HIGHER ORDER ELECTROWEAK **RADIATIVE CORRECTIONS USING COVARIANT APPROACH**



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July 22, 2024 Mahumm Ghaffar





2024 EIC User Group Early Career Workshop

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Motivation



- Motivation
- Introduction to covariant approach



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- Full electroweak e^-p scattering with tree level, NLO and NNLO level graphs



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- distinguishable target particle

Parity violating asymmetry calculations via leptonic/hadronic tensor for



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Results

Parity violating asymmetry calculations via leptonic/hadronic tensor for



MOTIVATION

- make predictions that match experiments to one part in ten billion.
- related to Higgs boson etc.
- no concrete evidence of BSM at 13 TeV centre-of-mass energy at LHC.
- directly accessible at existing high-energy colliders.

• The theory of Standard Model (SM) \rightarrow unifies Electromagnetic, Weak and Strong interactions \rightarrow can

• SM limitations \rightarrow don't include gravity, dark matter/dark energy existence, hierarchies of scale

• Theoretical door open for Beyond the SM (BSM) physics to be observed at TeV scale \rightarrow **but** till date

Low energy precision physics becomes important \rightarrow provides a way to reach mass scales not

• We are doing precision physics with full electroweak Parity Violating Asymmetry (A_{PV}) \rightarrow achieve by calculating the higher order corrections up to NNLO (α^4) using Covariant/ leptonic tensor approach.



PARITY VIOLATING ASYMMETRY

Formula:
$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$
, where:

- For QED, $|M_{\gamma R}| = |M_{\gamma L}|$, numerator contains just weak+electroweak cross terms.
- Denominator contains just QED terms as m_{Z} (90 GeV)> $m_{\rho-}(0.5 \text{ MeV})$

•
$$A_{PV} = \frac{|M_{ZZ}|_R^2 - |M_{ZZ}|_L^2 + |M_{\gamma Z}|_R^2 - |M_{\gamma Z}|_R^2}{|M_{\gamma \gamma}|_R^2 + |M_{\gamma \gamma}|_L^2}$$

$\sigma_R \propto |M_R|^2$ and $\sigma_L \propto |M_L|^2$

 $|M_{\gamma Z}|_L^2$



- particles in the loop.

• Elastic e^-p scattering is studied up to the NNLO level considering all SM

• A longitudinally polarized e^- scatters off an unpolarized proton target

Covariant Approach



- particles in the loop.



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 (M, p_2)

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Covariant Approach

 (M, p_2)







QED

$i\mathcal{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1)\frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1)$



QED

 $i\mathcal{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1)\frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1)$

 $\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$



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$$\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$$

 $F_1(t)$ and $F_2(t) \rightarrow$ Dirac and Pauli form factors depending on momentum transfer "t".



OFD

 $i\mathcal{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1)\frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1) \qquad i\mathcal{M} = \bar{u}(k_2)(-ie(a_v\gamma_{\mu} + a_p\gamma_{\mu}\gamma_5))u(k_1)\frac{-i}{t - M_z^2}\bar{u}(p_2)(-ie\Gamma^{\mu}_{Z-p})u(p_1)$

$$\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$$

 $F_1(t)$ and $F_2(t) \rightarrow \text{Dirac}$ and Pauli form factors depending on momentum transfer "t".



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 $i\mathcal{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1)\frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1) \qquad i\mathcal{M} = \bar{u}(k_2)(-ie(a_v\gamma_{\mu} + a_p\gamma_{\mu}\gamma_5))u(k_1)\frac{-i}{t - M_7^2}\bar{u}(p_2)(-ie\Gamma^{\mu}_{Z-p})u(p_1)$

$$\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$$

 $F_1(t)$ and $F_2(t) \rightarrow \text{Dirac}$ and Pauli form factors depending on momentum transfer "t".

 $\Gamma^{\mu}_{Z-p} = f_1^p(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}f_2^p(t) + g_1^p(t)\gamma^{\mu}\gamma_5$







OFD

 $i\mathscr{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1) - \frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1)$

$$\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$$

 $F_1(t)$ and $F_2(t) \rightarrow \text{Dirac}$ and Pauli form factors depending on momentum transfer t''.

