, not functions

Predictions for tracks in Energy Correlators

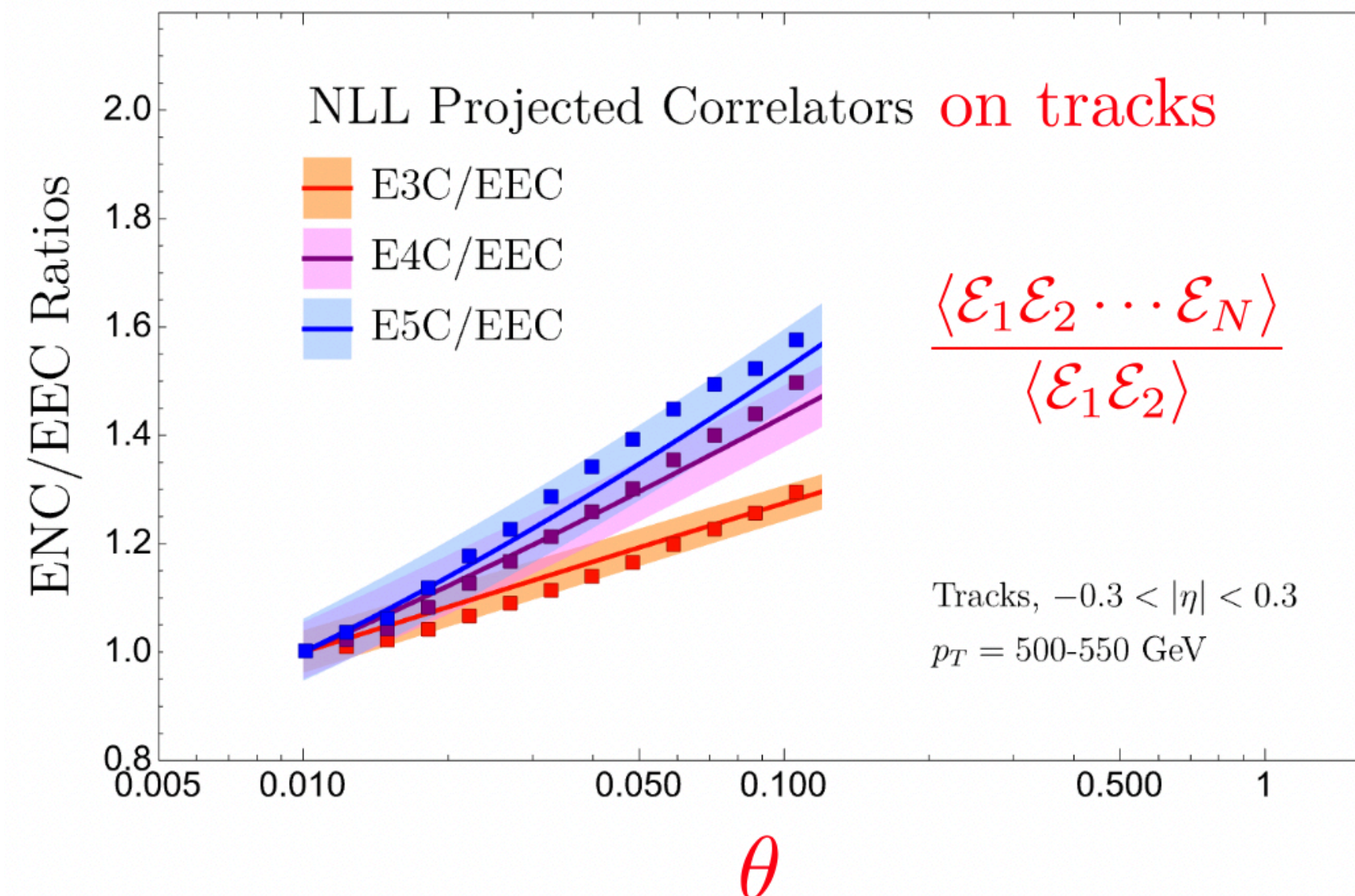
Chang, Procura, Thaler, Waalewijn `13

Jaarsma, Li, Mout, Waalewijn, Zhu et al `21, 22, 23

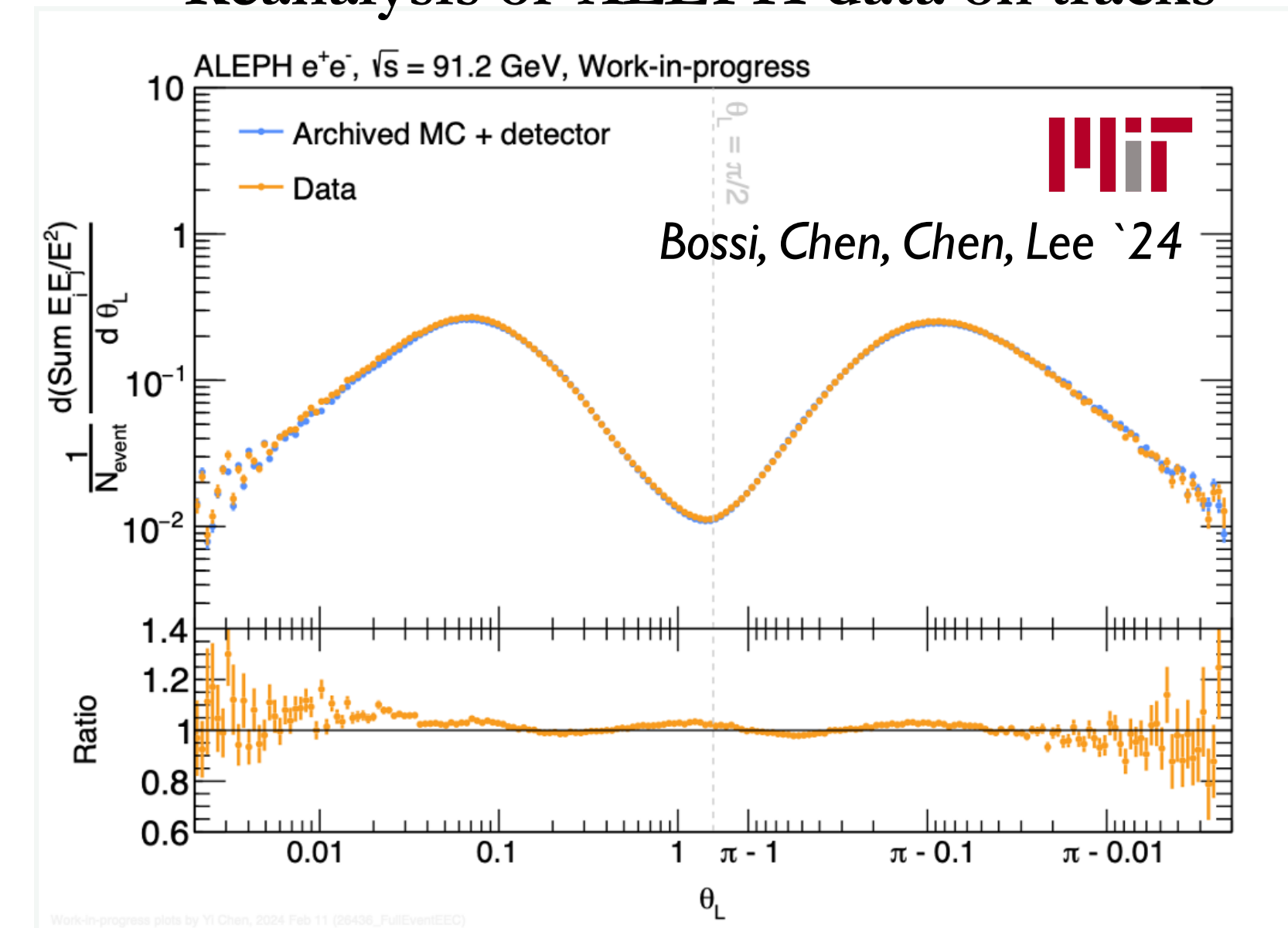
KL, Mout, Ringer, Waalewijn `23

KL, Mout `23

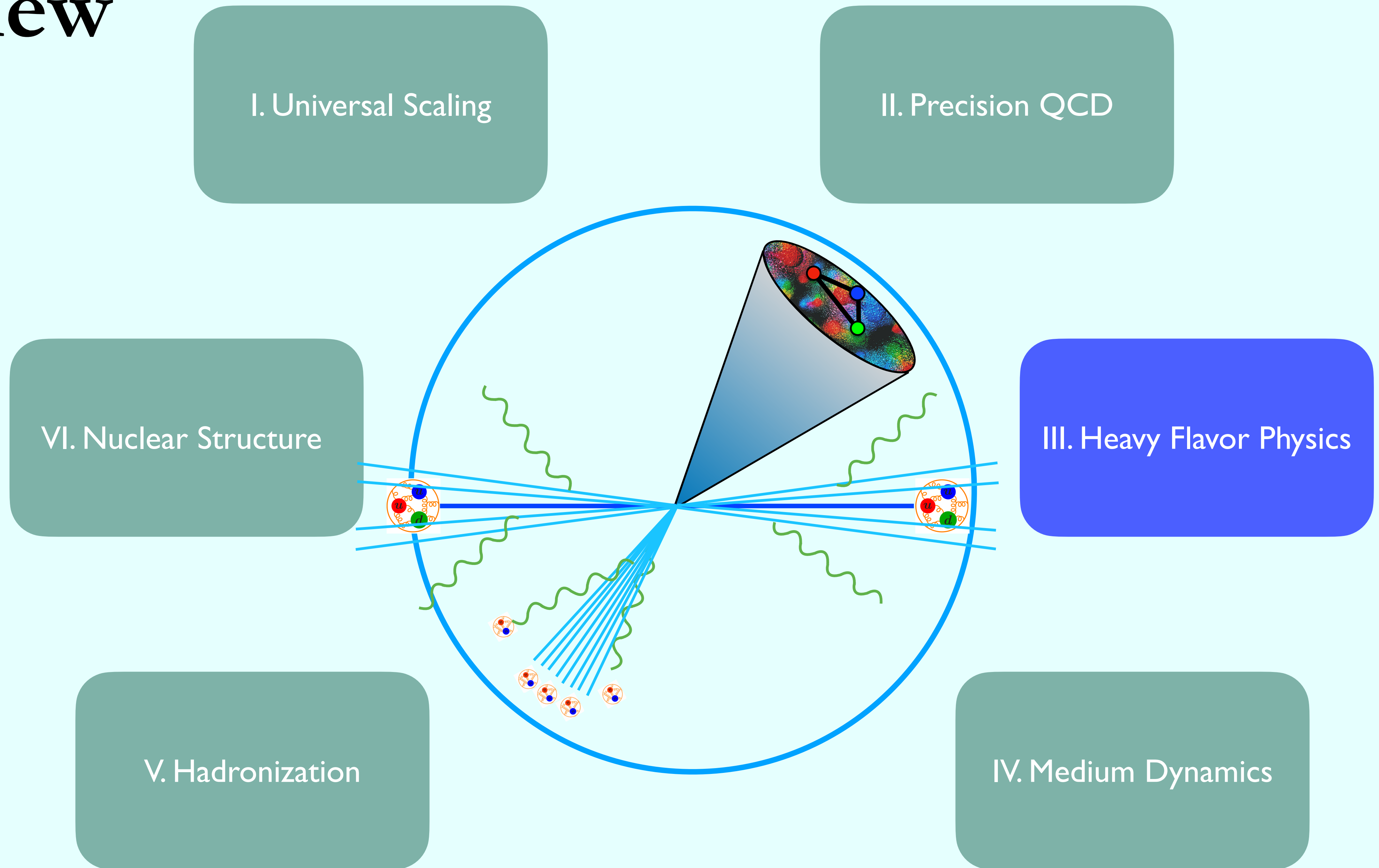
KL, Li, Mout, Waalewijn `In Progress



Reanalysis of ALEPH data on tracks

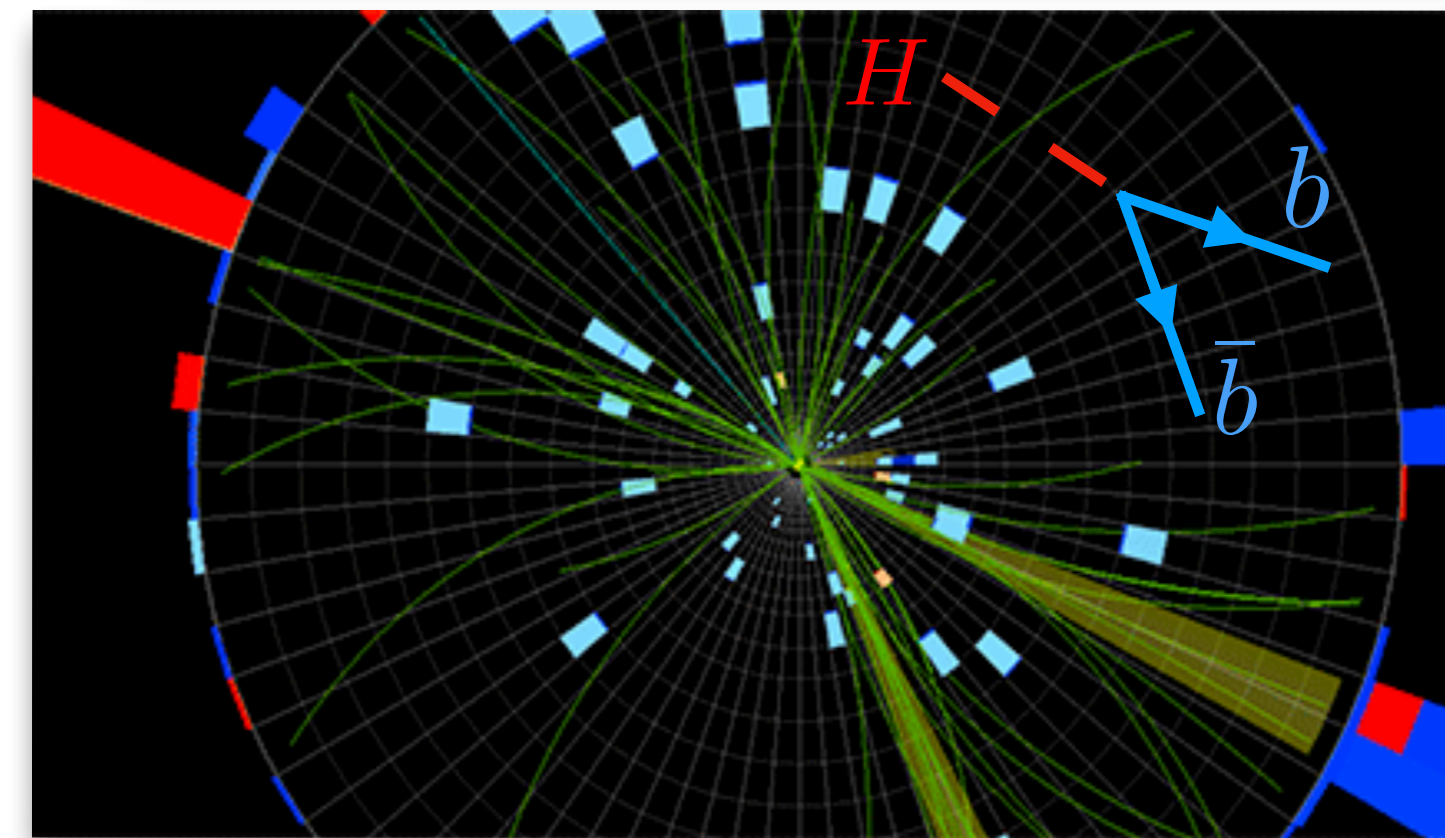
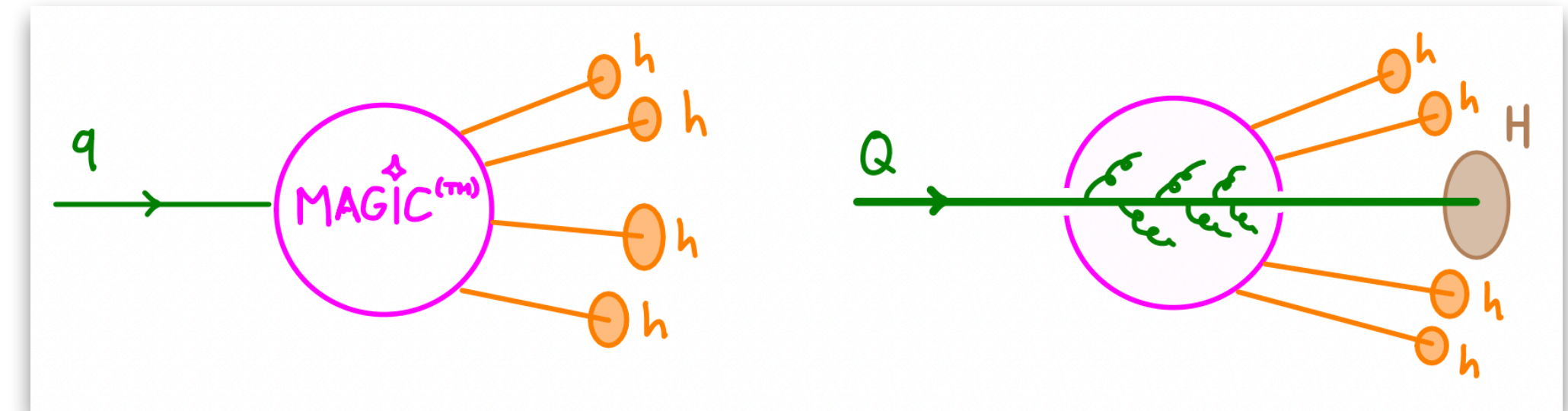
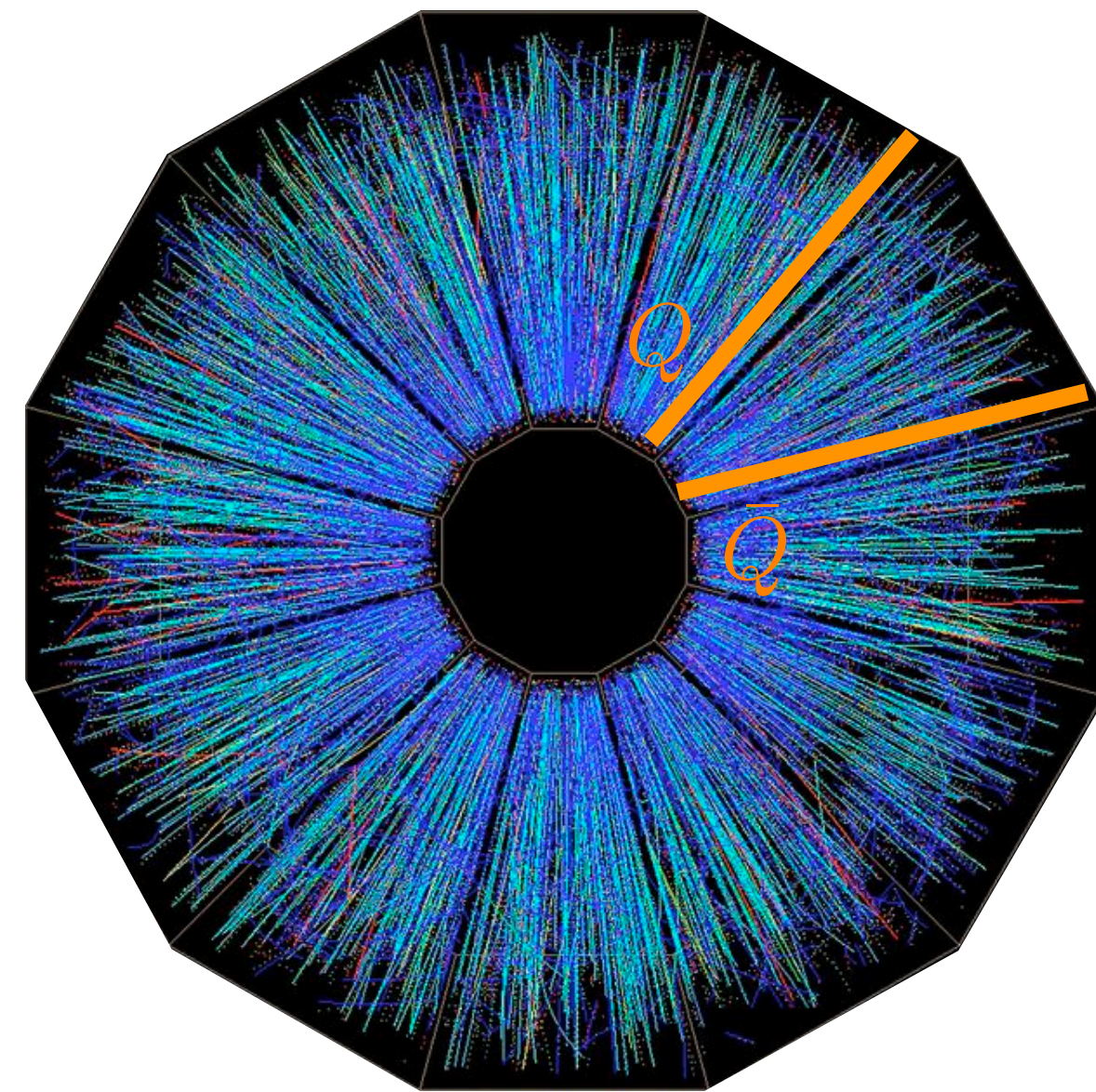


Overview



UNRAVELING HEAVY FLAVOR DYNAMICS

- **Heavy quark dynamics** are important for understanding **medium, hadronization, Higgs, BSM searches, flavor tagging, gluon structure, etc.**



Run 3 and sPHENIX will give us a lot more access to **heavy quarks with precise data!**

- Heavy quark introduces **new mass scale** m_Q
- **Jet substructure** allows us to precisely probe the dynamics from this new **heavy quark scale**

QUARK GLUON SCALING AND HADRONIZATION

- Energy correlators allow the **hadronization process to be directly imaged inside high energy jets: transition from interacting quarks and gluons to free hadrons is clearly visible!**