WEAK

$$i\mathcal{M} = \bar{u}(k_2)(-ie(a_v\gamma_{\mu} + a_p\gamma_{\mu}\gamma_5))u(k_1)\frac{-i}{t - M_Z^2}\bar{u}(p_2)(-ie\Gamma_{Z-\mu}^{\mu})u(k_1)\frac{-i}{t - M_Z^2}\bar{u}(p_2)(-ie\Gamma_{Z-\mu}^{\mu})u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u(k_1)u$$

$$\Gamma^{\mu}_{Z-p} = f_1^p(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}f_2^p(t) + g_1^p(t)$$

 $f_1^p(t), f_2^p(t)$ and $g_1^p(t) \rightarrow$ weak electric, magnetic and axial vector form factors.







OFD

 $i\mathscr{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1) - \frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1)$

$$\Gamma^{\mu}_{\gamma-p} = F^p_1(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F^p_2(t)$$

 $F_1(t)$ and $F_2(t) \rightarrow \text{Dirac}$ and Pauli form factors depending on momentum transfer t''.

WEAK

$$i\mathcal{M} = \bar{u}(k_2)(-ie(a_v\gamma_\mu + a_p\gamma_\mu\gamma_5))u(k_1)\frac{-i}{t - M_Z^2}\bar{u}(p_2)(-ie\Gamma_{Z-\mu}^\mu)u(k_1)\frac{-i}{t - M_Z^2}$$

$$\Gamma^{\mu}_{Z-p} = f_1^p(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}f_2^p(t) + g_1^p(t)$$

 $f_1^p(t), f_2^p(t)$ and $g_1^p(t) \rightarrow$ weak electric, magnetic and axial vector form factors.

$$a_{v} = \frac{I^{3} - 2\sin^{2}\theta_{W}Q_{f}}{2\sin\theta_{W}\cos\theta_{W}}, \quad a_{p} = \frac{I^{3}}{2\sin\theta_{W}\cos\theta_{W}}, \text{ where}$$
$$Q_{f} = -1(e^{-}) \text{ and } I_{3} = -\frac{1}{2}.$$







LEPTONIC TENSOR AND INTRODUCTION TO COVARIANT APPROACH



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 First introduced by Bardin and Shumeiko in 1976 (Nuclear Physics **B127**) to extract the infrared divergence from the lowestorder bremsstrahlung cross section.



LEPTONIC TENSOR AND INTRODUCTION TO COVARIANT APPROACH

- First introduced by Bardin and Shumeiko in 1976 (Nuclear Physics **B127**) to extract the infrared divergence from the lowestorder bremsstrahlung cross section.
- Recently used by Afanasev et al. (Phys.) Rev. D 66) to calculate QED radiative corrections in processes of exclusive Pion electroproduction.











































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COVARIANT APPROACH WITH LEPTONIC-HADRONIC TENSORS

 The differential cross section of general lepton-lepton/hadron scattering can be obtained by:

 $d\sigma \sim L^{\mu\nu}L_{\mu\nu}$

• where $W_{\mu\nu}$ is the hadronic tensor which in case of elastic e^-p scattering:

$$W_{\mu\nu} = H_1 g_{\mu\nu} + H_2 p_{1\mu} p_{1\nu} + H_3 p_{2\mu} p_{2\mu}$$

where p_1 and p_2 are incoming and outgoing protons momenta. H_1 , H_2 , H_3 , H_4 , H_5 and H_6 are the hadronic structure functions which can be extracted from experimental data.

$$_{\nu}$$
 or $d\sigma \sim L^{\mu\nu}W_{\mu\nu}$

$P_{\nu} + H_4 p_{1\mu} p_{2\nu} + H_5 p_{2\mu} p_{1\nu} + H_6 \epsilon_{\mu,\nu,p_1,p_2}$