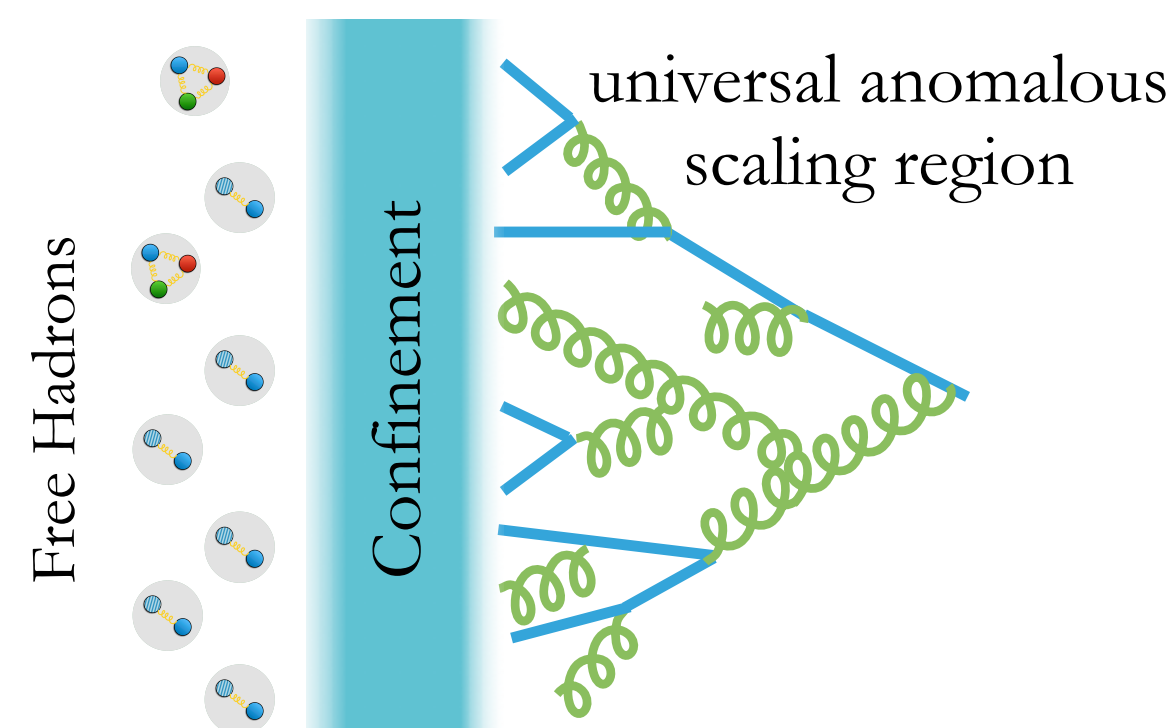
Free hadrons

$$\frac{d\sigma}{d\theta^2} = \text{const}$$

$$\frac{d\sigma}{d\theta} = \text{const} \times 2\theta$$

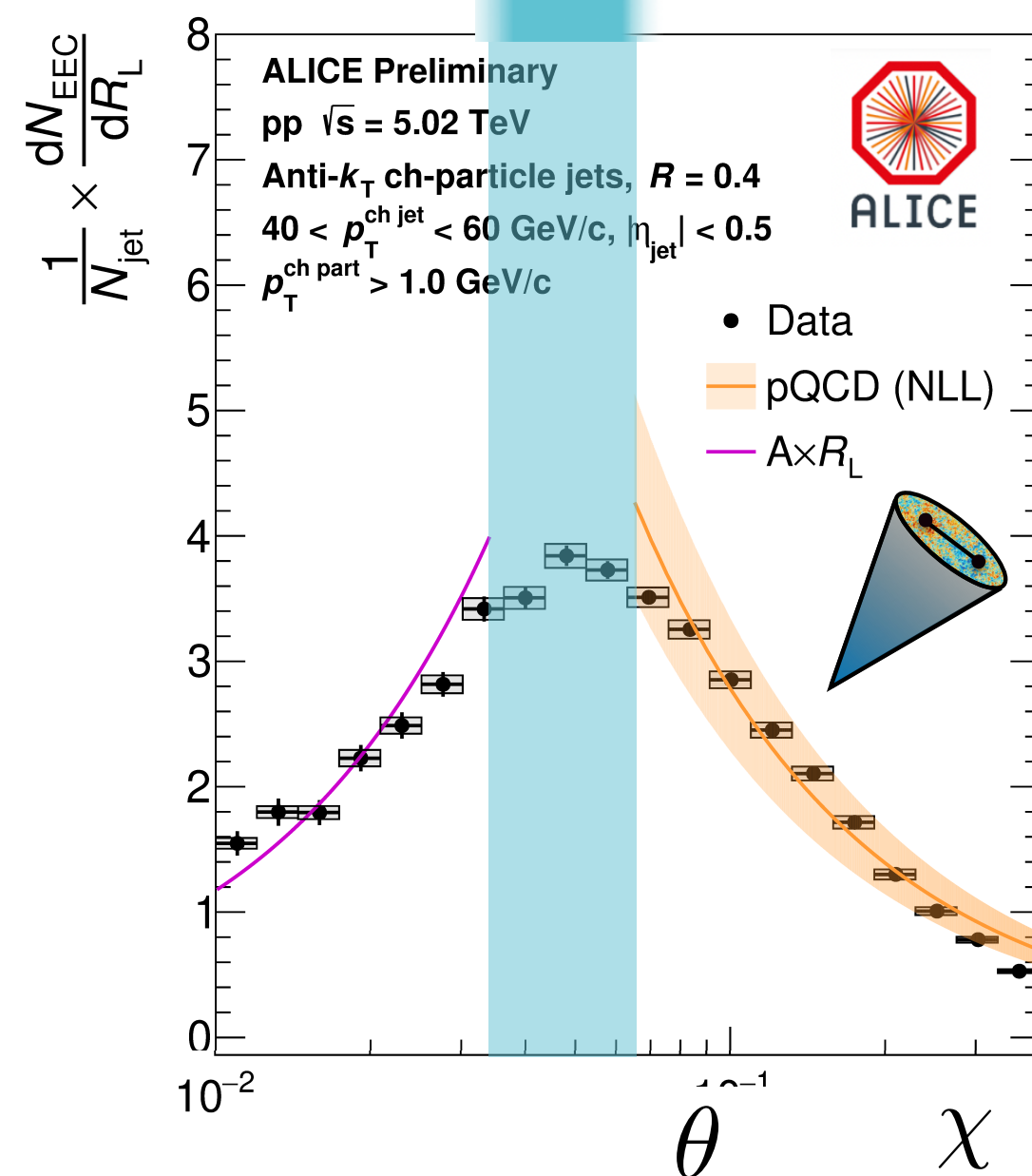
EEC gives angular scale

$$\mu \sim p_T \theta_{ij}$$



Interacting quarks and gluons

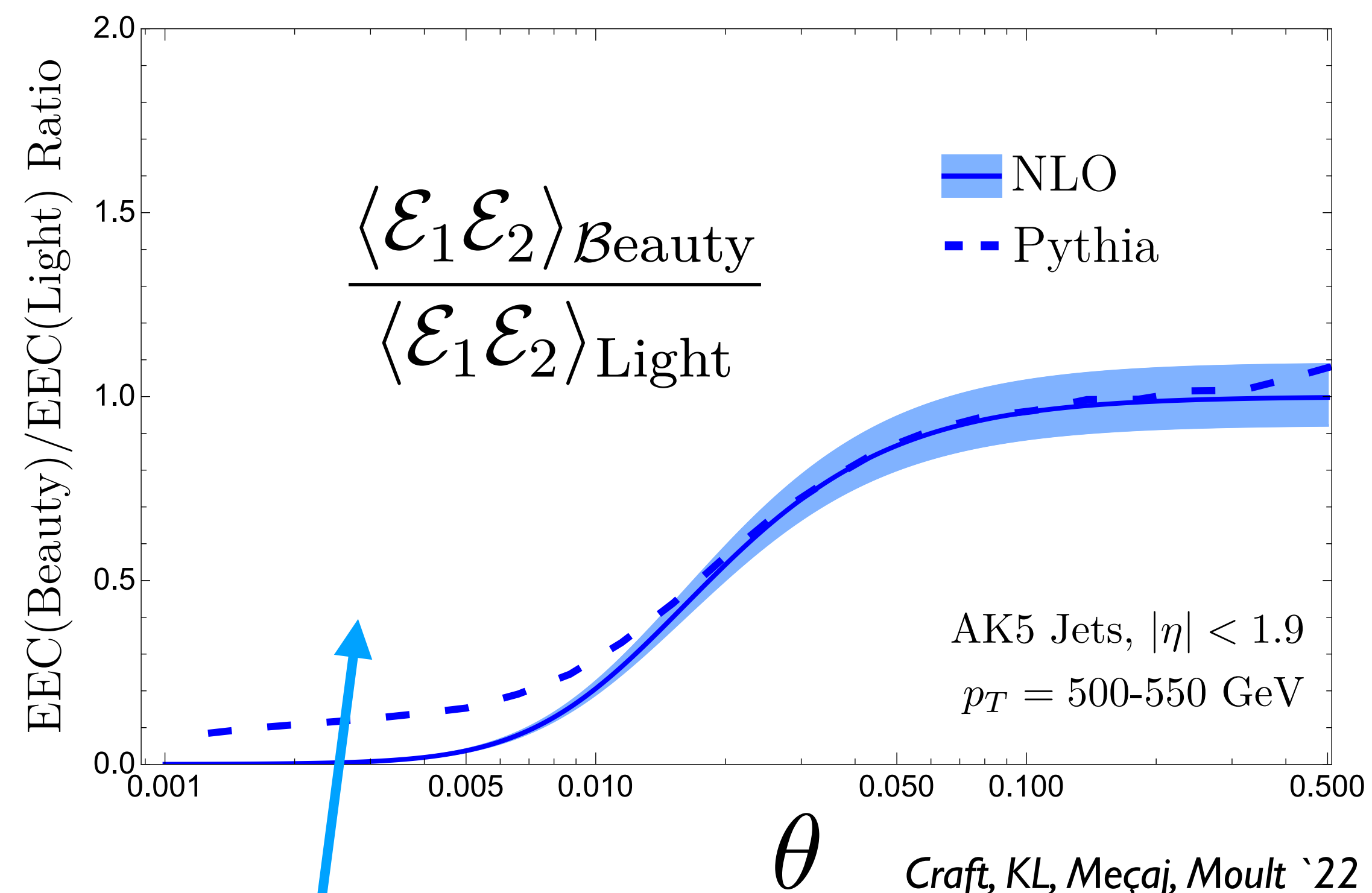
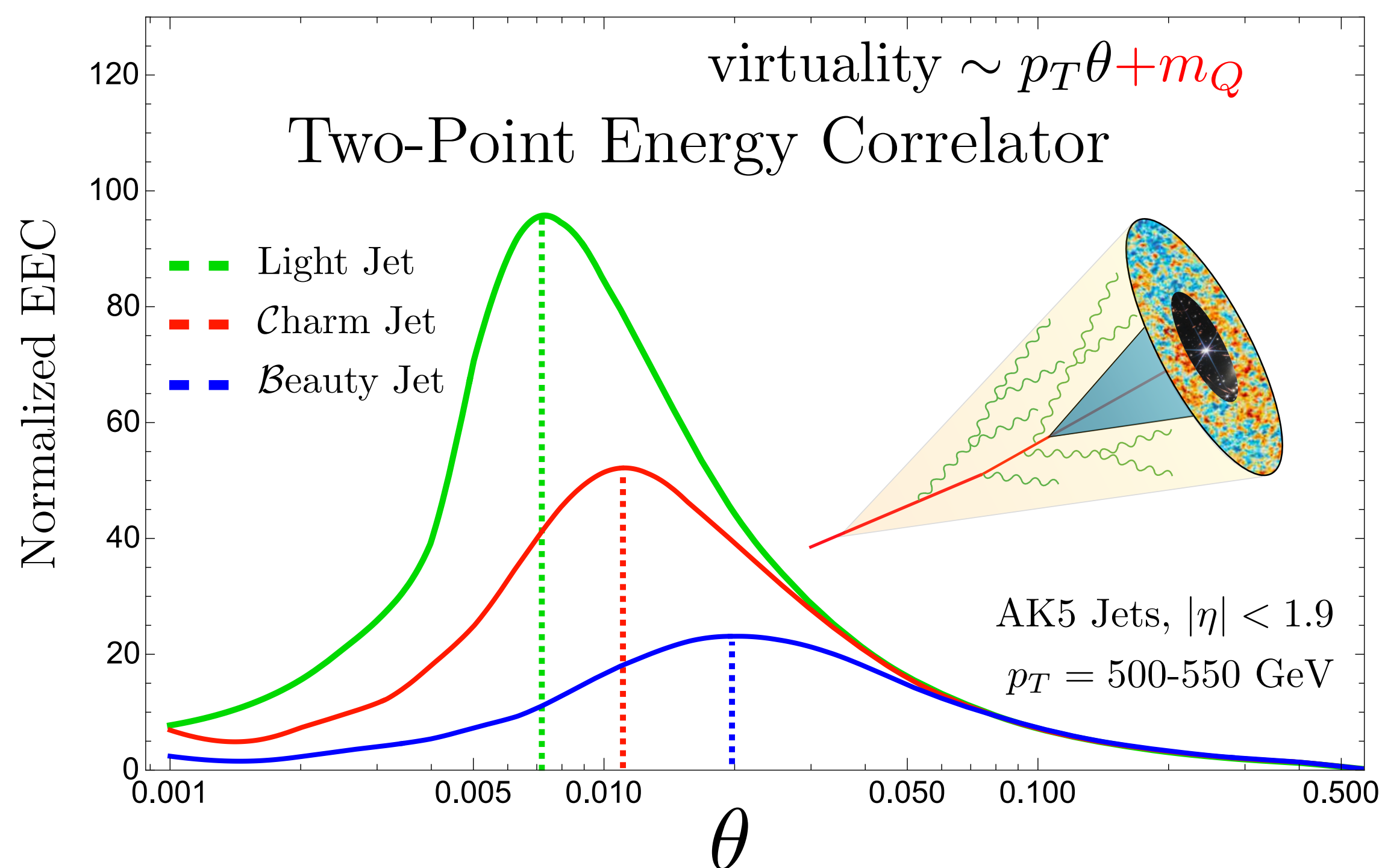
$$\mathcal{E}(\hat{n}_1) \mathcal{E}(\hat{n}_2) \sim \sum \theta^{\gamma(3)-2} \mathcal{O}_i(\hat{n}_1)$$



KL, Meçaj, Moul `22
 Komiske, Moul, Thaler, Zhu `22
 ALICE `23

IDENTIFYING THE INTRINSIC HEAVY QUARK SCALE

- **Two-point correlators** capture the effects of intrinsic mass, displaying earlier formation of heavy bound states due to their mass



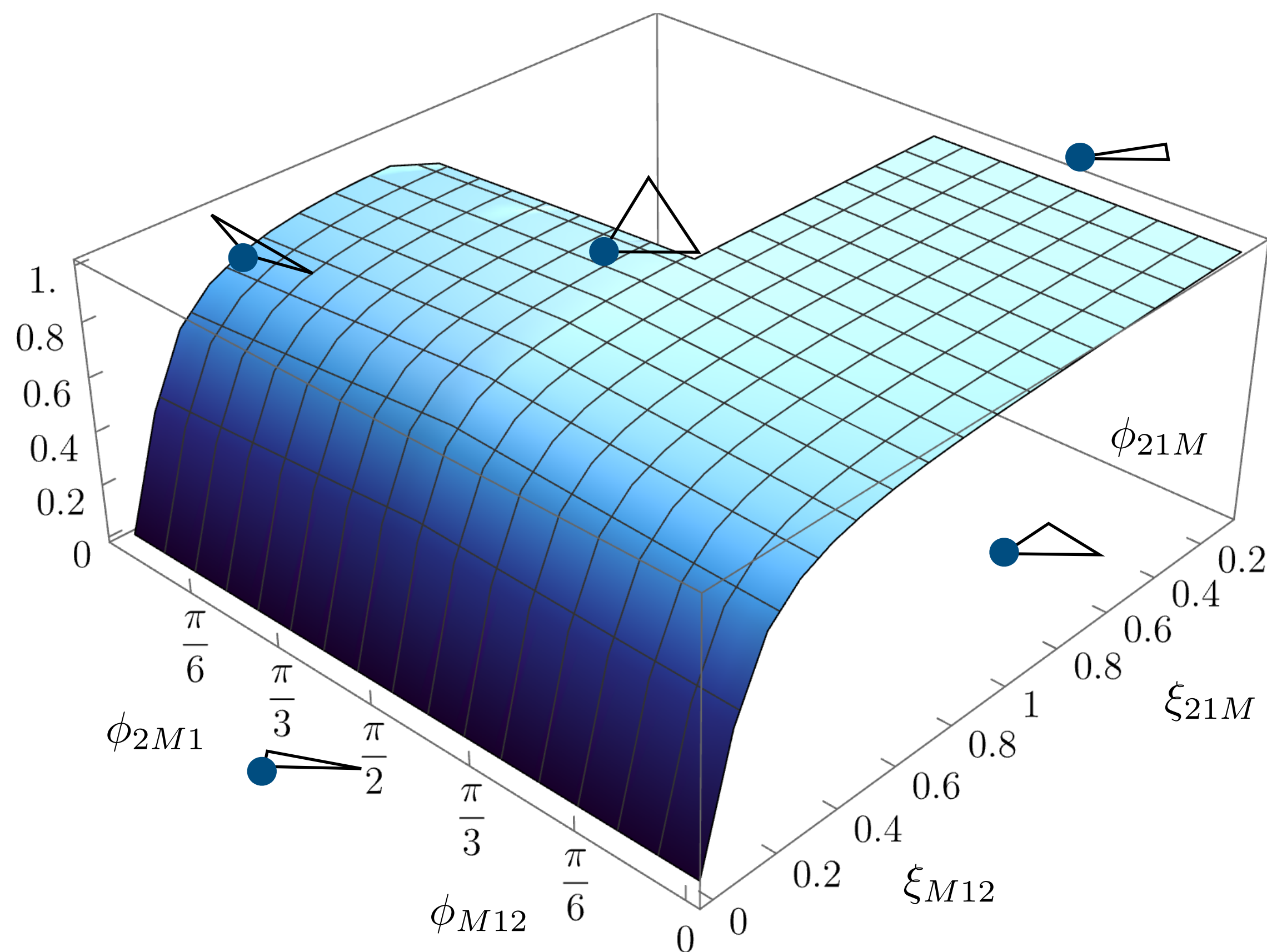
- Ratio of the **two-point correlators** clearly shows the **dead-cone region** around $\theta \lesssim \frac{m_Q}{E}$

PROBING THE DYNAMICS OF THE DEAD-CONE

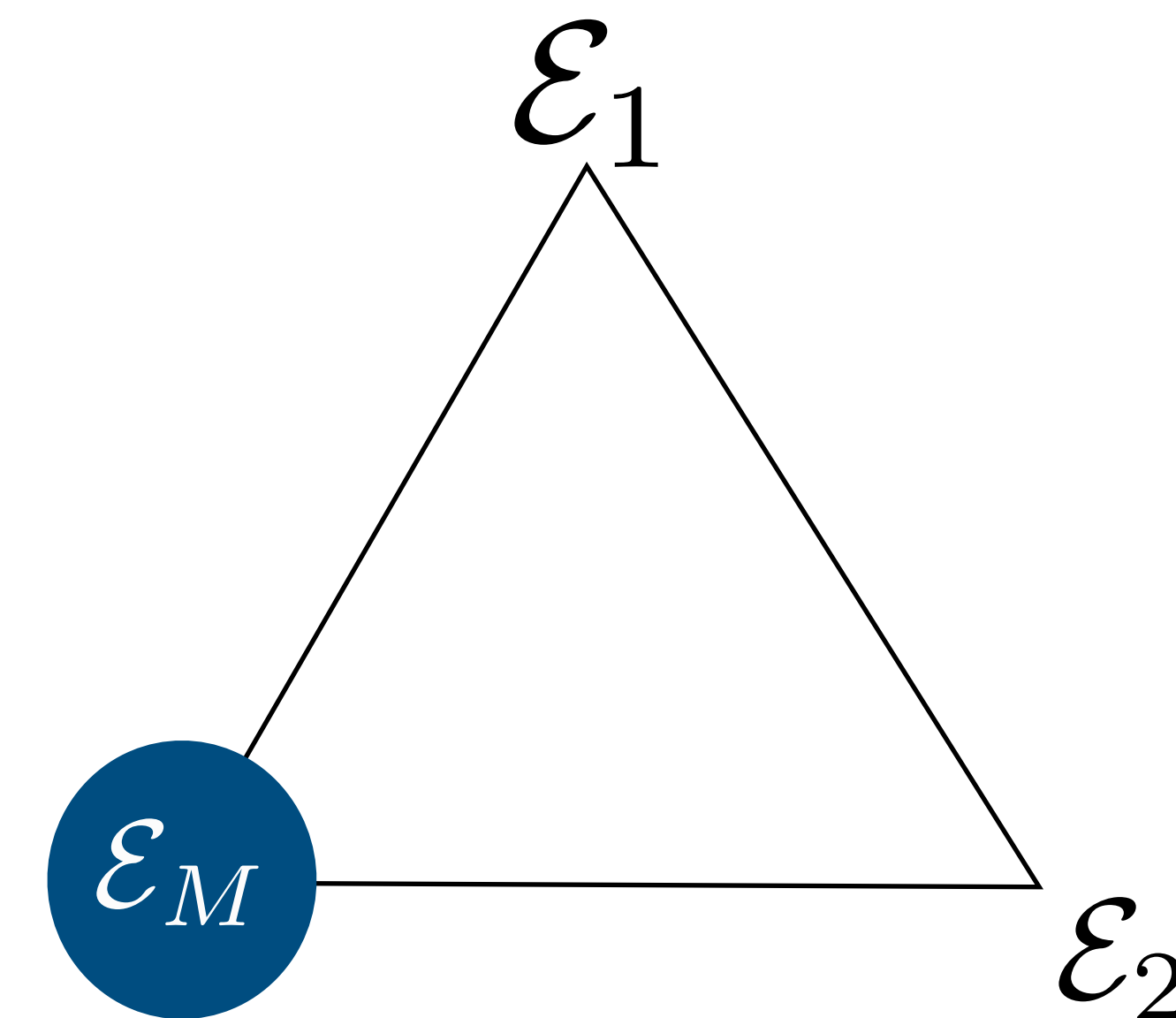
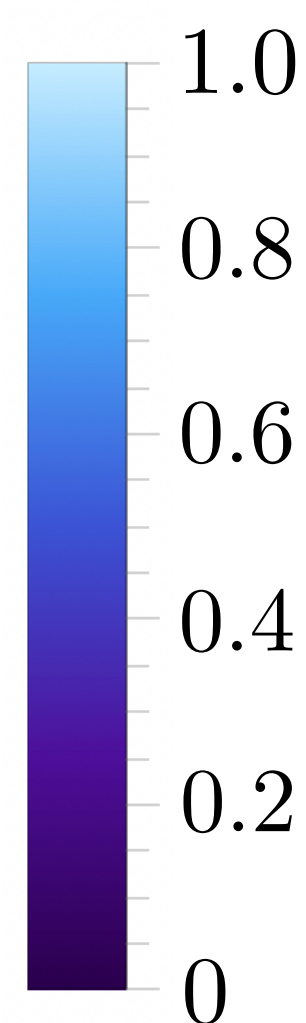
- **Application:** **three-point correlations** probe the non-trivial dynamics of the **dead-cone**

Craft, Gonzalez, KL, Meçaj, Moutl `In Progress

Ratio of Three-Point Massive Correlators



$$\frac{\langle \mathcal{E}_1 \mathcal{E}_2 \mathcal{E}_M \rangle}{\langle \mathcal{E}_1 \mathcal{E}_2 \mathcal{E}_{M=0} \rangle}$$



- **Achieve analytic calculation** using our **1 → 3 splitting functions**

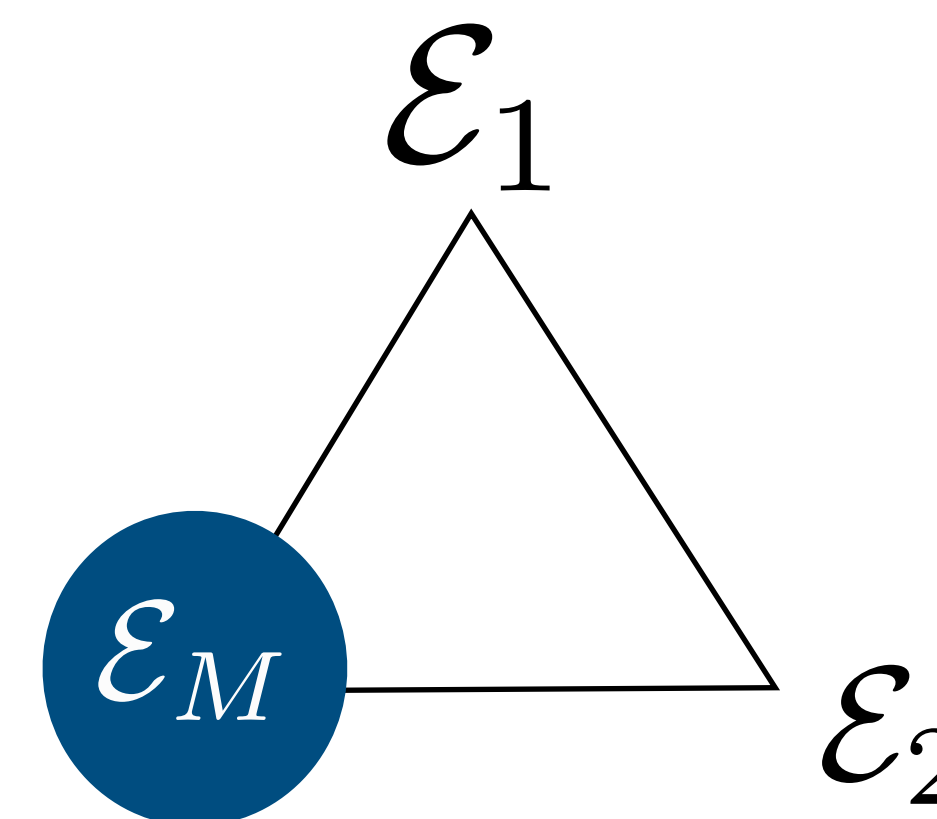
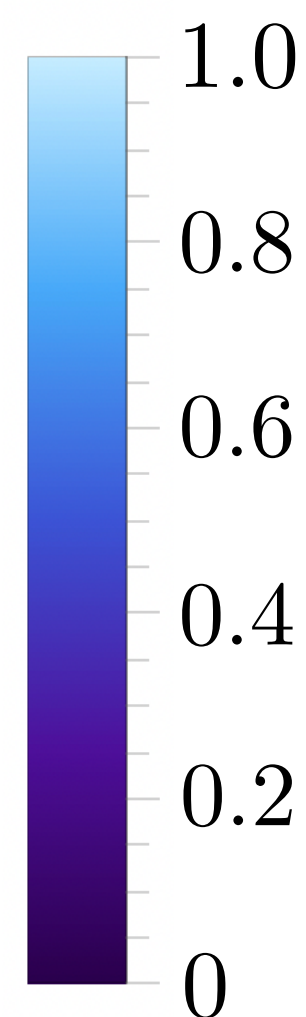
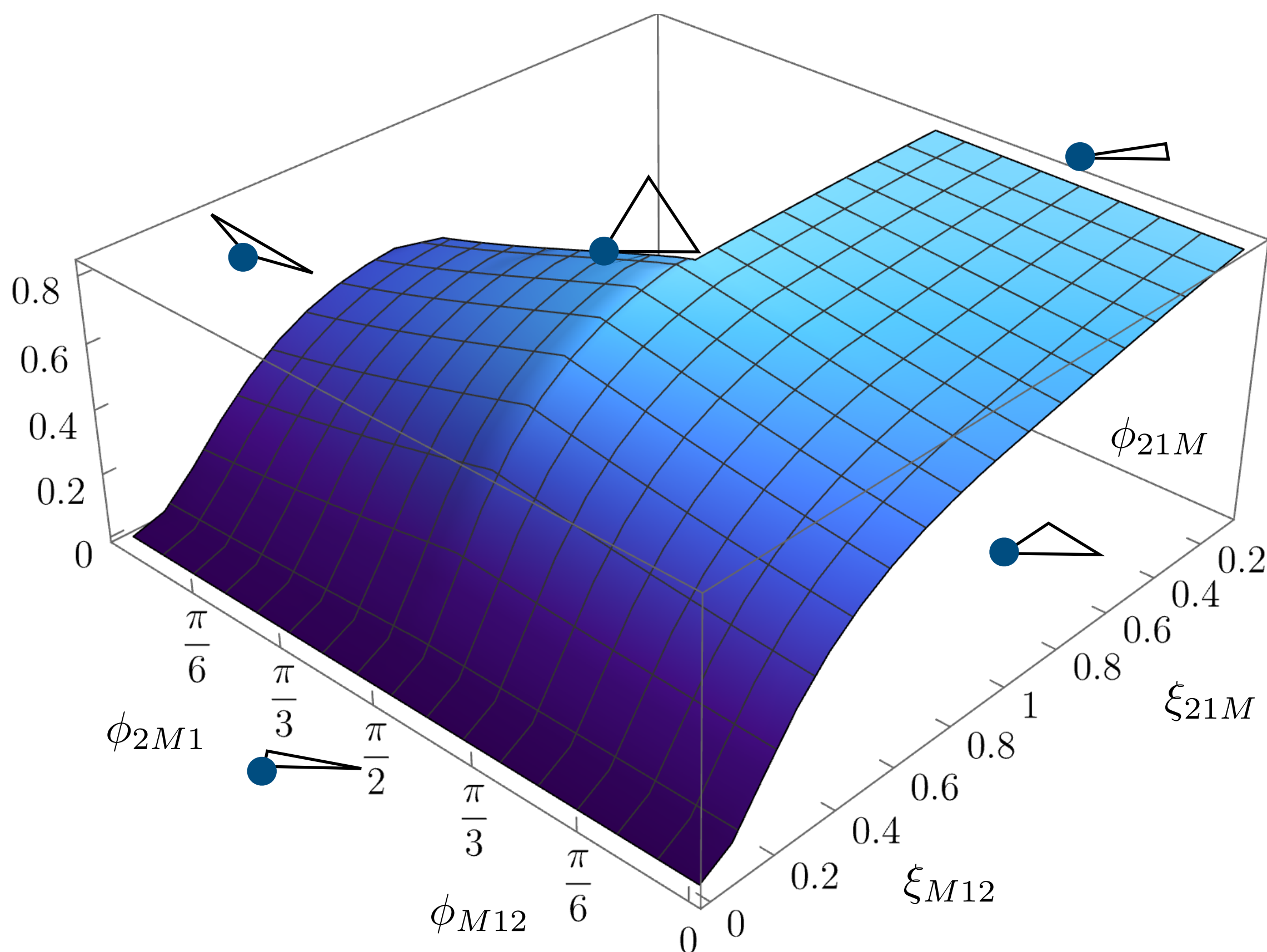
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Craft, Gonzalez, KL, Meçaj, Moutl `In Progress

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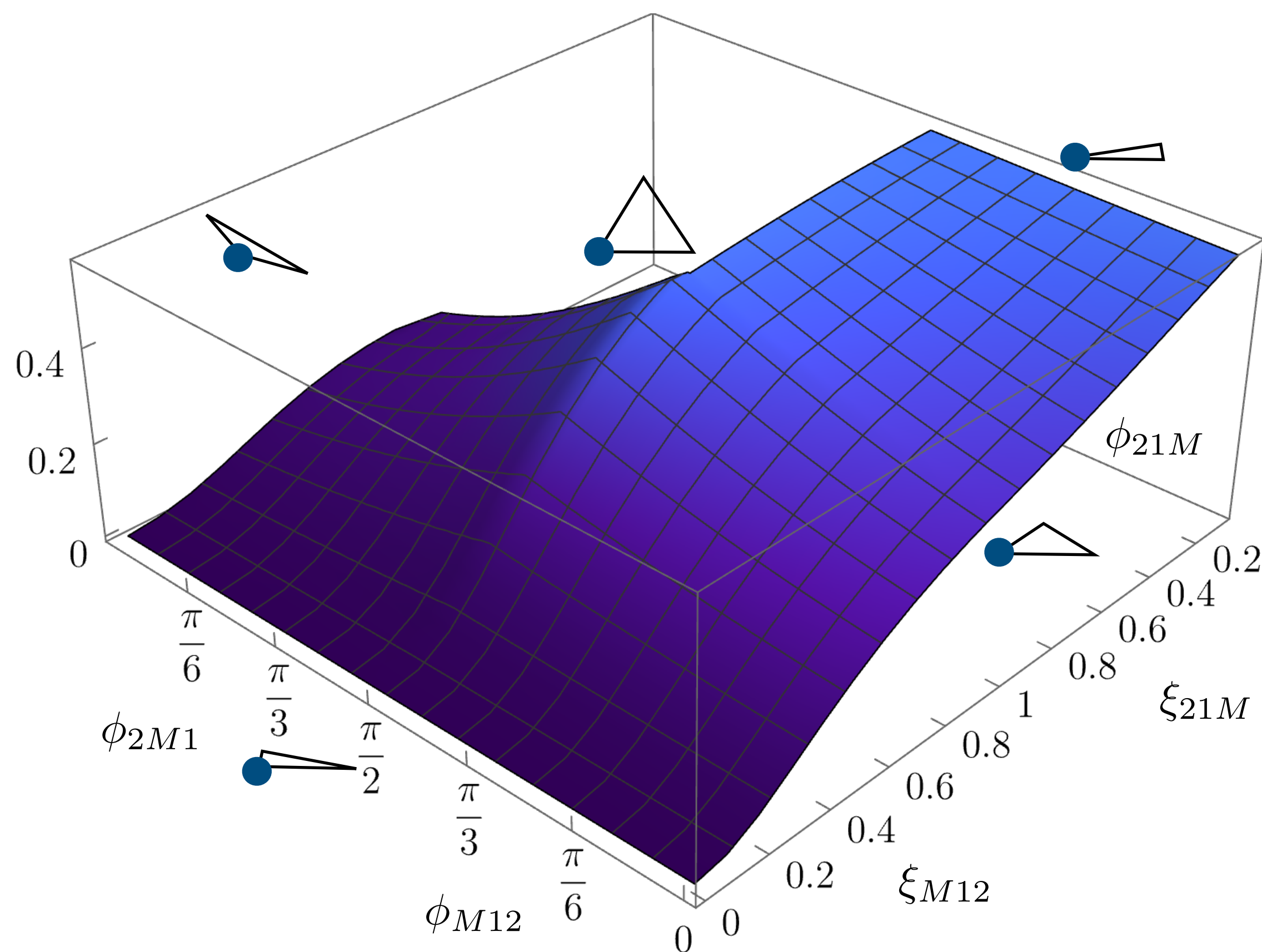
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PROBING THE DYNAMICS OF THE DEAD-CONE

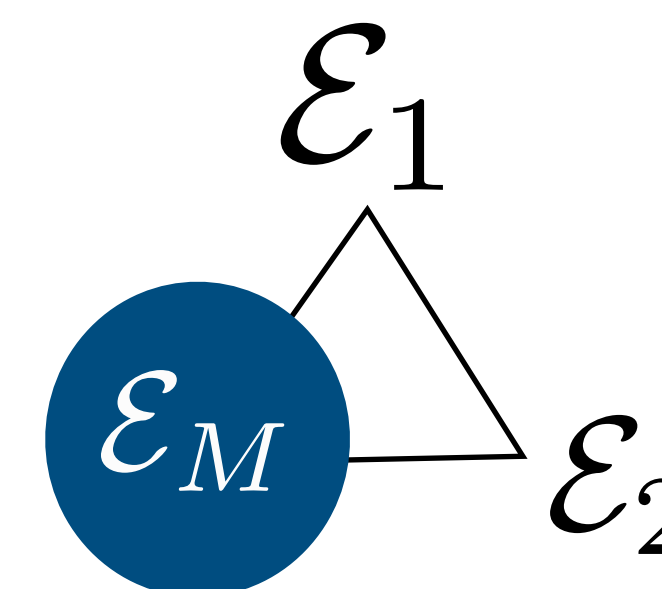
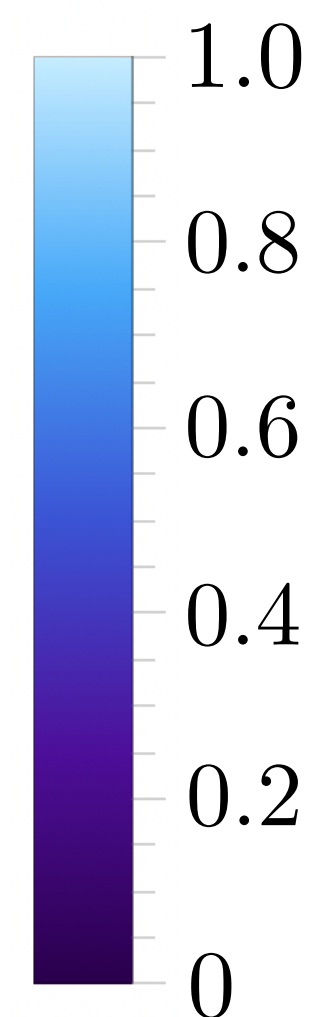
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Craft, Gonzalez, KL, Meçaj, Moutl `In Progress

Ratio of Three-Point Massive Correlators

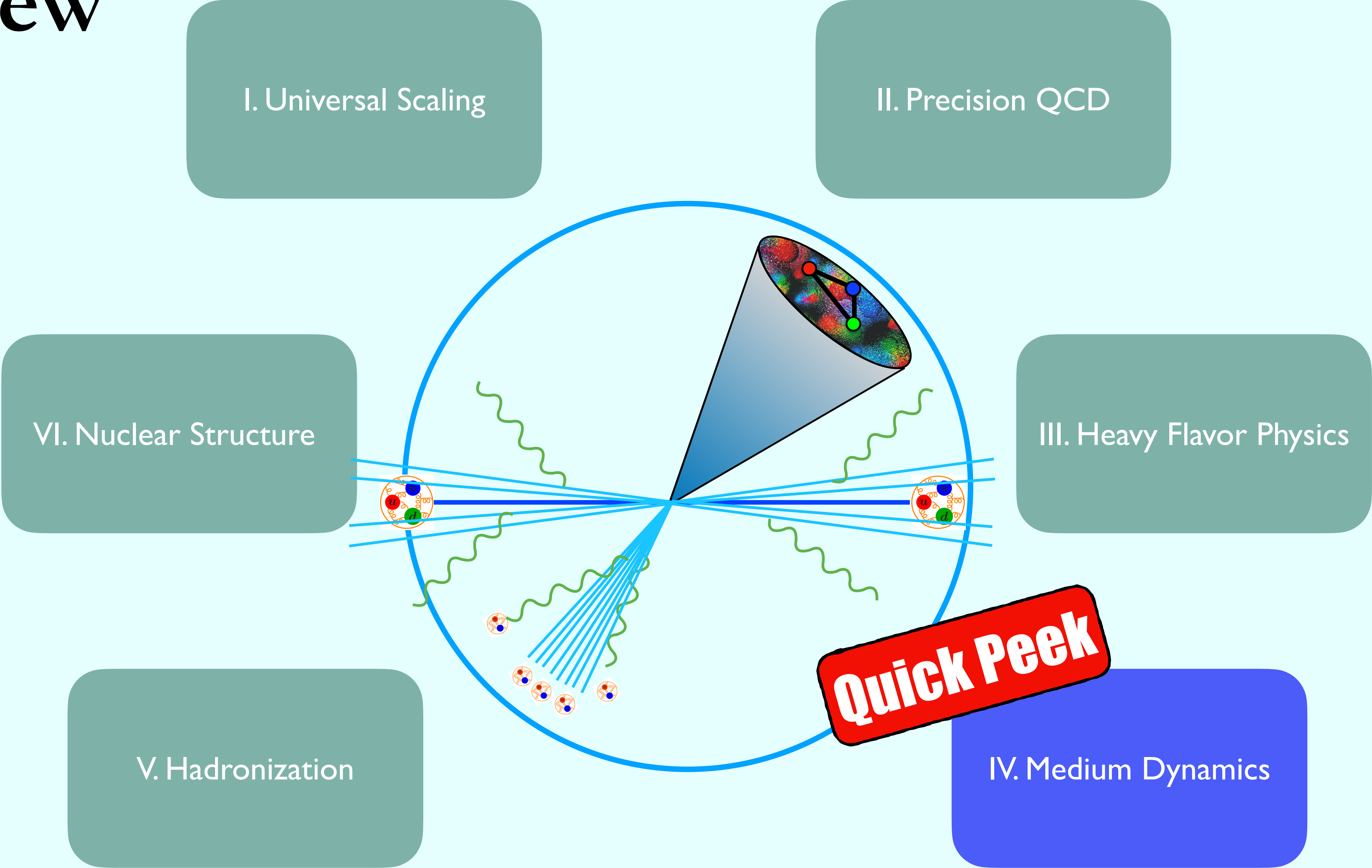


$$\frac{\langle \mathcal{E}_1 \mathcal{E}_2 \mathcal{E}_M \rangle}{\langle \mathcal{E}_1 \mathcal{E}_2 \mathcal{E}_{M=0} \rangle}$$



- **Achieve analytic calculation using our 1 → 3 splitting functions**

Overview

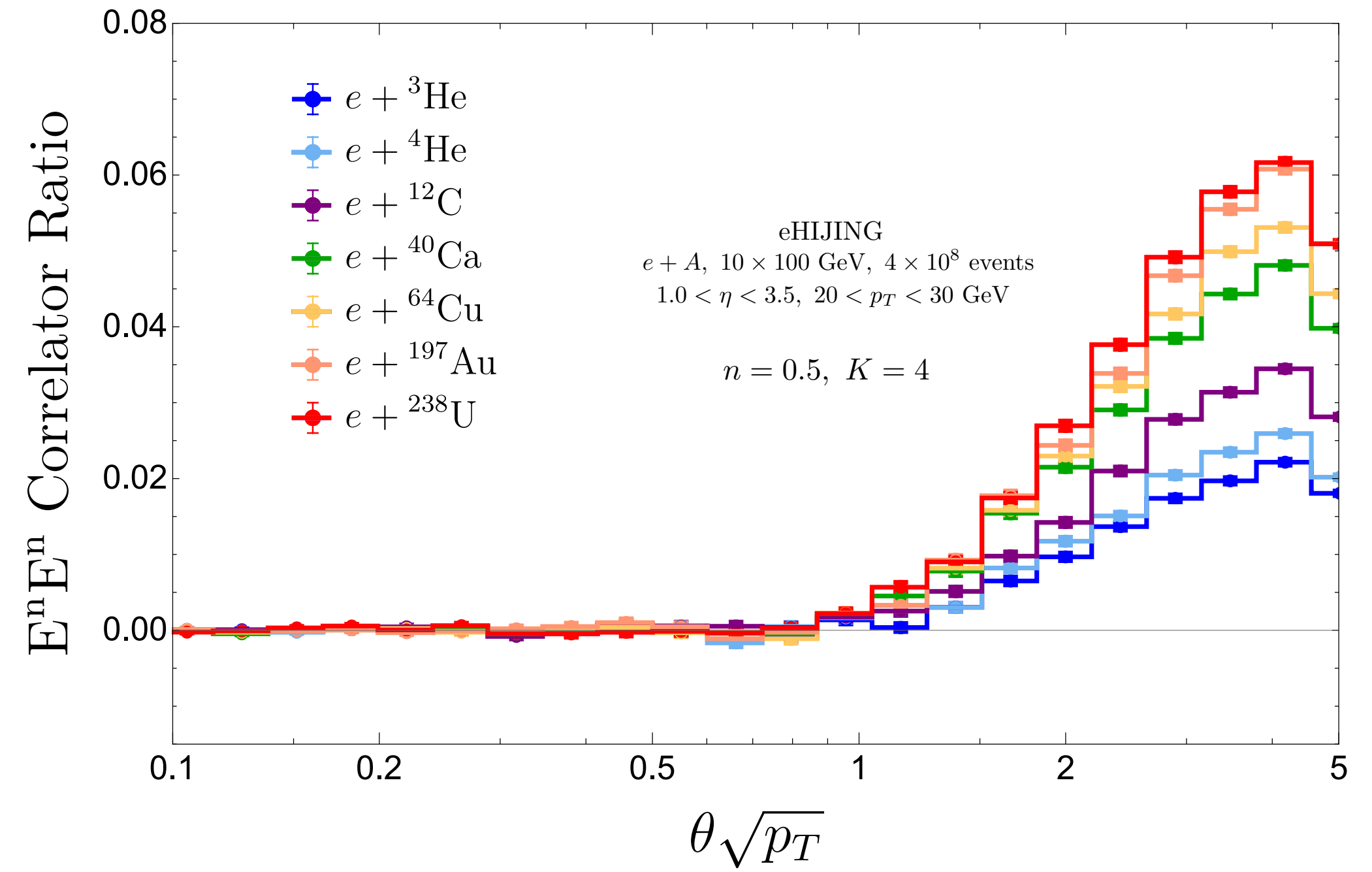
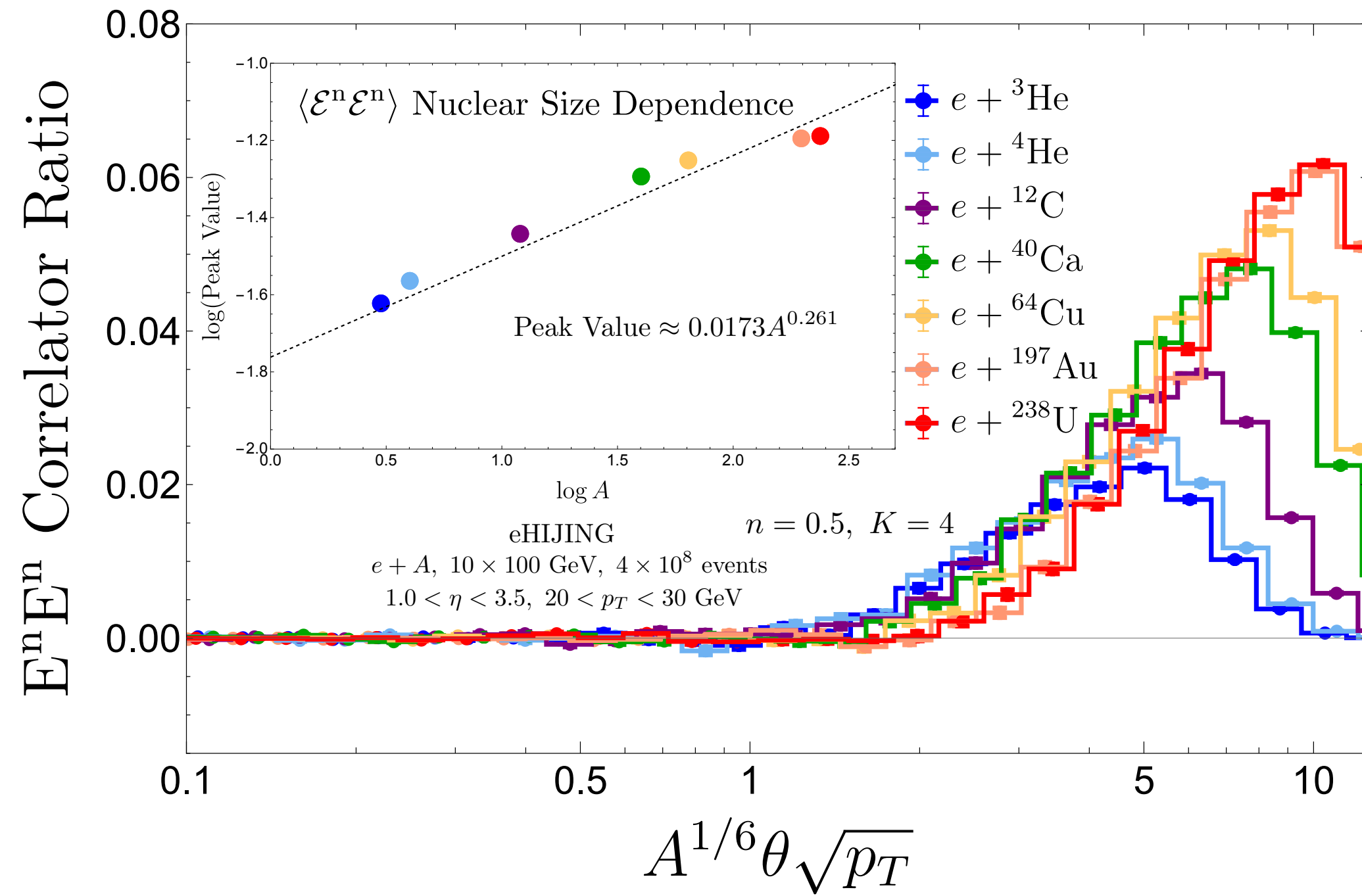
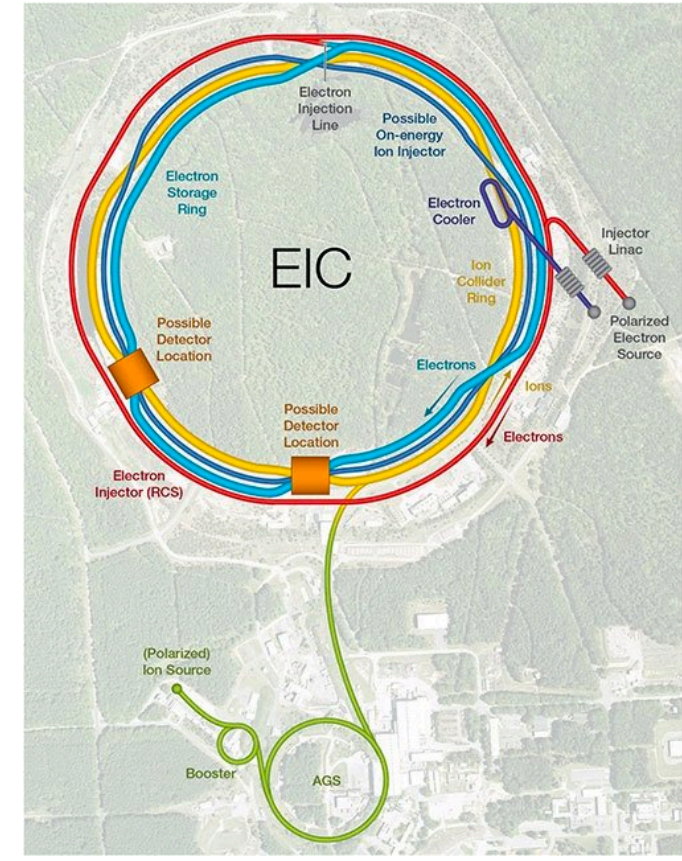
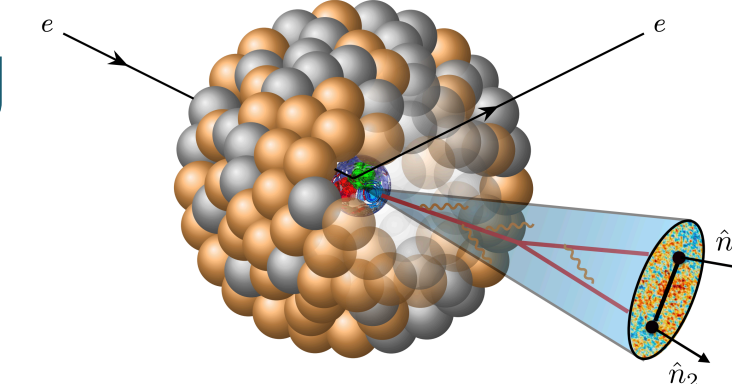