QED AND ELECTROWEAK HADRONIC COUPLINGS WITH FORM FACTORS

$$\begin{split} \Gamma^{\mu}_{\gamma-p}(q^2) &= ieCnp2\left(f2p\gamma^{\mu} + gp\gamma_L\gamma^{\mu}\omega_- + gp\gamma_R\gamma^{\mu}\omega_+ - \frac{f2p(p_1^{\mu} + p_2^{\mu})}{2m_p}\right) \\ \Gamma^{\mu}_{Z-p}(q^2) &= -ieCnp2\left(F2W\gamma^{\mu} + gpz_L\gamma^{\mu}\omega_- + gpz_R\gamma^{\mu}\omega_+ - \frac{F2W(p_1^{\mu} + p_2^{\mu})}{2m_p}\right) \end{split}$$

$$\Gamma^{\mu}_{Z-p}(q^2) = -ieCnp2\left(F2W\gamma^{\mu} + gpz_L\gamma^{\mu}\omega_{-} + gpz_L\gamma^$$

 $F2W = \frac{F2Vp - 4\sin^2\theta_W f2p}{4\cos\theta_W \sin\theta_W} \to \text{EW form factor}$

 $gp\gamma_L = gp\gamma_R = f1p(0) \rightarrow$ Electric form factor

$$Cnp2 = \left(\frac{\Lambda^2}{\Lambda^2 - t}\right)^2, \quad \Lambda = \sqrt{0.83} \ m_p^2$$

or
$$gpz_{(L,R)} = \frac{F1Vp - 4\sin^2\theta_W f1p \pm G1p}{4\sin\theta_W \cos\theta_W}$$

 $G1p = 1.267 \rightarrow$ Axial form factor

F(1,2)Vp = f(1,2)p - f(1,2)n



Full Electroweak Tree level Graphs

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Full Electroweak Tree level Graphs





Full Electroweak Tree level Graphs





Full Electroweak Tree level Graphs







Full Electroweak Tree level Graphs





Full Electroweak Tree level Graphs





TREE-LEVEL LEPTONIC TENSOR (α -ORDER)



 For tree-level upper part of the diagram (say e⁻p scattering), one can calculate leptonic tensor which is:

$$L^{0}_{\mu\nu} \propto 4\pi\alpha((l_1)g_{\mu\nu} + (l_2)k_{2\mu}k_{1\nu} + (l_3)k_{1\mu}k_{2\nu} + \dots)$$

where k_1 , k_2 are incoming and outgoing e^- momenta and $l_{1,2..}$ are tree level leptonic tensor structure functions.







































e











307 graphs

e



NEXT TO THE LEADING ORDER (NLO) LEPTONIC TENSOR (α^2 -ORDER)



• The NLO leptonic tensor can be obtained by multiplying tree-level upper diagram with the sum of oneloop level self energy (SE) and triangular diagrams.

$$L_{\mu\nu}^{NLO} = (m_1)g_{\mu\nu} + (m_2)k_{1\nu}k_{2\mu} + (m_3)k_{1\mu}k_{2\nu} + (m_4)k_{1\mu}k_{1\nu} + (m_5)k_{2\mu}k_{2\nu} + \dots$$

Where $m_{1,2,3...}$ are leptonic structure functions which depend on the momentum transfer "t" and written in terms of Passarino-Veltman integral functions. We used LoopTools Mathematica package to calculate them.

In total 307 graphs SE and triangular graphs.

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NEW RESULTS: OED AND ELECTROWEAK NNLO LEVEL LEPTONIC TENSOR (α^3 -ORDER)





Tree level, NLO and NNLO level A_{PV} for e^-p scattering versus θ_{CM} using QWEAK kinematics



 (e^-p) Tree level, NLO and NNLO level A_{PV} versus θ_{CM}

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Tree level, NLO and NNLO level A_{PV} for e^-p scattering versus θ_{CM} using QWEAK kinematics



 (e^-p) Tree level, NLO and NNLO level A_{PV} versus θ_{CM}

$$(E_{beam} = 1.16 \ GeV, \ \theta_{CM} = 14$$

---- Tree $A_{PV} \sim -298.9 \ p_{PV}$
---- NLO $A_{PV} \sim -239 \ pp_{PV}$
..... NNLO $A_{PV} \sim -230 \ p_{PV}$

QWEAK Measured

 $\sim -226.5 \pm 7.3$ (statistical) ± 5.8 (systematic) ppb

Phys. Rev. C 101, 055503





NLO AND NNLO LEVEL CORRECTION FACTORS

• The correction factors depend upon the scattering angle θ which appears in momentum transfer as