RESOLVING THE FEMTOSCALE IN JETS

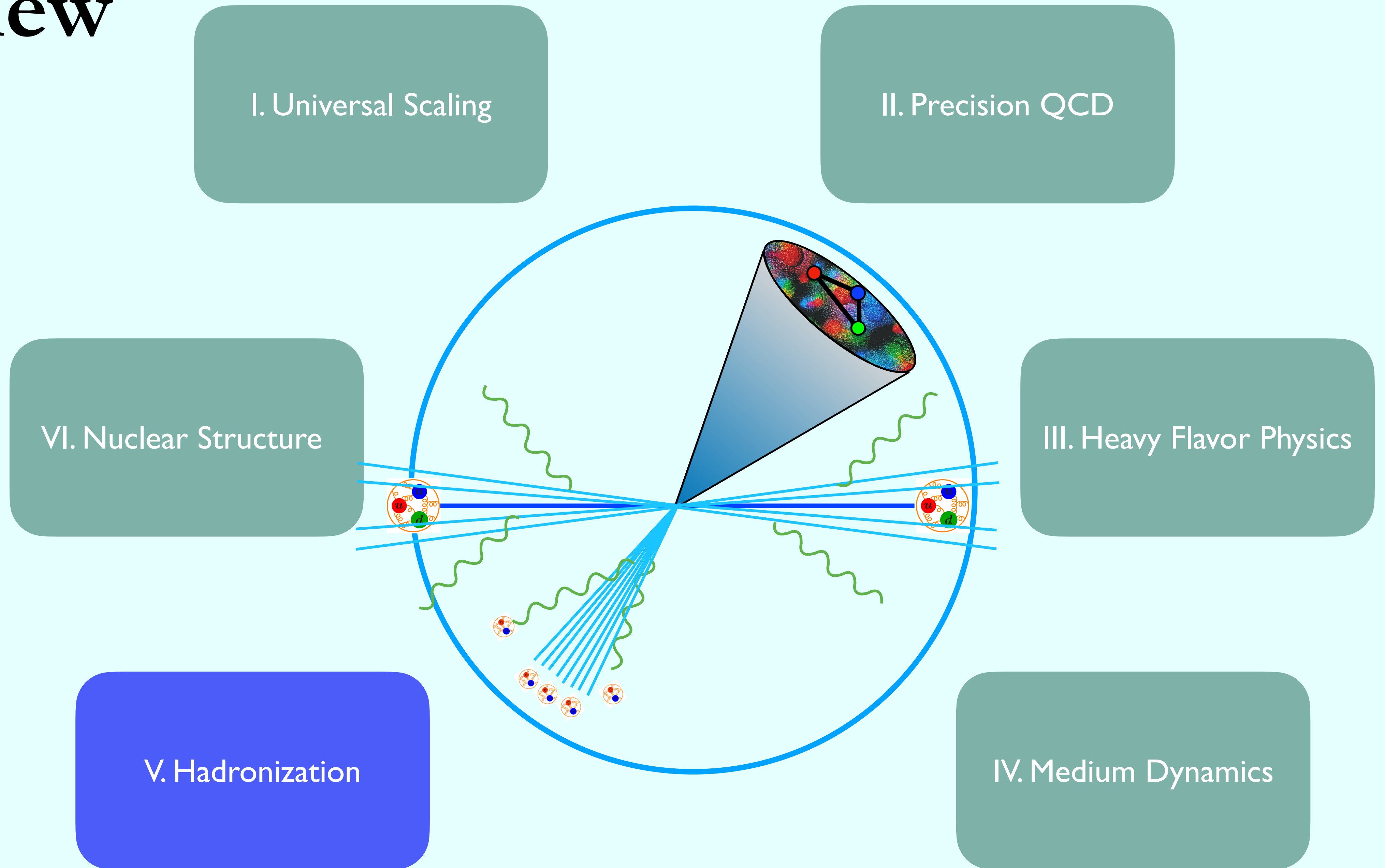
Devereaux, Fan, Ke, KL, Moul '23

- **Femtosecond nuclear size dependence** can be resolved within jets using **two-point correlations**

$$\theta_{\text{nucl}} \sim \frac{1}{\sqrt{p_T L}} \sim \frac{1}{\sqrt{p_T} A^{1/6}}$$



Overview



WHAT IS A DETECTOR?



- What constitutes a well-defined **field theory definition for a detector?**

Caron-Huot, Kologlu, Kravchuk, Meltzer, Simmons-Duffin `22

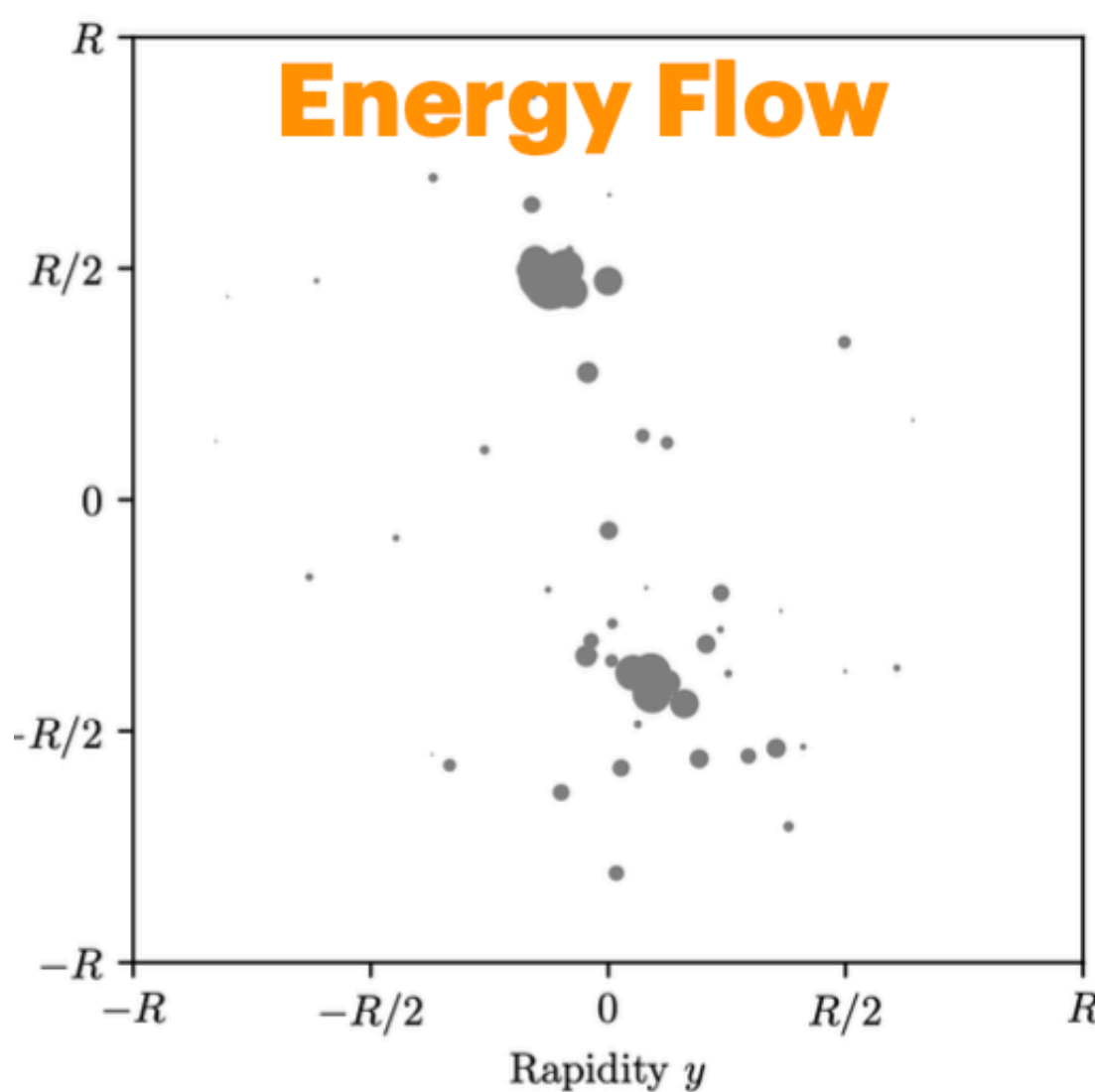
- Interesting measurements of energy flow can be made on a **restricted set of hadronic states, R** , for example, **charged hadrons (tracks)**

$$\mathcal{E}_R = \sum_{i \in R} \mathcal{E}_i$$

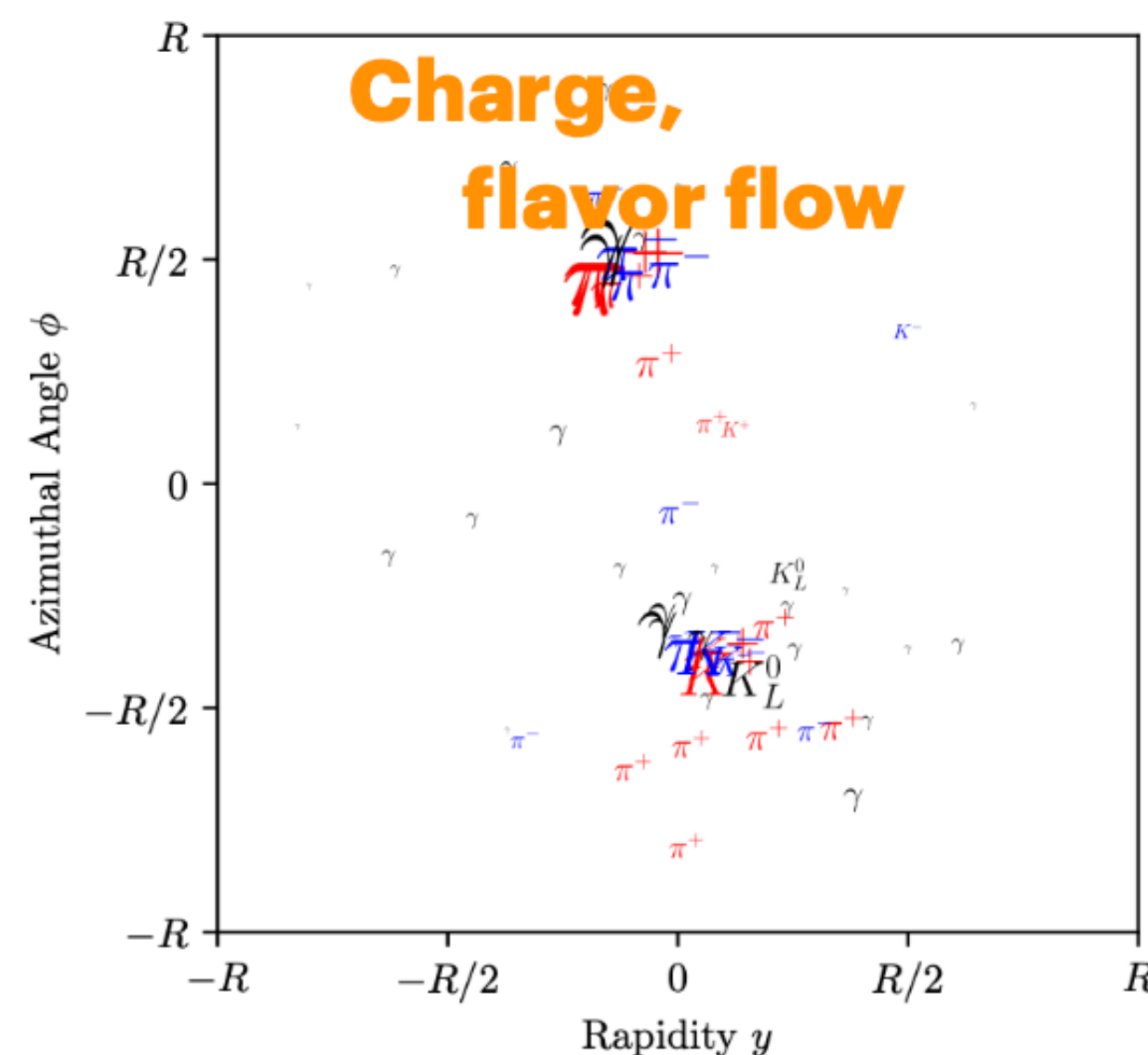
- Provides a **sharp link** between underlying field theory and observables

KL, Moulton `23

The **energy** flow is unpixelized and ignores charge/flavor information



Full event is a set of particles having momentum and charge/flavor



All observables



Well-defined detectors



\mathcal{E}_\pm

IRC safe

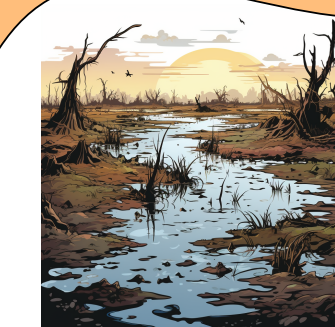
\mathcal{E}



\mathcal{E}_Q



\mathcal{E}_n



WHAT IS A DETECTOR?



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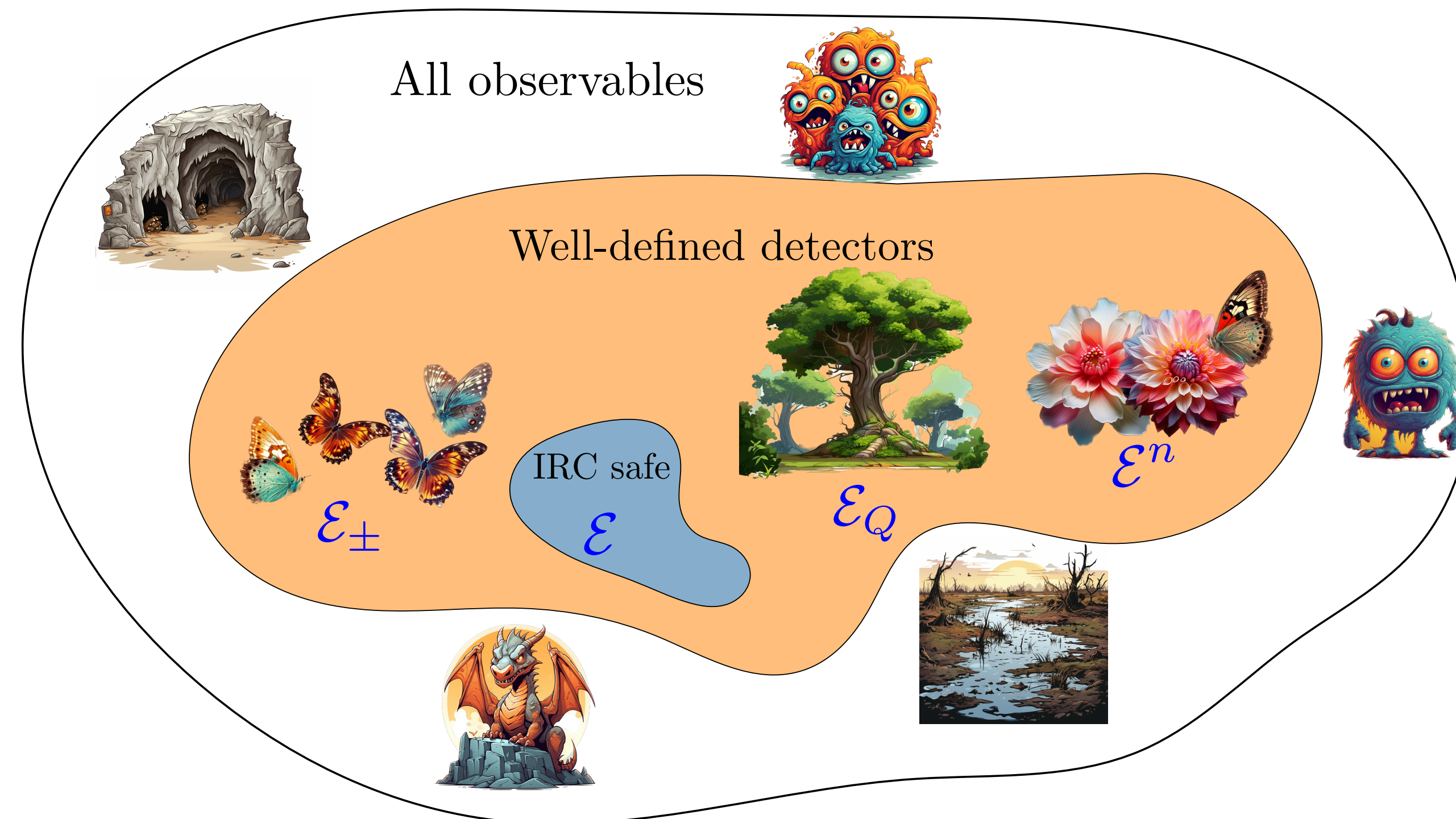
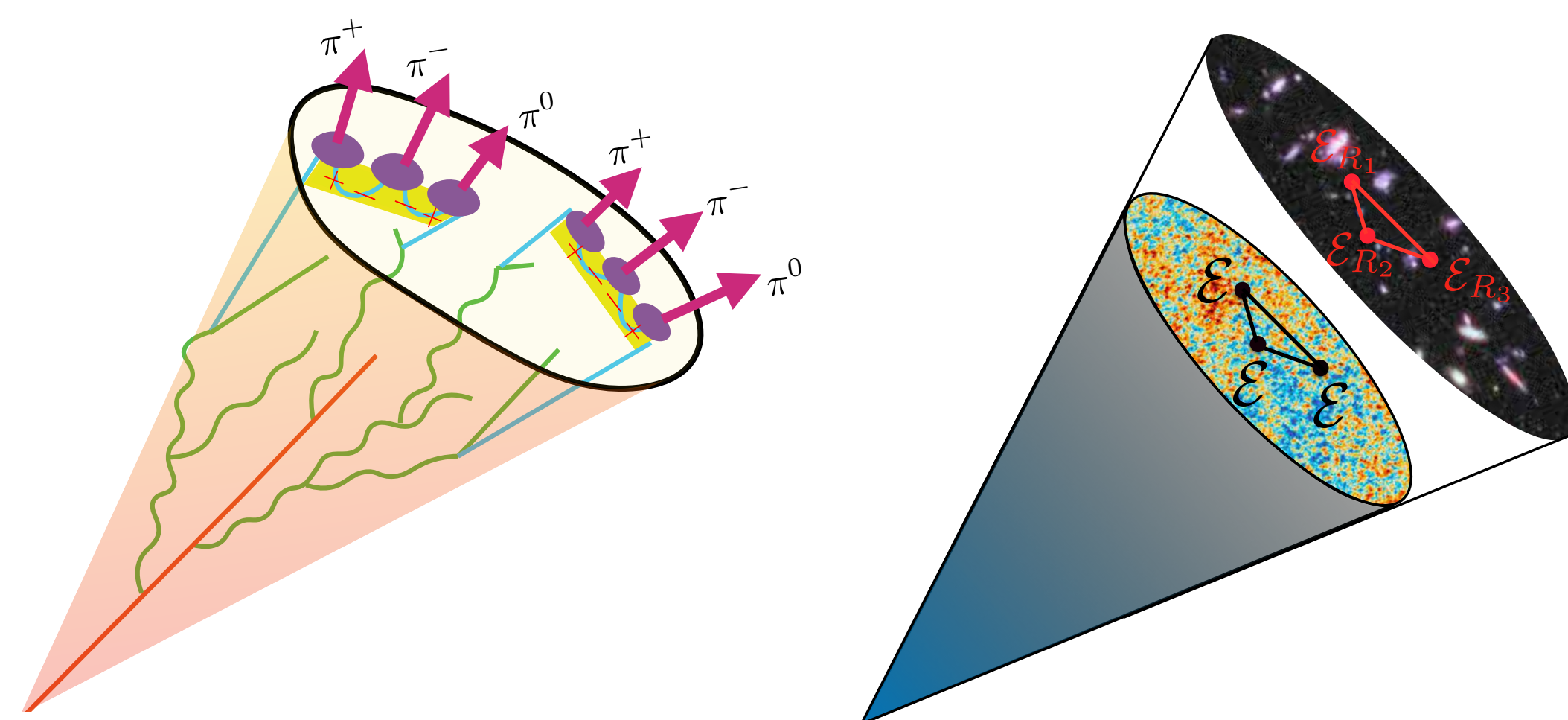
Caron-Huot, Kologlu, Kravchuk, Meltzer, Simmons-Duffin `22

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KL, Moulton `23

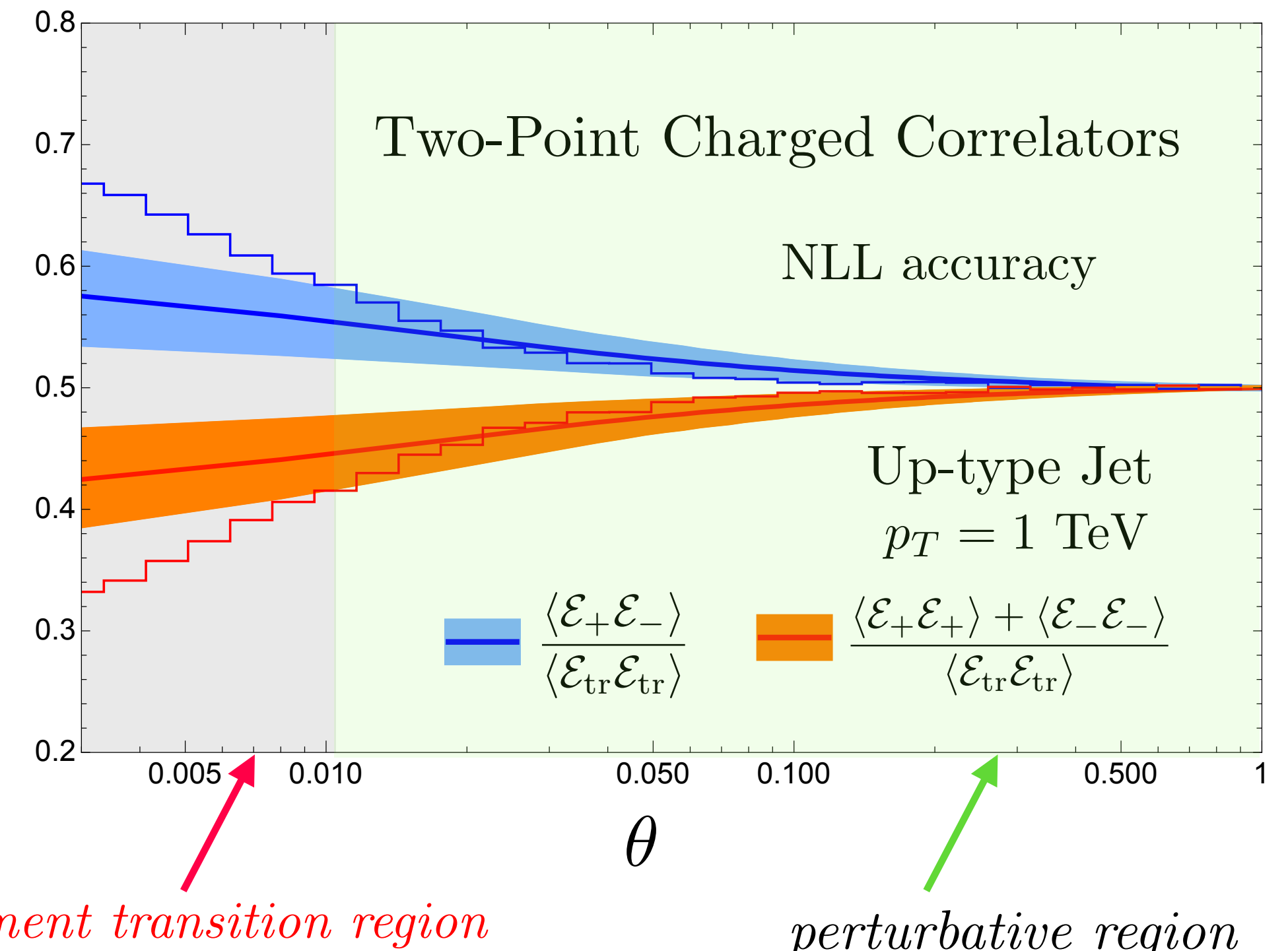
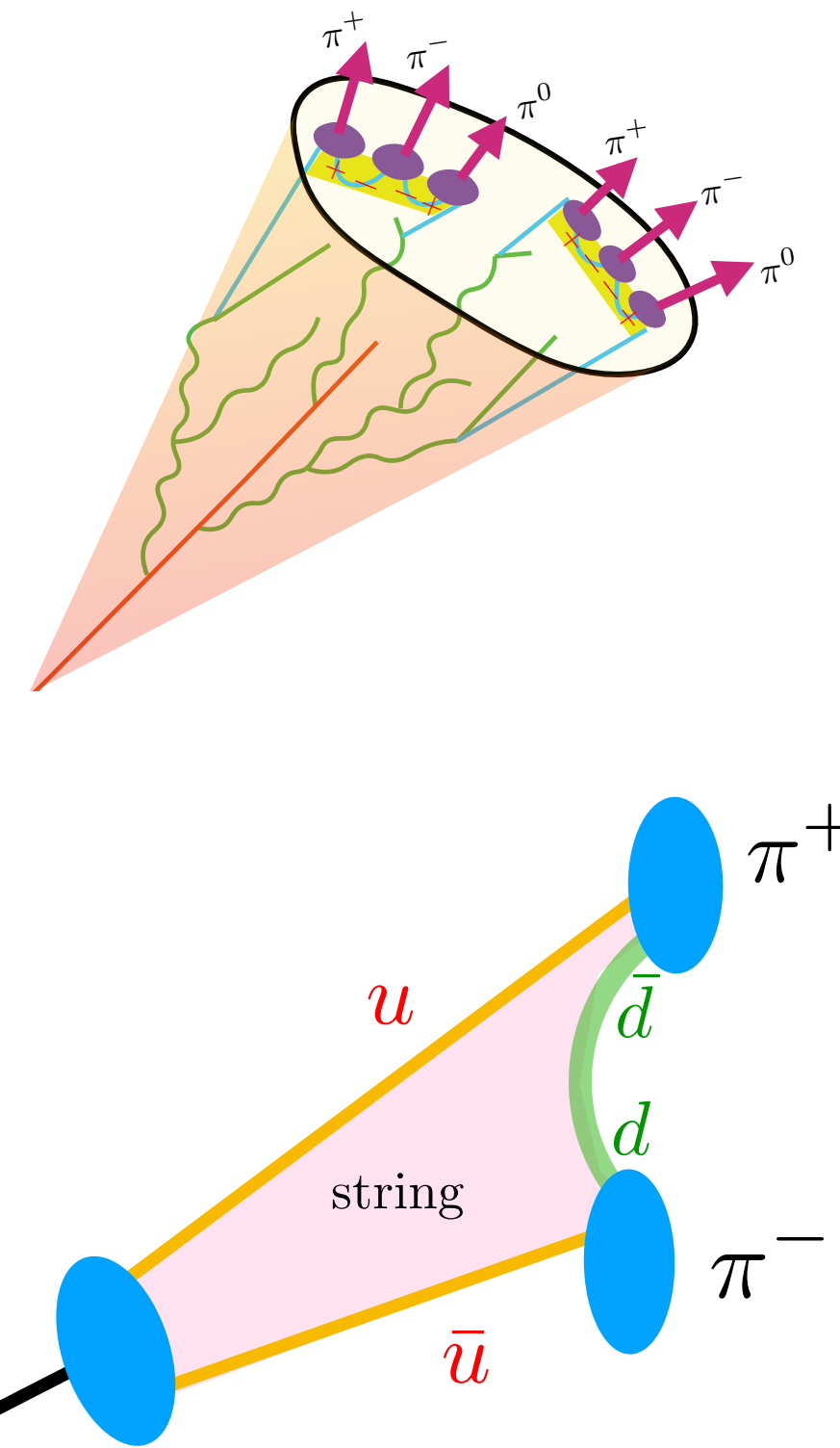
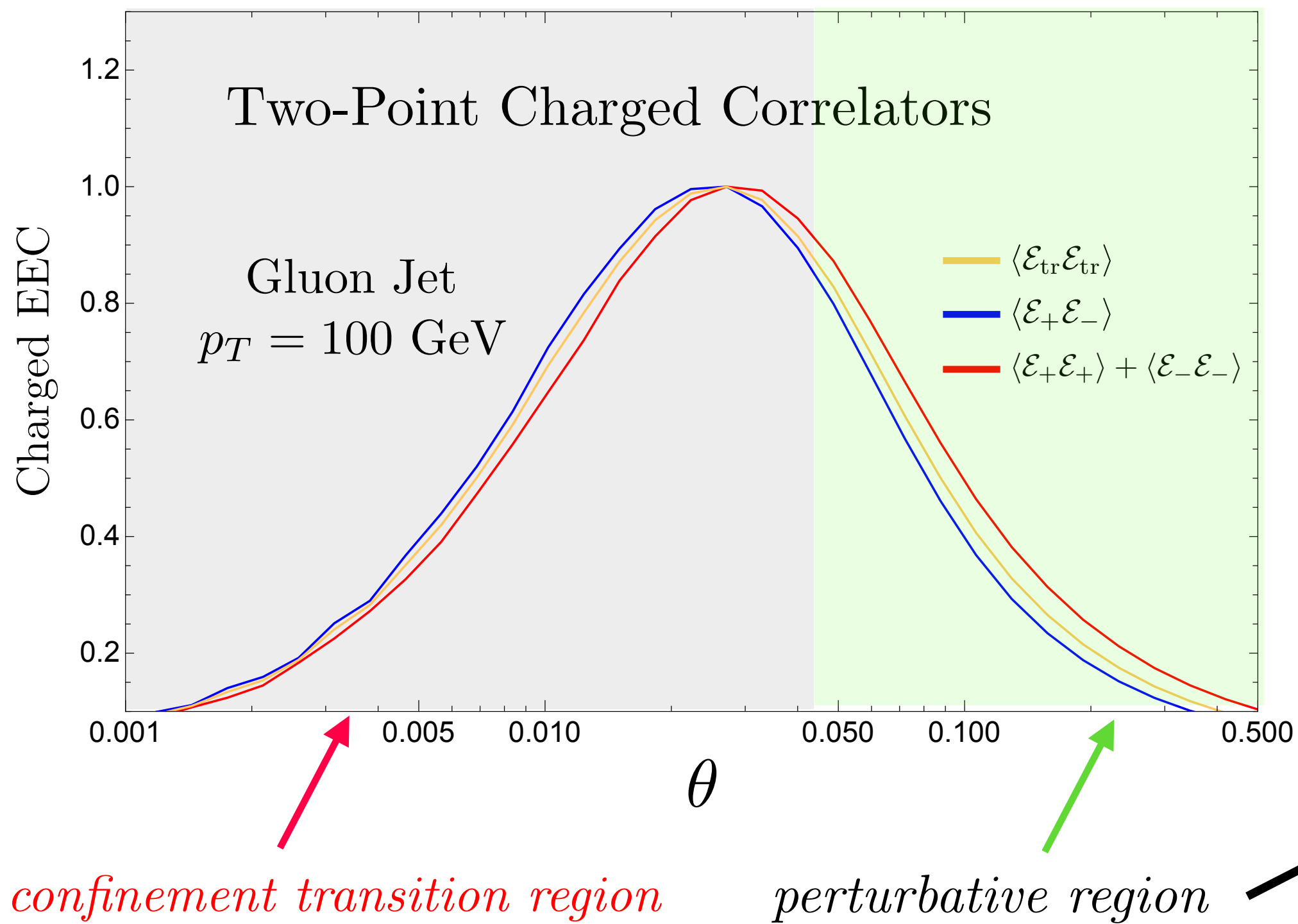


CORRELATION BETWEEN CHARGED HADRONS

KL, Moulton '23

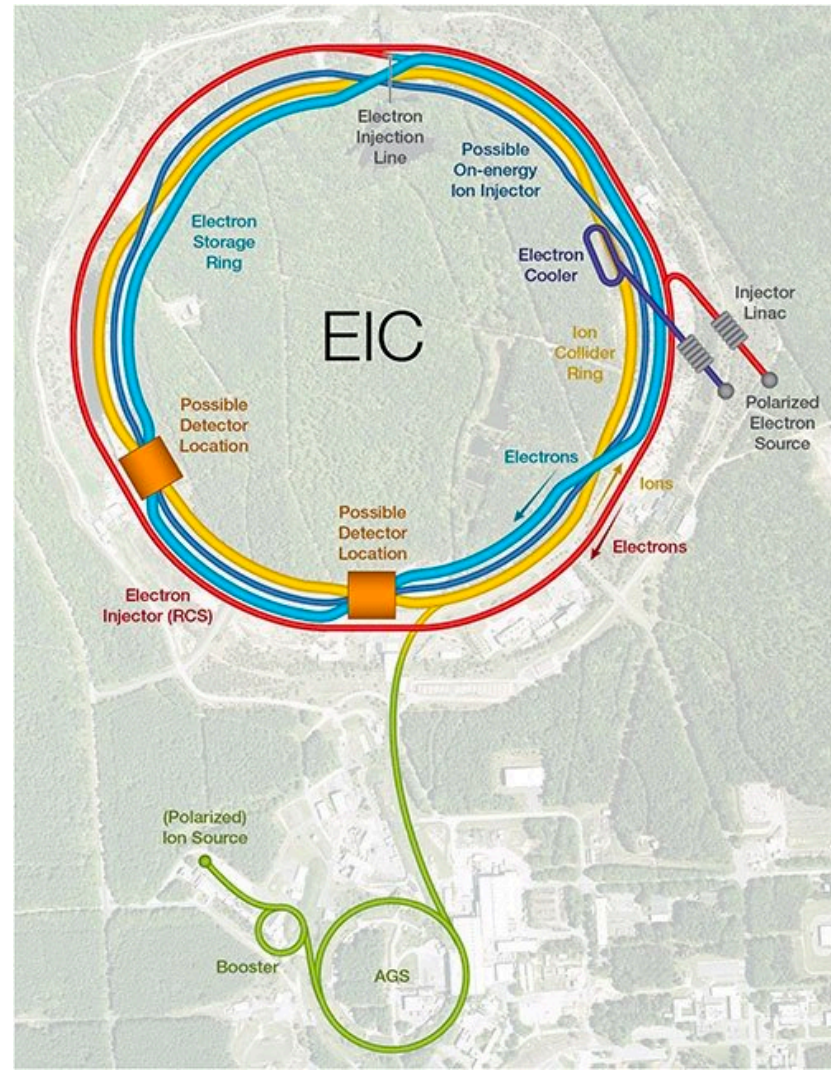
- Unlike-signed charged correlators are **correlated more** as the angle becomes smaller!

$$\langle \mathcal{E}_+ \mathcal{E}_- \rangle, \langle \mathcal{E}_+ \mathcal{E}_+ \rangle, \langle \mathcal{E}_- \mathcal{E}_- \rangle$$

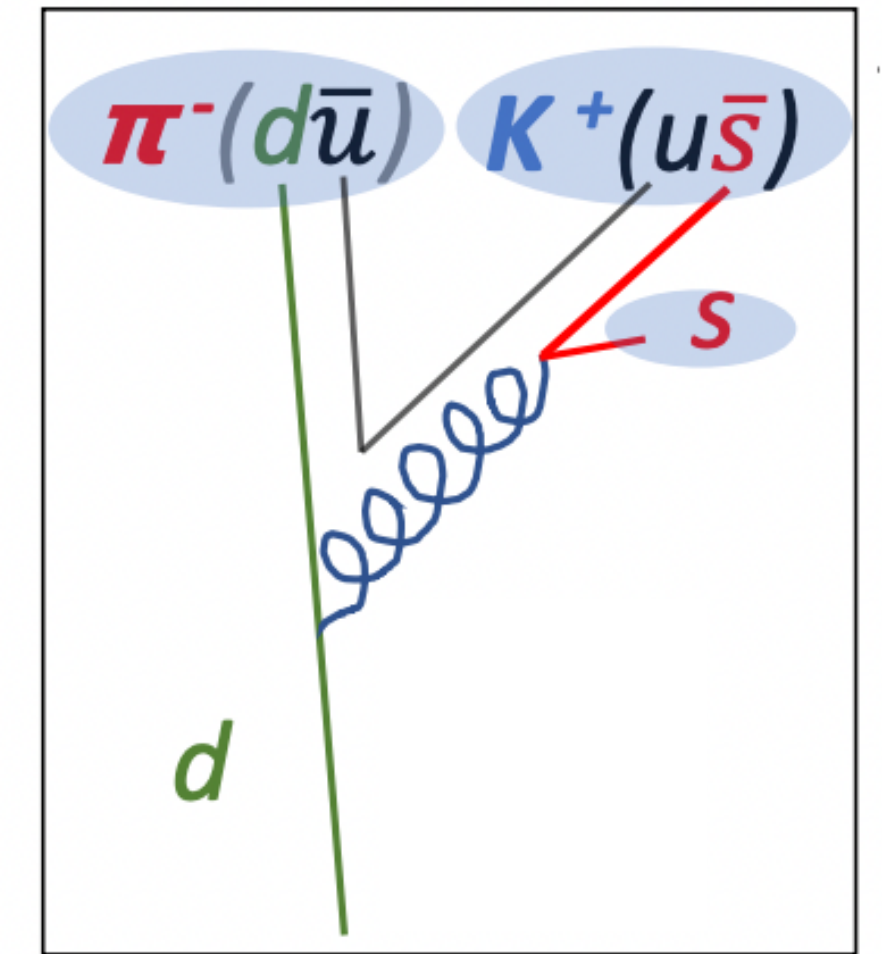
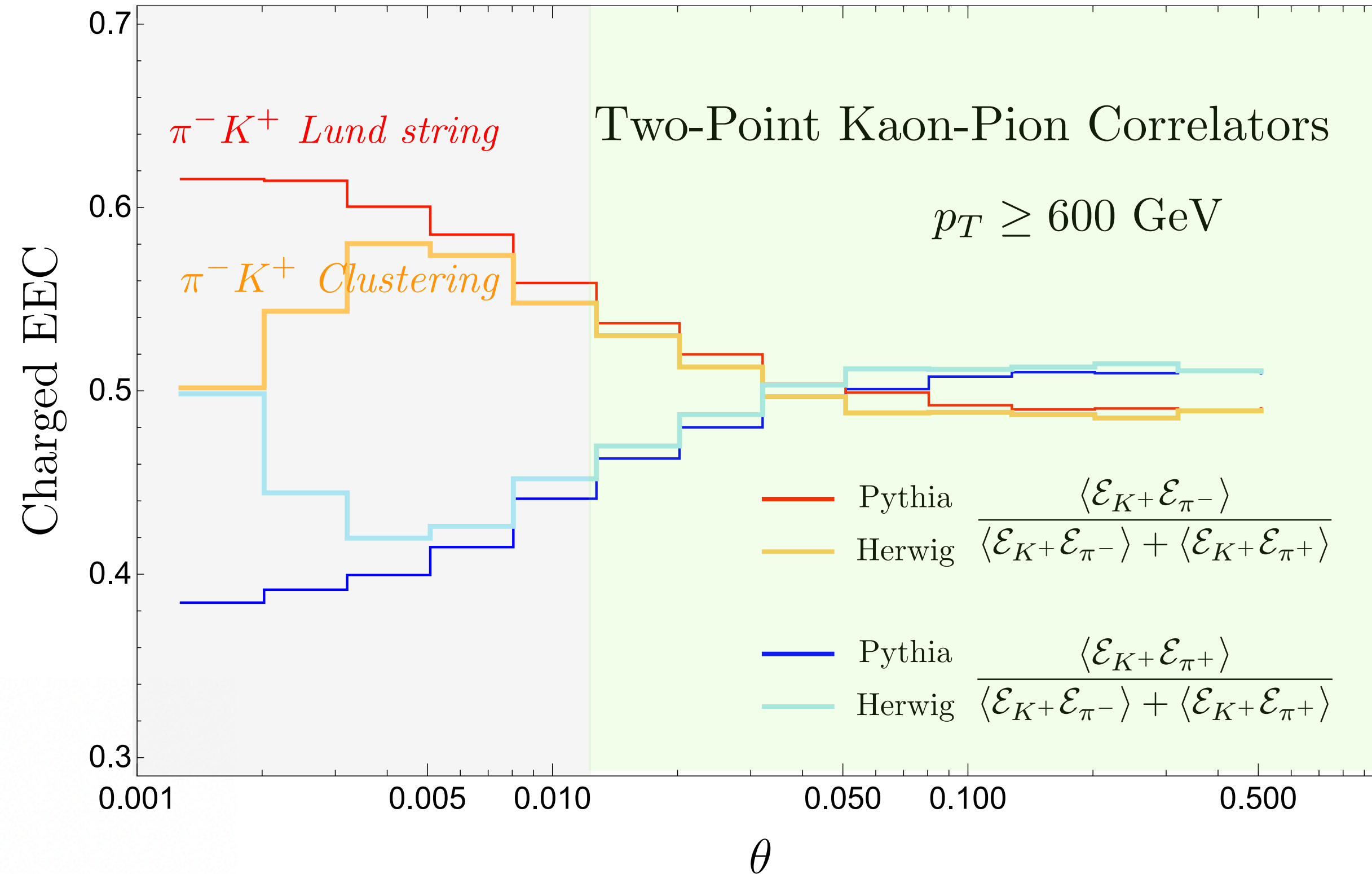


- The correlation between unlike-signed hadron pair is expected to grow in **string-like hadronization**

DISCRIMINATING HADRONIZATION MECHANISMS

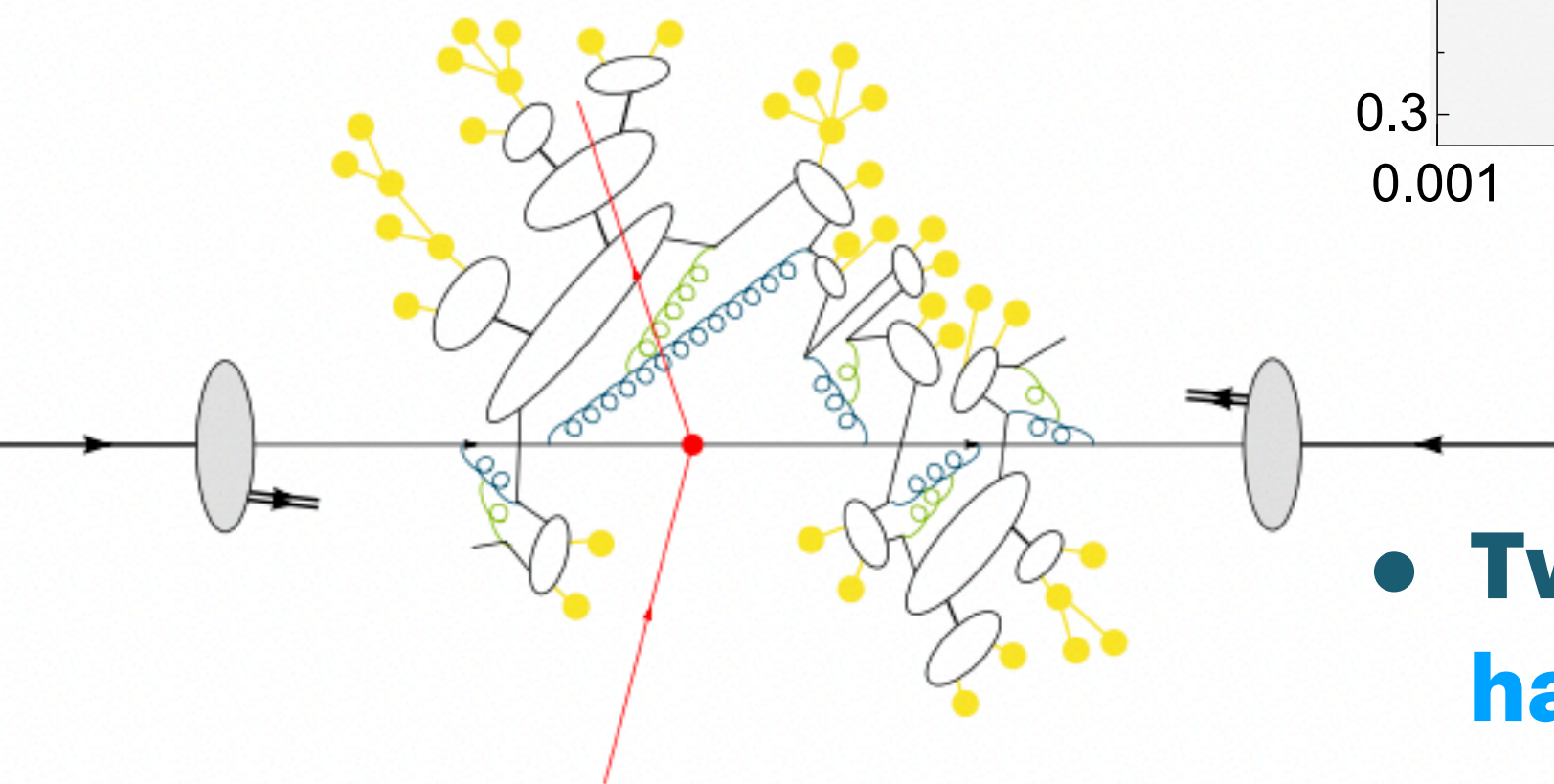


KL, Mout, Song, Sterman `In Progress



Lund string model (Pythia)

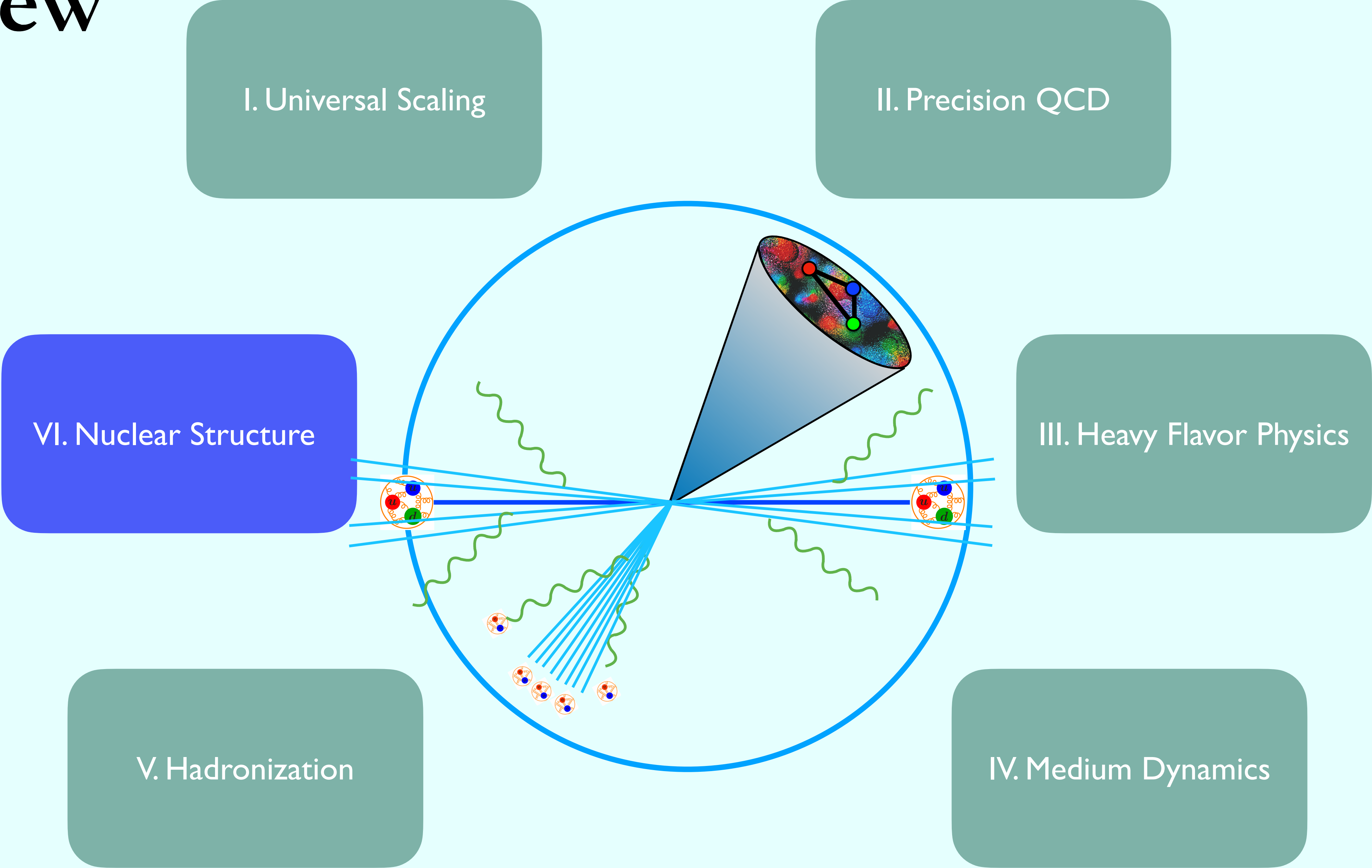
Clustering model (Herwig)



See also Chien, Deshpande, Mondal, Sterman

- Two-point charged correlators already **nontrivially probe** the two hadronization mechanisms by eye, and pave the path to go even beyond!