 $t = (k_2 - k_1)^2 =$

where,

 $p_{in} = p_{out}$

Corrected $A_{PV}\% = \left(\frac{Tree\ level}{1}\right)$

$$= -2 p_{in}^2 \left[1 - Cos(\theta)\right]$$

$$\left(\frac{A_{PV} - (NLO, NNLO)A_{PV}}{Tree \ level \ A_{PV}}\right) \times 100$$





NLO and NNLO level Corrected A_{PV} versus θ_{CM}

NLO and NNLO level Corrected A_{PV} Asymmetry with Correction factors in percentage for QWEAK kinematics

 $(\theta = 14.6^{\circ})$

NLO level Corrected $A_{PV} = 19.9 \%$ Quadratic level Corrected $A_{PV} = 22.5$ % --- Total Corrected $A_{PV} = 22.9\%$

Tree level, NLO and NNLO level A_{PV} for e^-p scattering versus θ_{CM} using P2 kinematics



 $(e^{-}p)$ Tree level, NLO and NNLO level A_{PV} versus θ_{CM}

 $(E_{beam} = 155 \ MeV, \ \theta_{CM} = 39.97^{0})$ --- Tree $A_{PV} \sim -95.6 \ ppb$ NNLO (Reducible Two loop) $A_{PV} \sim -71.6 \ ppb$



Tree level, NLO and NNLO level A_{PV} for e^-p scattering versus θ_{CM} using P2 kinematics



 $(e^{-}p)$ Tree level, NLO and NNLO level A_{PV} versus θ_{CM}

 $(E_{beam} = 155 \ MeV, \ \theta_{CM} = 39.97^{0})$ --- Tree $A_{PV} \sim -95.6 \ ppb$ NNLO (Reducible Two loop) $A_{PV} \sim -71.6 \ ppb$

P2 Proposed $A_{PV} \sim -67.34 \ ppb$



RESULTS:

- For completeness, work in progress to include soft and hard photon bremsstrahlung cross sections in the results.
- effects in electron-proton scattering.
- physics beyond the Standard Model at the precision frontier.

• Next goal is to consider the polarized proton target and study the A_{PV}

• We make predictions for the e^-p NNLO level radiative corrections. These theoretical predictions will be important for many experimental programs such as QWEAK, P2, EIC and MOLLER (background studies) searching for



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REFERENCES FOR BOX DIAGRAMS

[1] M. Gorchtein, Phys. Rev. C 73, 055201 (2006)

[2] Peter G. Blunden et al., Physical Review Letters 91(14)









Backup Slides

ELECTROWEAK LEPTONIC TENSOR STRUCTURE FUNCTIONS

• In case of tree level polarized e^-p scattering:

• With photon (γ) as a mediator \rightarrow **Five** leptonic tensors

 $g^{\mu\nu}$, $k_{2}^{\mu}k_{1}^{\nu}$, $k_{1}^{\mu}k_{2}^{\nu}$, $\epsilon^{s_{1}\mu\nu k_{1}}$, $\epsilon^{s_{1}\mu\nu k_{2}}$

where $s_1 \rightarrow$ helicity reference vector of the incoming electron.

• With Z boson or γ Z mixing \rightarrow **Eight** leptonic tensors

 $g^{\mu\nu}$, $k_2^{\mu}s_1^{\nu}$, $k_2^{\nu}s_1^{\mu}$, $k_1^{\mu}k_2^{\nu}$, $k_2^{\mu}k_1^{\nu}$, $\epsilon^{s_1\mu\nu k_1}$, $\epsilon^{s_1\mu\nu k_2}$, $\epsilon^{\mu\nu k_1k_2}$



NLO LEVEL AND QUADRATIC LEVEL LEPTONIC TENSOR STRUCTURE FUNCTIONS

- In case of one loop level polarized e^-p scattering:
- With photon (γ), Z boson or γ Z mixing \rightarrow **19** leptonic tensors
- In case of quadratic level polarized e^{-p} scattering:
- With photon (γ), Z boson or γ Z mixing \rightarrow **21** leptonic tensors