Overview

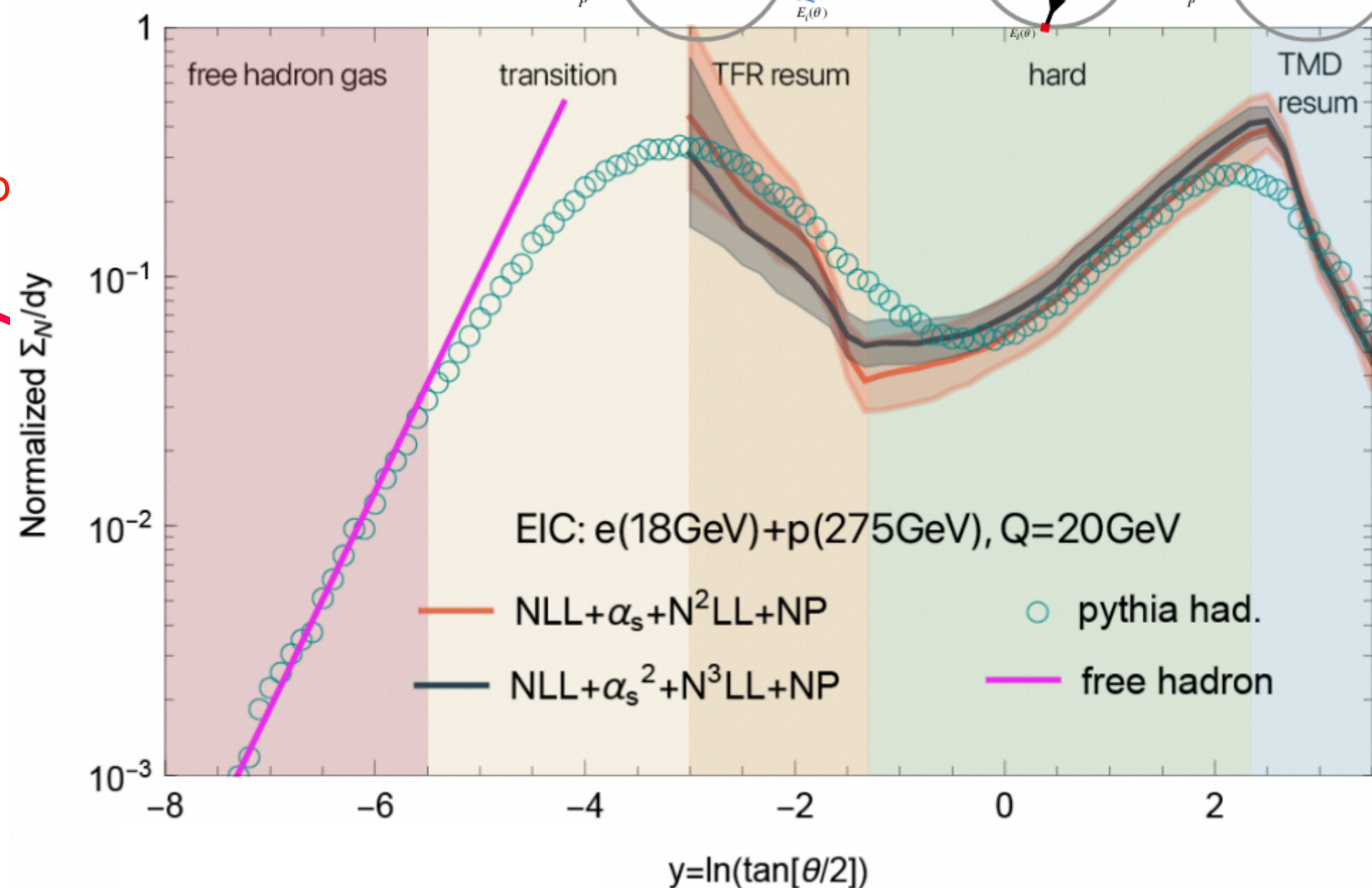


NUCLEON ENERGY CORRELATORS

- Recent proposal to measure the target remnants can probe nuclear structure directly!

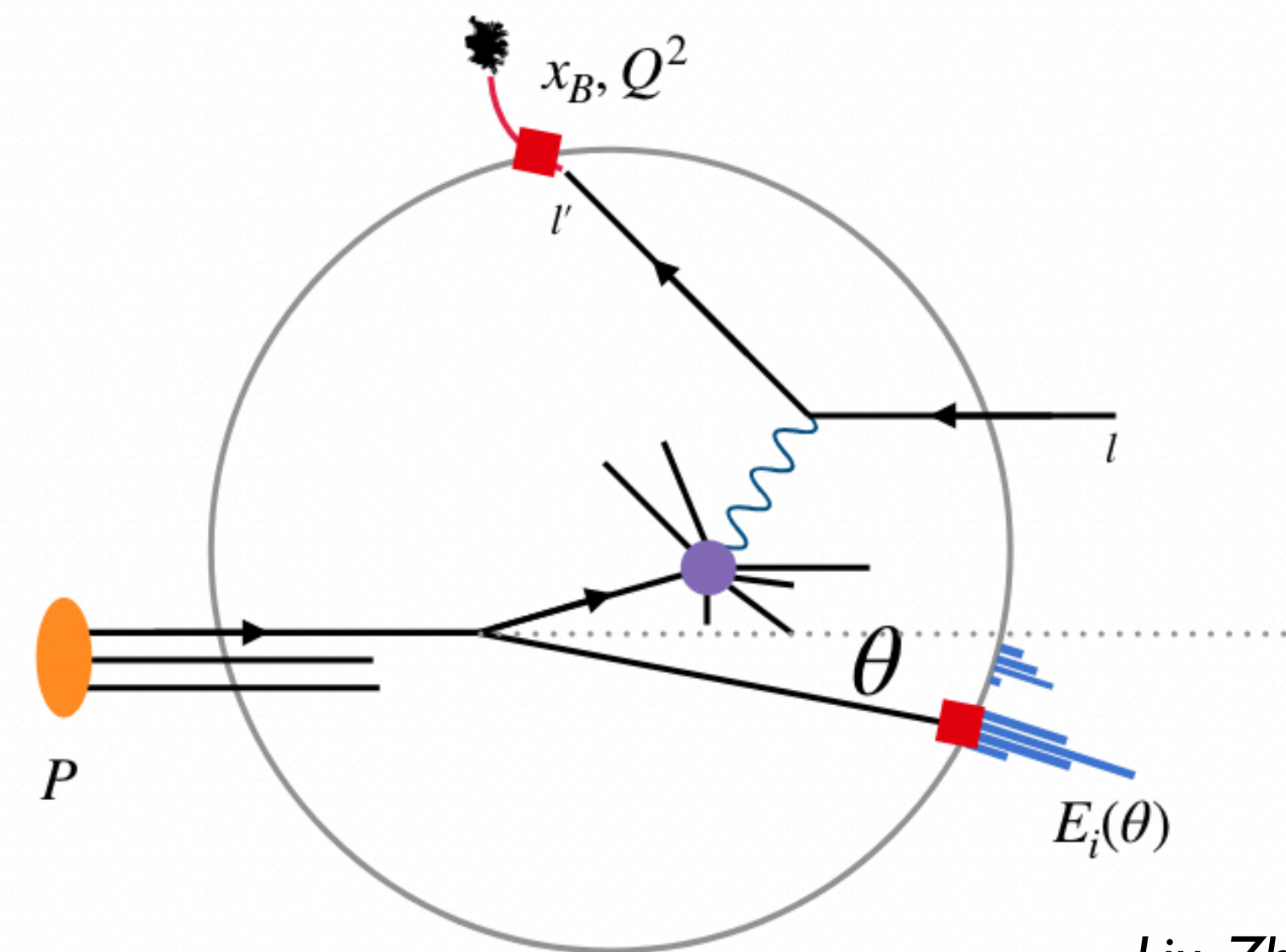
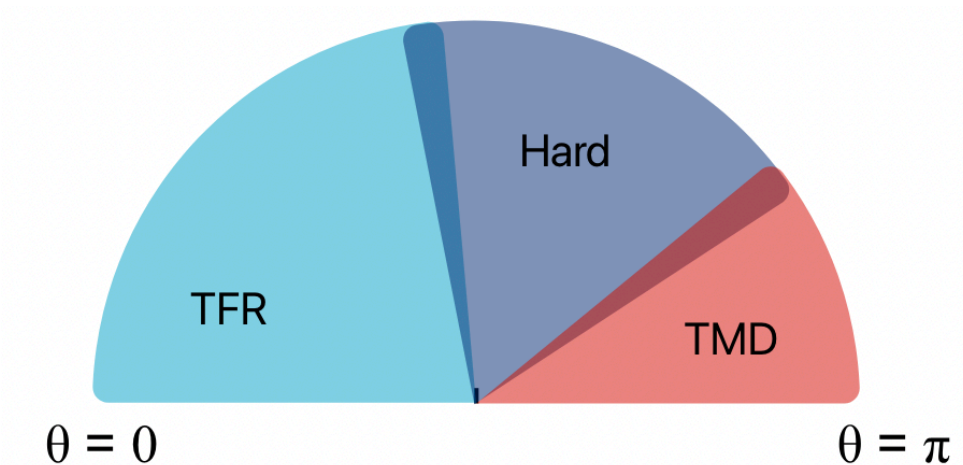
Collinear PDFs and Fracture Functions

NEEC



NP hadronizatio

Probes TMD structure

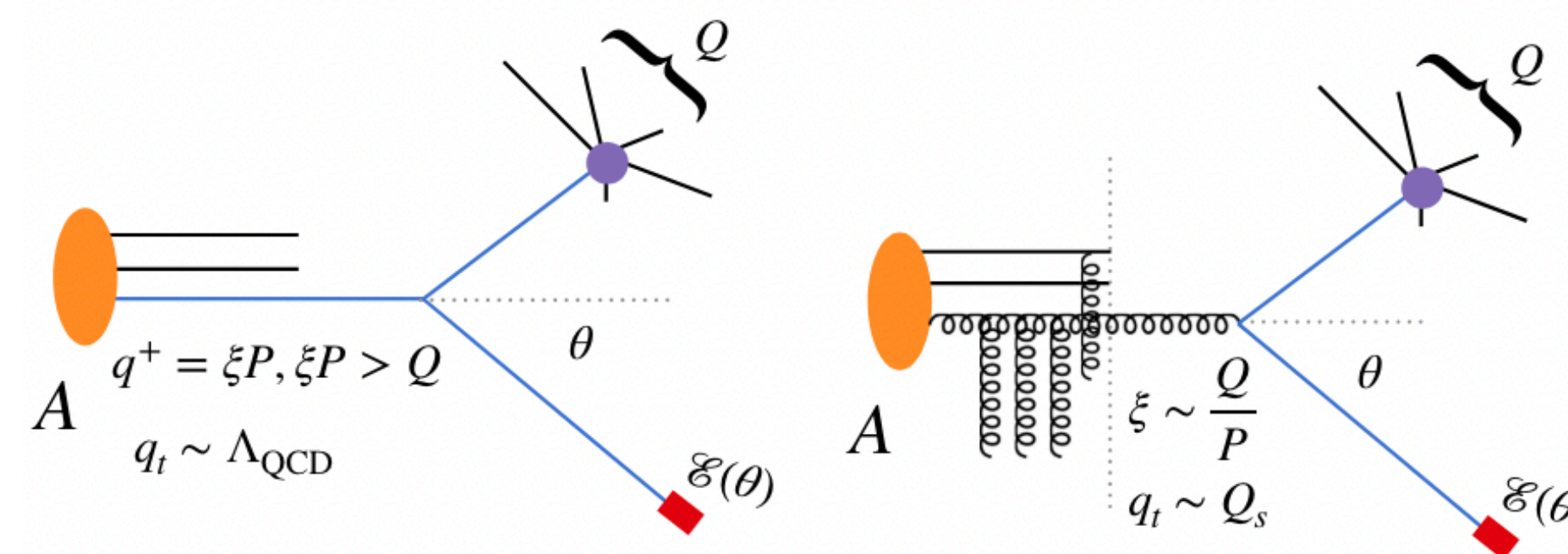
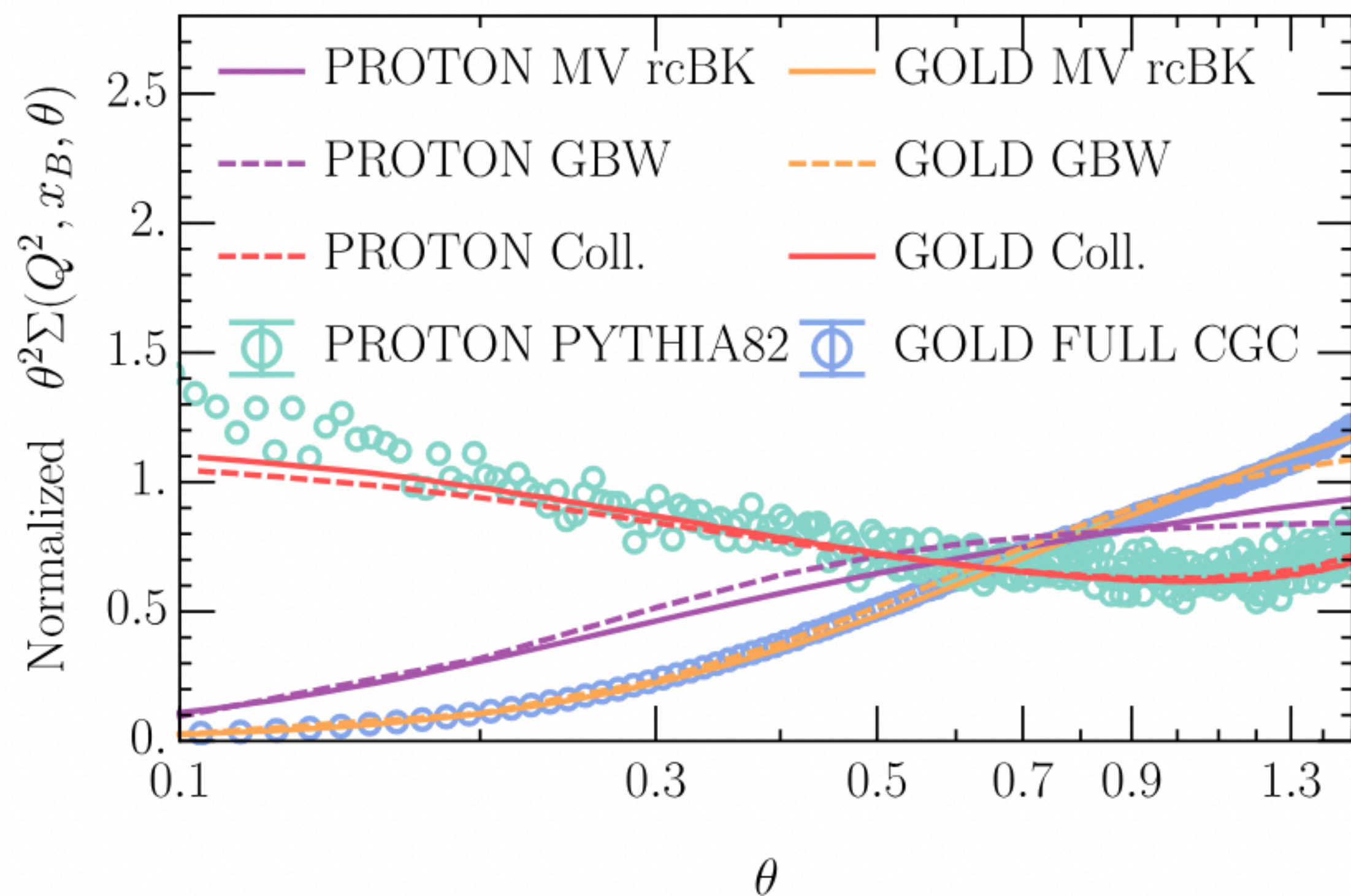


Liu, Zhu `24
 Can, Li, Mi `23
 Liu, Zhu `22

NUCLEON ENERGY CORRELATORS

- Deviation from collinear factorization showing sensitivity to the saturation scale Q_s !

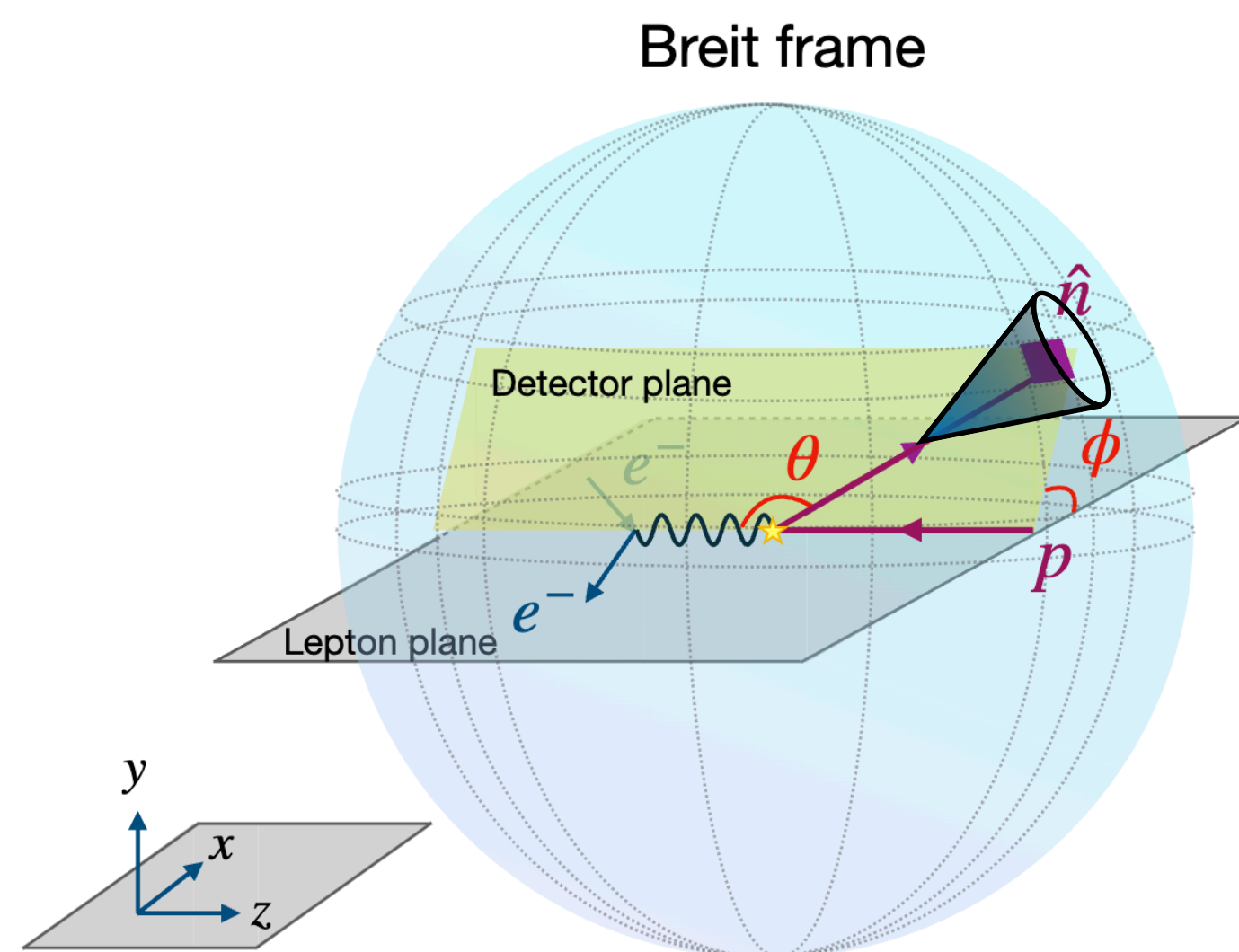
$$x_B = 3 \times 10^{-3}, Q^2 = 25\text{GeV}^2, \sqrt{s} = 105\text{GeV}$$



Liu, Liu, Pan, Yuan, Zhu '23

MOMENTUM IMBALANCE OF ENERGY FLOW

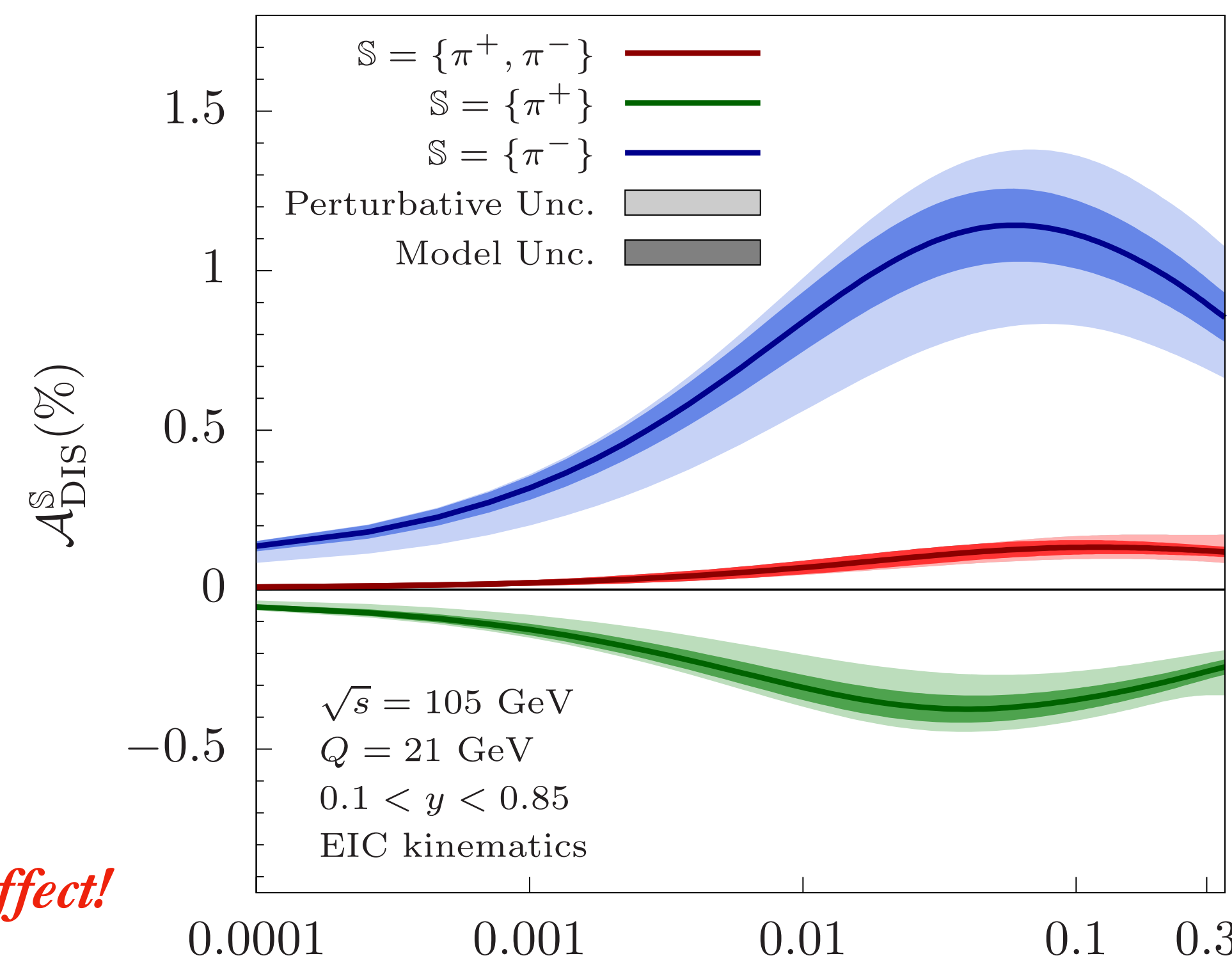
- One can further measure azimuthal angle dependence of the momentum imbalance of jets in order to study spin correlations



$$H_1^\perp = \text{Collins}$$

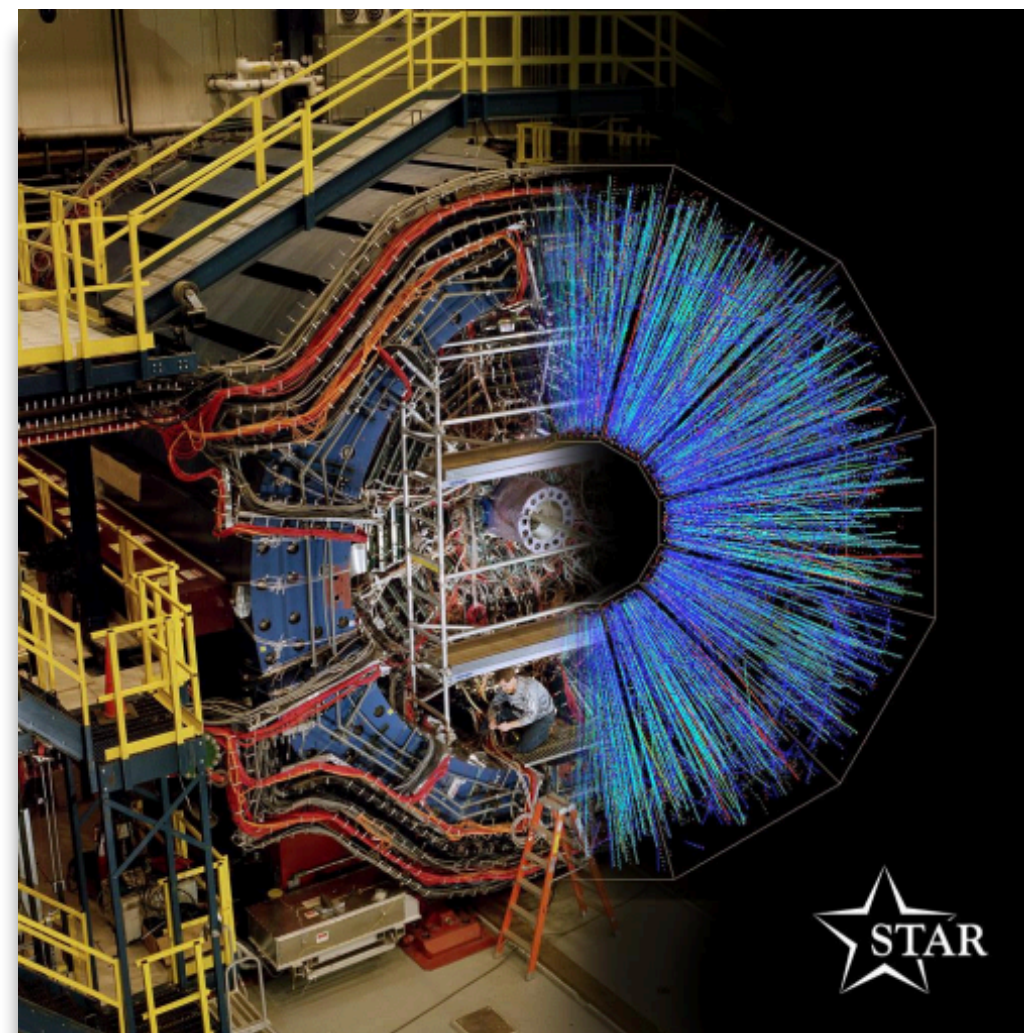
$$D_1 = \text{Unpolarized}$$

EIC can measure this collins effect!



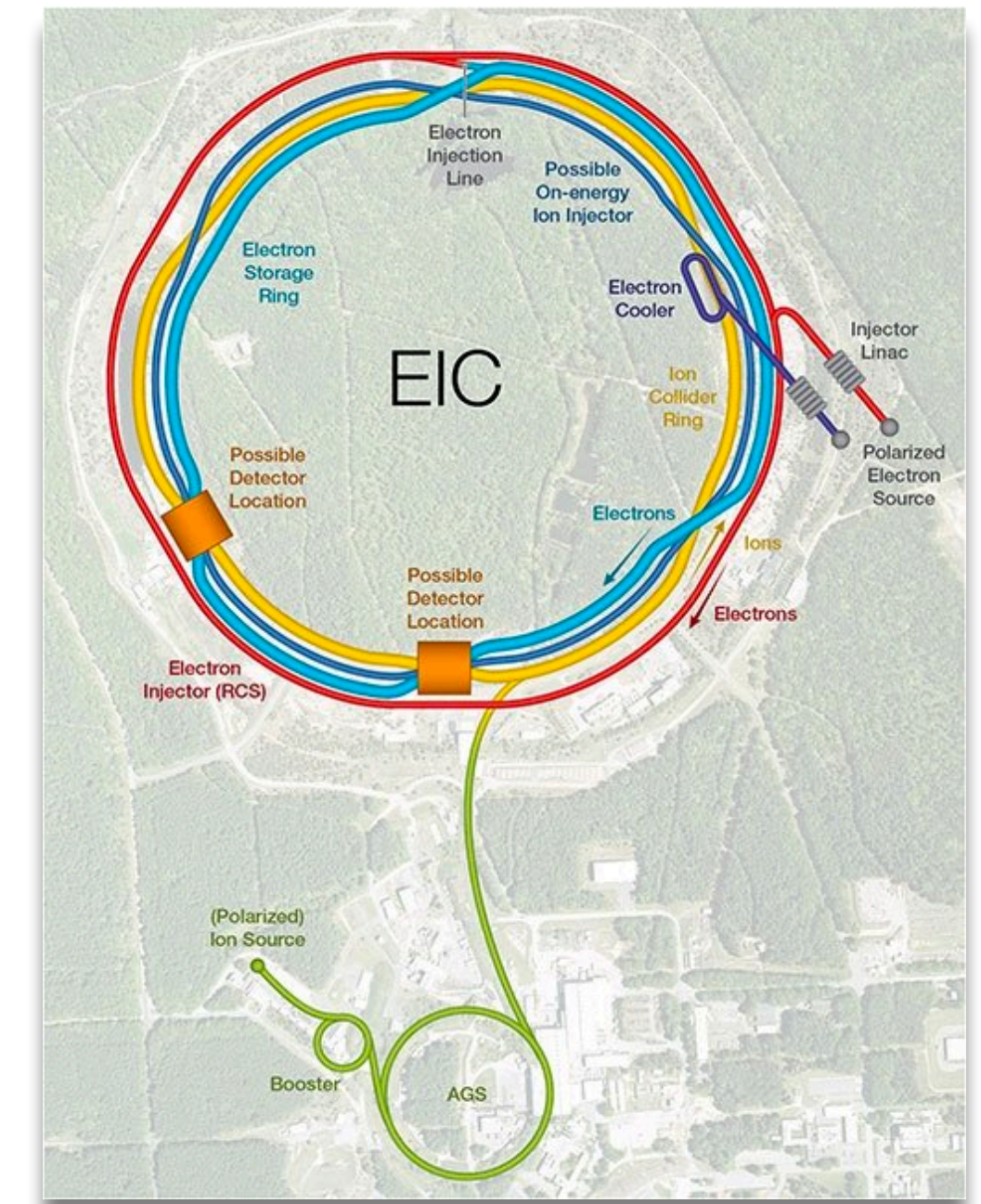
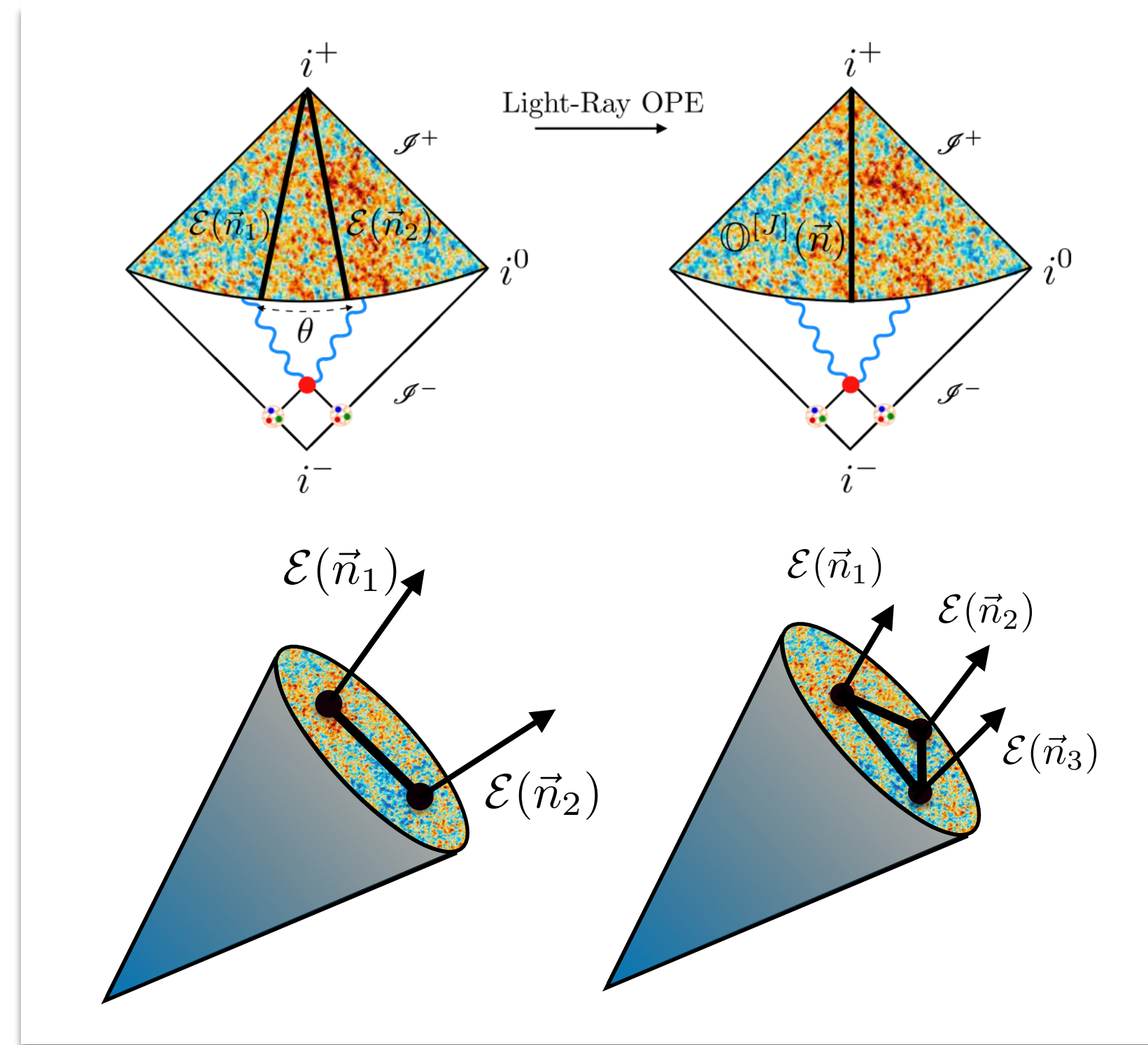


LHC 



RHIC 

Conformal Colliders meet Jets in Particle Colliders!

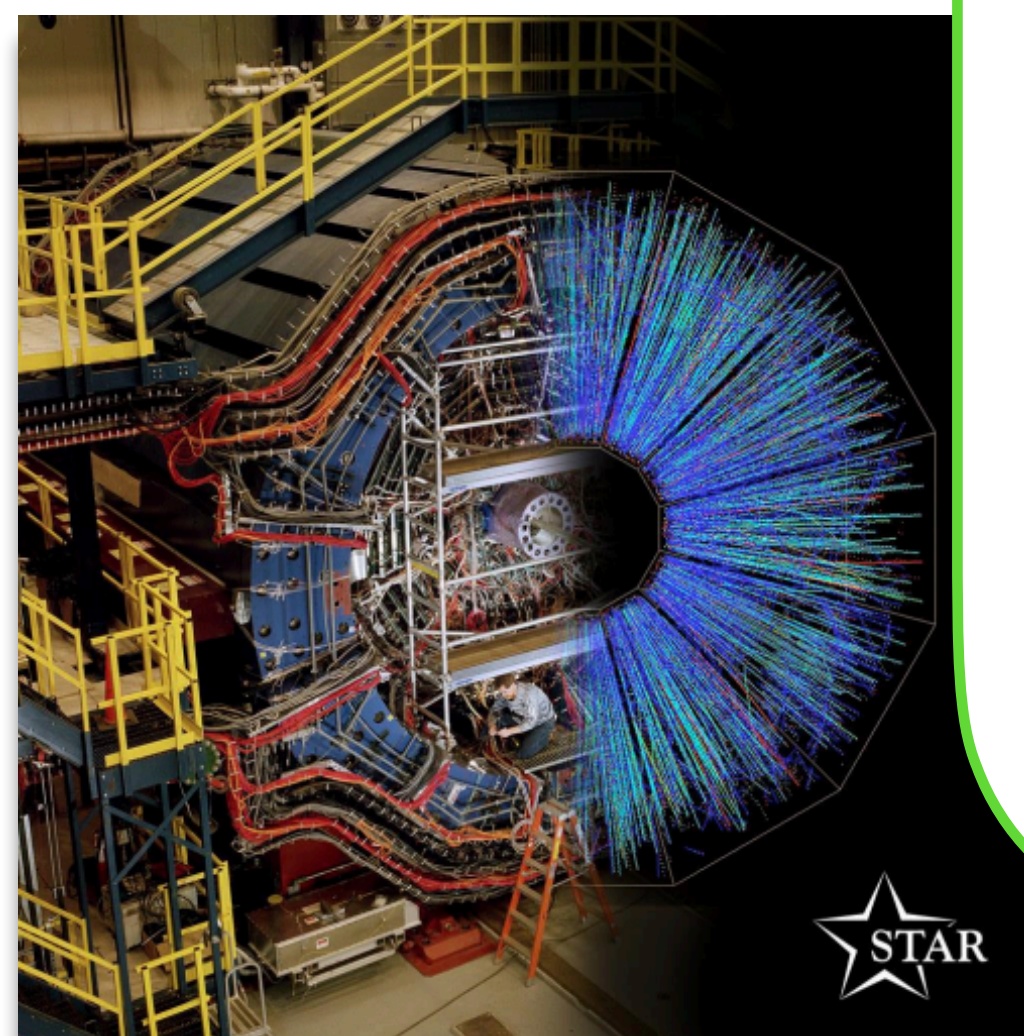


 **EIC**

Jets provide sharp link between underlying field theory and real world!



LHC



RHIC

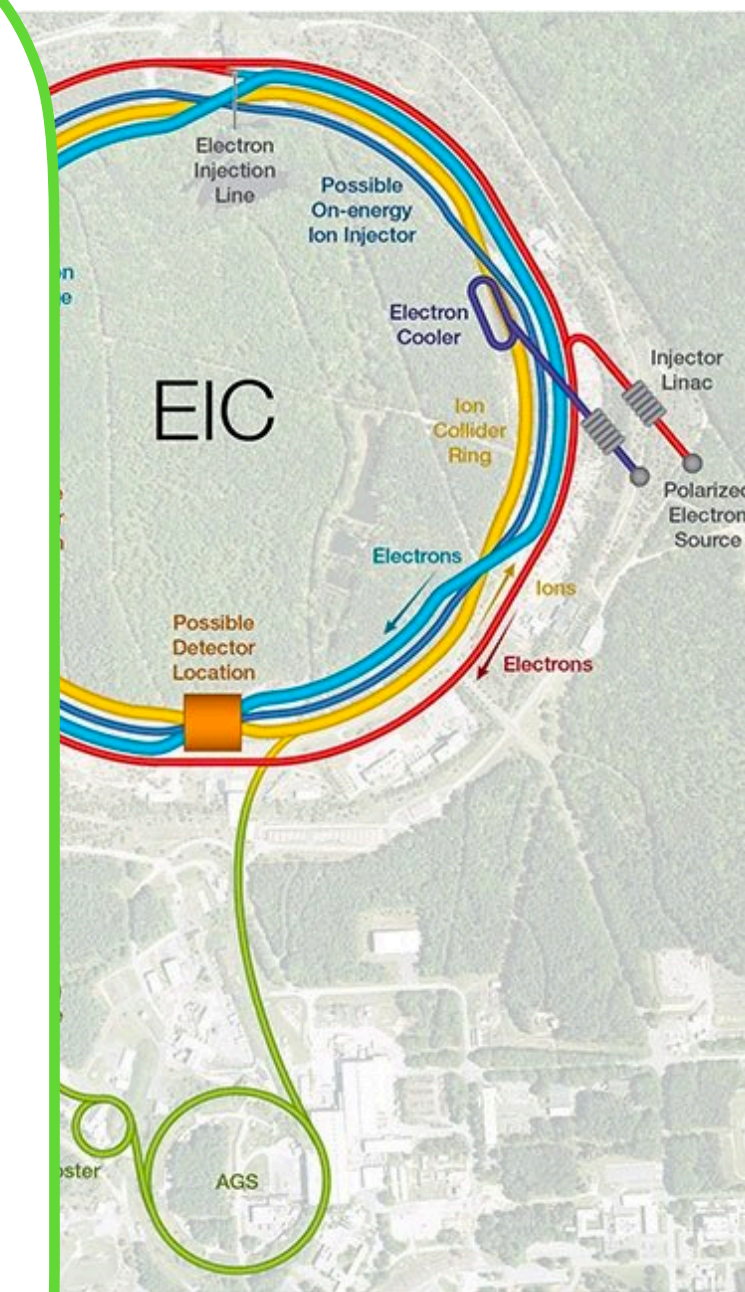


Conformal Colliders meet Jets in Particle Colliders!

Jets form a **universal language**, uniting studies across all colliders.

They create a **bridge** between the underlying field theory and the real world!

Jets provide sharp link between underlying field theory and real world!



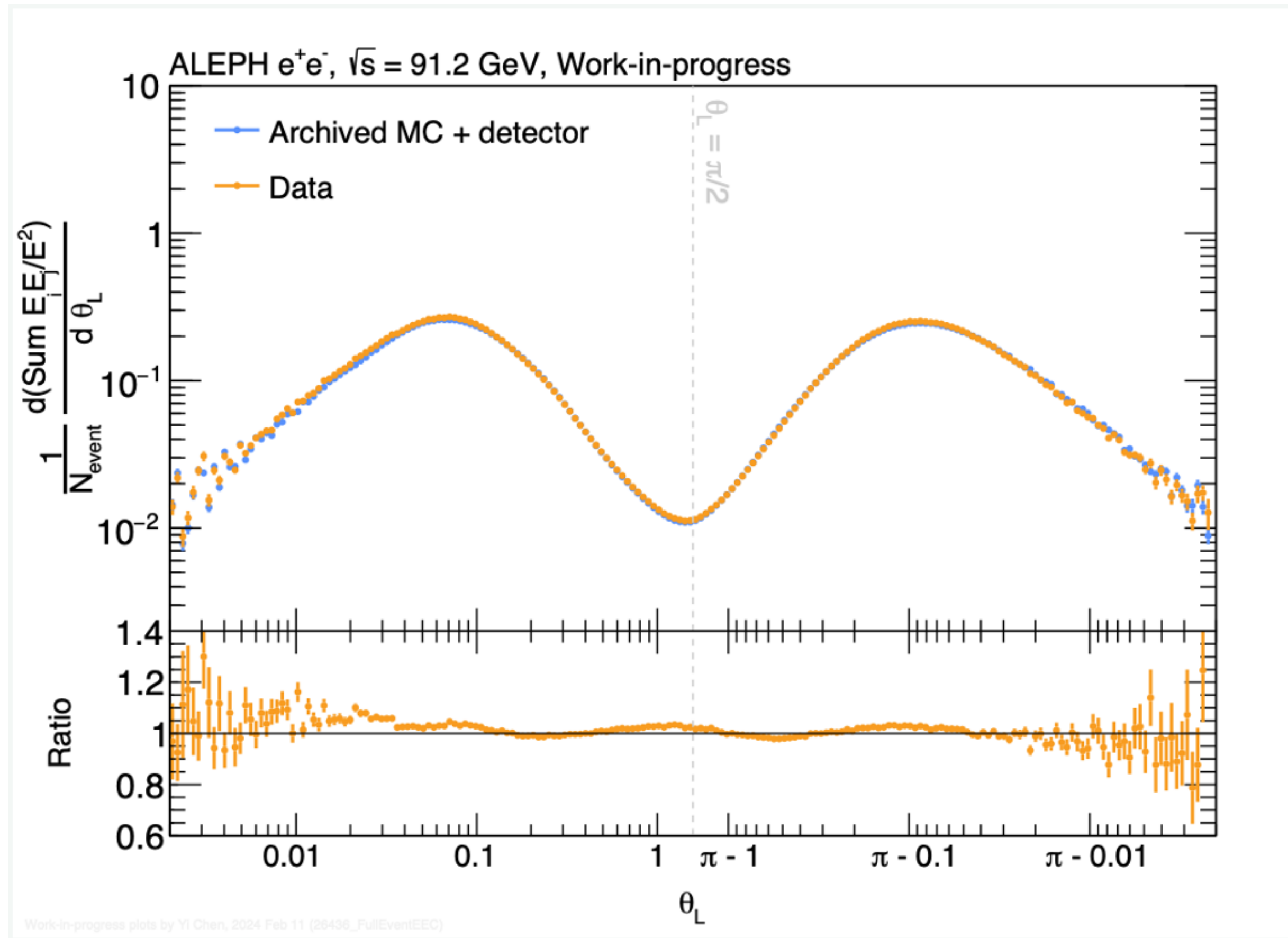
EIC

Backup slides

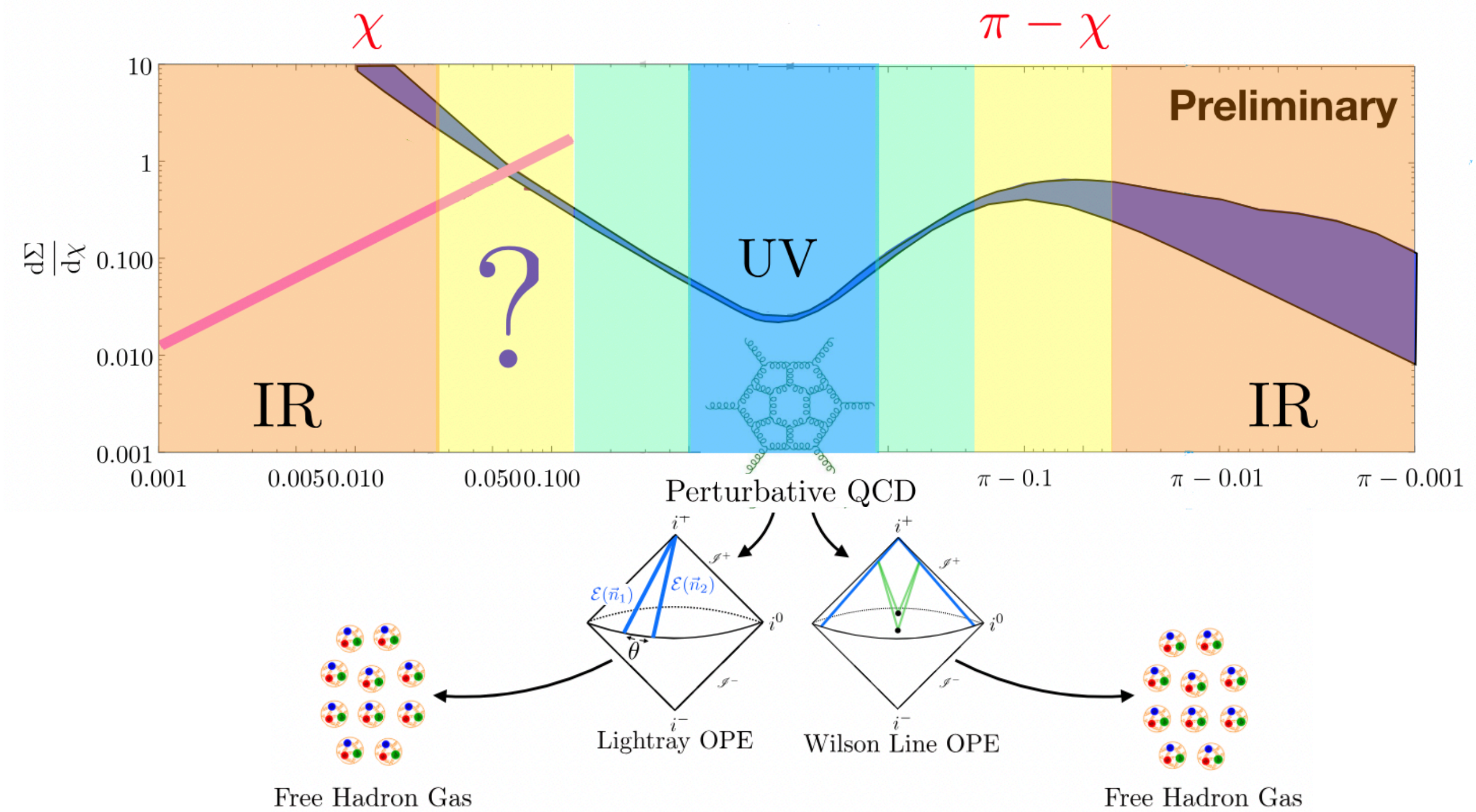
ENERGY CORRELATORS ON TRACK

1. Measurements on Tracks
2. Power corrections
3. Improved perturbative accuracy

Reanalysis of ALEPH data on tracks



Jaarsma, Li, Mout, Waalewijn, Zhu `In Progress

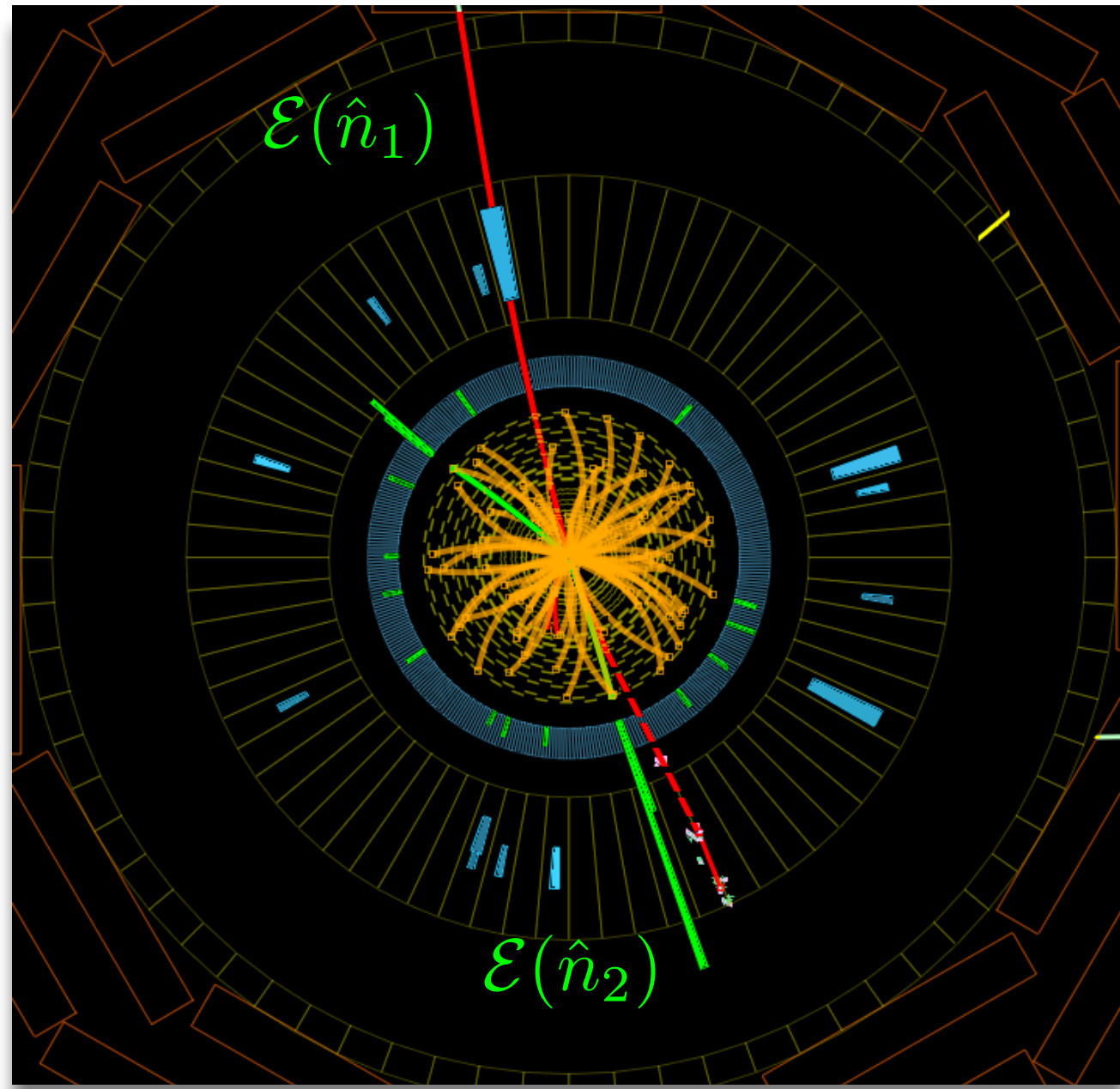


EEC on track for e^+e^- allows one to study event-wide correlations very precisely!

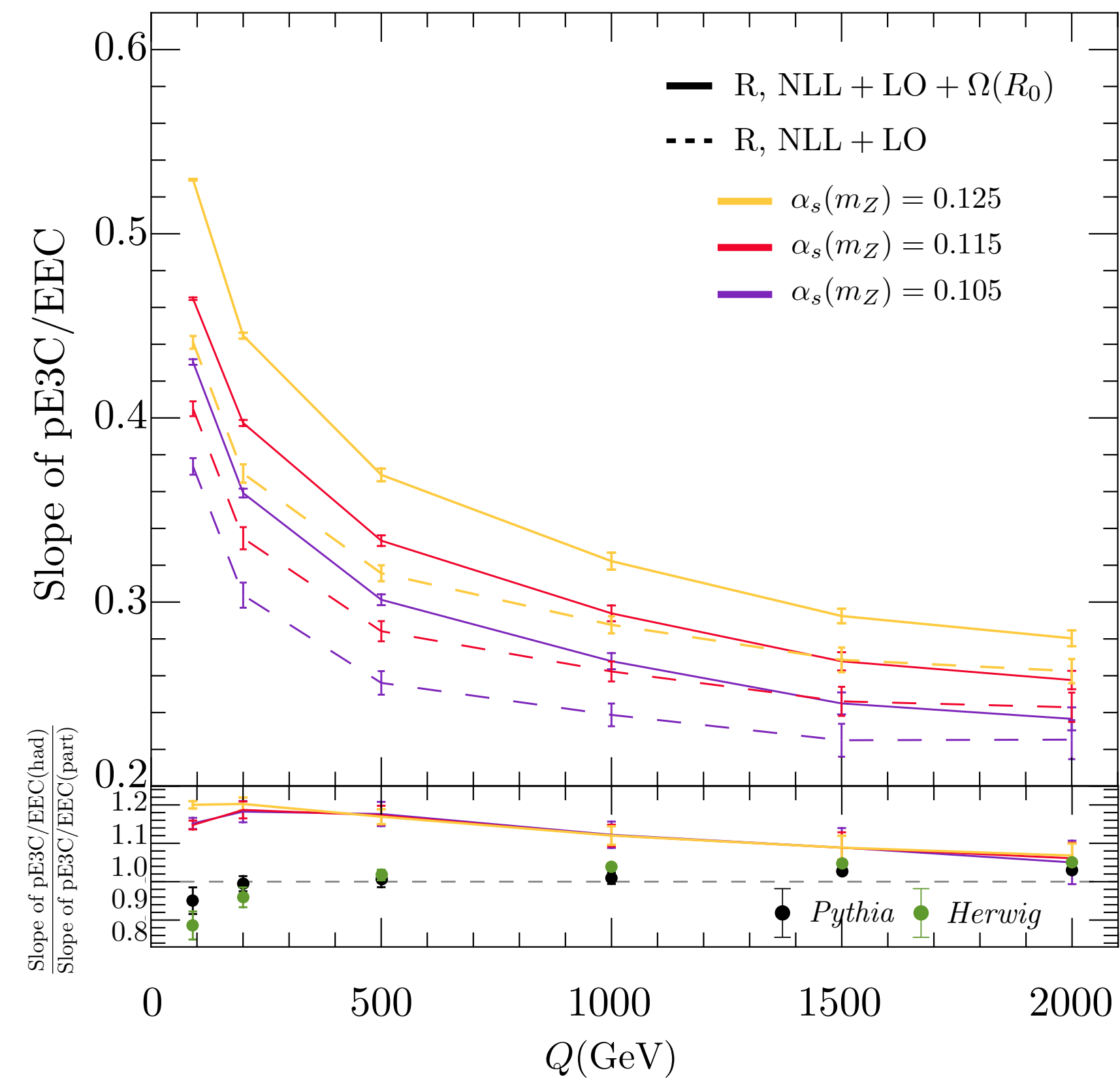
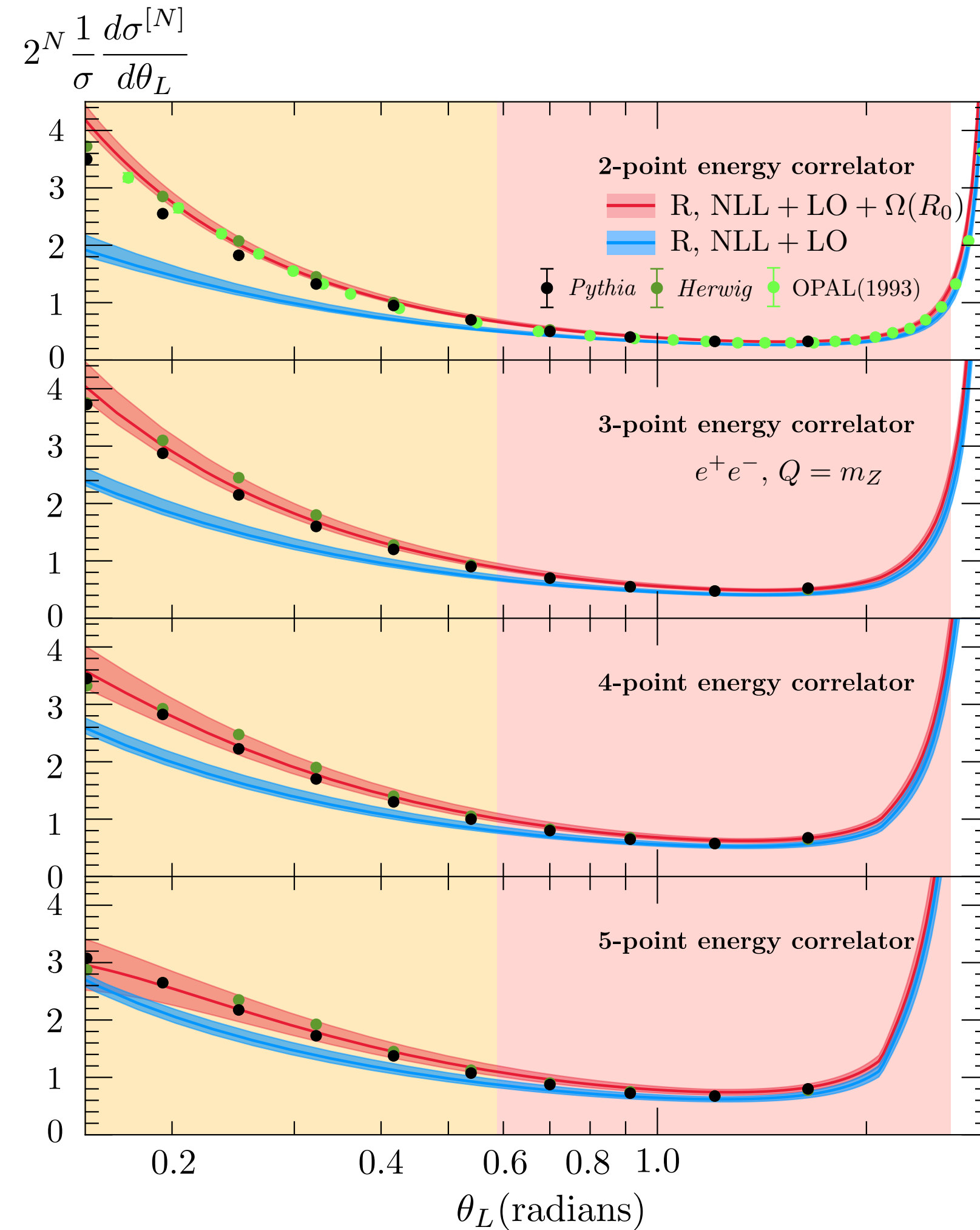
POWER CORRECTIONS

1. Measurements on Tracks
2. Power corrections
3. Improved perturbative accuracy

Schindler, Stewart, Sun '23
 KL, Pathak, Stewart, Sun '24
 Chen, Monni, Xu, Zhu '24



e^+e^- in the collinear limit exhibits same universal behavior as hadron jets



At $Q=1000$, 10% impact of power correction

$$\frac{1}{\sigma} \frac{d\sigma^{[N]}}{dx_L} = \frac{1}{\sigma} \frac{d\hat{\sigma}^{[N]}}{dx_L} + \frac{N}{2^N} \frac{\bar{\Omega}_{1q}}{Q (x_L (1-x_L))^{3/2}}$$

Universal Power Corrections

IMPROVING PERTURBATIVE ACCURACY

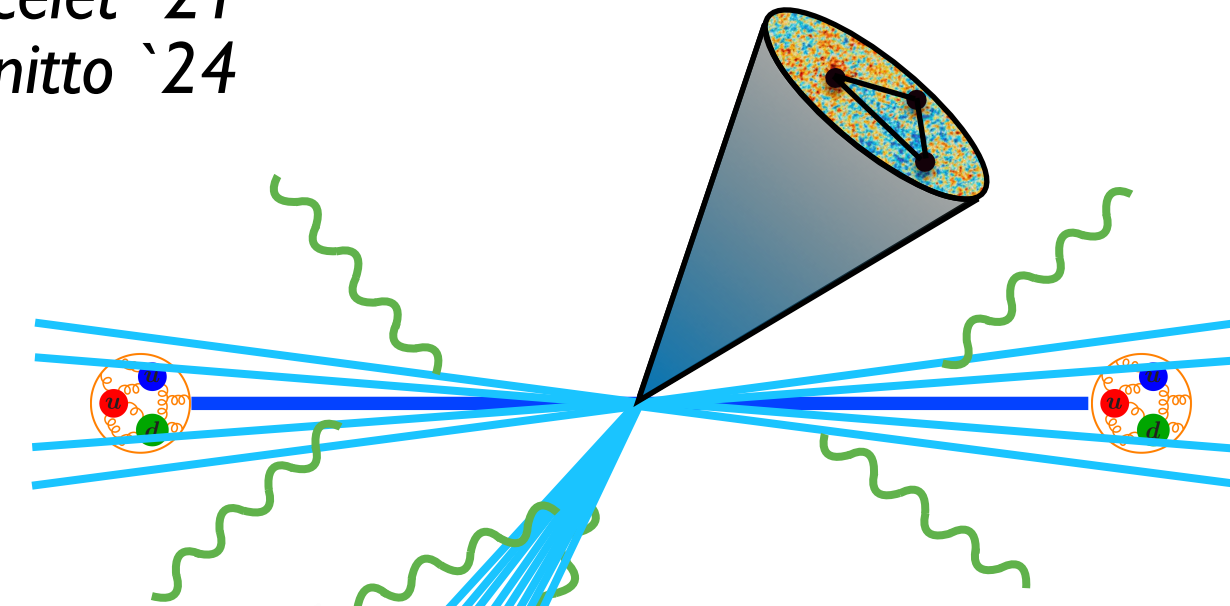
1. Measurements on Tracks
2. Power corrections
3. Improved perturbative accuracy

$$\langle \psi | \mathcal{E}(\vec{n}_1) \cdots \mathcal{E}(\vec{n}_{J-1}) | \psi \rangle$$

Czakon, Generet, Mitov, Poncelet '21
Bonino, Gehrmann, Stagnitto '24

$$\frac{d\sigma^{pp \rightarrow \text{jet}(\mathbf{N}\text{-proj})X}}{dp_T d\eta d\theta_L} = \sum_{a,b,c} f_{a/A} \otimes f_{b/B} \otimes H_{ab}^c \otimes \mathcal{G}_c^{\mathbf{N}\text{-proj}}(\theta_L)$$

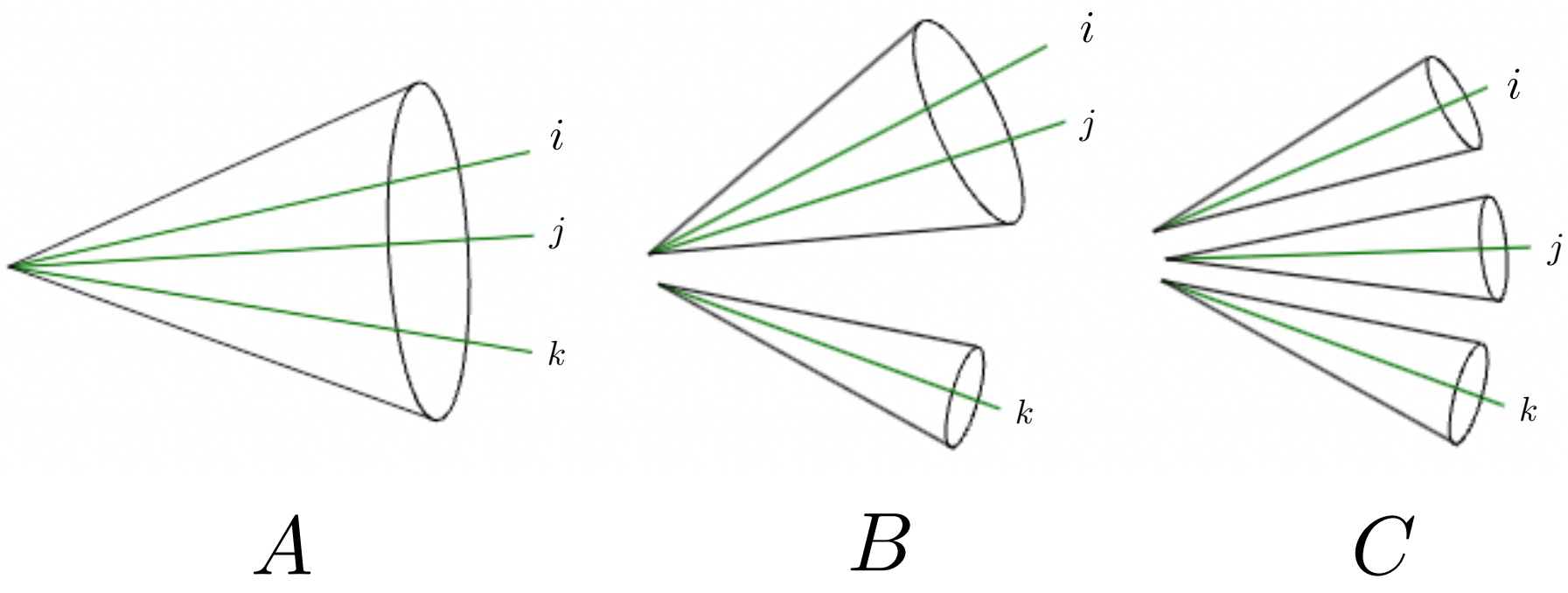
Λ_{QCD} p_T $p_T R$ $p_T \theta_L$



KL, Meçaj, Moutl '22
Kang, KL, Zhao '20
KL, Moutl, Zhang 'In Progress

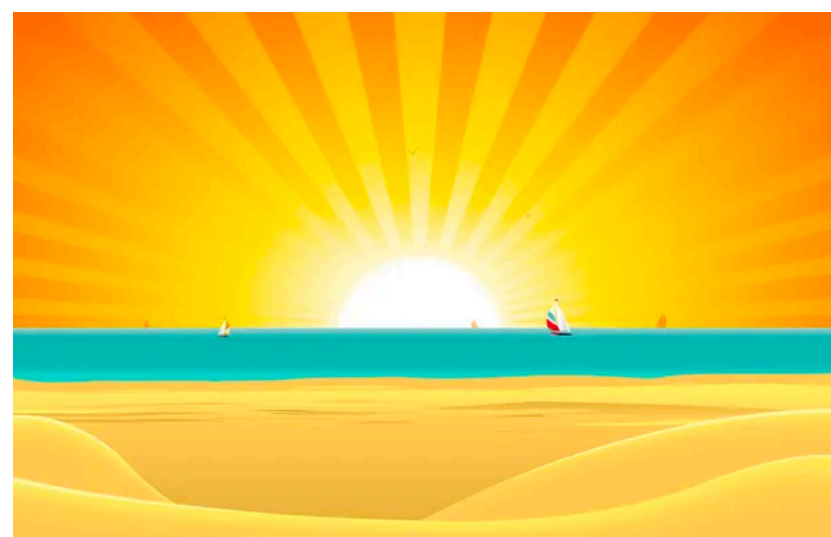
$$\mathcal{G}_c^{\mathbf{N}\text{-proj}}(z, R_L, p_T R, \mu) = \sum_j \int_0^1 dx x^N \mathcal{J}_{ij}(z, x, p_T R, \mu) \underbrace{J_{\text{EEC}}^{\mathbf{N}\text{-proj}}(R_L, x, \mu)}$$

Encodes complicated jet clustering algorithm details



\tilde{s}_{ik} = angle between i and k

z_i = momentum fraction of i



➤ **Unprecedented precision calculation of jet substructure on the horizon!**

